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Rational Engineering of the Dielectric Properties of Thin Silica Layers with a Single Plane of AgNPs

C. Rigoudy^{1, 2}, K. Makasheva¹, M. Belhaj², S. Dadouch², G. Teyssedre¹, and L. Boudou¹

¹LAPLACE, Université de Toulouse, CNRS, UPS, INPT, 118 route de Narbonne, 31062 Toulouse Cedex 9, France ²ONERA – The French Aerospace Lab, 31062 Toulouse Cedex 4, France

Abstract- Interface properties of dielectrics are of outmost importance when dealing with phenomena as charge injection or secondary electron emission. When the dielectric thickness scales down to the nm range, a deviation from the bulk material properties occurs; therefore, a possibility to tune the surface properties comes by modifying dielectrics at its sub-surface. This paper focuses on the precision of plasma-based processes in order to elaborate thin silica dielectric layers with a single plane of silver nanoparticles (AgNPs) at its sub-surface and thus to control their behavior under electrical stress. An illustration of the modulated properties of such nanostructured dielectrics is given through current measurements and electron emission.

I. INTRODUCTION

The dielectric charging phenomenon, *i.e.* the ability of dielectric materials to trap electrical charges, leads to modification of the local electric field. It causes change in the electrical behavior, and on a longer term degradation of the dielectric properties of materials under electrical stress [1,2]. The dielectric charging phenomenon is twofold: (i) it lays down the principle of operation of many applications, among which non-volatile memories [3] and electret based microphones [4], and (ii) it is considered the main cause of electrostatic sticking in MicroElectroMechanical Systems (MEMS) [5], charge build-up and breakdown in insulations [2,6,7] or being at the origin of potential gradients that may generate ElectroStatic Discharges (ESD), thus damage electronic devices and limit their lifetime, especially in space environment [8]. The later example is closely related to another physical phenomenon, namely electron emission from dielectrics under irradiation. In general, the emission efficiency is extremely sensitive to the surface properties of irradiated materials (composition, surface state, etc.), since the secondary electrons are typically emitted from the first 10 nm beneath the surface of the irradiated material. Main solutions to avoid ESD in those cases are oriented to geometric optimization of the devices and to design of surface coatings.

Even though the required outcome is on macroscopic level the solution is often existing at nanoscale and can be achieved by rational engineering of the dielectric properties of materials [6]. The lately developed nanometric dielectrics, or so-called nanodielectrics, have shown a considerable progress on many of the global electrical parameters of nanocomposite dielectrics (permittivity, losses, low field conductivity, space charge, etc.) since the pioneer works in the early 90s [9]. The Total Electron Emission Yield (TEEY) from dielectric materials under irradiation is one of the global parameters to control, when the featured application is intended for space environment. Therefore, sustainable solution calls for fine tuning of the material composition and structure, and a thorough analysis of the dielectric charging phenomenon.

This work aims to explore the impact of metal inclusions (silver nanoparticles, AgNPs), embedded in thin silica (SiO₂) layers, on the physical mechanisms responsible for dielectric charging and electron emission, in order to modulate them. Our strategy relies on synthesis of thin stratified coatings and is based on the principle of combined effect from metal/dielectric interfaces. An illustration of the modified properties of such nanostructured dielectrics is given through current measurements and electron emission in space environment.

Elaboration of stratified dielectric thin films requires facilities/methods permitting strict control of alignment of the metal inclusions at a given nanometric distance beneath the surface. Chemical or physical-based methods, or combination of both, can reply to the above requirements, however with different accuracy. Although well adapted for technological transfer the chemical (wet) methods for synthesis present various drawbacks as wasted materials (solvents), lack of precision in the metal NPs plane depth, uncertainties of the NPs distribution in the plane or different types of ligands that may preferentially trap electrical charges. Combined elaboration methods like dry dipping followed by plasma deposition [10], or physical methods like gas aggregation source [11], ion implantation [12], or plasma-based processes as the one employed in this work [6,7,13], should be privileged because of the accuracy in the process and the lack of uncontrolled chemical species surrounding the metal NPs.

II. PLASMA PROCESS FOR SAMPLE ELABORATION AND CHARACTERIZATION METHODS

The targeted structures consisted in a single plane of AgNPs embedded at 10 nm beneath the surface of 200 nm thick SiO₂ layer deposited on mirror polished stainless still and intrinsic Si-substrates. The searched total thickness of this stratified structure was of 220 nm. The selection of stainless still substrates ($3.3 \text{ cm} \times 3.3 \text{ cm}$) was related to the possibility to ensure good electrical contact on the back side for the charge current and electron emission measurements.

To achieve performance of the dielectric structure, the AgNPs should be well aligned in a plane, small enough so that their surface to volume ratio remains large, well distributed in the plane, with density that ensures charge transport regime far from the electrical percolation limit, however with interparticle distance allowing charge transport through tunneling or trap-assisted processes.

A. Plasma process and real-time monitoring

The above defined nanostructured thin dielectric layers have been elaborated in a plasma process successfully combining in the same reactor sputtering of a metallic target (physical vapor deposition - PVD) and plasma polymerization (plasma enhanced chemical vapor deposition - PECVD) [6,7,13].

The plasma was sustained in an argon axially-asymmetric radio-frequency (RF, f = 13.56 MHz) capacitively-coupled discharge at low pressure (p = 5.3 Pa). The RF-driven top electrode (smaller, 10 cm in diameter) was placed in front of a bottom electrode (larger, 12 cm in diameter) with a gap distance of 3.5 cm. The bottom electrode and the walls were grounded. The discharge powered electrode was an Ag-made target to bear the silver sputtering. The axially-asymmetric design of this RF discharge induces a self-bias voltage V_{dc} on the powered electrode, controlling in that way the sputtering. The self-bias voltage is at the origin of the acceleration of argon ions towards the silver target to achieve bombardment. The size and density of the deposited AgNPs were controlled by the injected in the plasma power P = 80 W ($V_{dc} = -950$ V) and the sputtering time ($t_s = 5$ s) for a given pressure.

The SiO₂-base and cover layers were deposited in the same reactor in the plasma sustained by P = 120 W ($V_{dc} = -820$ V). Optimal parameters for silica deposits have been identified in a previous work for the very same reactor [6,13]. Argon was the vector gas (Ar_{flow} = 2.8 sccm). The precursors used for SiO₂ deposition were hexamethyldisiloxane (HMDSO, $C_6H_{18}Si_2O$) and O_2 ($O_{2flow} = 1.25$ sccm). The HMDSO was injected by pulses ($t_{on} = 3.1$ s and period T = 5 s) giving an average HMDSO flow of 0.248 sccm. The total pressure was of p = 7.9 Pa. The thicknesses of the two silica layers were controlled by the deposition time $(t_{d-SiO2-base} = 25 \text{ min and } t_{d-SiO2-base} = 25 \text{ min and } t_{$ $_{SiO2-cover} = 90$ s). The use of oxygen has two purposes: to bring oxygen atoms needed for the SiO₂-deposition and to ensure plasma combustion of the C-content by creating by-products like CO, CO₂, COH₂ in order to evacuate the carbon atoms originating from the HMDSO decomposition.

The plasma state was monitored at each step of the deposition by measurements of the self-bias voltage and by Optical Emission Spectroscopy (OES) with a Princeton Instrument Acton Advanced Sp 2500A in the range 500 - 600 nm with resolution of 0.2 nm.

B. Structural characterization methods

Structural characterization of the resulting samples has been performed to determine the thickness and composition of the plasma silica matrix as well as to obtain the AgNPs size, shape, density and distribution and the total thickness of the structure.

The thicknesses of SiO_2 layers were determined by spectroscopic ellipsometry using a Semilab ellipsometer SE-2000 platform in the range 250–850 nm. Bruggeman's model was used to analyze the recorded spectra.

The types of molecular bonds within the silica layers were determined by Fourier Transform InfraRed (FTIR) spectroscopy using a VERTEX70 in transmission mode in the range 400 - 4000 cm⁻¹ with resolution of 2 cm⁻¹. The

transmission mode was achieved due to the transparency of intrinsic Si-substrates in the infrared domain.

Transmission Electron Microscopy (TEM) analysis of the elaborated structures in plane and cross-section views was performed with a JEM 2100F microscope operating at 200 kV.

C. Characterization of the electrical behavior of the structures

The electrical properties of the nanostructured layers have been studied through measurements of the polarization currents under strong electric fields and of the TEEY under electronic radiation conditions.

The study of polarization currents under strong electric fields was performed through a protocol of successive and increasing voltages [14]. Each voltage step lasts for 1000 s and is followed by a depolarization time at 0 V during 1000 s. The voltage was applied between the metallic substrate and an upper Au-electrode (150 nm thick and 1 cm in diameter).

The TEEY measurement were realized under high vacuum (10^{-6} Pa) in the energy range 0 - 2000 eV, using a 2 keV Kimball electron gun, according to the protocol described in details elsewhere [15].

III. RESULTS AND DISCUSSION

A. Importance of plasma stability during sample elaboration

The plasma stability during deposition was real-time monitored by OES and measurements of the self-bias voltage. The intensity of the emission lines is proportional to the electron density and the electron temperature in the plasma and the self-bias voltage is related to the power absorbed in the plasma. Thus, these two indicators reflect the stability of plasma deposition process and consequently, the precision and quality of the deposited layers. Typically, stabilization of the plasma process was achieved 30 min after plasma ignition.



Fig 1. OES spectra recorded during: (a) plasma silica deposition and (b) deposition of the single plane of AgNPs. The spectra are normalized to the intensity of Ar-line at 549.6 nm.

OES spectra of a stable plasma are shown on Fig. 1. For the silica deposition (Fig. 1a) only Ar-lines can be observed in the spectrum. A stable intensity of the Ar-line at 549.6 nm is a prerequisite to start deposition onto the substrates. Another condition is the absence of Ag-lines. It is achieved when the silver target is completely covered by SiO_2 and no Ag-atoms are sputtered into the plasma. For the AgNPs-deposition however, the Ag-lines are clearly observed (Fig. 1b). To correlate the amount of sputtered silver we use as indicator the ratio between the intensities of Ag- (at 546.5 nm) and Ar- (at 549.6 nm) lines. The same intensity ratio $(I_{Ag(546.5nm)}/I_{Ar(549.6nm)})$ is a criterion for reproducibility of the AgNPs-deposition.

Stability of the emission line intensities and of the self-bias voltage are the two key parameters in the deposition process, especially for long deposition times, as they give possibility to associate the deposition time to the composition (in particular the silica layer quality) and the thickness of deposited layers.

B. Quality of the plasma deposited SiO_2 layers and quantitative analysis of the single layer of AgNPs

The homogeneity of the SiO_2 -layer thickness was evaluated by ellipsometry measurements after mapping the entire surface of the stainless still substrates (10.9 cm²). The thickness map gives less than 10% difference between the thinnest and the thickest point.

According to the above given plasma conditions and deposition times the SiO_2 base layer is found 187.1 nm thick and the thickness of the SiO_2 cover layer is 9.1 nm.



Fig 2. FTIR spectrum for a sample of plasma SiO₂. Insert: Deconvolution of the asymmetrical stretching peak and shoulder.

FTIR measurements were performed systematically in order to control the structural properties of the elaborated SiO₂ layers. The three characteristic peaks of Si-O-Si vibrations are detected on the FTIR spectrum of SiO₂ layers (Fig. 2). Indeed, the rocking at 457 cm⁻¹, the symmetrical stretching at 814 cm⁻¹ and the asymmetrical stretching at 1063 cm⁻¹, including the shoulder at 1184 cm⁻¹ are visible on the spectrum. The plasma silica FTIR spectrum differs from a thermal SiO₂ one for the slightly larger full width at half maximum (FWHM) of the asymmetrical stretching: 87 cm^{-1} for the plasma SiO₂ compared to 78 cm^{-1} for thermal SiO₂. Although the used HMDSO precursor contains carbon, no carbon-traces are detectable on the FTIR spectrum. Presence of only SiO₂associated bonds means that the plasma process conditions and stability are met to grow a dielectric layer with composition and structure very near to that of thermal silica.

TEM imaging allows a direct observation of the elaborated structure (Fig. 3). The total thickness of the structure is measured from the cross section view (Fig. 3a). It is found 206 nm. The base SiO_2 layer is of 182 nm and the SiO_2 cover layer is of 9.5 nm. These values confirm the ones found by spectroscopic ellipsometry. Moreover, they are very close to

the targeted layer thicknesses 190 nm and 10 nm, respectively. These results confirm the importance of control during the plasma elaboration, and validate the chosen model to analyze the ellipsometric data. The insert in Fig. 3a shows well that for very thin SiO₂ cover layers (< 10 nm) the sample surface topography presents slight waviness reproducing the shape of AgNPs, due to the conformal deposition. Better planarity of the surface is achieved when increasing the depth of the AgNPs plane or when decreasing the size of the nanoparticles.



Fig 3. (a) Cross-section view TEM image of the SiO₂/AgNPs sample. Insert: enlargement of the near surface part. (b) Plane view TEM image of the SiO₂/AgNPs sample, (c) Statistical distribution of the AgNPs' major axis.

The plane view (Fig. 3b) of the sample shows that the AgNPs are well distributed in the plane and are far from percolation. Statistics on the AgNPs size are shown on Fig. 3c. The AgNP-distribution in size is bimodal with major axis centered at 4.1 nm and 17.8 nm for smaller and larger NPs. The AgNPs are of prolate spheroid shape, with eccentricity factor of 0.38 and 0.44 for the small and the large AgNPs respectively, in agreement with previous work [6,13]. Dark Field TEM imaging allowed us to conclude that the AgNPs are polycrystalline. No oxidation of the AgNPs was detected. It means that the SiO₂ cover layer protect the AgNPs plane from oxidation, thus their metallic character is preserved.

TEM imaging approves the quality of the plasma elaboration process. This verification is mandatory in order to accurately evaluate the impact of the nanostructuration, *i.e.* the nanometer scale parameters on the global electrical properties.

C. Current density under high electric field

As discussed before the dielectric charging leads to numerous complications over time, over use or under irradiation. Inclusion of AgNPs is a way to modify the electrical properties of a material, however preserving its dielectric behavior. The metallic nature of AgNPs increases the apparent relative permittivity of the nanocomposite. However, the AgNPs can also behave as artificial charge traps, which once occupied generate an internal electric field that opposes to the charges injection [7]. Comparison of the current density as a function of the electric field for a SiO_2 layer and the presented above nanostructured layer is given in Fig. 4. The precision in the elaboration process allowed us to compare the global response of the SiO_2 dielectric layers that differ at nanoscale level. One of the main differences found in conductivity data is that the current density at low field is 4 orders of magnitude higher with AgNPs than without.



Fig 4. Current density vs. electric field of thin SiO₂ layer (black triangle) and the SiO₂/AgNPs nanostructure (red points). The open red points represent instabilities in the current density measurements.

Present results show that a single plane of AgNPs within a plasma silica matrix deeply impacts the transport, trapping and release of electrical charges, which depends among others on the distance between the AgNPs plane and the sample surface. The influence on the electrical properties of the AgNPs size, and the AgNPs plane depth will be detailed in coming works.

D. Modifications in the Total Electron Emission Yield

Under irradiation, the TEEY is globally higher for dielectrics, being around 3 for a bulk SiO_2 , compared to metals for which it is around 2. Fig. 5 shows the TEEY of the thin SiO_2 layer alone and the nanostructured one with embedded at 9.5 nm AgNPs.



Fig 5. Total Electron Emission Yield of thin SiO₂ layers with AgNPs (red points) and without (black triangle).

The atypical shape of the TEEY is due to the small thickness of the dielectric layer [15]. However, one can see that addition of a single plane of AgNPs lowers the TEEY curve toward metal-like sample at low energy (Fig. 5). The global shape remains the same, the specific behavior of the thin dielectric is preserved. With this structure, the lowering can be modulated by varying the size as well as the depth of the AgNPs plane.

IV. CONCLUSION

In this study we explored the capacity of plasma-based deposition method to elaborate thin nanocomposite dielectric layers consisting of silica matrix with embedded single plane of AgNPs at a controlled distance from the surface. The precise control down to nm on the dielectric layer thickness, as well as on the dimensions of the AgNPs allows evaluation of the impact of nanoscale structuration on the global electrical properties of thin dielectrics and opens the avenue for rational engineering of the dielectric properties of thin silica layers.

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