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Exploring Phosphine Electronic Effects on

Molybdenum Complexes: A Combined

Photoelectron Spectroscopy and Energy

Decomposition Analysis Study

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ABSTRACT. In organometallic chemistry, and especially in the catalysis area, accessing the finest

tuning of a catalytic reaction pathway requires a detailed knowledge of the steric and electronic

influence of the ligands bound to the metal center. Usually, the M-L bond between a ligand and a

metal is depicted by the Dewar-Chatt-Duncanson model involving two opposite interactions, a σ-

donor and a π -acceptor effect of the ligand. The experimental evaluation of these effects is essential

and complementary to in-depth theoretical approaches that are able to provide a detailed

description of the M-L bond. In this work, we present a study of LMo(CO)₅ complexes with L

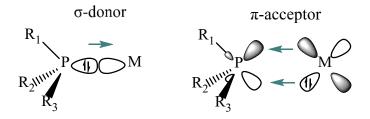
being various tertiary phosphines ligands by means of mass-selected high-resolution photoelectron

spectroscopy (PES) performed with synchrotron radiation, DFT and energy decomposition analyses (EDA) combined with the natural orbitals for chemical valence (NOCV) analysis. These methods enable a separated access of the σ -donor and π -acceptor effects of ligands by probing either the electronic configuration of the complex (PES) or the interaction of the ligand with the metal (EDA). Three series of PX₃ ligands with various electronic influence are investigated: the strong donating alkyl substituents (PMe₃, PEt₃ and PiPr₃), the intermediate PPh_xMe_(3-x) (x = 0-3) set and the $PPh_xPyrl_{(3-x)}$ set (x = 0-3 with Pyrl being the strong electron withdrawing pyrrolyl group C₄H₄N). For each complex, their adiabatic and vertical ionization energies (IEs) could be determined with a 0.03 eV precision. Experiment and theory show an excellent agreement, either for the IEs determination or for the electronic effect analysis. The ability to interpret the spectra is shown to depend on the character of the ligand. "Innocent" ligands provide the spectra the most straightforward to analyze whereas the "non-innocent" ligands (which are ionized prior to the metal center) render the analysis more difficult due to an increased number of molecular orbitals in the energy range considered. A very good linear correlation is finally found between the measured adiabatic ionization energies and the interaction energy term obtained by EDA for each of these two types of ligands which opens interesting perspective for the prediction of ligand characters.

Introduction

Nowadays, homogeneous organometallic catalysis has become a major tool in both academic and industrial fields. By surrounding the reactive metallic center, ligands play a crucial role through their electronic and steric properties.¹ They are indeed able to influence the net electron density on the metallic center and can therefore drive its behavior and reactivity. Custom-made ligands may thus lead to high selectivity in chemical transformation.² To access the finest tuning

of the catalytic reaction pathway, a detailed knowledge of the electronic influence of the ligands is required. Over the years, groundbreaking developments in theoretical and experimental approaches have made possible the increasing knowledge in transition metal coordination chemistry, including the effects of the ligands surrounding the metal center. The electronic interaction between a ligand L and a metal resulting in the M-L bond strength is commonly described by the Dewar Chatt Duncanson (DCD) model involving a σ -donor and π -acceptor capability of the ligand (see for instance the Scheme 1 for a phosphine ligand-metal interaction).¹, ³⁻⁵ Combination of these two effects leads to the global electronic effect of the ligand.



Scheme 1. Donor–acceptor interaction in transition metal complexes with phosphine as described by the DCD model.

Since several decades, many works have been devoted to the evaluation of these effects. The first and most known one refers to the seminal work of Tolman which defines the Tolman Electronic Parameter (TEP) descriptor based on the A_1 -symmetrical CO stretching frequency of $[(L)Ni(CO)_3]$ complexes.⁶ The TEP gives an experimental evaluation of the π -acceptor and net donor contribution of the ligands. From there, computational approaches were also developed to propose new descriptors issued from the TEP, as the calculated electronic parameter (CEP),⁷ the semiempirical electronic parameter (SEP)⁸ or a mode-decoupled local TEP (LTEP).⁹ These developments could show the great benefit brought by theory on this topic, but also how valuable complementary experimental and theoretical works are to describe ligands effects. Among the

various theoretical approaches that have been developed to describe the metal-ligand bonding and successfully applied for the description of ligand electronic effects, ¹⁰ one can for instance cite the energy decomposition analysis (EDA) combined to the natural orbitals for chemical valence (NOCV) analysis¹¹⁻¹⁴ or the natural bond orbital (NBO) approach.¹⁵⁻¹⁷ Using quantum mechanics/molecular mechanics (QM/MM) calculations, Suresh and coworkers were also able to determine molecular electrostatic potential (MESP) parameters from which electronic and steric descriptors can be obtained. 18-19 All these methods perform well but it is still interesting to confront them to experimental data to evaluate their performance. Experimentally, several approaches in addition to Tolman's have indeed been proposed to probe these electronic effects. Spectroscopy (NMR, 20-21 photoelectron spectroscopy, 22-23...) is one of the method of choice as it enables to directly survey the electronic and/or vibrational structure of the studied systems. For instance, Mampa et al. have demonstrated that it was possible to distinguish σ and π contributions for phosphine and phosphite ligands by joining IR and NMR data of iron-based complexes.²¹ This approach was successfully applied also for weak π -acceptor ligands such as P(4-XC₆H₄)₃. Ligand effects have also been probed by photoelectron spectroscopy (PES).²⁴ Lichtenberger was among the firsts to introduce this method in the field of transition metals in the 70's. 25-26 This technique enables to probe the electronic configuration of a system by measuring its ionization energies which are related to discrete energy levels of molecular orbitals.²⁷ According to the DCD model, the ligand electronic effect is due to the interaction of the metal orbitals with those of the ligand. PES is thus able to provide direct and critical information on the molecular orbitals (MOs) of the complex and therefore of both the σ -donor and the π -acceptor character of the ligand. Using a molecular orbital interaction scheme, Lichtenberger was able to interpret and analyze the electronic effect influence on numerous ligands, including phosphines. ^{25-26, 28-33} Another approach,

the quantitative analysis of ligand effects (QALE),³⁴⁻³⁵ consists in the analysis of several experimental thermochemical and spectroscopic data of the complex and the ligand to extract electronic parameters describing the stereoelectronic properties of the ligand. The QALE approach has demonstrated its ability for separating the σ and π contributions of the ligand, but it suffers of few limitations. For instance, the weak π -acceptor ligands are not well described with this methodology.³⁶

From these perspectives, it could be interesting to focus on descriptors that are able to build a bridge between experimental and theoretical approaches, i.e. that benefit from the in-depth analysis provided by theory but that can also be confronted to experimental data which are essential. One such example is provided by Coll et al. who evidenced the linear correlation existing between local ionization energies and TEP, providing thus a high-performance theoretical method to evaluate TEPs.³⁷ The study herein is based on such objective. As described above, PES has already been shown to be able to provide information on ligand electronic effects. To this end, the interpretation of the spectra was done using MOs interaction scheme which could be enriched today with the development of all the powerful quantitative theoretical methods. Our aim in this project was thus to confront PES analysis to a modern chemical bond description approach. In this perspective, we have performed a systematic study on LMo(CO)₅ complexes with L = CO and three series of L tertiary phosphine ligands, namely PR₃, with R = alkyl substituents, $PPh_xMe_{(3-x)}$ (x=1-3) and $PPh_xPyrl_{(3-x)}$ (x=1-3 and Pyrl being a pyrrolyl group (C₄H₄N)) (Scheme 2), using a state-of-the-art photoelectron spectroscopy setup combined with DFT and EDA/ETS-NOCV calculations. Our objective was to determine to which extent this experimental method may be correlated to theoretical data. The phosphorus(III) ligands were chosen because they belong to one of the most used ligand family in catalytic processes³⁸⁻⁴¹ and several experimental and/or theoretical works have already been devoted to the study of their electronic influence on metals.^{37, 42-47}

$$OC \sim CO$$

$$CO$$

$$CO$$

$$L = PPh_xMe_{(3-x)}$$

$$L = PPh_x Pyrl_{(3-x)}$$

Scheme 2. Structures of the complexes studied in this work.

After a description of the experimental and theoretical approaches, experimental results will be presented for each series of ligands and compared to theoretical data. A correlation between the ionization energies of the complexes and the interaction energy obtained from EDA is found and discussed at the end of this work.

Methods

Synthetic procedure

Molybdenum complexes were either commercially available (molybdenum hexacarbonyl, Sigma-Aldrich, St Quentin Fallavier, France) or synthetized from Mo(CO)₆ and the corresponding tertiary phosphines upon treatment with trimethylamine oxide in acetonitrile. Characterization details of the complexes are given in the Supplementary Material.

Photoelectron spectroscopy using synchrotron radiation

Experiments were carried out on the DESIRS beamline at the SOLEIL synchrotron facility in St. Aubin (France). 48 This undulator-based 49 VUV-beamline is equipped with a 6.65 m normal incidence monochromator 50 proposing several gratings depending on the required flux/resolution compromise. A 200 gr/mm grating was used, offering a photon flux in the 10^{12} - 10^{13} ph/s range with a moderate energy resolution (12 meV at 10 eV photon energy). The slits were set to deliver a photon energy bandwidth between 5 and 15 meV at 9 eV and the second order light was removed with an upstream gas filter filled with 0.20 mbar of Kr or 0.15 mbar of Xe (depending on the energy range considered). 51 Absolute energy calibration was obtained by using the $4s^24p^6 \rightarrow 4s^24p^5(^2P_{1/2})6s$ absorption line of Kr. 52

Ca. 1 g of compound was deposited on a glass wool nest in a home-made oven installed in the SAPHIRS vacuum chamber. ⁵³ The temperature was adjusted for each molecule, from 25°C up to 160°C. The generated vapor was transported with the help of 0.5 to 1 bar helium flow and expanded through a 70 μ m nozzle. After traversing a $\emptyset = 1$ mm double skimmer, the molecular beam is crossed at a right angle with the synchrotron radiation inside the interaction region of the versatile DELICIOUS III electron/ion imaging coincidence spectrometer. ⁵⁴ Electrons and ions

produced in the photoionization process were analyzed in coincidence by two velocity map imaging (VMI) analyzers. Photoelectrons were collected using a 176 V/cm extraction DC field corresponding to photoelectrons kinetic energies up to 6 eV with an energy resolution of $\Delta E/E =$ 5%. Typically, the 6.6-11 eV energy range was scanned with 20 meV or 40 meV energy steps. After a mathematical treatment of mass-selected photoelectron images,⁵⁵ 2D-photoelectron matrixes are plotted as a function of the photon energy and of the electron kinetic energy. Extracting photoelectrons with kinetic energies comprised between 0 and 150 meV gives the slow photon electron spectra (SPES). Error bars shown on the SPES have been obtained by assuming an initial Poisson distribution on the image pixels and by applying standard error propagation formulas. Appearance energies (AEs) and vertical ionization energies (IEvert) are given with an estimated error of 30 meV (taking into account the scan steps, the precision in the kinetic energies and in the photon energy). The photon flux (photodiode AXUV100 from IRD) was recorded at each scan point for normalization purposes and hence, all spectra are corrected from false coincidences and photon flux variation but not from the oven signal variation (due to the consumption of the sample). Finally, the ionization band envelops were fitted using an asymmetric Gaussian following a similar procedure to the one developed by Lichtenberger. 25, 56

Theoretical method

Two types of calculations were performed in this work: DFT-based calculations to access the electronic structure of the studied complexes and energy decomposition analysis to describe the L-Mo(CO)₅ bond.

Electronic structure calculations were performed using the GAUSSIAN 09 package.⁵⁷ Geometry optimizations and frequency calculations were obtained with the density functional theory (DFT)-

based PBE0 method⁵⁸⁻⁵⁹ coupled with the triple-ζ Ahlrichs-type basis set def2-TZVP.⁶⁰ Vibrational analysis served to locate minima (no imaginary frequency). A full Natural Bond Orbital (NBO, v3.1) analysis was also performed to obtain in particular the localized partial charges and NBOs on the ligand. To compute energies, we have chosen to use the long-range corrected hybrid LCωPBE0 (also known as LC-ωPBEh) functional⁶¹ using an IP-tuned range-separation parameter (ω = 0.3 Bohr⁻¹) combined with the def2-TZVP basis set. The IP-tuning procedure has been shown to provide very accurate estimations of IEs and is discussed in detail in the work of Gallandi et al. 62 Adiabatic and vertical ionization energies were accordingly calculated with this method from the energy difference between the ion with an optimized geometry and the ground-state neutral's geometry, respectively, and the ground-state neutral complex. The energies of the highest Kohn-Sham (KS) orbitals (HOMOs) of the neutral complexes were also computed and served for the interpretation of the photoelectron spectra pattern. Although it is known that KS orbitals eigenvalues predict IE_{vert} with very large absolute errors, 63 Zhan et al. 64 and Gengeliczki et al. 65 nevertheless showed that they may be used in the same way as Hartree-Fock (HF) orbitals with Koopman's theorem by using a correction factor. In practice, the negative of the calculated KS orbitals energies have to be shifted so that the HOMO corresponds to the first vertical ionization energy on the SPES.

Finally, the interaction between the phosphine ligand and its Mo(CO)₅ counterpart was probed by an energy decomposition analysis (EDA) following a Morokuma-like energy partition of the bond, ⁶⁶⁻⁶⁹ using the Amsterdam Density Functional (ADF) program⁷⁰ at the BP86-D3BJ/TZ2P(ZORA) level of theory. ⁷¹⁻⁷⁴ The EDA scheme enables to partition the interaction energy ΔE_{int} between the fragments into three terms: i) the repulsive Pauli interaction, ΔE_{Pauli} , which includes the destabilizing interactions between occupied orbitals, ii) the classical

electrostatic interaction, ΔE_{elstat} , between the fragments and iii) the orbital interaction, ΔE_{orb} , which is associated to the relaxation and the mixing of the orbital between the fragments. Combining the EDA approach with the natural orbitals for chemical valence (NOCV) theory, ^{12, 68, 75} enables to further decompose qualitatively and quantitatively the ΔE_{orb} term into the different deformation density contributions of the bonding ($\Delta E_{orb,\sigma}$, $\Delta E_{orb,\pi}$...). NOCVs with the same absolute energy eigenvalues (and opposite sign) can be grouped together to describe charge-transfer channels between the molecular fragments. The visualization of these pairwise NOCVs allows their assignment as donating and backdonating processes.

All calculations details are given in the Supplementary Information.

Results and Discussion

Molybdenum hexacarbonyl was the first complex to be studied to serve as the reference compound. Many gas phase experimental or theoretical works have already been undertaken on this molecule and its valence photoionization process is well-described in the literature. The SPES of Mo(CO)₆^{+•} was recorded with an oven temperature at 60°C and is presented in Figure 1.

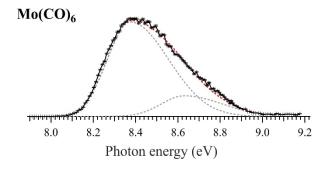


Figure 1. SPES of the Mo(CO)₆^{+•} ion recorded in the 7.9-9.2 eV energy range. Plain lines represent the experimental spectrum whereas dotted lines are the asymmetric Gaussian curves used to fit the data.

energy edge. In the 80's, Hubbard and Lichtenberger showed that this shoulder could be attributed to the vibrational progression of the v1 stretching mode of the ion. ⁷⁶ We have accordingly fitted the global band by taking into account this feature. Two asymmetric Gaussian curves have thus been necessary to describe the SPES, the first one related to the ionization of orbitals mainly with a d character (see below) and the second one to the v1 progression. Both are shown with dotted lines in Fig. 1 as well as the resulting envelops (red dotted line). This gives rise to an estimation of the vertical ionization energy of the complex: $IE_{vert}(Mo(CO)_6) = 8.37 \pm 0.03$ eV. This results agrees pretty well with the work of Hubbard and Lichtenberger⁷⁶ (8.45 eV) as well as other experimental^{45, 77} and theoretical⁷⁸⁻⁷⁹ ones. We also estimated the adiabatic IE of Mo(CO)₆ to be $IE_{ad}(Mo(CO)_6) = 8.13 \pm 0.03$ eV, very close from our calculated value of 8.16 eV and in satisfying accordance with the previous work of Chen et al. who measured a value of 8.21 ± 0.03 eV. 80 As already outlined, the band observed in the spectrum in Figure 1 is attributed to the ionization on the triply degenerate t_{2g} set of orbitals of the complex having a (Mo 4d + CO(π *)) character.⁷⁸ Changing one CO group by another ligand L decreases the symmetry of the molecule and leads to a splitting of the t2g set. One obtains then the doubly degenerate e and the b2 sets. This feature is visualized in our SPES and will be discussed for the different series of studied ligands. The first series of ligands studied is the PR_3 = alkyl set. SPES of $LMo(CO)_5^{+\bullet}$ ions with $L = PMe_3$, PEt₃ and PiPr₃ were obtained by heating the complexes from 50°C (PMe₃) to 95°C (PiPr₃) and are shown in Figure 2a)-c). In the 7.0 to 10.5 eV energy range, different noticeable features are present on the SPES as a function of alkyl substituents. The first observation from the SPES presented here concerns the measured adiabatic ionization energies which vary from 7.32 ± 0.03 eV for PMe₃

to 7.21 ± 0.03 for PEt₃ and 7.16 ± 0.03 eV for PiPr₃ within 0.06 eV agreements with calculated

A broad band is observed in the 8.0-9.0 eV region and exhibits a small shoulder on the high

Table 1. Experimental and calculated adiabatic and vertical ionization energies of the LMo(CO)₅ complexes (in eV).

	$\mathbf{IE_{ad}}$		IE _{vert}			
L	exp.a	calc.	exp. ^a this work	calc. this work	other works	
СО	8.13	8.16	8.37	8.45	see text	
PMe_3	7.32	7.26	7.55; 7.74; 9.84	7.53	7.64 (calc.) ⁸¹	
PEt ₃	7.21	7.17	7.45; 7.67; 9.52	7.46		
PiPr ₃	7.16	7.11	7.43; 7.67; 9.31	7.37		
PPh ₃	6.99	7.03	7.37; 7.57; ~8.7	7.30	7.30 (calc.) ⁸¹	
PPh_2Me	7.14	7.13	7.44; 7.62; ~9.1	7.38		
$PPhMe_2$	7.16	7.20	7.48; 7.67; ~9.2	7.44		
PPh ₂ Pyrl	7.25	7.33	~7.5; ~7.7	7.51		
$PPhPyrl_2$	7.42	7.50	not measured	7.69		
PPyrl ₃	7.59	7.67	not measured	7.95		

^a obtained with an estimated error of 0.03 eV

values (Table 1). These energies can be linked to the global interaction of the ligand, which includes orbital, steric, electrostatic effects. One can indeed expect that a stronger interacting ligand will tend to stabilize the metal within the complex and then increase its ionization energy. For instance, the interaction of occupied t_{2g} orbitals of the metal with low-lying empty π^* MOs of the ligand results in a bonding t_{2g} MO (d (metal) + π^* (ligand)), lower in energy than the initial t_{2g} orbital. In case the ejected electron is in the direct environment of the metal (i.e. the HOMO has a d (metal) character), it results thus in higher IEs. As the CO ligand shows a strong π -acceptor effect, it explains the high IE_{ad} of the Mo(CO)₆ complex (8.13 eV, Table 1). For the PR₃ set, the IE_{ad} trend indicates that the PMe₃ group is interacting more strongly with Mo compared to larger alkyls and has thus a greater net donor capability. This result may seem somewhat surprising as

Table 2. Hybridization of the P lone-pair and partial charges (q) of the ligand PX₃ in the neutral and charged LMo(CO)₅ complex obtained from by the NBO analysis performed at the PBE0/Def2-TZVP level of theory.

	P lone-pair	· hybridization	L partial charge	
L	% s(P) lp	% p(P) lp	LMo(CO)5	LMo(CO)5 ^{+•}
PMe ₃	50.14	49.85	0.649	0.770
PEt ₃	50.78	49.21	0.645	0.804
PiPr ₃	44.57	55.42	0.620	0.806
PPh ₃	45.31	54.67	0.577	0.889
PPh ₂ Me	47.16	52.83	0.605	0.886
PPhMe ₂	48.53	51.46	0.628	0.856
PPh ₂ Pyr	48.16	51.83	0.575	0.950
PPhPyr ₂	50.38	49.61	0.575	0.967
PPyr3	54.52	45.48	0.571	1.478

larger alkyl groups are expected to have a larger donor ability due to inductive effects. This trend was however confirmed by the decrease of the Mo-P bond length from PMe₃ to PiPr₃ (resp. from 2.51 to 2.61 Å) and the NBO analysis of the partial charges: the carried charge by the PR₃ ligands decreases from 0.649 for PMe₃ down to 0.620 for PiPr₃, showing a reduced electron deficiency on the ligand (Table 2) and thus, a smaller donating effect for larger alkyls. Further, two recent experimental works also reached the same conclusions. First, Kendall *et al.* studied the steric and electronic influences of Buchwald-type alkyl-JohnPhos ligands in LCr(CO)₅ complexes and showed using IR data that the net-donating ability of the phosphines was decreasing with the increasing size.⁸² They explained this finding by the influence that steric effects also have on the

L-M interaction and which thus favor smaller alkyl groups. Then, in a previous study, we carried out an investigation on L-Au⁺-CO complexes using mass spectrometry techniques to measure the LAu⁺-CO bond-dissociation energies. The same observations on the alkyl size effect arise and were discussed by means of electrostatic interaction and Pauli repulsion as determined by EDA.⁸³ In the present work, the global donor effect of PR₃ ligands on LMo(CO)₅ complexes is discussed further in the light of the EDA results, but first, the experimental results are analyzed in more details.

On the SPES in Fig. 2a)-c), a broad band, similar to that of Mo(CO)₆, is observed in the ca. 7.5 eV region. This band exhibits two local maxima which are attributed to ionization of the e and the b2 sets. The small shoulder corresponding to the v1 progression (see above) is also observed. Attribution of these bands can be done following Koopman's theorem. The energy position of the HOMO of the different complexes are displayed in Figure 2a)-c) (red sticks). The agreement between the bands position and the KS-orbital's energies is noteworthy and points out the efficiency of the IP-tuned LC-ωPBEh method. The first band observed can be linked to ionization from the HOMO, HOMO-1 and HOMO-2 orbitals. The shapes of these three KS-orbitals are displayed in Figure 3 and indicate that the two orbitals of the e set (HOMO and HOMO-1) both backbond to three CO and one phosphine whereas the b2 orbital backbonds to four CO. The difference in energy between these two sets of orbitals (e and b2) can thus be linked to the backbonding variation between a CO and a phosphine bound to Mo and gives an evaluation of the π -acceptor effect of the ligand. The experimental difference between the two maxima in the first bands is respectively 0.19, 0.22 and 0.24 eV from methyl to iso-propyl showing that the π stabilization is decreasing by increasing the alkyl size (Table 1). Note that the PR₃ π -acceptor

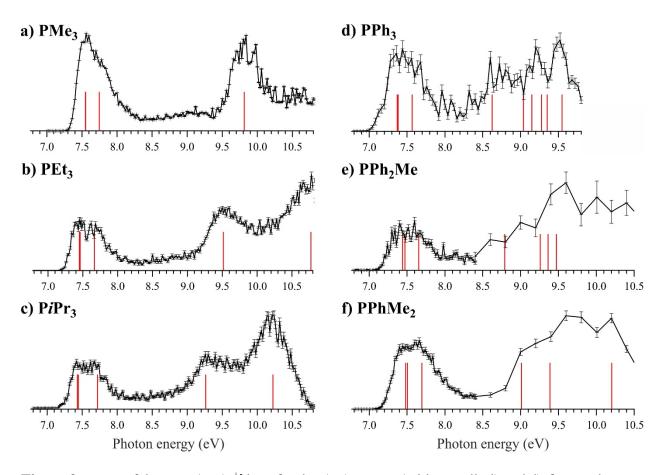


Figure 2. SPES of the LMo(CO)₅^{+•} ions for the a)-c) L=PR₃ (with R = alkyl) and d)-f) L=PPh_xMe_(3-x) series recorded in the 6.8-10.5 eV energy range. Red sticks represent the calculated energies of the first HOMOs of the neutral complexes. Because of a lack of experiment time, the spectra e) and f) were recorded in two steps: a first one from 6.6 to 8.4 eV with 20 meV steps and a second one from 8.4 to 10.5 eV with 200 meV steps.

orbital formally corresponds to a (σ^* (P-R) + 3d (P)) hybridized orbital (Scheme 1).⁸⁴ This experimental finding is supported by calculations which give also an increase of HOMO/HOMO-1 and HOMO-2 KS-orbitals energy difference by changing the alkyl size (resp. 0.20, 0.21 and 0.27 eV for Me to *i*Pr, see SI). This π -acceptor effect remains however much smaller than for a CO-group (as seen by the e set having a lower energy than the b2 set).

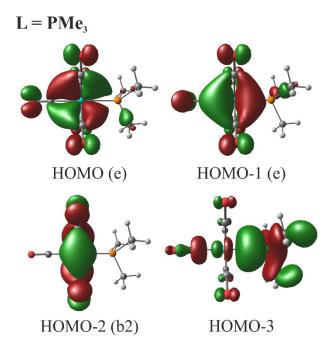


Figure 3. Kohn-Sham orbitals of Me₃PMo(CO)₅ obtained at the LC-ωPBEh/Def2-TZVP level of theory. The symmetry is indicated into brackets. Isosurface value: 0.03 a.u.

Furthermore, another broad band is present on the SPES in Fig. 2a)-c), at around 9.5 eV. It is attributed to ionization of the coordinated phosphine lone pair orbital which is involved in the σ bond with Mo. Yarbrough and Hall indeed compared the photoelectron spectra of several LMo(CO)₅ phosphine complexes with those of uncoordinated L ligands.⁸⁵ They could show that, under coordination, this P lone-pair band is the most destabilized because of its direct Mo-P interaction. Hence, the position of this band in the SPES is able to reflect the strength of the σ-donation of the phosphine ligand: the more destabilized it appears, the stronger is the σ-donation. In Fig. 2a)-c), the position of this band is shifted to lower energies by increasing the alkyl size: from 9.84 eV for PMe₃ to 9.52 eV for PEt₃ and 9.31 eV for PiPr₃. Note that the PMe₃ band position measured at 9.84 eV is in great accordance with the He(I) spectrum of the same Mo complex of Lichtenberger and Jatchko who measured a band at 9.87 eV.^{30, 86} These energies have to be compared with those of the P-lone pair in the uncoordinated ligand. Experimental data on the P

lone-pair band position could only be found for PMe₃ and PEt₃ in two works giving either 8.79 and 8.52 eV⁸⁵ or 8.60 and 8.27 eV⁸⁷ for PMe₃ and PEt₃, respectively. In both cases, the energy difference between these bands and the experimental bands observed in this work shows about the same difference (ca 1.0 or 1.2 eV). We have also evaluated the energy difference of these bands with theory as no experimental data were available for $PiPr_3$. The results are given in the SI and indicate also the same energy difference between PMe₃ and PEt₃, but a larger one for $PiPr_3$ suggesting thus the following trend for the σ -donor effect: PMe₃ \approx PEt₃ < $PiPr_3$. This observation is consistent with the known inductive effects of the alkyl groups: an increased σ -donation is indeed expected by increasing the alkyl size.

Table 3. EDA results of the L-Mo bond in LMo(CO)₅ complexes (in kcal/mol).

L	ΔE_{Pauli}	$\Delta E_{elstat}{}^{a}$	$\Delta E_{orb}{}^{a}$	ΔE_{int}
PMe ₃	96.28	-88.72	-48.26 (25.20/)	-40.70
PEt ₃	92.46	(64.8%) -84.44 (64.1%)	(35.2%) -47.26 (35.9%)	-39.25
PiPr ₃	87.96	-80.71 (63.9%)	-45.63 (36.1%)	-38.38
PPh ₃	81.97	-69.62 (60.4%)	-45.55 (39.6%)	-33.20
PPh ₂ Me	87.78	-76.58 (61.9%)	-47.19 (38.1%)	-35.99
PPhMe ₂	91.57	-81.87 (63.0%)	-48.08 (37.0%)	-38.39
PPyrlPh ₂	84.92	-69.99 (59.7%)	-47.31 (40.3%)	-32.38
PPyrl ₂ Ph	88.20	-70.68 (59.0%)	-49.18 (41.0%)	-31.66
PPyrl ₃	90.16	-70.31 (57.8%)	-51.23 (42.2%)	-31.38

 $[\]overline{\ }^a$ The values in parentheses give the contribution of $\Delta E_{elstat.}$ and $\Delta E_{orb.}$ to the total attractive interactions ($\Delta E_{elstat} + \Delta E_{orb}$).

The results of the energy decomposition analysis performed on the L-Mo(CO)₅ bond for the L = PR₃ (R = alkyl) are summarized in Table 3. The ΔE_{int} term appears to be the largest for PMe₃ (-40.70 kcal/mol) and decreases slightly for PEt₃ (-39.25 kcal/mol) and PiPr₃ (-38.38 kcal/mol). This result can be correlated with the decrease of IE_{ad} observed from PMe₃ to PiPr₃ (Table 1) that was rationalized by suggesting a stronger interaction between PMe₃ and the metal (see above). Decomposition of this term indicates that the total attractive interaction between the PR₃ and the Mo(CO)₅ fragments is dominated by the electrostatic interaction term ΔE_{elstat} (about 64%) whereas the covalent term introduced by ΔE_{orb} accounts only for 36%. Frenking rationalized this important electrostatic contribution by the large p character of the σ HOMO of PMe₃ which induces a good overlap between the phosphorous lone-pair electrons and the nucleus of the metal. 88-89 The same amount of p character is found for PEt₃ and PiPr₃ (above 50%, Table 2). Furthermore, the decomposition of ΔE_{orb} following the ETS-NOCV approach helps correlating the experimental results to the EDA. In Table 4, the ETS-NOCV results indicate that the deformation densities contributing to the orbital interaction on the L-M(CO)₅ bond have principally a σ -donor and a π acceptor character ($\Delta E_{\text{orb},\sigma}$, $\Delta E_{\text{orb},\pi 1}$ and $\Delta E_{\text{orb},\pi 2}$). The corresponding deformation densities $\Delta \rho_{\sigma/\pi}$ are presented in Figure 3 for L= PMe₃ (see the SI for the others ligands). $\Delta \rho_{\sigma}$ shows clearly the L \rightarrow M σ -donation whereas $\Delta \rho_{\pi 1}$ and $\Delta \rho_{\pi 2}$ represent the two components of the L \leftarrow M π backdonation. Comparison of the σ -donor and π -acceptor contributions of the L= PR₃ series confirms the experimental findings: the σ-donor contribution is larger for PiPr₃ than for PEt₃ and PMe₃ (76.4% of the total ($\sigma + \pi$) orbital contribution vs 74.1 and 73.6%, respectively, Table 4) whereas the π -acceptor character of the ligand is the highest for PMe₃. It represents indeed 26.3% of the $(\sigma + \pi)$ orbital contribution vs 25.9% for PEt₃ and 23.6% for PiPr₃ with NOCV eigenvalues for the total π contribution decreasing from 0.50 down to 0.43 along the series (Fig. 4 and S2). The

Table 4. ETS-NOCV of the L-Mo(CO)₅ complexes (in kcal/mol).

L	$\Delta E_{orb,\sigma}{}^{b}$	ΔEorb,π1 ^b	ΔEorb,π2 ^b
PMe ₃	-32.8	-5.9	-5.9
	(73.6%)	(13.3%)	(13.2%)
PEt ₃	-31.3	-5.5	-5.4
	(74.1%)	(13.1%)	(12.9%)
PiPr ₃	-30.1	-4.7	-4.6
	(76.4%)	(11.9%)	(10.1%)
PPh ₃	-28.6	-6.0	-5.6
	(71.2%)	(14.9%)	(13.9%)
PPh ₂ Me	-30.4	-6.3	-5.7
	(71.6%)	(14.9%)	(13.4%)
PPhMe2	-31.9	-6.2	-5.8
	(72.7%)	(14.1%)	(13.3%)
PPyrlPh ₂	-28.3	-7.7	-6.3
	(67.0%)	(18.1%)	(14.9%)
PPyrl ₂ Ph	-27.7	-8.7	-7.6
	(63.0%)	(19.8%)	(17.2%)
PPyrl3	-27.2	-9.5	-9.4
	(59.0%)	(20.6%)	(20.4%)

^a Only the major components are given here. A last contribution corresponding to a σ backdonation is also present but accounts for less than 2 % of the total orbital interaction ΔE_{orb} . This contribution was observed and discussed by Ardizzoia *et al.*¹³

Tolman analysis of these alkylphosphine ligands confirm these results: the C-O bond stretching frequency slightly varies in Ni(CO)₃L complexes from 2064 cm⁻¹ for PMe₃ to 2062 cm⁻¹ for PEt₃ and 2059 cm⁻¹ for P*i*Pr₃ thus revealing PMe₃ more π -acceptor than P*i*Pr₃ and globally more donating. Despite a smaller σ -donor and a larger π -acceptor effects of PMe₃ compared to the other PR₃ ligands, PMe₃ appears as the largest global donor ligand considering the trend in IE_{ad}, the NBO partial charges on PR₃ and the Δ E_{int} analysis. As already outlined above, this feature is confirmed by other experimental works. Se-83 The overall donating character of the ligand has

^b The contribution of each deformation density to the $(\Delta E_{orb,\sigma} + \Delta E_{orb,\pi1/\pi2})$ orbital interactions is indicated in parentheses.

indeed to be rationalize by considering all the energetic terms that account for the bonding. In the case of PMe₃, the attractive interaction largely overtakes that of PEt₃ and PiPr₃ (137.0 kcal/mol vs 131.7 and 126.3 kcal/mol, respectively, Table 3) and explains this non-intuitive effect.

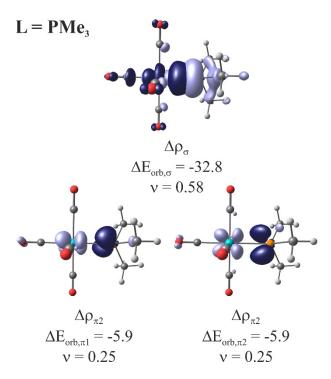


Figure 4. Contours of deformation density $\Delta \rho_{\sigma/\pi}$ connected with the pairs of interacting orbitals in the Me₃P-Mo(CO)₅ complex. Light purple surfaces represent a loss of electron density and dark blue surfaces, a gain of electron density. The associated orbital interaction energies ΔE_{orb} are given in kcal/mol. v is the eigenvalue indicating the size of the charge flow. Isosurface value: 0.0025 a.u.

The $L=PPh_xMe_{(3-x)}$ set of ligands was studied in the same way as the $L=PR_3$ sets and resulting SPES are shown in Fig. 2a), d)-f). Complexes with phenyl group were the most difficult to vaporize and required temperatures up to $160^{\circ}C$ ($L=PPh_3$) whereas an oven temperature of $130^{\circ}C$ was sufficient to observe the $L=PPh_2Me$ and $L=PPhMe_2$ complexes. Concerning the adiabatic IE, a clear energy decrease is observed by changing the number of phenyls from 0 (PMe₃, $IE_{ad}=7.32\pm0.03$ eV) to 3 (PPh₃, $IE_{ad}=6.99\pm0.03$ eV, see Table 1). The global interaction

between the $PPh_xMe_{(3-x)}$ ligands and Mo seems thus to be stronger by increasing the number of methyls. This finding is confirmed by the NBO partial charges analysis on the $PPh_xMe_{(3-x)}$ moiety: a net decrease of the charge is observed when increasing the number of phenyls $(q(PPh_xMe_{(3-x)})) =$

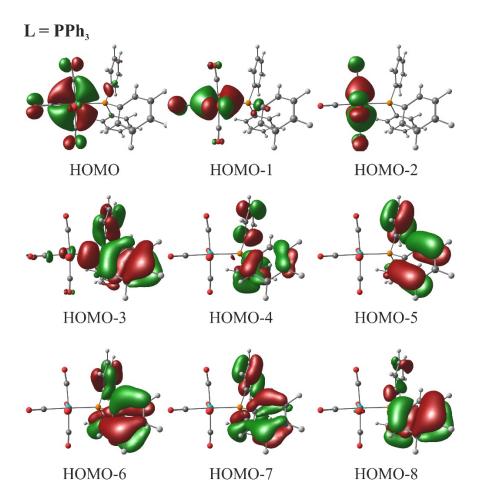


Figure 5. Kohn-Sham orbitals of Ph₃PMo(CO)₅ obtained at the LC-ωPBEh/Def2-TZVP level of theory. Isosurface value: 0.03 a.u.

0.649, 0.628, 0.605, 0.577 for x = 0 to 3, respectively (see Table 2)), indicating a decreased electron deficiency on -PPh_xMe_(3-x) and thus an increased global donor effect following PMe₃ > PPhMe₂ > PPh₂Me > PPh₃. This trend is confirmed by the behavior of PPhMe₂, PPh₂Me and PPh₃ observed in oxidative addition reactions of Pd(OAc)₂(PR₃)₂ complexes toward phenyl iodide: PPhMe₂ is

found to be more active than PPh₂Me and PPh₃ because of its higher donating ability to the metal.⁴⁶,

Further, as for $L = PR_3$, a common structure is observed for the $L = PPh_xMe_{(3-x)}$ SPES: a first intense band followed by a second one but which is much broader than for PR₃. MO analysis shows that this broadening is due to the contribution of the phenyl orbitals to the HOMOs which account in this energy region (see Figure 5 in the case of Ph₃PMo(CO)₅ and SI for the others). Interpretation of the first band in terms of π -acceptor effect remains nevertheless similar to that of the PR₃ set. The three first HOMOs of this series of complexes indeed possess a full d character (Fig. 5). The measured energy difference between the two first ionization energies reveal no noticeable change, respectively 0.19, 0.19, 0.18 and 0.20 eV for $PPh_xMe_{(3-x)}$ by varying x from 0 to 3 (Table 1). Similarly, the calculated HOMO/HOMO-1 and HOMO-2 energies show a difference of 0.20, 0.19, 0.18 and 0.19 eV respectively. The π -backdonation appears thus to be similar for this whole set of ligands. This observation is rather in agreement with the NBO analysis performed by Leyssens and coworkers showing that PMe₃ was a slightly better π-acceptor than PPh₃. ¹⁶ Concerning the higher-energy bands on the SPES on Fig. 2a),d)-f), their first contribution is also attributed to the P coordinated pair to Mo (e.g. HOMO-3 for PPh₃, see Fig. 5) and the others (up to 5 in the case of PPh₃) to the π orbitals of the phenyls. Due to the broadness of the band, a precise evaluation of the σ ionization energy is difficult to provide but one can still see some trend for the band position: the more phenyl groups are in the ligand, the lower in energy the band is: from 9.84 eV for PMe₃ to ca. 8.7 eV for PPh₃. These values have to be compared with the energies of the phosphorus lonepair in uncoordinated ligand. Ikuta et al. have provided the following estimation of the IEs for the set of $PPh_xMe_{(3-x)}$ ligands: $IE(PPh_xMe_{(3-x)}) = 8.62, 8.32, 8.23$ and 7.80 eV for x = 0 up to x = 3, respectively. 92 This indicates an energy difference between the P coordinated- and lone-pair MO

energies of 1.22, 0.88, 0.87 and 0.90 eV respectively, showing first the largest change upon coordination of the PMe₃ ligand and then a similar variation for PPhMe₂, PPh₂Me and PPh₃. This suggests a stronger σ -donor effect of the PMe₃ ligand and a similar effect for the others: PMe₃ > PPhMe₂ \approx PPh₂Me \approx PPh₃. Calculations confirms as well these energy differences (see SI) and thus the trend for this effect.

The EDA results for the $L = PPh_xMe_{(3-x)}$ sets of ligands are presented in Table 3. The interaction energy ΔE_{int} shows a decrease by increasing the number of phenyl groups from -40.70 kcal/mol for PMe₃ down to -33.20 kcal/mol for PPh₃, following the same trend as the adiabatic ionization energies. As expected with the structure of these ligands, the steric interaction ($\Delta E_{Pauli} + \Delta E_{elstat}$) between the Mo(CO)₅ and L fragments is playing a non-negligible role: from 7.35 kcal/mol for the small PMe₃ fragment, it raises up to 12.35 kcal/mol for PPh₃. This feature is mostly due to a decreasing ΔE_{elstat} term for large ligands: the overlap between the phosphorous lone-pair electrons and the nucleus of the metal cannot be as efficient as for small ligands (for steric reasons here but not because of the p character of the σ HOMO of PX₃ which is also high in PPh₃ (above 50%, see SI)). The orbital contribution to the total attractive bond interaction slightly increases with x (35.2% for x = 0 up to 39.6% for x = 3). The ETS-NOCV results are given in Table 4 (and the deformation densities in Fig.6 for PPh₃, in the SI for the other ligands). They confirm very well the experimental findings: PMe₃ appears to be a better σ-donor ligand compared to PPhMe₂, PPh₂Me and PPh₃, as seen by the slightly decreasing σ-donor contribution (73.6% vs 72.7%, 71.6% and 71.2%, respectively). On the other hand, the π -acceptor part on the total orbital interaction remains stable for all the series: between 13.2 and 13.9% of the $(\sigma + \pi)$ orbital contribution. The TEPs of these ligands match with these observations: the stretching frequencies v_{CO} in Ni(CO)₃L

complexes follow this trend: PMe₃ (2064 cm⁻¹) < PPhMe₂ (2065 cm⁻¹) < PPhMe₂ (2067 cm⁻¹) < PPh₃ (2069 cm⁻¹), giving PMe₃ slightly more σ-donating than PPh₃.⁹⁰

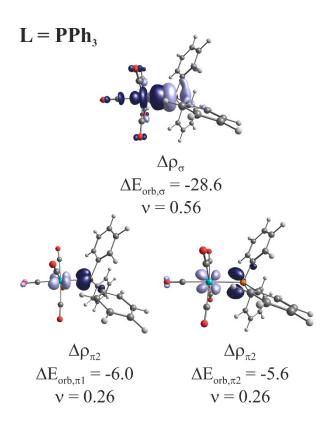


Figure 6. Contours of deformation density $\Delta \rho_{\sigma/\pi}$ connected with the pairs of interacting orbitals in the Ph₃P-Mo(CO)₅ complex. Light purple surfaces represent a loss of electron density and dark blue surfaces, a gain of electron density. The associated orbital interaction energies ΔE_{orb} are given in kcal/mol. ν is the eigenvalue indicating the size of the charge flow. Isosurface value: 0.0025 a.u.

The last set of ligands studied was the L (Pyrl = pyrrolyl, x = 0-3) series. The ion signal was optimized for a temperature between 100°C (PPyrl₃) and 150°C (PPh₃) and resulting SPES are presented in Fig. 7. Whereas the PPh₃ complex still has a SPES quite similar to the previous SPES shown (i.e. constituted by a first intense band at ca. 7.7 eV and a smaller one at ca. 8.7 eV), once the ligand contains at least one pyrrolyl, the SPES changes dramatically to a unique unstructured

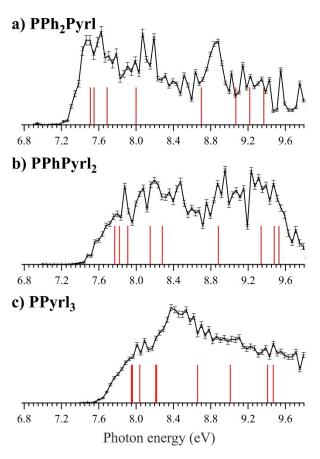


Figure 7. SPES of the LMo(CO)₅^{+•} ions for the PPh_xPyrl_(3-x) series recorded in the 6.8-9.7 eV energy range. Red lines represent the calculated energies of the first HOMOs of the neutral complexes

broad band. This feature is explained by considering the MOs involved in the signal: up to 9 HOMOs distributed all over the 7.5-9.8 eV range and presenting a mixing from the orbitals of the metal and of the ligands (see Fig. 8 to visualized the nine first HOMOs of Pyrl₃PMo(CO)₅ and in the SI for the other ligands). From HOMO to HOMO-5, the orbitals are characterized by a contribution of the pyrrolyl groups in addition to the metal d orbitals. This orbital mixing renders the interpretation of the electronic effects much more complex if not impossible. But reasoning on the adiabatic ionization variation energy still enables to get some insights on the ionization process.

A clear increase of the IE values is indeed observed by increasing the number of pyrrolyls: from 6.99 eV for PPh₃ to 7.59 eV for PPyrl₃. As the first HOMOs are not only related to metal orbitals but also to pyrrolyl's, it is thus not straightforward to conclude that the lowest the IE, the largest the global donor effect of the ligand. As stated for PR₃ ligands, this argument indeed holds true if the corresponding ejected electron is in the direct metal environment reflecting thus the interaction between the metal and the ligand. Considering the partial NBO charges on the PX₃ moiety in the neutral and in the ionized complex indicates an interesting feature. For PPh₃, only a slight increase of the electron deficiency on the ligand is observed upon ionization (+ 0.312) whereas it is very high for PPyrl₃ (+ 0.907) (Table 2). This shows that the charge is carried mostly by the ligand for PPyrl₃ on LMo(CO)₅^{+•} and on the metal for PPh₃, revealing a π (pyrrolyl) ionization and a d (Mo) ionization, respectively (assuming that no electronic rearrangement occurs after ionization). Thus, as the electron is ejected from the ligand (which is, in this case, termed as "non-innocent" ligand), 93-94 one can assume that the less interaction the ligand has with the metal, then the more density remains on it which accordingly increases its stability and therefore its ionization energy. This would thus indicate that PPyrl₃ is less globally donating than PPh₃ because its IE_{ad} is higher in energy. This trend is consistent with the known very strong electron-acceptor effect of pyrrolylcontaining phosphine ligands which counterbalances any possible donating effect, 95 and with the NBO charge analysis showing a slight decreased of the electron deficiency of the ligand by increasing the number of pyrrolyls (see SI). Considering now the MO shapes and energies in more details reveals that the orbitals accounting for the SPES can be roughly separated in three groups. The first one located below to 8.4 eV is constituted by orbitals with a combination of metal d and π (pyrrolyl) character (HOMO to HOMO-5 for PPyrl₃, Fig. 8). The second one concerns one orbital at 8.62, 8.70, 8.88 and 9.01 eV for ligand with zero to three pyrrolyls, respectively. This orbital is

the only one which exhibits in part a Mo-P σ character (HOMO-6 for PPyrl₃). And finally, the last group is composed by orbitals with a π (phenyl) and π (pyrrolyl) contributions (HOMO-7 and HOMO-8 for PPyrl₃). This analysis of the MOs repartition helps us to derive at least one property for this set of ligands: the calculated σ ionization is indeed distinctly decreasing by lowering the number of pyrrolyls in the PPh_xPyrl_(3-x) set and comparing these values with the calculated P lone pair ionization in uncoordinated ligand shows that the energy difference is decreasing with the number of pyrrolyls (see SI) giving thus PPyrl₃ less σ -donating than PPh₃. Furthermore, comparing the π -acceptor effect of the different ligands is not possible for this series of ligands because of the high mixing between the metal orbitals and the ligand orbitals.

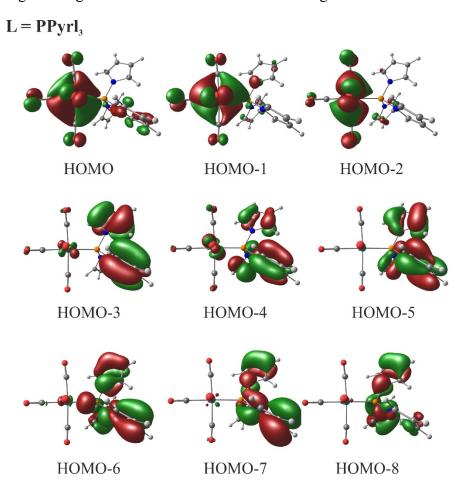


Figure 8. Kohn-Sham orbitals of Pyrl₃PMo(CO)₅ obtained at the LC-ωPBEh/Def2-TZVP level of theory. Isosurface value: 0.03 a.u.

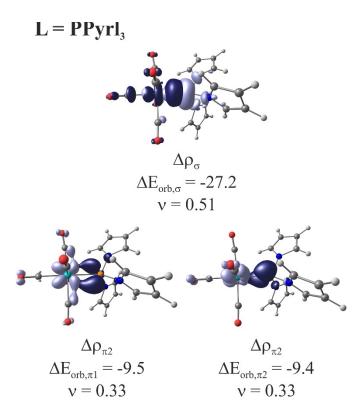


Figure 9. Contours of deformation density $\Delta \rho_{\sigma/\pi}$ connected with the pairs of interacting orbitals in the Pyrl₃P-Mo(CO)₅ complex. Light purple surfaces represent a loss of electron density and dark blue surfaces, a gain of electron density. The associated orbital interaction energies ΔE_{orb} are given in kcal/mol. v is the eigenvalue indicating the size of the charge flow. Isosurface value: 0.0025 a.u.

The EDA results for this series are displayed in Table 3. Not surprisingly, the electrostatic interaction remains the same (between -69.6 and -70.31 kcal/mol) along the series because of similar steric properties and p character on the σ MO of PX₃ (see SI). A small variation is observed for the interaction energy: it slightly decreases from PPh₃ to PPyrl₃. On the contrary to the other ligands sets studied in this work, the trend observed in ΔE_{int} is the opposite as that observed for IE_{ad} for this series. This feature will be discussed further. Only the Pauli interaction seems to be more influenced by the ligand change: it ranges from 81.97 kcal/mol, for PPh₃ up to 90.16 kcal/mol

for PPyrl₃. The deformation densities contributions to the orbital interaction are informative as they provides an evaluation of the σ -donor and π -acceptor effects of the ligands that were not obtained by experiment (Table 4 and Fig. 9 for the deformation densities of PPyrl₃, in the SI for the other ligands). The strong electron-withdrawing character of the N-pyrrolylphosphines appears with the large $\Delta E_{\text{orb},\pi 1}$ and $\Delta E_{\text{orb},\pi 2}$ contributions to the total orbital interaction (up to 41 % for PPyrl₃ with a significant NOCV eigenvalue of 0.66 vs 0.51 for PPh₃ (Fig. 9 and S2)), giving the following trend for the π -acceptor character of the ligand: PPyrl₃ > PPyrl₂Ph > PPyrlPh₂ > PPh₃. On the contrary, the σ -donor part is decreasing with the number of pyrrolyls increasing as confirmed by the MOs analysis. The Tolman analysis of N-pyrrolylphosphine ligands is consistent with these results, giving the following v_{co} stretching frequencies for RhCl(CO)L₂ complexes: 1965 cm⁻¹ (PPh₃), 1992 cm⁻¹ (PPyrlPh₂), 2007 cm⁻¹ (PPyrl₂Ph) and 2024 cm⁻¹ (PPyrl₃). 95 Finally, the last point that we would like to address in this study concerns the correlation that may exist between IE_{ad} and the interaction energy ΔE_{int} obtained from EDA. Whereas the latter indicates the strength of the interaction between the ligand and the metal, the former gives an assessment of the enrichment of the metal through the interaction with the ligand: the more net donating the ligand is, the higher IE_{ad} the complex has, because of the stabilization provided on this metal (assuming an ionization occurring on the d orbital of Mo). For the $L = PR_3$ and L =PPh_xMe_(3-x) sets of ligands, we have observed the same trend for both terms: for the most donating ligands, IE_{ad} and ΔE_{int} were the largest. On the other hand, this trend was reversed for the L = $PPh_xPyrl_{(3-x)}$ ligands: the less donating ligand exhibits the higher IE_{ad} and the smaller ΔE_{int} . In this case, we explain this discrepancy by the fact that ionization may also be occurring directly on the π orbitals of the ligand and not to the d MO of the metal ("non-innocent" ligands). As suggested by a reviewer, we have considered the average local ionization energy (ALIE) of the complexes

to reinforce this assumption. ALIE is defined as the sum over orbital electron densities times absolute orbital energies and divided by the total electron density. It helps unveiling molecular regions the most prone to lose an electron. In Fig. S3, the ALIE mapped molecular isosurfaces of the $L = PMe_3$, PPh_3 and $PPyrl_3$ complexes are displayed. The blue area where the electrons are the most weakly bond is concentrated for all three ligands in the metal region but it is also present in the ligand, principally for $L = PPyrl_3$ which confirms our experimental findings. The correlation obtained between the IE_{ad} and ΔE_{int} terms is illustrated in Figure 10. Two rather good linear correlation appear: the first one for the $L = PR_3$ and $L = PPh_xMe_{(3-x)}$ sets of ligands (red squares, $r^2 = 0.91$) and the second one for the $L = PPh_xPyrl_{(3-x)}$ ligands (blue diamonds, $r^2 = 0.94$). Although many factors certainly influence the interaction between a metal and a ligand, the correlation found here between these two parameters seems to us to be an interesting and promising trail to get an easy and rapid experimental evaluation of net donating ligand effects and also on the intrinsic character (innocent/non innocent) of the ligand.

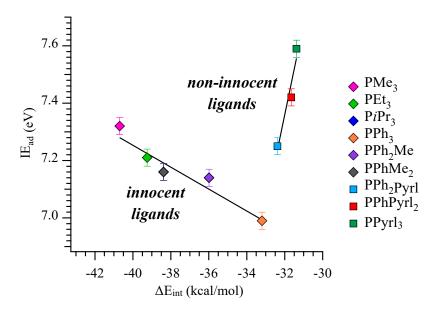


Figure 10. Correlation plots between IE_{ad} and ΔE_{int} (values given in Tables 1 and 2) for all ligands studied. Note that the points for $L = PiPr_3$ and $L = PPhMe_2$ are superimposed.

Conclusions

We have presented here a detailed study on the photoionization process of X₃PMo(CO)₅ complexes with various PX₃ ligands. Our interest was to determine if VUV spectroscopy was a suitable techniques for the experimental study of electronic properties of different phosphine ligands and separation of their π -acceptor and σ -donor effect. Experimental data obtained from a state-of-the-art experimental setup were confronted to the EDA method which is known to be very well adapted for ligands effect description. Three sets of ligands have been considered. The first one, the PR₃ set (R = alkyl), delivered SPES which could show that by increasing the alkyl size, the ligands have i) small decreasing π -acceptor properties, ii) increased σ -donor ability and iii-) an lower global donor electronic effect. For the second ligands set (PPh_xMe_(3-x)), SPES let us deduce that the π -acceptor effect remains unchanged in the series and that the same PPh₃ > PPh₂Me > PPhMe₂ > PMe₃ trend is observed in terms of σ -donor and global electronic effect. Finally, the last set to be studied was the PPh_xPyrl_(3-x) set whose SPES were dramatically different compared to the other ligands because of a high contribution of π (pyrrolyls) orbitals in the spectrum energy range. These spectra could thus not been interpreted in the same way as for the other ligands. We could however confirm the strong electron-acceptor and the poor donor ability of the PPyrl₃ ligand compared to PPh₃. An increased energy resolution could maybe help to observe better structured bands for these ligands and thus get more insights on the experimental evaluation of the σ -donor and π -acceptor effects of this set of ligands.

Comparison of these experimental results with EDA calculations shows an excellent agreement. The same trends and conclusions indeed arise from these two methods. The perspectives offered by such a joint experimental and theoretical work are very broad. Many very elaborated theoretical approaches have indeed been developed to describe ligand effects. But most of them also suffer

from the lack of experimental data to confirm their findings. On the other hand, the interpretation and analysis of the experimental spectra is greatly improved by using a theoretical approach such as EDA. The aim of the experimental part of this work was indeed not to deliver a complete description of the metal-ligand bonding as it can be done by theory. A description of the electronic structure of the complex was done and could thus give access to an evaluation of orbital description of the metal-ligand interaction. Other interaction terms (electrostatic, steric,...) terms were provided by theory.

Finally, a correlation could be found between the measure adiabatic ionization energies (IE_{ad}) and the interaction energies (ΔE_{int}) obtained from EDA. Depending on the ionization site of the complex (e.g. either on the d orbital of the metal or on the ligand's orbital), IE_{ad} and ΔE_{int} follow the same or the opposite trend. To our knowledge, it is the first time that such a correlation is evidenced. It opens interesting perspectives on the possibility to predict the innocent/non-innocent character of ligands which is known to orientate the chemistry of complexes. This correlation will be investigated further on other ligands and metals.

Supporting Information

Synthetic procedure; NMR spectra; computational data; DFT energies; ETS-NOCV; ALIE; cartesian coordinates

Acknowledgments

We would like to dedicated this work on the memory of Ridha Elattar, a brilliant young researcher of our group who passed away this year.

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TOC Graphic

