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Generalization of second order quasi-phase matching in whispering gallery mode resonators using Berry Phase

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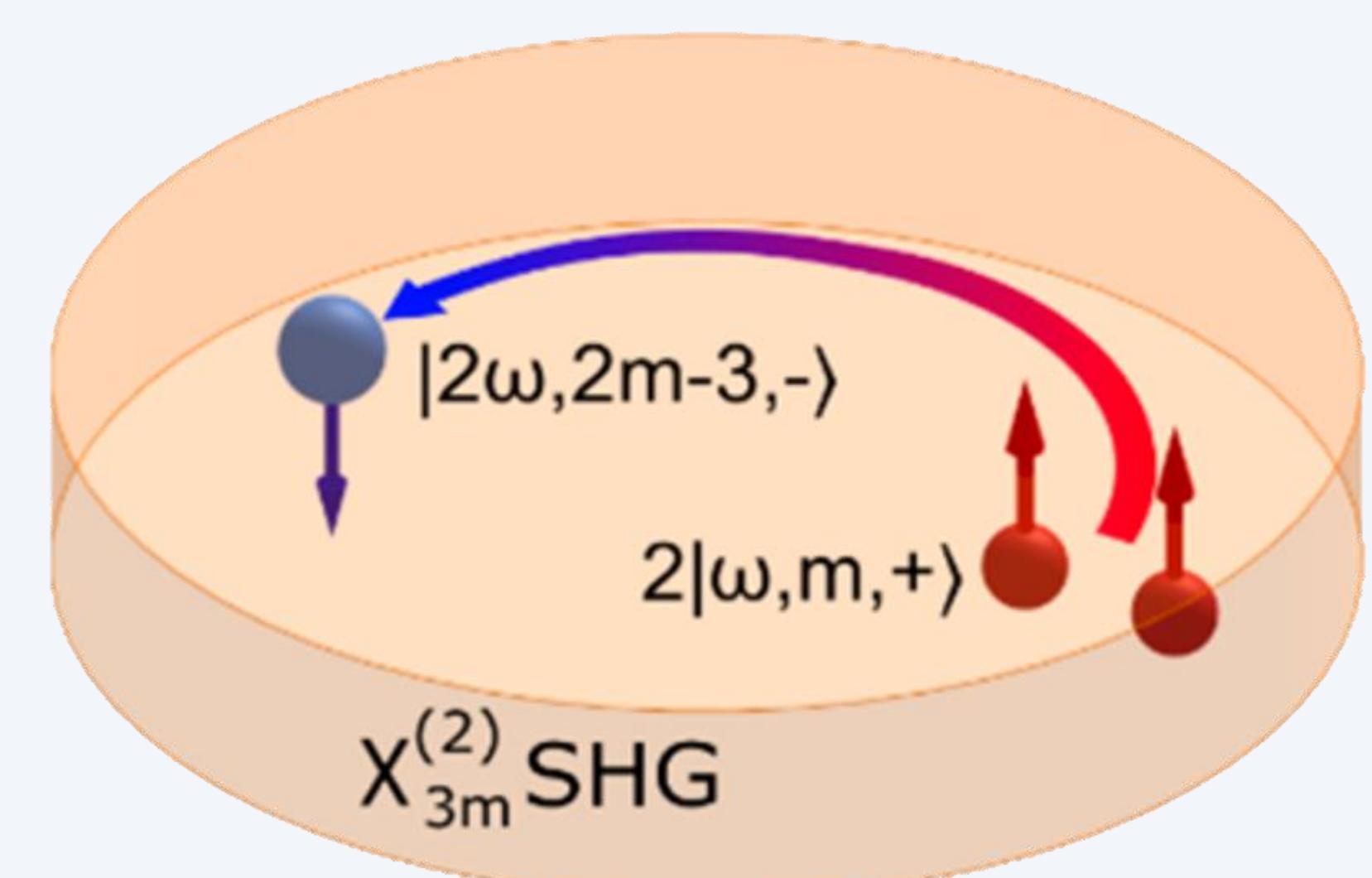
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Introduction

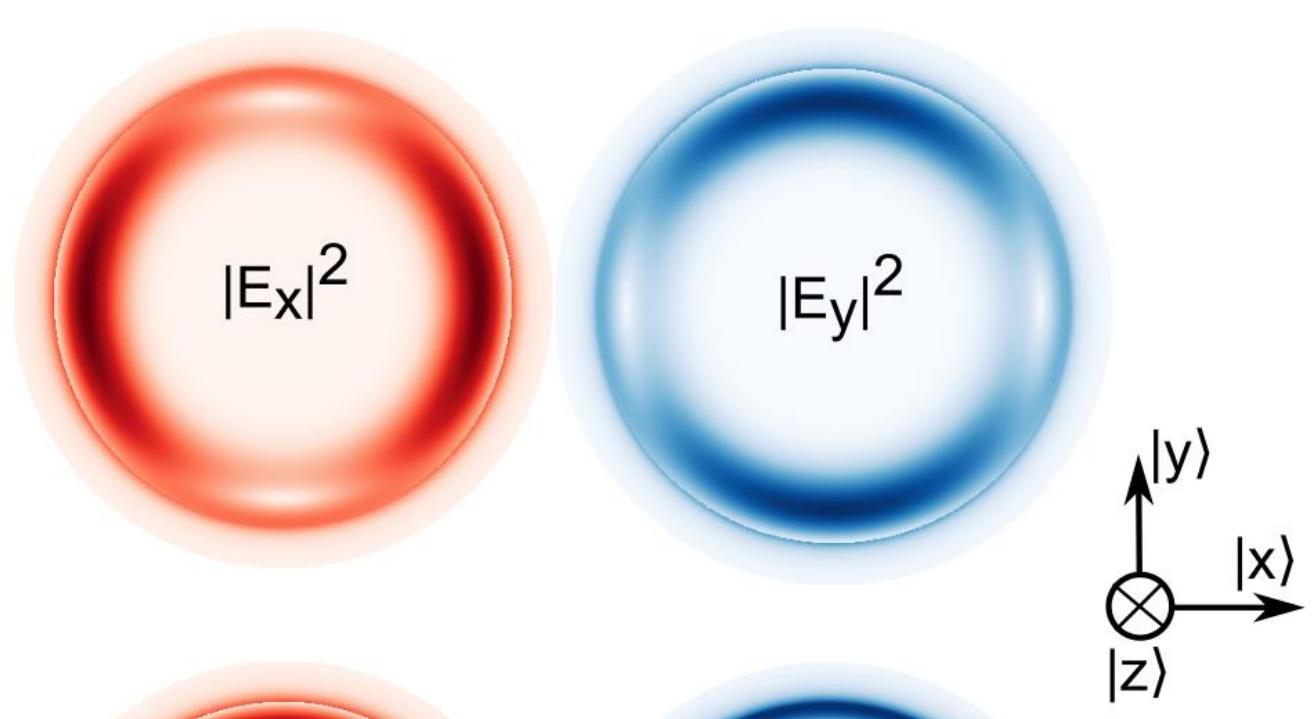
Second order processes in circular microresonators are highly investigated to design and fabricate high efficiency integrated devices:

- Need to conserve momentum conservation → Quasi-phase matching (QPM) → Azimuthal number selection rules (Δm)
- Natural QPM in WGM strongly depend on crystal symmetry
- Different additional quanta appear for different materials and crystallographic orientations
- No simple explanation of the physical source of these quanta in literature

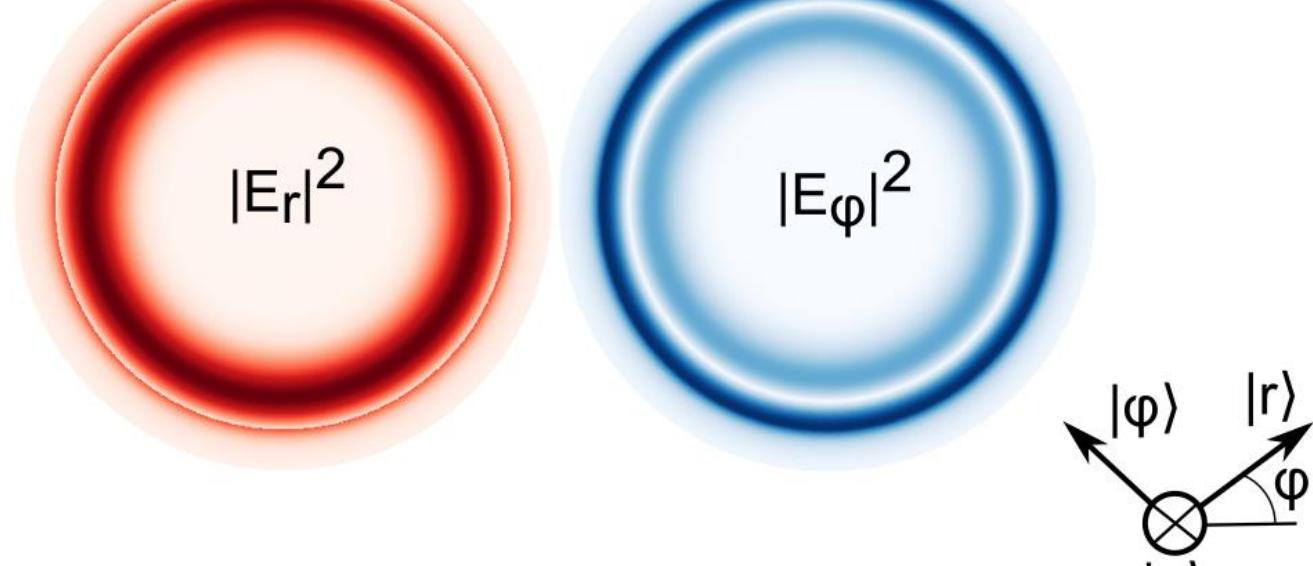


Transverse Spin Angular Momentum (TSAM) of photons: circular polarization basis

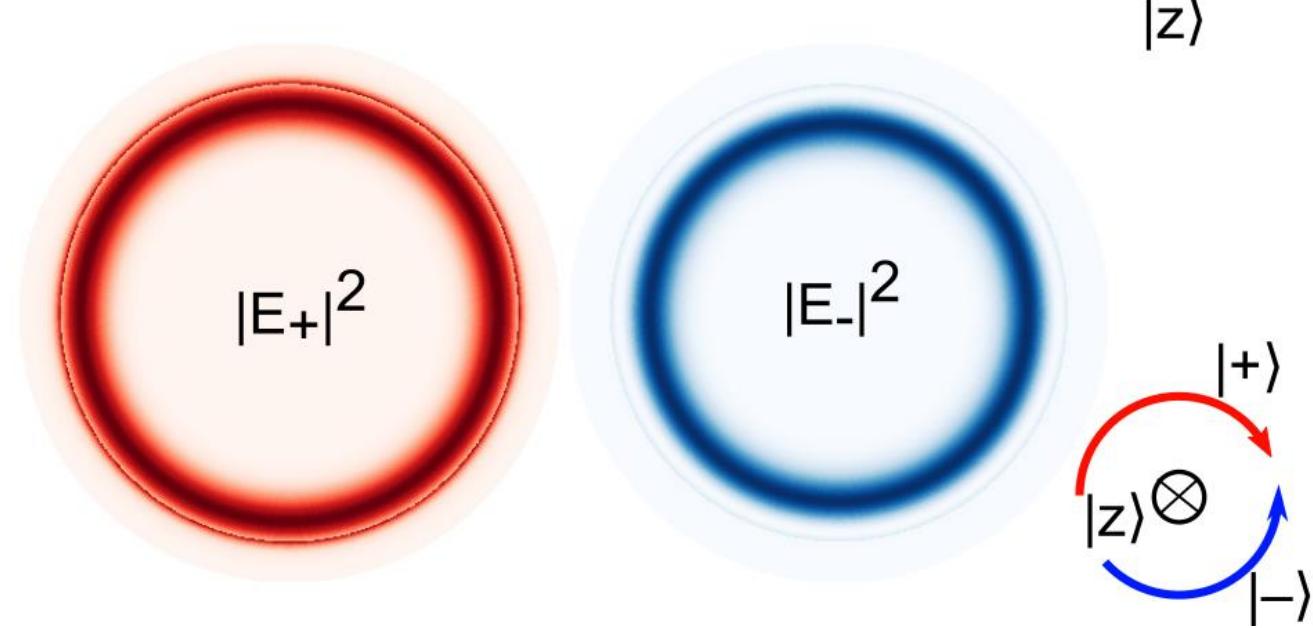
- Cartesian fixed basis



- Rotating polar basis



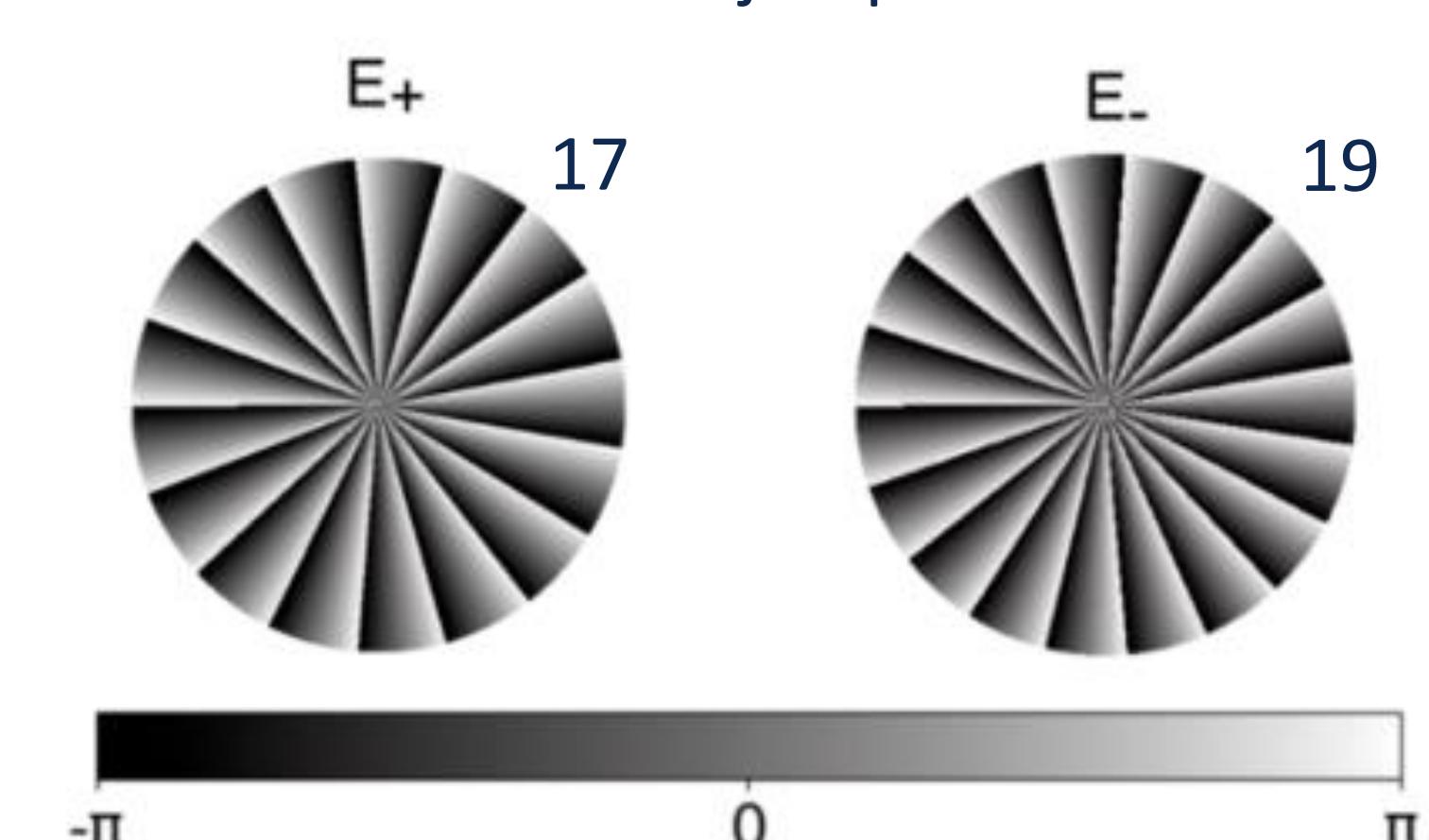
- Circular polarization basis (CP basis) along resonator axis → TSAM components⁽¹⁾



Berry phase in TSAM components: spin up & down polarization of photons

$$\begin{cases} |+\rangle = \frac{1}{\sqrt{2}}(|x\rangle + i|y\rangle) \\ |-\rangle = \frac{1}{\sqrt{2}}(|x\rangle - i|y\rangle) \end{cases} \xrightarrow{\quad} \begin{cases} E_+ = \frac{1}{\sqrt{2}}(S_m(r) + T_m(r))e^{i(m-1)\varphi} \\ E_- = \frac{1}{\sqrt{2}}(S_m(r) - T_m(r))e^{i(m+1)\varphi} \end{cases}$$

TE WGM m=18
Phase jumps:



- Opposite additional quanta on azimuthal dependence appears on both CP components
- This phenomenon is a direct consequence of the Berry phase experienced by each CP component during the revolution within the resonator⁽²⁾

Generalization of nonlinear susceptibility tensor and azimuthal selection rules

General form of the nonlinear tensor can be written in the circular polarization basis keeping the d_{ij} parameters known from the usual cartesian representation

$$\begin{pmatrix} P_+ \\ P_- \\ P_z \end{pmatrix} = \epsilon_0 \begin{pmatrix} \frac{\sqrt{2}(d_{11} + d_{12} + id_{16} + id_{22})}{4} & \frac{\sqrt{2}(d_{11} - 3d_{12} - 3id_{16} + id_{22})}{4} & \frac{\sqrt{2}(d_{13} - id_{23})}{2} & -id_{14} + \frac{d_{15}}{2} - \frac{d_{24}}{2} & \frac{d_{15}}{2} + \frac{d_{24}}{2} & \frac{\sqrt{2}(d_{11} + d_{12} - id_{16} - id_{22})}{4} \\ \frac{\sqrt{2}(d_{11} - 3d_{12} + 3id_{16} - id_{22})}{4} & \frac{\sqrt{2}(d_{11} + d_{12} - id_{16} - id_{22})}{4} & \frac{\sqrt{2}(d_{13} + id_{23})}{2} & \frac{d_{15}}{2} + \frac{d_{24}}{2} & id_{14} + \frac{d_{15}}{2} - \frac{d_{24}}{2} & \frac{\sqrt{2}(d_{11} + d_{12} + id_{16} + id_{22})}{4} \\ id_{14} + \frac{d_{15}}{2} - \frac{d_{24}}{2} & -id_{14} + \frac{d_{15}}{2} - \frac{d_{24}}{2} & d_{33} & \frac{\sqrt{2}(d_{13} - id_{23})}{2} & \frac{\sqrt{2}(d_{13} + id_{23})}{2} & \frac{d_{15}}{2} + \frac{d_{24}}{2} \\ \end{pmatrix} \begin{pmatrix} E_+ E_+ \\ E_- E_- \\ E_z E_z \\ 2E_- E_z \\ 2E_+ E_z \\ 2E_+ E_- \end{pmatrix}$$

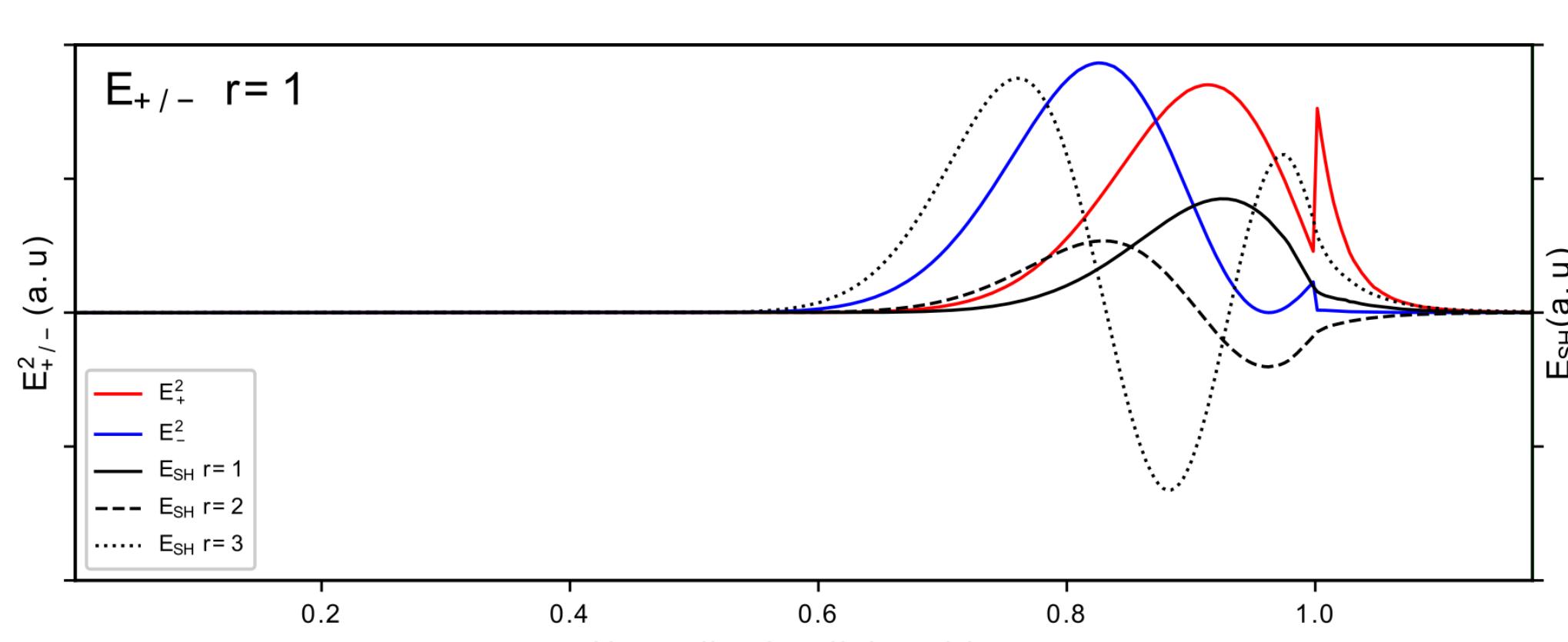
Each tensor element is associated to a single creation/annihilation process in CP basis with a unique azimuthal number mismatch

$$\begin{bmatrix} +1 & -3 & -1 & -2 & 0 & -1 \\ +3 & -1 & +1 & 0 & +2 & +1 \\ +2 & -2 & 0 & -1 & +1 & 0 \end{bmatrix}$$

Examples:

- Zinc-blende: TE-TE → TM $\Delta m = \pm 2$
- LnNbO₃ z-cut: TE+TE → TM $\Delta m = \pm 3$
- LnNbO₃ x-cut: TE+TE /TM → TM $\Delta m = \pm 1$
- Wurtzite: TE+TE → TM $\Delta m = 0$

Choosing the right WGM: towards better design of resonators for nonlinear processes



$$\beta_2 = \frac{i d_{14}}{4} \frac{\int d^3x \varepsilon E_{2z}^*(E_{1+}^2 - E_{1-}^2)}{\left[\int d^3x \varepsilon |E_1|^2 \right] \left[\int d^3x \varepsilon |E_2|^2 \right]^{\frac{1}{2}}}$$

- Azimuthal dependence can be inserted in the overlap integral⁽³⁾.
- A precise optimization can be made using higher radial orders
- Knowing exactly the Δm condition allows a precise tuning of integrated coupling systems

Conclusion

- TSAM description of WGMs to explicit the Berry phase dependence
- General expression to obtain the nonlinear susceptibility tensor in circular polarization basis
- Origin of the different azimuthal selection rules presents in second order processes in circular resonators
- Useful description to optimize and design highly efficient devices.
- Detailed work available at Lorenzo-Ruiz, Alejandro, and Yoan Léger. "Generalization of second order quasi-phase matching in whispering gallery mode resonators using Berry Phase." *ACS Photonics* (2020).

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