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# High Temperature Behaviour of Early Life Membrane Models

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10 **Abstract:** Origin of life scenarios generally assume an onset of cell formation in terrestrial hot springs or in the deep oceans close to hot vents, where energy was available for non-enzymatic reactions. Membranes of the protocells had therefore to withstand extreme conditions different from what is found on Earth's surface today. We present here an exhaustive study of temperature stability up to 80 °C of vesicles formed by a mixture of short-chain fatty acids and alcohols, which are  
15 plausible candidates for membranes permitting the compartmentalization of protocells. We confirm that the presence of alcohol has a strong structuring and stabilizing impact on the lamellar structures. Moreover and most importantly, at high temperature (> 60 °C) we observe a conformational transition in the vesicles, which results from vesicular fusion. Because all the most likely environments for the origin of life involve high temperature, our results imply the need to take into account such a  
20 transition and its effect when studying a protomembrane model's behaviour.

## Introduction

Compartmentalization of the protocell and thus the formation of the lipid boundary between the cell and its environment is central to the emergence of early life and to assure the required conditions for energy-dependent life processes. Indeed, uni- or multi-lamellar vesicles (ULV – MLV) are believed to  
25 be the ancestors of cell membranes <sup>1,2</sup> and to have provided a privileged environment to early biopolymers, enabling the formation of the earliest protomembranes and cells. To create such barriers, membranes were composed of self-assembling amphiphilic molecules. Contemporary membranes are all based on phospholipids, a group of highly functionalized, diverse and adaptable molecules, whose presence at the early stages of life is not probable. Another aspect that suggests the progressive  
30 transition between early amphiphiles and modern phospholipids is that the latter require a complex synthesis pathway and are unlikely to have been produced spontaneously, but rather elaborated in successive selective steps <sup>3</sup>.

In contrast, simple amphiphilic substrates were likely formed in the relatively harsh but potentially wet conditions of the early Earth by Fischer-Tropsch processes <sup>4,5</sup>. Photo-oxidation likely valorized  
35 alkanes by transforming them into oxygenated amphiphilic material such as fatty acids and fatty alcohols. This synthesis would have led to high abundance of short chain Fatty Acids (FA) ( $C < 12$ ), the abundance of which fall exponentially with the tail length <sup>1,6</sup>. FA have been proven to form vesicles when the chain length is eight or more carbons <sup>7</sup>.

Pioneering work has documented the ability of these amphiphiles to form membrane structures in  
40 solution, the stability of which is strongly dependent on environmental conditions <sup>8</sup>. The first requirement for vesicle formation is a solution pH equivalent to the FA dissociation constant ( $pK_a$ ), so that approximately half of the protonated molecules act as hydrogen bond donor to the adjacent ionized ones <sup>9,10</sup>. The second requirement is to reach the Critical Vesicle Concentration (CVC) <sup>9</sup>, below which these single chain amphiphiles exist only as monomers in solution or aggregated in  
45 micellar structures. The CVC of pure FA vesicles has been measured quantitatively as a function of the chain length <sup>7</sup>. A decrease in the CVC was observed with increasing chain length, as expected, because longer apolar hydrocarbon tails favour the more packed bilayered structures. Furthermore, the addition of small amounts of fatty alcohols (with same length as the FA, in a 10:1 FA : alcohol molar ratio) was found to dramatically lower the CVC and also to allow vesicle formation at pH values  
50 higher than the  $pK_a$  of the acid. This can be explained because the alcohol ability to form hydrogen bonds with the anionic FA balances the lack, at high pH, of hydrogenated FA that would serve the same purpose (as it happens in the pure FA samples around its  $pK_a$ ). However, FA vesicles do not withstand high salinity, especially when divalent cations such as  $Mg^{2+}$  or  $Ca^{2+}$  are involved <sup>7,11,12</sup>.

All these major limitations have led to a paradoxical situation where the molecules able to self-  
55 assemble into vesicles that were probably the most abundant in an origin of life scenario, namely the FA with  $C < 12$ , are also seen as the least promising ones because of their observed sensitivity to the environment.

Some of these apparent issues have been potentially solved by considering a mixture of the same chain alcohol and FA, as mentioned above, for its effect on the CVC and the pH of vesicle formation.  
60 On the other hand, the sensitivity towards salt concentration has also been shown to be lowered either by alcohol addition <sup>13</sup>, by using mixtures of different length FA <sup>14</sup> or by prebiotic amino acids binding the FA membrane <sup>15</sup>. These studies prove that ways to overcome the sensitivity towards high salinity exist while still considering plausible prebiotic systems and environmental conditions. Yet, previous studies on the thermal behaviour and stability of protomembranes have been focusing on longer  
65 aliphatic chains <sup>16,17</sup> driven by the observed environmental sensitivity at ambient temperature. Nevertheless, some surprisingly thermostable properties have been observed even for vesicles made with short chain amphiphiles. In particular, vesicles made with mixtures of decanoic “capric” acid -

decanol or “capric” acid – decanol - monocaprin (2:1 and 4:1:1 molar ratio, respectively) where shown to retain encapsulated oligonucleotides up to 70 °C for more than 1 h, or 100 °C for more than 70 1 min <sup>16</sup>. These findings suggest that the observed sensitivity of the membranes made by short aliphatic chains towards changes in pH, amphiphile and salt concentrations, do not necessarily imply a lack of stability towards temperature variations.

A first temperature-pressure study of the C10 mixture (capric acid+decanol 1:1 ratio) in buffered solution (bicine 0.2M pH 8.5) has been performed by Kapoor et al. <sup>18</sup>. They observed and assigned 75 two phase transitions experienced by the system: the first, at 10 °C, identified as “the appearance of disordered fluid-like chains”; the second, at  $\approx 60$  °C, explained as “morphological changes of lipid entities with overall fluid-like chains” (e.g. a shift in the micelle-vesicle equilibrium). Interestingly, they also observed an increase of the vesicle size beyond 50 °C (diameter increasing from 120 nm at 45 °C to 160 nm at 55 °C), although the temperature range investigated (5 – 55 °C) did not allow 80 them to study in more detail the phenomenon at higher temperatures.

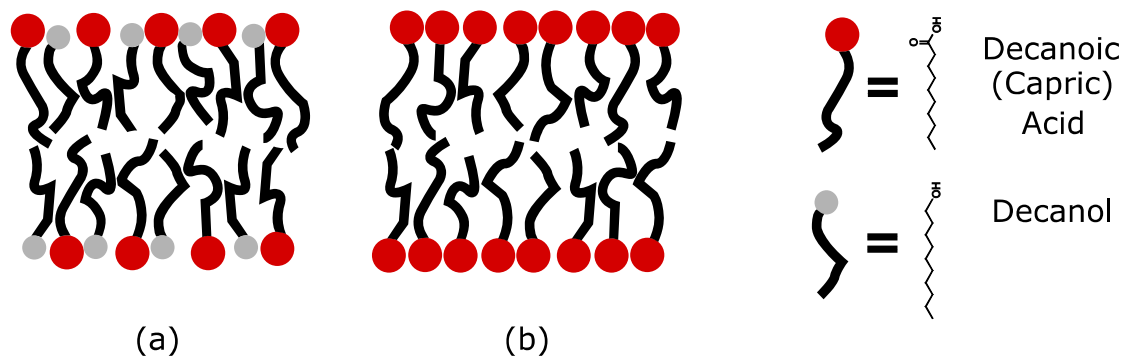
In this article, we extend the studies up to 80 °C and present a deeper and quantitative characterization of the membrane model composed of a capric acid:decanol 1:1 mixture (hereafter called C10 mix, Figure 1a) by following the properties and the evolution of the corresponding vesicles (partitioning, size distribution, lamellarity, membrane thickness, signs of micelle presence), as a promising model 85 of protocellular compartmentalization. We compare its characteristics with a similar model made of pure capric acid (Figure 1b). The stability with temperature is investigated for both systems, to understand the capability of the vesicles to withstand the high temperatures expected at hydrothermal vents/hot springs. It is of great interest to investigate the thermal stability of the protomembranes, as the high temperature would have allowed a number of important non-enzymatic reactions to take 90 place (e.g. DNA strand separation and reannealing) <sup>16</sup>.

We find that the alcohol insertion leads to a number of substantial differences in the vesicle behaviour:

- 1) It leads to a much larger proportion of vesicles, at the expense of other aggregates (micelles) or monomers in solution;
- 95 2) It protects vesicles from low temperature vesicle disruption, maintaining the multilamellar structures when present;
- 3) It triggers an abrupt conformational change above 60 °C. This leads to a transition from multilamellar vesicles to large unilamellar structures.

Such findings, in turn, could set a high T limit for the stability of vesicle systems to 60 °C at ambient 100 pressure, but could also open new alternative perspectives related to the observed changes in the

amphiphiles self-assembly. Our results point towards an alcohol induced, T triggered and slowly reversible vesicular fusion process.



**Figure 1.** The protomembrane models: (a) Bilayer made from a mixture of capric acid and decanol; (b) bilayer made of capric acid.

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## Experimental

### Sample preparation

Decanoic acid, 1-decanol, bicine buffer and D<sub>2</sub>O were purchased from Sigma Aldrich (Merck). The samples were prepared by dissolving the Decanoic acid (and the decanol, when needed) in a CHCl<sub>3</sub> solution, followed by drying with a flush of nitrogen (the Decanoic acid + 1-decanol mixture is liquid at ambient temperature). The samples were then placed in a dessicator and left under vacuum overnight. The sample weights were checked at each step, to make sure that all the chloroform was evaporated. At the final stage, a sample loss of  $\approx 5\%$  was measured for the C10 mix samples, probably due to a partial evaporation of the decanol. The final Decanoic acid : 1-decanol in the C10 mix samples was 1:1.

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The bicine buffer was prepared at a concentration of 0.2M in H<sub>2</sub>O. For the NMR experiments, the buffer was solubilized in <sup>2</sup>H depleted water. In the case of SANS experiments, the buffer solution was prepared by dissolving hydrogenated bicine in D<sub>2</sub>O. In all cases, the buffer was filtered with a 0.2  $\mu$ m millipore membrane before use.

120 The dried organic solutions were then suspended in buffer, titrated to desired pH/pD with aliquots of NaOH/NaOD, and vigorously vortexed, leading to final milky-foamy solutions.

For the Solid State NMR and SANS on D16 experiments, the total amphiphile concentration was set to 350 mM and the samples were measured as they were at this stage.

125 For DSC, SLS-DLS and SANS on D33, the total amphiphile concentration was 80 mM. Samples were extruded using the equipment provided by Avanti Polar Lipids, with the heating block regulated at 40 °C and using 100 nm polycarbonate membranes. The extrusion was always performed with 11 passes

per sample and less than 24h before the corresponding experiments. Light Scattering pre-characterization measurement ensured that the extruded vesicles were stable at ambient temperature for several days (data not shown).

### 130 Static - Dynamic Light Scattering

The multi-angle Static Light Scattering Measurements in the angular range 25-150° and the Dynamic Light scattering measurements at 150° (giving the results of Figure 2 in the text) were performed on an ALV CGS-3 instrument (ALV GmbH, Langen, Germany) at the Institut Laue Langevin (ILL) Grenoble with an incoming laser light of  $\lambda = 633\text{nm}$  wavelength. 1 ml of each sample was loaded on  
 135 round bottom borosilicate cuvettes after filtration through 0.45  $\mu\text{m}$  Millipore filters (0.2  $\mu\text{m}$  for the buffer), and left 10 min in the toluene bath for thermalisation before starting the measurements. A buffer measurement at the same conditions was performed to subtract its contribution, as well as toluene as a standard with known Rayleigh ratio ( $R_{\text{tol}} = 1.35 \cdot 10^{-5} \text{ cm}^{-1}$ ):

$$R_{\text{sample}} = R_{\text{tol}} * (I_{\text{sample}} - I_{\text{solvent}}) / I_{\text{toluene}} \quad (1)$$

140 where  $R$ , the Rayleigh ratio, is defined as the scattered intensity per unit of solid angle, scattering volume and incident intensity;  $I$  is the scattered intensity. The scattered intensity is corrected for the scattering volume (the sample volume visible for the detector) using  $I_{\text{sample/solv/tol}} = I_{\text{exp}} * \sin(\theta) / I_0$ , where  $\theta$  is the scattering angle and  $I_0$  the incoming intensity.

For the DLS measurements at 90° fixed angle, as a function of temperature, the instrument used was a  
 145 Malvern Zetasizer Nano ZS90. For the latter, 200  $\mu\text{l}$  quartz cells were used.

For the SLS measurements (ALV), the scattered intensity was collected for 10 s per angle. For DLS (ALV and Malvern), 3 runs of 30 s each were collected per angle for an error estimation of the diffusion coefficients.

The SLS scattering data were treated using the Guinier approximation<sup>19</sup>, whose validity was verified  
 150 for the 100 nm extruded vesicles. In fact, for the cases where  $q * R_g \leq 1$  ( $q$  being the scattering vector,  $q = 4\pi n \sin(\theta/2)$ ;  $R_g$  the gyration radius), one can approximate the scattering intensity (the Rayleigh ratio) as:

$$\ln R_{\text{sample}} \approx \ln(I_0) - q^2 R_g^2 / 3 \quad (2)$$

so that in a log-linear plot, the slope of the fitted line gives an estimation of  $R_g$ .

155 The Dynamic Light Scattering correlation functions were analysed following the method of the moments (a re-visitation of the method of cumulants) developed in<sup>20</sup>. The first two moments  $\mu_1$  (corresponding to the decay rate  $\Gamma$ ) and  $\mu_2$  were left as free parameters in the fit. For particles moving by Brownian diffusion, the decay rate is linked to the diffusion coefficient  $D$  by the relation  $\Gamma = D q^2$ .

In some cases, a sum of two populations, each one having its own  $\mu_1$  and  $\mu_2$ , was necessary for successful fits (indicating the existence of different families of scatterers of different sizes).

The Brownian diffusion coefficients  $D$  experimentally determined were then used to estimate the hydrodynamic radius  $R_h$  using the Stokes-Einstein equation:

$$D = k_B T / (6\pi R_h \mu) \quad (3)$$

with  $D$  the diffusion coefficient,  $k_B$  the Boltzmann constant,  $T$  the temperature,  $\mu$  the solvent viscosity (solvent viscosity and refraction index  $n$  were assumed the same as for  $H_2O$ ).

Fits of the DLS autocorrelation functions were performed using the Origin software (OriginLab Corp., Northampton, MA).

### Differential Scanning Calorimetry

For the DSC experiments, a microDSCIII calorimeter (Setaram, France) was used. Two Hastelloy C cells of 0.5 ml sample volume were used for both the samples and the buffer. The samples were measured against the buffer as a reference. The explored temperature range was  $-4 < T < 85$  °C, with scanning rate 0.5 °C /min. Three heating cycles were collected (the first for annealing, the last two to verify the reproducibility). No corrections or baseline subtractions were applied, the scans being used for qualitative comparison.

### 175 Solid State NMR

The Capric acid sample was prepared by hydrating 6mg of perdeuterated ( $^2H_{19}$ ) Capric acid with 100 $\mu$ l of 0.2M bicine buffer prepared in deuterium-depleted water (hydration of 94%). The C10 mix sample was prepared by mixing 3mg of perdeuterated Capric acid with decanol (1:1 molar ratio) and hydrated with 100 $\mu$ l of the same buffer (hydration of 97%). Samples were transferred into 100 $\mu$ l 4-mm zirconia rotors for NMR analyses.  $^2H$ -ssNMR experiments were performed at 76.77 MHz with a phase-cycled quadrupolar echo pulse sequence ( $90^\circ_x-t-90^\circ_y-t-acq$ )<sup>21</sup> and using a Bruker Avance III 500 MHz WB (11.75 T) spectrometer equipped with a solid state CPMAS 4mm H/X probe (IECB structural biophysics platform, Bordeaux, France). Acquisition parameters were as follows: spectral window of 500 kHz,  $\pi/2$  pulse width of 3.9  $\mu$ s, interpulse delays of 40  $\mu$ s, recycling delay of 1.3 s; number of scans from 500 to 3000. Samples were equilibrated for 30 min at each temperature before the acquisition. The acquisition time for Capric acid and C10 mix spectra was 10 and 65 min, respectively. Spectra were processed using a Lorentzian line broadening of 400 Hz before Fourier transformation from the top of the echo. Samples were equilibrated for 30 min at a given temperature before data acquisition. All spectra were processed and analyzed using Bruker Topspin 4.0.6 software.

### 190 Small Angle Neutron Scattering

SANS measurements were performed on two instruments, available at the ILL (Grenoble, France):

- D33 [<https://www.ill.eu/users/instruments/instruments-list/d33/>] <sup>22</sup>, with two neutron wavelength configurations ( $\lambda = 5 \text{ \AA}$  and  $\lambda = 14.5 \text{ \AA}$ ) and three sample-detector distances (2, 10 and 12 m) to cover the widest range of scattering vector ( $1.3 * 10^{-2} < q < 5.6$ )  $\text{nm}^{-1}$ ;
- 195 - D16 [<https://www.ill.eu/users/instruments/instruments-list/d16/>], with neutron wavelength  $\lambda = 4.5 \text{ \AA}$  and sample-detector distance 0.95 m, leading to the scattering vector range ( $2.6 * 10^{-1} < q < 5.3$ )  $\text{nm}^{-1}$ .

The sample suspensions, of 200  $\mu\text{l}$  each, were loaded in 1mm quartz cells (Hellma, Germany) on a sample holder with thermal regulation. For the start at low temperature during the temperature scans, samples were equilibrated for at least 20 mins. Then, for every T jump, the minimum equilibration time before starting the data acquisition of the first sample in the sample changer was 10 min; this time increased as samples were sequentially measured. We did not notice any inconsistencies due to the sample positions in the sample changer. All data were corrected for the scattering of the sample holder and the instrumental background. H<sub>2</sub>O was measured for detector efficiency calculation and scaling to absolute units ( $\text{cm}^{-1}$ ). A flat background was subtracted to account for the q-independent incoherent neutron scattering signal.

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The outcome from a SANS experiment is a series of curves of I vs. q, where

$$I(q) = |F(q)^2 * S(q)| = |P(q) * S(q)| \quad (4)$$

P(q) = F(q)<sup>2</sup> being the particle form factor and S(q) the structure factor. The form factor contains information about the shape of the scatterers at the different length scales (the q range) covered: in the case of vesicles this will be the vesicle size, shape, size polydispersity, lamellarity, bilayer thickness and corresponding polydispersity. The structure factor contains information about possible interaction between neighbour vesicles and, most importantly, the interaction between subsequent membranes forming the MLVs. The position of the first maximum of the correlation, due to the membrane S(q) in a SANS curve, can give a direct estimation of the membrane d-spacing through the relation  $d = 2\pi/q$  (with d: d-spacing; q: scattering vector).

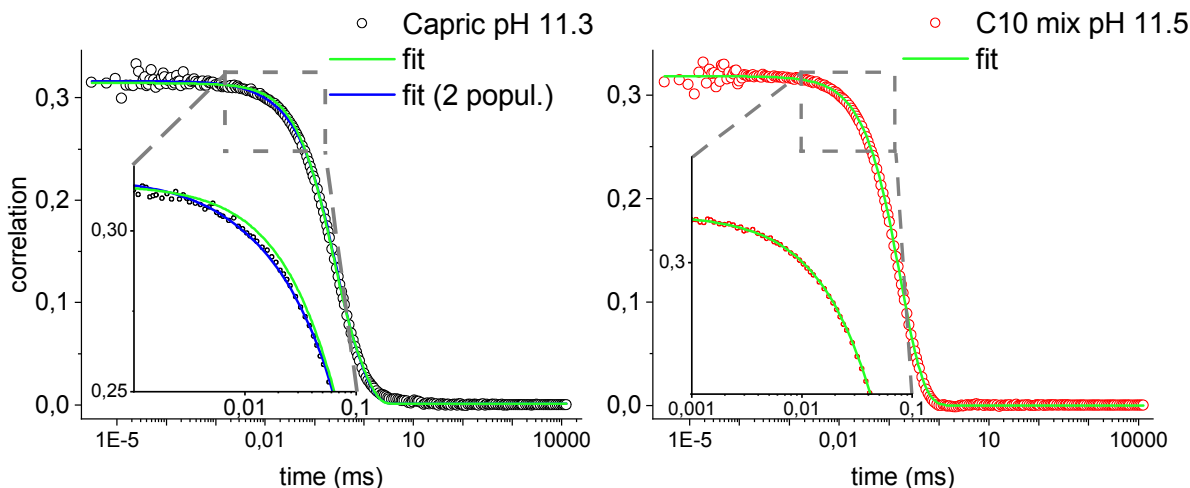
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The fitting of the SANS curves were performed using the SASView public software (<http://www.sasview.org/>). The equations for all the model form factors employed can be found in the “Material and methods” section (Supp. Info).

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225 **Figure 2.** Dynamic light scattering auto-correlation curves at high pH 11.4, and corresponding zooms on the short correlation times (insets) at 25 °C. Left: Capric acid 80 mM; right: C10 mix 80 mM. The solid lines are fits to the data as described in the “Experimental” section.

## Results and Discussion

### Vesicle stability at ambient temperature

230 First, we studied the two model systems, capric acid and C10 mix, at room temperature  $T = 25\text{ °C}$ . Because only the general features (vesicle presence and amount, bilayer thickness, micelle presence and amount) were investigated at this stage, the samples were extruded through a 100 nm polycarbonate membrane, in order to obtain a suspension of monodisperse vesicles. We verified that the use of bicine as buffer, commonly used in previous work with similar systems<sup>16,18,23,24</sup>, did not

235 affect vesicle stability and properties at ambient temperature. This was done by screening a number of concentrations and pH conditions of the samples in buffered solution with Static Light Scattering (SLS) (Figure 1S, Supp. Info). Our results show a well defined vesicle population only at the lowest measured pH (7.8) and the highest concentration (80mM) data for the capric acid sample, while in the C10 mix the same is found also at higher pH (up to 11.5) and lower concentration (down to 40mM).

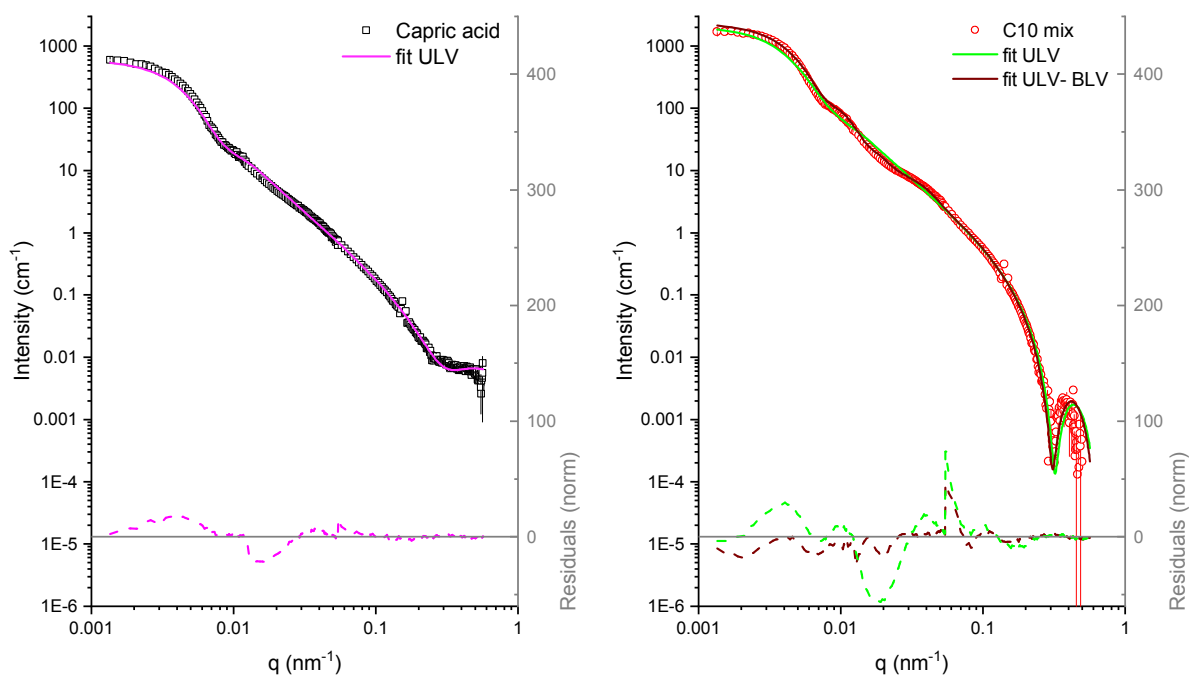
240 These are qualitatively consistent with what was previously found for the  $pK_a$  and the CVC of similar samples (pure capric acid and C10 mix 10:1 acid: alcohol ratio) that were suspended and titrated in water<sup>5,7</sup>. The bicine buffer did not lead to any significant variation in the behaviour of the model vesicles.

Dynamic Light Scattering (DLS) was then used as a tool to disentangle populations of very different

245 sizes (from tens of nanometers to micrometers), allowing for the detection of the presence of micelles and vesicles. By fitting the autocorrelation curves considering one (or two) populations of scatterers, we can show that under very alkaline conditions ( $pH \gg pK_a$ ), micelles are detected in pure capric

acid samples while for the C10 mix only vesicles are visible at every measured condition (pH, T) (Figure 2;

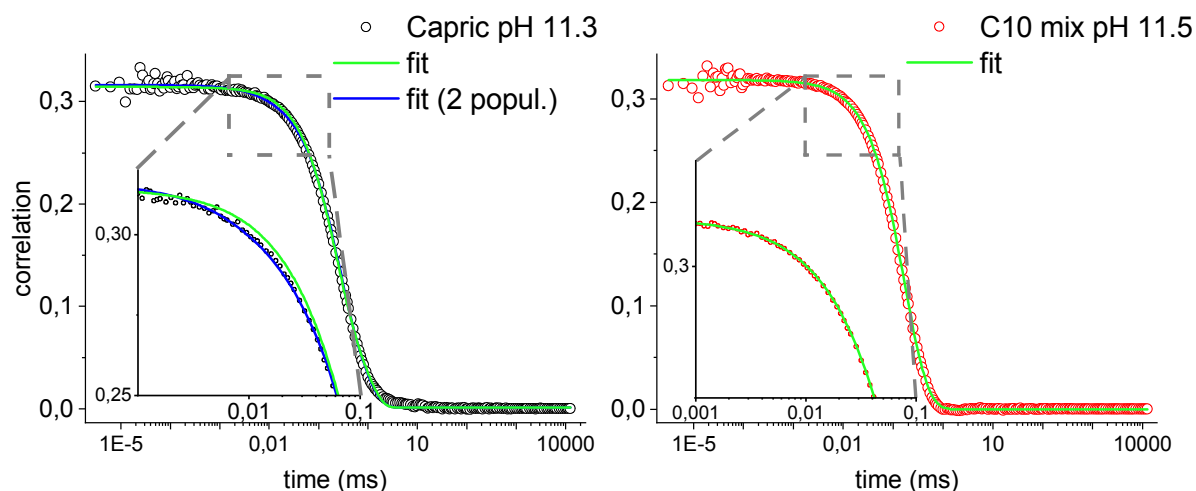
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**Figure 3.** SANS curves of capric acid sample 80 mM (left) and C10 mix 80 mM (right) at 25 °C collected on D33, and corresponding fits. Sample pH: 7.3 (capric acid) and 8.5 (C10 mix). For C10 mix, two models are compared and show that the pure unilamellar vesicle (ULV) model fails in describing correctly the features in the data, therefore a fraction of bilamellar vesicles (BLV) has to be added. The dotted lines are the normalised residuals in linear scale (right hand axis). The error bars are displayed only every 5 data points, for clarity.

data at pH 7.8 and 8.1 are shown in Figure 2S, Supp. Info). We found for the micelles a mean diameter of 1-2 nm, consistent with the expected size of capric acid micelles.



260 We thus kept a constant pH of 7.3 for the capric acid (to keep it nearby its  $pK_a$ <sup>25</sup>) and 8.5 for the C10

mix for the following Small-Angle Neutron Scattering (SANS) measurements. The vesicles made of pure capric acid present common features with those formed with the C10 mix, as found from our SANS results and fits to the data (Figure 3 and Table 1; see the “Materials and methods” section for the fitting model formulas and Table 1S for all used parameters, Supp. Info):

265 1. Both have vesicles with much smaller sizes than the extrusion pore diameter of 100 nm (59.6 nm and 80.6 nm diameter composed by 2 x (core radius + shell thickness)), although the C10 mix vesicles have a larger mean radius;

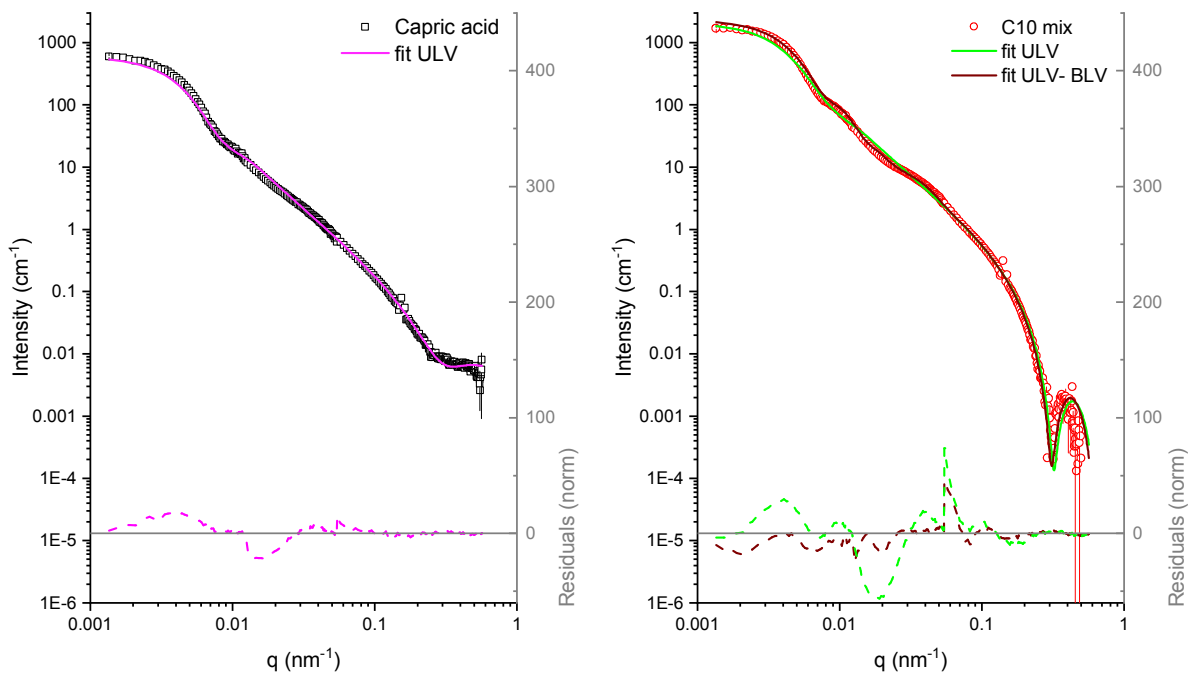
2. A membrane thickness of about 2 nm, as expected for a bilayer of amphiphiles with 1 nm chain length.

270 The following major differences are visible and inform about the substantial impact of the alcohol insertion. First, because the SANS data are very sensitive to the vesicle lamellarity, it is observed that, while the capric acid vesicles are always unilamellar, the C10 mix presents a fraction of bilamellar vesicles (BLV). Although the extrusion performed in this experiment could have interfered with the capability to form multilamellar structures in both samples, the data suggest a higher propensity of the  
275 latter (C10 mix) to form membrane stacks. This will be further confirmed also on non-extruded samples, as discussed in the next section.

The observed bilamellar vesicles differ substantially from ordinary phosphatidylcholine (PC) MLVs for what concerns the water layer thickness (11.8 nm for the C10 mix, Table 1) between the lamellae, and their characteristics resemble more the typical values found at maximum swelling for charged  
280 lipids in pure water <sup>29</sup>, while for PC lipids the solvent layer thickness is usually within 2-3 nm (DMPC, DPPC, DOPC, POPC) <sup>26-28</sup>.

As another main difference, according to the fit to the SANS data in absolute scale, we find that only  $\approx 30\%$  of the sample in the pure capric acid solution contributes to the vesicle form factor, even if it is close to its  $pK_a$ . This gives a second piece of indirect evidence about the main fraction of the FA ( $\approx$   
285 65-70%) remaining in solution as monomers or self-assembling into micelles. For the C10 mix sample, we observe that the totality of the sample (i.e. both the capric acid and the decanol),

assembles into the vesicles.



290 We conclude that the  
brings many  
vesicle self assembly  
been found previously  
of the FA CVC and the  
range of vesicle  
295 by <sup>7,18</sup>). The totality of  
mixture assembles into vesicles and leads to bigger, and in some cases multilayered, structures. No  
sign of micelles or monomers in solution (FA nor alcohol) is detected in the C10 mix, as opposed to  
the pure capric acid sample where only a small sample fraction aggregated into vesicles.

	Capric acid	C10 mix	
<i>scale (total)</i>	0.34	1	
<i>population fraction</i>	1	0.6 (ULV)	0.4 (BLV)
<i>core radius (nm)</i>	28.0	38.3 (ULV)	9.2 (BLV)
<i>shell thickness (nm)</i>	1.8	2.0 (ULV)	2.0 (BLV)
<i>solvent thickness (nm)</i>	-----	-----	11.8 (BLV)

addition of alcohol  
consequences on the  
other than what has  
(namely the decrease  
broadening of the pH  
existence as observed  
the FA/alcohol

300

**Table 1.** Main parameters extracted from the SANS data fitting at ambient temperature. Sample pH: 7.3 (capric acid) and 8.5 (C10 mix). Note, in particular, the scale factor of 0.34 calculated for the capric acid, a sign that only a fraction of sample contributes to the vesicle form factor. Full parameter list in Table 1S (Supp. Info).

305

	Capric acid	C10 mix	
<i>scale (total)</i>	0.34	1	
<i>population fraction</i>	1	0.6 (ULV)	0.4 (BLV)
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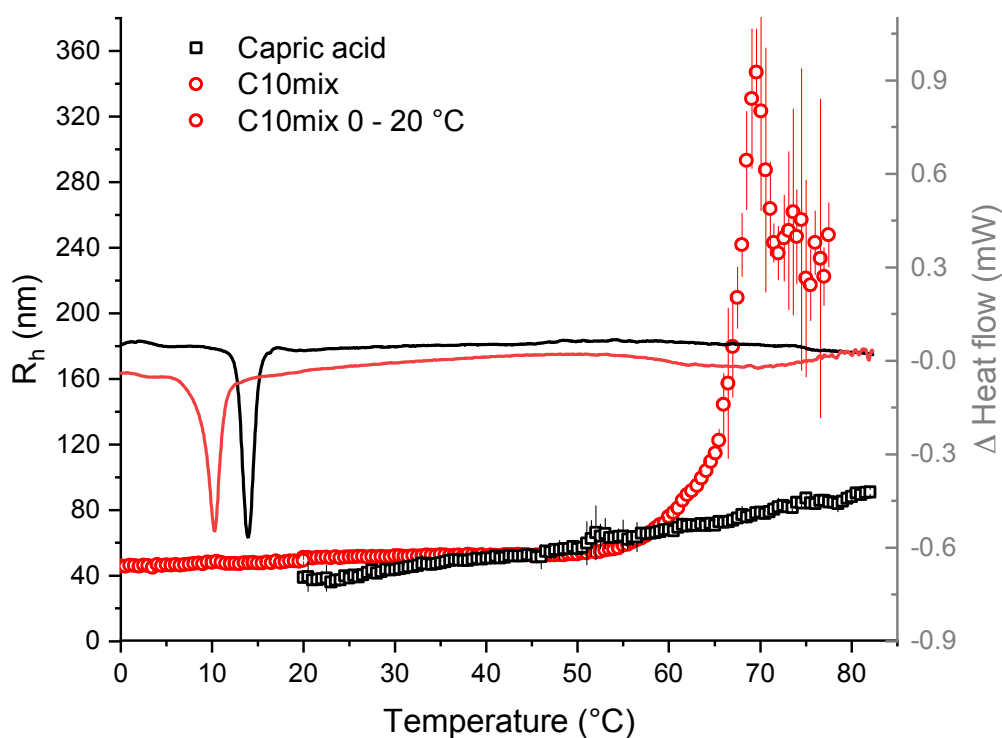
### Thermal stability

The vesicle thermal behaviour was explored by means of Differential Scanning Calorimetry (DSC, solid lines in Figure 4) and DLS (symbols in Figure 4) to explore membrane phase transitions (DSC) and follow vesicle dimension evolution (DLS). For all the experiments in this section, the pH of capric acid and C10 mix were set to 7.7 and 8.5, respectively. Because the focus here was to screen the mean vesicle size with temperature and to detect the gel-fluid membrane phase transition, the samples were extruded prior to the experiments.

The DSC thermogram of C10 mix vesicles shows a first sharp transition at 10.3 °C assigned to the gel-fluid membrane phase transition and a second, much broader one, detected above 60 °C. This result is consistent with what was observed previously<sup>18</sup>. The latter transition, for  $T \geq 60$  °C, was interpreted by<sup>18</sup> as a conformational micelle-vesicle equilibrium change. Here we additionally investigated pure capric acid vesicles and observed a single sharp transition at 13.9 °C but no sign of a second broad one.

For this reason, the high-temperature broad transition does not seem linked to a vesicle – micelle equilibrium shift as previously proposed.

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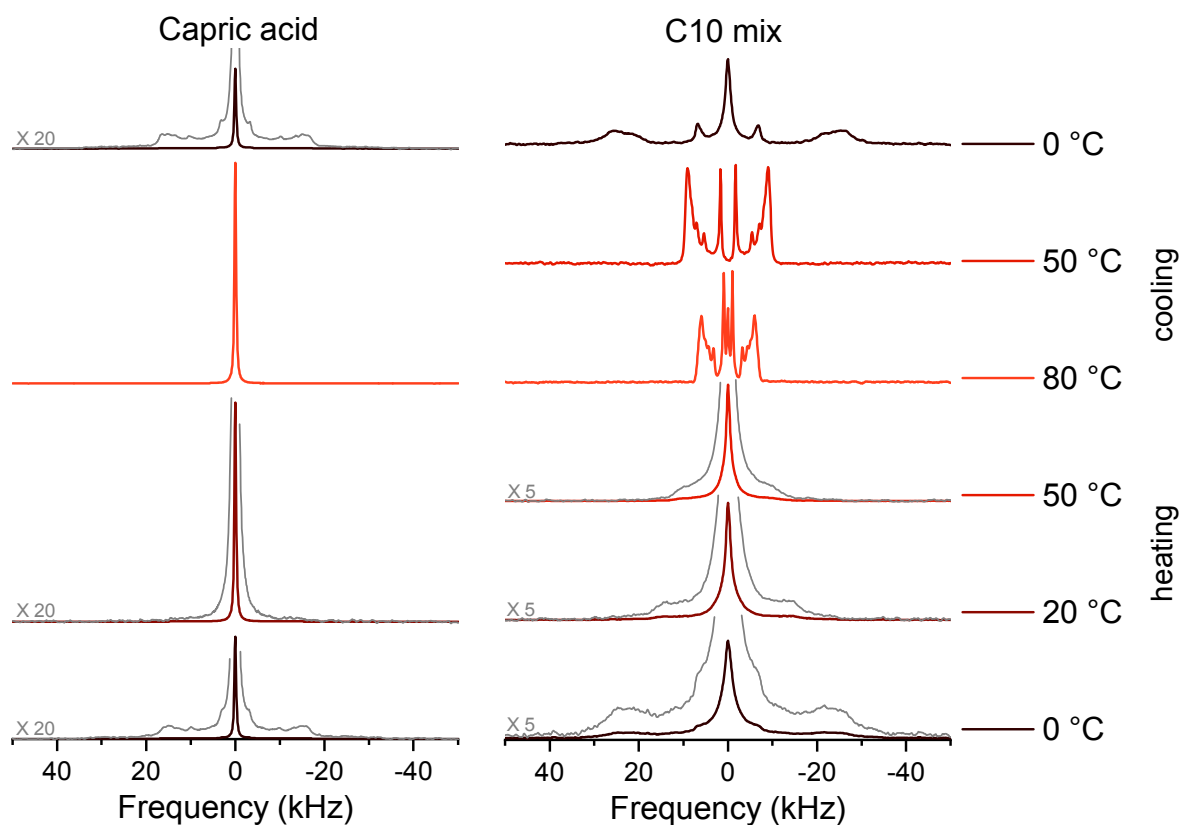


**Figure 4.** Left axis: Resulting  $R_h$  found from fitting the DLS autocorrelation functions at a fixed angle ( $90^\circ$ ) as function of the temperature, for the pure capric acid at 80 mM (black squares) and C10 mix at 80 mM (red circles). The additional dataset collected for C10 mix in the range 0 – 20 °C proves that there is no effect of the gel-fluid phase transition in the vesicle size, as expected. Right axis: DSC thermograms of the same samples (red and black lines).

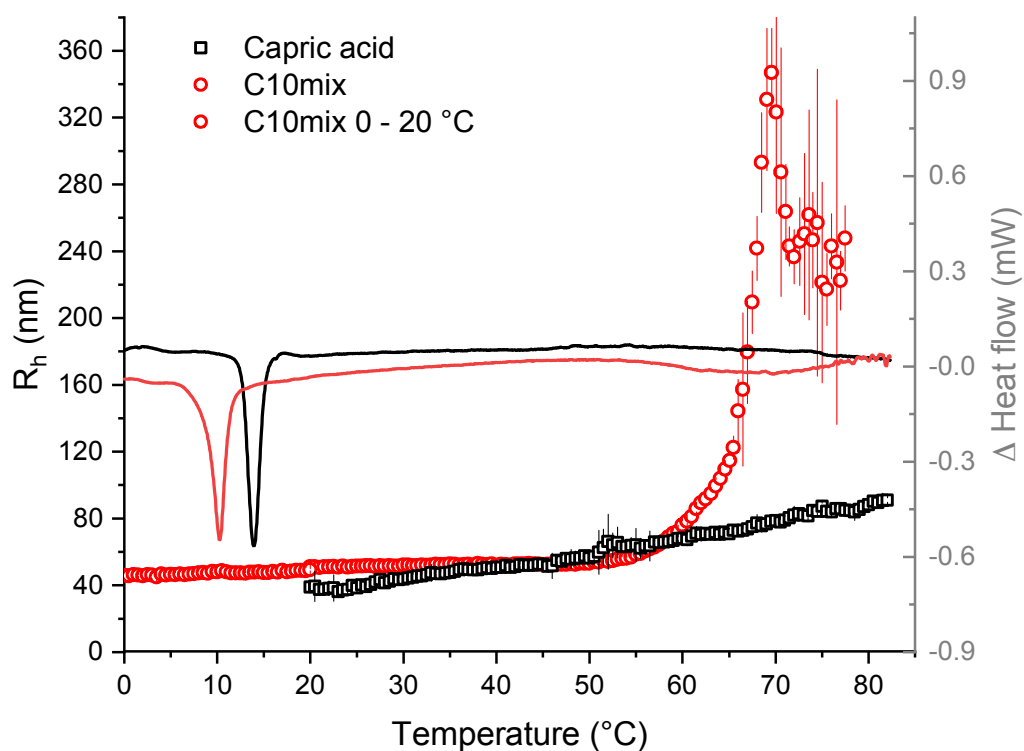
The capric acid sample is in fact the only one giving evidence for the presence of micelles (observed directly in our DLS data and indirectly in our SANS data as a lack of the expected vesicle signal, Figure 2 and Table 1). Furthermore, our DLS data vs T (Figure 4, symbols) show a corresponding, steep increase in the mean vesicle radius occurring at the same temperature and only for the C10 mix. This, together with the marked difference with the pure capric acid vesicles, suggests that the observed high temperature effect is an alcohol triggered phase transition towards bigger macrostructures.

The thermal behaviour of the samples was also addressed by performing solid-state (ss)  $^2\text{H}$ -NMR and SANS as a function of temperature for capric acid and C10 mix water dispersions. In this case, the samples were not extruded to avoid interferences in the samples lamellarity. ssNMR was used to bring quantitative structural and dynamic information about the dispersions, as has been accomplished long ago on liquid-crystal-like assemblies such as biological membranes, lipids, colloids and surfactants dispersed in water<sup>30–32</sup>. In our case, perdeuterated capric acid was used in order to probe the membrane dynamics (experimental details and ssNMR data analysis can be found in the Supp. Info).

Figure 5 shows some wide-line spectra recorded for a few temperatures for capric acid and C10 mix water dispersions.



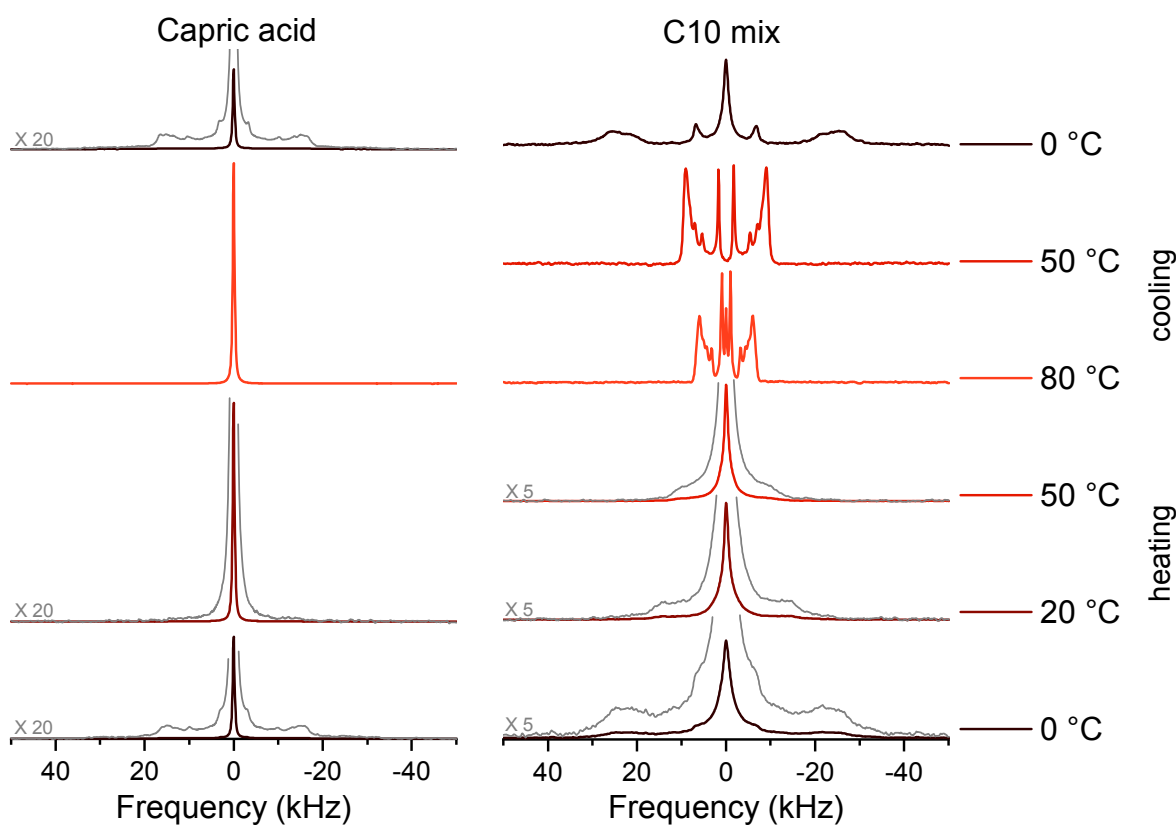
**Figure 5.**  $^2\text{H}_{19}$ -ssNMR spectra of  $^2\text{H}_{19}$ -capric acid 350 mM (left) and  $^2\text{H}_{19}$ -capric acid:decanol 350 mM (right) water dispersions (94% and 95% hydration, respectively) as a function of temperature, from bottom to top, on increasing the temperature from 0 to 80 °C and cooling back to 50 and 0 °C. Grey solid lines show details of the baseline by multiplying the vertical scale by 20 (left) and 5 (right).



In the capric acid data (left of Figure 5), an isotropic very intense Lorentzian line dominates most spectra from 0 °C to 80 °C indicative of small size objects (1-100 nm) tumbling very fast compared to the ssNMR time scale<sup>33</sup>. Close attention to the spectrum at 0 °C allows detection of a so-called “powder” (non-oriented) pattern indicating that slowly tumbling much larger objects (microns to submicron size) are also identified. At the other temperatures, there is no sign of large objects in the capric acid sample.

Conversely, in the case of the C10 mix (right set of curves in Figure 5) the coexistence of broad powder patterns and isotropic lines is observed at all temperatures, and the latter is much less abundant as was the case in the former sample. This points again towards a higher propensity of the alcohol-enriched vesicles to maintain the big multilayer structures. The hydrophobic bilayer thickness ( $b_H$ ) at  $T = 20$  °C, calculated from the NMR spectral simulations assuming no chain interdigitation, was  $b_H = 2 * L_{chain} = (1.66 \pm 0.08)$  nm (Table 3S, Supp. Info). This value is slightly lower than the one found from SANS fitting,  $b = (2.0 \pm 0.1)$  nm (Table 1 and Table 1S, Supp. Info), and the discrepancy may reflect twice the size of the small FA headgroup,  $(b - b_H)/2 = (2.0 - 1.66)$  nm/2  $\approx 0.17$  nm.





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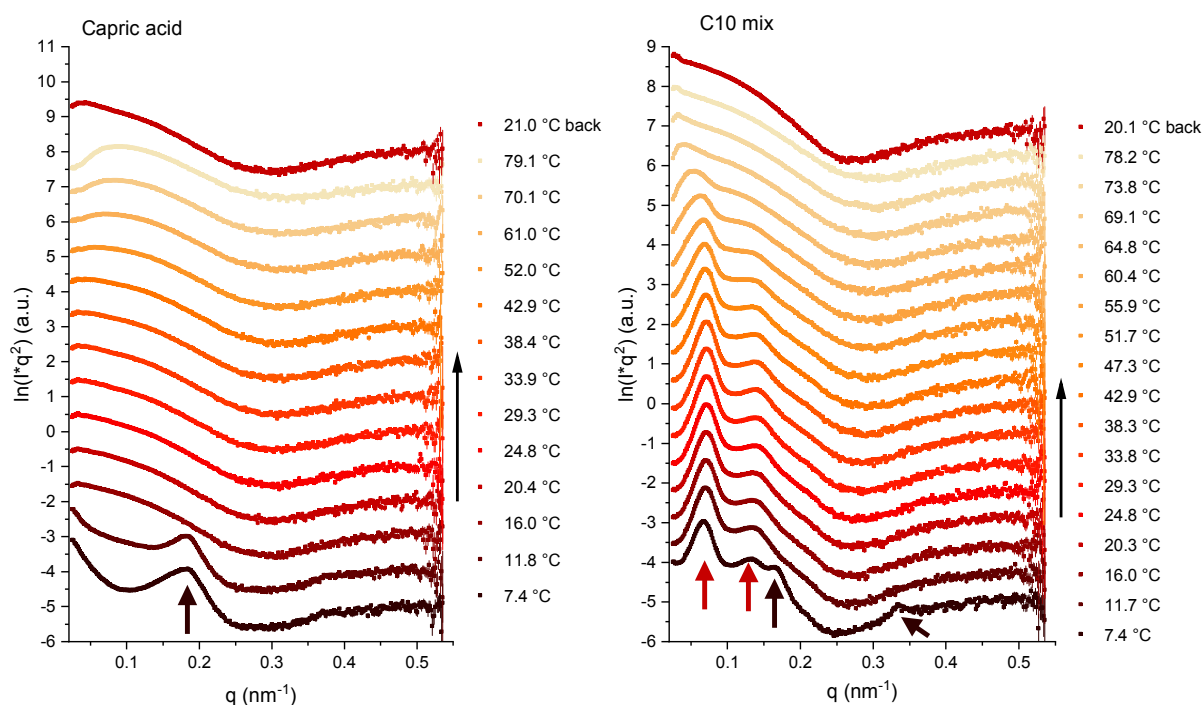
nally, the spectrum of the C10 mix at 80 °C (hence above the conformational transition seen at  $T \geq 60$  °C, Figure 4) shows an oriented-like pattern (bilayer normal perpendicular to the magnetic field of the spectrometer). Returning to 50 °C shows the same but wider oriented-like pattern, that differs considerably (the isotropic line is absent) from that observed on increasing temperature. Returning further to 0 °C shows a superimposition of powder pattern and isotropic lines with, however, a lower proportion (Figure 5S, Supp. Info). This clearly indicates a hysteresis related to the thermal history of the sample suggesting intense thermally induced sample reorganization.

375 We questioned whether the addition of decanol triggered a transition from a lamellar to a non-lamellar phase, e.g. hexagonal or cubic, as these can grow with a preferential direction and lead to NMR spectra such as the one observed for the C10 mix at 80 °C (inverted phases being improbable because of the low amphiphile concentration used). For that, we performed SANS experiments on the same samples.

380 At high ( $D_2O$ ) contrast, SANS is extremely sensitive to the presence and dimensions of MLVs and ULVs (LUVs and SUVs), and to small spherical or elongated micelles, as their dimensions are within the  $q$ -range of interest in this study ( $1.5 \times 10^{-3}$  to  $0.6 \text{ \AA}^{-1}$ , which corresponds in real space to dimensions of 420 to 1 nm). The multilamellar structures can be studied following the correlation between the adjacent membranes (given by the static structure factor  $S(q)$ )<sup>34</sup>. For instance, the relative position  
 385 between the oscillation' maxima, observed as broadened peaks in the SANS curves, will differ for

lamellar (positions  $n, 2n, 3n, \dots$ ,  $n$  being the order) or hexagonal ( $n, \sqrt{3}n, 2n, \dots$ ) phases<sup>35</sup>. Thus, this represents a convenient way to distinguish between lamellar and non-lamellar phases.

390 Additionally, even in the cases where there is no phase ordering and the correlations are absent (so that  $S(q) = 1$  over the entire  $q$ -range and no structure peak is observed), it is still possible to distinguish, from the decay of the curves (representing the form factor  $P(q)$ ), between planar lamellae and solid

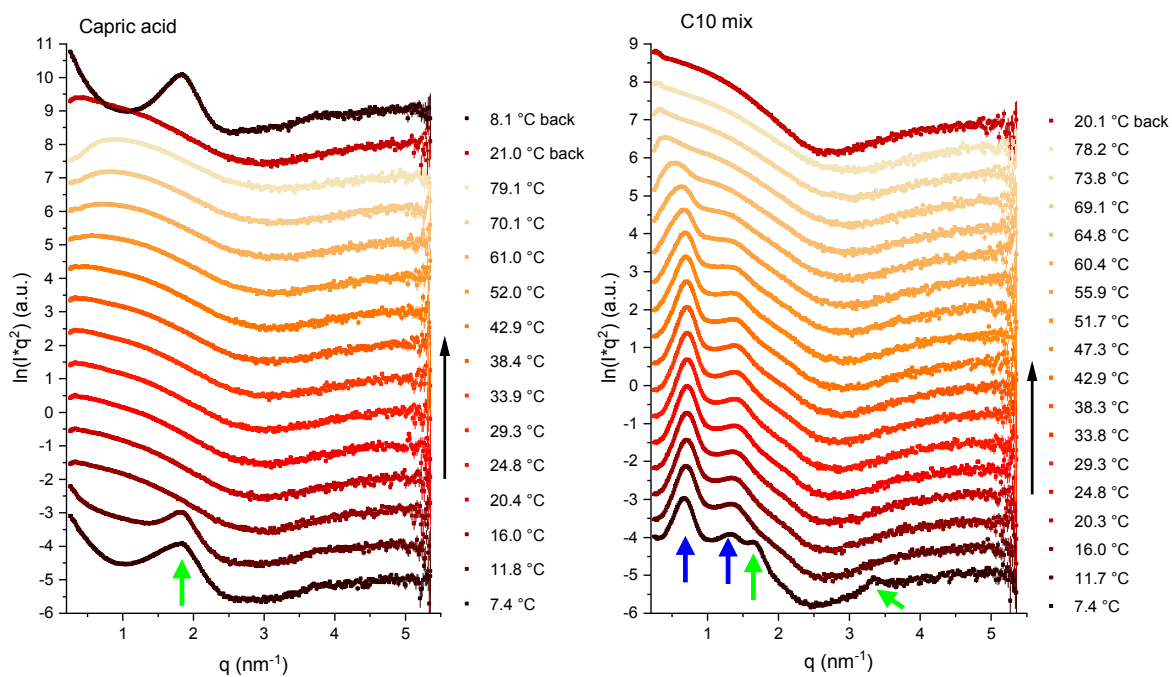


**Figure 6.** SANS curves at different temperatures collected on D16, shown in the representation  $\ln(I \cdot q^2)$  vs  $q$ , to highlight the presence of membrane correlations. This representation should flatten the curves that have the typical decay of vesicles ( $q^{-2}$ ) Left: Capric acid sample 350 mM. Right: C10 mix sample 350 mM. Arrows show the first two orders of the observed lamellar phases structure factor  $S(q)$ .

cylinders, characteristic of elongated micelles (the hexagonal phase "building blocks", diluted and free to move in solution)<sup>36</sup>. The planar lamellae would decay with a  $\approx q^{-2}$  power decay, while for cylinders the decay would follow a  $\approx q^{-1}$  decay.

400 The SANS curves were collected for both the capric acid and C10 mix samples as a function of temperature in the range 7 – 80 °C (Figure 6). All the curves show a decay close to a  $q^{-2}$  power decay, indicating the presence of vesicles. The capric acid vesicles (left plot of Figure 6) possess a correlation peak at  $q = (1.8 \pm 0.1) \text{ nm}^{-1}$  that corresponds to a  $d$ -spacing of  $(3.4 \pm 0.1) \text{ nm}$ . This correlation disappears at 16 °C and no other correlation is visible at higher temperatures. From 16 to 405 43 °C the curves are identical. From 52 °C on, a change is observed in the low  $q$  region. Remarkably, the fits performed on the capric acid curves for  $T \geq 16$  °C show the need to consider a model of vesicles coexisting with micelles to properly reproduce the experimental data (Table 4S, Supp. Info).

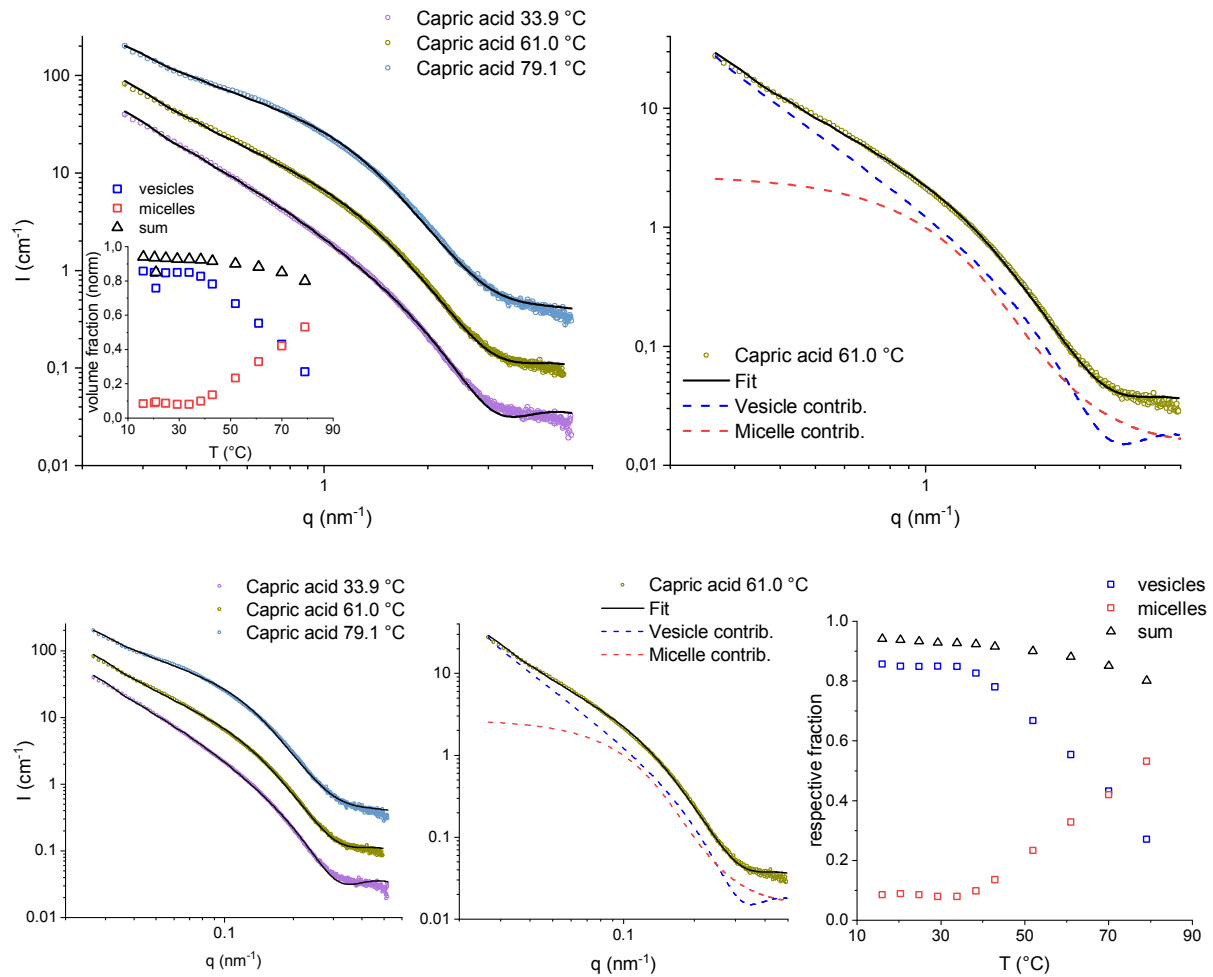
In Figure 7 some of the fits are shown, as well as the normalised respective fraction corresponding to each species (Figure 7, right panel). At  $T \leq 35$  °C,  $\approx 85\%$  of capric acid contributes to the vesicle signal while  $\approx 10\%$  goes into micelles.



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hen  $T \geq 35$  °C, a steep decrease in vesicle amount is observed, followed by a concomitant increase in the micelle signal. The sum of the two components also goes slightly down with the temperature

increase, probably due to FA solubilization.

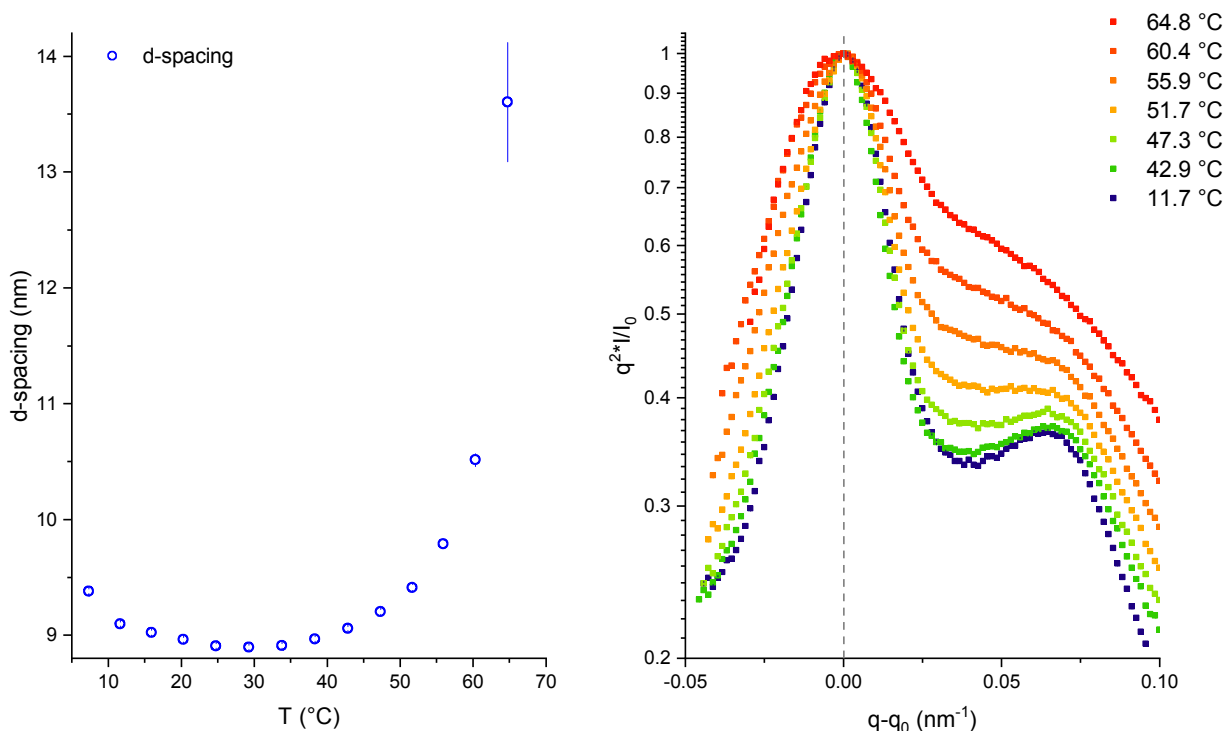


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**Figure 7.** Left: Examples of fits performed on the capric acid SANS data from D16. The error bars are displayed every 5 data points. Middle: Plot of the vesicle and micelle contributions and their sum for  $T = 61\text{ }^{\circ}\text{C}$ . Right: Respective fraction of vesicles and micelles.

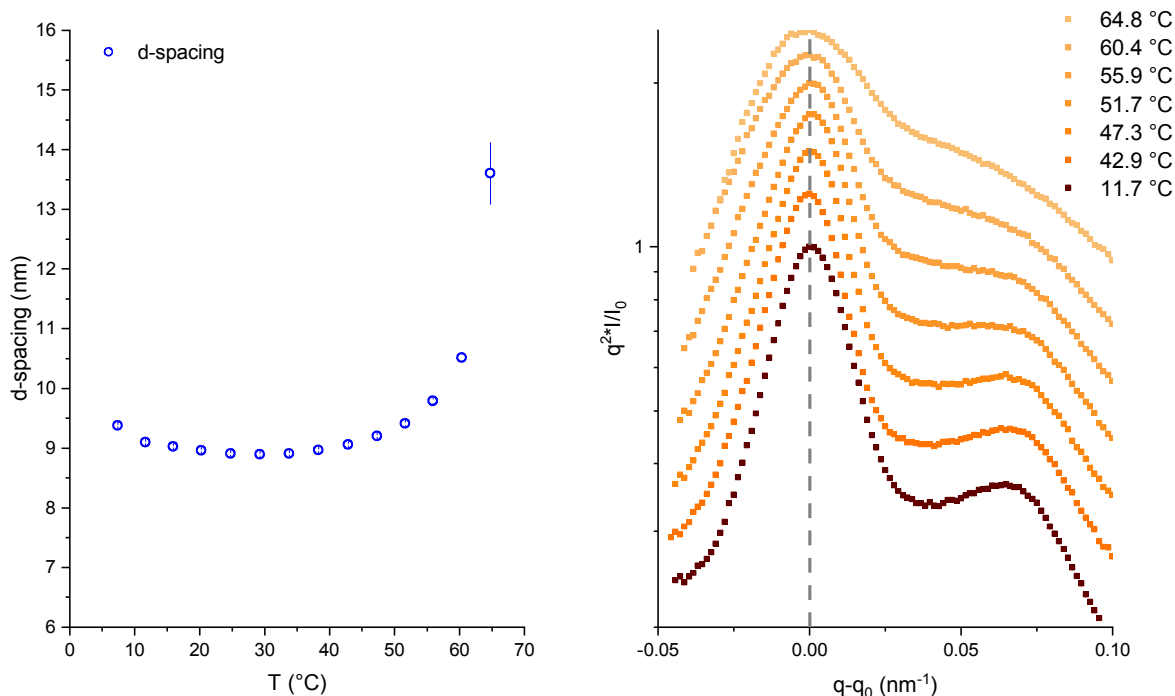
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Upon cooling back to  $20\text{ }^{\circ}\text{C}$ , the proportion of micelles and vesicles is recovered but the total amount of self-assembled material is diminished, which may be a sign of the occurrence of some phase separation. Finally, cooling down to  $8\text{ }^{\circ}\text{C}$  causes the complete recovery of the correlation peak (data not shown).



425 On the other hand, the correlation peaks observed in the C10 mix (right plot of Figure 6) show that the sample is biphasic at low temperature (7 °C): the first phase has a correlation at  $q = (0.7 \pm 0.1) \text{ nm}^{-1}$  (d-spacing of  $(9.4 \pm 0.1) \text{ nm}$ ), the second at  $q = (1.70 \pm 0.1) \text{ nm}^{-1}$  (d-spacing =  $(3.7 \pm 0.1) \text{ nm}$ ). Since the second order peaks are visible for both correlations at  $2 \cdot q_i$ , we are in the presence of two coexisting lamellar phases. Above 10 °C, only the swollen phase is visible, at  $q = (0.7 \pm 0.1) \text{ nm}^{-1}$  (d-spacing =  $(9.4 \pm 0.1) \text{ nm}$ ). The gel-fluid phase transition of the C10 mix, occurring at  $T = 10 \text{ °C}$  (DSC thermogram of Figure 4), can explain the slight decrease ( $\approx 0.5 \text{ nm}$ ) in the d-spacing of the swollen phase from 7 to 30 °C as a thinning of the bilayer upon melting (Figure 8, left). Remarkably, the d-spacing decrease of  $\Delta d = (0.42 \pm 0.03) \text{ nm}$  from 7 to 20 °C is in extremely good agreement with the decrease in hydrophobic bilayer thickness ( $\Delta b_H$ ) as calculated from the NMR spectral simulations,  $\Delta b_H = (0.44 \pm 0.11)$  (Supp. Info). Approaching the second phase change (above 50 - 60 °C), the membrane phase undergoes a fast swelling (now due to the water layer thickening between membranes) seen as a shift of the correlation towards lower  $q$ , until disappearing. This effect was observed in the literature and referred to as an unbinding transition<sup>29,37</sup>. Both the C10 mix curves at the highest temperature  $T = 78 \text{ °C}$  and after cooling back to  $T = 20 \text{ °C}$  fit extremely well with a bilayer form factor, with a  $\approx q^{-2}$  decay (Figure 9 and Table 5S, Supp. Info). This decay is much different from the one that would be observed in the case of elongated cylindrical micelles, with  $\approx 1 \text{ nm}$  radius (plotted as a grey dashed line in Figure 9). These results indicate that the high temperature structure of the C10 mix samples is still a unilamellar membrane. Interestingly, the two fitting results differ slightly in the bilayer thicknesses (1.8 nm at 85 °C, 2.1 nm at 20 °C, uncertainty 0.1 nm) and their polydispersity (0.27 at 85 °C, 0.15 at 20 °C).

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**Figure 8.** Left: d-spacing of the swollen C10 mix phase as function of T. Right: C10 mix SANS curves collected on D16, zoomed around the first lamellar correlation peak, in the representation  $q^2 \cdot I/I_0$  vs  $q - q_0$  to highlight the 1<sup>st</sup> and 2<sup>nd</sup> order peak broadenings upon heating. Some of the curves at low temperature were skipped, for clarity.

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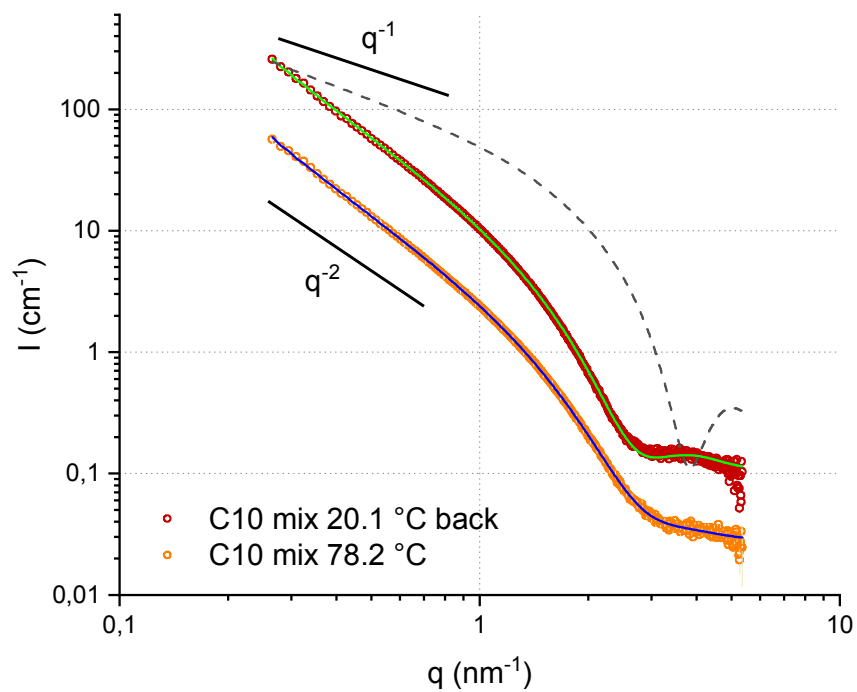
These are in line with the consideration that the molten amphiphiles' chains allow a wider range of bent/kinked configurations at higher temperatures. Most importantly, no scaling factor is required to fit the two curves, which means that no amphiphile solubilization or disaggregation happens while the temperature is increased, unlike with the pure capric acid where progressive solubilization takes place (Figure 7).

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The experimental results lead towards a high temperature, alcohol-triggered phase transition occurring above  $T = 60$  °C that leads to bigger structures still possessing a bilayer (Figure 9), which are oriented upon application of a magnetic field (Figure 5). Furthermore, the alcohol addition stabilizes the membrane by preventing a temperature induced vesicle dissolution into micelles.

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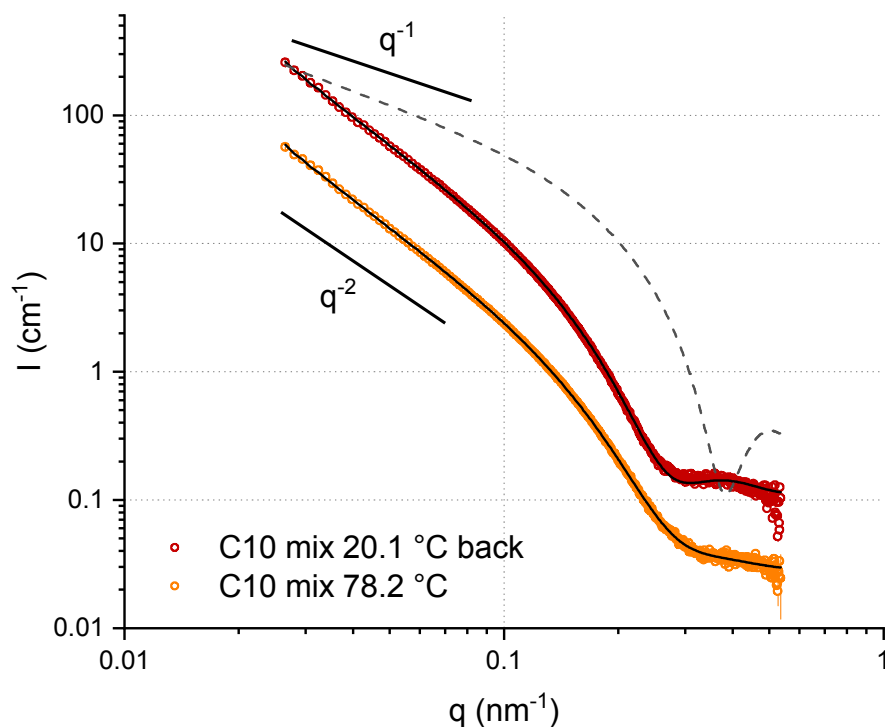
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Our data confirm that the

465 decanol addition has a marked stabilizing effect on the membrane at ambient temperature. In addition, we found additional measurable differences in the amphiphile behaviour that complement the

previous findings<sup>7,18</sup>, such as:



**Figure 9.** Results of fitting the SANS curves on D16 for C10 mix at 78.2 °C and cooled back to 20.1 °C, each with a bilayer model. The error bars are displayed every 5 data points, for clarity. The black dashed line shows the model of a system of elongated micelles with 1.0 nm radius and 100 nm length. The black straight lines show the  $q^{-1}$  and  $q^{-2}$  decay behaviours. Models considering vesicles or hollow cylinders, with core radius  $R > 20$  nm, could fit equally well the data in this  $q$  region.

- 470 a) A marked increase in the vesicle amount at the expense of micelles (Figure 2 and 3, Table 1);
- 475 b) A higher propensity to form bigger and multilamellar vesicles (Figure 3). Our data confirm that the decanol addition has a marked stabilizing effect on the membrane at ambient temperature. In addition, we found additional measurable differences in the amphiphile behaviour that complement the previous findings<sup>7,18</sup>, such as:

- a) A marked increase in the vesicle amount at the expense of micelles (Figure 2 and 3, Table 1);
- 480 b) A higher propensity to form bigger and multilamellar vesicles (Figure 3).



Apel and coworkers <sup>7</sup> interpreted the stabilizing effect of the alcohol, for the vesicle self-aggregation at ambient temperature, in terms of hydrogen bonding. In particular, they measured quantitatively that the aggregates of ionic FAs-alcohols (giving vesicles) are thermodynamically more favourable than the FA hydrogen bonds with water that induce micelle formation or favour monomers in solution.

485 This interpretation can explain our observed increase of the vesicle population at the expense of the micelles. However, this does not explain the higher propensity to form bigger, and even multilamellar structures. Our data imply that the alcohol must also affect considerably the amphiphile packing parameter <sup>38</sup>. Supposing the intrinsic geometry of the ionic FA as conic (because of the electrostatic repulsion of the anionic headgroups), the hydrogen bond with a neutral FA molecule leads to a

490 geometry which is approaching the cylindrical one and therefore can allow for the formation of flat or low curvature aggregates like vesicles. Nevertheless, it is the ionic FA-alcohol aggregates that lead most effectively to a cylindrical intrinsic geometry, so that planar structures are more easily formed and much bigger multilamellar vesicles can be produced.

It is in the thermal behaviour that the two systems, differing by the presence of the alcohol, show the

495 most remarkable distinctions, and some of them are found in the low temperature region. As shown in Figure 6, we observe for the capric acid at  $T < 16$  °C a single lamellar phase with a d-spacing =  $(3.4 \pm 0.1)$  nm, at a much lower value than that expected for the membrane correlation in a MLV. On the other hand, at  $T < 11$  °C the C10 mix shows a coexistence of such collapsed phase (with slightly different d-spacing =  $(3.7 \pm 0.1)$  nm) with the expected MLV swollen phase (d-spacing =  $(9.4 \pm 0.1)$

500 nm). We interpret the occurrence of the collapsed phase as the probable crossing of the corresponding FA Krafft temperature (Figure 6S, Supp. Info), known as the temperature below which the surfactant solubility is below its critical micelle concentration (CMC), causing flocculation. Therefore, the observed transition at low temperature of the capric acid (with its effect visible in the data in Figure 4, 5 and 6) is due to the flocculation/solubilization (as monomers, or assembled into small

505 vesicles/micelles) rather than to a gel-fluid membrane transition. The coexistence observed in Figure 6 (right) indicates that the decanol has the effect of partially preventing flocculation, so that a MLV phase is still present, with the amphiphiles in the gel phase so to give a thicker membrane (lowest temperatures in Figure 8 left). Upon chain melting at  $T \geq 10$  °C, the d-spacing decreases slightly ( $\leq 0.5$  nm), due to a decrease of the bilayer thickness when crossing the gel-fluid transition, as also

510 observed by <sup>17</sup>. From  $T = 30$  °C, the solvent induced swelling becomes more important and d-spacing starts to increase.

Most importantly in this study, the mixture containing decanol undergoes a high temperature phase transition, not observed in the pure FA system, which leads to much bigger (Figure 4), orientable (Figure 5), yet still bilayer (Figure 9) structures. High temperatures also have the effect of causing an

515 unbinding transition of the MLV swollen phase ( $T \approx 65 - 70$  °C) that leads to unilamellar membranes (Figure 6 and 9), only observed in the case of the C10 mix.

The peculiar oriented-like spectral shapes observed with NMR at 80 °C and 50 °C (Figure 5) have already been observed for bilayer-micelles (bicelles) when the system is under the form of small lipidic discs that align as a nematic phase in the magnetic field<sup>39-41</sup>. It has been also observed for model and natural membranes (MLVs) that are deformed by magnetic fields<sup>42,43</sup>. This happens when the elastic energy maintaining the vesicles with an average spherical shape is lower than the magnetic energy<sup>43</sup>. In such a situation the vesicle is deformed as a prolate ellipsoid along the magnetic field direction, the elastic modulus of the membrane being orders of magnitude lower than that of classical plasma membranes, which are not affected by magnetic fields. Because vesicles are being detected with neutrons we favour the latter explanation. As pointed out by the work of<sup>43</sup>, the extent of the huge deformation observed here (long-axis : short-axis ratio  $\approx$  8-10, see Supp. Info) is directly proportional to the vesicle size. Therefore the observed high temperature transition that leads to bigger structures has likely favoured the occurrence of such deformation. Nevertheless, the membrane packing of the C10 mix vesicles is also stable at the highest measured temperature (Figure 4S, Supp. Info).

Taken together, these results point towards the appearance of large, unilamellar vesicles at  $T \geq 60$  °C formed by fusion of the vesicles existing at low temperature. This interpretation is supported by the observed vesicle mean size increase (Figure 4), the MLV unbinding to give unilamellar vesicles (Figure 6, right) and the decrease of the small ( $< 100$  nm) vesicle fraction until complete disappearance (Figure 5, right and Figure 5S, Supp. Info). The second high-temperature broad transition observed in the DSC thermogram of C10 mix (Figure 4) is therefore a result of such fusion process, rather than a shift of the equilibrium between vesicles and micelles as previously reported<sup>18</sup>.

We note that the occurrence of this high temperature conformational transition has probably been observed with its effect in previous studies, although not directly identified:

(i) Jordan and coworkers<sup>14</sup>, by studying a different system made with mixtures of C10 – C15 FA/alcohols at strongly alkaline pH, high salt concentrations (NaCl, MgCl<sub>2</sub>, CaCl<sub>2</sub>) and high temperature  $T = 70$  °C with confocal microscopy, observed in several cases the occurrence of big elongated structures, where “the filaments appeared to be composed of chains of individual vesicles”. Although on different systems and in non-equilibrium conditions (the sample stage was not heated during measurements), their observation could be the effect of the high temperature transition described here. Although we observed oriented structures with ssNMR at high temperature, which could be because of an elongated shape of the bilayer structures, we could not conclude whether the *orientation* or also the *deformation* itself were triggered by the static magnetic field of the spectrometer. For this reason we refer to the high temperature products as large unilamellar vesicles, with a not necessarily spherical mean shape that is not the focus of the present study.

(ii) Mansy and Szostak<sup>16</sup> examined a mixture of capric acid : decanol in 2:1 ratio by measuring the vesicle capability to retain encapsulated oligonucleotides as a function of temperature. They observed

the occurrence of material leakage above 50 – 60 °C. In view of our results, we can argue that the observed leakage may not be due simply to a sudden increase in the protomembrane permeability, but a consequence of the morphological changes happening above that temperature threshold: the conformational transition and, perhaps more importantly in this case, the unbinding transition of the bilamellar vesicle fraction (probably present as the experimental buffer/pH/extrusion size conditions were the same as in our study). From this perspective, the actual permeability of the model protomembrane may still be sufficiently low even at higher temperatures, thus additional studies are required to address this point taking into account the MLV unbinding transition and the vesicle fusion processes that we identified.

The occurrence of this transition at  $T \geq 60$  °C could appear as unfavourable in terms of the C10 mix stability at the highest temperatures, the focus of our study given the implications for the origin of life. The pure FA vesicles, although less stable at room temperature, would not undergo the high T transition and hence appear more thermostable after all. Nevertheless, our results (Figure 7) show that the temperature increase induces a partial vesicle disassembling in the pure capric acid that leads to a higher amount of micelles. Conversely, in the C10 mix, no vesicle-to-micelle transition is observed, and the bilayer structuring (probed by SANS) and packing (probed by NMR) are preserved at the highest temperatures. Therefore, rather than protomembrane disruption, solubilization or micellar aggregation, the high temperature triggers the fusion of the small vesicles assembling in favour of large unilamellar bilayered structures.

## Conclusion

In summary, we found that the addition of same-length fatty alcohol to a short chain C10 FA suspension in equimolar ratio leads to several substantial effects. It improves the vesicles' stability at room temperature at the expense of the other amphiphile states (micelles, monomers). It partially prevents the vesicles from undergoing flocculation at temperatures lower than the expected FA Krafft point. Most importantly, it triggers a high temperature conformational change at  $T \geq 60$  °C that leads to vesicle fusion with stable membranes at temperature as high as 80 °C, the highest measured in this study. On one hand, this sets a warning on the interpretation of the results from all experiments of this kind (i.e. whether what is observed correlates with the occurrence of such phase transition at high temperature). In addition, as a perspective, it sets the basis for further studies: namely, whether such transition could serve as a prebiotic mean of solute exchange between different vesicles, as this would set a milestone in our understanding of origin of life related processes.

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This article is dedicated to our late colleague Isabelle Grillo

**Supporting Information.** Equations for the SANS model form factors; preliminary sample screening data; full list of parameters used for fits/simulation and annexes.

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