

Bounds for eigenvalues of the Dirichlet problem for the logarithmic Laplacian

Huyuan Chen, Laurent Véron

▶ To cite this version:

Huyuan Chen, Laurent Véron. Bounds for eigenvalues of the Dirichlet problem for the logarithmic Laplacian. 2020. hal-02977991v2

HAL Id: hal-02977991 https://hal.science/hal-02977991v2

Preprint submitted on 9 Nov 2020 (v2), last revised 4 Nov 2021 (v3)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Bounds for eigenvalues of the Dirichlet problem for the logarithmic Laplacian

Huyuan Chen¹

Department of Mathematics, Jiangxi Normal University, Nanchang, Jiangxi 330022, PR China

Laurent Véron²

Laboratoire de Mathématiques et Physique Théorique, Université de Tours, 37200 Tours, France

Abstract

We provide bounds for the sequence of eigenvalues $\{\lambda_i(\Omega)\}_i$ of the Dirichlet problem

$$L_{\Delta}u = \lambda u \text{ in } \Omega, \qquad u = 0 \text{ in } \mathbb{R}^N \setminus \Omega,$$

where L_{Δ} is the logarithmic Laplacian operator with Fourier transform symbol $2 \ln |\zeta|$. The logarithmic Laplacian operator is not positively definitive if the volume of the domain is large enough, hence the principle eigenvalue is no longer always positive. We also give asymptotic estimates of the sum of the first k eigenvalues. To study the principle eigenvalue, we construct lower and upper bounds by a Li-Yau type method and calculate the Rayleigh quotient for some particular functions respectively. Our results point out the role of the volume of the domain in the bound of the principle eigenvalue. For the asymptotic of sum of eigenvalues, lower and upper bounds are built by a duality argument and by Kröger's method respectively. Finally, we obtain the limit of eigenvalues and prove that the limit is independent of the volume of the domain.

Contents

1	Introduction and main results	1		
2	Upper bounds for the principle eigenvalue 2.1 Large domain: proof of Theorem 1.2-(i)			
3	Lower bounds	11		
4	Wey's limits 4.1 Upper bounds for the sum of eigenvalues	17		
K	eywords: Dirichlet eigenvalues; Logarithmic Laplacian.			
\mathbf{M}	MSC2010: 35P15; 35R09.			

1 Introduction and main results

Let L_{Δ} be the logarthmic Laplacian in \mathbb{R}^N , $N \geq 1$, defined by

$$L_{\Delta}u(x) = c_N \int_{\mathbb{R}^N} \frac{u(x)1_{B_1(x)}(y) - u(y)}{|x - y|^N} dy + \rho_N u(x), \tag{1.1}$$

¹chenhuyuan@yeah.net

²veronl@univ-tours.fr

where

$$c_N := \pi^{-N/2} \Gamma(N/2) = \frac{2}{\omega_{N-1}}, \qquad \rho_N := 2 \ln 2 + \psi(\frac{N}{2}) - \gamma,$$
 (1.2)

 $\omega_{N-1}:=H^{N-1}(S^{N-1})=\int_{S^{N-1}}dS,\,\gamma=-\Gamma'(1)$ is the Euler Mascheroni constant and $\psi=\frac{\Gamma'}{\Gamma}$ is the Digamma function.

The aim of this article is to provide estimates of the eigenvalues of the operator L_{Δ} in a bounded domain $\Omega \subset \mathbb{R}^N$ which are the real numbers λ such that there exists a solution to the Dirichlet problem

$$\begin{cases} L_{\Delta} u = \lambda u & \text{in } \Omega, \\ u = 0 & \text{in } \mathbb{R}^N \setminus \Omega. \end{cases}$$
 (1.3)

In recent years, there has been a renewed and increasing interest in the study of boundary value problems involving linear and nonlinear integro-differential operators. This growing interest is justified both seminal advances in the understanding of nonlocal phenomena from a PDE or a probabilistic point of view, see e.g. [3-6, 14, 15, 21, 32, 35, 36] and the references therein, and by important applications. Among nonlocal differential order operators, the simplest and most studied examples, are the fractional powers of the Laplacian which exhibit many phenomenological properties. Recall that, for $s \in (0,1)$, the fractional Laplacian of a function $u \in C_c^{\infty}(\mathbb{R}^N)$ is defined

$$\mathcal{F}((-\Delta)^s u)(\xi) = |\xi|^{2s} \widehat{u}(\xi)$$
 for all $\xi \in \mathbb{R}^N$,

where and in the sequel both \mathcal{F} and $\widehat{\cdot}$ denote the Fourier transform. Equivalently, $(-\Delta)^s$ can be written as a singular integral operator under the following form

$$(-\Delta)^{s} u(x) = c_{N,s} \lim_{\epsilon \to 0^{+}} \int_{\mathbb{R}^{N} \setminus B_{\epsilon}(x)} \frac{u(x) - u(y)}{|x - y|^{N + 2s}} dy, \tag{1.4}$$

where $c_{N,s}=2^{2s}\pi^{-\frac{N}{2}}s\frac{\Gamma(\frac{N+2s}{2})}{\Gamma(1-s)}$ and Γ is the Gamma function, see e.g. [36]. The fractional Laplacian has the following limiting properties when s approaches the values 0

and 1:

$$\lim_{s \to 1^{-}} (-\Delta)^{s} u(x) = -\Delta u(x) \quad \text{and} \quad \lim_{s \to 0^{+}} (-\Delta)^{s} u(x) = u(x) \quad \text{for } u \in C_{c}^{2}(\mathbb{R}^{N}),$$

see e.g. [14]. Recently, [8] shows a further expansion at s=0 that for $u\in C_c^2(\mathbb{R}^N)$ and $x\in\mathbb{R}^N$,

$$(-\Delta)^s u(x) = u(x) + sL_{\Delta}u(x) + o(s)$$
 as $s \to 0^+$

where, formally, the operator

$$L_{\Delta} := \frac{d}{ds} \Big|_{s=0} (-\Delta)^s \tag{1.5}$$

is given as a logarithmic Laplacian; indeed,

(i) for $1 , we have <math>L_{\Delta}u \in L^p(\mathbb{R}^N)$ and $\frac{(-\Delta)^s u - u}{s} \to L_{\Delta}u$ in $L^p(\mathbb{R}^N)$ as $s \to 0^+$;

(ii)
$$\mathcal{F}(L_{\Delta}u)(\xi) = 2 \ln |\xi| \, \widehat{u}(\xi)$$
 for a.e. $\xi \in \mathbb{R}^N$.

Note that the problems with integral-differential operators given by kernels with a singularity of order -N have received growing interest recently, as they give rise to interesting limiting regularity properties and Harnack inequalities without scaling invariance, see e.g. [24]. Another important domain of study consists in understanding the eigenvalues of the Dirichlet problem with zero exterior value [8]. We refer to [7,19] for more topics related to the logarithmic Laplacian and also [16,23] for general nonlocal operator and related embedding results. Let $\mathbb{H}(\Omega)$ denote the space of all measurable functions $u: \mathbb{R}^N \to \mathbb{R}$ with $u \equiv 0$ in $\mathbb{R}^N \setminus \Omega$ and

$$\iint_{\substack{x,y\in\mathbb{R}^N\\|x-y|\leq 1}} \frac{(u(x)-u(y))^2}{|x-y|^N} dxdy < +\infty.$$

As we shall see it, $\mathbb{H}(\Omega)$ is a Hilbert space under the inner product

$$\mathcal{E}(u, w) = \frac{c_N}{2} \iint_{\substack{x, y \in \mathbb{R}^N \\ |x - y| \le 1}} \frac{(u(x) - u(y))(w(x) - w(y))}{|x - y|^N} dx dy,$$

where c_N is given in (1.2), with associated norm $||u||_{\mathbb{H}(\Omega)} = \sqrt{\mathcal{E}(u,u)}$. By [13, Theorem 2.1], the embedding $\mathbb{H}(\Omega) \hookrightarrow L^2(\Omega)$ is compact. Throughout this article we identify $L^2(\Omega)$ with the space of functions in $L^2(\mathbb{R}^N)$ which vanish a.e. in $\mathbb{R}^N \setminus \Omega$. The quadratic form associated with L_{Δ} is well-defined on $\mathbb{H}(\Omega)$ by

$$\mathcal{E}_L: \mathbb{H}(\Omega) \times \mathbb{H}(\Omega) \to \mathbb{R}, \quad \mathcal{E}_L(u, w) = \mathcal{E}(u, w) - c_N \iint_{\frac{x, y \in \mathbb{R}^N}{|x - y| > 1}} \frac{u(x)w(y)}{|x - y|^N} dx dy + \rho_N \int_{\mathbb{R}^N} uw \, dx,$$

where ρ_N is defined in (1.2). A function $u \in \mathbb{H}(\Omega)$ is an eigenfunction of (1.3) corresponding to the eigenvalue λ if

$$\mathcal{E}_L(u,\phi) = \lambda \int_{\Omega} u\phi \, dx$$
 for all $\phi \in \mathbb{H}(\Omega)$.

Proposition 1.1. [8, Theorem 1.4] Let Ω be a bounded domain in \mathbb{R}^N . Then problem (1.3) admits a sequence of eigenvalues

$$\lambda_1(\Omega) < \lambda_2(\Omega) \le \cdots \le \lambda_i(\Omega) \le \lambda_{i+1}(\Omega) \le \cdots$$

and corresponding eigenfunctions ϕ_i , $i \in \mathbb{N}$ such that the following holds:

(a)
$$\lambda_i(\Omega) = \min\{\mathcal{E}_L(u, u) : u \in \mathbb{H}_i(\Omega) : ||u||_{L^2(\Omega)} = 1\}, \text{ where }$$

$$\mathbb{H}_1(\Omega) := \mathbb{H}(\Omega) \quad and \quad \mathbb{H}_i(\Omega) := \{ u \in \mathbb{H}(\Omega) : \int_{\Omega} u \phi_i \, dx = 0 \text{ for } i = 1, \dots, i-1 \} \quad for \ i > 1;$$

- (b) $\{\phi_i : i \in \mathbb{N}\}\ is\ an\ orthonormal\ basis\ of\ L^2(\Omega);$
- (c) ϕ_1 is positive in Ω . Moreover, $\lambda_1(\Omega)$ is simple, i.e., if $u \in \mathbb{H}(\Omega)$ satisfies (1.3) in weak sense with $\lambda = \lambda_1(\Omega)$, then $u = t\phi_1$ for some $t \in \mathbb{R}$;
- (d) $\lim_{i \to \infty} \lambda_i(\Omega) = +\infty$.

Due to lack of the homogenous property for the logarithmic Laplacian operator, the effect of the domain for the principle eigenvalue can't be expected as the Laplacian or fractional Laplacian, just by scaling the domain by their homogeneous property of such operators. Secondly, the logarithmic Laplacian operator is no longer positively definitive if $|\Omega|$ is to large, since it is proved in [8] that the positivity of the principle eigenvalue is equivalent to the comparison principle, which does not hold for balls with large radius. These properties of the logarithmic Laplacian operator enrich the asymptotics of the Dirichlet eigenvalues as we will see below, but also make more difficult the obtention of bounds for eigenvalues.

Observe that the inclusion $\mathbb{H}(O_1) \subset \mathbb{H}(O_2)$ implies that the mapping $O \mapsto \lambda_1(O)$ is nonincreasing, i.e. $\lambda_1(O_1) \geq \lambda_1(O_2)$ if $O_1 \subset O_2$. Our first results deal with upper and lower bounds on the principle eigenvalue and they are connected both with the measure and the distortion of the domain. We denote by H^k the k-dimensional Hausdorff measure in \mathbb{R}^N and for simplicity $H^N(E) = |E|$ for any Borel set $E \subset \mathbb{R}^N$. We also define the signed distance function to $\partial \Omega$ by

$$\rho(x) = \begin{cases} \operatorname{dist}(x, \partial \Omega) & \text{if } x \in \overline{\Omega}, \\ -\operatorname{dist}(x, \partial \Omega) & \text{if } x \in \overline{\Omega}^c \end{cases}$$

$$(1.6)$$

and, for $\nu > 0$, the internal and external foliations of $\partial \Omega$ by

$$T_{\nu}^{+} = \left\{ x \in \mathbb{R}^{N} : \rho(x) = \nu \right\} \subset \Omega \tag{1.7}$$

and

$$T_{\nu}^{-} = \left\{ x \in \mathbb{R}^{N} : \rho(x) = -\nu \right\} \subset \Omega^{c}, \tag{1.8}$$

respectively, and $T_{\nu} = T_{\nu}^{+} \cup T_{\nu}^{-}$.

Theorem 1.2. Let Ω be a bounded domain in \mathbb{R}^N and $\lambda_1(\Omega)$ be the principle eigenvalue of Dirichlet problem (1.3) obtained in Proposition 1.1.

(i) For R > 2, if we assume that $B_R \subset \Omega \subset B_{2R}$, and that there exists $c_0 > 1$ depending only on N such that for any $\nu \in [0, \frac{1}{2})$, there holds

$$\frac{1}{c_0}R^{N-1} \le H^{N-1}(T_{\nu}^+) \le c_0 R^{N-1} \quad \text{if} \quad N \ge 2.$$
 (1.9)

Then for $R \ge \max\left\{2, \frac{Nc_0}{2\omega_{N-1}}\right\}$, we have that

$$\lambda_1(\Omega) \le \omega_{N-1} \ln \frac{1}{R} + z_1(R), \tag{1.10}$$

where

$$z_1(R) = \rho_N + \omega_{N-1} \ln 2 + \frac{4c_0}{R} \left(1 + \frac{c_0}{2\omega_{N-1}R} \right).$$

(ii) For $R \in (0, \frac{1}{4})$, if we assume again that $B_R \subset \Omega \subset B_{2R}$, and that there exists $c_0 > 1$ such that for any $|\nu| \in [0, \frac{R}{4}]$, there holds

$$\frac{1}{c_0}R^{N-1} \le H^{N-1}(T_\nu) \le c_0 R^{N-1} \quad \text{if } N \ge 2.$$
 (1.11)

Then

$$\lambda_1(\Omega) \le 4 \ln \frac{1}{R} + c_1,\tag{1.12}$$

where $c_1 > 0$ independent of R.

Since the function z_1 is decreasing, estimate (1.10) indicates that there exists R^* such that $\lambda_1(\Omega) < 0$ when $R > R^*$. Furthermore,

$$\lambda_1(\Omega) \le \omega_{N-1} \ln \frac{2}{R} + \rho_N + O(R^{-1})$$
 as $R \to \infty$.

Our upper bounds are obtained by considering the Rayleigh quotient $\lambda_1(\Omega) \leq \frac{\mathcal{E}(u,u)}{\int_{\Omega} u^2 dx}$ with the particular function $u = w_{\sigma}(x) = \min\{\max\{\sigma\rho(x), 0\}, 1\}$, where ρ is defined in (1.6) and $\sigma > 0$. However, the dominating term of upper bounds arises from $c_N \iint_{|x-y| \geq 1} \frac{w_{\sigma}(x)w_{\sigma}(y)}{|x-y|^N} dx$ when R is large, while it does from $\mathcal{E}(w_{\sigma}, w_{\sigma})$ for R > 0 small enough.

Next we prove lower bounds of the principle eigenvalue.

Theorem 1.3. Assume that Ω is a bounded domain in \mathbb{R}^N and $\lambda_1(\Omega)$ is the principle eigenvalue of Dirichlet problem (1.3) obtained in Proposition 1.1. Let

$$d_N = \frac{2\omega_{N-1}}{N^2(2\pi)^N}. (1.13)$$

Then we have that

(i)

$$\lambda_1(\Omega) \geq -d_N|\Omega|$$
;

(ii) if
$$|\Omega| < \frac{2}{eNd_N}$$
,

$$\lambda_1(\Omega) > 0;$$

(iii) if
$$|\Omega| \leq \frac{2}{e^{e+1}Nd_N}$$
,

$$\lambda_1(\Omega) \geq \frac{2}{N} \left(\ln \left(\frac{2}{eNd_N|\Omega|} \right) - \ln \ln \left(\frac{2}{eNd_N|\Omega|} \right) \right),$$

where e is the Euler number.

We summarize our results by the following table of the main asymptotic term of principle eigenvalue with respect to the volume of domain Ω from the upper bound in Theorem 1.2 and the lower bound in Theorem 1.3 in the particular where $\Omega = B_R$ where $0 < r_0 < 1 < R_0 < +\infty$:

R	$(0, r_0)$	$(R_0, +\infty)$
Upper bound of $\lambda_1(\Omega)$	$4 \ln \frac{1}{R}$	$\omega_{N-1} \ln \frac{1}{R}$
Lower bound of $\lambda_1(\Omega)$	$2 \ln \frac{1}{R}$	$-\frac{2e\omega_{N-1}}{N^2d_N}R^N$

The main order asymptotic of principle eigenvalue with respect to R.

From above table we note that the lower bound of principle eigenvalue for large value R is rather unprecise.

The Hilbert-Pólya conjecture is to associate the zero of the Riemann Zeta function with the eigenvalue of a Hermitian operator. This quest initiated the mathematical interest for estimating the sum of Dirichlet eigenvalues of the Laplacian while in physics the question is related to count the number of bound states of a one body Schrödinger operator and to study their asymptotic distribution. In 1912, Weyl in [37] shows that the k-th eigenvalue $\mu_k(\Omega)$ of Dirichlet problem with the Laplacian operator has the asymptotic behavior $\mu_k(\Omega) \sim C_N(k|\Omega|)^{\frac{2}{N}}$ as $k \to +\infty$, where $C_N = (2\pi)^2 |B_1|^{-\frac{2}{N}}$. Later, Pólya [33] (in 1960) proved that

$$\mu_k(\Omega) \ge C \left(\frac{k}{|\Omega|}\right)^{\frac{2}{N}} \tag{1.14}$$

holds for $C = C_N$ and any "plane-covering domain" D in \mathbb{R}^2 , (his proof also works in dimension $N \geq 3$) and he also conjectured that (1.14) holds with $C = C_N$ for any bounded domain in \mathbb{R}^N . Rozenbljium [34] and independently Lieb [29] proved (1.14) with a positive constant C for general bounded domain. Li-Yau [28] improved the constant $C = \frac{N}{N+2}C_N$, and with that constant (1.14) is also called Berezin-Li-Yau inequality because this constant is achieved with the help of Legendre transform as in the Berezin's earlier paper [2]. The Berezin-Li-Yau inequality then is generalized in [11–13,25,29,31], for degenerate elliptic operators in [9,22,38] for the fractional Laplacian $(-\Delta)^s$ defined in (1.4) and the inequality reads

$$\mu_{s,k}(\Omega) \ge \frac{N}{N+2s} C_N \left(\frac{k}{|\Omega|}\right)^{\frac{2s}{N}}.$$
(1.15)

Due to the expression of the Fourier symbol of L_{Δ} , Berezin-Li-Yau method can not be applied to our problem (1.3). Our results are based on the appropriate estimates for the solutions of equations:

$$r \ln r = c$$
 and $\frac{r}{\ln r - \ln \ln r} = t$.

The estimates that we obtain provide a uniform lower bound of the sum of the first k-eigenvalues, independently of k, an estimate which has a particular interest when these eigenvalues are negative. More precisely, we have the following inequalities:

Theorem 1.4. Let $\Omega \subset \mathbb{R}^N$ be a bounded domain, $\{\lambda_i(\Omega)\}_{i\in\mathbb{N}}$ be the sequence of eigenvalues of problem (1.3) obtained in Proposition 1.1 and d_N be given in (1.13). Then there holds (i) for any $k \in \mathbb{N}^*$,

$$\sum_{i=1}^{k} \lambda_i(\Omega) \ge -d_N|\Omega|;$$

(ii) if $k > \frac{eNd_N}{2}|\Omega|$,

$$\sum_{i=1}^{k} \lambda_i(\Omega) > 0;$$

(iii) if $k \ge \frac{e^{e+1}Nd_N}{2}|\Omega|$,

$$\sum_{i=1}^{k} \lambda_i(\Omega) \ge \frac{2k}{N} \left(\ln k + \ln \left(\frac{2}{eNd_N|\Omega|} \right) - \ln \ln \left(\frac{2k}{eNd_N|\Omega|} \right) \right). \tag{1.16}$$

Using the monotonicity of the sequence of eigenvalues, we deduce the following lower bound for $\lambda_k(\Omega)$ from Theorem 1.4 part (iii).

Corollary 1.5. Under the assumption of Theorem 1.4, for $k \geq \frac{e^{e+1}Nd_N}{2}|\Omega|$, we have that

$$\lambda_k(\Omega) > 0$$

and

$$\lambda_k(\Omega) \ge \frac{2}{N} \left(\ln k + \ln \left(\frac{2}{eNd_N|\Omega|} \right) - \ln \ln \left(\frac{2k}{eNd_N|\Omega|} \right) \right). \tag{1.17}$$

Our goal is to provide an upper bound for the sum of eigenvalues. Motivated by Kröger's result for the Laplacian [25], we prove the following upper bound.

Theorem 1.6. Let $\Omega \subset \mathbb{R}^N$ be a bounded domain and $\{\lambda_i(\Omega)\}_{i\in\mathbb{N}}$ be the sequence of eigenvalues of problem (1.3). Then for $k > k_0 := \frac{1}{2}e^{N+1}d_N|\Omega|$,

$$\sum_{i=1}^{k} \lambda_i(\Omega) \le \frac{2k}{N} \left(\ln(k+1) + \ln\left(\frac{p_N}{|\Omega|}\right) + \frac{\omega_{N-1}}{\sqrt{|\Omega|}} \ln\ln\left(\frac{p_N(k+1)}{|\Omega|}\right) \right)$$
(1.18)

and

$$\lim_{k \to +\infty} (k \ln k)^{-1} \sum_{i=1}^{k} \lambda_i(\Omega) = \frac{2}{N}, \tag{1.19}$$

where $p_N = \frac{2(2\pi)^N N}{\omega_{N-1}}$.

Note that from (1.19) the limit of the sum of the first k-eigenvalues does not depend on the volume of Ω . Finally, we build the Wely's formula for the logarithmic Laplacian and indeed we have the following asymptotic estimate.

Theorem 1.7. Let $\Omega \subset \mathbb{R}^N$ be a bounded domain and $\{\lambda_i(\Omega)\}_{i\in\mathbb{N}}$ be the sequence of eigenvalues of problem (1.3). Then

$$\lim_{k \to +\infty} \frac{\lambda_k(\Omega)}{\ln k} = \frac{2}{N}.$$
(1.20)

It is worth noting that

(a) we have the same limits of $\frac{\lambda_k(\Omega)}{\ln k}$ and $(k \ln k)^{-1} \sum_{i=1}^k \lambda_i(\Omega)$ as $k \to +\infty$;

(b) Weyl's estimate (1.20) is derived by the lower bound and the upper bound of the first keigenvalues directly.

Usually, the asymptotic behavior of eigenvalues is derived by the counting functions. Inversely, the estimates of counting functions could be deduced by the asymptotic behavior of eigenvalues. Let $\mathcal{N}(t)$ be the counting function of $\{\lambda_k(\Omega)\}_{k\in\mathbb{N}}$, counts the number of eigenvalues below t>0, i.e.

$$\mathcal{N}(t) = \sum_{j \in \mathbb{N}} sgn_{+}(t - \lambda_{j}(\Omega)) = \sum_{j \in \mathbb{N}} (t - \lambda_{j}(\Omega))_{+}^{0}$$
(1.21)

Here $sgn_+(r)=1$ if r>0, $sgn_+(r)=0$ if $r\leq 0$ and $r_\pm=(|r|\pm r)/2$ denotes the positive and negative part of $x \in \mathbb{R}$. The counting function can also be expressed by introducing the trace of an operator

$$\mathcal{N}(t) = \operatorname{tr}(L_{\Lambda} - t)_{-}^{0}.$$

Note that from the bounds of $\lambda_k(\Omega)$, we can obtain estimates for the counting function

$$\lim_{t \to +\infty} \mathcal{N}(t)e^{-(\frac{N}{2}+\delta)t} = 0 \quad \text{and} \quad \liminf_{t \to +\infty} \mathcal{N}(t)e^{-(\frac{N}{2}-\delta)t} = +\infty.$$

By analyzing the asymptotic behavior of $\sum_{j\in\mathbb{N}}(t-\lambda_j(\Omega))_+$. The rest of this paper is organized as follows. In Section 2, we build the upper bound of the principle eigenvalue of Theorem 1.2 by considering particular test functions in the Rayleigh quotient. Section 3 is devoted to proving the lower bound by developing Li-Yau's method, and then we prove Theorem 1.3, Theorem 1.4 and Corollary 1.5. In Section 4, we show the upper bounds for the first k-eigenvalues in Theorem 1.6 and prove the Wely's limit of eigenvalues in Theorem 1.7. Finally, we obtain the estimates for the counting function.

$\mathbf{2}$ Upper bounds for the principle eigenvalue

2.1 Large domain: proof of Theorem 1.2-(i)

Set

$$\eta(t) = \min\{\max\{0, t\}, 1\}$$
 for all $t \in \mathbb{R}$,

and, for $\sigma > 0$,

$$w_{\sigma}(x) = \eta(\sigma^{-1}\rho(x)) \quad \text{for all } x \in \mathbb{R}^{N}.$$
 (2.1)

Note that $w_{\sigma} \in \mathbb{H}(\Omega)$, $w_{\sigma} \to 1$ in Ω as $\sigma \to 0^+$, and for $\sigma \in (0, 2R]$,

$$w_{\sigma}(x) = \frac{1}{\sigma}\rho(x)$$
 for all $x \in \Omega$.

Since $|\eta'| \leq 1$ and the signed distance function ρ is a contraction mapping, there always hols

$$|w_{\sigma}(x) - w_{\sigma}(y)| \le \frac{1}{\sigma} |x - y| \quad \text{for all } x, y \in \mathbb{R}^{N}.$$
 (2.2)

By definition of $\lambda_1(\Omega)$,

$$\lambda_1(\Omega) \le \inf_{\sigma>0} \frac{\mathcal{E}_L(w_\sigma, w_\sigma)}{\int_{\Omega} w_\sigma^2(x) dx}.$$
 (2.3)

If $\sigma \geq 2R$, we have

$$\int_{\Omega} w_{\sigma}^2(x) dx = \sigma^{-2} \int_{\Omega} \rho^2 dx,$$

while if $\sigma < 2R$, there holds

$$\begin{split} \int_{\Omega} w_{\sigma}^2(x) dx &= \sigma^{-2} \int_{\Omega_{\sigma}} \rho^2 dx + \int_{\Omega \backslash \Omega_{\sigma}} dx \\ &= \sigma^{-2} \int_{\Omega_{\sigma}} \rho^2 dx + |\Omega| - |\Omega_{\sigma}|, \end{split}$$

where

$$\Omega_{\sigma} = \left\{ x \in \mathbb{R}^N : 0 < \rho(x) < \sigma \right\} \subset \Omega. \tag{2.4}$$

Then

$$\int_{\Omega} w_{\sigma}^{2}(x)dx = \sigma^{-2} \int_{\Omega_{\sigma}} \rho^{2} dx + |\Omega| - |\Omega_{\sigma}|$$
$$\geq |\Omega| - |\Omega_{\sigma}|$$

Taking $\sigma \leq \frac{1}{2}$, using (1.9) and the co-area formula since $|\nabla \rho(x)| = 1$, we have that

$$|\Omega| - |\Omega_{\sigma}| = \int_{\Omega \setminus \Omega_{\sigma}} |\nabla \rho(x)| dx = |\Omega| - \int_{0}^{\sigma} H^{N-1}(T_{t}^{+}) dt$$

$$\geq |\Omega| - c_{0}\sigma R^{N-1} = |\Omega| - \frac{Nc_{0}\sigma}{\omega_{N-1}R} |B_{R}|,$$

since $N|B_1| = \omega_{N-1}$. Hence, under the assumption of Theorem 1.2-(i), we have that

$$\int_{\Omega} w_{\sigma}^{2}(x)dx \ge \left(1 - \frac{Nc_{0}\sigma}{\omega_{N-1}R}\right)|\Omega| \quad \text{for all } \sigma \in (0, \frac{1}{2}).$$
(2.5)

Concerning the term $\mathcal{E}_L(w_{\sigma}, w_{\sigma})$, we have the following upper bound.

Lemma 2.1. Under the assumption of Theorem 1.2-(i), let $\sigma = \frac{1}{4}$, then we have that

$$\mathcal{E}_L(w_{\sigma}, w_{\sigma}) \le \frac{4c_0}{R} |\Omega| + \left(\rho_N - \omega_{N-1} \ln \frac{R}{2}\right) \int_{\Omega} w_{\sigma}^2 dx. \tag{2.6}$$

Proof. We recall that

$$\mathcal{E}_L(w_{\sigma}, w_{\sigma}) = \mathcal{E}(w_{\sigma}, w_{\sigma}) - c_N \iint_{\substack{x,y \in \mathbb{R}^N \\ |x-y| > 1}} \frac{w_{\sigma}(x)w_{\sigma}(y)}{|x-y|^N} dxdy + \rho_N \int_{\Omega} w_{\sigma}^2 dx$$

and our proof is divided into two parts.

Step 1: Note that for $x, y \in \mathbb{R}^N$ such that $|x - y| \le 1$, we have that

$$|w_{\sigma}(x) - w_{\sigma}(y)| = 0, \quad \forall (x, y) \in (\Omega \setminus \Omega_{\sigma})^{2} \cup (\Omega^{c})^{2}.$$

Using (2.2), we have that

$$\mathcal{E}(w_{\sigma}, w_{\sigma}) = \frac{c_N}{2} \iint_{\substack{x,y \in \mathbb{R}^N \\ |x-y| \le 1}} \frac{(w_{\sigma}(x) - w_{\sigma}(y))^2}{|x-y|^N} dx dy$$

$$\leq c_N \sigma^{-2} \int_{\Omega_{\sigma}} \int_{B_1(x)} |x-y|^{2-N} dy dx$$

$$= \frac{c_N \omega_{N-1}}{2\sigma^2} |\Omega_{\sigma}|$$

$$\leq \frac{c_0}{R\sigma} |\Omega| = \frac{4c_0}{R} |\Omega|,$$

thanks to the identity $c_N \omega_{N-1} = 2$.

Step 2: We have that

$$c_N \iint_{\substack{x,y \in \mathbb{R}^N \\ |x-y| \ge 1}} \frac{w_{\sigma}(x)w_{\sigma}(y)}{|x-y|^N} dy dx = c_N \int_{\Omega} \int_{\Omega \cap B_1^c(x)} \frac{w_{\sigma}(x)w_{\sigma}(y)}{|x-y|^N} dy dx.$$

Note that for $\sigma = \frac{1}{4}$, we have that

$$\inf_{x \in \Omega} |\Omega \setminus (\Omega_{\sigma} \cup B_1(x))| \ge |\Omega| - \frac{1}{4} c_0 R^{N-1} - \frac{\omega_{N-1}}{N}.$$

Set

$$\mathcal{D}_1(\sigma, x) = (\Omega \setminus B_1(x)) \cap \Omega_{\sigma} = \Omega_{\sigma} \cap B_1^c(x)$$

and

$$\mathcal{D}_2(\sigma, x) = (\Omega \setminus B_1(x)) \cap (\Omega \setminus \Omega_{\sigma}) = \Omega \cap \Omega_{\sigma}^c \cap B_1^c(x).$$

Then $\mathcal{D}_1(\sigma, x) \cap \mathcal{D}_2(\sigma, x) = \emptyset$, $\mathcal{D}_1(\sigma, x) \cup \mathcal{D}_2(\sigma, x) = \Omega \cap B_1^c(x)$. If $x \in B_{\frac{R}{2}}$, then $B_{\frac{R}{3}}(x) \cap B_1^c(x) \subset \mathcal{D}_2(\sigma, x)$, which implies

$$\int_{\mathcal{D}_2(\sigma,x)} \frac{1}{|x-y|^N} dy \ge \int_{B_{\frac{R}{2}}(x) \cap B_1^c(x)} \frac{1}{|x-y|^N} dy = \omega_{N-1} \ln \frac{R}{2}$$

and

$$\int_{\Omega\cap B_1^c(x)}\frac{w_\sigma(y)}{|x-y|^N}dy=\int_{\mathcal{D}_1(\sigma,x)}\frac{\sigma\rho(y)}{|x-y|^N}dy+\int_{\mathcal{D}_2(\sigma,x)}\frac{1}{|x-y|^N}dy>\omega_{_{N-1}}\ln\frac{R}{2},$$

since the first term of the right-hand side is positive. Thus, we obtain

$$\int_{\Omega} \int_{\Omega \cap B_1^c(x)} \frac{w_{\sigma}(x)w_{\sigma}(y)}{|x-y|^N} dy dx \ge \omega_{N-1} \ln \frac{R}{2} \int_{\Omega} w_{\sigma}(x) dx$$
$$\ge \omega_{N-1} \ln \frac{R}{2} \int_{\Omega} w_{\sigma}^2(x) dx.$$

As a consequence and since $\sigma = \frac{1}{4}$, we infer that

$$\mathcal{E}_L(w_{\sigma}, w_{\sigma}) \le \frac{4c_0}{R} |\Omega| + \left(\rho_N - \omega_{N-1} \ln \frac{R}{2}\right) \int_{\Omega} w_{\sigma}^2 dx.$$

We complete the proof.

Proof of Theorem 1.2-(i). Since $R > \frac{Nc_0}{2\omega_{N-1}}$, we have that

$$\left(1 - \frac{Nc_0}{4\omega_{N-1}R}\right)^{-1} \le 1 + \frac{Nc_0}{2\omega_{N-1}R}$$

and

$$\begin{split} \lambda_{1}(\Omega) &\leq \frac{\mathcal{E}_{L}(w_{\sigma}, w_{\sigma})}{\int_{\Omega} w_{\sigma}^{2}(x) dx} \leq \frac{4Nc_{0}}{R} \Big(1 - \frac{c_{0}}{4\omega_{N-1}R} \Big)^{-1} + \rho_{N} - \omega_{N-1} \ln \frac{R}{2} \\ &\leq -\omega_{N-1} \ln R + \rho_{N} + \omega_{N-1} \ln 2 + \frac{4c_{0}}{R} \Big(1 + \frac{Nc_{0}}{2\omega_{N-1}R} \Big), \end{split}$$

which is the result.

2.2 Small domain: proof of Theorem 1.2-(ii)

The following upper estimate of $\mathcal{E}_L(w_{\sigma}, w_{\sigma})$ holds.

Lemma 2.2. Under the assumptions of Theorem 1.2-(ii), there exists $c_2 > 0$ independent of R such that for $\sigma = \min\left\{\frac{R}{4}, \frac{2R\omega_{N-1}}{Nc_0}\right\}$,

$$\mathcal{E}_L(w_\sigma, w_\sigma) \le \left(2\ln\frac{1}{R} + c_2\right)|\Omega| + \rho_N \int_{\Omega} w_\sigma^2 dx. \tag{2.7}$$

Proof. Since $\Omega \subset B_{2R}$ and $R < \frac{1}{4}$, there holds

$$\mathcal{E}_L(w_{\sigma}, w_{\sigma}) = \mathcal{E}(w_{\sigma}, w_{\sigma}) + \rho_N \int_{\Omega} w_{\sigma}^2 dx.$$

For r > 0, we denote

$$\Omega^{1,r} = \{x \in \mathbb{R}^N : |\rho(x)| < r\} \text{ and } \Omega^{2,r} = \{x \in \mathbb{R}^N : \rho(x) > -r\} = \Omega \cup \Omega^{1,r},$$

and for $r_1 > r_2 > 0$ set

$$\mathcal{A}^{r_1,r_2} = \Omega^{2,r_1} \setminus \Omega^{2,r_2} = \{ x \in \mathbb{R}^N : -r_2 \ge \rho(x) > -r_1 \}.$$

Note that for r > 2R, $\Omega \subset \Omega^{1,r} = \Omega^{2,r}$ since $\Omega \subset B_{2R}$. When $\sigma \in (0, \frac{R}{4}]$, we set

$$\mathcal{D} = \{(x,y) \in \mathbb{R}^N \times \mathbb{R}^N \text{ s.t. } |x-y| < 1 \text{ and } |w_{\sigma}(x) - w_{\sigma}(y)| > 0\}.$$

Then $\mathcal{D} \subset \mathcal{D}_1 \cup \mathcal{D}_2$, where

$$\mathcal{D}_1 = (\Omega^{1,\sigma} \times \Omega^{2,\sigma}) \cup (\Omega^{2,\sigma} \times \Omega^{1,\sigma}) \text{ and } \mathcal{D}_2 = (\mathcal{A}^{1,\sigma} \times \Omega) \cup (\Omega \times \mathcal{A}^{1,\sigma}).$$

Using (2.2), $|\Omega^{2,\sigma}| \leq |B_{2R+\sigma}(x)|$, we obtain from the definition of \mathcal{D}_1 and since $c_N \omega_{N-1} = 2$,

$$\frac{c_N}{2} \iint_{\mathcal{D}_1} \frac{(w_{\sigma}(x) - w_{\sigma}(y))^2}{|x - y|^N} dx dy \le c_N \sigma^{-2} \int_{\Omega^{1,\sigma}} \int_{\Omega^{2,\sigma}} |x - y|^{2-N} dy dx
\le c_N \sigma^{-2} \int_{\Omega^{1,\sigma}} \int_{B_{2R+\sigma}(x)} |x - y|^{2-N} dy dx
= \frac{c_N \omega_{N-1}}{2} (\frac{2R}{\sigma} + 1)^2 |\Omega^{1,\sigma}|
\le 2c_0 (\frac{2R}{\sigma} + 1)^2 R^{N-1} \sigma
\le \frac{Nc_N c_0 \sigma}{R} (\frac{2R}{\sigma} + 1)^2 |\Omega|.$$

On the other hand,

$$\frac{c_N}{2} \iint_{\mathcal{D}_2} \frac{(w_{\sigma}(x) - w_{\sigma}(y))^2}{|x - y|^N} dx dy \le c_N \int_{\mathcal{A}^{1,\sigma}} \int_{\Omega} |x - y|^{-N} dy dx
= c_N \int_{\Omega} \int_{\mathcal{A}^{1,\sigma}} |x - y|^{-N} dx dy
\le 2|\Omega| \ln \frac{2}{\sigma},$$

since for any $y \in \Omega$, $x \in \mathcal{A}^{1,\sigma} \Longrightarrow x \in B_2(y) \setminus B_{\sigma}(y)$, which implies

$$\int_{\mathcal{A}^{1,\sigma}} |x-y|^{-N} dx \le \int_{B_2(y)\setminus B_{\sigma}(y)} |x-y|^{-N} dx = \omega_{N-1} \ln \frac{2}{\sigma}.$$

Taking $\sigma = \varrho_0 R$ with

$$\varrho_0 = \min\left\{\frac{1}{4}, \frac{\omega_{N-1}}{2Nc_0}\right\},\,$$

we obtain that

$$\mathcal{E}(w_{\sigma}, w_{\sigma}) = \frac{c_N}{2} \iint_{\substack{x, y \in \mathbb{R}^N \\ |x-y| \le 1}} \frac{(w_{\sigma}(x) - w_{\sigma}(y))^2}{|x-y|^N} dx dy$$

$$\leq \left(\frac{Nc_N c_0 \sigma}{R} (\frac{2R}{\sigma} + 1)^2 + 2\ln\frac{2}{\sigma}\right) |\Omega|$$

$$< \left(2\ln\frac{1}{R} + c_2\right) |\Omega|,$$

where $c_2 = \frac{81Nc_0}{2\omega_{N-1}} + 4 \ln 2$, which yields (2.7).

End of the proof of Theorem 1.2-(ii). Estimate (2.5) is valid, hence, if $\sigma = \varrho_0 R$, we derive that $1 - \frac{Nc_0\sigma}{R\omega_{N-1}} \leq \frac{1}{2}$ and

$$\int_{\Omega} w_{\sigma}^{2}(x)dx \ge \frac{1}{2}|\Omega|.$$

Therefore,

$$\lambda_1(\Omega) \le \frac{\mathcal{E}_L(w_{\sigma}, w_{\sigma})}{\int_{\Omega} w_{\sigma}^2(x) dx} \le 2\left(2\ln\frac{1}{R} + c_2\right) + \rho_N = 4\ln\frac{1}{R} + 2c_2 + \rho_N,$$

which ends the proof.

3 Lower bounds

Let

$$q(r) = r \ln r$$
 for $r > 0$,

then g(e) = e, g(1) = 0 and $g(\frac{1}{e}) = -\frac{1}{e}$

Lemma 3.1. For $c \geq -\frac{1}{e}$, there exists a unique point $r_c \geq \frac{1}{e}$ such that

$$g(r_c) = c,$$

and we have that $r_c \leq 1 + c$. Furthermore,

(i) for $-\frac{1}{e} \le c \le 0$,

$$r_c \ge 1 + (e - 1)c \ge \frac{1}{e};$$

(ii) for $0 \le c \le e$,

$$r_c \geq 1 + \frac{e-1}{e}c;$$

(iii) for $c \geq e$,

$$r_c \ge 1 + \frac{e-1}{e}c$$

and

$$\frac{c}{\ln c} \le r_c \le \frac{c}{\ln c - \ln \ln c}.\tag{3.1}$$

Proof. The function g is increasing in $\left[\frac{1}{e}, +\infty\right)$ with value in $\left[-\frac{1}{e}, +\infty\right)$. Hence r_c is uniquely determined if $c \geq -\frac{1}{e}$, $c \mapsto r_c$ is increasing from $\left[-\frac{1}{e}, +\infty\right)$ onto $\left[\frac{1}{e}, +\infty\right)$, and g is convex.

For a > 0, we define $\psi_a(x) = (1 + ax) \ln(1 + ax) - x$ for $x > -\frac{1}{a}$. Then $\psi_a(x) > 0$ (resp. $\psi_a(x) < 0$) is equivalent to $1 + ax > r_x$ (resp. $1 + ax < r_x$). Note that $\psi'_a(x) = a(1 + \ln(1 + ax)) - 1$. Since $\psi_a'(-\frac{1}{a}) = -\infty$ and ψ_a' is increasing, $\psi_a'(0) = a - 1$ is the maximal (resp. minimal) value of ψ_a' on $(-\frac{1}{a},0]$ (resp. on $[0,\infty)$). Therefore, if a>1, ψ_a is positive on $(-\frac{1}{a},r_a^*)$ for some $r_a^*\in(-\frac{1}{a},0)$, negative on $(r_a^*,0)$ and positive on $(0,\infty)$. If 0< a<1, ψ_a is positive on $(-\frac{1}{a},0)$, negative on $(0,r_a^*)$ for some $r_a^*>0$ and positive on (r_a^*,∞) . If a=1, ψ_1 is positive on $[-\frac{1}{e},0)\cup(0,\infty)$ and vanishes only at 0. Then $\psi_1 \geq 0$ implies the first assertion.

Since e - 1 > 1 and $\psi_{e-1}(-\frac{1}{e}) = 0$, $\psi_{e-1}(x) < 0$ for $x \in (-\frac{1}{e}, 0)$. This gives (i). Since $0 < \frac{e-1}{e} < 1$, $\psi_{\frac{e-1}{e}}$ is negative on $(0, r_{\frac{e-1}{e}}^*)$ and positive on $(r_{\frac{e-1}{e}}^*, \infty)$. Since $\psi_{\frac{e-1}{e}}(e) = 0$, $r_{\underline{e-1}}^* = e$ and we get (ii) and (iii).

Since g is increasing on $[e, \infty)$, (3.1) is equivalent to

$$c - \frac{\ln c}{\ln \ln c} \le c \le \frac{c}{\ln c - \ln \ln c} \ln \left(\frac{c}{\ln c - \ln \ln c} \right) = c \frac{\ln c - \ln (\ln c - \ln \ln c)}{\ln c - \ln \ln c}.$$

Set $C = \ln c$, then

$$\frac{\ln c - \ln(\ln c - \ln \ln c)}{\ln c - \ln \ln c} = \frac{C - \ln(C - \ln C)}{C - \ln C} > 1 \quad \text{for } C > 1$$

and (3.1) follows.

Lemma 3.2. Let f be a real-valued function defined in \mathbb{R}^N with $0 \le f \le M_1$ and

$$2\int_{\mathbb{R}^N} \ln|z| f(z) dx = M_2.$$

Then (i)

$$M_2 \ge -\frac{2\omega_{N-1}}{N^2} M_1;$$

(ii)

$$\int_{\mathbb{R}^{N}} f(z) dz \leq \frac{M_{1} \omega_{_{N-1}}}{N} \left(e + \frac{N^{2} M_{2}}{2 M_{1} \omega_{_{N-1}}} \right) = \frac{e \omega_{_{N-1}}}{N} M_{1} + \frac{N}{2} M_{2};$$

(iii) assuming more that $\frac{M_2}{M_1} \ge \frac{2e^2\omega_{N-1}}{N^2}$, there holds

$$\int_{\mathbb{R}^{N}} f(z)dz \leq \frac{NM_{2}}{2} \left(\ln \left(\frac{N^{2}M_{2}}{2eM_{1}\omega_{N-1}} \right) - \ln \ln \left(\frac{N^{2}M_{2}}{2eM_{1}\omega_{N-1}} \right) \right)^{-1}.$$

Proof. We have

$$\begin{split} \frac{M_2}{2} &= \int_{B_1} \ln|z| f(z) dz + \int_{B_1^c} \ln|z| f(z) dz \\ &\geq M_1 \int_{B_1} \ln|z| dz + \int_{B_1^c} \ln|z| f(z) dz \geq -\frac{\omega_{N-1}}{N^2} M_1. \end{split}$$

Hence (i) holds.

For R > 0 we have that

$$(\ln |z| - \ln R)(f(z) - M_1 \mathbf{1}_{B_R}) \ge 0.$$

By integration over \mathbb{R}^N we get

$$\frac{M_2}{2} + \frac{M_1 \omega_{N-1} R^N}{N^2} \ge \ln R \int_{\mathbb{R}^N} f(z) dz.$$

The estimate from above of $\int_{\mathbb{R}^N} f(z)dz$ is obtained by

$$\int_{\mathbb{R}^N} f(z)dz \le \inf\Big\{A > 0 \text{ s.t. } \frac{M_2}{2} + \frac{M_1\omega_{N-1}R^N}{N^2} - A\ln R \ge 0 \text{ for all } R > 0\Big\}. \tag{3.2}$$

Set

$$\Theta_A(R) = \frac{M_2}{2} + \frac{M_1 \omega_{N-1} R^N}{N^2} - A \ln R.$$

then Θ_A achieves the minimum if

$$\frac{M_1 \omega_{N-1} R^N}{N} = A \Longleftrightarrow R = R_A := \left(\frac{NA}{\omega_{N-1} M_1}\right)^{\frac{1}{N}}.$$

Hence

$$\Theta_A(R_A) = \frac{M_2}{2} + \frac{A}{N} - \frac{A}{N} \ln \left(\frac{NA}{\omega_{N-1} M_1} \right). \tag{3.3}$$

Put $r = \frac{NA}{M_1\omega_{N-1}}$, then

$$\Theta_A(R_A) \ge 0 \Longleftrightarrow r \ln r - r \le \frac{N^2 M_2}{2M_1 \omega_{N-1}} \Longleftrightarrow g\left(\frac{r}{e}\right) \le \frac{N^2 M_2}{2e M_1 \omega_{N-1}}.$$
(3.4)

Then $\frac{r}{e} \leq r_c$ with $c = \frac{N^2 M_2}{2e M_1 \omega_{N-1}}$, inequality $r_c \leq 1 + c$ in Lemma 3.1 yields

$$r = \frac{NA}{M_1\omega_{N-1}} \le e + \frac{N^2M_2}{2M_1\omega_{N-1}} \Longrightarrow \int_{\mathbb{R}^N} f(z)dz \le \frac{M_1\omega_{N-1}}{N} \left(e + \frac{N^2M_2}{2M_1\omega_{N-1}}\right),$$

which is (ii).

Assuming now that $\frac{M_2}{M_1} \geq \frac{2e^2\omega_{N-1}}{N^2}$, we can apply Lemma 3.1-(iii) and get

$$\int_{\mathbb{R}^{N}} f(z)dz \leq \frac{NM_{2}}{2} \left(\ln \left(\frac{N^{2}M_{2}}{2eM_{1}\omega_{N-1}} \right) - \ln \ln \left(\frac{N^{2}M_{2}}{2eM_{1}\omega_{N-1}} \right) \right)^{-1},$$

which is (iii) and ends the proof.

Lemma 3.3. Let

$$\tilde{g}(r) = \frac{r}{\ln r - \ln \ln r}$$
 for $r > e$.

Then for $t > \frac{e^e}{e-1}$, there exists a unique point $r_t > e$ such that $\tilde{g}(r_t) = t$. Furthermore,

$$t(\ln t - \ln \ln t) \le r_t < t \ln t. \tag{3.5}$$

Proof. Since

$$\tilde{g}'(r) = \frac{1}{\ln r - \ln \ln r} - \frac{1 - (\ln r)^{-1}}{(\ln r - \ln \ln r)^2}$$

$$\geq \frac{1}{\ln r - \ln \ln r} \left(1 - \frac{1}{\ln r - \ln \ln r} \right) > 0,$$

the function \tilde{g} is increasing from $(e, +\infty)$ onto $(\frac{e^e}{e^{-1}}, +\infty)$. Setting $r_t^* = t(\ln t - \ln \ln t)$, then

$$\begin{split} \tilde{g}(r_t^*) &= \frac{t(\ln t - \ln \ln t)}{\ln t + \ln(\ln t - \ln \ln t) - \ln \ln(t(\ln t - \ln \ln t))} \\ &\leq \frac{t(\ln t - \ln \ln t)}{\ln t + \ln(\ln t - \ln \ln t) - \ln \ln(t \ln t)} \\ &= \frac{\ln t - \ln \ln t}{\ln t - \ln \frac{\ln(t \ln t)}{\ln t - \ln \ln t}} t \\ &\leq t, \end{split}$$

where the last inequality holds if

$$\frac{\ln(t\ln t)}{\ln t - \ln\ln t} \le \ln t,$$

which is equivalent to

$$\tilde{h}(\tau) := \tau^2 - (\ln \tau + 1)\tau - \ln \tau \ge 0, \quad \tau = \ln t.$$

Freezing the coefficient $\ln \tau$, $\tilde{h}(\tau) = (\tau - \tau_1)(\tau - \tau_2)$, where the τ_1 , τ_2 depend of τ , but $\tau_1 < 0 < \tau_2$, since $\tau_1 \tau_2 = -\ln \tau < 0$. Because $\tilde{h}(e) = e^2 - 2e - 1 = 0.9584 \pm 10^{-4}$, we have $e > \tau_1$. Hence $\tau > e$ implies $\tau > \tau_1$ which in turn implies $\tilde{h}(\tau) > 0$. Hence $r_t^* \le r_t$ using the monotonicity of \tilde{g} .

Let $s_t = t \ln t$, then

$$\tilde{g}(s_t) = \frac{t \ln t}{\ln t + \ln \ln t - \ln \ln(t \ln t)} < t$$

by the fact that

$$ln ln t - ln ln(t ln t) < 0 \quad \text{for } t > e.$$

Hence $s_t \geq r_t$, which ends the proof.

Proof of Theorem 1.4. Denote

$$\Phi_k(x,y) = \sum_{j=1}^k \phi_j(x)\phi_j(y), \qquad (x,y) \in \Omega \times \Omega,$$

and

$$\widehat{\Phi}_k(z,y) = (2\pi)^{-\frac{N}{2}} \int_{\mathbb{R}^N} \Phi_k(x,y) e^{ix \cdot z} dx,$$

where $\widehat{\Phi}_k$ is the Fourier transform with respect to x. Hence we have that

$$\int_{\mathbb{R}^N} \int_{\Omega} |\widehat{\Phi}_k(z,y)|^2 dz dy = \int_{\Omega} \int_{\Omega} |\Phi_k(x,y)|^2 dx dy = k$$

by the orthonormality of the $\{\phi_j\}_{i\in\mathbb{N}}$ in $L^2(\Omega)$. Furthermore, we note that

$$\int_{\Omega} |\widehat{\Phi}_{k}(z,y)|^{2} dy = \int_{\Omega} \left(\sum_{j=1}^{k} \widehat{\phi}_{j}(z) \phi_{j}(y) \right) \left(\sum_{j=1}^{k} \overline{\widehat{\phi}_{j}(z)} \phi_{j}(y) \right) dy$$

$$= \int_{\Omega} \left(\sum_{j,\ell=1}^{k} \widehat{\phi}_{j}(z) \overline{\widehat{\phi}_{\ell}(z)} \phi_{j}(y) \phi_{\ell}(y) \right) dy$$

$$= \sum_{j=1}^{k} |\widehat{\phi}_{j}(z)|^{2}.$$
(3.6)

Using again the orthonormality of the $\{\phi_j\}_{i\in\mathbb{N}}$ in $L^2(\Omega)$, we infer by the k-dim Pythagore theorem,

$$\int_{\Omega} |\widehat{\Phi}_{k}(z,y)|^{2} dy = (2\pi)^{-N} \int_{\Omega} \left| \sum_{j=1}^{k} \left(\int_{\Omega} e^{ix.z} \phi_{j}(x) dx \right) \phi_{j}(y) \right|^{2} dy$$

$$= (2\pi)^{-N} \sum_{j=1}^{k} \left| \int_{\Omega} e^{ix.z} \phi_{j}(x) dx \right|^{2}$$

$$\leq (2\pi)^{-N} |\Omega|.$$
(3.7)

We have, from the Fourier expression of L_{Δ} ,

$$\begin{split} \sum_{j=1}^k \lambda_j(\Omega) &= \int_{\Omega} \int_{\Omega} \Phi_k(x,y) L_{\Delta} \Phi_k(x,y) dy dx \\ &= 2 \sum_{j=1}^k \int_{\mathbb{R}^N} |\widehat{\phi}_j(z)|^2 \ln|z| dz \\ &= 2 \int_{\mathbb{R}^N} \left(\int_{\Omega} |\widehat{\Phi}_k(z,y)|^2 dy \right) \ln|z| dz. \end{split}$$

Now we apply Lemma 3.2 to the function

$$f(z) = \int_{\Omega} |\widehat{\Phi}_k(z, y)|^2 dy$$

with

$$M_1 = (2\pi)^{-N} |\Omega|$$
 and $M_2 = \sum_{j=1}^k \lambda_j(\Omega)$.

Part (i): By Lemma 3.2 (i),

$$\sum_{j=1}^{k} \lambda_j(\Omega) \ge -\frac{2\omega_{N-1}}{N^2(2\pi)^N} |\Omega| = -d_N |\Omega|,$$

where d_N is constant defined in (1.13).

Part (ii):

$$k = \int_{\mathbb{R}^N} f(z)dz \le \frac{e\omega_{N-1}|\Omega|}{N(2\pi)^N} + \frac{N}{2} \sum_{j=1}^k \lambda_j(\Omega),$$

which implies that

$$\sum_{j=1}^{k} \lambda_j(\Omega) \ge \frac{2k}{N} - \frac{2e\omega_{N-1}|\Omega|}{N^2(2\pi)^N}.$$

Part (iii): for $k \in \mathbb{N}$, if

$$\sum_{j=1}^{k} \lambda_j(\Omega) \ge \frac{2e^2 \omega_{N-1}}{N^2} |\Omega|,$$

then

$$k \le \frac{NM_2}{2} \left(\ln \left(\frac{N^2 M_2}{2e M_1 \omega_{N-1}} \right) - \ln \ln \left(\frac{N^2 M_2}{2e M_1 \omega_{N-1}} \right) \right)^{-1}.$$

Setting

$$r = \frac{N^2 M_2}{2e M_1 \omega_{N-1}}$$
 and $t = \frac{Nk}{e M_1 \omega_{N-1}} = \frac{(2\pi)^N Nk}{e \omega_{N-1} |\Omega|}$

we have from (3.5) that

$$r \ge r_t \ge t \left(\ln t - \ln \ln t \right), \tag{3.8}$$

for any $t > e^e$, i.e.

$$k > \frac{e^{e+1}\omega_{N-1}|\Omega|}{(2\pi)^N N} = \frac{e^{e+1}Nd_N}{2}|\Omega|.$$

This implies

$$\frac{N^2 M_2}{2e M_1 \omega_{N-1}} \ge \frac{(2\pi)^N Nk}{e \omega_{N-1} |\Omega|} \left(\ln \left(\frac{(2\pi)^N Nk}{e \omega_{N-1} |\Omega|} \right) - \ln \ln \left(\frac{(2\pi)^N Nk}{e \omega_{N-1} |\Omega|} \right) \right),$$

from what we infer

$$\sum_{j=1}^{k} \lambda_j(\Omega) \ge \frac{2k}{N} \left(\ln \left(\frac{2k}{eNd_N|\Omega|} \right) - \ln \ln \left(\frac{2k}{eNd_N|\Omega|} \right) \right), \tag{3.9}$$

which completes the proof.

Proof of Theorem 1.3 and Corollary 1.5. It is clear that Theorem 1.4 with k = 1 implies Theorem 1.3. From the inequality (3.9) we derive

$$\tilde{\lambda}_k(\Omega) := \frac{1}{k} \sum_{i=1}^k \lambda_i(\Omega) \ge \frac{2}{N} \left(\ln \left(\frac{2k}{eNd_N|\Omega|} \right) - \ln \ln \left(\frac{2k}{eNd_N|\Omega|} \right) \right).$$

Since $k \mapsto \lambda_k(\Omega)$ is nondecreasing, we conclude Corollary 1.5.

4 Wey's limits

4.1 Upper bounds for the sum of eigenvalues

For any bounded complex valued functions u, v defined on Ω , there holds

$$L_{\Delta}(uv)(x) = u(x)L_{\Delta}v(x) + c_N \int_{B_1(x)} \frac{u(x) - u(\zeta)}{|x - \zeta|^N} v(\zeta)d\zeta. \tag{4.1}$$

Lemma 4.1. For $z \in \mathbb{R}^N$, we denote

$$\mu_z(x) = e^{ix \cdot z}, \quad \forall x \in \mathbb{R}^N,$$

then

$$L_{\Delta}\mu_z(x) = (2\ln|z|)\mu_z(x), \quad \forall x \in \mathbb{R}^N.$$
(4.2)

Proof. Step 1: we claim that for

$$(-\Delta)^s \mu_z(x) = |z|^{2s} \mu_z(x), \quad \forall x \in \mathbb{R}^N.$$
(4.3)

Without loss of generality, it is enough to prove (4.3) with $z=te_1$, where t>0 and $e_1=(1,0,\cdots,0)\in\mathbb{R}^N$. For this, we write

$$v_t(x) = \mu_z(x_1) = e^{itx_1}, \quad x = (x_1, x') \in \mathbb{R} \times \mathbb{R}^{N-1}.$$

For $N \geq 2$ it implies by [7, Lemma 3.1] that

$$(-\Delta)^s v_t(x) = (-\Delta)^s_{\mathbb{R}} v_t(x_1).$$

Now we claim that

$$(-\Delta)^s_{\mathbb{R}} v_t(x_1) = t^{2s} v_t(x_1), \quad \forall x_1 \in \mathbb{R}.$$

$$(4.4)$$

Indeed, observe that $-\Delta_{\mathbb{R}} := -(v_t)_{x_1x_1} = t^2v_t$ in \mathbb{R} and then

$$(|\xi_1|^2 - t^2)\widehat{v_t}(\xi_1) = \mathcal{F}(-\Delta_{\mathbb{R}}v_t - t^2v_t)(\xi_1) = 0,$$

which implies that

$$\operatorname{supp}(\widehat{v_t}) \subset \{\pm t\},\$$

which in turn implies

$$(|\xi_1|^{2s} - t^{2s})\widehat{v_t}(\xi_1) = 0 = \mathcal{F}\left((-\Delta)_{\mathbb{R}}^s v_t - t^{2s} v_t\right)(\xi_1).$$

and finally

$$\left((-\Delta)_{\mathbb{R}}^s v_t - t^{2s} v_t \right) (\xi_1) = 0 \quad \text{in } \mathbb{R},$$

which yields

$$(-\Delta)^{s} v_{t}(x) = (-\Delta)^{s}_{\mathbb{R}} v_{t} = t^{2s} v_{t}(x), \quad \forall x \in \mathbb{R}^{N},$$

Step 2: we show (4.2). From the property (1.5) of L_{Δ} since μ_z is bounded,

$$0 = \frac{(-\Delta)^{s} \mu_{z}(x) - |z|^{2s} \mu_{z}(x)}{s}$$

$$= \frac{(-\Delta)^{s} \mu_{z}(x) - \mu_{z}(x)}{s} - \frac{|z|^{2s} - 1}{s} \mu_{z}(x)$$

$$\to L_{\Delta} \mu_{z}(x) - (2 \ln|z|) \mu_{z}(x) \quad \text{as } s \to 0^{+},$$

hence,

$$L_{\Delta}\mu_z(x) = (2\ln|z|)\mu_z(x), \quad \forall x \in \mathbb{R}^N,$$

which is the claim.

Next, let $\eta_0 \in C^1(\mathbb{R})$ be a nondecreasing real value function such that $\|\eta_0'\|_{L^\infty} \leq 2$ satisfying

$$\eta_0(t) = 1$$
 if $t \ge 1$, $\eta_0(t) = 0$ if $t \le 0$.

Since Ω is a bounded domain, there exists a C^1 domain $\mathcal{O} \subset \Omega$ such that $|\mathcal{O}| \geq \frac{3}{4}|\Omega|$. For $\sigma > 0$, we set again

$$w_{\sigma}(x) = \eta_0(\sigma^{-1}\bar{\rho}(x)), \quad \forall x \in \mathbb{R}^N.$$
(4.5)

where $\bar{\rho}(x) = \operatorname{dist}(x, \partial \mathcal{O})$. Observe that $w_{\sigma} \in \mathbb{H}_0(\Omega)$ and

$$w_{\sigma} \to 1$$
 in \mathcal{O} as $\sigma \to 0^+$.

Thus, there exists $\sigma_1 > 0$ such that for $\sigma \in (0, \sigma_1]$,

$$|\Omega| > \int_{\Omega} w_{\sigma} dx \ge \int_{\Omega} w_{\sigma}^2 dx > \frac{|\Omega|}{2}.$$

Lemma 4.2. Let

$$\mathcal{L}_z w_{\sigma}(x) = \int_{B_1(x)} \frac{w_{\sigma}(x) - w_{\sigma}(\zeta)}{|x - \zeta|^N} e^{i\zeta \cdot z} d\zeta,$$

then there holds

$$\left| \mathcal{L}_z w_{\sigma}(x) \right| \le \frac{2\omega_{N-1}}{\sigma} \quad \text{for } x \in \Omega.$$

Proof. Actually, if $x \in \Omega$, we have that

$$|w_{\sigma}(x) - w_{\sigma}(\zeta)| \le ||Dw_{\sigma}||_{L^{\infty}} |x - \zeta| \le \sigma^{-1} ||\eta_0'||_{L^{\infty}} |x - \zeta|,$$

then

$$\Big| \int_{B_1(x)} \frac{w_{\sigma}(x) - w_{\sigma}(\zeta)}{|x - \zeta|^N} e^{i\zeta \cdot z} d\zeta \Big| \le \frac{\|\eta_0'\|_{L^{\infty}}}{\sigma} \int_{B_1(x)} \frac{d\zeta}{|\zeta - x|^{N-1}} \le \frac{2\omega_{N-1}}{\sigma},$$

since $\|\eta_0'\|_{L^{\infty}} \leq 2$. This ends the proof.

4.2 Proof of Theorem 1.6 and 1.7

Proof of Theorem 1.6. We recall that $\Phi_k(x,y)$ and $\widehat{\Phi}_k(z,y)$ have been defined in the proof of Theorem 1.4. If we denote

$$\tilde{v}_{\sigma,z}(x) := v_{\sigma}(x,z) = w_{\sigma}(x)e^{\mathrm{i}x\cdot z},$$

the projection of v_{σ} onto the subspace of $L^{2}(\Omega)$ spanned by the ϕ_{j} for $1 \leq j \leq k$ can be written in terms of the Fourier transform of $w_{\sigma}\Phi_{k}$ with respect to the x-variable:

$$\int_{\Omega} v_{\sigma}(x,z) \Phi_k(x,y) dx = (2\pi)^{N/2} \mathcal{F}_x(w_{\sigma} \Phi_k)(z,y).$$

Put

$$v_{\sigma,k}(z,y) = v_{\sigma}(z,y) - (2\pi)^{N/2} \mathcal{F}_x(w_{\sigma}\Phi_k)(z,y)$$

and the Rayleigh-Ritz formula shows that

$$\lambda_{k+1}(\Omega) \int_{\Omega} |v_{\sigma,k}(z,y)|^2 dy \le \int_{\Omega} \overline{v_{\sigma,k}(z,y)} L_{\Delta,y} v_{\sigma,k}(z,y) dy$$

for any $z \in \mathbb{R}^N$ and $\sigma > 0$, where the right hand side is a real value

$$\int_{\Omega} \overline{v_{\sigma,k}(z,y)} L_{\Delta,y} v_{\sigma,k}(z,y) dy = \int_{\mathbb{R}^N} \overline{v_{\sigma,k}(z,y)} L_{\Delta,y} v_{\sigma,k}(z,y) dy = \int_{\mathbb{R}^N} 2 \ln |\xi| \Big| \mathcal{F}(v_{\sigma,k})(z,\xi) \Big|^2 d\xi,$$

although $v_{\sigma,k}$ is complex valued function. Then, integrating this last inequality with respect to z in $B_r \setminus B_1$, for r > 1, we obtain

$$\lambda_{k+1}(\Omega) \le \inf_{\sigma > 0} \frac{\int_{B_r \setminus B_1} \int_{\Omega} \overline{v_{\sigma,k}(z,y)} L_{\Delta,y} v_{\sigma,k}(z,y) dy dz}{\int_{B_r \setminus B_1} \int_{\Omega} |v_{\sigma,k}(z,y)|^2 dy dz}.$$

By Pythagore's theorem, we have that

$$\int_{\Omega} |v_{\sigma,k}(z,y)|^2 dy = \int_{\Omega} |v_{\sigma}(z,y)|^2 dy - (2\pi)^N \int_{\Omega} \sum_{j=1}^k |\mathcal{F}_x(w_{\sigma}\phi_i)(z)|^2 \phi_i(y)^2 dy,$$

integrating over $B_r \setminus B_1$ implies that

$$\int_{B_r \setminus B_1} \int_{\Omega} |v_{\sigma,k}(z,y)|^2 dy dz \ge \frac{\omega_{N-1} r^N}{N} \int_{\Omega} w_{\sigma}^2(y) dy - (2\pi)^N \sum_{j=1}^k \int_{B_r \setminus B_1} |\mathcal{F}_x(w_{\sigma} \phi_i)(z)|^2 dz.$$

On the other hand,

$$\int_{B_r \setminus B_1} \int_{\Omega} \overline{v_{\sigma,k}(z,y)} L_{\Delta,y} v_{\sigma,k}(z,y) dy dz = \int_{B_r \setminus B_1} \int_{\Omega} \overline{v_{\sigma}(z,y)} L_{\Delta,y} v_{\sigma}(z,y) dy dz - (2\pi)^N \int_{B_r \setminus B_1} \int_{\Omega} \overline{\mathcal{F}_x(w_{\sigma} \Phi_k)(z,y)} L_{\Delta,y} \mathcal{F}_x(w_{\sigma} \Phi_k)(z,y) dy dz,$$

where

$$\int_{B_r \setminus B_1} \int_{\Omega} \overline{\mathcal{F}_x(w_{\sigma} \Phi_k)(z, y)} L_{\Delta, y} \mathcal{F}_x(w_{\sigma} \Phi_k)(z, y) dy dz = \sum_{j=1}^k \lambda_j(\Omega) \int_{B_r \setminus B_1} |\mathcal{F}_x(w_{\sigma} \phi_j)(z)|^2 dz$$

and

$$\begin{split} \int_{B_r \backslash B_1} \int_{\Omega} \overline{v_{\sigma}(z,y)} L_{\Delta,y} v_{\sigma}(z,y) dy dz \\ & \leq \int_{B_r \backslash B_1} \int_{\Omega} w_{\sigma}^2(y) |L_{\Delta,y} e^{\mathrm{i}y.z}| dy dz + \int_{B_r \backslash B_1} \int_{\Omega} w_{\sigma}(y) |\mathcal{L}_z w_{\sigma}(y)| dy dz \\ & \leq \int_{B_r \backslash B_1} \int_{\Omega} w_{\sigma}^2(y) \ln|z| dy dz + \frac{2\omega_{N-1}}{\sigma} \int_{B_r \backslash B_1} \int_{\Omega} w_{\sigma}(y) dy dz \\ & = \frac{\omega_{N-1}}{N} \varrho_{2,\sigma} \left(r^N \ln r - \frac{1}{N} (r^N - 1) \right) + \frac{\omega_{N-1}^2}{N\sigma} \varrho_{1,\sigma} \left(r^N - 1 \right) \\ & \leq \frac{\omega_{N-1}}{N} \varrho_{2,\sigma} r^N \ln r + \frac{\omega_{N-1}^2}{N\sigma} \varrho_{1,\sigma} r^N \end{split}$$

with

$$\varrho_{1,\sigma} = \int_{\Omega} w_{\sigma}(y) dy$$
 and $\varrho_{2,\sigma} = \int_{\Omega} w_{\sigma}^{2}(y) dy$.

Because of Parseval's identity, there holds

$$\int_{B_r \setminus B_1} |\mathcal{F}_x(w_\sigma \phi_i)(z)|^2 dz \le \int_{\Omega} (w_\sigma \phi_i)^2 dx \le 1.$$

We fix a number $k_0 \in \mathbb{N}$ such that $\lambda_{k_0} \geq 1$. To this end, we can choose k_0 by (1.17) such that

$$\frac{2k_0}{eNd_N|\Omega|} \ge e$$

and

$$\ln\left(\frac{2k_0}{eNd_N|\Omega|}\right) - \ln\ln\left(\frac{2k_0}{eNd_N|\Omega|}\right) \ge \frac{N}{2}.$$

For simplicity, we may fix $k_0 := \frac{1}{2}e^{N+1}d_N|\Omega|$. And by increasing monotonicity, we have that $\lambda_k \ge 1$ for $k \ge k_0$.

For $k \geq k_0$, we choose r > 1 such that

$$\frac{2\omega_{N-1}}{N^2}r^N\ln r \geq \frac{\omega_{N-1}r^N}{N} \Longleftrightarrow r \geq e^{\frac{N}{2}} \quad \text{ and } \quad \frac{\omega_{N-1}r^N}{N}|\Omega| > 2(2\pi)^N k,$$

then we have that

$$\lambda_{k+1}(\Omega) \leq \frac{\frac{\omega_{N-1} r^N}{N} \left(\varrho_{2,\sigma} \frac{2}{N} \ln \frac{r^N}{e} + \varrho_{1,\sigma} \frac{\omega_{N-1}}{\sigma}\right) - (2\pi)^N \sum_{j=1}^k \lambda_j(\Omega) \int_{B_r} |\mathcal{F}_x(w_{\sigma}\phi_j)(z)|^2 dz}{\frac{\omega_{N-1} r^N}{N} \varrho_{2,\sigma} - (2\pi)^N \sum_{j=1}^k \int_{B_r} |\mathcal{F}_x(w_{\sigma}\phi_j)(z)|^2 dz}$$

Denote

$$A_1 = \frac{\omega_{N-1} r^N}{N} \left(\varrho_{2,\sigma} \frac{2}{N} \ln \frac{r^N}{e} + \varrho_{1,\sigma} \frac{\omega_{N-1}}{\sigma} \right) \quad \text{and} \quad A_2 = \frac{\omega_{N-1} r^N}{N} \varrho_{2,\sigma},$$

then

$$0 \leq \frac{A_{1} - (2\pi)^{N} \sum_{j=1}^{k} \lambda_{j}(\Omega) \int_{B_{r}} |\mathcal{F}_{x}(w_{\sigma}\phi_{j})(z)|^{2} dz}{A_{2} - (2\pi)^{N} \sum_{j=1}^{k} \int_{B_{r}} |\mathcal{F}_{x}(w_{\sigma}\phi_{j})(z)|^{2} dz} - \lambda_{k+1}(\Omega)$$

$$= \frac{\left(A_{1} - A_{2}\lambda_{k+1}(\Omega)\right) + (2\pi)^{N} \sum_{j=1}^{k} \left(\lambda_{k+1}(\Omega) - \lambda_{j}(\Omega)\right) \int_{B_{r}} |\mathcal{F}_{x}(w_{\sigma}\phi_{j})(z)|^{2} dz}{A_{2} - (2\pi)^{N} \sum_{j=1}^{k} \int_{B_{r}} |\mathcal{F}_{x}(w_{\sigma}\phi_{j})(z)|^{2} dz}$$

$$\leq \frac{\left(A_{1} - A_{2}\lambda_{k+1}(\Omega)\right) + (2\pi)^{N} \sum_{j=1}^{k} \left(\lambda_{k+1}(\Omega) - \lambda_{j}(\Omega)\right)}{A_{2} - (2\pi)^{N} k}$$

since $\lambda_{k+1}(\Omega) \ge \lambda_j(\Omega)$ for j < k+1 and $\int_{B_r} |\mathcal{F}_x(w_\sigma \phi_j)(z)|^2 dz \in (0,1)$. As a consequence, we obtain that

$$\lambda_{k+1}(\Omega) \le \frac{\frac{\omega_{N-1}r^N}{N} \left(\varrho_{2,\sigma} \frac{2}{N} \ln \frac{r^N}{e} + \varrho_{1,\sigma} \frac{\omega_{N-1}}{\sigma} \right) - (2\pi)^N \sum_{j=1}^k \lambda_j(\Omega)}{\frac{\omega_{N-1}r^N}{N} \varrho_{2,\sigma} - (2\pi)^N k}, \tag{4.6}$$

where

$$\frac{\omega_{N-1}r^N}{N}\varrho_{2,\sigma} - (2\pi)^N k > \frac{\omega_{N-1}r^N}{N} \frac{|\Omega|}{2} - (2\pi)^N k > 0.$$

We fix $\sigma = \sigma_1$ and first impose k, r > 1 such that

$$\frac{\omega_{N-1}r^N}{N}\varrho_{2,\sigma} = (2\pi)^N(k+1),$$

and take $r = k^{\frac{1}{N}}$ for $k \ge k_0 := \frac{eNd_N}{2}|\Omega|$, then we have that

$$\lambda_{k_0} \ge 0$$

and

$$(2\pi)^{N} \sum_{j=1}^{k+1} \lambda_{j}(\Omega) \leq \frac{2\omega_{N-1}}{N^{2}} \varrho_{2,\sigma} r^{N} \ln \frac{r^{N}}{e} + \frac{\omega_{N-1}^{2}}{N} \sigma^{-1} |\Omega|^{\frac{1}{2}} \sqrt{\varrho_{2,\sigma}} r^{N}$$

$$\leq (2\pi)^{N} \frac{2(k+1)}{N} \left(\ln(k+1) + \ln \frac{p_{N}}{\varrho_{2,\sigma}} + \frac{\omega_{N-1}}{\varrho_{2,\sigma}} \ln \ln \frac{p_{N}(k+1)}{\varrho_{2,\sigma}} \right)$$

$$\leq (2\pi)^{N} \frac{2(k+1)}{N} \left(\ln(k+1) + \ln \frac{2p_{N}}{|\Omega|} + \frac{2\omega_{N-1}}{\sqrt{|\Omega|}} \ln \ln \frac{2p_{N}(k+1)}{|\Omega|} \right),$$

where

$$p_N = \frac{2(2\pi)^N N}{\omega_N}$$
 and $\frac{|\Omega|}{2} \le \varrho_{2,\sigma} \le |\Omega|$.

Moreover, (1.19) follows by the lower bound (1.16) and the upper bound (1.18) directly. \Box

Proof of Theorem 1.7. Note that (1.17) reads as

$$\lambda_k(\Omega) \ge \frac{2}{N} \left(\ln k + \ln \left(\frac{2}{eNd_N|\Omega|} \right) - \ln \ln \left(\frac{2k}{eNd_N|\Omega|} \right) \right). \tag{4.7}$$

Using (1.18) and the monotonicity of $j \mapsto \lambda_j(\Omega)$, we take $m = \lfloor \frac{k}{\ln \ln k} \rfloor + 1$ and obtain

$$\lambda_{k+1}(\Omega) \leq \frac{1}{m} \left(\sum_{j=1}^{k+m} \lambda_j(\Omega) - \sum_{j=1}^{k} \lambda_j(\Omega) \right)$$

$$\leq \frac{2}{Nm} (k+m) \left(\ln(k+m) + 1 + \ln\left(\frac{p_N}{|\Omega|}\right) + \frac{\omega_{N-1}}{\sqrt{|\Omega|}} \ln\ln\left(\frac{p_N(k+1+m)}{|\Omega|}\right) \right)$$

$$- \frac{2}{Nm} k \left(\ln k + \ln\left(\frac{2}{eNd_N|\Omega|}\right) - \ln\ln\left(\frac{2k}{eNd_N|\Omega|}\right) \right)$$

$$\leq \frac{2}{N} \ln(k+m) + \frac{2}{N} (\ln\ln k) \ln\left(1 + \frac{1}{\ln\ln k}\right) + \delta_1 \ln\ln k + \delta_2 (\ln\ln k)^2$$

$$\leq \frac{2}{N} \ln k + \frac{4}{N} + \delta_3 (\ln\ln k)^2, \tag{4.8}$$

where δ_1 , δ_2 , $\delta_3 > 0$. Thus, we have that

$$\lim_{k \to +\infty} \frac{\lambda_k(\Omega)}{\ln k} = \frac{2}{N}.$$

We complete the proof.

4.3 Discussion about the counting function

From the asymptotic behavior (1.20) of $\lambda_k(\Omega)$, we obtain some asymptotic estimates for the counting function when $k \to \infty$.

Theorem 4.3. Let $\Omega \subset \mathbb{R}^N$ be a bounded domain and $\{\lambda_i(\Omega)\}_{i\in\mathbb{N}}$ be the sequence of eigenvalues of problem (1.3). Then for any $\delta > 0$

$$\lim_{t \to +\infty} \mathcal{N}(t)e^{-(\frac{N}{2}+\delta)t} = 0 \quad \text{and} \quad \liminf_{t \to +\infty} \mathcal{N}(t)e^{-(\frac{N}{2}-\delta)t} = +\infty. \tag{4.9}$$

Proof. From (1.20), for any $\epsilon \in (0, \frac{1}{4N})$, there exists $k_{\epsilon} > 0$ such that for $k \geq k_{\epsilon}$

$$(\frac{2}{N} - \epsilon) \ln k < \lambda_k(\Omega) \le (\frac{2}{N} + \epsilon) \ln k,$$

and let t > 0 be such that

$$(\frac{2}{N} - \epsilon) \ln k < \lambda_k(\Omega) \le t < \lambda_{k+1}(\Omega) \le (\frac{2}{N} + \epsilon) \ln(k+1),$$

which implies

$$e^{\frac{1}{\frac{2}{N} + \epsilon}t} - 1 < k < e^{\frac{1}{\frac{2}{N} - \epsilon}t}$$

This means

$$e^{\frac{1}{N+\epsilon}t} - 1 < \mathcal{N}(t) < e^{\frac{1}{N-\epsilon}t}$$

Therefore, we have that for any $\delta > 0$,

$$\lim_{t \to +\infty} \mathcal{N}(t) e^{(\frac{N}{2} + \delta)t} = 0 \quad \text{and} \quad \lim_{t \to +\infty} \mathcal{N}(t) e^{(\frac{N}{2} - \delta)t} = +\infty,$$

which ends the proof.

Remark 4.4. (i) From the upper bound (4.8) and the lower bound (1.16) there exists $\delta_4 > 0$ such that for $k > k_0$

$$\frac{2}{N}\ln k - \delta_4 \ln \ln k < \lambda_k(\Omega) \le t < \lambda_{k+1}(\Omega) \le (\frac{2}{N} + \epsilon) \ln(k+1) + \delta_4 \ln \ln k.$$

Similar arguments could be applied to improve the asymptotic behavior of N.

(ii) Another approach is to analyze the asymptotic behavior of $\sum_{j\in\mathbb{N}}(t-\lambda_j(\Omega))_+$.

Acknowledgements: This work is supported by NNSF of China, No: 12071189 and 12001252, by the Jiangxi Provincial Natural Science Foundation, No: 20202BAB201005, 20202ACBL201001.

References

- [1] W. Beckner, Pitt's inequality and the uncertainty principle, Proc. AMS 123(6), 1897-1905 (1995).
- [2] F. A. Berezin, Covariant and contravariant symbols of operators. *Izv. Akad. Nauk SSSR Ser. Mat.* 36, 1134-1167 (1972).
- [3] L. Caffarelli, S. Salsa and L. Silvestre, Regularity estimates for the solution and the free boundary to the obstacle problem for the fractional Laplacian, *Invent. Math.* 171, 425-461 (2008).
- [4] L. Caffarelli and L. Silvestre, Regularity theory for fully nonlinear integro-differential equations, *Comm. Pure Appl. Math.* 62(5), 597-638 (2009).
- [5] L. Caffarelli and L. Silvestre, Regularity results for nonlocal equations by approximation, *Arch. Ration. Mech. Anal.* 200(1), 59-88 (2011).
- [6] H-Y. Chen, P. Felmer and A. Quaas, Large solution to elliptic equations involving fractional Laplacian, *Annales de l'Institut Henri Poincaré-AN 32*, 1199-1228 (2015).
- [7] H-Y. Chen and L. Véron, Initial trace of positive solutions to fractional diffusion equations with absorption, *J. Funct. Anal.* 276, 1145-1200 (2019).
- [8] H-Y. Chen and T. Weth, The Dirichlet Problem for the Logarithmic Laplacian, *Comm. Part. Diff. Eq.* 44, 1100-1139 (2019).
- [9] H. Chen and A. Zeng, Universal inequality and upper bounds of eigenvalues for non-integer poly-Laplacian on a bounded domain, Calc. Var. Part. Diff. Eq. 56(5) (2017).
- [10] H. Chen and P. Luo, Lower bounds of Dirichlet eigenvalues for some degenerate elliptic operators, Calc. Var. Part. Diff. Eq. 54(3), 2831-2852 (2015).
- [11] Q. Cheng and H. Yang, Bounds on eigenvalues of Dirichlet Laplacian, Math. Ann. 337, 159-175 (2007).

- [12] Q. Cheng and G. Wei, A lower bound for eigenvalues of a clamped plate problem, Calc. Var. Part. Diff. Eq. 42(3/4), 579-590 (2011).
- [13] E. Correa and A. De Pablo, Nonlocal operators of order near zero, J. Math. Anal. Appl. 461, 837-867 (2018).
- [14] E. Di Nezza, G. Palatucci and E. Valdinoci, Hitchhiker's guide to the fractional Sobolev spaces, *Bull. Sci. Math.* 136(5), 521-573 (2012).
- [15] P. Felmer and A. Quaas, Fundamental solutions and Liouville type theorems for nonlinear integral operators, Adv. Math. 226, 2712-2738 (2011).
- [16] M. Felsinger, M. Kassmann and P. Voigt, The Dirichlet problem for nonlocal operators, Math. Zeit. 279, 779–809 (2015).
- [17] R. L. Frank and L. Geisinger, Refined semiclassical asymptotics for fractional powers of the Laplace operator, *Crelle's Journal* 712, 1-37 (2016).
- [18] R. L. Frank and L. Geisinger, Two-term spectral asymptotics for the Dirichlet Laplacian on a bounded domain, *Mathematical results in quantum physics*, World Sci. Publ., Hackensack, NJ, 2011, pp. 138–147.
- [19] RL. Frank, T. Konig and H. Tang, Classification of solutions of an equation related to a conformal log Sobolev inequality, arXiv: 2003.08135 (2020).
- [20] L. Geisinger, A short proof of Weyl's law for fractional differential operators, J. Math. Phys. 55, 011504 (2014).
- [21] S. Goyal and K. Sreenadh, On the Fučik spectrum of non-local elliptic operators, *Nonlinear Diff. Eq. Appl.* 21(4), 567-588 (2014).
- [22] E. M. Harrell II and S. Y. Yolcu, Eigenvalue inequalities for Klein-Gordon operators, *J. Funct. Anal.* 256(12), 3977-3995 (2009).
- [23] S. Jarohs and T. Weth, Local compactness and nonvanishing for weakly singular nonlocal quadratic forms, *Nonlinear Analysis*, Doi: 10.1016/j.na.2019.01.021 (2020).
- [24] M. Kassmann and A. Mimica, Intrinsic scaling properties for nonlocal operators, *J. Eur. Math. Soc.* 19(4), 983-1011 (2013).
- [25] P. Kröger, Estimates for sums of eigenvalues of the Laplacian, J. Funct. Anal. 126(1), 217-227 (1994).
- [26] A. Laptev, Dirichlet and Neumann eigenvalue problems on domains in Euclidean spaces, *J. Funct. Anal.* 151(2), 531-545 (1997).
- [27] A. Laptev and T. Weth, Spectral properties of the Logarithmic Laplacian, arXiv:2009.03395.
- [28] P. Li and S.-T.Yau, On the Schrödinger equation and the eigenvalue problem. *Commun. Math. Phys.* 88(3), 309-318 (1983).
- [29] E. Lieb, The number of bound states of one-body Schrödinger operators and the Weyl problem, *Proc. Sym. Pure Math.* 36, 241-252 (1980).
- [30] E. Lieb and M. Loss, *Analysis*, second ed., Graduate Studies in Mathematics, vol. 14, American Mathematical Society, Providence, RI, 2001.
- [31] A. Melas, A lower bound for sums of eigenvalues of the Laplacian, *Proc. Am. Math. Soc.* 131(2), 631-636 (2003).
- [32] R. Musina and A.I. Nazarov, On fractional Laplacians. Comm. Part. Diff. Eq. 39, 1780-1790 (2014).
- [33] G. Pólya, On the Eigenvalues of Vibrating Membranes (In Memoriam Hermann Weyl), *Proc. Lond. Math. Soc.* 3(1), 419-433 (1961).
- [34] G.V. Rozenbljum, Distribution of the discrete spectrum of singular operator. *Dokl. Akad. Nauk SSSR* 202, 1012-1015 (1972).
- [35] X. Ros-Oton and J. Serra, The Dirichlet problem for the fractional laplacian: regularity up to the boundary, J. Math. Pures Appl. 101, 275-302 (2014).
- [36] X. Ros-Oton and J. Serra, The Pohozaev identity for the fractional Laplacian, Arch. Ration. Mech. Anal. 213, 587-628 (2014).

- [37] H. Weyl, Das asymptotische Verteilungsgesetz der Eigenwerte linearer partieller Differentialgleichungen (mit einer Anwendung auf die Theorie der Hohlraumstrahlung). Math. Ann. 71(4), 441-479 (1912).
- [38] S.Y. Yolcu and T. Yolcu, Estimates for the sums of eigenvalues of the fractional Laplacian on a bounded domain, *Commun. Contemp. Math.* 15(3), 1250048 (2013).