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The early dark energy resolution to the Hubble tension in light of weak lensing surveys and lensing anomalies

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A constant Early Dark Energy (EDE) component contributing a fraction $f_{\text{EDE}}(z_c) \sim 10\%$ of the energy density of the universe around $z_c \simeq 3500$ and diluting as or faster than radiation afterwards, can provide a simple resolution to the Hubble tension, the $\sim 5\sigma$ discrepancy – in the Λ CDM context – between the H_0 value derived from early- and late-universe observations. However, it has been pointed out that including Large-Scale Structure (LSS) data, which are in $\sim 3\sigma$ tension with Λ CDM and EDE cosmologies, might break some parameter degeneracy and alter these conclusions. We reassess the viability of the EDE against a host of high- and low-redshift measurements, by combining LSS observations from recent weak lensing (WL) surveys with CMB, Baryon Acoustic Oscillation (BAO), growth function (FS) and Supernova Ia (SNIa) data. Introducing a model whose only parameter is $f_{\text{EDE}}(z_c)$, we report in agreement with past work a $\sim 2\sigma$ preference for non-zero $f_{\text{EDE}}(z_c)$ from PLANCK CMB data alone, while the tension with the local H_0 measurement from SH0ES is reduced below 2σ . Adding BAO, FS and SNIa does not affect this conclusion, while the inclusion of a prior on H_0 from SH0ES increase the preference for EDE over Λ CDM to the $\sim 3.6\sigma$ level. After checking the EDE non-linear matter power spectrum as predicted by standard semi-analytical algorithms via a dedicated set of N -body simulations, we test the 1-parameter EDE cosmology against WL data. We find that it does not significantly worsen the fit to the S_8 measurement as compared to Λ CDM, and that current WL observations do not exclude the EDE resolution to the Hubble tension. We also caution against the interpretation of constraints obtained from combining statistically inconsistent data sets within the Λ CDM cosmology. In light of the CMB lensing anomalies, we show that the lensing-marginalized CMB data also favor non-zero $f_{\text{EDE}}(z_c)$ at $\sim 2\sigma$, predicts H_0 in 1.4σ agreement with SH0ES and S_8 in 1.5σ and 0.8σ agreement with KiDS-VIKING and DES respectively. There still exists however a $\sim 2.5\sigma$ tension with the joint results from KiDS-VIKING and DES. With an eye on Occam’s razor, we finally discuss promising extensions of the EDE cosmology that could allow to fully restore cosmological concordance.

I. INTRODUCTION

In recent years, several tensions between probes of the early- and late-universe have emerged. First and foremost, there exists a strong mismatch between the prediction of the current expansion rate of the universe (known as Hubble constant) in the Λ cold dark matter (Λ CDM) model calibrated onto Planck CMB data, and its direct measurement using low redshift data (i.e., the classical distance ladder) [1, 2]. Originally, this “Hubble tension” was limited to the determination of the Hubble constant using type Ia supernovae by the SH0ES collaboration, whose latest *determination* is $H_0 = 74.03 \pm 1.42$ km/s/Mpc [3], while the *prediction* from the Λ CDM model inferred from Planck CMB data is $H_0 = 67.4 \pm 0.5$ km/s/Mpc [4]. In the last few years, tremendous progress have been made in measuring H_0 with alternative methods, such that nowadays there exist five other methods¹ to measure H_0 with few percent accuracy. Remarkably, various averages over these measurements (ex-

cluding correlated data) leads to H_0 values that ranges from 72.8 ± 1.1 and 74.3 ± 1.0 , in 4.5 to 6.3σ discrepancy with the prediction from Λ CDM [1, 2]. Similarly, it has been shown that the prediction from Λ CDM calibrated on any ‘early-universe’ data (e.g. BAO+SNIa+BBN [11–14], WMAP+SPT and/or ACT [12, 15]) is always in good agreement with that of Planck. A number of possible systematic effects affecting some of these measurements have been discussed (see e.g. [16–20]), yet the existence of several vastly different methods – none of which giving a value of H_0 smaller than ~ 70 km/s/Mpc – have triggered a wide range of theoretical activities to resolve the Hubble tension (see in particular [21] for a recent review). Indeed, this tension between different measurements of the Hubble constant could point to a major failure of the Λ CDM scenario, and hence to a new cosmological paradigm: that would be a new and unexpected breakthrough in cosmology.

CMB data do not provide an absolute measurement of H_0 . Rather, the value of H_0 is inferred within a given cosmological model from a measurement of the angular scale

¹ These include strong-lens time delays of quasars [5], Tip of the red giant branch from the ‘CCHP’ [6, 7] (and re-evaluation by the SH0ES team [8]), SNIa calibrated on Miras (an alternative to Cepheids) [9], water masers (sources of microwave stimulated

emission) in four galaxies at great distances [10] and Surface Brightness Fluctuations of distant galaxies [1].

of sound horizon $\theta_s \equiv r_s(z_*)/d_A(z_*)$, where $r_s(z_*)$ is the sound horizon at recombination and $d_A(z_*)$ is the angular diameter distance to recombination. The great challenge lies in that θ_s is nowadays measured at sub-percent-level accuracy with the latest CMB data [4]. This suggests two main ways of resolving the Hubble tension through new physics – based on the requirement to keep the key angular scale θ_s fixed – usually called *late*- and *early*-universe solutions. The first way boils down to changing the redshift evolution of the angular diameter distance in the late-universe, i.e. $z < z_*$, so as to force a higher H_0 , *without* changing $d_A(z_*)$ nor $r_s(z_*)$. The second way amounts in reducing $r_s(z_*)$ in the early-universe, which automatically requires to reduce $d_A(z_*)$ by the same amount to keep θ_s fixed, that is most naturally done by increasing the value of H_0 . A final, more subtle, way of resolving the H_0 tension comes from the fact that the position of the peaks receives an additional phase-shift from various effects, in particular from the gravitational pulling of CMB photons out of the potential wells by free-streaming neutrinos [22–24]. Suppressing this phase-shift can change the value of θ_s deduced from a CMB power spectra analysis and in turn significantly increase H_0 .

There have been many attempts to find extensions of the standard cosmological model, Λ CDM, which bring these estimates into agreement. However, theoretical explanations for the Hubble tension are not easy to come by. Late-time observables, especially BAO and luminosity-distance to SNIa, place severe limitations on late-time resolutions [12, 25–36]. On the other hand, early-time resolutions affect the physics that determines the fluctuations in the CMB. At first sight, given the precision measurements of the CMB from Planck, this might appear to be even more constraining than the late-time probes of the expansion rate. Excitingly, there are a few early-time resolutions which do not spoil the fit to current CMB temperature measurements [37–41], sometimes even *improving it* over Λ CDM.

EDE representing $\sim 10\%$ of the total energy density of the universe around matter-radiation equality and diluting faster than radiation afterwards has been shown to provide a very good resolution to this tension. However, taken at face value, this model triggers a number of questions. On the theoretical side, it suffers from a strong coincidence problem as the fluid needs to become dynamical around a key era of the universe. This is not without reminding the standard coincidence problem of DE that such models were originally introduced to resolve. However, this coincidence might be the sign of a very specific dynamics to be uncovered; in fact there exist models in which the field becomes dynamical precisely around matter-radiation equality, either because of a phase-transition triggered by some other process (e.g. the neutrino mass becoming of the order of the neutrino bath temperature [42]) or because of a non-minimal coupling to the Ricci curvature [43]. There exist also more concrete issues with Large-Scale Structure (LSS) observables, and in particular weak lensing

(WL) surveys, which this article aims to address. Indeed, a number of cosmic shear surveys (CFHTLenS [44], KiDS/Viking [45], DES [46], HSC [47]) have provided accurate measurements of the cosmological parameter $S_8 \equiv \sigma_8(\Omega_m/0.3)^{0.5}$ – where σ_8 measures the amplitude of fluctuations in a sphere of radius 8 Mpc/ h – which are systematically lower than the Λ CDM prediction. The significance of this “ S_8 tension” oscillates between 2 and 4σ depending on the experiments, such that the discrepancy cannot easily be attributed to a statistical fluke. The problem for EDE cosmologies lies in the fact that the prediction for the value of S_8 is somewhat higher than that of Λ CDM. Therefore, taken at face values, these experiments pose a challenge to EDE cosmologies, and could exclude these models as a resolution to the Hubble tension [48]. A similar conclusion was reached in Refs [49, 50] with the inclusion of BOSS data in the effective field theory (EFT) of LSS framework.

In this paper, we analyze the EDE cosmology resolving the Hubble tension in light of the latest PLANCK data (and the more precise polarization measurement) and confront it to the KiDS-VIKING measurement of the cosmic shear power spectrum [51] and the joint measurement of S_8 from KiDS-VIKING+DES². The KiDS-VIKING+DES measurements however rely on modelling the non-linear matter power spectrum on relatively small scales. This is done within numerical Einstein-Boltzmann solvers such as CAMB [53] or CLASS [54, 55], through the HALOFIT [56, 57] or HMCODE [58] algorithms, which have not been calibrated for EDE cosmologies. We thus check the predictions of these algorithms against the results of a set of dedicated cosmological N -body simulations, confirming that the qualitative departures from Λ CDM arising in the EDE cosmology are small enough to make use of these standard algorithms. We perform a series of Monte Carlo Markov Chain (MCMC) analyses with various combination of the latest Planck, BAO, growth factor and SNIa Luminosity distance measurements, the SH0ES measurement of H_0 , and KiDS/Viking/DES data, in order to assess whether current observations exclude the EDE resolution to the Hubble tension.

We find that, while the S_8 prediction from the best fit EDE cosmology is indeed $\sim 2.5\sigma$ higher than the measurement, KiDS data currently provide very little constraining power on the EDE parameters. Yet, it has been found in Ref. [51, 59] that a combination of KiDS and DES-Y1 (after re-calibration of the DES photo-metric redshifts) and yields $S_8 = 0.755^{+0.019}_{-0.021}$, a result that is in 3.2σ tension with Planck Λ CDM prediction³. At such a level of discrepancy, one should be cautious when inter-

² The re-analysis of BOSS data in the EFT of LSS is performed elsewhere [52].

³ The joint analysis of KiDS1000+BOSS+2dfLenS has determined $S_8 = 0.766^{+0.020}_{-0.014}$ [60] in 3σ tension with Planck. Making use of these data would not affect our conclusions.

preting results obtained from combining Planck and WL data, even within Λ CDM. Indeed, we show that despite the inclusion of a Gaussian S_8 likelihood, the resulting cosmological model yields a very bad fit to the S_8 data, while providing very strong constraints on any parameter correlated with S_8 (e.g. ω_{cdm} , A_s , $f_{\text{EDE}}(a_c)$). It is easily conceivable that the resolution to the S_8 tension lies elsewhere (whether new physics related – or not – to the EDE, or systematic effects), such that any constraints derived from these combined data are artificial.

It is also possible that this anomaly comes from CMB measurements. Interestingly, it has been found that there exists a lensing anomaly in Planck data (strengthened⁴ in the latest data released despite extensive efforts from the Planck team to pin down a possible systematic effect) [4, 62], which could be related to the S_8 tension [63, 64]. Indeed, the amplitude of the CMB lensing power spectrum deduced from the smoothing of the acoustic peaks at high- ℓ 's is higher than that predicted by the Λ CDM cosmology obtained from the ‘unlensed’ part of the CMB power spectrum⁵. Moreover, it is also higher than that deduced from the lensing reconstruction such that, while this anomaly looks like lensing, it cannot be attributed to some extra source of CMB lensing. Once marginalizing over the lensing information (which by itself is an interesting consistency check of the Λ CDM cosmology), it has been shown that the reconstructed cosmology has a smaller A_s and ω_{cdm} (as well as a higher H_0) and it shows no S_8 tension, but a remnant $\sim 3.5\sigma$ Hubble tension [63, 64]. Interestingly, the cosmology deduced once marginalizing over the lensing information is in better agreement with the recent results from the SPTPol [62, 65, 66], which shows no tension with the LSS measurement of S_8 , a weaker H_0 tension, and no lensing anomaly. Pin-pointing the source of such lensing anomaly (perhaps a simple statistical fluke, although quantifying is likelihood and how to treat it is complicated) is therefore of utmost importance to understand whether the S_8 tension derives from it.

Motivated by this fact, we perform an analysis of the Λ CDM and EDE cosmology against *Planck* and a prior on S_8 from the joint DES-Y1 and KiDS results, while marginalizing over the lensing information. We find that both the unlensed Λ CDM and EDE cosmology spectrum agrees better with LSS data, and that the presence of EDE does not affect the amount of anomalous lensing. This means that the anomalous lensing is not due to the EDE, and also that the success of EDE is not due to opening up a new degeneracy direction with some exotic lensing parameters.

The paper is structured as follows. In sec. II, we present the phenomenological EDE models studied in this work and compare it to other models from the literature.

In sec. III, we perform an ‘anatomy’ of the EDE resolution to the Hubble tension, to understand how each data set reacts to the presence of EDE. We additionally introduce the new baseline 1-parameter EDE model that is favored by PLANCK 2018 data. We present our N -Body simulations validating the HMcode prediction and confront the 1-parameter EDE model against WL data in sec. IV. We discuss the H_0 and S_8 tension in light of the lensing-marginalized CMB spectrum in sec. V and suggest some promising ways of restoring cosmological concordance in an extended EDE cosmology. Finally, we conclude in sec. VI.

II. THE PHENOMENOLOGICAL EARLY DARK ENERGY MODEL RESOLVING THE HUBBLE TENSION

The possible presence of a dark energy component before last-scattering has been studied for more than a decade [67, 68]. These alternative cosmological realizations have little to do with that under study here, as they typically assume tracking equation of state at early times. The idea of an anomalous era of expansion triggered by a frozen scalar field as a resolution to the Hubble tension was introduced in Ref. [69], where a background-only computation was shown to alleviate the Hubble tension. However, it is the work of Ref. [37] that showed through a fluid approximation the key role played by perturbations in the scalar field to allow for a resolution of the Hubble tension. Since this work, the treatment of the EDE component has been improved [70–72], and augmented to deal with alternative potentials and better motivated underlying fundamental models [39, 42, 70, 72–80]. In particular, it has been shown that *Planck* data not only provide a detection of the background dynamics of the EDE component, but also severely restricts the dynamics of perturbations [71, 72]. As such, *Planck* data allows for pinning down directly properties of the EDE, making the choice of model crucial. They favor either non-canonical kinetic term whereby the equation of state w is approximately equal to the effective sound speed c_s^2 [72], or potential that flattens close to the initial field value [71].

In this work, we study the modified axion potential introduced in Refs. [37, 69, 71, 81, 82],

$$V_n(\Theta) = m^2 f^2 [1 - \cos(\Theta)]^n, \quad (1)$$

where m represent the axion mass, f the decay constant and $\Theta \equiv \phi/f$ is a re-normalized field variable, so that $-\pi \leq \Theta \leq \pi$. We assume that the field always starts in slow-roll the background dynamics and without loss of generality we restrict $0 \leq \Theta_i \leq \pi$.

This potential is a phenomenological generalization of the well motivated axion-like potential (which can be recovered by setting $n = 1$) that arise generically in string theory [83–86]. Such a potential may be generated by higher-order instanton corrections [87], but taken at face

⁴ See also Ref. [61] for an independent re-analysis.

⁵ We will later on dub this the ‘unlensed’ Λ CDM cosmology for simplicity.

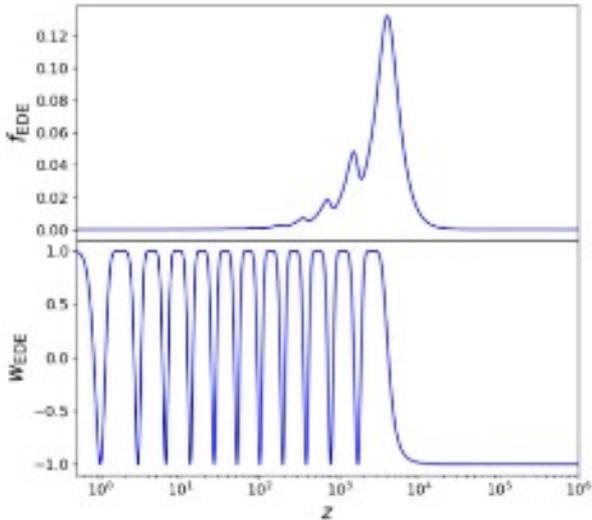


FIG. 1. Fractional energy density (upper) and equation of state (lower) in the EDE as a function of redshift. Cosmological parameters are set to the best-fit values from the PLANCK+BAO+SNIa+SH0ES analysis (see third column of Table I).

values would suffer from a strong fine-tuning issues necessary to the cancelling of the lowest orders instantons. Therefore, it should not be interpreted beyond a phenomenological description. We note that similar forms of potential, with power law minima and flattened “wings” have been used in the context of inflationary physics, as well as dark energy (see, e.g., Refs. [88–91]). Still, this form was devised to allow for flexibility in the background dynamics after the field becomes dynamical, and it also provides an excellent fit to both *Planck* and SH0ES data. In fact, to the best of our knowledge, it corresponds to the EDE scenario that leads to the best combined χ^2 of the cosmological data-sets under study (although the better theoretically motivated model studied in Ref. [39, 80] seems to perform equally well).

We refer to Refs. [71, 82] for all necessary details about the model. However, the key features can be summarized as follows: at early times the scalar field is frozen due to Hubble friction, until the Hubble rate drops below its mass value; the field then starts moving in the potential, and eventually oscillating around the minimum, at which point the energy density dilutes at a rate dictated by the asymptotic equation of state $w(n) = (n-1)/(n+1)$ (e.g., Refs. [82, 92, 93]). In Fig. 1 we show the redshift evolution for the fractional energy density and the equation of state in the EDE.

We can trade three out of the four model parameters $\{m, f, n, \Theta_i\}$ for phenomenological parameters: the first two of them describing the fractional energy density $f_{\text{EDE}}(z_c)$ at the critical redshift z_c where the field be-

comes dynamical and the asymptotic equation of states after the field becomes dynamical $w(n) = (n-1)/(n+1)$, respectively; the last degree of freedom lies in the dynamics of linear perturbations, whose phenomenology is captured by the effective sound speed c_s^2 . However, within the EDE scalar field scenario under study, such freedom is intrinsically encoded in the choice of the initial field value⁶ Θ_i , once the other phenomenological parameters have been fixed.

To perform our analyses, we use the modified version of the Einstein-Boltzmann code CLASS [54, 55] presented in Ref. [71]. The code is publicly available at <https://github.com/PoulinV/AxiCLASS> (the latest version, used for this study, can be found in the “devel” branch).

III. CONFRONTING EDE TO PLANCK 2018, BOSS AND PANTHEON DATA

In this section, we test the EDE scenario with various combinations of data-sets, in order to extract the cosmology that would resolve the Hubble tension, and compare with results from past literature making use of Planck 2015 data. We will test our phenomenological model against:

- PLANCK^{TTTEEE+ $\phi\phi$} : the high- ℓ TT,TE,EE, low- ℓ TT and EE data from *Planck* 2018 through the baseline PLIK, COMMANDER and SMALL likelihoods [94, 95], alone and combined with the lensing amplitude reconstruction (SIMCA likelihood); we make use of a Cholesky decomposition as implemented in MONTEPYTHON-V3 to handle the large number of nuisance parameters [96].
- BAO: the measurements from 6dFGS at $z = 0.106$ [25], SDSS MGS at $z = 0.15$ [26], and BOSS DR12 data at $z = 0.38, 0.51$ and 0.61 [31].
- FS: the measurements of the growth function $f\sigma_8(z)$ (FS) from the CMASS and LOWZ galaxy samples of BOSS DR12 at $z = 0.38, 0.51$, and 0.61 [31]. In practice, we make use of the “consensus” BAO and FS result that combines both in a single likelihood⁷.
- PANTHEON: the Pantheon SNIa catalogue, spanning redshifts $0.01 < z < 2.3$; we marginalize over the nuisance parameter \mathcal{M} describing the SNIa calibration.

⁶ In practice, it is the curvature of the potential, $\partial^2 V(\Theta)/\partial^2 \Theta$, close to the initial field value Θ_i that dictates the last of degree of freedom in the perturbation dynamics [71, 82].

⁷ We correct for a small mistake in the standard MONTEPYTHON-V3 implementation of the likelihood. This also explains why Refs. [37, 71], which did not correct for this mistake, reported slightly different constraints when including this likelihood.

- SH0ES: the SH0ES result, modelled with a Gaussian likelihood centered on $H_0 = 74.03 \pm 1.42$ km/s/Mpc [3]; however, choosing a different value that *combines* various direct measurements would not affect the result, given their small differences.

A. Baseline analysis: anatomy of the 3-parameter EDE model resolving the Hubble tension

Our baseline cosmology consists in the following combination of the six Λ CDM parameters $\{\omega_b, \omega_{\text{cdm}}, H_0, n_s, A_s, \tau_{\text{reio}}\}$, plus three parameters describing the EDE sector, namely $\{f_{\text{EDE}}(a_c), \text{Log}_{10}(a_c), \Theta_i\}$. We use wide flat priors on all these parameters. We follow the Planck convention and assume two massless neutrinos and one massive with $m_\nu = 0.06$ eV. We perform our MCMC analyses using MONTEPYTHON-V3, and consider chains to be converged with the Gelman-Rubin criterion⁸ $R - 1 < 0.1$ [97]. To extract best-fit parameters, we make use of the MINUIT algorithm [98] through the iMINUIT python package⁹. Starting from Planck only, we now discuss the impact of adding data-sets on the reconstructed EDE parameters. We compare the evolution of the χ^2 in the EDE cosmology as we add data-sets, to that of the Λ CDM model in the same combined fit, with and without SH0ES data. The results are presented in Tab. I and we show the 1D and 2D posterior distributions of $\{H_0, f_{\text{EDE}}(z_c), \Theta_i, \text{Log}_{10}(z_c), \omega_{\text{cdm}}, n_s, S_8\}$ in Fig. 2. All relevant χ^2 information is given in App. A.

Planck TT,TE,EE only: with PLANCK^{TTTEEE} data only and three free parameters, the EDE model under study is not detected. In agreement with Refs. [48, 80], we find that the fraction of EDE at z_c is limited to¹⁰ $f_{\text{EDE}}(z_c) < 0.088$, while $\text{Log}_{10}(z_c)$ and Θ_i are unconstrained. Interestingly, we also find that the best fit within PLANCK data only has $\{f_{\text{EDE}}(z_c) \sim 8.5\%$, $\text{Log}_{10}(z_c) \sim 3.56$, $\Theta_i \sim 2.8$, $H_0 \sim 70.5$ km/s/Mpc $\}$ and a $\Delta\chi_{\text{min}}^2 \equiv \chi_{\text{min}}^2(\Lambda\text{CDM}) - \chi^2(\text{EDE}) \simeq -5$ in favor of the EDE model¹¹. One can already note a curiosity: the best fit value of $f_{\text{EDE}}(z_c)$ is very close to the 2σ bound that we obtain. This, as we will discuss later, is due to

the choice of flat, uninformative prior on $\text{Log}_{10}(z_c)$ and Θ_i .

Planck TT,TE,EE+SH0ES: Once a prior on H_0 given by SH0ES is included in the analysis, the sampler explores more easily a part of parameter space with higher H_0 values, and the EDE is now well detected: $\{f_{\text{EDE}}(z_c) \simeq 0.11_{-0.031}^{+0.036}$, $\text{Log}_{10}(z_c) = 3.6_{-0.039}^{+0.14}$, $\Theta_i = 2.569_{-0.032}^{+0.36}\}$, with $H_0 = 71.4 \pm 1$ km/s/Mpc. This is in excellent agreement with results from past literature [37, 71, 72, 80]. Remarkably, the best fit values of both Θ_i and $\text{Log}_{10}(z_c)$ are in perfect agreement with that obtained PLANCK only. This is highly non-trivial, and seem to indicate that PLANCK does favor the region of the $\{\text{Log}_{10}(z_c), \Theta_i\}$ -space that resolves the Hubble tension. However, the best fit fraction reaches 15%, a value that one would naively consider to be strongly excluded by the PLANCK only analysis. In fact, that is not the case, as the fit to PLANCK data is barely affected by the additional H_0 prior, while one can get a perfect fit of SH0ES data. Concretely, the $\chi_{\text{min}}^2(\text{EDE})$ when fitting PLANCK+SH0ES increases by ~ 3 , such that even in this combined fit, the χ^2 of PLANCK data is smaller than that of Λ CDM fitted on PLANCK data only. This indicates that the limit on $f_{\text{EDE}}(z_c)$ derived in a PLANCK only analysis is not robust, as it is entirely driven by our choice of flat priors. This was also discussed in Refs. [71, 80], and the reason for that is clear¹²: there exists a strong χ^2 degeneracy in PLANCK data between Λ CDM and the EDE cosmology, that, given our choice of uninformative priors on Θ_i and $\text{Log}_{10}(z_c)$, leads to an artificially strong bound on $f_{\text{EDE}}(z_c)$. Indeed, once $f_{\text{EDE}}(z_c)$ drops below $\lesssim 4\%$ (as seen from the 2D posterior), its impacts on the power spectrum is not detectable given current measurement accuracy. As a result, the quantity $\text{Log}_{10}(z_c)$ and Θ_i have no impact on observables, such that any choice of $\text{Log}_{10}(z_c)$ and Θ_i leads to a cosmology indistinguishable from Λ CDM. Therefore, the sampler spends much more time exploring this degeneracy direction, rather than efficiently sampling the narrow degeneracy between $f_{\text{EDE}}(z_c)$ and H_0 , which requires a specific choice of $\text{Log}_{10}(z_c)$ and Θ_i to appear. Following Ref. [80], we will discuss a natural way to alleviate this issue in Section III B.

Planck TT,TE,EE+PP+BAO+SNIa+SH0ES: We now add to our analysis the lensing reconstruction PLANCK ^{$\phi\phi$} , the PANTHEON SNIa data-set, and the BAO data from BOSS. Strikingly, the addition of these three data-sets has almost no impact on the reconstructed posteriors, nor on the best fit. This is far from a trivial test to pass, as many of the suggested resolutions to the Hubble tension are strongly constrained by the

⁸ Most chains are in fact converged at the $R - 1 \sim 0.01$ level, this somewhat ‘loose’ but reasonable criterion was only used once including KiDS-VIKING data, which are much longer to converge.

⁹ <https://iminuit.readthedocs.io/>

¹⁰ Hereinafter, we quote 1-sided constraints at 95%C.L., and two-sided ones at 68%C.L.

¹¹ To guide the reader, we mention that a 1σ shift in the quality of the fit to PLANCK data roughly corresponds to a $\Delta\chi^2$ of ~ 6 (see the distribution of PLANCK’s χ^2 in the tables available at this link).

¹² Here, let us mention that Ref. [49] make the comment that such degeneracy does not exist. This is of course only true because they include the 3σ discrepant S_8 data to their analysis. The degeneracy is very clear within PLANCK data.

3-parameter EDE cosmology				
Parameter	PLANCK ^{TTTEEE}	+SH0ES	+PLANCK ^{φφ} +BAO+PANTHEON	+FS
H_0 [km/s/Mpc]	68.29(70.49) ^{+0.75} _{-1.3}	71.49(73.05) ± 1.2	71.34(72.41) ⁺¹ _{-1.1}	71.01(71.96) ^{+1.1} _{-1.1}
100 ω_b	2.252(2.270) ^{+0.019} _{-0.023}	2.284(2.281) ^{+0.022} _{-0.024}	2.282(2.292) ^{+0.021} _{-0.022}	2.28(2.285) ^{+0.021} _{-0.022}
ω_{cdm}	0.1232(0.1278) ^{+0.0019} _{-0.004}	0.13(0.135) ^{+0.0042} _{-0.004}	0.1297(0.1327) ^{+0.0036} _{-0.0039}	0.1289(0.1323) ± 0.0039
10 ⁹ A_s	2.116(2.124) ^{+0.035} _{-0.041}	2.153(2.160) ^{+0.036} _{-0.042}	2.152(2.183) ^{+0.031} _{-0.035}	2.144(2.135) ^{+0.032} _{-0.033}
n_s	0.9706(0.9829) ^{+0.0058} _{-0.0087}	0.9889(0.9966) ^{+0.0076} _{-0.0075}	0.9878(0.9963) ^{+0.0066} _{-0.007}	0.9859(0.9895) ^{+0.007} _{-0.0071}
τ_{reio}	0.0552(0.0524) ^{+0.0076} _{-0.0086}	0.0586(0.0558) ^{+0.0077} _{-0.0091}	0.0585(0.0633) ^{+0.007} _{-0.008}	0.0574(0.0528) ^{+0.007} _{-0.0079}
$f_{\text{EDE}}(z_c)$	< 0.088(0.085)	0.108(0.152) ^{+0.035} _{-0.028}	0.106(0.133) ^{+0.031} _{-0.028}	0.097(0.126) ^{+0.035} _{-0.029}
Log ₁₀ (z_c)	3.705(3.569) ^{+0.37} _{-0.22}	3.612(3.569) ^{+0.13} _{-0.049}	3.615(3.602) ^{+0.11} _{-0.029}	3.61(3.572) ^{+0.13} _{-0.054}
Θ_i	unconstrained (2.775)	2.604(2.756) ^{+0.33} _{-0.087}	2.722(2.759) ^{+0.17} _{-0.092}	2.557(2.705) ^{+0.37} _{-0.025}
100 θ_s	1.04165(1.04371) ^{+0.00039} _{-0.00034}	1.04131(1.04070) ^{+0.00039} _{-0.0004}	1.04143(1.04122) ^{+0.00036} _{-0.00039}	1.04145(1.04098) ^{+0.00038} _{-0.00039}
$r_s(z_{\text{rec}})$	142.8(140.1) ^{+1.9} _{-0.72}	138.8(136.4) ^{+1.7} _{-1.9}	139(137.5) ^{+1.7} _{-1.7}	139.4(137.8) ^{+1.7} _{-1.9}
S_8	0.839(0.834) ^{+0.018} _{-0.019}	0.838(0.842) ^{+0.018} _{-0.019}	0.838(0.846) ± 0.013	0.837(0.838) ± 0.013
Ω_m	0.314(0.304) ^{+0.0088} _{-0.0091}	0.3004(0.2969) ^{+0.0079} _{-0.0084}	0.301(0.2980) ^{+0.0051} _{-0.0055}	0.3022(0.3009) ^{+0.0053} _{-0.0054}
$\Delta\chi^2_{\text{min}}$ (ΛCDM w/ SH0ES)	—	-20.8	-19.1	-18.7
$\Delta\chi^2_{\text{min}}$ (ΛCDM w/o SH0ES)	-4.9	-1.5	-0.02	-0.6

TABLE I. The mean (best-fit) $\pm 1\sigma$ error of the cosmological parameters reconstructed from the combined analysis of various data sets (from left to right, each column adds a set of data to the previous one). We also report the $\Delta\chi^2_{\text{min}}$ with respect to a Λ CDM fit to the same data-sets, with and without a prior on H_0 from SH0ES.

addition of these data-sets [12, 25–35]. However, as noted in Refs. [37, 71, 80], we find that the reconstructed ω_{cdm} and n_s in the EDE cosmology are somewhat higher than in Λ CDM, such that the S_8 tension is slightly increased. As suggested in past literature [48–50], this opens up the possibility of constraining the EDE resolution using LSS data, and in particular the S_8 measurement from weak gravitational lensing surveys. However, combining KiDS+VIKING/HSC data with PLANCK to constrain extension to Λ CDM can be problematic as: i) they require the ability to predict the non-linear power spectrum at relatively small scales in models beyond Λ CDM; ii) the Λ CDM best fit model from PLANCK is not a good fit to these data.

All Data: As a starter, we add the ‘consensus’ $f\sigma_8$ BOSS likelihood to the analysis, which is consistent with the Λ CDM model from PLANCK; we find a mild $\sim 0.4\sigma$ decrease in the reconstructed mean, now being $H_0 \simeq 71 \pm 1$ km/s/Mpc and $f_{\text{EDE}} \simeq 0.1 \pm 0.03$. This is consistent with the fact that the $f\sigma_8$ measurements are sensibly lower than the Λ CDM prediction, while the EDE cosmology leads to slightly larger values. Still, the χ^2 of the FS data is barely affected; in fact, as before, Λ CDM provides a slightly worse fit to the joint dataset, even when the SH0ES prior is not included in the analysis. Before including weak lensing measurements to the analysis, we therefore conclude that the 3-parameter EDE model under study performs very well in resolving the Hubble tension, but future measurement of $f\sigma_8$ will certainly put the model under crucial tests.

B. Towards a 1-parameter resolution to the Hubble tension

Before turning to the inclusion of WL data, we show that the apparently tight bounds obtained when the SH0ES prior on H_0 is not included is due to our choice

of uninformative priors for Θ_i & $\text{Log}_{10}(z_c)$. In fact, one can strongly weaken the bound on f_{EDE} by *reducing* the EDE parameter space to a suitable choice of $\text{Log}_{10}(z_c)$ and Θ_i . This might sound counter-intuitive: in principle, one expects to relax constraints on a given parameter by enlarging the parameter space such as to introduce a new degeneracy. Here however, it is the poor prior choice which leads to a strong bound on $f_{\text{EDE}}(z_c)$ independently of the data combination. Fixing $\text{Log}_{10}(z_c)$ and Θ_i to some fiducial values surely rises the question of what values should one choose. In a realistic scenario, one might know these values a priori; one example is the scenario discussed in Ref. [42] in which a scalar field experiences a phase-transition around the redshift at which neutrinos becomes non-relativistic, such that the critical redshift is specified by the value of the neutrino mass, while Θ_i is set by the dynamics of the phase-transition (see also Refs. [43, 74, 77–80] for different EDE models with fewer free parameters). Here however, we have been considering a phenomenological model whose primary characteristics is to have enough freedom to extract information from the data to resolve the tension – we will therefore make use of that information and fix Θ_i & $\text{Log}_{10}(z_c)$ to their best fit value from PLANCK data only – which, we recall, are close-to-identical to that obtained in the combined fit. We report in Table II the reconstructed cosmological parameters from PLANCK only and from the combined fit of all data, with and without including SH0ES. We show the reconstructed 2D posteriors of $\{f_{\text{EDE}}(z_c), H_0, S_8\}$ in Figure 3. Notice how the degeneracy direction $f_{\text{EDE}}(z_c) - H_0$ clearly opens up. Furthermore, the mild $\Delta\chi^2$ preference in favor of the EDE cosmology now leads to reconstructing $f_{\text{EDE}}(z_c) = 0.082 \pm 0.037$, i.e., a $\sim 2\sigma$ preference for non-zero EDE *from PLANCK data only*. The inferred $H_0 = 70.1 \pm 1.4$ km/s/Mpc is now in agreement with the

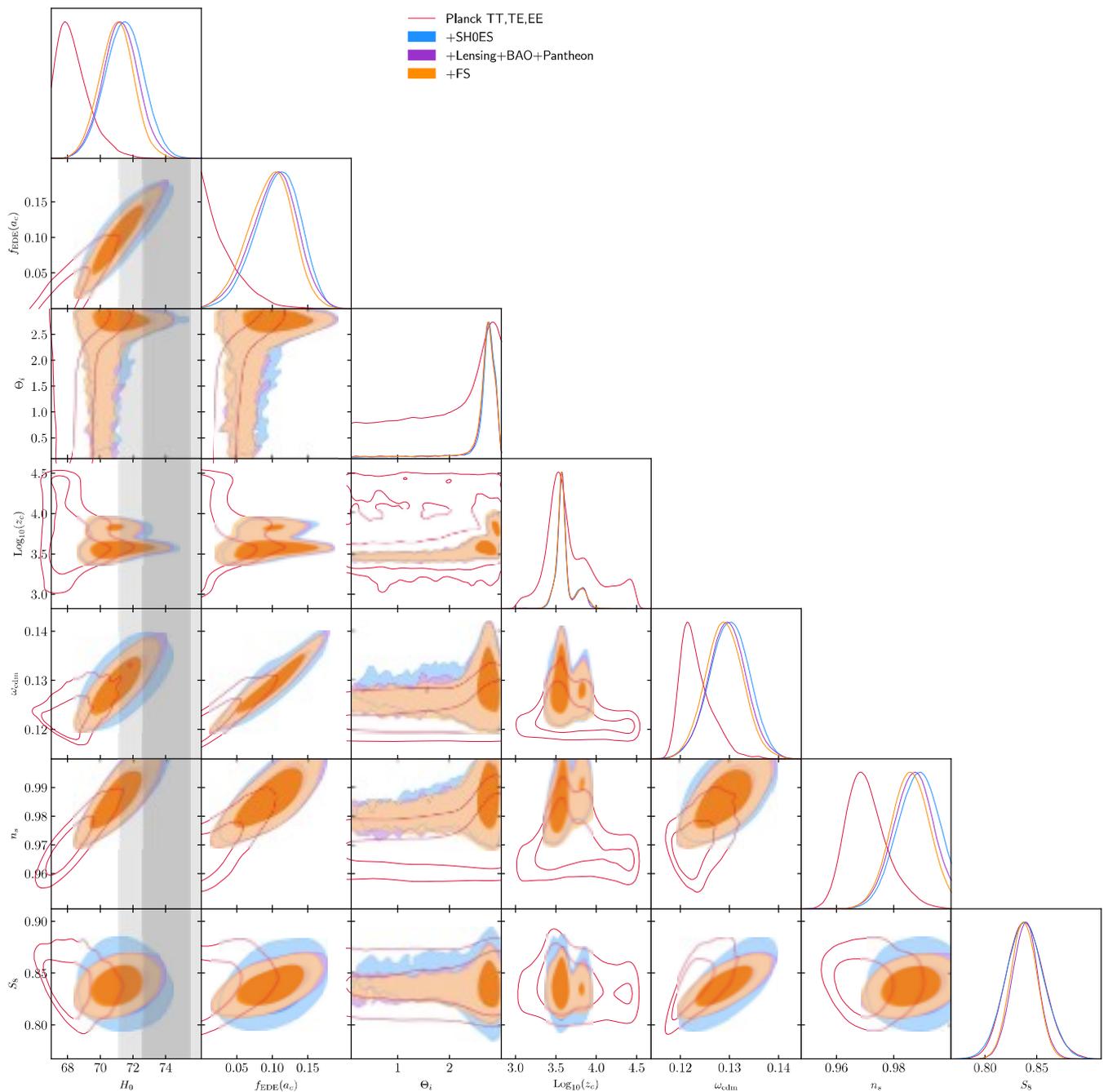


FIG. 2. Reconstructed 2D posterior distributions of a subset of parameters for various data set combinations (see legend) in the 3-parameters EDE cosmology.

SH0ES determination at better than¹³ 2σ . The addi-

¹³ From here on, we quote ‘tension’ and ‘agreement’ assuming Gaussian posteriors for simplicity. While this is surely a crude approximation (to be tested elsewhere [99] following Ref. [100]), we believe this is justified because the posterior of interest (H_0, S_8) are close to Gaussian, and since this is the approach followed

tion of BAO, FS and PANTHEON measurements has little impact; the reconstructed EDE fraction shifts downward by $\sim 0.3\sigma$, slightly degrading the success of the resolution to the Hubble tension, while the 2σ prefer-

by collaborations when quoting tensions in the Λ CDM context [3, 4].

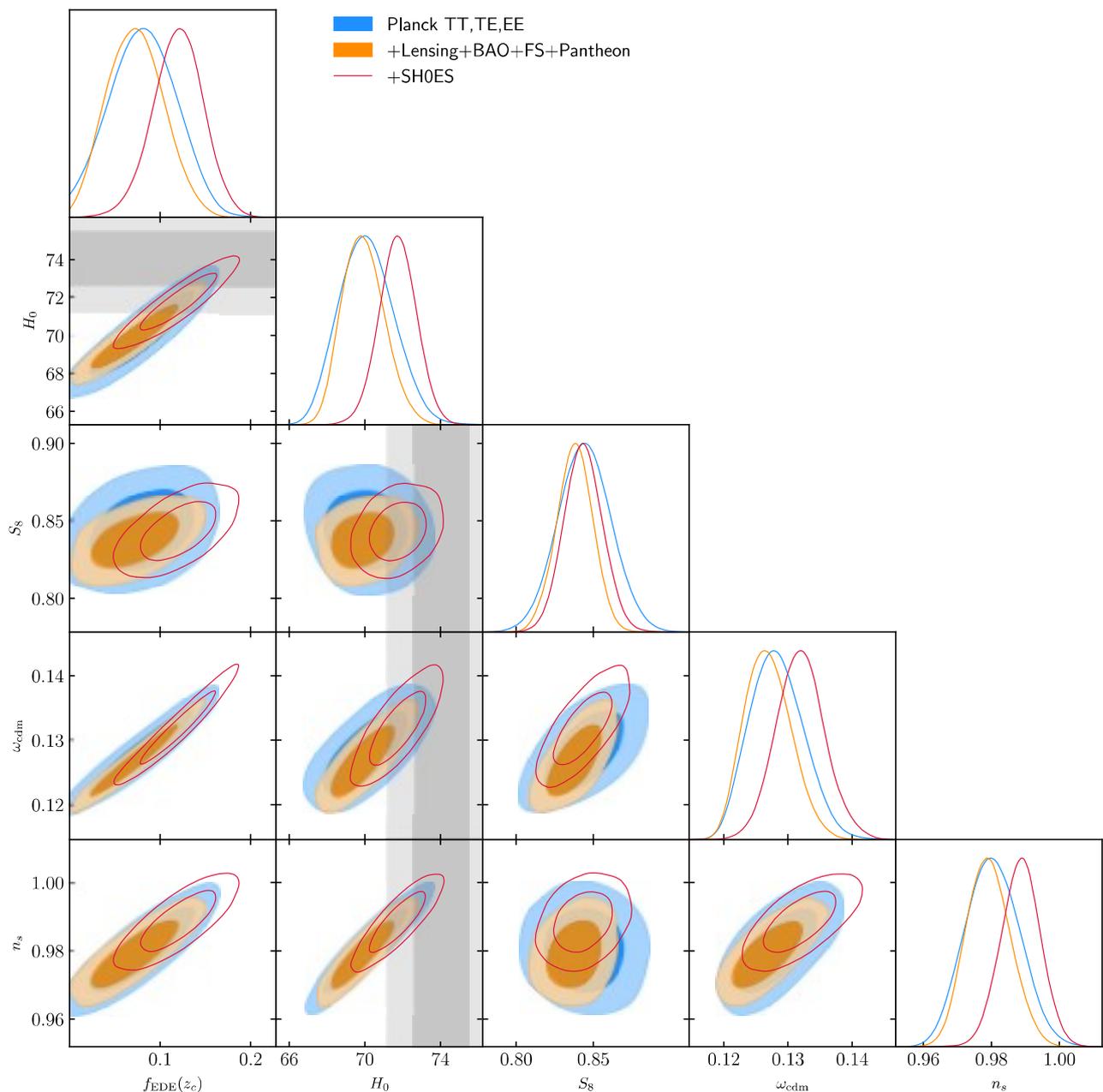


FIG. 3. Reconstructed 2D posterior distributions of a subset of parameters for various data set combinations (see legend) in the 1-parameter EDE cosmology.

ence for non-zero EDE is still present. These results are in excellent agreement with those presented in Ref. [80] for a different EDE model. Finally, the inclusion of a prior from SH0ES pulls up the fraction of EDE to $f_{\text{EDE}}(z_c) = 0.118 \pm 0.029$ and the value of $H_0 = 71.7 \pm 1$ km/s/Mpc, at the cost of a small degradation in χ^2_{min} ($\Delta\chi^2 \sim +6$). Yet, as before, the χ^2_{min} of the *combined* fit PLANCK+BAO+FS+PANTHEON+SH0ES in the 1-parameter EDE cosmology is slightly lower than a Λ CDM fit to PLANCK+BAO+FS+PANTHEON (no

SH0ES). This attests that, despite this small degradation in χ^2_{min} , the goodness of fit is still excellent. However, as discussed previously, the values of S_8 are in significant tension with weak lensing measurements, and one might expect that it is possible to strongly constrain the EDE model by including LSS data. We study this possibility in detail in Section IV, following Refs. [48–50].

1-parameter EDE cosmology			
Parameter	PLANCK ^{TTTEEE}	+PLANCK ^{$\phi\phi$} +BAO+PANTHEON+FS	+SH0ES
H_0 [km/s/Mpc]	70.10(70.83) ^{+1.4} _{-1.6}	70.00(69.84) ^{+0.98} _{-1.4}	71.71(72.21) ^{+1.0} _{-0.95}
100 ω_b	2.258(2.265) ^{+0.013} _{-0.018}	2.259(2.263) ^{+0.015} _{-0.016}	2.273(2.282) \pm 0.013
ω_{cdm}	0.1282(0.1306) ^{+0.0039} _{-0.0046}	0.1270(0.1265) ^{+0.0033} _{-0.0043}	0.1317(0.1310) ^{+0.0036} _{-0.0037}
10 ⁹ A_s	2.137(2.164) ^{+0.037} _{-0.042}	2.131(2.118) ^{+0.03} _{-0.036}	2.15(2.140) ^{+0.032} _{-0.031}
n_s	0.9803(0.9851) ^{+0.0079} _{-0.0085}	0.9795(0.9788) ^{+0.0063} _{-0.0074}	0.9884(0.9917) ^{+0.0062} _{-0.0057}
τ_{reio}	0.0558(0.0602) ^{+0.0079} _{-0.0085}	0.0553(0.0522) ^{+0.0069} _{-0.0075}	0.0561(0.0536) ^{+0.0071} _{-0.0076}
$f_{\text{EDE}}(z_c)$	0.082(0.104) ^{+0.037} _{-0.038}	0.074(0.070) ^{+0.03} _{-0.036}	0.118(0.122) ^{+0.029} _{-0.026}
100 θ_s	1.04147(1.0413) ^{+0.00036} _{-0.00035}	1.04153(1.04159) ^{+0.00035} _{-0.00032}	1.04157(1.04127) \pm 0.00034
$r_s(z_{\text{rec}})$	140.1(138.9) ^{+2.2} _{-2.0}	140.6(140.8) ^{+2.1} _{-1.6}	138.2(138.2) ^{+1.6} _{-1.8}
S_8	0.844(0.851) ^{+0.017} _{-0.018}	0.838(0.835) \pm 0.012	0.843(0.839) ^{+0.012} _{-0.013}
Ω_m	0.3084(0.3067) ^{+0.009} _{-0.0093}	0.3067(0.3071) ^{+0.0055} _{-0.0058}	0.3017(0.2962) ^{+0.0051} _{-0.0054}
$\Delta\chi_{\text{min}}^2$ (Λ CDM)	-5	-6	-18.5 (-0.5)

TABLE II. The mean (best-fit) $\pm 1\sigma$ error of the cosmological parameters reconstructed from the combined analysis of various data-sets (from left to right, each column adds a sets of data to the previous one). We also report the $\Delta\chi_{\text{min}}^2$ with respect to a Λ CDM fit to the same data-sets. In the last row, we also report the $\Delta\chi^2$ with respect to Λ CDM fit to the combined data without SH0ES in parenthesis.

IV. CONFRONTING EDE TO WEAK LENSING SURVEYS

In order to make use of weak gravitational lensing data to perform LSS analyses, one needs to accurately model the matter power spectrum in the late-time non-linear regime. To this purpose, one can adopt the HALOFIT semi-analytical prescription [56], as revised by [57], which has been shown to be accurate at 5% level in reproducing the non-linear power spectra of Λ CDM models up to wavenumbers $k \leq 10 h/\text{Mpc}$. However, the version developed by authors of Ref. [57] does not consider the impact of baryon feedback. A further improvement, dubbed as HMCODE, has been developed in Ref. [58], its main advantage being its flexibility to account for the effects of baryon physics on the small-scale clustering of matter, particularly important at very low redshifts. Both HALOFIT and HMCODE have been shown to be suitable to describe the Λ CDM scenario, as well as some common extensions beyond it, such as models with varying DE EoS or massive neutrinos [101]. In Section IV A we confront the non-linear matter power spectra produced by using HALOFIT/HMCODE in the EDE framework, against the outputs of dedicated cosmological N -Body simulations, to explicitly demonstrate the accuracy of our LSS data analyses. In Section IV B we discuss the results of our MCMC analysis against weak lensing data.

A. Non-linear matter power spectrum: a comparison with N -Body simulations

The goal of this Section is to show that the impact on the non-linear matter power spectrum do to the presence of EDE is mainly due to changes in the standard Λ CDM free parameters with respect to their reference values, and therefore such impact can be safely studied without further modifying or re-calibrating HALOFIT/HMCODE. To this end, we perform two sets of N -Body DM-only

simulations (one set for the EDE and one for the Λ CDM scenarios), as reported in Table III, by using the N -body code GADGET-3, a modified version of the publicly available numerical code GADGET-2 [102, 103]. The initial conditions have been produced by displacing the DM particles from a cubic Cartesian grid according to second-order Lagrangian Perturbation Theory, with the 2LPTIC public code [104], at redshift $z = 99$. The corresponding input linear matter power spectra, for both the EDE and Λ CDM cases, were computed with AXICLASS, the aforementioned modified version [71] of the publicly available code CLASS [55]. For all of the simulations, we kept the cosmological parameters fixed to their EDE best fit values from Ref. [71] (very close to ours), namely $H_0 = 72.81$, $\Omega_m = 0.2915$, $A_s = 2.191 \cdot 10^{-9}$, $n_s = 0.986$ for both cosmological scenarios; plus the additional parameters $\log_{10}(z_c) = 1.04106$, $f_{\text{EDE}}(z_c) = 0.132$, $\Theta_i = 2.72$, $n = 2.6$ for the EDE model. To bind together the matter power spectra extracted from simulations with different resolutions we adopt a splicing technique described in detail in Appendix B. Our results are summarized in Fig. 4, and we refer to Appendix B for a deeper technical discussion.

In the top panel of Fig. 4 we compare the matter power spectra extracted from our simulations, with the ones computed with HALOFIT/HMCODE, at redshift $z = 0$. The blue curves refer to the Λ CDM scenario – dubbed hereafter as Λ CDM “equivalent” – while the red ones refer to the EDE best fit model. As a reference, we also report the best fit Λ CDM case from Planck 2018. The spliced power spectra are denoted by thick dot-dashed lines. Symbols stand for the output power spectra of the “non-spliced” LB and HR simulations. The solid/dotted lines are the non-linear power spectra from HALOFIT/HMCODE, while the dashed lines are the corresponding linear power spectra used to set the initial conditions for the simulations. In the right panel, we adopt the same linestyle-code and color-code to show the ratio between the non-linear power spectra produced by

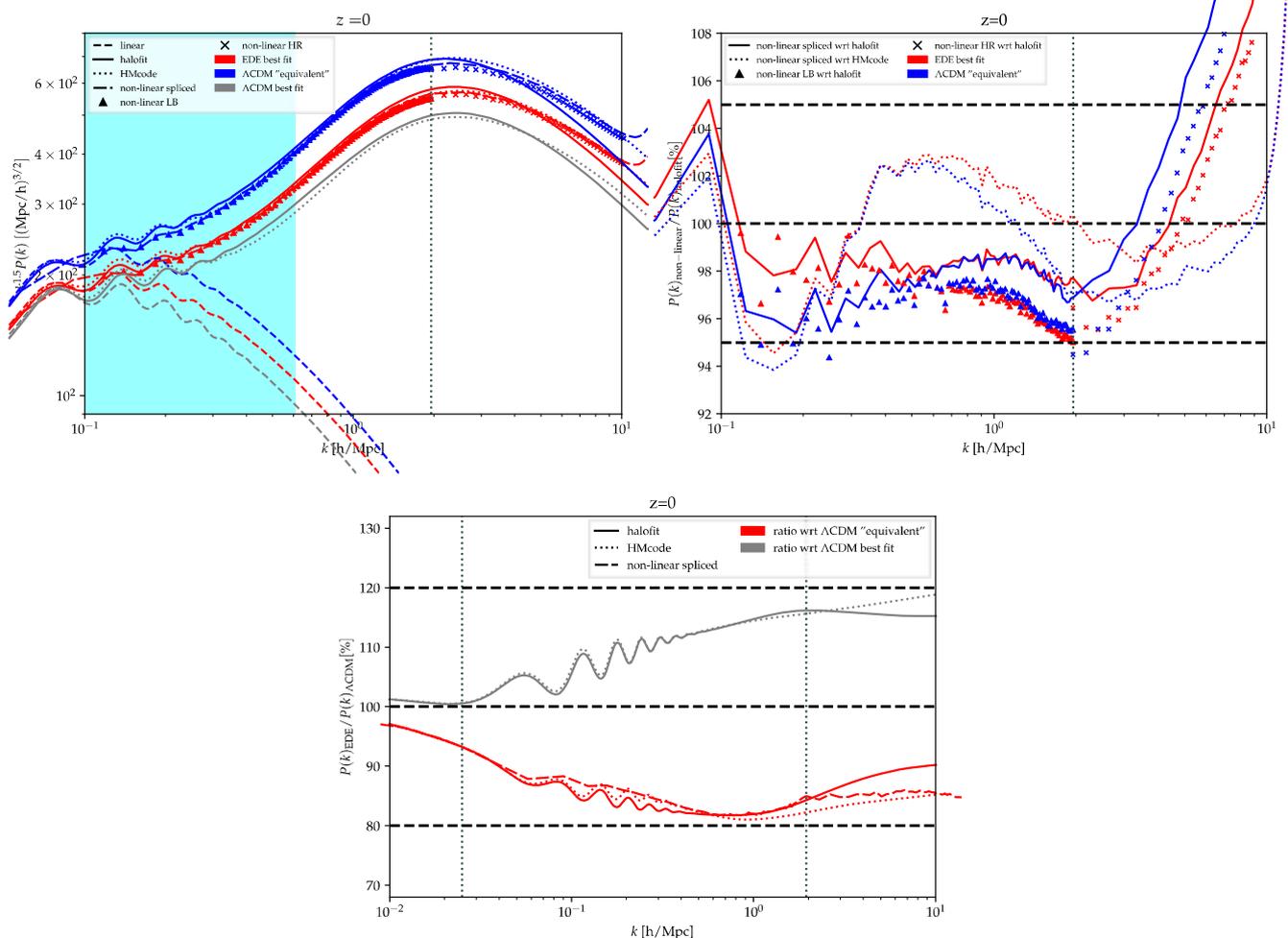


FIG. 4. In the top left panel we show the matter power spectra extracted from our simulations, and the ones computed with HALOFIT/HMCODE. The blue curves refer to the Λ CDM scenario, while the red ones refer to the EDE best fit model. We also report the best fit Λ CDM case from Planck 2018. The spliced power spectra are reported as thick dot-dashed lines. Symbols stand for the outputs of the LB and HR simulations. The solid/dotted lines are the non-linear power spectra from HALOFIT/HMCODE, whereas the dashed lines are the corresponding linear power spectra used to produce the initial conditions for the simulations. The cyan shaded band roughly corresponds to the scales probed by DES-Y1. In the top right panel, we show the ratio between the non-linear matter power spectra from our simulations and the ones computed with HALOFIT/HMCODE, for both the Λ CDM “equivalent” and the EDE best fit models, adopting the same linestyle-code and color-code. In the bottom right panel we compare departures from the Λ CDM model in terms of ratios of non-linear matter power spectra, adopting the same linestyle-code and color-code.

Model	Particles (N)	Box size (L)	Mass resolution	Label
Λ CDM/EDE	1024^3	$250 h^{-1}$ Mpc	$1.2 \cdot 10^9 h^{-1} M_{\odot}$	HR
Λ CDM/EDE	1024^3	$1000 h^{-1}$ Mpc	$7.5 \cdot 10^{10} h^{-1} M_{\odot}$	LB
Λ CDM/EDE	256^3	$250 h^{-1}$ Mpc	$7.5 \cdot 10^{10} h^{-1} M_{\odot}$	LR

TABLE III. Summary of the properties of the cosmological simulations used in this work. Notice that the Figures shown in this Section have been obtained by splicing together (for each redshift and model) the non-linear matter power spectra extracted from the first two simulations listed here, by using the third one to correct for finite-volume and resolution effects (see Appendix B for details). The labels listed in the last column stand for High Resolution, Large Box, and Low Resolution, respectively.

HALOFIT/HMCODE with respect to the ones extracted from our simulations. The thick horizontal lines highlight

$\pm 5\%$ deviations. In Appendix B we extend the analysis to three additional redshift bins – $z = 0.5, 1.5, 2$ – ob-

taining analogous results. We can thus conclude that the differences between HALOFIT/HMCODE predictions with respect to the outputs of our N -Body simulations are below 5% level, for scales $10^{-2} \lesssim k \lesssim 10 \text{ h/Mpc}$, at redshifts $0 \leq z \leq 2$, for both Λ CDM and EDE models. Whereas this is a very well established result for the Λ CDM paradigm, this is not often the case for alternative cosmological scenarios, such as the one considered in this work. Let us note that the exponential increase in the difference between the outputs by simulations and HALOFIT at $k \sim 10 \text{ h/Mpc}$ is absent when one compares the outputs from simulations with the predictions by HMCODE. As expected, the latter method is more accurate than HALOFIT in modelling the very small-scale and very low- z regime. In this work, we therefore make use of HMCODE to model the non-linear evolution of perturbations, following the approach adopted by the KiDS collaboration.

We also present our results in terms of ratios between the matter power spectra in the EDE and the Λ CDM models in Fig. 4 bottom panel. The comparison between the EDE best fit and the Λ CDM equivalent confirms that it is not the intrinsic presence of EDE that enhances the matter power spectrum on small scales, exacerbating the S_8 tension. Rather, the EDE reduces the growth of perturbations at fixed ω_{cdm} . As already pointed out, such power enhancement is instead due to variations in the standard Λ CDM parameters – mostly an increase in ω_{cdm} – induced to balance the EDE impact on the CMB. This suggests that the limitations of the EDE are not intrinsic to its presence, but rather to an accidental degeneracy that could be alleviated in an extended model. This will be the starting point of Section VB, where we will outline possible paths towards restoring the agreement with WL measurements in (extended) EDE cosmologies.

In view of these considerations, it is straightforward to conclude that LSS surveys constitute an ideal counterpart to CMB data, given the complementarity between the regimes that they probe. However, in section IV B we will show that currently available weak lensing data are not sensitive enough to unequivocally capture the signature of EDE. This will clearly not be the case when more precise data (e.g. from Euclid [105]) will become available. As also our results suggest, it will soon be necessary to go beyond the HALOFIT/HMCODE prescription for modelling the non-linear power spectrum (see e.g. [49, 50, 106, 107]). Furthermore, it might be already possible to test $\mathcal{O}(20\%)$ deviations in the small-scale power, as the ones shown in the bottom panel of Figure 4, with current Lyman- α forest flux power spectrum data [108–113] and the EFT of LSS data analysis of BOSS data (see e.g. Refs. [49, 50] for a recent analysis in the 3-parameter model). We leave these tasks for future work.

B. MCMC analysis against weak lensing data

In the following, we will focus on the 1-parameter EDE cosmology, fixing Θ_i and $\text{Log}_{10}(z_c)$ to their best fit values from PLANCK only. Firstly, we test the model against the KiDS-VIKING cosmic shear measurements. We follow the prescription described in Ref. [45] and make use of the HMCODE algorithm [58] (with 9 nuisance parameters) to model the non-linear matter power spectrum. Secondly, we perform an analysis trading KiDS-VIKING data for a split-normal likelihood on S_8 as inferred from the joint KiDS-VIKING+DES data using Complete Orthogonal Sets of E/B-Integrals (COSEBIs), namely¹⁴ $S_8 = 0.755^{+0.019}_{-0.021}$ [51]. We report results of MCMC analysis of Λ CDM and EDE against the KiDS-VIKING data and the joint KiDS-VIKING+DES data in table IV and V.

Results for Λ CDM: Starting with the Λ CDM cosmology, we find that combining PLANCK with KiDS-VIKING data leads to a mild degradation of the χ^2_{min} of the combined fit: While one might naively expect that the χ^2_{min} of the global fit should be roughly the sum of the χ^2_{min} of individual fits, we find that the global χ^2_{min} is degraded by $\sim +6.5$. Similarly, the inclusion of a tight Gaussian likelihood on S_8 as measured by KiDS-VIKING+DES leads to a degradation in the combined $\chi^2 \sim +15.5$, while one expects $\sim +1$ for a good fit. In Fig. 5, we show the reconstructed 2D posteriors of $\{H_0, S_8, 10^{-9}A_s, \omega_{\text{cdm}}, \Omega_m\}$ in the Λ CDM model. One can see that the degradation in χ^2_{min} is accompanied by shifts in the mean of any parameter correlated with S_8 , in particular A_s , ω_{cdm} and H_0 , without succeeding in getting a good fit to the WL data. We therefore stress that any of the combined results should be taken with a grain of salt, even in the Λ CDM framework. This joint analysis serves mostly to demonstrate that the EDE cosmology does not sensibly degrade the fit to the S_8 measurement as compared to Λ CDM, and that currently available WL measurements do not strongly constrain the EDE resolution to the Hubble tension.

Results for EDE against Planck+KiDS-VIKING: In Fig. 6 we show the reconstructed 2D posteriors of $\{f_{\text{EDE}}(z_c), H_0, S_8, \Omega_m\}$ in the 1-parameter EDE realization for various data combinations. We start by performing an analysis of EDE against KiDS-VIKING data only; as expected we find that the KiDS-VIKING data have no constraining power on the fraction of EDE. However, the reconstructed $S_8 = 0.738^{+0.041}_{-0.038}$ is $\sim 2.4\sigma$ discrepant with that obtained from previous analyses, suggesting a potential discordance between the cosmologies. For comparison, the prediction for S_8 in the Λ CDM model

¹⁴ We stress that S_8 is a model-dependent quantity, and it is particularly sensitive to the treatment of the neutrino mass. We therefore make use of the value that was derived following our convention, i.e. at fixed $\sum m_\nu = 0.06 \text{ eV}$.

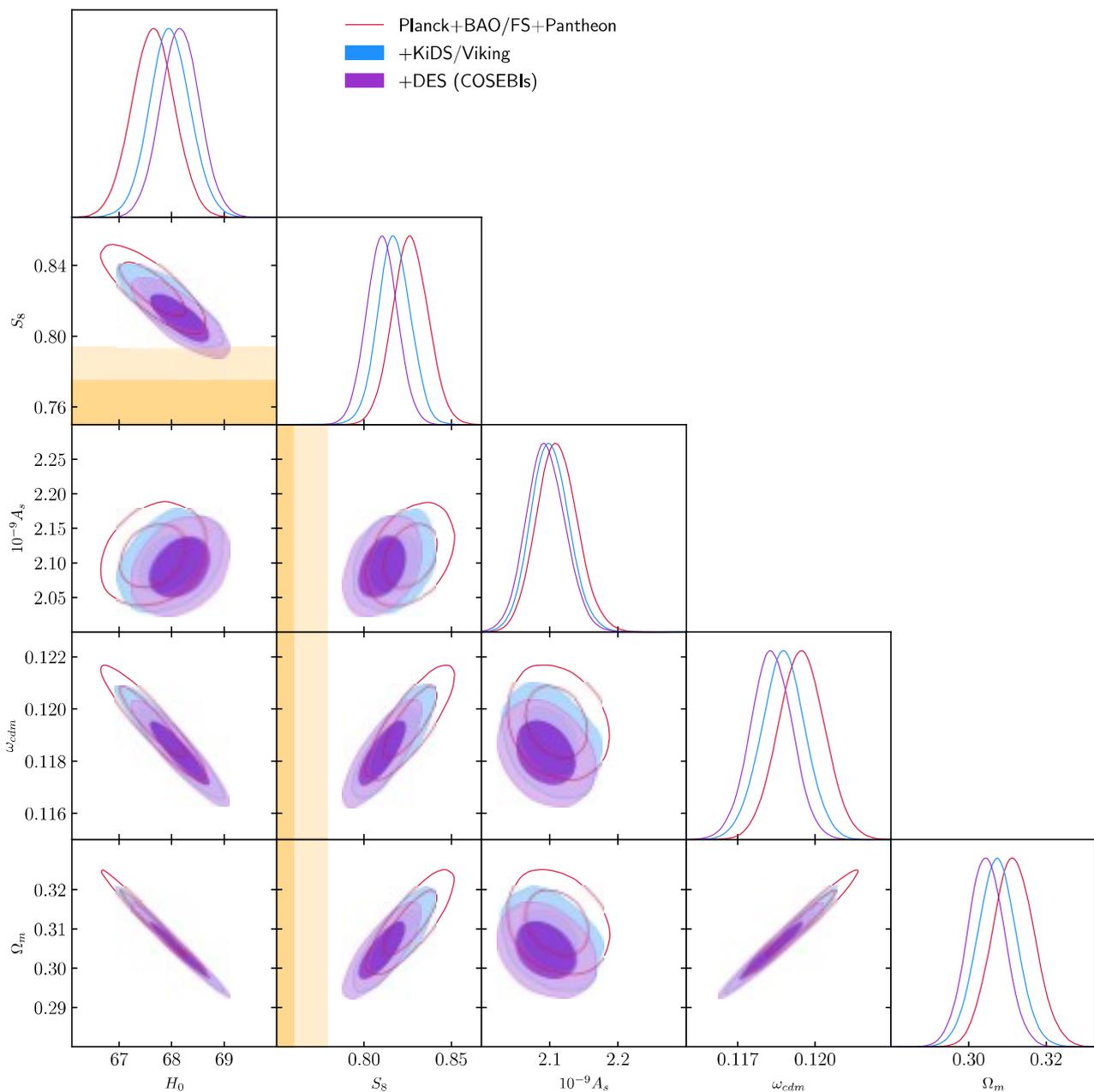


FIG. 5. Reconstructed 2D posterior distributions of a subset of parameters for various data set combinations (see legend) in the Λ CDM cosmology.

obtained from PLANCK data is 2.3σ discrepant with that from KIDS-VIKING data [45]. Therefore, although the mean value has increased, the level of the S_8 tension in the EDE cosmology is similar to that in Λ CDM because of larger error bars. Combining KIDS-VIKING to PLANCK+BAO+PANTHEON+FS, a non-zero EDE contribution is still favored at $\sim 1.5\sigma$, but the reconstructed mean fraction has moved downwards by $\sim 0.7\sigma$. This was expected, given the positive correlation between $f_{\text{EDE}}(z_c)$ and S_8 . In this cosmology,

PLANCK data are still slightly better fitted ($\Delta\chi^2_{\text{min}} \sim -6$) than in Λ CDM, while the fit to KIDS-VIKING data is degraded by $\sim +2$. Once a prior on H_0 from SH0ES is added, we find again $f_{\text{EDE}}(z_c) \sim 10 \pm 3\%$, at the cost of increasing the total $\chi^2_{\text{min}} \sim +5.5$. The increase in χ^2 is partly due to the inclusion of SH0ES ($\chi^2 \sim 1.62$, a reasonably good fit), and also to a mild degradation in the fit to PLANCK ($\sim +3$) and BAO ($\sim +1.6$). The reason is that the inclusion of KIDS-VIKING data reduces the degeneracy between $f_{\text{EDE}}(z_c)$ and the

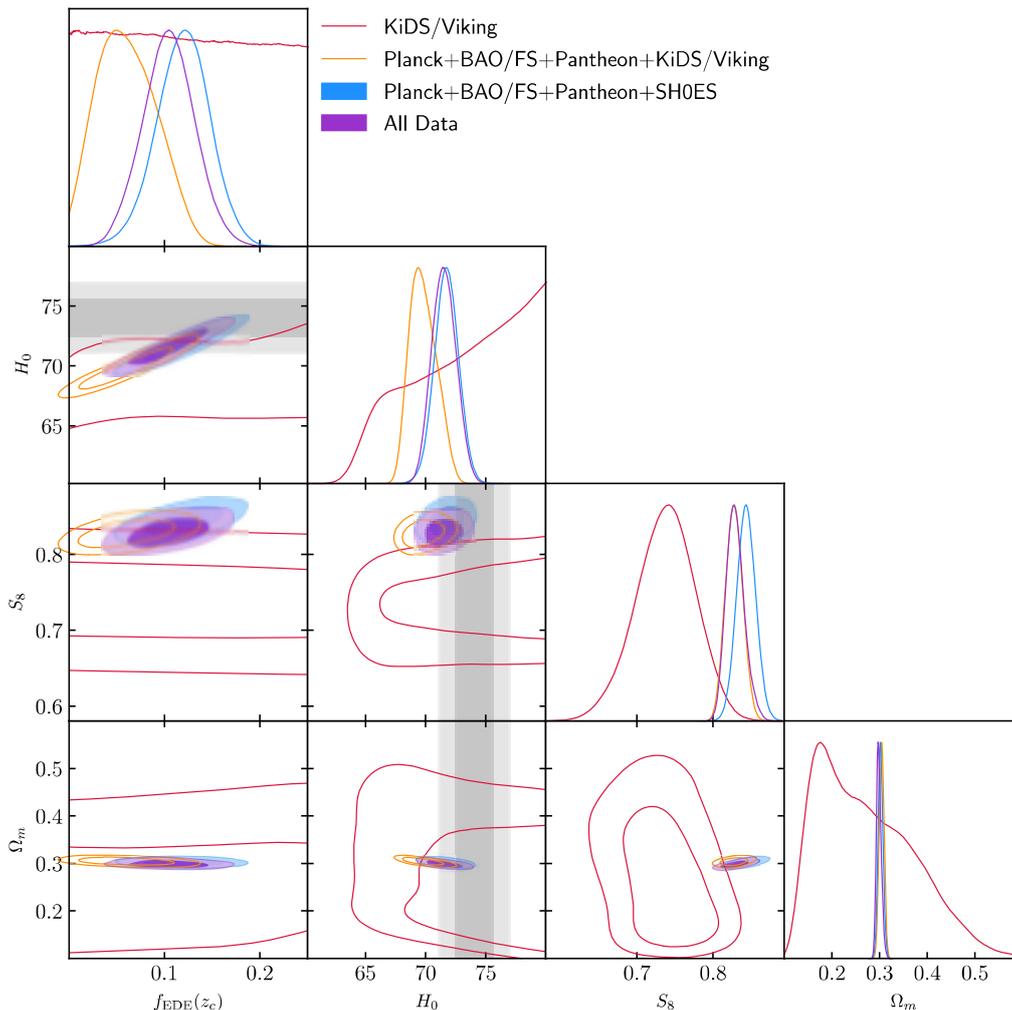


FIG. 6. Reconstructed 2D posterior distributions of a subset of parameters for various data set combinations (see legend) in the 1-parameter EDE cosmology.

Model	Λ CDM		EDE	
	Base+KiDS/VIKING	+SH0ES	Base+KiDS/VIKING	+SH0ES
H_0 [km/s/Mpc]	67.97 ± 0.38	68.4 ± 0.38	$69.75(68.95)^{+0.99}_{-1.1}$	$71.58(72.22)^{+1}_{-0.97}$
$100 \omega_b$	$2.248(2.248) \pm 0.013$	$2.257(2.256) \pm 0.013$	$2.261(2.253)^{+0.014}_{-0.015}$	$2.277(2.282)^{+0.013}_{-0.015}$
ω_{cdm}	$0.1187(0.1188) \pm 0.0009$	$0.1179(0.1180) \pm 0.0009$	$0.1245(0.1235)^{+0.0028}_{-0.0039}$	$0.1291(0.1310) \pm 0.0034$
$10^9 A_s$	$2.10(2.08) \pm 0.028$	$2.11(2.12)^{+0.03}_{-0.032}$	$2.117(2.116)^{+0.03}_{-0.033}$	$2.136(2.130)^{+0.03}_{-0.033}$
n_s	$0.9685(0.9667)^{+0.0038}_{-0.0036}$	$0.9708(0.9691)^{+0.0039}_{-0.0035}$	$0.9778(0.9740)^{+0.0061}_{-0.0068}$	$0.9872(0.9907)^{+0.0064}_{-0.0055}$
τ_{reio}	$0.0556(0.0520)^{+0.0069}_{-0.0066}$	$0.0585(0.0589)^{+0.0074}_{-0.0075}$	$0.0547(0.0559)^{+0.0067}_{-0.0074}$	$0.0552(0.0536)^{+0.0068}_{-0.0073}$
$f_{\text{EDE}}(z_c)$	—	—	$0.058(0.042)^{+0.028}_{-0.034}$	$0.104(0.122)^{+0.029}_{-0.025}$
$100 \theta_s$	$1.04198(1.04165) \pm 0.00028$	$1.04207(1.04210) \pm 0.00028$	$1.04165(1.04181)^{+0.00035}_{-0.00032}$	$1.04146(1.04130)^{+0.00031}_{-0.00034}$
S_8	$0.8172(0.8137)^{+0.009}_{-0.0096}$	$0.8092(0.8094)^{+0.0091}_{-0.0098}$	$0.826(0.831) \pm 0.011$	$0.829(0.828)^{+0.012}_{-0.011}$
Ω_m	$0.307(0.309) \pm 0.005$	$0.302(0.302) \pm 0.005$	$0.3037(0.3085)^{+0.0054}_{-0.0055}$	$0.2976(0.2962)^{+0.005}_{-0.0051}$
χ^2_{min}	3996.82	4011.16	3992.11	3997.67

TABLE IV. The mean (best-fit) $\pm 1\sigma$ error of the cosmological parameters reconstructed from the combined analysis of KIDS/Viking with other data. The ‘Base’ dataset refers to PLANCK+BAO/FS+PANTHEON. We also report the χ^2_{min} for each model and data set combination.

Λ CDM parameters, in particular the one with ω_{cdm} . Note that the goodness of PLANCK fit is not sensibly degraded as compared to Λ CDM, since the χ^2 stays better than that from Λ CDM fitted on PLANCK only. In fact, when compared to Λ CDM, the combined χ^2 is improved by ~ -13 (for one extra parameter), indicating a significant preference for EDE despite the presence of KiDS-VIKING data. Looking at the individual χ^2_{min} , we find indeed that the quality of the fit to KiDS-VIKING data in the EDE cosmology that resolves the Hubble tension is barely changed ($\Delta\chi^2 \sim +1.6$ for 195 data points [45]) compared to the Λ CDM fit to the same data set.

Results for EDE against Planck+KIDS+DES:

We now trade KiDS-VIKING data for a split-normal likelihood on S_8 as inferred from the joint KiDS-VIKING+DES data. We note that the tension between the value of this joint S_8 measurement and that predicted by our fiducial EDE model (obtained from the global fit of PLANCK+BAO+FS+PANTHEON+SH0ES) is at the $\sim 3.8\sigma$ level (slightly increased from 3.2σ tension in Λ CDM). It would be interesting to quantify the level of tension between these data sets in the Λ CDM framework using more robust statistical tools than the ‘difference in the mean’ used here, as done for instance in Refs. [100, 114]. However, we note that authors from Ref. [100] found that the less precise KiDS-VIKING data available at that time were already in significant statistical disagreement with the prediction from Λ CDM. We anticipate that this more robust approach would strengthen the case for a statistically significant discrepancy, even in Λ CDM, and therefore the need to apply caution when drawing conclusions from the combined analysis. We show the reconstructed 2D posteriors of $\{f_{\text{EDE}}(z_c), H_0, S_8, \Omega_m\}$ in the 1-parameter EDE model in Fig. 7. Without the SH0ES prior, $f_{\text{EDE}}(z_c)$ is compatible with 0 at 1σ , and we find an upper limit on $f_{\text{EDE}}(z_c) < 0.094$ at 95% C.L. This constraint is significantly weaker than that derived in Refs [48], despite the fact that we have *reduced* the parameter space. We have simply adopted a different ‘prior’ choice on Θ_i and $\text{Log}_{10}(z_c)$ (i.e. here we fix them), demonstrating that the current constraints from WL – besides being derived from statistically inconsistent data set – are not robust. Looking at χ^2_{min} , we find that the resulting best fit cosmology degrades the fit to Planck by $\sim +6$ while providing a poor fit to the S_8 likelihood ($\chi^2 = 8.3$ for a single data point). Still, the best-fit is marginally better than that of Λ CDM adjusted on the same sets of data ($\Delta\chi^2_{\text{min}} \sim -1.5$). Once we include the SH0ES prior, we find again $f_{\text{EDE}}(z_c)$ to be non-zero at more than 3σ , $f_{\text{EDE}}(z_c) \simeq 9 \pm 3\%$, with a global $\Delta\chi^2_{\text{min}} \simeq -11.6$. Looking at individual χ^2_{min} , we find that the fit to PLANCK^{TTTEEE}, PLANCK ^{$\phi\phi$} , BAO and FS data is somewhat degraded compared to the best fit EDE cosmology obtained without S_8 prior, as a consequence of the breaking of the $f_{\text{EDE}}(z_c) - \omega_{\text{cdm}}$ degeneracy. However, as expected, we note that the S_8 likelihood has

a $\chi^2 \simeq 9$, which is not particularly worst that the one obtained in the Λ CDM case without SH0ES ($\chi^2 \simeq 8.3$). This indicates that any constraint on the EDE derived from this combined analysis should be regarded with caution, as the cosmology reconstructed from the analysis does not provide a good fit to the S_8 data. This naturally impacts the reconstructed H_0 , which is $\sim 0.6\sigma$ lower than without the S_8 likelihood, although the fit to SH0ES is still reasonably good ($\chi^2 \simeq 2.4$). We therefore conclude that current S_8 measurements do not exclude the EDE resolution to the Hubble tension; however, they do call for new physics beyond EDE – or unknown systematics – to explain the intriguingly low measured S_8 values.

C. 1pEDE vs Λ CDM: a Bayesian view

In order to get a sense of whether 1pEDE is favored over Λ CDM, we perform a Bayesian analysis against 4 representative data combination, namely i) Base; ii) Base+SH0ES; iii) Base+KiDS/VIKING/DES; iv) Base+SH0ES+KiDS/VIKING/DES. We recall that our ‘Base’ dataset refers to PLANCK+BAO/FS+PANTHEON. We make use of the sampler MULTINEST [115], with 800 live points and a tolerance condition on the evidence for stopping the sampling equal to 0.1. We perform model comparison by calculating $\Delta\log B = \log B(\text{EDE}) - \log B(\Lambda\text{CDM})$ as reported in Table VI. As can be anticipated from the small $\Delta\chi^2$ values, we find that the 1pEDE is not favored in a Bayesian sense over Λ CDM, except in the case PLANCK+BAO/FS+PANTHEON+SH0ES, where the preference is ‘weak’ according to Jeffrey’s scale. We note that Ref. [37] reported ‘strong evidence’ in favor of EDE for the same data combination. We checked that the differences can be attributed to (in order of importance): i) the new Planck 2018 data; ii) a slightly less restrictive prior on f_{EDE} , taken to vary between $[0, 0.5]$ instead of $[0, 0.3]$; iii) the correct inclusion of BAO/FS data from BOSS DR12.

V. A COMMON RESOLUTION TO THE H_0 AND S_8 TENSIONS IN THE EDE COSMOLOGY?

A. EDE and the S_8 tension in light of Planck unlensed CMB spectrum

It has been noted that there exists a number of ‘curiosities’ in PLANCK that can potentially shed light on cosmological tensions. In particular, there is a residual oscillatory feature in the PLANCK TT data at $1100 \lesssim \ell \lesssim 2000$ compared to the best fit Λ CDM prediction [4, 62]. This feature can be captured by an extra source of smoothing of the acoustic peaks, as modelled by the ‘ $A_{\text{lens}}^{\phi\phi}$ ’ parameter which is used to re-scale the amplitude of the lensing potential power spectrum $C_l^{\phi\phi} \rightarrow A_{\text{lens}}^{\phi\phi} C_l^{\phi\phi}$, at every point in parameter space. However, the amplitude

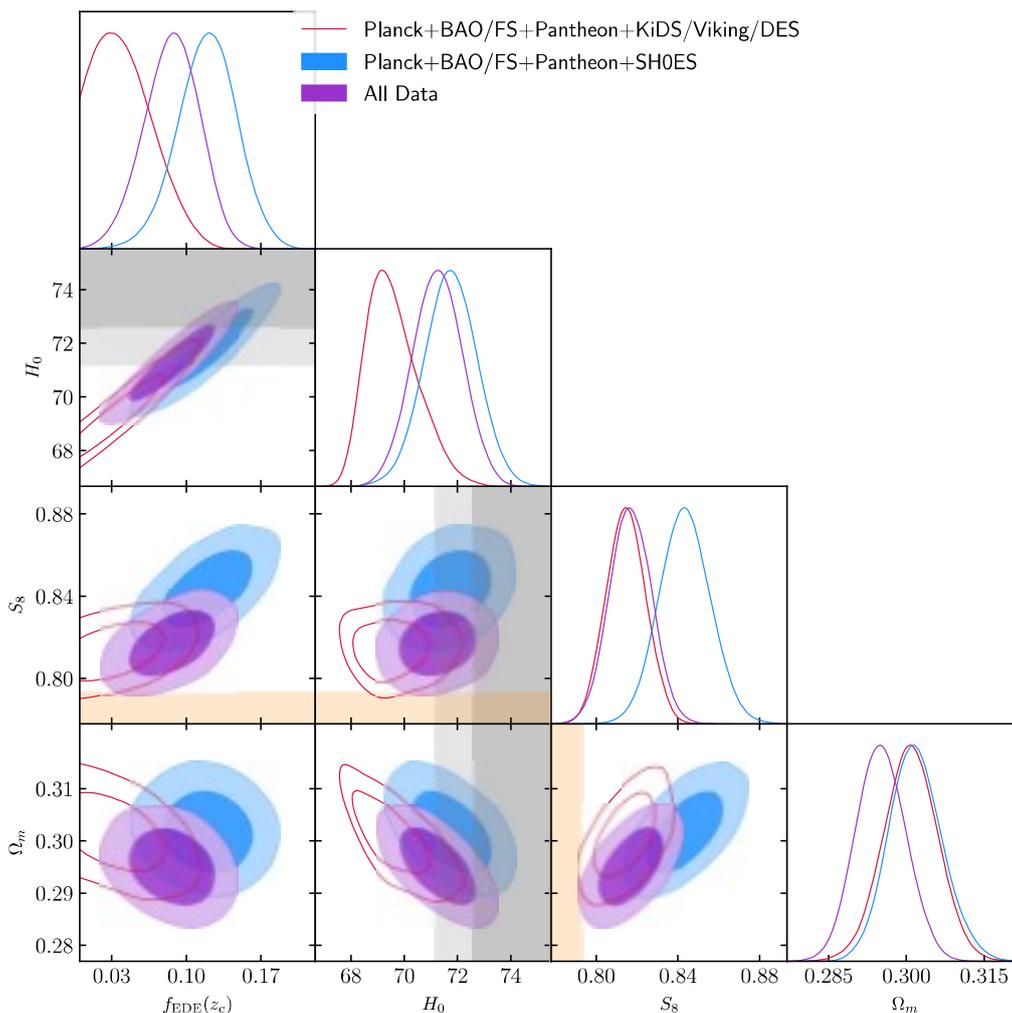


FIG. 7. Reconstructed 2D posterior distributions of a subset of parameters for various data set combinations (see legend) in the 1-parameter EDE cosmology.

Model	Λ CDM		EDE	
	Base+KiDS/VIKING/DES	+SH0ES	Base+KiDS/VIKING/DES	+SH0ES
H_0 [km/s/Mpc]	$68.16(68.15) \pm 0.38$	$68.56(68.69)^{+0.38}_{-0.39}$	$69.56(69.55)^{+0.72}_{-1.2}$	$71.29(71.81)^{+0.94}_{-0.94}$
$100 \omega_b$	$2.251(2.253) \pm 0.013$	$2.26(2.263)^{+0.013}_{-0.014}$	$2.262(2.270)^{+0.014}_{-0.015}$	$2.278(2.288) \pm 0.014$
ω_{cdm}	$0.1183(0.1183)^{+0.00084}_{-0.00082}$	$0.1175(0.1172)^{+0.00085}_{-0.00083}$	$0.1223(0.1199)^{+0.002}_{-0.0036}$	$0.1264(0.1270)^{+0.003}_{-0.0032}$
$10^9 A_s$	$2.094(2.091)^{+0.029}_{-0.03}$	$2.104(2.115)^{+0.029}_{-0.032}$	$3.046(2.107)^{+0.014}_{-0.015}$	$2.121(2.117) \pm 0.031$
n_s	$0.9691(0.9705) \pm 0.0037$	$0.9712(0.9731) \pm 0.0037$	$0.9765(0.9782)^{+0.0051}_{-0.0065}$	$0.9854(0.9892)^{+0.0055}_{-0.0057}$
τ_{reio}	$0.0546(0.0538)^{+0.0069}_{-0.0073}$	$0.0576(0.0602)^{+0.0069}_{-0.0077}$	$0.05339(0.0559)^{+0.0071}_{-0.0072}$	$0.05441(0.05254)^{+0.007}_{-0.0072}$
$f_{\text{EDE}}(z_c)$	—	—	$< 0.094(0.029)$	$0.087(0.097)^{+0.029}_{-0.024}$
$100 \theta_s$	$1.04198(1.04195)^{+0.00028}_{-0.00029}$	$1.04207(1.04209)^{+0.00028}_{-0.00029}$	$1.04178(1.04190)^{+0.00032}_{-0.00031}$	$1.04157(1.04149)^{+0.00033}_{-0.00032}$
S_8	$0.8043(0.8102)^{+0.0035}_{-0.0057}$	$0.8039(0.8023)^{+0.0036}_{-0.0058}$	$0.8145(0.8036)^{+0.0098}_{-0.01}$	$0.817(0.812)^{+0.01}_{-0.011}$
Ω_m	$0.3045(0.3046)^{+0.0048}_{-0.005}$	$0.2994(0.2978) \pm 0.0049$	$0.3008(0.2961)^{+0.0054}_{-0.0053}$	$0.2949(0.2919)^{+0.0047}_{-0.005}$
χ^2_{min}	3821.93	3837.98	3820.46	3826.35

TABLE V. The mean (best-fit) $\pm 1\sigma$ error of the cosmological parameters reconstructed from the combined analysis of the KIDS/Viking/DES data with other data discussed in the text, with and without a prior on H_0 from SH0ES. The ‘Base’ dataset refers to PLANCK+BAO/FS+PANTHEON. We also report the χ^2_{min} for each model and data set combination.

of the lensing potential power spectrum can also be estimated directly from the lensing-reconstruction and is

compatible with the Λ CDM expectation, such that while this extra smoothing looks like lensing, it cannot be at-

	Base	Base+SH0ES	Base+KiDS/VIKING/DES	Base+KiDS/VIKING/DES+SH0ES
$\Delta \log B$	+4	-0.7	+5.1	+2.2

TABLE VI. Difference in the Bayesian evidence $\Delta \log B = \log B(\text{EDE}) - \log B(\Lambda\text{CDM})$ between the 1pEDE and ΛCDM cosmology for the data combination as indicated by the row title. The ‘Base’ dataset refers to PLANCK+BAO/FS+PANTHEON.

tributed to actual gravitational lensing.

A thorough investigation of the lensing-like tensions in the Planck legacy release was performed in Refs. [62–64]. It has been noted in particular that, once marginalizing over the lensing information, the ‘unlensed’ CMB temperature and polarization power spectra favor a cosmology with a lower A_s and $\Omega_{\text{cdm}} h^2$. Indeed, these parameters are strongly correlated with the amplitude of the lensing power spectrum, such that the lensing-like anomaly pulls up these parameters. Additionally, since the acoustic feature of the CMB tightly constrains the parameter combination $\Omega_m h^3$, a lower $\Omega_{\text{cdm}} h^2$ is compensated by a higher H_0 . As a consequence, the unlensed ΛCDM cosmology shows no S_8 tension, and a milder (although still $> 3.5\sigma$ significant) H_0 tension. It was also pointed out that this ‘unlensed’ cosmology is in good agreement with the ΛCDM cosmology reconstructed from the SPTPol data [62, 65, 66].

It is therefore reasonable to ask what is the impact of such anomalies on extensions to ΛCDM like the EDE under study. To that end, we introduce two additional parameters A_{lens} and $A_{\text{lens}}^{\phi\phi}$ whose goal is to marginalize over the lensing information in Planck¹⁵. The latter parameter re-scales the amplitude of the theory lensing potential power spectrum, while the former only re-scales the amplitude of the acoustic peak smoothing. In practice, the amplitude of the acoustic peak smoothing is then determined by the product $A_{\text{lens}}^{\text{TTTTTEE}} \equiv A_{\text{lens}} \times A_{\text{lens}}^{\phi\phi}$. We first perform MCMC analysis of the ΛCDM and EDE cosmologies against PLANCK data. In a second step, we perform a global analysis combining all the data considered in this work. As before, the joint KiDS-VIKING+DES results is modelled via a split-normal likelihood distribution on S_8 . The results of these analysis are reported in Tab. VII and shown in Figs. 8 and 9.

Results for ΛCDM : We start by analyzing the ΛCDM cosmology in light of the ‘unlensed’ Planck spectra. We confirm the results of Refs. [4, 64]: we find that the amount of lensing determined from the peak smoothing $A_{\text{lens}}^{\text{TTTTTEE}}$ is $\sim 2.8\sigma$ higher than the expectation from the ΛCDM model deduced from the ‘unlensed’ CMB power spectrum. Moreover, the difference between the reconstructed $A_{\text{lens}}^{\phi\phi} \simeq 1.07 \pm 0.04$ and $A_{\text{lens}}^{\text{TTTTTEE}} \simeq 1.2 \pm 0.07$ illustrates the fact that this extra

smoothing component cannot be due to actual gravitational lensing. Still, this ΛCDM ‘unlensed’ cosmology is now in much better agreement with the S_8 measurements from KiDS and DES, as can be seen in Fig. 8. This is due to the fact that the reconstructed A_s and ω_{cdm} are lower than in the analysis including lensing information. We then perform a global analysis, including all data sets considered in this work. We find that the ‘unlensed’ ΛCDM cosmology can indeed accommodate a low S_8 , however this is at the cost of worsening somewhat the fit to BAO+FS data ($\Delta\chi^2 \simeq +3.5$), when compared to the ‘concordance’ ΛCDM model obtained from a fit to the full Planck data, BAO and FS (without S_8 priors). Additionally, we note that accommodating such a low S_8 requires a somewhat smaller ω_{cdm} and A_s (by a little less than 1σ), which are compensated for by pulling up the $A_{\text{lens}}^{\phi\phi}$ and $A_{\text{lens}}^{\text{TTTTTEE}}$ by a similar amount. The fit to SH0ES on the other hand is still very poor, $\chi_{\text{min}}^2 \simeq 10$, suggesting that the global unlensed cosmology is still in strong tension with SH0ES.

Results for EDE: Turning now to the 1-parameter EDE model, we wish to check whether the EDE cosmology deduced from ‘unlensed’ PLANCK spectra only is in better agreement with both S_8 and H_0 direct measurements¹⁶. As one can see from Fig. 9, the lensing-marginalized CMB data does favor non-zero $f_{\text{EDE}}(z_c)$ at $\sim 2\sigma$ ($\Delta\chi_{\text{min}}^2 \simeq -4$ with respect to ΛCDM) and predicts $H_0 \simeq 71.2 \pm 1.5$ and $S_8 \simeq 0.81 \pm 0.02$. Compared to the EDE cosmology reconstructed from the full Planck data, the ‘tension’ with H_0 and S_8 has therefore decreased by $\sim 1\sigma$ due to a shift in the mean of the reconstructed posterior in the unlensed cosmology. It is now in 1.4σ agreement with SH0ES but stays in mild ($\sim 2.7\sigma$) tension with the combined S_8 measurement. The S_8 prediction is however in very good agreement with the KiDS-VIKING and DES measurements when considered individually (an important note since the combined low S_8 value relies on a re-calibration of DES ‘photometric redshift by the KiDS-VIKING team). Additionally, the $A_{\text{lens}}^{\phi\phi}$ and $A_{\text{lens}}^{\text{TTTTTEE}}$ parameters are unchanged with respect to that reconstructed in the ΛCDM cosmology. Therefore, while the anomalous amount of lensing in PLANCK data is not an effect due to the presence of the EDE, these parameters do not correlate with a non-zero $f_{\text{EDE}}(z_c)$, i.e., they do not take values different from the ΛCDM ones to ‘hide’ the effect of the EDE.

¹⁵ An alternative, more thorough, way is to use CMB lensing principal components as introduced in Ref. [63, 64]. As we will show shortly, our reconstructed ‘unlensed’ cosmologies are in good agreement. Our approach follows that introduced in Refs. [116, 117].

¹⁶ A similar study was performed in Ref. [64] for N_{eff} . There, it was found that polarization and BAO data exclude N_{eff} as a resolution to the Hubble tension, even after marginalizing over the lensing anomaly in PLANCK.

Model	Λ CDM		EDE	
	PLANCK ^{TTTEEE} +PLANCK ^{$\phi\phi$}	All Data	PLANCK ^{TTTEEE} +PLANCK ^{$\phi\phi$}	All Data
H_0 [km/s/Mpc]	68.44(68.66) ^{+0.74} _{-0.72}	69.16(69.37) \pm 0.41	71.17(72.18) ^{+1.4} _{-1.6}	71.64(72.07) ^{+0.94} ₋₁
100 ω_b	2.262(2.269) \pm 0.018	2.277(2.284) \pm 0.014	2.284(2.292) \pm 0.02	2.292(2.294) \pm 0.015
ω_{cdm}	0.1179(0.1174) \pm 0.0016	0.1164(0.1159) \pm 0.00087	0.1253(0.1286) ^{+0.0037} _{-0.0045}	0.1248(0.1255) ^{+0.003} _{-0.0033}
10 ⁹ A_s	2.069(2.071) ^{+0.038} _{-0.035}	2.048(2.053) ^{+0.039} _{-0.032}	2.101(2.122) \pm 0.041	2.064(2.066) ^{+0.047} _{-0.033}
n_s	0.9718(0.9730) \pm 0.005	0.9755(0.9786) \pm 0.0037	0.9862(0.9925) ^{+0.008} _{-0.0088}	0.9884(0.9925) ^{+0.0059} _{-0.0056}
τ_{reio}	0.0494(0.0506) ^{+0.0089} _{-0.0079}	0.0464(0.0480) ^{+0.0093} _{-0.0074}	0.0507(0.0529) ^{+0.0087} _{-0.008}	0.0429(0.0431) ^{+0.012} _{-0.0071}
$A_{\text{lens}}^{\phi\phi}$	1.071(1.075) ^{+0.04} _{-0.043}	1.104(1.110) ^{+0.034} _{-0.038}	1.064(1.056) ^{+0.04} _{-0.043}	1.093(1.099) ^{+0.035} _{-0.039}
$A_{\text{lens}}^{\text{TTTEEE}}$	1.195(1.208) ^{+0.066} _{-0.07}	1.247(1.266) ^{+0.06} _{-0.066}	1.187(1.188) ^{+0.065} _{-0.07}	1.222(1.238) ^{+0.061} _{-0.067}
$f_{\text{EDE}}(z_c)$	—	—	0.078(0.108) ^{+0.035} _{-0.038}	0.082(0.092) \pm 0.027
100 θ_s	1.04205(1.04207) \pm 0.00031	1.04215(1.04214) \pm 0.00029	1.04165(1.04343) ^{+0.00036} _{-0.00035}	1.04165(1.04164) \pm 0.00034
S_8	0.800(0.795) ^{+0.019} _{-0.02}	0.780(0.776) \pm 0.011	0.801(0.812) \pm 0.02	0.794(0.793) \pm 0.013
Ω_m	0.302(0.297) ^{+0.009} _{-0.01}	0.2924(0.2883) ^{+0.0049} _{-0.0051}	0.2938(0.2870) ^{+0.0095} _{-0.01}	0.2891(0.2870) \pm 0.0052
χ_{min}^2 (Λ CDM)	2765.98	3816.23	2761.98	3808.40

TABLE VII. The mean (best-fit) $\pm 1\sigma$ error of the cosmological parameters reconstructed from the lensing-marginalized PLANCK data only and in combination with BAO/FS+PANTHEON+KiDS-VIKING-DES. We also report the χ_{min}^2 for each model and data set combination.

Once all data sets are included in the analysis, a non-zero fraction of EDE is favored at $\sim 3.5\sigma$. Interestingly, most of the reconstructed parameters do not shift by more than $\sim 0.5\sigma$; rather, the uncertainty on the reconstructed parameters tighten significantly, as one would expect from making use of additional data. However, similarly to what happens in the Λ CDM cosmology, the inclusion of the tight-and-low S_8 value does force a slightly ($\sim 1\sigma$) smaller A_s , that is compensated by slightly higher $A_{\text{lens}}^{\phi\phi}$ and $A_{\text{lens}}^{\text{TTTEEE}}$ parameters. The fit to SH0ES is good ($\chi^2 \sim 1.9$) and stable when compared to that obtained including Planck lensing information. On the other hand, as expected, the fit to the joint KiDS-VIKING+DES S_8 is better than in the ‘lensed’ cosmology ($\Delta\chi^2 \sim -4.3$), but its value is still somewhat poor ($\chi^2 \sim 4$). We emphasize again that the fit to individual S_8 measurements, on the other hand, is excellent. If future S_8 measurements stay low while becoming more precise, they will be in tension even with the ‘unlensed’ cosmology (whether Λ CDM or EDE), confirming the need for new physics beyond EDE (or an unknown systematic effect).

B. Extended cosmologies that could help resolving the S_8 tension

Given that both the Λ CDM cosmology and the EDE cosmology are in tension with S_8 measurements, it is reasonable to ask whether their could exist additional extensions that would help in accommodating the low S_8 value. Naturally one can argue that ‘Occam’s razor’ should prevail, and that the ‘true’ solution should be able to resolve both tension simultaneously. This might very well be the case, and it is without question that extending the parameter space until data fits is not a reasonable attitude. However, within a phenomenological framework, what is reasonable is to understand what aspects of a suggested resolution (EDE here) makes it at odds with

a certain data set (S_8 here). This question is especially interesting given that the fiducial Λ CDM cosmology is equally at odds with these data. In other words, this approach does not consist in ‘hiding’ bad effects of the EDE that compromises it with respect to Λ CDM in light of S_8 measurements. Rather, it consists in trying to understand what it would take (within a reasonable set of extensions) to accommodate current S_8 measurements in an EDE cosmology. If this can be achieved, the hope is then that it will lead to a set of predictions to be tested in the future, together with guidelines for model building which should make the extension less ‘ad-hoc’.

As was discussed extensively in this paper (see sec. IV A), it is interesting to note that at fixed ω_{cdm} , the EDE leads to a *decrease* in power at small scales that goes in the right direction to resolve the S_8 tension. However, the problem of EDE cosmologies is that they exploit a degeneracy with ω_{cdm} to counteract the effect of the EDE on the gravitational potential wells as seen in the CMB. Taken at face value, this poses both a concrete experimental problem –it is at odds with S_8 measurements– and a theoretical ‘tuning’ issue –why should these two apparently unrelated sectors ‘conspire’ to hide the EDE in CMB data? This logic should also be applied when considering sensible extension to a model. One of the less ‘theoretically costly’ possible explanation of the S_8 values is to invoke the fact that neutrinos are massive, and lead to a power suppression at small scales which decreases the value of σ_8 . Unfortunately it seems as though in practice the required sum of neutrino masses $\sum m_\nu \sim 0.3\text{eV}$ is excluded by PLANCK data. Moreover, neither DES nor KiDS-VIKING seems to have a preference for non-zero $\sum m_\nu$. However, it is interesting to note that constraints on the sum of neutrino masses can be strongly relaxed in extended cosmologies (e.g. [4, 118, 119]). In fact, in models attempting at resolving the Hubble tension with strongly interacting neutrinos [38, 120–122], it has been noted that Planck temperature 2015 data are in good agreement with $H_0 \simeq 72$ km/s/Mpc and neutrino masses $\sum m_\nu \sim 0.4\text{eV}$, but polarization data seems to

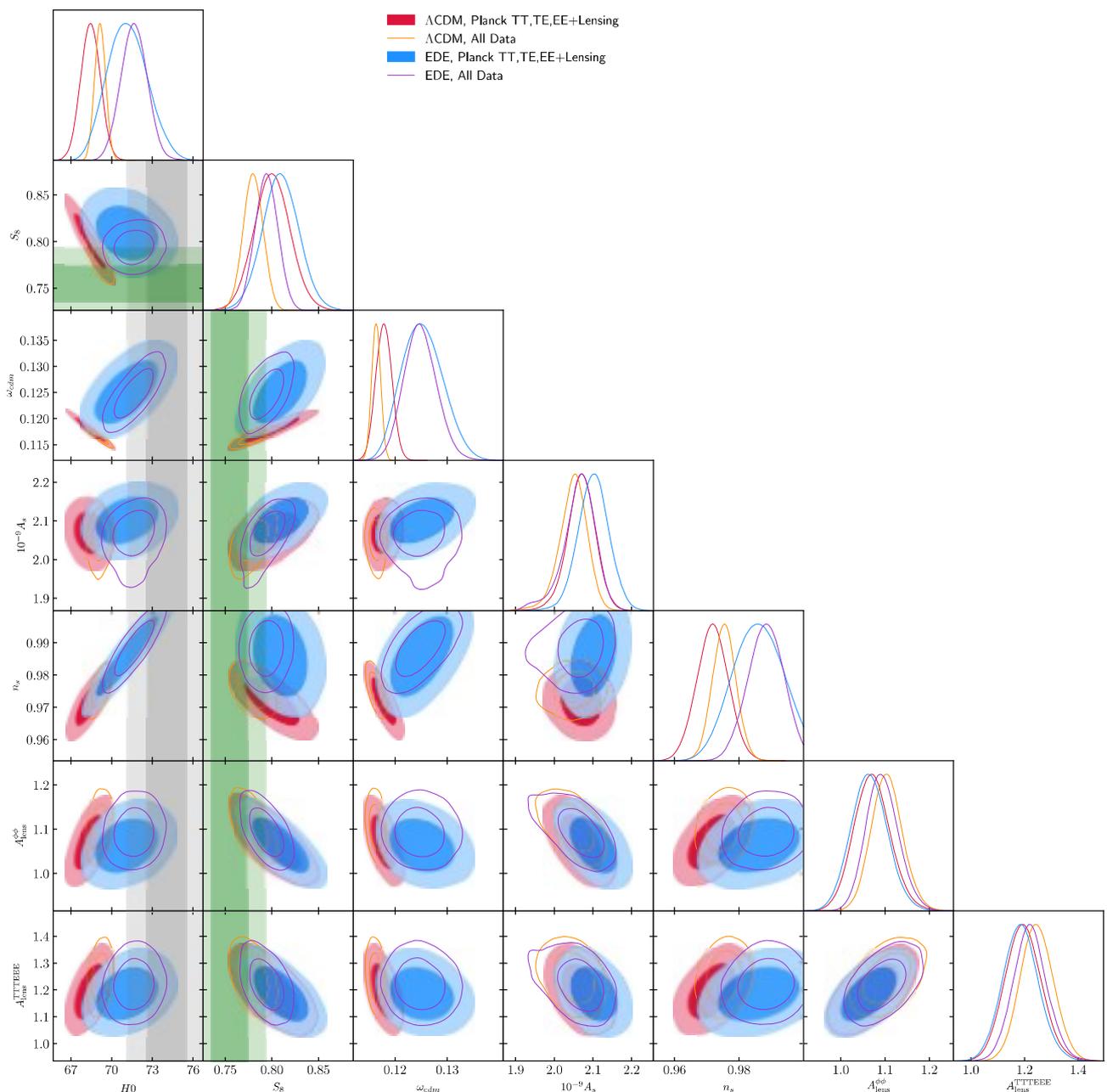


FIG. 8. Reconstructed 1- and 2D-posteriors of a subset of parameters in the Λ CDM and 1-parameter EDE cosmology for various data sets (see legend), once marginalizing over A_{lens} and $A_{lens}^{\phi\phi}$

restrict this resolution. In the EDE context, we have already mentioned that in Ref. [42], it was suggested that the non-relativistic transition of neutrinos could trigger a phase-transition in the EDE. It would therefore be very interesting in the future to study further the possible connection between EDE and neutrino masses. Along this idea, we have performed a MCMC run in the 1-parameter EDE model against all data-sets with the sum of neutrino masses let free to vary. We find that marginalizing over the neutrino mass does not affect the result presented

here. The S_8 value shifts downward only by $\sim 0.3\sigma$. This is because, while σ_8 does get suppressed, a non-zero $\sum m_\nu$ increases Ω_m , resulting only in a mild decrease in S_8 .

Another promising category of solutions to the S_8 tension invokes interaction between DM and an additional dark radiation (DR) component [123–125]. These models are particularly interesting because, by themselves and similarly to EDE, an additional radiation component leads to an increase in S_8 . This is due to the fact that

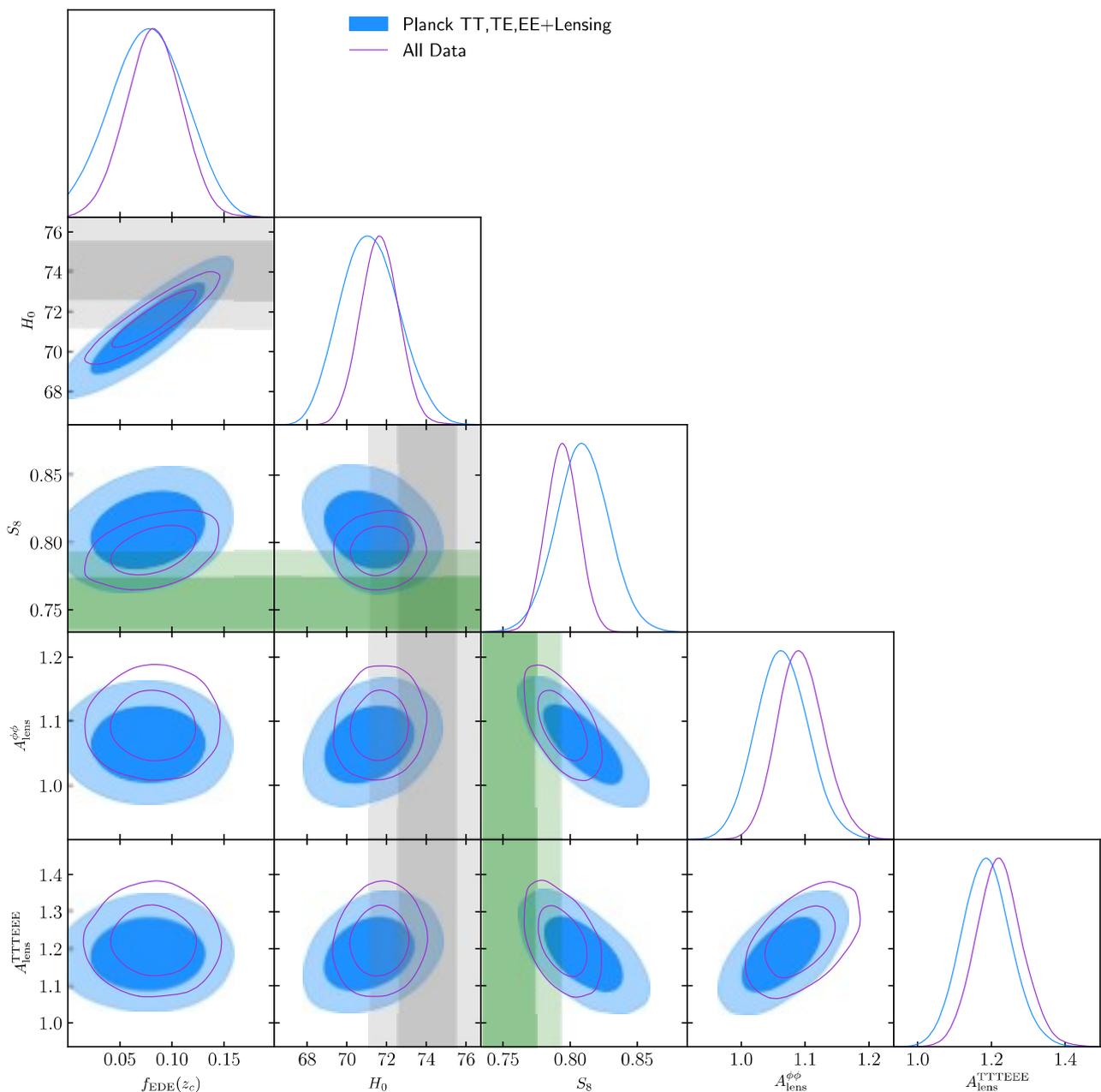


FIG. 9. Reconstructed 1- and 2D-posterior of a subset of parameters in the 1-parameter EDE cosmology for various data sets (see legend), once marginalizing over A_{lens} and A_{lens}^{ϕ} .

increasing the radiation density requires a simultaneous increase in the matter density to avoid a shift in matter radiation equality. As a consequence, Ω_m is higher in this model, leading to a higher S_8 . However, the introduction of an interaction between DM and DR leads to a power-suppression at small-scales and therefore to a smaller σ_8 , resulting in a net decrease in S_8 . Alternatively, there exists a number of model leading to interactions between DM and DE at late-times. Introducing an interaction between DM and EDE is a straight-forward extension of the

naive EDE model studied here. In fact, an axion EDE model whose dynamics is dictated by an interaction with dark gauge bosons was recently proposed in Ref. [74]. Interestingly, the parameter space is not enlarged in this model- rather, the critical redshift z_c at which the field starts to move is dictated by the ratio of the interaction rate over the Hubble rate. It will be interesting to generalize the model studied there (i.e., consider different type of interactions) and include linear cosmological perturbations, to check whether the presence of the ad-

ditional interaction could open a new degeneracy direction (alternative to the $f_{\text{ede}} - \omega_{\text{cdm}}$ one), which would prevent an increase in S_8 . Speculating further, if the presence of EDE is confirmed in the future, it is likely that it is connected to the existence of DE today, and perhaps even inflation. As a matter of fact, EDE models were introduced over a decade ago to alleviate the cosmological coincidence problem – the fact that the dark energy density and the matter density are very close from one another just today [81, 92]. Therefore, it is quite natural to ask whether the current epoch of accelerated expansion (and inflation) could be due to a dynamical scalar-field similar to the EDE, if more eras of such type can have occurred at other moments in the history of the Universe, and what would be the impact on the cosmic structure growth. In fact, in Ref. [101], it was shown that a time-evolving equation of state for DE is favored over Λ CDM from a combination of PLANCK, KiDS-450 and SH0ES data. However, these simple solutions are severely constrained by BAO and PANTHEON data. In future work, it will be interesting to check whether these constraints can be affected by the presence of EDE, and whether a more complete picture for early and late dark energy can help restoring cosmological concordance.

VI. CONCLUSIONS

In this paper, we have reassessed the viability of the EDE against a host of high- and low-redshift measurements, by combining LSS observations from recent weak lensing surveys KiDS-VIKING and DES with PLANCK 2018 CMB data, BOSS-DR12 BAO and growth function measurements, and the PANTHEON compilation of luminosity distance to SNIa. Our results can be summarized as follows:

- Within a phenomenological 3-parameters EDE model, we confirm that PLANCK+BAO+FS+PANTHEON+SH0ES favor $f_{\text{EDE}}(z_c) \simeq 0.1 \pm 0.03$, $z_c \simeq 4000_{-500}^{+1400}$ and $\Theta_i = 2.6_{-0.03}^{+0.4}$, with a $\Delta\chi^2 = -18.7$ compared to Λ CDM fitted on the same data set (i.e. a $\sim 3.6\sigma$ preference over Λ CDM¹⁷). The inclusion of the latest PLANCK data (and in particular the more precise polarization measurements) does not spoil the success of the EDE resolution to the Hubble tension. When compared to the ‘concordance’ Λ CDM model (i.e. obtained from analysis without SH0ES data), the EDE cosmology fits PLANCK+BAO+FS+PANTHEON equally well, but can additionally accommodate the high local H_0 values.

- Following the approach of Ref. [80], we have then shown that *reducing* the parameter space to a 1-parameter EDE model by fixing $\text{Log}_{10}(z_c)$ and Θ_i to their best fit values as obtained from a PLANCK data only analysis – which strikingly coincide with those from the combined analysis with SH0ES – leads to $\sim 2\sigma$ preference for non-zero EDE, namely $f_{\text{EDE}}(z_c) \simeq 0.08 \pm 0.04$ from PLANCK CMB data alone. In this cosmology, the inferred $H_0 \simeq 70 \pm 1.5$ km/s/Mpc is in agreement at better than 2σ with its local measurement from SH0ES. The addition of BAO, FS and PANTHEON data has no significant impact on the result. Including a prior on H_0 from SH0ES pulls up the reconstructed fraction to the $\sim 10\%$ level, with $H_0 \simeq 71.7 \pm 1$, while the fit to PLANCK is slightly better than in the concordance Λ CDM cosmology ($\Delta\chi^2 \sim -5$).
- To justify the inclusion of LSS data in our analyses, we have confronted the EDE non-linear matter power spectrum as predicted by standard semi-analytical algorithms against a dedicated set of N -body simulations. We have then tested the 1-parameter EDE cosmology against WL data, finding that it does not significantly worsen the fit to the S_8 measurements as compared to Λ CDM, and that current WL observations do not exclude the EDE resolution to the Hubble tension.
- We also caution against the interpretation of constraints obtained from combining PLANCK with KiDS-VIKING+DES. As we showed, the ‘compromise’ cosmology that is obtained is a poor fit to KiDS-VIKING+DES and degrades the fit to PLANCK data, even in Λ CDM. This illustrates that these data sets are statistically inconsistent in a Λ CDM framework, and it is easily conceivable that the resolution of this tension lies elsewhere (whether systematic effect or new physics).
- In light of the CMB lensing anomaly, we have shown that the lensing-marginalized CMB data favor non-zero EDE at $\sim 2\sigma$, predicts H_0 in 1.4σ agreement with SH0ES and S_8 in 1.5σ and 0.8σ agreement with KiDS-VIKING and DES, respectively. There still exists however a $\sim 2.5\sigma$ tension with the joint results from KiDS-VIKING and DES. Moreover, the presence of EDE does not affect the amount of anomalous lensing. This suggests that the anomalous lensing is not due to the presence of EDE, but also that the success of EDE is not due to opening up a new degeneracy direction with some exotic lensing parameters.
- With an eye on Occam’s razor, we finally discussed extensions of the EDE cosmology that could allow to accommodate the low S_8 values. In particular, we argue that EDE models which are coupled to

¹⁷ We assume Gaussian posteriors with 3 additional parameters for simplicity.

neutrinos; include interaction with an extra dark radiation bath or dark matter; or are connected to dynamical dark energy at late time (and perhaps inflation), are all worth exploring in future work as promising ways to fully restore cosmological concordance.

In another study [52], we confronted BOSS data to the 1-param EDE cosmology within the EFT of LSS framework [49, 50], finding that the constraints on f_{EDE} largely weaken and that LSS current data do not exclude the EDE resolution to the Hubble tension. An important follow-up to these studies will be to see whether the new ACT data [15], compatible with PLANCK (although see Ref. [126]), support – or restrict – the EDE resolution to the Hubble tension. Looking forward, future CMB experiment (such as Simons Observatory [127] and CMB-S4 [128]) and LSS data (from Euclid [105], LSST [129], JWST and DESI [130]) will be crucial in testing prediction of the EDE cosmology (and its potential extensions) [71, 107] and firmly confirm – or exclude – the presence of EDE.

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Appendix A: χ^2 tables

We report all χ_{min}^2 ’s obtained with the `MINUIT` algorithm [98] through the `MINUIT` python package for the various model and data-set combination considered in this work.

Appendix B: a closer look to N -Body simulations

The main systematical uncertainties in numerical simulations come from their limited box size and resolution,

as it has been thoroughly discussed in past literature (see, e.g., [131–134]). In order to minimize the missing large-scale modes, potentially affecting small-box simulations, and to overcome the impossibility of capturing the very non-linear scales in our large-box simulations, we adopted a splicing technique to bind together the matter power spectra extracted from simulations with different resolutions, for each redshift and model, as in Refs. [135, 136].

For both the Λ CDM model and the EDE best fit model from [71], we indeed performed: i) one Large Box (LB hereafter) simulation with $N = 1024^3$ DM particles and box size $L = 250 h^{-1}$ Mpc; ii) one High Resolution (HR hereafter) simulation with $N = 1024^3$ DM particles and box size $L = 1000 h^{-1}$ Mpc; iii) one Low Resolution (LR hereafter) simulation with the same box size of the HR one and the same resolution of the LB, namely $N = 256^3$ and $L = 250 h^{-1}$ Mpc, to be used as a transition simulation. The spliced non-linear matter power spectrum $P(k)$ is given by Ref. [136]

$$P(k) = \begin{cases} P_{\text{LB}}(k) \cdot \frac{P_{\text{HR}}(k_{\text{MIN}}^{250})}{P_{\text{LR}}(k_{\text{MIN}}^{250})}, & \text{if } k \leq k_{\text{MIN}}^{250} \\ P_{\text{LB}}(k) \cdot \frac{P_{\text{HR}}(k)}{P_{\text{LR}}(k)}, & \text{if } k_{\text{MIN}}^{250} < k < \frac{1}{2} k_{\text{Nyq}}^{\text{LB}} \\ P_{\text{HR}}(k) \cdot \frac{P_{\text{LB}}(0.5 \cdot k_{\text{Nyq}}^{\text{LB}})}{P_{\text{LR}}(0.5 \cdot k_{\text{Nyq}}^{\text{LB}})}, & \text{if } k \geq 0.5 \cdot k_{\text{Nyq}}^{\text{LB}} \end{cases} \quad (\text{B1})$$

where k_{MIN}^{250} is the minimum k -mode in our small-box simulations (HR and LR), while $k_{\text{Nyq}}^{\text{LB}}$ is the Nyquist wave-number of the LB one.

Besides the aforementioned systematical uncertainties, numerical simulations are also affected by two primary sources of statistical errors: the cosmic variance, affecting the large-scale part of the spectra, and the shot noise due to the discreteness of the DM particles, thereby affecting the smallest scales.

Concerning the shot noise term, its contribution to the power spectrum is simply given by $P_{\text{SN}} = (L/N)^3$. It is straightforward to see that it is largely subdominant at the scales and redshifts considered in this work, from Fig. 10, where we compare the matter power spectra extracted from our simulations with the ones computed with `HALOFIT/HMCODE`, in three different redshift bins from $z = 1.5$ to $z = 0.5$ – given that we have already discussed the $z = 0$ case in Section IV A. In Fig. 11 we plot the ratio between the power spectrum predicted by `HALOFIT` or `HMCODE` and that extracted from the numerical simulation in order to explicitly demonstrate that the differences are below 5% level, for scales $10^{-2} \lesssim k \lesssim 10 h/\text{Mpc}$, at redshifts $0.5 \leq z \leq 2$, for both Λ CDM and EDE models. This extends the $z = 0$ result presented in the main text to cover the full redshift range from `KIDS-VIKING`.

It is also informative to compare the prediction from algorithms with N -body at larger scales than that depicted in Figs. 10 and 11. Indeed, these are affected by higher statistical noise, due to cosmic variance, as one might already guess from the lower- k part of both figures. To

Λ CDM cosmology									
PLANCK high- ℓ TT,TE,EE	2347.86	2351.81	2347.02	2349.78	–	2351.53	2349.68	2351.24	2352.88
PLANCK low- ℓ EE	396.03	395.8	399.60	395.71	–	395.78	396.94	395.88	397.21
PLANCK low- ℓ TT	23.18	22.25	22.74	22.62	–	22.99	22.83	22.34	22.09
PLANCK lensing	–	–	8.65	9.56	–	9.49	9.05	9.93	10.05
PANTHEON	–	–	1026.83	1026.82	–	1026.84	1026.69	1026.72	1026.67
BAO FS BOSS DR12	–	–	6.25	6.11	–	6.23	5.88	5.86	6.18
BAO BOSS low- z	–	–	1.38	1.39	–	1.34	1.82	1.65	2.22
SH0ES	–	16.57	–	18.57	–	–	16.05	–	14.23
KIDS/VIKING	–	–	–	–	177.9	182.62	182.21	–	–
COSEBI	–	–	–	–	–	–	–	8.30	6.44
total	2767.07	2786.43	3812.47	3830.57	177.9	3996.82	4011.16	3821.93	3837.98

TABLE VIII. Best-fit χ^2 per experiment (and total) in the Λ CDM model.

3-parameter EDE cosmology					
PLANCK high- ℓ TT,TE,EE	2343.07	2350.24	2349.30	2347.73	–
PLANCK low- ℓ EE	397.47	396.20	398.19	395.88	–
PLANCK low- ℓ TT	21.54	20.80	20.56	21.09	–
PLANCK lensing	–	–	10.12	9.85	–
PANTHEON	–	–	1026.72	1026.68	–
BAO BOSS DR12	–	–	3.46	–	–
BAO BOSS low- z	–	–	2.06	1.81	–
BAO/FS BOSS DR12	–	–	–	6.73	–
SH0ES	–	0.47	1.38	2.13	–
total	2762.08	2767.72	2786.43	3811.89	–

TABLE IX. Best-fit χ^2 per experiment (and total) in the 3-parameter EDE model.

1-parameter EDE cosmology									
PLANCK high- ℓ TT,TE,EE	2345.02	2347.63	2344.98	2347.42	–	2345.16	2349.15	2350.22	2349.82
PLANCK low- ℓ EE	395.80	395.97	395.82	395.90	–	396.33	395.88	396.10	395.79
PLANCK low- ℓ TT	21.49	20.82	21.89	20.85	–	22.38	20.97	21.54	20.84
PLANCK lensing	–	–	9.39	10.00	–	9.07	10.04	10.22	10.91
PANTHEON	–	–	1026.80	1026.69	–	1026.84	1026.7	1026.69	1026.80
BAO FS BOSS DR12	–	–	6.44	7.18	–	6.43	7.08	6.47	7.70
BAO BOSS low- z	–	–	1.41	2.33	–	1.33	2.33	2.38	2.83
SH0ES	–	1.07	–	1.64	–	–	1.62	–	2.43
KIDS/VIKING	–	–	–	–	178.0	184.57	183.88	–	–
COSEBI	–	–	–	–	–	–	–	6.83	9.22
total	2762.31	2765.49	3806.74	3812.01	178.0	3992.11	3997.67	3820.46	3826.35

TABLE X. Best-fit χ^2 per experiment (and total) in the 1-parameter EDE model.

Model	Λ CDM		EDE cosmology	
PLANCK high- ℓ TT,TE,EE	2339.92	2340.96	2335.71	2336.12
PLANCK low- ℓ EE	395.67	395.87	395.80	397.01
PLANCK low- ℓ TT	21.93	21.03	20.65	20.44
PLANCK lensing	8.47	8.35	9.82	9.36
PANTHEON	–	1026.88	–	1026.99
BAO FS BOSS DR12	–	8.04	–	9.02
BAO BOSS low- z	–	3.20	–	3.48
SH0ES	–	10.67	–	1.91
COSEBI	–	1.2	–	4.07
total	2765.99	3816.23	2761.98	3808.40

TABLE XI. Best-fit χ^2 per experiment (and total) in Λ CDM and the 1-parameter EDE model when marginalizing over the lensing information in Planck.

beat down cosmic variance, one should run several statistical realizations of the same simulation, by producing initial conditions starting from different random seeds. To circumvent this issue and save computational time, we adopted the simple solution to run the two sets of simulations (EDE and Λ CDM) with identical random seeds for

the realization of their initial conditions, and to present our results in terms of ratios in the matter power spectra between the EDE and the Λ CDM models, in Figs. 12 and 13. Any scatter related to the cosmic variance is now removed, allowing us to go down one order of magnitude in terms of wave-numbers k 's. We show in both

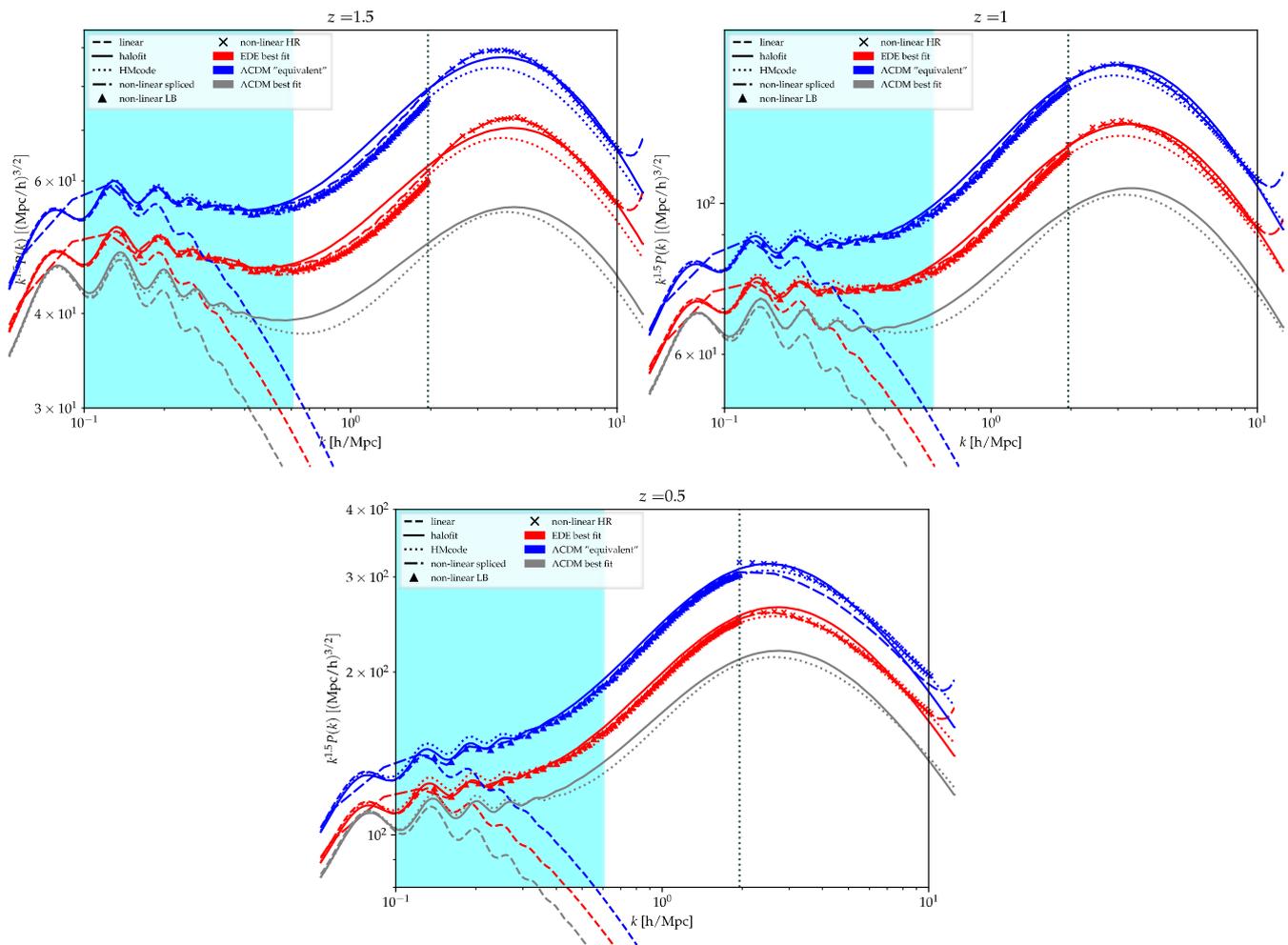


FIG. 10. Here we compare the matter power spectra extracted from our simulations, with the ones computed with HALOFIT/HMCODE, in three different redshift bins from $z = 1.5$ to $z = 0.5$. The blue curves refer to the Λ CDM scenario, whereas the red ones refer to the EDE best fit model. As a reference, we also report the best fit Λ CDM case from Planck 2018. The spliced power spectra are denoted by thick dot-dashed lines. Symbols stand for the output power spectra of the “non-spliced” LB and HR simulations. The solid/dotted lines are the non-linear power spectra from HALOFIT/HMCODE, while the dashed lines are the corresponding linear power spectra used to set the initial conditions for the simulations. The cyan shaded band approximately corresponds to the scales probed by DES-Y1.

figures by vertical dashed lines the scales corresponding to k_{MIN}^{250} and $k_{\text{Nyq}}^{\text{LB}}/2$. One can clearly see again that the EDE by itself lead to a decrease in power. However, the increase in ω_{cdm} , leads the EDE best-fit model to predict $\mathcal{O}(20\%)$ increase in power, when compared to the best-fit Λ CDM model. Note how the differences become even more manifest at higher redshift. This illustrates that high- z LSS measurements have the potential to put EDE under crucial tests [107].

Another way of presenting our results is in terms of the accuracy at which HALOFIT/HMCODE can predict *devi-*

ations in the non-linear power spectrum of EDE models with respect to the Λ CDM ‘equivalent’ case (as opposed to predicting the absolute power spectrum). This is what we show in Fig. 13, where we now compare the ratio between the EDE and Λ CDM power spectra from HALOFIT/HMCODE against the same ratio extracted from simulations. The thick horizontal lines highlight $\pm 5\%$ deviations. In light of all of this, we conclude that, in the EDE framework, HALOFIT/HMCODE predictions on Λ CDM departures are reliable at $\leq 5\%$ level with respect to the outputs of N -Body simulations, for scales $10^{-2} \lesssim k \lesssim 10$ h/Mpc , at redshifts $0 \leq z \leq 2$.

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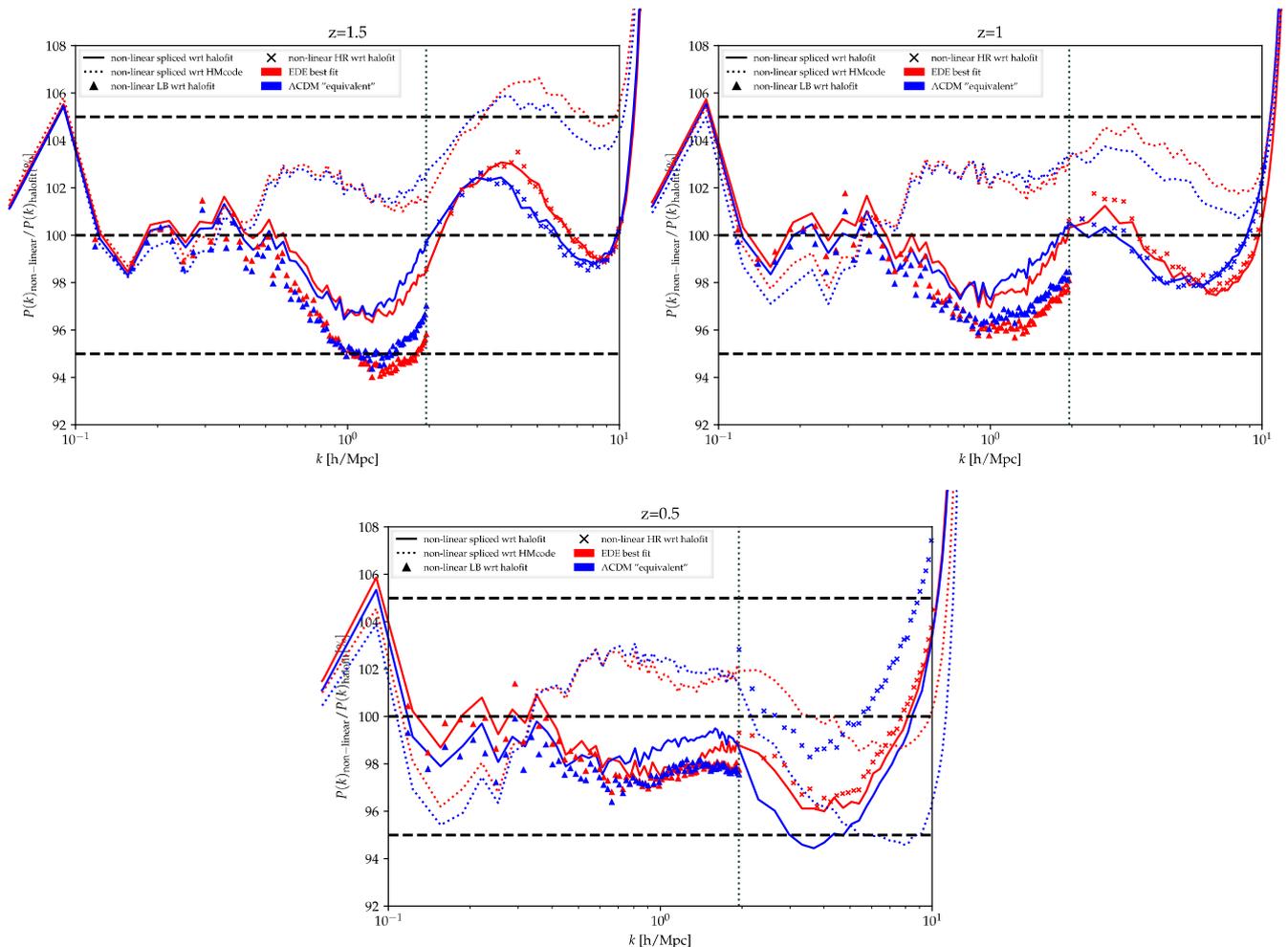


FIG. 11. Here we show the ratio between the non-linear matter power spectra from our simulations and the ones computed with HALOFIT/HMCODE, for both the Λ CDM equivalent and the EDE best-fit models. We have adopted the same linestyle-code and color-code of Figure 10.

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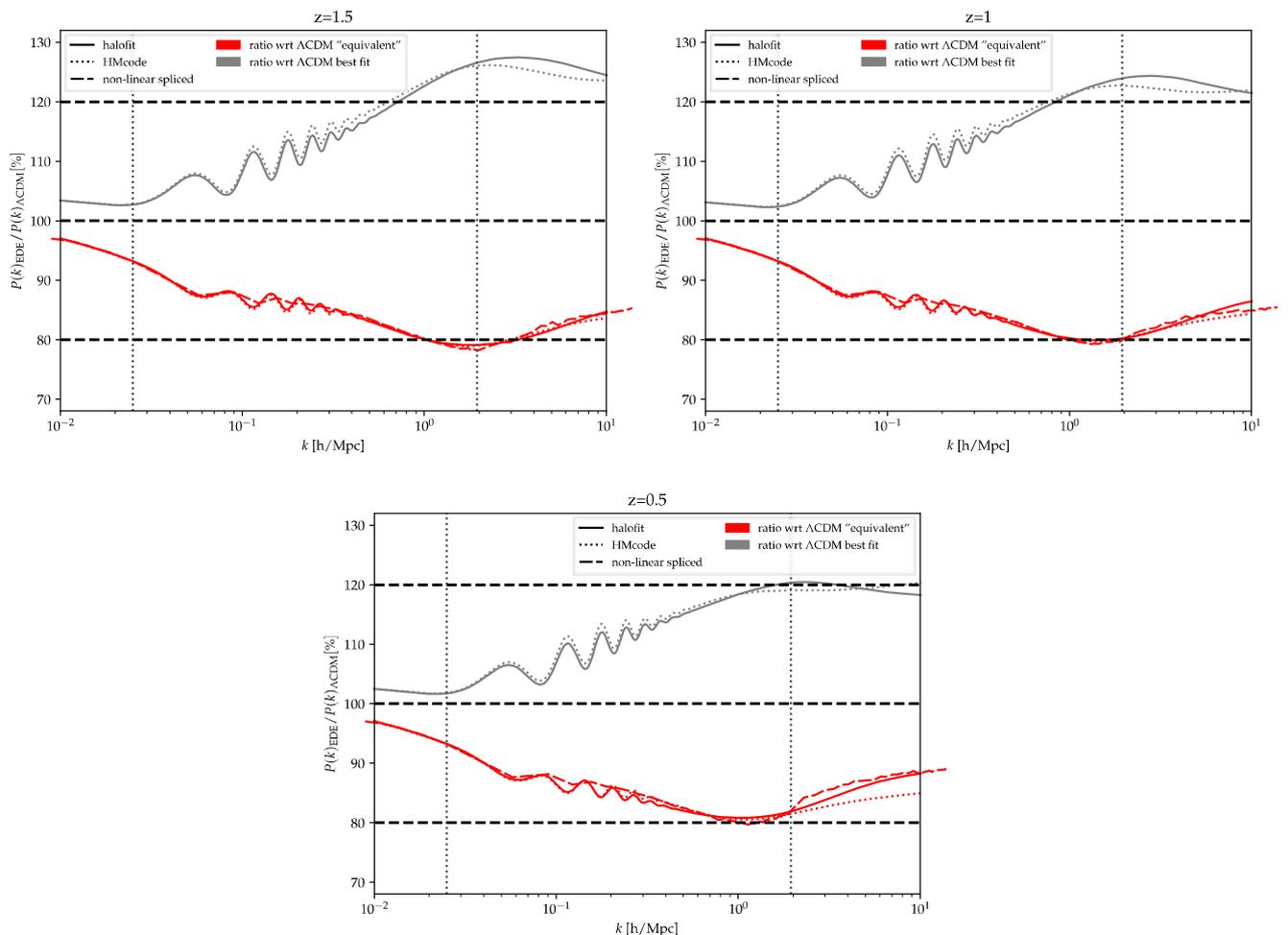


FIG. 12. Here we compare departures from the Λ CDM “equivalent” model in terms of ratios of non-linear matter power spectra. The EDE best-fit case is shown in red, while the gray lines refer to the Λ CDM best-fit model. Solid and dotted lines stand for the non-linear power spectra from HALOFIT and HMCODE, respectively. Dot-dashed lines refer to the outputs of our simulations. The cyan shaded band approximately corresponds to the scales probed by DES-Y1.

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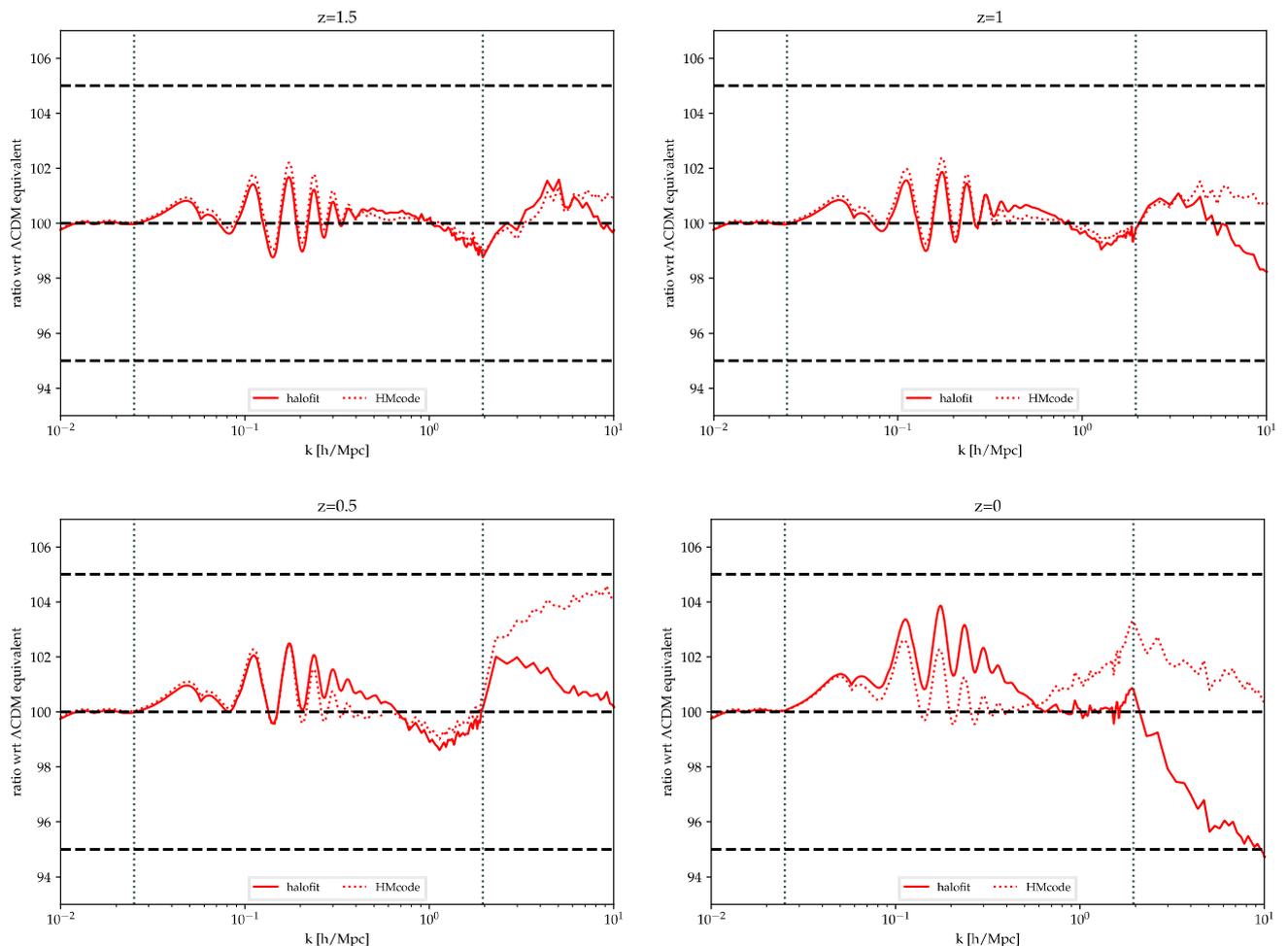


FIG. 13. Here we compare the outputs of our simulations with the HALOFIT/HMcode predictions, in terms of deviations in the ratios of the EDE best-fit power spectra over the Λ CDM “equivalent” ones, in four redshift bins from $z = 2$ to $z = 0$. Solid and dotted lines stand for HALOFIT and HMcode, respectively.

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