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1 Structural-scale modeling of the active confinement effect in the steel- 2 concrete bond for reinforced concrete structures

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7

8 Abstract

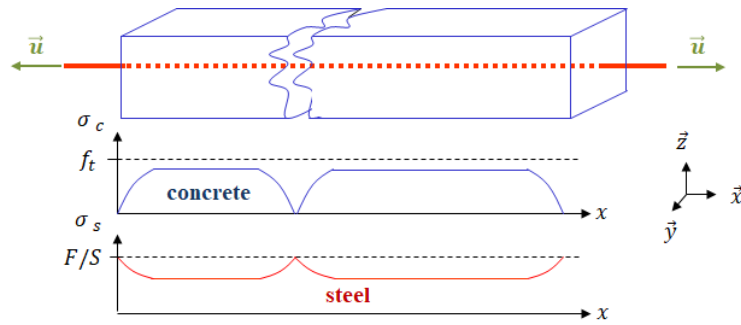
9 A numerical model to take into account the effect of the stress state on the bond behavior between
10 steel and concrete in reinforced concrete structures is proposed. It is based on a zero thickness
11 element, adapted to large-scale simulations and the use of 1D elements for steel bars. The proposed
12 model also assumes the definition of a bond stress – slip law which includes the confining pressure
13 around the steel bar as a parameter. The implementation of the model is presented and the
14 calibration of the bond law is discussed. A general equation is especially proposed. This evolution law
15 is validated through the comparison to 28 pullout tests. The model is able to reproduce the evolution
16 of the bond stress (especially the bond strength) as a function of the confinement pressure,
17 whatever the configuration (different concrete cover to steel diameter ratios). Finally, the effects at
18 the structural level are investigated on a reinforced concrete tie. The response for different confining
19 pressures is especially studied. It shows the capability of the model to reproduce the “expected”
20 tendencies with an increase of the initial elastic stiffness with increasing pressures and consequently
21 a higher number of cracks in the stabilized nonlinear regime. The “transfer length” is also shown to
22 decrease with increasing confining pressures.

23 1. Introduction

24 Steel is widely used in civil engineering applications to strengthen concrete in tension. These so-
25 called reinforced concrete structures, which present a more ductile behavior compared to plain
26 concrete, may nevertheless be subjected to cracking. In this case, when a crack initiates, stresses in
27 concrete drop to zero and the loading is totally supported by the reinforcement. They are then
28 responsible for stress transfer around the crack from steel to concrete. This progressive
29 redistribution, which can be easily demonstrated in the case of a reinforced concrete tie (Figure 1)
30 [1], is directly influenced by the bond properties [2]. That is why the influence of the steel-concrete
31 bond has to be carefully studied, especially when the crack properties, which are directly related to
32 this stress distribution, play a key role in the structural functions (failure mode and confinement [3]
33 for example).

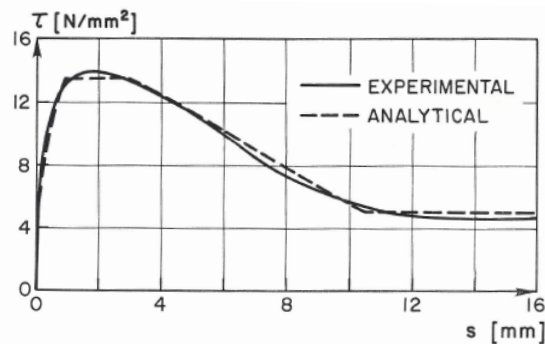
34 Experimentally, steel-concrete bond is generally described following three different steps [4]: a
35 perfect “chemical” bond (no slip), then a gradual degradation of concrete around the steel ribs,
36 followed by crack propagation (associated with a steel-concrete slip), and finally a total degradation
37 of the interface with only a residual friction. It generally leads to the definition of an adhesion law
38 that gives the evolution of the bond stress as a function of the slip at the interface (Figure 2). The
39 influent parameters on this adhesion law have been widely studied in the literature.

40



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42
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Figure 1. Principle of the distribution of steel and concrete stresses in a reinforced concrete tie after the first crack (σ_c and σ_s stand for the stresses in concrete and steel respectively) [1].



44
45

Figure 2 : Example of experimental bond stress-slip law [5].

46 Steel and concrete properties (relative rib area [6], steel diameter [7], concrete compressive and
47 tensile strength [8]-[9] and concrete cover to steel diameter ratio [10]) can be considered as the main
48 impacting parameters. They lead to potential complex formulations for the adhesion law, including a
49 distinction between splitting and pullout failure ([11] among others).

50 Besides these material or geometrical parameters, the stress state around the reinforcement may
51 have also an impact on the adhesion law. Especially, the confinement inside concrete may increase
52 the bond strength. This confinement effect can be induced either directly (by the application of an
53 external loading, like a pressure or a prestress) or indirectly (through the presence of secondary
54 reinforcements which prevents concrete cracking in certain directions [12]). Experimental studies,
55 generally performed on pullout tests, conclude on a positive influence of the concrete compression
56 stress state, whose range is dependent on the geometry of the specimen (concrete cover) [13] [14].
57 Eligehausen et al. [5] especially showed that the maximum bond stress increases with the imposed
58 lateral pressure. Malvar [15] obtained around twice the initial bond strength by applying a lateral
59 pressure from 3.5 to 31 MPa. Verderame et al. [16] and Jin et al. [17] demonstrated that the active
60 confinement had a significant effect on the cyclic bond behavior.

61 Some empirical formula were then proposed for the adhesion law to take into account this effect.
62 Based on the experimental observations from [5], [15] and [18], Lowes et al. [19] proposed a
63 relationship between the ultimate bond strength and the confining pressure, which explicitly
64 includes the lateral pressure. Xu et al. [20], Zhang et al. [21], Wu et al. [22] also included the positive
65 effect of lateral compression and the negative effect of lateral tension [23] in the bond stress – slip
66 law.

67 Even if the influence of the stress state on the bond properties has been experimentally observed, its
68 inclusion in a finite element model, compatible with engineering computations, is rather rare in the
69 literature, despite the attempts from Lowes et al. [19] for example. In this contribution, a dedicated

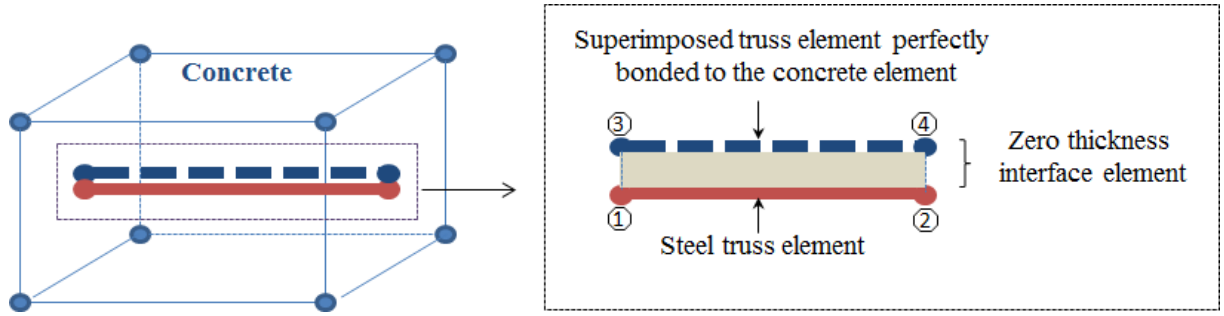
70 finite element model is thus proposed, based on the initial contribution from Mang et al. [24]. It is
71 improved to take into account the stress state in the bond behavior. It supposes the definition of a
72 bond stress – slip law, which is dependent on the stress state. A law is proposed and discussed by a
73 comparison to experimental results on pullout tests. Finally, the simulation of a reinforced concrete
74 tie subjected to different lateral pressures illustrates the impact of the confinement at a structural
75 level. It is especially shown how the lateral pressure can affect the transfer length between steel and
76 concrete.

77 **2. Steel concrete bond model including “confinement” effect**

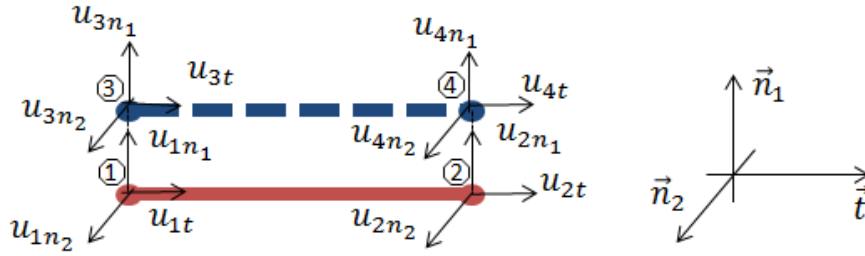
78 2.1. Presentation of the interface element

79 When reinforced concrete structures are considered, one of the most usual hypotheses, especially
80 for engineering computations, is to model the steel reinforcement as truss elements and to consider
81 a no-slip perfect relation between steel and concrete. This perfect relation is generally applied
82 through kinematic relations between both models, using the shape functions of each element.
83 However, it may have consequences, especially when the crack properties (spacing and openings) are
84 studied, as the steel – concrete bond directly influences their evolutions ([25] for example). To take
85 into account the interfacial behavior between steel and concrete in a more appropriate manner,
86 different models exist. They range from analytical or semi-analytical approaches (tension-stiffness
87 effect in uniaxial tension [26], [27]) to more complex simulation methods (including fracture
88 mechanics [28]). In the frame of finite element method and continuum mechanics, Ngo and Scordelis
89 [29] proposed a spring element, associated with a linear law, to relate concrete and steel nodes. To
90 improve the description of the bond behavior, joint elements have been developed. These zero
91 thickness elements, introduced at the interface between steel and concrete, allow the use of a
92 nonlinear law ([30], [31] among others). Special finite elements can also be used to enclose, in a
93 same element, the material behavior (steel or/and concrete) and the bond effects [32].
94 Ibrahimbegovic et al. [33]), among others, also proposed embedded elements whose principle is to
95 describe the steel-concrete bond behavior through an enrichment of the degrees of freedom. Even if
96 these solutions give appropriate results, one of their main drawbacks, in the context of industrial
97 applications, is the need to explicitly consider the interface between steel and concrete. It may
98 impose meshing difficulties and heavy computational cost which are not compatible with large scale
99 simulations.

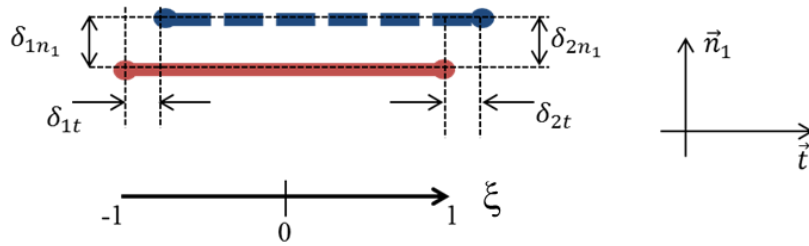
100 To overcome these difficulties, alternative solutions exist. For example, in [34], the slippage is
101 accounted for in an indirect manner through damage factors and the method is applied successfully
102 to full-scale RC structures. However, when the values of the slip are needed (for example, to capture
103 the position and the opening of the cracks [35]), the slippage has to be explicitly computed. To do so,
104 Lykidis [36] proposed a link element using 1D rebar elements embedded within 20-noded hexahedral
105 element. An alternative approach has also been developed in [1] then [24] to represent bond effects
106 between steel, modeled with truss elements, and the surrounding 3D concrete through a “1D-3D”
107 interface element. This type of developments can be seen as a macroscopic representation of local
108 effects at the interface between steel bar ribs and surrounding concrete. It is to be noted that these
109 local effects could be also considered at a very local scale using only adapted constitutive laws for
110 concrete and steel. However, the resulting approach would not be compatible with structural scale
111 computations, contrary to the proposed strategy.



112
113 Figure 3. Principle of the interface element between steel and concrete [24].



114
115 Figure 4. Degrees of freedom of the interface element [24].



116
117 Figure 5. Definition of the slip between steel and concrete in the interface element in the (\vec{t}, \vec{n}_1) plane [24].

118 In this contribution, the formulation of the element, initially developed in [24], is improved to take
119 into account the influence of the concrete stress state.

120 The principle of the interface element is first briefly recalled. It is a zero thickness four node element
121 which relates each steel truss element with an associated superimposed segment, perfectly bonded
122 to the surrounding concrete (Figure 3), through additional kinematic relations. The nodal unknowns
123 of the reinforcement bar are thus retained and a relative slip between steel and concrete becomes
124 possible through the interface element. Each node of the interface element has three degrees of
125 freedom (nodal displacements) (Figure 4). The relation between the generalized slip in the local
126 direct frame $\{\delta(\xi)\}$ (Figure 5) and the nodal displacements $\{u\}$ is written in the following form:

127
$$\{\delta(\xi)\} = \{\delta_t(\xi) \quad \delta_{n_1}(\xi) \quad \delta_{n_2}(\xi)\}^T = \bar{\bar{B}}(\xi)\{u\} \quad (1)$$

128 with

129
$$\bar{\bar{B}}(\xi) = [\bar{\bar{B}}_1(\xi) \quad \bar{\bar{B}}_2(\xi) \quad -\bar{\bar{B}}_1(\xi) \quad -\bar{\bar{B}}_2(\xi)] \quad (2)$$

130 and

131
$$\begin{aligned} \bar{\bar{B}}_1(\xi) &= 0.5(1 - \xi)\bar{\bar{l}}_3 \\ \bar{\bar{B}}_2(\xi) &= 0.5(1 + \xi)\bar{\bar{l}}_3 \end{aligned} \quad (3)$$

132 where $\underline{I}_3 = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$ and $-1 \leq \xi \leq 1$ (Figure 5).

133 Constitutive laws needs then to be defined between the bond stress $\{\sigma(\xi)\} = \begin{Bmatrix} \sigma_t(\xi) \\ \sigma_{n_1}(\xi) \\ \sigma_{n_2}(\xi) \end{Bmatrix}$ and the slip
 134 $\{\delta(\xi)\}$. In the tangential direction, the tangential stress σ_t is computed from the tangential slip:

$$135 \quad \sigma_t(\xi) = f(\delta_t(\xi)) \quad (4)$$

136 In the normal directions, a linear relation is assumed between the stresses σ_{n_1} and σ_{n_2} and the
 137 corresponding normal slips:

$$138 \quad \begin{Bmatrix} \sigma_{n_1}(\xi) \\ \sigma_{n_2}(\xi) \end{Bmatrix} = k_n \begin{Bmatrix} \delta_{n_1}(\xi) \\ \delta_{n_2}(\xi) \end{Bmatrix} \quad (5)$$

139 For the sake of simplicity, the value of the normal stiffness k_n is chosen high enough to be
 140 representative of a perfect bond in the normal directions ($k_n = 10^{15}$ Pa.m⁻¹ in the following). An
 141 improvement could be to take into account the effect of the slip behavior in the normal direction in a
 142 more appropriate way. One solution would be to consider a unilateral contact if normal stress is in
 143 tension, or to include a simple damage model. However, in classical configuration, the mechanical
 144 degradation of the concrete elements related to the interface element is supposed to be enough to
 145 capture the overall behavior correctly.

146 This first version of the model is improved to consider the effect of the confinement pressure in the
 147 tangential bond behavior. Eq. (4) is thus replaced by:

$$148 \quad \sigma_t(\xi) = f(\delta_t(\xi), p_{lat}) \quad (6)$$

149 with p_{lat} the concrete confinement pressure. It is computed using the same definition as in [20] [21]
 150 and based, by hypothesis, on the mean value of the normal concrete stresses:

$$151 \quad p_{lat} = \frac{\sigma_{n_1n_1} + \sigma_{n_2n_2}}{2} \quad (7)$$

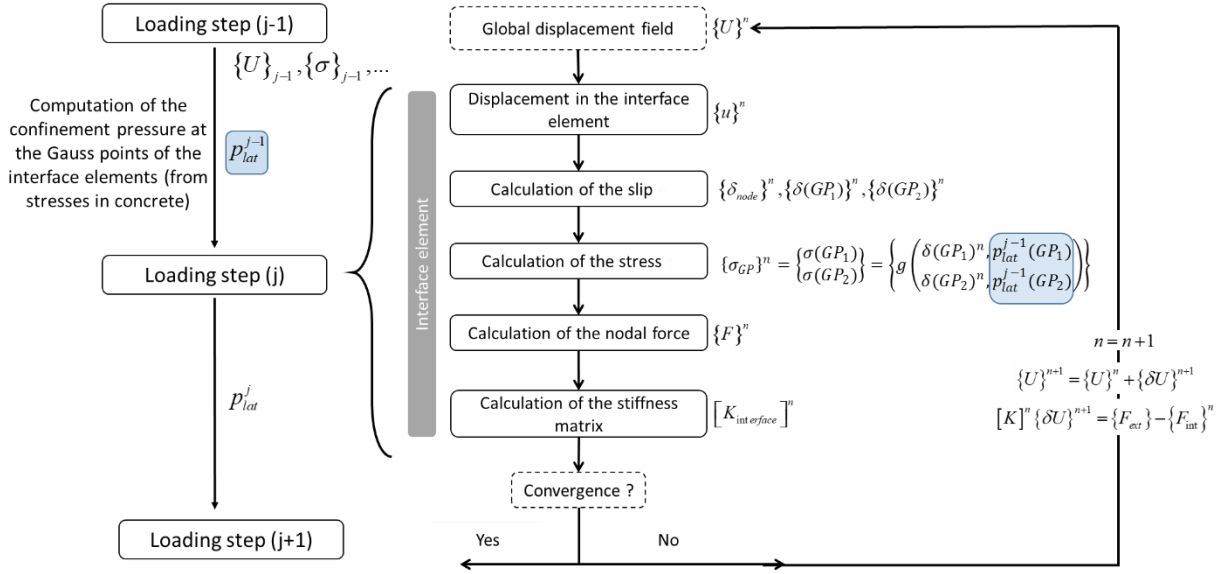
152 where $\sigma_{n_1n_1}$ and $\sigma_{n_2n_2}$ are the calculated concrete stresses in n_1 and n_2 directions. Eq. (6) is
 153 computed at each of the two integration points of the interface element to obtain the bond nodal
 154 forces after an analytical integration [24]. The concrete stresses are thus needed to be calculated at
 155 the position of these integration points, using the shape functions of the concrete elements.

156 In this contribution, the distribution of the confinement pressure p_{lat} is calculated at the end of each
 157 converged loading step. It is then used at the following loading step, especially for equation (6).
 158 There is no update during each internal iterative loop. This kind of consideration may delay the
 159 response of the active confinement on the bond but can be easily compensated by using sufficiently
 160 small calculation steps. Moreover, regarding the quasi-static evolution of the studied systems and
 161 the expected loading history (confinement pressure generally applied in one-step), this simplification
 162 is considered valid.

163 The general algorithm is summarized in Figure 6. The overall convergence is obtained with a
 164 tolerance equal to 10^{-4} .

165

166



167

168 Figure 6. Principle of the resolution for the interface element. j stands for the loading step, n for the iteration
 169 number. U and u are the global displacement field and the displacement in the interface element respectively. δU
 170 is the increment in the displacement for each iteration. δ is the slip, σ the stress and GP₁ and GP₂ are the
 171 positions of the two integration points for each interface element. F are the nodal forces and K the resolution
 172 matrix.

173 It is to be noted that one advantage of the proposed method is the possibility to mesh 1D steel bars
 174 and 3D concrete volumes can be meshed independently. For large industrial structures with a high
 175 number of rebars ([37] for example), an efficient mesh generation method can also be considered
 176 ([38]).

177 2.2. Proposition of a tangential bond law including confinement

178 As previously mentioned, taking into account the influence of the concrete stresses on the bond
 179 properties supposes the definition of an appropriate bond stress – slip law in the tangential direction.
 180 This should observe the following experimental statements:

- 181 • Zhang, et al. [14], Lowes, et al. [19], Robins and Standish [39], among others, reported that
 182 the active confinement effect is not significant on the shape of the bond law but can be considered
 183 only on the value of the bond stress for a given slip.
- 184 • Tension and compression stress states have different effects on the bond behavior. The bond
 185 strength increases with increasing lateral compression ([20], [39]), while it decreases with increasing
 186 lateral tension [22].
- 187 • Bond properties increase with the ratio of lateral pressure over the compressive strength f_c :

188
$$\sigma_t = f \left(\frac{P_{lat}}{f_c} \right) \text{ ([5], [21], [40])}$$

- 189 • Finally, the lower the concrete cover c to steel diameter d_s ratio, the greater the effect of
 190 lateral pressure ([20], [41]). This transition in the behavior is recognized in [42]. The confining

191 pressure is considered to be only able to enhance splitting behavior (small c/d_s ratio) whereas the
 192 pullout limit state (where the concrete fails in shear – higher c/d_s ratio) is not enhanced significantly
 193 by confinement. This ratio has already been highlighted in case of no “active” confinement ($p_{lat} = 0$)
 194 [11].

195 Based on these four main considerations, the influence of the lateral pressure is proposed to be
 196 considered in Eq. (6) through Eq. (8). For sake of simplicity, in the following, τ will stand for the
 197 tangential bond stress σ_t .

$$198 \quad \tau(p_{lat}, \delta_t) = \tau_0(\delta_t) \left(1 - \text{sgn}(p_{lat}) \cdot \alpha \sqrt{\left| \frac{p_{lat}}{f_c} \right|} \right) \quad (8)$$

199 τ_0 is the bond stress – slip law for $p_{lat} = 0$ and α is a parameter. sgn stands for the sign of p_{lat} (>0 in
 200 tension and <0 in compression) in order to represent the different effect of the confinement in
 201 tension and in compression. α is defined as a function on the concrete cover to steel bar diameter
 202 ratio. A transition value for c/d_s is especially considered, below which the influence of lateral
 203 pressure is quite strong and above which its influence is weak [20] [42]. This transition value is
 204 chosen from [11]:

$$205 \quad \left(\frac{c}{d_s}\right)_{tr} = 0.39 \frac{f_c}{f_t} - 0.24 \quad (9)$$

206 with f_t the concrete tensile strength.

207 α is near 1 for very small c/d_s ratios and tends toward 0 for large values, with a quite strong drop
 208 around the transition value. A continuous expression is proposed based on exponentials:

$$209 \quad \alpha = 1 - e^{-\beta \left(\frac{c}{d_s}\right)^a} \quad \text{if } \frac{c}{d_s} < \left(\frac{c}{d_s}\right)_{tr} \quad (10)$$

$$210 \quad \alpha = e^{-\gamma \left(\frac{c}{d_s}\right)^b} \quad \text{if } \frac{c}{d_s} > \left(\frac{c}{d_s}\right)_{tr} \quad (11)$$

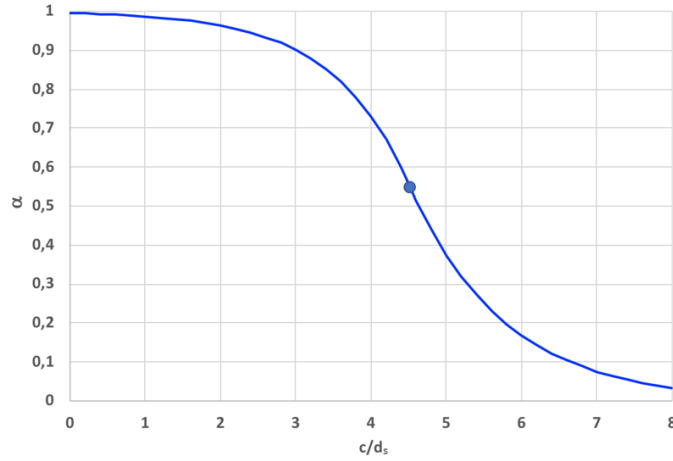
211 β and γ are parameters to control the shape of the exponentials, while a and b are adjusted to assure
 212 the continuity of α and its derivative at the transition point $\left(\frac{c}{d_s}\right)_{tr}$. In the following $\beta=1$ and $\gamma=0.8$. An

213 example of the evolution of α as a function of $\left(\frac{c}{d_s}\right)$ for $\left(\frac{c}{d_s}\right)_{tr} = 4.5$ ($a=5.31$ and $b=3.77$) is given in

214 Figure 7. The resulting bond stress – slip laws for $\left(\frac{c}{d_s}\right) = 2$ and $\left(\frac{c}{d_s}\right) = 7$ and different p_{lat} values are

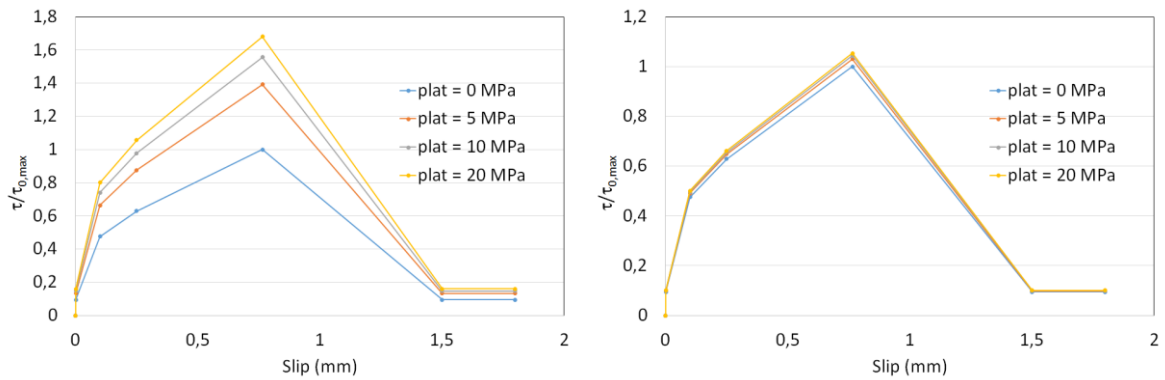
215 given in Figure 8 (f_c chosen equal to 30 MPa). In Figure 8, a piecewise linear curve has been
 216 considered as input data for $\tau_0(g)$ but some more complex evolutions could also have been chosen
 217 (nonlinear curve for example).

218



219

220 Figure 7. Evolution of α as a function of c/d_s for $(c/d_s)_{tr} = 4.5$.



221

222 Figure 8. Bond stress – slip laws for $(c/d_s)_{tr} = 4.5$ and $f_c = 30$ MPa and different confinement pressures. Left,
223 $(c/d_s) = 2$, right $(c/d_s) = 7$. $\tau_{0, \max}$ stands for the bond strength at $p_{lat} = 0$.

224 The calibration of the bond stress – slip law can be summarized in four main steps:

225 • Definition of the initial interface law $\tau_0(\delta_t)$ from either experimental or empirical evolutions.
226 In particular, the experimental bond stress – slip law obtained from a pullout test without any
227 confining pressure ($P_{conf} = 0$ MPa) can be used [43]. $\tau_0(\delta_t)$ may be slightly different from the
228 experimental pullout curve at $P_{conf} = 0$ MPa, to take into account the effect of lateral pressure $p_{lat} (\neq 0)$
229 around the steel bar during the pullout, even for a zero applied confinement pressure (structural
230 stress related to the type of test). If the experimental bond stress – slip law at $P_{conf} = 0$ MPa is not
231 available, it is possible to build the expected curve using the methodology proposed in [11], solely
232 from the material and geometric parameters.

233 • Given the geometrical and material parameters $(\frac{c}{d_s})$, f_c and f_t , calculation of $(\frac{c}{d_s})_{tr}$ from Eq. (9)

234 • Given $(\frac{c}{d_s})_{tr}$, calculation of parameters a and b to ensure the continuity of Eq. (10) and Eq. (11)

235 and its derivatives. Definition of $\alpha(\frac{c}{d_s})$.

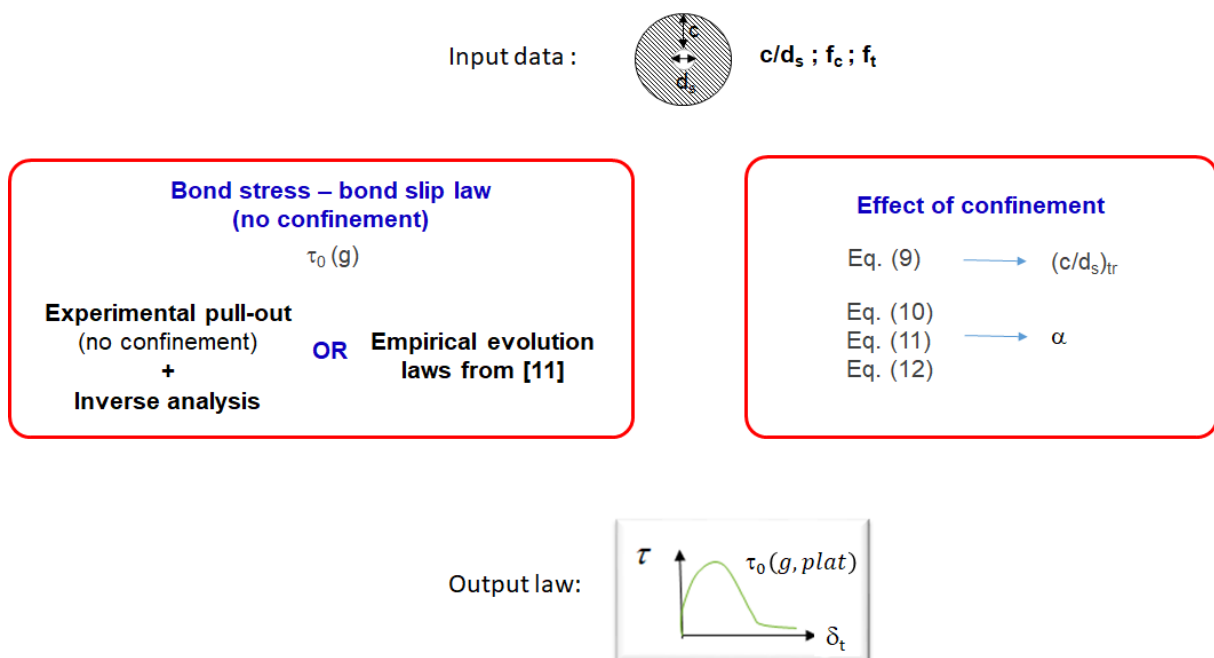
236 • Given $\alpha(\frac{c}{d_s})$, $(\frac{c}{d_s})$ and $\tau_0(\delta_t)$, definition of the bond stress – slip law $\tau(p_{lat}, \delta_t)$ from Eq. (7)

237 and Eq. (8).

238 As a summary, for a given geometry (associated to a given $(\frac{c}{d_s})$), and a given concrete (associated to
 239 a compressive strength and a tensile strength), the definition of the adhesion law only requires the
 240 initial bond stress – slip law $\tau_0(\delta_t)$ (for $p_{lat} = 0$) that can be obtained either experimentally by inverse
 241 analysis or numerically (Figure 9).

242 Compared to other evolution laws that are found in literature, the proposed methodology can be
 243 viewed as more general. For example, in [44], two evolution laws were proposed for two steel
 244 diameters, without any generalization. In [20], [21], [22], the shape and the parameters of the laws
 245 were functions of the loading (compression, tension, tension – compression respectively). Moreover,
 246 these former approaches were not expected to distinguish splitting and pullout failures as it is the
 247 case for the present methodology through the comparison to the transition value of $(\frac{c}{d_s})$. It can also
 248 be considered as easy to use, especially compared to more complex approaches in literature. For
 249 example, in [18], relations between the bond strength, the radial stress, the slip and the crack
 250 opening have been proposed, whose extension to a single independent relationship between
 251 concrete confinement and peak bond strength and radial force is not trivial [45]. Finally, as the
 252 proposed relation is function of the mean normal stresses around the bar, it is expected to reproduce
 253 both “active” and indirect confinements (through an applied pressure or transverse reinforcements
 254 respectively). Finally, the application to cyclic loadings could be considered, provided an adaptation
 255 of the bond stress-slip to alternative load (from [5] for example, combined with additional
 256 experimental data) and an adapted constitutive law for concrete.

257



258

259 Figure 9. Summary of the calibration of the bond stress – slip law, including the effect of the stress state.

260

261

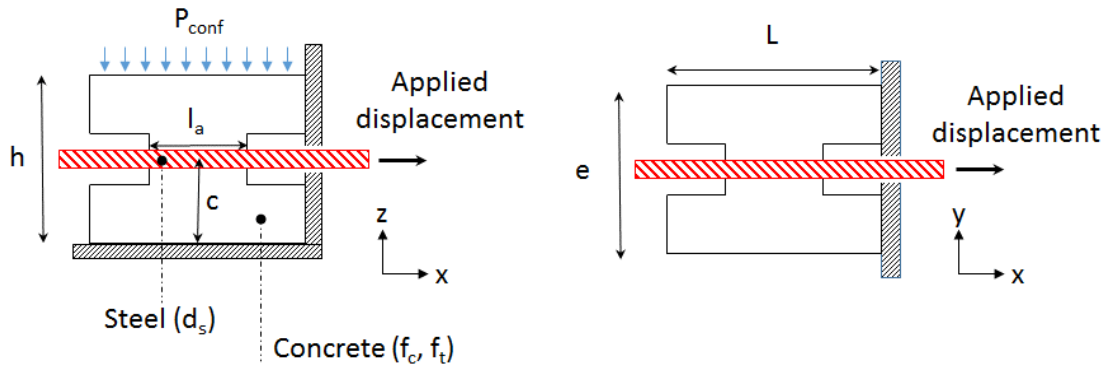
262 **3. Validation of the evolution law on pullout tests**

263 To validate the proposed bond stress – slip law, pullout tests from [11], [13] and [41] are simulated.

264 3.1 Test specimen

265 The generic geometry of pullout tests is presented in Figure 10. The detailed geometry, properties
 266 (concrete compressive strength f_c , tensile strength f_t and Young modulus E_c , steel Young modulus E_s
 267 and yielding stress f_e) and loadings (applied uniaxial or biaxial compressive confinement pressure
 268 P_{conf}) are given in Table 1 and Table 2. Concrete properties are either obtained from the given
 269 average values of experimental uniaxial compressive and tensile tests when available or evaluated
 270 using the equations provided in [2] ($E_c = 22000(\frac{f_c}{10})^{0.3}$), $f_t = 0.3(f_c)^{2/3}$).

271 The proposed validation includes different concrete properties, steel diameters and confining
 272 pressures to represent a wide enough range of configurations. In particular, it encloses
 273 configurations in which the ratio of the concrete cover over the steel diameter is higher and smaller
 274 than the transition value (Eq. (9)). It means that pullout and splitting failures are both considered in
 275 the validation process.



276
 277 Figure 10. General geometry of pull-out tests.

278

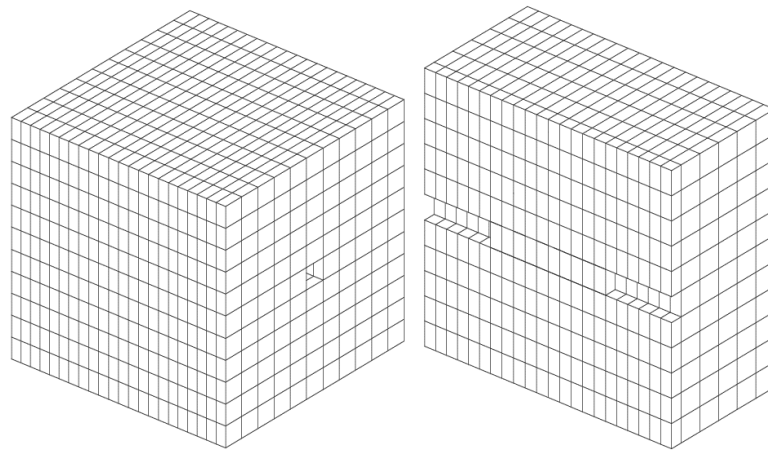
Reference of the simulated pull-out tests	d_s (mm)	L (mm)	h (mm)	e (mm)	l_a (mm)	c/d_s
(Xu et al., 2012) [13]	16	150	150	150	$5d_s$	4.19
	22					2.91
(Shang et al., 2017)[41]	14	150	150	150	$5d_s$	4.86
	18					3.67
	22					2.91
(Torre-Casanova et al., 2013)[11]	12	180	180	180	60	7

279 Table 1. Geometries of simulated pullout tests.

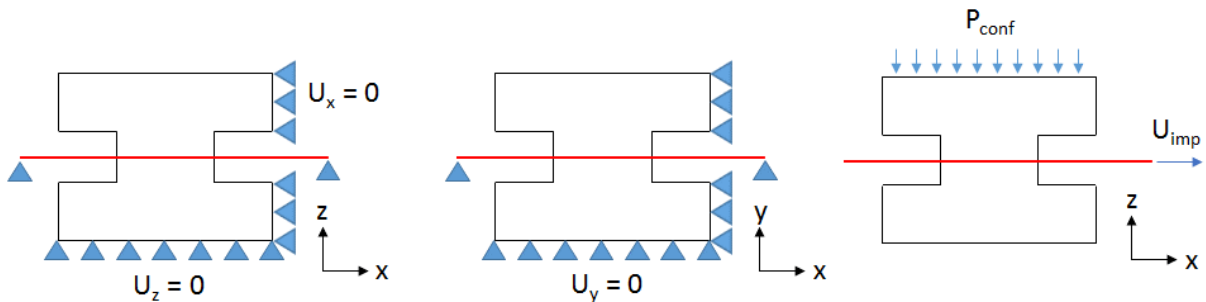
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Reference of the simulated pull-out tests	d_s (mm)	E_s (GPa)	f_e (MPa)	f_c (MPa)	E_c (GPa)	f_t (MPa)	P_{conf}	
							Type	Value (MPa)
(Xu et al., 2012) [13]	16	200	400	35	31.7	3.21	uniaxial	0 → 18
	22			30.25		2.91		
(Shang et al., 2017) [41]	14	200	560	33.93	31.7	3.14	uniaxial	0 → 20
	18							
	22							
(Torre-Casanova et al., 2013) [11]	12	200	400	36.6	28	3.12	biaxial	0 → 10

281 Table 2. Material properties and loadings for the pullout tests.



282
283 Figure 11. Example of concrete mesh for the simulations of pullout tests from [41] – $d_s = 14$ mm. On the left,
284 entire mesh, on the right, cross section along the position of the steel bar.



285
286 Figure 12. Boundary conditions (left and middle) and loadings (right) for the simulation of pullout tests – case of
287 uniaxial confinement.

288 3.2. Presentation of the simulations

289 The simulations are performed in the finite element code Cast3M [46] using cubic elements for
290 concrete, 1D truss elements for steel and the steel-concrete bond model presented in section 2.1. An
291 illustration of the concrete mesh is proposed in Figure 11. The loading and boundary conditions are
292 illustrated in Figure 12. They include zero displacement conditions on appropriate faces and imposed
293 displacements at the loaded end of the steel bar. The confinement is modeled through a uniform

294 pressure applied in one-step on the adequate face(s) at the beginning of the loading (before the
 295 imposed displacement).
 296 Concrete behavior is simulated using a damage model, which includes irreversible strains [47].
 297 Damage is represented by two independent variables d^+ and d^- which have respectively an influence
 298 in tension and compression. The stress σ is evaluated from the following relation:

$$299 \quad \sigma = [1 - d^+(\sigma'^+)]\sigma'^+ + [1 - d^-(\sigma'^-)]\sigma'^- \quad (12)$$

300 where σ'^+ and σ'^- correspond respectively to the positive and the negative parts of the effective
 301 stress σ' :

$$302 \quad \sigma' = \mathbf{C}(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}^p) \quad (13)$$

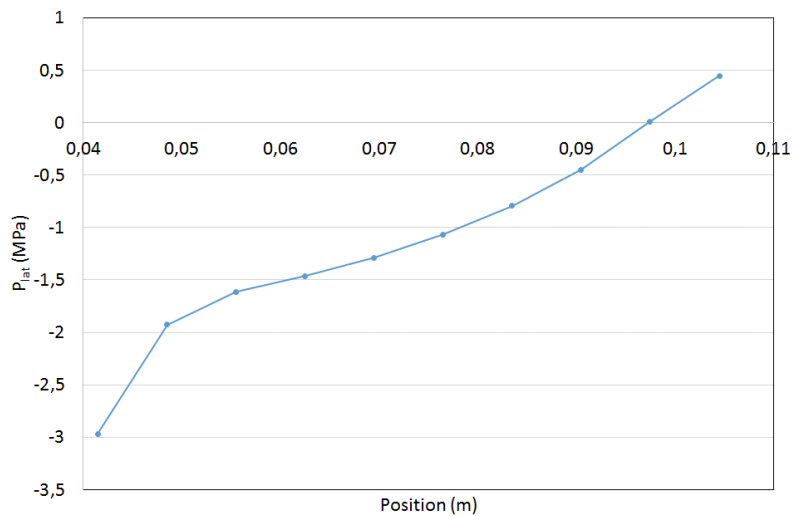
303 In this relation, \mathbf{C} is the tensor of elasticity and $\boldsymbol{\varepsilon}$ represents the total strain. $\boldsymbol{\varepsilon}^p$ stands for the
 304 irreversible strains which are governed by the damage evolution in compression:

$$305 \quad \dot{\boldsymbol{\varepsilon}}^p = \beta_p E_c H(\dot{d}^-) \langle \boldsymbol{\sigma}' : \dot{\boldsymbol{\varepsilon}} \rangle / (\boldsymbol{\sigma}' : \boldsymbol{\sigma}') \mathbf{C}^{-1} : \boldsymbol{\sigma}' \quad (14)$$

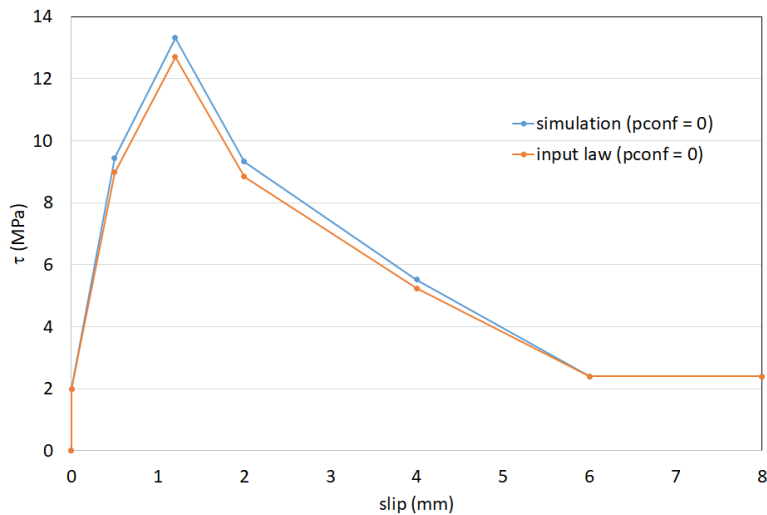
306 where β_p is a model parameter and H is the Heaviside function. $\langle . \rangle$ represents the positive part of the
 307 tensor. The tensile part of the model is regularized using the Hillerborg concept of fracture energy
 308 that guarantees a constant energy release, independently from the mesh size [48]. The model
 309 parameters are chosen to be as representative as possible to the experimental uniaxial concrete
 310 properties.

311 For sake of simplicity, the 1D steel bar follows an elastic- perfect plastic law using the steel properties
 312 provided in Table 2. It was checked that the yielding stress was not reached during the simulation (no
 313 influence of the plastic part of the model).

314 The bond stress - slip law is calibrated following the methodology described in Figure 9 from the
 315 experimental result at a zero confinement pressure. The obtained simulated p_{lat} distribution along
 316 the bar is given in Figure 13 at $P_{conf} = 0$ MPa. The difference between the input law and the simulated
 317 bond behavior, which results from the inverse analysis and the effect of the non-zero p_{lat} distribution,
 318 is illustrated in Figure 14.



319
 320 Figure 13. Simulation of a pullout test from [36] - $d_s = 14$ mm. p_{lat} distribution along the steel bar for $P_{conf} = 0$
 321 MPa at the applied displacement corresponding to the bond strength.



322

323 Figure 14. Simulation of a pullout test from [41] - $d_s = 14$ mm. Input calibrated bond law and simulated bond
 324 stress – slip law for $P_{conf} = 0$ MPa.

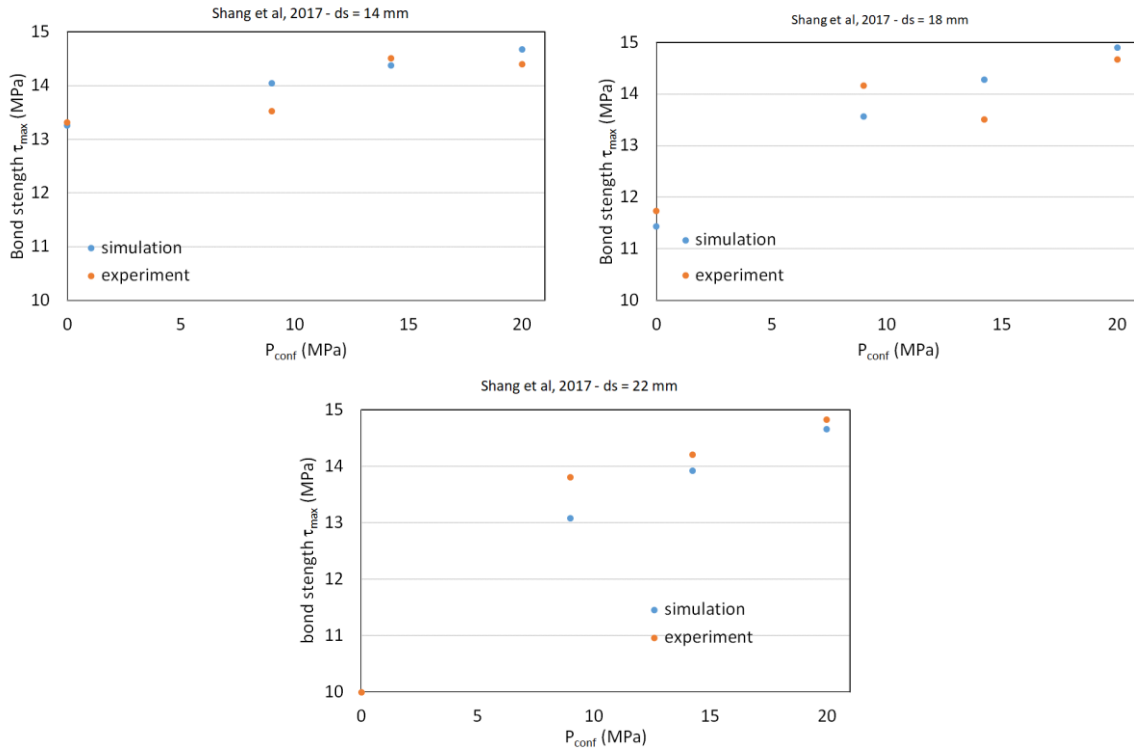
325 **3.3. Results**

326 The simulated maximum bond stresses are compared to the experimental ones for different levels of
 327 applied confinement pressures in Figure 15 ([41]), Figure 16 ([13]) and Figure 17 ([11]). It has been
 328 chosen to perform the comparison on the bond strength, as it is the main mean characteristic, which
 329 is experimentally studied. Moreover, the existence of mean values (which may be the only quantity
 330 given in the experimental results - [36] for example) enable to reduce the experimental discrepancy
 331 associated to this type of tests, compared to single bond stress – slip curves.

332 The simulated bond strength is calculated using the same equation as for the experiment:

333
$$\tau_{max} = \frac{F_{max}}{\pi d_s l_a} \quad (15)$$

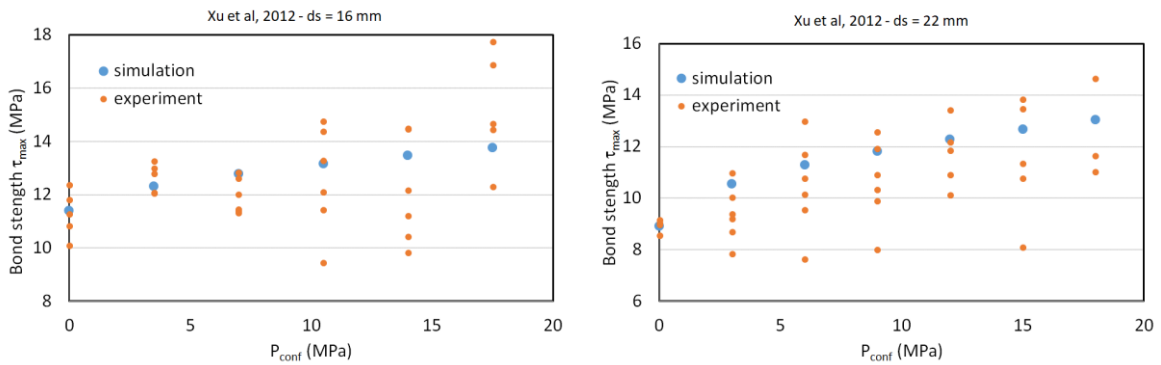
334 where F_{max} is the maximum applied force at the loaded end of the steel bar. Eq. (15) thus represents
 335 a “mean” measure of the bond along the steel bar.



336

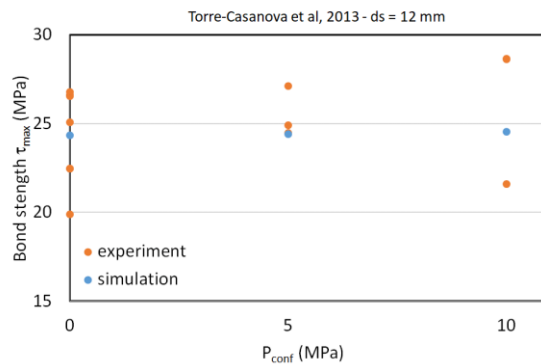
337

338 Figure 15. Evolution of the bond strength as a function of the applied confinement pressure – experiments from
339 [41] (only mean values were experimentally provided).



340

341 Figure 16. Evolution of the bond strength as a function of the applied confinement pressure – experiments from
342 [13].



343

344 Figure 17. Evolution of the bond strength as a function of the applied confinement pressure – experiments from
345 [11].

346

Applied confining pressure (MPa)	$d_s = 16 \text{ mm}$			$d_s = 22 \text{ mm}$		
	Simulated bond strength (MPa)	Mean experimental bond strength (MPa)	Mean error (%)	Simulated bond strength (MPa)	Mean experimental bond strength (MPa)	Mean error (%)
0	11.39	11.28	5.9	8.92	8.92	2.1
$0.1 f_c$	12.31	12.64	4.2	10.55	9.36	12.6
$0.2 f_c$	12.78	11.93	6	11.29	10.45	13.5
$0.3 f_c$	13.17	12.57	11.9	11.81	10.60	12.6
$0.4 f_c$	13.48	12.09	16.7	12.28	11.69	8.4
$0.5 f_c$	13.76	15.20	14.7	12.68	11.5	15.3
$0.6 f_c$	-	-	-	13.04	12.44	12.7

347 Table 3. Simulated and experimental bond strength and mean error. Pullout tests from [13].

Applied confining pressure (MPa)	$d_s = 12 \text{ mm}$		
	Simulated bond strength (MPa)	Mean experimental bond strength (MPa)	Mean error (%)
0	24.36	24.27	9.3
5	24.42	25.5	4.4
10	24.54	26.30	15.1

348 Table 4. Simulated and experimental bond strength and mean error. Pullout tests from [11].

Applied confining pressure (MPa)	$d_s = 14 \text{ mm}$			$d_s = 18 \text{ mm}$		
	Simulated bond strength (MPa)	Mean experimental bond strength (MPa)	Mean error (%)	Simulated bond strength (MPa)	Mean experimental bond strength (MPa)	Mean error (%)
0	13.26	13.32	0.4	11.43	11.73	2.6
9	14.05	13.53	3.7	13.57	14.16	4.3
14	14.38	14.51	0.9	14.24	13.51	5.1
20	14.68	14.4	1.9	14.9	14.67	1.5

Applied confining pressure (MPa)	$d_s = 22 \text{ mm}$		
	Simulated bond strength (MPa)	Mean experimental bond strength (MPa)	Mean error (%)
0	9.98	10	0.2
9	13.08	13.81	5.5
14	13.92	14.21	2.0
20	14.66	14.83	1.1

349 Table 5. Simulated and experimental bond strength and mean error. Pullout tests from [41].

350 The results are summarized in Table 3, Table 4 and Table 5 with a comparison on the mean values of
351 the simulated and experimental bond strength in each configuration. For a given configuration
352 (same reference, same steel diameter and same applied confinement pressure), the mean error is
353 also computed as the mean difference between the simulated bond strength and each measured
354 experimental bond strength, over the simulated bond strength.

355 A good agreement is generally obtained for the different values of the applied pressures and in the
356 different geometrical and material configurations (especially for different $(\frac{c}{d_s})$). It is to be noted that

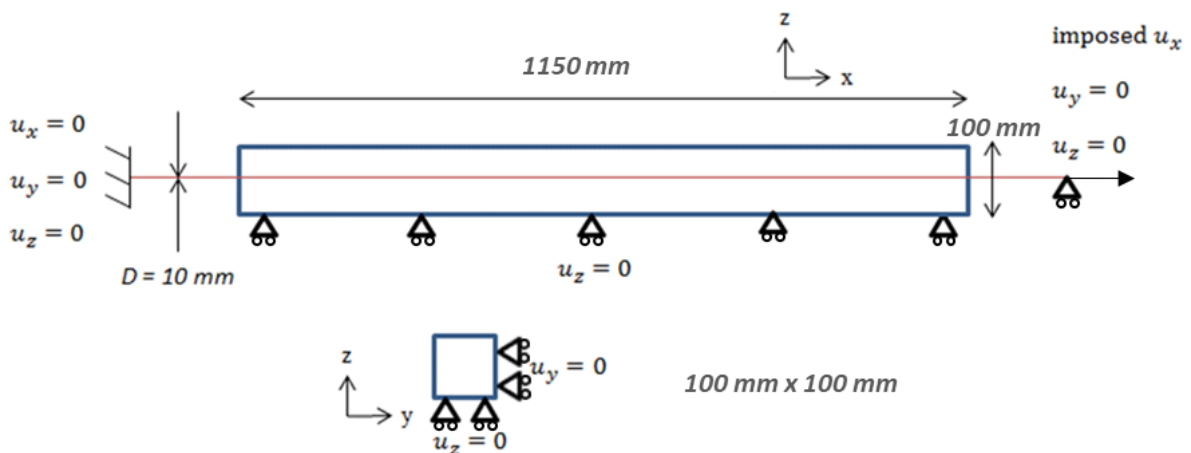
357 the experimental results may show an expected certain dispersion. In this case, the validation of the
 358 simulation is obtained if the simulation results (which are by definition deterministic in our case) are
 359 in the range of the experimental ones (with a maximum mean error equal to 15%, which is quite
 360 reasonable here). As a conclusion, the proposed comparison validates the proposed methodology
 361 and especially the evolution law for the bond stress (at least regarding the bond strength).

362 4. Structural consequences of the confinement effect – structural case

363 The previous section was dedicated to the validation of the bond stress – slip law as a function of the
 364 lateral pressure along the steel bar. In this section, the consequences at the structural level is going
 365 to be investigated by the simulation of a reinforced concrete tie loaded both in tension and in
 366 confinement. The consequences are studied in terms of global and local behaviors.

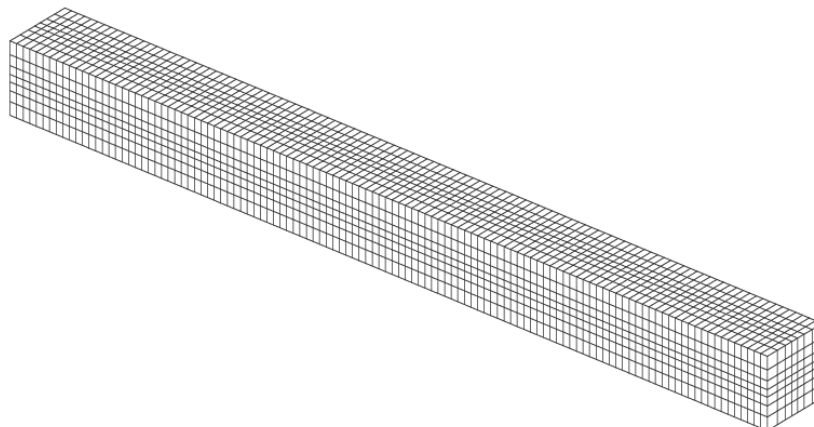
367 4.1. Presentation of the test

368 A reinforced concrete tie (length L equal to 1.15 m, square section S_c equal to 0.01 m^2), crossed by a
 369 steel bar (diameter d_s equal to 10 mm for a section S_s equal to 78.5 mm^2) is considered (Figure 18).
 370 Concrete and steel are meshed using solid elements for concrete and truss elements for steel,
 371 respectively (Figure 19).



372

373 Figure 18. Presentation of the reinforced concrete tie. Geometry, boundary conditions and loading.



374

375 Figure 19. Concrete mesh for the reinforced concrete tie.

376

377 At each end of the steel bar, one element is added to apply the boundary conditions (no
 378 displacement at one end) and the loading (imposed horizontal displacement at the other end) (Figure
 379 18). This particular structural test was chosen because it was successfully modeled in [24] and
 380 experimentally studied in [49].

381 The steel bar is modeled using an elastic-plastic law with linear hardening. Concrete follows the
 382 damage constitutive law presented in the previous section from [47]. Parameters given in Table 6
 383 and Table 7 are chosen, in agreement with the experimental data. A random distribution of the
 384 tensile strength is introduced in order to localize the damage during loading (Figure 20). This
 385 Gaussian, stationary random scalar field is generated using the turning band method. It has an
 386 isotropic exponential covariance, whose matrix is computed from the standard deviation and the
 387 correlation lengths. The mean value, the standard deviation and the correlation length are chosen
 388 equal to the mean tensile strength (2.9 MPa), 5% and 3 cm respectively.

389 For the bond model, only the bond stress – slip law $\tau_0(\delta_t)$ has to be provided. A piecewise linear
 390 curve is chosen, following the recommendations from [11]. The parameters are given in Table 8. The
 391 effect of the confinement is taken into account using the methodology presented in the previous
 392 sections (calculation of α especially).

393 A monotonic increasing displacement is applied to the loaded end of the steel bar. To evaluate the
 394 influence of the active confinement, a compressive pressure (P_{conf}) can be applied, through a surface
 395 pressure, on both lateral faces (normal to the direction of the steel bar) before the loading in
 396 displacement. The pressure is then kept constant during the imposed displacement. The lateral
 397 pressures are chosen to be equal to 5, 10 or 20 MPa.

398

Young modulus	Poisson ratio	Yielding stress	Hardening modulus
E_s	ν_s	σ_s^e	E_h
200 GPa	0.3	500 MPa	3245 MPa

399 Table 6. Steel parameters.

Young modulus	Poisson ratio	Tensile strength	Compressive strength	Fracture energy
E_c	ν_c	f_t	f_c	G_f
30.2 GPa	0.2	2.9 MPa	56.9 MPa	150 N/m

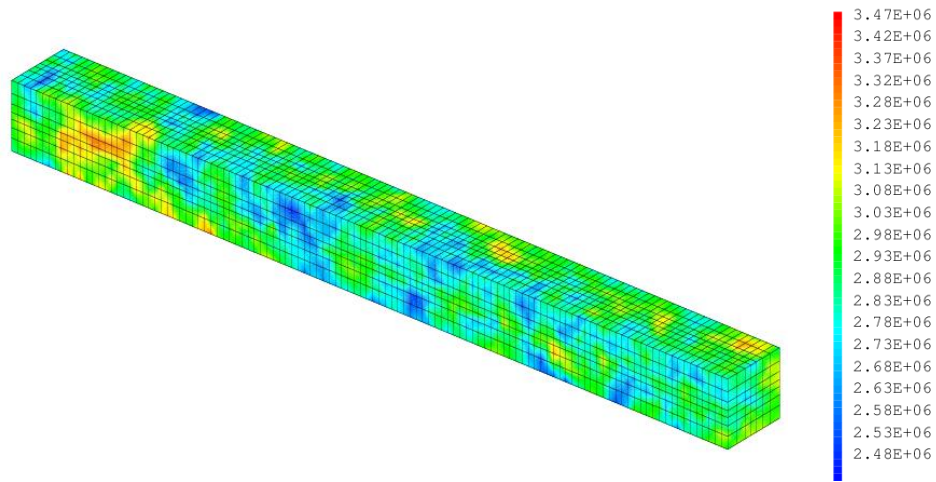
400 Table 7. Concrete parameters.

Bond stress (MPa)	2	10	13.2	21	2	2
Slip (mm)	0.002	0.1	0.25	0.765	1.5	1.8

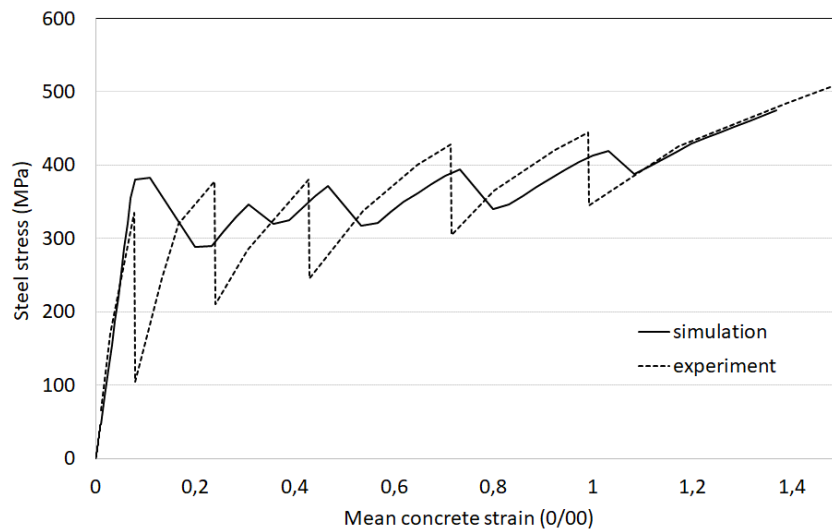
401 Table 8. Parameters for the bond stress – slip law $\tau_0(\delta_t)$.

402

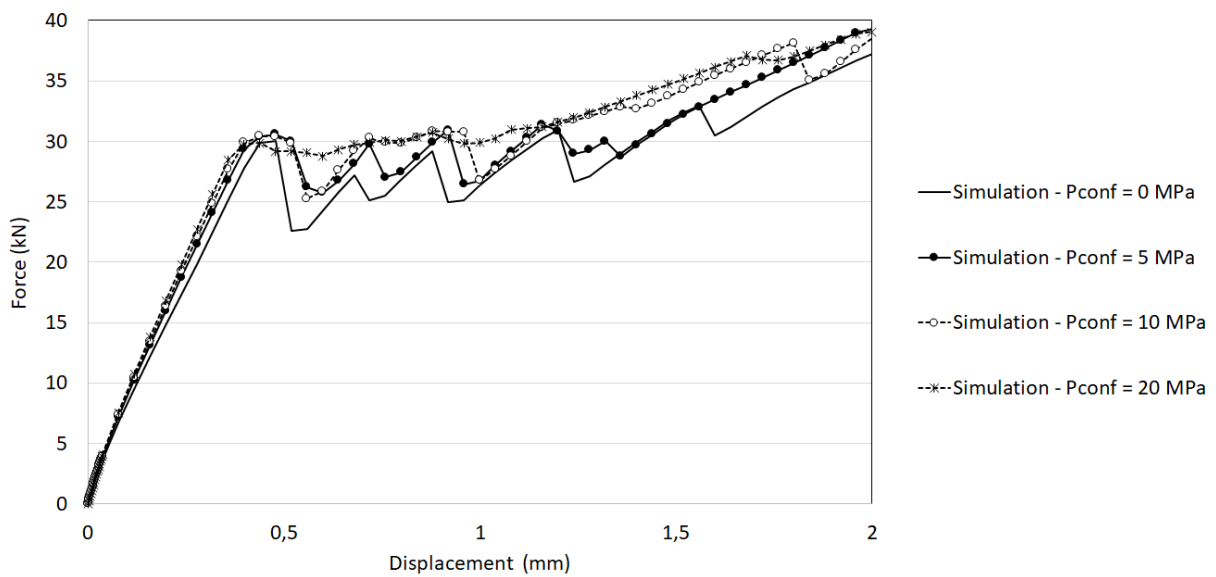
403



404 Figure 20. Distribution of the initial tensile strength (in Pa) in the reinforced concrete tie.



405
406 Figure 21. Steel stress as a function of the mean concrete strain. Comparison between the experimental result and
407 the simulation at $P_{\text{conf}} = 0$ MPa.



408 Figure 22. Force – displacement curves for the reinforced concrete tie considering different applied confinement
409 pressures.

410

4.2. Results

411 Numerical results are compared in terms of global force – displacement curves and local evolutions
 412 (mechanical degradation especially). The experimental comparison for $P_{\text{conf}} = 0$ MPa is first provided.
 413 Figure 21 gives the evolution of the steel stress at the end of the bar as a function of the mean strain
 414 in concrete, as considered in [49]. A good agreement is obtained, which validates the methodology
 415 for the simulation. The expected differences between the simulation and the experimental curve
 416 have already been discussed in [24].

417 Figure 22 shows the evolution of the force as a function of the applied displacement for different
 418 external pressures. In every case, the expected global evolution is obtained in three main steps: a
 419 linear regime in which concrete and steel behave elastically, then a nonlinear regime where concrete
 420 is gradually damaged (active cracking) and finally a stage where the number of cracks in concrete
 421 does not evolve any more (stabilized cracks). Some unloading zones are also observed when new
 422 localized damaged zones appear. The effect of the confinement pressure is especially observed on
 423 these local unloading phases, which are less significant when the confinement pressure is high.

424 Moreover, the first loading stage of the force-displacement curve is found to be stiffer as the
 425 confinement pressure increases, a phenomenon that was expected as the confinement pressure
 426 tends to stiffen the bond stress – slip law (Eq. (8)). It is also to be noted that the force F_{nl} which
 427 corresponds to the first unloading (at about 31 kN) is almost the same whatever the confining
 428 pressure. It can be demonstrated ([1]), if the length of the tie is enough compared to the transfer
 429 length (which is the case here, see Figure 26), that F_{nl} can be estimated through the following
 430 equation:

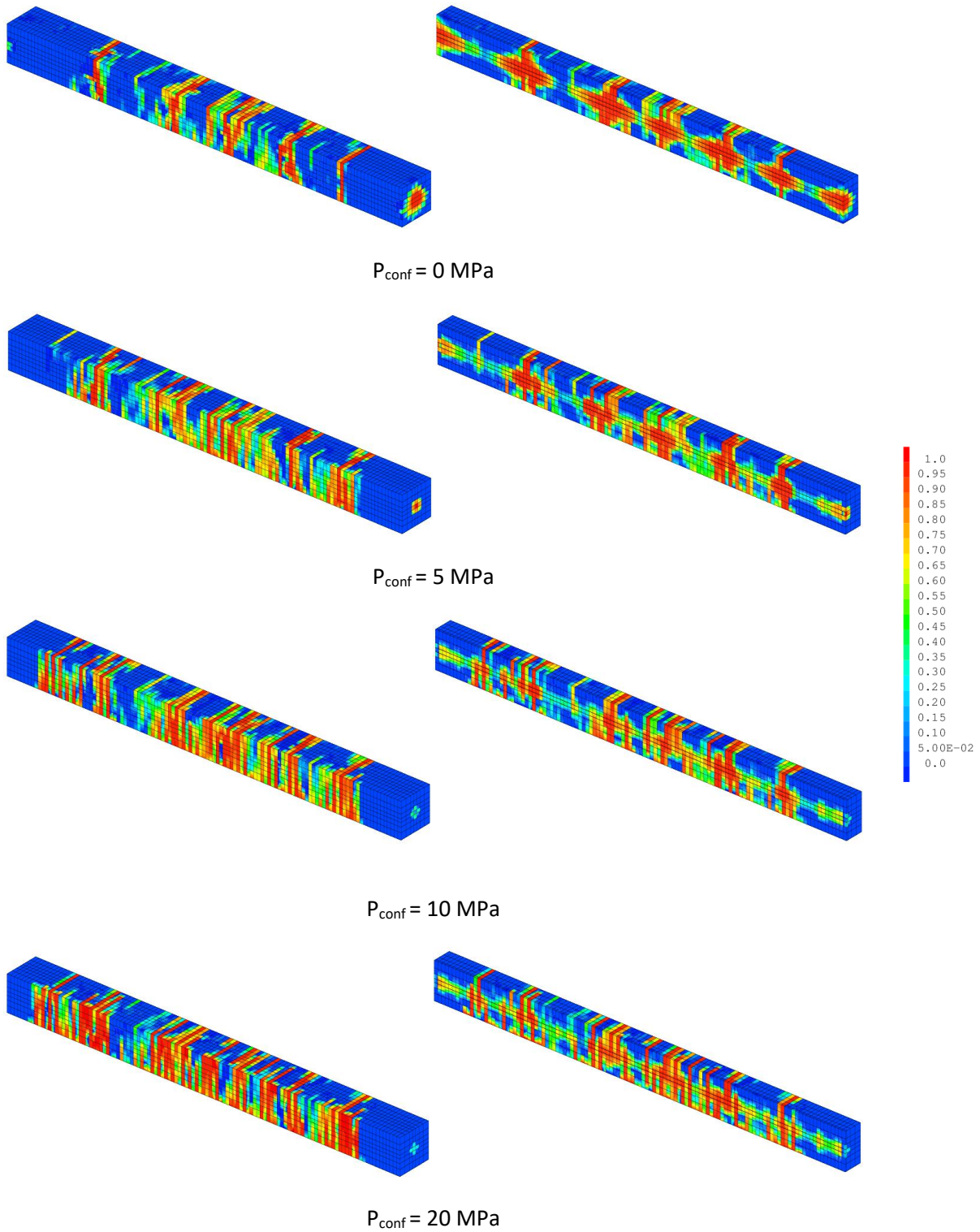
$$F_{nl} = \frac{f_t \cdot (E_s S_s + E_c S_c)}{E_c} \quad (16)$$

431 This equation is not dependent on the bond properties, as it is observed numerically.

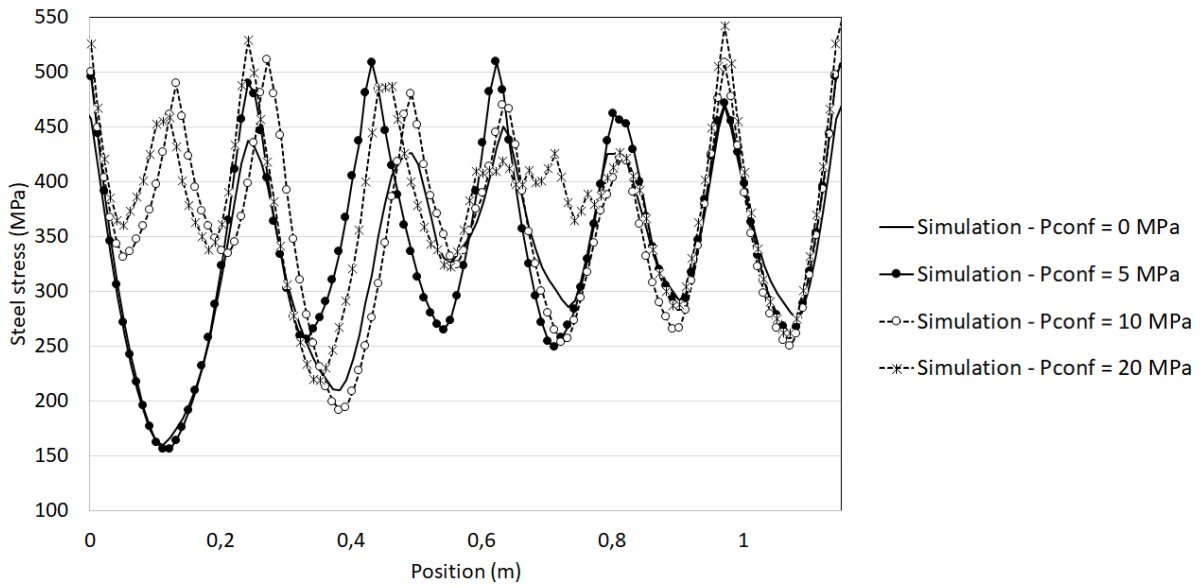
432 The damage distributions in concrete at the last displacement increment are provided in Figure 23
 433 (entire concrete tie and cross section at the position of the steel bar). As mentioned in [24], for a zero
 434 applied confinement pressure, localized damages zones appear along the tie. The higher the
 435 confinement pressure, the higher the number of localized damaged zones. For high values of the
 436 confinement pressure (20 MPa), as the bond is stronger, damage is also observed along the steel bar,
 437 almost continuously. The confining pressure also affects the value of the damage at the exit point of
 438 the rebar with a decreasing value of the mechanical degradation with an increasing confinement. The
 439 constitutive model for concrete may explain this effect, with an increasing resistance to local shear
 440 with increasing compressive stresses.

441 Figure 24 illustrates the distribution of the stress along the steel bar at the end of the loading. Peaks
 442 in the stress are related to the position of the cracks, which can also be located from the change in
 443 the sign of the slip as illustrated in Figure 25 [35]. It enables to characterize, even using a damage
 444 model, both the position and the opening of “equivalent cracks” using an additional post-processing
 445 step on the relative displacement [35]. From these results, it is shown that the increase in the
 446 confinement pressure is responsible for an increase in the number of cracks. For $P_{\text{conf}} = 20$ MPa, the
 447 distribution of the stress along the rebar (between $x = 0.6$ and $x = 0.8$ m) is disrupted by the
 448 continuous damage localization.

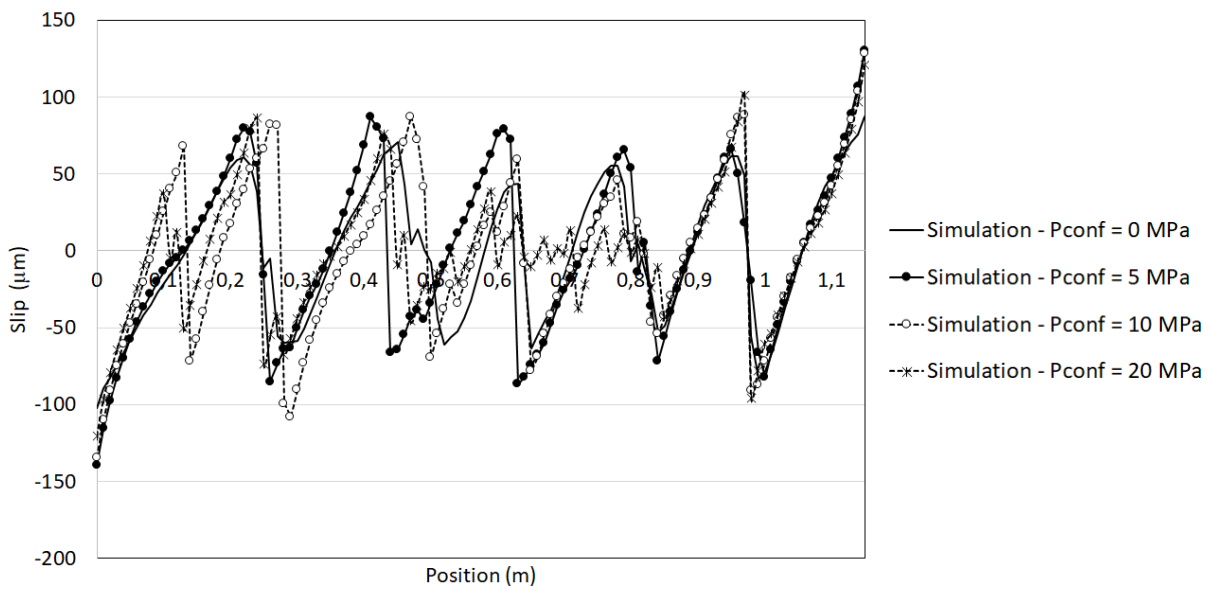
449 From the slip distribution, the number of cracks and average crack spacing are finally computed and
450 summarized in Table 9. As expected from the steel stress distribution, the higher the confinement,
451 the higher the number of cracks. This is essentially due to the stiffening of the bond with the increase
452 in the confinement pressure, which decreases the transfer length and enables the apparition of
453 closer “cracks” (decrease in crack spacing).
454



455 Figure 23. Damage distributions at the end of the loading for different applied confinement pressures.



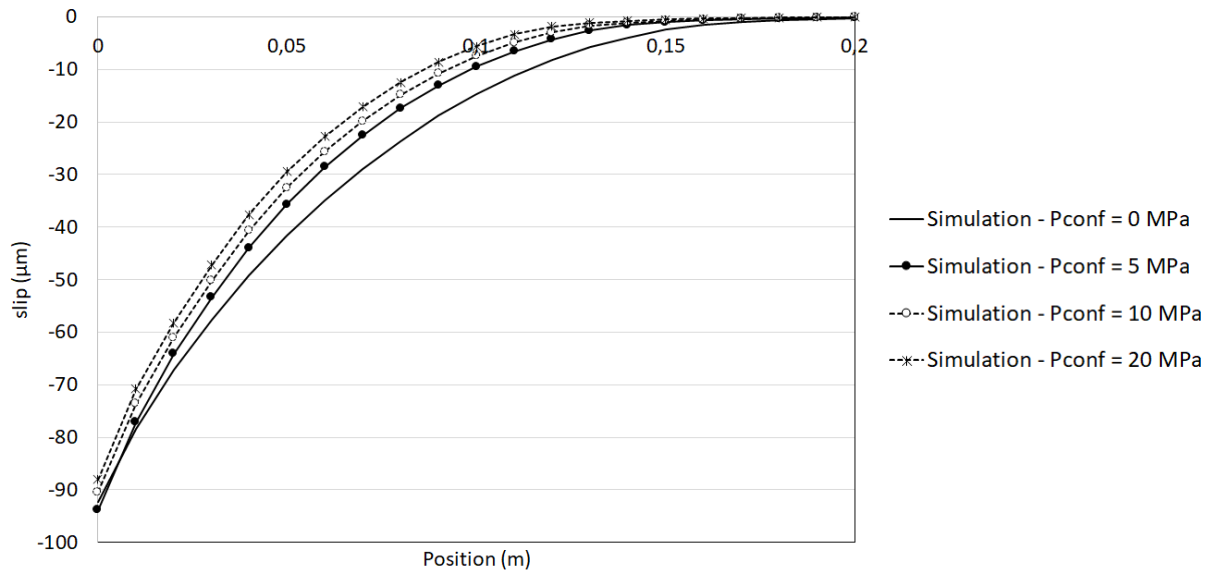
456 Figure 24. Evolution of the stress along the steel bar for different confinement pressures – Last displacement
 457 increment.



458 Figure 25. Evolution of the slip along the steel bar for different confinement pressures. Last displacement
 459 increment.

	$P_{\text{conf}} = 0 \text{ MPa}$	$P_{\text{conf}} = 5 \text{ MPa}$	$P_{\text{conf}} = 10 \text{ MPa}$	$P_{\text{conf}} = 20 \text{ MPa}$
Number of cracks	5	5	6	~6*
Average spacing (mm)	183	182	168	170
Minimum spacing (mm)	150	160	140	100
Maximum spacing (mm)	240	190	240	250

460 Table 9. Number of cracks and average crack spacing for different P_{conf} . *due to the damage localization along
 461 the steel bar, the crack position is estimated from the local maximum in the steel stress distribution.



462 Figure 26. Evolution of the slip along the steel bar in the elastic phase for different confinement pressures.

Confinement pressure (MPa)	0	5	10	20
Transfer length (cm)	14.1	12.8	11.5	10.7

463 Table 10. Computed transfer length.

464 This effect is also underlined in Figure 26, which provides the evolution of the slip along the rebar
 465 during the elastic regime. From this evolution, it is possible to evaluate the corresponding transfer
 466 lengths, considering that the transfer length is the length to reach a 5μm slip from the end of the bar.
 467 The computed values are summarized in Table 10, which confirms the decrease in the transfer length
 468 with an increasing confinement pressure.

469 5. Conclusions

470 In this contribution, a numerical model to consider the effect of the stress state on the bond
 471 behavior between steel and concrete was proposed. It is based on a zero thickness element, adapted
 472 to large-scale simulations and the use of 1D elements for steel bars. It supposes the definition of a
 473 bond stress – slip law that includes the confining pressure around the steel bar as a parameter. The
 474 implementation of the model was presented and the calibration of the bond law was discussed. A
 475 general equation was especially proposed. Based on experimental observations from the state-of-
 476 the-art, it includes a different behavior in confining tension or compression and two different phases
 477 depending on the concrete cover to steel diameter ratio.

478 The evolution law was validated through the comparison to 28 experimental pullout tests. The model
 479 was able to reproduce the evolution of the bond stress (especially the bond strength) as a function of
 480 the applied confinement pressure, whatever the configuration. A methodology was proposed,
 481 especially to take into account the non-zero lateral pressure, even in the case of a zero applied
 482 confinement pressure (structural effect due to pullout configuration). Considering the full range of
 483 simulated pullouts, the maximum mean error between the simulated and experimental bond
 484 strength reaches 15 %, which is reasonable regarding the wide variability in the experimental results.

485 Finally, the consequences at the structural level were investigated on a reinforced concrete tie. The
 486 response for different levels of confining pressures was studied. It shows the capability of the model

487 to reproduce the “expected” tendencies with especially an increase of the initial elastic stiffness with
488 increasing pressures and consequently a higher number of cracks in the stabilized nonlinear regime.
489 The “transfer length” was also shown to decrease with increasing confining pressures.

490 As a conclusion, the proposed model is able to simulate the mechanical behavior of reinforced
491 concrete at the structural scale, including the confinement effect. However, additional experimental
492 validations are necessary to further evaluate the correctness and the validity of the proposed slip
493 model. It would suppose the proposition of adapted experimental campaigns to focus on the
494 confinement effect. The applicability of the model to more full-scale applications (including for
495 example the applications studied in [50], [51]) could then be demonstrated.

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