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ACCEPTED MANUSCRIPT

## Low temperature epitaxial growth of GaP on Si by atomic-layer deposition with plasma activation

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# Low temperature epitaxial growth of GaP on Si by atomic-layer deposition with plasma activation

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*Keywords: atomic layer deposition, gallium phosphide, plasma*

## Abstract

An approach for epitaxial growth of GaP layers on Si substrates at low temperature (380°C) by plasma-enhanced atomic layer deposition (PEALD) is explored. A significant improvement of the crystalline properties of the GaP layers is obtained using additional in-situ Ar plasma treatment. The epitaxial growth for the first 20-30 nm of GaP on Si is demonstrated from transmission electron microscopy. Moreover, the use of in-situ Ar plasma treatment during the PEALD process allows one to increase the growth rate per cycle from  $0.9 \pm 0.1$  Å/cycle to  $1.9 \pm 0.1$  Å/cycle and reduce the RMS roughness from 3.76 nm to 1.88 nm. The effect of Ar plasma treatment on the electronic properties of the GaP/Si interface is studied by deep level transient spectroscopy (DLTS). A defect level at  $(0.33 \pm 0.03)$  eV below the conduction band is observed in the subsurface layer of Si for the GaP/Si structure grown under Ar plasma treatment. However, the defect response observed by DLTS vanishes after rapid thermal annealing at 500 °C in nitrogen ambient.

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## 1. Introduction

Gallium phosphide (GaP) epitaxially grown on Si is of great interest for the integration of III-V and Si technologies. GaP being a III-V semiconductor with a band gap of 2.26 eV and having a lattice mismatch with Si of less than 0.4% may be used as a nucleation layer for further growth of III-V compounds or as a window/emitter layer for silicon based solar cells [1-3].

Epitaxial growth of GaP layers on Si substrates may be achieved by metal-organic chemical vapor epitaxy (MOVPE) [4] or molecular beam epitaxy (MBE) [5]. Several important problems should be solved for the growth of polar III-V materials on non-polar Si substrates. The nucleation phase of III-V materials may give rise to bonds between two group-III or two group-V atoms. These bonds between similar atoms act as a boundary between two single crystalline domains opposite in phase (antiphase boundaries), which could propagate along the direction of growth. To improve the crystalline properties of the GaP layer the time modulated procedure was proposed recently, where deposition of a Ga monolayer is followed by that of a P one for the GaP nucleation stage. The terms atomic layer epitaxy (ALE) for MOVPE [6-8] and mobility enhanced epitaxy (MEE) for MBE [8] were used to describe this procedure, which could significantly enhance the quality of the epitaxy. However, high temperatures, above 600 °C, used for Si surface deoxidation and reconstruction as well as for the growth process lead to a significant decrease in the lifetime of minority charge carriers in silicon substrates, which negatively affects the resulting efficiency of solar cells [9, 10].

We propose to use a plasma deposition technique to provide epitaxial growth of III-V compounds on Si at low temperature. This is possible due to the fact that plasma energy is locally scattered by surface atoms and can enhance the rate of surface reactions and the surface diffusion of adatoms [11]. To gain the advantage of the time modulated process (ALE or MEE) developed for nucleation of the epitaxial layers the atomic layer deposition (ALD) approach is suggested to be used in combination with plasma stimulation. Plasma-enhanced atomic layer deposition (PEALD) has several advantages such as high-quality uniform deposition on large

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3 areas, conformal layer growth on textured surfaces, trenches and holes with a high aspect ratio.  
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5 Moreover, PEALD can be used to grow thin layers of material over large areas with high  
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7 throughput, which is one of the main advantages for photovoltaic applications.  
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10 The possibility to obtain microcrystalline GaP by a PEALD-like method with a  
11 continuous hydrogen plasma was previously shown [12]. However, hydrogen plasma of high  
12 power used in those processes to increase the crystallinity of GaP leads to damages in the Si  
13 subsurface layer, which results in a drop of the photovoltaic performance [13]. Thus, a way to  
14 provide epitaxial growth of GaP at low temperature without deteriorating the electronic  
15 properties of the silicon substrate is an important issue. Recently, the use of in-situ annealing  
16 in argon plasma was shown to improve the crystalline quality to the epitaxial level of aluminum  
17 nitride layers obtained by PEALD [14]. Here we explore the possibility to obtain thin epitaxial  
18 GaP layers using PEALD at temperatures below 400 °C using in-situ annealing in Ar plasma.  
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## 33 **2. Experiment details**

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36 The process of ALD growth of GaP was realized using the Oxford Plasmalab 100  
37 PECVD setup with a capacitively coupled plasma reactor. Trimethylgallium (TMG) and  
38 phosphine (PH<sub>3</sub>) were used as the precursors of gallium and phosphorus, respectively.  
39 Hydrogen (H<sub>2</sub>) was used as a carrier gas for TMG providing a precursor flow into the chamber  
40 using a bubbling system. The growth surface was alternatively submitted to TMG and PH<sub>3</sub>  
41 flows. To provide PH<sub>3</sub> decomposition at low temperature an RF (13.56 MHz) plasma was used  
42 at the phosphorous deposition step. After exposure to each precursor, the chamber was purged  
43 with argon to avoid mixing precursors and parasitic CVD growth. Additionally, a modified  
44 process was also explored, where an Ar plasma treatment step was introduced immediately  
45 before the TMG step. The main deposition parameters are presented in Table 1.  
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**Table 1.** Process parameters of GaP deposition by PEALD

Parameter\Step	Gallium	Purge	Phosphorus	Purge	Ar plasma
Gas composition	5% TMG/H <sub>2</sub>	Ar	PH <sub>3</sub> /H <sub>2</sub>	Ar	Ar
Time, s	5	10	3	10	0-15
Dose, nanomoles/cm <sup>2</sup>	5-50	-	0.5-3.7	-	-
Pressure, mTorr	350	0	350	0	350
RF power density, mW/cm <sup>2</sup>	0	0	90	0	90

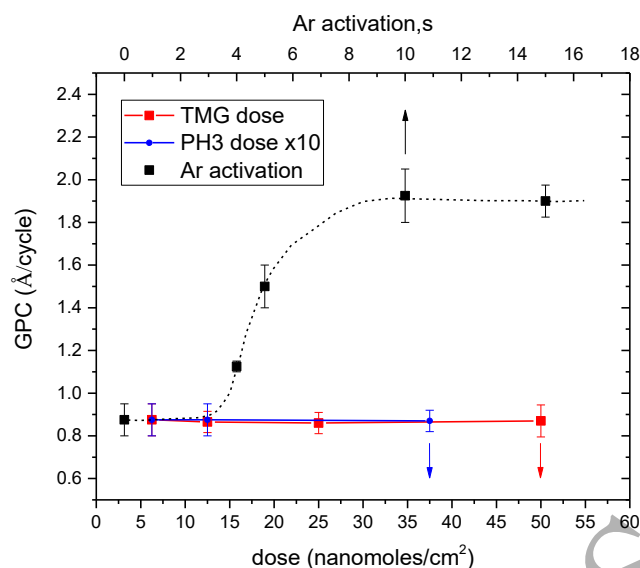
GaP layers were grown on Si (100) substrates with a 4° misorientation toward [110] as conventionally used for epitaxy [15]. Si substrates were cleaned by the *Shiraki* method followed by an HF-dip removal of oxide prior to the GaP growth [16]. The layers were grown at the temperature of 380 °C. The choice of the process temperature is determined by the requirement for the evaporation of excess phosphorus from the growth surface, which occurs at temperatures above 300 °C. In this case the growth of only one monolayer of phosphorous atoms, which are bonded to Ga, is expected during one cycle. On the other hand, TMG starts to thermally decompose at about 280°C and keeps at least one methyl group up to 480°C [17]. The deposition rate was estimated by determining the resulting layer thickness using laser ellipsometry at a wavelength of 632.8 nm (Horiba PZ2000) and step profilometry (Ambios XP) on selective chemical etching steps of GaP on the silicon surface. The structural properties of the samples were investigated by transmission electron microscopy (TEM) and electron diffraction (JEOL JEM 2100F) with 200 kV acceleration voltage. Capacitance deep-level transient spectroscopy (DLTS) was applied to study the influence of plasma treatment and in particular the possible creation of defects in the subsurface region of the silicon wafer to explore possible defects in the silicon near to the GaP/n-Si interface. For this characterization technique, special structures with Schottky barriers were fabricated by vacuum evaporation of gold on the top of GaP layers through a mask with circular holes (diameter of 1 mm). DLTS measurements were carried out

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3 using a Boonton 7200 capacitance meter and a Janis VPF 100 liquid nitrogen cryostat in the  
4 temperature range of 80–360 K at the following conditions:  $V_{\text{init}} = -2$  V,  $V_{\text{pulse}} = +2$  V,  $t_{\text{pulse}} = 50$   
5 ms. The time rate window was varied in the range of 10–1000 s<sup>-1</sup>.  
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### 10 11 12 **3. Results and discussions**

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14 For the ALD process without Ar plasma activation two series of 200 cycles were  
15 performed. In the first series the TMG dose was varied from 5 to 50 nmol/cm<sup>2</sup> at a constant PH<sub>3</sub>  
16 dose of 1.25 nmol/cm<sup>2</sup>. During the second one the PH<sub>3</sub> dose was varied from 0.5 to 3.7  
17 nmol/cm<sup>2</sup> at a constant TMG dose of 50 nmol/cm<sup>2</sup>. For both series the growth rate of GaP does  
18 not depend on the dose of TMG and PH<sub>3</sub> and is equal to 0.9±0.1 Å/cycle (Figure 1) being more  
19 than 2.5 times lower compared to one monolayer per cycle. The saturation of the growth rate,  
20 which is observed in the whole range of TMG and PH<sub>3</sub> dose, means that a self-limitation occurs  
21 for both precursors, indicating that the ALD growth mode was achieved.  
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33 Next the effect of Ar plasma activation before the TMG step in each cycle on the deposition  
34 rate per cycle was investigated. The treatment time in Ar plasma was varied in the range of 0–  
35 15 s at constant doses of TMG and PH<sub>3</sub> of 12.5 nmol/cm<sup>2</sup> and 1.25 nmol/cm<sup>2</sup>, respectively. The  
36 dependence of the growth rate on the duration of the argon plasma step is presented in Figure  
37 1. First an increase of growth rate per cycle (GPC) with the argon plasma step duration ( $t_{\text{Ar}}$ ) is  
38 observed. Then for  $t_{\text{Ar}} \geq 10$  s the GPC reaches saturation at the value of 1.9±0.1 Å/cycle, which  
39 does not exceed one monolayer per cycle (2.73 Å/cycle), being in compliance with the ALD  
40 regime. However, introduction of an additional Ar plasma step (200 W, 15 s) before the PH<sub>3</sub>  
41 step does not lead to an increase of the growth rate. This is probably due to the fact that PH<sub>3</sub>  
42 decomposition occurs at high RF plasma power (200 W), which provides significant energy to  
43 phosphorous atoms during this step. An influence of Ar plasma treatment at the previous step  
44 seems to be less important for phosphorous deposition.  
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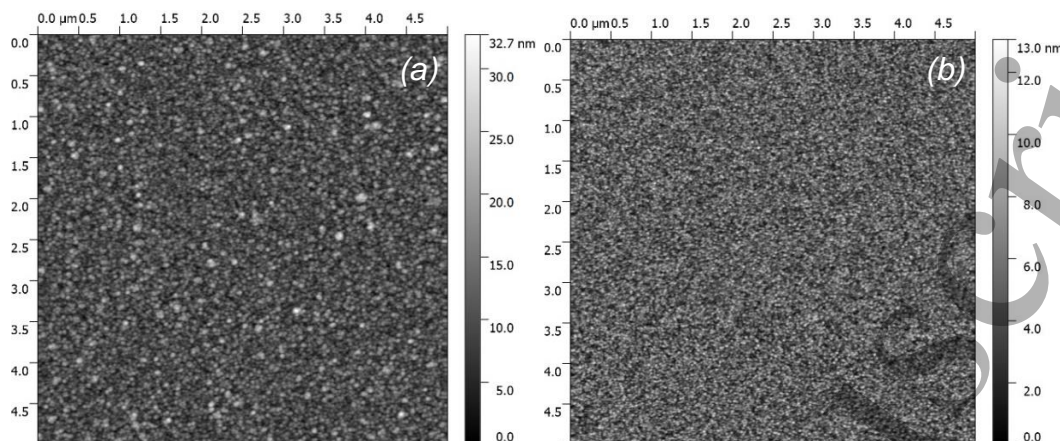
**Figure 1.** Growth rate per cycle (GPC) as a function of TMG and PH<sub>3</sub> doses and dependence on Ar plasma step duration.

Along with the increase of the deposition rate due to the Ar plasma treatment, a significant decrease of the surface roughness of the layers was observed by atomic force microscopy (AFM) (Figure 2). AFM measurements were first performed for samples with the same thickness. Thus, 400 cycles of PEALD without Ar plasma and 200 cycles with Ar plasma step duration of 15 s were used to obtain two GaP layers of about 40 nm. The RMS roughness decreases from  $3.76 \pm 0.1$  nm for the layer grown without Ar plasma to  $1.88 \pm 0.1$  nm for the sample prepared with Ar plasma. However, the roughness of the deposited films depends on the thickness. The GaP layer deposited without Ar plasma during 200 cycles (thickness of about 20 nm) exhibits RMS roughness of  $1.55 \pm 0.1$  nm. For very thin layers, both modes produce similar smooth surface. Indeed, for samples processed using 20 cycles with Ar plasma treatment and samples processed using 40 cycles without Ar plasma, which both have a thickness of about 4 nm, the RMS roughness was equally found in the range of  $0.12 \pm 0.1$  nm indicating smooth surface for initial growth.

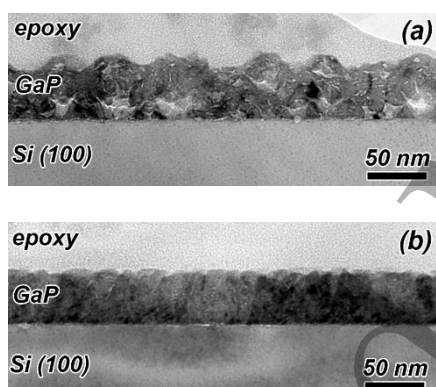
Obviously, the Ar plasma treatment affects the structure of the GaP layers. Bright field TEM images presented in Figure 3 confirm the improved RMS roughness when the Ar plasma



step is introduced in the PEALD process. Moreover, a strong difference in the structure of the GaP films can also be observed. Indeed, the layer obtained with the additional Ar plasma step is much denser and uniform compared to the one fabricated without Ar plasma treatment.



**Figure 2.** AFM images of GaP layers grown on Si by PE-ALD without and with 15 s Ar plasma treatment step.

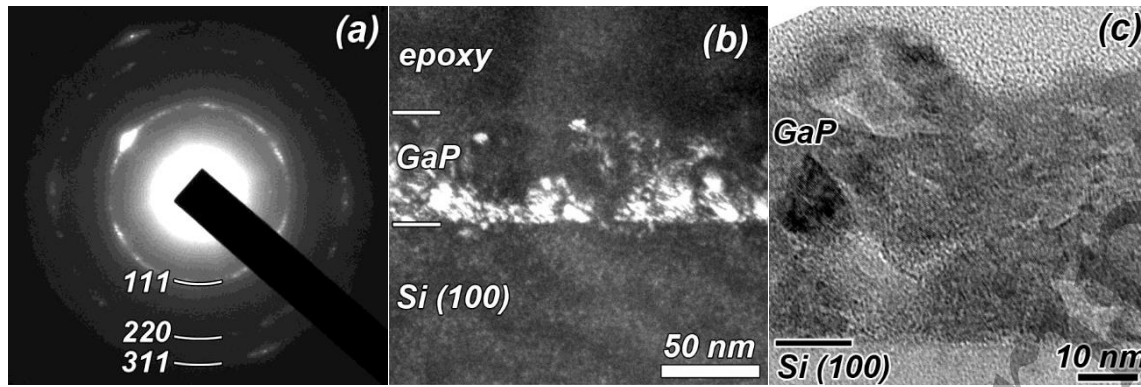


**Figure 3.** Cross section bright field TEM image of GaP on Si grown by PEALD without Ar plasma step (a) and with 15 s Ar plasma step (b).

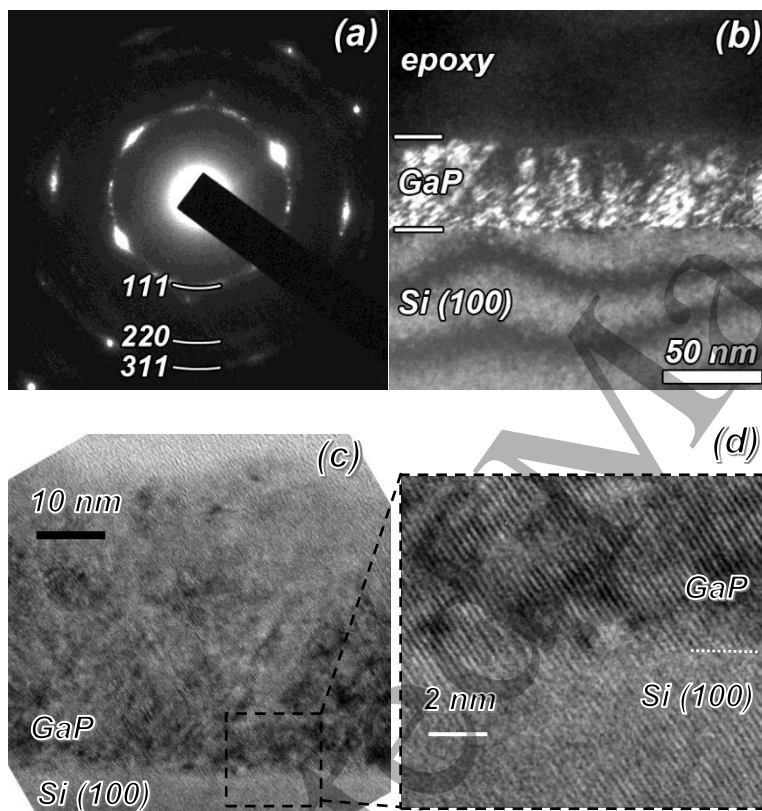
Detailed studies of structural properties were performed using selective area electron diffraction (SAED) and high-resolution transmission electron microscopy (HRTEM). The SAED pattern is a superposition of the diffraction from a small part of single-crystal Si substrate, GaP layer and a small part of amorphous epoxy glue. The diffraction from the monocrystal Si substrate is a spot grid, while if the GaP layer is microcrystalline, the diffraction looks as a set of solid rings. The more the rings look like spots, the closer the GaP structure is to a single crystal.

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3 For the GaP layer obtained without Ar plasma the SAED pattern consists of almost solid  
4 circles, which indicates that the layer is microcrystalline with crystallites of different  
5 orientations (Figure 4a). The Dark field (DF) image obtained in part of  $g = (111)$  ring indicates  
6 that the GaP crystallites of this orientation (bright contrast) are concentrated at the GaP/Si  
7 interface (Figure 2b), which corresponds to initial epitaxial growth on the Si substrate.  
8 However, DF images taken from the other parts of the (111) plane (not presented) show separate  
9 GaP crystallites located in the middle-upper part of the GaP layer and having a different  
10 orientation. The presence of 10 nm size GaP crystallites with various orientations could also be  
11 observed in the HRTEM image presented in Figure 4c. A GaP epitaxial layer of a few  
12 nanometers could be observed by HRTEM at the GaP/Si interface, while further layer growth  
13 leads to deviations in the crystal orientation. Moreover, the thickness of the GaP epitaxial layer  
14 is nonuniform along GaP/Si interface.  
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30 A detailed TEM study for the GaP/Si structures obtained by PEALD with the Ar plasma  
31 step demonstrates that the GaP layer is rather epitaxial (Figure 5). The diffraction pattern of  
32 epitaxial GaP layer, which exhibits marked spots, is clearly visible in the SAED pattern of the  
33 layer (Figure 5a). According to the dark-field image obtained in the part of  $g = (111)$  GaP-ring  
34 epitaxial growth of GaP is observed within the whole layer (Figure 5b). However, crystallites  
35 with different orientations could be observed 20-30 nm far from the GaP/Si interface, which is  
36 also confirmed by the rings in the diffraction pattern. The HRTEM image (Figure 5c) confirms  
37 epitaxial growth of GaP layer on the Si substrate. Substrate orientation is clearly reproduced by  
38 subsequent GaP growth (Figure 5d). Despite the high concentration of lattice defects like twins,  
39 misfits and threading dislocations, which could be observed by TEM in the epitaxial layer, a  
40 significant improvement of the crystalline properties was achieved using the in-situ Ar plasma  
41 treatment.  
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**Figure 4.** SAED pattern (a), DF image obtained in the part of  $g = (111)$  ring (b), and HRTEM image (c) of GaP layer grown by PE-ALD without Ar plasma.



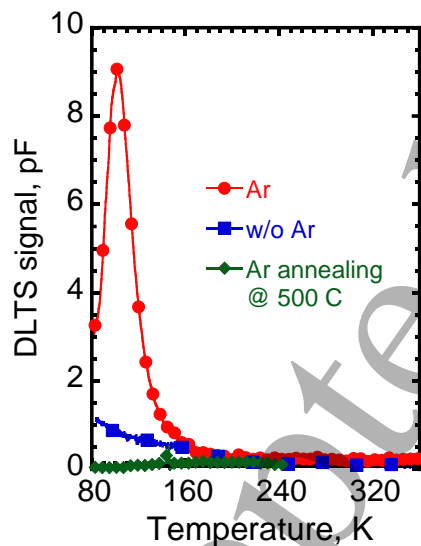
**Figure 5.** SAED pattern (a), DF image obtained in the part of  $g = (111)$  ring (b), and HRTEM images (c,d) of GaP layer grown by PE-ALD without Ar plasma.

The Ar plasma treatment introduced immediately before the TMG decomposition step influences both the growth rate and the structure of the layer. Both properties can be affected by the ion bombardment during Ar plasma treatment. Recently, the importance of substrate bombardment ion energy and flux on the film properties was demonstrated for the PEALD

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3 process [18]. The influence of the energy and dose of ions controlled by the substrate bias on  
4 the growth rate as well as on the crystalline properties was shown. In another work performed  
5 in the same Plasmalab setup as used here, it was shown that both Ar ion energy and flux increase  
6 with the RF power [19]. Energy of this ion flux bombardment during the Ar plasma treatment  
7 provides local heating of the growing surface as was proposed in Ref.14. Indeed, an increase of  
8 the surface temperature could lead to more effective TMG decomposition. According to ref. 17  
9 in the temperature range from 380°C to 480°C TMG starts to lose the second methyl group.  
10 Methyl groups of TMG provide steric hindrance of the growth rate, which is a main reason why  
11 in practice the GPC is normally lower than one monolayer per cycle [20]. The observed  
12 saturation of the GPC with the increase of Ar plasma treatment duration (Figure 1a) could  
13 indicate that a threshold temperature is reached when TMG loses the second methyl group. In  
14 addition, the local heating of the surface leads to enhanced surface migration of adatoms and,  
15 therefore, to the improvement of the structural properties as was observed for the growth of  
16 AlN by PEALD with an Ar plasma treatment [14].

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35 However, any plasma treatment is known to be a potential source of radiation-induced defects,  
36 which in our case could be formed in the Si substrate. Previously, for the GaP growth performed  
37 with high power hydrogen plasma treatment, a damage of the Si substrate was observed by  
38 TEM and defects were detected at the GaP/Si interface by electrical measurements [13]. Here  
39 no damage of Si is observed by TEM (Figure 1b) at the GaP/Si interface, despite the short (3 s)  
40 treatment in H<sub>2</sub> plasma during the PH<sub>3</sub>/H<sub>2</sub> step and the long (15 s) Ar plasma treatment.  
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However, to have a better assessment of the electronic properties of the GaP/Si interface, deep  
level transient spectroscopy (DLTS) measurements were performed on GaP/n-Si structures  
with a top Schottky barrier. DLTS spectra obtained for a rate window of 50 s<sup>-1</sup> are presented in  
Figure 6. For the GaP/Si structure grown without Ar plasma treatment no responses are  
detected, while a clear defect response was observed in the range of 80–140 K for the sample  
fabricated with additional Ar plasma. Using the classical DLTS data treatment (see

supplementary material) we could identify the defect level to be at  $0.33 \text{ eV} \pm 0.03 \text{ eV}$  below the conduction band. Ar plasma is known to be a source of radiation-induced defects in Si during magnetron sputtering of ITO on a-Si:H/c-Si heterostructures due to hard UV radiation. However, thermal annealing at a temperature approximately one hundred degrees higher than the deposition temperature allows one to recover the carrier lifetime in the Si wafer [21]. Our GaP/Si structure grown with the Ar plasma treatment step was annealed at  $500 \text{ }^\circ\text{C}$  during 1 minute in nitrogen ambient. To verify that annealing of GaP layer does not affect the structural properties they were additionally monitored by Raman spectroscopy. No changes were observed for microcrystalline GaP films after annealing at the temperature of  $500^\circ\text{C}$  (see supplementary material). After annealing no defect response could be detected by DLTS (Figure 6). Thus, argon plasma leads to the formation of defects in the sub-surface region of the silicon wafer, however their concentration could be drastically reduced by rapid thermal annealing at  $500 \text{ }^\circ\text{C}$ .



**Figure 6.** DLTS spectra (rate window of  $50 \text{ s}^{-1}$ ) for GaP layers grown with or without Ar plasma step, and also for a layer grown with the Ar plasma step and annealed at  $500 \text{ }^\circ\text{C}$  during 1 minute in nitrogen ambient.

## 5. Conclusion

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3 In conclusion, GaP epitaxial layers on silicon substrates were obtained by PEALD at a  
4 temperature of 380 °C. It was shown that introducing an additional Ar plasma step during  
5 growth increases the growth rate per cycle and reduces the surface roughness due to an increase  
6 in the surface migration of adsorbed atoms. This effect has a saturation, which also indicates  
7 the compliance with the ALD regime and the absence of uncontrollable volumetric chemical  
8 reaction or re-evaporation of material from the chamber walls in the Ar plasma. The Ar plasma  
9 step also allows one to significantly improve the quality of epitaxial layers of GaP on Si  
10 substrates grown by PEALD. Finally, it leads to the formation of defects detected by DLTS in  
11 Si close to the GaP/Si interface, however these defects are suppressed by rapid thermal  
12 annealing at 500 °C.  
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### 38 **References**

- 39  
40 [1] *III-V Compound Semiconductors: Integration with Silicon-Based Microelectronics*, (Eds:  
41 T. Li, M. Mastro, A. Dadgar), CRC Press **2016**.  
42  
43 [2] C. Zhang, N.N. Faleev, L. Ding, M. Boccard, M. Bertoni, Z. Holman, R.R. King, C.B.  
44 Honsberg, in *2016 IEEE 43rd Photovolt. Spec. Conf.*, 1950–1953, IEEE **2016**.  
45  
46 [3] M. Feifel, J. Ohlmann, J. Benick, T. Rachow, S. Janz, M. Hermle, F. Dimroth, J. Belz, A.  
47 Beyer, K. Volz, D. Lackner, *IEEE J. Photovoltaics*, **2017**, 7, 502.  
48  
49 [4] J.M. Olson, M.M. Al-Jassim, A. Kibbler, K.M. Jones, *J. Cryst. Growth*, **1986**, 77, 515.  
50  
51 [5] Mahdad Sadeghi, Shumin Wang, *J. Cryst. Growth*, **2001**, 227–228, 279.  
52  
53  
54 [6] Y. Sakuma, K. Kodama, M. Ozeki, *Appl. Phys. Lett.* **1990**, 56, 827.  
55  
56  
57  
58  
59  
60

- [7] J.R. Gong, S. Nakamura, M. Leonard, S.M. Bedair, N.A. El-Masry, *J. Electron. Mater.* **1992**, *21*, 965.
- [8] T. Tsuji, H. Yonezu, M. Yokozeki, Y. Takagi, Y. Fujimoto, N. Ohshima, *Jpn. J. Appl. Phys. A*, **1997**, *36*, 5431.
- [9] R. Varache, M. Darnon, M. Descazeaux, M. Martin, T. Baron, D. Muñoz, *Energy Procedia*, **2015**, *77*, 493.
- [10] L. Ding, C. Zhang, T. U. Nærland, N. Faleev, C. Honsberg, M. I. Bertoni, *Energy Procedia* **2016**, *92*, 617.
- [11] A.J.M. Mackus, S.B.S. Heil, E. Langereis, H.C.M. Knoop, M.C.M. van de Sanden, W.M.M. Kessels, *J. Vac. Sci. Technol. A Vacuum*, **2010**, *28*, 77.
- [12] A.S. Gudovskikh, A. V. Uvarov, I.A. Morozov, A.I. Baranov, D.A. Kudryashov, K.S. Zelentsov, A.S. Bukatin, K.P. Kotlyar, *J. Vac. Sci. Technol. A*, **2018**, *36*, 02D408.
- [13] A. S. Gudovskikh, A. V. Uvarov, I. A. Morozov, A. I. Baranov, D. A. Kudryashov, K.S. Zelentsov, A. Jaffre, S. Le Gall, A. Darga, A. Brezard-Oudot, J.P. Kleider, *Phys. Status Solidi A*, **2019**, *216*, 1800617.
- [14] H.-Y. Shih, W.-H. Lee, W.-C. Kao, Y.-C. Chuang, R.-M. Lin, H.-C. Lin, M. Shiojiri, M.-J. Chen, *Sci. Rep.* **2017**, *7*, 39717.
- [15] T.J. Grassman, M.R. Brenner, S. Rajagopalan, R. Unocic, R. Dehoff, M. Mills, H. Fraser, S.A. Ringel, *Appl. Phys. Lett.* **2009**, *94*, 232106.
- [16] A. Ishizaka, Y. Shiraki, *J. Electrochem. Soc.* **1986**, *133*, 666.
- [17] F. Lee, T.R. Gow, R.I. Masel, *J. Electrochem. Soc.* **1989**, *136*, 2640.
- [18] T. Faraz, H. C. M. Knoop, M. A. Verheijen, C. A. A. van Helvoirt, S. Karwal, A. Sharm Beladiya, A. Szeghalmi, D. M. Hausmann, Jon Henri, Mariadriana Creatore, and W. M. M. Kessels *ACS Applied Materials & Interfaces* **2018**, *10*, 13158.
- [19] D. Gahan, S. Daniels, C. Hayden, P. Scullin, D. O'Sullivan, Y.T. Pei, M.B. Hopkins, *Plasma Sources Science and Technology* **2012**, *21*, 024004.

1  
2  
3 [20] *Chemical Vapour Deposition. Precursors, Processes and Applications* (Eds: A.C. Jones,  
4 M.L. Hitchman). Royal Society of Chemistry, Cambridge **2008**.

5  
6  
7 [21] B. Demarex, St. De Wolf, A. Descoedres, Z. Ch. Holman, Ch.e Ballif, *Appl. Phys. Lett.*  
8 **2012**, *101*, 171604.  
9  
10  
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