



Duality and bicrystals on infinite binary matrices

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DUALITY AND BICRYSTALS ON INFINITE BINARY MATRICES

THOMAS GERBER AND CÉDRIC LECOUEVEY

ABSTRACT. The set of finite binary matrices of a given size is known to carry a finite type A bicrystal structure. We first review this classical construction, explain how it yields a short proof of the equality between Kostka polynomials and one-dimensional sums together with a natural generalisation of the $2M - X$ Pitman transform. Next, we show that, once the relevant formalism on families of infinite binary matrices is introduced, this is a particular case of a much more general phenomenon. Each such family of matrices is proved to be endowed with Kac-Moody bicrystal and tricrystal structures defined from the classical root systems. Moreover, we give an explicit decomposition of these multicrystals, reminiscent of the decomposition of characters yielding the Cauchy identities.

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1. INTRODUCTION

Crystals are fundamental combinatorial objects introduced by Kashiwara and Littelmann at the beginning of the nineties. They can be interpreted as the combinatorial skeletons of certain modules for complex Lie algebras and their infinite-dimensional analogues: the Kac-Moody algebras. More precisely, crystals are oriented graphs with colored arrows defined from renormalisations (called the Kashiwara operators) of the Chevalley generators of the Kac-Moody algebra considered. In Kashiwara's approach, they are obtained via "crystallisation" at $q = 0$ of representations of the corresponding quantum group, a q -deformation of the Kac-Moody algebra introduced by Jimbo. Since its introduction, crystal theory has revealed numerous fruitful interactions with modern particle physics theory and integrable systems. We refer the reader to [3] and the references therein for a recent exposition.

The present paper is concerned with generalisations of crystals where multimodule structures are considered rather than just ordinary module structures. This means that we consider analogues of crystals for complex vector spaces endowed with commuting actions of several Lie algebras (or Kac-Moody algebras). In fact, most of the combinatorial structures (crystals, Fock spaces, one-dimensional sums, Pitman transforms) that we shall consider in the sequel were defined to solve problems connected to theoretical physics. We expect similar interactions with the results we establish here.

The prototypical example of an interesting bicrystal is obtained by starting from the $\mathfrak{gl}_n \times \mathfrak{gl}_\ell$ -module $\mathbb{C}^n \otimes \mathbb{C}^\ell$ and considering the associated symmetric and antisymmetric $(\mathfrak{gl}_n \times \mathfrak{gl}_\ell)$ -modules $S(\mathbb{C}^n \otimes \mathbb{C}^\ell)$ and $\Lambda(\mathbb{C}^n \otimes \mathbb{C}^\ell)$. Using two sets of indeterminates $\{x_1, \dots, x_n\}$ and $\{y_1, \dots, y_\ell\}$, one can check that their characters are given by the formulas

$$\text{char } S(\mathbb{C}^n \otimes \mathbb{C}^\ell) = \prod_{\substack{1 \leq i \leq n \\ 1 \leq j \leq \ell}} \frac{1}{1 - x_i y_j} \quad \text{and} \quad \text{char } \Lambda(\mathbb{C}^n \otimes \mathbb{C}^\ell) = \prod_{\substack{1 \leq i \leq n \\ 1 \leq j \leq \ell}} (1 + x_i y_j).$$

The classical Cauchy identities then gives the decomposition of each characters in terms of the Schur functions. More precisely, recall that the irreducible finite-dimensional \mathfrak{gl}_n -modules are parametrised by partitions of length n (that is, nonincreasing sequences $\lambda = (\lambda_1, \dots, \lambda_n)$ of nonnegative integers) and the character of the irreducible module parametrised by λ is the Schur symmetric polynomial $s_\lambda(x)$ in the variables x_1, \dots, x_n . We then have

$$\begin{aligned} \text{char } S(\mathbb{C}^n \otimes \mathbb{C}^\ell) &= \prod_{\substack{1 \leq i \leq n \\ 1 \leq j \leq \ell}} \frac{1}{1 - x_i y_j} = \sum_{\substack{\lambda \text{ partition of} \\ \text{length } \max(n, \ell)}} s_\lambda(x) s_\lambda(y) \quad \text{and} \\ \text{char } \Lambda(\mathbb{C}^n \otimes \mathbb{C}^\ell) &= \prod_{\substack{1 \leq i \leq n \\ 1 \leq j \leq \ell}} (1 + x_i y_j) = \sum_{\substack{\lambda \text{ contained in the} \\ \text{rectangle } n \times \ell}} s_\lambda(x) s_{\lambda^{\text{tr}}}(y) \end{aligned}$$

where λ^{tr} is the transpose of the partition λ . Both identities admit an elegant combinatorial proof based on the Robinson-Schensted-Knuth (RSK) correspondence, see [6]. The idea is first to observe that

$$\prod_{\substack{1 \leq i \leq n \\ 1 \leq j \leq \ell}} \frac{1}{1 - x_i y_j} = \sum_{(m_{ij}) \in \mathcal{N}} \prod_{\substack{1 \leq i \leq n \\ 1 \leq j \leq \ell}} (x_i y_j)^{m_{ij}}$$

where \mathcal{N} is the set of $n \times \ell$ matrices with nonnegative integer entries. Next, the RSK correspondence yields a bijection between \mathcal{N} and the set of pairs of semistandard tableaux with the same partition shape. The identity then follows from the fact that s_λ is the generating function of the set of semistandard tableaux with shape λ for the evaluation map on tableaux. Similarly, one has

$$\prod_{\substack{1 \leq i \leq n \\ 1 \leq j \leq \ell}} (1 + x_i y_j) = \sum_{(m_{ij}) \in \mathcal{M}} \prod_{\substack{1 \leq i \leq n \\ 1 \leq j \leq \ell}} (x_i y_j)^{m_{ij}}$$

where \mathcal{M} is the set of binary $n \times \ell$ matrices, and the second Cauchy identity comes from an adapted version of the RSK correspondence. This construction admits numerous extensions, notably based on generalisations of the previous Cauchy identities due to Littlewood, and involving the characters of simple modules corresponding to the orthogonal or symplectic Lie algebras. We refer the reader to [7] and the references therein for a more detailed presentation, and for a generalisation to the Demazure characters. It is also worth mentioning that the previous constructions have interesting applications not only in combinatorics and algebra but also in problems related to physics and percolation models [2].

We now turn out to the main topic of this article. One can define two commuting families of crystal operators (one for \mathfrak{gl}_n and another for \mathfrak{gl}_ℓ) directly on the set of matrices \mathcal{N} and \mathcal{M} yielding the structure of a $(\mathfrak{gl}_n \times \mathfrak{gl}_\ell)$ -crystal, or type $(A_{n-1} \times A_{\ell-1})$ -crystal. This approach, explained in detail in [43], see also [15, Chapter 11], permits to overcome the RSK construction. One gets the decomposition

$$\mathcal{N} = \bigoplus_{\substack{\lambda \text{ partition of} \\ \text{length } \max(n, \ell)}} B(\lambda) \otimes \dot{B}(\lambda) \quad \text{and} \quad \mathcal{M} = \bigoplus_{\substack{\lambda \text{ contained in the} \\ \text{rectangle } n \times \ell}} B(\lambda) \otimes \dot{B}(\lambda^{\text{tr}})$$

where $B(\lambda)$ (respectively $\dot{B}(\lambda)$) is the crystal of the \mathfrak{gl}_n -module (respectively \mathfrak{gl}_ℓ -module) parametrised by λ . This immediately implies the Cauchy identities. In the following sections, we shall present different extensions of this construction of multicrystal structures on sets of binary matrices (possibly infinite). They arise naturally from the notion of combinatorial Fock spaces $F(s)$, where s is a integer. The A_{n-1} and $A_{\ell-1}$ crystal structures of $F(s)$ corresponds to the \mathfrak{gl}_n and \mathfrak{gl}_ℓ modules

$$\bigoplus_{s_1 + \dots + s_\ell = s} \Lambda^{s_1}(\mathbb{C}^n) \otimes \dots \otimes \Lambda^{s_\ell}(\mathbb{C}^n), \quad \text{and} \quad \bigoplus_{\dot{s}_1 + \dots + \dot{s}_n = s} \Lambda^{\dot{s}_1}(\mathbb{C}^\ell) \otimes \dots \otimes \Lambda^{\dot{s}_n}(\mathbb{C}^\ell).$$

In affine type A , there is a similar construction due to Uglov [42] where the ordinary wedge products are replaced by their thermodynamical limits. It plays a central role in the representation theory of Cherednik algebras, see [39], and in the construction of some representations of the Virasoro algebra [21]. We get the structure of an $(\widehat{\mathfrak{sl}}_n \times \mathcal{H} \times \widehat{\mathfrak{sl}}_\ell)$ -module where \mathcal{H} is a Heisenberg algebra, and therefore a tricrystal structure on the affine Fock space. This structure has been made completely explicit in [13].

A key tool in our approach is to exploit a combinatorial duality which permits to easily switch between the different combinatorial actions defined on $F(s)$. In finite type A , this coincides with the transposition of binary matrices, and enables us to bypass the RSK correspondence. This is particularly convenient because the insertion algorithm (on which the RSK correspondence is based) in other types is more complicated and less well understood. This point of view enables us to unify and extend the existing constructions in affine type, by considering a block transposition on appropriate infinite matrices. Through the previous duality, each combinatorial object or Dynkin diagram automorphism for one structure admits a natural counterpart for the other one. For example, the cyclage operation introduced by Lascoux and Schützenberger on the semistandard tableaux of type A_{n-1} will correspond to the promotion operator on the $A_{\ell-1}$ -highest weight vertex in $F(s)$. This permits us to give a short proof of the equality between Kostka polynomials and one-dimensional sums established by Nakayashiki and Yamada in [37]. This observation also provides an algebraic interpretation of the Pitman transform $M - 2X$ introduced in [38] to study conditional Brownian motions in the spirit of [1] where its dual version is shown to map each Littelmann path on its associated highest path. Also, by considering subsets of the combinatorial Fock spaces $F(s)$ invariant under Dynkin diagram automorphisms of type A (affine or not), we get various bicrystal (or tricrystal) structures of classical types on sets of binary matrices. In a connected direction, based on the results of [31], we also establish that it is also possible to define a combinatorial Fock space with an $(X_\infty \times A_{\ell-1})$ -structure (with X of type B, C or D) on some infinite binary matrices similar to the finite type A construction. We expect that this bicrystal to be related to the charge statistics defined for type C in [28], once the appropriate duality relating both crystal structures is discovered.

As already mentioned, the combinatorial Fock spaces that we shall study can be defined in terms of (infinite) binary matrices (this corresponds to the second Cauchy identity). This is indeed the natural context corresponding to the existing mathematical material (such as Kashiwara-Nakashima columns tableaux and Uglov's Fock space), but it would be interesting to get analogous results for infinite matrices with nonnegative integer coefficients.

In the present paper, we have chosen to illustrate our results by numerous examples. Quite often, they can be established by adapting proofs existing in the literature to the unified formalism that we propose. We then give precise references rather than complete proofs. The next sections are organised as follows. Section 2 is devoted to a re-exposition of the type $(A_{n-1} \times A_{\ell-1})$ -crystal structure on \mathcal{M} . This serves as a basis for the various generalisations of the subsequent sections. In particular, we quickly reach a simple proof that the two crystals commute in Theorem 2.21, recovering the results of [43] and [15]. Also, thanks to the connection between the cyclage and the promotion operator, we are able to give a short proof of the relation between the charge and the energy function originally proved in [37]. In Section 3, the classical $(A_{n-1} \times A_{\ell-1})$ -crystal structure on \mathcal{M} is made compatible with an affine similar construction given in [10]. The highest weight vertices for the different possible (simple, double an triple) actions are described. Also, a new combinatorial interpretation of the (e, s) -cores introduced in [20] for describing the blocks of cyclotomic Hecke algebras is proposed. Section 4 describes

a type $X_\infty \times A_{\ell-1}$ analogue to the previous bicrystal. This means that we define a type $A_{\ell-1}$ -crystal structure on products of X_∞ -columns. This is done directly in terms of sliding (Jeu de Taquin) operations. Nevertheless, due to the lack of a straightforward duality, the results of Section 2 are needed to prove that this indeed yields the desired $A_{\ell-1}$ -crystal structure. The results of Section 5 focus on the vertices of the combinatorial type A (of both finite and affine type) Fock spaces fixed under the action of Dynkin diagram automorphisms. By using results of Naito and Sagaki [36], they are proved to have various bicrystal (or tricrystal) structures of classical types. Finally in Section 6, we relate Pitman's $2M - X$ transform on the line to its dual version (the $X - 2M$ transform) and the promotion operator on tensor products of type $A_1^{(1)}$ Kirillov-Reshetikhin crystals. This permits notably to show that iterations of this transform on any trajectory will eventually tend to the trivial one (that is, with all steps equal to 1). Using this time $A_\ell^{(1)}$ -crystals, a highest dimensional generalisation of the $2M - X$ transform which shares the same convergence behavior is defined. Its probabilistic properties will be studied elsewhere.

2. FINITE TYPE A DUALITY

In the rest of the paper, fix $n, \ell \in \mathbb{Z}_{\geq 2}$.

2.1. Products of type A columns. Let P be the weight lattice for the Lie algebra \mathfrak{sl}_n , with basis $\{\omega_1, \dots, \omega_{n-1}\}$, where $\omega_1, \dots, \omega_{n-1}$ are the fundamental weights for \mathfrak{sl}_n . Each partition $\lambda = (\lambda_1, \dots, \lambda_n) \in \mathbb{Z}_{\geq 0}^n$ shall be identified with the A_{n-1} -dominant weight $\lambda = \sum_{i=1}^{n-1} a_i \omega_i$ where for any $i = 1, \dots, n-1$ the integer a_i is the number of columns of height i in the Young diagram of the partition λ . Observe that the contribution of the columns of height n is thus equal to zero and we have a one-to-one correspondence between the type A_{n-1} -dominant weights and the partitions with at most $n-1$ parts. In what follows, it thus makes sense to use the symbol λ as a partition with at most n parts or a dominant weight of type A_{n-1} .

Example 2.1. Let $n = 3$ and $\lambda = 3\omega_1 + \omega_2$. Then the corresponding partition is $(4, 1, 0) = \begin{array}{|c|c|c|c|} \hline & & & \\ \hline & & & \\ \hline & & & \\ \hline \end{array}$.

A column of type A_{n-1} is a subset c of $\{1, \dots, n\}$ such that $|c| \leq n$, which we identify with the semistandard Young tableau of shape $\omega_{|c|} = (1, \dots, 1)$ containing the elements of c .

Example 2.2. The set $\{1, 3, 4\} = \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline 4 \\ \hline \end{array}$ is a column of type A_3 .

Definition 2.3. Let c_1, \dots, c_ℓ be columns of type A_{n-1} . The symbol $b = c_\ell \otimes \dots \otimes c_1$ is called

- (1) a *tableau* if the top-aligned juxtaposition $c_1 \dots c_\ell$ yields a semistandard Young tableau.
- (2) an *antitablaeu* if the bottom-aligned juxtaposition $c_1 \dots c_\ell$ yields a semistandard skew Young tableau.

The *shape* of a tableau (respectively of an antitablaeu) is the partition $(|c_1|, \dots, |c_\ell|)^{\text{tr}}$ (respectively $(|c_\ell|, \dots, |c_1|)^{\text{tr}}$).

Example 2.4. Let $\ell = 3$ and $n = 2$

- (1) The product $\begin{array}{|c|} \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array}$ is a tableau which we identify with $\begin{array}{|c|c|c|} \hline 1 & 1 & 2 \\ \hline 3 & & \\ \hline \end{array}$.
- (2) The product $\begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array}$ is an antitablaeu which we identify with $\begin{array}{|c|c|c|} \hline & 1 & 1 \\ \hline 2 & 2 & 3 \\ \hline \end{array}$.
- (3) The product $\begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array}$ is both a tableau and an antitablaeu, which we identify in with $\begin{array}{|c|c|} \hline 1 & 1 \\ \hline 2 & 3 \\ \hline \end{array}$.
- (4) The products $\begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array}$ and $\begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array}$ are neither tableaux nor antitablaeux.

For the next definition, define first the *word* $w(b)$ of a product $b = c_\ell \otimes \dots \otimes c_1$ to be the concatenation of the elements of c_ℓ (in increasing order), then $c_{\ell-1}$, and so on.

Definition 2.5. The element b is called *Yamanouchi*¹ if every prefix of $w(b)$ contains at least as many letters i as $i+1$, for all $i = 1, \dots, n-1$.

Example 2.6. Let $\ell = 4$ and $n = 3$, and take $b = \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array}$. Then $w(b) = 12132$, and the different prefixes are 1, 12, 121, 1213, 12132 and we see that b is Yamanouchi.

Clearly, for all $\lambda \in P_+$, there is a unique Yamanouchi tableau (respectively antitablaeu) of shape λ .

¹ Note that the usual convention is to use suffixes instead of prefixes. In fact, our definition coincides with the notion of *lattice word* in the literature.

Example 2.7. The Yamanouchi tableau and antitabou of shape $(4, 3, 1)$ are respectively given by

$$\begin{array}{|c|c|c|c|} \hline 1 & 1 & 1 & 1 \\ \hline 2 & 2 & 2 & \\ \hline 3 & & & \\ \hline \end{array} \quad \text{and} \quad \begin{array}{|c|c|c|} \hline & & 1 \\ \hline & 1 & 1 & 2 \\ \hline 1 & 2 & 2 & 3 \\ \hline \end{array}.$$

2.2. Crystal structures. From now on, fix $s \in \mathbb{Z}_{\geq 0}$. Further, for $\mathbf{s} = (s_1, \dots, s_\ell) \in R(n, \ell) = \{0, \dots, n\}^\ell$ and $\dot{\mathbf{s}} = (\dot{s}_\ell, \dots, \dot{s}_1) \in R(\ell, n) = \{0, \dots, \ell\}^n$, denote $|\mathbf{s}| = \sum_{j=1}^\ell s_j$ and $|\dot{\mathbf{s}}| = \sum_{i=1}^n \dot{s}_i$. For $p = n, \ell$ and for all $A \subseteq \mathbb{Z}^p$, write Finally, we denote

$$\mathcal{S}(s) = \{\lambda = (\lambda_1, \dots, \lambda_n) \vdash s \mid \lambda_1 \leq \ell\}.$$

The elements introduced in the previous section appear as vertices of certain tensor products of crystal graphs, which we start by recalling. For all $j = 1, \dots, \ell$, the crystal of the irreducible highest weight \mathfrak{sl}_n -module of highest weight ω_{s_j} (with the convention $\omega_0 = \omega_n = 0$) can be realised using columns of height s_j [17, Chapter 7]. More precisely, $B(\omega_{s_j})$ is the A_{n-1} -crystal with vertices the columns of height s_j and arrows i from c to c' when c' is obtained from c by changing i into $i + 1$. Observe that the trivial crystal of highest weight 0 can so be realised as the graph with a unique vertex: the empty column or the column containing all the integers $1, \dots, n$.

Definition 2.8. Let $\mathbf{s} = (s_1, \dots, s_\ell) \in R(n, \ell)$. The combinatorial Fock space associated to \mathbf{s} is the A_{n-1} -crystal

$$F(\mathbf{s}) = B(\omega_{s_\ell}) \otimes \dots \otimes B(\omega_{s_1})$$

By classical crystal theory, the elements of $F(\mathbf{s})$ can be realised as tensor products of ℓ columns with entries in $\{1, \dots, n\}$. Let us recall the rule for computing $F(\mathbf{s})$, following [17, Section 4.4]. Fix $i \in \{1, \dots, n-1\}$ and let $b = c_\ell \otimes \dots \otimes c_1 \in F(\mathbf{s})$. consider the subword $w_i(b)$ of $w(b)$ obtained by keeping only letters i and $i + 1$, and encode each i by a symbol $+$ and each $i + 1$ by a symbol $-$. Deleting all factors $+-$ recursively yields a word called the i -signature of b .

Theorem 2.9. The action of the A_{n-1} -crystal operators on $F(\mathbf{s})$ is given by the following rule

- (1) The raising crystal operator e_i acts on $b \in F(\mathbf{s})$ by changing the entry $i + 1$ corresponding to the rightmost $-$ in the i -signature of b into i if it exists; and by 0 otherwise.
- (2) The lowering crystal operator f_i acts on $b \in F(\mathbf{s})$ by changing the entry i corresponding to the leftmost $+$ in the i -signature into $i + 1$ if it exists; and by 0 otherwise.

We define similarly the combinatorial Fock space $\dot{F}(\dot{\mathbf{s}}) = B(\omega_{\dot{s}_1}) \otimes \dots \otimes B(\omega_{\dot{s}_n})$ for $\dot{\mathbf{s}} = (\dot{s}_1, \dots, \dot{s}_n) \in \mathbb{Z}^n$, and the rule for computing $\dot{F}(\dot{\mathbf{s}})$ is the same as for $F(\mathbf{s})$, expect that the role of n and ℓ have been swapped².

Example 2.10. Let $\ell = 4, n = 3, \mathbf{s} = (2, 2, 1, 2), b = \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} = c_4 \otimes c_3 \otimes c_2 \otimes c_1$ and choose $i = 1$. Then $w(b) = 1213112$, so that $w_1(b) = 121112$, which gives the encoded word $+-++++-$. Thus, the i -signature of b is $++$, whose leftmost $+$ corresponds to the entry 1 of c_3 . Therefore, we have $f_1 b = \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} = c_4 \otimes c_3 \otimes c_2 \otimes c_1$. In other terms, in the crystal $F(\mathbf{s})$, we have an arrow

$$\begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \xrightarrow{1} \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array}.$$

Let us now explain how the tableaux and Yamanouchi elements naturally appear in the context of crystals. In the following, we set

$$F(s) = \bigoplus_{\mathbf{s} \in R(n, \ell)(s)} F(\mathbf{s}) \quad \text{and} \quad \dot{F}(s) = \bigoplus_{\dot{\mathbf{s}} \in R(\ell, n)(s)} \dot{F}(\dot{\mathbf{s}}).$$

The following results are well-known, see for instance [34].

Theorem 2.11.

- (1) The set of tableaux in $F(s)$ is closed under the crystal operators, and tableaux of a given shape $\lambda \in \mathcal{S}(s)$ form a connected component of $F(s)$ denoted $B(\lambda)$. Moreover, for any $b \in F(s)$, there is a unique tableau $P(b) \in F(s)$ such that the induced map $b \mapsto P(b)$ is an A_{n-1} -crystal isomorphism.
- (2) An element $b \in F(s)$ is a highest weight vertex in the crystal if and only if b is Yamanouchi.

² Note that the order in which the components of \mathbf{s} and $\dot{\mathbf{s}}$ are enumerated is reversed. Though this seems artificial at this point, this will be crucial in Section 2.3

Remark 2.12. In fact, for all $\lambda \in \mathcal{S}(s)$, $B(\lambda)$ is the crystal of the irreducible highest weight module with highest weight λ .

Theorem 2.11 also holds for $\dot{F}(s)$, replacing tableaux by antit tableaux. We denote similarly $\dot{F}(\mu)$ the connected component of $\dot{F}(s)$ consisting of all antit tableaux of shape $\mu \in \mathcal{S}(s) = \{\mu \vdash s \mid \mu^{\text{tr}} \in \mathcal{S}(s)\}$.

Example 2.13. Take $n = 3$, $\ell = 3$ and $\mathbf{s} = (s_3, s_2, s_1) = (2, 1, 2)$. Then one checks that $b = \begin{smallmatrix} 1 \\ 2 \end{smallmatrix} \otimes \begin{smallmatrix} 3 \end{smallmatrix} \otimes \begin{smallmatrix} 1 \\ 2 \end{smallmatrix} \in F(\mathbf{s})$ is Yamanouchi. We compute the connected component of $F(\mathbf{s})$ containing b :

$$\begin{smallmatrix} 1 \\ 2 \end{smallmatrix} \otimes \begin{smallmatrix} 3 \end{smallmatrix} \otimes \begin{smallmatrix} 1 \\ 2 \end{smallmatrix} \xrightarrow{2} \begin{smallmatrix} 1 \\ 2 \end{smallmatrix} \otimes \begin{smallmatrix} 3 \end{smallmatrix} \otimes \begin{smallmatrix} 1 \\ 3 \end{smallmatrix} \xrightarrow{1} \begin{smallmatrix} 1 \\ 2 \end{smallmatrix} \otimes \begin{smallmatrix} 3 \end{smallmatrix} \otimes \begin{smallmatrix} 2 \\ 3 \end{smallmatrix}.$$

This is isomorphic to the crystal $B(\omega_2)$, which we can compute:

$$\begin{array}{|c|c|} \hline 1 & 1 \\ \hline 2 & 2 \\ \hline 3 & \\ \hline \end{array} \xrightarrow{2} \begin{array}{|c|c|} \hline 1 & 1 \\ \hline 2 & 3 \\ \hline 3 & \\ \hline \end{array} \xrightarrow{1} \begin{array}{|c|c|} \hline 1 & 2 \\ \hline 2 & 3 \\ \hline 3 & \\ \hline \end{array}.$$

This means that $P(b) = \begin{array}{|c|c|} \hline 1 & 1 \\ \hline 2 & 2 \\ \hline 3 & \\ \hline \end{array}$, and so on. Alternatively, we can use an isomorphic realisation of this crystal by antit tableaux:

$$\begin{array}{|c|c|} \hline 1 \\ \hline 1 & 2 \\ \hline 2 & 3 \\ \hline \end{array} \xrightarrow{2} \begin{array}{|c|c|} \hline 1 \\ \hline 1 & 2 \\ \hline 3 & 3 \\ \hline \end{array} \xrightarrow{1} \begin{array}{|c|c|} \hline 1 \\ \hline 2 & 2 \\ \hline 3 & 3 \\ \hline \end{array}.$$

In general, $P(b)$ can be computed by carrying out one of the following procedures:

- performing Schensted's insertion on the word $w(b)$ [6, Section 1.1],
- performing the Jeu de Taquin on the skew Young tableau corresponding to b [6, Section 1.2],
- applying a sequence of plactic relation to $w(b)$ [6, Section 2.1].

Example 2.14. Let $\ell = 2$, $n = 3$ and let $b = \begin{smallmatrix} 1 \\ 2 \\ 3 \end{smallmatrix} \otimes \begin{smallmatrix} 2 \end{smallmatrix}$. Let us compute $P(b)$ by using the first two methods. The word associated to b is $w(b) = 12413$, and Schensted's insertion yields the following sequence of tableaux

$$\begin{smallmatrix} 1 \end{smallmatrix}, \begin{smallmatrix} 1 \\ 2 \end{smallmatrix}, \begin{smallmatrix} 1 \\ 2 \\ 3 \end{smallmatrix}, \begin{smallmatrix} 1 & 2 \\ 2 & 3 \end{smallmatrix} = P(b).$$

Now, the minimal skew Young tableau³ associated to b is

$$\begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 2 & 3 \\ \hline \end{array}$$

and the Jeu de Taquin corresponds to the following two sliding operations

$$\begin{array}{|c|} \hline 1 \\ \hline \bullet & 2 \\ \hline 2 & 3 \\ \hline \end{array} \rightarrow \begin{array}{|c|c|} \hline 1 \\ \hline 2 & 2 \\ \hline \bullet & 3 \\ \hline \end{array} \rightarrow \begin{array}{|c|c|} \hline 1 \\ \hline 2 & 2 \\ \hline 3 & \bullet \\ \hline \end{array}, \begin{array}{|c|} \hline \bullet & 1 \\ \hline 2 & 2 \\ \hline 3 & \\ \hline \end{array} \rightarrow \begin{array}{|c|} \hline 1 \\ \hline 2 & \bullet \\ \hline 2 & 2 \\ \hline 3 & \\ \hline \end{array} \rightarrow \begin{array}{|c|c|} \hline 1 & 2 \\ \hline 2 & \bullet \\ \hline 2 & 2 \\ \hline 3 & \\ \hline \end{array} = P(b).$$

2.3. The duality. There is a duality

$$\begin{array}{ccc} F(s) & \longleftrightarrow & \dot{F}(s) \\ b & \longleftrightarrow & b^* \end{array}$$

defined as follows. If $b = c_\ell \otimes \cdots \otimes c_1$ is a tensor product of columns, then for each $i = 1, \dots, n$, let $d(i)$ be the column with letters in the set

$$\{j \in \{1, \dots, \ell\} \mid i \in c_j\}.$$

Then we set $b^* = d(1) \otimes \cdots \otimes d(n) \in \dot{F}(s)$.

³ This skew tableau is minimal in the sense that its skew shape is minimal (for the inclusion of skew shapes) among all the possible shapes of the skew tableaux associated to b .

Example 2.15. Let $\ell = 5$ and $n = 4$. Take

$$b = \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array}.$$

Then

$$d(1) = \{1, 2, 3, 5\}, \quad d(2) = \{2, 5\}, \quad d(3) = \{1, 5\}, \quad d(4) = \{2, 3, 4\},$$

so that

$$b^* = \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline \end{array}.$$

Remark 2.16. As mentioned in the introduction, we can use binary matrices to represent elements in $F(s)$ and $\dot{F}(s)$, which yields an easy description of the duality $*$. More precisely, encode $b = c_\ell \otimes \cdots \otimes c_1$ by the $n \times \ell$ matrix M defined by

$$M_{i,j} = \begin{cases} 1 & \text{if } i \in c_j \\ 0 & \text{otherwise} \end{cases}$$

Then b^* is the element of $\dot{F}(s)$ encoded by M^{tr} , the transpose of M . For instance, take b as in Example 2.15. Then b and b^* are respectively encoded by the following matrices (remember that we read the columns of b starting from the right)

$$M = \begin{bmatrix} 1 & 1 & 1 & 0 & 1 \\ 0 & 1 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 & 1 \\ 0 & 1 & 1 & 1 & 0 \end{bmatrix} \quad \text{and} \quad M^{\text{tr}} = \begin{bmatrix} 1 & 0 & 1 & 0 \\ 1 & 1 & 0 & 1 \\ 1 & 0 & 0 & 1 \\ 0 & 0 & 0 & 1 \\ 1 & 1 & 1 & 0 \end{bmatrix}.$$

Therefore, we recover the crystal skew Howe duality of [15, Section 11.2], also studied in [43].

The duality $*$ intertwines several classical notions. We already observe some occurrences of this phenomenon now, and will give more results in the upcoming sections. For all map $\varphi : \dot{F}(s) \rightarrow \dot{F}(s) \sqcup \{0\}$, denote $\varphi^* : F(s) \rightarrow F(s) \sqcup \{0\}$ the map determined by the formula

$$(\varphi(b))^* = \begin{cases} \varphi(b^*)^* & \text{if } \varphi(b^*) \in \dot{F}(s) \\ 0 & \text{if } \varphi(b^*) = 0. \end{cases}$$

that is, φ^* is the conjugation of φ by the duality $*$. Similarly, for all map $\psi : F(s) \rightarrow F(s) \sqcup \{0\}$, denote ${}^*\psi : \dot{F}(s) \rightarrow \dot{F}(s) \sqcup \{0\}$ the map determined by the formula

$${}^*\psi(b^*) = \begin{cases} (\psi(b))^* & \text{if } \psi(b) \in F(s) \\ 0 & \text{if } \psi(b) = 0. \end{cases}$$

Note that we have $\psi = \varphi^*$ if and only if ${}^*\psi = \varphi$.

Proposition 2.17. *The element b is a tableau (respectively Yamanouchi) if and only if b^* is Yamanouchi (respectively an antitableau).*

Proof. Let $b \in F(s)$ be a tableau. Then $s_j \geq s_{j+1}$ for all $j = 1, \dots, \ell - 1$, which means that the number of j 's in b^* is greater than or equal to the number of $j + 1$'s, which is a necessary condition to be Yamanouchi. In fact, the k first components of b correspond by duality to the subtableau of b with letters less or equal to k . Since all such subtableau is semistandard, the vertex b^* has the Yamanouchi property. The converse holds by the same observation, and the analogue statement with antitableaux holds similarly. \square

Example 2.18.

(1) Let $\ell = 2$, $n = 4$, and $b = \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 4 \\ \hline \end{array}$, so that b is a tableau. Then $b^* = \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array}$, which is Yamanouchi.

(2) Let $\ell = 5$, $n = 3$, and $b = \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array}$, so that b is Yamanouchi. Then $b^* = \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array}$, which is an antitableau.

We have already recalled the Jeu de Taquin procedure for computing the tableau in Theorem 2.11. Let us denote J_j the map $F(s) \rightarrow F(s) \sqcup \{0\}$ where $J_j(b)$ is obtained from b by an elementary horizontal Jeu de Taquin slide from column $j+1$ to column j if possible, and $J_j(b) = 0$ otherwise. Let us moreover recall that there is a unique isomorphism of A_{n-1} -crystals $B(\omega_i) \otimes B(\omega_{i'}) \xrightarrow{\sim} B(\omega_{i'}) \otimes B(\omega_i)$ called the *combinatorial R-matrix*, which we can compute by a simple combinatorial procedure, see [40, Section 4.8] and the references therein. It induces an isomorphism

$$R_{j,j'} : F(s) \longrightarrow F(s)$$

permuting the components $B(\omega_{s_j})$ and $B(\omega_{s_{j'}})$, which is the composition of R -matrices of the form $R_{j,j+1}$, which we denote R_j for simplicity. Finally, recall that the Weyl group (here the symmetric groups on n letters) acts on the crystal $F(s)$ by letting the Coxeter generators σ_i , $i = 1 \dots, n-1$ act by reversing each i -string [3, Definition 2.35], see also [40, Section 2.1].

All of the above maps have counterparts defined on $\dot{F}(s)$ (defined using antitableaux instead of tableaux when needed). We now prove that the duality $*$ intertwines crystal operators with elementary Jeu de Taquin slides, as well as the Weyl group action with the R -matrix.

Theorem 2.19. *For all $j = 1, \dots, \ell-1$, we have*

$$(1) \quad \dot{e}_j^* = J_j \quad \text{and} \quad (2) \quad \dot{\sigma}_j^* = R_j.$$

Proof.

- (1) Assume first that b^* is a highest weight vertex. Then $\dot{e}_j b^* = 0$, so we have by definition $\dot{e}_j^* b = 0 = J_j(b)$. Note that this simply means that the Jeu de Taquin is not authorised for b , which makes sense because b is a tableau by Theorem 2.11 and Proposition 2.17. Now, assume that $b^* = d_1 \otimes \dots \otimes d_n$ is not a highest weight vertex, so that $\dot{e}_j b^* \in \dot{F}(s)$. Let i denote the index of the column containing the entry of b^* affected by \dot{e}_j , see Theorem 2.9. Then by definition of the duality $*$, \dot{e}_j^* acts on $b = c_\ell \otimes \dots \otimes c_1$ by sliding entry i from c_{j+1} to c_j . Now, consider the Jeu de Taquin between columns j and $j+1$ of b . The procedure used to determine i ensures us that all entries of c_j that are smaller than or equal to i are matched to an element of c_{j+1} . Therefore, the first entry that slides from c_{j+1} to c_j is i . In other terms, $J_j(b)$ is obtained from b by sliding i from c_{j+1} to c_j . Thus $\dot{e}_j^* b = J_j(b)$.
- (2) Set $N = |c_{j+1}| - |c_j|$, where $b = c_\ell \otimes \dots \otimes c_1$. If $N \geq 0$, then the element $R_j(b)$ is obtained by using N Jeu de Taquin slides between columns j and $j+1$ of b . Therefore, by Part (1), $(R_j(b))^* = \dot{e}_j^N b^*$. Now, by definition of the duality $*$, we have $N = \varepsilon_j(b^*) - \varphi_j(b^*)$, therefore $\dot{e}_j^N b^* = \dot{\sigma}_j(b^*)$ and the claim is proved. The case $N < 0$ is proved similarly by using the lowering operators \dot{f}_j instead. □

Example 2.20. (1) Take $n = 6$ and $\ell = 4$, and let

$$b^* = \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array}, \quad \text{so that} \quad b = \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 6 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 4 \\ \hline 5 \\ \hline \end{array}.$$

Then one can check that

$$\dot{e}_1 b^* = \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array}, \quad \text{whose dual is} \quad b' = \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 6 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array}.$$

On the other hand, the horizontal slide from column 2 to column 1 of b is achieved as

$$\begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 1 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \begin{array}{|c|} \hline \bullet \\ \hline \end{array} \rightarrow \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \begin{array}{|c|} \hline \bullet \\ \hline \end{array} \rightarrow \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \begin{array}{|c|} \hline \bullet \\ \hline \end{array} \rightarrow \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \begin{array}{|c|} \hline \bullet \\ \hline \end{array} = \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 6 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array},$$

and we see that we recover b' .

(2) Take $n = 5$, $\ell = 4$, and

$$b^* = \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 3 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \quad \text{so that} \quad b = \begin{array}{|c|} \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 3 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline 5 \\ \hline \end{array}.$$

Let us look at $j = 2$. One checks that $\varepsilon_j(b^*) = 3$ and $\varphi_j(b^*) = 1$, (i.e. there are 3 incoming and 1 outgoing arrows with color 2 at vertex $b^* \in \hat{F}(s)$), so that

$$s_2 b^* = \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array}, \quad \text{whose dual is} \quad b' = \begin{array}{|c|} \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline 5 \\ \hline \end{array}.$$

Now, we can compute $R_2(b)$, by doing the following two Jeu de Taquin slides between columns 2 and 3

$$\begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \begin{array}{|c|} \hline 3 \\ \hline 4 \\ \hline \end{array} \rightarrow \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \begin{array}{|c|} \hline 3 \\ \hline 3 \\ \hline 4 \\ \hline \end{array} \rightarrow \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \begin{array}{|c|} \hline 3 \\ \hline 3 \\ \hline 4 \\ \hline \end{array}, \quad \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \begin{array}{|c|} \hline 3 \\ \hline 3 \\ \hline 4 \\ \hline \end{array} \rightarrow \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline \end{array} \rightarrow \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline \end{array} = \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array},$$

from which we recover b' .

Of course, Theorem 2.19 also holds when switching $F(s)$ and $\hat{F}(s)$ and taking the dual versions of the different maps.

2.4. Bicrystal structure. We are ready to prove that the A_{n-1} and $A_{\ell-1}$ crystal structures commute (modulo the duality $*$). This is best stated as follows.

Theorem 2.21. *For all $j = 1, \dots, \ell - 1$, the restriction of \dot{f}_j^* to any connected component of $F(s)$ is either 0 or an A_{n-1} -crystal isomorphism. Similarly, for all $i = 1, \dots, n - 1$, the restriction of *f_i to any connected component of $\hat{F}(s)$ is either 0 or an $A_{\ell-1}$ -crystal isomorphism.*

Proof. Note that by symmetry, both statements are equivalent, and it is enough to prove that for all $j = 1, \dots, \ell - 1$, the map \dot{f}_j^* is an A_{n-1} -crystal isomorphism. Fix $j \in \{1, \dots, \ell - 1\}$. First of all, if $b \in F(s)$ is a highest weight vertex, then b is Yamanouchi by Theorem 2.11(2), and b^* is an antitableau by Proposition 2.17. By Theorem 2.11(1), $\dot{f}_j b^*$ is again an antitableau, so $\dot{f}_j^* b$ is Yamanouchi, i.e. a highest weight vertex. This shows that \dot{f}_j^* maps highest weight vertices to highest weight vertices with the same weight. It remains to show that it commutes with the lowering crystal operators f_i . We have, for all $i = 1, \dots, n - 1$,

$$\begin{aligned} \dot{f}_j^* f_i &= J_j^{-1} f_i \quad \text{by Theorem 2.19(1)} \\ &= f_i J_j^{-1} \quad \text{by Theorem 2.11(1)} \\ &= f_i \dot{f}_j^*. \end{aligned}$$

□

This yields an $(A_{n-1} \times A_{\ell-1})$ -crystal structure on $F(s)$. For $\underline{i} = (i_1, \dots, i_r) \in \{1, \dots, n\}^r$ (respectively $\underline{j} = (j_1, \dots, j_t) \in \{1, \dots, \ell\}^t$), denote $f_{\underline{i}} = f_{i_r} \cdots f_{i_1}$ (respectively $\dot{f}_{\underline{j}} = \dot{f}_{j_t} \cdots \dot{f}_{j_1}$). The following corollary is immediate from Proposition 2.17 and Theorem 2.21. For a given $\lambda \in P_+$, denote b_λ the Yamanouchi tableau of shape λ .

Corollary 2.22. *Each connected component of the $(A_{n-1} \times A_{\ell-1})$ -crystal $F(s)$ has a unique source vertex. The sources are exactly the Yamanouchi tableaux. In other terms, we have*

$$F(s) = \bigoplus_{\substack{\lambda = \omega_{s_1} + \dots + \omega_{s_\ell} \\ s_1 + \dots + s_\ell = s}} \dot{f}_{\underline{j}}^* f_{\underline{i}} b_\lambda,$$

where the sum runs over all possible \underline{i} and \underline{j} .

Remark 2.23. Accordingly, the $(A_{n-1} \times A_{\ell-1})$ -crystal structure can be considered on $\hat{F}(s)$. In this case, the sources are the Yamanouchi antitableaux.

Example 2.24. Take $\ell = 4, n = 3$. The element

$$\begin{array}{|c|} \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ 2 \\ 3 \\ \hline \end{array}$$

is a Yamanouchi tableau, i.e. a source in the $(A_{n-1} \times A_{\ell-1})$ -crystal. Its dual is the following Yamanouchi antitableau

$$\begin{array}{|c|} \hline 1 \\ 2 \\ 3 \\ 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ 2 \\ 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline \end{array}.$$

By Corollary 2.22, for all $b \in F(s)$, there is a unique Yamanouchi tableau \bar{b} such that $b = f_{\underline{j}}^* f_{\underline{i}} \bar{b}$ for some \underline{i} and \underline{j} . Set

$$P(b) = f_{\underline{i}} \bar{b} \quad \text{and} \quad Q(b) = f_{\underline{j}}^* \bar{b}^*.$$

Theorem 2.25. *For all $b \in F(s)$, $P(b)$ is a tableau and $Q(b)$ is an antitableau of transpose shape. Moreover, the assignment*

$$\Phi : b \mapsto (P(b), Q(b))$$

yields a bijection $F(s) \rightarrow \Phi(F(s))$ called the crystal RSK correspondence. In particular, we have the decomposition

$$F(s) \simeq \bigoplus_{\lambda \in \mathcal{S}(s)} B(\lambda) \otimes \dot{B}(\lambda^{\text{tr}}).$$

Proof. By Theorem 2.21, the two crystals commute, so since \bar{b} is a highest weight vertex for the $A_{\ell-1}$ -crystal, $P(b) = f_{\underline{i}} \bar{b}$ is also a highest weight vertex for the $A_{\ell-1}$ -crystal, i.e. $P(b)^*$ is Yamanouchi. By Proposition 2.17, $P(b)$ is a tableau. Similarly, we get that $Q(b)$ is an antitableau. Moreover, Φ is clearly injective, so that $\Phi : F(s) \rightarrow \Phi(F(s))$ is a bijection. \square

Accordingly, for $b^* \in \dot{F}(s)$, we will write $P(b^*) = Q(b)$ and $Q(b^*) = P(b)$.

2.5. Crystal structure on self-dual elements. An element $b = c_{\ell} \otimes \cdots \otimes c_1 \in F(s)$ is called *self-dual* if $b^* = c_1 \otimes \cdots \otimes c_{\ell}$. In particular, self-dual elements exist only if $n = \ell$. Alternatively, if M denotes the binary matrix associated to b , then b is self-dual if and only if $M^{\text{tr}} = M$. We denote $F(s)^*$ the set of self-dual elements of $F(s)$. Now, for all $i = 1, \dots, n-1$, set

$$f_i^* = f_i^* f_i.$$

The operators $f_i^*, i = 1, \dots, n-1$, induce an A_{n-1} -crystal structure on $F(s)^*$. More precisely, we get the decomposition

$$F(s)^* \simeq \bigoplus_{\substack{\lambda \in \mathcal{S}(s) \\ \lambda^{\text{tr}} = \lambda}} B(\lambda).$$

2.6. Keys and bikeys. Keys are certain tableaux introduced by Lascoux and Schützenberger [26] that are used to compute Demazure crystals [3]. In this section, we generalise the notion of keys using the bicrystal structure of Section 2.4.

Definition 2.26. *An element $b = c_{\ell} \otimes \cdots \otimes c_1 \in F(s)$ is called a key (respectively an antikey) if b is a tableau (respectively an antitableau) and if $c_{\ell} \subseteq \cdots \subseteq c_1$ (respectively $c_1 \subseteq \cdots \subseteq c_{\ell}$).*

For $b \in F(s)$, let $\mathcal{O}_{S_n}(b)$ be the orbit of b under the action of the Weyl group S_n . The following proposition is easy to establish by induction on the length of the elements of S_n .

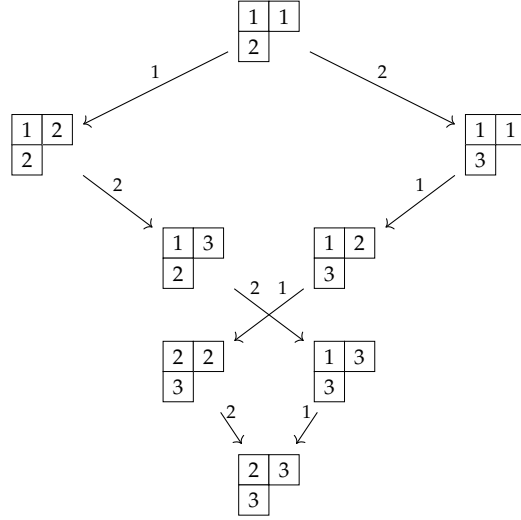
Proposition 2.27. *The set of all keys of given shape $\lambda \in \mathcal{S}(s)$ is equal to $\mathcal{O}_{S_n}(b_{\lambda})$.*

Therefore, we have the following characterisation of keys (which could also be proved directly by using the definition of $*$).

Proposition 2.28. *Let $b \in F(s)$. Then b is a key if and only if b^* is a product of Yamanouchi columns.*

Proof. We know that b_{λ}^* is the Yamanouchi antitableau of shape λ^{tr} . In particular, b_{λ}^* is an antikey. Now by Theorem 2.19(2) and Proposition 2.27, b is a key if and only if b^* is obtained from b_{λ}^* by applying a sequence of R_j 's, which simply permute the columns since each column of b_{λ}^* is included in the next. \square

Example 2.29. Take $n = 3, \ell = 2$. Then the crystal of highest weight $\lambda = \omega_1 + \omega_2$ is realised by tableaux as follows



The keys of shape λ are obtained by keeping all vertices at the extremity of each i -string (for $i = 1, 2$), namely

$$\begin{array}{|c|c|} \hline 1 & 1 \\ \hline 2 & \\ \hline \end{array}, \begin{array}{|c|c|} \hline 1 & 2 \\ \hline 2 & \\ \hline \end{array}, \begin{array}{|c|c|} \hline 1 & 1 \\ \hline 3 & \\ \hline \end{array}, \begin{array}{|c|c|} \hline 1 & 3 \\ \hline 3 & \\ \hline \end{array}, \begin{array}{|c|c|} \hline 2 & 2 \\ \hline 3 & \\ \hline \end{array}, \begin{array}{|c|c|} \hline 2 & 3 \\ \hline 3 & \\ \hline \end{array},$$

whose respective dual are the following products of Yamanouchi columns

$$\begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array} \otimes \emptyset, \begin{array}{|c|} \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \emptyset, \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \emptyset \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array}, \begin{array}{|c|} \hline 2 \\ \hline \end{array} \otimes \emptyset \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array}, \emptyset \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array}, \emptyset \otimes \begin{array}{|c|} \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array}.$$

We are ready to introduce the generalised notion of key.

Definition 2.30. An element $b = c_1 \otimes \cdots \otimes c_\ell \in F(s)$ is called a *bikey* if there exists $\dot{w} \in S_\ell$ such that $c_{\dot{w}(1)} \subseteq \cdots \subseteq c_{\dot{w}(\ell)}$.

Write \mathcal{K} for the set of bikeys in $F(s)$. Observe that for bikeys, the distinction between tableaux and antit tableaux disappears. More precisely, when b is a bikey, for any $j = 1, \dots, \ell - 1$ we have the inclusion $c_j \subseteq c_{j+1}$ or $c_{j+1} \subseteq c_j$. It thus follows from Theorem 2.19(2) that $\dot{w}^*(b) = c_{\dot{w}(1)} \otimes \cdots \otimes c_{\dot{w}(\ell)}$. Moreover, $\dot{w}^*(b)$ is a tableau (and an antit tableau), thus for a bikey b we must have $P(b) = \dot{w}^*(b)$. For all $\lambda \in \mathcal{S}(s)$, write

$$\mathcal{K}(\lambda) = \{b \in \mathcal{K} \mid P(b) \text{ has shape } \lambda\}.$$

The next result justifies the terminology of the previous definition.

Proposition 2.31. For all $\lambda \in \mathcal{S}(s)$, we have $\mathcal{K}(\lambda) = \mathcal{O}_{S_n \times S_\ell}(b_\lambda)$.

Proof. We have already observed that $\dot{w}^*(b)$ is a tableau. Since $c_{\dot{w}(1)} \subseteq \cdots \subseteq c_{\dot{w}(\ell)}$, it is a key of shape λ . Thus, there exists $w \in S_n$ such that $w\dot{w}^*(b) = b_\lambda$ which proves that $\mathcal{K}(\lambda) \subseteq \mathcal{O}_{S_n \times S_\ell}(b_\lambda)$. Now for any $w \in S_n$ and $\dot{w} \in S_\ell$, the vertex $b = w\dot{w}^*(b_\lambda) = \dot{w}^*w(b_\lambda)$ is a bikey since $w(b_\lambda)$ is a key on which the action of \dot{w}^* (which is a combination of R -matrices) reduces to a permutation of the columns. Hence we get $\mathcal{K}(\lambda) = \mathcal{O}_{S_n \times S_\ell}(b_\lambda)$. \square

Finally, we prove that the duality $*$ maps bikeys to bikeys. For this, consider the set $\dot{\mathcal{K}}$ of bikeys in $\dot{F}(s)$, and for $\mu \vdash s$ let

$$\dot{\mathcal{K}}(\mu) = \{\dot{b} \in \dot{\mathcal{K}} \mid P(\dot{b}) \text{ has shape } \mu\}.$$

Proposition 2.32. We have $b \in \mathcal{K}(\lambda)$ if and only if $b^* \in \dot{\mathcal{K}}(\lambda^{\text{tr}})$.

Proof. It suffices to observe that for any $b \in \mathcal{K}$, there exist a partition λ and $(w, \dot{w}) \in S_n \times S_\ell$ such that $b = w\dot{w}^*(b_\lambda)$. Then $b^* = {}^*w\dot{w}(b_\lambda^*) = {}^*w\dot{w}(b_{\lambda^{\text{tr}}})$ belongs to $\dot{\mathcal{K}}$. \square

Let \mathcal{H} be the set of binary matrices corresponding to the bikeys in \mathcal{K} . For any $\lambda \in \mathcal{S}(s)$, recall that the binary matrix M_λ corresponding to b_λ is the $n \times \ell$ matrix whose entries equal to 1 form a pattern corresponding to the Young diagram of λ . Also, we have a right S_ℓ -action (respectively a left S_n -action) on the set of $n \times \ell$ matrices where the action of $\dot{w} \in S_\ell$ (respectively of $w \in S_n$) is given by the right (respectively left) multiplication by the permutation matrix $P_{\dot{w}}$ (respectively P_w) associated to \dot{w} (respectively w). Since permuting the columns of b (respectively of b^*) is equivalent to permuting the columns (respectively the rows) of the corresponding binary matrix, we get the following corollary.

Corollary 2.33. *We have*

$$\mathcal{K} = \bigsqcup_{(w, \bar{w}) \in S_n \times S_\ell} P_w M_\lambda P_{\bar{w}}$$

and

$$\sum_{(m_{i,j}) \in \mathcal{K}} \prod_{\substack{1 \leq i \leq n \\ 1 \leq j \leq \ell}} (x_i y_j)^{m_{i,j}} = \sum_{\lambda \in \mathcal{S}(s)} m_\lambda(x) m_{\lambda^{\text{tr}}}(y).$$

2.7. Cyclage and promotion. In this section, we show that the duality $*$ intertwines two important maps, namely the cyclage and the promotion operators. This will be used in Section 2.8 to study the relationship between charge and energy.

Definition 2.34.

(1) Let $b = c_\ell \otimes \cdots \otimes c_1 \in F(s)$. The cyclage of b is the element of $F(s)$ defined by

$$\zeta(b) = c_{\ell-1} \otimes \cdots \otimes c_1 \otimes c_\ell.$$

(2) Let $b^* = d_1 \otimes \cdots \otimes d_n \in \dot{F}(s)$. The promotion of b^* is the element of $\dot{F}(s)$ defined by $\text{pr}(b^*) = \text{pr}(d_1) \otimes \cdots \otimes \text{pr}(d_n)$ where, for all $i = 1, \dots, n$,

$$\text{pr}(d_i) = \{k + 1 \bmod \ell; k \in d_i\}$$

Remark 2.35. The promotion operator permits to endow $\dot{F}(s)$ with the structure of an affine type $A_{\ell-1}^{(1)}$ -crystal that we shall denote by $\dot{F}(s)^{\text{aff}}$. This structure is obtained by considering the additional crystal operators $\dot{f}_0 = \text{pr}^{-1} f_1 \circ \text{pr}$ and $\dot{e}_0 = \text{pr}^{-1} e_1 \circ \text{pr}$. The crystal $\dot{F}(s)^{\text{aff}}$ splits into affine connected components. Each such component is isomorphic to a tensor product of n column Kirillov-Reshetikhin crystals. Observe that these components are not highest weight crystals.

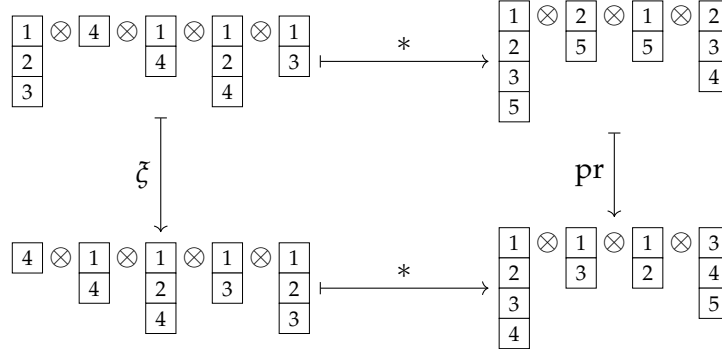
The following proposition is immediate from the definitions.

Proposition 2.36.

(1) For all $b \in F(s)$, we have $\zeta(b)^* = \text{pr}(b^*)$.

(2) For all $b^* \in \dot{F}(s)$, we have $\text{pr} \circ R(b^*) = R \circ \text{pr}(b^*)$

Example 2.37. Take the same values as in Example 2.15. Then we have



In order for the cyclage to be relevant for the description of the Kostka polynomials, we need to restrict to some cases, which we call *authorised*.

Definition 2.38. Let $b = c_\ell \otimes \cdots \otimes c_1 \in F(s)$.

(1) We say that the cyclage is authorised for b if either

(a) b is of dominant evaluation, i.e. b has no more $k + 1$ than k for all $k = 1, \dots, n - 1$, and $1 \notin c_\ell$.

(b) b is a tableau and there is no $k \in \{1, \dots, n\}$ appearing in every column c_j .

(2) Suppose that there exists $k \in \{1, \dots, n\}$ such that $k \in c_j$ for all $j = 1, \dots, \ell$. In particular, b is not authorised. Let k_0 be minimal with this property. The reduction of b is the element $\text{red}(b)$ obtained by deleting all occurrences of k_0 in b , and replacing each $k > k_0$ by $k - 1$.

Remark 2.39. In particular, if b verifies (2), there exists m such that the cyclage is authorised for the tableau $\text{red}^m(b)$. We denote $\bar{\zeta} : b \mapsto \zeta(\text{red}^m(b))$.

Example 2.40. Take $\ell = 2, n = 3$ and $b = \begin{bmatrix} 1 \\ 2 \\ 3 \end{bmatrix} \otimes \begin{bmatrix} 2 \\ 3 \end{bmatrix}$. We have $k_0 = 2$ and $\text{red}(b) = \begin{bmatrix} 1 \\ 2 \end{bmatrix} \otimes \begin{bmatrix} 2 \end{bmatrix}$.

It can be read off the dual if the cyclage is authorised, and how reduction acts.

Proposition 2.41.

- (1) Let $b \in F(s)$ and write $b^* = d_1 \otimes \cdots \otimes d_n$. The cyclage is authorised for b if and only if either
 - (a) $|d_1| \geq \cdots \geq |d_n|$ and $\ell \notin d_1$, or
 - (b) b^* is Yamanouchi and for all $i = 1, \dots, n$, $d_i \neq \{1, \dots, \ell\}$.
- (2) $\text{red}(b)^*$ is obtained from $\text{red}(b)$ by deleting the leftmost column $\{1, \dots, \ell\}$ in b^*

Proof. For (1), one checks from the definition that an element b is of dominant evaluation if and only if its dual b^* has columns of non-increasing size, and (a) follows. Point (b) is a direct consequence of Proposition 2.17. Point (2) is clear from the definition of $*$. \square

2.8. Charge and energy. Charge and energy are two classic statistics defined on $F(s)$ and $\dot{F}(s)$. The goal of this section is to establish that the duality $*$ intertwines these two notions.

Lemma 2.42. For all tableau $b \in F(s)$, there exists $m > 1$ such that $P(\bar{\xi}^m(b))$ is a column.

Definition 2.43. The charge is defined as follows. Let $b = c_\ell \otimes \cdots \otimes c_1 \in F(s)$.

- (1) If b is a column, set $\text{ch}(b) = 0$.
- (2) If b is a tableau, define $\text{ch}(b)$ by induction by setting

$$\text{ch}(b) = \text{ch}(\xi(b)) + |c_\ell| \quad \text{and} \quad \text{ch}(\text{red}(b)) = \text{ch}(b)$$

and using Lemma 2.42.

- (3) In general, set $\text{ch}(b) = \text{ch}(P(b))$.

We have seen in 2.2 how one can associate to any product $b^* = d_1 \otimes d_2 \in \dot{F}(s)$ of two columns a unique skew tableau T of minimal shape. Write $\text{leg}(b^*)$ for the number of boxes in the left column of T having no box to their right. The following definition of the *local energy* H is slightly different but equivalent to the original one (see [37] and also [40] for an equivalent definition in the case of rows and affine crystals).

Definition 2.44. Let $\ell = 2$ and $b^* = d_1 \otimes d_2 \in \dot{F}(s)$.

$$H(b^*) = \begin{cases} \text{leg}(b^*) & \text{if } |d_1| \geq |d_2| \\ H(R_1(b^*)) & \text{if } |d_1| < |d_2|. \end{cases}$$

Example 2.45. The minimal skew tableau corresponding to

$$b^* = \begin{bmatrix} 1 \\ 3 \\ 4 \\ 6 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 7 \\ 8 \end{bmatrix} \quad \text{is} \quad \begin{array}{cc} & 1 \\ & 3 \\ & 4 \\ 1 & 6 \\ 7 & \\ 8 & \end{array}, \quad \text{therefore we get } H(b^*) = 2.$$

Let $b^* = d_1 \otimes \cdots \otimes d_n \in \dot{F}(s)$. For all $1 \leq i < j \leq n$, denote

$$d_1 \otimes \cdots \otimes d_i \otimes d_{i+1}^{(j)} \otimes \cdots \otimes d_n$$

the element obtained from b^* by applying successively the R -matrices R_{j-1}, \dots, R_{i+1} .

Definition 2.46. The energy D is defined as follows. For $b^* = d_1 \otimes \cdots \otimes d_n \in \dot{F}(s)$, we set

$$D(b^*) = \sum_{1 \leq i < j \leq n} H(d_i \otimes d_{i+1}^{(j)}).$$

Remark 2.47. If $\dot{s}_1 = \cdots = \dot{s}_n$, then the R -matrices act as the identity, thus

$$D(b^*) = \sum_{1 \leq i \leq n-1} (n-i) H(d_i \otimes d_{i+1}).$$

Lemma 2.48. Let $b^* = d_1 \otimes \cdots \otimes d_n \in \dot{F}(s)$ be Yamanouchi such that $d_i = \{1, \dots, \ell\}$ for some $i = 1, \dots, n$. Then $D(b^*) = D(b^\times)$, where b^\times is the element obtained from b^* by deleting its leftmost column $\{1, \dots, \ell\}$.

Proof. Since D is invariant under the action of the combinatorial R -matrices, we can assume that $d_1 = \{1, \dots, \ell\}$. Then

$$D(b^*) = D(b^\times) + \sum_{i=1}^{n-1} (n-i) H(d_1 \otimes d_{i+1}).$$

But by definition of H , we have $H(d_1 \otimes d_{i+1}) = 0$ for any $i = 1, \dots, n-1$. \square

We are now going to prove the following proposition.

Proposition 2.49. *Let $b^* = d_1 \otimes \dots \otimes d_n \in \dot{F}(s)$ be Yamanouchi such that $d_i \neq \{1, \dots, \ell\}$ for all $i = 1, \dots, n$. Then $D(\text{pr}(b^*)) = D(b^*) - N$, where N is the number of ℓ in b^* .*

To do this, let us set for any $b^* = d_1 \otimes \dots \otimes d_n \in \dot{F}(s)$

$$\Delta_n(b^*) = D(b^*) - D(\text{pr}(b^*)).$$

When $n = 2$, one can use Definition 2.44 to compute Δ_2 . This immediately gives:

Lemma 2.50. *Assume $b^* = d_1 \otimes d_2$ with $h(d_1) \geq h(d_2)$. Then*

$$\Delta_2(d_1 \otimes d_2) = \begin{cases} -1 & \text{if } \ell \in d_1 \text{ and } \ell \notin d_2, \\ 1 & \text{if } \ell \notin d_1 \text{ and } \ell \in d_2, \\ 0 & \text{otherwise.} \end{cases}$$

Proof of Proposition 2.49. Since the combinatorial R -matrices preserve the energy, we can and shall assume that $h(d_1) \geq \dots \geq h(d_n)$. For any Yamanouchi vertex $b^* = d_1 \otimes \dots \otimes d_n \in \dot{F}(s)$, we have

$$\Delta_n(b^*) = \Delta_n(d_1 \otimes \dots \otimes d_n) = \Delta_{n-1}(d_1 \otimes \dots \otimes d_{n-1}) + \sum_{i=1}^{n-1} \Delta_2(d_i \otimes d_{i+1}^{(n)}).$$

because the promotion operator commutes with the R -matrices. Since b^* is Yamanouchi, the vertex $d_1 \otimes \dots \otimes d_{n-1}$ is also Yamanouchi. By induction on n , we can assume that $\Delta_{n-1}(d_1 \otimes \dots \otimes d_{n-1})$ is equal to the number of letters ℓ in $d_1 \otimes \dots \otimes d_{n-1}$. This shows that the proposition is in fact equivalent to the assertion

$$S_n(b^*) := \sum_{i=1}^{n-1} \Delta_2(d_i \otimes d_{i+1}^{(n)}) = \begin{cases} 1 & \text{if } \ell \in d_n, \\ 0 & \text{otherwise.} \end{cases}$$

for any Yamanouchi vertex $b^* = d_1 \otimes \dots \otimes d_n \in \dot{F}(s)$. Here again, we proceed by induction on n . When $n = 2$, this follows from Lemma 2.50. Note that we cannot have $\ell \in d_1$ in this case because we have assumed that $d_1 \neq \{1, \dots, \ell\}$. When $n \geq 3$, set

$$b^*(n-1) = d_1 \otimes \dots \otimes d_{n-2} \otimes d_{n-1}^{(n)}.$$

Then $b^*(n-1)$ is Yamanouchi and

$$S_n(b^*) = S_{n-1}(b^*(n-1)) + \Delta_2(d_{n-1} \otimes d_n).$$

Assume first $\ell \notin d_n$ and $\ell \notin d_{n-1}$. Then, we have $\ell \notin d_{n-1}^{(n)}$ and by the induction hypothesis $S_{n-1}(b^*(n-1)) = 0$. Also, $\Delta_2(d_{n-1} \otimes d_n) = 0$ and we thus get $S_n(b^*) = 0$ as desired. The arguments are similar for the three remaining cases: $\ell \in d_n$ and $\ell \notin d_{n-1}$, $\ell \notin d_n$ and $\ell \in d_{n-1}$, $\ell \in d_n$ and $\ell \in d_{n-1}$. \square

Combining Definition 2.43, the combinatorial description of the duality $*$, Lemma 2.48 and Proposition 2.49, we get the expected following theorem.

Theorem 2.51. *For all $b \in F(s)$, we have $\text{ch}(b) = D(b^*)$.*

3. AFFINE TYPE A DUALITY

The study of the combinatorics of highest weight representations for the quantum groups of affine type A has been initiated in [21] in the context of solvable lattice models. There, the affine Fock space plays a crucial role. It has been subsequently considered in various settings, ranging from mathematical physics to symmetric functions or finite group representation theory, see [27] for a good review. In this section, we examine the crystal combinatorics of the affine Fock space, and show that the results of Section 2 generalise to this setting.

3.1. Infinite columns. In this affine setting, columns are semi-infinite, i.e. of the form

$$c = \{x_1, \dots, x_k\} \sqcup \{x \in \mathbb{Z} \mid x < x_1\} \quad \text{and represented as } c = \begin{array}{|c|} \hline x_1 \\ \hline \vdots \\ \hline x_k \\ \hline \end{array},$$

for some $k \in \mathbb{Z}_{\geq 1}$ and some $x_1 < \dots < x_k \in \mathbb{Z}$. Note that this expression is not unique. For instance, the values $k = 2, x_1 = 1, x_2 = 3$ and $k = 3, x_1 = 0, x_2 = 1, x_3 = 3$ both yield the semi-infinite column $\{\dots, -2, -1, 0, 1, 3\}$. If we impose that x_1 is minimal in \mathbb{Z} with the condition that $x_1 + 1 \notin c$, then this expression becomes unique: we call it the *standard form* of c . Finally, we set

$$s(c) = x_1 + k - 1.$$

3.2. Affine combinatorial Fock space. Following [8, Chapter 6], highest weight crystals of type $A_{n-1}^{(1)}$ and level ℓ can be realised by using the combinatorics of symbols (or equivalently of abaci). By analogy with type A , we identify symbols with tensor products of columns (and abaci with semi-infinite binary matrices as in Remark 2.16).

Definition 3.1. Let $\mathbf{s} = (s_1, \dots, s_\ell) \in \mathbb{Z}^\ell$. The affine combinatorial Fock space of type $A_{n-1}^{(1)}$ and level ℓ associated to \mathbf{s} is

$$\widehat{F}(\mathbf{s}) = \{b = c_\ell \otimes \dots \otimes c_1 \mid s(c_j) = s_j \text{ for all } j = 1, \dots, \ell\}.$$

For all $b \in \widehat{F}(\mathbf{s})$, the tuple \mathbf{s} is called the *shape* of b ⁴.

We say that $b \in \widehat{F}(\mathbf{s})$ is in *standard form* if all columns have the same smallest entry and (at least) one of them is in standard form. For convenience, we represent negative entries $-x$ by \bar{x} .

Example 3.2. The representation

$$\begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 3 \\ \hline \end{array}$$

is the standard form of the element $b = c_3 \otimes c_2 \otimes c_1$ where $c_3 = \{\dots, -2, -1, 0, 1\}$, $c_2 = \{\dots, -2, -1, 2\}$, $c_1 = \{\dots, -2, -1, 0, 3\}$.

We now define affine analogues of tableaux in a natural way, by requesting semistandardness on a cylinder. We

use the following notation: if $c = \begin{array}{|c|} \hline x_1 \\ \hline \vdots \\ \hline x_k \\ \hline \end{array}$ is a column, set $c^+ = \begin{array}{|c|} \hline x_1 \\ \hline \vdots \\ \hline x_1 + n - 1 \\ \hline x_1 + n \\ \hline \vdots \\ \hline x_k + n \\ \hline \end{array}.$

Definition 3.3. An element $b = c_\ell \otimes \dots \otimes c_1 \in \widehat{F}(\mathbf{s})$ is called

- (1) *cylindric* if the top-aligned juxtaposition $c_\ell^+ c_1 \dots c_\ell$ is a semistandard Young tableau (with entries in \mathbb{Z}).
- (2) *anticylindric* if the bottom-aligned juxtaposition $c_1 \dots c_\ell c_1^+$ is a semistandard skew Young tableau (with entries in \mathbb{Z}).

Example 3.4. Take $\ell = 3, n = 2$.

- (1) The element

$$\begin{array}{|c|} \hline \bar{1} \\ \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 1 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 2 \\ \hline \end{array}$$

is cylindric because

$$\begin{array}{|c|c|c|c|} \hline \bar{1} & \bar{1} & \bar{1} & \bar{1} \\ \hline 0 & 0 & 1 & 1 \\ \hline 1 & 2 & 4 & \\ \hline 3 & & & \\ \hline \end{array}$$

is a semistandard Young tableau.

⁴ The element \mathbf{s} is usually called the (multi)charge in the literature, see [21]. We choose this alternative terminology for two reasons : by analogy with the finite case (Section 2.1), since \mathbf{s} plays a similar role to the partition λ (in fact, \mathbf{s} can itself be represented by a partition as in the original approach [21, Definition 3.1]); and because “charge” has already been introduced with a different meaning in Section 2.8.

(2) The element

$$\begin{array}{|c|} \hline \bar{2} \\ \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{2} \\ \hline 0 \\ \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{2} \\ \hline \bar{1} \\ \hline 0 \\ \hline 2 \\ \hline \end{array}$$

is anticylindric because

$$\begin{array}{|c|c|c|c|} \hline & & & \bar{2} \\ \hline & & \bar{2} & \bar{1} \\ \hline \bar{2} & \bar{2} & \bar{1} & 0 \\ \hline \bar{1} & 0 & 0 & 1 \\ \hline 0 & 1 & 1 & 2 \\ \hline 2 & 3 & 3 & 4 \\ \hline \end{array}$$

is a semistandard skew Young tableau.

Remark 3.5. Note that the elements of $\widehat{F}(\mathbf{s})$ can alternatively be described as *charged multipartitions*: they form the standard basis of the (algebraic) Fock space, see [8, Chapter 6], [42]. More precisely, tensor products of columns are the representation of charged multipartitions by their β -sets, also known as Lusztig's symbols [8, Chapter 5]. In turn, tensor products that are cylindric correspond to *cylindric multipartitions*, which are important objects in algebraic combinatorics. They were first introduced by Gessel and Krattenthaler [14], and have since then been used in combinatorics, representation theory, and mathematical physics, see [5] and [12] for more details.

Further, we give an affine analogue of the Yamanouchi property. For this, given an element $b = c_\ell \otimes \cdots \otimes c_1 \in \widehat{F}(\mathbf{s})$, we define the first n -period of b to be the sequence $P_1 = (x, x-1, \dots, x-n+1)$ of entries of b such that x is maximal in b and such that for all $i = 0, \dots, n-1$, $x-i \in c_{k_i}$ where $k_0 \leq \dots \leq k_{n-1}$ with k_i maximal. One then defines the r -th n -period of b by induction, by setting P_r to be the first period of the element $b^{(r-1)}$ obtained from b by removing P_1, \dots, P_{r-1} if they all exist.

Definition 3.6. An element $b \in \widehat{F}(\mathbf{s})$ is called *totally n -periodic* if the r -th n -period of b exists for all $r \in \mathbb{Z}_{\geq 1}$.

Example 3.7. Take $n = 3$ and $\ell = 4$. Then the following element is totally periodic

$$\begin{array}{|c|} \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 0 \\ \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 0 \\ \hline 2 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 0 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline \end{array},$$

where we have highlighted the first period in blue, the second in purple and the third in yellow (and the periods P_r for $r > 3$ then obviously exist).

3.3. Affine crystal structures. In the following, for $\mathbf{s} = (s_\ell, \dots, s_1) \in \mathbb{Z}^\ell$ and $\dot{\mathbf{s}} = (\dot{s}_1, \dots, \dot{s}_n) \in \mathbb{Z}^n$, denote

$$b_{\mathbf{s}} = \boxed{s_\ell} \otimes \cdots \otimes \boxed{s_1} \quad \text{and} \quad \dot{b}_{\dot{\mathbf{s}}} = \boxed{\dot{s}_1} \otimes \cdots \otimes \boxed{\dot{s}_n}.$$

Moreover, define analogues of the sets $\mathcal{S}(s)$ and $\dot{\mathcal{S}}(s)$ of Section 2 by considering

$$\mathcal{D}(s) = \left\{ (s_\ell, \dots, s_1) \in \mathbb{Z}^\ell(s) \mid s_\ell \leq \cdots \leq s_1 \leq s_\ell + n \right\} \quad \text{and} \quad \dot{\mathcal{D}}(s) = \left\{ (\dot{s}_1, \dots, \dot{s}_n) \in \mathbb{Z}^\ell(s) \mid \dot{s}_1 \leq \cdots \leq \dot{s}_n \leq \dot{s}_1 + \ell \right\},$$

and set

$$\mathcal{D} = \bigsqcup_{s \in \mathbb{Z}} \mathcal{D}(s) \quad \text{and} \quad \dot{\mathcal{D}} = \bigsqcup_{s \in \mathbb{Z}} \dot{\mathcal{D}}(s).$$

As in Section 2.2, the combinatorial Fock space $\widehat{F}(\mathbf{s})$ can be endowed with the structure of an $A_{n-1}^{(1)}$ -crystal, where the action of the lowering crystal operators f_i is given by changing a certain entry x such that $x \bmod n = i$ to $x+1$, see for instance [8, Chapter 6]. Also, the connected component containing $b_{\mathbf{s}}$ is isomorphic to the crystal $\widehat{B}(\mathbf{s})$ of the irreducible highest weight module of highest weight $\omega_{\mathbf{s}} = \omega_{s_\ell} + \cdots + \omega_{s_1}$ where $\omega_0, \dots, \omega_{n-1}$ are the fundamental weights of type $A_{n-1}^{(1)}$. Let us recall the crystal structure on $\widehat{F}(\mathbf{s})$, due to [21], [4], [42]. We construct the word $w_i(b)$ as in Section 2.2, by the following procedure:

- (1) For all $k \in \mathbb{Z}$, let $w_i^{(k)}(b)$ be the word obtained by reading the entries $i+kn$ and $i+kn+1$ in b from left to right (i.e., going through the columns c_ℓ, \dots, c_1).
- (2) Set $w_i(b)$ to be the concatenation of the words $w_i^{(k)}(b)$ ordered by decreasing values of k .

Remark 3.8. Because of the semi-infinite form of each column, it is enough to consider only the integers k such that $i + kn$ appears in the standard form of b . That way, we construct $w_i(b)$ from finitely many subwords $w_i^{(k)}$.

Example 3.9. For $n = 4, i = 3, \ell = 3$ and $b = \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline 4 \\ \hline 7 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 4 \\ \hline 6 \\ \hline 7 \\ \hline \end{array} = c_3 \otimes c_2 \otimes c_1$, where we highlighted the entries with

relevant residues with boldface, we have $w_3^{(1)}(b) = 77$ and $w_3^{(0)}(b) = 3344$, therefore $w_3(b) = 773344$.

As in type A , this induces a word in the symbols $+$ and $-$ by encoding each $i + kn$ by $+$ if and each $i + kn + 1$ by $-$. Deleting all factors $+-$ recursively yields a word called the i -signature of b . With these conventions, one checks that the following result is equivalent to [8, Theorem 6.2.12].

Theorem 3.10. *The set $\hat{F}(\mathbf{s})$ is endowed with an $A_{n-1}^{(1)}$ -crystal structure, where*

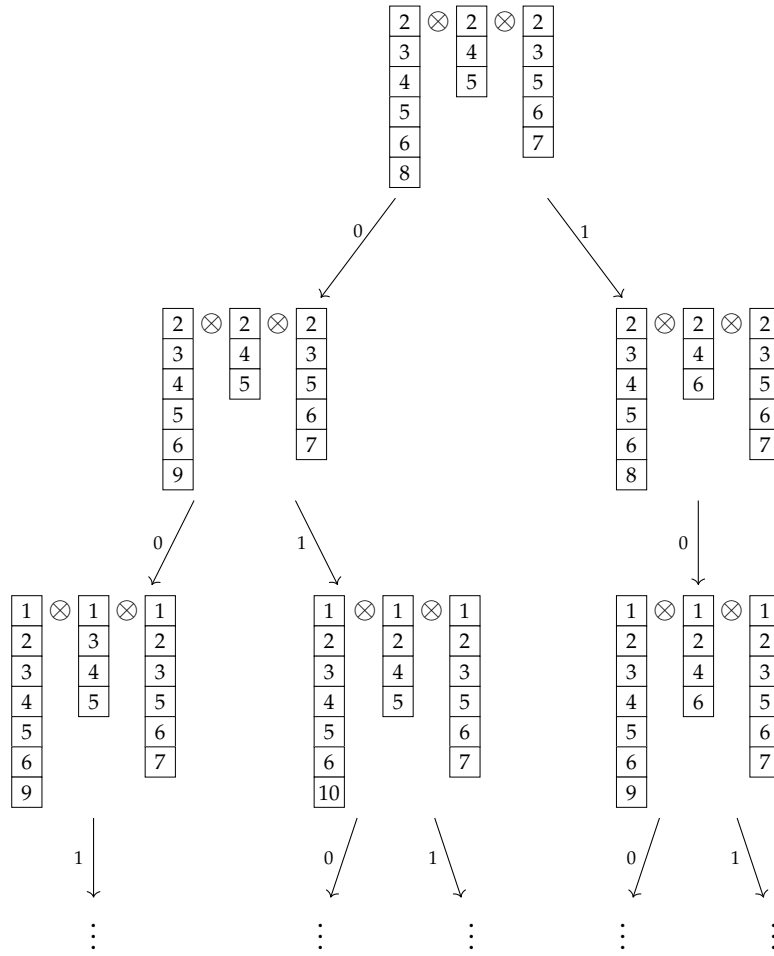
- (1) *The lowering crystal operators f_i act on $b \in \hat{F}(\mathbf{s})$ by changing the entry x corresponding to the leftmost $+$ in the i -signature of b into $x + 1$ if it exists; and by 0 otherwise.*
- (2) *The raising crystal operators f_i act on $b \in \hat{F}(\mathbf{s})$ by changing the entry $x + 1$ corresponding to the rightmost $-$ in the i -signature of b into x if it exists; and by 0 otherwise.*

Example 3.11. In the previous example, we first get the sequence $++++--$, and the signature is $++$. The leftmost $+$ corresponds to the 7 in c_2 , and therefore f_3 acts on b by changing that 7 into an 8, i.e.

$$f_3 b = \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline 8 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 1 \\ \hline 4 \\ \hline 6 \\ \hline 7 \\ \hline \end{array}.$$

Example 3.12. Take $\ell = 3, n = 2, \mathbf{s} = (s_3, s_2, s_1) = (7, 4, 6)$ and $b = c_3 \otimes c_2 \otimes c_1 = \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline 6 \\ \hline 8 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline 5 \\ \hline 6 \\ \hline 7 \\ \hline \end{array}$. One checks that

$e_0 b = e_1 b = e_2 b = 0$, i.e. b is a highest weight vertex, and that the beginning of the connected component of the $A_{n-1}^{(1)}$ -crystal containing b is equal to



Remark 3.13. If n is sufficiently large, any column (in standard form) has at most one entry with residue i or $i + 1$, and therefore one recovers the finite type A crystal rule (forgetting the arrows indexed by 0).

In fact, we have

$$\hat{F}(\mathbf{s}) \simeq \hat{F}((s_\ell)) \otimes \cdots \otimes \hat{F}((s_1)).$$

Remark 3.14. Note that the rule used to compute the affine crystal above is not the tensor product rule, which means that the above isomorphism is not an equality in general. However, we recover the tensor product rule (and hence the equality) if the differences $s_{j+1} - s_j$ are sufficiently large.

Similarly, for any n -tuple of integers $\dot{\mathbf{s}} = (\dot{s}_1, \dots, \dot{s}_n)$, there is a level n $A_{\ell-1}^{(1)}$ -crystal structure on the set of tensor products of n columns

$$\hat{F}(\dot{\mathbf{s}}) = \{a = d_1 \otimes \cdots \otimes d_n \mid s(d_i) = \dot{s}_i \text{ for all } i\}.$$

In order for these crystals structures to commute (see Section 3.6), we need to use a slightly different rule for computing the $A_{\ell-1}^{(1)}$ -crystal. More precisely, we construct a word $\dot{w}_j(a)$ by the following procedure:

- (1) For all $k \in \mathbb{Z}$, let $\dot{w}_j^{(k)}(a)$ be the word obtained by reading the entries $j + k\ell$ and $j + k\ell + 1$ in a from left to right (i.e., going through the columns d_1, \dots, d_n).
- (2) Set $\dot{w}_j(a)$ to be the concatenation of the words $\dot{w}_j^{(k)}$ ordered by increasing values of k .

Now, we compute the j -signature by the exact same procedure as for the $A_{n-1}^{(1)}$ -crystal, and the action of the crystal operators \dot{f}_j and \dot{e}_j is given by the same rule as in Theorem 3.10. We obtain again

$$\hat{F}(\mathbf{s}) \simeq \hat{F}((\dot{s}_1)) \otimes \cdots \otimes \hat{F}((\dot{s}_n)).$$

Example 3.15. For $\ell = 3, j = 0, n = 4$ and $a = \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline \end{array} = d_1 \otimes d_2 \otimes d_3 \otimes d_4$, where we highlighted

$$\begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline \end{array}$$

the entries with relevant residues with boldface, we have $\dot{w}_0^{(0)}(a) = 0101001$ and $\dot{w}_0^{(1)}(a) = 3434$, therefore $\dot{w}_0(a) = 01010013434$. This yields the sequence $+ - + - + + - + -$, so the 0-signature is $+$, in which the leftmost $+$ corresponds to the 0 in d_3 . Therefore, \dot{f}_0 acts on a by changing that 0 into a 1, i.e.

$$\dot{f}_0 a = \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline \end{array}.$$

In other terms, in the crystal $\hat{F}(s)$, we have an arrow

$$\begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline \end{array} \xrightarrow{0} \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 4 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline 4 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline \end{array}.$$

Let us set $\hat{F}(s) = \bigoplus_{\mathbf{s} \in \mathbb{Z}^\ell(s)} \hat{F}(\mathbf{s})$ and $\hat{F}(s) = \bigoplus_{\mathbf{s} \in \mathbb{Z}^n(s)} \hat{F}(\mathbf{s})$. We can now state an analogue of Theorem 2.11.

Theorem 3.16.

- (1) The set of cylindric elements of $\hat{F}(\mathbf{s})$ is stable under the crystal operators $f_i, e_i, i = 0, \dots, n-1$. Moreover, for any $b \in \hat{F}(\mathbf{s})$, there is a cylindric element $\mathcal{P}(b) \in \hat{F}(\mathbf{s})$ such that $\mathcal{P}(b)$ and b appear at the same place in two isomorphic connected components of $\hat{F}(\mathbf{s})$.
- (2) An element $b \in \hat{F}(s)$ is a highest weight vertex in the $A_{n-1}^{(1)}$ -crystal if and only if b is totally periodic.

Proof. One checks that Definition 3.3 is equivalent to the definition of cylindricity used in [41] and [9]. Therefore, the first statement of (1) translates to [41, Section 3], see also [9, Proposition 4.10]. A constructive proof of the second statement of (1) can be derived from [9, Section 4.3]. Part (2) was obtained in [19, Theorem 5.9]. \square

Remark 3.17. Note that $\mathcal{P}(b)$ is not unique, unlike the tableau $P(b)$ of Theorem 2.11. There is uniqueness by putting some extra constraints on $\mathcal{P}(b)$, see [9] for more details.

3.4. Uglov's duality. There is a duality

$$\begin{array}{ccc} \hat{F}(s) & \longleftrightarrow & \hat{F}(s) \\ b & \longleftrightarrow & b^* \end{array}$$

defined in [42, Remark 4.2], see also [44, Section 1.1.5]. It is best explained in terms of binary matrices. We first encode $b = c_\ell \otimes \dots \otimes c_1 \in \hat{F}(s)$ by the $\infty \times \ell$ matrix M defined by

$$M_{i,j} = \begin{cases} 1 & \text{if } i \in c_j \\ 0 & \text{otherwise} \end{cases}$$

Now, for all $k \in \mathbb{Z}$, consider the submatrices $M^{(k)}$ of size $n \times \ell$ of M defined by, for all $1 \leq i \leq n$,

$$M_{i,j}^{(k)} = M_{(k-1)n+i,j}.$$

Set $N^{(k)} = M^{(k)\text{tr}}$. Then the $\infty \times n$ matrix N defined by, for all $j \in \mathbb{Z}$, say $(k-1)\ell + 1 \leq j \leq k\ell$,

$$N_{j,i} = N_{j-(k-1)\ell,i}^{(k)}$$

decodes to an element b^* of $\hat{F}(s)$.

Example 3.18. Take $\ell = 2, n = 3$, and $b = \begin{bmatrix} \overline{2} \\ 0 \\ 1 \\ 3 \end{bmatrix} \otimes \begin{bmatrix} \overline{2} \\ \overline{1} \\ 2 \end{bmatrix}$. Then we have

$$\begin{array}{c} \vdots \\ -5 \\ -4 \\ -3 \\ -2 \\ -1 \\ 0 \\ 1 \\ 2 \\ 3 \\ 4 \\ 5 \\ 6 \\ \vdots \end{array} \begin{array}{c} \vdots \\ \begin{bmatrix} 1 & 1 \end{bmatrix} \\ \begin{bmatrix} 1 & 1 \end{bmatrix} \\ \begin{bmatrix} 1 & 1 \end{bmatrix} \\ \begin{bmatrix} 1 & 1 \end{bmatrix} \\ \begin{bmatrix} 1 & 0 \end{bmatrix} \\ \begin{bmatrix} 0 & 1 \end{bmatrix} \\ \begin{bmatrix} 0 & 1 \end{bmatrix} \\ \begin{bmatrix} 1 & 0 \end{bmatrix} \\ \begin{bmatrix} 1 & 0 \end{bmatrix} \\ \begin{bmatrix} 0 & 1 \end{bmatrix} \\ \begin{bmatrix} 0 & 0 \end{bmatrix} \\ \begin{bmatrix} 0 & 0 \end{bmatrix} \\ \begin{bmatrix} 0 & 0 \end{bmatrix} \\ \vdots \end{array} \begin{array}{c} \\ M^{(-1)} \\ \\ \\ M^{(0)} \\ \\ M^{(1)} \\ \\ M^{(2)} \\ \\ \\ \end{array} \quad \text{and we get} \quad N = \begin{array}{c} \vdots \\ -3 \\ -2 \\ -1 \\ 0 \\ 1 \\ 2 \\ 3 \\ 4 \\ \vdots \end{array} \begin{array}{c} \vdots \\ \begin{bmatrix} 1 & 1 & 1 \end{bmatrix} \\ \begin{bmatrix} 1 & 1 & 1 \end{bmatrix} \\ \begin{bmatrix} 1 & 1 & 0 \end{bmatrix} \\ \begin{bmatrix} 1 & 0 & 1 \end{bmatrix} \\ \begin{bmatrix} 0 & 1 & 0 \end{bmatrix} \\ \begin{bmatrix} 1 & 0 & 1 \end{bmatrix} \\ \begin{bmatrix} 0 & 0 & 0 \end{bmatrix} \\ \begin{bmatrix} 0 & 0 & 0 \end{bmatrix} \\ \vdots \end{array} \begin{array}{c} \\ M^{(-1)\text{tr}} = N^{(-1)} \\ M^{(0)\text{tr}} = N^{(0)} \\ M^{(1)\text{tr}} = N^{(1)} \\ M^{(2)\text{tr}} = N^{(2)} \\ \\ \end{array} ,$$

where we indicated the rows and column indices. This gives $b^* = \begin{bmatrix} \overline{2} \\ \overline{1} \\ 0 \\ 2 \end{bmatrix} \otimes \begin{bmatrix} \overline{2} \\ \overline{1} \\ 1 \\ 2 \end{bmatrix} \otimes \begin{bmatrix} \overline{2} \\ 0 \\ 2 \end{bmatrix}$.

Remark 3.19. Assume that each column $c = \{x_1, \dots, x_k\}$ of b written in standard form verifies $0 \leq x_1 < \dots < x_k \leq n$.

- (1) Then b is determined by $M^{(1)}$, and the duality \star is just transposing this single matrix. We recover precisely the duality \star , see Remark 2.16. For instance, for $\ell = 4, n = 3$ and $b = \begin{bmatrix} 0 \\ 2 \\ 3 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 3 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \\ 3 \end{bmatrix}$. The corresponding

matrix is

$$\begin{bmatrix} 1 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 1 & 1 \end{bmatrix}, \quad \text{whose transpose is} \quad \begin{bmatrix} 1 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 1 \end{bmatrix},$$

$$\text{so we get } b^* = \begin{bmatrix} 0 \\ 1 \\ 2 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 4 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \\ 3 \\ 4 \end{bmatrix}.$$

- (2) Moreover, the orders \prec and \rightarrow used to compute the crystals coincide, in analogy with the finite case again.

The first result concerning the duality \star is an analogue of Proposition 2.17. The following result can be checked purely combinatorially. We omit its proof because it is essentially equivalent to [12, Theorem 3.3] in the convention of the present paper.

Proposition 3.20. *Let $b \in \widehat{F}(\mathfrak{s})$. Then b is totally periodic if and only if b^* is anticylindric.*

Remark 3.21. Similarly, there is a notion of antiperiods that yield a characterisation of highest weight vertices in $\widehat{F}(\mathfrak{s})$. We get that b is cylindric if and only if b^* is totally antiperiodic, completing the previous proposition

Now, we would like to establish an analogue of Theorem 2.19. Let \widehat{S}_n be the affine symmetric group on n elements, which is the Weyl group of type $A_{n-1}^{(1)}$. We denote by $\sigma_i, i = 0, \dots, n-1$ the involutions generating \widehat{S}_n , subject to the usual braid relations modulo n . The group \widehat{S}_n acts on the crystal $\widehat{F}(\mathfrak{s})$ as in finite type. Similarly, we denote $\sigma_j, j = 0, \dots, \ell-1$ the generators of \widehat{S}_ℓ . First of all, we focus on the maps \dot{e}_j and $\dot{\sigma}_j$ but only in the case $j = 1, \dots, \ell-1$. The case $j = 0$ will be treated in upcoming Section 3.5 using the promotion operator. As in

Section 2.3, for $j = 1, \dots, \ell - 1$, we denote by J_j the elementary Jeu de Taquin map on $\hat{F}(s)$ between columns j and $j + 1$, and by R_j the combinatorial R -matrix on $\hat{F}(s)$ realising the isomorphism of crystals

$$F((s_1, \dots, s_j, s_{j+1}, \dots, s_\ell)) \xrightarrow{\sim} F((s_1, \dots, s_{j+1}, s_j, \dots, s_\ell)).$$

Theorem 3.22. *For all $j = 1, \dots, \ell - 1$, we have*

$$(1) \quad \dot{e}_j^* = J_j \quad \text{and} \quad (2) \quad \dot{\sigma}_j^* = R_j.$$

Proof. The proof is analogous to that of Theorem 2.19. In particular, (1) is immediate. For (2), an explicit formula for R_j was given in [18, Proposition 5.2.2]. This formula coincides with the formula for the R -matrix in finite type which can be found in [40, Example 4.10]. Therefore, we conclude using Theorem 2.19. \square

3.5. Affine cyclage and promotion. We now show that the results of Section 2.7 generalise to the affine case. Consider the *affine cyclage* map

$$\begin{aligned} \hat{\xi} : \quad \hat{F}(s) &\longrightarrow F(s + n) \\ c_\ell \otimes \dots \otimes c_1 &\longmapsto c_{\ell-1} \otimes \dots \otimes c_1 \otimes c_\ell^+ \end{aligned}$$

where c_ℓ^+ at the beginning of 3.2. By [18, Proposition 5.2.1], $\hat{\xi}$ is an isomorphism of $A_{n-1}^{(1)}$ -crystals. Now, consider the promotion map

$$\begin{aligned} \text{pr} : \quad \hat{F}(s) &\longrightarrow \hat{F}(s + n) \\ d_1 \otimes \dots \otimes d_n &\longmapsto \text{pr}(d_1) \otimes \dots \otimes \text{pr}(d_n) \end{aligned}$$

where $\text{pr}(d_i) = \{k + 1; k \in d_i\}$ for all $i = 1, \dots, n$. Exactly as in Proposition 2.36, we have the following result, which is immediate by definition of \star .

Proposition 3.23. *For all $b \in \hat{F}(s)$,*

$$\hat{\xi}(b)^\star = \text{pr}(b^\star).$$

As a direct consequence, we get a description of the action of the maps \dot{e}_0 and $\dot{\sigma}_0$ directly on $\hat{F}(s)$. This completes the statement of Theorem 3.22.

Corollary 3.24.

$$(1) \quad \dot{e}_0^* = \hat{\xi}^{-1} \circ J_1 \circ \hat{\xi} \quad \text{and} \quad (2) \quad \dot{\sigma}_0^* = \hat{\xi}^{-1} \circ R_1 \circ \hat{\xi}.$$

By analogy with Theorem 3.22, we denote $R_0 = \dot{\sigma}_0^*$. If c_j is a column, denote by c_j^- the column such that $(c_j^-)^+ = c_j$. Since $\hat{\xi}^{-1}(c_\ell \otimes \dots \otimes c_1) = c_1^{-1} \otimes c_\ell \otimes \dots \otimes c_2$, we have

$$R_0(c_\ell \otimes \dots \otimes c_1) = c_1^- \otimes c_{\ell-1} \otimes \dots \otimes c_2 \otimes c_\ell^+.$$

3.6. Triple crystal structure. The affine duality \star is more complex than its finite type counterpart $*$, and this is related to the existence of third crystal structure on $\hat{F}(s)$ [10]. This arises from the action of a Heisenberg algebra \mathcal{H} , and the resulting crystal is of type A_∞ . More precisely, each $b \in \hat{F}(s)$ writes uniquely on the form $b = a_\kappa(\bar{b})$ for some partition κ and some highest weight vertex \bar{b} for the \mathcal{H} -structure. Here a_κ is the so-called *Heisenberg crystal operator* corresponding to the partition κ . The explicit action of a_κ is described in [11] in terms of translating (generalised) n -periods. By writing $\kappa(b) = \kappa$, the map $b \mapsto \kappa(b)$ yields a bijection between the connected component containing b and the Young lattice \mathcal{Y} . Now, the Young lattice carries a crystal structure corresponding to the basic representation of type A_∞ . More precisely, there is an arrow $\lambda \xrightarrow{k} \mu$ in \mathcal{Y} with $k \in \mathbb{Z}$ if and only if μ/λ has only one box of content k . This endows $\hat{F}(s)$ with the structure of an A_∞ -crystal. The following result is [10, Theorem 6.17].

Theorem 3.25. *The three crystals commute⁵, i.e. for all $i = 1, \dots, n - 1$, for all $j = 1, \dots, \ell - 1$ and for all partition κ , we have*

$$f_i \dot{f}_j^* = \dot{f}_j^* f_i, \quad a_\kappa f_i = f_i a_\kappa, \quad \text{and} \quad \dot{f}_j^* a_\kappa = a_\kappa \dot{f}_j^*.$$

Recall that we have introduced the sets $\mathcal{D}(s)$, $\mathcal{D}(s)$ and the notation b_s, \dot{b}_s for $\mathbf{a} \in \mathcal{D}(s), \dot{\mathbf{s}} \in \mathcal{D}(s)$ in Section 3.3. Write respectively $\dot{\omega}_0, \dots, \dot{\omega}_{\ell-1}$ and $\dot{\delta}$ the fundamental weights and the null root for the root system of type $A_{\ell-1}^{(1)}$. The definition of \star implies the following important property.

⁵ Note that in [10], the commutation was proved for a twisted version of this duality. This is accounted for here by enumerating columns in the reverse order in $\hat{F}(s)$, and using the different ordering.

Lemma 3.26. *Let $b \in \widehat{F}(s)$. We have $b = b_s$ for some $s \in \mathcal{D}(s)$ if and only if $b^* = \dot{b}_{s^*}$ for some $s^* \in \dot{\mathcal{D}}(s)$. Moreover, the $A_{\ell-1}^{(1)}$ -dominant weight corresponding to \dot{b}_{s^*} has the form*

$$(n - s_1 + s_\ell)\dot{\omega}_0 + \sum_{j=1}^{\ell-1} (s_j - s_{j+1})\dot{\omega}_j + k\delta$$

where k is an integer.

Example 3.27. Take $\ell = 4, n = 3$ and $s = (-2, -1, -1, 1) \in \mathcal{D}(-3)$. Then

$$b_s = \begin{bmatrix} 2 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 2 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 1 \end{bmatrix}, \quad \text{i.e.} \quad b_s = \begin{bmatrix} 2 \\ 2 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 2 \\ 2 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 2 \\ 2 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 2 \\ 1 \\ 0 \\ 1 \end{bmatrix}.$$

The corresponding matrix is

$$M = \begin{matrix} & \vdots & \\ \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 0 \\ 1 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} & \begin{matrix} M^{(0)} \\ \\ M^{(1)} \end{matrix} & \text{which gives} & M^* = \begin{bmatrix} 1 & 1 & 1 \\ 1 & 1 & 0 \\ 1 & 1 & 0 \\ 1 & 0 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} & \begin{matrix} M^{(0)\text{tr}} \\ \\ M^{(1)\text{tr}} \end{matrix} \end{matrix},$$

therefore

$$b_s^* = \begin{bmatrix} 3 \\ 2 \\ 1 \\ 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 3 \\ 2 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 3 \\ 2 \\ 1 \end{bmatrix} = \dot{b}_{s^*} \quad \text{with } s^* = (1, -1, -3) \in \dot{\mathcal{D}}(-3).$$

Remark 3.28. We see that the map $s \mapsto s^*$ is an analogue of the transposition of partitions used in Section 2.

The commutation of the three crystals in Theorem 3.25 induces an $(A_{n-1}^{(1)} \times A_\infty \times A_{\ell-1}^{(1)})$ -crystal structure on $\widehat{F}(s)$, and as in the finite case (see Corollary 2.22), each connected component of the tricrystal of $\widehat{F}(s)$ has a unique source vertex. The following corollary is the translation of [10, Theorem 6.19] in our terminology.

Corollary 3.29. *Let $b \in \widehat{F}(s)$. Then b is a highest weight vertex in the $(A_{n-1}^{(1)} \times A_\infty \times A_{\ell-1}^{(1)})$ -crystal if and only if $b = b_s$ for some $s \in \mathcal{D}(s)$.*

To be complete, we give a characterisation of the highest weight vertices in the different bicrystals. In order to do this, we consider the vertices of $\widehat{B}(s)$, the connected component of the $A_{n-1}^{(1)}$ -crystal containing b_s , where $s \in \mathcal{D}(s)$. These are called n -FLOTW elements, and they have an explicit combinatorial description, see [8, Definition 5.7.8]⁶. In particular, FLOTW elements are cylindric, which is immediate from Theorem 3.16 because b_s is.

Theorem 3.30.

- (1) b is a highest weight vertex in the $(A_{n-1}^{(1)} \times A_{\ell-1}^{(1)})$ -crystal if and only if b is cylindric and totally periodic.
- (2) b is a highest weight vertex in the $(A_{n-1}^{(1)} \times A_\infty)$ -crystal if and only if b^* is ℓ -FLOTW.

Proof. Assertion (1) follows from the characterisation of the highest weight vertices in the $A_{n-1}^{(1)}$ -crystal and the $A_{\ell-1}^{(1)}$ -crystal given in Theorem 3.16(2) and Proposition 3.20. Assertion (2) is proved in [10]. \square

⁶ Note that FLOTW elements are originally defined for a more constrained condition on s , but it is easy to see that our condition induces the same combinatorial characterisation

Remark 3.31. Note that it is more challenging to give a simple description of the $(A_{n-1}^{(1)} \times A_\infty)$ -highest weight vertices. Nevertheless, a characterisation of the A_∞ -highest weight vertices has been given in [13, Theorem 5.1], but this does not yield to an analogue of the previous theorem.

By Theorem 3.25 and Corollary 3.29, for each $b \in \widehat{F}(s)$, there is a unique $\mathbf{s} \in \mathcal{D}(s)$ such that $b = f_{\underline{i}}^* a_\kappa f_{\underline{j}} b_{\mathbf{s}}$ for some $\underline{i}, \underline{j}$ and some partition κ . Set

$$\mathcal{P}(b) = f_{\underline{i}} b_{\mathbf{s}}, \quad \kappa(b) = \kappa, \quad \text{and} \quad \mathcal{Q}(b) = f_{\underline{j}} b_{\mathbf{s}}^*.$$

Note that the elements $\mathcal{P}(b)$ and $\mathcal{Q}(b)$ are FLOTW by definition. Therefore, we get an analogue of Theorem 2.25, yielding an affine crystal version of the RSK correspondence.

Theorem 3.32. *The assignment*

$$\widehat{\Phi} : b \mapsto (\mathcal{P}(b), \kappa(b), \mathcal{Q}(b))$$

yields a bijection $\widehat{F}(s) \rightarrow \widehat{\Phi}(\widehat{F}(s))$. In particular, we have the decomposition

$$\widehat{F}(s) \simeq \bigoplus_{\mathbf{s} \in \mathcal{D}(s)} \widehat{B}(\mathbf{s}) \otimes \mathcal{Y} \otimes \widehat{B}(\mathbf{s}^*).$$

Remark 3.33. Another way to express the affine crystal RSK correspondence is to consider the bijection

$$b \longleftrightarrow (\mathcal{P}(b), \mathcal{Q}(b))$$

where $\mathcal{Q}(b)$ is defined as before, and $\mathcal{P}(b) = a_\kappa f_{\underline{i}} b_{\mathbf{s}}$ (in particular $\mathcal{P}(b)$ is cylindric). In a symmetric fashion, one can establish a bijection

$$b \longleftrightarrow (\mathcal{P}(b), \mathcal{Q}(b))$$

where $\mathcal{Q}(b)$ is anticylindric. It would be interesting to compare this with the results of [25].

3.7. Bicrystal structure on self-dual elements. As in Section 2.5, let us consider the set of self-dual elements $\widehat{F}(s)^*$, that is, the set of all $b = c_1 \otimes \cdots \otimes c_\ell \in F(s)$ such that $b^* = c_\ell \otimes \cdots \otimes c_1$. Again, self-dual elements exist only if $n = \ell$. Moreover, by [10, Proposition 5.2], for all $b \in \widehat{F}(s)$, we have $\kappa(b^*) = \kappa(b)^{\text{tr}}$, therefore if $b \in \widehat{F}(s)^*$, we have

$$\kappa(b) = \kappa(b)^{\text{tr}}.$$

Now, by [23], self-transpose partitions realise the crystal graph $\widetilde{\mathcal{Y}}$ of the basic representation of type B_∞ , by setting

$$\kappa \xrightarrow{k} \kappa' \text{ in } \widetilde{\mathcal{Y}} \quad \text{if and only if} \quad \kappa \xrightarrow{k} \kappa'' \xrightarrow{-k} \kappa' \text{ in } \mathcal{Y}.$$

If one prefers, one can also consider the vertices of $\widetilde{\mathcal{Y}}$ as shifted Young diagrams since self-conjugate partitions are in bijection with strict partitions. Further, for all $i = 1, \dots, n-1$, set again

$$f_i^* = f_i^* f_i.$$

Therefore, the operators f_i^* for $i = 1, \dots, n-1$ and a_κ for $\kappa = \kappa^{\text{tr}}$ induce an $(A_{n-1}^{(1)} \times B_\infty)$ -crystal structure on $\widehat{F}(s)^*$. More precisely, we get the decomposition

$$\widehat{F}(s)^* \simeq \bigoplus_{\substack{\mathbf{s} \in \mathcal{D}(s) \\ \mathbf{s}^* = \mathbf{s}}} \widehat{B}(\mathbf{s}) \otimes \widetilde{\mathcal{Y}}.$$

Other multicrystal structures on fixed points sets will be studied in Section 5.

3.8. Affine keys and bikeys. We now construct affine analogues of keys and bikeys as defined in Section 2.6.

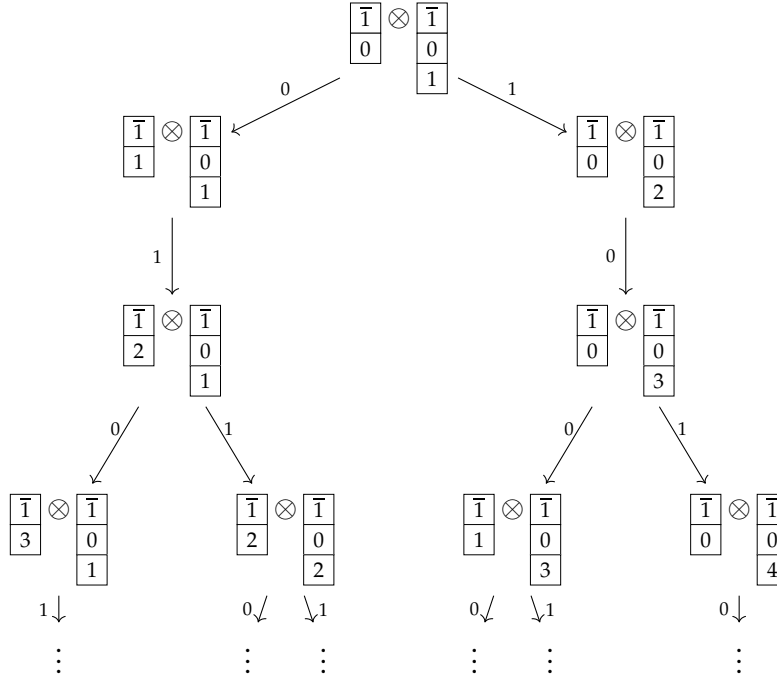
Definition 3.34. *An element $b = c_\ell \otimes \cdots \otimes c_1 \in \widehat{F}(s)$ is called an affine key if b is cylindric and $c_\ell \subseteq \cdots \subseteq c_1 \subseteq c_\ell^+$.*

Remark 3.35. Note that [20] uses the terminology (n, \mathbf{s}) -cores instead of affine bikeys. Indeed, it is observed that if $b = c_\ell \otimes \cdots \otimes c_1$ is an affine key, then each c_j is in particular the beta-set of a partition which is an n -core.

In particular, if $b \in \widehat{F}(s)$ is an affine key, b is cylindric and therefore $s(b) \in \mathcal{D}(s)$. Now, for an element $b \in \widehat{F}(s)$, denote $\mathcal{O}_{\widehat{S}_n}(b)$ the orbit of b under the action of \widehat{S}_n . The following result is an analogue of Proposition 2.27 and is a reformulation of [20, Proposition 5.14].

Proposition 3.36. *Let $\mathbf{s} \in \mathcal{D}(s)$. The set of all affine keys in $\widehat{F}(\mathbf{s})$ is equal to $\mathcal{O}_{\widehat{S}_n}(b_{\mathbf{s}})$.*

Example 3.37. Take $\ell = n = 2$ and $\mathbf{s} = (r_1, r_2) = (1, 0)$, so that $b_{\mathbf{s}} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \end{bmatrix}$. Let us compute the beginning of the connected component of $\widehat{F}(\mathbf{s})$ containing $b_{\mathbf{s}}$.



Note that $b_{\mathbf{s}}$ is cylindric, and therefore all vertices appearing are cylindric. In the picture, only

$$b_{\mathbf{s}} = \begin{bmatrix} \bar{1} \\ 0 \end{bmatrix} \otimes \begin{bmatrix} \bar{1} \\ 0 \end{bmatrix}, \quad \begin{bmatrix} \bar{1} \\ 1 \end{bmatrix} \otimes \begin{bmatrix} \bar{1} \\ 0 \end{bmatrix} \quad \text{and} \quad \begin{bmatrix} \bar{1} \\ 0 \end{bmatrix} \otimes \begin{bmatrix} \bar{1} \\ 2 \end{bmatrix}$$

are in $\mathcal{O}_{\widehat{S}_n}(b_{\mathbf{s}})$. One checks that these are affine keys in the sense of Definition 3.34, and that the others are not, which illustrates Proposition 3.36.

We can now give a characterisation of the dual of an affine key, in the spirit of Proposition 2.28. It will be convenient to consider the action of \widehat{S}_n on \mathbb{Z}^n determined by the formulas

$$w(z_1, \dots, z_n) = (z_{w(1)}, \dots, z_{w(n)}) \text{ for } w \in S_n \quad \text{and} \quad \sigma_0(z_1, \dots, z_n) = (z_1 - \ell, z_2, \dots, z_{n-1}, z_n + \ell).$$

Observe that the set $\mathcal{D}(\mathbf{s})$ is a fundamental domain for this action.

Proposition 3.38. *Let $b \in \widehat{F}(\mathbf{s})$. Then b is an affine key if and only if $b^* = \dot{b}_{\dot{\mathbf{s}}}$ for some $\dot{\mathbf{s}} \in \mathbb{Z}^n(\mathbf{s})$.*

Proof. It suffices to observe that $b_{\mathbf{s}}^* = \dot{b}_{\dot{\mathbf{s}}}$ for some $\dot{\mathbf{s}} \in \mathcal{D}(\mathbf{s})$. The rest of the proof is analogous to that of Proposition 2.28. The maps $R_i, i = 1, \dots, n-1$ simply permute the columns of $\dot{b}_{\dot{\mathbf{s}}}$, and by Corollary 3.24 ${}^*\sigma_0(\begin{bmatrix} \bar{r}_1 \\ \vdots \\ \bar{r}_n \end{bmatrix} \otimes \dots \otimes \begin{bmatrix} \bar{r}_1 \\ \vdots \\ \bar{r}_n \end{bmatrix}) = \begin{bmatrix} \bar{r}_1 - \ell \\ \vdots \\ \bar{r}_2 \end{bmatrix} \otimes \dots \otimes \begin{bmatrix} \bar{r}_{n-1} \\ \vdots \\ \bar{r}_n + \ell \end{bmatrix}$, i.e. ${}^*\sigma_0(\dot{b}_{\dot{\mathbf{s}}}) = \dot{b}_{\sigma_0(\dot{\mathbf{s}})}$. Thus, for all $w \in \widehat{S}_n$, we have ${}^*w(\dot{b}_{\dot{\mathbf{s}}}) = \dot{b}_{w(\dot{\mathbf{s}})}$. Now, we have

$$\begin{aligned} b \text{ is an affine key} &\Leftrightarrow b \in \mathcal{O}_{\widehat{S}_n}(b_{\mathbf{s}}) \text{ for some } \mathbf{s} \in \mathcal{D}(\mathbf{s}) \\ &\Leftrightarrow b^* \in \mathcal{O}_{\widehat{S}_n}(\dot{b}_{\dot{\mathbf{s}}}) \text{ for some } \dot{\mathbf{s}} \in \mathcal{D}(\mathbf{s}) \\ &\Leftrightarrow \text{there exists } w \in \widehat{S}_n \text{ such that } b^* = {}^*w(\dot{b}_{\dot{\mathbf{s}}}) \\ &\Leftrightarrow \text{there exists } w \in \widehat{S}_n \text{ such that } b^* = \dot{b}_{w(\dot{\mathbf{s}})} \end{aligned}$$

□

Let $\mathbf{s} \in \mathcal{D}(\mathbf{s})$. By analogy with Proposition 2.31, we say that $b \in \widehat{F}(\mathbf{s})$ is an *affine bikey* of shape \mathbf{s} if b is in the orbit of $b_{\mathbf{s}}$ under the action of $\widehat{S}_n \times \widehat{S}_{\ell}$. We denote $\widehat{\mathcal{K}}(\mathbf{s})$ the set of bikeys of shape \mathbf{s} . Similarly, given $\dot{\mathbf{s}} \in \mathcal{D}(\mathbf{s})$, let $\widehat{\mathcal{K}}(\dot{\mathbf{s}}) \subseteq \widehat{F}(\dot{\mathbf{s}})$ be the set of bikeys of shape $\dot{\mathbf{s}}$. By Lemma 3.26, for all $\mathbf{s} \in \mathcal{D}(\mathbf{s})$, there exists $\dot{\mathbf{s}} \in \mathcal{D}(\mathbf{s})$ such that $b_{\mathbf{s}}^* = \dot{b}_{\dot{\mathbf{s}}}$. Therefore, $b \in \widehat{\mathcal{K}}(\mathbf{s})$ if and only if $b^* \in \widehat{\mathcal{K}}(\dot{\mathbf{s}})$, which gives an analogue to Proposition 2.32.

Finally, we can describe affine bikeys directly, in a slightly less explicit way than in finite type. First, as in Section 2.6 the maps R_j for $j = 1, \dots, \ell$ act on affine keys by permuting their columns. Therefore, the elements of $\hat{\mathcal{K}}(\mathbf{s})$ are obtained from affine keys by combining permutations of their columns and applications of R_0 .

Example 3.39. Take $\ell = n = 2$ and $\mathbf{s} = (r_1, r_2) = (1, 0)$ as in Example 3.37. The orbit of $b_{\mathbf{s}}$ under the action of \hat{S}_{ℓ} consists of the following elements

$$b_{\mathbf{s}} = \begin{array}{|c|} \hline 0 \\ \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 0 \\ \hline 1 \\ \hline \end{array}, \quad \begin{array}{|c|} \hline 0 \\ \hline 1 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline 0 \\ \hline 1 \\ \hline \end{array}, \quad \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline \end{array}, \quad \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline \end{array}, \quad \begin{array}{|c|} \hline \bar{2} \\ \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{2} \\ \hline \bar{1} \\ \hline 0 \\ \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array}, \quad \dots$$

The same procedure applied to each affine key in $\hat{F}(\mathbf{s})$ yields $\hat{\mathcal{K}}(\mathbf{s})$.

4. BICRYSTAL STRUCTURES INVOLVING INFINITE RANK CLASSICAL ROOT SYSTEMS

In this section, we denote by X_{∞} any infinite Dynkin diagram of type $A_{+\infty}, B_{\infty}, C_{\infty}$ or D_{∞} . We refer to [31] and the references therein for the notations and definition used in this section. In particular, the nodes of the Dynkin diagrams of types B_{∞}, C_{∞} or D_{∞} are parametrised by the integers of $\mathbb{Z}_{\geq 0}$ so that one gets a Dynkin diagram of type $A_{+\infty}$ by removing the 0-node. Our aim in this section, is to explain how the $A_{n-1} \times A_{\ell-1}$ bicrystal structure described in Section 2.4 can be generalised to obtain $X_{\infty} \times A_{\ell-1}$ bicrystal structures.

4.1. Recollection on extremal crystals of type X_{∞} . By [31], one can associate to each partition λ an extremal weight crystal⁷ $B_{\infty}(\lambda)$ of type X_{∞} which can be regarded as the direct limit of the finite type X_n -crystal $B_n(\lambda)$ associated to the dominant weight λ when n tends to infinity. To be more precise let introduce the infinite alphabets

$$\mathcal{A}_{X_{\infty}} = \begin{cases} \{1 < \dots < n < \dots\} \text{ for } X = A, \\ \{\dots < \bar{n} < \dots < \bar{1} < 0 < 1 < \dots < n < \dots\} \text{ for } X = B, \\ \{\dots < \bar{n} < \dots < \bar{1} < 1 < \dots < n < \dots\} \text{ for } X = C, \\ \{\dots < \bar{n} < \dots < \bar{2} < \bar{1}, 1 < 2 < \dots < n < \dots\} \text{ for } X = D. \end{cases}$$

The crystal $B_{\infty}(\lambda)$ admits a convenient realisation in terms of Kashiwara-Nakashima tableaux of type X_{∞} defined exactly as their finite rank counterpart by relaxing the admissibility condition of the columns. The general picture is identical to the finite type A . First, the crystals $B_{\infty}(1)$ for each type $A_{+\infty}, B_{\infty}, C_{\infty}$ and D_{∞} are respectively

$$\begin{array}{c} \begin{array}{c} \boxed{1} \xrightarrow{1} \boxed{2} \xrightarrow{2} \dots \xrightarrow{n-2} \boxed{n-1} \xrightarrow{n-1} \boxed{n} \xrightarrow{n} \dots \\ \dots \xrightarrow{n} \boxed{\bar{n}} \xrightarrow{n-1} \boxed{\overline{n-1}} \xrightarrow{n-2} \dots \xrightarrow{2} \boxed{\bar{2}} \xrightarrow{1} \boxed{\bar{1}} \xrightarrow{0} \boxed{0} \xrightarrow{0} \boxed{1} \xrightarrow{1} \boxed{2} \xrightarrow{2} \dots \xrightarrow{n-2} \boxed{n-1} \xrightarrow{n-1} \boxed{n} \xrightarrow{n} \dots \\ \dots \xrightarrow{n} \boxed{\bar{n}} \xrightarrow{n-1} \boxed{\overline{n-1}} \xrightarrow{n-2} \dots \xrightarrow{2} \boxed{\bar{2}} \xrightarrow{1} \boxed{\bar{1}} \xrightarrow{0} \boxed{1} \xrightarrow{1} \boxed{2} \xrightarrow{2} \dots \xrightarrow{n-2} \boxed{n-1} \xrightarrow{n-1} \boxed{n} \xrightarrow{n} \dots \end{array} \\ \begin{array}{c} \dots \xrightarrow{n} \boxed{\bar{n}} \xrightarrow{n-1} \boxed{\overline{n-1}} \xrightarrow{n-2} \dots \xrightarrow{2} \boxed{\bar{2}} \xrightarrow{1} \boxed{\bar{1}} \xrightarrow{0} \boxed{1} \xrightarrow{1} \boxed{2} \xrightarrow{2} \dots \xrightarrow{n-2} \boxed{n-1} \xrightarrow{n-1} \boxed{n} \xrightarrow{n} \dots \\ \begin{array}{c} \boxed{1} \\ \uparrow 1 \\ \boxed{\bar{2}} \downarrow 0 \\ \boxed{\bar{1}} \end{array} \end{array} \end{array}$$

Next, one introduces for any $k \geq 1$, column tableaux c of height k as the vertices of the crystal $B_{\infty}(1)^{\otimes k}$ with extremal vertex the column $c(k) = 1 \otimes \dots \otimes k$. Finally, for any partition λ with ℓ -columns of heights $k_1 \geq \dots \geq k_{\ell}$ the tableaux of shape λ are identified with the vertices $b = c_{\ell} \otimes \dots \otimes c_1$ in the connected component of the crystal $B_{\infty}(1^{k_{\ell}}) \otimes \dots \otimes B_{\infty}(1^{k_1})$ with extremal vertex $c(k_{\ell}) \otimes \dots \otimes c(k_1)$.

In type C_{∞} , one can show that a tableau $T = c_1 \dots c_{\ell}$ of shape λ is a filling of the Young diagram λ by letters of $\mathcal{A}_{C_{\infty}}$ so that the filling of each column c_i of λ is strictly increasing from top to bottom and the split form of T is semistandard. The split form of T is obtained by splitting each column c_i according to the following procedure. Given a column c , its split form is the pair (lc, rc) of columns contained no pair of letters (z, \bar{z}) with z unbarred defined as follows. Let $I = \{z_1 < \dots < z_r\}$ the set of unbarred letters z such that the pair (z, \bar{z}) occurs in c . Define the set $J = \{t_1 < \dots < t_r\}$ of unbarred letters such that:

⁷The crystal $B_{\infty}(\lambda)$ is not a highest weight crystal in types $C_{\infty}, B_{\infty}, D_{\infty}$.

- t_1 is the lowest unbarred letter satisfying: $t_1 > z_1, t_1 \notin c$ and $\bar{t}_1 \notin c$,
- for $i = 2, \dots, r$, t_i is the lowest unbarred letter satisfying: $t_i > \min(t_{i-1}, z_i), t_i \notin c$ and $\bar{t}_i \notin c$.

Then write:

rc for the column obtained by changing in c , \bar{z}_i into \bar{t}_i for each letter $z_i \in I$ and by reordering if necessary,
 lc for the column obtained by changing in c , z_i into t_i for each letter $z_i \in I$ and by reordering if necessary.

Now T is a tableau of type C_∞ when its split form $\text{spl}(T) = lc_1rc_1 \cdots lc_rrc_r$ is semistandard. Observe this implies that T itself is semistandard but a semistandard tableau on \mathcal{A}_{C_∞} is not always of type C_∞ . Also, the tableaux of type $A_{+\infty}$ which are the semistandard tableaux on the alphabet $\mathcal{A}_{A_{+\infty}}$ coincide with the tableaux of type C_∞ with no barred letter.

Example 4.1. For

$$T = \begin{array}{|c|c|c|} \hline \bar{2} & \bar{1} & 2 \\ \hline 1 & 2 & \\ \hline 2 & 3 & \\ \hline \end{array} \quad \text{one gets} \quad \text{spl}(T) = \begin{array}{|c|c|c|c|c|c|} \hline \bar{3} & \bar{2} & \bar{1} & \bar{1} & 2 & 2 \\ \hline 1 & 1 & 2 & 2 & & \\ \hline 2 & 3 & 3 & 3 & & \\ \hline \end{array}$$

therefore T is a tableau of type C_∞ , but

$$T' = \begin{array}{|c|c|c|} \hline \bar{2} & \bar{2} & 2 \\ \hline 1 & 2 & \\ \hline 2 & 3 & \\ \hline \end{array} \quad \text{with} \quad \text{spl}(T') = \begin{array}{|c|c|c|c|c|c|} \hline \bar{3} & \bar{2} & \bar{4} & \bar{2} & 2 & 2 \\ \hline 1 & 1 & 2 & 2 & & \\ \hline 2 & 3 & 3 & 4 & & \\ \hline \end{array}$$

is not, although it is semistandard on \mathcal{A}_{C_∞} .

The tableaux of types B_∞ and D_∞ can be described in a similar way, through a splitting operation on their columns. Nevertheless, in type B_∞ (respectively D_∞), blocks of the form $\begin{bmatrix} 0 \\ 0 \end{bmatrix}$ (respectively $\begin{bmatrix} \bar{1} \\ 1 \end{bmatrix}$) can appear in the same column. This slightly modifies the procedure for computing the splitting form of the tableaux. A juxtaposition $T = c_1 \cdots c_l$ of columns with decreasing heights will be a tableau of type B_∞ (respectively D_∞) if its split form $\text{spl}(T) = lc_1rc_1 \cdots lc_lrc_l$ is semistandard (respectively is semistandard and each two columns tableau $rc_i lc_{i+1}$ avoids a particular pattern π_D). We refer the reader to [30] for a complete description.

There exists a convenient notion of weight on the crystals $B_\infty(1)^{\otimes m}$, $m \geq 0$ (and thus also on any tensor product $B_\infty(1^{k_\ell}) \otimes \cdots \otimes B_\infty(1^{k_1})$). For any $b = x_1 \otimes \cdots \otimes x_m$ in $B(1)^{\otimes m}$, the weight $\text{wt}(b)$ is the sequence $(\gamma_i)_{0 \leq i}$ where for any $0 \leq i$ the integer γ_i is equal to the number of letters \bar{i} in b minus its number of letters i .

Given λ and μ two partitions, it was proved in [31] that the crystal $B_\infty(\lambda) \otimes B_\infty(\mu)$ decomposes as the disjoint union of extremal weight crystals and this decomposition is independent of the type considered.

Theorem 4.2. *For the type X_∞ extremal crystals $B_\infty(\lambda)$ and $B_\infty(\mu)$ we have*

$$B_\infty(\lambda) \otimes B_\infty(\mu) \simeq \bigoplus_v B_\infty(v)^{c_{\lambda,\mu}^v}$$

where $c_{\lambda,\mu}^v$ is the usual Littlewood-Richardson coefficients associated to λ, μ and v .

Remark 4.3.

- (1) The theorem extends to the tensor products $B_\infty(\lambda^1) \otimes B_\infty(\lambda^2) \otimes \cdots \otimes B_\infty(\lambda^\ell)$ associated to any sequence $(\lambda^1, \dots, \lambda^\ell)$ of partitions. This tensor product contains a unique component isomorphic to $B_\infty(\lambda^1 + \cdots + \lambda^\ell)$ that we will refer as the principal component (here $\lambda^1 + \cdots + \lambda^\ell$ is the partition whose k -th part is the sum of the k -th part of each $\lambda^a, a = 1, \dots, \ell$).
- (2) For any $m \geq 0$, any connected component B of $B(1)^{\otimes m}$ contains a unique *distinguished extremal vertex*, namely a vertex $b_{\text{ext}} = x_1 \otimes \cdots \otimes x_m$ such that x_1, \dots, x_m are unbarred letters and $x_1 \cdots x_m$ is a Yamanouchi word. When $b_{\text{ext}} = b_\lambda$ is the tableau of shape λ whose i -th row contains only letters i , $B(b_\lambda)$ coincides with the set of tableaux of shape λ . In the general case, we get that $\text{wt}(b) = \lambda$ is a partition and there is a unique crystal isomorphism from $B(b_{\text{ext}}) = B$ to $B(b_\lambda)$ sending b_{ext} on b_λ .

4.2. Jeu de Taquin on two columns and A_1 -crystal structure. For any integer $u \geq 0$, write $B_\infty(\omega_u)$ for the type X_∞ -crystal extremal associated to the partition 1^u . It contains exactly the column tableaux of type X_∞ and shape 1^u . More generally, we identify the partition λ with a_u columns of height u with the formal weight $\lambda = \sum_u a_u \omega_u$. By Theorem 4.2, we get for any integers $u, v \geq 1$

$$B_\infty(\omega_u) \otimes B_\infty(\omega_v) \simeq \bigoplus_{t=0}^{\min(u,v)} B_\infty(\omega_t + \omega_{u+v-t}) = B_\infty(\omega_{u+v}) \oplus B_\infty(\omega_1 + \omega_{u+v-1}) \oplus \cdots \oplus B_\infty(\omega_{\min(u,v)} + \omega_{\max(u,v)}).$$

When $u \leq v$, the principal component $B_\infty(\omega_u + \omega_v)$ of $B_\infty(\omega_u) \otimes B_\infty(\omega_v)$ contains exactly the tableaux of type X_∞ and shape $\omega_u + \omega_v$. When $u > v$, the vertices of the principal component $B_\infty(\omega_v + \omega_u)$ of $B_\infty(\omega_u) \otimes B_\infty(\omega_v)$ are called antit tableaux of type X_∞ . As in Section 2.4, they can be easily described by using the splitting operation given in Section 4.1.

Consider a vertex $c_2 \otimes c_1$ in $B_\infty(\omega_u) \otimes B_\infty(\omega_v)$ which is not a tableau. This means we have either $u > v$, or $u \leq v$ and the connected component $B_\infty(c_2 \otimes c_1)$ containing $c_2 \otimes c_1$ is isomorphic to a crystal $B_\infty(\omega_t + \omega_{u+v-t})$ with $t \leq u - 1$. In both cases the above crystal decomposition implies that the crystal $B_\infty(\omega_{u-1}) \otimes B_\infty(\omega_{v+1})$ contains a unique component isomorphic to $B_\infty(c_2 \otimes c_1)$. Write $\dot{e}(c_2 \otimes c_1)$ for the vertex in $B_\infty(\omega_{u-1}) \otimes B_\infty(\omega_{v+1})$ matched with $c_2 \otimes c_1$ by this isomorphism. When $c_2 \otimes c_1$ is a tableau, we set $\dot{e}(c_2 \otimes c_1) = 0$.

Similarly, for any vertex $c_2 \otimes c_1$ in $B_\infty(\omega_u) \otimes B_\infty(\omega_v)$ which is not an antit tableau, there exists a unique vertex $\dot{f}(c_2 \otimes c_1)$ in $B_\infty(\omega_{u+1}) \otimes B_\infty(\omega_{v-1})$ such that the components $B_\infty(c_2 \otimes c_1)$ and $B_\infty(\dot{f}(c_2 \otimes c_1))$ are isomorphic and $\dot{f}(c_2 \otimes c_1)$ is matched with $c_2 \otimes c_1$ by this isomorphism. When $c_2 \otimes c_1$ is an antit tableau, we set $\dot{f}(c_2 \otimes c_1) = 0$. For any $c_2 \otimes c_1$ in $B_\infty(\omega_u) \otimes B_\infty(\omega_v)$, write $\dot{B}_\infty(c_2 \otimes c_1)$ for the set obtained by applying operators \dot{e} and \dot{f} to $c_2 \otimes c_1$. We get immediately the following proposition.

Proposition 4.4. *The set $\dot{B}_\infty(c_2 \otimes c_1)$ has the structure of a highest weight A_1 -crystal.*

The previous proposition can be regarded as an analogue of the Jeu de Taquin procedure on skew tableaux of type A with two columns. In fact, such a notion of skew tableaux also exists in type B_∞, C_∞ and D_∞ . A skew tableau of type X_∞ is defined as the filling of a skew Young diagram by letters of \mathcal{A}_{X_∞} whose duplicated form is semistandard (and avoid the pattern π_D in type D_∞). To any skew tableau $c_1 \cdots c_\ell$, one associates the tensor product of columns $c_\ell \otimes \cdots \otimes c_1$. Conversely, any tensor product $c_2 \otimes c_1$ of two columns can be regarded as a minimal skew tableau as in Example 2.14. The operators \dot{e} and \dot{f} can then be interpreted as horizontal sliding operations on this skew tableau. In type $A_{+\infty}$, one so recovers the usual Jeu de Taquin procedure, see Example 2.14, in types B_∞ and C_∞ this corresponds to the sliding operations described in [30].

Example 4.5. In type C_∞ , for

$$c_1 \otimes c_2 = \begin{array}{|c|} \hline \bar{2} \\ \hline \bar{1} \\ \hline 2 \\ \hline 5 \\ \hline \end{array} \otimes \begin{array}{|c|} \hline \bar{4} \\ \hline \bar{2} \\ \hline 3 \\ \hline 5 \\ \hline \end{array}, \quad \text{the corresponding minimal skew tableau is } \begin{array}{|c|c|} \hline \bar{2} & \\ \hline \bar{4} & \bar{1} \\ \hline \bar{2} & 2 \\ \hline 3 & 5 \\ \hline 5 & \\ \hline \end{array}.$$

By using the sliding procedure in type C_∞ or the previous crystal isomorphisms, one gets

$$\dot{e} \left(\begin{array}{|c|c|} \hline & \bar{2} \\ \hline \bar{4} & \bar{1} \\ \hline \bar{2} & 2 \\ \hline 3 & 5 \\ \hline 5 & \\ \hline \end{array} \right) = \begin{array}{|c|c|} \hline \bar{4} & \bar{2} \\ \hline \bar{1} & \bar{1} \\ \hline 1 & 5 \\ \hline 3 & \\ \hline 5 & \\ \hline \end{array} \quad \text{and} \quad \dot{f} \left(\begin{array}{|c|c|} \hline & \bar{2} \\ \hline \bar{4} & \bar{1} \\ \hline \bar{2} & 2 \\ \hline 3 & 5 \\ \hline 5 & \\ \hline \end{array} \right) = \begin{array}{|c|c|} \hline & \bar{2} \\ \hline \bar{4} & \bar{1} \\ \hline \bar{2} & 3 \\ \hline 5 & 5 \\ \hline \end{array}.$$

4.3. Bicrystal structure on product of columns of type X_∞ . Consider $n, \ell \in \mathbb{Z}_{\geq 2}, s = (s_1, \dots, s_\ell) \in \mathbb{Z}_{\geq 0}^\ell$ and $s \in \mathbb{Z}_{\geq 0}$. With the notation of the previous paragraph, write $B_\infty(s) = B_\infty(\omega_{s_1}) \otimes \cdots \otimes B_\infty(\omega_{s_\ell})$ and set $B_\infty(s) = \bigoplus_{|s|=s} B_\infty(s)$. By definition $B_\infty(s)$ is a type X_∞ -crystal. We are going to show that it admits in fact the structure of a type $X_\infty \times A_{\ell-1}$ -crystal. First, for any $j = 1, \dots, \ell - 1$ define the operators \dot{e}_j and \dot{f}_j on each $B_\infty(s)$ as the action of \dot{e} and \dot{f} on its j and $j + 1$ -th components, that is for any $c_1 \otimes \cdots \otimes c_\ell \in B_\infty(s)$

$$\dot{e}_j = c_1 \otimes \cdots \otimes \dot{e}(c_j \otimes c_{j+1}) \otimes \cdots \otimes c_\ell \quad \text{and} \quad \dot{f}_j = c_1 \otimes \cdots \otimes \dot{f}(c_j \otimes c_{j+1}) \otimes \cdots \otimes c_\ell.$$

By definition, the operators \dot{e} and \dot{f} commute with the action of the X_∞ -crystal operators, therefore \dot{e}_j and \dot{f}_j also commute with the operators e_i and f_i with $i \in X_\infty$.

Proposition 4.6. *For any $b = c_1 \otimes \cdots \otimes c_\ell \in B_\infty(s)$, the colored and oriented graph $\dot{B}(b)$ defined by applying the operators \dot{e}_j and \dot{f}_j to b has the structure of a type $A_{\ell-1}$ crystal.*

Proof. Assume first that b only contains unbarred letters. Then the result follows from Theorem 2.21. In the general case, there exists a path in $B_\infty(s)$ from b to the distinguished extremal vertex b_{ext} defined in Remark 4.3. This path corresponds to a composition of operators e_i and $f_i, i \in X_\infty$ which commute with the operators \dot{e}_j and $\dot{f}_j, j \in \{1, \dots, \ell\}$. Therefore the graph $\dot{B}(b)$ and $\dot{B}(b_{\text{ext}})$ are isomorphic and we are done because b_{ext} only contains unbarred letters. \square

Denote by $\dot{f}_{\underline{j}}$ any finite composition of operators \dot{f}_j with $j \in \{1, \dots, \ell\}$ and by $k_{\underline{i}}$ any finite composition of operators e_i, f_i with $i \geq 0$.

Theorem 4.7.

- (1) The crystal $B_\infty(s)$ has the structure of an $(X_\infty \times A_{\ell-1})$ -crystal. This bicrystal is extremal for the X_∞ -structure and of highest weight for the $A_{\ell-1}$ -structure.
- (2) The tableaux b_λ associated to the partitions of rank s are the unique vertices in $B_\infty(s)$ both distinguished extremal for X_∞ and of highest weight for $A_{\ell-1}$.
- (3) We have the decomposition

$$B_\infty(s) = \bigoplus_{|\lambda|=s} \dot{f}_{\underline{j}} k_{\underline{i}} b_\lambda$$

where the sum runs over all the possible \underline{i} and \underline{j} . In particular

$$B_\infty(s) \simeq \bigoplus_{|\lambda|=s} B_\infty(\lambda) \otimes B(\lambda^{\text{tr}})$$

Proof. Assertion (1) is just a reformulation of the previous proposition. To prove Assertion (2), observe first that the vertices b_λ are clearly both distinguished extremal for X_∞ and of highest weight for $A_{\ell-1}$. Consider now $b = c_\ell \otimes \dots \otimes c_1$ with this property. Since b only contains unbarred letters, we can apply Corollary 2.22 and conclude that $b = b_\lambda$ for $\lambda \vdash s$. Assertion (3) easily follows from 1 et 2. \square

Corollary 4.8. The $A_{\ell-1}$ -weight highest vertices in $B_\infty(s)$ are exactly the tableaux of type X_∞ .

Proof. Consider b a $A_{\ell-1}$ -weight vertices in $B_\infty(s)$. Then, by Assertion (3) of the previous theorem, there should exist a partition λ of rank s and sequence $k_{\underline{i}}$ such that $b = k_{\underline{i}} b_\lambda$. This means that b is a tableau of type X_∞ . Conversely, any tableau b can be written $b = k_{\underline{i}} b_\lambda$ and for $j = 1, \dots, \ell - 1$, we have $\dot{e}_j(b) = \dot{e}_j(k_{\underline{i}} b_\lambda) = k_{\underline{i}} \dot{e}_j(b_\lambda) = 0$ since both crystal structures commute. \square

We conclude this section with three important remarks.

Remark 4.9.

- (1) Contrary to Section 2.4, we did not introduced a duality $b \rightarrow b^*$ from which the $A_{\ell-1}$ -crystal structure on $B_\infty(s)$ can easily be made explicit. Although the $A_{\ell-1}$ -highest vertices of $B_\infty(s)$ correspond to Kashiwara-Nakashima tableaux of type X_∞ , we cannot use the same duality as in type $A_{\ell-1}$. Indeed, in types B_∞ and D_∞ , these tableaux are not semistandard on \mathcal{A}_{X_∞} in general. Moreover, even in type C_∞ , the action of a crystal operator \dot{e}_j and \dot{f}_j does not coincide with a horizontal sliding of the ordinary Jeu de Taquin of type A . This problem, also related to the cyclage operation on tableaux defined in [28] and [29] will be considered elsewhere.
- (2) The previous corollary can also be obtained without referring to Section 2.4. For $b = c_1 \otimes \dots \otimes c_\ell$ a $A_{\ell-1}$ -highest weight in $B_\infty(s)$, the equality $\dot{e}_j(b) = 0$ for any $j = 1, \dots, \ell - 1$ means that $c_j \otimes c_{j+1}$ is a tableau of two columns for any $j = 1, \dots, \ell - 1$. This implies that b is a tableau of type X_∞ because the tableaux are characterised locally by conditions on their pairs of consecutive columns.
- (3) The results of this section can also be regarded as bicrystal structures on particular matrix sets. In type $A_{+\infty}$ or C_∞ , each vertices $b = c_1 \otimes \dots \otimes c_\ell$ of $B_\infty(s)$ can indeed be encoded in an infinite matrix $M = (m_{i,j})$ with rows indexed by $\mathbb{Z} \setminus \{0\}$ and columns indexed by $\{1, \dots, \ell\}$, where $m_{i,j} = 1$ if c_j contains the letter $i \in \mathbb{Z} \setminus \{0\}$ and $m_{i,j} = 0$ otherwise. In type B_∞ we proceed similarly, except that the integers in the rows $m_{0,j}$ may be greater than 1. Finally in type D_∞ , the coefficient $m_{0,j}$ counts the number of blocks $\begin{smallmatrix} \bar{1} \\ 1 \end{smallmatrix}$ in c_j .

5. FIXED POINTS IN BICRYSTALS

In this section, we show that bicrystals structures of classical types can be obtained by considering fixed points sets under involutions defined on our combinatorial Fock spaces. We shall detail mainly the classical case, the methods and arguments being similar in the affine situation.

5.1. Dynkin diagram involutions in finite type A. In type A_{n-1} , let θ be the Dynkin diagram automorphism mapping the vertex i to $n - i$, for all $i = 1, \dots, n - 1$. Then θ induces a map on the set of dominant weights of type A_{n-1} that we denote by the same symbol

$$\theta : P_+ \rightarrow P_+, \quad \sum_{i=1}^{n-1} a_i \omega_i \mapsto \sum_{i=1}^{n-1} a_i \omega_{n-i}.$$

We can also define the map θ on the partitions of length at most n : θ replaces each column of height $0 < i < n$ by a column of height $n - i$ and fixes the columns of height 0 or n .

For any partition λ , recall that $B(\lambda)$ is the crystal of highest weight vertex b_λ . It can be realised by using semistandard tableaux of shape λ , and in this case b_λ is the tableau of shape λ with only k 's in row k [17, Section 7.3]. The map θ induces a map on $B(\lambda)$ that we denote by the same symbol

$$\theta : B(\lambda) \longrightarrow B(\theta(\lambda)),$$

which maps b_λ to $b_{\theta(\lambda)}$ and which is a θ -isomorphism of crystals, i.e. for all $b \in B(\lambda)$, and for all $i = 1, \dots, n - 1$,

$$\theta(f_i b) = f_{\theta(i)}(\theta(b)) = f_{n-i}(\theta(b)).$$

Example 5.1. Take $n = 3$ and let $\lambda = \omega_1 = \square$, so that $\theta(\lambda) = \omega_2 = \square$. We get the following crystals.

$$B(\lambda) = \boxed{1} \xrightarrow{1} \boxed{2} \xrightarrow{2} \boxed{3} \quad \text{and} \quad B(\theta(\lambda)) = \boxed{\begin{smallmatrix} 1 \\ 3 \end{smallmatrix}} \xrightarrow{2} \boxed{\begin{smallmatrix} 1 \\ 3 \end{smallmatrix}} \xrightarrow{1} \boxed{\begin{smallmatrix} 2 \\ 3 \end{smallmatrix}}.$$

5.2. Crystal structure on fixed points sets. Denote

$$P_+^\theta = \{\lambda \in P_+ \mid \theta(\lambda) = \lambda\}$$

the set of dominant weights fixed under θ , so that

$$\lambda \in P_+^\theta \Leftrightarrow \exists a_1, \dots, a_{\lfloor n/2 \rfloor} \in \mathbb{Z} \text{ s.t. } \lambda = \begin{cases} \sum_{i=1}^{n/2-1} a_i(\omega_i + \omega_{n-i}) + a_{n/2}\omega_{n/2} & \text{if } n-1 \text{ is odd,} \\ \sum_{i=1}^{(n-1)/2} a_i(\omega_i + \omega_{n-i}) & \text{if } n-1 \text{ is even.} \end{cases}$$

Definition 5.2. For any weight λ in P^θ , set $\lambda' = \sum_{i=1}^{\lfloor n/2 \rfloor} a_i \omega'_i$ where ω'_i is the i -th fundamental weight of type $X_{\lfloor n/2 \rfloor}$, with

$$X = B \text{ if } n-1 \text{ is odd} \quad \text{and} \quad X = C \text{ if } n-1 \text{ is even}$$

For any partition λ with at most n parts and such that $\theta(\lambda) = \lambda$, there exists a unique crystal involution on $B(\lambda)$ such that

$$\theta(\tilde{f}_{i_1} \cdots \tilde{f}_{i_r}(b_\lambda)) = \tilde{f}_{\theta(i_1)} \cdots \tilde{f}_{\theta(i_r)}(b_\lambda)$$

for any sequence i_1, \dots, i_r in $\{1, \dots, n - 1\}$ of arbitrary length. This permits to consider the set

$$B(\lambda)^\theta = \{b \in B(\lambda) \mid \theta(b) = b\}.$$

In [35, 36] Naito and Sagaki established the following theorem. Consider a partition λ with at most n parts and such that $\theta(\lambda) = \lambda$.

Theorem 5.3. For $i = 1, \dots, \lfloor n/2 \rfloor$, define the modified crystal operators

(1) if $n - 1$ is odd,

$$f_i^\theta = \begin{cases} f_i f_{n-i} & \text{if } i = 1, \dots, n/2 - 1 \\ f_i & \text{if } i = n/2. \end{cases}$$

(2) if $n - 1$ is even,

$$f_i^\theta = \begin{cases} f_i f_{n-i} & \text{if } i = 1, \dots, (n-1)/2 - 1 \\ f_i f_{i+1}^2 f_i & \text{if } i = (n-1)/2. \end{cases}$$

Then every $b \in B(\lambda)^\theta$ is obtained by applying a sequence of modified crystal operators to the highest weight vertex b_λ of $B(\lambda)$. This induces a crystal structure on $B(\lambda)^\theta$. Moreover, we have

$$B(\lambda)^\theta \simeq B^{(X_{\lfloor n/2 \rfloor})}(\lambda'),$$

the crystal of classical type $X_{\lfloor n/2 \rfloor}$ with highest weight λ' . In each case, the action of the modified crystal operators is mapped to the action of the classical crystal operators.

Example 5.4.

(1) Take $n = 4$, and $\lambda = \omega_1 + \omega_3 = \begin{smallmatrix} \square & \square \\ \square & \end{smallmatrix}$, so that $\lambda \in P_+^\theta$. Then we have

$$B(\lambda)^\theta = \begin{smallmatrix} \square & \square \\ 1 & 1 \\ 2 & \\ 3 & \end{smallmatrix} \xrightarrow{1} \begin{smallmatrix} \square & \square \\ 1 & 2 \\ 2 & \\ 4 & \end{smallmatrix} \xrightarrow{2} \begin{smallmatrix} \square & \square \\ 1 & 2 \\ 3 & \\ 4 & \end{smallmatrix} \xrightarrow{2} \begin{smallmatrix} \square & \square \\ 1 & 3 \\ 3 & \\ 4 & \end{smallmatrix} \xrightarrow{1} \begin{smallmatrix} \square & \square \\ 2 & 4 \\ 3 & \\ 4 & \end{smallmatrix} \simeq B^{(B_2)}(\lambda')$$

with $\lambda' = \omega'_1$.

(2) Take $n = 5$, and $\lambda = \omega_1 + \omega_4 = \begin{smallmatrix} \square & \square \\ \square & \\ \square & \\ \square & \end{smallmatrix}$, so that $\lambda \in P_+^\theta$. Then we have

$$B(\lambda)^\theta = \begin{smallmatrix} \square & \square \\ 1 & 1 \\ 2 & \\ 3 & \\ 4 & \end{smallmatrix} \xrightarrow{1} \begin{smallmatrix} \square & \square \\ 1 & 2 \\ 2 & \\ 3 & \\ 5 & \end{smallmatrix} \xrightarrow{2} \begin{smallmatrix} \square & \square \\ 1 & 4 \\ 3 & \\ 4 & \\ 5 & \end{smallmatrix} \xrightarrow{1} \begin{smallmatrix} \square & \square \\ 2 & 5 \\ 3 & \\ 4 & \\ 5 & \end{smallmatrix} \simeq B^{(C_2)}(\lambda')$$

with $\lambda' = \omega'_1$.

Case (1) is a particular occurrence of a phenomenon called “similarity of crystal bases” by Kashiwara [23, Theorem 5.1]. In fact, in the case where $n - 1$ is odd, we can also exhibit a type $C_{n/2}$ crystal structure on a subset of $B(\lambda)^\theta$ as follows.

Assume $n - 1$ is odd and let λ is a partition with at most n part and with associated weight of the form $\lambda = \sum_{i=1}^{n/2} a_i(\omega_i + \omega_{n-i})$. Let $\eta : B(\lambda) \rightarrow B(\lambda)$ be the involution defined by $\eta(b_\lambda) = b_\lambda$ and

$$\eta(\tilde{f}_{i_1} \cdots \tilde{f}_{i_r}(b_\lambda)) = \tilde{f}_{n-i_1}^{a_{i_1}} \cdots \tilde{f}_{n-i_r}^{a_{i_r}}(b_\lambda) \text{ with } a_{i_k} = \begin{cases} 1 & \text{if } i_k \neq n/2, \\ 2 & \text{otherwise.} \end{cases}$$

for any sequence i_1, \dots, i_r in $\{1, \dots, n - 1\}$ of arbitrary length. Accordingly, we write

$$P_+^\eta = \left\{ \lambda \in P_+ \mid \lambda = \sum_{i=1}^{n/2} a_i(\omega_i + \omega_{i+1}) \right\} \subset P_+^\theta,$$

and for $\lambda \in P_+^\eta$, let $\lambda^\dagger = \sum_{i=1}^{n/2} a_i \omega_i^\dagger$ where ω_i^\dagger denotes the i -th fundamental weight of type $C_{n/2}$. Let

$$B(\lambda)^\eta = \{b \in B(\lambda) \mid \eta(b) = b\}.$$

Theorem 5.5. Assume $n - 1$ is odd and let λ be a partition with associated weight in P_+^η . Consider the modified crystal operators

$$f_i^\eta = f_i f_{n-i} \quad \text{for all } i = 1, \dots, n/2.$$

Then every $b \in B(\lambda)^\eta$ is obtained by applying a sequence of modified crystal operators f_i^η to the highest weight vertex b_λ of $B(\lambda)$. This induces a crystal structure on $B(\lambda)^\eta$. Moreover, we have

$$B(\lambda)^\eta \simeq B^{(C_{n/2})}(\lambda^\dagger),$$

the crystal of type $C_{n/2}$ with highest weight λ^\dagger , where the action of the modified crystal operators is mapped to the action of the classical crystal operators.

Example 5.6. Take $n = 4$, and $\lambda = \omega_1 + \omega_3 = \begin{smallmatrix} \square & \square \\ \square & \end{smallmatrix}$, so that $\lambda \in P_+^\eta$. Then we have

$$B(\lambda)^\eta = \begin{smallmatrix} \square & \square \\ 1 & 1 \\ 2 & \\ 3 & \end{smallmatrix} \xrightarrow{1} \begin{smallmatrix} \square & \square \\ 1 & 2 \\ 2 & \\ 4 & \end{smallmatrix} \xrightarrow{2} \begin{smallmatrix} \square & \square \\ 1 & 3 \\ 3 & \\ 4 & \end{smallmatrix} \xrightarrow{1} \begin{smallmatrix} \square & \square \\ 2 & 4 \\ 3 & \\ 4 & \end{smallmatrix} \simeq B^{(C_2)}(\lambda')$$

with $\lambda' = \omega'_1$.

Remark 5.7. For $\lambda \in P_+$, one might also consider the set $B(\lambda)_\theta = \{b \in B(\lambda) \mid \theta(\text{wt}(b)) = \text{wt}(b)\} \supset B(\lambda)^\theta$. Nevertheless, the modified crystal operators of Theorem 5.3 and Theorem 5.5 do not endow $B(\lambda)_\theta$ with the structure of a crystal. For example in type A_3 for $\lambda = (4, 3, 3)$, $B(\lambda)_\theta$ would so have a source vertex of weight $(4, 3, 2, 1)$ whose associated connected component is not a type B_2 -crystal.

5.3. Bicrystal structure on fixed points sets: the finite case. In this section, we combine the results from Section 2.4 and Section 5.2 to exhibit bicrystals structures of mixed types $A \times B$, $A \times C$ and $B \times C$. Recall that we have fixed $s \in \mathbb{Z}_{\geq 0}$ and considered

$$F(s) = \bigoplus_{\mathbf{s} \in \mathbb{Z}^\ell(s)} F(\mathbf{s}) = \bigoplus_{\mathbf{s} \in \mathbb{Z}^\ell(s)} B(\omega_{s_1}) \otimes \cdots \otimes B(\omega_{s_n}).$$

By Corollary 2.22 and Theorem 2.25, we get

$$F(s) = \bigoplus_{\lambda \in \mathcal{S}(s)} f_{\underline{i}}^* f_{\underline{i}} b_\lambda, \quad \text{which implies} \quad F(s) \simeq \bigoplus_{\lambda \in \mathcal{S}(s)} B(\lambda) \otimes B(\lambda^{\text{tr}})$$

where $\mathcal{S}(s)$ is the set of partitions of s with Dynkin diagram contained in the rectangle $n \times \ell$. Now, for all $b = f_{\underline{i}}^* f_{\underline{i}} b_\lambda \in F(s)$ such that $\theta(\lambda) = \lambda$, we set

$$\theta(b) = f_{\underline{i}}^* \theta(f_{\underline{i}} b_\lambda)$$

where $\theta(f_{\underline{i}} b_\lambda) = f_{n-i_r} \cdots f_{n-i_1} b_\lambda$ if $\underline{i} = (i_1, \dots, i_r)$, and we consider

$$F(s)^\theta = \left\{ b \in F(s) \mid b = f_{\underline{i}}^* f_{\underline{i}} b_\lambda \text{ with } \theta(\lambda) = \lambda \text{ and } \theta(b) = b \right\}.$$

By Theorem 5.3, we then get a bicrystal structure.

Theorem 5.8. *The set $F(s)^\theta$ has an $(X_{\lfloor n/2 \rfloor} \times A_{\ell-1})$ -crystal structure, where*

$$X = \begin{cases} B & \text{if } n-1 \text{ is odd,} \\ C & \text{if } n-1 \text{ is even.} \end{cases}$$

This yields the decomposition

$$F(s)^\theta \simeq \bigoplus_{\lambda \in \mathcal{S}(s), \theta(\lambda)=\lambda} B^{(X_{\lfloor n/2 \rfloor})}(\lambda') \otimes B(\lambda^{\text{tr}}).$$

In fact, we can consider fixed points on either side, and on both sides simultaneously. To see this, we need to consider the automorphism $\dot{\theta}$ of the Dynkin diagram of type $A_{\ell-1}$. Similarly to θ , the map $\dot{\theta}$ induces an involution on the set of partitions with at most ℓ columns flipping rows of length j and $\ell-j$ for $0 < j < \ell$ and fixing rows of length 0 or ℓ . Then, for any partition λ fixed by $\dot{\theta}$, we can define $\dot{\theta}$ on $b = f_{\underline{i}}^* f_{\underline{i}} b_\lambda \in F(s)$ by setting

$$\dot{\theta}(b) = f_{\underline{i}} (\dot{\theta} f_{\underline{j}})^* b_\lambda$$

where $\dot{\theta}(f_{\underline{j}}^* b_\lambda) = f_{\ell-j_r}^* \cdots f_{\ell-j_1}^* b_\lambda$ if $\underline{j} = (j_1, \dots, j_r)$, and also consider

$$F(s)^{\dot{\theta}} = \left\{ b \in F(s) \mid b = f_{\underline{i}}^* f_{\underline{i}} b_\lambda \text{ with } \dot{\theta}(\lambda) = \lambda \text{ and } \dot{\theta}(b) = b \right\}.$$

Theorem 5.9. *The set $F(s)^{\dot{\theta}}$ has an $(A_{n-1} \times \dot{X}_{\lfloor \ell/2 \rfloor})$ -crystal structure, where*

$$\dot{X} = \begin{cases} B & \text{if } \ell-1 \text{ is odd} \\ C & \text{if } \ell-1 \text{ is even} \end{cases}$$

This yields the decomposition

$$F(s)^{\dot{\theta}} \simeq \bigoplus_{\lambda \in \mathcal{S}(s), \dot{\theta}(\lambda)=\lambda} B(\lambda) \otimes B^{(\dot{X}_{\lfloor \ell/2 \rfloor})}((\lambda^{\text{tr}})').$$

Now set

$$\mathcal{S}(s)^{\theta, \dot{\theta}} = \{ \lambda \in \mathcal{S}(s) \mid \theta(\lambda) = \lambda \text{ and } \dot{\theta}(\lambda) = \lambda \}.$$

Then the set

$$F(s)^{\theta, \dot{\theta}} = F(s)^\theta \cap F(s)^{\dot{\theta}}.$$

is the subset of $F(s)$ obtained from the double highest weight vertices b_λ with $\lambda \in \mathcal{S}(s)^{\theta, \dot{\theta}}$ by applying the modified crystal operators f_i^θ and $f_j^{\dot{\theta}}$.

Theorem 5.10. *The set $F(s)^{\theta, \check{\theta}}$ has the structure of an $(X_{[n/2]} \times \check{X}_{[\ell/2]})$ -crystal where*

$$X_{[n/2]} = \begin{cases} B & \text{if } n-1 \text{ is odd,} \\ C & \text{if } n-1 \text{ is even,} \end{cases} \quad \text{and } \check{X}_{[\ell/2]} = \begin{cases} B & \text{if } \ell-1 \text{ is odd,} \\ C & \text{if } \ell-1 \text{ is even.} \end{cases}$$

Moreover, we have the decomposition

$$F(s)^{\theta, \check{\theta}} \simeq \bigoplus_{\lambda \in \mathcal{S}(s)^{\theta, \check{\theta}}} B^{(X_{[n/2]})}(\lambda') \otimes B^{(\check{X}_{[\ell/2]})}((\lambda^{\text{tr}})').$$

Remark 5.11.

- (1) It would be interesting to have a combinatorial description of the set $F(s)^{\theta, \check{\theta}}$ as tensor products of columns or binary matrices.
- (2) By applying the relevant weight functions on the previous decompositions obtained for $F(s)^{\theta}$, $F(s)^{\check{\theta}}$ and $F(s)^{\theta, \check{\theta}}$, one can get analogues of Cauchy identities in our bicrystal context. For example, for any vertex b in $F(s)^{\theta, \check{\theta}}$, $\text{wt}(b)'$ is a weight of type $X_{[n/2]}$ whereas $\text{wt}(b^*)'$ is a weight of type $\check{X}_{[\ell/2]}$. One obtains

$$\sum_{b \in F(s)^{\theta, \check{\theta}}} x^{\text{wt}(b)'} y^{\text{wt}(b^*)'} = \sum_{\lambda \in \mathcal{S}(s)^{\theta, \check{\theta}}} s_{\lambda'}^{(X_{[n/2]})}(x) s_{(\lambda^{\text{tr}})'}^{(\check{X}_{[\ell/2]})}(y)$$

where for any partition ν of length at most m (respectively at most p), $s_{\nu}^{(X_m)}$ (respectively $s_{\nu}^{(\check{X}_p)}$) stands for the Weyl character of type X_m (respectively \check{X}_p) associated to ν .

- (3) The results of this paragraph have analogues when the map θ is replaced by the map η defined in Section 5.2 and P^{θ} by P^{η} .
- (4) In [24], King established the interesting Cauchy type formula

$$\prod_{i=1}^n \prod_{j=1}^{\ell} (x_i + x_i^{-1} + y_j + y_j^{-1}) = \sum_{\lambda \subset n \times \ell} s_{\lambda}^{(C_n)}(x) s_{[\lambda]}^{(C_{\ell})}(y)$$

where $[\lambda]$ is the transposed of the rectangular complement of λ in $n \times \ell$. Although the right-hand side looks similar to the sum appearing in our results when $n-1$ and $\ell-1$ are even, it is not obvious to relate both.

To conclude this paragraph, let us introduce a last natural involution $\check{\theta}$ on the set $\mathcal{S}(s)$. Here for any $\lambda \in \mathcal{S}(s)$, the Young diagram of $\check{\theta}(\lambda)$ is obtained from that of λ by changing each column of height $0 \leq i \leq n$ in a column of height $n-i$. In other words, $\check{\theta}(\lambda)$ is the π -rotation of the complement of λ in the $n \times \ell$ -rectangle as illustrated in Example 5.12 below. In particular θ and $\check{\theta}$ coincide on the partitions with no column of height n or 0 but this is not true in general.

Example 5.12. Take $\ell = 5, n = 3$ and $\lambda = \begin{array}{|c|c|c|c|c|} \hline & & & & \\ \hline & & & & \\ \hline & & & & \\ \hline \end{array}$. Then λ and $\check{\theta}(\lambda)$ fit in the $n \times \ell$ -rectangle as follows

$$\lambda \longrightarrow \begin{array}{|c|c|c|c|c|} \hline & & & & \\ \hline & & & & \\ \hline & & & & \\ \hline \end{array} \longrightarrow \pi\text{-rotation of } \check{\theta}(\lambda) \quad \text{so we have} \quad \check{\theta}(\lambda) = \begin{array}{|c|c|c|c|c|} \hline & & & & \\ \hline & & & & \\ \hline & & & & \\ \hline \end{array}.$$

One may also observe that $\check{\theta}(\lambda) = \lambda$ if and only if $\check{\theta}(\lambda)$ is obtained from λ by changing each row of length $0 \leq j \leq \ell$ in a row of length $\ell-j$. In fact, we see that $\check{\theta}(\lambda) = \lambda$ if and only if each box of the Young diagram of λ is paired with a box of $n \times \ell$ outside λ .

Also, the weights of type A_{n-1} and $A_{\ell-1}$ associated to the partitions in the set

$$\mathcal{S}^{\check{\theta}}(s) = \{\lambda \in \mathcal{S}(s) \mid \check{\theta}(\lambda) = \lambda\}$$

are fixed simultaneously by θ and $\check{\theta}$ (since columns of height n and rows of length ℓ do not contribute to A_{n-1} and $A_{\ell-1}$ weights, respectively). Thus

$$F(s)^{\check{\theta}} = \{b \in F(s)^{\theta, \check{\theta}} \mid b = f_i f_i^* b_{\lambda} \text{ with } \check{\theta}(\lambda) = \lambda\}$$

has a bicrystal structure exactly as in Theorem 5.10. Also, note that when ℓ and n are both odd, the set $\mathcal{S}^{\check{\theta}}(s)$ is empty.

Example 5.13. Assume $n = 4$ and $\ell = 3$. Then the set $\mathcal{S}^{\hat{\theta}}(s)$ contains exactly the partitions $(2, 2, 1, 1), (3, 3, 0, 0), (3, 2, 1, 0)$. We have

$$\sum_{b \in F(s)^{\hat{\theta}}} x^{\text{wt}(b)'} y^{\text{wt}(b^*)'} = s_{(1,1)}^{(B_2)}(x) s_{(2)}^{(\dot{C}_1)} + s_{(3,3)}^{(B_2)}(x) s_{(0)}^{(\dot{C}_1)}(y) + s_{(2,1)}^{(B_2)}(x) s_{(1)}^{(\dot{C}_1)}(y).$$

5.4. Bicrystal structure on fixed points sets: the affine case. Consider the type $A_{n-1}^{(1)}$ Dynkin diagram automorphism $\theta : i \mapsto -i \bmod n$, for all $i = 0, \dots, n-1$. It induces an involution on the cone of dominant weights P_+ sending each fundamental weight $\omega_i, i = 0, \dots, n-1$ on $\omega_{-i \bmod n}$. Let P_+^{θ} be the subset of P_+ of dominant weights fixed by θ . When $n = 2m$ is even (respectively $n = 2m-1$ is odd), there is a bijection $\lambda \mapsto \lambda'$ between the sets P_+^{θ} and the set $P_+^{(D_{m+1}^{(2)})}$ (respectively $P_+^{(A_{2(m-1)}^{(2)})}$) of dominant weights for the root system $D_{m+1}^{(2)}$ (respectively $A_{2(m-1)}^{(2)}$) (in Kac's classification of affine Dynkin diagrams [22]).

For any $\lambda \in P_+$, similarly to the classical case, we have a crystal anti-isomorphism also denoted by θ from $B(\lambda)$ to $B(\theta(\lambda))$ which flips the labels i and $-i \bmod n$ of the arrows. Assuming that $\theta(\lambda) = \lambda$, one gets an involution on $B(\lambda)$ (also denoted by θ) and it makes sense to set $B^{\theta}(\lambda) = \{b \in B(\lambda) \mid \theta(b) = b\}$. Let us define some modified crystal operators by

$$\begin{aligned} f_i^{\theta} &= f_i f_{2m-i}, i = 1, \dots, n-1, & f_0^{\theta} &= f_0 \text{ and } & f_m^{\theta} &= f_m \text{ when } n = 2m, \\ f_i^{\theta} &= f_i f_{2m-1-i}, i = 1, \dots, n-1, & f_0^{\theta} &= f_0 \text{ and } & f_m^{\theta} &= f_m f_{m-1}^2 f_m \text{ when } n = 2m-1. \end{aligned}$$

In [36], it was proved that when $n = 2m-1$ (respectively $n = 2m$) these operators stabilize $B(\lambda)^{\theta}$ and the crystal structure obtained in this way is isomorphic to $B^{(D_{m+1}^{(2)})}(\lambda')$ (respectively to $B^{(A_{2m}^{(2)})}(\lambda')$). Write for short $B^{(X)}(\lambda')$ the crystal obtained in both cases.

By Lemma 3.26, for any $\mathbf{s} \in \mathcal{D}(s)$, the image of the vertex $b = b_{\mathbf{s}}$ by the duality \star is $b_{\mathbf{s}^{\star}}$ with $\mathbf{s}^{\star} \in \mathcal{D}(s)$. Recall that we had denoted $\omega_{\mathbf{s}} = \omega_{s_{\ell}} + \dots + \omega_{s_1}$ for $\mathbf{s} \in \mathcal{D}(s)$ and $\dot{\omega}_{\dot{\mathbf{s}}} = \dot{\omega}_{\dot{s}_n} + \dots + \dot{\omega}_{\dot{s}_1}$ for $\dot{\mathbf{s}} \in \mathcal{D}(s)$. By mimicking the construction in Section 5.3, one can define, for any combinatorial Fock space $\hat{F}(s)$, the crystal $\hat{F}(s)^{\theta}$ as the crystal generated from the triple highest weight vertices $b_{\mathbf{s}}$ with $\mathbf{s} \in \mathcal{D}(s)^{\theta}$ where

$$\mathcal{D}(s)^{\theta} = \{\mathbf{s} \in \mathcal{D}(s) \mid \theta(\omega_{\mathbf{s}}) = \omega_{\mathbf{s}}\}.$$

Then, $\hat{F}(s)^{\theta}$ has the structure of an $(X \times A_{\infty} \times A_{\ell-1}^{(1)})$ -crystal. Similarly, by considering $\dot{\theta}$ the Dynkin diagram automorphism of type $A_{\ell-1}^{(1)}$, one can define the crystal $\hat{F}(s)^{\dot{\theta}}$, which has an $(A_{n-1}^{(1)} \times A_{\infty} \times \dot{X})$ -crystal structure. Finally, write

$$\mathcal{D}(s)^{\theta, \dot{\theta}} = \{\mathbf{s} \in \mathcal{D}(s) \mid \theta(\omega_{\mathbf{s}}) = \omega_{\mathbf{s}} \text{ and } \dot{\theta}(\dot{\omega}_{\mathbf{s}^{\star}}) = \dot{\omega}_{\mathbf{s}^{\star}}\}$$

Example 5.14. Assume $\mathbf{s} = (s_1, \dots, s_{\ell})$ belong to $\mathcal{D}(s)$ and is such that $0 \leq s_1 \leq s_1 \leq \dots \leq s_{\ell} < n$ with $s_j = s_{\ell-j+1}$ for any $j = 1, \dots, \ell$. Then we clearly have $\theta(\mathbf{s}) = \mathbf{s}$. But by Lemma 3.26 we also get $\dot{\theta}(\mathbf{s}) = \mathbf{s}$ since for any $j = 1, \dots, \ell-1$ we have $s_{j+1} - s_j = (n - s_j) - (n - s_{j+1}) = s_{\ell-j+1} - s_{\ell-j}$. In particular $\mathbf{s} \in \mathcal{D}(s)^{\theta, \dot{\theta}}$.

Now define $\hat{F}(s)^{\theta, \dot{\theta}}$ as the crystal generated by the operators f_i^{θ} and $(f_i^{\dot{\theta}})^{\star}$ applied on the triple highest weight vertex $b_{\mathbf{s}}$ with $\mathbf{s} \in \mathcal{D}(s)^{\theta, \dot{\theta}}$. Then $\hat{F}(s)^{\theta, \dot{\theta}}$ has the structure of an $(X \times A_{\infty} \times \dot{X})$ -crystal. The crystals $\hat{F}(s)^{\theta}, \hat{F}(s)^{\dot{\theta}}$ and $\hat{F}(s)^{\theta, \dot{\theta}}$ can then be regarded as combinatorial Fock spaces carrying a crystal structure other than type A .

When $n = 2m$ is even, one can also define the subset P_+^{η} (respectively P_+^{ζ}) of P_+^{θ} of weights with an even ω_0 -coordinate (respectively with even ω_0 and ω_m -coordinates). Then, there is a bijective map $\omega_{\mathbf{s}} \mapsto \omega_{\mathbf{s}}^{\dagger}$ between P_+^{η} and $P_+^{(\tilde{A}_{2m}^{(2)})}$ the set of dominant weights for the affine root system $\tilde{A}_{2m}^{(2)}$. We also have a bijective map $\omega_{\mathbf{s}} \mapsto \omega_{\mathbf{s}}^{\ddagger}$ between P_+^{ζ} and $P_+^{(C_m^{(1)})}$ the set of dominant weights for the affine root system $C_m^{(1)}$. Let us define some modified crystal operators by

$$\begin{aligned} f_i^{\eta} &= f_i f_{2m-i}, i = 1, \dots, n-1, & f_0^{\eta} &= f_0^2 \text{ and } & f_m^{\eta} &= f_m, \\ f_i^{\zeta} &= f_i f_{2m-i}, i = 1, \dots, n-1, & f_0^{\zeta} &= f_0^2 \text{ and } & f_m^{\zeta} &= f_m^2. \end{aligned}$$

Write $\hat{F}(s)^{\eta}$ (respectively $\hat{F}(s)^{\zeta}$) for the subcrystal of $\hat{F}(s)$ obtained by applying the operators f_i^{η} (respectively the operators f_i^{ζ}) to vertices of the form $b_{\mathbf{s}}$ with \mathbf{s} fixed by η (respectively ζ). By the results of [23], one gets that $\hat{F}(s)^{\eta}$

and $\widehat{F}(s)^\zeta$ are isomorphic to the crystals $F^{\tilde{A}_{2m}^{(2)}}(s^+)$ and $F^{C_m^{(1)}}(s^\pm)$ of type $\tilde{A}_{2m}^{(2)}$ and $C_m^{(1)}$, respectively. From any combinatorial Fock space $\widehat{F}(s)$, one can then define the combinatorial Fock spaces $\widehat{F}(s)^\eta$, $F^\eta(s)$, $\widehat{F}(s)^\zeta$ and $\widehat{F}(s)^\zeta$ as previously and also get triple structures of crystal with one structure of affine type other than A .

In general, one can define sets

$$\mathcal{D}(s)^{\sharp, \flat} = \{\mathbf{s} \in \mathcal{D}(s) \mid \omega_{\mathbf{s}} \in P_+^\sharp \text{ and } \dot{\omega}_{\mathbf{s}} \in \dot{P}_+^\flat\}$$

where the symbols \sharp and \flat belong to the set $\{\theta, \eta, \zeta\}$. We can then define $\widehat{F}(s)^{\sharp, \flat}$ as the crystal generated by the operators f_i^\sharp and $(f_i^\flat)^*$ applied on the triple highest weight crystal b_s with $\mathbf{s} \in \mathcal{D}(s)^{\sharp, \flat}$. It admits the structure of an $(X \times A_\infty \times \check{X})$ -crystal described by the table below.

	$\theta, n = 2m - 1$	$\theta, n = 2m$	$\eta, n = 2m$	$\zeta, n = 2m$
$\dot{\theta}, \ell = 2p - 1$	$A_{2(m-1)}^{(2)} \times A_\infty \times \dot{A}_{2(p-1)}^{(2)}$	$D_{m+1}^{(2)} \times A_\infty \times \dot{A}_{2(p-1)}^{(2)}$	$\tilde{A}_{2m}^{(2)} \times A_\infty \times \dot{A}_{2(p-1)}^{(2)}$	$C_m^{(1)} \times A_\infty \times \dot{A}_{2(p-1)}^{(2)}$
$\dot{\theta}, \ell = 2p$	$A_{2(m-1)}^{(2)} \times A_\infty \times \dot{D}_{p+1}^{(2)}$	$D_{m+1}^{(2)} \times A_\infty \times \dot{D}_{p+1}^{(2)}$	$\tilde{A}_{2m}^{(2)} \times A_\infty \times \dot{D}_{p+1}^{(2)}$	$C_m^{(1)} \times A_\infty \times \dot{D}_{p+1}^{(2)}$
$\dot{\eta}, \ell = 2p$	$A_{2(m-1)}^{(2)} \times A_\infty \times \check{A}_{2p}^{(2)}$	$D_{m+1}^{(2)} \times A_\infty \times \check{A}_{2p}^{(2)}$	$\tilde{A}_{2m}^{(2)} \times A_\infty \times \check{A}_{2p}^{(2)}$	$C_m^{(1)} \times A_\infty \times \check{A}_{2p}^{(2)}$
$\dot{\zeta}, \ell = 2p$	$A_{2(m-1)}^{(2)} \times A_\infty \times \dot{C}_p^{(1)}$	$D_{m+1}^{(2)} \times A_\infty \times \dot{C}_p^{(1)}$	$\tilde{A}_{2m}^{(2)} \times A_\infty \times \dot{C}_p^{(1)}$	$C_m^{(1)} \times A_\infty \times \dot{C}_p^{(1)}$

In each case, $\widehat{F}(s)^{\sharp, \flat}$ can be regarded as a combinatorial Fock space carrying a triple crystal structure, two of them being of affine type other than A . Observe that this gives rises to all possible classical affine crystal structures except those corresponding to Dynkin diagrams containing a sub-Dynkin diagram of classical type D .

6. PROMOTION OPERATOR AND A GENERALISATION OF PITMAN'S 2M-X TRANSFORM

In this section, we first relate the Pitman transform 2M-X to the affine crystals $A_1^{(1)}$ and show how the energy can be used to prove that the successive iterations of this transform on trajectories on \mathbb{Z} with steps ± 1 tend to the trivial trajectory, all of whose steps are equal to 1. We next define a transformation analogue in higher dimension and establish that it also yields a natural convergence of trajectories.

6.1. Affine type $A_1^{(1)}$ -crystal and Pitman's 2M-X transform. In this paragraph, we shall consider tensor products $B^{\otimes n}$ of the affine Kirillov-Reshetikhin crystal of type $A_1^{(1)}$ where

$$B : 1 \xrightleftharpoons[1]{0} 2.$$

There is a straightforward bijection between the vertices $b^* = \varepsilon_1 \otimes \cdots \otimes \varepsilon_n \in B^{\otimes n}$ and the trajectories π of length n on the set \mathbb{Z} of integers starting at 0 with steps $+1$ or -1 defined by $\pi(k) = \pi_+(k) - \pi_-(k)$ for any $k = 0, \dots, n$ where $\pi_+(k)$ (respectively $\pi_-(k)$) is the numbers of letters 1 (respectively of letters 2) in $\varepsilon_1 \otimes \cdots \otimes \varepsilon_k$. In the sequel, we shall abuse the notation and identify the vertices b^* with their associated path π . This corresponds to the Littelman path model for A_1 and it is easy to check that b^* is a Yamanouchi word if and only if $\pi(k) \geq 0$ for any $k = 0, \dots, n$.

We now define the two Pitman transforms \mathcal{P}_{\min} and \mathcal{P}_{\max} on the trajectories $\pi \in B^{\otimes n}$ by

$$\mathcal{P}_{\min}(\pi)(k) = \pi(k) - 2 \min_{0 \leq a \leq k} \pi(a) \text{ and } \mathcal{P}_{\max}(\pi)(k) = 2 \max_{0 \leq a \leq k} \pi(a) - \pi(k)$$

for any $0 \leq k \leq n$. The following properties are easy to check.

- (1) The image by \mathcal{P}_{\min} or \mathcal{P}_{\max}^* of any trajectory $\pi \in B^{\otimes n}$ is a trajectory which always remains nonnegative.
- (2) The trajectory $\mathcal{P}_{\min}(\pi)$ corresponds to the highest weight vertex associated to π in $B^{\otimes n}$ for the A_1 -structure obtained by deleting the 0-arrows.
- (3) The nonnegative trajectories are fixed by the transformation \mathcal{P}_{\min} (since we then have $\inf_{0 \leq a \leq k} \pi(a) = 0$ for any $0 \leq a \leq n$). This is not true for the transformation \mathcal{P}_{\max} .
- (4) We have $\mathcal{P}_{\min}(\pi) = \pi$ if and only if $\min_{0 \leq a \leq k} \pi(a) = 0$ for any $0 \leq k \leq n$, that is π remains nonnegative.
- (5) We have $\mathcal{P}_{\max}(\pi) = \pi$ if and only if $\max_{0 \leq a \leq k} \pi(a) = \pi(k)$ for any $0 \leq k \leq n$. This means that $\pi(k) = k$ for any $0 \leq k \leq n$, i.e. π is the trivial trajectory π_0 whose all steps are equal to 1.

In the particular case $A_1^{(1)}$, the promotion operator pr acts on each trajectory π just by flipping the steps $+1$ and -1 . Thus, the path $\text{pr}(\pi)$ is obtain by reflecting the path π i.e. we have $\text{pr}(\pi)(k) = -\pi(k)$ for any $0 \leq k \leq n$. This implies that

$$\mathcal{P}_{\max} = \mathcal{P}_{\min} \circ \text{pr}.$$

By using the results of Section 2.7, we get the following proposition

Proposition 6.1.

- For any π in $B^{\otimes n}$, we have $D(\mathcal{P}_{\min}(\pi)) = D(\pi)$: the Pitman transform \mathcal{P}_{\min} preserves the energy D .
- For any nonnegative trajectory π in $B^{\otimes n}$, we have $D(\text{pr}(\pi)) = D(\pi) - \pi_-$ where π_- is equal to the number of steps -1 in π : the promotion operator makes decrease the energy of a trajectory as the number of its negative steps.
- For any trajectory π in $B^{\otimes n}$, there exists an integer m_0 such that for any $m \geq m_0$, we have $\mathcal{P}_{\max}^{(m)}(\pi) = \pi_0$.

Proof. The first claim follows from the fact that $\mathcal{P}_{\min}(\pi)$ is the highest weight vertex of the A_1 -connected component containing π and D is constant on classical components. The second one is a consequence of Lemma 2.50. Finally, the sequence of integers $D(\pi^{(m)})$ is nonnegative, strictly decreasing while $\pi^{(m)}$ contains at least a step -1 . It will eventually becomes equal to zero for m sufficiently large. Then $\pi^{(m)} = \pi_0$ because π_0 is the unique nonnegative trajectory such that $D(\pi_0) = 0$ (or the unique fixed point by the transform \mathcal{P}_{\max}). \square

Example 6.2. Starting with $\pi = 112121$, we get

$$\begin{aligned} \text{pr}(\pi) &= 221212, & \mathcal{P}_{\max} &= \mathcal{P}_{\min} \circ \text{pr}(\pi) = 111212, & \text{pr}\mathcal{P}_{\max}(\pi) &= 222121, \\ \mathcal{P}_{\max}^2(\pi) &= 111121, & \text{pr} \circ \mathcal{P}_{\max}^2(\pi) &= 222212, & \mathcal{P}_{\max}^3(\pi) &= 111112, \\ \text{pr} \circ \mathcal{P}_{\max}^3(\pi) &= 222221, & \mathcal{P}_{\max}^4(\pi) &= 111111. \end{aligned}$$

Remark 6.3. In his seminal article [38], Pitman proves that the image by the transforms \mathcal{P}_{\max} and \mathcal{P}_{\min} of a one-dimensional Brownian motion is a 3-dimensional Bessel process (i.e. the norm of a 3-dimensional Brownian motion). One can replace this Brownian motion by a random walk with transitions $+1$ and -1 and related probabilities p_1, p_{-1} where $p_1 + p_{-1} = 1$. Its image by \mathcal{P}_{\min} or \mathcal{P}_{\max} yields a Markov chain on $\mathbb{Z}_{\geq 0}$. Later, it was observed by Biane, Bougerol and O’Connell [1] that \mathcal{P}_{\min} can be interpreted in Littelmann’s path theory as the transform associating to each path its corresponding highest weight path. The previous one-dimensional results then admit higher dimensional generalisations (see for example [1] and [32]). As far as we are aware, our interpretation of \mathcal{P}_{\max} in terms of affine crystals is new.

6.2. Generalisation to higher dimension. The generalised Pitman transform \mathcal{P}_{\min} introduced in [1] is defined for any finite root system R with Dynkin diagram I . It associates to each Littelmann path, its highest weight path. Here, one can first define a Pitman transform $\mathcal{P}_{i,\min}$ for any node $i \in I$ as in Section 6.1: it just computes the highest weight path for the A_1 -crystal corresponding to the node i . Then one has $\mathcal{P}_{\min} = \mathcal{P}_{\min,i_1} \cdots \mathcal{P}_{\min,i_r}$ where $s_{i_1} \cdots s_{i_r}$ is a decomposition of the longest element w_0 of the Weyl group of R as a product of elementary reflections $s_i, i \in I$. In particular \mathcal{P}_{\min} does not depend on the reduced decomposition considered.

Proposition 6.1 suggests to look for a natural higher dimensional generalisation of the results of Section 2.4. Let B be the $A_{\ell-1}^{(1)}$ -crystal with set of vertices $\{1, 2, \dots, \ell\}$ such that for any $j = 0, \dots, \ell - 1$ we have

$$\dot{f}_j(k) = \begin{cases} k + 1 \bmod \ell & \text{if } k = j \bmod \ell, \\ 0 & \text{otherwise.} \end{cases}$$

This is the simplest nontrivial Kirillov-Reshetikhin crystal of type $A_{\ell-1}^{(1)}$ associated to the rectangle 1×1 . For any integer n , the vertices of $\dot{B}_{n,\ell} = B^{\otimes n}$ coincide with the tensor product $b^* = d_1 \otimes \cdots \otimes d_n$ of n columns with height 1 on the alphabet $\{1, \dots, \ell\}$. With the notation of Section 2.4, the corresponding vectors $b \in F(n)$ are the tensor products $b = c_\ell \otimes \cdots \otimes c_1$ of ℓ -columns in which each letter of $\{1, \dots, n\}$ appears exactly once. Set $b_{\text{fin}}^* = 1^{\otimes n}$. Then $b_{\text{fin}} = \emptyset \otimes \emptyset \otimes \cdots \otimes c_{1,\dots,n}$ where $c_{1,\dots,n}$ is the column containing exactly the letters $\{1, \dots, n\}$.

One might first define a transform \mathcal{P}'_{\max} as in [1] by setting $\mathcal{P}'_{\max} = \mathcal{P}_{\max,i_1} \cdots \mathcal{P}_{\max,i_r}$ where $w_0 = s_{i_1} \cdots s_{i_r}$ is a reduced decomposition of w_0 . But then, as illustrated by the following example, \mathcal{P}'_{\max} would depend on the chosen reduced decomposition.

Example 6.4. Assume $\ell = 3$ and $n = 5$. To apply $\mathcal{P}_{\max,1}$ (respectively $\mathcal{P}_{\max,2}$) to a vertex b^* in $\dot{B}_{5,3}$, we have first to flip the letters 1 and 2 (respectively 2 and 3) and next compute the source vertex of the 1-chain (respectively

the 2-chain) corresponding to the vertex so obtained. For the vertex $b^* = 2 \otimes 1 \otimes 2 \otimes 3 \otimes 2 = 21232$ (we omit the symbol \otimes for short), we get

$$\begin{aligned} 21232 &\xrightarrow{\mathcal{P}_{\max,1}} 12131 \xrightarrow{\mathcal{P}_{\max,2}} 12121 \xrightarrow{\mathcal{P}_{\max,1}} 11212, \\ 21232 &\xrightarrow{\mathcal{P}_{\max,2}} 21223 \xrightarrow{\mathcal{P}_{\max,1}} 12113 \xrightarrow{\mathcal{P}_{\max,2}} 12112. \end{aligned}$$

Thus $\mathcal{P}_{\max,1}\mathcal{P}_{\max,2}\mathcal{P}_{\max,1} \neq \mathcal{P}_{\max,2}\mathcal{P}_{\max,1}\mathcal{P}_{\max,2}$.

In order to generalize Proposition 6.1, we rather set

$$\mathcal{P}_{\max} : \begin{cases} \dot{B}_{n,\ell} \rightarrow \dot{B}_{n,\ell} \\ b^* \rightarrow \text{pr} \circ \mathcal{P}_{\min}(b^*) \end{cases}$$

where \mathcal{P}_{\min} is the generalised Pitman transform of [1], that is $\mathcal{P}_{\min}(b^*)$ is the highest weight vertex of the $A_{\ell-1}$ -connected component containing b^* .

Theorem 6.5.

- (1) For any vertex b^* in $\dot{B}_{n,\ell}$, we have $D(\mathcal{P}_{\min}(b^*)) = D(b^*)$.
- (2) For any vertex b^* in $\dot{B}_{n,\ell}$, there exists an integer m_0 such that for any $m \geq m_0$, we have $\mathcal{P}_{\max}^{(m)}(b^*) = b_{\text{fin}}^*$.

Proof. Assertion (1) follows from the fact that the energy D is constant over classical components and $\mathcal{P}_{\min}(b^*)$ is the highest weight vertex associated to b^* . For Assertion (2), set $b_m^* = \text{pr}^{-1} \circ \mathcal{P}_{\max}^{(m)}(b^*)$ for any $m \geq 0$. Then $b_1^* = \mathcal{P}_{\min}(b^*)$ and thus b_1 is a tableau. More generally, by using the results of Section 2.7, we get that the sequence $b_m, m \geq 1$ coincides with the sequence of tableaux $\zeta^{m-1}(b_1)$, that is with the cyclage sequence defined from the standard tableau b_1 . Since all these tableaux are standard, this sequence is indeed well-defined and eventually ends on the column $c_{1,\dots,n}$. This means that there exists an integer m_0 such that $\mathcal{P}_{\max}^{(m_0)}(b^*) = b_{\text{fin}}^*$. Since $\mathcal{P}_{\max}^{(a)}(b_{\text{fin}}^*) = b_{\text{fin}}^*$ for any $a \geq 0$, we get $\mathcal{P}_{\max}^{(m)}(b^*) = b_{\text{fin}}^*$ for any $m \geq m_0$. \square

Remark 6.6.

- (1) By slightly generalizing the notion of an authorised cyclage operation, it is possible to define an analogue of \mathcal{P}_{\max} on any tensor product of columns (not only for columns of height 1) which yields a similar convergence property.
- (2) In [32], random walks are defined from tensor products of crystals. In particular, one can endow the crystal $\dot{B}_{n,\ell}$ with a probability distribution compatible with the $A_{\ell-1}$ -weight graduation (i.e.; two vertices with the same $A_{\ell-1}$ -weight have the same probability). This permits to define a random walk on the weight lattice of type $A_{\ell-1}$ whose image by \mathcal{P}_{\min} is a Markov chain in the Weyl chamber. We thought that our transform \mathcal{P}_{\max} also admits interesting probabilistic properties that we aim to study later.
- (3) For the other classical affine root systems, it is also possible to define an analogue of the transformation \mathcal{P}_{\max} by using the results of [33] which essentially reduces their study (and notably the computation of the energy) to affine type A crystals by using relevant Dynkin diagram automorphisms.

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