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COURANT-SHARP EIGENVALUES OF COMPACT FLAT SURFACES: KLEIN BOTTLES AND CYLINDERS

PIERRE BÉRARD, BERNARD HELFFER, AND ROLA KIWAN

ABSTRACT. The question of determining for which eigenvalues there exists an eigenfunction which has the same number of nodal domains as the label of the associated eigenvalue (Courant-sharp property) was motivated by the analysis of minimal spectral partitions. In previous works, many examples have been analyzed corresponding to squares, rectangles, disks, triangles, tori, Möbius strips, A natural toy model for further investigations is the flat Klein bottle, a non-orientable surface with Euler characteristic 0, and particularly the Klein bottle associated with the square torus, whose eigenvalues have higher multiplicities. In this note, we prove that the only Courant-sharp eigenvalues of the flat Klein bottle associated with the square torus (resp. with square fundamental domain) are the first and second eigenvalues. We also consider the flat cylinders $(0, \pi) \times \mathbb{S}_r^1$ where $r \in \{0.5, 1\}$ is the radius of the circle \mathbb{S}_r^1 , and we show that the only Courant-sharp Dirichlet eigenvalues of these cylinders are the first and second eigenvalues.

1. INTRODUCTION

Given a compact Riemannian surface (M, g) , we write the eigenvalues of the Laplace-Beltrami operator $-\Delta_g$,

$$(1.1) \quad \lambda_1(M) < \lambda_2(M) \leq \lambda_3(M) \leq \dots ,$$

in nondecreasing order, starting from the label 1, with multiplicities accounted for. If the boundary ∂M of M is non-empty, we consider Dirichlet eigenvalues.

Courant's nodal domain theorem (1923) states that any eigenfunction associated with the eigenvalue λ_k has at most k nodal domains (connected components of the complement of the zero set of u). The eigenvalue λ_k is called *Courant-sharp* if there exists an associated eigenfunction with precisely k nodal domains. It follows from Courant's theorem that the eigenvalues λ_1 and λ_2 are Courant-sharp, and that $\lambda_{k-1} < \lambda_k$ whenever λ_k is Courant-sharp. Courant-sharp eigenvalues appear naturally in the context of partitions, [11].

Back in 1956, Pleijel proved that there are only finitely many Courant-sharp eigenvalues. More precisely, Pleijel's original proof [24] applied to Dirichlet eigenvalues of bounded domains in \mathbb{R}^2 . It was later adapted to more general domains ([23]) and to general closed Riemannian manifolds ([6]).

Courant-sharp eigenvalues of flat tori and Möbius strips were studied in [13, 19, 4] and [5] respectively. In view of the classification of complete, flat surfaces (see [18, p. 222-223], or [26, Chap. 2.5]) it is natural to investigate the Courant-sharp property

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for the other compact flat surfaces, Klein bottles and cylinders. The purpose of the present paper is to prove the following theorems.

Theorem 1.1. *Let \mathbb{K}_1 denote the flat Klein bottle associated with the square torus, resp. \mathbb{K}_2 the flat Klein bottle whose fundamental domain is a square. Then the only Courant-sharp eigenvalues of \mathbb{K}_c , $c \in \{1, 2\}$, are λ_1 and λ_2 .*

Theorem 1.2. *Let \mathcal{C}_r denote the flat cylinder $(0, \pi) \times \mathbb{S}_r^1$, with the product metric. Here \mathbb{S}_r^1 denotes the circle with radius r . Then, for $r \in \{\frac{1}{2}, 1\}$, the only Courant-sharp Dirichlet eigenvalues of \mathcal{C}_r , are λ_1 and λ_2 .*

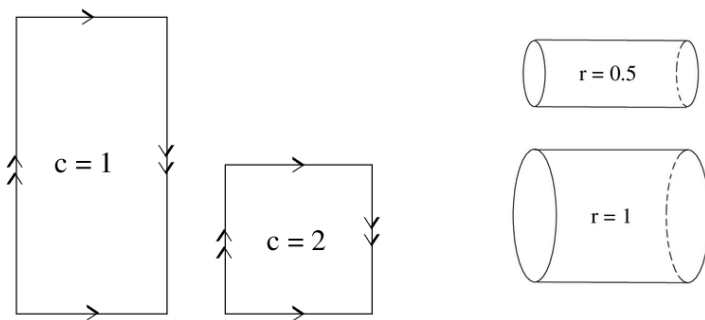


FIGURE 1.1. The Klein bottles \mathbb{K}_c (left) and the cylinders \mathcal{C}_r (right).

Courant-sharp eigenvalues have previously been determined for several compact surfaces. We refer to the following papers and their bibliographies.

- ◇ Closed surfaces: round 2-sphere and projective plane, [22]; flat tori, [13, 19, 4].
- ◇ Plane domains and compact surfaces with boundary (different boundary conditions might be considered): square, [24, 3, 16, 10, 9]; equilateral triangle, [4]; 2-rep-tiles, [1]; thin cylinders, [12]; Möbius strips, [5].

There are also a few results in higher dimensions, see for example [15, 20, 14].

Most of the papers mentioned above adapt the method introduced by Pleijel in [24] to the example at hand.

The paper is organized as follows. In Section 2, we recall the main lines of Pleijel's method. More precisely, we show how a lower bound on the ratio $\frac{\lambda_k(M)}{k}$ and a lower bound on the Weyl counting function can be used to restrict the search for Courant-sharp eigenvalues to a finite set of eigenvalues. To actually determine the Courant-sharp eigenvalues one then needs to analyze the nodal patterns of eigenfunctions in a finite set of eigenspaces. In Section 3, we recall some basic facts concerning Klein bottles. In Sections 4 and 5, we adapt Pleijel's method (Section 2) to the flat Klein bottles \mathbb{K}_1 and \mathbb{K}_2 . In Section 6, we adapt Pleijel's method (Section 2) to the flat cylinders \mathcal{C}_r , with $r \in \{\frac{1}{2}, 1\}$.

2. PLEIJEL'S METHOD SUMMARIZED

In order to prove that the number of Courant-sharp (Dirichlet) eigenvalues of a compact Riemannian surface (M, g) is finite, one only needs two ingredients,

- (1) the classical Weyl asymptotic law, $W_M(\lambda) = \frac{|M|}{4\pi} \lambda + O(\sqrt{\lambda})$, for the counting function $W_M(\lambda) = \#\{j \mid \lambda_j(M) < \lambda\}$, where $|M|$ denotes the area of M ;
- (2) an inequality à la Faber-Krahn for domains $\omega \subset M$ with small enough area,

$$(2.1) \quad \delta_1(\omega) \geq (1 - \varepsilon)^2 \frac{\pi j_{0,1}^2}{|\omega|} \quad \text{provided that } |\omega| \leq C(M, \varepsilon) |M|,$$

for some $\varepsilon \in [0, 1)$, and for some constant $C(M, \varepsilon)$. Here, $\delta_1(\omega)$ denotes the least Dirichlet eigenvalue of ω , and $j_{0,1}$ is the first positive zero of the Bessel function J_0 ($j_{0,1} \approx 2.404825$).

Note that the right-hand side of (2.1) is equal to $(1 - \varepsilon)^2 \delta_1(\omega^*)$, where ω^* denotes a disk in \mathbb{R}^2 with area equal to $|\omega|$. The existence of such an inequality (for a given $\varepsilon \in (0, 1)$) follows from the asymptotic isoperimetric inequality proved in [6, Appendice C], and from the usual symmetrization argument to translate the isoperimetric inequality into a Faber-Krahn type inequality for the first Dirichlet eigenvalue, [6, Lemme 15].

In order to give more quantitative information on the Courant-sharp eigenvalues, one needs a lower bound on the Weyl function, in the form,

$$(2.2) \quad W_M(\lambda) \geq \frac{|M|}{4\pi} \lambda - A(M) \sqrt{\lambda} - B(M), \quad \text{for all } \lambda \geq 0$$

for some constants $A(M)$ and $B(M)$ depending on the geometry of (M, g) .

Lemma 2.1. *Assume that (M, g) satisfies (2.1), with $\varepsilon < 1 - \frac{2}{j_{0,1}}$, and (2.2). Let $\lambda_k(M)$ be a Courant-sharp eigenvalue with $k \geq \frac{1}{C(M, \varepsilon)}$. Then,*

$$(2.3) \quad \frac{\lambda_k(M)}{k} \geq (1 - \varepsilon)^2 \frac{\pi j_{0,1}^2}{|M|},$$

$$(2.4) \quad \frac{|M|}{(1 - \varepsilon)^2 \pi j_{0,1}^2} \lambda_k(M) \geq k \geq \frac{|M|}{4\pi} \lambda_k(M) - A(M) \sqrt{\lambda_k(M)} - B(M) + 1,$$

In particular, $F_{M, \varepsilon}(\lambda_k(M)) \leq 0$, where

$$(2.5) \quad F_{M, \varepsilon}(\lambda) = \frac{|M|}{4\pi} \left(1 - \frac{4}{(1 - \varepsilon)^2 j_{0,1}^2} \right) \lambda - A(M) \sqrt{\lambda} - B(M) + 1,$$

so that $\sqrt{\lambda_k(M)} \leq D(M, \varepsilon)$, where $D(M, \varepsilon)$ is the largest root of the equation $F_{M, \varepsilon}(\lambda) = 0$.

Proof. To prove (2.3), choose any eigenfunction associated with $\lambda_k(M)$ and having precisely k nodal domains. Since $k \geq \frac{1}{C(M, \varepsilon)}$, one of them, call it ω , has area $|\omega| \leq \frac{|M|}{k} \leq C(M, \varepsilon) |M|$, so that its first Dirichlet eigenvalue satisfies (2.1). Then, use the fact that $\lambda_k(M) = \delta_1(\omega)$.

To prove (2.4), use the fact that $k - 1 = W_M(\lambda_k(M))$ (this is because the eigenvalue is Courant-sharp), and apply inequalities (2.2) and (2.3).

Since $j_{0,1} > 2$, choosing $\varepsilon < 1 - \frac{2}{j_{0,1}}$, the coefficient of the leading term in $F_{M, \varepsilon}$ is positive, and the function tends to infinity when λ tends to infinity. \square

Corollary 2.2. *For $\varepsilon < 1 - \frac{2}{j_{0,1}}$, in order to be Courant-sharp, the eigenvalue $\lambda_k(M)$ must satisfy $\lambda_{k-1}(M) < \lambda_k(M)$, and*

$$(2.6) \quad \begin{cases} \text{either } k < \frac{1}{C(M,\varepsilon)}, \\ \text{or } k \geq \frac{1}{C(M,\varepsilon)}, \sqrt{\lambda_k(M)} \leq D(M,\varepsilon), \text{ and } \frac{\lambda_k(M)}{k} \geq (1-\varepsilon)^2 \frac{\pi j_{0,1}^2}{|M|}. \end{cases}$$

To conclude whether the eigenvalue $\lambda_k(M)$ is actually Courant-sharp, it remains to determine the maximum number of nodal domains of an eigenfunction in the eigenspace $\mathcal{E}(\lambda_k(M))$.

Remarks 2.3.

- (1) In Sections 4 and 6, we will use the fact that we can choose $\varepsilon = 0$ in the isoperimetric inequality (2.1) for the flat Klein bottles and for the flat cylinders. This follows from [17], and was already used in [19, 4, 5] for the flat torus and for the Möbius band.
- (2) When $\partial M \neq \emptyset$, and for Neumann or Robin eigenfunctions of (M, g) , the inequality (2.1) can only be applied to a nodal domain which does not touch the boundary ∂M . To prove a result à la Pleijel for Neumann or Robin eigenfunctions, as in [25, 21, 16, 10, 9], it is necessary to take care of this difficulty.
- (3) The proof of Lemma 2.1 also yields the following inequality. Let $\kappa(k)$ denote the maximal number of nodal domains of an eigenfunction associated with λ_k . Then,

$$\limsup_k \frac{\kappa(k)}{k} \leq \gamma(2) := \left(\frac{2}{j_{0,1}} \right)^2 < 1.$$

This inequality, which generalizes Pleijel's inequality [24] valid for plane domains, is actually a particular case of a general result valid for any closed Riemannian manifold, [6]. The interesting feature is that the upper bound $\gamma(n) < 1$ only depends on the dimension.

3. PRELIMINARIES ON KLEIN BOTTLES

3.1. Klein bottles. In this note, we are interested in the flat Klein bottles. More precisely, given $a, b > 0$, we consider the isometries of \mathbb{R}^2 given by

$$(3.1) \quad \begin{cases} \tau_1 : (x, y) \mapsto (x, y + b), \\ \tau_2 : (x, y) \mapsto (x + a, y), \\ \tau : (x, y) \mapsto (x + \frac{a}{2}, b - y). \end{cases}$$

We denote by G_2 (resp. G) the group generated by τ_1 and τ_2 (resp. by τ_1 and τ). These groups act properly and freely by isometries on \mathbb{R}^2 equipped with the usual scalar product. Since $\tau^2 = \tau_2$, the group G_2 is a subgroup of index 2 of the group G . We denote by $\mathbb{T}_{a,b}$ (resp. $\mathbb{K}_{a,b}$) the torus \mathbb{R}^2/G_2 (resp. the Klein bottle \mathbb{R}^2/G). We equip $\mathbb{T}_{a,b}$ and $\mathbb{K}_{a,b}$ with the induced flat Riemannian metrics.

A fundamental domain for the action of G_2 (resp. G) on \mathbb{R}^2 is the rectangle $\mathcal{T}_{a,b} = (0, a) \times (0, b)$ (resp. the rectangle $\mathcal{K}_{a,b} = (0, \frac{a}{2}) \times (0, b)$, see Figure 3.1 (A)). The horizontal sides of $\mathcal{K}_{a,b}$ are identified with the same orientation, the vertical sides are identified with the opposite orientations.

The geodesics of the Klein bottle are the images of the lines in \mathbb{R}^2 under the Riemannian covering map $\mathbb{R}^2 \rightarrow \mathbb{K}_{a,b}$ (see [8]). They can be looked at in the fundamental domain $\mathcal{K}_{a,b}$, taking into account the identifications $(x, 0) \sim (x, b)$ and $(0, y) \sim (\frac{a}{2}, b - y)$. Among them, we have some special geodesics, see Figure 3.1 (B)-(C),

- ◇ $t \mapsto (t, 0)$ and $t \mapsto (t, \frac{b}{2})$, for $0 \leq t \leq \frac{a}{2}$, which are periodic geodesics of length $\frac{a}{2}$;
- ◇ for $0 < y_0 < \frac{b}{2}$, $\gamma_{y_0} : t \mapsto \begin{cases} (t, y_0), & 0 \leq t \leq \frac{a}{2}, \\ (t - \frac{a}{2}, b - y_0), & \frac{a}{2} \leq t \leq a, \end{cases}$ which is a periodic geodesic of length a ; the two horizontal lines in blue in Figure 3.1 (C) yield a periodic geodesic of the Klein bottle;
- ◇ for $0 \leq x_0 \leq \frac{a}{2}$, $t \mapsto (x_0, t)$, with $0 \leq t \leq b$, is a periodic geodesic of length b .

Remark 3.1. *The description of geodesics of the Klein bottles as projected lines implies that the shortest, nontrivial, periodic geodesic of $\mathbb{K}_{a,b}$ has length $\min\{\frac{a}{2}, b\}$.*

Remark 3.2. *Scissoring the Klein bottle along the blue geodesic $t \mapsto \gamma_{y_0}(t)$, with $0 < y_0 < \frac{b}{2}$ and $0 \leq t \leq a$, divides the surface into two Möbius strips whose center lines are the geodesics $t \mapsto (t, 0)$ and $t \mapsto (t, \frac{b}{2})$, see Figure 3.1 (D).*

The isometry τ of \mathbb{R}^2 induces an isometry on the torus $\mathbb{T}_{a,b}$ so that we can identify $\mathbb{K}_{a,b}$ with the quotient $\mathbb{T}_{a,b}/\{Id, \tau\}$. It follows that the eigenfunctions of the Klein bottle $\mathbb{K}_{a,b}$ are precisely the eigenfunctions of the torus $\mathbb{T}_{a,b}$ which are invariant under the map τ . Because τ is orientation reversing, the surface $\mathbb{K}_{a,b}$ is non-orientable with orientation double cover $\mathbb{T}_{a,b}$.

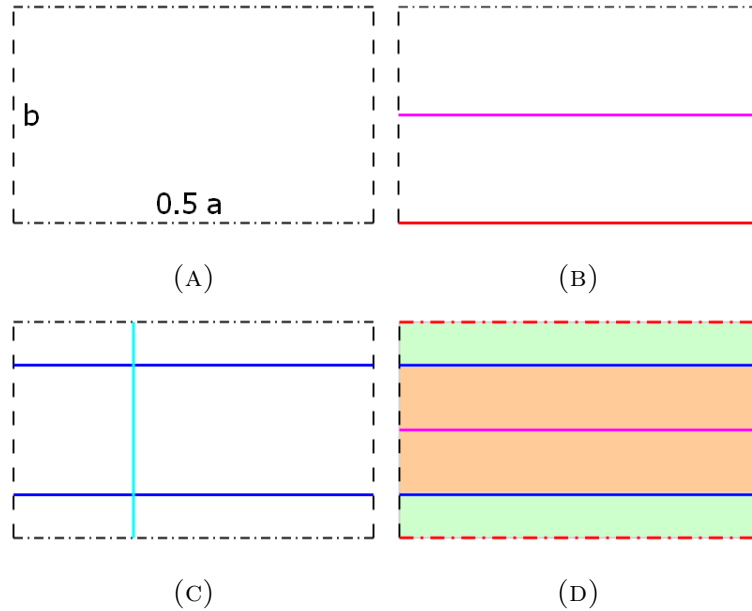


FIGURE 3.1. Fundamental domain, geodesics, partition into Möbius strips

3.2. The spectrum of Klein bottles. A complete family of (complex) eigenfunctions of the flat torus $\mathbb{T}_{a,b}$ is

$$(3.2) \quad f_{m,n}(x, y) = \exp\left(i\frac{2\pi mx}{a}\right) \exp\left(i\frac{2\pi ny}{b}\right), \quad m, n \in \mathbb{Z},$$

with associated eigenvalues $\hat{\lambda}(m, n) = 4\pi^2\left(\frac{m^2}{a^2} + \frac{n^2}{b^2}\right)$. Given some eigenvalue λ of $\mathbb{T}_{a,b}$, we introduce the set,

$$(3.3) \quad \mathcal{L}_\lambda = \left\{ (m, n) \in \mathbb{Z}^2 \mid \hat{\lambda}(m, n) = \lambda \right\}$$

A general (complex) eigenfunction of $\mathbb{T}_{a,b}$, with eigenvalue λ , is of the form

$$(3.4) \quad \phi = \sum_{(m,n) \in \mathcal{L}_\lambda} \alpha_{m,n} f_{m,n},$$

with $\alpha_{m,n} \in \mathbb{C}$.

The function ϕ is invariant under τ , $\phi = \phi \circ \tau$, if and only if

$$\sum_{(m,n) \in \mathcal{L}_\lambda} \alpha_{m,n} f_{m,n}(x, y) = \sum_{(m,n) \in \mathcal{L}_\lambda} \alpha_{m,n} (-1)^m f_{m,-n}(x, y),$$

or, equivalently, if and only if

$$(3.5) \quad \alpha_{m,-n} = (-1)^m \alpha_{m,n}, \quad \forall (m, n) \in \mathcal{L}_\lambda.$$

We can rewrite a τ -invariant eigenfunction ϕ as,

$$(3.6) \quad \phi = \sum_{(m,0) \in \mathcal{L}_\lambda, m \text{ even}} \alpha_{m,0} f_{m,0} + \sum_{(m,n) \in \mathcal{L}_\lambda, n > 0} \alpha_{m,n} (f_{m,n} + (-1)^m f_{m,-n}).$$

The following lemma follows readily.

Lemma 3.3 ([7]). *A complete family of real eigenfunctions of the flat Klein bottle $\mathbb{K}_{a,b}$ is given by the following functions.*

$$(3.7) \quad \begin{cases} \text{For } m = 0, n \in \mathbb{N} : & \cos\left(\frac{2\pi ny}{b}\right); \\ \text{for } m \in \mathbb{N}^\bullet \text{ even}, n \in \mathbb{N} : & \cos\left(\frac{2\pi mx}{a}\right) \cos\left(\frac{2\pi ny}{b}\right); \sin\left(\frac{2\pi mx}{a}\right) \cos\left(\frac{2\pi ny}{b}\right); \\ \text{for } m \in \mathbb{N}^\bullet \text{ odd}, n \in \mathbb{N}^\bullet : & \cos\left(\frac{2\pi mx}{a}\right) \sin\left(\frac{2\pi ny}{b}\right); \sin\left(\frac{2\pi mx}{a}\right) \sin\left(\frac{2\pi ny}{b}\right). \end{cases}$$

Here, \mathbb{N} denotes the set of non-negative integers, and \mathbb{N}^\bullet the set of positive integers.

Remark 3.4. *If $\mathcal{L}_\lambda \cap (\{0\} \times \mathbb{Z}) \neq \emptyset$, the multiplicity of λ is odd; if $\mathcal{L}_\lambda \cap (\{0\} \times \mathbb{Z}) = \emptyset$, the multiplicity of λ is even.*

3.3. Choices for a and b . In this paper, we restrict our attention to the case $a = b = 2\pi$, i.e. to the flat Klein bottle $\mathbb{K}_1 := \mathbb{K}_{2\pi, 2\pi}$, whose fundamental domain is the rectangle $\mathcal{K}_1 = (0, \pi) \times (0, 2\pi)$, and to the case $a = 2\pi, b = \pi$, i.e., to the flat Klein bottle $\mathbb{K}_2 := \mathbb{K}_{2\pi, \pi}$, whose fundamental domain is the square $\mathcal{K}_2 = (0, \pi) \times (0, \pi)$. As in [13] for flat tori, we could consider other values of the pair (a, b) , in particular we could look at what happens when a is fixed, b tends to zero, and vice-versa.

We denote the associated square flat tori by \mathbb{T}_1 and \mathbb{T}_2 respectively, with corresponding fundamental domains $\mathcal{T}_1 = (0, 2\pi) \times (0, 2\pi)$ and $\mathcal{T}_2 = (0, 2\pi) \times (0, \pi)$.

As points of the spectrum, the eigenvalues of the flat Klein bottle $\mathbb{K}_c, c \in \{1, 2\}$ are the numbers of the form $\hat{\lambda}(p, q) = p^2 + c^2 q^2$, with $p, q \in \mathbb{N}$, and the extra condition

that p is even when $q = 0$. As usual, the eigenvalues of \mathbb{K}_c are listed in nondecreasing order, multiplicities accounted for, starting from the label 1,

$$(3.8) \quad 0 = \lambda_1(\mathbb{K}_c) < \lambda_2(\mathbb{K}_c) \leq \lambda_3(\mathbb{K}_c) \leq \dots$$

For $\lambda \geq 0$, we introduce the *Weyl counting function*,

$$(3.9) \quad W_{\mathbb{K}_c}(\lambda) = \#\{j \mid \lambda_j(\mathbb{K}_c) < \lambda\}.$$

Weyl's asymptotic law tells us that

$$(3.10) \quad W(\lambda) = \frac{|\mathbb{K}_c|}{4\pi} \lambda + O(\sqrt{\lambda}) = \frac{\pi}{2c} \lambda + O(\sqrt{\lambda}),$$

where $|\mathbb{K}_c|$ denotes the area of the Klein bottle, namely $|\mathbb{K}_c| = \frac{2\pi^2}{c}$.

For later purposes, we also introduce the set

$$(3.11) \quad \mathcal{L}_c(\lambda) := \{(m, n) \in \mathbb{N}^2 \mid m^2 + c^2 n^2 < \lambda\}$$

and the *counting function*,

$$(3.12) \quad L_c(\lambda) = \#\left(\{(m, n) \in \mathbb{N}^2 \mid m^2 + c^2 n^2 < \lambda\}\right).$$

4. COURANT-SHARP EIGENVALUES OF THE KLEIN BOTTLES $\mathbb{K}_c, c \in \{1, 2\}$

The purpose of this section is to determine the Courant-sharp eigenvalues of the Klein bottles \mathbb{K}_c , following Pleijel's method, Section 2.

Lemma 4.1. *Let $\mathbb{K}_c, c \in \{1, 2\}$ be the flat Klein bottles introduced in Section 3. Let $\lambda_k(\mathbb{K}_c)$ be a Courant-sharp eigenvalue of \mathbb{K}_c , with label $k \geq \frac{2\pi}{c}$ (i.e. $k \geq 7$ when $c = 1$; $k \geq 4$ when $c = 2$).*

When $c = 1$,

$$(4.1) \quad \frac{\lambda_k(\mathbb{K}_1)}{k} \geq \frac{j_{0,1}^2}{2\pi} \geq 0.920422,$$

$$(4.2) \quad \frac{2\pi}{j_{0,1}^2} \lambda_k(\mathbb{K}_1) \geq k \geq \frac{\pi}{2} \lambda_k(\mathbb{K}_1) - 2\sqrt{\lambda_k(\mathbb{K}_1)} - 2.$$

When $c = 2$,

$$(4.3) \quad \frac{\lambda_k(\mathbb{K}_2)}{k} \geq \frac{j_{0,1}^2}{\pi} \geq 1.840844,$$

$$(4.4) \quad \frac{\pi}{j_{0,1}^2} \lambda_k(\mathbb{K}_2) \geq k \geq \frac{\pi}{4} \lambda_k(\mathbb{K}_2) - \frac{3}{2}\sqrt{\lambda_k(\mathbb{K}_2)} - 1.$$

In particular, $\lambda_k(\mathbb{K}_1) < 25$ and $\lambda_k(\mathbb{K}_2) \leq 47$.

Proof. By Remark 3.1, the shortest closed geodesic of $\mathbb{K}_c, c \in \{1, 2\}$, has length π . According to [17, § 7], any domain $\omega \subset \mathbb{K}_c$, with area $|\omega| \leq \pi$, satisfies the Euclidean isoperimetric inequality $|\partial\omega|^2 \geq 4\pi|\omega|$ (where $|\partial\omega|$ denotes the length of the boundary $\partial\omega$). It follows that

$$\delta_1(\omega) \geq \frac{\pi j_{0,1}^2}{|\omega|}.$$

If $\lambda_k(\mathbb{K}_c)$ is Courant-sharp with $k \geq \frac{2\pi}{c}$, then, for any associated eigenfunction u with precisely k nodal domains, there exists at least one nodal domain ω with $|\omega| \leq \frac{\mathbb{K}_c}{k} \leq \pi$, and we have

$$\lambda_k(\mathbb{K}_c) = \delta_1(\omega) \geq \frac{\pi j_{0,1}^2}{|\omega|} \geq k \frac{c j_{0,1}^2}{2\pi}.$$

This proves inequalities (4.1) and (4.3).

To prove the other inequalities, consider the set

$$\mathcal{E}_c(\lambda) := \left\{ (x, y) \mid 0 \leq x, y, x^2 + c^2 y^2 < \lambda \right\}.$$

Then,

$$\mathcal{E}_c(\lambda) \subset \bigcup_{(m,n) \in \mathcal{L}_c(\lambda)} [m, m+1] \times [n, n+1],$$

and hence

$$(4.5) \quad L_c(\lambda) = \#(\mathcal{L}_c(\lambda)) \geq |\mathcal{E}_c(\lambda)| = \frac{\pi}{4c} \lambda.$$

From the description of the spectrum of \mathbb{K}_c (Lemma 3.3), the contribution of a pair $(m, n) \in \mathcal{L}_c(\lambda)$, to $W_{\mathbb{K}_c}(\lambda)$ is

- a) 2 if $m, n \geq 1$;
- b) 1 if $m = 0$, and $n \geq 0$;
- c) 2 if $m \geq 2$ is even, and $n = 0$;
- d) 0 otherwise.

Recall that multiplicities also arise from the number of solutions of the equation $m^2 + c^2 n^2 = m_0^2 + c^2 n_0^2$. It follows that,

$$\begin{aligned} W_{\mathbb{K}_c}(\lambda) &= 2 \left(L_c(\lambda) - \lfloor \sqrt{\lambda} \rfloor - \lfloor \frac{\sqrt{\lambda}}{c} \rfloor - 1 \right) + \left(\lfloor \frac{\sqrt{\lambda}}{c} \rfloor + 1 \right) + 2 \lfloor \frac{\sqrt{\lambda}}{2} \rfloor, \\ &= 2 L_c(\lambda) - 2 \lfloor \sqrt{\lambda} \rfloor - \lfloor \frac{\sqrt{\lambda}}{c} \rfloor + 2 \lfloor \frac{\sqrt{\lambda}}{2} \rfloor - 1, \end{aligned}$$

where $\lfloor x \rfloor$ denotes the integer part of $x \geq 0$. It follows that, for all $\lambda \geq 0$,

$$W_{\mathbb{K}_c}(\lambda) \geq \begin{cases} \frac{\pi}{2} \lambda - 2\sqrt{\lambda} - 3 & \text{when } c = 1, \\ \frac{\pi}{4} \lambda - \frac{3}{2}\sqrt{\lambda} - 2 & \text{when } c = 2. \end{cases}$$

Inequalities (4.2) and (4.4) follow from the previous inequalities. \square

Table 4.1 displays the eigenvalues of \mathbb{K}_c less than or equal to 25 ($c = 1$), resp. 47 ($c = 2$), the corresponding labeled eigenvalues, and the ratio $\frac{\lambda_{k_{\min}}(\mathbb{K}_c)}{k_{\min}}$ which should be larger than or equal to

$$\begin{aligned} \frac{j_{0,1}^2}{2\pi} &\approx 0.92042 & \text{if } \lambda_{k_{\min}}(\mathbb{K}_1) \text{ is Courant-sharp,} \\ \frac{j_{0,1}^2}{\pi} &\approx 1.84084 & \text{if } \lambda_{k_{\min}}(\mathbb{K}_2) \text{ is Courant-sharp.} \end{aligned}$$

Since the eigenvalues $\lambda_1(\mathbb{K}_c), \lambda_2(\mathbb{K}_c)$ are Courant-sharp, we conclude from Table 4.1 that Theorem 1.1 is proved in the case $c = 2$. To finish the proof of the theorem

\mathbb{K}_1	$\lambda_{k_{\min}}(\mathbb{K}_1)$	$\lambda_{k_{\max}}(\mathbb{K}_1)$	$\frac{\lambda_{k_{\min}}(\mathbb{K}_1)}{k_{\min}}$	\mathbb{K}_2	$\lambda_{k_{\min}}(\mathbb{K}_2)$	$\lambda_{k_{\max}}(\mathbb{K}_2)$	$\frac{\lambda_{k_{\min}}(\mathbb{K}_2)}{k_{\min}}$
1	λ_2	λ_2	—	4	λ_2	λ_4	—
2	λ_3	λ_4	—	5	λ_5	λ_6	1
4	λ_5	λ_7	—	8	λ_7	λ_8	1.1429
5	λ_8	λ_{11}	0.625	13	λ_9	λ_{10}	1.4444
8	λ_{12}	λ_{13}	0.6666	16	λ_{11}	λ_{13}	1.4545
9	λ_{14}	λ_{14}	0.6428	17	λ_{14}	λ_{15}	1.2143
10	λ_{15}	λ_{18}	0.6666	20	λ_{16}	λ_{19}	1.2500
13	λ_{19}	λ_{22}	0.6842	25	λ_{20}	λ_{21}	1.2500
16	λ_{23}	λ_{25}	0.6956	29	λ_{22}	λ_{23}	1.3182
17	λ_{26}	λ_{29}	0.6538	32	λ_{24}	λ_{25}	1.3333
18	λ_{30}	λ_{31}	0.6	36	λ_{26}	λ_{28}	1.3846
20	λ_{32}	λ_{35}	0.625	37	λ_{29}	λ_{30}	1.2759
25	λ_{36}	λ_{40}	0.6944	40	λ_{31}	λ_{34}	1.2903
				41	λ_{35}	λ_{36}	1.1714
				45	λ_{37}	λ_{38}	1.2162
				52	λ_{39}	λ_{42}	1.3333

TABLE 4.1. The first eigenvalues of \mathbb{K}_1 (left) and \mathbb{K}_2 (right)

when $c = 1$, it remains to investigate $\lambda_3(\mathbb{K}_1)$ and $\lambda_5(\mathbb{K}_1)$. This is done in the next section.

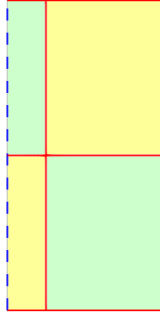
5. THE EIGENVALUES $\lambda_3(\mathbb{K}_1)$ AND $\lambda_5(\mathbb{K}_1)$ ARE NOT COURANT-SHARP

In order to finish the proof of Theorem 1.1 for the Klein bottle \mathbb{K}_1 , we investigate the eigenvalues $\lambda_3(\mathbb{K}_1) = \hat{\lambda}(1, 1)$ and $\lambda_5(\mathbb{K}_2) = \hat{\lambda}(2, 0) = \hat{\lambda}(0, 2)$.

5.1. The eigenvalue $\lambda_3(\mathbb{K}_1)$ is not Courant-sharp. A general eigenfunction associated with λ_3 has the form $(A \cos(x) + B \sin(x)) \sin(y)$. It is sufficient to look at eigenfunctions of the form $\sin(x - \alpha) \sin(y)$. These eigenfunctions have exactly two nodal domains in \mathbb{K}_1 . It follows that λ_3 is not Courant-sharp, see Figure 5.1.

5.2. The eigenvalue $\lambda_5(\mathbb{K}_1)$ is not Courant-sharp. A general eigenfunction associated with λ_5 has the form $A \cos(2x) + B \sin(2x) + C \cos(2y)$. Up to multiplication by a scalar, it suffices to consider the family $\cos \theta \cos(2x - \alpha) + \sin \theta \cos(2y)$, with $\theta \in [0, \pi)$ and $\alpha \in [0, \pi)$. Choosing the fundamental domain appropriately, we can assume that $\alpha = 0$. Changing y to $y + \frac{\pi}{2}$, we see that it suffices to consider $\theta \in [0, \frac{\pi}{2}]$. The nodal sets are known explicitly when $\theta = 0$ or $\frac{\pi}{2}$. We now consider the family

$$(5.1) \quad \phi_\theta(x, y) = \cos \theta \cos(2x) + \sin \theta \cos(2y), \quad \theta \in (0, \frac{\pi}{2}).$$

FIGURE 5.1. Nodal domains of $\sin(x - \alpha) \sin(y)$

The critical zeros (points at which both the function and its differential vanish) satisfy the system,

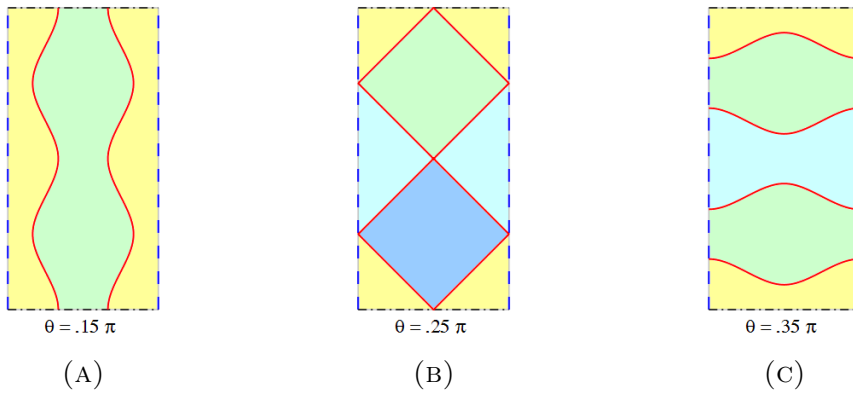
$$(5.2) \quad \begin{cases} \cos \theta \cos(2x) + \sin \theta \cos(2y) = 0, \\ \sin(2x) = 0, \\ \sin(2y) = 0. \end{cases}$$

It follows that critical zeros only occur for $\theta = \frac{\pi}{4}$, so that the nodal set of ϕ_θ is a regular curve when $\theta \neq \frac{\pi}{4}$. The nodal set of the eigenfunction

$$(5.3) \quad \phi_{\frac{\pi}{4}}(x, y) = \cos(2x) + \cos(2y) = 2 \cos(x + y) \cos(x - y)$$

is explicit, see Figure 5.2 (B). An analysis à la Stern, see [5, Section 5] or [2], shows that the nodal sets of ϕ_θ for $0 < \theta < \frac{\pi}{4}$ and $\frac{\pi}{4} < \theta < \frac{\pi}{2}$ are given by Figures 5.2 (A) and (C) respectively. More precisely, we first note that the common zeros to ϕ_0 and $\phi_{\frac{\pi}{2}}$ are common zeros to all ϕ_θ . Since $\theta \in (0, \frac{\pi}{2})$, and hence $\sin \theta \cos \theta > 0$, it follows that, except for the common zeros, the nodal set of ϕ_θ is contained in the set $\cos(2x) \cos(2y) < 0$. Finally, depending on the sign of $\frac{\pi}{4} - \theta$, we can use the nodal set of $\cos(2x)$ or the nodal set of $\cos(2y)$ as barrier to obtain the behaviour described in the figures.

It follows that an eigenfunction associated with λ_5 has at most 4 nodal domains in \mathbb{K}_1 , and hence that $\lambda_5(\mathbb{K}_1)$ is not Courant-sharp.

FIGURE 5.2. Nodal sets for $\cos \theta \cos(2x) + \sin \theta \cos(2y)$

This completes the proof of Theorem 1.1 in the case $c = 1$.

6. COURANT-SHARP DIRICHLET EIGENVALUES OF FLAT CYLINDERS

6.1. **Preliminaries.** Given $r > 0$, let $\mathcal{C}_r = (0, \pi) \times \mathbb{S}_r^1$ denote the cylinder with radius r and length π , equipped with the flat product metric. We consider the Dirichlet boundary condition on $\partial\mathcal{C}_r = \{0, \pi\} \times \mathbb{S}_r^1$. A complete family of Dirichlet eigenfunctions is given by,

$$(6.1) \quad \begin{cases} \sin(mx) \cos\left(\frac{ny}{r}\right), & \text{for } m \in \mathbb{N}^\bullet, n \in \mathbb{N}, \\ \sin(mx) \sin\left(\frac{ny}{r}\right), & \text{for } m \in \mathbb{N}^\bullet, n \in \mathbb{N}^\bullet, \end{cases}$$

with associated eigenvalues the numbers $\hat{\lambda}_r(m, n) = m^2 + \frac{n^2}{r^2}$.

For $m \in \mathbb{N}^\bullet$, the point $(m, 0)$ contributes one eigenfunction $\sin(mx)$; for $m, n \in \mathbb{N}^\bullet$, the point (m, n) contributes two functions, $\sin(mx) \cos\left(\frac{ny}{r}\right)$ and $\sin(mx) \sin\left(\frac{ny}{r}\right)$. Multiplicities of the eigenvalues also occur when the equation $m^2 + \frac{n^2}{r^2} = m_0^2 + \frac{n_0^2}{r^2}$ has multiple solutions within the above range.

As usual, we arrange the Dirichlet eigenvalues of \mathcal{C}_r in non-decreasing order, starting from the label 1, multiplicities accounted for,

$$(6.2) \quad 0 < \lambda_1(\mathcal{C}_r) < \lambda_2(\mathcal{C}_r) \leq \lambda_3(\mathcal{C}_r) \leq \dots$$

The purpose of this section is to prove Theorem 1.2, i.e. to determine the Courant-sharp eigenvalues in two specific cases $r \in \left\{\frac{1}{2}, 1\right\}$ whose eigenvalues have higher multiplicities. The corresponding cylinder are the orientation covers of the Möbius band we studied in [5]. As in [12, 13], we could also consider other values of r , in particular r close to zero or r very large.

Remark 6.1. *As pointed out in the introduction, $\lambda_1(\mathcal{C}_r)$ and $\lambda_2(\mathcal{C}_r)$ are always Courant-sharp, and any Courant-sharp eigenvalue $\lambda_k(\mathcal{C}_r)$ satisfies the inequality $\lambda_k(\mathcal{C}_r) > \lambda_{k-1}(\mathcal{C}_r)$.*

Following Pleijel’s method, Section 2, we introduce Weyl’s counting function,

$$(6.3) \quad W_{\mathcal{C}_r}(\lambda) = \#\{j \mid \lambda_j(\mathcal{C}_r) < \lambda\}.$$

According to Weyl’s law,

$$(6.4) \quad W_{\mathcal{C}_r}(\lambda) = \frac{|\mathcal{C}_r|}{4\pi} \lambda + O(\sqrt{\lambda}) = \frac{r\pi}{2} \lambda + O(\sqrt{\lambda}),$$

where $|\mathcal{C}_r|$ denotes the area of the cylinder.

6.2. **Courant-sharp eigenvalues of \mathcal{C}_r , $r \in \left\{\frac{1}{2}, 1\right\}$.** According to [17, § 6], a domain $\omega \subset \mathcal{C}_r$, with area $|\omega| \leq 4\pi r^2$ satisfies a Euclidean isoperimetric inequality, and hence its least Dirichlet eigenvalue satisfies

$$(6.5) \quad \delta_1(\omega) \geq \frac{\pi j_{0,1}^2}{|\omega|}.$$

Lemma 6.2. *Let $r \in \left\{\frac{1}{2}, 1\right\}$. Let $\lambda_k(\mathcal{C}_r)$ be a Courant-sharp eigenvalue of \mathcal{C}_r , with $k \geq \frac{\pi}{2r}$. The following inequalities hold.*

$$(6.6) \quad \frac{\lambda_k(\mathcal{C}_r)}{k} \geq \frac{j_{0,1}^2}{2\pi r} \geq \begin{cases} \frac{j_{0,1}^2}{\pi} \approx 1.840844 & \text{if } r = \frac{1}{2}, \\ \frac{j_{0,1}^2}{2\pi} \approx 0.920422 & \text{if } r = 1, \end{cases}$$

$$(6.7) \quad k - 1 = W_{\mathcal{C}_r}(\lambda_k(\mathcal{C}_r)) \geq \frac{\pi r}{2} \lambda_k(\mathcal{C}_r) - (2r + 1) \sqrt{\lambda_k(\mathcal{C}_r)} - 2.$$

$$(6.8) \quad \text{For } r = \frac{1}{2}, \quad \frac{\pi}{4} \left(1 - \frac{4}{j_{0,1}^2}\right) \lambda_k(\mathcal{C}_{\frac{1}{2}}) - 2\sqrt{\lambda_k(\mathcal{C}_{\frac{1}{2}})} - 1 \leq 0,$$

$$(6.9) \quad \lambda_k(\mathcal{C}_{\frac{1}{2}}) \leq 76.25,$$

$$(6.10) \quad \text{For } r = 1, \quad \frac{\pi}{4} \left(1 - \frac{4}{j_{0,1}^2}\right) \lambda_k(\mathcal{C}_1) - \frac{3}{2} \sqrt{\lambda_k(\mathcal{C}_1)} - \frac{1}{2} \leq 0,$$

$$(6.11) \quad \lambda_k(\mathcal{C}_1) \leq 42.40$$

Proof. The arguments are similar to those used in the proof of Lemma 4.1. Inequality (6.6) follows from (6.5). The inequality in (6.7) follows from the description of the Dirichlet spectrum of \mathcal{C}_r . The other inequalities follow from (6.7) and (6.6). \square

The following tables give the eigenvalues, the corresponding range of labels, and the ratios $\frac{\lambda_{k_{min}}}{k_{min}}$ for $\mathcal{C}_{\frac{1}{2}}$ (left) and \mathcal{C}_1 (right). For a Courant-sharp eigenvalue, this ratio should be greater than 1.840844 for $\mathcal{C}_{\frac{1}{2}}$, and greater than 0.920422 for \mathcal{C}_1 .

According to Remark 6.1, the eigenvalues λ_1 and λ_2 are always Courant-sharp.

The table for $\mathcal{C}_{\frac{1}{2}}$ (left) shows that it only remains to analyze the eigenvalue $5 = \lambda_3(\mathcal{C}_{\frac{1}{2}}) = \lambda_4(\mathcal{C}_{\frac{1}{2}})$. The associated eigenspace is generated by the eigenfunctions $\sin(x) \cos(2y)$ and $\sin(x) \sin(2y)$. Functions in this eigenspace have 2 nodal domains. This finishes the proof of Theorem 1.2 in the case $r = \frac{1}{2}$.

The table for \mathcal{C}_1 (right) shows that it only remains to analyze two eigenvalues 4 and 5. The eigenvalue $4 = \lambda_4(\mathcal{C}_1)$ is simple, with associated eigenfunction $\sin(2x)$ which has two nodal domains. The eigenvalue $5 = \lambda_5(\mathcal{C}_1) = \dots = \lambda_8(\mathcal{C}_1)$ has multiplicity 4. The corresponding eigenspace is generated by the eigenfunctions $\sin(2x) \cos(x)$, $\sin(2x) \sin(x)$ and $\sin(x) \cos(2x)$, $\sin(x) \sin(2y)$ which, according to [5], turn out to span the second eigenspace $\mathcal{E}(\lambda_2(M_1))$ of the square Möbius strip, whose orientation cover is \mathcal{C}_1 . Since the eigenfunctions in $\mathcal{E}(\lambda_2(M_1))$ have two nodal domains, it follows that the eigenfunctions in the eigenspace $\mathcal{E}(\lambda_5(\mathcal{C}_1))$ have at most four nodal domains. This finishes the proof of Theorem 1.2 in the case $r = 1$. Note that one can give a precise analysis of nodal patterns in the eigenspace $\mathcal{E}(\lambda_5(\mathcal{C}_1))$ by using the same arguments as in [5, Section 4].

$\lambda(\mathcal{C}_{1/2})$	k_{min}	k_{max}	$\frac{\lambda_{k_{min}}(\mathcal{C}_{1/2})}{k_{min}}$
1	1	1	–
4	2	2	–
5	3	4	–
8	5	6	1,60000
9	7	7	1,28571
13	8	9	1,62500
16	10	10	1,60000
17	11	12	1,54545
20	13	16	1,53846
25	17	19	1,47059
29	20	21	1,45000
32	22	23	1,45455
36	24	24	1,50000
37	25	26	1,48000
40	27	30	1,48148
41	31	32	1,32258
45	33	34	1,36364
49	35	35	1,40000
52	36	39	1,44444
53	40	41	1,32500
61	42	43	1,45238
64	44	44	1,45455
65	45	48	1,44444
68	49	52	1,38776
72	53	54	1,35849
73	55	56	1,32727
80	57	60	1,40351

$\lambda(\mathcal{C}_1)$	k_{min}	k_{max}	$\frac{\lambda_{k_{min}}(\mathcal{C}_1)}{k_{min}}$
1	1	1	–
2	2	3	1,00000
4	4	4	1,00000
5	5	8	1,00000
8	9	10	0,88889
9	11	11	0,81818
10	12	15	0,83333
13	16	19	0,81250
16	20	20	0,80000
17	21	24	0,80952
18	25	26	0,72000
20	27	30	0,74074
25	31	35	0,80645
26	36	39	0,72222
29	40	43	0,72500
32	44	45	0,72727
34	46	49	0,73913
36	50	50	0,72000
37	51	54	0,72549
40	55	58	0,72727
41	59	62	0,69492
45	63	66	0,71429
49	67	67	0,73134
50	68	73	0,73529

TABLE 6.1. The first eigenvalues of $\mathcal{C}_{\frac{1}{2}}$ (left) and \mathcal{C}_1 (right)

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