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# Formation of SiP<sub>2</sub> nanocrystals embedded in SiO<sub>2</sub> from phosphorus rich SiO<sub>1.5</sub> thin films

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## Abstract

We investigate the structural, vibrational and optical properties of phosphorus-rich SiO<sub>1.5</sub> thin films annealed at 1100°C. For phosphorus (P) contents larger than 3 at. %, high resolution transmission electron microscopy characterizations reveal the presence of both spherical shaped SiP<sub>2</sub> nanoparticles crystallizing in an orthorhombic structure and bumps in epitaxy with the underlying Si substrate. Energy dispersive spectroscopy measurements confirm the SiP<sub>2</sub> stoichiometry. Moreover, electron energy loss spectroscopy characterizations allow us to determine the exact location of P and Si atoms. Apart from SiP<sub>2</sub> nanoparticles, P atoms are found to be located in the bumps and in the Si substrate to a level of 1 at. %, which is explained by P diffusion during annealing. The vibrational properties determined by Raman spectroscopy are found to be in excellent agreement with density functional theory calculations of the vibration modes for the SiP<sub>2</sub> alloy. Finally, the quenching of the photoluminescence with the increasing P content is explained on the basis of the structural data.

Si nanocrystals (Si-NCs) have been extensively studied during the last two decades due to their original electronic and optical properties. Additional functionalities can be expected if the nanocrystals can be effectively doped. This is, however, a challenging issue which is related with several physical and technological barriers. In particular, the increasing formation energy for substitutional dopants with the decreasing size can lead to a “self-purification” mechanism for small nanocrystals [1]. Nevertheless, both P and B doped Si nanocrystals embedded in a SiO<sub>2</sub> matrix have been successfully obtained by high temperature phase separation of Si rich SiO<sub>x</sub> thin films deposited by e-beam evaporation [2], plasma enhanced chemical vapor deposition [3], ion implantation [4] or by using a combination of the two last techniques [5]. Another important question concerned the localization of the dopants. König et al. [6] used atom probe tomography to demonstrate that in the case of P-doped Si nanocrystals embedded in SiO<sub>2</sub> prepared by plasma enhanced chemical vapor deposition, only 15 % of the P atoms are found within the nanocrystals while 30 % are localized at the interface with the SiO<sub>2</sub> matrix and 55% in the silica matrix. Alternatively, ion beam synthesis by silicon/dopant co-implantation appeared to be a viable route to obtain a high concentration of P atoms inside Si nanocrystals [4]. Very recently, it has been shown both theoretically [7] and experimentally [8-11] that either phosphorus or boron doped Si-NCS can be of interest for tunable localized surface plasmon resonances (LSPR). When the dopant concentration is very high, the solid solubility limit could be exceeded and alloying could become possible. Indeed, the Si-P equilibrium phase diagram established so far predicts that two stable compounds may form (SiP and SiP<sub>2</sub>) depending on the relative amount of phosphorus [12-14]. Both orthorhombic and pyrite type bulk SiP<sub>2</sub> alloys were first synthesized by Wadsten in 1967 [15,16]. Donohue et al. [17] refined the structure of pyrite type SiP<sub>2</sub> by using X-ray powder diffraction. Based on galvanomagnetic and reflectance measurements, these authors claim that pyrite-type SiP<sub>2</sub> could be a semimetal. Farberovich et al. [18] came to the same

conclusion by calculating the band structure using the orthogonalized plane wave method. Vogt et al. [19] characterized the vibrational properties of pyrite type  $\text{SiP}_2$  and obtained the frequencies of all Raman active phonons at room temperature. Meier et al. [20] performed *ab initio* calculations of the fundamental vibration frequencies of pyrite type  $\text{SiP}_2$ . A good agreement was found with both Raman and infrared studies. Bachhuber et al. [21] further used first principle calculations to investigate the structure, the bonding and the vibrational properties of pyrite type  $\text{SiP}_2$ . By using the vapor transport method, SpringThorpe [22] succeeded to prepare single crystals of orthorhombic  $\text{SiP}_2$ . This compound was shown to be a p-type semiconductor with an optical bandgap of about 1.89 eV. By using the Sn flux method with Gd as a mineralizer in a fused silica tube, Zhang et al. [23] identified by means of X-ray diffraction another orthorhombic phase of  $\text{SiP}_2$  crystallizing in the *Pnma* group. Diffuse reflectance measurements as well as preliminary density functional theory (DFT) calculations suggest a semiconducting behavior with a bandgap close to 1.4 eV. Moreover,  $\text{SiP}_2$  microparticles prepared by ball milling [24] but also nanocomposites consisting of  $\text{SiP}_2$  and amorphous carbon as well as carbon nanotube wrapped  $\text{SiP}_2$  [25] emerged as promising anode materials for high performance Li- and Na-ion batteries [26,27].

Very recently,  $\text{Si}_x\text{P}_y$  alloys ( $x, y = 1, 2$ ) have gained a renewed interest since they exhibit a lamellar structure. This property might be of interest for engineering novel Si-based 2D materials, which could allow filling the gap between silicene and phosphorene. Density functional theory (DFT) calculations have shown that 2D- $\text{SiP}_x$  ( $x = 1, 2$ ) are direct bandgap semiconductors [28-31], which might be of particular interest for future optoelectronic applications. Moreover, according to its band structure, 2D- $\text{SiP}_2$  appears to be promising for photocatalytic water splitting [31].

Till now, there are no reports available in the literature concerning the growth of orthorhombic  $\text{SiP}_2$  either in thin films or nanocrystals. In this paper, we demonstrate the

formation of orthorhombic SiP<sub>2</sub> nanocrystals embedded in silicon dioxide thin films obtained by annealing of phosphorus rich SiO<sub>1.5</sub> thin films at 1100°C. We provide a careful structural and vibrational study of SiP<sub>2</sub> nanocrystals. The experimental results are found to be in excellent agreement with density functional theory calculations of the vibration modes for the SiP<sub>2</sub> alloy. The optical properties of P-rich SiO<sub>1.5</sub> thin films are further discussed in relation with the obtained structural results.

Phosphorus-rich SiO<sub>1.5</sub> thin films were prepared by co-evaporation on Si(001) substrates held at room temperature of SiO and SiO<sub>2</sub> from e-beam guns and P from a GaP decomposition source [32]. The SiO and SiO<sub>2</sub> fluxes were controlled by a quartz microbalance. The P content, which is defined by  $[P]/([Si]+[O]+[P])$ , was obtained by comparing the relative intensities of the P 2*p*, Si 2*p* and O 1*s* core levels measured using X-ray photoelectron spectroscopy (XPS). The obtained values were cross-checked by secondary ion mass spectroscopy (SIMS) measurements. The total film thickness was about 200 nm for all samples considered in this study. The samples were post-grown annealed in a rapid thermal annealing furnace at 1100°C during 5 minutes under N<sub>2</sub> atmosphere. The structural properties of the samples were investigated by high resolution transmission electron microscopy (HRTEM), scanning transmission electron microscopy (STEM), energy dispersive spectroscopy (EDS) and electron energy loss spectroscopy (EELS) using a JEOL ARM 200 F probe corrected with cold field emission gun. EELS spectrum images were recorded with a Gatan GIF quantum ER filter. To avoid any phase precipitation due to ion or electron beam damages, the samples were mechanically thinned up to the electron transparency by the tripod method and the TEM was operated at 80 kV with a probe current lower than 70 pA for EDS and EELS spectrum images. For EELS investigations, zero-loss and core-loss spectrum images were recorded simultaneously to allow advanced data post processing (correction of energy drift, thickness correction). To separate overlapping Si edges

and P edges and chemical states of Si and extract quantitative data, multiple linear least square fitting (MLS) was used using reference spectra recorded on pure Si and on pure thermal silica grown on a Si wafer. The reference spectra were recorded with the same STEM and spectrometer parameters than that of the sample investigated. Raman scattering measurements were performed at RT by using a 532 nm laser diode as excitation source. The scattered light is analyzed by a monochromator equipped with a 1800 lines/mm grating and a CCD detector. The steady state photoluminescence was excited with a 325 nm He-Cd laser and subsequently analyzed by a monochromator equipped with a 600 grooves/mm grating and measured by a photomultiplier tube operating in the visible-IR range. The spectra were corrected from the response of the detector. Density functional theory (DFT) calculations were performed with the CRYSTAL14 [33] periodic computation code using the B3LYP functional [33, 34] and a triple  $\zeta$  localized (atom centered) basis set [35]. Both the atomic positions and the cell parameters of SiP<sub>2</sub> were optimized. The harmonic vibration frequencies were calculated at the  $\Gamma$  point and the corresponding intensities were evaluated through a coupled perturbed Kohn-Sham (CPKS) approach.

Figure 1 (a) shows a HRTEM image of a SiO<sub>1.5</sub> thin film containing 10 at. % of phosphorus after annealing for 5 minutes at 1100°C. Two kinds of features can be distinguished. Firstly, the interface with the substrate is not sharp but one can identify bumps which are in epitaxy with the substrate. A high resolution STEM-high angle annular dark field (HAADF) image of the interface between a typical bump and the Si substrate is shown in Fig. 1(b). The zone axis of the Si substrate is  $\langle 110 \rangle$ . The interface is indicated by a red dash-dotted line. One can clearly identify the dumbbell of Si  $\langle 110 \rangle$  both in the bump and in the substrate showing that the bumps are in epitaxy with the underlying Si substrate. Secondly, in the thin film, one can distinguish large, spherical shaped nanoparticles having typical sizes between 30 and 40 nm. One can clearly distinguish the atomic planes demonstrating almost perfect crystallinity. We

note, however, that these nanocrystals often contain a high density of twins and stacking faults. Similar particles can also be identified in samples containing 3 and 6 at. % of phosphorus (not shown here). Their size increases from about 12 nm in the film containing 3 at. % of phosphorus to 31 nm in the film containing about 10 at. % of phosphorus. Figure 1(c) displays a Fourier transform of the spherical nanocrystal identified in Figure 1(a). A simulation of the electron diffraction pattern of  $\text{SiP}_2$  with *Pbam* orthorhombic structure (JCPDS file 00-044-1124) was compared to the Fourier transform. A good agreement is found for the zone axis [014]. This is the first experimental observation of  $\text{SiP}_2$  nanocrystals crystallizing in an orthorhombic structure. In order to characterize the chemical composition of both types of nanostructures, we have performed EDS spectrum images (SI). Figure 2 shows a HAADF image (Fig. 2(a)) together with the chemical semi-quantitative maps for Si (Fig. 2(b)), O (Fig. 2(c)) and P (Fig. 2(d)) extracted from EDS-SI. The EDS spectra extracted from EDS-SI in three different regions (delimited in Figs. 2(b), 2(c) and 2(d)) are shown in Figure 2(e). Three different peaks can be identified at 0.5, 1.7 and 2.0 keV corresponding to the  $K\alpha$  emission lines of oxygen, of silicon and of phosphorus, respectively. In the substrate, one can find about 1 at. % of phosphorus indicating that P diffusion has occurred during annealing at 1100°C. Moreover, the EDS spectrum extracted from the bump area shows that the bumps contain also about 1 at. % of phosphorus. The oxygen related peak at 0.5 keV originates from  $\text{SiO}_2$  which is present both on top and below the bump. Finally, for the large nanocrystals, the Si-related peak at 1.7 keV originates from both Si atoms located in  $\text{SiO}_2$  and in  $\text{SiP}_2$ . Indeed, under the electron probe, X-rays are emitted from the whole thickness of the sample. Consequently, Si atoms in  $\text{SiO}_2$  present both on top and below the nanocrystals will also emit. To obtain the real Si content in the nanoparticles, the contribution resulting from Si atoms present in  $\text{SiO}_2$  have to be subtracted. This can be done easily by assuming that all oxygen atoms originate from the stoichiometric  $\text{SiO}_2$  compound. A Si content of about 32.8

at. % and a phosphorus content of 67.2 at. % was then obtained, which is compatible with the presence of a  $\text{SiP}_2$  compound. However EDS measurements do not allow to distinguish clearly between silicon atoms in  $\text{SiP}_2$  nanoparticles, in silicon bumps and in the silicon dioxide matrix. Moreover, the EDS signal is known to be not fully localized. In order to get a more precise chemical mapping at the nanoscale and to be able to distinguish between  $\text{Si}^{(0)}$  and  $\text{Si}^{(4+)}$ , we have performed further STEM-EELS spectrum images. Figure 3(a) shows a STEM-HAADF image of an annealed film containing 10 at. % of phosphorus annealed at  $1100^\circ\text{C}$ . One can clearly distinguish both the nanoparticles and the bumps. Figure 3(b) shows an EELS spectrum integrated over the whole area shown in Fig. 3(a) (black solid line). To get a chemical map, we have performed a minimum least square fitting of this spectrum by using reference spectra for  $\text{Si}^{(0)}$  and  $\text{Si}^{(4+)}$  (from a lamella of thermal  $\text{SiO}_2$  grown by oxidation of  $\text{Si}(100)$ ) measured in the same experimental conditions (thickness, accelerating voltage, probe current, convergence and collection angle, etc.). Moreover, the fitting leads to a residual spectrum (orange solid line in figure 3(b) having an edge at 129 eV, which corresponds to phosphorus [36]. By considering the reference spectra for both  $\text{Si}^{(0)}$  and  $\text{Si}^{(4+)}$  and the residual spectrum for phosphorus, we were able to reconstruct the chemical map for each element with a new MLS fitting of each pixel of the SI. Figure 3(c)-(d) show the chemical maps for both  $\text{Si}^{(0)}$  and P. One can recognize that Si is present in the substrate, in the bumps and in the nanocrystals. P atoms are localized in the nanocrystals and at the interface between the substrate and the film, which originates from P diffusion during annealing. It is noteworthy that the maps for Si and P superimpose perfectly in the thin films indicating that both Si and P atoms are located in the nanocrystals. Full EELS chemical maps showing  $\text{Si}^{(0)}$ ,  $\text{Si}^{(4+)}$  and P as well as the superposition of maps is presented in the Supplementary Information (Fig. S1). Finally, one can discuss briefly the formation of the bumps at the interface with the substrate. In case of pure  $\text{SiO}_{1.5}$  films, they are not observed. Their size and height increases as the P



content in the films increases. During annealing at high temperatures, both P atoms and Si atoms in excess which do not participate in either  $\text{SiP}_2$  or  $\text{SiO}_2$  will diffuse towards the substrate. The diffusion will be even enhanced for larger P contents thus leading to formation of larger bumps. The Si atoms eventually accumulate into holes created in the oxide at the interface thus leading to the formation of bumps in epitaxy with the  $\text{Si}(001)$  substrate.

All the previous characterizations were done at the nanometer scale. To obtain a spectroscopic signature of the  $\text{SiP}_2$  nanocrystals, we have further performed Raman scattering measurements. Figure 4 (a) shows Raman spectra measured for  $\text{SiO}_{1.5}$  thin films with various P contents ranging from 0.3 at.% to 10 at. % and annealed at  $1100^\circ\text{C}$ . The films were deposited onto silica substrates in order to get rid of the huge signal originating from the Si-TO line at  $520.5\text{ cm}^{-1}$ . For P contents below 3 at. %, the spectra are essentially dominated by the signal originating from the silica substrate. However, for P contents of about 10 at. %, new and sharp peaks can be observed suggesting the formation of a new crystalline phase. According to all structural results from this work, these peaks could possibly be a signature of the  $\text{SiP}_2$  compound. To clarify the origin of these peaks, we have performed DFT calculations of the vibrational modes for the orthorhombic  $\text{SiP}_2$  structure. Figure 4 (b) shows the experimental spectrum (blue solid line) as well as the calculated Raman spectra for the orthorhombic  $\text{SiP}_2$  alloy (magenta solid line). It is obvious that a good agreement is obtained between the experimental and the calculated Raman spectrum thus further confirming that the observed nanocrystals consist indeed of  $\text{SiP}_2$ . Table I summarizes the position of both calculated and measured vibrational wavenumbers. Some differences appear nevertheless between the theoretical and the experimental values, which could be due either to the approximations used for the DFT calculations or to the fact that a non-negligible strain may be exerted by the silica matrix onto the  $\text{SiP}_2$  nanocrystals. The different vibrational modes are

also identified next to the calculated vibrational wavenumbers. Finally, our structural investigations help to shed new light on the evolution of the photoluminescence of  $\text{SiO}_{1.5}:\text{P}$  thin films with increasing P contents. Figure 5 displays the room temperature PL spectra measured for  $\text{SiO}_{1.5}$  thin films with various P contents ranging from 0.3 at.% to 10 at. %. For the lowest P content (i.e. 0.3 at.%), the PL signal is characterized by a rather broad band extending from about 700 nm to 1000 nm. In this particular sample,  $\text{SiP}_2$  nanocrystals could not be detected. The observed luminescence thus originates from radiative recombinations between charge carriers confined in pure Si nanocrystals. For the sample containing about 3 at. % of P, a strong decrease of the PL intensity is observed. Similar evolutions have been reported in the literature and are generally ascribed to the Auger effect [37] in Si-NCs containing P atoms which could then act as dopants. Spatially resolved EELS measurements performed on the same sample (see Fig.S1, Supplementary Informations) revealed the coexistence of both  $\text{SiP}_2$  nanocrystals having typical sizes of about 15 nm and Si nanocrystals with sizes ranging between 4 and 6 nm. Only few of them have sizes below the Bohr radius of the exciton in bulk Si and can contribute to the observed photoluminescence thus explaining the strong intensity decrease when the P content increases from 0.3 at. % to 3 at. % of P. Moreover, the PL peak shifts to larger wavelength thus suggesting an increase of the Si nanocrystal size, which is indeed observed. Finally, no PL could be detected for the sample having the highest P content (i.e. 10 at. %). In this case, all Si atoms are either involved in the  $\text{SiP}_2$  nanocrystals or in the bumps at the interface with the substrate.

To summarize, we provided clear evidence of the formation of  $\text{SiP}_2$  nanocrystals in P-rich  $\text{SiO}_{1.5}$  thin films annealed at 1100°C. The latter are found to crystallize in an orthorhombic structure. Elemental chemical mapping confirmed the  $\text{SiP}_2$  stoichiometry and allowed us to determine the exact location of both P and Si atoms. Apart from the  $\text{SiP}_2$  nanocrystals, P atoms are located in Si bumps, which are in epitaxy with the substrate, and at

the substrate surface to a level of about 1 at. %. The vibrational properties are found to be in excellent agreement with DFT calculations of the vibrational modes for the SiP<sub>2</sub> alloy. Finally, the photoluminescence properties of SiO<sub>1.5</sub> thin films with various P contents in the range 0.3-10 at. % can be understood by considering the structural properties. Our results shed new light into the complex evolution of ternary Si-O-P alloys during high temperature annealing.

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Figure captions:

Figure 1: (a) HRTEM image of a  $\text{SiO}_{1.5}$  thin film containing 10 at. % of P annealed at  $1100^\circ\text{C}$  for 5 minutes. (b) HRSTEM-HAADF image of the interface between a typical bump and the Si substrate. One can identify the [110] dumbbell of Si. The substrate is oriented along the  $\langle 110 \rangle$  zone axis for all images. The interface is indicated by a red dash-dotted line. The intensity difference between the substrate and the bump is due to the presence of silica above and below the bump. (c) Fourier transform of the nanocrystal delimited by a red dash-dotted line in (a).

Figure 2: (a) STEM-HAADF image of a  $\text{SiO}_{1.5}$  thin film containing 10 at. % of P annealed at  $1100^\circ\text{C}$  for 5 minutes. Quantitative chemical maps drawn from Si K (b), O K (c) and P K (d) signals. (e) EDS spectra extracted in SI from the different square zones indicated in (a)-(d).

Figure 3: (a) Spectrum image of a  $\text{SiO}_{1.5}$  thin film containing 10 at. % of P annealed at  $1100^\circ\text{C}$  for 5 minutes. (b) EELS spectrum of the sample (black solid line) together with reference spectra relative to  $\text{Si}^{4+}$  (cyan),  $\text{Si}^0$  (red) and P (orange). K edges maps for  $\text{Si}^{(0)}$  (c) and P (d) resulting from the MLS fitting

Figure 4: (a) Raman spectra of  $\text{SiO}_{1.5}$  thin films with different P contents annealed at  $1100^\circ\text{C}$  for 5 minutes. The films were grown onto silica substrates. (b) Experimental Raman spectrum of a  $\text{SiO}_{1.5}$  thin film containing 10 at. % of P (blue solid line) and calculated spectrum for  $\text{SiP}_2$  (magenta solid line)

Table I: Measured ( $\omega_{\text{EXP}}$ ) and calculated ( $\omega_{\text{DFT}}$ ) vibrational wavenumbers of orthorhombic  $\text{SiP}_2$ . The symmetry species of the vibrational modes are indicated in brackets.

Figure 5: Room temperature photoluminescence spectra of  $\text{SiO}_{1.5}$  thin films with different P contents ranging between 0.3 at. % and 10 at. %. The spectra relative to the samples containing 3 at. % and 10 at. % of P have been multiplied by 10.

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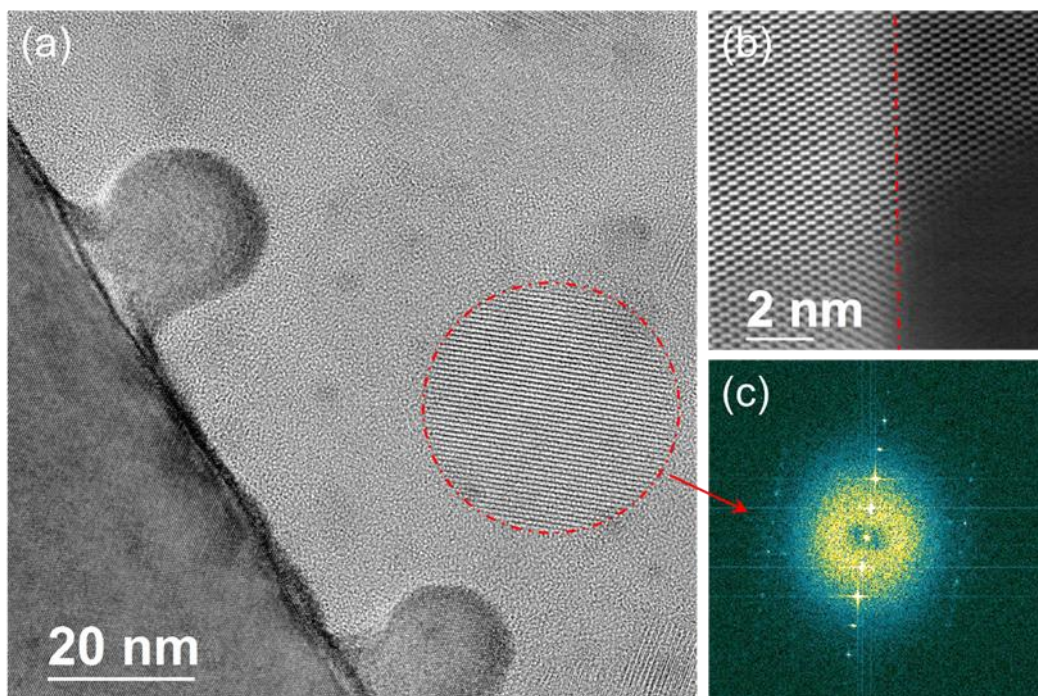


Fig.1

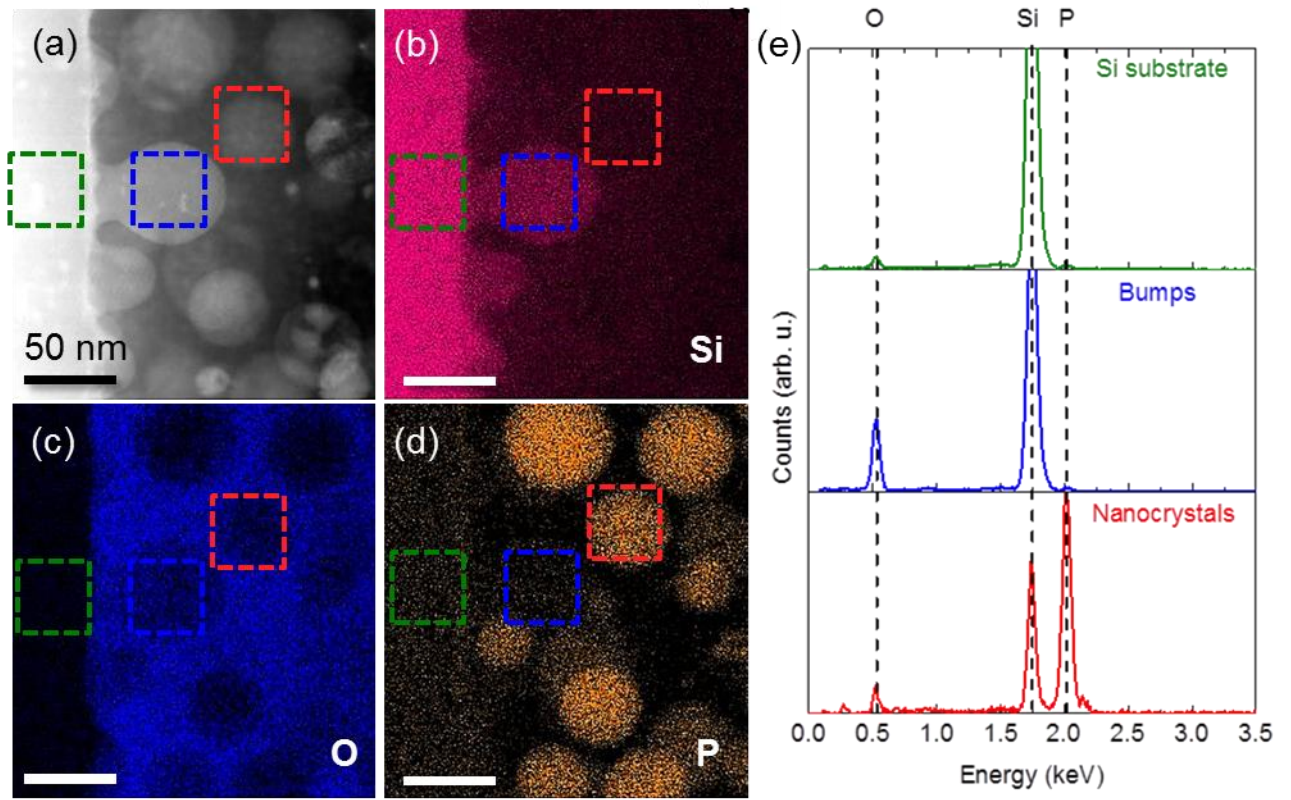


Fig.2

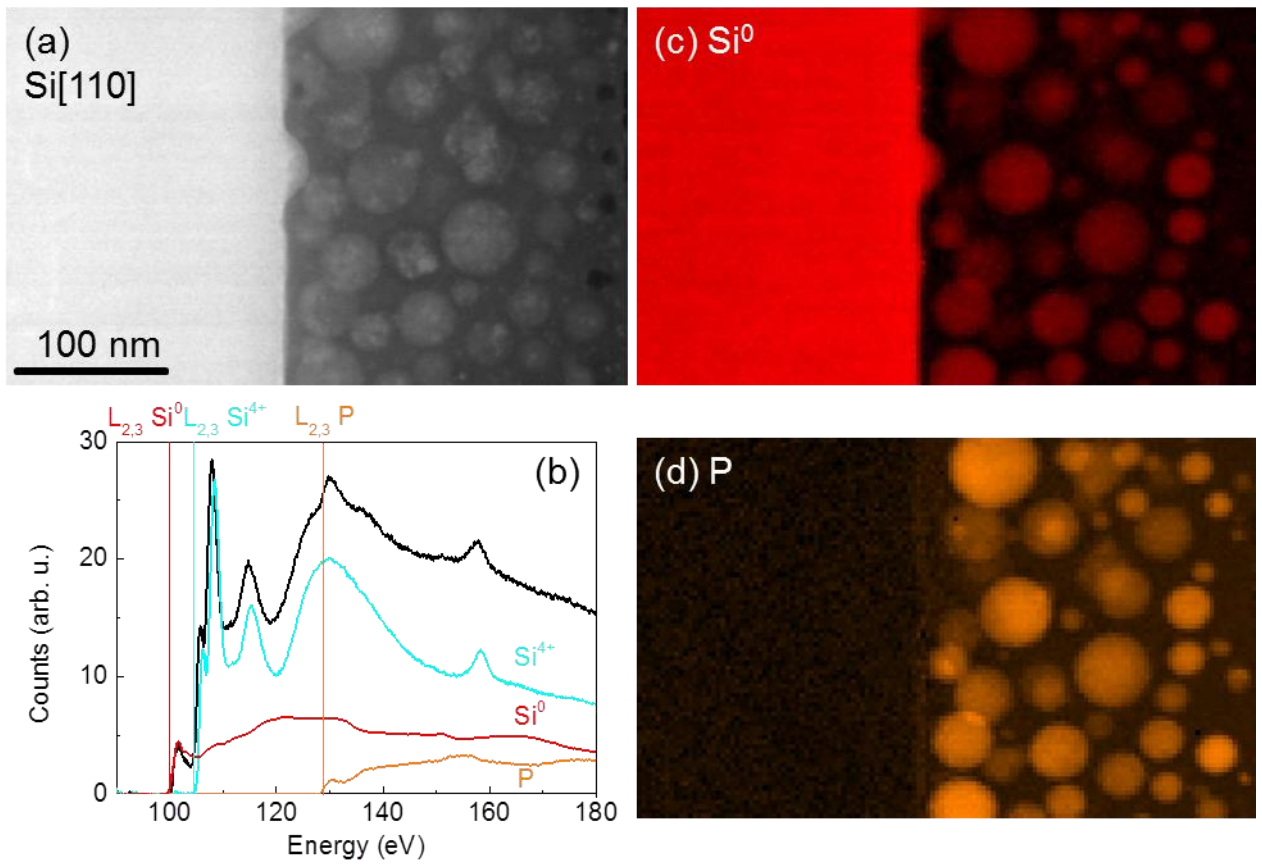


Fig.3

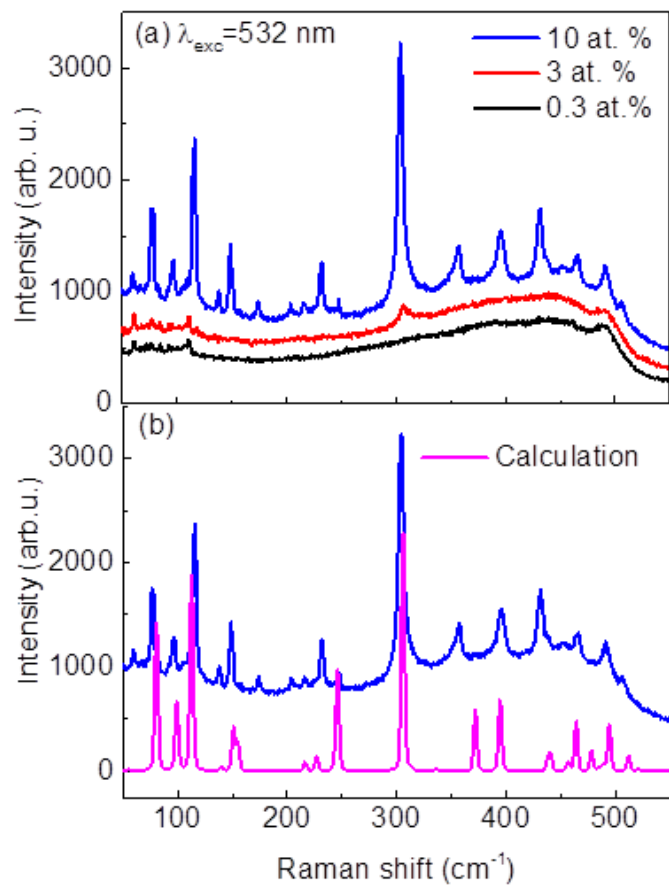


Fig.4

$\omega_{\text{DFT}} (\text{cm}^{-1})$	$\omega_{\text{EXP}} (\text{cm}^{-1})$	$\omega_{\text{DFT}} (\text{cm}^{-1})$	$\omega_{\text{EXP}} (\text{cm}^{-1})$
80 (Ag)	77	306 (Ag)	304
99 (B1g)	96	371 (Ag)	356
112 (Ag)	115	394 (B3g)	395
139 (B1g)	138	439 (Ag)	431
150 (B3g)	149	456 (B1g)	452
187 (B2g)	173	464 (Ag)	465
216 (B1g)	203	477 (Ag)	478
226 (Ag)	216	494 (Ag)	490
245 (B1g)	231	511 (B1g)	506

Table I

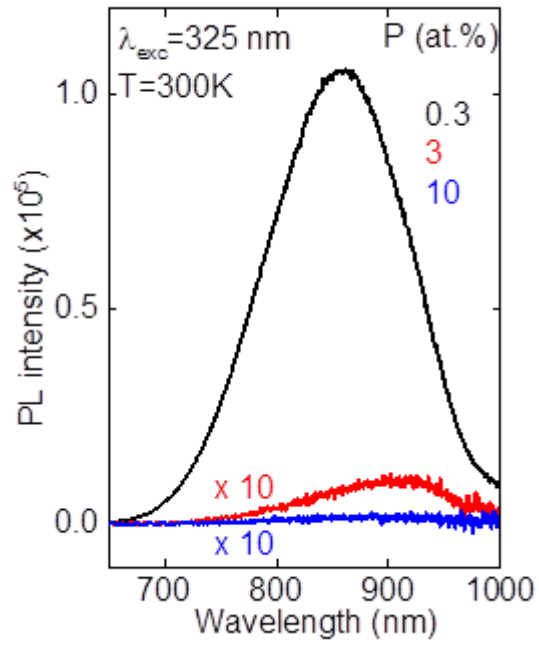


Fig. 5