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Submillimeter-wave spectroscopy of and interstellar search for thioacetaldehyde

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Abstract

We present a new study of the millimeter and submillimeter wave spectra of the thioacetaldehyde molecule, CH₃CHS, a sulfur-bearing analog of the ubiquitous interstellar molecule acetaldehyde (CH₃CHO). Here, we present a laboratory investigation aimed at determining accurate spectroscopic parameters for CH₃CHS to enable astronomical searches for this molecule using radio telescope arrays at millimeter and submillimeter wavelengths. New laboratory measurements have been carried out between 150 and 660 GHz using the spectrometer of PhLAM in Lille (France). Rotational transitions up to $J = 60$ were assigned in the ground, first and second excited torsional states of thioacetaldehyde and fit using the RAM Hamiltonian model. The final fit includes 62 parameters to give an overall weighted root-mean-square deviation of 0.9. On the basis of our spectroscopic results, CH₃CHS was searched for, but not detected, in data from the Atacama Large Millimeter/submillimeter Array (ALMA) between 84 GHz and 114 GHz toward the hot molecular core Sgr B2(N2). The non-detection implies that thioacetaldehyde is at least 90 times less abundant than acetaldehyde in this source. We also searched for but found no evidence of thioacetaldehyde in a number of prestellar and protostellar sources targeted by the Astrochemical Surveys at IRAM (ASAI) Large Program in nearby low-mass star forming regions.

Keywords: Microwave spectrum, torsional large amplitude motion, thioacetaldehyde, ISM

1. Introduction

More than 200 molecules have thus far been detected in the interstellar medium (ISM) [1]. Of these, ~70 and ~20 are oxygen-bearing and sulfur-bearing chemical compounds, respectively. Sulfur is the tenth most abundant element in the galaxy and to date, almost all sulfur compounds detected in the ISM have the corresponding oxygen derivatives also observed in the ISM. Existing problems with the systematic understanding of interstellar sulfur chemistry [2, 3, 4, 5] have motivated recent searches for new sulfur-bearing interstellar molecules. To that purpose a number of studies of sulfur-bearing analogs, not yet detected in the ISM, of well known oxygen-bearing interstellar molecules have been performed in recent years: thioacetic acid (CH₃COSH) [6], thioacetamide (CH₃CSNH₂) [7], dimethylsulfide ((CH₃)₂S) [8], O- and S-thiomethylformates (CH₃SCHO, CH₃OCHS) [9]. In this work, we present the results of a new spectroscopic study for the thioacetaldehyde molecule (CH₃CHS). With the corresponding oxygen derivative, acetaldehyde [10] as well as the unsubstituted derivative, thioformaldehyde, [11] being already detected in the ISM,

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thioacetaldehyde is a reasonable candidate for detection in the ISM. Except for the tentatively detected ethanethiol, [12] which contains 9 atoms, all the detected sulfur compounds have less than 7 atoms. If detected, thioacetaldehyde would therefore be one of the largest sulfur-bearing species in the ISM.

Several prior spectroscopic and synthetic studies have been devoted to thioacetaldehyde. As early as 1974, Kroto et al. [13] reported the photoelectron and microwave spectrum of this compound generated by vacuum flash pyrolysis (VFP) of the corresponding trimer: 1,3,5-trimethyl *s*-trithiane. The microwave spectrum showed evidence of hindered internal rotation [13] and it was detailed a few years later together with a study of eight isotopic variants of thioacetaldehyde [14]. The measurements were done up to 40 GHz and transitions up to $J = 35$ in the ground torsional state were assigned [14]. Information on dipole moment as well as methyl top internal rotation barrier heights was also obtained. Using the same approach, visible, UV spectra and phosphorescence excitation and emission spectra were observed [15, 16]. The VFP of alkyl thiosulfinates [17] or thiocyanohydrines [18] was also used to characterize or to trap the formed thioacetaldehyde. The photolysis of 2-methylthietane and 2,4-dimethylthietane in argon matrices allowed recording matrix infrared spectra [19]. Thioacetaldehyde was also prepared by matrix photolysis of ethanesulfonyl chloride or thiirane and by VFP of allylethylsulfide [20] and generated by photolysis of phenacylethylsulfide or by thermolysis of thiabicycloheptene [21].

In this paper, we present the results of our study of the torsion-rotational spectrum of thioacetaldehyde. The aim of the study is to obtain accurate spectroscopic parameters of thioacetaldehyde to support astronomical observations by radio telescope arrays, in particular at millimeter and submillimeter wavelengths. The new measurements were carried out between 150 GHz and 660 GHz and analyzed using the so-called rho-axis method (RAM). The analysis of the spectrum was performed up to $J = 60$ and up to the second excited torsional state, and a fit within experimental error was obtained. Predictions based on our final fit are used to search for CH₃CHS with the Atacama Large Millimeter/submillimeter Array (ALMA) toward the hot molecular core Sgr B2(N2) as well as in a number of observational datasets from the Astrochemical Surveys at IRAM (ASAI) Large Program.

The rest of the paper is organized as follows. In Section 2 we give in brief experimental details for the new measurements used in this work. In Section 3 we describe the assignment process and our fitting results. Section 4 describes our astronomical observations and the results of our search for thioacetaldehyde. Section 5 summarizes the results obtained in this study.

2. Experimental Section

2.1. Synthesis

Thioacetaldehyde was synthesized by flash vacuum thermolysis of ethyl-2-propynyl sulfide [22] at 750°C under 0.1 mbar. A stoichiometric amount of allene, a compound without permanent dipolar moment, was formed during the reaction. The gaseous flow was directly introduced in the cell of the spectrometer.

2.2. Lille - submillimeter spectra

The measurements in the frequency range under investigation were performed using the Lille spectrometer [23], equipped with a fast-scan mode [24]. The frequency ranges 150–330, and 400–660 GHz were covered with various active and passive frequency multipliers from VDI Inc. and an Agilent synthesizer (12.5–18.25 GHz) was used as the reference source of radiation. The absorption cell was a stainless-steel tube (6 cm diameter, 220 cm long). During measurements, the sample was at a pressure of about 10 Pa and at room temperature, the linewidth was limited by Doppler broadening. Estimated uncertainties for measured line frequencies are 30 kHz, 50 kHz, 100 kHz, and 200 kHz depending on the observed S/N and the frequency range.

3. Spectroscopic analysis and results

The so-called RAM torsion-rotation Hamiltonian used in this work is based on the works [25], [26], and [27]. We employed the RAM36 (rho-axis-method for 3 and 6-fold barriers) code that realizes the RAM

approach for the molecules with the C_{3v} top attached to a molecular frame of C_s or C_{2v} symmetry and having 3- or 6-fold barrier to internal rotation respectively [28, 29]. This code was successfully applied to the oxygen-bearing analog of thioacetaldehyde, acetaldehyde [30], which has comparable and rather large coupling between internal and overall rotations ($\rho = 0.33$ for CH_3CHO , $\rho = 0.26$ for CH_3CHS) as well as V_3 value ($V_3 = 407.6 \text{ cm}^{-1}$ for CH_3CHO , $V_3 = 556.6 \text{ cm}^{-1}$ for CH_3CHS). A general expression for the RAM Hamiltonian implemented in this code as well as further details of the used theoretical approach may be found in Refs. [28, 29]. The fits for thioacetaldehyde were performed with 21 torsional basis functions in the first diagonalization step and 9 torsional basis functions in the second diagonalization step.

We started our analysis from the results of [14], and as the first step we have refit the dataset of [14] with the RAM36 program [28, 29]. The initial predictions obtained from this refit were good enough for starting assignments in the 150–330 GHz range. Further assignments were performed in a usual bootstrap manner, with numerous cycles of refinement of the parameter set while the new data were gradually added. It should be noted that initially we aimed at studying only the ground torsional state with a number of low K series from the first excited torsional state included for stabilizing the torsional part of the Hamiltonian, since we were mainly interested in reliable predictions for astronomical use. But it appeared that for this dataset we were not able to get a fit within experimental error (the weighted root mean square (rms) deviation was about 4.5). One of our assumptions was that the obtained solution corresponded to some local minima and in order to try to get out of it we decided to add the second excited torsional state to our analysis. Following several cycles of adding new $v_t=1$ and $v_t=2$ data to the fit and refinement of the parameter set confirmed correctness of our assumption and finally let us get a fit within experimental error for $v_t=0,1,2$ dataset.

The final dataset of the current study includes both our new measurements and the data from the literature [14] (6 lines out of 93 were excluded from Ref. [14] dataset because of large obs.-cal. values). The dataset contains 8333 A - and E -type transitions with $J \leq 60$ and $K_a \leq 25$ that due to blending correspond to 6908 measured lines. The final weighted rms deviation was 0.9 with all groups of data fitted within their measurement uncertainties. The final set of molecular parameters is presented in Table 1. It contains 62 parameters, which are distributed between the $n_{op} = 2, 4, 6, 8, 10$ orders as 7, 21, 25, 6, 3, respectively. This is consistent with the total numbers of determinable parameters of 7, 22, 50, 95, and 161 for those orders, as calculated from the differences between the total number of symmetry-allowed Hamiltonian terms of order n_{op} and the number of symmetry-allowed contact transformation terms of order $n_{op} - 1$ [31]. There were no convergence problems observed for the final fit.

Table 1: Fitted parameters of the RAM Hamiltonian for the thioacetaldehyde molecule.

n_{tr} ^a	Parameter ^b	Operator ^c	Value ^d
2 ₂₀	F	p_α^2	6.809672(26)
2 ₂₀	V_3	$\frac{1}{2}(1 - \cos 3\alpha)$	556.6140(24)
2 ₁₁	ρ	$J_z p_\alpha$	0.26038433(72)
2 ₀₂	A_{RAM}	J_z^2	1.61784741(70)
2 ₀₂	B_{RAM}	J_x^2	0.19945055(20)
2 ₀₂	C_{RAM}	J_y^2	0.176923079(33)
2 ₀₂	D_{zx}	$\{J_z, J_x\}$	-0.1047031(14)
4 ₄₀	F_m	p_α^4	$-0.39768(66) \times 10^{-3}$
4 ₄₀	V_6	$\frac{1}{2}(1 - \cos 6\alpha)$	-22.7413(28)
4 ₃₁	ρ_m	$J_z p_\alpha^3$	$-0.66956(86) \times 10^{-3}$
4 ₂₂	F_J	$J^2 p_\alpha^2$	$0.18497(48) \times 10^{-5}$
4 ₂₂	F_K	$J_z^2 p_\alpha^2$	$-0.80501(43) \times 10^{-3}$
4 ₂₂	F_{zx}	$\frac{1}{2} p_\alpha^2 \{J_z, J_x\}$	$0.14856(53) \times 10^{-4}$
4 ₂₂	V_{3J}	$J^2(1 - \cos 3\alpha)$	$0.253960(94) \times 10^{-3}$
4 ₂₂	V_{3K}	$J_z^2(1 - \cos 3\alpha)$	$-0.192064(64) \times 10^{-1}$
4 ₂₂	V_{3zx}	$\frac{1}{2}(1 - \cos 3\alpha) \{J_z, J_x\}$	$0.17186(19) \times 10^{-2}$
4 ₂₂	V_{3xy}	$(J_x^2 - J_y^2)(1 - \cos 3\alpha)$	$0.127485(86) \times 10^{-3}$

Table 1: continued.

ntr^a	Parameter ^b	Operator ^c	Value ^d
4 ₂₂	D_{3zy}	$\frac{1}{2} \sin 3\alpha \{J_z, J_y\}$	$-0.8237(15) \times 10^{-2}$
4 ₂₂	D_{3xy}	$\frac{1}{2} \sin 3\alpha \{J_x, J_y\}$	$0.2399(38) \times 10^{-4}$
4 ₁₃	ρ_J	$J^2 J_z p_\alpha$	$0.85246(25) \times 10^{-5}$
4 ₁₃	ρ_K	$J_z^3 p_\alpha$	$-0.46829(14) \times 10^{-3}$
4 ₁₃	ρ_{zx}	$\frac{1}{2} p_\alpha \{J_z^2, J_x\}$	$0.8797(57) \times 10^{-5}$
4 ₀₄	D_{zxJ}	$\frac{1}{2} \{J_z, J_x\} J^2$	$0.75397(24) \times 10^{-6}$
4 ₀₄	Δ_J	$-J^4$	$0.1210501(89) \times 10^{-6}$
4 ₀₄	Δ_{JK}	$-J^2 J_z^2$	$-0.447150(52) \times 10^{-5}$
4 ₀₄	Δ_K	$-J_z^4$	$0.108273(22) \times 10^{-3}$
4 ₀₄	δ_J	$-2J^2 (J_x^2 - J_y^2)$	$0.287961(44) \times 10^{-7}$
4 ₀₄	δ_K	$-\{J_z^2, (J_x^2 - J_y^2)\}$	$-0.688(11) \times 10^{-7}$
6 ₆₀	V_9	$\frac{1}{2} (1 - \cos 9\alpha)$	$0.3553(22)$
6 ₄₂	F_{mJ}	$J^2 p_\alpha^4$	$0.317(10) \times 10^{-8}$
6 ₄₂	V_{6J}	$J^2 (1 - \cos 6\alpha)$	$0.37406(66) \times 10^{-4}$
6 ₄₂	V_{6K}	$J_z^2 (1 - \cos 6\alpha)$	$-0.5154(70) \times 10^{-3}$
6 ₄₂	D_{6zy}	$\frac{1}{2} \sin 6\alpha \{J_z, J_y\}$	$0.4466(27) \times 10^{-3}$
6 ₃₃	ρ_{mK}	$J_z^3 p_\alpha^3$	$0.897(28) \times 10^{-7}$
6 ₃₃	ρ_{mxy}	$\frac{1}{2} p_\alpha^3 \{J_z, (J_x^2 - J_y^2)\}$	$-0.1371(54) \times 10^{-8}$
6 ₃₃	ρ_{3xy}	$\frac{1}{2} \{J_z, J_x, J_y, p_\alpha, \sin 3\alpha\}$	$0.1617(41) \times 10^{-5}$
6 ₂₄	F_{JK}	$J^2 J_z^2 p_\alpha^2$	$-0.3857(50) \times 10^{-8}$
6 ₂₄	F_{KK}	$J_z^4 p_\alpha^2$	$0.1421(19) \times 10^{-6}$
6 ₂₄	V_{3JJ}	$J^4 (1 - \cos 3\alpha)$	$-0.25678(69) \times 10^{-8}$
6 ₂₄	V_{3KK}	$J_z^4 (1 - \cos 3\alpha)$	$0.2616(28) \times 10^{-5}$
6 ₂₄	V_{3xyJ}	$J^2 (1 - \cos 3\alpha) (J_x^2 - J_y^2)$	$-0.12861(61) \times 10^{-8}$
6 ₂₄	V_{3x2y2}	$\frac{1}{2} \cos 3\alpha \{J_x^2, J_y^2\}$	$0.2712(46) \times 10^{-8}$
6 ₂₄	V_{3zxK}	$\frac{1}{2} (1 - \cos 3\alpha) \{J_z^3, J_x\}$	$-0.9077(71) \times 10^{-6}$
6 ₂₄	D_{3xyK}	$\frac{1}{2} \sin 3\alpha \{J_z^2, J_x, J_y\}$	$0.2067(73) \times 10^{-6}$
6 ₁₅	ρ_{JK}	$J^2 J_z^3 p_\alpha$	$-0.3414(21) \times 10^{-8}$
6 ₁₅	ρ_{KK}	$J_z^5 p_\alpha$	$0.7595(49) \times 10^{-7}$
6 ₀₆	D_{zxJK}	$\frac{1}{2} J^2 \{J_z^3, J_x\}$	$0.389(17) \times 10^{-10}$
6 ₀₆	Φ_J	J^6	$0.19952(55) \times 10^{-12}$
6 ₀₆	Φ_{JK}	$J^4 J_z^2$	$-0.825(19) \times 10^{-11}$
6 ₀₆	Φ_{KJ}	$J^2 J_z^4$	$-0.8072(37) \times 10^{-9}$
6 ₀₆	Φ_K	J_z^6	$0.14083(59) \times 10^{-7}$
6 ₀₆	ϕ_J	$2J^4 (J_x^2 - J_y^2)$	$0.8980(27) \times 10^{-13}$
6 ₀₆	ϕ_{JK}	$J^2 \{J_z^2, (J_x^2 - J_y^2)\}$	$-0.2730(65) \times 10^{-11}$
8 ₆₂	V_{9K}	$J_z^2 (1 - \cos 9\alpha)$	$-0.699(33) \times 10^{-4}$
8 ₅₃	ρ_{6xy}	$\frac{1}{2} \{J_z, J_x, J_y, p_\alpha, \sin 6\alpha\}$	$-0.405(13) \times 10^{-6}$
8 ₂₆	F_{KKK}	$J_z^6 p_\alpha^2$	$0.1347(37) \times 10^{-10}$
8 ₂₆	V_{3JJJ}	$J^6 (1 - \cos 3\alpha)$	$0.825(27) \times 10^{-14}$
8 ₀₈	L_{JK}	$J^4 J_z^4$	$0.1473(71) \times 10^{-14}$
8 ₀₈	L_{KKJ}	$J^2 J_z^6$	$-0.458(21) \times 10^{-13}$
10 ₆₄	V_{9JJ}	$J^4 (1 - \cos 9\alpha)$	$-0.4033(48) \times 10^{-10}$
10 ₆₄	D_{9xyK}	$\frac{1}{2} \sin 9\alpha \{J_z^2, J_x, J_y\}$	$-0.793(34) \times 10^{-7}$
10 ₄₆	V_{6KKK}	$J_z^6 (1 - \cos 6\alpha)$	$-0.1877(97) \times 10^{-9}$

Table 1: continued.

ntr^a	Parameter ^b	Operator ^c	Value ^d
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^a $n = t + r$, where n is the total order of the operator, t is the order of the torsional part and r is the order of the rotational part, respectively.

^b Parameter nomenclature based on the subscript procedures of ref. [32].

^c $\{A, B\} = AB + BA$, $\{A, B, C\} = ABC + CBA$, etc. The product of the operator in the third column of a given row and the parameter in the second column of that row gives the term actually used in the torsion-rotation Hamiltonian of the program, except for F , ρ and A_{RAM} , which occur in the Hamiltonian in the form $F(p_\alpha - \rho J_z)^2 + A_{RAM} J_z^2$.

^d All values are in cm^{-1} (except ρ which is unitless). Statistical uncertainties are shown as one standard uncertainty in the units of the last two digits.

On the basis of the parameters of our final fit, we calculated a list of transitions for astronomical use. This list includes information on transition quantum numbers, transition frequencies, calculated uncertainties, lower state energies, and transition strengths. The predictions are made up to 700 GHz for $v_t \leq 2$ and $J \leq 60$, and we limit our predictions to the transitions with calculated uncertainties lower than 0.1 MHz. The RAM36 code labels energy levels by free rotor quantum number m , overall rotational angular momentum quantum number J , and K_a , K_c numbers. Table 2 summarizes relations between torsional v_t and free rotor m quantum numbers, also giving energies of the A/E torsional substates. Lower state energies in Table 2 (as well as in the list of transitions for astronomical use) are given referenced to the $J = 0$ A-type $v_t = 0$ level. This level was calculated to be 82.4851 cm^{-1} above the bottom of the torsional potential well. Thus $v_t=2$ E substate is approximately 168 cm^{-1} below the top of the barrier. The line strengths have been calculated using the values $\mu_a = 2.26 \text{ D}$ and $\mu_b = 0.56 \text{ D}$ [14], which were transformed to the RAM axis system of the current study. In addition, we provide the rotation-torsion part of the partition function $Q_{rt}(\text{T})$ of thioacetaldehyde calculated via direct summation over the rotational-torsion states. The states up to $J = 100$ and $v_t = 8$ were taken into account in $Q_{rt}(\text{T})$ calculation. The predictions, partition function as well as line list of the data set treated in our study are available in the Supplementary material associated with this article.

Table 2: Energies of the $J = 0$ levels for $v_t = 0, 1, 2$ states of thioacetaldehyde (in cm^{-1}).

v_t	A-type	m	E-type	m
0	0.0	0	0.0098	1
1	159.5637	-3	159.2535	-2
2	301.8678	3	305.8351	4

4. Radioastronomical Observations

4.1. Search toward Sgr B2(N2) with ALMA

We used the spectroscopic predictions derived in this work for thioacetaldehyde to search for this molecule toward the hot molecular core Sgr B2(N2). We used the imaging spectral line survey EMOCA (Exploring Molecular Complexity with ALMA) conducted toward Sgr B2(N) with the Atacama Large Millimeter/submillimeter Array (ALMA) in its observing cycles 0 and 1. Sgr B2(N) is a protocluster forming high-mass stars located in the giant molecular cloud Sgr B2, at about 100 pc in projection from the Galactic center and a distance of 8.2 kpc [33]. Details about the survey and the data calibration and imaging were reported in ref. [34]. In short, the survey was performed with a median angular resolution of $1.6''$ and a phase center located at $(\alpha, \delta)_{\text{J2000}} = (17^{\text{h}}47^{\text{m}}19^{\text{s}}87, -28^{\circ}22'16''.0)$. The survey covers the frequency range from 84.1 to 114.4 GHz with a spectral resolution of 488 kHz, which corresponds to 1.7 to 1.3 km s^{-1} over this frequency range. Here we analyze the spectrum at the peak position of Sgr B2(N2) located at $(\alpha, \delta)_{\text{J2000}} = (17^{\text{h}}47^{\text{m}}19^{\text{s}}86, -28^{\circ}22'13''.4)$. The spectrum is modeled with the Weeds software [35] under the

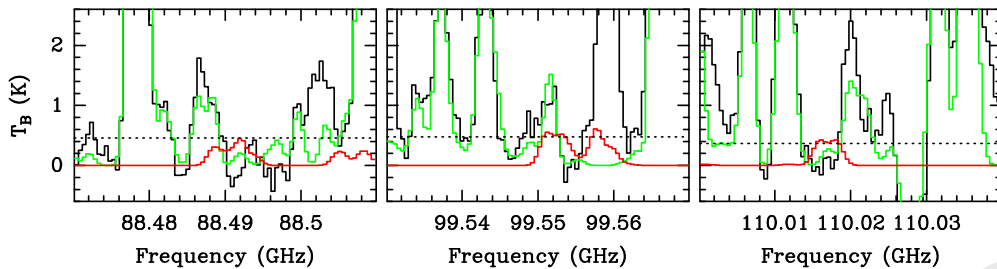


Figure 1: Synthetic LTE spectrum of thioacetaldehyde (in red) used to derive the upper limit to its column density, overlaid on the ALMA spectrum of Sgr B2(N2) (in black) and the synthetic spectrum that contains the contributions of all the species (but not thioacetaldehyde) that we have identified so far in this source on the basis of the EMOCA survey (in green). The dotted line in each panel indicates the 3σ noise level. Other transitions of thioacetaldehyde that are expected to be weaker than 3σ and/or heavily contaminated by other species are not shown.

assumption of local thermodynamic equilibrium (LTE), which is adequate given the high densities at the scales probed with EMOCA in Sgr B2(N2) ($> 1 \times 10^7 \text{ cm}^{-3}$, see [36]).

120 We did not detect any transition of thioacetaldehyde in the EMOCA spectrum of Sgr B2(N2). Figure 1 displays three frequency ranges that illustrate the non-detection. The LTE synthetic spectrum used to derive the upper limit to the column density of thioacetaldehyde is displayed in red and was computed assuming the same parameters (temperature, emission size, linewidth, and velocity offset with respect to the systemic velocity) as the ones we derived for acetaldehyde with the EMOCA survey [37]. These parameters are listed in Table 3, along with the parameters derived previously for methanol and methyl mercaptan in Sgr B2(N2) using the same survey [38]. We find that thioacetaldehyde is at least ~ 90 times less abundant than acetaldehyde in Sgr B2(N2). For comparison, methyl mercaptan is about 120 times less abundant than methanol in Sgr B2(N2). The non-detection of thioacetaldehyde is thus not surprising if the abundance ratio of oxygen-bearing to sulfur-bearing molecules is the same for these two families of organic molecules
125
130 in this source.

Table 3: Parameters of our best-fit LTE model of methanol, methyl mercaptan, acetaldehyde^a, and column density upper limit for thioacetaldehyde, toward Sgr B2(N2) on the basis of the EMOCA survey.

Molecule	Status ^b	N_{det} ^c	Size ^d ($''$)	T_{rot} ^e (K)	N ^f (cm^{-2})	F_{vib} ^g	ΔV ^h (km s^{-1})	V_{off} ⁱ (km s^{-1})	$\frac{N_{\text{ref}}}{N}$ ^j
CH_3OH , $v = 0^*$	d	41	1.4	160	4.0 (19)	1.00	5.4	-0.5	1
CH_3SH , $v = 0$	d	12	1.4	180	3.4 (17)	1.00	5.4	-0.5	118
CH_3CHO , $v = 0^*$	d	19	1.1	160	5.3 (17)	1.08	5.6	0.0	1
CH_3CHS , $v = 0$	n	0	1.1	160	< 6.2 (15)	1.03	5.6	0.0	> 86

^a The parameters for methanol and methyl mercaptan were published in ref. [38], and for acetaldehyde in ref. [37].

^b d: detection, n: non-detection.

^c Number of detected lines (conservative estimate, see Sect. 3 of ref.[34]). One line of a given species may mean a group of transitions of that species that are blended together.

^d Source diameter ($FWHM$).

^e Rotational temperature.

^f Total column density of the molecule. x (y) means $x \times 10^y$.

^g Correction factor that was applied to the column density to account for the contribution of vibrationally excited states, in the cases where this contribution was not included in the partition function of the spectroscopic predictions.

^h Linewidth ($FWHM$).

ⁱ Velocity offset with respect to the assumed systemic velocity of Sgr B2(N2), $V_{\text{sys}} = 74 \text{ km s}^{-1}$.

^j Column density ratio, with N_{ref} the column density of the previous reference species marked with a $*$.

4.2. Search using Astrochemical Surveys at IRAM Large Program

We have searched for thioacetaldehyde in a number of observational datasets from the Astrochemical Surveys at IRAM Large Program, which are publicly available and have been previously published elsewhere.

The ASAI spectra cover a range of source types from cold, dark clouds to Class 0/1 protostars and shocked outflows, with observational details available in Ref. [39]. Because the observational details of these sources are described elsewhere; only brief overview is provided here. For each source, because these are non-detections, we have had to make assumptions about the physical conditions in which thioacetaldehyde might be found. For all sources considered in this subsection, these parameters were chosen to match those used in a similar analysis of cyanoketene [40] (since many of these sources do not have acetaldehyde detections in them). These parameters and the literature sources they were obtained from are provided in Table 4.

Table 4: Source parameters assumed for thioacetaldehyde in each of the ASAI sets of observations.

Source	Telescope	θ_s^a ($''$)	T_{bg}^b (K)	ΔV (km s^{-1})	T_b^\dagger (mK)	T_{ex}^b (K)	Refs.
Barnard 1	IRAM	–	2.7	0.8	2.5	10	[41], [42]
IRAS 4A	IRAM	–	2.7	5.0	2.3	21	[41], [43]
L1157B1	IRAM	–	2.7	8.0	1.4	60	[44]
L1157mm	IRAM	–	2.7	3.0	2.8	60	[44]
L1448R2	IRAM	–	2.7	8.0	2.4	60	[45]
L1527	IRAM	–	2.7	0.5	1.9	12	[45], [46]
L1544	IRAM	–	2.7	0.5	2.3	10	[47], [48]
SVS13A	IRAM	0.3	2.7	3.0	7.5	80	[41], [43]
TMC1	IRAM	–	2.7	0.3	6.5	7	[49], [50]

^aExcept where noted, the source is assumed to fill the beam.

^b T_{bg} is the background temperature. T_{ex} is the excitation temperature.

[†]Taken either as the 3σ RMS noise level at the location of the target line, or for line confusion limited spectra, the reported 3σ RMS noise of the observations.

We find no evidence of thioacetaldehyde in any of our studied sources. We derive upper limits to the column density in each source using the formalisms outlined in Ref. [51], which assume the molecules are well described by a single excitation temperature and include corrections for optical depth. Frequencies, energy levels, degeneracies, and line strengths were obtained from the laboratory spectroscopy described in this work. The partition function, which includes a correction for the lowest torsional vibrational mode of the molecule, is given in Supplementary Material associated with this article. In warmer source ($T > 100$ K), vibrational corrections from higher lying modes must also be considered. To that end, we have carried out quantum chemical calculations using Gaussian 09 [52] at the wB97XD/6-311++G(d,p) level of theory and basis set to determine the harmonic frequencies of thioacetaldehyde. The harmonic frequencies are also given in the Supplementary Material.

For each source, we simulated a spectrum using the physical conditions for the source and the line parameters measured in this work. Then, the 3σ upper limit to the column density was derived, using the line that gave the most rigorous constraint (i.e. the line that would have the highest signal-to-noise ratio in the event of a detection). All lines used to derive upper limits to the column density belong to the ground torsional state of thioacetaldehyde. These lines are provided in Table 5 along with the resulting upper limits.

5. Conclusion

A new study of the rotational spectrum of thioacetaldehyde CH_3CHS was carried out in the frequency range of 150–660 GHz in order to provide accurate predictions for astronomical searches. The rotational transitions of the ground, first and second excited torsional states were assigned up to $J = 60$ and fit using a RAM Hamiltonian within experimental accuracy. The overall weighted rms deviation of the global fit of 6908 line frequencies of CH_3CHS is 0.9, indicating that we have an appropriate set of parameters to provide reliable predictions for astronomical observations. Unfortunately, our first attempt to search for CH_3CHS with ALMA data between 84 GHz and 114 GHz toward the hot molecular core Sgr B2(N2) was not successful. The non-detection implies that thioacetaldehyde is at least 90 times less abundant than

Table 5: 3σ upper limits to thioacetaldehyde and the line parameters used to calculate them in each of the ASAI sets of observations.

Source	Frequency (MHz)	Transition ($S: J''_{K_a, K_c} - J''_{K_a, K_c}$)	E_u (K)	$S_{ij}\mu^2$ (Debye ²)	Q (Q_{tr}, Q_{vib}) ^a	N_T (cm ⁻²)	$N(H_2)$ (cm ⁻²)	X_{H_2}	Refs.
Barnard 1	101408.2	$E: 9_{1,8} - 8_{1,7}$	22.4	45.4	267 (267, 1.00)	$\leq 6.9 \times 10^{11}$	1.5×10^{23}	$\leq 6 \times 10^{-12}$	[42]
IRAS 4A	108228.9	$A: 10_{1,10} - 9_{1,9}$	30.7	50.6	861 (861, 1.00)	$\leq 2.5 \times 10^{12}$	3.7×10^{23}	$\leq 6 \times 10^{-12}$	[42]
L1157B1	108228.9	$A: 10_{1,10} - 9_{1,9}$	30.7	50.6	861 (861, 1.00)	$\leq 3.6 \times 10^{12}$	1×10^{21}	$\leq 3 \times 10^{-9}$	[42]
L1157mm	108228.9	$A: 10_{1,10} - 9_{1,9}$	30.7	50.6	861 (861, 1.00)	$\leq 4.2 \times 10^{12}$	6×10^{21}	$\leq 6 \times 10^{-10}$	[42]
L1448R2	97429.9	$A: 9_{1,9} - 8_{1,8}$	25.5	45.4	4161 (4161, 1.00)	$\leq 7.2 \times 10^{12}$	3.5×10^{23}	$\leq 2 \times 10^{-11}$	[45]
L1527	86623.7	$A: 8_{1,8} - 7_{1,7}$	20.8	40.2	364 (364, 1.00)	$\leq 3.0 \times 10^{11}$	2.8×10^{22}	$\leq 1 \times 10^{-11}$	[45]
L1544	99110.2	$E: 9_{0,9} - 8_{0,8}$	23.8	46.0	267 (267, 1.00)	$\leq 2.9 \times 10^{11}$	5×10^{21}	$\leq 6 \times 10^{-11}$	[53]
SVS13A	255531.7	$A: 23_{3,20} - 22_{3,19}$	165.4	116.0	6666 (6661, 1.00)	$\leq 1.0 \times 10^{16}$	3×10^{24}	$\leq 3 \times 10^{-9}$	[54]
TMC1	135082.5	$E: 12_{1,11} - 11_{1,10}$	44.3	60.9	129 (129, 1.00)	$\leq 1.0 \times 10^{13}$	1×10^{22}	$\leq 9 \times 10^{-10}$	[42]

^aCalculated at the excitation temperature assumed for the source. See Table 4.

165 acetaldehyde in this source. We also find no evidence of thioacetaldehyde in any of the prestellar and protostellar sources targeted by the ASAI Large Program. Based on our results, frequency predictions for astrophysical use were produced up to 700 GHz. The predictions, as well as data set treated in our study, are available as supplementary material to this article.

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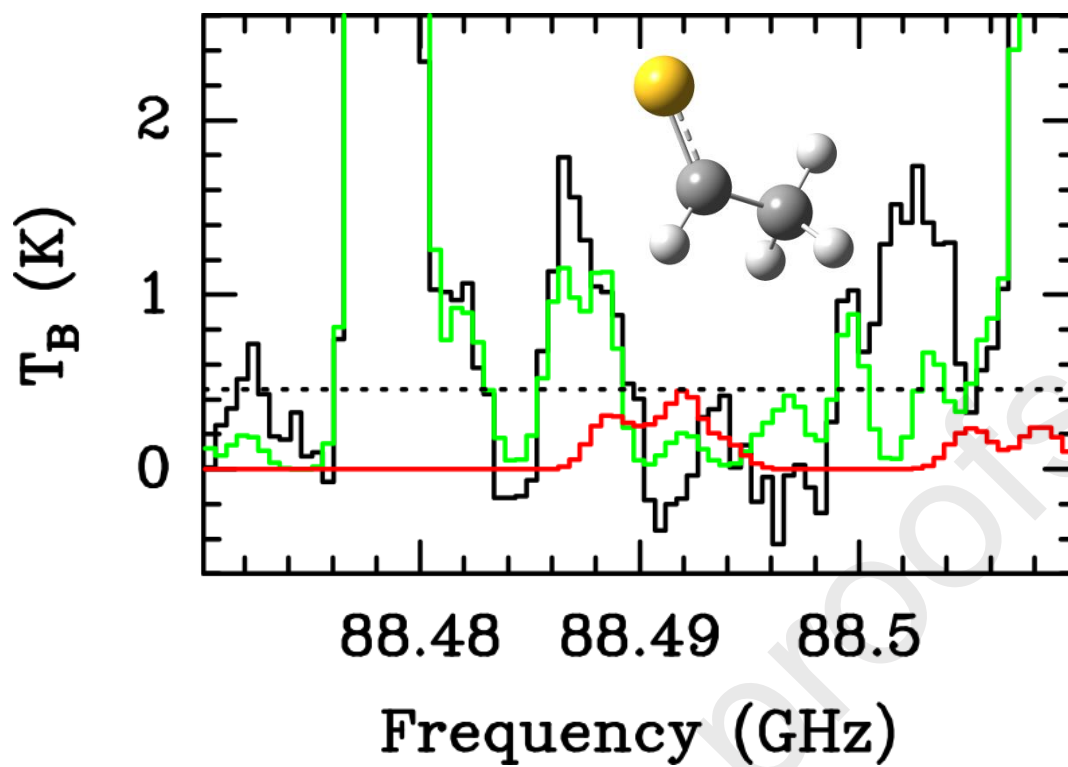
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A new study of the CH₃CHS spectrum up to 660 GHz has been carried out. The rotation transitions in $v_t=0,1,2$ states with J up to 60 have been assigned. The new data were fit within experimental error using the Rho axis method. Thioacetaldehyde was searched for, but not detected, in the ISM.

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Declaration of interests

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

The authors declare the following financial interests/personal relationships which may be considered as potential competing interests: