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► **To cite this version:**

Teddy Pichard, Nina Aguillon, Bruno Després, Edwige Godlewski, Michael Ndjinga. Existence and uniqueness of generalized solutions to hyperbolic systems with linear fluxes and stiff sources. *Journal of Hyperbolic Differential Equations*, 2021, 18 (3), pp.653-700. 10.1142/S021989162150020X. hal-02920155

HAL Id: hal-02920155

<https://hal.science/hal-02920155v1>

Submitted on 24 Aug 2020

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EXISTENCE AND UNIQUENESS OF GENERALIZED SOLUTIONS TO HYPERBOLIC SYSTEMS WITH LINEAR FLUXES AND STIFF SOURCES

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Received (Day Mth. Year)

Revised (Day Mth. Year)

Communicated by [editor]

Abstract. Motivated by the modelling of boiling two-phase flows, we study systems of balance laws with a source term defined as a discontinuous function of the unknown. Due to this discontinuous source term, the classical theory of partial differential equations (PDE) is not sufficient here. Restricting to a simpler system with linear fluxes, a notion of generalized solution is developed. An important point in the construction of a solution is that the curve along which the source jumps, which we call the boiling curve, must never be tangent to the characteristics. This leads to exhibit sufficient conditions which ensure the existence and uniqueness of a solution in two different situations: first when the initial data is smooth and such that the boiling curve is either overcharacteristic or

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subcharacteristic; then with discontinuous initial data in the case of Riemann problems. A numerical illustration is given in this last case.

Keywords: Balance laws; Stiff sources; Generalized solutions.

1. Introduction

The starting point is a system of balance laws of the form

$$\partial_t U + \partial_x F(U) = \begin{cases} S^- & \text{if } h(U) \leq h^b, \\ S^+ & \text{if } h(U) > h^b. \end{cases} \quad (1.1)$$

For instance, a system of the form (1.1) was obtained as a drift-flux model ([21, 19,20,27,3,26]) for homogenized boiling two-phase flows. The discontinuity of the source term approximates the boiling phenomenon, which is much faster than the time scale of the flow. In this context, the unknown U and the flux $F(U)$ are given by

$$U = (\alpha\rho_v, \rho, \rho u, \rho e), \quad F(U) = (\alpha\rho_v u, \rho u, \rho u^2 + p, \rho e + pu)$$

where $\alpha\rho_v$ is the partial density of vapor, and $(\rho, \rho u, \rho e)$ are the total density, the total momentum and the total energy of the mixture of vapor and liquid. Here, $p(U)$ is some pressure law, and $h(U)$ is an enthalpy function. The liquid is assumed to boil only when the enthalpy $h(U)$ is above the boiling threshold h^b . The source is chosen of the form

$$S^- = (0, 0, 0, \phi), \quad S^+ = (K\phi, 0, 0, \phi),$$

which corresponds to imposing a source $\phi > 0$ of energy everywhere and a creation $K\phi > 0$ of vapor only above the boiling threshold $h(U) > h^b$. See e.g. [9] for further thermodynamics consideration on the construction of such boiling models.

Most of the results from the "classical theory" require the source term $S(U)$ to be Lipschitz-continuous, and are therefore unavailable here. For instance, in a simple ordinary differential equation (ODE) framework (consider e.g. a 0D version of (1.1) or see examples in Section 2), the Picard-Lindelöf theorem cannot be applied and neither uniqueness nor existence of a solution to such an ODE is ensured. Even the notion of "classical solution" is unclear and one typically works with the integral notion of Carathodory solutions ([11]). In the literature, several generalizations of the Picard-Lindelöf theorem provide similar results under weaker regularity requirements on the source term, e.g. under directional continuity or bounded variations ([25,5,7,18]). Other approaches provide results in a more general context, e.g. in Filippov theory ([13,2,8]), the ODE is replaced by a differential inclusion, and a wider set of solutions is sought. However, these results mainly apply to ODE problems and there are few generalizations to PDE, to the authors' knowledge.

This work focuses on a system with linear fluxes oversimplified compared

to (1.1), however presenting already interesting features:

$$\partial_t u - \mu \partial_x u = \sigma(u, v), \quad u(0, x) = u_0(x), \quad (1.2a)$$

$$\partial_t v + \nu \partial_x v = \varsigma(u, v), \quad v(0, x) = v_0(x), \quad (1.2b)$$

with source terms defined by

$$\sigma(u, v) = \begin{cases} a & \text{if } h(u, v) \leq 0, \\ b & \text{if } h(u, v) > 0, \end{cases} \quad \varsigma(u, v) = \begin{cases} c & \text{if } h(u, v) \leq 0, \\ d & \text{if } h(u, v) > 0, \end{cases} \quad (1.2c)$$

and where the function h is linear

$$h(u, v) = h_1 u + h_2 v. \quad (1.2d)$$

For simplicity, the constants are chosen strictly positive

$$h_1 > 0, \quad h_2 > 0, \quad \mu > 0, \quad \nu > 0,$$

though the results below can be generalized. Also, one remarks that the sources σ and ς in the two equations are discontinuous at the same (x, t) -positions since they both jump when $h(u, v) = 0$.

System (1.2) retains the discontinuity of the source with respect to the unknown, but the fluxes and the enthalpy are simplified into linear functions and the system is only composed of two equations. The linearity of h and the reduction of the number of equations is only meant to simplify the mathematical analysis, and the results below could be generalized under weaker but more technical requirements. On the contrary, the hypothesis on the linearity of the fluxes is necessary in the present work. It allows to use the method of characteristics and it prevents the formation of shock and rarefaction waves. Extending the present work with nonlinear fluxes would naturally require to consider such nonlinear waves.

The main results of this paper consist in constructing a solution to (1.2) in three different cases, i.e. subject to various inequalities involving the initial data u_0 and v_0 and the source values a, b, c and d . The zone of transition Γ of the source terms plays an important role in the construction below. This is defined by

$$\Gamma = \{(x, t) \in \mathbb{R} \times \mathbb{R}^+ \text{ s.t. } h(u, v)(x, t) = 0\}. \quad (1.3)$$

In all the cases described in the theorems below, the set Γ is a line, and is afterward called the "boiling line". The three sets of conditions required for the well-posedness of (1.2) correspond to imposing that

- for regular initial conditions,
 - either Γ is an overcharacteristic line, i.e. it remains outside the cones generated by the two characteristics passing at every point $(x, t) \in \Gamma$ (see Fig. 1 top left). A precise definition is given in Section 3.1 and the precise requirements and the result are stated in Theorem 3.3 and the proof is detailed in Section 4;

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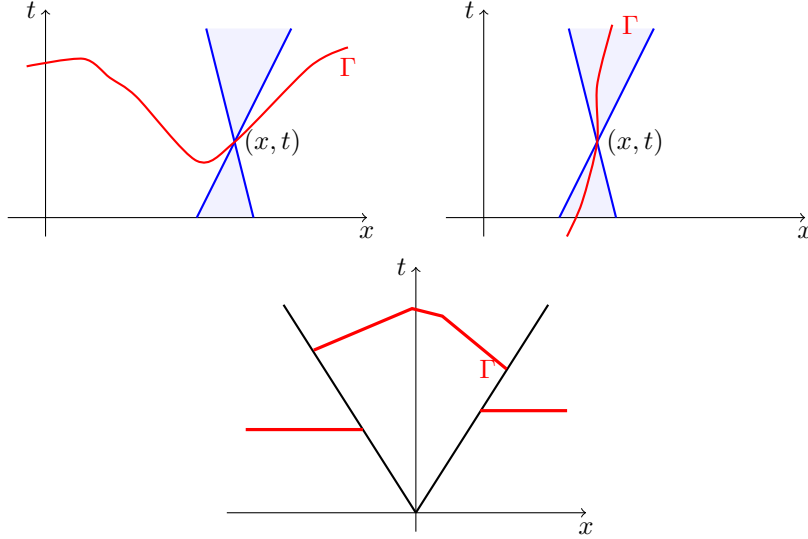


Fig. 1. Examples of boiling lines Γ (red) that are overcharacteristic (top left), subcharacteristic (top right) or in the case of Riemann problems (bottom), with the two characteristic curves (blue) passing at a point $(x, t) \in \Gamma$.

- or Γ is a subcharacteristic line, i.e. it remains inside the cones generated by the two characteristics passing at every point $(x, t) \in \Gamma$ (see Fig. 1 top right). Similarly the precise definition is given in Section 3.1, the statement in Theorem 3.5 and the proof in Section 5;
- for initial conditions u_0 and v_0 that are discontinuous at $x = 0$ and constant elsewhere, i.e. in the case of Riemann problems, the set Γ is required to be a line which is overcharacteristic in the middle region (see Fig. 1 bottom), the statement of the result is given by Theorem 3.7 and the proof in Section 6.

We give sufficient conditions for existence and uniqueness of a generalized solution. However, the conditions are not necessary and, as illustrated in a simplified framework in Section 2.1 (the first two cases of this test case), it remains unclear what happens when these conditions are violated.

The paper is organized as follows. In the next section, the main problem and basic ideas are presented, especially the notion of *generalized solution* is introduced. In Section 3, the notion of non-characteristic boiling line is introduced and the main theorems are presented. The following three sections are devoted to prove these results. Section 7 gathers the concluding comments.

2. Generalized solutions

First, a simplified ODE framework is considered. Then, an extension to PDE's is described.

2.1. Carathodory solutions for ODE

When considering discontinuous sources in an ODE framework, it is possible to construct solutions when convenient inequalities on the data hold ([23,18,24]).

Consider the scalar ODE

$$\frac{du}{dt}(t) = s(u(t), t), \quad u(0) = u_0. \quad (2.1)$$

If s is discontinuous at some point (u, t) , then the derivative $\frac{du}{dt}$ is not defined in a classical sense at this point.

Definition 2.1 ([11,23]). Consider $u_0 \in \mathbb{R}$ and $s \in L^1(\mathbb{R} \times]0, T[)$. A function $u \in W^{1,1}(]0, T[)$ is a Carathodory solution to (2.1) on $]0, T[$ if it satisfies

$$\forall t \in]0, T[, \quad u(t) = u_0 + \int_0^t s(u(\tau), \tau) d\tau. \quad (2.2)$$

Remark that when s is continuous, then a C^1 function satisfying (2.1) for all $t \in \mathbb{R}^+$ also satisfies (2.2). Thus, this definition generalizes the notion of solution to ODE with discontinuous RHS.

Consider now that the source $s = s(u)$ depends only on u and is given by

$$s(u) = \begin{cases} a & \text{if } u \leq 0, \\ b & \text{otherwise.} \end{cases} \quad (2.3)$$

The space of parameters $(a, b) \in \mathbb{R}^2$ can be decomposed into three subsets. Each of these subsets corresponds to a particular behaviour of the solutions. This can be summarized by the geometrical representation of the solutions on Fig. 2 and by:

- $b \leq 0 \leq a$ and $a \neq b$

The Carathodory solution u is attracted toward $u = 0$ (left plot on Fig. 2).

The function

$$u(t) = \begin{cases} u_0 + at & \text{if } u_0 \leq 0, \\ u_0 + bt & \text{if } u_0 \geq 0, \end{cases} \quad (2.4)$$

is the unique Carathodory solution to (2.1-2.3) on $t \in [0, T_{\max}[$ with

$$T_{\max} = \begin{cases} \frac{-u_0}{a} & \text{if } a \neq 0 \text{ and } u_0 \leq 0, \\ \frac{-u_0}{b} & \text{if } b \neq 0 \text{ and } u_0 \geq 0, \\ +\infty & \text{otherwise.} \end{cases}$$

This solution (2.4) cannot be prolonged for times $t > T_{\max}$.

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- $a \leq 0 \leq b$

The Carathodory solution u is repulsed away from $u = 0$ (right plot on Fig. 2).

If $u_0 \neq 0$, the function

$$u(t) = \begin{cases} u_0 + at & \text{if } u_0 < 0, \\ u_0 + bt & \text{if } u_0 > 0, \end{cases} \quad (2.5)$$

is the unique Carathodory solution to (2.1-2.3) on \mathbb{R}^+ .

If $u_0 = 0$, both $t \mapsto at$ and $t \mapsto bt$ are Carathodory solutions to (2.1-2.3) on $t \in \mathbb{R}^+$.

- $ab > 0$

The Carathodory solution u crosses the discontinuity line $u = 0$.

Proposition 2.2. *Suppose that $ab > 0$. Then, for all $u_0 \in \mathbb{R}$, there exists a unique Carathodory solution to (2.1-2.3) on \mathbb{R}^+ .*

This ensures **existence and uniqueness of a Carathodory solution for all $t \in \mathbb{R}^+$** . This is a straightforward consequence of various more general results, e.g. [25,5,7,24,17]. The proof is detailed since its main idea is afterward extended in the PDE framework.

Proof. Suppose that $a > 0$ and $b > 0$ (similar computations hold if $a < 0$ and $b < 0$). The function (right plot on Fig. 2).

$$u(t) = \begin{cases} u_0 + bt & \text{if } u_0 > 0, \\ u_0 + at & \text{if } u_0 \leq 0 \text{ and } t \leq \frac{-u_0}{a}, \\ b \left(t + \frac{u_0}{a} \right) & \text{if } u_0 \leq 0 \text{ and } t > \frac{-u_0}{a}, \end{cases} \quad (2.6)$$

is absolutely continuous and satisfies (2.1-2.3) for all $t \in \mathbb{R}^+$. Therefore, it is a Carathodory solution to (2.1-2.3) on \mathbb{R}^+ .

Since $s > 0$, any Carathodory solution \tilde{u} is strictly increasing. Thus, there exists a unique switch time t_s such that $\tilde{u}(t_s) = 0$. Thus the source has the form $s(\tilde{u}(t)) = \tilde{s}(t)$ with

$$\tilde{s}(t) = \begin{cases} a & \text{if } t \leq t_s, \\ b & \text{if } t > t_s. \end{cases} \quad (2.7)$$

Injecting (2.7) in (2.2) leads to a unique $t_s = \frac{-u_0}{a}$, and thus a unique $\tilde{u} = u$. □

The next sections focus on a problem corresponding to a PDE version of the last case.

2.2. Generalized solutions in a PDE framework

The PDE framework is formalized in the scalar case, then extended to the case of system (1.2).

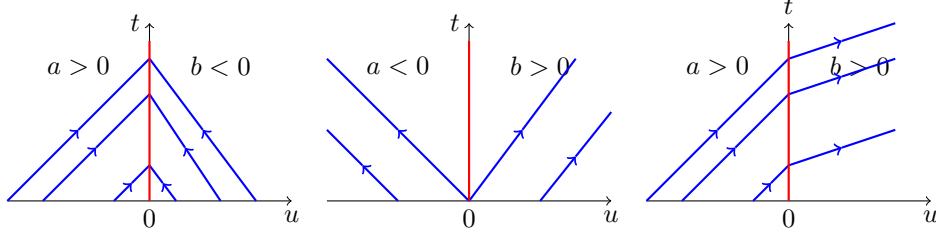


Fig. 2. Representation in phase space (u, t) of the functions (2.4; left), (2.5; middle) and (2.6; right) for different initial conditions u_0 .

2.2.1. Scalar case

Consider the scalar equation

$$(\partial_t u + \lambda \partial_x u)(x, t) = s(u(x, t), x, t), \quad u(0, x) = u_0(x), \quad (2.8)$$

where $\lambda \in \mathbb{R}$ is some given constant. First, the method of characteristics is formalized.

Definition 2.3. The line $C_\lambda(x)$, resp. $C_\lambda(x, t)$, is defined as the characteristic curve issuing from $(x, 0)$, resp. passing at (x, t) . These read

$$C_\lambda(x) := \{(x + \lambda\tau, \tau), \quad \tau \in \mathbb{R}^+\}, \quad (2.9a)$$

$$C_\lambda(x, t) := \{(x - \lambda\eta, t - \eta), \quad \eta \in]-\infty, t]\} \equiv C_\lambda(x - \lambda t), \quad (2.9b)$$

where (2.9a), resp. (2.9b), is written for convenience in the forward, resp. backward, direction.

Following Definition 2.1, a generalized solution is defined as an integral solution along the characteristics.

Definition 2.4. For a function $u \in L^1_{loc}(\mathbb{R} \times \mathbb{R}^+)$, denote

$$u_{\lambda, x} : \begin{cases} \mathbb{R}^+ \rightarrow \mathbb{R} \\ \tau \mapsto u(x + \lambda\tau, \tau) \end{cases}$$

Define the set

$$E_\lambda := \left\{ u \in L^1_{loc}(\mathbb{R} \times \mathbb{R}^+) \quad s.t. \quad \forall x \in \mathbb{R}, \quad u_{\lambda, x} \in W^{1, \infty}_{loc}(\mathbb{R}^+) \right\}. \quad (2.10)$$

Definition 2.5. Consider $u_0 \in L^1_{loc}(\mathbb{R})$ and $s \in L^\infty(\mathbb{R} \times \mathbb{R} \times \mathbb{R}^+)$ function of the variables (u, x, t) . A function $u \in E_\lambda$ is a **generalized solution** to (2.8) if it satisfies for all $(x, t) \in \mathbb{R} \times \mathbb{R}^+$

$$u(x, t) = u_0(x - \lambda t) + \int_0^t s(u(x - \lambda\eta, t - \eta), x - \lambda\eta, t - \eta) d\eta. \quad (2.11)$$

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Again, when s is continuous, a function u that is C^1 and satisfies (2.8) for all $(x, t) \in \mathbb{R} \times \mathbb{R}^+$ also satisfies (2.11). Thus, this definition generalizes the notion of classical solution. However, solutions to hyperbolic equations are rarely expected to have such a strong regularity, and another notion of generalized solution was introduced in [1], the present one extends it with discontinuous sources for 1D problems.

Remark 2.6. To simplify the analysis in the next sections, two choices about the regularity of the different functions were performed in Definitions 2.4 and 2.5:

- Since all the source terms in our applications are bounded, we have chosen to fix the regularity of s to L^∞ in Definition 2.5 and thus we required a $W^{1,\infty}$ regularity in Definition 2.4 along the characteristics. However, weaker integrability of s could have been sufficient to define the integral in (2.11).

On the contrary in the ODE case, for generality (as it is easier to define), the source s was only L^1 in Definition 2.1, and the regularity of the Carathodory solutions was fixed to $W^{1,1}$.

- The spatial domain was chosen to be the full real line to avoid considering boundary conditions. By extension, the time variable can also go to infinity. However, in most cases we consider below, the integral in (2.11) becomes infinite when $t \rightarrow +\infty$. Then the integrability of the solutions $u \in E_\lambda$ can only be local. By extension, as the solution is only locally integrable, there was no point of restricting the initial condition u_0 with global integrability.

Of course, one could restrict the (x, t) -domain into a bounded set (with boundary conditions) and restrict the solution to be in L^1 such that $u_{\lambda,x}$ is $W^{1,\infty}$. Such a restriction was considered in the ODE in cases as the time domain $]0, T[$ can be bounded.

1- Case s independent of u : Suppose here that

$$s(u, x, t) = s(x, t),$$

with $s \in L^\infty(\mathbb{R} \times \mathbb{R}^+)$, i.e. suppose that the source depends only on the variables (x, t) and not on the unknown u .

With such a source, the integral on the RHS of (2.11) is well-defined and for all $u_0 \in L^1_{loc}(\mathbb{R})$, there exists a unique generalized solution to (2.8) on $(x, t) \in \mathbb{R} \times \mathbb{R}^+$.

To shorten the writings, a notation for this solution is introduced.

Notation 2.7. Given $\lambda \in \mathbb{R}$, $u_0 \in L^1_{loc}(\mathbb{R})$ and $s \in L^\infty(\mathbb{R} \times \mathbb{R}^+)$, denote

$$r_\lambda^{u_0}(s) : \begin{cases} \mathbb{R} \times \mathbb{R}^+ \rightarrow \mathbb{R}, \\ (x, t) \mapsto u_0(x - \lambda t) + \int_0^t s(x - \lambda\eta, t - \eta) d\eta. \end{cases} \quad (2.12)$$

2- Case s independent of (x, t) : Suppose here that $s = s(u)$ with

$$s(u) = \begin{cases} a & \text{if } u \leq 0, \\ b & \text{otherwise,} \end{cases} \quad (2.13)$$

where $(a, b) \in \mathbb{R}^2$ are given constants.

One easily obtains existence and uniqueness of such a generalized solution.

Corollary 2.8. *Suppose that $ab > 0$. Then for all $u_0 \in L^1_{loc}(\mathbb{R})$, there exists a unique generalized solution to (2.8, 2.13) on $\mathbb{R} \times \mathbb{R}^+$.*

Proof. By definition, u is a generalized solution to (2.8) if for all $x \in \mathbb{R}$, the function $t \mapsto u(x + \lambda t, t)$ is a Carathodory solution to (2.1-2.3), where x acts as a parameter. Proposition 2.2 provides the result. \square

2.2.2. Case of a system

Notation 2.9. The dependent variables associated to a system are denoted with a capital letter, e.g. $U = (u, v)$, similarly $\Sigma = (\sigma, \varsigma)$ (also with subscripts or superscripts), while those associated to the sub-equations are denoted with small letters. We will write identically $f(U) = f(u, v)$.

Consider the system of equations in $U = (u, v)$

$$(\partial_t u - \mu \partial_x u)(x, t) = \sigma(u(x, t), v(x, t), x, t), \quad u(x, 0) = u_0(x), \quad (2.14a)$$

$$(\partial_t v + \nu \partial_x v)(x, t) = \varsigma(u(x, t), v(x, t), x, t), \quad v(x, 0) = v_0(x). \quad (2.14b)$$

where μ, ν are given constants and Σ given sources.

Notation 2.10. Denote

$$E = E_{-\mu} \times E_{\nu}.$$

Definition 2.11. *Consider $U_0 \in (L^1_{loc}(\mathbb{R}))^2$ and $\Sigma = (\sigma, \varsigma) \in (L^\infty(\mathbb{R}^2 \times \mathbb{R} \times \mathbb{R}^+))^2$ depending on (U, x, t) . Then, a function $U = (u, v) \in E$ is a **generalized solution of** (2.14) if U satisfies for all $(x, t) \in \mathbb{R} \times \mathbb{R}^+$*

$$u(x, t) = u_0(x + \mu t) + \int_0^t \sigma(u(x + \mu\eta, t - \eta), v(x + \mu\eta, t - \eta), x + \mu\eta, t - \eta) d\eta, \quad (2.15a)$$

$$v(x, t) = v_0(x - \nu t) + \int_0^t \varsigma(u(x - \nu\eta, t - \eta), v(x - \nu\eta, t - \eta), x - \nu\eta, t - \eta) d\eta. \quad (2.15b)$$

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1- Case Σ independent of U : Suppose here that

$$\forall(U, x, t) \in \mathbb{R}^2 \times \mathbb{R} \times \mathbb{R}^+, \quad \Sigma(U, x, t) = \tilde{\Sigma}(x, t),$$

where $\tilde{\Sigma} \in (L^\infty(\mathbb{R} \times \mathbb{R}^+))^2$ depends only on $(x, t) \in \mathbb{R} \times \mathbb{R}^+$ and not on the unknown U .

With such a source, the integrals on the RHS of (2.15) are well-defined. Thus, for all $U_0 \in (L^1_{loc}(\mathbb{R}))^2$, there exists a unique generalized solution U to (2.8) on $(x, t) \in \mathbb{R} \times \mathbb{R}^+$.

Notation 2.12. Given $U_0 = (u_0, v_0) \in (L^1_{loc}(\mathbb{R}))^2$ and $\tilde{\Sigma} = (\tilde{\sigma}, \tilde{\zeta}) \in (L^\infty(\mathbb{R} \times \mathbb{R}^+))^2$, using the notation 2.7, denote

$$R(\tilde{\Sigma}) = (r_{-\mu}^{u_0}(\tilde{\sigma}), r_{\nu}^{v_0}(\tilde{\zeta})). \quad (2.16)$$

2- Case Σ independent of (x, t) : Coming back to (1.2) with a source $\Sigma = (\sigma, \varsigma)$ of the form (1.2c), i.e. depending only on the unknown U , then the integral equation (2.15) rewrites

$$U = R(\Sigma(U)). \quad (2.17)$$

Notation 2.13. Emphasizing the dependence on U , note $\Gamma(U)$ the set Γ defined in (1.3).

In the present work, only cases where $\Gamma(U)$ is a line are considered.

Notation 2.14. Suppose that $U = (u, v) \in E$ is such that $\Gamma(U)$ is a line. The (x, t) -space is decomposed into two zones separated by $\Gamma(U)$

$$Z^+(U) := \{(x, t) \in \mathbb{R} \times \mathbb{R}^+ \text{ s.t. } h(u, v)(x, t) > 0\}, \quad Z^-(U) := \mathbb{R} \times \mathbb{R}^+ \setminus Z^+(U).$$

In such a case, the boiling line satisfies $\Gamma(U) = \overline{Z^+(U)} \cap Z^-(U)$.

When Σ depends U , the two equations of (2.15) are coupled. Indeed, the evaluation of both integrals in the right-hand side require the value of both $U = (u, v)$ in all the so-called cone of dependence ([6] (in pale blue on Fig. 3) delimited by $C_{-\mu}(x, t)$ and $C_{\nu}(x, t)$, and no straightforward application of Corollary 2.8 provides existence and uniqueness of a generalized solution. This problem requires further investigations which are described below.

2.3. Construction of a generalized solution

First, the method is presented as a prediction-correction technique, then it is illustrated on a numerical example.

2.3.1. Prediction-correction method

A generalized solution $U \in E$ of (1.2) is constructed based on a prediction \tilde{U} which is corrected into \bar{U} , then sufficient conditions are exhibited which ensure that \bar{U} is solution.

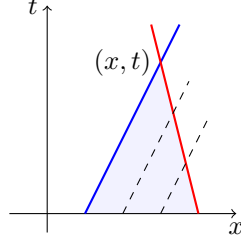


Fig. 3. Representation of a cone of dependence (pale blue) at a point (x, t) delimited by $C_{-\mu}(x, t)$ (red) and $C_{\nu}(x, t)$ (dark blue); and $C_{\nu}(\tilde{x}, \tilde{t})$ (dashed) for different $(\tilde{x}, \tilde{t}) \in C_{-\mu}(x, t)$.

Step 1: Define the function $\tilde{\Sigma} = (\tilde{\sigma}, \tilde{\zeta}) \in (L^{\infty}(\mathbb{R} \times \mathbb{R}^+))^2$ by

$$\tilde{\sigma}(x, t) = \sigma(U_0)(x + \mu t) = \begin{cases} a & \text{if } h(U_0)(x + \mu t) \leq 0, \\ b & \text{if } h(U_0)(x + \mu t) > 0, \end{cases}$$

$$\tilde{\zeta}(x, t) = \varsigma(U_0)(x - \nu t) = \begin{cases} c & \text{if } h(U_0)(x - \nu t) \leq 0, \\ d & \text{if } h(U_0)(x - \nu t) > 0. \end{cases}$$

The source $\tilde{\Sigma}$ can be interpreted as the source Σ at small time, before the switch indicator $h(U)$ has changed sign. Indeed, the source $\tilde{\sigma}$, resp. $\tilde{\zeta}$, has in (x, t) the value of σ , resp. ς , at the foot of the characteristics $C_{-\mu}(x, t)$, resp. $C_{\nu}(x, t)$.

Step 2: Solve

$$\partial_t \tilde{u} - \mu \partial_x \tilde{u} = \tilde{\sigma}(x, t), \quad \tilde{u}(x, 0) = u_0(x), \quad (2.18a)$$

$$\partial_t \tilde{v} + \nu \partial_x \tilde{v} = \tilde{\zeta}(x, t), \quad \tilde{v}(x, 0) = v_0(x), \quad (2.18b)$$

or equivalently compute

$$\tilde{U} = R(\tilde{\Sigma}), \quad (2.18c)$$

where $\tilde{\Sigma} = (\tilde{\sigma}, \tilde{\zeta})$ is independent of the unknown $\tilde{U} = (\tilde{u}, \tilde{v})$. Similarly, \tilde{U} is a prediction of the generalized solution at small times. Indeed, consider that (x, t) is such that $h(U)$ has not changed sign along $C_{-\mu}(x, t)$ and $C_{\nu}(x, t)$, then \tilde{U} satisfies (2.15) at (x, t) and is a generalized solution.

Step 3: Solve

$$\partial_t \bar{u} - \mu \partial_x \bar{u} = \sigma(\tilde{U}(x, t)), \quad \bar{u}(x, 0) = u_0(x), \quad (2.19a)$$

$$\partial_t \bar{v} + \nu \partial_x \bar{v} = \varsigma(\tilde{U}(x, t)), \quad \bar{v}(x, 0) = v_0(x). \quad (2.19b)$$

or equivalently compute

$$\bar{U} = R \circ \Sigma(\tilde{U}). \quad (2.19c)$$

This is a new decoupled system over $\bar{U} = (\bar{u}, \bar{v})$ where the source is computed using $\tilde{U} = (\tilde{u}, \tilde{v})$. Here, \bar{U} corresponds to a correction of (2.18) at larger times.

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There remains to check whether \bar{U} is a generalized solution to (1.2):

Step 4: Check that

$$\text{sign}\left(h(\tilde{U})\right) = \text{sign}\left(h(\bar{U})\right). \quad (2.20)$$

By definition, this implies that $Z^\pm(\tilde{U}) = Z^\pm(\bar{U})$, and especially this leads to

$$\Sigma(\tilde{U}) = \Sigma(\bar{U}).$$

Re-injecting it in (2.19) reads the original system (1.2) and thus \bar{U} is a generalized solution.

In fact, this last step holds only under conditions over the data. Some sufficient conditions are provided in Section 3.

The uniqueness of this generalized solution is also studied *a posteriori*, once \bar{U} is constructed.

Remark 2.15. An alternative to this approach based on a fixed point method is proposed in [6] with smooth source terms $\Sigma = (\sigma, \varsigma)$. It consists in solving iteratively

$$\partial_t u^{i+1} - \mu \partial_x u^{i+1} = \sigma(u^i, v^i), \quad u^{i+1}(0, x) = u_0(x), \quad (2.21a)$$

$$\partial_t v^{i+1} + \nu \partial_x v^{i+1} = \varsigma(u^i, v^i), \quad v^{i+1}(0, x) = v_0(x), \quad (2.21b)$$

or equivalently computing $U^{i+1} = R \circ \Sigma(U^i)$, until convergence (if convergence occurs). The present method consists in providing $U^0 = \tilde{U}$ (in Step 1) in this fixed point such that it converges at the first iteration, i.e. $U^1 = U^2 = \dots = U^\infty$.

2.3.2. Example with numerical values

Due to the technicality of the computations in the proofs below, the steps of the last subsection are first illustrated on a numerical example. The data are chosen to satisfy the conditions described in the case of Section 3.2 below, i.e. of Theorem 3.3.

Consider the Cauchy problem (1.2) with

$$\begin{aligned} h_1 = h_2 = \mu = \nu = 1, \quad u_0(x) = -10, \quad v_0(x) = x, \\ a = 10, \quad b = 20, \quad c = 30, \quad d = 40, \end{aligned}$$

with these data, (1.2) reads

$$\begin{aligned} \partial_t u - \partial_x u = \begin{cases} 10 & \text{if } u + v \leq 0, \\ 20 & \text{if } u + v > 0, \end{cases} \quad \partial_t v + \partial_x v = \begin{cases} 30 & \text{if } u + v \leq 0, \\ 40 & \text{if } u + v > 0, \end{cases} \quad (2.22) \\ u_0(x) = -10, \quad v_0(x) = x. \end{aligned}$$

Step 1: Define for all $(x, t) \in \mathbb{R} \times \mathbb{R}^+$

$$\tilde{\sigma}(x, t) = \begin{cases} 10 & \text{if } -10 + x + t \leq 0, \\ 20 & \text{if } -10 + x + t > 0, \end{cases} \quad \tilde{\varsigma}(x, t) = \begin{cases} 30 & \text{if } -10 + x - t \leq 0, \\ 40 & \text{if } -10 + x - t > 0, \end{cases}$$

At initial time, $h(U_0)(x) = 0$ only for $x = 10$. Thus, $\tilde{\sigma}$ switches value along $C_{-1}(10)$ and $\tilde{\zeta}$ along $C_1(10)$ (see Fig. 4).

Step 2: Integrating these sources along the characteristics provides

$$\forall (x, t) \in \mathbb{R} \times \mathbb{R}^+, \quad \tilde{u}(x, t) = \begin{cases} -10 + 10t & \text{if } -10 + x + t \leq 0, \\ -10 + 20t & \text{if } -10 + x + t > 0, \end{cases}$$

$$\tilde{v}(x, t) = \begin{cases} x - t + 30t & \text{if } -10 + x - t \leq 0, \\ x - t + 40t & \text{if } -10 + x - t > 0. \end{cases}$$

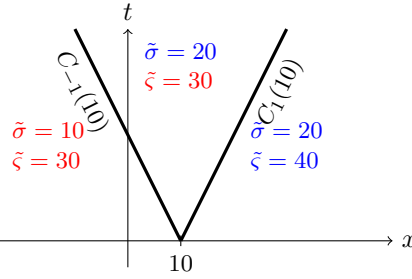


Fig. 4. Representation of the three zones where the source $\tilde{\Sigma}$ has different constant values.

Step 3: Computing $h(\tilde{U})$ yields

$$\forall (x, t) \in \mathbb{R} \times \mathbb{R}^+, \quad h(\tilde{u}, \tilde{v})(x, t) = \begin{cases} -10 + x + 39t & \text{if } x \leq 10 - t, \\ -10 + x + 49t & \text{if } 10 - t < x \leq 10 + t, \\ -10 + x + 59t & \text{otherwise} \end{cases}$$

Computing where $h(\tilde{U}(x, t)) = 0$ in $\mathbb{R} \times \mathbb{R}^+$ provides the line $\Gamma(\tilde{U})$

$$\Gamma(\tilde{U}) = \left\{ \left(x, \frac{10-x}{39} \right), \quad x \in]-\infty, 10] \right\}.$$

Then $h(\tilde{U})$ is positive above this line (for all $t > T(x) = \frac{10-x}{39}$ and also for all $x > 10$) and negative under. The sign of $h(\tilde{U})$ provides the value $\Sigma(\tilde{U})$, and integrating along the characteristics gives (see Fig. 5)

$$\bar{u}(x, t) = \begin{cases} -10 + 10t & \text{if } x \leq 10 - 39t & (I), \\ -10 + 20t & \text{if } x > 10 - t & (II), \\ -10 + 10t_s + 20(t - t_s) & \text{otherwise (III),} \end{cases} \quad t_s = \frac{-10 + x + t}{40}, \quad (2.23a)$$

$$\bar{v}(x, t) = \begin{cases} (x - t) + 30t & \text{if } x \leq 10 - 39t & (IV), \\ (x - t) + 40t & \text{if } x > 10 + t & (V), \\ (x - t) + 30\tau_s + 40(t - \tau_s) & \text{otherwise (VI),} \end{cases} \quad \tau_s = \frac{-10 + x + t}{38}. \quad (2.23b)$$

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Step 4: Computing $h(\bar{u}, \bar{v})$ reads for all $(x, t) \in \mathbb{R} \times \mathbb{R}^+$

$$h(\bar{u}, \bar{v})(x, t) = \begin{cases} -10 + x + 39t & \text{if } x \leq 10 - 39t & (I) \cap (IV), \\ -10 + x + 59t & \text{if } x > 10 + t & (II) \cap (V), \\ -10 + x + 59t - 10\tau_s & \text{if } 10 + t \geq x > 10 - t & (II) \cap (VI), \\ -10 + x + 59t - 10(t_s + \tau_s) & \text{otherwise} & (III) \cap (VI). \end{cases}$$

One verifies that $h(\bar{U})$ is negative for $x < 10 - 39t$, *i.e.* below $\Gamma(\tilde{U})$ and is positive for $x > \min(10 - 39t, 0)$, *i.e.* above and on the right of $\Gamma(\tilde{U})$. Thus $\bar{U} = (\bar{u}, \bar{v})$ is a generalized solution.

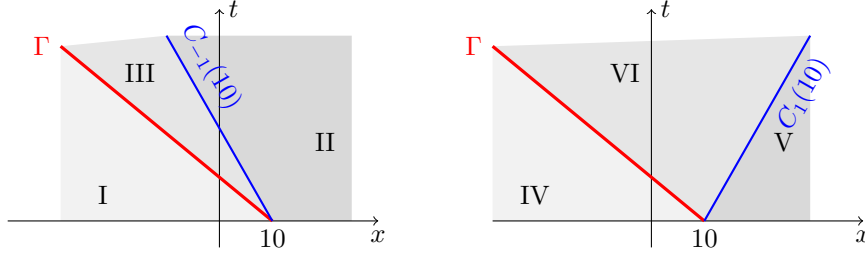


Fig. 5. Zones (I), (II) or (III) (resp. (IV), (V) or (VI)) in which \bar{u} (left), resp. \bar{v} (right), takes the value given in (2.23).

3. Main results

First, the notion of non-characteristic curve is presented. Then, conditions are exhibited to ensure that $\Gamma(\tilde{U})$ is non-characteristic, then we obtain existence and uniqueness of a generalized solution.

3.1. Non-characteristic curves

The expression *non-characteristic* is understood here, for a curve Γ , as the curve being non-tangent to $C_{-\mu}(x, t)$ and to $C_{\nu}(x, t)$ at all point $(x, t) \in \Gamma$. Therefore, the slopes of such tangents to Γ are compared to the velocities $-\mu$ and ν .

Two notions of non-characteristic curves are used. The following definitions are adapted from the literature (see e.g. [12,14,28]) to the present problem.

Definition 3.1 (Overcharacteristic curve). *The curve $G = \{(x, T(x)), x \in \mathbb{R}\}$ with $T \in W^{1,\infty}(\mathbb{R})$ is overcharacteristic if*

$$\frac{1}{-\mu} < \inf T', \quad \text{and} \quad \sup T' < \frac{1}{\nu}. \quad (3.1)$$

This imposes that G does not enter in the cone defined by the characteristics at (x, t) (see Fig. 6, left).

Imposing that $\Gamma(U)$ is overcharacteristic corresponds to imposing that the boiling front propagates faster than the flow, i.e. $X' < -\mu$ or $X' > \nu$.

Definition 3.2 (Subcharacteristic curve). *The curve $G = \{(X(t), t), t \in \mathbb{R}^+\}$ with $X \in W^{1,\infty}(\mathbb{R}^+)$ is subcharacteristic if*

$$-\mu < \inf X', \quad \text{and} \quad \sup X' < \nu. \quad (3.2)$$

This imposes that the curve G stays inside the cone defined by the characteristics at (x, t) (see Fig. 6, right).

Imposing that $\Gamma(U)$ is subcharacteristic corresponds to imposing that the boiling front propagates with a speed between $-\mu$ and ν .

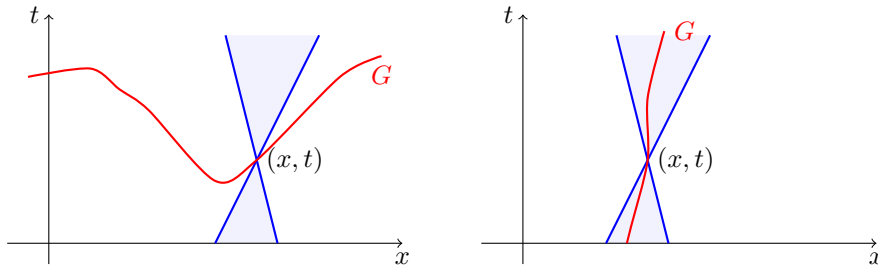


Fig. 6. Examples of an overcharacteristic (3.1; left) and subcharacteristic (3.2; right) curve G in red, and $C_{-\mu}(x, t)$ and $C_{\nu}(x, t)$ passing at $(x, t) \in G$ in blue.

The requirements in the next paragraphs are obtained (after computation) by imposing that $\Gamma(\tilde{U})$ is non-characteristic, where \tilde{U} is the small time solution given in Section (2.3.1). The cases where $\Gamma(\tilde{U})$ is tangent to a characteristic are rejected as it might affect the existence or the uniqueness of the solution for similar reasons as the case $ab \leq 0$ was rejected in the ODE framework of Section 2.1. Indeed, assuming e.g. that $\Gamma(\tilde{U})$ is tangent to $C_{-\mu}(x, t)$ (on a non-empty time interval), then fixing the value of Σ in (x, t) to (a, c) or to (b, d) could provide two different U satisfying (2.15).

Theorems 3.3, 3.5 and 3.7 below could be proved using e.g. [5,8]. These theorems are presented in a simplified form for the sake of readability, but more general versions of these well-posedness results are given in Lemmas 4.2, 5.2 and 6.1 which do not fit within the scope of the above cited literature. Furthermore, the present approach provides some additional notions about the solution and its boiling line $\Gamma(U)$.

3.2. Overcharacteristic case

Under some specified requirements, (1.2) has a unique generalized solution.

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Theorem 3.3. *Suppose that $a \leq b$ and $c \leq d$. Then, for all $(u_0, v_0) \in (W^{1,\infty}(\mathbb{R}))^2$ satisfying*

$$\inf u'_0 > -\frac{h(a, c)}{h_1(\mu + \nu)}, \quad \sup v'_0 < \frac{h(a, c)}{h_2(\mu + \nu)}, \quad (3.3)$$

there exists a unique generalized solution to (1.2).

This corresponds to the case when the sources in the two equations have higher values when $h(u, v) > 0$ than when $h(u, v) < 0$. Especially this holds for the application (1.1) we have in mind.

Remark 3.4. In Section 4, the boiling line $\Gamma(U)$ associated to this unique solution U is proved to be globally overcharacteristic under condition (3.3).

In practice, the solution in this case is $(W_{loc}^{1,\infty}(\mathbb{R} \times \mathbb{R}^+))^2$ which is a stronger regularity than the requirement to be in E .

3.3. Subcharacteristic case

Under some specified requirements, (1.2) has a unique generalized solution.

Theorem 3.5. *Suppose that $b \leq a$ and $d \leq c$. Then, for all $(u_0, v_0) \in (W^{1,\infty}(\mathbb{R}))^2$ satisfying*

$$\inf u'_0 > \frac{-h(b, d)}{h_1(\mu + \nu)}, \quad \inf v'_0 > \frac{h(a, c)}{h_2(\mu + \nu)}, \quad (3.4)$$

there exists a unique generalized solution to (1.2).

This corresponds to the case when the sources in the two equations have lower values when $h(u, v) > 0$ than when $h(u, v) < 0$.

Remark 3.6. In Section 5, the boiling line $\Gamma(U)$ associated to this unique solution U is proved to be subcharacteristic under condition (3.4). This happens e.g. for a steady solution since the boiling front is stationary with a corresponding boiling curve of the form $\Gamma = \{(X, t), t \in \mathbb{R}\}$ with constant X .

In practice, the solution in this case is also $(W_{loc}^{1,\infty}(\mathbb{R} \times \mathbb{R}^+))^2$.

3.4. Riemann problems

Riemann problems are often studied as they offer a simple setup involving discontinuities in the solution. Furthermore, approximating their solutions leads to the construction of the well-known Godunov-type schemes. Consider initial conditions of the form

$$u(0, x) = u_0(x) = \begin{cases} u_L^0 & \text{if } x \leq 0, \\ u_R^0 & \text{otherwise,} \end{cases} \quad v(0, x) = v_0(x) = \begin{cases} v_L^0 & \text{if } x \leq 0, \\ v_R^0 & \text{otherwise.} \end{cases} \quad (3.5)$$

Again, we specify some requirements ensuring existence and uniqueness of a generalized solution.

Theorem 3.7. *Suppose that*

$$a < b, \quad c < d \quad \text{and} \quad h(a, c) > 0. \quad (3.6)$$

Then, for all $(U_L^0, U_R^0) \in \mathbb{R}^4$, there exists a unique generalized solution to (1.2, 3.5).

Since a generalized solution $U = (u, v)$ is sought in E , then u is Lipschitz continuous along $C_{-\mu}(x, t)$ and v along $C_\nu(x, t)$ for all $(x, t) \in \mathbb{R} \times \mathbb{R}^+$. However, as for a Riemann problem without source term, u jumps across $C_{-\mu}(0)$ and v across $C_\nu(0)$. Thus the solution is constructed differently in the three subsets of $\mathbb{R} \times \mathbb{R}^+$ separated by $C_{-\mu}(0)$ and $C_\nu(0)$ (see Fig. 7).

$$D_L = \{(x, t) \in \mathbb{R} \times \mathbb{R}^+ \text{ s.t. } x \leq -\mu t\}, \quad (3.7a)$$

$$D_R = \{(x, t) \in \mathbb{R} \times \mathbb{R}^+ \text{ s.t. } \nu t < x\}, \quad (3.7b)$$

$$D^* = \{(x, t) \in \mathbb{R} \times \mathbb{R}^+ \text{ s.t. } -\mu t < x \leq \nu t\}. \quad (3.7c)$$

In Section 6, the generalized solution is constructed in the following way:

In D_L and D_R : The generalized solution $U = (u, v)$ depends only on time

$$\forall (x, t) \in D_L, \quad U(x, t) = U_L(t), \quad \forall (x, t) \in D_R, \quad U(x, t) = U_R(t).$$

Furthermore, U_L , and U_R are continuous piecewise affine functions of t . Thus, the restrictions of the boiling line in D_L and D_R are just horizontal (see Fig. 7)

$$\Gamma_L := \{(y, \tau_L), \quad y \in]-\infty, y_L]\} \subset D_L,$$

$$\Gamma_R := \{(y, \tau_R), \quad y \in [y_R, +\infty[)\} \subset D_R,$$

where τ_L , τ_R , y_L and y_R are constants to determine. By definition, such lines Γ_L and Γ_R are overcharacteristic.

In D^* : the solution $U = (u, v)$ is also continuous piecewise affine and the source jumps along a boiling line Γ^* of the form

$$\Gamma^* := \{(y, T(y)), \quad y \in [y_L^*, y_R^*]\} \subset D^*,$$

with a function T continuous piecewise affine over $[y_L^*, y_R^*]$ (to determine).

In practice, the function T is composed of at most three pieces such that Γ^* is overcharacteristic (see Fig. 7). The restriction on the source term leads to reject all cases where Γ^* is subcharacteristic.

Remark 3.8.

- The condition $h(a, c) > 0$ (or (6.1a) below) is imposed to ensure that the jump of the source term (σ, ς) does not impact the strict monotonicity of $h(u, v)$ along the characteristics inside D_L , D^* and D_R (in the spirit of Proposition 2.2).

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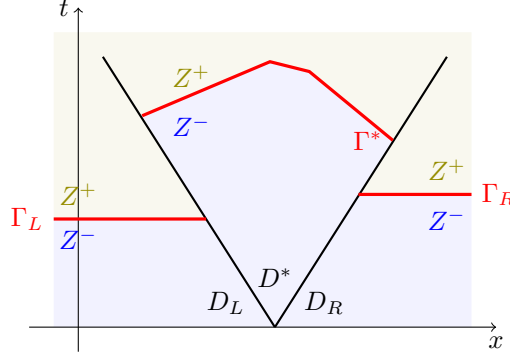


Fig. 7. Example of a boiling curve (in red) in the case of a Riemann problem in Theorem 6.1.

- The conditions $a < b$ and $c < d$ (or (6.1b) and (6.1c) below) are imposed to ensure that the boiling curve Γ^* in D^* is overcharacteristic.

4. Solution in the regular overcharacteristic case (Theorem 3.3)

The proof of Theorem 3.3 is based on Lemma 4.2 described in a first subsection. This lemma is proved in the remaining subsections using the technique of Section 2.3.1.

4.1. Notations and precise statement

In order to shorten the writings, the following notations are used.

Notation 4.1. We set

$$\begin{aligned} \sigma_0(x) &:= \sigma(U_0)(x), & \varsigma_0(x) &:= \varsigma(U_0)(x), \\ \Theta(y, z) &:= \frac{-[h_1 u'_0(y) + h_2 v'_0(z)]}{h_1(\mu u'_0 + \sigma_0)(y) + h_2(-\nu v'_0 + \varsigma_0)(z)}, \\ m &:= \inf_{(y,z) \in \mathbb{R}^2} \Theta(y, z), & M &:= \sup_{(y,z) \in \mathbb{R}^2} \Theta(y, z), \\ \theta &:= \frac{-(\mu + \nu)m}{1 - \nu m}, & \vartheta &:= \frac{(\mu + \nu)M}{1 + \mu M}. \end{aligned} \quad (4.1)$$

Lemma 4.2. Suppose that $U_0 \in (W^{1,\infty}(\mathbb{R}))^2$ is such that

$$-\frac{1}{\mu} < m, \quad M < \frac{1}{\nu}, \quad (4.2a)$$

$$h_1(\mu + \nu) \sup_x (-u'_0(x)) < h_1 \min(a, b, b + (a - b)\theta) + h_2 \min(c, d), \quad (4.2b)$$

$$h_2(\mu + \nu) \sup_x (v'_0(x)) < h_1 \min(a, b) + h_2 \min(c, d, d + (c - d)\vartheta). \quad (4.2c)$$

Then, there exists a unique generalized solution to (1.2).

The method of Section 2.3.2 is followed to prove Lemma 4.2 in the next subsections: first, a function $\bar{U} \in E$ is constructed; second, it is proved to be a generalized solution to (1.2) under the conditions (4.2); third, it is proved to be the only generalized solution.

Remark 4.3. The coefficients m and M are lower and upper bounds of the derivative T' of the function T involved in the boiling line $\Gamma(\bar{U})$, where \bar{U} is defined in (2.18).

Assuming Lemma 4.2 holds, then Theorem 3.3 is proved:

Proof of Theorem 3.3. First, let us verify that m and M are well-defined under condition (3.3). Write

$$Den(y, z) := h_1(\mu u'_0 + \sigma_0)(y) + h_2(-\nu v'_0 + \varsigma_0)(z), \quad Num(y, z) := -(h_1 u'_0(y) + h_2 v'_0(z)),$$

such that $\Theta = \frac{Num}{Den}$. Injecting (3.3) in $Den(y, z)$ yields

$$\inf_{(y,z) \in \mathbb{R}^2} Den(y, z) > -h(a, c) + \inf_{(y,z) \in \mathbb{R}^2} h_1 \sigma_0(y) + h_2 \varsigma_0(z) \geq 0.$$

Now, fixing $\sigma_0 = a$ and $\varsigma_0 = c$ in Θ , one observes that the function

$$(X, Y) \mapsto \frac{-(h_1 X + h_2 Y)}{h_1(\mu X + a) + h_2(-\nu Y + c)},$$

is strictly monotonic w.r.t. $X \in \left] -\frac{h(a,c)}{h_1(\mu+\nu)}, +\infty \right[$ and to $Y \in \left] -\infty, \frac{h(a,c)}{h_2(\mu+\nu)} \right[$.

Thus, the extreme values of Θ can be found at the extreme values $\sup u'_0$, $\sup v'_0$, $\inf u'_0$ and $\inf v'_0$. Computing Θ with these values of u'_0 and v'_0 , one finds that (4.2a) holds. Similar computations hold with the other values of σ_0 and ς_0 .

Now, with these bounds on m and M , one obtains that $\theta \in \left] -\infty, 1 \right[$ and $\vartheta \in \left] -1, +\infty \right[$. With these bounds on θ and ϑ , one verifies that (4.2) holds. \square

4.2. Construction of the solution \bar{U}

Here, for writing purposes, the computation of \bar{U} is decomposed into three steps: 1-construct $\tilde{U} = (\tilde{u}, \tilde{v})$; 2-compute $Z^-(\tilde{U})$, $Z^+(\tilde{U})$ and $\Gamma(\tilde{U})$; 3-compute $\bar{U} = (\bar{u}, \bar{v})$.

Solution in short time \tilde{U}

As in Section 2.3.2, a source $\tilde{\Sigma} = (\tilde{\sigma}, \tilde{\varsigma})$ constant along the (respective) characteristics is constructed, i.e. $\tilde{\sigma}$ (respectively $\tilde{\varsigma}$) is constant along $C_{-\mu}(x)$ (resp. $C_{-\nu}(x)$) for all x , and equal to the value of σ (resp. ς) at the foot of the characteristic $C_{-\mu}(x, t)$ (resp. $C_{\nu}(x, t)$). This yields

$$\tilde{\sigma}(x, t) = \sigma_0(x + \mu t), \quad \tilde{\varsigma}(x, t) = \varsigma_0(x - \nu t).$$

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Then, a field $\tilde{U} = (\tilde{u}, \tilde{v}) = R(\tilde{\Sigma})$ is constructed

$$\tilde{u}(x, t) = u_0(x + \mu t) + \sigma_0(x + \mu t)t, \quad (4.3a)$$

$$\tilde{v}(x, t) = v_0(x - \nu t) + \varsigma_0(x - \nu t)t. \quad (4.3b)$$

Computation of the boiling curve $\Gamma(\tilde{U})$

With this \tilde{U} , we compute T , the graph of which is $\Gamma(\tilde{U})$, and verify that $\Gamma(\tilde{U})$ is overcharacteristic.

Lemma 4.4. *Suppose (4.2) holds. Then, there exists a unique $T \in W^{1,\infty}(\mathbb{R})$ such that for all $(x, t) \in \mathbb{R} \times \mathbb{R}^+$*

$$\begin{cases} h(\tilde{U})(x, t) < 0 & \text{if and only if} & t < T(x), \\ h(\tilde{U})(x, t) = 0 & \text{if and only if} & t = T(x), \\ h(\tilde{U})(x, t) > 0 & \text{otherwise.} \end{cases} \quad (4.4)$$

Furthermore, T satisfies a.e. in $x \in \mathbb{R}$

$$m \leq T'(x) \leq M. \quad (4.5)$$

Especially, T satisfies (3.1) and its graph $\Gamma(\tilde{U})$ is overcharacteristic.

Proof. This proof is in three steps: 1-compute all possible pointwise values of T considering σ and ς constant; 2-verify that this T satisfies (3.1); 3-compute $\Gamma(\tilde{U})$.

(1) Denote $\sigma^- = a$, $\sigma^+ = b$, $\varsigma^- = c$ and $\varsigma^+ = d$ and define

$$u^\pm(x, t) = u_0(x + \mu t) + t\sigma^\pm, \quad v^\pm(x, t) = v_0(x - \nu t) + t\varsigma^\pm.$$

Since $U_0 \in (W^{1,\infty}(\mathbb{R}))^2$ and since h is linear, $h(u^\star, v^\otimes) \in W^{1,\infty}(\mathbb{R} \times \mathbb{R}^+)$ for all $\star = \pm$ and $\otimes = \pm$. Using an implicit function theorem (see e.g. [10] for a weak version) provides

$$T^{\star,\otimes} \in W^{1,\infty}(\mathbb{R}), \text{ s.t. } h(u^\star, v^\otimes)(x, T^{\star,\otimes}(x)) = 0, \quad \star = \pm, \quad \otimes = \pm.$$

(2) Differentiating a.e. the function $x \mapsto h(u^\star, v^\otimes)(T^{\star,\otimes}(x), x)$ leads to

$$\begin{aligned} \frac{dT^{\star,\otimes}}{dx}(x) &= -\frac{\partial_x h(u^\star, v^\otimes)(x, T^{\star,\otimes}(x))}{\partial_t h(u^\star, v^\otimes)(x, T^{\star,\otimes}(x))} \\ &= \frac{-[h_1 u'_0(x + \mu T^{\star,\otimes}(x)) + h_2 v'_0(x - \nu T^{\star,\otimes}(x))]}{h_1 [\mu u'_0(x + \mu T^{\star,\otimes}(x)) + \sigma^\star] + [-\nu v'_0(x - \nu T^{\star,\otimes}(x)) + \varsigma^\otimes]}. \end{aligned} \quad (4.6)$$

The denominator in (4.6) is non-zero according to (4.2). As in the proof of Proposition 3.3,

$$(X, Y) \mapsto -\frac{h_1 X + h_2 Y}{h_1(\mu X + \sigma^\star) + h_2(-\nu Y + \varsigma^\otimes)}$$

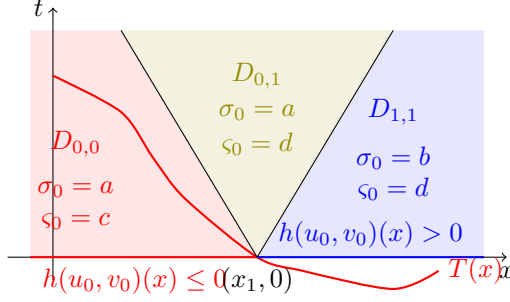


Fig. 8. Example of a function T such that $T(x_1) = 0$ and domains $D_{0,0}$, $D_{0,1}$ and $D_{1,1}$ separated by $C_{-\mu}(x_1)$ and $C_\nu(x_1)$.

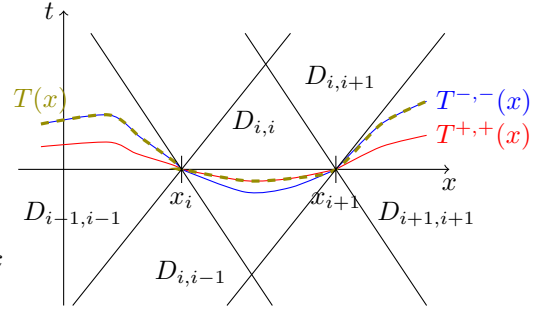


Fig. 9. Example of domains $D_{i,j}$ and functions T^{++} , T^{--} and T .

is strictly monotonic w.r.t. $X \in [\inf u'_0, \sup u'_0] \subset]\sigma^*/\mu, +\infty[$ and to $Y \in [\inf v'_0, \sup v'_0] \subset]-\infty, \varsigma^\otimes/\nu[$. With the bounds (4.2b) on u'_0 and (4.2c) on v'_0 , the derivatives (4.6) satisfy (4.5).

- (3) Suppose that the initial source $\Sigma_0 = (\sigma_0, \varsigma_0)$ does not jump, then Σ_0 is constant equal to (a, c) or (b, d) , and T is either $T^{+,+}$ or $T^{-,-}$, which both satisfy (4.5).

Suppose that Σ_0 jumps at the points $(x_i)_{i \in J}$ (potentially infinite numbered) such that $x_i < x_{i+1}$. By construction, $h(\tilde{U})$ is discontinuous across $C_{-\mu}(x_i)$ and $C_\nu(x_i)$ and $W^{1,\infty}$ elsewhere. Then, decompose the (x, t) -space into (see Fig. 8 and 9)

$$D_{i,j} = \left\{ \bigcup_{x \in]x_i, x_{i+1}[} C_{-\mu}(x) \right\} \cap \left\{ \bigcup_{x \in]x_j, x_{j+1}[} C_\nu(x) \right\}. \quad (4.7)$$

where $x_{\min J-1} = -\infty$ if $\min x_i > -\infty$ and $x_{\max J+1} = +\infty$ if $\max x_i < +\infty$.

Inside $D_{i,j}$, the source $\tilde{\Sigma} = (\tilde{\sigma}, \tilde{\varsigma})$ is constant, and $h(\tilde{U}) \in W^{1,\infty}(D_{i,j})$ (from (4.3)). Since the $T^{*,\otimes}$ satisfy (4.5), their graphs are overcharacteristic. And they pass at the points $(x_i, 0)$ for all $i \in J$. This provides for all $\star = \pm$ and $\otimes = \pm$

$$\{(x, T^{*,\otimes}(x)) \mid x \in \mathbb{R}\} \in \bigcup_i D_{i,i}.$$

In each $D_{i,i}$, the value of \tilde{U} is either U^+ or U^- , then T is also either $T^{+,+}$ or $T^{-,-}$ over $]x_i, x_{i+1}[$. They both satisfy (4.5) in each $]x_i, x_{i+1}[$ then T satisfies (4.5) over \mathbb{R} . And $h(\tilde{U}) \neq 0$ for all $(x, t) \in D_{i,j}$ for $i \neq j$.

Furthermore, using (4.2b) and (4.2c), $h(u^*, v^\otimes)$ is strictly increasing along $C_{-\mu}(x)$ and $C_\nu(x)$ for all $x \in \mathbb{R}$. Thus it is negative for all $t < T^{*,\otimes}(x)$ and positive for all $t > T^{*,\otimes}(x)$. One deduces that $h(\tilde{U})$ has opposite signs on each side of $\Gamma(\tilde{U})$. \square

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Correction of the solution for long time \bar{U}

Computing $\bar{U} = R(\Sigma(\bar{U}))$ from (2.16,2.12) gives

$$\bar{u}(x, t) = u_0(x + \mu t) + a\mathcal{T}(x, t) + b(t - \mathcal{T}(x, t)), \quad (4.8a)$$

$$\bar{v}(x, t) = v_0(x - \nu t) + c\mathfrak{T}(x, t) + d(t - \mathfrak{T}(x, t)), \quad (4.8b)$$

$$\mathcal{T}(x, t) = \int_0^t \mathbf{1}_{\mathbb{R}^-} (h(\tilde{u}, \tilde{v})(x + \mu(t - \tau), \tau)) d\tau,$$

$$\mathfrak{T}(x, t) = \int_0^t \mathbf{1}_{\mathbb{R}^-} (h(\tilde{u}, \tilde{x})(x - \nu(t - \tau), \tau)) d\tau.$$

where the integrals can be computed explicitly

$$\mathcal{T}(x, t) = \begin{cases} 0 & \text{if } \mathfrak{t}(x + \mu t) < 0, \\ t & \text{if } 0 < t < \mathfrak{t}(x + \mu t), \\ \mathfrak{t}(x + \mu t) & \text{otherwise,} \end{cases} \quad (4.9a)$$

$$\mathfrak{T}(x, t) = \begin{cases} 0 & \text{if } \tau(x - \nu t) < 0, \\ t & \text{if } 0 < t < \tau(x - \nu t), \\ \tau(x - \nu t) & \text{otherwise.} \end{cases} \quad (4.9b)$$

Here, $\mathfrak{t}(x)$, resp. $\tau(x)$, is the time of intersection between $C_{-\mu}(x)$, resp. $C_\nu(x)$, and $\Gamma(\bar{U})$. It yields

$$\mathfrak{t}(x) = T((Id + \mu T)^{-1}(x)), \quad \tau(x) = T((Id - \nu T)^{-1}(x)). \quad (4.9c)$$

Remark that if T' is bounded by (3.1), then $Id + \mu T$ and $Id - \nu T$ are strictly monotonic, thus bijective, and (4.9c) is well-defined.

For $\bar{U} = (\bar{u}, \bar{v})$ to be a generalized solution, there remains to prove that it satisfies (2.15).

Lemma 4.5. *Suppose that (4.2) holds. Then, $\bar{U} = (\bar{u}, \bar{v})$ given by (4.8-4.9) satisfies*

$$(\partial_t - \mu \partial_x)h(\bar{U}) > 0, \quad (\partial_t + \nu \partial_x)h(\bar{U}) > 0 \quad \text{a.e. in } (x, t) \in \mathbb{R} \times \mathbb{R}^+.$$

Proof. Compute

$$(\partial_t - \mu \partial_x)h(\bar{U})(x, t) = h_1(\partial_t - \mu \partial_x)\bar{u}(x, t) + h_2(\partial_t - \mu \partial_x)\bar{v}(x, t), \quad (4.10a)$$

$$(\partial_t + \nu \partial_x)h(\bar{U})(x, t) = h_1(\partial_t + \nu \partial_x)\bar{u}(x, t) + h_2(\partial_t + \nu \partial_x)\bar{v}(x, t), \quad (4.10b)$$

where

$$(\partial_t - \mu \partial_x)\bar{u}(x, t) = \begin{cases} a & \text{if } (x, t) \in Z^-(\bar{U}), \\ b & \text{if } (x, t) \in Z^+(\bar{U}), \end{cases} \quad (4.11a)$$

$$(\partial_t + \nu \partial_x)\bar{v}(x, t) = \begin{cases} c & \text{if } (x, t) \in Z^-(\bar{U}), \\ d & \text{if } (x, t) \in Z^+(\bar{U}). \end{cases} \quad (4.11b)$$

Note that all the identities hold a.e. Using (4.9), the other terms read

$$(\partial_t - \mu \partial_x)\bar{v}(x, t) = -(\mu + \nu)v'_0(x - \nu t) + d + (c - d)(\partial_t - \mu \partial_x)\mathfrak{T}(x, t), \quad (4.12a)$$

$$(\partial_t + \nu \partial_x)\bar{u}(x, t) = (\mu + \nu)u'_0(x + \mu t) + b + (a - b)(\partial_t + \nu \partial_x)\mathcal{T}(x, t), \quad (4.12b)$$

$$\begin{aligned}
(\partial_t + \nu\partial_x)\mathcal{T}(x, t) &= \begin{cases} 0 & \text{if } \mathfrak{t}(x + \mu t) < 0, \\ 1 & \text{if } 0 < t < \mathfrak{t}(x + \mu t), \\ (\mu + \nu)\mathfrak{t}'(x + \mu t) & \text{otherwise.} \end{cases} \\
(\partial_t - \mu\partial_x)\mathfrak{T}(x, t) &= \begin{cases} 0 & \text{if } \tau(x - \nu t) < 0, \\ 1 & \text{if } 0 < t < \tau(x - \nu t), \\ -(\mu + \nu)\tau'(x - \nu t) & \text{otherwise,} \end{cases} \\
\mathfrak{t}'(x) &= \frac{T'}{1 + \mu T'} ((Id + \mu T)^{-1}(x)), & \tau'(x) &= \frac{T'}{1 - \nu T'} ((Id - \nu T)^{-1}(x)), \\
&= \frac{1}{\frac{1}{T'} + \mu} ((Id + \mu T)^{-1}(x)), & &= \frac{1}{\frac{1}{T'} - \nu} ((Id - \nu T)^{-1}(x)),
\end{aligned}$$

From (4.5), $1 + \mu T' > 0$ and $1 - \nu T' > 0$ a.e. Then, where $T'(x) = 0$, one has $\mathfrak{t}'(x) = 0$ and $\tau'(x) = 0$, elsewhere

$$\begin{aligned}
\mathfrak{t}'(x) &= \frac{1}{\frac{1}{T'} + \mu} ((Id + \mu T)^{-1}(x)) \geq \frac{1}{\frac{1}{\sup T'} + \mu} \geq \frac{M}{1 + \mu M} \quad \text{a.e.} \\
\tau'(x) &= \frac{1}{\frac{1}{T'} - \nu} ((Id - \nu T)^{-1}(x)) \leq \frac{1}{\frac{1}{\inf T'} - \nu} \leq \frac{m}{1 - \nu m} \quad \text{a.e.}
\end{aligned}$$

Injecting these estimates in (4.10) leads to

$$\begin{aligned}
(\partial_t - \mu\partial_x)h(\bar{u}, \bar{v}) &\geq h_1 \min(a, b) + h_2 [-(\mu + \nu) \sup v'_0 + \min(d, c, d + (c - d)\vartheta)], \\
(\partial_t + \nu\partial_x)h(\bar{u}, \bar{v}) &\geq h_1 [(\mu + \nu) \inf u'_0 + \min(b, a, b + (a - b)\vartheta)] + h_2 \min(c, d),
\end{aligned}$$

the right-hand sides are strictly positive according to (4.2b) and (4.2c), which gives the result. \square

From this lemma, one deduces that (2.20) holds and thus that \bar{U} is a generalized solution.

Lemma 4.6. *Suppose that (4.2) holds. Then, $\tilde{U} = (\tilde{u}, \tilde{v})$ given by (4.3) and $\bar{U} = (\bar{u}, \bar{v})$ given by (4.8-4.9) satisfy*

$$\text{sign}(h(\tilde{U})) = \text{sign}(h(\bar{U})).$$

Proof. Consider first a sub-domain $D_{i,i}$ (from (4.7)), such that $h(U_0)(x) < 0$ for all $(x, t) \in D_{i,i}$. In this set, $\text{sign}(h(\tilde{U}))$ is given by (4.4).

Consider $(x, t) \in D_{i,i}$ such that $0 < t < T(x)$ (see Fig. 9). Since $\Gamma(\tilde{U})$ is overcharacteristic, we have

$$\forall \tau \in [0, T(x)[, \quad h(\tilde{U})(x + \mu\tau, t - \tau) < 0 \quad \text{and} \quad h(\tilde{U})(x - \nu\tau, t - \tau) < 0.$$

Comparing (4.3) and (4.8) gives $\tilde{U}(x, t) = \bar{U}(x, t)$, and $\text{sign}(h(\tilde{U})(x, t)) = \text{sign}(h(\bar{U})(x, t))$.

Now, use Lemma 4.5 for all other points. That is first for all $(x, t) \in D_{i,i}$ such that $t > T(x)$, then for all $(x, t) \in D_{i,i}$ such that $h(u_0, v_0)(x) > 0$ and eventually for all $(x, t) \in D_{i,j}$ with $i < j$, one has $h(\bar{U})(x, t) > 0$. Especially, $\text{sign}(h(\tilde{U})(x, t)) = \text{sign}(h(\bar{U})(x, t))$. \square

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4.3. Uniqueness of the generalized solution

The proof of uniqueness mimics the one of Lemma 4.6, but with the constraint (4.2) on the generalized solution U instead of assuming that U is given by the construction (4.8-4.9).

First let us state a regularity result.

Lemma 4.7. *Consider $T \in W^{1,\infty}(]z, z + h[)$ such that for all $x \in]z, z + h[$,*

$$\frac{-1}{\mu} < \lim_{\substack{y \rightarrow x \\ y < x}} T'(y) < \frac{1}{\nu}, \quad \frac{-1}{\mu} < \lim_{\substack{y \rightarrow x \\ y > x}} T'(y) < \frac{1}{\nu}, \quad (4.13)$$

and a generalized solution $U \in E$ such that (see Fig. 10)

$$\forall (x, t) \in]z, z + h[\times]0, T(x)], \quad U(x, t) = \tilde{U}(x, t) \quad \text{and} \quad h(U)(x, t) \leq 0.$$

Then, $h(U)$ is continuous at $(x, T(x))$ for all $x \in]z, z + h[$.

Suppose furthermore that (4.2) holds and

$$m \leq T'(x) \leq M \quad \text{for almost all } x \in \mathbb{R}, \quad (4.14)$$

where m and M are defined in (4.1). Then, for all $x \in]z, z + h[$ and all $\epsilon > 0$ (small enough)

$$\begin{aligned} h(U)(x - \mu\epsilon, T(x) + \epsilon) &> h(U)(x, T(x)), \\ h(U)(x + \nu\epsilon, T(x) + \epsilon) &> h(U)(x, T(x)). \end{aligned}$$

Remark 4.8. Since $T \in W^{1,\infty}(]z, z + h[)$, the limits in (4.13) are well-defined. Furthermore, the requirement (4.13) is weaker than (3.1) as the sup of T' could equal $\frac{1}{\nu}$ in $z + h$.

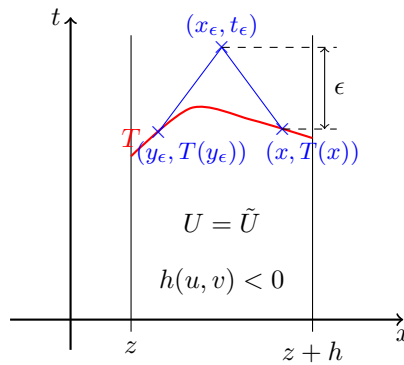


Fig. 10. Configuration for Lemma 4.7.

Proof. By hypothesis, $Id - \nu T$ and $Id + \mu T$ are bijective over $]z, z + h[$.

Consider $x \in [z, z + h]$ and $\epsilon > 0$ small enough. Define $t_\epsilon = T(x) + \epsilon$ and $(x_\epsilon, y_\epsilon) \in]z, z + h]^2$ such that $(x_\epsilon, t_\epsilon) \in C_\nu(y_\epsilon, T(y_\epsilon)) \cap C_{-\mu}(x, T(x))$. This reads

$$\begin{cases} x_\epsilon = x - \mu\epsilon = y_\epsilon + \nu\epsilon, \\ t_\epsilon = T(x) + \epsilon = T(y_\epsilon) + \epsilon, \end{cases} \Rightarrow y_\epsilon = (Id - \nu T)^{-1}((Id - \nu T)(x) - (\mu + \nu)\epsilon). \quad (4.15)$$

Thus, we impose the bound $0 < \epsilon < \frac{(z-x) + \nu T(x)}{\mu + \nu}$ on ϵ in order that $y_\epsilon > z$.

Compute the integrals along $C_{-\mu}(x_\epsilon, t_\epsilon)$,

$$\begin{aligned} u(x_\epsilon, t_\epsilon) &= u(x, T(x)) + \int_{T(x)}^{t_\epsilon} \sigma(u, v)(x - \mu\tau, \tau) d\tau \\ &= \tilde{u}(x, T(x)) + \check{\sigma}\epsilon, \end{aligned}$$

where

$$\check{\sigma} := \frac{1}{\epsilon} \int_{T(x)}^{T(x)+\epsilon} \sigma(u, v)(x - \mu\tau, \tau) d\tau \in [a, b].$$

Similarly, along $C_\nu(x_\epsilon, t_\epsilon)$,

$$\begin{aligned} v(x_\epsilon, t_\epsilon) &= v(y_\epsilon, T(y_\epsilon)) + \int_{T(y_\epsilon)}^{T(x)+\epsilon} \varsigma(u, v)(y_\epsilon + \nu\tau, \tau) d\tau \\ &= \tilde{v}(y_\epsilon, T(y_\epsilon)) + \check{\zeta}\epsilon + \hat{\zeta} \int_{y_\epsilon}^x T'(\xi) d\xi, \end{aligned}$$

where

$$\check{\zeta} := \frac{1}{\epsilon} \int_{T(x)}^{T(x)+\epsilon} \varsigma(u, v)(y_\epsilon + \nu\tau, \tau) d\tau, \quad \hat{\zeta} := \frac{1}{T(x) - T(y_\epsilon)} \int_{T(y_\epsilon)}^{T(x)} \varsigma(u, v)(y_\epsilon + \nu\tau, \tau) d\tau,$$

such that $(\check{\zeta}, \hat{\zeta}) \in [c, d]^2$. Here, one rewrites

$$\begin{aligned} \tilde{v}(y_\epsilon, T(y_\epsilon)) &= \tilde{v}(x, T(x)) - \int_{y_\epsilon}^x \frac{d}{d\xi} \tilde{v}(\xi, T(\xi)) d\xi \\ &= \tilde{v}(x, T(x)) - \int_{y_\epsilon}^x [(1 - \nu T'(\xi))v'_0(\xi - \nu T(\xi)) + cT'(\xi)] d\xi. \end{aligned}$$

This leads to

$$\begin{aligned} D_\epsilon &:= \frac{h(u, v)(x_\epsilon, t_\epsilon) - h(u, v)(x, T(x))}{\epsilon} \\ &= h(\check{\sigma}, \check{\zeta}) - \frac{h_2}{\epsilon} \int_{y_\epsilon}^x [(1 - \nu T'(\xi))v'_0(\xi - \nu T(\xi)) + (c - \hat{\zeta})T'(\xi)] d\xi. \end{aligned} \quad (4.16)$$

Then, one obtains from (4.15)

$$\begin{aligned} \lim_{\epsilon \rightarrow 0} D_\epsilon &= h(\check{\sigma}, \check{\zeta}) - \frac{h_2(\mu + \nu)}{1 - \nu T'(x^-)} [(1 - \nu T'(x^-))v'_0(x^- - \nu T(x^-)) + (c - \hat{\zeta})T'(x^-)] \\ &= h(\check{\sigma}, \check{\zeta}) - h_2(\mu + \nu) \left[v'_0(x^- - \nu T(x^-)) + (c - \hat{\zeta}) \frac{T'(x^-)}{1 - \nu T'(x^-)} \right], \end{aligned}$$

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where $f(x^-) = \lim_{\substack{y \rightarrow x \\ y < x}} f(y)$. By hypothesis, these derivatives exist a.e. in $x \in]z, z + h[$.

Thus these limits exist, and $1 - \nu T'(x^-) > 0$. Especially, $\lim_{\epsilon \rightarrow 0} D_\epsilon$ exists and is finite. Thus $h(U)$ is continuous at $(x, T(x))$ for all $x \in \mathbb{R}$.

Assume furthermore that (4.2) and (4.14) hold. Rewrite

$$\epsilon = \frac{1}{\mu + \nu} \int_{(Id - \nu T)(x) - (\mu + \nu)\epsilon}^{(Id - \nu T)(x)} dy = \frac{1}{\mu + \nu} \int_{y_\epsilon}^x (1 - \nu T')(\xi) d\xi.$$

Then

$$\begin{aligned} D_\epsilon &= \int_{y_\epsilon}^x \frac{(1 - \nu T')(\xi)}{\mu + \nu} \left[h(\bar{\sigma}, \bar{\zeta}) - h_2(\mu + \nu) \left(v'_0(\xi - T(\xi)) + (c - \bar{\zeta}) \frac{T'}{1 - \nu T'}(\xi) \right) \right] d\xi \\ &\geq \int_{y_\epsilon}^x \frac{(1 - \nu T')(\xi)}{\mu + \nu} h(\bar{\sigma}, \bar{\zeta} - (\mu + \nu) \sup v'_0 + (c - \bar{\zeta})\vartheta) d\xi > 0. \end{aligned}$$

Similar computations hold on the other characteristics. \square

Definition 4.9. Following [6], define the cone of dependence (see Fig. 3)

$$C(x, t) := \{(y, \tau) \in \mathbb{R}^+ \times \mathbb{R} \quad \text{s.t.} \quad \tau \in [0, t] \text{ and } y \in [x - \nu(t - \tau), x + \mu(t - \tau)]\}.$$

Considering a generalized solution $U = (u, v) \in E$, define (see Fig. 11)

$$D^-(U) := \{(x, t) \in \mathbb{R} \times \mathbb{R}^+ \quad \text{s.t.} \quad h(u, v)(y, \tau) \leq 0 \quad \text{a.e. in } (y, \tau) \in C(x, t)\} \subset Z^-(U). \quad (4.17)$$

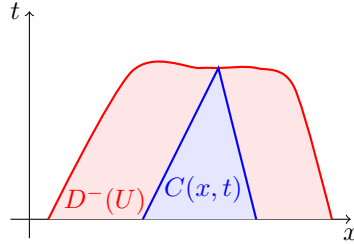


Fig. 11. Example of a set $D^-(U)$ (pale red) and a cone of dependence $C(x, t)$ (pale blue) for a point $(x, t) \in D^-(U)$.

Now, let us show that U is equal \bar{U} in $Z^-(\bar{U})$.

Lemma 4.10. Suppose that (4.2) holds and consider a generalized solution $U \in E$. Then,

$$D^-(U) = Z^-(\bar{U}).$$

Especially, $U(x, t) = \bar{U}(x, t)$ for all $(x, t) \in Z^-(\bar{U})$.

Proof. Any generalized solution $U = (u, v) \in E$ is of the form (2.15) which rewrites (4.8) for some functions $(\mathcal{T}, \mathfrak{T})$. Since $U_0 \in (W^{1,\infty}(\mathbb{R}))^2$ and $U \in E$, we obtain that $(\mathcal{T}, \mathfrak{T}) \in E$.

Since U satisfies (2.15), then $U(x, t) = \tilde{U}(x, t)$ for almost all $(x, t) \in D^-(U)$. Thus

$$D^-(U) \subset Z^-(\tilde{U}). \quad (4.18)$$

By construction,

$$D^-(U) = \{(x, t) \in \mathbb{R} \times \mathbb{R}^+ \text{ s.t. } x \in [x_{2i}, x_{2i+1}] \text{ and } t \in [0, T_i(x)]\}, \quad (4.19)$$

where the x_j are such that $h(U_0)(x) \leq 0$ a.e. on $[x_{2i}, x_{2i+1}]$ and the $T_i \in W^{1,\infty}(\mathbb{R})$ are such that

$$0 \leq T_i(x) \quad \forall x \in [x_{2i}, x_{2i+1}] \quad \text{and} \quad \frac{-1}{\mu} \leq T_i'(x) \leq \frac{1}{\nu} \text{ for almost all } x \in [x_{2i}, x_{2i+1}],$$

and $(x, T_i(x) + \epsilon) \notin D^-(U)$ for all $\epsilon > 0$. According to (4.18), $T_i \leq T$ a.e.

First, the graph of T_i is shown to be overcharacteristic in a subset of $[x_{2i}, x_{2i+1}]$, then in the entire interval $[x_{2i}, x_{2i+1}]$, and finally it is shown to be $\Gamma(\tilde{U})$, i.e. $T_i(x) = T(x)$ for all $x \in [x_{2i}, x_{2i+1}]$.

- (1) Suppose first that $T_i'(x_{2i}) = \frac{1}{\nu}$ (see Fig. 12). Then, $T_i'(x_{2i}) > T'(x_{2i})$ and by continuity of T_i and T , there exists $\epsilon > 0$ such that $T_i(x) > T(x)$ over $x \in [x_{2i}, x_{2i} + \epsilon]$. By definition of T_i ,

$$\forall x \in [x_{2i}, x_{2i} + \epsilon], \quad (x, T_i(x)) \in C(x, T_i(x)) \subset D^-(\tilde{U}).$$

However, one verifies geometrically that

$$\forall x \in [x_{2i}, x_{2i} + \epsilon], \quad C(x, T_i(x)) \setminus Z^-(\tilde{U}) \neq \emptyset,$$

which contradicts (4.18). Then $T_i'(x_{2i}) < \frac{1}{\nu}$. Thus there exists $\epsilon > 0$ such that $0 < T_i'(x) < \frac{1}{\nu}$ for all $x \in [x_{2i}, x_{2i} + \epsilon]$. Similarly, one finds $0 > T_i'(x) > \frac{-1}{\mu}$ for all $x \in [x_{2i+1} - \epsilon, x_{2i+1}]$.

- (2) Suppose now that $T_i'(x) < \frac{1}{\nu}$ for all $x \in [x_{2i}, y[$ and $T_i'(y) = \frac{1}{\nu}$ (see Fig. 12). Especially, $Id - \nu T$ is bijective over $[x_{2i}, y]$.

Since T' cannot reach such a value and $T_i \leq T$, there exists $\epsilon > 0$ such that $T_i(x) < T(x)$ for all $x \in [y - \epsilon, y]$.

Lemma 4.7 provides the continuity of $h(U)$ along the graph of T_i . Especially

$$h(U)(x, T_i(x)) = h(\tilde{U})(x, T_i(x)) < 0 \quad \forall x \in [y - \epsilon, y].$$

By continuity, one can find a neighbourhood of $(x, T_i(x))$ on which $h(U)$ is strictly negative (in gray on Fig. 12). Inside this neighbourhood, one may find a point (z, τ) with $\tau > T_i(z)$ such that $h(U)$ is strictly negative over all $C(z, \tau)$. This implies that $C(z, \tau) \subset D^-(U)$, which contradicts (4.19).

Thus, $T_i'(x) < \frac{1}{\nu}$ for all $x \in [x_{2i}, x_{2i+1}]$. Similar computations give $T_i'(x) > \frac{-1}{\mu}$ for all $x \in [x_{2i}, x_{2i+1}]$ and thus T_i has an overcharacteristic graph.

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- (3) Finally, suppose that there exists $y \in [x_{2i}, x_{2i+1}]$ such that $T_i(x) < T(x)$. Applying again Lemma 4.7 along the graph of T_i provides the existence of a neighbourhood of $(y, T_i(y))$ on which $h(u, v) < 0$. This neighbourhood contains $C(z, \tau) \setminus D^-(U)$ for some (z, τ) satisfying $\tau > T_i(z)$ which contradicts the definition of $D^-(U)$. Thus $T(x) = T_i(x)$ for all $x \in [x_{2i}, x_{2i+1}]$. \square

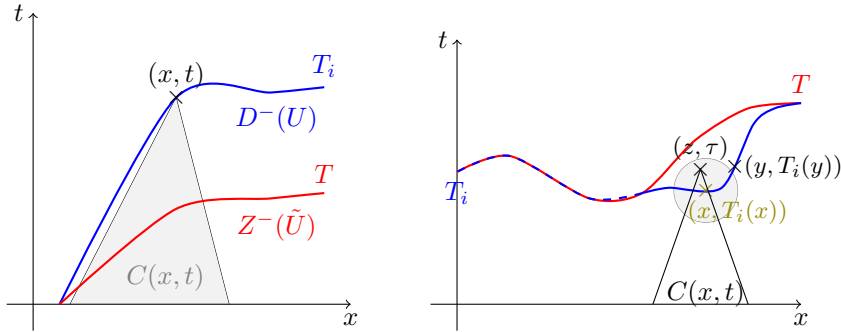


Fig. 12. Configuration at steps 1 (left) and 2 (right) of the proof of Lemma 4.10.

Finally, uniqueness is obtained in all the (x, t) -space: from Lemma 4.10, \bar{U} is the unique generalized solution in $Z^-(\tilde{U})$. Furthermore,

$$Z^-(\tilde{U}) = \{(x, t) \in \mathbb{R} \times \mathbb{R}^+ \text{ s.t. } t \leq T(x)\},$$

where the graph of $T \in W^{1,\infty}(\mathbb{R})$ is overcharacteristic. Then Lemma 4.7 provides the positivity of $h(U)$ for all $(x, t) \notin Z^-(\tilde{U})$, and one identifies \bar{U} to U through (4.8).

One easily verifies that the unique generalized solution \bar{U} depends continuously on a, b, c, d and U_0 inside the set defined by (4.2). Especially, the problem (2.15) is well-posed in the sense of Hadamard under these constraints.

5. Solution in the regular subcharacteristic case (Theorem 3.5)

Again, the proof of Theorem 3.5 is based on Lemma 5.2 below.

5.1. Notations and precise statement

Let us introduce notations similar to those of the preceding section.

Notation 5.1. Define

$$\Xi(y, z) := \frac{-[h_1(\mu u'_0 + \sigma_0)(y) + h_2(-\nu v'_0 + \varsigma_0)(z)]}{h_1 u'_0(y) + h_2 v'_0(z)} = \frac{1}{\Theta(y, z)},$$

$$\begin{aligned} \mathfrak{m} &:= \inf_{(y,z) \in \mathbb{R}^2} \Xi(y, z), & \mathfrak{M} &:= \sup_{(y,z) \in \mathbb{R}^2} \Xi(y, z), \\ \theta &:= \frac{\mu + \nu}{\mu + \mathfrak{m}}, & \vartheta &:= \frac{\mu + \nu}{\nu - \mathfrak{M}}. \end{aligned} \quad (5.1)$$

Lemma 5.2. *Suppose that $U_0 \in (W^{1,\infty}(\mathbb{R}))^2$ is such that*

$$-\mu < \mathfrak{m}, \quad \mathfrak{M} < \nu, \quad (5.2a)$$

$$h_1(\mu + \nu) \sup_x (-u'_0) < h_1 \min(a, b, a + (b - a)\theta) + h_2 \min(c, d), \quad (5.2b)$$

$$h_2(\mu + \nu) \inf_x (v'_0) > h_1 \max(a, b) + h_2 \max(c, d, d + (c - d)\vartheta). \quad (5.2c)$$

Then, there exists a unique generalized solution $U \in E$ to (1.2)

Remark 5.3. Again, the coefficients \mathfrak{m} and \mathfrak{M} correspond to the bounds on the derivative X' of the function X the graph of which is $\Gamma(\tilde{U})$.

Proof of Theorem 3.5. This is obtained by a straightforward adaptation of the proof of Theorem 3.3. \square

5.2. Construction of the generalized solution

The function $\tilde{U} = (\tilde{u}, \tilde{v})$ has the same definition (4.3) as in the last section.

Boiling curve $\Gamma(\tilde{U})$

One verifies that $\Gamma(\tilde{U})$ is subcharacteristic.

Lemma 5.4. *Suppose (5.2) holds. Then, there exists a unique $X \in W^{1,\infty}(\mathbb{R})$ such that for all $(x, t) \in \mathbb{R} \times \mathbb{R}^+$*

$$\begin{cases} h(\tilde{U})(x, t) < 0 & \Leftrightarrow x < X(t), \\ h(\tilde{U})(x, t) = 0 & \Leftrightarrow x = X(t), \\ h(\tilde{U})(x, t) > 0 & \Leftrightarrow x > X(t). \end{cases} \quad (5.3)$$

Furthermore, X satisfies for almost all $t \in \mathbb{R}^+$

$$\mathfrak{m} \leq X'(t) \leq \mathfrak{M}. \quad (5.4)$$

Especially, X satisfies (3.2) and its graph $\Gamma(\tilde{U})$ is subcharacteristic.

Remark 5.5. This case corresponds to the cases $T'(x) > \frac{1}{\nu}$ and $T'(x) < \frac{-1}{\mu}$ (see Fig. 13).

Proof. First, using the hypothesis (5.2), one computes

$$\begin{aligned} \partial_x h(U_0)(x) &= h_1 u'_0(x) + h_2 v'_0(x) \\ &> \frac{h_1(\max(a, b) - \min(a, b)) + h_2(\max(c, d) - \min(c, d))}{\mu + \nu} > 0. \end{aligned}$$

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Thus there exists one unique point where $h(U_0)$ changes sign.

Reproducing the computations of the proof of Lemma 4.4 (with the same notations) leads to

$$\begin{aligned} (\partial_t - \mu \partial_x)h(u^*, v^\otimes)(x, t) &\leq -h_2(\mu + \nu) \inf v'_0 + h_1 \max(a, b) + h_2 \max(c, d) < 0, \\ (\partial_t + \nu \partial_x)h(u^*, v^\otimes)(x, t) &\geq h_1(\mu + \nu) \inf u'_0 + h_1 \min(a, b) + h_2 \min(c, d) > 0. \end{aligned}$$

The implicit function theorem provides again the existence of functions $X^{*,\otimes} \in W^{1,\infty}(\mathbb{R}^+)$ such that $h(\bar{u}^*, \bar{v}^\otimes)(X^{*,\otimes}(t), t) = 0$. Differentiating it with respect to t leads to

$$X^{*,\otimes}'(t) = \frac{-[h_1(\mu u'_0(X^{*,\otimes}(t) + \mu t) + \sigma^*) + h_2(-\nu v'_0(X^{*,\otimes}(t) - \nu t) + \varsigma^\otimes)]}{h_1 u'_0(X^{*,\otimes}(t) + \mu t) + h_2 v'_0(X^{*,\otimes}(t) - \nu t)}. \quad (5.5)$$

The inf and sup values are obtained for the inf and sup values of u'_0 and v'_0 and these estimates show that $X^{*,\otimes}$ satisfies (5.4) for all $\star = \pm$ and $\otimes = \pm$.

Finally, one identifies $X = X^{-,+}$ matching the signs of h at $t = 0$. \square

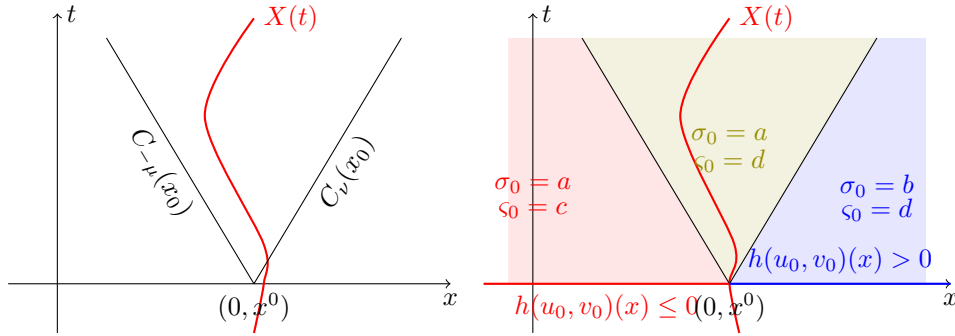


Fig. 13. Representation of $\Gamma(\bar{U})$ (left), and of different zones where $h(\bar{u}, \bar{v})$ is linear (right).

Correction of the solution for long time \bar{U}

From (4.4), $Z^+(\bar{U})$ is on the right of $Z^-(\bar{U})$, then \bar{U} is given by (4.8) where \mathcal{T} and \mathfrak{T} are

$$\mathcal{T}(x, t) = \begin{cases} t & \text{if } \mathfrak{t}(x + \mu t) < 0, \\ 0 & \text{if } 0 < t < \mathfrak{t}(x + \mu t), \\ t - \mathfrak{t}(x + \mu t) & \text{otherwise,} \end{cases} \quad (5.6a)$$

$$\mathfrak{T}(x, t) = \begin{cases} 0 & \text{if } \tau(x - \nu t) < 0, \\ t & \text{if } 0 < t < \tau(x - \nu t), \\ \tau(x - \nu t) & \text{otherwise,} \end{cases} \quad (5.6b)$$

where $\mathfrak{t}(x)$, resp. $\tau(x)$, is the time of intersection between $C_{-\mu}(x)$, resp. C_ν , and $\Gamma(\tilde{U})$. It reads

$$\mathfrak{t}(x) = (X + \mu Id)^{-1}(x), \quad \tau(x) = (X - \nu Id)^{-1}(x). \quad (5.6c)$$

Again, $h(\tilde{U})$ is shown to have the same sign as $h(\tilde{U})$ by exploiting its monotonicity along the characteristics.

Lemma 5.6. *Suppose that (5.2) holds. Then, $\bar{U} = (\bar{u}, \bar{v})$ given by (4.8,5.6) satisfies*

$$(\partial_t - \mu \partial_x)h(\bar{U}) < 0, \quad (\partial_t + \nu \partial_x)h(\bar{U}) > 0 \quad \text{a.e. in } (x, t) \in \mathbb{R} \times \mathbb{R}^+,$$

Proof. The formula (4.10) holds in the present case. Then,

$$(\partial_t - \mu \partial_x)\mathfrak{I}(x, t) = \begin{cases} 0 & \text{if } \tau(x - \nu t) < 0, \\ 1 & \text{if } 0 < t < \tau(x - \nu t), \\ -(\mu + \nu)\tau'(x - \nu t) & \text{otherwise,} \end{cases} \quad (5.7a)$$

$$(\partial_t + \nu \partial_x)\mathcal{T}(x, t) = \begin{cases} 1 & \text{if } \mathfrak{t}(x + \mu t) < 0, \\ 0 & \text{if } 0 < t < \mathfrak{t}(x + \mu t), \\ 1 - (\mu + \nu)\mathfrak{t}'(x + \mu t) & \text{otherwise.} \end{cases} \quad (5.7b)$$

Using (5.6), one has

$$\mathfrak{t}'(x) = \frac{1}{X'((X + \mu Id)^{-1}(x)) + \mu}, \quad \tau'(x) = \frac{1}{X'((X - \nu Id)^{-1}(x)) - \nu}.$$

re-injecting this in (4.10) and using the estimates (5.4) on X' and those (5.2b-5.2c) on U_0 provides the result. \square

5.3. Uniqueness of the generalized solution

Again there remains to identify the unique solution U to \bar{U} given by (4.8,5.6). This is obtained by adapting the proof of Section 5.3.

6. Solution of the Riemann problem (Theorem 3.7)

The proof of Theorem 3.7 is based on Lemma 6.1 below. The solution is easily found in the left domain $D_L = \{(x, t) \in \mathbb{R} \times \mathbb{R}^+ \text{ s.t. } x \leq -\mu t\}$ and in the right domain $D_R = \{(x, t) \in \mathbb{R} \times \mathbb{R}^+ \text{ s.t. } x > \nu t\}$. In the middle domain, its construction is adapted from the steps of Subsection 2.3.1.

6.1. Notations and precise statement

The following requirements are slightly less restrictive than (3.6) and sufficient to successfully construct a solution that provides a boiling line Γ^* overcharacteristic in the middle domain D_* .

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Lemma 6.1. *Suppose that*

$$\min_{\tilde{\sigma}=a,b} \min_{\tilde{\zeta}=c,d} h(\tilde{\sigma}, \tilde{\zeta}) > 0, \quad (6.1a)$$

$$\min_{\tilde{\sigma}=a,b} \frac{h(\tilde{\sigma}, d)}{h(\tilde{\sigma}, c)} \min_{\tilde{\zeta}=c,d} h(a, \tilde{\zeta}) > h_1(a - b), \quad (6.1b)$$

$$\min_{\tilde{\sigma}=a,b} h(\tilde{\sigma}, c) \min_{\tilde{\zeta}=c,d} \frac{h(b, \tilde{\zeta})}{h(a, \tilde{\zeta})} > h_2(c - d). \quad (6.1c)$$

Then, for all $(U_L^0, U_R^0) \in \mathbb{R}^4$, there exists a unique generalized solution to (1.2,3.5).

This result implies Theorem 3.7.

Proof of Theorem 3.7. Recall that $h_1 > 0$ and $h_2 > 0$.

If (3.6) holds, then

$$\min_{\tilde{\sigma}=a,b} \min_{\tilde{\zeta}=c,d} h(\tilde{\sigma}, \tilde{\zeta}) = h(a, c) > 0.$$

Furthermore, $h_1(a - b) < 0$ and $h_2(c - d) < 0$, while

$$\begin{aligned} \min_{\tilde{\sigma}=a,b} \frac{h(\tilde{\sigma}, d)}{h(\tilde{\sigma}, c)} \min_{\tilde{\zeta}=c,d} h(a, \tilde{\zeta}) &> \frac{h(a, d)}{h(b, c)} h(a, c) > 0 > h_1(a - b), \\ \min_{\tilde{\sigma}=a,b} h(\tilde{\sigma}, c) \min_{\tilde{\zeta}=c,d} \frac{h(b, \tilde{\zeta})}{h(a, \tilde{\zeta})} &> h(a, c) \frac{h(b, c)}{h(a, d)} > 0 > h_2(c - d). \end{aligned}$$

Thus, (6.1) holds and Lemma 6.1 provides the result. \square

The proof of Lemma 6.1 is restricted to the case where the Riemann data satisfy

$$h(u_L^0, v_L^0) < 0, \quad h(u_R^0, v_R^0) < 0, \quad h(u_R^0, v_L^0) < 0. \quad (6.2)$$

However it extends easily to general Riemann data. As illustrated below, (6.2) corresponds to the case where $h(U)$ changes sign in every domains D_L , D^* and D_R , especially

$$\begin{aligned} \sigma(u_L^0, v_L^0) &= \sigma(u_R^0, v_R^0) = \sigma(u_R^0, v_L^0) = a, \\ \varsigma(u_L^0, v_L^0) &= \varsigma(u_R^0, v_R^0) = \varsigma(u_R^0, v_L^0) = c. \end{aligned}$$

6.2. Solution in the left and right domains D_L and D_R

Denote $S = L, R$ the generic index associated to a side domain D_L or D_R . A space independent solution is computed in a straightforward way.

Lemma 6.2. *Suppose that $h(a, c)h(b, d) > 0$. Then, there exists a unique generalized solution $U = (u, v)$ over D_S . Furthermore, U does not depend on x and satisfies*

for all $(x, t) \in D_S$

$$u(x, t) = u_S(t) = \begin{cases} u_S^0 + at & \text{if } 0 < t < \tau_S, \\ u_S^0 + a\tau_S + b(t - \tau_S) & \text{if } \tau_S < t, \end{cases} \quad (6.3a)$$

$$v(x, t) = v_S(t) = \begin{cases} v_S^0 + ct & \text{if } 0 < t < \tau_S, \\ v_S^0 + c\tau_S + d(t - \tau_S) & \text{if } \tau_S < t. \end{cases}, \quad (6.3b)$$

$$\tau_S = \frac{-h(u_S^0, v_S^0)}{h(a, c)}, \quad (6.3c)$$

where u_S and v_S are defined by (6.4) below.

Proof. Define u_S and v_S such that

$$u'_S(t) = \sigma(u_S, v_S)(t), \quad v'_S(t) = \varsigma(u_S, v_S)(t), \quad u_S(0) = u_S^0, \quad v_S(0) = v_S^0. \quad (6.4)$$

One computes

$$\frac{d}{dt}h(U_S)(t) = \begin{cases} h(a, c) & \text{if } h(U_S)(t) \leq 0, \\ h(b, d) & \text{if } h(U_S)(t) > 0, \end{cases} \quad h(U_S)(0) = h(U_S^0),$$

which has a unique Carathodory solution according to Proposition 2.2 and (6.1a). One verifies that $(x, t) \mapsto U_S(t)$ is a generalized solution over D_S . The uniqueness is obtained by adapting the proof of Lemma 4.10. \square

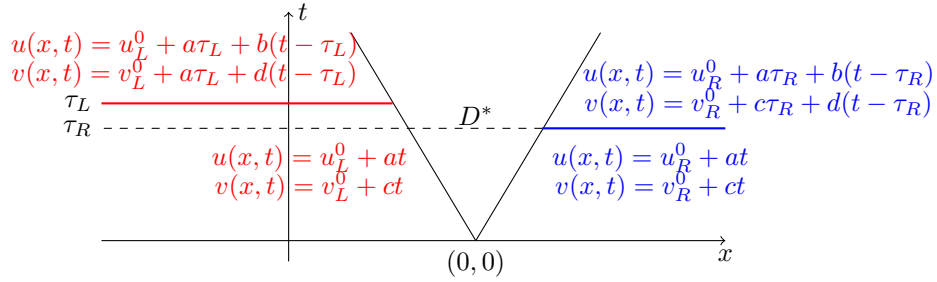


Fig. 14. Example of a solution $U = (u, v)$ in D_L and D_R .

6.3. Solution in the middle domain D^*

Now, the solution (6.3) is prolonged in the middle domain D^* . As illustrated on Fig. 1 (bottom) and in the numerical example of Section 6.5 below, $h(u, v)$ may switch sign several times along a characteristic curve, due to the discontinuities of u , resp. v , across $C_{-\mu}(0)$, resp. $C_\nu(0)$. For this purpose, and since U is already known in D_L and D_R , the construction of \tilde{U} is restricted inside D^* and is constructed using continuity of u along $C_\nu(0)$ and of v along $C_{-\mu}(0)$.

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The construction of the solution follows the same guidelines as before, but is adapted to take into account boundary conditions along $C_{-\mu}(0)$ and $C_\nu(0)$.

*Solution in short time \tilde{U} in D^**

Consider the short time solution $\tilde{U} = (\tilde{u}, \tilde{v})$ inside D^* . The solution \tilde{u} is continuous through $C_\nu(0)$ and along the characteristics $C_{-\mu}(x, t)$ the feet of which lie on $\{x > 0\}$. At such positions, $\tilde{u}(x, 0) = u_R^0$. Similarly, \tilde{v} is continuous through $C_{-\mu}(0)$ and along the characteristics $C_\nu(x, t)$ the feet of which lie on $\{x < 0\}$. At such positions, $\tilde{v}(x, 0) = v_L^0$. Since $h(u_R^0, v_L^0) < 0$ is strictly negative and exploiting the regularity of the solution, $h(\tilde{U})$ is expected to remain negative for small enough time.

For these reasons, we fix the source term $\tilde{\Sigma} = (a, c)$ and replace the initial conditions by conditions along $C_\nu(0)$ and $C_{-\mu}(0)$. In other words, we replace the problem (2.18) for $\tilde{U} = (\tilde{u}, \tilde{v})$ by

$$\begin{aligned} \partial_t \tilde{u} - \mu \partial_x \tilde{u} &= a, & \tilde{u}(\nu t, t) &= u_R(t), \\ \partial_t \tilde{v} + \nu \partial_x \tilde{v} &= c, & \tilde{v}(-\mu t, t) &= v_L(t). \end{aligned}$$

Let $(x, t) \in D_*$. The (backward) characteristic $C_{-\mu}(x, t)$ crosses the (forward) characteristic $C_\nu(0)$ at time $T_R(x, t) = \frac{\mu t + x}{\mu + \nu}$. The characteristic $C_\nu(x, t)$ crosses the (forward) characteristic $C_{-\mu}(0)$ at time $T_L(x, t) = \frac{\nu t - x}{\mu + \nu}$. This provides

$$\begin{aligned} \tilde{u}(x, t) &= u_R(T_R(x, t)) + a(t - T_R(x, t)) & (6.5a) \\ &= u_R^0 + \begin{cases} at & \text{if } T_R(x, t) < \tau_R, \\ a\tau_R + b(T_R(x, t) - \tau_R) + a(t - T_R(x, t)) & \text{otherwise} \end{cases} \end{aligned}$$

$$\begin{aligned} \tilde{v}(x, t) &= v_L(T_L(x, t)) + c(t - T_L(x, t)) & (6.5b) \\ &= v_L^0 + \begin{cases} ct & \text{if } T_L(x, t) < \tau_L, \\ c\tau_L + d(T_L(x, t) - \tau_L) + c(t - T_L(x, t)) & \text{otherwise} \end{cases} \end{aligned}$$

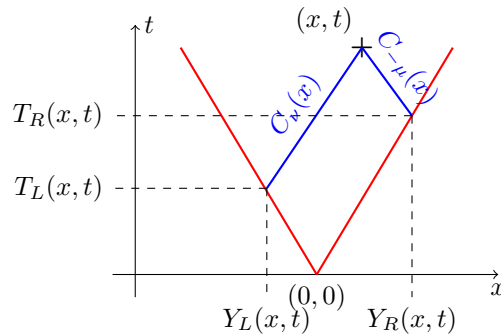


Fig. 15. Example of a configuration to compute $T_L(x, t)$ and $T_R(x, t)$ from a point $(x, t) \in D^*$.

Computation of the boiling curve $\Gamma(\tilde{U})$ in D^*

We now prove that $\{(x, t) \in D_* : h(\tilde{U})(x, t) = 0\}$ is overcharacteristic.

Lemma 6.3. *Suppose that (6.1) holds. Then there exists a unique $\tau_L^* > 0$ and a unique $\tau_R^* > 0$ such that $h(\tilde{U})(-\mu\tau_L^*, \tau_L^*) = 0$ and $h(\tilde{U})(\nu\tau_R^*, \tau_R^*) = 0$. Moreover, there exists a unique function $T \in W^{1,\infty}(I)$ such that for all $(x, t) \in D^*$*

$$\begin{cases} h(\tilde{U})(x, t) < 0 & \text{if and only if } t < T(x), \\ h(\tilde{U})(x, t) = 0 & \text{if and only if } t = T(x), \\ h(\tilde{U})(x, t) > 0 & \text{otherwise.} \end{cases}$$

Furthermore, T is continuous piecewise affine with at most three pieces and satisfies (3.1), i.e. its graph is overcharacteristic.

Proof. Decompose D^* into (see Fig. 16)

$$\begin{aligned} R_1 &= \{(x, t) \in D^*, \text{ s.t. } T_R(x, t) \leq \tau_R, \text{ and } T_L(x, t) \leq \tau_L\}, \\ R_2 &= \{(x, t) \in D^*, \text{ s.t. } T_R(x, t) > \tau_R, \text{ and } T_L(x, t) \leq \tau_L\}, \\ R_3 &= \{(x, t) \in D^*, \text{ s.t. } T_R(x, t) \leq \tau_R, \text{ and } T_L(x, t) > \tau_L\}, \\ R_4 &= \{(x, t) \in D^*, \text{ s.t. } T_R(x, t) > \tau_R, \text{ and } T_L(x, t) > \tau_L\}. \end{aligned}$$

From the expressions (6.5) of \tilde{u} and \tilde{v} we get

$$\begin{aligned} \forall (x, t) \in R_1, \quad h(\tilde{U})(x, t) &= h(u_R^0, v_L^0) + h(a, c)t, \\ \forall (x, t) \in R_2, \quad h(\tilde{U})(x, t) &= h(u_R^0, v_L^0) + th\left(\frac{a\nu + b\mu}{\mu + \nu}, c\right) + h_1(a - b)\left(\tau_R - \frac{x}{\mu + \nu}\right), \\ \forall (x, t) \in R_3, \quad h(\tilde{U})(x, t) &= h(u_R^0, v_L^0) + th\left(a, \frac{d\nu + c\mu}{\mu + \nu}\right) + h_2(c - d)\left(\tau_L + \frac{x}{\mu + \nu}\right), \\ \forall (x, t) \in R_4, \quad h(\tilde{U})(x, t) &= h(u_R^0, v_L^0) + th\left(\frac{a\nu + b\mu}{\mu + \nu}, \frac{d\nu + c\mu}{\mu + \nu}\right) \\ &\quad + h\left((a - b)\left(\tau_R - \frac{x}{\mu + \nu}\right), (c - d)\left(\tau_L + \frac{x}{\mu + \nu}\right)\right). \end{aligned}$$

As a consequence, $h(\tilde{U})(x, t) = 0$ if and only if

$$\text{in } R_1, \quad t = \frac{-h(u_R^0, v_L^0)}{h(a, c)}, \quad (6.6a)$$

$$\text{in } R_2, \quad t = \frac{-h\left(u_R^0 + (a - b)\left(\tau_R - \frac{x}{\mu + \nu}\right), v_L^0\right)}{h\left(\frac{a\nu + b\mu}{\mu + \nu}, c\right)}, \quad (6.6b)$$

$$\text{in } R_3, \quad t = \frac{-h\left(u_R^0, v_L^0 + (c - d)\left(\tau_L + \frac{x}{\mu + \nu}\right)\right)}{h\left(a, \frac{c\mu + d\nu}{\mu + \nu}\right)}, \quad (6.6c)$$

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$$\text{in } R_4, \quad t = \frac{-h \left(u_R^0 + (a-b) \left(\tau_R - \frac{x}{\mu+\nu} \right), v_L^0 + (c-d) \left(\tau_L + \frac{x}{\mu+\nu} \right) \right)}{h \left(\frac{a\nu+b\mu}{\mu+\nu}, \frac{c\mu+d\nu}{\mu+\nu} \right)}. \quad (6.6d)$$

Remark that \tilde{u} and \tilde{v} are continuous in D^* , thus $h(\tilde{U})$ is continuous. Then (6.6) defines a function T continuous piecewise affine. Using (6.1), one verifies that the x -derivative of (6.6) satisfies (4.5). \square

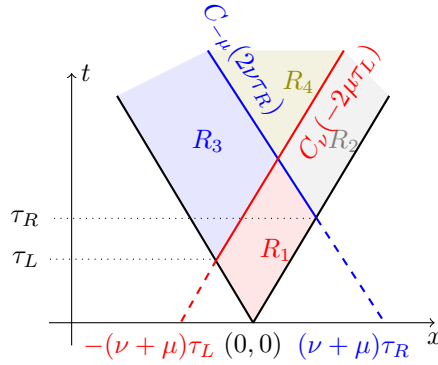


Fig. 16. Representation of the four regions of D^* in which \tilde{u} and \tilde{v} are linear.

*Correction of the solution for long time \bar{U} in D^**

Now, we correct the short time solution \tilde{U} so that the source term is (b, d) above the boiling curve $\Gamma(\tilde{U})$ as it should be. Define

$$\bar{u}(x, t) = u_R(T_R(x, t)) + a\mathcal{T}(x, t) + b((t - T_R(x, t)) - \mathcal{T}(x, t)), \quad (6.7a)$$

$$\bar{v}(x, t) = v_L(T_L(x, t)) + c\mathfrak{I}(x, t) + d((t - T_L(x, t)) - \mathfrak{I}(x, t)), \quad (6.7b)$$

$$\mathcal{T}(x, t) = \int_{T_R(x, t)}^t \mathbb{1}_{\mathbb{R}^-} (h(\tilde{u}, \tilde{v})(x + \mu(t - \tau), \tau)) d\tau,$$

$$\mathfrak{I}(x, t) = \int_{T_L(x, t)}^t \mathbb{1}_{\mathbb{R}^-} (h(\tilde{u}, \tilde{v})(x - \nu(t - \tau), \tau)) d\tau.$$

The characteristic $C_{-\mu}(x, t)$ in D^* consists of two parts, one below the boiling line $\Gamma(\tilde{U})$ in the time interval $t \in [T_R(x, t), \mathcal{T}(x, t)]$ where the source term is a , the other above the boiling line in the time interval $t \in [T_R(x, t), \mathcal{T}(x, t)]$ and where the source is b . The two graphs intersect at most once since T is overcharacteristic. Similarly, $\mathfrak{I}(x, t)$ plays the same role along the characteristic $C_{\nu}(x, t)$ (see Fig. 17).

Explicit computations (Appendix Appendix C) give

$$\mathcal{T}(x, t) = \begin{cases} 0 & \text{if } \tau_R^* \leq T_R(x, t), \\ t - T_R(x, t) & \text{if } t \leq T(x), \\ \mathfrak{t}(T_R(x, t)) - T_R(x, t) & \text{otherwise,} \end{cases} \quad (6.8a)$$

$$\mathfrak{T}(x, t) = \begin{cases} 0 & \text{if } \tau_L^* \leq T_L(x, t), \\ t - T_L(x, t) & \text{if } t \leq T(x), \\ \tau(T_L(x, t)) - T_L(x, t) & \text{otherwise,} \end{cases} \quad (6.8b)$$

where τ_R^* and τ_L^* are the times at which ends and starts Γ^* , i.e. the times spent in $Z^-(\tilde{U})$ along $C_\nu(0)$ and along $C_{-\mu}(0)$. It yields (see Appendix Appendix C)

$$\tau_R^* = T((Id - \nu T)^{-1}(0)), \quad \tau_L^* = T((Id + \mu T)^{-1}(0)). \quad (6.8c)$$

And $\mathfrak{t}(t)$, resp. $\tau(t)$, are the times spent in $Z^-(\tilde{U})$ along $C_{-\mu}(\nu t, t)$, resp. $C_\nu(-\mu t, t)$, i.e. starting at the point $(\nu t, t) \in C_\nu(0)$, resp. $(-\mu t, t) \in C_{-\mu}(0)$. It yields (see Appendix Appendix C)

$$\mathfrak{t}(t) = T((Id + \mu T)^{-1}((\mu + \nu)t)), \quad \tau(t) = T((Id - \nu T)^{-1}(-(\mu + \nu)t)). \quad (6.8d)$$

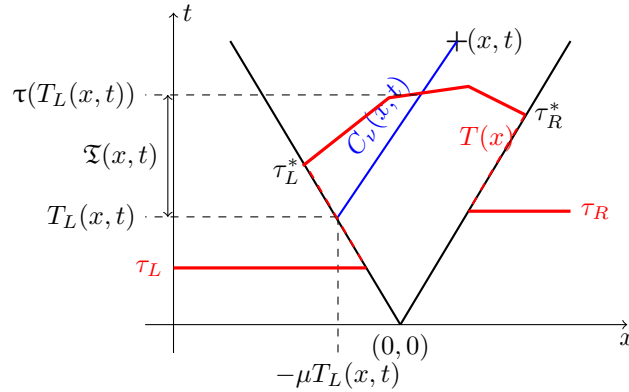


Fig. 17. Example of configuration for the computation of \mathfrak{T} .

Conclusion

We now verify that \bar{U} is indeed a solution to the Riemann problem. The key point is that $h(\bar{U})$ is strictly increasing along the characteristics, thus it crosses the critical value 0 at most once. Before this time, the source term is (a, c) and $\bar{U} = \tilde{U}$; after this time, the source term is (b, d) as in the definition of \bar{U} . The proof is an adaptation of Section 4.2.

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Lemma 6.4. *Suppose that (6.1) holds. Then, $\bar{U} = (\bar{u}, \bar{v})$ given by (6.7-6.8) satisfies*

$$(\partial_t - \mu\partial_x)h(\bar{U}) > 0, \quad (\partial_t + \nu\partial_x)h(\bar{U}) > 0 \quad \text{a.e. in } (x, t) \in D^*.$$

Proof. The PDE given by (4.10-4.11) and (4.12) is modified into

$$(\partial_t - \mu\partial_x)\bar{v}(x, t) = (\partial_t - \mu\partial_x)T_L(x, t) [v'_L(T_L(x, t)) - d] + (\partial_t - \mu\partial_x)\mathfrak{F}(x, t)(c - d) + d, \quad (6.9a)$$

$$(\partial_t + \nu\partial_x)\bar{u}(x, t) = (\partial_t + \nu\partial_x)T_R(x, t) [u'_R(T_R(x, t)) - b] + (\partial_t + \nu\partial_x)\mathcal{T}(x, t)(a - b) + b. \quad (6.9b)$$

One verifies that

$$\begin{aligned} (\partial_t - \mu\partial_x)T_R(x, t) &= 0, & (\partial_t + \nu\partial_x)T_R(x, t) &= 1, \\ (\partial_t + \nu\partial_x)T_L(x, t) &= 0, & (\partial_t - \mu\partial_x)T_L(x, t) &= 1, \end{aligned}$$

and the other derivatives read

$$\begin{aligned} (\partial_t + \nu\partial_x)\mathcal{T}(x, t) &= \begin{cases} 0 & \text{if } \tau_R^* < T_R(x, t), \\ 0 & \text{if } t < T(x), \\ \mathfrak{t}'(T_R(x, t)) - 1 & \text{otherwise,} \end{cases} \\ (\partial_t - \mu\partial_x)\mathfrak{F}(x, t) &= \begin{cases} 0 & \text{if } \tau_L^* < T_L(x, t), \\ 0 & \text{if } t < T(x), \\ \mathfrak{t}'(T_L(x, t)) - 1 & \text{otherwise,} \end{cases} \\ \mathfrak{t}'(t) &= \frac{(\mu + \nu)T'}{1 + \mu T'} ((Id + \mu T)^{-1}((\mu + \nu)t)), \\ \mathfrak{t}'(t) &= \frac{-(\mu + \nu)T'}{1 - \nu T'} ((Id - \nu T)^{-1}(-(\mu + \nu)t)). \end{aligned}$$

Using the definition $h(x, T(x)) = 0$, one obtains

$$T'(x) = \frac{\partial_x h(\bar{u}, \bar{v})(x, T(x))}{\partial_t h(\bar{u}, \bar{v})(x, T(x))} = \frac{[(\partial_t - \mu\partial_x) - (\partial_t + \nu\partial_x)] h(\bar{u}, \bar{v})(x, T(x))}{[\nu(\partial_t - \mu\partial_x) + \mu(\partial_t + \nu\partial_x)] h(\bar{u}, \bar{v})(x, T(x))},$$

From (6.5), one computes

$$\begin{aligned} (\partial_t - \mu\partial_x)h(\bar{U})(x, t) &= h(a, \varsigma_L(x, t)), & \varsigma_L(x, t) &= \varsigma(u_L, v_L)(T_L(x, t)) \\ (\partial_t + \nu\partial_x)h(\bar{U})(x, t) &= h(\sigma_R(x, t), c), & \sigma_R(x, t) &= \sigma(u_R, v_R)(T_R(x, t)) \end{aligned}$$

which provides

$$T'(x) = \frac{h(a, \varsigma_L(x, T(x))) - h(\sigma_R(x, T(x)), c)}{\nu h(a, \varsigma_L(x, T(x))) + \mu h(\sigma_R(x, T(x)), c)},$$

and leads to

$$\begin{aligned} \mathfrak{t}'(t) &= 1 - \frac{h(\sigma_R(x_R(t), t_R(t)), c)}{h(a, \varsigma_L(x_R(t), t_R(t)))}, & \mathfrak{t}'(t) &= 1 - \frac{h(a, \varsigma_L(x_L(t), t_L(t)))}{h(\sigma_R(x_L(t), t_L(t)), c)}, \\ x_R(t) &= (Id + \mu T)^{-1}((\mu + \nu)t), & x_L(t) &= (Id - \nu T)^{-1}(-(\mu + \nu)t), \\ t_R(t) &= T(x_R(t)), & t_L(t) &= T(x_L(t)). \end{aligned}$$

Thus we get

$$(\partial_t - \mu \partial_x)h(\bar{U})(x, t) = h(\sigma(\bar{U})(x, t), \varsigma_L(x, t) + (d - c)\alpha(x, t)), \quad (6.10a)$$

$$(\partial_t + \nu \partial_x)h(\bar{U})(x, t) = h(\sigma_R(x, t) + (b - a)\beta(x, t), \varsigma(\bar{U})(x, t)), \quad (6.10b)$$

with the coefficients

$$\alpha(x, t) = \begin{cases} 0 & \text{if } t < T(x) \text{ or } \tau_L^* < T_L(x, t), \\ \frac{h(a, \varsigma_L(x, t))}{h(\sigma_R(x, t), c)} & \text{otherwise,} \end{cases} \quad (6.10c)$$

$$\beta(x, t) = \begin{cases} 0 & \text{if } t < T(x) \text{ or } \tau_R^* < T_R(x, t), \\ \frac{h(\sigma_R(x, t), c)}{h(a, \varsigma_L(x, t))} & \text{otherwise.} \end{cases} \quad (6.10d)$$

One concludes on the strict positivity from hypothesis (6.1):

- Either $t < T(x)$ or $\tau_R^* < T_R(x, t)$, resp. $\tau_L^* < T_L(x, t)$, then $\alpha = 0$, resp. $\beta = 0$, and (6.1a) provides the result.
- Or, $t > T(x)$ and $\Sigma(\bar{U})(x, t) = (b, d)$, and (6.10) turns into

$$\begin{aligned} (\partial_t - \mu \partial_x)h(\bar{U})(x, t) &= h\left(b, \varsigma_L(x, t) + (d - c)\frac{h(a, \varsigma_L(x, t))}{h(\sigma_R(x, t), c)}\right) \\ &= \frac{h(a, \varsigma_L(x, t))}{h(\sigma_R(x, t), c)} \left(\frac{h(\sigma_R(x, t), c)}{h(a, \varsigma_L(x, t))}h(b, \varsigma_L(x, t)) + h_2(d - c)\right), \\ (\partial_t + \nu \partial_x)h(\bar{U})(x, t) &= h\left(\sigma_R(x, t) + (b - a)\frac{h(\sigma_R(x, t), c)}{h(a, \varsigma_L(x, t))}, d\right) \\ &= \frac{h(\sigma_R(x, t), c)}{h(a, \varsigma_L(x, t))} \left(\frac{h(a, \varsigma_L(x, t))}{h(\sigma_R(x, t), c)}h(\sigma_R(x, t), d) + h_1(b - a)\right), \end{aligned}$$

which are strictly positive by hypothesis (6.1b) and (6.1c). \square

Finally, \bar{U} is a generalized solution.

Lemma 6.5. *Suppose that (6.1) holds. Then, (2.20) holds in D^* , and $U \in E$ defined by*

$$U(x, t) = U_L(x, t)\mathbf{1}_{D_L}(x, t) + \bar{U}(x, t)\mathbf{1}_{D^*}(x, t) + U_R(x, t)\mathbf{1}_{D_R}(x, t) \quad (6.11)$$

is a generalized solution.

Proof. This is a straightforward adaptation of Lemma 4.6: $h(\bar{U}) = h(\tilde{U}) < 0$ below Γ^* , then from the monotonicity of $h(\bar{U})$ along the characteristics and since Γ^* is overcharacteristic, one obtains that $h(\bar{U}) > 0$ above Γ^* . \square

6.4. Uniqueness of the generalized solution

The uniqueness is again obtained by adapting the computations of Section 4, i.e. when (6.1) holds, then (6.7-6.8) is the unique generalized solution in D^* and (6.11) is the unique generalized solution in $\mathbb{R} \times \mathbb{R}^+$.

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6.5. Numerical illustration

The solution \bar{U} is illustrated through the numerical example

$$\partial_t u - \partial_x u = \begin{cases} 0.2 & \text{if } u + v \leq 0, \\ 1 & \text{otherwise,} \end{cases} \quad u(0, x) = \begin{cases} 0.1 & \text{if } x \leq 2, \\ -0.4 & \text{otherwise,} \end{cases} \quad (6.12a)$$

$$\partial_t v + \partial_x v = \begin{cases} 1 & \text{if } u + v \leq 0, \\ 2 & \text{otherwise,} \end{cases} \quad v(0, x) = \begin{cases} -0.4 & \text{if } x \leq 2, \\ 0.2 & \text{otherwise.} \end{cases} \quad (6.12b)$$

over $(x, t) \in [0, 4] \times [0, 1]$. These data satisfy (6.1).

This numerical solution is obtained using the upwind scheme

$$u_i^{n+1} = u_i^n \left(1 - \frac{\Delta t}{\Delta x}\right) + u_{i+1}^n \frac{\Delta t}{\Delta x} + \begin{cases} 0.2\Delta t & \text{if } u_i^n + v_i^n \leq 0, \\ \Delta t & \text{otherwise,} \end{cases} \quad (6.13a)$$

$$v_i^{n+1} = v_i^n \left(1 - \frac{\Delta t}{\Delta x}\right) + v_{i-1}^n \frac{\Delta t}{\Delta x} + \begin{cases} \Delta t & \text{if } u_i^n + v_i^n \leq 0, \\ 2\Delta t & \text{otherwise.} \end{cases} \quad (6.13b)$$

For general purposes, the use of an adapted discretization, especially of well-balanced techniques such as upwinding the source ([4,15,22,16]), is recommended to avoid numerical artifacts. The study of such adapted schemes for discontinuous sources is postponed to future work and (6.13) was found sufficient here and the qualitative results presented below are converged.

The spatial domain $x \in [0, 4]$ is meshed with 800 cells and the time step is defined using a CFL of 0.95. The numerical solution (u, v) is plotted on Fig. 18 at the times $t = 0.1$, $t = 0.2$, $t = 0.3$, $t = 0.6$ and $t = 1$. For each t , the solution $U = (u, v)$ is discontinuous at two points that move at velocity $-\mu$ and ν . In D_L and D_R , $U = U_L$ and $U = U_R$ is space independent, while it is piecewise linear in D^* . Though, the number of pieces in this center domain and their slopes vary in time.

This change in time is due to the switch of sign of $h(U)$. Where $h(U)$ changes sign, the derivatives of the solution U become discontinuous. Then, the location of this change of regularity is transported along the characteristics. This phenomenon happens on the boundary of D^* at the times τ_L and τ_R when $h(U_L)$ and $h(U_R)$ change sign, but also inside D^* . These lines where the solution loses its C^1 regularity are represented on Fig. 18 (bottom left) in green over the colormap of $h(U)$, and the yellow line is the line of discontinuity $\Gamma(U)$.

7. Conclusion, discussions and outlooks

This section gathers a summary of the results and a discussion about the constraints on the set of data u_0 , v_0 , a , b , c and d .

7.1. Conclusion

In this paper we consider a 2×2 hyperbolic system with linear fluxes and a source term that is discontinuous in the unknown as a first attempt to study the numerical

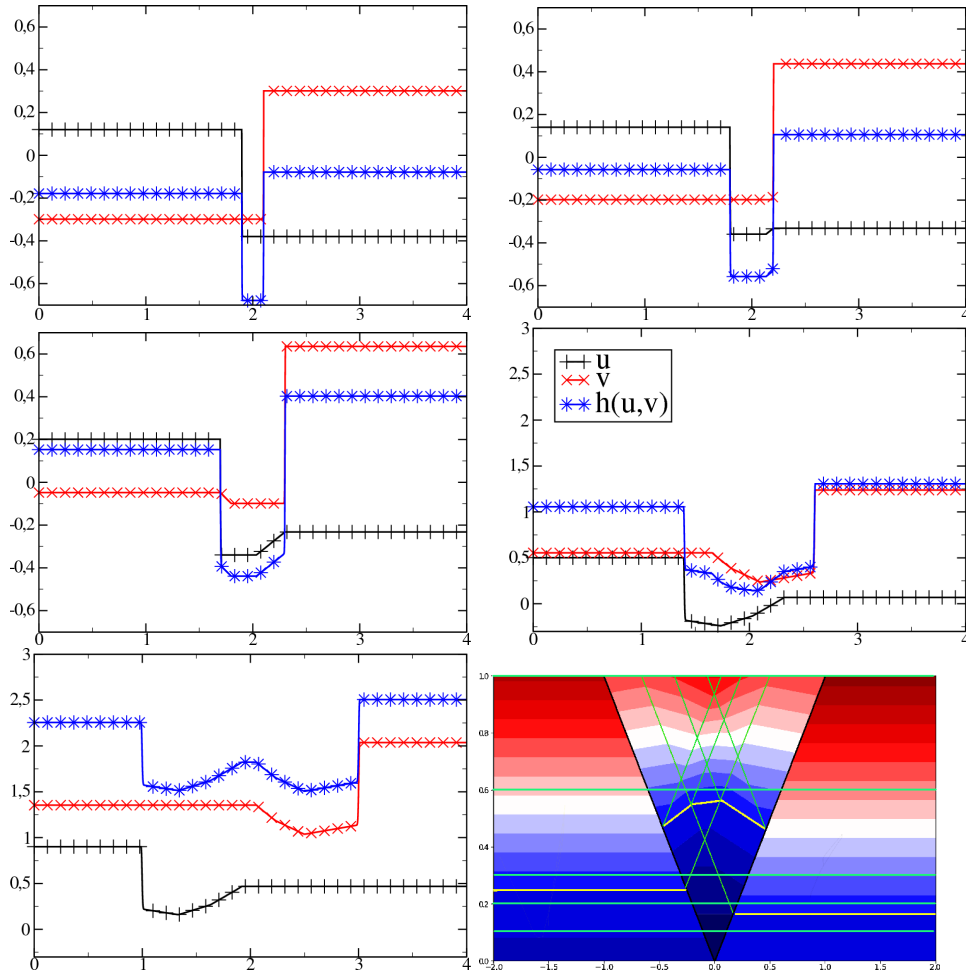


Fig. 18. Numerical solution $U = (u, v)$ to (6.12) at times $t = 0.1$ (top left), $t = 0.2$ (top right), $t = 0.3$ (middle left), $t = 0.6$ (middle right) and $t = 1$ (bottom left), and representation (bottom right) of these times (in cyan) over the colormap of $h(u, v)$ in (x, t) -space.

modeling of flows with fast evolving sources, such as a boiling two-phase flow. Two important notions are introduced to study such problems. First, the solution is understood in an integral sense along the characteristics. Second, the source term jumps along a line in the (x, t) -space that is non-tangent (in a strong sense) to all the characteristic curves. These notions lead to prove the existence and uniqueness of a solution under different sets of restrictions over the data (initial conditions and value of the source) according to the regularity of the data: first in the case of regular initial conditions with either an overcharacteristic or a subcharacteristic line of discontinuity, then for Riemann problems.

Two relations between these sets of data are now discussed:

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- the relation between a Riemann problem (1.2, 3.5) in the limit $(\mu, \nu) \rightarrow (0, 0)$ and a system of two ODEs with discontinuous RHS,
- the relation between the PDE (1.2) when the initial condition (satisfying (4.2)) tends to a Heaviside function and the resulting Riemann problem.

7.2. *The Riemann problem in the zero velocity limit*

Formally, when $(\mu, \nu) \rightarrow 0_{\mathbb{R}^2}$, the Riemann problem (1.2,3.5) degenerates toward a system of two ODE's of the form (2.1), one for $x \in \mathbb{R}^-$ and one for $x \in \mathbb{R}^+$. However, the requirement in Theorem 6.1 degenerates in a condition more restrictive than in Lemma 6.2 for the existence and uniqueness of a solution. Recall that (6.1) is sufficient but not necessary for the existence of a solution.

Indeed, one computes the solution U from (1.2,3.5) in the limit $(\mu, \nu) \rightarrow 0_{\mathbb{R}^2}$, and U locally tends to

$$u(x, t) = \begin{cases} u_L(t) & \text{if } x \leq 0, \\ u_R(t) & \text{otherwise,} \end{cases} \quad v(x, t) = \begin{cases} v_L(t) & \text{if } x \leq 0, \\ v_R(t) & \text{otherwise,} \end{cases}$$

where u_L, v_L, u_R and v_R are defined in (6.3). By construction, these functions solve the ODE (6.4) and they exist under the condition $h(a, c)h(b, d) > 0$.

For the existence of a solution to the Riemann problem, one remarks that the intensity of the velocities $\mu > 0$ and $\nu > 0$ have no influence on condition (6.1).

7.3. *The overcharacteristic problem in the limit of Heaviside initial condition*

The three theorems 3.3, 3.5 and 3.7 provide solutions under particular assumptions on the data u_0, v_0, a, b, c, d .

One remarks that the sets of data used in the different theorems are disjoint. For instance, the condition (4.2) on u_0 of Theorem 4.2 implies that $\inf u'_0 > -\infty$ and $\sup v'_0 < \infty$. In particular, using this overcharacteristic framework, we may never consider the limit

$$(u_0, v_0) \rightarrow (u_L, v_L) + (u_R - u_L, v_R - v_L)\mathbb{1}_{\mathbb{R}^+}$$

with $u_L > u_R$ or $v_L < v_R$, corresponding to a Riemann problem. However, such a problem can be studied in the framework of Theorem 3.7 under condition (6.1).

This also implies that we may not expect the possibility of approximating such a problem with stiff initial data by a more regular one with the conditions that we have exhibited.

7.4. *Outlooks*

The main outlook to the present work consists in constructing numerical schemes for (1.1), possibly based on the method used to construct solutions to (1.2). A naive discretization of these equations may produce numerical artefacts on certain test

cases, such as artificial oscillations, non-physical wave speeds, loss of admissibility (non-positive pressure or energy) or lack of preservation of equilibrium states. This impacts both the precision and the stability of the method. Using the scheme (6.13), we have been able to trigger such artefacts with a careful choice of the parameters a, b, c, d, u_0 and v_0 . Note that the parameters (6.12) of Section 6.5 were also chosen carefully to avoid these. There is a large literature discussing the discretization of source terms in balance laws including the well-balanced schemes (see e.g. [4,15,22, 16]), but it is mainly restricted to regular source terms. Extending these schemes to discontinuous source terms presents new difficulties, e.g. to obtain the existence of discrete steady states for (1.1). The present construction provides a solution to (1.2) and its behaviour in certain regimes. One may exploit the construction of Section 4 and 5 to construct well-balanced finite difference schemes or the one of Section 6 to construct well-balanced Godunov schemes for PDE with discontinuous sources.

Other extensions of the present work toward the general PDE (1.1) may be considered. The generalization from (1.2) to a system with N equations with linear fluxes and non-linear enthalpy is expected to be only technical and to present no real difficulty. However, the extension to non-linear fluxes is much more complicated. Indeed, only contact discontinuities were considered here, and the present method may misbehave when shock and rarefaction waves are forming. Especially, the notion of integral solutions along the characteristics can be generalized with non-linear fluxes when the solution is smooth, but such a construction needs further work in the presence of shocks.

Acknowledgement

This work was funded by CEA and LJLL (Sorbonne University and CNRS) through LRC MANON.

Appendix A. Computation of T_L and T_R in the Riemann problem

Consider $(x, t) \in D^*$ and define $(T_R(x, t), Y_R(x, t))$, resp. $(T_L(x, t), Y_L(x, t))$, the point of intersection between $C_{-\mu}(x, t)$ and $C_\nu(0)$, resp. $C_\nu(x, t)$ and $C_{-\mu}(0)$. By definition, these points satisfy

$$\begin{cases} Y_R + \mu T_R = x + \mu t, \\ Y_R - \nu T_R = 0, \end{cases} \Rightarrow \begin{cases} T_R(x, t) = \frac{\mu t + x}{\mu + \nu}, \\ Y_R(x, t) = \frac{\nu(\mu t + x)}{\mu + \nu}, \end{cases} \quad (\text{A.1a})$$

$$\begin{cases} Y_L + \mu T_L = 0, \\ Y_L - \nu T_L = x - \nu t, \end{cases} \Rightarrow \begin{cases} T_L(x, t) = \frac{\nu t - x}{\mu + \nu}, \\ Y_L(x, t) = \frac{-\mu(\nu t - x)}{\mu + \nu}. \end{cases} \quad (\text{A.1b})$$

Appendix B. Computation of $\tilde{\Sigma}$ in the Riemann problem

Any generalized solution $U = (u, v) \in E$, especially u remains continuous along $C_{-\mu}(x)$ and v along $C_{\nu}(x)$ for all $x \neq 0$. The value of the sources $\tilde{\Sigma}$ in D^* is chosen constant based on the following property.

Lemma Appendix B.1. *Any generalized solution $U = (u, v)$ satisfies*

$$\lim_{\substack{(x,t) \rightarrow (0,0) \\ (x,t) \in D^*}} h(U)(x, t) = h(u_R^0, v_L^0).$$

Proof. Using the method of characteristics leads to define the solution as

$$\begin{aligned} u(x, t) &= u_R(T_R(x, t)) + \int_{T_R(x, t)}^t \sigma(u, v)(x + \mu(t - \tau), \tau) d\tau, \\ v(x, t) &= v_L(T_L(x, t)) + \int_{T_L(x, t)}^t \varsigma(u, v)(x - \nu(t - \tau), \tau) d\tau, \end{aligned}$$

where u_R and v_L are given by (6.3). One computes

$$\begin{aligned} h(U)(x, t) &= h(u_R(T_R(x, t)), v_L(T_L(x, t))) \\ &+ h\left(\int_{T_R(x, t)}^t \sigma(u, v)(x + \mu(t - \tau), \tau) d\tau, \int_{T_L(x, t)}^t \varsigma(u, v)(x - \nu(t - \tau), \tau) d\tau\right). \end{aligned}$$

Since σ is bounded by a and b and ς by c and d , the integrals in this formula are bounded. By continuity of T_R , T_L and h , computing $h(U)$ in the limit $D^* \ni (x, t) \rightarrow (0, 0)$ provides the result. \square

Appendix C. Details on the computation of \mathcal{T} and \mathfrak{T} in D^*

Since Γ^* is overcharacteristic, for all $(x, t) \in D^*$, the characteristic curves $C_{-\mu}(x, t)$ and $C_{\nu}(x, t)$ cross Γ^* at most once.

- Either $t < T(x)$ then the part of the characteristics in D^* and before time t (see Fig. 17) are below Γ^* and thus entirely in $Z^-(U)$. Then the term in the integral (6.7) is constant equal to one and

$$\mathcal{T}(x, t) = t - T_R(x, t) \quad \text{if } t \leq T(x), \quad \mathfrak{T}(x, t) = t - T_L(x, t) \quad \text{if } t \leq T(x).$$

- Denote τ_R^* and τ_L^* the points such that

$$(-\mu\tau_L^*, \tau_L^*) \in C_{-\mu}(0) \cap \Gamma^*, \quad (\nu\tau_R^*, \tau_R^*) \in C_{\nu}(0) \cap \Gamma^*,$$

i.e. the starting and ending points of Γ^* in D^* .

Then, consider a point $(x, t) \in D^*$ such that $T_R(x, t) \geq \tau_R^*$ and $T_L(x, t) \geq \tau_L^*$. Again since Γ^* is overcharacteristic, the part of the characteristics $C_{-\mu}(x, t)$ and $C_{\nu}(x, t)$ in D^* and before time t are entirely in $Z^+(\tilde{U})$. Then the term in the integral (6.7) is constant equal to zero and

$$\mathcal{T}(x, t) = 0 \quad \text{if } T_R(x, t) \geq \tau_R^*, \quad \mathfrak{T}(x, t) = 0 \quad \text{if } T_L(x, t) \geq \tau_L^*.$$

The values of τ_L^* and τ_R^* satisfy

$$\begin{cases} Y_L^* + \mu\tau_L^* = 0, \\ T(Y_L^*) = \tau_L^*, \end{cases} \Rightarrow \begin{cases} Y_L^* = (Id + \mu T)^{-1}(0), \\ \tau_L^* = T((Id + \mu T)^{-1}(0)), \end{cases}$$

$$\begin{cases} Y_R^* - \nu\tau_R^* = 0, \\ T(Y_R^*) = \tau_R^*, \end{cases} \Rightarrow \begin{cases} Y_R^* = (Id - \nu T)^{-1}(0), \\ \tau_R^* = T((Id - \nu T)^{-1}(0)). \end{cases}$$

These values are obtained along $C_{-\mu}(0^+)$ and $C_\nu(0^-)$ when $h(\tilde{U})$ reaches the value 0. Along these lines, one computes

$$\begin{aligned} \forall (x, t) \in C_{-\mu}(0^+) \cap Z^-(\tilde{U}), \quad h(\tilde{U})(x, t) &= h(u_R^0 + at, v_L(t)) \\ &= \begin{cases} h(u_R^0, v_L^0) + h(a, c)t & \text{if } t \leq \tau_L, \\ h(u_R^0, v_L^0 + (c-d)\tau_L) + h(a, d)t & \text{otherwise,} \end{cases} \\ \forall (x, t) \in C_\nu(0^-) \cap Z^-(\tilde{U}), \quad h(\tilde{U})(x, t) &= h(u_R(t), v_L^0 + ct) \\ &= \begin{cases} h(u_R^0, v_L^0) + h(a, c)t & \text{if } t \leq \tau_R, \\ h(u_R^0 + (a-b)\tau_R, v_L^0) + h(b, c)t & \text{otherwise,} \end{cases} \end{aligned}$$

which provides

$$\begin{aligned} h(\tilde{U})(-\mu t, t) = 0 &\Leftrightarrow t = \tau_L^* = \begin{cases} \frac{-h(u_R^0, v_L^0)}{h(a, c)} & \text{if } \frac{-h(u_R^0, v_L^0)}{h(a, c)} < \tau_L, \\ \frac{-h(u_R^0, v_L^0 + (c-d)\tau_L)}{h(a, d)} & \text{otherwise,} \end{cases} \\ h(\tilde{U})(\nu t, t) = 0 &\Leftrightarrow t = \tau_R^* = \begin{cases} \frac{-h(u_R^0, v_L^0)}{h(a, c)} & \text{if } \frac{-h(u_R^0, v_L^0)}{h(a, c)} < \tau_R, \\ \frac{-h(u_R^0 + (a-b)\tau_R, v_L^0)}{h(b, c)} & \text{otherwise.} \end{cases} \end{aligned}$$

- Otherwise, at $(x, t) \in D^*$, the characteristic curve $C_{-\mu}(x, t)$, or $C_\nu(x, t)$, has crossed once Γ^* . Then the term in the integral (6.7) is one before \mathfrak{t} , resp. τ , and 0 after, then these integrals rewrite

$$\begin{aligned} \mathcal{T}(x, t) &= \mathfrak{t}(T_R(x, t)) - T_R(x, t) & \text{if } t < T(x) \text{ and } T_R(x, t) \leq \tau_R^*, \\ \mathcal{X}(x, t) &= \tau(T_L(x, t)) - T_L(x, t) & \text{if } t < T(x) \text{ and } T_L(x, t) \leq \tau_L^*. \end{aligned}$$

Here, $\mathfrak{t}(T_R(x, t))$, resp. $\tau(T_L(x, t))$, is the time spent in $Z^-(\tilde{U})$ along the characteristic $C_{-\mu}(x, t) = C_{-\mu}(\nu T_R(x, t), T_R(x, t))$, resp. $C_\nu(x, t) = C_\nu(-\mu T_L(x, t), T_L(x, t))$. Especially, these times satisfy

$$(\mathfrak{t}, \mathfrak{t})(t) \in C_{-\mu}(\nu t, t) \cap \Gamma^*, \quad (\xi, \tau)(t) \in C_\nu(-\mu t, t) \cap \Gamma^*,$$

that rewrites

$$\begin{aligned} \begin{cases} \mathfrak{r}(t) + \mu\mathfrak{t}(t) = (\nu + \mu)t, \\ T(\mathfrak{r}(t)) = \mathfrak{t}(t), \end{cases} &\Leftrightarrow \begin{cases} \mathfrak{r}(t) = (Id + \mu T)^{-1}((\nu + \mu)t), \\ \mathfrak{t}(t) = T((Id + \mu T)^{-1}((\nu + \mu)t)), \end{cases} \\ \begin{cases} \xi(t) - \nu\tau(t) = -(\nu + \mu)t, \\ T(\xi(t)) = \tau(t), \end{cases} &\Leftrightarrow \begin{cases} \xi(t) = (Id - \nu T)^{-1}(-(\nu + \mu)t), \\ \tau(t) = T((Id - \nu T)^{-1}(-(\nu + \mu)t)), \end{cases} \end{aligned}$$

As for τ_L^* and τ_R^* , \mathfrak{t} and τ can be rewritten based on u_R and v_L and a, b, c and d .

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