# Liouville results and asymptotics of solutions of a quasilinear elliptic equation with supercritical source gradient term <br> Marie-Françoise Bidaut-Veron 

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# Liouville results and asymptotics of solutions of a quasilinear elliptic equation with supercritical source gradient term 

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#### Abstract

We consider the elliptic quasilinear equation $-\Delta_{m} u=u^{p}|\nabla u|^{q}$ in $\mathbb{R}^{N}$ with $q \geq m$ and $p>0,1<m<N$. Our main result is a Liouville-type property, namely, all the positive $C^{1}$ solutions in $\mathbb{R}^{N}$ are constant. We also give their asymptotic behaviour : all the solutions in an exterior domain $\mathbb{R}^{N} \backslash B_{r_{0}}$ are bounded. The solutions in $B_{r_{0}} \backslash\{0\}$ can be extended as a continuous functions in $B_{r_{0}}$. The solutions in $\mathbb{R}^{N} \backslash\{0\}$ has a finite limit $l \geq 0$ as $|x| \rightarrow \infty$. Our main argument is a Bernstein estimate of the gradient of a power of the solution, combined with a precise Osserman's type estimate for the equation satisfied by the gradient.


Key Words: Liouville property, Bernstein method, Keller-Osserman estimates.
MSC2010: 35J92.

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## 1 Introduction

In this paper we study local and global properties of positive solutions of the equation

$$
\begin{equation*}
-\operatorname{div}\left(|\nabla u|^{m-2} \nabla u\right):=-\Delta_{m} u=u^{p}|\nabla u|^{q} \tag{1.1}
\end{equation*}
$$

in $\mathbb{R}^{N},(N \geq 1,1<m<N$ and $p>0)$ in the supercritical case

$$
\begin{equation*}
q \geq m \tag{1.2}
\end{equation*}
$$

We are concerned by the Liouville property in $\mathbb{R}^{N}$, which is wether all the positive $C^{1}$ solutions are constant. We also study the asymptotic behaviour of any solution of (1.1) near a singularity in the punctured ball $B_{r_{0}} \backslash\{0\}$, in $\mathbb{R}^{N} \backslash\{0\}$ or in an exterior domain $\mathbb{R}^{N} \backslash B_{r_{0}}$.

In the case $q=0$, equation (1.1) reduces to the classical Lane-Emden-Fowler equation

$$
\begin{equation*}
-\Delta_{m} u=u^{p} \tag{1.3}
\end{equation*}
$$

which has already been the subject of countless publications. One of the questions solved is that the Liouville property holds if and only if

$$
p<p_{m}^{*}:=\frac{N(m-1)+m}{N-m}
$$

Note that $p_{m}^{*}$ is the Sobolev exponent. Since it is impossible to quote all the articles on the subject, we only mention here the pioneering works and references therein. Gidas and Spruck [21] first showed the nonexistence of positive solutions in $\mathbb{R}^{N}$ for $m=2$ and $p<p_{2}^{*}$. They combine the Bernstein technique applied in the equation satisfed by the gradient of a suitable power of $u$, with delicate integral estimates ensuring the Harnack inequality, see also [6]. Then the complete behaviour up to the case $p=p_{2}^{*}$ was obtained by moving plane methods by [14], see also [16]. In the general case $m>1$, the nonexistence of nontrivial solutions for $p<p_{2}^{*}$ was proved in a beautiful article of Serrin and Zou [33], then the extension to the case $p=p_{2}^{*}$ was done by [30] for $m<2$, then [36] for $1<m<2$, and finally [29] for any $m>1$.

When $p=0,(1.1)$ reduces to the Hamilton-Jacobi equation

$$
-\Delta_{m} u=|\nabla u|^{q}
$$

The Liouville property was proved in [24] for $m=2$, and in [9] for any $m>1$, using the Bernstein technique. In that case the nonexistence holds for any $q>m-1$, without any sign condition on the solution. Estimates of the gradient for more general problems can be found in [23].

For the general case of equation (1.1), consider the range of exponents

$$
p>0, \quad p+q+1-m>0
$$

As in the case $q=0$, there exists a "first subcritical case", where

$$
p<\frac{N(m-1)}{N-m}-\frac{(N-1) q}{N-m}
$$

for which any supersolution in $\mathbb{R}^{N}$ of equation (1.1) is constant, from [19]. Beyond this case, a second critical case appears when $0 \leq q<m-1$ : indeed there exist radial positive nonconstant solutions of (1.4) whenever $p \geq p_{m, q}^{*}$, where

$$
p_{m, q}^{*}=\frac{N(m-1)+m}{N-m}-\frac{q((N-1) q-N(m-1)+m)}{(N-m)(m+1-q)},
$$

see [15] and [8].
When $m=2<N$ and $p>0$, equation

$$
\begin{equation*}
-\Delta u=u^{p}|\nabla u|^{q} \tag{1.4}
\end{equation*}
$$

was studied in [8] for $0<q \leq 2$. The case $q=2$ could be solved explicitely by a change of the unknown function, showing that the Liouville property holds for any $p>0$. Using a direct Bernstein technique we obtained a first range of values of $(p, q)$ for which the Liouville property holds, in particular it holds when $p+q-1<\frac{4}{N-1}$, covering the first subcritical case. Using an integral Bernstein technique in the spirit of [21] we obtained a wider range of $(p, q)$ ensuring the Liouville property, recovering Gidas and Spruck result $p<\frac{N+2}{N-2}$ when $q=0$. However some deep questions remained unsolved: Does the property hold for any $p<p_{2, q}^{*}$ when $q<1$ ? Does it hold for any $p>0$ when $1 \leq q<2$ ?

In a recent article, Filippucci, Pucci and Souplet [20, Theorem 1.1] considered the case $m=2$, $q>2$, of a superquadratic growth in the gradient, a case which was not covered by [8]. They proved the following:
Theorem [20, Theorem 1.1] Any classical positive and bounded solution of equation (1.4) in $\mathbb{R}^{N}$ with $q \geq 2$ and $p>0$ is constant.

In this article, we prove that the Liouville property holds true not only for (1.4) but for the quasilinear equation (1.1) without the assumption of boundedness on the solution. Our main result is the following

Theorem 1.1 Let $u$ be any positive $C^{1}\left(\mathbb{R}^{N}\right)$ solution of equation (1.1), with $1<m<N$ and

$$
\begin{equation*}
q \geq m, \quad p \geq 0 \tag{1.5}
\end{equation*}
$$

Then $u$ is constant.
We show that the case $q=m$ can still be solved explicitely, giving the complete behaviour of the solutions of the equation, see Theorem 2.1. Next we assume $q>m$. We prove that all the solutions in an exterior domain are bounded, and we give the asymptotic behaviour ( $|x| \rightarrow 0$ and $|x| \rightarrow \infty)$ of the solutions in $\mathbb{R}^{N} \backslash\{0\}$ :

Theorem 1.2 Assume $1<m<N, q>m, p \geq 0$. Then any positive $C^{1}$ solution $u$ of (1.1) in $\mathbb{R}^{N} \backslash B_{r_{0}}$ is bounded. If $u$ is a non-constant solution, then $|\nabla u|$ does not vanish for $|x|>r_{0}$. Moreover any positive solution $u$ in $\mathbb{R}^{N} \backslash\{0\}$ satisfies

$$
\begin{equation*}
\lim _{|x| \rightarrow \infty} u(x)=l \geq 0 \tag{1.6}
\end{equation*}
$$

If $l>0$, there exist constants $C_{1}, C_{2}>0$ such that for $|x|$ large enough,

$$
\begin{equation*}
C_{1}|x|^{\frac{N-m}{m-1}} \leq|u(x)-l| \leq C_{2}|x|^{\frac{N-m}{m-1}} . \tag{1.7}
\end{equation*}
$$

Concerning the solutions in $B_{r_{0}} \backslash\{0\}$ and in particular in $\mathbb{R}^{N} \backslash\{0\}$ we proved an estimate of the gradient, showing that the solution is continuous up to 0 but the gradient is singular at 0 :

Theorem 1.3 Assume $1<m<N, q>m, p \geq 0$. Any positive solution $u$ in $B_{r_{0}} \backslash\{0\}$ is bounded near 0 , it can be extended as a continuous function in $B_{r_{0}}$, such that $u(0)>0$, and for any $x \in B_{\frac{r_{0}}{2}} \backslash\{0\}$

$$
\begin{equation*}
|\nabla u(x)| \leq C|x|^{-\frac{1}{q-m+1}} \tag{1.8}
\end{equation*}
$$

where $C=C(N, p, q, m, u)$. Finally

$$
\begin{equation*}
|u(x)-u(0)| \leq C|x|^{\frac{q-m}{q-m+1}}, \tag{1.9}
\end{equation*}
$$

near 0 , where $C=C(N, p, q, m, u(0))$. Moreover, if $u$ is defined in $\mathbb{R}^{N} \backslash\{0\}$, then $u(x) \leq u(0)$ in $\mathbb{R}^{N} \backslash\{0\}$.

Note that the exponent involved in (1.8) is independent of $p$, actualy the solution behaves like a solution of the Hamilton-Jacobi equation

$$
\begin{equation*}
-\Delta_{m} u=c|\nabla u|^{q} \tag{1.10}
\end{equation*}
$$

with $c=u^{p}(0)$.
Finally we make an exhaustive study of the radial solutions for $q>m$, showing the sharpness of the nonradial results. We reduce the study to the one of an autonomous quadratic polynomial system of order 2, following the technique introduced in [10]. Compared to other classical techniques, it provides a complete description of all the positive solutions, in particular the global ones, without questions of regularity. We prove the following:

Theorem 1.4 Assume $1<m<N, q>m, p \geq 0$ and $u$ is any positive non constant radial solution $r \mapsto u(r)$ of (1.1) in an interval $(a, b) \subseteq(0, \infty)$.
(i) If $a=0$, then $u$ is bounded, decreasing and singular:

$$
\begin{equation*}
\lim _{r \rightarrow 0} u=u_{0}>0, \quad \lim _{r \rightarrow 0} r\left|u^{\prime}\right|^{q-m+1}=\frac{a_{m, q}}{u_{0}^{p}}, \quad a_{m, q}=\frac{(N-1) q-N(m-1)}{q+1-m} \tag{1.11}
\end{equation*}
$$

And for given $u_{0}>0$, there exist infinitely many such solutions;
(ii) If $b=\infty$, then $u$ admits a limit a limit $l \geq 0$ at infinity and

$$
\begin{equation*}
\lim _{r \rightarrow \infty} r^{\frac{N-m}{m-1}}|u(r)-l|=k>0 \tag{1.12}
\end{equation*}
$$

Furthermore, for given $l>0, c \neq 0$ there exists a unique local solution near $\infty$, such that

$$
\begin{equation*}
\lim _{r \rightarrow \infty} r^{\frac{N-m}{m-1}}(u(r)-l)=c \tag{1.13}
\end{equation*}
$$

(iii) For any $u_{0}>0$, there exist infinitely many solutions in $(0, \infty)$, decreasing, such that $\lim _{r \rightarrow 0} u=$ $u_{0}$, but a unique one, satisfying

$$
\begin{equation*}
\lim _{r \rightarrow 0} u=u_{0} \quad \text { and } \quad \lim _{r \rightarrow \infty} u=0 \tag{1.14}
\end{equation*}
$$

There exist infinitely many solutions defined on an interval $(0, \rho)$, such that $\lim _{r \rightarrow \rho} u=0$, and an infinity such that $\lim _{r \rightarrow \rho} u^{\prime}=-\infty$. Finally, there exist an infinity of solutions in $(\rho, \infty)$ such that $\lim _{r \rightarrow \rho} u=0$, and an infinity of solutions such that $\lim _{r \rightarrow \rho} u^{\prime}=\infty$.

Note that Theorems 1.2 and 1.3 lead to the following natural question: are all the solutions in $\mathbb{R}^{N} \backslash\{0\}$ radially symmetric? This is still an open problem, even in the case $p=0$ of the Hamilton-Jacobi equation.

To conclude this paper, we improve another result of [20], where it was noticed that [20, Theorem 1.1] was still valid for $p<0, q \geq 2$. We prove here a much more general result covering the case $p=0$.

Theorem 1.5 Assume $1<m<N, p \leq 0$ and $p+q+1-m>0$. Then there exists a constant $C=C(N, p, q, m)>0$ such that for any positive $C^{1}$ solution $u$ of (1.1) in a bounded domain $\Omega$,

$$
|\nabla u(x)| \leq C \operatorname{dist}(x, \partial \Omega)^{-\frac{1}{q+1-m}}, \quad \forall x \in \Omega
$$

If $\Omega=\mathbb{R}^{N}$, then $u$ is constant.

Let us give a brief comment on the analogous equation with an absorption term:

$$
\begin{equation*}
-\Delta_{m} u+u^{p}|\nabla u|^{q}=0 \tag{1.15}
\end{equation*}
$$

In the case $m=2,0<q<2$, a complete classification of the solutions with isolated singularities was performed in [17]. A main contribution was recently given by the same authors in [18] where they obtained optimal estimates of the gradient for any $1<m \leq N, p, q \geq 0, p+q-m+1>0$, still by the Bernstein method.

Our paper is organized as follows. In Section 2 we first treat the case $q=m$. In Section 3 we give the main ideas of our proofs when $q>m=2$, and we introduce some tools for the general case $q>m>1$. Our main theorems are proved in Section 4, and Section 5 is devoted to the radial case. The extension to the case $p \leq 0$ is given in Section 6 .

## 2 The case $q=m$

If $q=m$ we can express explicitely the solutions of (1.1). We prove the following:
Theorem 2.1 Let $1<m<N, p \geq 0, q=m$. Then
(i) any $C^{1}$ positive solution in $\mathbb{R}^{N}$ is constant;
(ii) any nonconstant positive solution in $\mathbb{R}^{N} \backslash B_{r_{0}}$ has a limit $l$ at $\infty$ and

$$
\lim _{|x| \rightarrow \infty}|x|^{\frac{m-N}{m-1}}|u-l|=c>0
$$

(iii) any positive solution in $B_{r_{0}} \backslash\{0\}$ extends as a continuous function in $B_{r_{0}}$, or satisfies

$$
\begin{equation*}
\lim _{x \rightarrow 0} \frac{u^{p+1}}{|\ln | x| |}=\frac{(N-m)(p+1)}{m-1} \tag{2.1}
\end{equation*}
$$

(iv) any positive solution in $\mathbb{R}^{N} \backslash\{0\}$ is radial.

Proof. We use a change of variable already considered in [1]: the equation takes the form

$$
\begin{equation*}
-\Delta_{m} u=\beta(u)|\nabla u|^{m}, \text { with } \beta(u)=u^{p} \tag{2.2}
\end{equation*}
$$

We set $\gamma(\tau)=\int_{0}^{\tau} \beta(\theta) d \theta=\frac{\tau^{p+1}}{p+1}$, and

$$
\begin{equation*}
U(x)=\Psi(u(x))=\int_{0}^{u(x)} e^{\frac{\gamma(\theta)}{m-1}} d \theta:=\int_{0}^{u(x)} e^{\frac{\theta^{p+1}}{(p+1)(m-1)}} d \theta . \tag{2.3}
\end{equation*}
$$

A function $u$ is a solution of (1.1) if and only if the above function $U$ satisfies

$$
-\Delta_{m} U=0,
$$

and if $u$ is nonnegative not identically $0, U$ is $m$-harmonic and positive. Conversely, $u$ is derived from $U$ by

$$
\begin{equation*}
u(x)=\Psi^{-1}(U(x))=\int_{0}^{U(x)} \frac{d s}{1+g(s)} \text { where } g(s)=\int_{0}^{s} \beta(\Psi(w)) d w=\int_{0}^{s} \Psi^{p}(w) d w \tag{2.4}
\end{equation*}
$$

(i) If $u$ is a solution in $\mathbb{R}^{N}$ of (2.2), it is constant. Indeed any nonnegative $m$-harmonic functions $U$ defined in $\mathbb{R}^{N}$ is constant, from the Harnack inequality, see [28], [31] and [33, Theorem II].
(ii) If $u$ is defined in $\mathbb{R}^{N} \backslash B_{r_{0}}$, then $U$ is bounded, it admits a limit $L$ at $\infty$ and there holds $|U(x)-L| \leq C|x|^{\frac{p-N}{p-1}}$ near $\infty$, see [3] for more general results. Clearly the same properties hold for $u$ (with another limit).
(iii) If $u$ is defined in $B_{r_{0}} \backslash\{0\}$, it follows from [31] that, either $U$ extends as a continuous $m$-harmonic function in $B_{r_{0}}$, or it behaves like $k|x|^{\frac{p-N}{p-1}}$ near 0 , so (2.1) holds.
(iv) If $u$ is a solution in $\mathbb{R}^{N} \backslash\{0\}$, it is proved in [27] that $U$ is radial and endows the form

$$
U(x)=k|x|^{\frac{m-N}{m-1}}+\lambda \quad \text { with } k, \lambda \geq 0
$$

Then $u$ is radial, and, using (2.4), it has the expression

$$
u(x)=\int_{0}^{\lambda} \frac{d s}{1+g(s)}+\int_{0}^{k|x|^{\frac{m-N}{m-1}}} \frac{d s}{1+g(s-\lambda)}
$$

## 3 Main arguments of the proofs

### 3.1 Ideas in the case $m=2$

Before detailling the proof of Theorem 1.1. for $q>m$, we give an overview of it in the simple case of equation (1.4), with $m=2, p>0, q>2$. We set $u=v^{b}$, with $b \in(0,1)$, and obtain

$$
-\Delta v=(b-1) \frac{|\nabla v|^{m}}{v}+b^{q-1} v^{s}|\nabla v|^{q},
$$

with $s=1-q+b(p+q-1)$. Next we explicit the equation satisfied by $z=|\nabla v|^{2}$. Taking in account the Böchner formula and Cauchy-Schwarz inequality in $\mathbb{R}^{N}$,

$$
\left.-\frac{1}{2} \Delta z+\frac{1}{N}(\Delta v)^{2}+<\nabla(\Delta v), \nabla v\right) \leq-\frac{1}{2} \Delta z+(H e s s v)^{2}+\langle\nabla(\Delta v), \nabla v>=0
$$

we get an estimate of the form, with universal constants $C_{i}>0$,

$$
-\Delta z+C_{1} v^{2 s} z^{q} \leq C_{2} \frac{z^{2}}{v^{2}}+C_{3} \frac{1}{v}\langle\nabla z, \nabla v\rangle+C_{4} v^{s} z^{\frac{q-2}{2}}\langle\nabla z, \nabla v\rangle
$$

then

$$
\begin{equation*}
-\Delta z+C_{5} v^{2 s} z^{q} \leq C_{6} \frac{z^{2}}{v^{2}}+C_{7} \frac{|\nabla z|^{2}}{z} . \tag{3.1}
\end{equation*}
$$

Using the Hölder inequality we deduce,

$$
-\Delta z+C_{8} v^{2 s} z^{q} \leq C_{9} v^{-\frac{2(q+2 s)}{q-2}}+C_{7} \frac{|\nabla z|^{2}}{z}
$$

The crucial step is an estimate of Osserman's type in a ball $B_{\rho}$ valid for functions satisfying the inequality

$$
-\Delta z+\alpha(x) z^{k} \leq \beta(x)+d \frac{|\nabla z|^{2}}{z} \quad \text { in } B_{\rho},
$$

where $k>1$. This is proved in Lemma 3.1 below, and it asserts that

$$
z(x) \leq C(N, k, d)\left(\frac{1}{\rho^{2}} \max _{B_{\rho}} \frac{1}{\alpha}\right)^{\frac{1}{k-1}}+\left(\max _{B_{\rho}} \frac{\beta}{\alpha}\right)^{\frac{1}{k}} \quad \text { in } B_{\frac{\rho}{2}} .
$$

Then we take $b=\frac{q-2}{p+q-1}$, in the same spirit as in [20], so that $\frac{B}{\alpha}$ is constant and $\alpha^{-1}(x)=v^{2}(x)$. We obtain an estimate

$$
\max _{\bar{B}_{\frac{p}{2}}}|\nabla v| \leq C\left(\left(\frac{\max _{B_{\rho}} v}{\rho}\right)^{\frac{1}{q-1}}+1\right)
$$

But any solution in $\mathbb{R}^{N}$ satisfies for any $\rho \geq 1$

$$
\begin{equation*}
\max _{B_{\rho}} v \leq v(0)+C \rho \max _{B_{\rho}}|\nabla v| \leq C \rho\left(1+\max _{B_{\rho}}|\nabla v|\right), \tag{3.2}
\end{equation*}
$$

which yields

$$
\max _{B_{\frac{\rho}{2}}}|\nabla v| \leq C\left(\left(\max _{B_{\rho}}|\nabla v|\right)^{\frac{1}{q-1}}+1\right)
$$

Using the bootstrap method developped in [11] and [9] based upon the fact that $\frac{1}{q-1}<1$, we deduce that $|\nabla v| \in L^{\infty}\left(\mathbb{R}^{N}\right)$. Note that the boundness of $|\nabla v|$ had been obtained in [20] but under the extra assumption $u \in L^{\infty}\left(\mathbb{R}^{N}\right)$, an assumption that we get rid of. Returning to $u=v^{b}$, it means that

$$
-\Delta u=u^{p}|\nabla u|^{q} \leq C \frac{|\nabla u|^{2}}{u},
$$

and the same happens for $u-l$, where $l=\inf _{\mathbb{R}^{N}} u$. It implies that $w_{l}=(u-l)^{\sigma}$ is subharmonic for $\sigma$ large enough. Then from [9], see also Lemma 3.3 below, and since $u$ is superharmonic,

$$
\sup _{B_{R}} w_{l} \leq C\left(\frac{1}{\left|B_{2 R}\right|} \int_{B_{2 R}} w^{\frac{1}{\sigma}}\right)^{\sigma}=C\left(\frac{1}{\left|B_{2 R}\right|} \int_{B_{2 R}}(u-l)\right)^{\sigma} \leq C^{\prime}\left(\inf _{B_{R}} u-l\right)^{\sigma} .
$$

Since $C^{\prime}$ is independent of $R$, it follows that $\sup _{\mathbb{R}^{N}} w_{l}=0$, thus $u \equiv l$.
Next we consider a solution in an exterior domain and we replace (3.2) by a more precise comparison estimate between $v(x)$ and its infimum on a sphere of radius $|x|$, and use the fact that this infimum is bounded as $r \rightarrow \infty$. Then we can show that $u$ is still bounded, and obtain the behaviour near $\infty$ by a careful study of $u$ and $w_{l}$. Finally we study the behaviour in $B_{r_{0}} \backslash\{0\}$ by the Bernstein technique, not relative to $v$ but directly to $u$, that means we take $b=1$, so that $s=p$. From (3.1) the function $\xi=|\nabla u|^{2}$ satisfies

$$
-\Delta \xi+C_{5} u^{2 p} \xi^{q} \leq C_{6} \frac{\xi^{2}}{u^{2}}+C_{7} \frac{|\nabla z|^{2}}{z}
$$

and $k=\inf _{B_{\frac{r_{0}}{2}} \backslash\{0\}} u$ is positive by the strong maximum principle, thus

$$
-\Delta \xi+C_{8} \xi^{q} \leq C_{9} \xi^{2}+C_{7} \frac{|\nabla z|^{2}}{z} \leq \frac{C_{8}}{2} \xi^{q}+C_{11}+C_{7} \frac{|\nabla z|^{2}}{z}
$$

from what we deduce the estimates of $\xi$.

### 3.2 Some tools

In the sequel we use the Bernstein method. In the case $p=0$, it appeared that the square of the gradient is a subsolution of an elliptic equation with absorption, for which one can find estimates from above of Osserman's type. In the case of equation (1.1), the problem is more difficult, but such upper estimates were also a main step in study of [8] of equation (1.4) for $q<2$. Here also they constitue a crucial step of our proofs below. The following Lemma gives an Osserman's type property of such equations, extending of [8, Lemma 2.2], see also used in [7, Proposition 2.1].

Lemma 3.1 Let $\Omega$ be a domain of $\mathbb{R}^{N}$, and $z \in C(\Omega) \cap C^{2}(G)$, where $G=\{x \in \Omega: z(x) \neq 0\}$. Let $w \mapsto \mathcal{A} w=-\sum_{i, j=1}^{N} a_{i j} \frac{\partial^{2} w}{\partial x_{i} \partial x_{j}}$ be a uniformly elliptic operator in the open set $G$ :

$$
\begin{equation*}
\theta|\xi|^{2} \leq \sum_{i, j=1}^{N} a_{i j} \xi_{i} \xi_{j} \leq \Theta|\xi|^{2}, \quad \theta>0 \tag{3.3}
\end{equation*}
$$

Suppose that for any $x \subset G$,

$$
\mathcal{A}(z)+\alpha(x) z^{k} \leq \beta(x)+d \frac{|\nabla z|^{2}}{z},
$$

with $k>1$, and $d=d(N, p, q)$, and $\alpha, \beta$ are continuous in $\Omega$ and $\alpha$ is positive. Then there exists $c=c(N, p, q, k)>0$ such that for any ball $\bar{B}\left(x_{0}, \rho\right) \subset \Omega$ there holds

$$
z\left(x_{0}\right) \leq c\left(\frac{1}{\rho^{2}} \max _{B_{\rho}\left(x_{0}\right)} \frac{1}{\alpha}\right)^{\frac{1}{k-1}}+\left(\max _{B_{\rho}\left(x_{0}\right)} \frac{\beta}{\alpha}\right)^{\frac{1}{k}} .
$$

Proof. Let $\bar{B}\left(x_{0}, \rho\right) \subset \Omega$. We can assume that $z\left(x_{0}\right) \neq 0$. Let $r=\left|x-x_{0}\right|$. Let $w$ be the function defined in $B_{\rho}\left(x_{0}\right)$ by

$$
w(x)=\lambda\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}}+\mu
$$

where $\lambda, \mu>0$. Let $G_{1}$ be a connected component of $\left\{x \in B_{\rho}\left(x_{0}\right) ; z(x)>w(x)\right\}$. Then $G_{1} \subset G$ and $\overline{G_{1}} \subset \bar{B}\left(x_{0}, \rho\right) \subset G$. We define $\mathcal{L} w$ in $B_{\rho}\left(x_{0}\right)$ by

$$
\mathcal{L}(w)=\mathcal{A}(w)+\alpha(x) w^{k}-\beta(x)-d \frac{|\nabla w|^{2}}{w} .
$$

Then

$$
\begin{gathered}
w_{x_{i}}=\frac{4 \lambda}{k-1}\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}-1} x_{i}, \\
w_{x_{i} x_{j}}= \\
\frac{4 \lambda}{k-1}\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}-1} \delta_{i j}+\frac{4 \lambda(k+1)}{(k-1)^{2}}\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}-2} x_{i} x_{j}, \\
\mathcal{A}(w)=-\sum_{i, j=1}^{N} a_{i j} w_{x_{i} x_{j}}=\frac{4 \lambda}{k-1}\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}-1}\left(-\sum_{i, j=1}^{N} a_{i j} \delta_{i j}\right) \\
\\
+\frac{4 \lambda(k+1)}{(k-1)^{2}}\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}-2}\left(-\sum_{i, j=1}^{N} a_{i j} x_{i} x_{j}\right) \\
\geq \\
=-\Theta\left(\frac{4 \lambda N}{k-1}\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}-1}+\frac{4 \lambda(k+1)}{(k-1)^{2}}\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}-2} r^{2}\right. \\
= \\
=-\Theta\left(\frac{4 \lambda N}{k-1}\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}-2}\left(N\left(\rho^{2}-r^{2}\right)+\frac{k+1}{k-1} r^{2}\right)\right. \\
|\nabla w|^{2}=\frac{16 \lambda^{2}}{(k-1)^{2}}\left(\rho^{2}-r^{2}\right)^{-\frac{4}{k-1}-2} r^{2} \Longrightarrow \frac{|\nabla w|^{2}}{w} \leq \frac{16 \lambda}{(k-1)^{2}}\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}-2} r^{2},
\end{gathered}
$$

and

$$
w^{k}=\left(\lambda\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}}+\mu\right)^{k} \geq \mu^{k}+\lambda^{k}\left(\rho^{2}-r^{2}\right)^{-\frac{2 k}{k-1}}=\mu^{k}+\lambda^{k}\left(\rho^{2}-r^{2}\right)^{-\frac{2}{k-1}-2} .
$$

We deduce from this series of inequalities,

$$
\begin{aligned}
\mathcal{L}(w) \geq & \alpha(x) \mu^{k}-\beta(x)+\lambda\left(\rho^{2}-r^{2}\right)^{-\frac{2 k}{k-1}}\left(\lambda^{k-1} C(x)\right. \\
& -\Theta\left(\frac{4}{k-1}\left(N \rho^{2}+\left(\frac{k+1}{k-1}-N\right) r^{2}\right)-\frac{16 d r^{2}}{(k-1)^{2}}\right. \\
\geq & \alpha(x) \mu^{k}-\beta(x)+\lambda\left(\rho^{2}-r^{2}\right)^{-\frac{2 k}{k-1}}\left(\lambda^{k-1} C(x)-c^{\prime} \rho^{2}\right),
\end{aligned}
$$

where $c^{\prime}=\Theta\left(\frac{4}{k-1}\left(2 N+\frac{k+1}{k-1}\right)+\frac{16 d}{(k-1)^{2}}=c^{\prime}(N, p, q, k)\right.$. We deduce that $\mathcal{L}(w) \geq 0$ if we impose

$$
\mu^{k} \geq \max _{B_{\rho}\left(x_{0}\right)} \frac{\beta}{\alpha} \text { and } \lambda^{k-1} \geq c^{\prime} \rho^{2} \max _{B_{\rho}\left(x_{0}\right)} \frac{1}{\alpha} .
$$

If $x_{1} \in G_{1}$ is such that $z\left(x_{1}\right)-w\left(x_{1}\right)=\max _{G_{1}}(z-w)>0$, then $\nabla z\left(x_{1}\right)=\nabla w\left(x_{1}\right)$, and $\mathcal{A}(z-$ $w)\left(x_{1}\right) \geq 0$. Therefore

$$
\left.0 \geq \mathcal{L}(z-w)\left(x_{1}\right)\right)=\mathcal{A}(z-w)\left(x_{1}\right)+\alpha(x)\left(z^{k}-w^{k}\right)\left(x_{1}\right)+d\left(\frac{|\nabla w|^{2}}{w}-\frac{|\nabla z|^{2}}{z}\right)
$$

Since the last term is positive, it is a contradiction. Then $z \leq w$ in $B_{\rho}\left(x_{0}\right)$. In particular $z\left(x_{0}\right) \leq$ $w\left(x_{0}\right)$.

We also use a bootstrap argument, initialy used in [11, Lemma 2.2], and then in [9] in more general form.

Lemma 3.2 Let $d, h \in \mathbb{R}$ with $d \in(0,1)$ and $y$ be a positive nondecreasing function on some interval $\left(r_{1}, \infty\right)$. Assume that there exist $K>0$ and $\varepsilon_{0}>0$ such that, for any $\varepsilon \in\left(0, \varepsilon_{0}\right]$ and $r>r_{1}$,

$$
y(r) \leq K \varepsilon^{-h} y^{d}(r(1+\varepsilon))
$$

Then there exists $C=C\left(K, d, h, \varepsilon_{0}\right)$ such that $\sup _{\left(r_{1}, \infty\right)} y \leq C$.
Proof. Consider the sequence $\left\{\varepsilon_{n}\right\}=\left\{\varepsilon_{0} 2^{-n}\right\}_{n \geq 1}$. Since the series $\sum \varepsilon_{n}$ is convergent, the sequence $\left\{P_{m}\right\}:=\left\{\prod_{i=1}^{m}\left(1+\varepsilon_{i}\right)\right\}_{m \geq 1}$ is convergent too, with limit $P>0$. Then there holds for any $r>r_{1}$

$$
y(r) \leq K \varepsilon_{1}^{-h} y^{d}\left(r\left(1+\varepsilon_{1}\right)\right)=K \varepsilon_{1}^{-h} y^{d}\left(r P_{1}\right) .
$$

We deduce by induction,

$$
y(r) \leq K^{1+d+. .+d^{m}}\left(\varepsilon_{1}^{-h} \varepsilon_{2}^{-h d} . . \varepsilon_{m}^{-h d^{m-1}}\right) y^{d^{m}}\left(P_{m} r\right)=\left(K \varepsilon_{0}^{-h}\right)^{1+d+. .+d^{m}}\left(2^{h\left(1+2 d+. .+m d^{m-1}\right)}\right) y^{d^{m}}\left(P_{m} r\right),
$$

and $r P_{m} \rightarrow r P, d^{m} \rightarrow 0$, thus $\left(y\left(P_{m} r\right)\right)^{d^{m}} \rightarrow 1$. Therefore we deduce that for any $r>r_{1}$,

$$
y(r) \leq\left(K \varepsilon_{0}^{-h}\right)^{\sum_{m=1}^{\infty} d^{m}} 2^{\sum_{m=1}^{\infty} m d^{m-1}}=\left(K \varepsilon_{0}^{-h}\right)^{\frac{1}{1-d}} 2^{\frac{d}{(1-d)^{2}}} .
$$

We also mention below a property of $m$-subharmonic functions given in [9, Lemma 2.1]. It's proof is also based upon a boostrap method and is valid for more general quasilinear operators:

Lemma 3.3 Let $u \in W_{\text {loc }}^{1, m}(\Omega)$ be nonnegative, m-subharmonic function in a domain $\Omega$ of $\mathbb{R}^{N}$. Then for any $\tau>0$, there exists a constant $C=C(N, m, \tau)>0$ such that for any ball $B_{2 \rho}\left(x_{0}\right) \subset \Omega$ and any $\varepsilon \in\left(0, \frac{1}{2}\right]$,

$$
\sup _{B_{\rho}\left(x_{0}\right)} u \leq C \varepsilon^{-\frac{N m^{2}}{\tau^{2}}}\left(\frac{1}{\left|B_{(1+\varepsilon) \rho}\left(x_{0}\right)\right|} \int_{B_{(1+\varepsilon) \rho}\left(x_{0}\right)} u^{\tau}\right)^{\frac{1}{\tau}}
$$

Finally we use some simple properties of mean value on spheres of $m$-superharmonic functions, in the same spirit as the ones given in [2, Lemmas 3.7, 3.8, 3.9] for mean values on annulus, and in [13] for $m=2$. For the sake of completeness we recall their proofs.

Lemma 3.4 Let $u \in C^{1}(\Omega)$ be nonnegative, $m$-superharmonic in $\Omega$.
(i) If $\Omega=\mathbb{R}^{N} \backslash B_{r_{0}}$, then

$$
r \mapsto \mu(r):=\inf _{|x|=r} u,
$$

is bounded in $\left(r_{0}, \infty\right)$, and strictly monotone or constant for large $r$. (ii) If $\Omega=B_{r_{0}} \backslash\{0\}$, then $r \mapsto \mu(r)$ is nonincreasing in $\left(0, r_{0}\right)$.

Proof. (i) Let $r>r_{0}$ be fixed. The function $f(x)=\mu(r)\left(1-\left(\frac{|x|}{r_{0}}\right)^{\frac{m-N}{m-1}}\right)$ is $m$-harmonic, and $u \geq f$ on $\partial B_{r} \cup \partial B_{r_{0}}$, therefore $u \geq f$ in $\overline{B_{r}} \backslash B_{r_{0}}$. Let $k>0$ large enough such that $1-k^{\frac{p-N}{p-1}} \geq \frac{1}{2}$. If we take $r>k r_{0}$ and any $x$ such that $|x|=k r_{0}$ we obtain

$$
u(x) \geq \mu\left(k r_{0}\right) \geq f(x)=\mu(r)\left(1-k^{\frac{p-N}{p-1}}\right) \geq \frac{1}{2} \mu(r)
$$

so $\mu(r)$ is bounded for $r>k r_{0}$. For any $r_{2}>r_{1}>r_{0}, \varphi\left(r_{1}, r_{2}\right):=\inf _{\overline{B_{r_{2}}} \backslash B_{1}} u=\min \left(\mu\left(r_{1}\right), \mu\left(r_{2}\right)\right)$ from the maximum principle. Then $\varphi$ is nonincreasing in $r_{2}$ and nondecreasing in $r_{1}$. If $\mu$ has a strict local minimum at some point $r$, then for $0<\delta<\delta_{0}$ small enough, $\mu(r)<\varphi\left(r-\delta_{0}, r+\delta_{0}\right) \leq$ $\varphi(r-\delta, r+\delta)$, which yields a contradiction as $\delta \rightarrow 0$. Then $\mu$ is monotone. If it is constant on two intervals ( $a, b$ ) and ( $a^{\prime}, b^{\prime}$ ) with $b<a^{\prime}$ and non-constant on ( $b, a^{\prime}$ ) it follows by Vazquez's maximum principle [35] that $u$ is constant on $\bar{B}_{b} \backslash B_{a}$ and on $\bar{B}_{b^{\prime}} \backslash B_{a^{\prime}}$ but non constant on $B_{a^{\prime}} \backslash \bar{B}_{b}$. It means, always by Vazquez's maximum principle,

- either $\min \left\{\mu(r): a<r<b^{\prime}\right\}=\mu(a)$ (if $\mu$ is nondecreasing) and the minimum of $u$ in $\bar{B}_{b^{\prime}} \backslash B_{a}$ is achieved in any point in $B_{b} \backslash \bar{B}_{a}$, hence $u$ is constant in $\bar{B}_{b^{\prime}} \backslash B_{a}$,
- or $\min \left\{\mu(r): a<r<b^{\prime}\right\}=\mu(a)$ (if $\mu$ is nonincreasing) and the minimum of $u$ in $\bar{B}_{b^{\prime}} \backslash B_{a}$ is achieved in any point in $B_{b^{\prime}} \backslash \bar{B}_{a^{\prime}}$, hence $u$ is constant in $\bar{B}_{b^{\prime}} \backslash B_{a}$.
In both case we obtain a contradiction. Hence $\mu$ is either strictly monotone for $r$ large enough, or it is constant, and so is $u$.
(ii) For given $r_{1}<r_{0}$, and $\delta>0$, there exists $\varepsilon_{\delta} \leq r_{1}$ such that for $0<\varepsilon<\varepsilon_{\delta}$, such that $\delta \varepsilon^{\frac{m-N}{m-1}} \geq \mu\left(r_{1}\right)$. Let $h(x)=\mu\left(r_{1}\right)-\delta|x|^{\frac{m-N}{m-1}}$. Then $u \geq h$ on $\partial B_{r_{1}} \cup \partial B_{\varepsilon}$, then $u \geq h$ in $\overline{B_{r_{1}}} \backslash B_{\varepsilon}$. Making $\varepsilon \rightarrow 0$ and then $\delta \rightarrow 0$, one gets $u \geq \mu\left(r_{1}\right)$ in $B_{r_{1}} \backslash\{0\}$, thus $\mu(r) \geq \mu\left(r_{1}\right)$ for $r<r_{1}$.


## 4 Proof of the main results

### 4.1 Proof of the Liouville property for $q>m$

We first give a general Bernstein estimate for solutions of equation (1.1):
Lemma 4.1 Let $u$ be any $C^{1}$ positive solution of ((1.1) in a domain $\Omega$, with $m>1$ and $p, q$ arbitrary real numbers. Let $G=\{x \in \Omega:|\nabla u(x)| \neq 0\}$. Let $u=v^{b}$ with $b \in \mathbb{R} \backslash\{0\}$ and $z=|\nabla v|^{2}$. Then the operator

$$
\begin{equation*}
w \mapsto \mathcal{A}(w)=-\Delta w-(m-2) \frac{D^{2} w(\nabla v, \nabla v)}{|\nabla v|^{2}}=-\sum_{i, j=1}^{N} a_{i j} v_{x_{i} x_{j}} \tag{4.1}
\end{equation*}
$$

with coefficients $a_{i, j}$ depending on $\nabla v$, is uniformly elliptic in $G$, and for any $\varepsilon>0$, there exists $C_{\varepsilon}=C_{\varepsilon}(N, m, p, q, b, \varepsilon)$ such that

$$
\begin{align*}
-\frac{1}{2} \mathcal{A}(z) & +\left(\frac{1-\varepsilon}{N}(b-1)^{2}(m-1)^{2}-(1-b)(m-1)\right) \frac{z^{2}}{v^{2}}+\frac{1-2 \varepsilon}{N}|b|^{2(q-m+1)} v^{2 s} z^{q+2-m}  \tag{4.2}\\
& +\left(\frac{1}{N} 2(b-1)(m-1)-s\right)|b|^{q-m} b v^{s-1} z^{\frac{q+4-m}{2}} \leq C_{\varepsilon} \frac{|\nabla z|^{2}}{z}
\end{align*}
$$

Proof. The following identities hold if $u=v^{b}: \nabla u=b v^{b-1} \nabla v$,

$$
\begin{gathered}
|\nabla u|^{m-2} \nabla u=|b|^{m-2} b v^{(b-1)(m-1)}|\nabla v|^{m-2} \nabla v, \\
\Delta_{m} u=|b|^{m-2} b\left(v^{(b-1)(m-1)} \Delta_{m} v+(b-1)(m-1) v^{(b(m-1)-m}|\nabla v|^{m}\right), \\
-v^{(b-1)(m-1)} \Delta_{m} v=(b-1)(m-1) v^{(b(m-1)-m}|\nabla v|^{m}+|b|^{q} v^{b p+(b-1) q}|\nabla v|^{q},
\end{gathered}
$$

and finally

$$
\begin{equation*}
-\Delta_{m} v=(b-1)(m-1) \frac{|\nabla v|^{m}}{v}+|b|^{q-m} b v^{s}|\nabla v|^{q} \tag{4.3}
\end{equation*}
$$

with

$$
\begin{equation*}
s=m-1-q+b(p+q-m+1) \tag{4.4}
\end{equation*}
$$

We set $z=|\nabla v|^{2}$. Then in $G$,

$$
-\Delta_{m} v=f \Longleftrightarrow-\Delta v-(m-2) \frac{D^{2} v(\nabla v, \nabla v)}{|\nabla v|^{2}}=f|\nabla v|^{2-m},
$$

from which identity we infer

$$
-\Delta v=(m-2) \frac{D^{2} v(\nabla v, \nabla v)}{|\nabla v|^{2}}+(b-1)(m-1) \frac{|\nabla v|^{2}}{v}+|b|^{q-m} b v^{s}|\nabla v|^{q+2-m}
$$

where

$$
<H e s s ~ v(\nabla v), \nabla v>=D^{2} v(\nabla v, \nabla v)=\frac{1}{2}\langle\nabla z, \nabla v\rangle
$$

We recall the Böchner formula combined with Cauchy-Schwarz inequality,

$$
\left.-\frac{1}{2} \Delta z+\frac{1}{N}(\Delta v)^{2}+<\nabla(\Delta v), \nabla v\right) \leq-\frac{1}{2} \Delta z+(\text { Hess } v)^{2}+\langle\nabla(\Delta v), \nabla v>=0
$$

Since

$$
-\Delta v=\frac{m-2}{2} \frac{\langle\nabla z, \nabla v>}{z}+(b-1)(m-1) \frac{z}{v}+|b|^{q-m} b v^{s} z^{\frac{q+2-m}{2}},
$$

we deduce

$$
\begin{aligned}
<\nabla(\Delta v), \nabla v> & =-\frac{m-2}{2}<\nabla \frac{<\nabla z, \nabla v>}{z}, \nabla v>+(1-b)(m-1)<\nabla \frac{z}{v}, \nabla v> \\
& -|b|^{q-m} b\left(s v^{s-1} z^{\frac{q+4-m}{2}}+\frac{q+2-m}{2} v^{s} z^{\frac{q-m}{2}}<\nabla z, \nabla v>\right) .
\end{aligned}
$$

we observe that

$$
<\nabla \frac{z}{v}, \nabla v>=\frac{\langle\nabla z, \nabla v>}{v}-\frac{z^{2}}{v^{2}} \quad \text { and } \quad \frac{\left\langle\nabla z, \nabla v>^{2}\right.}{z^{2}} \leq \frac{|\nabla z|^{2}}{z}
$$

thus

$$
\begin{aligned}
-\frac{m-2}{2}<\nabla \frac{<\nabla z, \nabla v>}{z}, \nabla v> & =-\frac{m-2}{2}\left(\frac{D^{2} z(\nabla v, \nabla v)}{z}+\frac{1}{2} \frac{|\nabla z|^{2}}{z}-\frac{\left\langle\nabla z, \nabla v>^{2}\right.}{z^{2}}\right) \\
& \geq-\frac{m-2}{2} \frac{D^{2} z(\nabla v, \nabla v)}{z}-|m-2| \frac{|\nabla z|^{2}}{z} .
\end{aligned}
$$

We define the operator $\mathcal{A}$ by (4.1); it satisfies (3.3) with $\theta=\min (1, m-1)$ and $\Theta=\max (1, m-1)$, so it is uniformly elliptic in $G$. Therefore

$$
\begin{align*}
-\frac{1}{2} \mathcal{A}(z) & +\frac{1}{N}(\Delta v)^{2}-(1-b)(m-1) \frac{z^{2}}{v^{2}}-|b|^{q-m} b s v^{s-1} z^{\frac{q+4-m}{2}} \\
& \leq(b-1)(m-1) \frac{<\nabla z, \nabla v>}{v}+\frac{(q+2-m)|b|^{q-m} b}{2} v^{s} z^{\frac{q-m}{2}}<\nabla z, \nabla v>  \tag{4.5}\\
& +|m-2| \frac{|\nabla z|^{2}}{z}
\end{align*}
$$

For $\varepsilon>0$ there holds by Hölder's inequality,

$$
\begin{aligned}
& \frac{q+2-m}{2} v^{s} z^{\frac{q-m}{2}}<\nabla z, \nabla v>\leq \frac{\varepsilon}{N}|b|^{2(q-m+1)} v^{2 s} z^{q+2-m}+C_{\varepsilon} \frac{|\nabla z|^{2}}{z}, \\
(\Delta v)^{2}= & \left(\frac{m-2<\nabla z, \nabla v>}{2}+(b-1)(m-1) \frac{z}{v}+|b|^{q-m} b v^{s} z^{\frac{q+2-m}{2}}\right)^{2} \\
\geq & (b-1)^{2}(m-1)^{2} \frac{z^{2}}{v^{2}}+|b|^{2(q-m+1)} v^{2 s} z^{q+2-m}+2(b-1)(m-1)|b|^{q-m} b v^{s-1} z^{\frac{q+4-m}{2}} \\
& -(m-2) \frac{|\nabla z|}{\sqrt{z}}\left(|b-1|(m-1) \frac{z}{v}+|b|^{q-m+1} v^{s} z^{\frac{q+2-m}{2}}\right),
\end{aligned}
$$

and for any $\varepsilon>0$,

$$
\begin{gathered}
(m-2) \frac{|\nabla z|}{\sqrt{z}}\left(|b-1|(m-1) \frac{z}{v} \leq \frac{\varepsilon}{N}(b-1)^{2}(m-1)^{2} \frac{z^{2}}{v^{2}}+C_{\varepsilon} \frac{|\nabla z|^{2},}{z}\right. \\
(m-2) \frac{|\nabla z|}{\sqrt{z}}|b|^{q-m+1} v^{s} z^{\frac{q+2-m}{2}} \leq \frac{\varepsilon}{N}|b|^{2(q-m+1)} v^{2 s} z^{q+2-m}+C_{\varepsilon} \frac{|\nabla z|^{2}}{z}
\end{gathered}
$$

thus (4.2) follows.
Proof of Theorem 1.1. We use Lemma 4.1 with $b \in(0,1)$, combined with the estimate

$$
\left(\frac{1}{N} 2(b-1)(m-1)-s\right)|b|^{q-m} b v^{s-1} z^{\frac{q+4-m}{2}} \leq \frac{\varepsilon}{N}|b|^{2(q-m+1)} v^{2 s} z^{q+2-m}+C_{\varepsilon} \frac{z^{2}}{v^{2}}
$$

Then there exist constants $C_{i}>0$ depending only on $m, b, N, p, q$, such that

$$
\frac{1}{2} \mathcal{A}(z)+C_{1} v^{2 s} z^{q+2-m} \leq C_{2} \frac{z^{2}}{v^{2}}+C_{3} \frac{|\nabla z|^{2}}{z}
$$

Next we choose $s=-1$ in (4.4), thus

$$
b(p+q-m+1)=q-m
$$

which is positive because $q>m$. We deduce using Hölder inequality,

$$
\mathcal{A}(z)+\frac{C_{4} z^{q+2-m}-C_{5}}{v^{2}} \leq \mathcal{A}(z)+\frac{C_{1} z^{q+2-m}-C_{2} z^{2}}{v^{2}} \leq C_{3} \frac{|\nabla z|^{2}}{z}
$$

If we apply Lemma 3.1 with

$$
\alpha(x)=\frac{C_{4}}{v^{2}(x)}, \quad \beta(x)=\frac{C_{5}}{v^{2}(x)}, \quad k=q+2-m
$$

we deduce that any solution in $\bar{B}_{\rho}\left(x_{0}\right) \rho>0$, satisfies

$$
\left.z\left(x_{0}\right) \leq C_{6}\left(\frac{1}{\alpha \rho^{2}}\right)^{\frac{1}{k-1}}+\left(\frac{C_{5}}{C_{4}}\right)^{\frac{1}{k}} \leq C_{7}\left(\frac{\max _{B_{\rho}\left(x_{0}\right)} v}{\rho}\right)^{\frac{2}{q+1-m}}+1\right)
$$

which yields

$$
\begin{equation*}
\left|\nabla v\left(x_{0}\right)\right| \leq C_{8}\left(\left(\frac{\max _{\bar{B}_{\rho}\left(x_{0}\right)} v}{\rho}\right)^{\frac{1}{q+1-m}}+1\right) \tag{4.6}
\end{equation*}
$$

where we observe that $\frac{1}{q+1-m}<1$, since $q>m$. Let $\varepsilon \in\left(0, \frac{1}{2}\right]$. As a consequence, for any solution in $B_{2 R}$, (or even $\bar{B}_{\frac{3 R}{2}}$ ) considering any $x_{0} \in \bar{B}_{R}$ and taking $\rho=R \varepsilon$, we get

$$
\begin{gather*}
\max _{\bar{B}_{R}}|\nabla v| \leq c\left(\left(\frac{\max _{\bar{B}_{R(1+\varepsilon)}} v}{\varepsilon R}\right)^{\frac{1}{q+1-m}}+1\right) \leq c \varepsilon^{-\frac{1}{q+1-m}}\left(\left(\frac{\max _{\bar{B}_{R(1+\varepsilon)}} v}{R}\right)^{\frac{1}{q+1-m}}+1\right)  \tag{4.7}\\
\max _{\bar{B}_{R(1+\varepsilon)}} v \leq v(0)+R(1+\varepsilon) \max _{\bar{B}_{R(1+\varepsilon)}}|\nabla v| \\
\frac{\max _{\bar{B}_{R(1+\varepsilon)}} v}{R} \leq \frac{1+v(0)}{R}+(1+\varepsilon) \max _{\bar{B}_{R(1+\varepsilon)}}|\nabla v| \leq c_{0}\left(\frac{1}{R}+\max _{\bar{B}_{R(1+\varepsilon)}}|\nabla v|\right)
\end{gather*}
$$

where $c_{0}=2+v(0)$ depends on $v(0)$. If $R \geq 1$,

$$
\left.\left(\frac{\max _{\bar{B}_{R(1+\varepsilon)}} v}{R}\right)^{\frac{1}{q+1-m}}+1 \leq c_{0}^{\frac{1}{q+1-m}}\left(1+\max _{\bar{B}_{R(1+\varepsilon)}}|\nabla v|\right)^{\frac{1}{q+1-m}}\right)+1 \leq c_{1}\left(1+\max _{\bar{B}_{R(1+\varepsilon)}}|\nabla v|\right)^{\frac{1}{q+1-m}}
$$

Then from (4.7),
$y(R):=1+\max _{\bar{B}_{R}}|\nabla v| \leq 1+c_{2} \varepsilon^{-\frac{1}{q+1-m}}\left(1+\max _{\bar{B}_{R(1+\varepsilon)}}|\nabla v|\right)^{\frac{1}{q+1-m}} \leq c_{3} \varepsilon^{-\frac{1}{q+1-m}}\left(1+\max _{\bar{B}_{R(1+\varepsilon)}}|\nabla v|\right)^{\frac{1}{q+1-m}}$.

Using the definition of $y$, this is

$$
y(R) \leq c \varepsilon^{-\frac{1}{q+1-m}}(y((1+\varepsilon) R))^{\frac{1}{q+1-m}},
$$

where $c$ depends on $v(0)$. Therefore $y(R)$ is bounded as a consequence of Lemma 3.2. Thus $|\nabla v|$ is bounded and using the definition of $v$ with the value of $b$,

$$
|\nabla v|^{q-m}=u^{p+1}|\nabla u|^{q-m} \in L^{\infty}\left(\mathbb{R}^{N}\right) .
$$

Next we consider any $l \geq 0$ such that $u-l>0$. The function $u_{l}=u-l$ satisfies

$$
0 \leq-\Delta_{m} u_{l} \leq C_{\infty} \frac{|\nabla u|^{m}}{u} \leq C_{\infty} \frac{\left|\nabla u_{l}\right|^{m}}{u_{l}}
$$

with $C_{\infty}=\left\|u^{p+1}|\nabla u|^{q-m}\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}$. Then the function $w_{l}=u_{l}^{\sigma}$ with $\sigma>1$ to be specified below, satisfies

$$
\begin{aligned}
-\Delta_{m} w_{l} & =\sigma^{m-1} u_{l}^{(\sigma-1)(m-1)}\left(-\Delta_{m} u_{l}+(\sigma-1)(m-1) \frac{\left|\nabla u_{l}\right|^{m}}{u_{l}}\right) \\
& \leq \sigma^{m-1}\left((\sigma-1)(m-1)-C_{\infty}\right) u_{l}^{\sigma(m-1)-m}\left|\nabla u_{l}\right|^{m}
\end{aligned}
$$

Therefore $w_{l}$ is $m$-subharmonic for $\sigma$ large enough.
We first take $l=0$, so $w=u^{\sigma}$. By Lemma 3.3, for any $\tau>0$, there exists a constant $C_{\tau}=$ $C_{\tau}(N, m, \tau)$ such that

$$
\begin{equation*}
\sup _{B_{R}} w \leq C_{\tau}\left(\frac{1}{\left|B_{2 R}\right|} \int_{B_{2 R}} w^{\tau}\right)^{\frac{1}{\tau}}=C_{\tau}\left(\frac{1}{\left|B_{2 R}\right|} \int_{B_{2 R}} u^{\tau \sigma}\right)^{\frac{1}{\tau}}, \tag{4.8}
\end{equation*}
$$

and since $u$ is $m$-superharmonic, there holds for any $\theta \in\left(0, \frac{N(m-1)}{N-m}\right),[34]$

$$
\begin{equation*}
\inf _{B_{R}} u \geq c_{\theta}\left(\frac{1}{\left|B_{2 R}\right|} \int_{B_{2 R}} u^{\theta}\right)^{\frac{1}{\theta}} \tag{4.9}
\end{equation*}
$$

Taking $\tau=\frac{\theta}{\sigma}$, we deduce

$$
\begin{equation*}
\sup _{B_{R}} u=\left(\sup _{B_{R}} w\right)^{\frac{1}{\sigma}} \leq C_{\tau}^{\frac{1}{\sigma}}\left(\frac{1}{\left|B_{2 R}\right|} \int_{B_{2 R}} u^{\tau \sigma}\right)^{\frac{1}{s \sigma}} \leq \frac{C_{\tau}^{\frac{1}{\sigma}}}{c_{\theta}} \inf _{B_{R}} u . \tag{4.10}
\end{equation*}
$$

This means that $u$, and also $w$, satisfies the Harnack inequality in $\mathbb{R}^{N}$ :

$$
\sup _{B_{R}} w \leq \frac{C_{\tau}}{c_{\theta}^{\sigma}} \inf _{B_{R}} w
$$

But $r \mapsto \mu(r)=\inf _{|x|=r} u=\inf _{B_{r}} u$ from the maximum principle, is nonincreasing, so it has a limit $L \geq 0$ as $r \rightarrow \infty$. This implies that $u$ is bounded and $l=\inf _{\mathbb{R}^{N}} u \geq 0$. If we replace $u$ by $u_{l}$ and $w$ by $w_{l}$, then (4.8) holds with $w$ and $u$ replaced respectively by $w_{l}$ and $u_{l}$ since $w_{l}$ is m-subharmonic, but also (4.9) holds with $u$ replaced by $u_{l}$ since $u_{l}$ is m-superharmonic. Thus

$$
\sup _{B_{R}} w_{l} \leq C\left(\inf _{B_{R}} u_{l}\right)^{\sigma} .
$$

Therefore $\sup _{B_{R}} w_{l}$ tends to 0 as $R \rightarrow \infty$. Then $w_{l} \equiv 0$, thus $u \equiv l$.

### 4.2 Asymptotic behaviour near $\infty$

In this section we consider the behaviour of solutions defined in an exterior domain.
Proof of Theorem 1.2. Consider a nonnegative solution $u=v^{b}(0<b<1)$ of (1.1) in $\mathbb{R}^{N} \backslash B_{r_{0}}$. From (4.6) the function $v$ satisfies in $\bar{B}_{\rho}\left(x_{0}\right)(\rho>0)$,

$$
\begin{equation*}
\left|\nabla v\left(x_{0}\right)\right| \leq C\left(\left(\frac{\max _{\bar{B}_{\rho}\left(x_{0}\right)} v}{\rho}\right)^{\frac{1}{q+1-m}}+1\right) \tag{4.11}
\end{equation*}
$$

with $C=C(N, p, q, m)$. Here we denote by $c_{i}$ some positive constants depending on $r_{0}, N, p, q, m$. Let $R>4 r_{0}$ and $0<\varepsilon \leq \frac{1}{4}$. Applying (4.11) with $\rho=\varepsilon R$, we get

$$
\left|\nabla v\left(x_{0}\right)\right| \leq c_{1}\left(\left(\frac{\max _{\bar{B}_{\varepsilon R}\left(x_{0}\right)} v}{\varepsilon R}\right)^{\frac{1}{q+1-m}}+1\right) \leq c_{1} \varepsilon^{-\frac{1}{q+1-m}}\left(\left(\frac{\max _{\bar{B}_{\varepsilon R}\left(x_{0}\right)} v}{R}\right)^{\frac{1}{q+1-m}}+1\right)
$$

then

$$
\begin{gathered}
\max _{|x|=R}|\nabla v| \leq c_{2}\left(\left(\frac{\max _{R\left(1-\frac{\varepsilon}{2}\right) \leq|x| \leq R\left(1+\frac{\varepsilon}{2}\right) v}}{\varepsilon R}\right)^{\frac{1}{q+1-m}}+1\right) \leq c_{3}\left(\left(\frac{\max _{\frac{R}{1+\varepsilon} \leq|x| \leq R(1+\varepsilon)^{v}} v}{\varepsilon R}\right)^{\frac{1}{q+1-m}}+1\right), \\
\max _{\frac{R}{2} \leq|x| \leq 2 R}|\nabla v| \leq c_{4}\left(\left(\frac{\max \frac{R}{2(1+\varepsilon)} \leq|x| \leq 2 R(1+\varepsilon)^{v}}{\varepsilon R}\right)^{\frac{1}{q+1-m}}+1\right),
\end{gathered}
$$

and finally,

$$
1+\max _{\frac{R}{2} \leq|x| \leq 2 R}|\nabla v| \leq c_{5} \varepsilon^{-\frac{1}{q+1-m}}\left(\left(\frac{\max \frac{R}{2(1+\varepsilon)} \leq|x| \leq 2 R(1+\varepsilon)^{v}}{R}\right)^{\frac{1}{q+1-m}}+1\right)
$$

From Lemma 3.4-(i), $\mu(r)=\inf _{|x|=r} u=\left(\inf _{|x|=r} v\right)^{b}$ is bounded : let $M=\max _{r \geq r_{0}} \mu(r)$. Note that $M$ depends on $u$. Now for any $x$ such that $|x|=\rho$, there exists at least one point $x_{\rho}$ where $v\left(x_{\rho}\right)=\inf _{|x|=\rho} v$. We can join any point $x \in S_{\rho}$ to $x_{\rho}$ by a connected chain of balls of radius $\varepsilon \rho$ with at points $x_{i} \in S_{\rho}$ and this chain can be constructed so that it has at most $\frac{\pi}{\varepsilon}$ elements. Considering one ball containing $x$ and joining it to a ball containing $x_{\rho}$, we get that

$$
v(x) \leq v\left(x_{\rho}\right)+C_{N} \varepsilon^{-1} \rho \max _{\frac{\rho}{1+\varepsilon} \leq|x| \leq \rho(1+\varepsilon)}|\nabla v| \leq M^{\frac{1}{b}}+C_{N} \varepsilon^{-1} \rho \max _{\frac{\rho}{1+\varepsilon} \leq|x| \leq \rho(1+\varepsilon)}|\nabla v| .
$$

Then

$$
\begin{aligned}
\max _{\frac{R}{2(1+\varepsilon)} \leq|x| \leq 2 R(1+\varepsilon)} v & \leq c_{M}^{1}\left(1+\varepsilon^{-1} R \max _{\frac{R}{2(1+3 \varepsilon)} \leq|x| \leq 2 R(1+3 \varepsilon)}|\nabla v|\right) \\
& \leq c_{M}^{2} \varepsilon^{-1} R\left(1+\max _{\frac{R}{2(1+3 \varepsilon)} \leq|x| \leq 2 R(1+3 \varepsilon)}|\nabla v|\right)
\end{aligned}
$$

and

$$
\frac{1}{\varepsilon R} \max _{\frac{R}{2(1+\varepsilon)} \leq|x| \leq 2 R(1+\varepsilon)} v \leq c_{M}^{3} \varepsilon^{-2}\left(1+\max _{\frac{R}{2(1+3 \varepsilon)} \leq|x| \leq 2 R(1+3 \varepsilon)}|\nabla v|\right)
$$

Using estimate (4.7) we obtain

$$
\begin{equation*}
1+\max _{\frac{R}{2} \leq|x| \leq 2 R}|\nabla v| \leq c_{M}^{4} \varepsilon^{-\frac{2}{q+1-m}}\left(1+\max _{\frac{R}{2(1+3 \varepsilon)} \leq|x| \leq 2 R(1+3 \varepsilon)}|\nabla v|\right)^{\frac{1}{q+1-m}} . \tag{4.12}
\end{equation*}
$$

Let $\left\{\varepsilon_{n}\right\}_{n \geq 1}$ be a positive decreasing sequence such that $P_{n}:=\prod_{j=1}^{n}\left(1+\varepsilon_{j}\right) \rightarrow 2$ and $\Theta_{n}:=$ $\prod_{j=1}^{n} \varepsilon_{j+1}^{d^{j}} \rightarrow \Theta>0$ when $n \rightarrow \infty$. It is easy to find such sequences such that $\varepsilon_{j} \sim 2^{-j}$. For $\frac{R}{2} \leq a<2 R \leq b$ we set

$$
y(a, b)=1+\max _{a \leq|x| \leq b}|\nabla v| .
$$

Then (4.12) with $(a, b)=\left(\frac{R}{2}, 2 R\right)$ and $\varepsilon_{1}=3 \varepsilon$ asserts that

$$
y\left(\frac{R}{2}, 2 R\right) \leq c_{5} \varepsilon_{1}^{-h}\left(y\left(\frac{R}{2\left(1+\varepsilon_{1}\right)}, 2 R\left(1+\varepsilon_{1}\right)\right)\right)^{d} \quad \text { with } h=\frac{2}{q+1-m} \text { and } d=\frac{1}{q+1-m} \in(0,1) .
$$

Applying (4.12) with $(a, b)=\left(\frac{R}{2 P_{n}}, 2 R P_{n}\right)$ we obtain

$$
y\left(\frac{R}{2 P_{n}}, 2 R P_{n}\right) \leq c_{5} \varepsilon_{n+1}^{-h}\left(y\left(\frac{R}{2 P_{n+1}}, 2 R P_{n+1}\right)\right)^{d} .
$$

By induction, we deduce

$$
\begin{equation*}
y\left(\frac{R}{2}, 2 R\right) \leq c_{5}^{1+d+d^{2}+\ldots+d^{n}} \Theta_{n}^{-h}\left(y\left(\frac{R}{2 P_{n+1}}, 2 R P_{n+1}\right)\right)^{d^{n+1}} \tag{4.13}
\end{equation*}
$$

Since $y\left(\frac{R}{2 P_{n+1}}, 2 R P_{n+1}\right) \rightarrow y\left(\frac{R}{4}, 4 R\right)$, we obtain that

$$
\begin{equation*}
1+\max _{\frac{R}{2} \leq|x| \leq 2 R}|\nabla v| \leq c_{5}^{\frac{d}{1-d}} \Theta^{-h}:=C(M, N, p, q, m) . \tag{4.14}
\end{equation*}
$$

Then we conclude again that $|\nabla v|$ is bounded for $|x| \geq 4 r_{0}$, then in $\mathbb{R}^{N} \backslash B_{r_{0}}$ since we have assumed that $u \in C^{1} \mathbb{R}^{N} \backslash B_{r_{0}}$. We consider again the function $w=u^{\sigma}$, for $\sigma$ depending of $r_{0}$, large enough so that $(\sigma-1)(m-1) \geq\left\|u^{p+1}|\nabla u|^{q-m}\right\|_{L^{\infty}\left(\mathbb{R}^{N} \backslash B_{r_{0}}\right)}$. As in the proof of Theorem 1.1 we conclude that $w$ is $m$-subharmonic in $\mathbb{R}^{N} \backslash \bar{B}_{r_{0}}$. Hence $u$ satisfies the Harnack inequality, using the estimate (4.10). Therefore, for any $R>2 r_{0}$,

$$
\sup _{\frac{R}{2} \leq|x| \leq 3 \frac{R}{2}} u \leq C \inf _{\frac{R}{2} \leq|x| \leq 3 \frac{R}{2}} u .
$$

Since $u$ is is $m$-superharmonic, it follows by the strong maximum principle, that it cannot have any local minimum in $\mathbb{R}^{N} \backslash \bar{B}_{r_{0}}$. Since $u^{\sigma}$ is m-subharmonic it cannot have any local maximum too, and $u$ shares this property. As a consequence $|\nabla u|$ does not vanish in $\mathbb{R}^{N} \backslash \bar{B}_{r_{0}}$. The function $r \mapsto \mu$ is bounded by Lemma 3.4, hence $u$ is also bounded by the above Harnack inequality. Finally $\mu(r)$ is monotone for large $r$, so it admits a limit $l \geq 0$ when $r \rightarrow \infty$.

If $\mu(r)$ is nonincreasing for $r \geq r_{1}>r_{0}$, then $u-l \geq 0$, so we can consider the function $w_{l}$ instead of $w$. Then

$$
\max _{R \leq|x| \leq 2 R} w_{l} \leq C\left(\inf _{R \leq|x| \leq 2 R}(u-l)\right)^{\sigma},
$$

Then $w_{l}$ tends to 0 , thus $u$ tends to $l$ as $|x| \rightarrow \infty$. Since $u-l$ is $m$-superharmonic in $\mathbb{R}^{N} \backslash B_{r_{0}}$, then there holds

$$
u(x)-l \geq C|x|^{\frac{m-N}{m-1}}
$$

with $C=C\left(r_{0}, N, m, u\right)$, see for example [5, Proposition 2.6], [33, Lemma 2.3]. It is the case in particular when $u$ is a solution in $\mathbb{R}^{N} \backslash\{0\}$. Note that the radial solutions such that $\mu$ is nonincreasing are precisely defined in $(0, \infty)$.

Now, it follows from the upper estimate of $y(R)$, that the function $u$ satisfies

$$
-\Delta_{m} u=u^{p}|\nabla u|^{q} \leq C \frac{|\nabla u|^{m}}{u},
$$

in $\mathbb{R}^{N} \backslash B_{r_{0}}$. Next suppose that $l>0$. Then

$$
-\Delta_{m} u \leq C^{\prime}|\nabla u|^{m} .
$$

The function $U$ (still used in case $q=m$ ), defined by

$$
U=(m-1)\left(e^{\frac{u-l}{m-1}}-1\right)
$$

satisfies $-\Delta_{m} U \leq 0$ and $U$ tend to 0 at $\infty$. Then there exists $R_{\varepsilon}>0$ such that $U(x) \leq \varepsilon$ for $|x| \geq R_{\varepsilon}$. For $R>R_{\varepsilon}$, the function $x \mapsto \omega(x):=\varepsilon+\left(\sup _{|z|=r_{0}} U(z)\right)\left(\frac{|x|}{r_{0}}\right)^{\frac{m-N}{m-1}}$ is a $m$-harmonic in $B_{R} \backslash B_{r_{0}}$, hence it is larger than $U$. Letting $\varepsilon \rightarrow 0$ we get $U \leq C|x|^{\frac{m-N}{m-1}}$ near $\infty$; and $U$ has the same behaviour as $u-l$, so we deduce the estimate from above,

$$
u(x)-l \leq C|x|^{\frac{m-N}{m-1}} .
$$

Then we get the estimate (1.7).
Remark 4.2 (i) In case $u$ is defined in $\mathbb{R}^{N} \backslash\{0\}$ and $l=0$, we obtain the estimates

$$
C_{1}|x|^{\frac{m-N}{m-1}} \leq u(x) \leq C_{2}|x|^{\frac{1}{\sigma} \frac{m-N}{m-1}} .
$$

It would be interesting to improve the estimate from above.
(ii) If $u$ is defined in $\mathbb{R}^{N} \backslash B_{r_{0}}$ and if $\mu$ is nonincreasing, we have proved that $u$ has a limit $l \geq 0$ as $|x| \rightarrow \infty$. If $\mu$ is nondecreasing, we only obtain that $\mu(r)=\inf _{|x|=r} u(x)$ has a limit $l$, and $\sup _{|x|=r} u(x)$ has a limit $\lambda \geq l$. Indeed the function $w$ is $m$-subharmonic positive and bounded, so the function $r \mapsto \sup _{|x|=r} w=\left(\sup _{|x|=r} u\right)^{\sigma}$ is also monotone for large $r$ and has a limit $\lambda^{\sigma}$. We have $w=u^{\sigma} \leq \lambda^{\sigma}$, so $\sup _{|x|=r} u$ is also nondecreasing. But we cannot prove that $\lambda=l$.

### 4.3 Behaviour near an isolated singularity

In this section we study the behaviour of solutions with an isolated singularity at the origin.
Proof of Theorem 1.3 Let $u$ be a nonnegative solution $u$ of (1.1) in $B_{r_{0}} \backslash\{0\}$. We apply directly the Bernstein method to $u$ : we obtain by Lemma 4.1 with $b=1$, and then $s=p$. Setting $\xi=|\nabla u|^{2}$, we get

$$
\frac{1}{2} \mathcal{A}(\xi)+C_{1} u^{2 p} z^{q+2-m} \leq C_{2} \frac{\xi^{2}}{u^{2}}+C_{3} \frac{|\nabla \xi|^{2}}{\xi}
$$

By the strong maximum principle, there exists a constant $a_{0}>0$ depending on $r_{0}$ and $N, p, q$, such that $u \geq a_{0}$ in $B_{\frac{r_{0}}{2}} \backslash\{0\}$. Therefore, there holds

$$
\frac{1}{2} \mathcal{A}(\xi)+C_{1}^{2 p} a_{0}^{2 p} z^{q+2-m} \leq C_{2} \frac{\xi^{2}}{a_{0}^{2}}+C_{3} \frac{|\nabla \xi|^{2}}{\xi}
$$

in $B_{\frac{r_{0}}{2}} \backslash\{0\}$. Then from Lemma 3.1, we deduce the inequality

$$
z\left(x_{0}\right) \leq c\left(\left(\frac{1}{a_{0}^{2 p} \rho^{2}}\right)^{\frac{1}{q+1-m}}+\left(\frac{1}{a_{0}^{2(p+1)}}\right)^{\frac{1}{q+2-m}}\right) \leq c_{0}^{2}\left(\frac{1}{\rho^{\frac{2}{q+1-m}}}+1\right)
$$

for any ball $\bar{B}_{\rho}\left(x_{0}\right) \subset B_{\frac{r_{0}}{2}} \backslash\{0\}$, with $c=c(N, p, q, m)$ and $c_{0}^{2}=c\left(a_{0}^{-\frac{p}{q+1-m}}+a_{0}^{-\frac{p+1}{q+2-m}}\right)$. Hence for any $x \in \bar{B}_{R} \backslash\{0\}$, with $R \leq \min \left(1, \frac{r_{0}}{8}\right)$,

$$
|\nabla u(x)| \leq \frac{2 c_{0}}{|x|^{\frac{1}{q+1-m}}} \leq \frac{2 c_{0}}{|x|^{\frac{1}{q+1-m}}}
$$

As a consequence, considering any $x_{R}$ such that $x, x^{\prime} \in \bar{B}_{\frac{R}{2}} \backslash\{0\}$, there holds

$$
\left|u\left(x^{\prime}\right)-u(x)\right| \leq 2 c_{0} R^{\frac{q-m}{q+1-m}} .
$$

Since $q>m, u$ is bounded near 0 . Then, with constants $C>0$ depending on $a_{0}$,

$$
-\Delta_{m} u=f \leq C|\nabla u|^{q} \leq C|x|^{-\frac{q}{q+1-m}} .
$$

Then $f \in L_{l o c}^{\frac{N}{m}+\varepsilon}\left(B_{\frac{R}{4}}\right)$, since $N-\frac{N}{m} \frac{q}{q+1-m}=\frac{N(m-1)(q-m)}{m(q+1-m)}>0$. Thus from [31] $u$ can be extended as a continuous function, solution of the equation in the sense of distributions. Then we deduce that for any $x \in \bar{B}_{\frac{R}{2}} \backslash\{0\}$,

$$
|u(0)-u(x)| \leq c_{0}|x|^{\frac{q-m}{q+1-m}}
$$

Moreover, replacing $r_{0}$ by $\rho>0$ small enough such that $u(x) \geq \frac{u(0)}{2}$ in $B_{\rho}$, then $a_{0} \geq \frac{u(0)}{2}$, hence $c_{0} \leq C(N, p, q, m, u(0))$, and for $|x| \leq \min \left(1, \frac{\rho}{8}\right)$, we infer

$$
|u(0)-u(x)| \leq C|x|^{\frac{q-m}{q+1-m}} .
$$

Next assume that $u$ is defined in $\mathbb{R}^{N} \backslash\{0\}$ and is not constant. Then $u$ is bounded, since it is bounded near 0 and $\infty$. Then $r \mapsto \mu(r)=\inf _{|x|=r} u$ is nonincreasing, thus $\mu(r) \leq u(0)$ : indeed $\forall \varepsilon>0$, we have $\mu(|x|) \leq u(x) \leq u(0)+\varepsilon$ for any $|x| \leq r_{\varepsilon}$, then from the monotone decreasingness, $\mu(r) \leq u(0)+\varepsilon$ for any $r>0$. From Theorem $1.2 \lim _{|x| \rightarrow \infty} u=l=\lim _{r \rightarrow \infty} \mu$. and then necessarily $l \leq u(0)$. Suppose that there exists $x \neq 0$ such that $u(x)>u(0)$; then $u$ has a maximum in $\mathbb{R}^{N} \backslash\{0\}$, but $|\nabla u|$ cannot vanish in $\mathbb{R}^{N} \backslash\{0\}$ by Theorem 1.2 , so we get a contradiction.

## 5 Radial case

If $u$ is a positive radial solution of (1.1), and if we denote for simplicity $u(r)=u(x)$ with $r=|x|$, then $u$ satisfies the following o.d.e.

$$
\begin{equation*}
\left(\left|u^{\prime}\right|^{m-2} u^{\prime}\right)^{\prime}+\frac{N-1}{r}\left|u^{\prime}\right|^{m-2} u^{\prime}+u^{p}\left|u^{\prime}\right|^{q}=r^{1-N}\left(r^{N-1}\left|u^{\prime}\right|^{m-2} u^{\prime}\right)^{\prime}+u^{p}\left|u^{\prime}\right|^{q}=0 . \tag{5.1}
\end{equation*}
$$

We begin with a simple observation about the set of zeros of $u^{\prime}$. We have shown above that any solution of the exterior problem is either constant, or its gradient does not vanish. In the radial case, the proof is elementary:

Proposition 5.1 Assume $q>m-1, p \geq 0$. Then any nonnegative radial solution of (5.1) on a segment $\left[r_{1}, r_{2}\right] \subset(0, \infty)$ is constant, or strictly monotone.

Proof. By the strong maximum principle [35], we can assume that $u>0$ on $\left(r_{1}, r_{2}\right)$. The function

$$
r \mapsto W(r):=r^{N-1}\left|u^{\prime}(r)\right|^{m-2} u^{\prime}(r)
$$

is nonincreasing. Suppose that $u^{\prime}$ has two zeros $\rho_{1}$ and $\rho_{2}$ in $\left(r_{1}, r_{2}\right)$, then by integrating $W^{\prime}$ and using the equation, we deduce that $u^{\prime} \equiv 0$ on $\left[\rho_{1}, \rho_{2}\right]$, hence $u$ is constant therein, therefore we can assume that $\left[\rho_{1}, \rho_{2}\right]$ is the maximal subinterval of $\left[r_{1}, r_{2}\right]$ where $u^{\prime}$ vanishes. If $\left[r_{1}, \rho_{1}\right] \neq\left[r_{1}, r_{2}\right]$, for example $r_{1}<\rho_{1}$, then $u^{\prime}>0$ on $\left(r_{1}, \rho_{1}\right)$ where $u^{\prime}(r)=\left(r^{1-N} W(r)\right)^{\frac{1}{m-1}}$. By (5.1),

$$
\begin{equation*}
\frac{m-1}{m-1-q}\left(W^{\frac{m-1-q}{m-1}}\right)^{\prime}=|W|^{-\frac{q}{m-1}} W^{\prime}=-r^{N-1-(N-1) \frac{q}{m-1}} u^{p} \tag{5.2}
\end{equation*}
$$

on $\left(r_{1}, \rho_{1}\right)$ and $\left.\lim _{r \rightarrow \rho_{1}} u^{\prime} r\right)=0$. Since $m-1-q<$ this implies $\left.\lim _{r \rightarrow \rho_{1}} W^{\frac{m-1-p}{m-1}} r\right)=\infty$, a contradiction since $u$ is bounded on $\left[r_{1}, \rho_{1}\right]$. We proceed similarly if $r_{1}=\rho_{1}$ but $\rho_{2}<r_{2}$ or if $\rho_{1}=\rho_{2}$. Hence either $u$ is constant or it is strictly monotone.

Next we make a complete description of the radial solutions for $p \geq 0, q>m$.

### 5.1 The case $p=0$

This case $p=0$ of the Hamilton-Jacobi equation is well known, since equation (5.2) can be directly integrated, so the solutions are explicit, and are a Ariadne's thread for studying the case $p>0$. We find different types of nonconstant solutions according to the sign of $u^{\prime}$ :

$$
\left\{\begin{array}{c}
u^{\prime}=r^{\frac{1-N}{m-1}}\left(C_{1}-a_{m, q}^{-1} r^{\left.-a_{m, q}\right)}\right)^{-\frac{1}{q-m+1}} \\
u^{\prime}=-r^{\frac{1-N}{m-1}}\left(C_{2}+a_{m, q}^{-1} r^{\left.-a_{m, q}\right)}\right)^{-\frac{1}{q-m+1}}
\end{array}\right.
$$

where $a_{m, q}=\frac{(N-1) q-N(m-1)}{m-1}>0$ since $q>\frac{N(m-1)}{N-1}>m-1$; and the value of $u$ follows by integration, with the requirement that $u>0$. The solutions such that $C_{1}>0$ satisfy $\lim _{r \rightarrow 0} r^{\frac{1}{q-m+1}} u^{\prime}(r)=$ $-a_{m, q}^{\frac{1}{q-m+1}}$, then $\lim _{r \rightarrow 0} u(r)=u_{0}>0$, since $q>m$. The conclusions of theorem 1.4 follow in that case.

### 5.2 The case $p>0$

Equation (5.1) can be reduced to an autonomous system, since it is invariant by the transformation $u \mapsto T_{\lambda} u(\lambda>0)$ given by

$$
\begin{equation*}
T_{\lambda} u(x)=\lambda^{-\frac{q-m}{p+q+1-m}} u(\lambda x) . \tag{5.3}
\end{equation*}
$$

Here we perform a change of unknown, introduced in [10], which consists in a differentiation of the equation, as in the Bernstein technique. We set

$$
\begin{equation*}
X(t)=-r \frac{u^{\prime}(r)}{u(r)}, \quad Z(t)=-r u^{p}\left|u^{\prime}\right|^{q-m} u^{\prime}, \quad t=\ln r, \tag{5.4}
\end{equation*}
$$

and obtain the following quadratic system of Kolmogorov type, valid any point $t$ where $u^{\prime}(t) \neq 0$, and any reals $m, p, q$,

$$
\left\{\begin{array}{ccc}
X_{t} & = & X\left(X-\frac{N-m}{m-1}+\frac{Z}{m-1}\right)  \tag{5.5}\\
Z_{t} & = & Z\left(N-\frac{N-1}{m-1} q-p X+\frac{q+1-m}{m-1} Z\right)
\end{array}\right.
$$

in the region $Q=\left\{(X, Z) \in \mathbb{R}^{2}: X Z>0\right\}$. Note that the trajectories $X=0$ and $Z=0$ are not admissible in our study. Since $p+q \neq m-1$, we can recover $u$ and $u^{\prime}$ by

$$
\begin{equation*}
u=\left(r^{q-m}|Z||X|^{m-1-q}\right)^{\frac{1}{p+q-m+1}}, \quad u^{\prime}=-\frac{X u}{r}=\left(r^{-(p+1)}|Z||X|^{p}\right)^{\frac{1}{p+q-m+1}} \operatorname{sign}(-X) . \tag{5.6}
\end{equation*}
$$

The fixed points of the system in $\bar{Q}$ are

$$
N_{0}=\left(0, a_{m, q}\right)=\left(0, \frac{(N-1) q-N(m-1)}{q+1-m}\right), \quad O=(0,0), \quad A_{0}=\left(\frac{N-m}{m-1}, 0\right) .
$$

We begin by a local study of the different points and the correponding results for the solutions of (5.1):

Lemma 5.2 (i) The point $N_{0}$ is a source, with eigenvalues $0<\lambda_{1}=\frac{q-m}{q+1-m}<\lambda_{2}=\frac{(N-1) q-N(m-1)}{m-1}$ and eigenvectors $v_{1}=\left(1, c_{m, q}\right)$ with $c_{m, q}>0$ and $v_{2}=(0,1)$. Then there exist infinitely many singular decreasing solutions $u$ of (5.1) defined near 0 , satisfying (1.11).
(ii) The point $O$ is a sink, with eigenvalues $0>\xi_{1}=-\frac{N-m}{m-1}>\xi_{2}=N-\frac{N-1}{m-1} q$, and eigenvectors $u_{1}=(1,0)$ and $u_{2}=(0,1)$. Then there exist infinitely many solutions $u$ of (1.1) defined for large $r$ and either increasing or decreasing near $\infty$, satisfying (1.12) and (1.13)
(iii) The point $A_{0}$ is a saddle point, with eigenvalues $\mu_{1}=-\frac{(N-m) p+(N-1) q-(m-1) N}{m-1}<0<\mu_{2}=$ $\frac{N-m}{m-1}$ and eigenvectors $w_{1}=\left(1,-d_{m, q}\right)$ with $d_{m, q}>0$ and $w_{2}=(1,0)$. Then for any $c>0$ there exists a unique solution $u$ of (5.1), defined at least for large $r$, such that $\lim _{r \rightarrow \infty} r^{\frac{N-m}{m-1}} u=c>0$.

Proof. (i) We perform the linearization at $N_{0}$ : setting $Z=a_{m, q}+\bar{Z}$, we get, with $X>0$ and $Z>0$,

$$
\left\{\begin{array}{ccc}
X_{t} & = & \frac{q-m}{q+1-m} X \\
\bar{Z}_{t} & = & a_{m, q}\left(-p X+\frac{q+1-m}{m-1} \bar{Z}\right)
\end{array}\right.
$$

which gives the eigenvalues $\lambda_{1}, \lambda_{2}$ and their respective eigenvectors, with the value of $c_{m, q}$

$$
c_{m, q}=\frac{p(N-1) q-N)(m-1)}{(N-1) q^{2}-2 N(m-1) q+(m-1)(N(m-1)+m)} .
$$

So $N_{0}$ is a source; the particular trajectory $X=0$ associated to $\lambda_{2}$ is not admissible. There exists an infinity of trajectories starting from $N_{0}$ as $t \rightarrow-\infty$, associated to the eigenvalue $\lambda_{1}$; the solutions $(X, Z)$ satisfy $X>0, \lim _{t \rightarrow-\infty} e^{-\frac{q-m}{q+1-m} t} X=C_{0}$, where $C_{0}>0$ is arbitrary and $\lim _{t \rightarrow-\infty} Z=a_{m, q}$; then from (5.6) and the definition of $Z$, there exist infinitely many decreasing singular solutions $u$ of (5.1) defined near 0 , satisfying (1.11).
(ii) The linearisation at $O$ gives the system

$$
\left\{\begin{array}{rlc}
X_{t} & = & -\frac{N-m}{m-1} X \\
Z_{t} & = & \left(N-\frac{N-1}{m-1} q\right) Z,
\end{array}\right.
$$

with admits the eigenvalues $\xi_{1}, \xi_{2}$. So $O$ is a sink, two particular trajectories are the axis $X=0$ and $Z=0$ which not admissible. There is an infinity of trajectories converging to $O$ as $t \rightarrow \infty$, tangent to the axis $Z=0$, associated to the eigenvalue $\xi_{1}$, with either $X, Z>0$, or $X, Z<0$. They satisfy

$$
\begin{equation*}
X \sim_{t \rightarrow \infty} C_{1} e^{-\frac{N-m}{m-1} t}, \quad Z=\sim_{t \rightarrow \infty} C_{2} e^{\left(N-\frac{N-1}{m-1} q\right) t}, \text { with } C_{1}, C_{2}>0 . \tag{5.7}
\end{equation*}
$$

The corresponding solutions $u$ of (5.1) are defined for large $r$, and either decreasing or increasing; from (5.6), we obtain $\lim _{r \rightarrow \infty} u=\left(C_{1}^{m-1-q} C_{2}\right)^{\frac{1}{p+q-m+1}}=l>0$ and $\lim _{r \rightarrow \infty} r^{\frac{N-1}{m-1}} u^{\prime}=-l C_{1}$, thus $\lim _{r \rightarrow \infty} r^{\frac{N-m}{m-1}}(u-l)=-l C_{1}$. Thus (1.12) and (1.13) follow. The uniqueness property follows from the uniqueness of a trajectory satisfying (5.7) for given $C_{1}, C_{2}$, see also Remark 5.3 below.
(iii) Linearisation at $A_{0}$ : setting $X=\frac{N-m}{m-1}+\bar{X}$, we get

$$
\left\{\begin{array}{ccc}
\bar{X}_{t} & = & \frac{N-m}{m-1}\left(\bar{X}+\frac{Z}{m-1}\right) \\
Z_{t} & = & -\frac{(N-m) p+(N-1)-(m-1) N}{m-1} Z
\end{array}\right.
$$

which admits the eigenvalues $\mu_{1}<0<\mu_{2}$ and the eigenvectors, with

$$
d_{m, q}=\frac{m-1}{N-m}\left(N-m+\left|\mu_{1}\right|(m-1)\right.
$$

It is a saddle point. The trajectory $X=0$ associated to $\mu_{2}$ is not admissible. Then a unique trajectory $\mathcal{T}_{A_{0}}$ converging to $A_{0}$ as $t \rightarrow \infty$. By the scaling (5.3), we deduce the uniqueness property for $u$.

Next we a complete description of the local and global solutions in the phase-plane leading to the conclusions of Theorem 1.4:

Proof of Theorem 1.4 when $p>0$. We consider the sets

$$
\begin{gathered}
\mathcal{L}_{X}=\left\{(X, Z) \in Q: X_{t}=0\right\}=\left\{(X, Z) \in Q: X-\frac{N-m}{m-1}+\frac{Z}{m-1}=0\right\}, \\
\mathcal{L}_{Z}=\left\{(X, Z) \in Q: Z_{t}=0\right\}=\left\{(X, X) \in Q: N-\frac{N-1}{m-1} q-p X+\frac{q+1-m}{m-1} Z\right\} .
\end{gathered}
$$

The straight line $\mathcal{L}_{X}$ has an extremity at $A_{0}$, with slope $-(m-1)$, and the slope of $\mathcal{T}_{A_{0}}$ is $-d_{m, q}<$ $-(m-1)$, so $\mathcal{T}_{A_{0}}$ is above $\mathcal{L}_{X}$ near $t=\infty$. The line $\mathcal{L}_{X}$ has an extremity at $N_{0}$, with slope $\frac{p(m-1)}{q+1-m}$,
is located above $\mathcal{L}_{X}$ for $X>0$. The trajectories issued from $N_{0}$ have the slope $c_{m, q}$, and we check that it is greater than $\frac{p(m-1)}{q+1-m}$ because $q>1-m$; so they start above $\mathcal{L}_{Z}$.
(i) The trajectory $\mathcal{T}_{A_{0}}$ stays in the region $\mathcal{R}=\left\{0<X<\frac{N-m}{m-1} ; X-\frac{N-m}{m-1}+\frac{Z}{m-1}>0\right\}$ which is negatively invariant. Then $X_{t}$ stays positive, thus $X$ is increasing, hence bounded. Either $\mathcal{T}_{A_{0}}$ stays under $\mathcal{L}_{Z}$, then $Z_{t}<0$ and $Z$ is bounded, thus $\mathcal{T}_{A_{0}}$ converges to $N_{0}$, or it crosses the line $\mathcal{L}_{Z}$ at time $t_{0}$, and for $t<t_{0}$ there holds $Z_{t}>0$ so that $Z$ stays bounded, and $\mathcal{T}_{A_{0}}$ still converges to $N_{0}$; in fact the second eventuality holds, because of the slope of the eigenvector at $N_{0}$. So the trajectory, $\mathcal{T}_{A_{0}}$ joins $N_{0}$ to $A_{0}$. By scaling, for any $u_{0}>0$ there exists a unique solution $u$ defined in $(0, \infty)$ satisfying (1.14).
(ii) All the trajectories with one point in the bounded invariant region $\mathcal{R}^{\prime}$ delimitated by the axis $X=0, Z=0$ and $\mathcal{T}_{A_{0}}$, join $N_{0}$ to $O$, and the corresponding solutions $u$ are positive on $(0, \infty)$, decreasing, and satisfy (1.11). The trajectories with one point in the region $\mathcal{R}^{\prime \prime} \subset Q$ above $\mathcal{T}_{A_{0}}$ converge to $N_{0}$ as $t \rightarrow-\infty$, and satisfy $X_{t}>0$, since $\mathcal{T}_{A_{0}}$ is above $\mathcal{L}_{X}$, and cannot be bounded, since there is no fixed point in this region. They can be of two types:

- Either they cross $\mathcal{L}_{Z}$, then after crossing $Z$ is decreasing, necessarily to 0 ; then from (5.6), u is defined in a maximal interval $(0, \rho)$ with $u(\rho)=0$. Such solutions exists because by any point on $\mathcal{L}_{Z}$ passes a trajectory.
- Or they stay above $\mathcal{L}_{Z}$, thus $Z$ increases to $\infty$; in this case from (5.6) $u$ is defined in a maximal interval $(0, \rho)$ with $\lim _{r \rightarrow \rho} u^{\prime}=-\infty$; Let us show the existence of such solutions: For given $c>0$, we define

$$
\mathcal{L}_{c}=\left\{(X, Z) \in Q: X>0, Z=c X+a_{m, q}\right\} .
$$

We compute the field on this line, and show that it is entering the region above $\mathcal{L}_{c}$ for $c$ large enough: indeed we obtain,

$$
\begin{aligned}
\frac{Z_{t}}{X_{t}}-c & =\frac{Z\left(\frac{q+1-m}{m-1}\left(Z-a_{m, q}\right)-p X\right)}{X\left(X-\frac{N-m}{m-1}+\frac{c X+a_{m, q}}{m-1}\right)}-c \\
& =\frac{Z\left(c \frac{q+1-m}{m-1}-p\right)}{X-\frac{N-m}{m-1}+\frac{c X+a_{m, q}}{m-1}}-c>\frac{(c(q+1-m)-p(m-1))\left(c X+a_{m, q}\right)}{(m-1+c) X+a_{m, q}-(N-m)}-c,
\end{aligned}
$$

and

$$
\begin{aligned}
& (c(q+1-m)-p(m-1))\left(c X+a_{m, q}\right)-c\left((m-1+c) X+a_{m, q}-(N-m)\right) \\
& =c X\left(c(q-m)-(p+1)(m-1)-\left(p+a_{m, q}(m-1)\right)+c(q-m) a_{m, q}+N-m\right)-p(m-1) a_{m, q}
\end{aligned}
$$

is positive for large $c$, since $q-m>0$. All the solutions with one point above $\mathcal{L}_{c}$ stay in this region, so above $\mathcal{L}_{Z}$, which proves the existence.
(iii) All the trajectories with one point in $\{(X, Z) \in Q: X<0\}$ satisfy $X_{t}>0$ from (5.5). Then $X$ increases necessarily up to 0 , and then $Z_{t}>0$ for large $t$, thus $(X, Z)$ converges to $O$, and $u$ is defined for $r$ large enough, increasing and $\lim _{r \rightarrow \infty} u=l>0$.

- Either they cross $\mathcal{L}_{Z}$, then before crossing $Z$ is decreasing, necessarily to 0 ; then from (5.6), $u$ is defined in a maximal interval $(0, \rho)$ with $u(\rho)=0$. Such solutions exist as above.
- Or they stay under $\mathcal{L}_{Z}$, thus $X$ and $Z$ decrease to $-\infty$; in this case from (5.6) $u$ is defined in a maximal interval $(0, \rho)$ with $\lim _{r \rightarrow \rho} u^{\prime}=-\infty$. Let us show their existence: for given $k>0$ we
compute the field on the line $\mathcal{L}^{k}=\{(X, Z) \in Q: X<0, Z=k X\}$. On this line $X_{t}>0$ and

$$
\begin{aligned}
Z_{t}-k X_{t} & =Z\left(N-\frac{N-1}{m-1} q-p X+\frac{q+1-m}{m-1} Z-X+\frac{N-m}{m-1}-\frac{Z}{m-1}\right) \\
& =Z\left(\left(\frac{q-m}{m-1}-\frac{p+1}{k}\right) Z-\frac{N-1}{m-1}(q-m)\right)=Z^{2}\left(\frac{q-m}{m-1}-\frac{p+1}{k}\right)+\frac{N-1}{m-1}(q-m)|Z|
\end{aligned}
$$

is positive for large $k$. The region below $\mathcal{L}^{k}$ is therefore negatively invariant, then the existence is folllows. This conclude the proof.

Remark 5.3 The change of variable $u(r)=\tilde{u}(s), s=r^{\frac{m-N}{m-1}}$, introduced in [22], and also used in [13] in case $m=2<N$, leads to the equation

$$
\begin{equation*}
\left(\left|\tilde{u}_{s}\right|_{s}^{m-2} \tilde{u}\right)_{s}+\left(\frac{m-N}{m-1}\right)^{q-m} s^{\frac{N-1}{N-m}(q-m)} \tilde{u}^{p}\left|\tilde{u}_{s}\right|^{q}=0 \tag{5.8}
\end{equation*}
$$

Hence if $u$ is not constant $\tilde{u}_{s}$ does not vanish, from Remark 5.1, and (5.8) is equivalent to

$$
\begin{equation*}
(m-1) \tilde{u}_{s s}+\left(\frac{m-N}{m-1}\right)^{q-m} s^{\frac{N-1}{N-m}(q-m)} \tilde{u}^{p}\left|\tilde{u}_{s}\right|^{q-m+2}=0 . \tag{5.9}
\end{equation*}
$$

In particular we find again the existence and uniqueness of local solutions near $\infty$, satisfying (1.13) for given $l \geq 0$ and $c \neq 0(c>0$ if $l=0)$; indeed the problem reduces to the equation (5.9) with the initial conditions $\tilde{u}(0)=l$ and $\tilde{u}_{s}(0)=c$.

## 6 The case $p<0$

Proof of Theorem 1.5. We still consider $u=v^{b}$, with $b>0$ : we recall that from (4.3) (4.4)

$$
-\Delta_{m} v=(b-1)(m-1) \frac{|\nabla v|^{m}}{v}+b^{q-m+1} v^{s}|\nabla v|^{q}
$$

with $s=1-q+m+b(p+q-m+1)$. Next we take

$$
b=\frac{q+1-m}{p+q-m+1},
$$

thus here $b \geq 1$ and $s=0$, so,

$$
\begin{equation*}
-\Delta_{m} v=(b-1)(m-1) \frac{|\nabla v|^{m}}{v}+b^{q-m+1}|\nabla v|^{q} \tag{6.1}
\end{equation*}
$$

where the two terms have the same sign. Then $z=|\nabla v|^{2}$ satisfies

$$
\mathcal{A}(v)=-\Delta v-\frac{m-2}{2} \frac{<\nabla z, \nabla v>}{z}=(b-1)(m-1) \frac{z}{v}+b^{q-m+1} z^{\frac{q+2-m}{2}} .
$$

Setting, we get from (4.5)

$$
\begin{aligned}
& -\frac{1}{2} \mathcal{A}(z)+\frac{1}{N}(\Delta v)^{2}+(b-1)(m-1) \frac{z^{2}}{v^{2}} \\
\leq & (b-1)(m-1)\left(\frac{<\nabla z, \nabla v>}{v}+\frac{q+2-m}{2} b^{q-m+1} z^{\frac{q-m}{2}}<\nabla z, \nabla v>,\right.
\end{aligned}
$$

where now the term in $\frac{z^{2}}{v^{2}}$ has a positive coefficient. Since $b \geq 1$, we get an estimate of the form

$$
\mathcal{A}(z)+C_{1} z^{q+2-m} \leq C_{3} \frac{|\nabla z|^{2}}{z} .
$$

Since $q+2-m>1$, we deduce the estimate in any ball $B_{\rho}\left(x_{0}\right)$,

$$
\left|\nabla v\left(x_{0}\right)\right| \leq C\left(\frac{1}{\rho}\right)^{\frac{1}{q+1-m}},
$$

from Lemma 3.1, where $C$ is a universal constant, which leads to the conclusions.

## References

[1] Abdelhamid H. and Bidaut-Véron M.F., On the connection between two quasilinear elliptic problems with source terms of order 0 or 1, Comm. Contemporary Math. 12 (2010), 727-788.
[2] Armstrong S.M. and Sirakov B., Nonexistence of positive supersolutions of elliptic equations via the maximum principle, Comm. Part. Diff. Equ. 36 (2011), 2011-2047.
[3] Avila A. and Brock F., Asymptotics at infinity of solutions for p-Laplace equations in exterior domains, Nonlinear Anal. 69 (2008), 1615-1628.
[4] Bidaut-Véron M.F., Local and global behavior of solutions of quasilinear Emden-Fowler type equations, Arch. Rat. Mech. Anal. 107 (1989), 293-324.
[5] Bidaut-Véron M.F. and Pohozaev S., Nonexistence results and estimates for some nonlinear elliptic problems, J. Analyse Mathématiques 84 (2001), 1-49.
[6] Bidaut-Véron M.F. and Véron L., Nonlinear elliptic equations on compact Riemannian manifolds and asymptotics of Emden equations, Inventiones Math. 106 (1991), 489-539.
[7] Bidaut-Véron M.F., Garcia-Huidobro M. and Véron L., Local and global properties of solutions of quasilinear Hamilton-Jacobi equations, J. Funct. Anal. 267 (2014), 3294-3331.
[8] Bidaut-Véron M.F., Garcia-Huidobro M. and Véron L, Estimates of solutions of elliptic equations with a source reaction term involving the product of the function and its gradient, Duke Math. J. 168 (2019), 1487-1537.
[9] Bidaut-Véron M.F., Garcia-Huidobro M. and Yarur C., Keller-Osserman estimates for some quasilinear elliptic systems, Comm. Pure Applied Anal. 12 (2013), 1547-1568.
[10] Bidaut-Véron M.F. and Giacomini H., A new dynamical approach of Emden-Fowler equations and systems, Advances Diff. Equ. 15 (2010), 1033-1082.
[11] Bidaut-Véron M.F. and Grillot P., Singularities in Hamiltonian elliptic systems with absorption terms, Ann. Scuola Norm. Sup. Pisa 28 (1999), 229-271.
[12] Bidaut-Véron M.F. and Grillot P., Asymptotic behaviour of elliptic systems with mixed absorption and source terms, Asymptotic Analysis 19 (1999), 117-147.
[13] Burgos-Perez M.A., Garcia-Melian J. and Quass A., Classification of supersolutions and Liouville theorems for some nonlinear elliptic problems, Discrete Contin. Dyn. Systems 36 (2016), 4703-4721.
[14] Caffarelli L., Gidas G. and Spruck J., Asymptotic symmetry and local behavior os semilinear elliptic equations with critical Sobolev growth, Comm. Pure Applied Math. 42 (1989), 271-297.
[15] Caristi G. and Mitidieri E., Nonexistence of solutions of quasilinear elliptic equations, Adv. Diff. Equ. 2,3 (1997), 319-359.
[16] Chen W. and Li C., Classification of solutions of some nonlinear elliptic equations, Duke Math. J. 63 (1991), 615-622.
[17] Ching J., Cirstea F., Existence and classification of singular solutions to nonlinear elliptic equations with a gradient term Anal. Part. Diff. Equ. 8 (2015), 1931-1962
[18] Ching J. and Cirstea F., Gradient estimates for nonlinear elliptic equations with a gradientdependent nonlinearity, Proc. Royal Soc. Edinburg Sect. Ser. A 150 (2020), 1361-1376.
[19] Filippucci R., Nonexistence of positive weak solutions of elliptic inequalities, Nonlinear Anal. 70 (2009), 2903-2916.
[20] Filippucci R., Pucci P. and Souplet P., A Liouville-type theorem for an elliptic equation with superquadratic growth in the gradient, Adv. Nonlinear Studies 20 (2020), 245-251.
[21] Gidas B. and Spruck J. Global and local behavior of positive solutions of nonlinear elliptic equations, Comm. Pure Appl. Math. 34 (1981), 525-598.
[22] Guedda M. and Véron L., Local and Global Properties of Solutions of Quasilinear Elliptic Equations, J. Diff. Equ. 76 (1988), 159-189.
[23] Leonori T. and Porretta A., Large solutions and gradient bounds for quasilinear elliptic equations, Comm. Part. Diff. Equ. 41 (2016), 952-998.
[24] Lions P.L., Quelques remarques sur les problèmes elliptiques quasilinéaires du second ordre, J. Anal. math. 45 (1985), 234-254.
[25] Marcus M. and Nguyen PhuocT., Elliptic equations with nonlinear absorption depending on the solution and its gradient, Proc. Lond. Math. Soc. 111 (2015), 205-239.
[26] Mitidieri E. and Pohozaev S., A priori estimates and blow-up of solutions to nonlinear par-tial differential equations and inequalities, Trudy Mat. Instit. Steklov 234 (2001), 1-383.
[27] Kichenassamy S. and Véron L. Singular solutions of the p-Laplace equation, Math. Annal. 275 (1986), 599-615.
[28] Reshetniak, Y. G., Mappings with bounded deformation as extremals of Dirichlet type integrals, Sibirsk Math. Zh. 9 (1966), 652-666.
[29] Sciunzi B., Classification of the $D^{1, p}\left(R^{N}\right)$-solutions of the critical p-Laplace equation in $R^{N}$, Adv. in Math. 291 (2016), 12-23.
[30] Sciunzi B., Merchan S., Montoro L., Radial symmetry and applications for a problem involving the $\Delta_{p}($.$) operator and critical nonlinearity in R^{N}$, Adv. in Math. 265 (2014), 313-335.
[31] Serrin J. Local behaviour of solutions of quasi-linear equations, Acta Math. 111 (1964), 247302.
[32] Serrin J. Isolated singularities of solutions of quasi-linear equations, Acta Math. 113 (1965), 219-240
[33] Serrin J. and Zou H., Cauchy-Liouville and universal boundedness theorems for quasilinear elliptic equations and inequalities, Acta Math. 189 (2002), 79-142.
[34] Trudinger N. S., On harnack type inequality and their applications to quasilinear elliptic equations, Comm. Pure Appl. Math. 20 (1967), 721-747.
[35] Vàzquez J. L., A Strong Maximum Principle for Some Quasilinear Elliptic Equations, Appl. Math. Opt. 12 (1984), 191-202.
[36] Vétois J., A priori estimates and application to the symmetry of solutions for critical p-Laplace equations, J. Diff. Equ. 260 (2016), 149-161.


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