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RIEMANN'S NON-DIFFERENTIABLE FUNCTION AND THE BINORMAL CURVATURE FLOW

VALERIA BANICA AND LUIS VEGA

ABSTRACT. We make a connection between a famous analytical object introduced in the 1860s by Riemann, as well as some variants of it, and a nonlinear geometric PDE, the binormal curvature flow. As a consequence this analytical object has a non-obvious non-linear geometric interpretation. We recall that the binormal flow is a standard model for the evolution of vortex filaments. We prove the existence of solutions of the binormal flow with smooth trajectories that are as close as desired to curves with a multifractal behavior. Finally, we show that this behavior falls within the multifractal formalism of Frisch and Parisi, which is conjectured to govern turbulent fluids.

1. INTRODUCTION

In this article we construct the graph of Riemann's non-differentiable function, and variants of it, by using the binormal curvature flow, a geometric flow of curves in three dimensions that is related to the evolution of vortex filaments. We also make a rigorous connection between the binormal flow and the multifractal formalism of Frisch and Parisi.

1.1. Riemann's function and the multifractal formalism. A classical problem of mathematical analysis is finding real variable functions that are continuous but not differentiable at any point. Although it seems that the first example is due to Bolzano, traditionally the reference names in this matter are Riemann and Weierstrass, the latter attributing to Riemann the example

$$(1) \quad \varphi_R(t) = \sum_{j=1}^{\infty} \frac{\sin(tj^2)}{j^2}.$$

In fact Weierstrass, faced with the impossibility of proving that the previous function is not differentiable at any point, proposes his own examples which are later known as Weierstrass' functions. After fundamental contributions from Hardy in 1915 [30], the problem was not solved until 1960 by Gerver, who proved in [26] and [27] that the function φ_R is not differentiable except at times $t_{p,q} = \pi p/q$, with p and q odd numbers, in which case the derivative is precisely $-1/2$. Later Duistermaat [20] studied the self-similarity properties of the complex function, intimately associated with the Riemann's function φ_R , defined as:

$$(2) \quad \varphi_D(t) = \sum_{j=1}^{\infty} \frac{e^{itj^2}}{ij^2}.$$

He drew attention to the apparent fractal properties of the graph generated by it. Finally, Jaffard proved using the wavelet transform in [32] that in fact Riemann's function φ_R , and analogously the complex version studied by Duistermaat φ_D is a multifractal function that moreover satisfies what is known as the multifractal formalism of Frisch and Parisi. The motivation of the latter notion has its roots in the theory developed by Frisch and Parisi to explain certain data obtained in [1] by Antonia, Hopfinger, Gagne and Anselmet on the velocity structure functions in turbulent shear flows, that exit the homogeneous and isotropic framework of Kolmogorov 41's theory of turbulence.

More concretely, Jaffard's result in [32] is about the spectrum of singularities, that is the function $d_f(\beta)$ which associates β with the Hausdorff dimension of the sets of points t_0 where f has pointwise Hölder regularity of exponent β . This Hölder exponent is defined as the supremum of $\{\alpha : f \in \mathcal{C}^\alpha(t_0)\}$. Here $\mathcal{C}^\alpha(t_0)$ stands for the functions f for which there exists a polynomial P of order at most α such that locally at t_0

$$|f(t) - P(t - t_0)| < C|t - t_0|^\alpha.$$

For instance Weierstrass' functions $W_{a,b}(x) = \sum_{n \in \mathbb{N}^*} a^n \cos(b^n x)$ with $a < 1 < ab$ are nowhere differentiable, but have constant Hölder exponent $\alpha = -\log a / \log b$. Thus they belong to the class of monofractal functions characterized by the fact that their spectrum support are reduced to one point, encoding, despite of the fractal appearance of its graph, a sort of disciplined irregularity. However, the points where the Hölder exponent is reached might be a fractal set, and actually this is the reason of the monofractal label. The devil's staircase is a famous example of monofractal function, as it has only one finite Hölder exponent, reached on the Cantor's triadic set. In turn, multifractal functions are those whose Hölder exponent takes at least two finite values. The most complex such functions are those with spectrum positive at least on a whole interval. This encodes the fact that the regularity varies roughly between close points. For more details on these notions one can consult [33]. In [32] it is proved that for $\beta \in [1/2, 3/4]$,

$$(3) \quad d_{\varphi_R}(\beta) = 4\beta - 2.$$

It was also shown in [32] that (3) fits with what Frisch and Parisi conjecture in [25]:

$$(4) \quad d_f(\beta) = \inf_p (\beta p - \eta_f(p) + 1),$$

where $\eta_f(p)$ is defined in terms of Besov regularity:

$$(5) \quad \eta_f(p) = \sup\{s, f \in B_p^{\frac{s}{p}, \infty}\}.$$

We refer the reader to §8.5.3 of [24] and p.443 of [32] for the details on this multifractal formalism. Also, it was proved recently in [9] that Riemann's function is intermittent. The results in [32] and [9] are analytical in nature, and no direct connection is established between Riemann's function and turbulence. The aim of this article is to make a connection between Riemann's function and the time evolution of vortex filaments.

1.2. Vortex filaments: the binormal flow model and particular solutions. The vortex filaments are present in 3-D fluids having vorticity concentrated along a curve, and are a key element of quantum and classical fluid turbulent dynamics. This low regularity framework is difficult to analyze through Euler and Navier-Stokes equation. It is however at the heart of current investigations (see for instance [34],[12]). In this article we consider the binormal flow equation (BF), a classical reduced model for vortex filament dynamics. This model was formally derived by truncating the integral given by Biot-Savart's law ([16],[38],[2],[13]). Recently a rigorous argument, but still under some strong assumptions, has been given by Jerrard and Seis in [34]. If the vorticity concentrates along a curve $\chi(t, x)$, where t stands for the temporal variable and x is the arclength parameter, the BF evolution is

$$(6) \quad \partial_t \chi = \partial_x \chi \wedge \partial_x^2 \chi.$$

Using the Frenet system is immediate to see that (6) is also written as

$$(7) \quad \partial_t \chi = \kappa b,$$

where κ stands for the curvature of the curve and b for the binormal vector. By differentiation with respect to x , the tangent vector T of a BF solution solves the Schrödinger map with values in the unit sphere \mathbb{S}^2 , that is the classical continuous Heisenberg model used in ferromagnetics

$$(8) \quad \partial_t T = T \wedge \partial_x^2 T.$$

Finally, and thanks to the Hasimoto transformation:

$$\psi(t, x) = \kappa(t, x) e^{i \int_0^x \tau(t, s) ds},$$

with τ denoting the torsion of the curve, one gets that the function $\psi(t)$, called filament function of $\chi(t)$, satisfies the 1-D focusing cubic Schrödinger equation

$$(9) \quad i\psi_t + \psi_{xx} + \frac{1}{2} (|\psi|^2 - A(t)) \psi = 0,$$

for some real function $A(t)$ ([31]). Conversely, from a solution of the 1-D cubic Schrödinger equation one can construct a solution of the binormal flow solving either the Frenet equations, or better through the construction of a parallel frame $(T, e_1, e_2)(t, x)$. This type of frame fits better with our needs because it is not necessary to suppose that $\kappa > 0$ (for details on this construction see for instance §2 of [3]). It is important to note that with this construction, the binormal flow solution obtained from $\psi(t, x)$ solution of (9) is the same as the one obtained from $\psi(t, x) e^{-i \frac{\Phi(t)}{2}}$ solution of (9) with $A(t)$ replaced by $A(t) - \Phi(t)$. Thus one can always reduce to the usual cubic nonlinear Schrödinger equation, i.e., (9) with $A(t) = 0$.

Simple examples that can be obtained by this construction are:

- the straight line; $\psi(t, x) = 0$ and $A(t) = 0$;
- the circle; $\psi(t, x) = c > 0$ and $A(t) = c^2$;
- the helix; $\psi(t, x) = c e^{ix\omega_0 - it\omega_0^2}$ with $c > 0$ and $A(t) = c^2$;
- the self-similar solutions; $\psi(t, x) = \frac{c}{\sqrt{t}} e^{i \frac{x^2}{4t}}$ with $c > 0$ and $A(t) = \frac{c^2}{t}$.

After integrating the frame system one gets solutions of the non-linear equations (8) and (6). For doing that one needs to know the trajectory in time of one point. This is rather easy for the first three examples but is more delicate for the last one. As a matter of fact, it is better to solve directly (8) and (6) to get the four examples mentioned above, instead of using (9).

For instance for the selfsimilar solutions it is enough to look for solutions of the type $\chi(t, x) = \sqrt{t}G(x/\sqrt{t})$. Then it is easy to get that G has to solve the non-linear ode

$$(10) \quad \frac{1}{2}G - \frac{x}{2}G' = G' \wedge G''.$$

From this is rather simple to conclude that G is determined by the fact that the curvature has to be a constant c and the torsion has to be $\tau(s) = s/2$, see [11]. This means that $\chi(t)$ has curvature $\kappa(t, x) = \frac{c}{\sqrt{t}}$ and torsion $\tau(t, x) = \frac{x}{2t}$, and that $\chi(t)$ tends to two different lines at $x \pm \infty$ that are the same at all times. As a consequence $\chi(t)$ is a smooth function for $t > 0$ that becomes at $t = 0$ a polygonal line with one corner located at $x = 0$. These selfsimilar solutions were characterized in [29]. We will make a very strong use of this characterization in this paper. In particular the angle θ of the corner is related to c by the formula

$$(11) \quad \sin \frac{\theta}{2} = e^{-\frac{\pi}{2}c^2}.$$

Recall that if χ_0 is a polygonal line with just one corner of angle θ located at $x = 0$, then its curvature is given by $\kappa(0, x) = (\pi - \theta)\delta(x)$. Nevertheless, for constructing the solution of the binormal flow for that χ_0 , one has to solve (9) with initial data $\psi(0, x) = c\delta(x)$ and c as in (11). We will also need to know what is the relation between $G(0)$ and the two asymptotic lines of G at infinity and the plane that contains them, see [29]. Observe that from (10) we get that the trajectory of the corner is

$$\chi(t, 0) = \sqrt{t}G(0) = 2\frac{c}{\sqrt{t}}b(0),$$

with $(T(0), n(0), b(0))$ the Frenet frame at $x = 0$ of the profile curve G , which can be taken any orthonormal matrix due to the rotation invariance of (10). In this paper we will follow [29] and take $(T(0), n(0), b(0))$ the canonical orthonormal basis of \mathbb{R}^3 .

Similarly, the straight line, characterized by $\kappa = 0$, is a trivial solution of (6), and the circle and the helix can be easily obtained by looking at traveling solutions of (8). This immediately gives the dynamics of these solutions, and the particular fact that they conserve their shapes. Indeed, the circle moves with a constant speed along the axis perpendicular to the plane where it is contained, with direction depending on the initial orientation given by the arclength parametrization. The helix evolves by screwing up or down, also depending on the initial orientation. At this point it is important to recall that vortex filaments with the shapes of straight lines, circles, and helices do exist, both in experiments and as solutions of Euler equations. Also the selfsimilar solutions are very reminiscent of the flow behind a delta wing jet and in the reconnection process of helium superfluid. We refer the reader to [4] for the corresponding references.

It is worth mentioning that the helix can be obtained from the circle using one of the symmetries of the set of solutions of (9). These are the Galilean transformations: if $\psi(t, x)$ solves (9) with a constant $A(t)$, so does

$$\psi_{\omega_0}(t, x) = e^{ix\tau_0 - it\omega_0^2} \psi(t, x - 2\omega_0 t)$$

for any $\omega_0 \in \mathbb{R}$.

1.3. Numerical evidence about the connection between Riemann's function and the line, circle, and helix filaments. In [18] the Galilean transformations are used to look for solutions of (6) that are initially a planar regular polygon of M sides. The reason is simply because, in view of the construction of self-similar solutions, it is natural to look for solutions of the cubic Schrödinger equation with initial data

$$(12) \quad \psi_M(0, x) = c_M \sum_{j \in \mathbb{Z}} \delta(x - \frac{2\pi}{M}j) = c_M \frac{M}{2\pi} \sum_{j \in \mathbb{Z}} e^{ixMj},$$

with $c_M > 0$ related to the angle $\theta_M = \frac{M-2}{M}\pi$ by the relation (11), and δ denoting Dirac's delta function. The last equality uses Poisson's summation formula

$$(13) \quad \sum_{j \in \mathbb{Z}} f(j) = \sum_{j \in \mathbb{Z}} \hat{f}(2\pi j) = \sum_{j \in \mathbb{Z}} \int e^{-i2\pi j y} f(y) dy.$$

Then, it follows immediately from (12) that the $\psi_M(0, x)$ is invariant under the discrete subgroup of the Galilean transformations given by $\omega_0 \in \mathbb{Z}$. As a consequence, if uniqueness holds, it is proved in [18] that then

$$(14) \quad \psi_M(t, x) = \tilde{c}_M(t) \sum_{j \in \mathbb{Z}} e^{ixMj - it(Mj)^2},$$

with $\tilde{c}_M(t)$ a real function which is determined by geometric means. Later on it was showed in [5] the existence of a formal conservation law whose validity implies that $\tilde{c}_M(t)$ should indeed be a constant so

$$\tilde{c}_M^2(t) = -\frac{M^2}{4\pi^2} \ln(\cos \frac{\pi}{M}).$$

Notice that for all M we have

$$\frac{1}{\tilde{c}_M} \psi_M(\frac{t}{M^2}, \frac{x}{M}) = \sum_{j=-\infty}^{\infty} e^{ixj - itj^2},$$

and that $\lim_{M \rightarrow \infty} \tilde{c}_M = \frac{1}{4\pi}$. Hence,

$$(15) \quad \lim_{M \rightarrow \infty} 4\pi \psi_M(\frac{t}{M^2}, \frac{x}{M}) = \sum_{j=-\infty}^{\infty} e^{ixj - itj^2},$$

which is the solution of the linear Schrödinger equation with periodic boundary conditions,

$$(16) \quad \begin{aligned} i\psi_t + \psi_{xx} &= 0 \\ \psi(0, x) &= \sum_{j \in \mathbb{Z}} \delta(x - 2\pi j). \end{aligned}$$

Moreover, this solution describes the Talbot effect in Optics, see [7]. In [18] the consequences of this effect in (6) with initial data given by regular polygons were considered, suggesting a possible connection with the turbulent dynamics observed in non-circular jets (see on this subject for instance [28]). Let us mention also that at a less singular level, the Talbot effect for the linear and nonlinear Schrödinger equations on the torus with initial data given by functions with bounded variation has been largely studied ([39],[42],[45],[22],[15]). Let us mention also that the fractal behavior of one corner has been observed numerically in the context of the architecture of aortic valve fibers in [41] and [43].

Immediately we obtain that fixing x at the origin and integrating in time the limit in (15) we obtain¹,

$$(17) \quad \int_0^t \sum_{j=-\infty}^{\infty} e^{-i\tau j^2} d\tau = \sum_{j \in \mathbb{Z}} \frac{e^{itj^2} - 1}{ij^2} =: \mathfrak{R}(t).$$

In fact

$$\mathfrak{R}(t) = 2\varphi_D(t) + t + \sum_{j \in \mathbb{Z}^*} \frac{1}{j^2},$$

where φ_D is the function defined in (2) and studied by Duistermaat. It follows from the same arguments given by Gerber that at the points where $\varphi_D(t)$ is differentiable the derivative is precisely $-1/2$. Therefore at those points $\mathfrak{R}'(t) = 0$. It turns out that when one looks at the trajectory in the complex plane of $\mathfrak{R}(t)$, see Figure 1, there is no tangent at

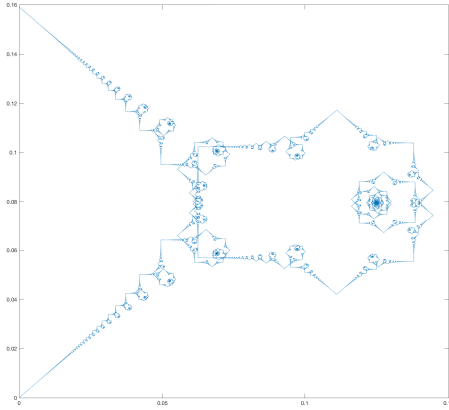


FIGURE 1. Graph of $\mathfrak{R}(t) = \int_0^t \sum_j e^{i\tau j^2} d\tau = \sum_{j \in \mathbb{Z}} \frac{e^{itj^2} - 1}{ij^2}$

those points due to the fact that the curve spirals around them. In fact, it has been recently

¹The term $j = 0$ is understood to be t .

proved by Eceizabarrena in [21] that the trajectory, although continuous, and contrary to what happens with Riemann's function, does not have a tangent at any point.

Therefore, it was very natural to ask what is the trajectory in time of any of the corners of the M -regular polygon. In Figure 2 we see the examples obtained in [18] for $M = 3$

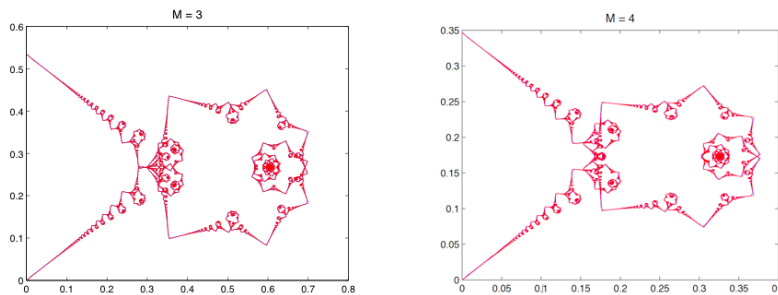


FIGURE 2. Trajectory that at time $t = 0$ starts in a vertex of an equilateral triangle (left) and of a square (right)

and $M = 4$. Similar pictures can be obtained for any M and it becomes evident after looking at them, that they converge to the one of $\mathfrak{R}(t)$ given in Figure 1. After doing an appropriate renormalization, this convergence is proved numerically in [18]. Later on, it is also proved numerically in [19] the convergence of the Fourier coefficients of the time derivative of the trajectory. These results have been extended in [17] to the case of helical regular polygons that converge to either a helix or to straight line. In the case of the helices, and depending on its pitch, different versions of \mathfrak{R} are obtained². The results in this paper prove analytically the aforementioned convergence using an approximation by non-closed polygonal lines.

1.4. Presentation of the results. Our main statement asserts the existence of various families of solutions $\{\chi_n\}_{n \in \mathbb{N}}$ of the binormal flow such that the trajectory of the corner $\chi_n(t, 0)$ near $t = 0$ is governed by the modified version of Riemann's function \mathfrak{R} as n goes to infinity.

Theorem 1.1. *Let $n \in \mathbb{N}^*$, $\nu \in]0, 1]$, $\Gamma > 0$. There exist $T > 0$, independent of n , and smooth solutions $\chi_n(t)$ of the binormal flow on $(-T, T) \setminus \{0\}$, weak solutions on $(-T, T)$, that at time $t = 0$ become polygonal lines $\chi_n(0)$ with corners located at $j \in \mathbb{Z}$ with $|j| \leq n^\nu$, of same torsion $\omega_0 \in \pi\mathbb{Q}$ and angles θ_n such that*

$$(18) \quad \lim_{n \rightarrow \infty} n(\pi - \theta_n) = \Gamma,$$

²The sequence of the squares that appears in \mathfrak{R} has to be changed into the squares of any arithmetic progression with integer coefficients, see [17].

and

$$(19) \quad \chi_n(0, 0) = (0, 0, 0), \quad \partial_x \chi_n(0, 0^\pm) = \left(\sin \frac{\theta_n}{2}, \pm \cos \frac{\theta_n}{2}, 0 \right).$$

For these solutions we have the following description of the trajectory of the corner $\chi_n(t, 0)$:

$$(20) \quad n(\chi_n(t, 0) - \chi_n(0, 0)) - (0, \mathfrak{R}(\tilde{\mathfrak{R}}(t)), \Im(\tilde{\mathfrak{R}}(t))) \xrightarrow{n \rightarrow \infty} 0,$$

uniformly on $(0, T)$. The function $\tilde{\mathfrak{R}}$ is multifractal, its spectrum of singularities satisfies (3) and the multifractal formalism formula (4). In the torsion-free case $\tilde{\mathfrak{R}}(t) = -\Gamma \frac{\mathfrak{R}(4\pi^2 t)}{4\pi^2}$, with \mathfrak{R} given in (17). The expressions for the cases with a non-trivial torsion are given in (71) and (72).

In the torsion-free case the polygonal lines in Theorem 1.1 can be chosen to approach the following special cases:

- the straight line; by taking $\nu < 1$. Indeed, the total variation angle of $\chi_n(0, x)$ as x varies from $-\infty$ to ∞ is

$$\theta_n^{total} := (\pi - \theta_n)(2 \lfloor n^\nu \rfloor + 1) \xrightarrow{n \rightarrow \infty} \frac{\Gamma}{n}(2 \lfloor n^\nu \rfloor + 1),$$

so if $\nu < 1$ we get convergence of θ_n^{total} to zero as n goes to infinity.

- a regular polygonal loop; by taking $\nu = 1$ and $\theta_n = \frac{(2n-1)\pi}{2n+1}$. Indeed, this means that the shape of $\chi_n(0)$ is composed of a regular closed polygon with $2n + 1$ edges of size 1, for $|x| \leq n$, and of two half-lines as $|x| \geq n$.
- a regular polygonal multi-loop; by increasing the number of corners of the regular polygonal loop to $x \in \{j \in \mathbb{Z}, |j| \leq mn^\nu\}$, for $m \in \mathbb{N}^*$. The proof of the conclusion of Theorem 1.1 goes the same, for times T_m of size $\frac{1}{m}$.

To approach other natural special cases we recall that the binormal flow is invariant under scaling: if χ is a solution then $\lambda\chi(\frac{t}{\lambda^2}, \frac{x}{\lambda})$ is a solution also for $\lambda > 0$. Thus from Theorem 1.1 we get for $\mu \in \mathbb{R}$ solutions of the binormal flow

$$\tilde{\chi}_n(t, x) = \frac{1}{n^\mu} \chi_n(n^{2\mu}t, n^\mu x).$$

For times smaller than $\frac{T}{n^{2\mu}}$ the convergence (20) becomes

$$(21) \quad n^{1+\mu}(\tilde{\chi}_n(\frac{t}{n^{2\mu}}, 0) - \tilde{\chi}_n(0, 0)) - (0, \mathfrak{R}(\tilde{\mathfrak{R}}(t)), \Im(\tilde{\mathfrak{R}}(t))) \xrightarrow{n \rightarrow \infty} 0,$$

uniformly on $(0, T)$. This convergence is for instance valid for polygonal lines that tend to two lines at infinity and that locally approach the following curves::

- a circular loop; by rescaling the regular polygonal loop above with $\mu = 1$. Indeed, $\tilde{\chi}_n(0)$ is composed by a regular closed polygon with $2n + 1$ edges of size $\frac{1}{n}$, thus inscribed in a circle of radius $\frac{1}{n \sin \frac{\pi}{2n+1}}$ for $|x| \leq n$ and two half-lines for $|x| \geq n$. In particular the polygon, as n goes to infinity, converges to a circle of size $\frac{2}{\pi}$.

- a circular multi-loop; by rescaling the regular polygonal multi-loop above with $\mu = 1$. This example confirms the numerical simulations of [17] that can be seen in the video <https://www.youtube.com/watch?v=bwbpKvqGk-o&feature=youtu.be>.
- the self-similar solution; by proceeding in the following way. Denote by θ the angle of a self-similar solution, and choose $\theta_n = \pi - \frac{\theta}{2n+1}$ and $\nu = 1$ so that $\theta_n^{total} = \theta$. Then, the shape of $\chi_n(0)$ is composed by a polygonal line with $2n + 1$ corners with the same angle and edges of size 1 that is inscribed in a circular sector of radius of size $2n + 1$ for $|x| \leq n$, and of two half-lines for $|x| \geq n$. By rescaling with $\mu > 1$, we get that $\tilde{\chi}_n(0)$ is composed of a polygonal line with $2n + 1$ corners with the same angle θ_n , and edges of size $\frac{1}{n^\mu}$ inside a circular sector of radius of size $\frac{1}{n^{\mu-1}}$ for $|x| \leq 1$, and of two half-lines for $|x| \geq 1$. Moreover the angle between the half-lines is precisely θ .

We note that in the above configurations the loops imply the existence of self-intersections, something that can not happen in a vortex filament. Nevertheless, they are relevant from a theoretical point of view as an analytical approximation to a real dynamics. Observe that the number of loops, although fixed, can be arbitrary large.

In the non-trivial torsion case the above examples give families of helicoidal polygonal lines. This way we have a non-planar approximation of the straight line, as well as, after rescaling, an approximation of a helical shape with as many turns as desired.

In order to explain the proof of Theorem 1.1 we have to recall some previous work done in [4] regarding the evolution through the binormal flow of non-closed polygonal lines $\chi_0(x)$ that tend as $x \rightarrow \pm\infty$ to two lines. These polygonal lines are characterized modulo a translation and a rotation by the fact that the corners are located at the integers $k \in \mathbb{Z}$ and by the curvature angles θ_k and torsion angles at the corners. For constructing the evolution of χ_0 according to the binormal flow, we define first a sequence of complex numbers $\{\alpha_k\}$ in terms of the curvature and torsion angles of χ_0 . In particular, the identity (11) has to be satisfied. The identity involving the torsion angles is more complicated, and is detailed in §6. We impose that for $s > 1/2$ and $p = 2$, the sequence $\{\alpha_k\}$ has to belong to $l^{p,s}$, the space of sequences of complex numbers which is determined by the condition

$$\|\alpha_k\|_{l^{p,s}}^p := \sum_{k \in \mathbb{Z}} |\alpha_k|^p (1 + |k|)^{2s} < +\infty.$$

Then, we solve on $t > 0$ the equation (9) with $A(t) = \frac{\sum_{k \in \mathbb{Z}} |\alpha_k|^2}{2\pi t}$ and with

$$(22) \quad \psi(t, x) = \sum_{k \in \mathbb{Z}} e^{-i \frac{|\alpha_k|^2}{4\pi} \log \sqrt{t}} (\alpha_k + R_k(t)) e^{it\Delta} \delta(x - k),$$

such that

$$(23) \quad \sup_{0 < t < T} \frac{1}{t^\gamma} \|\{R_k(t)\}\|_{l^{2,s}} + t \|\{\partial_t R_k(t)\}\|_{l^{2,s}} < C.$$

Here T and C depend only on $\|\{\alpha_k\}\|_{l^{2,s}}$ and $\gamma \in (0, 1)$. It is a remarkable fact that we have the mass conservation

$$(24) \quad \sum_{k \in \mathbb{Z}} |\alpha_k + R_k(t)|^2 = \sum_{k \in \mathbb{Z}} |\alpha_k|^2, \quad \forall 0 \leq t \leq T.$$

Then from this solution ψ we obtain in [4] a smooth solution $\chi(t)$ of the binormal flow (6) on $(-T, T) \setminus \{0\}$, which is a weak solution on $(-T, T)$, that at time $t = 0$ becomes the desired polygonal line χ_0 .

Thus to construct the solutions χ_n in Theorem 1.1 we consider for $n \in \mathbb{N}^*$ sequences $\{\alpha_{n,k}\}_{k \in \mathbb{Z}}$ satisfying

$$(25) \quad \alpha_{n,k} = \begin{cases} c_n e^{ik\omega_0}, & |k| \leq n^\nu, \\ 0, & |k| > n^\nu, \end{cases}$$

where $c_n > 0$ is defined by the identity (11) in terms of the angle θ_n from the statement of Theorem 1.1. Note that we have

$$(26) \quad c_n = \sqrt{-\frac{2}{\pi} \arccos \frac{\pi - \theta_n}{2}} \stackrel{n \rightarrow \infty}{\approx} \frac{\Gamma}{2\sqrt{\pi n}}.$$

Then the above results proved in [4] insure us the existence of BF solutions that at time $t = 0$ are polygonal lines with $2n^\delta + 1$ corners located at $x \in \{-\lfloor n^\nu \rfloor, \dots, -1, 0, 1, \dots, \lfloor n^\nu \rfloor\}$, all of them with the same curvature angle θ_n and torsion ω_0 . Having in mind that the binormal flow is invariant under translations and rotations we can consider $\chi_n(0)$ such that (19) holds.

Concerning the first part of Theorem 1.1 we have to observe that if we directly use the results in [4] we obtain solutions (22) of the Schrödinger equation (9), constructed by a Picard iteration procedure that is valid for times T_n that depend on the weighted norm of $\|\{\alpha_{n,k}\}\|_{l^{2,s}}$, $s > 1/2$. In particular, we obtain that T_n vanishes as n goes to infinity for $\delta = 1$. In this article we shall improve the iteration procedure on $\{R_{n,k}\}_{k \in \mathbb{Z}}$ by using l^1 -based spaces. This allows us to consider times T that depend only on $\|\alpha_{n,k}\|_{l_k^1}$, so that in our case T can be chosen independent of n . Another important remark is that the construction of the solution of the binormal flow (6) done in [4], and that is based on solutions of (9), depends just on the l^1 norm of $\{\alpha_{n,k}\}_{k \in \mathbb{Z}}$ and of $\{R_{n,k}\}_{k \in \mathbb{Z}}$. Therefore, the construction of the solutions of the binormal flow from the Schrödinger ones is assured by the results in [4].

Regarding the second part of the statement of the theorem, we will prove it by splitting the trajectory of the corner into three parts. One part will disappear in the limit due to an improved decay in n that we obtained for $R_{n,k}$ in l^1 and $l^{1,1}$. Another part will be shown to be negligible by a fine analysis done in the construction of the parallel frame. It will be based on repeated integration by parts on some oscillatory integrals that naturally appear. They include some problematic resonant terms that eventually disappear thanks to the specific value of $A(t)$. Finally, the last part, and in the torsion-free case, includes Riemann's function that appears thanks to Poisson summation formula. For helicoidal

polygonal lines, the proof goes the same except that we end up with some variants of \mathfrak{R} . In this case we prove, following the approach of Chamizo and Ubis in [14], that their spectrum of singularities is the same as the one of Riemann's function in (3). Also, we show that the multifractal formalism of Frish and Parisi (4) is satisfied, by using exponential sums estimates.

The paper is organized as follows. The proof of the first part of Theorem 1.1 is done in §2. The second part of Theorem 1.1 in the torsion-free case is proved in §3 using §2, and two results related to the convergence of the normal vectors proved in §4 and §5 respectively. In the last section §6 we shall treat the nontrivial torsion cases of Theorem 1.1. This involves proving the results by Jaffard in [32] in the more general setting of the squares of arithmetic progressions.

2. PROOF OF THE FIRST PART OF THEOREM 1.1

In view of Proposition 2.1 below, for each sequence $\{\alpha_{n,k}\}_{k \in \mathbb{Z}}$ defined in (25) we get a solution of the Schrödinger equation (9) with $A(t) = \frac{\sum_{k \in \mathbb{Z}} |\alpha_{n,k}|^2}{2\pi t}$, of type (22), with time of existence T independent of n . As mentioned above in the Introduction, the construction of the corresponding BF solutions involves only the l^1 norms of $\{\alpha_{n,k}\}_{k \in \mathbb{Z}}$ and $\{R_{n,k}\}_{k \in \mathbb{Z}}$, so the first part of Theorem 1.1 follows.

In the following Proposition we improve the fixed point argument on $\{R_{n,k}(t)\}_{k \in \mathbb{Z}}$ in a way it suits our purposes here.

Proposition 2.1. *Let $n \in \mathbb{N}$, $\gamma \in (0, 1)$, $q > 1$, $C > 0$ and $\{\alpha_{n,k}\}_{k \in \mathbb{Z}}$ a sequence such that $|\alpha_{n,k}| \leq \frac{C}{n}$ for $|k| \leq n$ and $\alpha_{n,k} = 0$ otherwise. There exist $T \in (0, 1)$ depending only on γ and q and a unique solution written as*

$$\sum_{k \in \mathbb{Z}} e^{-i \frac{|\alpha_{n,k}|^2}{4\pi} \log \sqrt{\tau}} (\alpha_{n,k} + R_{n,k}(t)) e^{it\Delta} \delta_k(x),$$

of the equation

$$i\psi_t + \psi_{xx} + \frac{1}{2} \left(|\psi|^2 - \frac{\sum_k |\alpha_{n,k}|^2}{2\pi t} \right) \psi = 0,$$

with the property:

$$(27) \quad \sup_{t \in (0, T)} \|t^{-\gamma} R_{n,k}(t)\|_{l_k^1} \leq C(\gamma, q) \frac{1}{n^{2-\frac{2}{q}}},$$

$$\sup_{\tau \in (0, T)} \|t \partial_t R_{n,k}(t)\|_{l_k^1} \leq C(\gamma, q), \quad \sup_{\tau \in (0, T)} \|t \partial_t R_{n,k}(t)\|_{l_k^q} \leq C(\gamma, q) \frac{1}{n^{1-\frac{1}{q}}},$$

and

$$(28) \quad \sup_{t \in (0, T)} \|t^{-\gamma} R_{n,k}(t)\|_{l_k^{1,1}} \leq C(\gamma, q) \frac{1}{n^{1-\frac{2}{q}}}.$$

Proof. We follow the argument in [4], so that we have to find a fixed point for the application $\Phi(\{R_j\}) = \{\Phi_k(\{R_j\})\}$ given by

$$(29) \quad \Phi_k(\{R_j\})(t) = -i \int_0^t f_k(\tau) d\tau + i \int_0^t \frac{1}{8\pi\tau} (|\alpha_k + R_k(\tau)|^2 - |\alpha_k|^2) (\alpha_k + R_k(\tau)) d\tau,$$

where

$$(30) \quad f_k(t) = \frac{1}{8\pi t} \sum_{(j_1, j_2, j_3) \in NR_k} e^{-i \frac{\Delta_{k, j_1, j_2, j_3}}{4\tau}} e^{-i\omega_{k, j_1, j_2, j_3} \log \sqrt{t}} (\alpha_{j_1} + R_{j_1}(t)) \overline{(\alpha_{j_2} + R_{j_2}(t))} (\alpha_{j_3} + R_{j_3}(t)),$$

and

$$\omega_{k, j_1, j_2, j_3} = \frac{|\alpha_k|^2 - |\alpha_{j_1}|^2 + |\alpha_{j_2}|^2 - |\alpha_{j_3}|^2}{4\pi}, \quad \Delta_{k, j_1, j_2, j_3} = k^2 - j_1^2 + j_2^2 - j_3^2,$$

$$NR_k = \{(j_1, j_2, j_3), j_1 - j_2 + j_3 = k, \Delta_{k, j_1, j_2, j_3} \neq 0\},$$

see for instance (24) in [4]; for simplicity we have omitted the n -subindex. We shall perform the fixed point argument in the ball

$$X^{\gamma, q, n} := \{\{M_k\} \in \mathcal{C}^1((0, T), l^1) \cap \mathcal{C}((0, T), l^{1,1}), \|\{M_k\}\|_{X^{\gamma, q, n}} < \delta\},$$

where

$$\begin{aligned} \|\{M_k\}\|_{X^{\gamma, q, n}} &:= n^{2-\frac{2}{q}} \sup_{t \in (0, T)} \|t^{-\gamma} R_k(t)\|_{l^1} + \sup_{1 \leq \tilde{q} \leq q} n^{1-\frac{1}{\tilde{q}}} \sup_{t \in (0, T)} \|t \partial_t R_k(t)\|_{l^{\tilde{q}}} \\ &\quad + n^{1-\frac{2}{q}} \sup_{t \in (0, T)} \|t^{-\gamma} R_{n, k}(t)\|_{l^{1,1}}, \end{aligned}$$

and $T \in (0, 1)$ will be specified later.

Let $\{R_j\} \in X^{\gamma, q, n}$. We start with the estimates of $\|\Phi_k(\{R_j\})(t)\|_{l^1}$. To estimate the first term in the expression (29) we shall perform as in (36) in [4] an integration by parts in time to get advantage of the non-resonant phase $\Delta_{k, j_1, j_2, j_3}$ and to obtain integrability in time:

$$(31) \quad \begin{aligned} i \int_0^t f_k(\tau) d\tau &= t \sum_{(j_1, j_2, j_3) \in NR_k} \frac{e^{-i \frac{\Delta_{k, j_1, j_2, j_3}}{4\tau}} e^{-i\omega_{k, j_1, j_2, j_3} \log \sqrt{t}}}{2\pi \Delta_{k, j_1, j_2, j_3}} (\alpha_{j_1} + R_{j_1}(t)) \overline{(\alpha_{j_2} + R_{j_2}(t))} (\alpha_{j_3} + R_{j_3}(t)) \\ &\quad - \int_0^t \sum_{(j_1, j_2, j_3) \in NR_k} \frac{e^{-i \frac{\Delta_{k, j_1, j_2, j_3}}{4\tau}}}{2\pi \Delta_{k, j_1, j_2, j_3}} \partial_\tau (e^{-i\omega_{k, j_1, j_2, j_3} \log \sqrt{t}} \tau (\alpha_{j_1} + R_{j_1}(\tau)) \overline{(\alpha_{j_2} + R_{j_2}(\tau))} (\alpha_{j_3} + R_{j_3}(\tau))) d\tau. \end{aligned}$$

We shall exploit the decay given by $\Delta_{k,j_1,j_2,j_3} = 2(j_1 - j_2)(j_3 - j_2)$ on NR_k yielding for $1 \leq q < \infty$ the estimate

$$(32) \quad \left\| \sum_{(j_1,j_2,j_3) \in NR_k} \frac{M_{j_1} N_{j_2} P_{j_3}}{\Delta_{k,j_1,j_2,j_3}} \right\|_{l_k^1} \leq \sum_{j_2 \notin \{j_1,j_3\}} \left| \frac{M_{j_1} N_{j_2} P_{j_3}}{(j_1 - j_2)(j_3 - j_2)} \right| \leq C_q \|M_j\|_{l^q} \|N_j\|_{l^q} \|P_j\|_{l^1},$$

obtained by performing Hölder estimates in the j_1, j_2 variables. Similarly, we get also as upper-bounds $C\|M_j\|_{l^q} \|N_j\|_{l^1} \|P_j\|_{l^q}$ and $C\|M_j\|_{l^1} \|N_j\|_{l^q} \|P_j\|_{l^q}$. Therefore

$$(33) \quad \begin{aligned} \left\| \int_0^t f_k(\tau) d\tau \right\|_{l^1} &\leq Ct(1 + \|\alpha_j\|_{l^\infty}^2) (\|\alpha_j\|_{l^q}^2 \|\alpha_j\|_{l^1} + \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^1} (\|\alpha_j\|_{l^1}^2 + \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^1}^2)) \\ &\quad + Ct(\|\alpha_j\|_{l^q}^2 + \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^q}^2) \sup_{\tau \in (0,T)} \|\tau \partial_\tau R_j(\tau)\|_{l^1}. \end{aligned}$$

The second term in (29) contains only cubic terms with at least a power of R_k so we conclude that for all $1 \leq q < \infty$, as $l^q \subset l^1$,

$$(34) \quad \begin{aligned} \|\Phi_k(\{R_j\})(t)\|_{l^1} &\leq Ct(1 + \|\alpha_j\|_{l^\infty}^2) (\|\alpha_j\|_{l^q}^2 \|\alpha_j\|_{l^1} + \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^1} (\|\alpha_j\|_{l^1}^2 + \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^1}^2)) \\ &\quad + Ct(\|\alpha_j\|_{l^q}^2 + \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^1}^2) \sup_{\tau \in (0,T)} \|\tau \partial_\tau R_j(\tau)\|_{l^1} \\ &\quad + Ct^\gamma \sup_{\tau \in (0,T)} \|\tau^{-\gamma} R_j(\tau)\|_{l^1} (\|\alpha_j\|_{l^\infty}^2 + \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^\infty}^2). \end{aligned}$$

In particular, as $\|\alpha_j\|_{l^\infty} \leq \frac{C}{n}$, $\|\alpha_j\|_{l^1} \leq \frac{C}{n^1}$, $\|\alpha_j\|_{l^q} \leq \frac{C}{n^{1-\frac{1}{q}}}$, we have

$$(35) \quad \begin{aligned} \sup_{\tau \in (0,T)} \|\tau^{-\gamma} \Phi_k(\{R_j\})(\tau)\|_{l^1} &\leq CT^{1-\gamma} \left(\frac{1}{n^{2(1-\frac{1}{q})}} + \frac{T^\gamma \delta}{n^{2(1-\frac{1}{q})}} \left(1 + \frac{T^{2\gamma} \delta^2}{n^{4(1-\frac{1}{q})}} \right) \right) \\ &\quad + CT^{1-\gamma} \left(\frac{1}{n^{2(1-\frac{1}{q})}} + \frac{T^{2\gamma} \delta^2}{n^{2(1-\frac{1}{q})}} \right) \delta + C \frac{\delta}{n^{2(1-\frac{1}{q})}} \left(\frac{1}{n^2} + \frac{T^{2\gamma} \delta^2}{n^{4(1-\frac{1}{q})}} \right). \end{aligned}$$

So for T and δ less than a constant depending only on γ and q we have

$$n^{2(1-\frac{1}{q})} \sup_{\tau \in (0,T)} \|\tau^{-\gamma} \Phi_k(\{R_j\})(\tau)\|_{l^1} \leq \frac{\delta}{3}.$$

Now we shall get estimates on $\partial_\tau \Phi_k(\{R_j\})(\tau)$. As we have for all $1 \leq \tilde{q}$

$$\left\| \sum_{(j_1,j_2,j_3) \in NR_k} M_{j_1} N_{j_2} P_{j_3} \right\|_{l_k^{\tilde{q}}} \leq \|\{M_j\} \star \{N_j\} \star \{P_j\}(k)\|_{l_k^{\tilde{q}}} \leq C \|M_j\|_{l^1} \|N_j\|_{l^1} \|P_j\|_{l^{\tilde{q}}},$$

we get from (29)

$$\sup_{\tau \in (0,T)} \|\tau \partial_\tau \Phi_k(\{R_j\})(\tau)\|_{l^{\tilde{q}}} \leq C \|\alpha_j + R_j(\tau)\|_{l^1}^2 \|\alpha_j + R_j(\tau)\|_{l^{\tilde{q}}}$$

$$+C\|R_j(t)\|_{l^{\tilde{q}}}(\|\alpha_j\|_{l^\infty}^2 + \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^\infty}^2).$$

Hence for $1 \leq \tilde{q}$, as $l^{\tilde{q}} \subset l^1$,

(36)

$$\sup_{\tau \in (0,T)} \|\tau \partial_\tau \Phi_k(\{R_j\})(\tau)\|_{l^{\tilde{q}}} \leq C(1 + \frac{T^{2\gamma}\delta^2}{n^{4(1-\frac{1}{q})}})(\frac{1}{n^{1-\frac{1}{\tilde{q}}}} + \frac{T\gamma\delta}{n^{2(1-\frac{1}{q})}}) + C\frac{T\gamma\delta}{n^{2(1-\frac{1}{q})}}(\frac{1}{n^2} + \frac{T^{2\gamma}\delta^2}{n^{4(1-\frac{1}{q})}}).$$

Therefore, again for T and δ less than a constant depending only on γ and q we have

$$\sup_{1 \leq \tilde{q} \leq q} n^{1-\frac{1}{\tilde{q}}} \sup_{\tau \in (0,T)} \|\tau \partial_\tau \Phi_k(\{R_j\})(\tau)\|_{l^{\tilde{q}}} \leq \frac{\delta}{3}.$$

Finally, the control of the weighted norm $\|\{\Phi_k(\{R_j\})(\tau)\|_{l^{1,1}}$ is obtained similarly, by using weighted estimates of type

$$\begin{aligned} & \left\| \sum_{(j_1, j_2, j_3) \in NR_k} \frac{M_{j_1} N_{j_2} P_{j_3}}{\Delta_{k, j_1, j_2, j_3}} \right\|_{l^{1,1}} = \sum_k \left| \sum_{(j_1, j_2, j_3) \in NR_k} \frac{M_{j_1} N_{j_2} P_{j_3}}{\Delta_{k, j_1, j_2, j_3}} k \right| \\ & \leq \sum_{j_2 \notin \{j_1, j_3\}} \left| \frac{M_{j_1} N_{j_2} P_{j_3} (j_1 - j_2 + j_3)}{(j_1 - j_2)(j_3 - j_2)} \right| \\ & \leq C\|M_j\|_{l^q}\|N_j\|_{l^q}\|P_j\|_{l^{1,1}} + C\|M_j\|_{l^1}\|N_j\|_{l^q}\|P_j\|_{l^1}, \end{aligned}$$

for all $1 \leq \tilde{q} < \infty$. We get in the same way:

$$\begin{aligned} (37) \quad & \sup_{\tau \in (0,T)} \|\tau^{-\gamma} \Phi_k(\{R_j\})(\tau)\|_{l^{1,1}} \leq CT^{1-\gamma}(1 + \|\alpha_j\|_{l^\infty}^2)(\|\alpha_j\|_{l^q}^2 \|\alpha_j\|_{l^{1,1}} + \|\alpha_j\|_{l^q} \|\alpha_j\|_{l^1}^2) \\ & + \sup_{\tau \in (0,T)} (\|R_j(\tau)\|_{l^{1,1}} + \|R_j(\tau)\|_{l^1})(\|\alpha_j\|_{l^1}^2 + \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^1}^2) \\ & + CT^{1-\gamma}(\|\alpha_j\|_{l^q} \|\alpha_j\|_{l^{1,1}} \sup_{\tau \in (0,T)} \|\tau \partial_\tau R_j(\tau)\|_{l^q} + \|\alpha_j\|_{l^q} \|\alpha_j\|_{l^1} \sup_{\tau \in (0,T)} \|\tau \partial_\tau R_j(\tau)\|_{l^1}) \\ & + ((\|\alpha_j\|_{l^{1,1}} + \|\alpha_j\|_{l^1}) \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^1} + \|\alpha_j\|_{l^1} \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^{1,1}}) \sup_{\tau \in (0,T)} \|\tau \partial_\tau R_j(\tau)\|_{l^1} \\ & + (\sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^1}^2 + \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^1} \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^{1,1}}) \sup_{\tau \in (0,T)} \|\tau \partial_\tau R_j(\tau)\|_{l^1} \\ & + C \sup_{\tau \in (0,T)} \|\tau^{-\gamma} R_j(\tau)\|_{l^{1,1}} (\|\alpha_j\|_{l^\infty}^2 + \sup_{\tau \in (0,T)} \|R_j(\tau)\|_{l^\infty}^2). \end{aligned}$$

Thus again for T and δ less than a constant depending only on γ and q we have

$$n^{1-\frac{2}{q}} \sup_{\tau \in (0,T)} \|\tau^{-\gamma} \Phi_k(\{R_j\})(\tau)\|_{l^{1,1}} \leq \frac{\delta}{3}.$$

Summarizing, we have obtained the existence of T and δ less than a constant depending only on γ and q such that the stability estimate holds : if $\{R_k\} \in X^{\gamma, q, n}$ then $\{\Phi_k(\{R_j\})\} \in X^{\gamma, q, n}$. Thus to end the fixed point argument we need only the contraction estimates, that can be obtained in the same way. \square

3. PROOF OF THE SECOND PART OF THEOREM 1.1 IN THE PLANAR CASE

Let us first recall from §4.3 in [4] that for constructing the solutions of BF using the parallel frame (T, e_1, e_2) the following equations have to be solved:

$$(38) \quad T_x = \Re u e_1 + \Im u e_2 = \Re(\bar{u} N),$$

$$(39) \quad N_x = e_{1x} + i e_{2x} = -\Re u T - i \Im u T = -u T,$$

$$(40) \quad T_t = -\Im u_x e_1 + \Re u_x e_2 = \Im(\bar{u}_x N),$$

$$(41) \quad N_t = -i u_x T + i \left(\frac{|u|^2}{2} - \frac{\sum_{k \in \mathbb{Z}} |\alpha_k|^2}{2t} \right) N,$$

$$(42) \quad \chi_t = T \wedge T_x = T \wedge \Re(\bar{u} N) = \Im(\bar{u} N).$$

Above we have taken

$$(43) \quad u(t, x) = \sum_j e^{-i(|\alpha_j|^2 - \sum_{k \in \mathbb{Z}} |\alpha_k|^2) \log \sqrt{t}} (\alpha_j + R_j(t)) \frac{e^{i \frac{(x-j)^2}{4t}}}{\sqrt{t}},$$

$N = e_1 + i e_2$, and we have omitted the n -subindices for simplicity. We note that the ansatz (43) comes from the one given in Proposition 2.1 applied to the sequence $\{\sqrt{4\pi i} \alpha_n\}$ instead of $\{\alpha_n\}$, and thus the notation R_j in (43) comes from the remainder term in Proposition 2.1 divided by $\sqrt{4\pi i}$. Thus the remainder term in (43) enjoys the same decay properties (27)-(28).

In view of (42) we can write the evolution of the corner located at $x = 0$ as

$$\begin{aligned} \chi_n(t, 0) - \chi_n(0, 0) &= \int_0^t \Im(\bar{u}_n N_n(\tau, 0)) d\tau \\ &= \Im \int_0^t \sum_j e^{i(|\alpha_{n,j}|^2 - \sum_{k \in \mathbb{Z}} |\alpha_{n,k}|^2) \log \sqrt{\tau}} (\alpha_{n,j} + R_{n,j}(\tau)) \frac{e^{-i \frac{j^2}{4\tau}}}{\sqrt{\tau}} N_n(\tau, 0) d\tau. \end{aligned}$$

By using Proposition 2.1 the term involving $R_{n,j}(\tau)$ yields decay in n . Therefore

$$\chi_n(t, 0) - \chi_n(0, 0) = \Im \int_0^t \sum_{|j| \leq n^\nu} \frac{\alpha_{n,j}}{\sqrt{\tau}} e^{-i \frac{j^2}{4\tau}} e^{-i \sum_{k \neq j} |\alpha_{n,k}|^2 \log \sqrt{\tau}} N_n(\tau, 0) d\tau + r_n(t),$$

with

$$(44) \quad |r_n(t)| \leq \frac{C}{n^{2^-}}, \quad \forall t \in (0, T).$$

Here 2^- means any number smaller than 2, on which the constant depend.

From Lemma 4.5 in [4] we get the existence of the limit

$$\lim_{t \rightarrow 0} e^{i \sum_{j \neq x} |\alpha_{n,j}|^2 \log \frac{|x-j|}{\sqrt{t}}} N_n(t, x) =: \tilde{N}_n(0, x) \in \mathbb{S}^2 + i\mathbb{S}^2.$$

Hence we can write:

$$(45) \quad \begin{aligned} \chi_n(t, 0) - \chi_n(0, 0) &= \Im(e^{-ic_n^2 \sum_{1 \leq |j| \leq n^\nu} \log |j|} \tilde{N}_n(0, 0)) \int_0^t \sum_{|j| \leq n^\nu} \frac{\overline{\alpha_{n,j}} e^{-i \frac{j^2}{4\tau}}}{\sqrt{\tau}} d\tau \\ &+ \Im(e^{-ic_n^2 \sum_{1 \leq |j| \leq n^\nu} \log |j|} \int_0^t \sum_{|j| \leq n^\nu} \frac{\overline{\alpha_{n,j}} e^{-i \frac{j^2}{4\tau}}}{\sqrt{\tau}} g_n(\tau) d\tau) + r_n(t), \end{aligned}$$

where

$$(46) \quad g_n(t) = e^{i\Phi_n(t)} N_n(t, 0) - \tilde{N}_n(0, 0), \quad \Phi_n(t) = \sum_{j \in \mathbb{Z}} |\alpha_{n,j}|^2 \log \frac{|j|}{\sqrt{t}} = c_n^2 \sum_{1 \leq |j| \leq n^\nu} \log \frac{|j|}{\sqrt{t}}.$$

As we are in the case $\alpha_{n,j} = c_n$ for $|j| \leq n^\nu$, the first term makes appear the Riemann's function as follows.

Lemma 3.1.

$$\left| \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i \frac{j^2}{4\tau}}}{\sqrt{\tau}} d\tau - \frac{e^{-i \frac{\pi}{4}}}{2\pi\sqrt{\pi}} \Re(4\pi^2 t) \right| \leq \frac{C}{n^\nu},$$

uniformly on $(0, T)$.

Proof. We first replace the summation in j over the whole set of integers. Indeed, by integration by parts we have

$$\left| \int_0^t \sum_{|j| > n^\nu} \frac{e^{-i \frac{j^2}{4\tau}}}{\sqrt{\tau}} d\tau \right| = \left| \left[- \sum_{|j| > n^\nu} \frac{i4\tau\sqrt{\tau} e^{-i \frac{j^2}{4\tau}}}{j^2} \right]_0^t + \int_0^t \sum_{|j| > n^\nu} \frac{i6\sqrt{\tau} e^{-i \frac{j^2}{4\tau}}}{j^2} d\tau \right| \leq \frac{C}{n^\nu}.$$

Now we shall use Poisson's summation formula $\sum_{j \in \mathbb{Z}} f(j) = \sum_{j \in \mathbb{Z}} \hat{f}(2\pi j)$:

$$\begin{aligned} \sum_{j \in \mathbb{Z}} e^{i4\pi^2 t j^2} &= \sum_{j \in \mathbb{Z}} \int e^{-i2\pi x j + i4\pi^2 t x^2} dx = \frac{1}{\sqrt{4\pi^2 t}} \sum_{j \in \mathbb{Z}} \int e^{-iy \frac{j}{\sqrt{t}} + iy^2} dy \\ &= \frac{1}{2\pi\sqrt{t}} \sum_{j \in \mathbb{Z}} \widehat{e^{i \cdot^2}} \left(\frac{j}{\sqrt{t}} \right) = \frac{e^{i \frac{\pi}{4}}}{2\sqrt{\pi}\sqrt{t}} \sum_{j \in \mathbb{Z}} e^{-i \frac{j^2}{4t}}, \end{aligned}$$

and the statement follows after integration in time. \square

In view of this result, of (45) and (44) we obtain, as $0 < \nu \leq 1$,

$$n(\chi_n(t, 0) - \chi_n(0, 0)) - \Im(nc_n e^{-ic_n^2 \sum_{1 \leq |j| \leq n^\nu} \log |j|} \tilde{N}_n(0, 0)) \frac{e^{-i \frac{\pi}{4}}}{2\pi\sqrt{\pi}} \Re(4\pi^2 t)$$

$$-\Im(nc_n e^{-ic_n^2 \sum_{1 \leq |j| \leq n^\nu} \log |j|} \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} g_n(\tau) d\tau) \xrightarrow{n \rightarrow \infty} 0,$$

uniformly on $(0, T)$. We use Proposition 4.1 to get

$$n(\chi_n(t, 0) - \chi_n(0, 0)) - \Im(nc_n e^{-ic_n^2 \sum_{1 \leq |j| \leq n^\nu} \log |j|} \tilde{N}_n(0, 0) \frac{e^{-i\frac{\pi}{4}}}{2\pi\sqrt{\pi}} \Re(4\pi^2 t)) \xrightarrow{n \rightarrow \infty} 0.$$

By using the convergence (26) of c_n and the convergence of $\tilde{N}_n(0, 0)$ obtained in Proposition 5.1:

$$\lim_{n \rightarrow \infty} \tilde{N}_n(0, 0) = \left(0, \frac{1-i}{\sqrt{2}}, \frac{-1-i}{\sqrt{2}}\right),$$

we have

$$n(\chi_n(t, 0) - \chi_n(0, 0)) - \Im\left(\Gamma(0, -i, -1) \frac{\Re(4\pi^2 t)}{4\pi^2}\right) \xrightarrow{n \rightarrow \infty} 0,$$

and Theorem 1.1 follows.

4. A CONVERGENCE ESTIMATE FOR THE NORMAL VECTORS

Proposition 4.1. *Let g_n be as defined in (46). Then*

$$\int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} g_n(\tau) d\tau \xrightarrow{n \rightarrow \infty} 0,$$

uniformly on $(0, T)$.

Proof. In view of (41) and (43) we have, by omitting the subindices n for simplicity, except for c_n ,

$$\begin{aligned} \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} g(\tau) d\tau &= \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \left(-iu_x T + i \left(\frac{|u|^2}{2} - \frac{\sum_{j \in \mathbb{Z}} |\alpha_j|^2}{2s} \right) N + i\Phi_s N \right) e^{i\Phi}(s, 0) ds d\tau \\ &= - \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau e^{i[n^\nu]c_n^2 \log s} \sum_{k \neq 0} (\alpha_k + R_k(s)) \frac{e^{i\frac{k^2}{4s}}}{2s\sqrt{s}} k T(s, 0) e^{i\Phi(s)} ds d\tau \\ &\quad + i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{k^2 \neq l^2} (\alpha_k + R_k(s)) (\overline{\alpha_l + R_l(s)}) \frac{e^{i\frac{k^2-l^2}{4s}}}{2s} N(s, 0) e^{i\Phi(s)} ds d\tau \\ &\quad + i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \left(\sum_k (\alpha_k + R_k(s)) (\overline{\alpha_{-k} + R_{-k}(s)}) \frac{1}{2s} + \Phi_s \right) N(s, 0) e^{i\Phi(s)} ds d\tau \end{aligned}$$

$$(47) \quad =: I + J + K.$$

We start with the second term J .

Lemma 4.2.

$$J = i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{k^2 \neq l^2} (\alpha_k + R_k(s)) \overline{(\alpha_l + R_l(s))} \frac{e^{i\frac{k^2-l^2}{4s}}}{2s} N(s, 0) e^{i\Phi(s)} ds d\tau \xrightarrow{n \rightarrow \infty} 0.$$

Proof. Let us first observe that when $R_j(s)$ appears, it insures integration in s . Moreover there is also enough decay in n in view of (27). On the remaining term

$$\tilde{J} := i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{k^2 \neq l^2, |k|, |l| \leq n^\nu} \alpha_k \overline{\alpha_l} \frac{e^{i\frac{k^2-l^2}{4s}}}{2s} N(s, 0) e^{i\Phi(s)} ds d\tau,$$

we shall perform an integration by parts

$$\begin{aligned} \tilde{J} &= - \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \sum_{k^2 \neq l^2, |k|, |l| \leq n^\nu} \alpha_k \overline{\alpha_l} \frac{e^{i\frac{k^2-l^2}{4\tau}}}{k^2 - l^2} 2\tau N(\tau, 0) e^{i\Phi(\tau)} d\tau \\ &+ 2 \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{k^2 \neq l^2, |k|, |l| \leq n^\nu} \alpha_k \overline{\alpha_l} \frac{e^{i\frac{k^2-l^2}{4s}}}{k^2 - l^2} (2sN(s, 0) e^{i\Phi(s)})_s ds d\tau. \end{aligned}$$

We also notice that

$$(48) \quad \left| \sum_{k^2 \neq l^2, |k|, |l| \leq n^\nu} \frac{1}{k^2 - l^2} \right| \leq C(\log n^\nu)^2,$$

and that in view of (41) and of the definition (46) of Φ ,

$$(49) \quad |(2sN(s, 0) e^{i\Phi(s)})_s| \leq C \frac{n}{\sqrt{s}}.$$

Therefore we have got integrability in s and τ . Then, we have convergence to zero as n goes to infinity for the boundary term and for the integral term with $j = 0$. On the remaining integral terms with $j \neq 0$ we shall perform an integration by parts in τ to get summability in j without loss in n :

$$\begin{aligned} & \int_0^t \sum_{1 \leq |j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{k^2 \neq l^2, |k|, |l| \leq n^\nu} \alpha_k \overline{\alpha_l} \frac{e^{i\frac{k^2-l^2}{4s}}}{k^2 - l^2} (2sN(s, 0) e^{i\Phi(s)})_s ds d\tau \\ &= \sum_{1 \leq |j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4t}}}{ij^2} 4t\sqrt{t} \int_0^t \sum_{k^2 \neq l^2, |k|, |l| \leq n^\nu} \alpha_k \overline{\alpha_l} \frac{e^{i\frac{k^2-l^2}{4s}}}{k^2 - l^2} (2sN(s, 0) e^{i\Phi(s)})_s ds \\ &- \int_0^t \sum_{1 \leq |j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{ij^2} \left(4\tau\sqrt{\tau} \int_0^\tau \sum_{k^2 \neq l^2, |k|, |l| \leq n^\nu} \alpha_k \overline{\alpha_l} \frac{e^{i\frac{k^2-l^2}{4s}}}{k^2 - l^2} (2sN(s, 0) e^{i\Phi(s)})_s ds \right)_\tau d\tau. \end{aligned}$$

Again by using (48) and (49) we get the convergence to zero as n goes to infinity. \square

We consider now the first term and last terms in the decomposition (47):

Lemma 4.3.

$$I + K = - \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau e^{i[n^\nu]c_n^2 \log s} \sum_{k \neq 0} (\alpha_k + R_k(s)) \frac{e^{i\frac{k^2}{4s}}}{2s\sqrt{s}} k T(s, 0) e^{i\Phi(s)} ds d\tau$$

$$+ i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \left(\sum_k (\alpha_k + R_k(s)) (\overline{\alpha_{-k} + R_{-k}(s)}) \frac{1}{2s} + \Phi_s \right) N(s, 0) e^{i\Phi(s)} ds d\tau \xrightarrow{n \rightarrow \infty} 0.$$

Proof. The terms involving $\{R_k(s)\}$ in the first integral for $j = 0$ and in the second integral converge all to zero as n goes to 0 by using (27) with $\gamma > \frac{1}{2}$. On the remaining terms, that involve $\{R_k(s)\}$ and $j \neq 0$ in the first integral, we integrate by parts in τ to get summability in j . Eventually using again (27) with $\gamma > \frac{1}{2}$ we get convergence to zero as n goes to 0.

Thus we have to show

$$(50) \quad \tilde{I} + \tilde{K} = - \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau e^{i[n^\nu]c_n^2 \log s} \sum_{1 \leq |k| \leq n^\nu} \alpha_k \frac{e^{i\frac{k^2}{4s}}}{2s\sqrt{s}} k T(s, 0) e^{i\Phi(s)} ds d\tau$$

$$+ i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \left(\sum_{|k| \leq n^\nu} \alpha_k \overline{\alpha_{-k}} \frac{1}{2s} - \sum_{|k| \leq n^\nu} |\alpha_k|^2 \frac{1}{2s} \right) N(s, 0) e^{i\Phi(s)} ds d\tau \xrightarrow{n \rightarrow \infty} 0.$$

In the case $\alpha_{n,k} = c_n$ for $|k| \leq n^\nu$ we have $\alpha_k = \alpha_{-k}$ and the term \tilde{K} vanishes. Otherwise \tilde{K} will cancel with a piece of \tilde{I} , as we shall see later. The term \tilde{I} involves a bad power in s for integration, so we need to integrate by parts in the s variable:

$$\tilde{I} = -2i \int_0^t \sum_{|j| \leq n^\nu, 1 \leq |k| \leq n^\nu} \alpha_k \frac{e^{i\frac{k^2-j^2}{4\tau}}}{k} e^{i[n^\nu]c_n^2 \log \tau} T(\tau, 0) e^{i\Phi(\tau)} d\tau$$

$$+ 2i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{1 \leq |k| \leq n^\nu} \alpha_k \frac{e^{i\frac{k^2}{4s}}}{k} (e^{i[n^\nu]c_n^2 \log s} \sqrt{s} T(s, 0) e^{i\Phi(s)})_s ds d\tau.$$

We first treat the boundary term. When $j^2 = k^2$ we get a $\frac{\log n}{n}$ decay. When $j^2 \neq k^2$ we perform an integration by parts and get decay in n for all terms except the one coming from T_τ by using the estimate

$$(51) \quad \alpha_n \sum_{|j| \leq n^\nu, 1 \leq |k| \leq n^\nu, j^2 \neq k^2} \left| \frac{1}{k(k^2 - j^2)} \right| \leq \frac{C}{n^{1-\epsilon}}.$$

We are left with the term involving T_τ :

$$I_b := -2i \int_0^t \sum_{|j| \leq n^\nu, 1 \leq |k| \leq n^\nu, j^2 \neq k^2} \alpha_k \frac{e^{i\frac{k^2-j^2}{4\tau}}}{k(k^2 - j^2)} e^{i[n^\nu]c_n^2 \log \tau} \Im(\overline{u_x} N(\tau, 0)) e^{i\Phi(\tau)} \tau^2 d\tau$$

$$\begin{aligned}
&= -i \int_0^t \sum_{|j| \leq n^\nu, 1 \leq |k| \leq n^\nu, j^2 \neq k^2} \alpha_k \frac{e^{i \frac{k^2 - j^2}{4\tau}}}{k(k^2 - j^2)} e^{i \lfloor n^\nu \rfloor c_n^2 \log \tau} \\
&\quad \times \Im \left(i \sum_l e^{-i \lfloor n^\nu \rfloor c_n^2 \log \tau} \overline{(\alpha_l + R_l(\tau))} l e^{-i \frac{l^2}{4\tau}} N(\tau, 0) \right) e^{i\Phi(\tau)} \sqrt{\tau} d\tau.
\end{aligned}$$

In view of (51) and (28) we obtain the convergence to zero as n goes to infinity for the term involving R_l . So it remains to estimate the integral

$$\begin{aligned}
\tilde{I}_b &:= -i \int_0^t \sum_{|j| \leq n, 1 \leq |k| \leq n, j^2 \neq k^2} \alpha_k \frac{e^{i \frac{k^2 - j^2}{4\tau}}}{k(k^2 - j^2)} e^{i \lfloor n^\nu \rfloor c_n^2 \log \tau} \\
&\quad \times \Im \left(i \sum_{1 \leq |l| \leq n^\nu} e^{-i \lfloor n^\nu \rfloor c_n^2 \log \tau} \overline{\alpha_l} l e^{-i \frac{l^2}{4\tau}} N(\tau, 0) \right) e^{i\Phi(\tau)} \sqrt{\tau} d\tau.
\end{aligned}$$

For $l^2 = \pm(k^2 - j^2)$ the summation in l disappears and we get again by (51) a $\frac{\log^2 n}{n}$ decay. For $l^2 \neq \pm(k^2 - j^2)$ we perform an integration by parts and for the worse term involving the $u_x T$ in N_τ we use

$$(52) \quad c_n^2 \sum_{|j| \leq n^\nu, 1 \leq |k|, |l| \leq n^\nu, j^2 \neq k^2, l^2 \neq \pm(k^2 - j^2), p} \frac{|(\alpha_p + R_p(\tau)) p|}{|k(k^2 - j^2)(k^2 - j^2 \pm l^2)|} \leq \frac{C}{n^{1-}},$$

having in mind (28) for the part concerning the R_p terms.

Finally we treat the last remaining term, namely the integral term in \tilde{I} :

$$\begin{aligned}
&2i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i \frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{1 \leq |k| \leq n^\nu} \alpha_k \frac{e^{i \frac{k^2}{4s}}}{k} (e^{i \lfloor n^\nu \rfloor c_n^2 \log s} \sqrt{s} T(s, 0) e^{i\Phi(s)})_s ds d\tau \\
&= 2i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i \frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{1 \leq |k| \leq n^\nu} \alpha_k \frac{e^{i \frac{k^2}{4s}}}{k} (e^{i \lfloor n^\nu \rfloor c_n^2 \log s} \sqrt{s} e^{i\Phi(s)})_s T(s, 0) e^{i\Phi(s)} ds d\tau \\
&\quad + 2i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i \frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{1 \leq |k| \leq n^\nu} \alpha_k \frac{e^{i \frac{k^2}{4s}}}{k} e^{i \lfloor n^\nu \rfloor c_n^2 \log s} \sqrt{s} T_s(s, 0) e^{i\Phi(s)} ds d\tau.
\end{aligned}$$

For the first integral we observe that

$$|(e^{i \lfloor n^\nu \rfloor c_n^2 \log s} \sqrt{s} e^{i\Phi(s)})_s| \leq C \left(\frac{1}{n\sqrt{s}} + \frac{1}{\sqrt{s}} + \frac{1}{ns} \right).$$

Then, if $j = 0$, we get $\frac{\log n}{n}$ decay. For the terms $j \neq 0$ we integrate by parts in τ to get summability in j , so that we get a $\frac{\log n}{n}$ bound. Therefore we are left with the second

integral

$$\begin{aligned}
& 2i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{1 \leq |k| \leq n^\nu} \alpha_k \frac{e^{i\frac{k^2}{4s}}}{k} e^{i2[n^\nu]c_n^2 \log \sqrt{s}} \sqrt{s} \Im(\overline{u_x} N(s, 0)) e^{i\Phi(s)} ds d\tau \\
&= i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{1 \leq |k| \leq n} \alpha_k \frac{e^{i\frac{k^2}{4s}}}{k} e^{i2[n^\nu]c_n^2 \log \sqrt{s}} \\
&\quad \times \Im\left(i \sum_{l \neq 0} e^{-i2n[n^\nu]c_n^2 \log \sqrt{s}} (\alpha_l + R_l(s)) \frac{l}{s} e^{-i\frac{l^2}{4s}} N(s, 0)\right) e^{i\Phi(s)} ds d\tau.
\end{aligned}$$

For the R_l terms we use (28) to get integrability in s and a $\frac{1}{n^{1-}}$ decay. We are left with

$$\begin{aligned}
I_i &:= i \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{1 \leq |k| \leq n^\nu} \alpha_k \frac{e^{i\frac{k^2}{4s}}}{k} e^{i[n^\nu]c_n^2 \log s} \\
&\quad \times \Im\left(i \sum_{1 \leq |l| \leq n^\nu} e^{-i[n^\nu]c_n^2 \log s} \frac{l}{\alpha_l s} e^{-i\frac{l^2}{4s}} N(s, 0)\right) e^{i\Phi(s)} ds d\tau \\
&= \frac{i}{2} \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{1 \leq |k|, |l| \leq n^\nu} \alpha_k \overline{\alpha_l} \frac{l}{k} e^{i\frac{k^2-l^2}{4s}} \frac{1}{s} N(s, 0) e^{i\Phi(s)} ds d\tau \\
&\quad + \frac{i}{2} \int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{1 \leq |k|, |l| \leq n^\nu} e^{i[n^\nu]c_n^2 \log s} \alpha_k \alpha_l \frac{l}{k} e^{i\frac{k^2+l^2}{4s}} \frac{1}{s} \overline{N(s, 0)} e^{i\Phi(s)} ds d\tau.
\end{aligned}$$

In the case $\alpha_{n,k} = c_n$ for $|k| \leq n^\nu$, in the first integral the terms with $k = l$ and the terms with $k = -l$ cancel. Otherwise they cancel with \tilde{K} in (50). For all the remaining terms we need to perform an integration by parts to settle the integration in s . The terms not involving the $u_x T$ term of N_s are converging to zero as n goes to infinity since

$$(53) \quad c_n^2 \sum_{1 \leq |k|, |l| \leq n^\nu, k^2 \neq l^2} \frac{|l|}{|k(k^2 \pm l^2)|} \leq \frac{C}{n^{2-}}.$$

It goes the same by using (28) for the R_p -part coming from the $u_x T$ term of N_s . Therefore the last terms to treat are

$$\int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{1 \leq |k|, |l|, |p| \leq n^\nu} \alpha_k \overline{\alpha_l} \alpha_p \frac{lp}{k(k^2 - l^2)} e^{i\frac{k^2-l^2+p^2}{4s}} \frac{e^{i[n^\nu]c_n^2 \log s}}{\sqrt{s}} T(s, 0) e^{i\Phi(s)} ds d\tau,$$

and

$$\int_0^t \sum_{|j| \leq n^\nu} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} \int_0^\tau \sum_{1 \leq |k|, |l|, |p| \leq n^\nu} \alpha_k \alpha_l \overline{\alpha_p} \frac{lp}{k(k^2 + l^2)} e^{i\frac{k^2+l^2-p^2}{4s}} \frac{e^{i[n^\nu]c_n^2 \log s}}{\sqrt{s}} T(s, 0) e^{i\Phi(s)} ds d\tau.$$

For $j = 0$ we get again by (53) a $\frac{1}{n^{1-}}$ decay. For $j \neq 0$ we perform a last integration by parts in τ to get summability in j without loss in n and use again (53) to get a $\frac{1}{n^{1-}}$ decay. \square

In view of (47) and the last two lemmas the Proposition 4.1 is proved. \square

5. CONVERGENCE OF THE MODULATED NORMAL VECTORS AT $(t, x) = (0, 0)$

Proposition 5.1. *The following convergence holds:*

$$\lim_{n \rightarrow \infty} \tilde{N}_n(0, 0) = \left(0, \frac{1-i}{\sqrt{2}}, \frac{-1-i}{\sqrt{2}}\right).$$

Proof. Recall that from Lemma 4.5 in [4] we have

$$\lim_{t \rightarrow 0} e^{i \sum_{j \neq x} |\alpha_{n,j}|^2 \log \frac{|x-j|}{\sqrt{t}}} N_n(t, x) =: \tilde{N}_n(0, x),$$

with

$$|e^{i \sum_j |\alpha_{n,j}|^2 \log \frac{|x-j|}{\sqrt{t}}} N_n(t, x) - \tilde{N}_n(0, x)| \leq nC \frac{\sqrt{t}}{x}, \quad x \in \left(0, \frac{1}{4}\right),$$

$$|e^{i \sum_{j \neq 0} |\alpha_{n,j}|^2 \log \frac{|j|}{\sqrt{t}}} N_n(t, 0) - \tilde{N}_n(0, 0)| \leq nC \sqrt{t}.$$

The growth in n comes from $\|\{\alpha_j\}\|_{l^{1,1}}$. Moreover, Lemma 4.6 in [4] insures us that

$$\tilde{N}_n(0, x_1) = \tilde{N}_n(0, x_2), \quad x_1, x_2 \in \left(0, \frac{1}{4}\right).$$

We have for $t \in (0, T)$ and $x \in \left(0, \frac{1}{4}\right)$:

$$\begin{aligned} |\tilde{N}_n(0, 0) - \tilde{N}_n(0, 0^+)| &\leq |\tilde{N}_n(0, 0) - e^{i \sum_{j \neq 0} |\alpha_{n,j}|^2 \log \frac{|j|}{\sqrt{t}}} N_n(t, 0)| + \\ &|N_n(t, 0) - N_n(t, x)| + |e^{i \sum_{j \neq 0} |\alpha_{n,j}|^2 \log \frac{|j|}{|x-j|}} - 1| |N_n(t, x)| \\ &+ |e^{i |\alpha_{n,0}|^2 \log \frac{x}{\sqrt{t}}} - 1| |N_n(t, x)| + |e^{i \sum_j |\alpha_{n,j}|^2 \log \frac{|x-j|}{\sqrt{t}}} N_n(t, x) - \tilde{N}_n(0, x)|. \end{aligned}$$

In view of the above estimates we get for $t \in (0, T)$ and $x = \frac{1}{8}$:

$$(54) \quad |\tilde{N}_n(0, 0) - \tilde{N}_n(0, 0^+)| \leq nC \sqrt{t} + C \frac{\log n}{n^2} + C \frac{\log t}{n^2} + |N_n(t, 0) - N_n(t, x)|.$$

Now for the last term, we use (39) to write

$$N_n(t, 0) - N_n(t, x) = \int_0^x (-uT)(t, y) dy = - \int_0^x \sum_j e^{i \lfloor n^\nu \rfloor c_n^2 \log t} (\alpha_{n,j} + R_{n,j}(t)) \frac{e^{-i \frac{(y-j)^2}{4t}}}{\sqrt{t}} T_n(t, y) dy.$$

By using Proposition 2.1 we obtain that the term involving $R_{n,j}$ is controlled by $\frac{C}{n^2}$. On the remaining part we perform an integration by parts for $j \neq 0$:

$$\begin{aligned} \frac{e^{i[n^\nu]c_n^2 \log t}}{\sqrt{t}} \int_0^x \sum_{1 \leq |j| \leq n^\nu} \alpha_j e^{-i\frac{(y-j)^2}{4t}} T_n(t, y) dy &= \left[e^{i[n^\nu]c_n^2 \log t} 2i\sqrt{t} \sum_{1 \leq |j| \leq n^\nu} \alpha_j \frac{e^{-i\frac{(y-j)^2}{4t}}}{y-j} T_n(t, y) \right]_0^x \\ &\quad - e^{i[n^\nu]c_n^2 \log t} 2i\sqrt{t} \int_0^x \sum_{1 \leq |j| \leq n^\nu} \alpha_j e^{-i\frac{(y-j)^2}{4t}} \left(\frac{T_n(t, y)}{y-j} \right)_y dy. \end{aligned}$$

As $\partial_y T_n$ is upper-bounded by $\frac{C}{\sqrt{t}}$ we get that all this part is controlled by $C \frac{\log n}{n}$. We are left with the term for $j = 0$:

$$\alpha_0 \frac{e^{i[n^\nu]c_n^2 \log t}}{\sqrt{t}} \int_0^x e^{-i\frac{y^2}{4t}} T_n(t, y) dy.$$

The integral on $y \in (0, \sqrt{t})$ is upper-bounded by $\frac{C}{n}$. On the remaining region of integration we perform an integration by parts

$$\begin{aligned} \alpha_0 \frac{e^{i[n^\nu]c_n^2 \log t}}{\sqrt{t}} \int_{\sqrt{t}}^x e^{-i\frac{y^2}{4t}} T_n(t, y) dy &= \left[\alpha_0 e^{i[n^\nu]c_n^2 \log t} 2i\sqrt{t} \frac{e^{-i\frac{y^2}{4t}}}{y} T_n(t, y) \right]_{\sqrt{t}}^x \\ &\quad - \alpha_0 e^{i[n^\nu]c_n^2 \log t} 2i\sqrt{t} \int_{\sqrt{t}}^x e^{-i\frac{y^2}{4t}} \left(\frac{T_n(t, y)}{y} \right)_y dy. \end{aligned}$$

Again as $\partial_y T_n$ is upper-bounded by $\frac{C}{\sqrt{t}}$ we get that this part is controlled by $C \frac{\log t}{n}$.

Summarizing we have obtained that for $t \in (0, T)$

$$(55) \quad |\tilde{N}_n(0, 0) - \tilde{N}_n(0, 0^+)| \leq nC\sqrt{t} + C \frac{\log n}{n} + C \frac{\log t}{n}.$$

By choosing t small enough with respect to n we obtain that

$$N_n(0, 0) - \tilde{N}_n(0, 0^+) \xrightarrow{n \rightarrow \infty} 0,$$

and the Proposition follows by using the next lemma.

Lemma 5.2. *The following convergence holds:*

$$\lim_{n \rightarrow \infty} \tilde{N}_n(0, 0^+) = \left(0, \frac{1-i}{\sqrt{2}}, \frac{-1-i}{\sqrt{2}} \right).$$

Proof. We recall the results at hand of $\tilde{N}_n(0, 0^\pm)$: in view Lemma 4.4 and Lemma 4.7 of [4] there exist rotations $\tilde{\Theta}_{n,k}$ such that

$$\begin{aligned} T_n(0, k^+) &= \tilde{\Theta}_{n,k}(A_{|\alpha_{n,k}|}^+), \quad T_n(0, k^-) = \tilde{\Theta}_{n,k}(-A_{|\alpha_{n,k}|}^-), \\ \tilde{N}_n(0, k^\pm) &= e^{i \sum_{j \neq k} |\alpha_{n,k}|^2 \log |k-j|} e^{-i \arg(\alpha_{n,k})} \tilde{\Theta}_{n,k}(B_{|\alpha_{n,k}|}^\pm), \end{aligned}$$

and

$$T_n(0, 0^\pm) = \left(\sin \frac{\theta_n}{2}, \pm \cos \frac{\theta_n}{2}, 0 \right),$$

to fit the directions of $\partial_x \chi_n(0, 0^\pm)$ in (19). Therefore with respect to the notations in Lemma 4.4 and Lemma 4.7 of [4], the rotation $\tilde{\Theta}_{n,0}$ is determined by its values

$$\tilde{\Theta}_{n,0}(A_{c_n}^\pm) = \left(\sin \frac{\theta_n}{2}, \pm \cos \frac{\theta_n}{2}, 0 \right).$$

In particular

$$\tilde{\Theta}_{n,0}\left(\frac{A_{c_n}^+ \wedge A_{c_n}^-}{|A_{c_n}^+ \wedge A_{c_n}^-|}\right) = (0, 0, 1).$$

It follows that $\tilde{N}_n(0, 0^\pm)$ is determined in terms of $\alpha_n, A_{c_n}^\pm$ and $B_{c_n}^\pm$. More precisely, we decompose

$$(56) \quad \begin{cases} \Re(B_{c_n}^+) = a_{n,1}A_{c_n}^+ + a_{n,2}A_{c_n}^- + a_{n,3}\frac{A_{c_n}^+ \wedge A_{c_n}^-}{|A_{c_n}^+ \wedge A_{c_n}^-|}, \\ \Im(B_{c_n}^+) = b_{n,1}A_{c_n}^+ + b_{n,2}A_{c_n}^- + b_{n,3}\frac{A_{c_n}^+ \wedge A_{c_n}^-}{|A_{c_n}^+ \wedge A_{c_n}^-|}. \end{cases}$$

We have then

$$(57) \quad \begin{aligned} \tilde{N}_n(0, 0^+) &= e^{ic_n^2 2 \log(|n^\nu|!)} e^{-i \arg(\alpha_{n,0})} (\tilde{\Theta}_{n,0}(\Re(B_{c_n}^+)) + i \tilde{\Theta}_{n,0}(\Im(B_{c_n}^+))) \\ &= e^{ic_n^2 2 \log(|n^\nu|!)} e^{-i \arg(\alpha_{n,0})} \\ &\times \left(((a_{n,1} + a_{n,2}) \sin \frac{\theta_n}{2}, (a_{n,1} - a_{n,2}) \cos \frac{\theta_n}{2}, a_{n,3}) + i((b_{n,1} + b_{n,2}) \sin \frac{\theta_n}{2}, (b_{n,1} - b_{n,2}) \cos \frac{\theta_n}{2}, b_{n,3}) \right). \end{aligned}$$

Now to conclude we have to check if $\tilde{N}_n(0, 0^+)$ has a limit. On one hand θ_n converges to π , and on the other hand we have $\arg(\alpha_{n,0}) = 0$. Thus for having a limit for $\tilde{N}_n(0, 0^+)$ it is enough to show that $a_{n,1} + a_{n,2}, (a_{n,1} - a_{n,2}) \cos \frac{\theta_n}{2}, a_{n,3}$ are convergent, and similarly for the b -coefficients.

We recall that from Theorem 1 (iii)-(iv) in [29] we know that the unitary vectors $\Re(B_{c_n}^\pm), \Im(B_{c_n}^\pm), A_{c_n}^\pm$ satisfy:

$$(58) \quad B_{c_n}^\pm \perp A_{c_n}^\pm, \quad (A_{c_n,1}^+, A_{c_n,2}^+, A_{c_n,3}^+) = (A_{c_n,1}^-, A_{c_n,2}^-, A_{c_n,3}^-), \quad A_{c_n,1}^+ = e^{-\frac{\pi}{2}c_n^2},$$

and in particular

$$(59) \quad \frac{A_{c_n}^+ \wedge A_{c_n}^-}{|A_{c_n}^+ \wedge A_{c_n}^-|} = \frac{(0, A_{c_n,3}^+, -A_{c_n,2}^+)}{\sqrt{1 - e^{-\pi c_n^2}}}, \quad \langle A_{c_n}^-, A_{c_n}^+ \rangle = 2e^{-\pi c_n^2} - 1.$$

We also have from formulas (55), (47), (48), (69) and (56) in [29]:

$$(60) \quad B_{c_n}^+ \stackrel{n \rightarrow \infty}{\approx} \left(-c_n \sqrt{\frac{\pi}{2}}, -1, 0\right) + i \left(c_n \sqrt{\frac{\pi}{2}}, 0, -1\right), \quad A_{c_n}^+ \stackrel{n \rightarrow \infty}{\approx} \left(e^{-\frac{\pi}{2}c_n^2}, -c_n \sqrt{\frac{\pi}{2}}, c_n \sqrt{\frac{\pi}{2}}\right).$$

First, by looking at the first two coordinates of the decomposition (56) we get

$$(61) \quad (\Re(B_{c_n}^+))_1 = (a_{n,1} + a_{n,2})e^{-\frac{\pi}{2}c_n^2},$$

and

$$(62) \quad (\Re(B_{c_n}^+))_2 = (a_{n,1} - a_{n,2})A_{c_n,2}^+ + a_{n,3}\frac{A_{c_n,3}^+}{\sqrt{1 - e^{-\pi c_n^2}}},$$

and similarly for b_n using $\Im(B_{c_n}^+)$.

Secondly, from the orthogonality relation in (58) we have

$$\begin{cases} 0 = \langle \Re(B_{c_n}^+), A_{c_n}^+ \rangle = a_{n,1} + a_{n,2} \langle A_{c_n}^-, A_{c_n}^+ \rangle = a_{n,1} + a_{n,2}(2e^{-\pi c_n^2} - 1), \\ 0 = \langle \Im(B_{c_n}^+), A_{c_n}^+ \rangle = b_{n,1} + b_{n,2} \langle A_{c_n}^-, A_{c_n}^+ \rangle = b_{n,1} + b_{n,2}(2e^{-\pi c_n^2} - 1), \end{cases}$$

thus by using also (61),

$$(63) \quad \begin{cases} a_{n,1} = a_{n,2}(1 - 2e^{-\pi c_n^2}), & a_{n,2} = \frac{\Re(B_{c_n}^+)e^{\frac{\pi}{2}c_n^2}}{2(1 - e^{-\pi c_n^2})}, \\ b_{n,1} = b_{n,2}(1 - 2e^{-\pi c_n^2}), & b_{n,2} = \frac{\Im(B_{c_n}^+)e^{\frac{\pi}{2}c_n^2}}{2(1 - e^{-\pi c_n^2})}. \end{cases}$$

Finally, by using (60) we obtain

$$(64) \quad a_{n,1} \stackrel{n \rightarrow \infty}{\approx} \frac{1}{c_n 2\sqrt{2\pi}}, \quad a_{n,2} \stackrel{n \rightarrow \infty}{\approx} -\frac{1}{c_n 2\sqrt{2\pi}}, \quad b_{n,1} \stackrel{n \rightarrow \infty}{\approx} -\frac{1}{c_n 2\sqrt{2\pi}}, \quad b_{n,2} \stackrel{n \rightarrow \infty}{\approx} \frac{1}{c_n 2\sqrt{2\pi}}.$$

Then, in view of (60) and (62) we get

$$(65) \quad a_{n,3} \xrightarrow{n \rightarrow \infty} -\frac{1}{\sqrt{2}}, \quad b_{n,3} \xrightarrow{n \rightarrow \infty} -\frac{1}{\sqrt{2}}.$$

In particular

$$(a_{n,1} - a_{n,2}) \cos \frac{\theta_n}{2} \xrightarrow{n \rightarrow \infty} \frac{1}{\sqrt{2}}, \quad (b_{n,1} - b_{n,2}) \cos \frac{\theta_n}{2} \xrightarrow{n \rightarrow \infty} -\frac{1}{\sqrt{2}},$$

and from (61) and (60) we have

$$a_{n,1} + a_{n,2} \xrightarrow{n \rightarrow \infty} 0, \quad b_{n,1} + b_{n,2} \xrightarrow{n \rightarrow \infty} 0.$$

Then, from (57), the convergence and the explicit limit of $\tilde{N}_n(0, 0^+)$ as n goes to infinity follow. □

□

6. PROOF OF THEOREM 1.1 IN THE NONTRIVIAL TORSION CASE

First we shall recall the construction of the solutions to the binormal flow which at an initial time are given by non-planar polygonal lines. In [4] we showed that given a sequence of the complex numbers $\{\alpha_n\}_{n \in \mathbb{Z}}$, with at least two non-trivial values, starting from the Schrödinger solution in (22) we obtain the BF evolution of a polygonal line $\chi(0)$ fully characterized modulo translation and rotation by the following description of its curvature and torsion angles.

We denote $n_0, n_1 \in \mathbb{Z}$ two consecutive locations where the sequence is non-trivial: $\alpha_{n_0} \neq 0, \alpha_{n_1} \neq 0, \alpha_k = 0, \forall n_0 < k < n_1$. We consider the ordered set of integers $\mathcal{L} = \{n_k\}_{k \in L}$ where the sequence does not vanish, containing in particular n_0 and n_1 , so L stand for a finite or infinite set of consecutive integers including 0 and 1. Then \mathcal{L} represents the locations of the corners of the polygonal line $\chi(0)$. At a corner located at $n_k \in \mathcal{L}$, the

curvature angle is the one of the selfsimilar BF solution $\chi_{|\alpha_{n_k}|}$, that is θ_k given by formula (11):

$$\sin \frac{\theta_k}{2} = e^{-\frac{\pi}{2}|\alpha_{n_k}|^2}.$$

For $k, k+1 \in L, k \geq 0$ the torsion is determined by the identities (66)

$$\begin{cases} \frac{T_{k-1} \wedge T_k}{|T_{k-1} \wedge T_k|} \cdot \frac{T_k \wedge T_{k+1}}{|T_k \wedge T_{k+1}|} = -\cos(f_{|\alpha_{n_k}|} - f_{|\alpha_{n_{k+1}}|} + \text{Arg}(\alpha_{n_k}) - \text{Arg}(\alpha_{n_{k+1}})), \\ \frac{T_{k-1} \wedge T_k}{|T_{k-1} \wedge T_k|} \wedge \frac{T_k \wedge T_{k+1}}{|T_k \wedge T_{k+1}|} = -\text{sgn}(\sin(f_{|\alpha_{n_k}|} - f_{|\alpha_{n_{k+1}}|} + \text{Arg}(\alpha_{n_k}) - \text{Arg}(\alpha_{n_{k+1}})))T_k, \end{cases}$$

where we denoted $T_k = \partial_x \chi(0, x)$ for $x \in (n_k, n_{k+1})$, and f is a function described in [4] that we will not explicit here as we will work with sequences $\{\alpha_n\}_{n \in \mathbb{Z}}$ of same modulus. For negative subindices $k < 0$ the torsion is defined in a similar way.

In particular, by considering the sequence of complex numbers in (25) we obtain indeed the BF evolution of the helicoidal polygonal lines $\chi_n(0)$ in the statement of Theorem 1.1.

The proof of Theorem 1.1 in the nontrivial torsion case goes the same as for the planar case treated in the previous subsections, except we get a modification of the term yielding Riemann's function in Lemma 3.1. More precisely, we obtain

$$(67) \quad n(\chi_n(t, 0) - \chi_n(0, 0)) - \frac{\Gamma}{2\sqrt{\pi}} \Im \left((0, \frac{1-i}{\sqrt{2}}, \frac{-1-i}{\sqrt{2}}) \int_0^t \sum_{|j| \leq n^\nu} e^{-ij\omega_0} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} d\tau \right) \xrightarrow{n \rightarrow \infty} 0,$$

uniformly on $(0, T)$. We shall treat the integral as in the proof of Lemma 3.1. First the summation can be extended to the whole set of integers as by integrating by parts

$$\left| \int_0^t \sum_{|j| > n^\nu} e^{-ij\omega_0} \frac{e^{-i\frac{j^2}{4\tau}}}{\sqrt{\tau}} d\tau \right| \leq \frac{C}{n^\nu}.$$

Thus we have to analyze

$$(68) \quad I(t) = \int_0^t \sum_{j \in \mathbb{Z}} e^{-ij\omega_0} e^{-i\frac{j^2}{4\tau}} \frac{d\tau}{\sqrt{\tau}} = \int_0^t e^{i\tau\omega_0^2} \sum_{j \in \mathbb{Z}} e^{-i\frac{(j+2\tau\omega_0)^2}{4\tau}} \frac{d\tau}{\sqrt{\tau}}.$$

Now we use again Poisson's summation formula $\sum_{j \in \mathbb{Z}} f(j) = \sum_{j \in \mathbb{Z}} \hat{f}(2\pi j)$ to get

$$\begin{aligned} \sum_{j \in \mathbb{Z}} e^{-i2\pi jr} e^{i4\pi^2 t j^2} &= \sum_{j \in \mathbb{Z}} \int e^{-i2\pi x(j+r) + i4\pi^2 t x^2} dx = \frac{1}{2\pi\sqrt{t}} \sum_{j \in \mathbb{Z}} \int e^{-ix\frac{j+r}{\sqrt{t}} + ix^2} dx \\ &= \frac{1}{2\pi\sqrt{t}} \sum_{j \in \mathbb{Z}} \widehat{e^{i \cdot^2}} \left(\frac{j-r}{\sqrt{t}} \right) = \frac{e^{i\frac{\pi}{4}}}{2\sqrt{\pi}\sqrt{t}} \sum_{j \in \mathbb{Z}} e^{-i\frac{(j-r)^2}{4t}}. \end{aligned}$$

By using this formula with $r = 2\tau\omega_0$ we obtain

$$(69) \quad I(t) = 2\sqrt{\pi} e^{-i\frac{\pi}{4}} \int_0^t e^{i\tau\omega_0^2} \sum_{j \in \mathbb{Z}} e^{-i4\pi\tau\omega_0 j} e^{i4\pi^2 \tau j^2} d\tau = 2\sqrt{\pi} e^{-i\frac{\pi}{4}} \int_0^t \sum_{j \in \mathbb{Z}} e^{i\pi^2 \tau (2j - \frac{\omega_0}{\pi})^2} d\tau.$$

We recall now that ω_0 is of rational type, $\omega_0 = \frac{a}{b}\pi$, with a, b with no common divisors, thus

$$I(t) = 2\sqrt{\pi}e^{-i\frac{\pi}{4}} \int_0^t \sum_{j \in \mathbb{Z}} e^{i\pi^2 \tau (2j - \frac{a}{b})^2} d\tau = \frac{2b^2}{\pi\sqrt{\pi}} e^{-i\frac{\pi}{4}} \sum_{j \in \mathbb{Z}} \frac{e^{i\pi^2 \frac{t}{b^2} (2bj - a)^2} - 1}{i(2bj - a)^2}.$$

In view of (67)-(68) we obtain

$$(70) \quad n(\chi_n(t, 0) - \chi_n(0, 0)) - \Im((0, -i, -1) \Gamma \sum_{j \in \mathbb{Z}} \frac{e^{i4\pi^2 t (j - \frac{a}{2b})^2} - 1}{i4\pi^2 (j - \frac{a}{2b})^2}) \xrightarrow{n \rightarrow \infty} 0.$$

The nontrivial torsion case of Theorem 1.1 with

$$(71) \quad \tilde{\mathfrak{R}}(t) = -\Gamma \sum_{j \in \mathbb{Z}} \frac{e^{i4\pi^2 t (j - \frac{a}{2b})^2} - 1}{i4\pi^2 (j - \frac{a}{2b})^2},$$

follows from the following Propositions 6.1-6.3.

Proposition 6.1. *Let $n \in \mathbb{N}, m \in \mathbb{N}^*$. The spectrum of singularities of the function*

$$(72) \quad \mathfrak{R}_{n,m}(t) = \sum_{j \in \mathbb{Z}} \frac{e^{i2\pi t (mj - n)^2} - 1}{(mj - n)^2},$$

enjoys the same property (3) of Riemann's function:

$$d_{\mathfrak{R}_{n,m}}(\beta) = 4\beta - 2, \quad \forall \beta \in \left[\frac{1}{2}, \frac{3}{4}\right].$$

Proof. We first notice that Chamizo and Ubis prove in Theorem 2.3 of [14] that the above sum with denominator j^2 instead of $(mj - n)^2$ has the same spectrum as Riemann's function. The proof we give below follows very closely their method and the one by Oskholkov and Chakhkiev [40].

We start with finding asymptotics of $\mathfrak{R}_{n,m}$ near rational points.

Lemma 6.2. *Let $p \in \mathbb{N}, q \in \mathbb{N}^*$ with no common divisors. There exist a positive constants C such that we have the following estimate ³:*

$$(73) \quad \left| \mathfrak{R}_{n,m}\left(\frac{p}{q} + h\right) - \mathfrak{R}_{n,m}\left(\frac{p}{q}\right) - \frac{z_0}{m} \frac{\tau_0}{q} \sqrt{h} \right| \leq C \min(|h|\sqrt{q}, |hq|^{3/2}),$$

where $z_0 \in \mathbb{C}^*$ and

$$\tau_0 = \sum_{r=0}^{q-1} e^{i2\pi \frac{p(mr-n)^2}{q}}.$$

For q odd $\tau_0 = \sqrt{q}$.

³ $\sqrt{-1} = e^{i\frac{\pi}{2}}$.

Proof. We decompose $j = ql + r$, with $l \in \mathbb{Z}$ and $0 \leq r < q$ and write

$$\mathfrak{R}_{n,m}\left(\frac{p}{q}\right) = \sum_{j \in \mathbb{Z}} \frac{e^{i2\pi \frac{p}{q}(mj-n)^2} - 1}{(mj-n)^2} = \sum_{r=0}^{q-1} \sum_{l \in \mathbb{Z}} \frac{e^{i2\pi \frac{p}{q}(mr-n)^2} - 1}{(m(ql+r)-n)^2}.$$

Then, assuming for simplicity that $h > 0$,

$$\mathfrak{R}_{n,m}\left(\frac{p}{q} + h\right) - \mathfrak{R}_{n,m}\left(\frac{p}{q}\right) = \sum_{r=0}^{q-1} e^{i2\pi \frac{p}{q}(mr-n)^2} \sum_{l \in \mathbb{Z}} \frac{e^{i2\pi h(m(ql+r)-n)^2} - 1}{(m(ql+r)-n)^2}.$$

We use Poisson's summation formula $\sum_{l \in \mathbb{Z}} f(l) = \sum_{l \in \mathbb{Z}} \hat{f}(2\pi l)$ in the second summation:

$$\mathfrak{R}_{n,m}\left(\frac{p}{q} + h\right) - \mathfrak{R}_{n,m}\left(\frac{p}{q}\right) = \sum_{r=0}^{q-1} e^{i2\pi \frac{p}{q}(mr-n)^2} \sum_{l \in \mathbb{Z}} \int \frac{e^{i2\pi h(m(qx+r)-n)^2} - 1}{(m(qx+r)-n)^2} e^{-i2\pi lx} dx.$$

By changing variable $y = m(qx+r) - n$, and $s = \sqrt{h}y$, we have

$$\begin{aligned} \mathfrak{R}_{n,m}\left(\frac{p}{q} + h\right) - \mathfrak{R}_{n,m}\left(\frac{p}{q}\right) &= \frac{1}{dq} \sum_{r=0}^{q-1} e^{i2\pi \frac{p}{q}(mr-n)^2} \sum_{l \in \mathbb{Z}} \int \frac{e^{i2\pi hy^2} - 1}{y^2} e^{-i2\pi l \frac{y-(mr-n)}{dq}} dy \\ &= \frac{1}{dq} \sum_{r=0}^{q-1} e^{i2\pi \frac{p}{q}(mr-n)^2} \sum_{l \in \mathbb{Z}} e^{i2\pi l \frac{(mr-n)}{dq}} \int \frac{e^{i2\pi hy^2} - 1}{y^2} e^{-i2\pi l \frac{y}{mq}} dy \\ &= \frac{\sqrt{h}}{mq} \sum_{r=0}^{q-1} e^{i2\pi \frac{p}{q}(mr-n)^2} \sum_{l \in \mathbb{Z}} e^{i2\pi l \frac{(mr-n)}{mq}} \int \frac{e^{i2\pi s^2} - 1}{s^2} e^{-i2\pi l \frac{s}{\sqrt{hmq}}} ds \\ &= \frac{\sqrt{h}}{mq} \sum_{r=0}^{q-1} e^{i2\pi \frac{p}{q}(mr-n)^2} \sum_{l \in \mathbb{Z}} e^{i2\pi l \frac{(mr-n)}{mq}} J\left(\frac{2\pi l}{\sqrt{hmq}}\right), \end{aligned}$$

where

$$(74) \quad J(x) = \int \frac{e^{i2\pi s^2} - 1}{s^2} e^{-isx} ds.$$

We note that $J(0) \neq 0$ is well-defined and by integrating by parts twice we get for $|x| > 1$, see [40]

$$(75) \quad |J(x)| \leq \frac{C}{x^2}.$$

Thus

$$\begin{aligned} \mathfrak{R}_{n,m}\left(\frac{p}{q} + h\right) - \mathfrak{R}_{n,m}\left(\frac{p}{q}\right) &= \frac{\sqrt{h}}{mq} \sum_{l \in \mathbb{Z}} J\left(\frac{2\pi l}{\sqrt{hmq}}\right) \sum_{r=0}^{q-1} e^{i2\pi \frac{p}{q}(mr-n)^2} e^{i2\pi l \frac{(mr-n)}{mq}} \\ &= \frac{\sqrt{h}}{mq} \sum_{l \in \mathbb{Z}} J\left(\frac{2\pi l}{\sqrt{hmq}}\right) \tau_l, \end{aligned}$$

with

$$\tau_l = e^{i2\pi l \frac{n}{mq}} \sum_{r=0}^{q-1} e^{i2\pi \frac{p(mr-n)^2 + 2lr}{q}}.$$

Now we recall that the classical bounds on Gauss sums yields $|\tau_l| \leq \sqrt{2q}$, with $|\tau_l| = \sqrt{q}$ for q qdd. Then by using the estimate (75) on $J(\frac{2\pi l}{\sqrt{hmq}})$ for all l that gets summability in l , we obtain the upper-bound $|hq|^{3/2}$ in (73) with $z_0 = J(0)$. Moreover, for $\sqrt{t}q > e$ upper-bound $J(\frac{2\pi l}{\sqrt{hmq}})$ by a constant for $|l| \leq \sqrt{h}q$ and the estimate (75) for the remaining l 's. This yields the upper-bound $|h|\sqrt{q}$ in (73) with $z_0 = J(0)$, so the proof of the Lemma is complete. \square

Based on these estimates around rational points one can follow the proof of Theorem 2.3 in [14] to get that the spectrum of singularities of $\mathfrak{R}_{n,m}$ is given by (3). For the sake of completeness we give here shortly the argument, based on approximations by continued fractions (a.b.c.f.). Consider for $2 \leq r \leq \infty$

$$A_r = \{x \in [0, 1] \setminus \mathbb{Q} \text{ such that } |x - \frac{p_k}{q_k}| = \frac{1}{q_k} \text{ with } \frac{p_k}{q_k} \text{ the a.b.c.f of } x, \limsup r_k = r\},$$

and A_r^* the same set with the extra-condition that for a subsequence q_{k_m} are odd numbers and r_{k_m} converges to r . Concerning the Hausdorff dimensions of these sets one has from Jarník's theorems ([23]):

$$(76) \quad \dim_{\mathcal{H}} A_r = \dim_{\mathcal{H}} A_r^* = \dim_{\mathcal{H}} \cup_{s \geq r} A_s = \frac{2}{r}.$$

For $x \in A_r$ and a small $h \neq 0$ there exists k such that

$$|x - \frac{p_k}{q_k}| = \frac{1}{q_k} \leq |h| < \frac{1}{q_{k-1}} = |x - \frac{p_{k-1}}{q_{k-1}}|.$$

As for a.b.c.f. $|x - \frac{p_k}{q_k}| \leq \frac{1}{q_k q_{k+1}}$ we obtain

$$|h|^{-\frac{1}{r_k}} \leq q_k < |h|^{-1 + \frac{1}{r_{k-1}}}.$$

Then, combining with (73), we have

$$(77) \quad \begin{aligned} |\mathfrak{R}_{n,m}(x+h) - \mathfrak{R}_{n,m}(x)| &\leq |\mathfrak{R}_{n,m}(x+h) - \mathfrak{R}_{n,m}(\frac{p_k}{q_k}) - \frac{z_0 \tau_0}{m q_k} \sqrt{x+h - \frac{p_k}{q_k}}| \\ &+ |\mathfrak{R}_{n,m}(x) - \mathfrak{R}_{n,m}(\frac{p_k}{q_k}) - \frac{z_0 \tau_0}{d q_k} \sqrt{x - \frac{p_k}{q_k}}| + |\frac{z_0 \tau_0}{d q_k} \sqrt{x+h - \frac{p_k}{q_k}} - \frac{z_0 \tau_0}{m q_k} \sqrt{x - \frac{p_k}{q_k}}| \\ &\leq Ch\sqrt{q_k} + C \frac{\sqrt{h}}{\sqrt{q_k}} \leq Ch^{\frac{1}{2} + \frac{1}{2r_{k-1}}} + Ch^{\frac{1}{2} + \frac{1}{2r_k}}. \end{aligned}$$

Finally we note that for q_k odd, by taking in (73) the value $h_k = x - \frac{p_k}{q_k}$ so that $|h_k| = \frac{1}{q_k}$, we have, as $r_k \geq 2$,

$$|\mathfrak{R}_{n,m}(x) - \mathfrak{R}_{n,m}(x - h_k)| \geq Ch_k^{\frac{1}{2} + \frac{1}{2r_k}} - Ch_k^{\frac{3}{2} - \frac{3}{2r_k}} \geq Ch_k^{\frac{1}{2} + \frac{1}{2r_k}}.$$

Then at $x \in A_r^*$ the function $\mathfrak{R}_{n,m}$ has local Hölder exponent $\frac{1}{2} + \frac{1}{2r}$ and by (76) we get

$$d_{\mathfrak{R}_{n,m}}\left(\frac{1}{2} + \frac{1}{2r}\right) \geq \dim_{\mathcal{H}} A_r^* = \frac{2}{r}.$$

We obtain the converse inequality by noting that if the function $\mathfrak{R}_{n,m}$ has local Hölder exponent $\frac{1}{2} + \frac{1}{2r}$ at some $x \in [0, 1]$, then from (77) the point x is either in \mathbb{Q} or in $\cup_{s \geq r} A_s$, of which Hausdorff dimension is given in (76). By varying $r \in [2, \infty]$ we recover the spectrum of $\mathfrak{R}_{n,m}$ to be the one given in (3). □

Proposition 6.3. *Let $n \in \mathbb{N}, m \in \mathbb{N}^*$. The function*

$$(78) \quad \mathfrak{R}_{n,m}(t) = \sum_{j \in \mathbb{Z}} \frac{e^{i2\pi t(mj-n)^2} - 1}{(mj-n)^2},$$

satisfies the multifractal formalism formula (4).

Proof. In view of Proposition 6.1 and of the multifractal formalism formula (4) we have to compute

$$\eta_{\mathfrak{R}_{n,m}}(p) = \sup\{s, \mathfrak{R}_{n,m} \in B_p^{\frac{s}{p}, \infty}\}.$$

Recall that Riemann's function φ_D is defined in (2). Jaffard proved in [32] the multifractal formalism formula (4) for φ_D using sharp bounds of the L^p norms of the partial sums of φ'_D from [47]. In the following we shall obtain similar bounds for $\mathfrak{R}'_{n,m}$.

Up to rescaling the variable t by $\frac{1}{m}$ we are thus interested in L^p norms of dyadic partial sums of

$$\sum_{j \in \mathbb{Z}} e^{i2\pi(tj^2 - t\frac{2n}{m}j)}.$$

We shall need the following general lemma on partial sums of exponential sums. It is a classical type of result, initially motivated by the study of Vinogradov's mean value conjecture. The short proof we give is based on the explicit expression at rational times of the fundamental solution of the linear Schrödinger equation with periodic boundary conditions given in (15). As it is well known this fundamental solution is intimately linked to the Talbot effect.

Lemma 6.4. *Let $N \in \mathbb{N}$ and $\sigma_N(x) = \sigma(\frac{x}{N})$ where σ is a smooth real positive function with compact support. Let $a, b, q \in \mathbb{N}$, $1 \leq a < q \leq N$, $(a, q) = 1$ and $0 \leq b < q$.*

i) For q odd there exists $\delta > 0$, depending just on σ , such that

$$(79) \quad C \frac{N}{\sqrt{q}} \leq \left| \sum_{j \in \mathbb{Z}} \sigma_N(j) e^{i2\pi(tj^2 - xj)} \right| \leq \tilde{C} \frac{N}{\sqrt{q}},$$

if

$$\left| t - \frac{a}{q} \right| \leq \frac{\delta}{N^2}, \quad \left| x - \frac{b}{q} \right| \leq \frac{\delta}{N}.$$

ii) We have the upper-bound

$$(80) \quad \left| \sum_{j \in \mathbb{Z}} \sigma_N(j) e^{i2\pi(tj^2 - xj)} \right| \leq C \frac{N}{\sqrt{q}(1 + N\sqrt{|t - \frac{a}{q}|})},$$

if

$$\left| t - \frac{a}{q} \right| < \frac{1}{qN}.$$

Remark 6.5. The bounds in (79) can be found in the literature for q varying up to \sqrt{N} , with $\delta = 1$ (see Lemma 1 in page 56 of [46], or see for instance (2.46)-(2.47) in [10]). A different proof of the upper-bound (80) for $q < N$ can be found in Lemma 3.18 in [10].

Proof. We denote $\tilde{t} = \frac{t}{2\pi}$. For $\tilde{t} = \frac{1}{2\pi} \frac{a}{q} + h$, we see the (conjugated) sum as a linear Schrödinger solution:

$$\begin{aligned} S_N(t, x) &:= \sum_{j \in \mathbb{Z}} \sigma_N(j) e^{-i4\pi^2 \tilde{t} j^2 + i2\pi x} = (e^{ih\Delta} e^{i\frac{a}{q}\Delta} (\sum_{j \in \mathbb{Z}} \sigma_N(j) e^{i2\pi \cdot j}))(x) \\ &= (e^{ih\Delta} e^{i\frac{a}{q}\Delta} (\mathcal{F}^{-1}(\sigma_N) \star \mathcal{F}^{-1}(\sum_{j \in \mathbb{Z}} e^{i2\pi \cdot j}))(x) = (e^{ih\Delta} e^{i\frac{a}{q}\Delta} (\mathcal{F}^{-1}(\sigma_N) \star (\sum_{j \in \mathbb{Z}} e^{i2\pi \cdot j}))(x) \\ &= (e^{ih\Delta} (\mathcal{F}^{-1}(\sigma_N) \star e^{i\frac{1}{2\pi} \frac{a}{q} \Delta} (\sum_{j \in \mathbb{Z}} e^{i2\pi \cdot j}))(x). \end{aligned}$$

Now we use the fact that (for instance choosing $M = 2\pi$ in formulas (37) combined with (42) from [18]):

$$e^{i\frac{1}{2\pi} \frac{a}{q} \Delta} (\sum_{j \in \mathbb{Z}} e^{i2\pi \cdot j})(x) = \sum_{j \in \mathbb{Z}} \tau_j \delta(x - \frac{j}{q}),$$

with the coefficients τ_j given in terms of Gauss sums. In particular $|\tau_j| = \frac{1}{\sqrt{q}}$ if q is odd and $|\tau_j| \leq \frac{\sqrt{2}}{\sqrt{q}}$. Then we have

$$\begin{aligned} S_N(t, x) &= (e^{ih\Delta} (\sum_{j \in \mathbb{Z}} \tau_j \mathcal{F}^{-1}(\sigma_N)(\cdot - \frac{j}{q}))(x) = \sum_{j \in \mathbb{Z}} \tau_j (e^{ih\Delta} (\mathcal{F}^{-1}(\sigma_N)))(x - \frac{j}{q}) \\ &= \sum_{j \in \mathbb{Z}} \tau_j \int e^{i(x - \frac{j}{q})\xi} e^{-ih\xi^2} \sigma_N(\xi) d\xi = N \sum_{j \in \mathbb{Z}} \tau_j \int e^{i(x - \frac{j}{q})N\xi} e^{-ihN^2\xi^2} \sigma(\xi) d\xi \\ &= N \sum_{j \in \mathbb{Z}} \tau_j \int e^{i(x - \frac{b}{q} + \frac{b-j}{q})N\xi} e^{-ihN^2\xi^2} \sigma(\xi) d\xi \\ &= N c_b \int e^{i(x - \frac{b}{q})N\xi} e^{-ihN^2\xi^2} \sigma(\xi) d\xi \\ &\quad + N \sum_{j \neq b} \tau_j \int e^{i(x - \frac{b}{q} + \frac{b-j}{q})N\xi} e^{-ihN^2\xi^2} \sigma(\xi) d\xi =: I_1 + I_2. \end{aligned}$$

The first term I_1 gives the right growth and is not zero if q is odd as $|hN^2| < \delta$ and $|x - \frac{b}{q}| < \frac{\delta}{N}$. The last sum I_2 can be upper-bounded by two integrations by parts from the ξ -linear phase, that insures summation in j and the upper-bound $C \frac{N}{\sqrt{q}} (\frac{q}{N})^2$. Thus (79) follows.

For getting the upper-bound (80) we write:

$$S_N(t, x) = N \sum_{j=qx} c_j \int e^{i(x-\frac{j}{q})N\xi} e^{-ihN^2\xi^2} \sigma(\xi) d\xi \\ + N \sum_{j \neq qx} \tau_j \int e^{i(x-\frac{j}{q})N\xi} e^{-ihN^2\xi^2} \sigma(\xi) d\xi =: \tilde{I}_1 + \tilde{I}_2.$$

In \tilde{I}_2 we use again integrations by parts yield the upper-bound $C \frac{N}{\sqrt{q}}$ for \tilde{I}_2 . For $|t - \frac{a}{q}| \leq \frac{\delta}{N^2}$ the upper-bound (80) follows by using the fact that the integrals in \tilde{I}_1 is bounded by a constant depending only on σ . For $\frac{\delta}{N^2} \leq |t - \frac{a}{q}| \leq \frac{1}{qN}$ we get (80) by using the dispersion inequality on \tilde{I}_1 yielding the upper-bound $C \frac{N}{\sqrt{q}} \frac{1}{\sqrt{hN^2}}$. \square

For $N \in \mathbb{N}$ we consider the L^p norm of the dyadic partial sum of $\mathfrak{R}'_{n,m}$:

$$I_{N,p} := \int_0^1 \left| \sum_{j \in \mathbb{Z}} \sigma_N(j) e^{i2\pi(tj^2 - t\frac{2n}{m}j)} \right|^p dt,$$

with σ a smooth real positive function valued 1 on $1 < |x| < 2$ and vanishing on $|x| < \frac{1}{2}$ and $|x| > 4$

First we get a lower bound by applying (79) with $x = t\frac{2n}{m}$, $a = m\tilde{a}$, $b = 2n\tilde{a}$, with q, N large with respect to n, m , with $(m\tilde{a}, q) = 1$. Indeed then the condition on x is satisfied as for these choices of parameters we have

$$\left| t - \frac{m\tilde{a}}{q} \right| \leq \frac{\delta}{N^2} \implies \left| x - \frac{2n\tilde{a}}{q} \right| = \left| t\frac{2n}{m} - \frac{2n\tilde{a}}{q} \right| \leq \frac{2n}{m} \frac{\delta}{N^2} + \left| \frac{m\tilde{a}}{q} \frac{2n}{m} - \frac{2n\tilde{a}}{q} \right| = \frac{2n}{m} \frac{\delta}{N^2} \leq \frac{\delta}{N}.$$

Thus we get the following lower bound by integrating on one region $|t - \frac{m\tilde{a}}{q}| \leq \frac{\delta}{N^2}$ for instance for $q = 2$ by applying (79):

$$I_{N,p} \geq \int_{|t - \frac{m\tilde{a}}{q}| \leq \frac{\delta}{N^2}} \left| \sum_{j \in \mathbb{Z}} \sigma_N(j) e^{i2\pi(tj^2 - \frac{2tn}{m}j)} \right|^p dt \geq CN^{p-2},$$

that will suit our purposes for $p > 4$. For $p = 4$ we shall need to improve this lower bound by integrating on the union of all the disjoint regions $|t - \frac{m\tilde{a}}{q}| \leq \frac{\delta}{N^2}$ for all $q \leq N$. We shall use rough estimates that are enough for our purposes:

$$I_{N,p} \geq CN^2 \sum_{q \leq N} \frac{\#\{1 \leq m\tilde{a} < q, (m\tilde{a}, q) = 1\}}{q^2} \geq CN^2 \sum_{q \leq N, q \text{ prime}} \frac{\lfloor \frac{q}{m} \rfloor - 1}{q^2}$$

$$\begin{aligned}
&\geq C(m)N^2 \sum_{q \leq N, q \text{ prime}} \frac{1}{q} \geq C(m)N^2 \sum_{j=1}^{\lfloor \log N \rfloor} \sum_{2^{j-1} \leq q < 2^j, q \text{ prime}} \frac{1}{q} \\
&\geq C(m)N^2 \sum_{j=1}^{\lfloor \log N \rfloor} \frac{1}{2^j} \#\{q, 2^{j-1} \leq q < 2^j, q \text{ prime}\}.
\end{aligned}$$

We use now the law of distribution of prime numbers, i.e. that the number of prime numbers less than a given number x grows as $\frac{x}{\log x}$ to get

$$I_{N,p} \geq C(m)N^2 \sum_{j=1}^{\lfloor \log N \rfloor} \frac{1}{j} \geq C(m)N^2 \log(\log N).$$

For the remaining case $p < 4$ we do the above calculation but just in the region $N/2 < q \leq N$. Thus we get the lower-bound:

$$(81) \quad I_{N,p} \geq C \begin{cases} N^{p-2}, & p > 4, \\ N^2 \log(\log N), & p = 4, \\ \frac{N^{\frac{p}{2}}}{\log N}, & 0 < p < 4. \end{cases}$$

Finally we shall get upper-bounds almost similar to the above lower-bound for $I_{N,p}$. From Dirichlet principle we have the following covering of $[0, 1]$: for all $t \in [0, 1]$, there is $a \in \mathbb{N}, q \in \mathbb{N}^*, (a, q) = 1, q \leq N$ such that

$$\left| t - \frac{a}{q} \right| < \frac{1}{qN}.$$

Then from (80),

$$\begin{aligned}
I_{N,p} &\leq C_p \int_0^1 \left| \sum_{j \in \mathbb{Z}} \sigma_N(j) e^{i2\pi(tj^2 - t\frac{2n}{m}j)} \right|^p dt \\
&\leq C_{m,p} \sum_{q=1}^N \sum_{a=0}^{q-1} \int_{|t - \frac{a}{q}| < \frac{1}{qN}} \left| \sum_{j \in \mathbb{Z}} \sigma_N e^{i2\pi(tj^2 - t\frac{2n}{m}j)} \right|^p dt \\
&\leq C_{m,p} \sum_{q=1}^N q \left(\frac{N^{p-2}}{q^{\frac{p}{2}}} + \frac{1}{q^{\frac{p}{2}}} \int_{\frac{1}{N^2} < |y| < \frac{1}{qN}} \frac{1}{|y|^{\frac{p}{2}}} dy \right).
\end{aligned}$$

For the first term we have the following estimate

$$\sum_{q=1}^N q \frac{N^{p-2}}{q^{\frac{p}{2}}} \leq C \begin{cases} N^{p-2}, & p > 4, \\ N^2 \log N, & p = 4, \\ N^{\frac{p}{2}}, & 0 < p < 4, \end{cases}$$

For the second term we use:

$$\int_{\frac{1}{N^2} < |y| < \frac{1}{qN}} \frac{1}{|y|^{\frac{p}{2}}} dy \leq C \begin{cases} N^{p-2}, & p > 2, \\ \log \frac{N}{q}, & p = 2, \\ \frac{1}{(qN)^{1-\frac{p}{2}}}, & 0 < p < 2, \end{cases}$$

that implies

$$\sum_{q=1}^N \frac{q}{q^{\frac{p}{2}}} \int_{\frac{1}{N^2}|y| < \frac{1}{qN}} \frac{1}{|y|^{\frac{p}{2}}} dy \leq C \begin{cases} N^{p-2}, & p > 4, \\ N^2 \log N, & p = 4, \\ N^{\frac{p}{2}}, & 0 < p < 4, \end{cases}$$

and thus the same upper-bounds for $I_{N,p}$. Combining with (81) we have obtained (82)

$$C \begin{cases} N^{p-2}, & p > 4, \\ N^2 \log(\log N), & p = 4, \\ \frac{N^{\frac{p}{2}}}{\log N}, & 0 < p < 4, \end{cases} \leq \int_0^1 |\sigma_N \sum_{j \in \mathbb{Z}} e^{i2\pi t(j^2 - \frac{2n}{m}j)}|^p dt \leq C \begin{cases} N^{p-2}, & p > 4, \\ N^2 \log N, & p = 4, \\ N^{\frac{p}{2}}, & 0 < p < 4. \end{cases}$$

Now we note that the function is m -periodic, with frequencies $j(j - \frac{2n}{m})$. As $D^k \leq |j(j - \frac{2n}{m})| < D^{k+1}$ implies $C_{n,m,D} D^{\frac{k}{2}} < |j| < \tilde{C}_{n,m,D} D^{\frac{k+1}{2}}$ we can use (82), where all the powers of N in the equivalent are positive, to obtain for large D -adic blocs Δ_k :

$$C \begin{cases} D^{\frac{k(p-2)}{2p}}, & p > 4, \\ D^{\frac{k}{2}} \log^{\frac{1}{2}} \log D^k, & p = 4, \\ \frac{D^{\frac{k}{4}}}{\log D^k}, & 0 < p < 4, \end{cases} \leq \|\Delta_k \sum_{j \in \mathbb{Z}} e^{i2\pi t(j - \frac{n}{m})^2}\|_{L^p} \leq C \begin{cases} D^{\frac{k(p-2)}{2p}}, & p > 4, \\ D^{\frac{k}{2}} \log^{\frac{1}{2}} D^k, & p = 4, \\ D^{\frac{k}{4}}, & 0 < p < 4. \end{cases}$$

Therefore we get that $\mathfrak{R}'_{n,m}$ belongs to $B_p^{-\frac{1}{4}, \infty}$ for $0 < p < 4$ and to $B_p^{-\frac{1}{2} + \frac{1}{p}, \infty}$ for $p > 4$, and these are optimal, yielding

$$\eta_{\mathfrak{R}_{n,m}}(p) = \sup\{s, \mathfrak{R}_{n,m} \in B_p^{s, \infty}\} = \begin{cases} 1 + \frac{p}{2}, & p \geq 4, \\ \frac{3p}{4}, & 0 < p < 4. \end{cases}$$

Then the relation (4) follows as for $\beta \in [\frac{1}{2}, \frac{3}{4}]$ we have

$$\begin{aligned} & \inf_p (\beta p - \eta_{\mathfrak{R}_{n,m}}(p) + 1) - d_{\mathfrak{R}_{n,m}}(\beta) \\ &= \inf(\{(\beta - \frac{1}{2})(p - 4), p \geq 4\} \cup \{(p - 4)(\beta - \frac{3p}{4}), 0 < p < 4\}) = 0. \end{aligned}$$

□

Remark 6.6. *Riemann's function (17) is the integral of the Jacobi theta function of one variable $\vartheta(t) = \sum_{j \in \mathbb{Z}} e^{i\pi t j^2}$, and it follows from Jaffard's results in [32] that it is a multifractal function satisfying the multifractal formalism of Frish and Parisi (4). From Propositions 6.1-6.3 it follows that these properties are also valid for the integrals of the Jacobi theta companions functions, for instance $\tilde{\vartheta}(t) = \sum_{j \in \mathbb{Z}} e^{i\pi t(j + \frac{1}{2})^2}$.*

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REFERENCES

- [1] R.A. Antonia, E.J. Hopfinger, Y. Gagne, and F. Anselmet, Temperature structure functions in turbulent shear flows, *Phys. Rev. A*, **30** (1984), 2704–2707.
- [2] R.J. Arms and F.R. Hama, Localized-induction concept on a curved vortex and motion of an elliptic vortex ring, *Phys. Fluids* **8** (1965), 553–560.
- [3] V. Banica and L. Vega, Stability of the selfsimilar dynamics of a vortex filament, *Arch. Ration. Mech. Anal.* **210** (2013), 673–712.
- [4] V. Banica and L. Vega, Evolution of polygonal lines by the binormal flow, *Ann. PDE*, **6** (2020), Paper No. 6, 53 pp.
- [5] V. Banica and L. Vega, On the energy of critical solutions of the binormal flow, *Comm. PDE*, **45** (2020), 820–845.
- [6] M. Berry, I. Marzoli and W. Schleich, Quantum carpets, carpets of light, *Physics World*, **14** (2001), 39–46.
- [7] M.V. Berry, Quantum fractals in boxes, *J. Phys. A: Math. Gen.* **29** (1996), 6617–6629.
- [8] M.V. Berry and S. Klein, Integer, fractional and fractal Talbot effects, *J. Mod. Opt.* **43** (1996), 2139–2164.
- [9] A. Boritchev, D. Eceizabarrena and V. Vilaça da Rocha, Riemann's non-differentiable function is intermittent, *ArXiv 1910.13191*.
- [10] J. Bourgain, Fourier transform restriction phenomena for certain lattice subsets and applications to nonlinear evolution equations, Part I: Schrödinger Equations, *Geom. Funct. Anal.* **3** (1993), 107–156.
- [11] T.F. Buttke, A numerical study of superfluid turbulence in the Self Induction Approximation, *J. of Compt. Physics* **76** (1988), 301–326.
- [12] J. Bedrossian, P. Germain et B. Harrop-Griffiths, Vortex filament solutions of the Navier-Stokes equations, *ArXiv 1809.04109*.
- [13] A.J. Callegari and L. Ting, Motion of a curved vortex filament with decaying vertical core and axial velocity, *SIAM. J. Appl. Math.* **35** (1978), 148–175.
- [14] F. Chamizo and A. Ubis, Multifractal behaviour of polynomial Fourier series, *Adv. Math.* **250** (2014), 1–34.
- [15] V. Chousionis, B. Erdogan and N. Tzirakis, Fractal solutions of linear and nonlinear dispersive partial differential equations, *Proc. London Math. Soc.* **110** (2015), 543–564.
- [16] L.S. Da Rios, On the motion of an unbounded fluid with a vortex filament of any shape, *Rend. Circ. Mat. Palermo* **22** (1906), 117–135.
- [17] F. de La Hoz, S. Kumar and L. Vega, On the evolution of the vortex filament equation for regular M-polygons with nonzero torsion, *SIAM J. Appl. Math.* (2020).
- [18] F. de la Hoz et L. Vega, *Vortex filament equation for a regular polygon*, *Nonlinearity* **27** (2014), 3031–3057.
- [19] F. de la Hoz et L. Vega, On the relationship between the one-corner problem and the M-corner problem for the vortex filament equation, *J. Nonlinear Sci.* **28** (2018), 2275–2327.
- [20] J.J. Duistermaat, Selfsimilarity of Riemann's nondifferentiable function, *Nieuw Arch. Wisk.* **9** (1991), 303–337.
- [21] D. Eceizabarrena, Geometric differentiability of Riemann's non-differentiable function, *Adv. Math.* **366** (2020).
- [22] M.B. Erdogan and N. Tzirakis, Talbot effect for the cubic non-linear Schrödinger equation on the torus, *Math. Res. Lett.* **20** (2013), 1081–1090.
- [23] K. Falconer, *Fractal Geometry: Mathematical Foundations and Applications*, second ed. John Wiley & Sons Inc., Hoboken, NJ, 2003.
- [24] U. Frisch, *Turbulence*, Cambridge University Press, Cambridge, 1995. The legacy of A. N. Kolmogorov.

- [25] U. Frisch and G. Parisi, On the singularity structure of fully developed turbulence, *Proc. Enrico Fermi International Summer School in Physics* (1985), 84–88. Appendix to Fully developed turbulence and intermittency, by U. Frisch.
- [26] J. Gerver, The differentiability of the Riemann function at certain rational multiples of π , *Amer. J. Math.* **92** (1970), 33–55
- [27] J. Gerver, More on the differentiability of the Riemann function, *Amer. J. Math.* **93** (1971), 33–41.
- [28] F.F. Grinstein, E. Gutmark and T. Parr, Near field dynamics of subsonic free square jets. A computational and experimental study, *Phys. Fluids* **7** (1995), 1483–1497.
- [29] S. Gutiérrez, J. Rivas and L. Vega, Formation of singularities and self-similar vortex motion under the localized induction approximation, *Comm. Part. Diff. Eq.* **28** (2003), 927–968.
- [30] G.H. Hardy, Weierstrass’ non-differentiable function, *Trans. Amer. Math. Soc.* **17** (1915), 301–325.
- [31] H. Hasimoto, A soliton in a vortex filament, *J. Fluid Mech.* **51** (1972), 477–485.
- [32] S. Jaffard, The spectrum of singularities of Riemann’s function, *Rev. Mat. Iberoam.* bf 2 (1996), 441–460.
- [33] S. Jaffard, Wavelet techniques in multifractal analysis, *Fractal Geometry and Applications: A Jubilee of Benot Mandelbrot*, M. Lapidus and M. van Frankenhuysen Eds., Proceedings of Symposia in Pure Mathematics, AMS, Vol. 72, Part 2, 91–152 (2004).
- [34] R. L. Jerrard et C. Seis, On the vortex filament conjecture for Euler flows, *Arch. Ration. Mech. Anal.* **224** (2017), 135–172.
- [35] R. L. Jerrard et D. Smets, On the motion of a curve by its binormal curvature, *J. Eur. Math. Soc.* **17** (2015), 1148–1515.
- [36] N. Kita, Mode generating property of solutions to the nonlinear Schrödinger equations in one space dimension, *Nonlinear dispersive equations, GAKUTO Internat. Ser. Math. Sci. Appl., Gakkotosho, Tokyo* **26** (2006), 111–128.
- [37] N. Koiso, Vortex filament equation and semilinear Schrödinger equation, *Nonlinear Waves, Hokkaido University Technical Report Series in Mathematics* **43** (1996), 221–226.
- [38] Y. Murakami, H. Takahashi, Y. Ukita and S. Fujiwara, On the vibration of a vortex filament, *Appl. Phys. Colloquium* (1937), 1–5.
- [39] K.I Oskolkov, A class of I.M. Vinogradov’s series and its applications in harmonic analysis, Springer Ser. Comput. Math. 19, Springer, New York, 1992, 353–402.
- [40] K.I Oskolkov, M. A. Chakhkiev, On the “nondifferentiable” Riemann function and the Schrödinger equation. *Proc. Steklov Inst. Math.* **280** (2013), 248–262.
- [41] C. S. Peskin and D. M. McQueen, Mechanical equilibrium determines the fractal fiber architecture of aortic heart valve leaflets. *Am. J. Physiol. (Heart Circ. Physiol. 35)* **266** (1994), 319–328.
- [42] I. Rodnianski, Fractal solutions of the Schrödinger equation, *Contemp. Math.* **255** (2000), 181–187.
- [43] J. Stern and C. Peskin, Fractal dimension of an aortic heart valve leaflet, *Fractals* **2** (1994), 461–464.
- [44] H.F. Talbot, Facts related to optical science, No. IV. *Philos. Mag.* **9** (1836), 401–407.
- [45] M. Taylor, The Schrödinger equation on spheres, *Pacific J. Math.* **209** (2003), 145–155.
- [46] I.M. Vinogradov, The method of trigonometrical series in the theory of numbers, translated by K. F. Roth, Anne Davenport, Interscience, London, 1954
- [47] Z. Zalcwasser, Sur les polynomes associés aux fonctions modulaires ϑ , *Studia Math.* **7** (1938), 16–35.

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