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Densities of Active species in N₂/H₂ RF and HF afterglows: application to surface nitriding of TiO₂ nanocrystals[★]

André Ricard^{1,*}, Jean-Philippe Sarrette¹, Yunfei Wang², and Yu-Kwon Kim²

¹ LAPLACE, Université de Toulouse, CNRS, INPT, UPS, 118 route de Narbonne, 31062 Toulouse, France

² Department of Chemistry and Department of Energy System Research, Ajou University, Suwon 443-749, South Korea

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Abstract. N₂/0–5% H₂ flowing afterglows from Radio Frequency (RF) and High Frequency (HF) sources have been analyzed by optical emission spectroscopy. In similar conditions (pressure 5–6 Torr, flow rate 0.5 slm and power 100 W), it is found in pure N₂ a nearly constant N-atom density from the pink to the late afterglow, which is higher in HF than in RF: (1–2) and $0.4 \times 10^{15} \text{ cm}^{-3}$, respectively. With a N₂/2% H₂ gas mixture, the early afterglows is changed to a late afterglow with about the same N-atom density for both RF and HF cases: $(8\text{--}9) \times 10^{14} \text{ cm}^{-3}$. Anatase TiO₂ nanocrystals and Atomic Layer Deposition-grown films were exposed to the RF afterglows at room temperature. XPS analysis of the samples has shown that the highest N/Ti ratio of 0.24 can be achieved with the pure N₂ late afterglow. In the HF pure N₂ late afterglow, however, the N/Ti coverage was limited to 0.04 in spite of higher N-atom density. Such differences in the N content between the two RF and HF cases are attributed to the presence of a high O-atom impurity of $2 \times 10^{13} \text{ cm}^{-3}$ in HF as compared to that ($8 \times 10^{11} \text{ cm}^{-3}$) in RF.

1 Introduction

N₂ plasmas have been frequently used for the surface nitridation of metals and oxides for various applications [1,2]. For the case of TiO₂, N₂ plasma can be used to make oxynitride layers within the TiO₂ matrix for enhanced visible light absorption and subsequent enhancement in photocatalytic activity [2]. However, the high temperature and the high density of ions and electrons in the plasmas may induce undesirable defects within the TiO₂ matrix that can induce a negative effect in the resulting photocatalytic performances. As an alternative damage-free strategy in the nitridation of TiO₂, afterglows or the post-discharge regions of N₂ plasmas can be used for a selective nitridation of TiO₂ surface since the densities of ions and electrons are much lower in the afterglows than in the plasmas [3].

With such a purpose in mind, we compare the afterglows of N₂ plasmas generated by two different excitation sources of High Frequency 2450 MHz (HF) and Radio Frequency 13.6 MHz (RF) in this study. Such a comparison was started in N₂ and Ar–N₂ flowing plasmas [4] showing that the RF discharge is more vibrationally excited than the HF one. As a consequence, it is produced more high N₂(X,v) vibrational levels in RF than in HF and then a more intense pink afterglow in RF.

The two setups A and B used in this work are similar with about same discharge tube inner dia. 6 (A) and 5 (B) mm and afterglow tube inner dia. 21 (A) and 18 (B) mm, same gas pressure (5–8 Torr), flow rate (0.5–1 slm), power

(100–150 W). The most significant difference between the setup A and B is in the oxygen impurity as demonstrated hereafter. Previous results have been obtained in N₂ and N₂–H₂ flowing RF [5] and HF afterglows [6,7]. By analyzing the radial distribution of the N₂ 1st pos (580 nm) system in the RF afterglow [5], it has been shown that the early afterglow changed from a pink to a late afterglow by introducing very small amounts of H₂ (from 10^{–5} to 10^{–3}) into N₂.

The kinetic reactions in the early and late afterglows involving the N atoms, the N₂(X,v) vibrationally excited ground state, the N₂(A) metastable molecules and the N₂⁺ ions have been previously studied in pure N₂ [8–10]. In RF and HF flowing afterglows, N-atoms densities have been measured by NO titration. The concentrations of O-atoms (due to the presence of oxygen species in the impurities of the gas tank), of N₂(A) and N₂(X,v > 13) metastable molecules and of N₂⁺ ions were determined from band-ratio intensities [11]. Concentrations of N-atoms, O-atoms in impurity, N₂(X,v > 13) and N₂(A) metastable molecules and N₂⁺ ions were monitored in N₂/ < 5% H₂ RF and HF afterglows.

2 Experimental setup

The experimental set-ups of the RF and HF plasmas and afterglows are reproduced in Figure 1 [4]. The discharge quartz tubes have nearly the same inner diameters (ID) of 6 mm for the RF plasma and 5 mm for the HF plasma and have a length of 30 cm in RF and 16 cm in HF.

In the post-discharge region, the RF discharge tube is connected to a straight quartz tube of 21 mm ID (Fig. 1a) while the HF discharge tube is connected to a bent quartz tube of 18 mm ID itself connected to a 5 L Pyrex chamber (Fig. 1c).

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* e-mail: ricard@laplace.univ-tlse

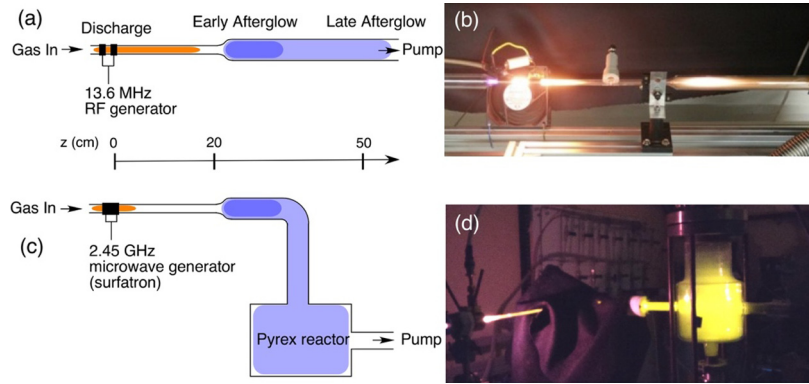


Fig. 1. Schematic diagrams and pictures of the RF (a and b) and the HF (c and d) set-ups A and B. In the RF (HF) set-up, the inner diameter of the quartz tube is 6 (5) mm in the discharge region and 21 (18) mm in the afterglow region.

The RF plasma is produced between two rings separated by 2 cm in the tube of 6 mm ID. The HF plasma is produced by a surfatron cavity [5]. In RF, with a N_2 flow rate $Q = 1$ slm, a pressure $p = 8$ Torr and an incident power of 100 W, a pink (early) afterglow extends on 15 cm in the 21 mm i.d. tube (from $z = 32$ cm to $z = 47$ cm, see Fig. 1a and b), followed by a late afterglow (after $z = 48$ cm). In HF, a pink afterglow is observed in the bent part of the 18 mm tube with a N_2 flow rate of 0.5 slm, a pressure $p = 5$ Torr and an incident power of 100 W. The NO titration of N-atoms in the RF and HF late afterglows was performed by introducing an Ar-1.5%NO flow after the pink afterglows. At 8 Torr, 0.5 slm and 100 W, it was previously obtained N-atom densities of 1.0 and 2.0, 10^{15} cm^{-3} , for respectively, HF and RF N_2 afterglow conditions [4].

The emission spectroscopy was performed by means of an optical fiber connected, in the RF set-up, to a Monera 500 spectrometer with 500 mm focal length, grating blazed at 500 nm, slits of 0.5 mm, resolution of 0.8 nm and a PMT (Hamamatsu R928 at 900 V). In the HF set-up, an optical fiber is connected to an Acton Spectra Pro 2500i spectrometer (grating 600 g/mm) equipped with a Pixis 256E CCD detector (front illuminated 1024×256 pixels).

The plasmas and the afterglows are characterized by the emission of the N_2 1st pos system in the red part of the spectrum between 500 and 1100 nm and of the N_2 2nd pos and N_2^+ 1st neg systems in the UV-blue part between 300 and 400 nm. The NO_β (320 nm) and NH (336 nm) emissions are also detected in the afterglows.

3 N_2 pink and late afterglows

As reported in [11], densities of N-atoms, O-atoms in impurity, $N_2(X, v > 13)$ and $N_2(A)$ metastable molecules and N_2^+ ions can be obtained by the line-ratio intensity method after calibration of the N-atom density by NO titration. The results have been recently published [12] for the RF afterglow at 6 Torr, 0.6 slm, 100 W. The N-atom density kept a nearly constant value of $4 \times 10^{14} \text{ cm}^{-3}$ from the pink to the late afterglow as the density of $N_2(X, v > 13)$ molecules and the N_2^+ ions decreased from 8 to 1 (10^{13} cm^{-3}) and from 2 to 0.4 (10^{10} cm^{-3}), respectively. The O-atom density remained at (8–11) (10^{11} cm^{-3}).

With the HF afterglow, it has been analyzed the afterglow from $z = 25$ to 50 cm in the tube of 18 mm i.d. (Fig. 1c). A pink afterglow was detected in the shoulder of the 18 mm i.d. tube between $z = 17$ and $z = 25$ cm. It is reported in Table 1 the N_2 active species densities versus the distance z in the HF afterglow at 5 Torr, 0.5 slm, 100 W.

The value of N-atom density is indicated in Table 1. It has been found to be nearly constant from the pink to the late HF afterglow.

Comparing with the densities measured in the RF afterglow [12], the concentrations of N, O, $N_2(A)$ and N_2^+ are higher in the HF afterglow. The $N_2(A)$ and N_2^+ density are found to decrease from the pink to the late HF afterglow.

4 $N_2/<5\%H_2$ plasma and early afterglow

As previously reported in [5], the RF pink afterglow disappeared when a few percent of H_2 is added to N_2 . The rate coefficients of NH(A) excitation in the afterglow are unknown and therefore the line ratio intensity method [11] cannot be applied to determine the NH and H densities.

The P1/P2 intensity ratio between two rotational subbands of the N_2 1st pos (2–0) vibrational transition at 775 nm can be related with the N_2 rotational temperature T_R , assumed to be equal to the gas temperature [4]. At 8 Torr, 0.5 slm and 100 W, it is found in the RF plasma an increase of T_R from 700 to 900 K when the percentage of H_2 introduced into N_2 is increased from 0.2% to 5%.

Concentrations of the active species densities in the RF afterglow at $z = 26$ cm are given in reference [12]. The N_2^+ density is not reported since the N_2^+ ions disappeared with the H_2 inlet.

The density of N-atoms remain almost constant when the $H_2\%$ was increased from 0.2% to 5% in the $N_2/<5\%H_2$ afterglows: $[N] = (7-8) (10^{14} \text{ cm}^{-3})$ but the O-atom density decreased from 8 to 2 (10^{13} cm^{-3}). Otherwise, the $N_2(X, v > 13)$ density decreased significantly from 8 to 1 (10^{13} cm^{-3}) as H_2 was introduced into N_2 , explaining the disappearance of the pink afterglow in the $N_2/<5\%H_2$ gas mixtures where they are produced by high $N_2(X, v > 13)$ vibrational Penning collisions [5,9].

In the HF afterglow, the condition of $N_2/2\%H_2$, 5 Torr, 0.5 slpm, 100 W has been chosen. The active species densities were measured in the afterglow for 3 positions

Table 1. N₂ active species densities determined in the HF afterglow of dia. 18 mm tube at 5 Torr, 0.5 slm, 100 W.

z (cm)	25 (end of the pink afterglow)	27	29	31	50 (late)
a_{N+N}	0.2	0.3	0.45	0.6–0.7	1.0
$[N]$ (10^{15} cm^{-3})	0.9	1.0	1.2	1.4–1.8	2.0
$[O]$ (10^{13} cm^{-3})	3.3	4.0	4.0	4.0–7.0	9.0
$[N_2(A)]$ (10^{11} cm^{-3})	2.3–4.0	2.4	1.7	1.3–1.6	
$[N_2(X, v > 13)]$ (10^{14} cm^{-3})	0.6–1.0	0.7	0.7	0.6–0.7	
$[N_2^+]$ (10^{10} cm^{-3})	20.0	2.0	0.5	0.4–0.5	

Table 2. N₂/2%H₂ active species densities determined in the HF afterglow at different z positions for 5 Torr, 0.5 slpm, 100 W.

z (cm)	17	25	50	5 L reactor
a_{N+N}	0.40	0.50	1.00	1.00
$[N]$ (10^{14} cm^{-3})	8.0	9.0	11.0	3.0
$[O]$ (10^{13} cm^{-3})	2.0	1.4	0.4	0.2
$[N_2(A)]$ (10^{11} cm^{-3})	5.0	2.0	0.7	–
$[N_2(X, v > 13)]$ (10^{13} cm^{-3})	1.3	3.0	–	–
$[N_2^+]$ (10^{10} cm^{-3})	4.5	0.4	–	–

along the 18 mm i.d. tube $z=17$, 25 and 50 cm and inside the 5 L reactor. The results are reported in Table 2. The N₂(A) density could not be measured in the 5 L reactor as the I_{316} intensity was too weak. Moreover, the N₂(X, $v > 13$) metastable and N₂⁺ ion concentrations could not be obtained for $z=50$ cm and in the 5 L reactor since the a_{N+N} factor was found equal to 1 (pure late afterglow).

As observed in RF afterglows, the O-atom density strongly decreases by adding 2%H₂ to N₂. The decrease from 2 to $0.2 \times 10^{13} \text{ cm}^{-3}$ appears to be sensitive. Thus, one of the interest of the N₂/ $<5\%$ H₂ gas mixture is for a surface treatment with reduced O-atom impurity.

5 TiO₂/Si substrate treatments

By XPS analysis of TiO₂ nanocrystals on Si samples treated by the N₂ RF afterglow at 6 Torr, 0.6 slpm, 100 W, we obtained a N/Ti ratio of 0.22–0.24 when the sample was treated at $z=61$ cm for 5 min in the N₂ late afterglow. The ratio does not increase after a longer treatment time up to 60 min. In these conditions, the measured active species densities are $[N]=4 \times 10^{14} \text{ cm}^{-3}$, $[O]=8 \times 10^{11} \text{ cm}^{-3}$, $[N_2(A)]=0.5 \times 10^{11} \text{ cm}^{-3}$, $[N_2(X, v > 13)]=1 \times 10^{13} \text{ cm}^{-3}$ and $[N_2^+]=0.4 \times 10^{10} \text{ cm}^{-3}$. On the other hand, in the N₂ HF late afterglow at $z=35$ cm, $Q=0.5$ slm, $p=7$ Torr and 150 W, the N/Ti coverage was limited to 0.04 in spite of higher N-atom density: $(1-2) \times 10^{15} \text{ cm}^{-3}$. Such a dramatic difference in the nitriding performance between the RF and HF afterglows can be explained by higher O-atom impurity in HF: $2 \times 10^{13} \text{ cm}^{-3}$. The O/N ratios in the RF and HF afterglows were 2×10^{-3} and $(1-2) \times 10^{-2}$, respectively. Thus the high O/N ratio in the HF afterglows is related to the low N/Ti ratio.

In the pure N₂ RF pink afterglow ($a_{N+N}=0$) at $z=26$ cm, the N content was only 3–4%. The pink afterglow is characterized by a high N₂⁺ density. With $a_{N+N}=0$, the N-atoms cannot be detected by the $a_{N+N} I_{580}$ intensity and the line-ratio method cannot be used. At the end of the pink afterglow, a N-atom density of $(3-4) \times 10^{14} \text{ cm}^{-3}$ and a density of N₂⁺ ions of $2 \times 10^{10} \text{ cm}^{-3}$ were measured [12]. It appears that a high N₂⁺ density is detrimental to achieving a high N coverage.

In the RF N₂/ $<5\%$ H₂ gas mixtures, the N content of 2–4% were obtained from the samples treated at $z=26$ cm for the treatment times up to 60 min. At this point, the measured active species densities are $[N]=7-8 \times 10^{14} \text{ cm}^{-3}$, $[O]=2 \times 10^{11} \text{ cm}^{-3}$, $[N_2(A)]=0.6 \times 10^{11} \text{ cm}^{-3}$ and $[N_2(X, v > 13)]=7-8 \times 10^{13} \text{ cm}^{-3}$. Compared with the pure N₂ afterglow, densities of active species are higher in N₂/ $<5\%$ H₂, except for the O-atoms. The addition of H₂ in the plasma could produce NH radicals and H-atoms which appeared to be not effective for an efficient N/Ti coverage in spite of higher N-atom density.

Under the present treatment conditions with the N₂ and N₂–H₂ afterglows, it can be concluded that treatment time of 5 min is sufficient for the saturation N coverage at room temperature (RT) for both RF and HF treatment conditions.

6 Conclusion

The RF N₂ early afterglow in a tube of 21 mm i.d. was transformed into a pink with a part a_{N+N} of N+N recombination between 0 and 0.1 to a late with $a_{N+N}=0.8$ when H₂ was introduced between 0.2% and 5% into N₂ at

8 Torr, 0.5 slm, 100 W. The pink afterglow was also observed after the HF N₂ plasma in a tube of 18 mm i.d. at 5 Torr, 0.5 slm, 100 W which disappeared as 2% of H₂ was introduced into N₂.

In the N₂–(0–5%)H₂ gas mixtures of RF plasma, the N-atoms, N₂(X, $v > 13$) and N₂(A) metastable molecules kept a nearly constant value of $(7-8) \times 10^{14} \text{ cm}^{-3}$, $(7-12) \times 10^{13} \text{ cm}^{-3}$ and $6 \times 10^{10} \text{ cm}^{-3}$, respectively. The emission of the N₂⁺ ions was too weak to be detected. The density of O-atoms in impurity decreased from 8 to $2 \times 10^{11} \text{ cm}^{-3}$. The N₂/H₂ gas mixtures appeared to be of interest to reduce the O-atom impurity in the afterglow for a surface treatment.

TiO₂ samples in forms of anatase nanocrystals and Atomic Layer Deposition-grown thin films were exposed to the RF and HF afterglows at RT. XPS analysis has shown that TiO₂ treated in the late afterglow of N₂ RF plasma has the highest N content (N/Ti ratio of 0.24). The high N content achieved is attributed to the high N-atom density in the afterglow. Under the very similar operating conditions, the treatment in the HF N₂ late afterglow resulted in low N/Ti ratios of 0.02–0.04 which was attributed to high O-atom impurity: O/N = 2×10^{-3} in RF and $(1-2) \times 10^{-2}$ in HF. Thus, it can be concluded that the densities of N and O atoms in the late afterglows play an important role in the surface nitriding performance of TiO₂ for applications in photocatalysis.

However, the present conclusion applies only to the case when the substrate is just introduced to the flowing afterglows operating at RT. Under this condition, the substrate temperature remains close to RT throughout the whole treatment process. We are in the process of further in-depth XPS analysis on the surface nitriding perfor-

mances of TiO₂ with HF N₂ and N₂/ $<2\%$ H₂ afterglows at elevated substrate temperatures, which is beyond the scope of the present study.

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