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► **To cite this version:**

Hedy Attouch, Radu Ioan Boț, Robert Csetnek. Fast optimization via inertial dynamics with closed-loop damping. 2020. hal-02912177

HAL Id: hal-02912177

<https://hal.science/hal-02912177>

Preprint submitted on 5 Aug 2020

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FAST OPTIMIZATION VIA INERTIAL DYNAMICS WITH CLOSED-LOOP DAMPING

HEDY ATTOUCH*, RADU IOAN BOT[†], AND ERNÖ ROBERT CSETNEK[‡]

Abstract. In a Hilbert space \mathcal{H} , in order to develop fast optimization methods, we analyze the asymptotic behavior, as time t tends to infinity, of inertial continuous dynamics where the damping acts as a closed-loop control. The function $f : \mathcal{H} \rightarrow \mathbb{R}$ to be minimized (not necessarily convex) enters the dynamic through its gradient, which is assumed to be Lipschitz continuous on the bounded subsets of \mathcal{H} . This gives autonomous dynamical systems with nonlinear damping and nonlinear driving force. We first consider the case where the damping term $\partial\phi(\dot{x}(t))$ acts as a closed-loop control of the velocity. The damping potential $\phi : \mathcal{H} \rightarrow \mathbb{R}_+$ is a convex continuous function which achieves its minimum at the origin. We show the existence and uniqueness of a global solution to the associated Cauchy problem. Then, we analyze the asymptotic convergence properties of the trajectories generated by this system. To do this, we use techniques from optimization, control theory, and PDE's: Lyapunov analysis based on the decreasing property of an energy-like function, quasi-gradient and Kurdyka-Lojasiewicz theory, monotone operator theory for wave-like equations. Convergence rates are obtained based on the geometric properties of the data f and ϕ . We put forward minimal hypotheses on the damping term guaranteeing the convergence of trajectories, thus showing the dividing line between strong and weak damping. When f is strongly convex, we give general conditions on the damping potential ϕ which provide exponential convergence rates. Then, we extend the results to the case where an additional Hessian-driven damping enters the dynamic, which reduces the oscillations. Finally, we consider an inertial system with a closed-loop damping involving jointly the velocity $\dot{x}(t)$ and the gradient $\nabla f(x(t))$. This study naturally leads to similar results for the proximal-gradient algorithms which can be derived by temporal discretization, some of them are studied in the article. In addition to its original results, this work surveys the numerous works devoted in recent years to the interaction between continuous damped inertial dynamics and numerical algorithms for optimization, with the emphasis on autonomous systems, closed-loop adaptive procedures, and convergence rates.

Key words. closed-loop damping; convergence rates; damped inertial gradient systems; Hessian damping; quasi-gradient systems; Kurdyka-Lojasiewicz inequality; maximally monotone operators.

AMS subject classifications. 37N40, 46N10, 49M30, 65K05, 65K10, 90B50, 90C25.

1. Introduction. Let \mathcal{H} be a real Hilbert space endowed with the scalar product $\langle \cdot, \cdot \rangle$ and norm $\| \cdot \|$. Let $f : \mathcal{H} \rightarrow \mathbb{R}$ be a differentiable function (not necessarily convex), whose gradient $\nabla f : \mathcal{H} \rightarrow \mathcal{H}$ is Lipschitz continuous on the bounded subsets of \mathcal{H} , and such that $\inf_{\mathcal{H}} f > -\infty$ (when considering the Hessian of f , we will assume that f is twice differentiable). Our objective is to study from the optimization point of view the Autonomous Damped Inertial Gradient Equation

$$(ADIGE) \quad \ddot{x}(t) + \mathcal{G}\left(\dot{x}(t), \nabla f(x(t)), \nabla^2 f(x(t))\right) + \nabla f(x(t)) = 0,$$

where the damping term $\mathcal{G}\left(\dot{x}(t), \nabla f(x(t)), \nabla^2 f(x(t))\right)$ acts as a closed-loop control. Under suitable assumptions, this term will induce dissipative effects, which tend to stabilize asymptotically (*i.e.* as $t \rightarrow +\infty$) the trajectories to critical points of f (minimizers in the case where f is convex). Since we work with autonomous (dissipative) systems, we can take an arbitrary initial time t_0 . As is usual, we take $t_0 = 0$, and hence work on the time interval $[0, +\infty[$. Our study aims to deepen

*IMAG, Université Montpellier, CNRS, Place Eugène Bataillon, 34095 Montpellier CEDEX 5, France. E-mail: hedy.attouch@umontpellier.fr. Supported by COST Action: CA16228

[†]University of Vienna, Faculty of Mathematics, Oskar-Morgenstern-Platz 1, A-1090 Vienna, Austria. E-mail: radu.bot@univie.ac.at.

[‡]University of Vienna, Faculty of Mathematics, Oskar-Morgenstern-Platz 1, A-1090 Vienna, Austria. E-mail: robert.csetnek@univie.ac.at. Research supported by FWF (Austrian Science Fund), Project P 29809-N32.

the understanding of the rich interactions between dissipative dynamic systems and optimization algorithms. Indeed, our first motivation was to understand if the closed-loop damping can do as well (and possibly improve) the fast convergence properties of the accelerated gradient method of Nesterov. Closely related questions concern the search for minimum hypotheses on the damping term guaranteeing rapid convergence of values and convergence of trajectories. This is a wide subject which concerns continuous optimization, as well as the study of the stabilization of oscillating systems in mechanics and physics. Due to the highly nonlinear characteristics of (ADIGE) (non-linearity occurs both in the damping term and in the gradient of f), our convergence analysis will mainly rely on the combination of the quasi-gradient approach for inertial systems initiated by Bégout–Bolte–Jendoubi [41] with the theory of Kurdyka–Lojasiewicz. The price to pay is that some of the results are only valid in finite dimensional Hilbert spaces.

1.1. Presentation of the results.

1.1.1. (ADIGE-V). Our study mainly concerns the case where \mathcal{G} only depends on the velocity. So, we consider the differential inclusion

$$(1.1) \quad (\text{ADIGE-V}) \quad 0 \in \ddot{x}(t) + \partial\phi(\dot{x}(t)) + \nabla f(x(t)),$$

where $\phi : \mathcal{H} \rightarrow \mathbb{R}$ is a convex continuous function that achieves its minimum at the origin, and the operator $\partial\phi : \mathcal{H} \rightarrow 2^{\mathcal{H}}$ is its convex subdifferential. The suffix V refers to the velocity as a control variable. This model encompasses several classical situations:

- The case $\phi(u) = \frac{\gamma}{2}\|u\|^2$ corresponds to the Heavy Ball with Friction method

$$(1.2) \quad (\text{HBF}) \quad \ddot{x}(t) + \gamma\dot{x}(t) + \nabla f(x(t)) = 0$$

introduced by B. Polyak [77, 78] and further studied by Attouch–Goudou–Redont [26] (exploration of local minima), Alvarez [7] (convergence in the convex case), Haraux–Jendoubi [64, 65] (convergence in the analytic case), Bégout–Bolte–Jendoubi [41] (convergence based on the Kurdyka–Lojasiewicz property), to cite part of the rich literature devoted to this subject.

- The case $\phi(u) = r\|u\|$ corresponds to the dry friction effect. Then, (ADIGE-V) is a differential inclusion, which, when the velocity is not equal to zero, writes as follows

$$\ddot{x}(t) + r \frac{\dot{x}(t)}{\|\dot{x}(t)\|} + \nabla f(x(t)) = 0.$$

The importance of this case in optimization comes from the finite time stabilization property of the trajectories, which is satisfied generically with respect to the initial data. The rigorous mathematical treatment of this case has been considered by Adly–Attouch–Cabot [5], Amann–Diaz [10], see Adly–Attouch [2, 3, 4] for recent developments.

- Taking $\phi(u) = \frac{1}{p}\|u\|^p$ with $p \geq 1$ allows to treat these questions in a unifying way. We will pay particular attention to the role played by the parameter p in the asymptotic convergence analysis. For $p > 1$ the dynamic writes

$$\ddot{x}(t) + \|\dot{x}(t)\|^{p-2}\dot{x}(t) + \nabla f(x(t)) = 0.$$

We will see that the case $p = 2$ separates the weak damping ($p > 2$) from the strong damping ($p < 2$), hence the importance of this case.

1.1.2. (ADIGE-H). Then, we will extend the previous results to the differential inclusion

$$\text{(ADIGE-H)} \quad \ddot{x}(t) + \partial\phi(\dot{x}(t)) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) \ni 0,$$

which, besides a damping potential ϕ as above, also involves a geometric damping driven by the Hessian of f , hence the terminology. The inertial system

$$\text{(DIN)}_{\gamma,\beta} \quad \ddot{x}(t) + \gamma\dot{x}(t) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) = 0,$$

was introduced in [9]. In the same spirit as (HBF), the dynamic $(\text{DIN})_{\gamma,\beta}$ contains a *fixed* positive viscous friction coefficient $\gamma > 0$. The introduction of the Hessian-driven damping allows to damp the transversal oscillations that might arise with (HBF), as observed in [9] in the case of the Rosenbrock function. The need to take a geometric damping adapted to f had already been observed by Alvarez [7] who considered the inertial system

$$\ddot{x}(t) + \Gamma\dot{x}(t) + \nabla f(x(t)) = 0,$$

where $\Gamma : \mathcal{H} \rightarrow \mathcal{H}$ is a linear positive anisotropic operator (see also [45]). But still this damping operator is fixed. For a general convex function, the Hessian-driven damping in $(\text{DIN})_{\gamma,\beta}$ performs a similar operation in a closed-loop adaptive way. The terminology (DIN) stands shortly for Dynamic Inertial Newton system. It refers to the natural link between this dynamic and the continuous Newton method, see Attouch-Svaiter [35]. Recent studies have been devoted to the study of the inertial dynamic

$$\ddot{x}(t) + \frac{\alpha}{t}\dot{x}(t) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) = 0,$$

which combines asymptotic vanishing damping with Hessian-driven damping. The corresponding algorithms involve a correcting term in the Nesterov accelerated gradient method which reduces the oscillatory aspects, see Attouch-Peypouquet-Redont [33], Attouch-Chbani-Fadili-Riahi [22], Shi-Du-Jordan-Su [80].

1.1.3. (ADIGE-VH). Finally, we will consider the new dynamical system

$$\text{(ADIGE-VH)} \quad \ddot{x}(t) + \partial\phi\left(\dot{x}(t) + \beta\nabla f(x(t))\right) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) \ni 0,$$

where the damping term $\partial\phi\left(\dot{x}(t) + \beta\nabla f(x(t))\right)$ involves both the velocity vector and the gradient of the potential function f . The parameter $\beta \geq 0$ is attached to the geometric damping induced by the Hessian. As previously considered, ϕ is a damping potential function. Assuming that f is convex and ϕ is a sharp function at the origin, that is $\phi(u) \geq r\|u\|$ for some $r > 0$, we will show that, for each trajectory generated by (ADIGE-VH), the following properties are satisfied:

- i) $x(\cdot)$ converges weakly as $t \rightarrow +\infty$, and its limit belongs to $\text{argmin}_{\mathcal{H}} f$.
- ii) $\dot{x}(t)$ and $\nabla f(x(t))$ converge strongly to zero as $t \rightarrow +\infty$.

1.2. Contents. The paper is organized in accordance with the above presentation. In Section 2, we first recall some classical facts concerning the Heavy Ball with Friction method and the Su-Boyd-Candès dynamic approach to the Nesterov method. Then, we successively examine each of the cases considered above: Sections 3, 4, 5, 6, 7 are devoted to the closed-loop control of the velocity, which is the main part of our study. We show the existence and uniqueness of a global solution for

the Cauchy problem, the exponential convergence rate in the case f strongly convex, the effect of weak damping, and finally analyze the convergence under the Kurdyka–Lojasiewicz property (KL). Section 8 considers some first related algorithmic results. Section 9 is devoted to the closed-loop damping with Hessian driven damping. Section 10 is devoted to the closed-loop damping involving the velocity and the gradient. We conclude by mentioning several lines of research for the future.

2. Classical facts. Let's recall some classical facts concerning the (HBF) system, the Su-Boyd-Candès dynamic approach to the Nesterov method, and the introduction of the Hessian-driven damping into these dynamics. These results will serve as comparison tools for our study.

2.1. (HBF) dynamic system. The Heavy Ball with Friction system

$$(\text{HBF})_r \quad \ddot{x}(t) + r\dot{x}(t) + \nabla f(x(t)) = 0,$$

was introduced by B. Polyak [77, 78]. It involves a fixed viscous friction coefficient $r > 0$. Assuming that f is a convex function such that $\text{argmin}_{\mathcal{H}} f \neq \emptyset$, we know, by Alvarez's theorem [7], that each trajectory of $(\text{HBF})_r$ converges weakly, and its limit belongs to $\text{argmin}_{\mathcal{H}} f$. In addition, we have the following convergence rates, the proof of which (see [18]) is based on the decrease property of the following Lyapunov functions

$$\mathcal{E}(t) := \frac{1}{r^2}(f(x(t)) - \min_{\mathcal{H}} f) + \frac{1}{2}\|x(t) - x^* + \frac{1}{r}\dot{x}(t)\|^2,$$

where $x^* \in \text{argmin}_{\mathcal{H}} f$.

THEOREM 2.1. *Let $f : \mathcal{H} \rightarrow \mathbb{R}$ be a convex function of class \mathcal{C}^1 such that $\text{argmin} f \neq \emptyset$, and let r be a positive parameter. Let $x(\cdot) : [0, +\infty[\rightarrow \mathcal{H}$ be a solution trajectory of $(\text{HBF})_r$. Set $x(0) = x_0$ and $\dot{x}(0) = x_1$. Then, we have*

$$(i) \quad \int_0^{+\infty} (f(x(t)) - \min_{\mathcal{H}} f) dt < +\infty, \quad \int_0^{+\infty} t\|\dot{x}(t)\|^2 dt < +\infty.$$

$$(ii) \quad f(x(t)) - \min_{\mathcal{H}} f \leq \frac{C(x_0, x_1)}{t}, \quad \|\dot{x}(t)\| \leq \frac{\sqrt{2C(x_0, x_1)}}{\sqrt{t}}, \text{ where}$$

$$C(x_0, x_1) =: \frac{3}{2r} \left(f(x_0) - \min_{\mathcal{H}} f \right) + r \text{dist}(x_0, \text{argmin} f)^2 + \frac{5}{4r} \|x_1\|^2.$$

$$(iii) \quad f(x(t)) - \min_{\mathcal{H}} f = o\left(\frac{1}{t}\right) \quad \text{and} \quad \|\dot{x}(t)\| = o\left(\frac{1}{\sqrt{t}}\right) \quad \text{as } t \rightarrow +\infty.$$

Let us now consider the case of a strongly convex function. Recall that a function $f : \mathcal{H} \rightarrow \mathbb{R}$ is μ -strongly convex for some $\mu > 0$ if $f - \frac{\mu}{2}\|\cdot\|^2$ is convex. We have the following exponential convergence rate, whose proof relies on the decrease property of the following Lyapunov function

$$\mathcal{E}(t) := f(x(t)) - \min_{\mathcal{H}} f + \frac{1}{2}\|\sqrt{\mu}(x(t) - x^*) + \dot{x}(t)\|^2,$$

where x^* is the unique minimizer of f .

THEOREM 2.2. *Suppose that $f : \mathcal{H} \rightarrow \mathbb{R}$ is a function of class \mathcal{C}^1 which is μ -strongly convex for some $\mu > 0$. Let $x(\cdot) : [0, +\infty[\rightarrow \mathcal{H}$ be a solution trajectory of*

$$(2.1) \quad \ddot{x}(t) + 2\sqrt{\mu}\dot{x}(t) + \nabla f(x(t)) = 0.$$

Set $x(0) = x_0$ and $\dot{x}(0) = x_1$. Then, the following property is satisfied: for all $t \geq 0$

$$f(x(t)) - \min_{\mathcal{H}} f \leq Ce^{-\sqrt{\mu}t},$$

where $C := f(x_0) - \min_{\mathcal{H}} f + \mu \text{dist}(x_0, S)^2 + \|x_1\|^2$.

A recent account on the best tuning of the damping coefficient can be found in Aujol-Dossal-Rondepierre [37]. The above results show the important role played by the geometric properties of the data in the convergence rates. Apart from the convex case, the first convergence result for (HBF) was obtained by Haraux–Jendoubi [64] in the case where $f : \mathbb{R}^n \rightarrow \mathbb{R}$ is a real-analytic function. They have shown the central role played by Lojasiewicz’s inequality in these questions (see also [55]). Then, on the basis of Kurdyka’s work in real algebraic geometric, Lojasiewicz’s inequality was extended in [42] by Bolte–Daniilidis–Ley–Mazet to a large class of tame functions, possibly nonsmooth. This is the Kurdyka–Lojasiewicz inequality, to which we will briefly refer (KL). The convergence of first and second-order proximal-gradient dynamical systems in the context of the (KL) property was obtained by Bot–Csetnek [44] and Bot–Csetnek–László [46]. The (KL) property will be a key tool for obtaining convergence rates based on the geometric properties of the data. Note that this theory only works in the finite dimensional setting¹ (the infinite dimensional setting is a difficult topic which is the subject of current research), and only for autonomous systems. This explains why working with autonomous systems is very important, since it allows us to use the powerful theory (KL).

2.2. Su-Boyd-Candès dynamic approach to Nesterov accelerated gradient method.

The following non-autonomous system

$$(\text{AVD})_{\alpha} \quad \ddot{x}(t) + \frac{\alpha}{t}\dot{x}(t) + \nabla f(x(t)) = 0,$$

will serve as a reference to compare our results with the open-loop damping approach. It was introduced in the context of convex optimization by Su–Boyd–Candès in [81]. As a specific feature, the viscous damping coefficient $\frac{\alpha}{t}$ vanishes (tends to zero) as t goes to infinity, hence the terminology ”Asymptotic Vanishing Damping”. This contrasts with (HBF) where the viscous friction coefficient is fixed, which prevents obtaining fast convergence of the values for general convex functions. Recall the main results concerning the asymptotic behaviour of the trajectories generated by $(\text{AVD})_{\alpha}$.

- For $\alpha \geq 3$, each trajectory $x(\cdot)$ of $(\text{AVD})_{\alpha}$ satisfies the asymptotic convergence rate of the values $f(x(t)) - \inf_{\mathcal{H}} f = \mathcal{O}(1/t^2)$ as $t \rightarrow +\infty$.
- For $\alpha > 3$, each trajectory converges weakly to a minimizer of f , see [23]. It is shown in [32] and [72] that the asymptotic convergence rate of the values is actually $o(1/t^2)$.
- For $\alpha \leq 3$, the convergence rate of the objective values $\mathcal{O}(t^{-\frac{2\alpha}{3}})$, see [11] and [25].
- $\alpha = 3$ is a critical value. It corresponds to the historical case studied by Nesterov [73, 74].

The convergence of the trajectories is an open question in this case.

Based on the dynamic approach above, many recent studies have been devoted to the convergence properties of the associated algorithms, which have led to a better understanding and improvement of Nesterov’s accelerated gradient algorithm; see [11], [19], [22], [23], [25], [32], [54], [81].

These rates are optimal, that is, they can be reached, or approached arbitrarily close, as shown by the following example from [23]. Let us show that $\mathcal{O}(1/t^2)$ is the worst possible case for the rate

¹In the field of PDE’s, the Lojasiewicz–Simon theory [57] makes it possible to deal with certain classes of particular problems, such as semi-linear equations.

of convergence of the values for the $(AVD)_\alpha$ trajectories, when $\alpha \geq 3$. It is attained as a limit in the following example. Take $\mathcal{H} = \mathbb{R}$ and $f(x) = c|x|^\gamma$, where c and γ are positive parameters. We look for nonnegative solutions of $(AVD)_\alpha$ of the form $x(t) = \frac{1}{t^\theta}$, with $\theta > 0$. This means that the trajectory is not oscillating, it is a completely damped trajectory. Let us determine the values of c , γ and θ that provide such solutions. On the one hand,

$$\ddot{x}(t) + \frac{\alpha}{t}\dot{x}(t) = \theta(\theta + 1 - \alpha)\frac{1}{t^{\theta+2}}.$$

On the other hand, $\nabla f(x) = c\gamma|x|^{\gamma-2}x$, which gives $\nabla f(x(t)) = c\gamma\frac{1}{t^{\theta(\gamma-1)}}$. Thus, $x(t) = \frac{1}{t^\theta}$ is solution of $(AVD)_\alpha$ if, and only if,

- i) $\theta + 2 = \theta(\gamma - 1)$, which is equivalent to $\gamma > 2$ and $\theta = \frac{2}{\gamma-2}$; and
- ii) $c\gamma = \theta(\alpha - \theta - 1)$, which is equivalent to $\alpha > \frac{\gamma}{\gamma-2}$ and $c = \frac{2}{\gamma(\gamma-2)}(\alpha - \frac{\gamma}{\gamma-2})$.

We have $\min f = 0$ and

$$f(x(t)) = \frac{2}{\gamma(\gamma-2)}\left(\alpha - \frac{\gamma}{\gamma-2}\right)\frac{1}{t^{\frac{2\gamma}{\gamma-2}}}.$$

The speed of convergence of $f(x(t))$ to 0 depends on the parameter γ . The exponent $\frac{2\gamma}{\gamma-2}$ is greater than 2, and tends to 2 when γ tends to $+\infty$. This limiting situation is obtained by taking a function f which becomes very flat around the set of its minimizers. Therefore, without any other geometric assumptions on f , we cannot expect a convergence rate better than $\mathcal{O}(1/t^2)$. This means that it is not possible to obtain a rate $\mathcal{O}(1/t^r)$ with $r > 2$, which holds for all convex functions. Hence, when $\alpha \geq 3$, $\mathcal{O}(1/t^2)$ is sharp. This is not contradictory with the rate $o(t^{-2})$ obtained when $\alpha > 3$.

2.3. Hessian-driven damping. The inertial system

$$(DIN)_{\gamma,\beta} \quad \ddot{x}(t) + \gamma\dot{x}(t) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) = 0,$$

was introduced in [9]. In line with (HBF), it contains a *fixed* positive viscous friction coefficient γ . The introduction of the Hessian-driven damping makes it possible to neutralize the oscillations likely to occur with (HBF), a key property for numerical optimization purpose.

Another motivation for the study of $(DIN)_{\gamma,\beta}$ comes from mechanics, and the modeling of damped shocks. In [29], Attouch-Maingé-Redont consider the inertial system with Hessian-driven damping

$$(2.2) \quad \ddot{x}(t) + \gamma\dot{x}(t) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) + \nabla g(x(t)) = 0,$$

where $g : \mathcal{H} \rightarrow \mathbb{R}$ is a smooth real-valued function. An interesting property of this system is that, after the introduction of an auxiliary variable y , it can be equivalently written as a first-order system involving only the time derivatives $\dot{x}(t)$, $\dot{y}(t)$ and the gradient terms $\nabla f(x(t))$, $\nabla g(x(t))$. More precisely, the system (2.2) is equivalent to the following first-order differential equation

$$(2.3) \quad \begin{cases} \dot{x}(t) + \beta\nabla f(x(t)) + ax(t) + by(t) = 0, \\ \dot{y}(t) - \beta\nabla g(x(t)) + ax(t) + by(t) = 0, \end{cases}$$

where a and b are real numbers such that: $a + b = \gamma$ and $\beta b = 1$. Note that (2.3) is different from the classical Hamiltonian formulation, which would still involve the Hessian of f . In contrast, the formulation (2.3) uses only first-order information from the function f (no occurrence of the Hessian of f). Replacing ∇f by ∂f in (2.3) allows us to extend the analysis to the case of a convex lower

semicontinuous function $f : \mathcal{H} \rightarrow \mathbb{R} \cup \{+\infty\}$, and so to introduce constraints in the model. When $f = \delta_K$ is the indicator function of a closed convex set $K \subset \mathcal{H}$, the subdifferential operator ∂f takes account of the contact forces, while ∇g takes account of the driving forces. In this setting, by playing with the geometrical damping parameter β , one can describe nonelastic shock laws with restitution coefficient (for more details we refer to [29] and references therein). Combination of dry friction ($\phi(u) = r\|u\|$) with Hessian damping has been considered by Adly–Attouch [2], [4]. To accelerate this system, several studies considered the case where the viscous damping is vanishing. As a model example, which is based on the Su–Boyd–Candès continuous model for the Nesterov accelerated gradient method, we have

$$\ddot{x}(t) + \frac{\alpha}{t}\dot{x}(t) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) = 0.$$

Considering this system, let us quote Attouch–Peypouquet–Redont [33], Attouch–Chbani–Fadili–Riahi [22], Bot–Csetnek–László [47], Castera–Bolte–Févotte–Pauwels [53], Kim [69], Lin–Jordan [71], Shi–Du–Jordan–Su [80]. At first glance, the presence of the Hessian may seem to cause numerical difficulties. However, this is not the case because the Hessian intervenes in the above ODE in the form $\nabla^2 f(x(t))\dot{x}(t)$, which is nothing other than the derivative with respect to time of $\nabla f(x(t))$. Thus, the temporal discretization of this dynamic provides first-order algorithms which, by comparison with the accelerated gradient method of Nesterov, contain a correction term which is equal to the difference of the gradients at two consecutive steps. While preserving the convergence properties of $(AVD)_\alpha$, the above system provides fast convergence to zero of the gradients, namely $\int_{t_0}^\infty t^2 \|\nabla f(x(t))\|^2 dt < +\infty$ for $\alpha \geq 3$ and $\beta > 0$, and reduces the oscillatory aspects.

2.4. Inertial dynamics with dry friction. Although dry friction (also called Coulomb friction) plays a fundamental role in mechanics, its use in optimization has only recently been developed. Due to the nonsmooth character of the associated damping function $\phi(u) = r\|u\|$, the dynamics is a differential inclusion, which, when the speed is not equal to zero, is given by

$$\ddot{x}(t) + r \frac{\dot{x}(t)}{\|\dot{x}(t)\|} + \nabla f(x(t)) = 0.$$

In this case, the energy estimate gives $\int_0^{+\infty} \|\dot{x}(t)\| dt < +\infty$. Therefore, the trajectory has finite length, and it converges strongly. The limit x_∞ of the trajectory $x(\cdot)$ satisfies

$$\|\nabla f(x_\infty)\| \leq r.$$

Thus, x_∞ is an “approximate” critical point of f . In practice, for optimization purpose, we choose a small $r > 0$. This amounts to solving the optimization problem $\min_{\mathcal{H}} f$ with the variational principle of Ekeland, instead of the Fermat rule. The importance of this case in optimization comes from the finite time stabilization property of the trajectories, which is satisfied generically with respect to the initial data. The rigorous mathematical treatment of this case has been considered by Adly–Attouch–Cabot [5], see Adly–Attouch [2, 3, 4] for recent developments. Corresponding PDE’s results have been obtained by Amann–Diaz [10] for the nonlinear wave equation, and by Carles–Gallo [52] for the nonlinear Schrödinger equation.

2.5. Closed-loop versus open-loop damping. In the strongly convex case, the autonomous system (HBF) provides an exponential rate of convergence. On the other hand, the $(AVD)_\alpha$ system provides a convergence rate of order $1/t^\alpha$. Thus, in this case, the closed-loop damping behaves

better than the open-loop damping. For general convex functions (*i.e.* in the worst case), we have the opposite situation. $(AVD)_\alpha$ provides a convergence rate $1/t^2$, while (HBF) gives only $1/t$. In this paper, we will study the impact of the closed-loop damping on the rate of convergence. A related question is: using closed-loop damping, can we obtain for general convex functions, a convergence rate of order $1/t^2$, *i.e.* as good as the Nesterov accelerated gradient method? As we will see, to answer these questions, we will have to study different types of closed-loop damping, and rely on the geometric properties of the data. These questions fall within the framework of an active research current, to quote some recent works, Apidopoulos–Aujol–Dossal–Rondepierre [12] (geometrical properties of the data), Iutzeler–Hendricx [68] (online acceleration), Lin–Jordan [71] (control perspective on high-order optimization), Poon–Liang [79] (geometry of first-order methods and adaptive acceleration).

3. Damping via closed-loop velocity control, existence and uniqueness. In this section, we will successively introduce the notion of damping potential, then prove the existence and uniqueness of the solution of the corresponding Cauchy problem.

3.1. Damping potential. We consider the differential inclusion

$$(ADIGE-V) \quad 0 \in \ddot{x}(t) + \partial\phi(\dot{x}(t)) + \nabla f(x(t)),$$

where ϕ is a convex damping potential, which is defined below.

DEFINITION 3.1. *A function $\phi : \mathcal{H} \rightarrow \mathbb{R}_+$ is a damping potential if it satisfies the following properties:*

- (i) ϕ is a nonnegative convex continuous function;
- (ii) $\phi(0) = 0 = \min_{\mathcal{H}} \phi$;
- (iii) the minimal section of $\partial\phi$ is bounded on the bounded sets, that is, for any $R > 0$

$$\sup_{\|u\| \leq R} \|(\partial\phi)^0(u)\| < +\infty.$$

In the above, $(\partial\phi)^0(u)$ is the element of minimal norm of the closed convex non empty set $\partial\phi(u)$, see [49, Proposition 2.6]. Note that, when \mathcal{H} is finite dimensional, then property (iii) is automatically satisfied. Indeed, in this case, $\partial\phi$ is bounded on the bounded sets, see [40, Proposition 16.17].

The concept of damping potential is flexible, and allows to cover various situations. For example,

$$\phi_1(u) = \frac{\gamma}{2}\|u\|^2 + r\|u\|, \quad \phi_2(u) = \max\left\{\frac{\gamma}{2}\|u\|^2; r\|u\|\right\}$$

are damping potentials which combine dry friction with viscous damping, see [4].

3.2. Existence and uniqueness results. In this section, we study the existence and the uniqueness of the solution of the Cauchy problem associated with (ADIGE-V), where ϕ is a convex damping potential. No convexity assumption is made on the function f , which is supposed to be differentiable. Let us precise the notion of strong solution.

DEFINITION 3.2. *The trajectory $x : [0, +\infty[\rightarrow \mathcal{H}$ is said to be a strong global solution of (ADIGE-V) if it satisfies the following properties:*

- (i) $x \in \mathcal{C}^1([0, +\infty[; \mathcal{H})$,
- (ii) $\dot{x} \in \text{Lip}(0, T; \mathcal{H})$, $\ddot{x} \in L^\infty(0, T; \mathcal{H})$ for all $T > 0$,
- (iii) For almost all $t > 0$, $0 \in \ddot{x}(t) + \partial\phi(\dot{x}(t)) + \nabla f(x(t))$.

Note that, since $\dot{x} \in \text{Lip}(0, T; \mathcal{H})$, it is absolutely continuous on the bounded time intervals, its distributional derivative coincides with its derivative almost everywhere (which exists). Thus, the acceleration \ddot{x} belongs to $L^\infty(0, T; \mathcal{H})$ for all $T > 0$, but it is not necessarily continuous, see [49, Appendix] for further details on vector-valued Lebesgue and Sobolev spaces.

Let's prove the following existence and uniqueness result for the associated Cauchy problem.

THEOREM 3.3. *Let $f : \mathcal{H} \rightarrow \mathbb{R}$ be a differentiable function whose gradient is Lipschitz continuous on the bounded subsets of \mathcal{H} , and such that $\inf_{\mathcal{H}} f > -\infty$. Let $\phi : \mathcal{H} \rightarrow \mathbb{R}_+$ be a damping potential (see Definition 3.1). Then, for any $x_0, x_1 \in \mathcal{H}$, there exists a unique strong global solution $x : [0, +\infty[\rightarrow \mathcal{H}$ of (ADIGE-V) such that $x(0) = x_0$ and $\dot{x}(0) = x_1$, that is*

$$\begin{cases} 0 \in \ddot{x}(t) + \partial\phi(\dot{x}(t)) + \nabla f(x(t)) \\ x(0) = x_0, \dot{x}(0) = x_1. \end{cases}$$

Proof. We successively consider the case where ∇f is Lipschitz continuous over the whole space, then the case where it is supposed to be Lipschitz continuous only on the bounded sets. In both cases, the idea is to mix the existence results for ODEs which are respectively based on the Cauchy–Lipschitz theorem, and on the theory of maximally monotone operators. We treat the two cases independently because the proof is much simpler in the first case.

Case a) ∇f is Lipschitz continuous on the whole space. The Hamiltonian formulation of (ADIGE-V) gives the equivalent first-order differential inclusion in the product space $\mathcal{H} \times \mathcal{H}$:

$$(3.1) \quad 0 \in \dot{z}(t) + \partial\Phi(z(t)) + F(z(t)),$$

where $z(t) = (x(t), \dot{x}(t)) \in \mathcal{H} \times \mathcal{H}$, and

- $\Phi : \mathcal{H} \times \mathcal{H} \rightarrow \mathbb{R}$ is the convex function defined by $\Phi(x, u) = \phi(u)$
- $F : \mathcal{H} \times \mathcal{H} \rightarrow \mathcal{H} \times \mathcal{H}$ is defined by $F(x, u) = (-u, \nabla f(x))$.

Since ∇f is Lipschitz continuous on the whole space \mathcal{H} , we immediately get that F is a Lipschitz continuous mapping on $\mathcal{H} \times \mathcal{H}$. So, we can apply a result related to evolution equations governed by Lipschitz perturbations of convex subdifferentials [49, Proposition 3.12] in order to conclude that (3.1) has a unique strong global solution with initial data $z(0) = (x_0, x_1)$.

Case b) ∇f is Lipschitz continuous on the bounded sets. The major difficulty in (ADIGE-V) is the presence of the term $\partial\phi(\dot{x}(t))$, which involves a possibly nonsmooth operator $\partial\phi$. A natural idea is to regularize this operator, and thus obtain a classical evolution equation. To this end, we use the Moreau–Yosida regularization. Let us recall some basic facts concerning this regularization procedure. For any $\lambda > 0$, the Moreau envelope of ϕ of index λ is the function $\phi_\lambda : \mathcal{H} \rightarrow \mathbb{R}$ defined by: for all $u \in \mathcal{H}$,

$$\phi_\lambda(u) = \min_{\xi \in \mathcal{H}} \left\{ \phi(\xi) + \frac{1}{2\lambda} \|u - \xi\|^2 \right\}.$$

The function ϕ_λ is convex, of class $\mathcal{C}^{1,1}$, and satisfies $\inf_{\mathcal{H}} \phi_\lambda = \inf_{\mathcal{H}} \phi$, $\text{argmin}_{\mathcal{H}} \phi_\lambda = \text{argmin}_{\mathcal{H}} \phi$. One can consult [16, section 17.2.1], [36], [40], [49] for an in-depth study of the properties of the Moreau envelope in a Hilbert framework. In our context, since $\phi : \mathcal{H} \rightarrow \mathbb{R}$ is a damping potential, we can easily verify that ϕ_λ is still a damping potential. In particular $\phi_\lambda(0) = \inf_{\mathcal{H}} \phi_\lambda = 0$. According to the subdifferential inequality for convex functions, this implies that, for all $u \in \mathcal{H}$

$$(3.2) \quad \langle \nabla \phi_\lambda(u), u \rangle \geq 0.$$

This inequality will be useful later. We will also use the following inequality, see [49, Proposition 2.6]: for any $\lambda > 0$, for any $u \in \mathcal{H}$

$$(3.3) \quad \|\nabla\phi_\lambda(u)\| \leq \|(\partial\phi)^0(u)\|.$$

So, for each $\lambda > 0$, we consider the approximate evolution equation

$$(3.4) \quad \ddot{x}_\lambda(t) + \nabla\phi_\lambda(\dot{x}_\lambda(t)) + \nabla f(x_\lambda(t)) = 0, \quad t \in [0, +\infty[.$$

We will first prove the existence and uniqueness of a global classical solution x_λ of (3.4) satisfying the Cauchy data $x_\lambda(0) = x_0$ and $\dot{x}_\lambda(0) = x_1$. Then, we will prove that the filtered sequence (x_λ) converges uniformly as $\lambda \rightarrow 0$ over the bounded time intervals towards a solution of (ADIGE-V).

According to the Hamiltonian formulation of (3.4), it is equivalent to consider the first-order (in time) system

$$\begin{cases} \dot{x}_\lambda(t) - u_\lambda(t) = 0; \\ \dot{u}_\lambda(t) + \nabla\phi_\lambda(u_\lambda(t)) + \nabla f(x_\lambda(t)) = 0, \end{cases}$$

with the Cauchy data $x_\lambda(0) = x_0$, $u_\lambda(0) = x_1$. Set

$$Z_\lambda(t) = (x_\lambda(t), u_\lambda(t)) \in \mathcal{H} \times \mathcal{H}.$$

The above system can be written equivalently as

$$\dot{Z}_\lambda(t) + F_\lambda(Z_\lambda(t)) = 0, \quad Z_\lambda(0) = (x_0, x_1),$$

where $F_\lambda : \mathcal{H} \times \mathcal{H} \rightarrow \mathcal{H} \times \mathcal{H}$, $(x, u) \mapsto F_\lambda(x, u)$ is defined by

$$F_\lambda(x, u) = \left(0, \nabla\phi_\lambda(u)\right) + \left(-u, \nabla f(x)\right).$$

Hence F_λ splits as follows

$$F_\lambda(x, u) = \nabla\Phi_\lambda(x, u) + G(x, u)$$

where

$$\Phi(x, u) = \phi(u), \quad \Phi_\lambda(x, u) = \phi_\lambda(u), \quad G(x, u) = \left(-u, \nabla f(x)\right).$$

Therefore, it is equivalent to consider the first-order differential inclusion with Cauchy data

$$(3.5) \quad \dot{Z}_\lambda(t) + \nabla\Phi_\lambda(Z_\lambda(t)) + G(Z_\lambda(t)) = 0, \quad Z_\lambda(0) = (x_0, x_1).$$

According to the Lipschitz continuity of $\nabla\Phi_\lambda$, and the fact that G is Lipschitz continuous on the bounded sets, we have that the sum operator $\nabla\Phi_\lambda + G$ which governs (3.5) is Lipschitz continuous on the bounded sets. As a consequence, the existence of a local solution to (3.5) follows from the classical Cauchy–Lipschitz theorem. To pass from a local solution to a global solution, we use a standard energy argument, and the following a priori estimate on the solutions of (3.4). After taking the scalar product of (3.4) with \dot{x}_λ , and using (3.2), we get that the global energy

$$(3.6) \quad \mathcal{E}_\lambda(t) := f(x_\lambda(t)) - \inf_{\mathcal{H}} f + \frac{1}{2}\|\dot{x}_\lambda(t)\|^2,$$

is a decreasing function of t . According to the Cauchy data, and f minorized, this implies that, on any bounded time interval, the filtered sequences of functions (x_λ) and (\dot{x}_λ) are bounded. According to the property (3.3) of the Yosida approximation, and the property (iii) of the damping potential ϕ , this implies that

$$\|\nabla\phi_\lambda(x_\lambda(t))\| \leq \|(\partial\phi)^0(x_\lambda(t))\|$$

is also bounded uniformly with respect to $\lambda > 0$ and t bounded. According to the constitutive equation (3.4), this in turn implies that the filtered sequence (\dot{x}_λ) is also bounded. This implies that if a maximal solution is defined on a finite time interval $[0, T[$, then the limits of $x_\lambda(t)$ and $\dot{x}_\lambda(t)$ exist, as $t \rightarrow T$. Then, we can apply the local existence result, which gives a solution defined on a larger interval, thus contradicting the maximality of T .

To prove the uniform convergence of the filtered sequence (Z_λ) on the bounded time intervals, we proceed in a similar way as in the proof of Brézis [49, Theorem 3.1], see also Adly-Attouch [2] in the context of damped inertial dynamics. Take $T > 0$, and $\lambda, \mu > 0$. Consider the corresponding solutions of (3.5) on $[0, T]$

$$\begin{aligned} \dot{Z}_\lambda(t) + \nabla\Phi_\lambda(Z_\lambda(t)) + G(Z_\lambda(t)) &= 0, & Z_\lambda(0) &= (x_0, x_1) \\ \dot{Z}_\mu(t) + \nabla\Phi_\mu(Z_\mu(t)) + G(Z_\mu(t)) &= 0, & Z_\mu(0) &= (x_0, x_1). \end{aligned}$$

Let's make the difference between the two equations, and take the scalar product with $Z_\lambda(t) - Z_\mu(t)$. We obtain

$$(3.7) \quad \begin{aligned} \frac{1}{2} \frac{d}{dt} \|Z_\lambda(t) - Z_\mu(t)\|^2 + \langle \nabla\Phi_\lambda(Z_\lambda(t)) - \nabla\Phi_\mu(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle \\ + \langle G(Z_\lambda(t)) - G(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle = 0. \end{aligned}$$

We now use the following ingredients:

a) According to the general properties of the Yosida approximation (see [49, Theorem 3.1]), we have

$$\langle \nabla\Phi_\lambda(Z_\lambda(t)) - \nabla\Phi_\mu(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle \geq -\frac{\lambda}{4} \|\nabla\Phi_\mu(Z_\mu(t))\|^2 - \frac{\mu}{4} \|\nabla\Phi_\lambda(Z_\lambda(t))\|^2.$$

Since the filtered sequences (x_λ) and (\dot{x}_λ) are uniformly bounded on $[0, T]$, there exists a constant C_T such that, for all $0 \leq t \leq T$

$$\|Z_\lambda(t)\| \leq C_T.$$

According to (3.3), and the fact that ϕ is a damping potential (property (iii) of Definition (3.1)), we deduce that

$$\|\nabla\Phi_\lambda(Z_\lambda(t))\| \leq \sup_{\|\xi\| \leq C_T} \|(\partial\phi)^0(\xi)\| = M_T < +\infty.$$

Therefore

$$\langle \nabla\Phi_\lambda(Z_\lambda(t)) - \nabla\Phi_\mu(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle \geq -\frac{1}{4} M_T (\lambda + \mu).$$

b) According to the local Lipschitz assumption on ∇f , the mapping $G : \mathcal{H} \times \mathcal{H} \rightarrow \mathcal{H} \times \mathcal{H}$ is Lipschitz continuous on the bounded sets. Using again that the sequence (Z_λ) is uniformly bounded on $[0, T]$, we deduce that there exists a constant L_T such that, for all $t \in [0, T]$, for all $\lambda, \mu > 0$

$$\|G(Z_\lambda(t)) - G(Z_\mu(t))\| \leq L_T \|Z_\lambda(t) - Z_\mu(t)\|.$$

Combining the above results, and using Cauchy–Schwarz inequality, we deduce from (3.7) that

$$\frac{1}{2} \frac{d}{dt} \|Z_\lambda(t) - Z_\mu(t)\|^2 \leq \frac{1}{4} M_T(\lambda + \mu) + L_T \|Z_\lambda(t) - Z_\mu(t)\|^2.$$

We now proceed with the integration of this differential inequality. Using that $Z_\lambda(0) - Z_\mu(0) = 0$, elementary calculus gives

$$\|Z_\lambda(t) - Z_\mu(t)\|^2 \leq \frac{M_T}{4L_T}(\lambda + \mu)(e^{2L_T t} - 1).$$

Therefore, the filtered sequence (Z_λ) is a Cauchy sequence for the uniform convergence on $[0, T]$, and hence it converges uniformly. This means the uniform convergence of x_λ and \dot{x}_λ to x and \dot{x} respectively.

To go to the limit on (3.4), let us write it as follows

$$(3.8) \quad \nabla \phi_\lambda(\dot{x}_\lambda(t)) = \xi_\lambda(t)$$

where

$$\xi_\lambda(t) := -\ddot{x}_\lambda(t) - \nabla f(x_\lambda(t)).$$

We now rely on the variational convergence properties of the Yosida approximation. Since (ϕ_λ) converges increasingly to ϕ as $\lambda \downarrow 0$, the sequence of integral functionals

$$\Psi^\lambda(\xi) := \int_0^T \phi_\lambda(\xi(t)) dt$$

converges increasingly to

$$\Psi(\xi) = \int_0^T \phi(\xi(t)) dt.$$

Therefore (Ψ^λ) Mosco-converges to Ψ in $\mathbb{L}^2(0, T; \mathcal{H})$. According to the theorem which makes the link between the Mosco convergence of a sequence of convex lower semicontinuous functions and the graph convergence of their subdifferentials, see Attouch [13, Theorem 3.66]), we have that

$$\partial \Psi^\lambda \rightarrow \partial \Psi$$

with respect to the topology strong $-\mathbb{L}^2(0, T; \mathcal{H}) \times$ weak $-\mathbb{L}^2(0, T; \mathcal{H})$. According to (3.8) we have

$$\xi_\lambda = \nabla \Psi^\lambda(\dot{x}_\lambda).$$

Since $\dot{x}_\lambda \rightarrow \dot{x}$ strongly in $\mathbb{L}^2(0, T; \mathcal{H})$ and ξ_λ converges weakly in $\mathbb{L}^2(0, T; \mathcal{H})$ to ξ given by

$$(3.9) \quad \xi(t) = -\ddot{x}(t) - \nabla f(x(t)),$$

we deduce that $\xi \in \partial \Psi(\dot{x})$, that is

$$\xi(t) \in \partial \phi(\dot{x}(t)).$$

According to the formulation (3.9) of ξ , we finally obtain that x is a solution of (ADIGE-V).

The uniqueness of the solution of the Cauchy problem is obtained exactly in the same way as in the case of the global Lipschitz assumption. \square

Remark 3.4. The above existence and uniqueness result uses as essential ingredient that the potential function f to be minimized is a differentiable function, whose gradient is locally Lipschitz continuous. The introduction of constraints into f via the indicator function would lead to solutions involving shocks when reaching the boundary of the constraint. In this case, existence can be still obtained in finite dimension, but uniqueness may not be satisfied, see Attouch–Cabot–Redont [21].

4. Closed-loop velocity control, preliminary convergence results . Let $x : [0, +\infty[\rightarrow \mathcal{H}$ be a solution trajectory of (ADIGE-V).

4.1. Energy estimates. Define the global energy at time t as follows:

$$(4.1) \quad \mathcal{E}(t) := f(x(t)) - \inf_{\mathcal{H}} f + \frac{1}{2} \|\dot{x}(t)\|^2.$$

Take the scalar product of (ADIGE-V) with $\dot{x}(t)$. According to the derivation chain rule, we get

$$(4.2) \quad \frac{d}{dt} \mathcal{E}(t) + \langle \partial\phi(\dot{x}(t)), \dot{x}(t) \rangle = 0.$$

The convex subdifferential inequality, and $\phi(0) = 0$, gives, for all $u \in \mathcal{H}$

$$(4.3) \quad \langle \partial\phi(u), u \rangle \geq \phi(u).$$

Combining the two above inequalities, we get

$$(4.4) \quad \frac{d}{dt} \mathcal{E}(t) + \phi(\dot{x}(t)) \leq 0.$$

Since ϕ is nonnegative, this implies that the global energy is non-increasing. Since f is minorized, this implies that the velocity $\dot{x}(t)$ is bounded over $[0, +\infty[$. Precisely,

$$(4.5) \quad \sup_{t \geq 0} \|\dot{x}(t)\| \leq R_1 := \sqrt{2\mathcal{E}(0)}.$$

To go further, suppose that the trajectory $x(\cdot)$ is bounded (this is verified for example if f is coercive), and set

$$(4.6) \quad \sup_{t \geq 0} \|x(t)\| \leq R_2.$$

Let us now establish a bound on the acceleration. For this, we rely on the approximate dynamics

$$(4.7) \quad \ddot{x}_\lambda(t) + \nabla\phi_\lambda(\dot{x}_\lambda(t)) + \nabla f(x_\lambda(t)) = 0, \quad t \in [0, +\infty[.$$

A similar estimate as above gives $\sup_{t \geq 0} \|\dot{x}_\lambda(t)\| \leq R_1 := \sqrt{2\mathcal{E}(0)}$. According to property (iii) of the damping potential, we obtain

$$\|\nabla\phi_\lambda(\dot{x}_\lambda(t))\| \leq \sup_{\|u\| \leq R_1} \|(\partial\phi)^0(u)\| = M_1 < +\infty.$$

According to the local Lipschitz continuity property of ∇f

$$\|\nabla f(x_\lambda(t))\| \leq \sup_{\|x\| \leq R_2} \|\nabla f(x)\| = M_2 < +\infty.$$

Combining the two above inequalities with (4.7), we get that for all $\lambda > 0$, and all $t \geq 0$

$$\|\ddot{x}_\lambda(t)\| \leq M_1 + M_2.$$

Since $\ddot{x}_\lambda(t)$ converges weakly to $\ddot{x}(t)$ as $\lambda \rightarrow 0$ (see the proof of Theorem 3.3), we obtain

$$(4.8) \quad \sup_{t \geq 0} \|\ddot{x}(t)\| < +\infty.$$

Moreover, by integrating (4.4), we immediately obtain $\int_0^{+\infty} \phi(\dot{x}(t)) dt < +\infty$. Let us summarize the above results, and complete them, in the following proposition.

PROPOSITION 4.1. *Let $x : [0, +\infty[\rightarrow \mathcal{H}$ be a solution trajectory of (ADIGE-V). Then, the global energy $\mathcal{E}(t) = f(x(t)) - \inf_{\mathcal{H}} f + \frac{1}{2} \|\dot{x}(t)\|^2$ is non-increasing, and*

$$\sup_{t \geq 0} \|\dot{x}(t)\| < +\infty, \quad \int_0^{+\infty} \phi(\dot{x}(t)) dt < +\infty.$$

Suppose moreover that x is bounded. Then

$$(4.9) \quad \sup_{t \geq 0} \|\ddot{x}(t)\| < +\infty.$$

Suppose moreover that there exists $p \geq 1$, and $r > 0$ such that, for all $u \in \mathcal{H}$, $\phi(u) \geq r\|u\|^p$. Then

$$(4.10) \quad \lim_{t \rightarrow +\infty} \|\dot{x}(t)\| = 0.$$

Proof. We just need to prove the last point. From $\int_0^{+\infty} \phi(\dot{x}(t)) dt < +\infty$ and $\phi(u) \geq r\|u\|^p$, we get $\int_0^{+\infty} \|\dot{x}(t)\|^p dt < +\infty$. This estimate, combined with $\sup_{t \geq 0} \|\ddot{x}(t)\| < +\infty$ classically implies that $\lim_{t \rightarrow +\infty} \|\dot{x}(t)\| = 0$. \square

Let us complete the above result by examining the convergence of the acceleration towards zero. To get this result, we need additional assumptions on the data f and ϕ .

PROPOSITION 4.2. *Let $x : [0, +\infty[\rightarrow \mathcal{H}$ be a bounded solution trajectory of (ADIGE-V). Suppose that f is a C^2 function, and that ϕ is a C^2 function which satisfies:*

(i) (local strong convexity) there exists positive constants $\gamma > 0$, and $\rho > 0$ such that for all u in \mathcal{H} with $\|u\| \leq \rho$ the following inequality holds

$$\langle \nabla^2 \phi(u) \xi, \xi \rangle \geq \gamma \|\xi\|^2 \quad \text{for all } \xi \in \mathcal{H};$$

(ii) (global growth) there exist $p \geq 1$ and $r > 0$ such that $\phi(u) \geq r\|u\|^p$ for all $u \in \mathcal{H}$. Then

$$(4.11) \quad \lim_{t \rightarrow +\infty} \|\ddot{x}(t)\| = 0.$$

Proof. Let us derivate (ADIGE-V), and set $w(t) := \ddot{x}(t)$. We obtain

$$\dot{w}(t) + \nabla^2 \phi(\dot{x}(t)) w(t) = -\nabla^2 f(x(t)) \dot{x}(t).$$

Take the scalar product of the above equation with $w(t)$. We get

$$\frac{1}{2} \frac{d}{dt} \|w(t)\|^2 + \langle \nabla^2 \phi(\dot{x}(t)) w(t), w(t) \rangle = -\langle \nabla^2 f(x(t)) \dot{x}(t), w(t) \rangle.$$

According to Proposition 4.1, we have $\lim_{t \rightarrow +\infty} \|\dot{x}(t)\| = 0$. From the local strong convexity assumption (i), and the Cauchy-Schwarz inequality, we deduce that for t sufficiently large, say $t \geq t_1$, we have

$$\frac{1}{2} \frac{d}{dt} \|w(t)\|^2 + \gamma \|w(t)\|^2 \leq \|\nabla^2 f(x(t)) \dot{x}(t)\| \|w(t)\|.$$

Since $x(\cdot)$ is bounded and ∇f is Lipschitz continuous on the bounded sets, we obtain, that for some $C > 0$

$$\frac{1}{2} \frac{d}{dt} \|w(t)\|^2 + \gamma \|w(t)\|^2 \leq C \|\dot{x}(t)\| \|w(t)\| \text{ for all } t \geq t_1.$$

After multiplication by $e^{2\gamma t}$, and integration from t_1 to t , we get

$$\frac{1}{2} (e^{\gamma t} \|w(t)\|)^2 \leq \frac{1}{2} (e^{\gamma t_1} \|w(t_1)\|)^2 + C \int_{t_1}^t e^{\gamma \tau} \|\dot{x}(\tau)\| (e^{\gamma \tau} \|w(\tau)\|) d\tau.$$

According to the Gronwall Lemma (see [49, Lemma A.5]) we obtain

$$e^{\gamma t} \|w(t)\| \leq e^{\gamma t_1} \|w(t_1)\| + C \int_{t_1}^t e^{\gamma \tau} \|\dot{x}(\tau)\| d\tau.$$

Therefore

$$\|\ddot{x}(t)\| \leq \|\ddot{x}(t_1)\| e^{-\gamma(t-t_1)} + C e^{-\gamma t} \int_{t_1}^t e^{\gamma \tau} \|\dot{x}(\tau)\| d\tau.$$

Since $\lim_{t \rightarrow +\infty} \|\dot{x}(t)\| = 0$, we have $\lim_{t \rightarrow +\infty} e^{-\gamma t} \int_{t_1}^t e^{\gamma \tau} \|\dot{x}(\tau)\| d\tau = 0$. Therefore, by passing to the limit on the above inequality we get $\lim_{t \rightarrow +\infty} \|\ddot{x}(t)\| = 0$. \square

COROLLARY 4.3. *Let us make the assumption of Proposition 4.2 and suppose that the trajectory $x(\cdot)$ is relatively compact. Then for any sequence $x(t_n) \rightarrow x_\infty$ with $t_n \rightarrow +\infty$ we have $\nabla f(x_\infty) = 0$. Set $S = \{x \in \mathcal{H} : \nabla f(x) = 0\}$. Therefore,*

$$\lim_{t \rightarrow +\infty} d(x(t), S) = 0.$$

Remark 4.4. a) Without geometric assumption on the function f , the trajectories of (ADIGE-v) may fail to converge. In [26] Attouch–Goudou–Redont exhibit a function $f : \mathbb{R}^2 \rightarrow \mathbb{R}$ which is C^1 , coercive, whose gradient is Lipschitz continuous on the bounded sets, and such that the (HBF) system admits an orbit $t \mapsto x(t)$ which does not converge as $t \rightarrow +\infty$. The above result expresses that in such situation, the attractor is the set $S = \{x \in \mathcal{H} : \nabla f(x) = 0\}$.

b) It is necessary to assume that ϕ is a smooth function in order to get the conclusion of Corollary 4.3. In fact in the case of the dry friction, that is $\phi(u) = r\|u\|$, there is convergence of the orbits to points satisfying $\|\nabla f(x_\infty)\| \leq r$, and which are not in general critical point of f .

4.2. Model example. Consider the case $\phi(u) = \frac{r}{p}\|u\|^p$ with $p > 1$, in which case the dynamic (ADIGE-V) writes

$$(4.12) \quad \ddot{x}(t) + r \|\dot{x}(t)\|^{p-2} \dot{x}(t) + \nabla f(x(t)) = 0.$$

Therefore, for $p > 2$, the viscous damping coefficient $\gamma(\cdot)$ that enters equation (4.12), and which is equal to

$$(4.13) \quad \gamma(t) := r \|\dot{x}(t)\|^{p-2}$$

tends to zero as $t \rightarrow +\infty$. So, we are in the setting of the inertial dynamics with vanishing damping coefficient. Consequently, in the associated inertial gradient algorithms, the extrapolation

coefficient tends to 1, and we can expect fast asymptotic convergence results. To summarize, in the case of (4.12), and for f coercive, we have obtained that, for $p > 2$

$$(4.14) \quad \lim_{t \rightarrow +\infty} \gamma(t) = 0, \quad \gamma(\cdot) \in L^{\frac{p}{p-2}}(0, +\infty).$$

Are these informations sufficient to derive the convergence rate of the values, and obtain similar convergence properties as for the $(AVD)_\alpha$ system?

To give a first answer to this question, we rely on the results of Cabot-Engler-Gaddat [51], Attouch-Cabot [18] and Attouch-Chbani-Riahi [24] which concern the asymptotic stabilization of inertial gradient dynamics with general time-dependent viscosity coefficient $\gamma(t)$. In the case of a vanishing damping coefficient, the key property which insures the asymptotic minimization property is that

$$\int_0^{+\infty} \gamma(t) dt = +\infty.$$

This means that the coefficient $\gamma(t)$ can go to zero as $t \rightarrow +\infty$, but not too fast in order to dissipate the energy enough. On the positive side, $\gamma(t) = \frac{\alpha}{t}$ does satisfy the conditions (4.14) for any $p > 0$, which does not exclude the Nesterov case. On the negative side, we can easily find $\gamma(t)$ such that

$$(4.15) \quad \lim_{t \rightarrow +\infty} \gamma(t) = 0, \quad \gamma(\cdot) \in L^{\frac{p}{p-2}}(0, +\infty) \text{ and } \int_0^{+\infty} \gamma(t) dt < +\infty.$$

So, without any other hypothesis, we cannot conclude from this information alone. At this point, the idea is to introduce additional information, assuming a geometric property on the function f to minimize. In the next two sections, we successively consider the case where f is a strongly convex function, then the case of the functions f satisfying the Kurdyka-Lojasiewicz property.

5. The strongly convex case: exponential convergence rate. We will study the asymptotic behavior of the system (ADIGE-V) when f is a strongly convex function. Recall that $f : \mathcal{H} \rightarrow \mathbb{R}$ is said to be μ -strongly convex (with $\mu > 0$) if $f - \frac{\mu}{2} \|\cdot\|^2$ is convex. Then, we will consider the particular case where f is strongly convex and quadratic. Finally, we will give numerical illustrations in dimension one.

5.1. General strongly convex function f .

THEOREM 5.1. *Let $f : \mathcal{H} \rightarrow \mathbb{R}$ be a differentiable function which is μ -strongly convex for some $\mu > 0$, and whose gradient is Lipschitz continuous on the bounded sets. Let \bar{x} be the unique minimizer of f .*

Let $\phi : \mathcal{H} \rightarrow \mathbb{R}_+$ be a damping potential (see Definition 3.1) which is differentiable, and whose gradient is Lipschitz continuous on the bounded subsets of \mathcal{H} . Suppose that ϕ satisfies the following growth conditions:

- (i) (local) there exists positive constants α , and $\rho > 0$ such that for all u in \mathcal{H} with $\|u\| \leq \rho$

$$\langle \nabla \phi(u), u \rangle \geq \alpha \|u\|^2.$$

- (ii) (global) there exists $p \geq 1$, $c > 0$, such that for all u in \mathcal{H} , $\phi(u) \geq c \|u\|^p$.

Then, for any solution trajectory $x : [0, +\infty[\rightarrow \mathcal{H}$ of (ADIGE-V), we have exponential convergence rate to zero as $t \rightarrow +\infty$ for $f(x(t)) - f(\bar{x})$, $\|x(t) - \bar{x}\|$ and the velocity $\|\dot{x}(t)\|$.

Proof. We will use the following inequalities which are attached to the strong convexity of f :

$$(5.1) \quad f(\bar{x}) - f(x(t)) \geq \langle \nabla f(x(t)), \bar{x} - x(t) \rangle + \frac{\mu}{2} \|x(t) - \bar{x}\|^2$$

$$(5.2) \quad f(x(t)) - f(\bar{x}) \geq \frac{\mu}{2} \|x(t) - \bar{x}\|^2.$$

Let us consider the global energy (introduced in (4.1), in the preliminary estimates)

$$\mathcal{E}(t) := \frac{1}{2} \|\dot{x}(t)\|^2 + f(x(t)) - f(\bar{x}).$$

By Proposition 4.1, the velocity vector $\dot{x}(t)$ is bounded on \mathbb{R}_+ . Moreover, $\mathcal{E}(\cdot)$ is non-increasing, and hence bounded from above. By definition of $\mathcal{E}(t)$, this implies that $f(x(t))$ is bounded from above. Since f is strongly convex, it is coercive, which implies that $x(\cdot)$ is bounded. Since $x(\cdot)$ and $\dot{x}(\cdot)$ are bounded, and the vector fields ∇f and $\nabla \phi$ are locally Lipschitz continuous, we deduce from the constitutive equation

$$\ddot{x}(t) = -\nabla \phi(\dot{x}(t)) - \nabla f(x(t))$$

that $\ddot{x}(\cdot)$ is also bounded. According to the preliminary estimates established in Proposition 4.1, we have $\int_0^{+\infty} \phi(\dot{x}(t)) dt < +\infty$. Combining this property with the global growth assumption (ii) on ϕ , we deduce that there exists $p \geq 1$ such that

$$\int_0^{+\infty} \|\dot{x}(t)\|^p dt < +\infty.$$

Since $\ddot{x}(\cdot)$ is bounded, this implies that $\dot{x}(t) \rightarrow 0$ as $t \rightarrow +\infty$. So, for t sufficiently large, say $t \geq t_1$

$$\|\dot{x}(t)\| \leq \rho.$$

Time derivation of $\mathcal{E}(\cdot)$, together with the constitutive equation (ADIGE-V), gives for $t \geq t_1$

$$(5.3) \quad \begin{aligned} \dot{\mathcal{E}}(t) &= \langle \dot{x}(t), -\nabla \phi(\dot{x}(t)) - \nabla f(x(t)) \rangle + \langle \dot{x}(t), \nabla f(x(t)) \rangle \\ &= -\langle \dot{x}(t), \nabla \phi(\dot{x}(t)) \rangle \\ &\leq -\alpha \|\dot{x}(t)\|^2, \end{aligned}$$

where the last inequality comes from the growth condition (i) on ϕ , and $\|\dot{x}(t)\| \leq \rho$ for $t \geq t_1$.

Since $\dot{x}(\cdot)$ is bounded, let L be the Lipschitz constant of $\nabla \phi$ on a ball that contains the velocity vector $\dot{x}(t)$ for all $t \geq 0$. Since $\nabla \phi(0) = 0$ we have, for all $t \geq 0$

$$(5.4) \quad \|\nabla \phi(\dot{x}(t))\| \leq L \|\dot{x}(t)\|.$$

Using successively (ADIGE-V), (5.4) and (5.1), we obtain

$$(5.5) \quad \begin{aligned} \frac{d}{dt} \left(\langle x(t) - \bar{x}, \dot{x}(t) \rangle \right) &= \|\dot{x}(t)\|^2 + \langle x(t) - \bar{x}, -\nabla \phi(\dot{x}(t)) - \nabla f(x(t)) \rangle \\ &\leq \|\dot{x}(t)\|^2 + L \|x(t) - \bar{x}\| \|\dot{x}(t)\| - \langle x(t) - \bar{x}, \nabla f(x(t)) \rangle \\ &\leq \|\dot{x}(t)\|^2 + \frac{L^2}{2\mu} \|\dot{x}(t)\|^2 + \frac{\mu}{2} \|x(t) - \bar{x}\|^2 + \langle \bar{x} - x(t), \nabla f(x(t)) \rangle \\ &\leq \left(1 + \frac{L^2}{2\mu} \right) \|\dot{x}(t)\|^2 + f(\bar{x}) - f(x(t)). \end{aligned}$$

Take now $\epsilon > 0$ (we will specify below how it should be chosen), and define

$$h_\epsilon(t) := \mathcal{E}(t) + \epsilon \langle x(t) - \bar{x}, \dot{x}(t) \rangle.$$

Time derivation of h_ϵ , together with (5.3) and (5.5), gives for $t \geq t_1$

$$\dot{h}_\epsilon(t) \leq - \left(\alpha - \epsilon \left(1 + \frac{L^2}{2\mu} \right) \right) \|\dot{x}(t)\|^2 - \epsilon (f(x(t)) - f(\bar{x})).$$

Choose $\epsilon > 0$ such that $\alpha - \epsilon \left(1 + \frac{L^2}{2\mu} \right) > 0$, and take $C_1 := \min\{\alpha - \epsilon \left(1 + \frac{L^2}{2\mu} \right); \epsilon\}$. We deduce that

$$(5.6) \quad \dot{h}_\epsilon(t) \leq -C_1 \left(\|\dot{x}(t)\|^2 + f(x(t)) - f(\bar{x}) \right).$$

Further, from (5.2) and the Cauchy–Schwarz inequality we easily obtain

$$\begin{aligned} h_\epsilon(t) &\leq \frac{1}{2} \|\dot{x}(t)\|^2 + f(x(t)) - f(\bar{x}) + \frac{\epsilon}{2} \|x(t) - \bar{x}\|^2 + \frac{\epsilon}{2} \|\dot{x}(t)\|^2 \\ &\leq \left(\frac{1}{2} + \frac{\epsilon}{2} \right) \|\dot{x}(t)\|^2 + \left(1 + \frac{\epsilon}{\mu} \right) (f(x(t)) - f(\bar{x})) \\ &\leq \left(1 + \epsilon \left(\frac{1}{2} + \frac{1}{\mu} \right) \right) \left(\|\dot{x}(t)\|^2 + f(x(t)) - f(\bar{x}) \right). \end{aligned}$$

Combining this inequality with (5.6), we obtain

$$\dot{h}_\epsilon(t) + C_2 h_\epsilon(t) \leq 0,$$

with $C_2 := \frac{C_1}{1 + \epsilon \left(\frac{1}{2} + \frac{1}{\mu} \right)} > 0$. Then, the Gronwall inequality classically implies

$$(5.7) \quad h_\epsilon(t) \leq h_\epsilon(0) e^{-C_2 t}.$$

Finally, from (5.2) and the Cauchy–Schwarz inequality we have

$$\begin{aligned} h_\epsilon(t) &\geq \frac{1}{2} \|\dot{x}(t)\|^2 + f(x(t)) - f(\bar{x}) - \frac{\epsilon}{2} \|x(t) - \bar{x}\|^2 - \frac{\epsilon}{2} \|\dot{x}(t)\|^2 \\ &\geq \left(\frac{1}{2} - \frac{\epsilon}{2} \right) \|\dot{x}(t)\|^2 + \left(1 - \frac{\epsilon}{\mu} \right) (f(x(t)) - f(\bar{x})). \end{aligned}$$

Therefore, by taking ϵ small enough, we obtain the existence of $C_3 > 0$ such that

$$h_\epsilon(t) \geq C_3 \left(\|\dot{x}(t)\|^2 + f(x(t)) - f(\bar{x}) \right).$$

Combining this inequality with (5.7) and (5.2), we obtain an exponential convergence rate to zero for $f(x(t)) - f(\bar{x})$, $\|x(t) - \bar{x}\|$ and the velocity $\|\dot{x}(t)\|$. \square

Remark 5.2. In section 7.4, as a consequence of the convergence analysis based on the Kurdyka–Lojasiewicz theory, we will extend the above results to the case where we only assume a quadratic growth assumption

$$f(x) - \inf_{\mathcal{H}} f \geq c \operatorname{dist}(x, \operatorname{argmin} f)^2.$$

5.2. Case f convex quadratic positive definite. Let us make precise the previous results in the case $f(x) = \frac{1}{2} \langle Ax, x \rangle$, where $A : \mathcal{H} \rightarrow \mathcal{H}$ is a linear continuous positive definite self-adjoint operator. Then $\nabla f(x) = Ax$, and (ADIGE-V) is written

$$(5.8) \quad \ddot{x}(t) + \partial\phi(\dot{x}(t)) + A(x(t)) \ni 0.$$

Let us first prove that for a general damping potential ϕ , there is ergodic convergence of the trajectories.

THEOREM 5.3. *Let $x : [0, +\infty[\rightarrow \mathcal{H}$ be a solution trajectory of (5.8), where ϕ is a damping potential, and $A : \mathcal{H} \rightarrow \mathcal{H}$ is a linear continuous positive definite self adjoint operator. Then, we have the following ergodic convergence result for the weak topology: as $t \rightarrow +\infty$,*

$$\frac{1}{t} \int_0^t x(\tau) d\tau \rightharpoonup x_\infty,$$

where the limit x_∞ satisfies

$$(5.9) \quad 0 \in \partial\phi(0) + Ax_\infty.$$

When ϕ is differentiable at the origin, we have $Ax_\infty = 0$, that is $x_\infty = 0$.

When $\phi(x) = r\|x\|$, we have $\|Ax_\infty\| \leq r$.

Proof. The Hamiltonian formulation of (5.8) gives the equivalent first-order differential inclusion in the product space $\mathcal{H} \times \mathcal{H}$:

$$(5.10) \quad 0 \in \dot{z}(t) + \partial\Phi(z(t)) + F(z(t)),$$

where $z(t) = (x(t), \dot{x}(t)) \in \mathcal{H} \times \mathcal{H}$, and

- $\Phi : \mathcal{H} \times \mathcal{H} \rightarrow \mathbb{R}$ is the convex continuous function defined by $\Phi(x, u) = \phi(u)$
- $F : \mathcal{H} \times \mathcal{H} \rightarrow \mathcal{H} \times \mathcal{H}$ is defined by $F(x, u) = (-u, Ax)$.

The trick is to renorm the product space $\mathcal{H} \times \mathcal{H}$ as follows:

The mapping $(x, y) \mapsto \langle Ax, y \rangle$ defines a scalar product on \mathcal{H} , which is equivalent to the initial one. Accordingly, let us equip the product space $\mathcal{H} \times \mathcal{H}$ with the scalar product

$$\langle\langle (x_1, u_1), (x_2, u_2) \rangle\rangle := \langle Ax_1, x_2 \rangle + \langle u_1, u_2 \rangle$$

With respect to this new scalar product, let us observe that:

- $F : \mathcal{H} \times \mathcal{H} \rightarrow \mathcal{H} \times \mathcal{H}$ is a linear continuous skew symmetric operator. Since A is self-adjoint

$$\begin{aligned} \langle\langle F(x, u), (x, u) \rangle\rangle &= \langle\langle (-u, Ax), (x, u) \rangle\rangle \\ &= -\langle Au, x \rangle + \langle Ax, u \rangle = 0. \end{aligned}$$

- The subdifferential of Φ is unchanged that is $\partial\Phi(x, u) = (0, \partial\phi(u))$.

Therefore, the differential inclusion (5.10) is governed by the sum of two maximally monotone operators, one of them is the subdifferential of a convex continuous function, the other one is a monotone skew symmetric operator. By the classical Rockafellar theorem (see [40, Corollary 24.4]), their sum is still maximally monotone. Consequently, we can apply the theory concerning the semi groups generated by general maximally monotone operators, and conclude that $z(t)$ converges weakly and in an ergodic way towards a zero $z_\infty = (x_\infty, u_\infty)$ of $\partial\Phi + F$. This means

$$(0, \partial\phi(u_\infty)) + (-u_\infty, Ax_\infty) = (0, 0).$$

Equivalently $u_\infty = 0$ and $\partial\phi(0) + Ax_\infty \ni 0$. □

In the case of the wave equation, this type of argument was developed by Haraux in [63, Lecture 12, Theorem 45]. A recent account on these questions can be found in Haraux-Jendoubi [65] and Alabau Boussouira-Privat-Trélat [6].

5.3. Numerical illustrations. Finding explicit solutions in a closed form of nonlinear oscillators is an important subject, which has direct applications in various fields. In the one-dimensional case, the corresponding second-order differential equation

$$\ddot{x}(t) + d(x(t), \dot{x}(t))\dot{x}(t) + g(x(t)) = 0$$

is known as the Levinson-Smith equation. It is reduced to the Liénard equation when d depends only on x . There is a vast literature on this difficult subject. One can consult [60, 59] for recent reports on the subject and the description of some of the different techniques developed to resolve these questions. In our setting, we will provide some insight on this question by combining energetic and topological arguments.

5.3.1. A numerical one-dimensional example. Consider the case $\mathcal{H} = \mathbb{R}$, $f(x) = \frac{1}{2}|x|^2$, and $\phi(u) = \frac{1}{p}|u|^p$ with $p > 1$. Then, (ADIGE-V) writes

$$(5.11) \quad \ddot{x}(t) + |\dot{x}(t)|^{p-2}\dot{x}(t) + x(t) = 0.$$

It is a linear oscillator with nonlinear damping. According to the previous results,

- For $p = 2$, according to the strong convexity of the potential function $f(x) = \frac{1}{2}|x|^2$ and Theorem 5.1, we have convergence at an exponential rate of $x(t)$ and $\dot{x}(t)$ toward 0. Indeed, $p = 2$ is the only value of p for which the hypothesis of Theorem 5.1 are satisfied. For $p > 2$ the local hypothesis (i) is not satisfied, and for $p < 2$ the gradient of ϕ fails to be Lipschitz continuous on the bounded sets containing the origin.

- For $p > 1$, let us first show that $\lim_{t \rightarrow +\infty} \dot{x}(t) = 0$. This results from Proposition 4.1, and the fact that the trajectory is bounded. This last property results from the fact that the global energy $\mathcal{E}(t) = \frac{1}{2}|\dot{x}(t)|^2 + \frac{1}{2}|x(t)|^2$ is non-increasing, and hence convergent (and bounded from above).

Let us show that $x(t)$ tends to zero. Since $\lim_{t \rightarrow +\infty} \dot{x}(t) = 0$, and $\lim_{t \rightarrow +\infty} \mathcal{E}(t)$ exists, we have

$$(5.12) \quad \lim_{t \rightarrow +\infty} |x(t)|^2 = \lim_{t \rightarrow +\infty} \mathcal{E}(t) \text{ exists.}$$

Since the identity operator clearly satisfies the assumptions of Theorem 5.3, we have the following ergodic convergence result $\lim_{t \rightarrow +\infty} \frac{1}{t} \int_0^t x(\tau) d\tau = 0$. There are two possibilities:

a) For t sufficiently large, $x(t)$ has a fixed sign. According to (5.12) this implies that

$$\lim_{t \rightarrow +\infty} x(t) := x_\infty \text{ exists.}$$

The convergence implies the ergodic convergence. Therefore, $\lim_{t \rightarrow +\infty} \frac{1}{t} \int_0^t x(\tau) d\tau = x_\infty$. But we know that the ergodic limit is zero, hence $x_\infty = 0$.

b) The trajectory changes sign an infinite number of times as $t \rightarrow +\infty$. This means that there exists sequences s_n and t_n which tend to infinity such that $x(t_n)x(s_n) < 0$. Since the trajectory is continuous, by the mean value theorem, this implies the existence of $\tau_n \in [s_n, t_n]$ such that $x(\tau_n) = 0$. Hence $x(\tau_n)^2 = 0$ for all $n \in \mathbb{N}$, with $\tau_n \rightarrow +\infty$. Since $\lim_{t \rightarrow +\infty} |x(t)|^2$ exists, this implies that $\lim_{t \rightarrow +\infty} |x(t)|^2 = 0$. Clearly, this implies that $\lim_{t \rightarrow +\infty} x(t) = 0$.

So, for $p > 1$, for any solution trajectory of (5.11), we have:

$$(5.13) \quad \lim_{t \rightarrow +\infty} x(t) = 0 \quad \text{and} \quad \lim_{t \rightarrow +\infty} \dot{x}(t) = 0.$$

Now let's analyze how the trajectories and their speeds go to zero. In fact, the situation differs notably according to $p > 2$ or $p < 2$. As we shall see, the case $p > 2$ corresponds to a weak damping, while the case $p < 2$ corresponds to a strong damping.

Case $p > 2$. Since the speed $|\dot{x}(t)|$ tends to zero, we have $\gamma(t) := |\dot{x}(t)|^{p-2} \rightarrow 0$ as $t \rightarrow +\infty$. The viscous damping coefficient $\gamma(t)$ becomes asymptotically small. As a result, the damping effect

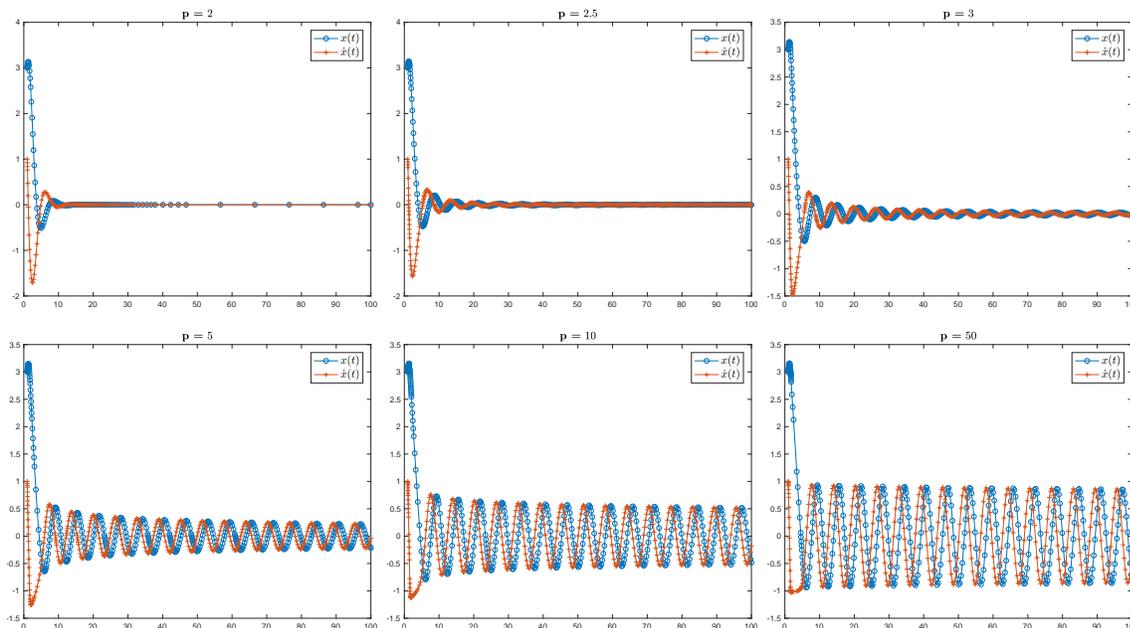


Fig. 5.1: The evolution of the trajectories $x(t)$ (blue line) and $\dot{x}(t)$ (red line) of the dynamical system (5.11) for different values of $p \geq 2$.

also becomes weak (that's what we call weak damping). As p increases, the damping effect tends to diminish, the trajectory tends to oscillate more and more, and the convergence rate deteriorates. We illustrate this in Figure 5.1, where we can see the evolution of the trajectory $x(t)$ (blue line) and its derivative $\dot{x}(t)$ (red line) of the dynamical system (5.11) with starting point $(x(0), \dot{x}(0)) = (3, 1)$ and for different values of p greater than or equal to 2. The trajectory and the velocity tend towards zero, however, the oscillations become stronger as p increases, and the convergence towards zero becomes very slow. For large p , the oscillatory aspect conforms to the ergodic convergence of the trajectory to 0 (indeed in dimension one the trajectory tends towards zero, but we can expect that in higher dimensions there is only ergodic convergence towards zero). Note also that for $p > 2$, and p close to 2, the trajectory is close to that corresponding to $p = 2$, and therefore enjoys excellent convergence properties. It would be interesting to study this situation, because it is a natural candidate to obtain convergence rates similar to the accelerated gradient method of Nesterov.

Case $1 < p < 2$. According to (5.13) we have $\lim_{t \rightarrow +\infty} \dot{x}(t) = 0$, and $\lim_{t \rightarrow +\infty} x(t) = 0$. Since $\lim_{t \rightarrow +\infty} \dot{x}(t) = 0$ and $2 - p > 0$, we have that the viscous damping coefficient

$$\gamma(t) := \frac{1}{|\dot{x}(t)|^{2-p}} \rightarrow +\infty \text{ as } t \rightarrow +\infty.$$

We are in the setting of a strong damping effect. This situation was analyzed in the following result of [17], which we reproduce here. It concerns the asymptotic behaviour of

$$(\text{IGS})_\gamma \quad \ddot{x}(t) + \gamma(t)\dot{x}(t) + \nabla f(x(t)) = 0.$$

PROPOSITION 5.4. *Let $f : \mathcal{H} \rightarrow \mathbb{R}$ be a function of class \mathcal{C}^1 such that ∇f is bounded on the bounded subsets of \mathcal{H} . Given $r > 0$ and $\theta > 1$, assume that $\gamma(t) = r t^\theta$ for every $t \geq t_0 \geq 0$. Then each bounded solution trajectory $x(\cdot)$ of $(\text{IGS})_\gamma$ satisfies $\int_{t_0}^{+\infty} \|\dot{x}(t)\| dt < +\infty$, and hence converges strongly toward some $x^* \in \mathcal{H}$.*

According to this result, we can obtain some information about the convergence rate of the velocity to zero. We have two cases: either $\int_{t_0}^{+\infty} \|\dot{x}(t)\| dt < +\infty$, or $\int_{t_0}^{+\infty} \|\dot{x}(t)\| dt = +\infty$. In this last case, according to Proposition 5.4, we cannot have $\gamma(t) = \frac{1}{|\dot{x}(t)|^{2-p}}$ of order $r t^\theta$ with $\theta > 1$. This excludes the possibility to have $|\dot{x}(t)|$ of order $\frac{1}{t^{2-p}}$ with $\theta > 1$. So, the best that we can expect is, $|\dot{x}(t)| \sim \frac{1}{t^{\frac{1}{2-p}}}$ as $t \rightarrow +\infty$. This estimate is in accordance with the exponential decay when $p = 2$, and the finite length property when $p = 1$. We emphasize the fact that the above argument is not a rigorous proof, it just gives an indication of the type of convergence rate that we can expect.

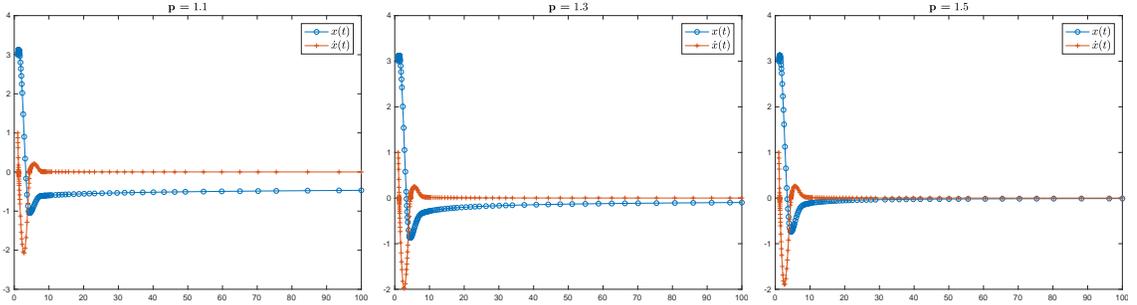


Fig. 5.2: Evolution of $x(t)$ (blue) and $\dot{x}(t)$ (red) for different values of $1 < p < 2$.

In Figure 5.2 we can see the evolution of the trajectory $t \mapsto x(t)$ (blue line) and of its derivative $t \mapsto \dot{x}(t)$ (red line) of the dynamical system (5.11) with starting point $(x(0), \dot{x}(0)) = (3, 1)$ for different values of $1 < p < 2$. Because of the strong damping, the trajectories exhibit small oscillations, and the velocity converges fastly to zero. By contrast, the convergence of $x(t)$ to zero highly depends on the parameter p . When p is close to 1, the convergence of the trajectory to zero is poor, however, already a slight increase of p concisely improves the convergence of the trajectory. Indeed, when p becomes large the convergence of the trajectory improves.

6. Weak damping: from slow convergence to attractor effect. As we already noticed, even in the case of a strongly convex function f , when the damping effect becomes too weak, then the convergence property is degraded. In the case of the damping $\|\dot{x}(t)\|^{p-2}\dot{x}(t)$, this corresponds to situations where $p > 2$. In this section, we give examples showing that in the case of a general convex function, the situation is even worse, and the trajectory may not converge in the case of weak damping. In this case, one has to replace the convergence notions by the concept of attractor, a central subject for the theory of dynamic systems, and PDE's, see Hale [61], Haraux [62] for seminal contributions to the subject in the case of gradient systems (*i.e.* systems for which there exists a Lyapunov function). For optimization purposes, this is a promising research topic, largely to be explored in the case of a general damping function. In the next section, we take a convex function f with a continuum of minimizers, and examine the lack of convergence when the damping becomes too weak. In fact, as we have already underlined, convergence depends both on the geometric properties of the damping potential and on the potential function f to be minimized. The corresponding geometric aspects concerning f will be examined a little later.

6.1. An example where convergence fails to be satisfied. The following example is based on Haraux [62, section 5.1]. Take $\mathcal{H} = \mathbb{R}$, and $f : \mathbb{R} \rightarrow \mathbb{R}$ a convex function of class \mathcal{C}^1 which achieves its minimal value on the line segment $[a, b]$, with $a < b$. We suppose that f is coercive, *i.e.* $\lim_{|x| \rightarrow +\infty} f(x) = +\infty$. Its graph is shown in the Figure 6.1 below, and looks like a bowl with a flat bottom.

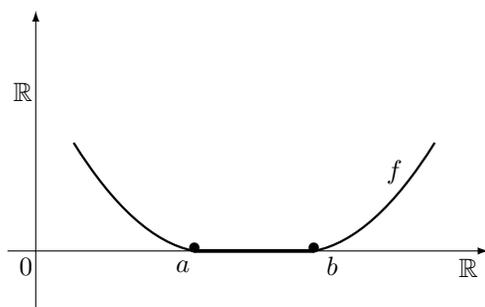


Fig. 6.1: Counterexample to convergence, $p \geq 3$.

Consider the evolution equation with closed-loop damping

$$(6.1) \quad \ddot{x}(t) + |\dot{x}(t)|^{p-2}\dot{x}(t) + \nabla f(x(t)) = 0.$$

Let us discuss, according to the value of p , the convergence properties of the trajectories of this system. We will need the following elementary lemma, see [62, Lemma 5.1.3].

LEMMA 6.1. *Let $v \in \mathcal{C}^2(\mathbb{R}_+)$ which satisfies, for some $c > 0$*

$$\dot{v}(0) > 0; \quad \ddot{v}(t) \geq -c\dot{v}(t)^2 \text{ for all } t \geq 0.$$

Then, v is an increasing function, and

$$\lim_{t \rightarrow +\infty} v(t) = +\infty.$$

Proof. As long as $\dot{v}(t) > 0$, by integration of the differential inequality $\ddot{v}(t) + c\dot{v}(t)^2 \geq 0$, we obtain

$$\dot{v}(t) \geq \frac{1}{ct + \frac{1}{\dot{v}(0)}}.$$

This immediately implies that $\dot{v}(t) > 0$ for all $t \geq 0$. By integrating the above inequality, we obtain

$$v(t) \geq v(0) + \int_0^t \frac{1}{c\tau + \frac{1}{\dot{v}(0)}} d\tau$$

which implies $\lim_{t \rightarrow +\infty} v(t) = +\infty$. \square

PROPOSITION 6.2. *Suppose that $p \geq 3$. Then, any solution trajectory of (6.1) which is not constant, passes an infinitely of times through the points a and b .*

Proof. According to Proposition 4.1, the trajectory $x(\cdot)$ is bounded and satisfies

$$(6.2) \quad \lim_{t \rightarrow +\infty} \|\dot{x}(t)\| = 0 \quad \text{and} \quad \sup_{t \geq 0} \|\ddot{x}(t)\| < +\infty.$$

Let us argue by contradiction, and assume that there exists some $t_1 > 0$ such that $x(t) \geq a$ for all $t \geq t_1$. We can distinguish two cases:

• **First case:** $\dot{x}(t) \geq 0$ for all $t \geq t_1$. Then, $t \mapsto x(t)$ is increasing and bounded, hence converges to some $x_\infty \in \mathbb{R}$. From the constitutive equation (6.1), $\lim_{t \rightarrow +\infty} \|\dot{x}(t)\| = 0$, and the continuity of ∇f , we deduce that $\lim_{t \rightarrow +\infty} \ddot{x}(t) = -\nabla f(x_\infty)$. Using again that $\lim_{t \rightarrow +\infty} \|\dot{x}(t)\| = 0$, we deduce that $\nabla f(x_\infty) = 0$. Since $x \mapsto \nabla f(x)$ is an increasing function, and $\nabla f(x(t_1)) \geq 0$, we obtain that

$$\nabla f(x(t)) = 0 \quad \text{for all } t \geq t_1.$$

Returning to the constitutive equation (6.1), we get

$$\ddot{x}(t) + |\dot{x}(t)|^{p-2} \dot{x}(t) = 0 \quad \text{for all } t \geq t_1.$$

Since $\lim_{t \rightarrow +\infty} \|\dot{x}(t)\| = 0$, and $p \geq 3$, we have for t sufficiently large, say $t \geq t_2 \geq t_1$ $|\dot{x}(t)|^{p-1} \leq |\dot{x}(t)|^2$. Therefore, for all $t \geq t_2$

$$\ddot{x}(t) + |\dot{x}(t)|^2 \geq 0.$$

Since $x(\cdot)$ is not constant, there exists some $t_3 \geq t_2$ such that $\dot{x}(t_3) > 0$. According to Lemma 6.1, we have $\lim_{t \rightarrow +\infty} x(t) = +\infty$, a clear contradiction with the convergence of $x(t)$.

• **Second case:** there exists $t_2 \geq t_1$ such that $\dot{x}(t_2) < 0$. From the constitutive equation (6.1) and $x(t) \geq a$ we get, for all $t \geq t_2$

$$\ddot{x}(t) + |\dot{x}(t)|^{p-2} \dot{x}(t) = -\nabla f(x(t)) \leq 0$$

This implies, for t large enough

$$\ddot{x}(t) \leq |\dot{x}(t)|^2.$$

Let's apply Lemma 6.1 to $-x(\cdot)$. Since $-\dot{x}(t_2) > 0$, we obtain $\lim_{t \rightarrow +\infty} x(t) = -\infty$, a contradiction. Therefore, there is an infinite number of times such that $x(t) = a$. A similar argument gives the same type of result with b , which shows that the trajectory oscillates indefinitely between a and b . \square

By contrast, if the damping effect is sufficiently important, there is convergence. In our situation, this corresponds to the case $2 < p < 3$, as shown in the following proposition.

PROPOSITION 6.3. *Suppose that $2 < p < 3$. Then, any solution trajectory of (6.1) converges, and its limit belongs to $[a, b]$.*

Proof. We sketch the main lines of the proof, whose details can be found in Haraux-Jendoubi [65, Theorem 9.2.1], which deals with a slightly more general situation. Take x a solution trajectory of (6.1), and denote by $\omega(x)$ its limit set, that is the set of all its limit points as $t \rightarrow +\infty$ (limits of sequences $x(t_n)$ for $t_n \rightarrow +\infty$). By classical argument, this set is a connected subset of $\{\nabla f = 0\}$, that is $\omega(x) \subset [a, b]$. If $\omega(x)$ is reduced to a singleton, the proof is finished. Let us therefore examine the complementary case

$$\omega(x) = [c, d] \subset [a, b], \text{ with } c < d,$$

and show that this leads to a contradiction. Set $l := \frac{1}{2}(c + d)$. Let us prove that $\lim_{t \rightarrow +\infty} x(t) = l$, which gives $\omega(x) = \{l\}$, a clear contradiction with $\omega(x) = [c, d]$, $c \neq d$. First, since l belongs to the interior of $\omega(x)$, according to the intermediate value property, there exists a sequence (t_n) with $t_n \rightarrow +\infty$ such that $x(t_n) = l$. By continuity of x , and since l belongs to the interior of $[c, d]$, for each $n \in \mathbb{N}$ there exists $\delta_n > 0$ such that

$$x(t) \in [c, d] \text{ for all } t \in [t_n, t_n + \delta_n].$$

Let's prove that for n large enough we can take $\delta_n = +\infty$. Set

$$\theta_n = \inf \{t > t_n : x(t) \notin [c, d]\},$$

and assume $\theta_n < +\infty$. So for all $t \in [t_n, \theta_n]$ we have $\nabla f(x(t)) = 0$, and (6.1) reduces to

$$\ddot{x}(t) + |\dot{x}(t)|^{p-2} \dot{x}(t) = 0.$$

After multiplying by $\dot{x}(t)$, we get for all $t \in [t_n, \theta_n]$

$$\frac{d}{dt} |\dot{x}(t)|^2 + 2|\dot{x}(t)|^p = 0.$$

After integration from t_n to $t \in [t_n, \theta_n]$, we get for all $t \in [t_n, \theta_n]$

$$|\dot{x}(t)| = \left(|\dot{x}(t_n)|^{-p+2} + (p-2)(t-t_n) \right)^{\frac{-1}{p-2}}.$$

After further integration, we get for all $t \in [t_n, \theta_n]$

$$\begin{aligned} |x(t) - l| &= |x(t) - x(t_n)| \leq \int_{t_n}^t |\dot{x}(s)| ds \\ &= \frac{1}{p-3} \left(|\dot{x}(t_n)|^{-p+2} + (p-2)(t-t_n) \right)^{\frac{p-3}{p-2}} + \frac{1}{3-p} |\dot{x}(t_n)|^{3-p} \\ &\leq \frac{1}{3-p} |\dot{x}(t_n)|^{3-p}, \end{aligned}$$

where, to obtain the last inequality, we use the hypothesis $2 < p < 3$. Since $\dot{x}(t_n)$ converges to zero as $t_n \rightarrow +\infty$, for n large enough we have that $\theta_n = +\infty$. This means that for n large enough

$$x(t) \in [c, d] \text{ and } |x(t) - l| \leq \frac{1}{3-p} |\dot{x}(t_n)|^{3-p} \quad \forall t \in [t_n, +\infty),$$

which implies that $x(t)$ converges to l as $t \rightarrow +\infty$. \square

Figure 6.2 illustrates the attractor effect when the damping becomes too weak. Take $f : \mathbb{R} \rightarrow \mathbb{R}$ $f(x) = \frac{1}{2}(x+1)^2$ for $x \leq -1$, $f(x) = 0$ for $|x| < 1$, and $f(x) = \frac{1}{2}(x-1)^2$ for $x \geq 1$.

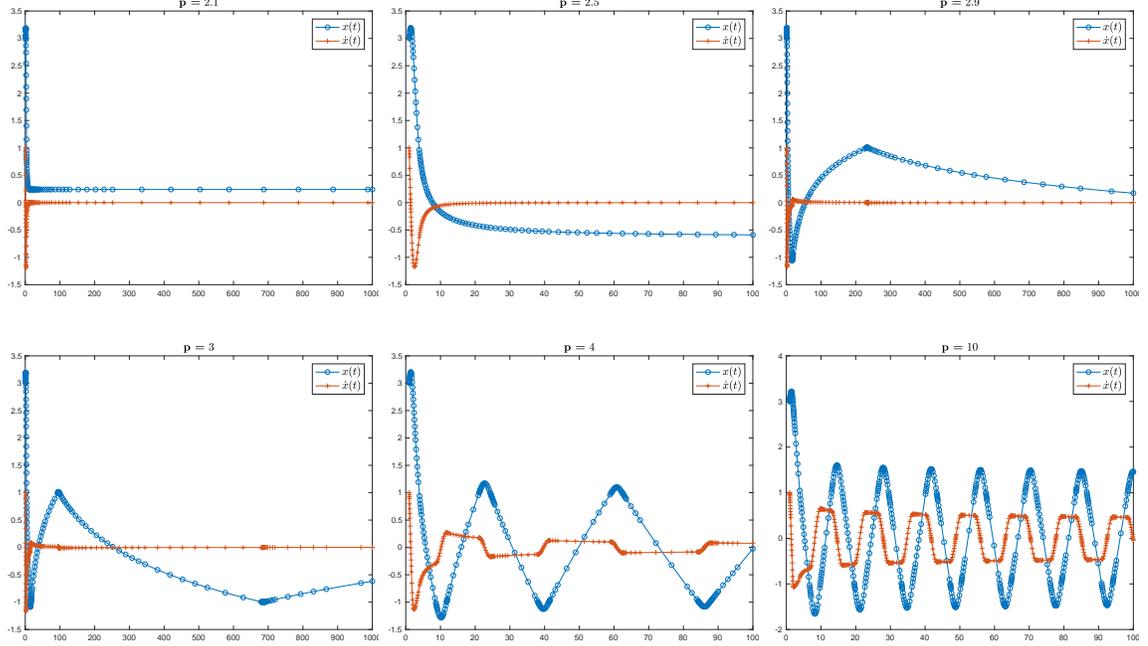


Fig. 6.2: Evolution of the trajectories $x(t)$ (blue) and $\dot{x}(t)$ (red) of the dynamical system (5.11) for $p \geq 2$.

For $p = 3$ there is a radical change in trajectory behavior. For $p \geq 3$, they do not converge asymptotically, and exhibit very oscillating behavior by steadily passing through the points -1 and $+1$. For $2 < p < 3$, there is numerical evidence that the trajectories converge, which confirms the conclusion of Proposition 6.3.

6.2. An explicit one-dimensional example. Take $\mathcal{H} = \mathbb{R}$ and $f(x) = c|x|^\gamma$, where c and γ are positive parameters. Let us look for solutions of

$$(6.3) \quad \ddot{x}(t) + r|\dot{x}(t)|^{p-2}\dot{x}(t) + \nabla f(x(t)) = 0,$$

when $p > 1$. Precisely, we look for nonnegative solutions of the form $x(t) = \frac{1}{t^\theta}$, with $\theta > 0$. This means that the trajectory is not oscillating, it is a completely damped trajectory. We proceed by identification, and determine the values of the parameters c, γ, r, p and θ which provide such solutions. On the one hand,

$$\ddot{x}(t) + r|\dot{x}(t)|^{p-2}\dot{x}(t) = \frac{\theta(\theta+1)}{t^{\theta+2}} - \frac{\theta^{p-1}r}{t^{(\theta+1)(p-1)}}.$$

On the other hand, $\nabla f(x) = c\gamma|x|^{\gamma-2}x$, which gives

$$\nabla f(x(t)) = \frac{c\gamma}{t^{\theta(\gamma-1)}}.$$

Thus, $x(t) = \frac{1}{t^\theta}$ is solution of (6.3) if, and only if,

$$\frac{\theta(\theta + 1)}{t^{\theta+2}} - \frac{\theta^{p-1}r}{t^{(\theta+1)(p-1)}} + \frac{c\gamma}{t^{\theta(\gamma-1)}} = 0.$$

This is equivalent to solve the following system:

- i)* $\theta + 2 = \theta(\gamma - 1)$;
- ii)* $\theta + 2 = (\theta + 1)(p - 1)$;
- iii)* $c\gamma = \theta^{p-1}r - \theta(\theta + 1)$;
- iv)* $\theta > 0, c > 0$

Solving *i)* and *ii)* with respect to γ gives $2 < p < 3$, $\gamma > 2$, and the following values of θ and p

$$\theta = \frac{2}{\gamma - 2}, \quad p = \frac{3\gamma - 2}{\gamma}.$$

Condition *iii)* gives $c = \frac{\theta}{\gamma} (\theta^{p-2}r - (\theta + 1))$ and the nonnegativity condition *iv)* gives $r > \frac{\theta+1}{\theta^{p-2}}$. We have $\min f = 0$ and

$$f(x(t)) = c \frac{1}{t^{\frac{2\gamma}{\gamma-2}}} = \frac{c}{t^{\frac{2}{p-2}}}.$$

To summarize, we have shown that when taking $2 < p < 3$, and $f(x) = c|x|^{\frac{2}{3-p}}$ there exists a solution of

$$\ddot{x}(t) + r\|\dot{x}(t)\|^p \dot{x}(t) + \nabla f(x(t)) = 0,$$

of the form $x(t) = \frac{1}{t^{\frac{3-p}{p-2}}}$, for which

$$f(x(t)) - \min f = \frac{c}{t^{\frac{2}{p-2}}}.$$

As expected, the speed of convergence of $f(x(t))$ to 0 depends on the parameter p . Therefore, without other geometric assumptions on f , for $2 < p < 3$, we cannot expect a convergence rate better than $\mathcal{O}\left(\frac{1}{t^{\frac{2}{p-2}}}\right)$. When $p \rightarrow 3$ from below, the function $f(x) = c|x|^{\frac{2}{3-p}}$ becomes very flat around its minimum (the origin) and the convergence rate of $x(t) = \frac{1}{t^{\frac{3-p}{p-2}}}$ to the origin becomes very slow.

7. Damping via closed-loop velocity control, quasi-gradient and (KL). In this section, $\mathcal{H} = \mathbb{R}^N$ is the finite dimensional Euclidean space. This will allow us to use the Kurdyka–Lojasiewicz property, which we briefly designate by (KL). Unless otherwise indicated, no convexity assumption is made on the function f to minimize, which will be assumed to satisfy (KL). To obtain the convergence of the orbits, the need for a geometric assumption on the function f to be minimized has long been recognized. As for the steepest descent, without additional geometric assumptions on the potential function f , the bounded orbits of the heavy ball with friction dynamic (HBF) may not converge. Let's recall the result from [26] where it is shown a function $f : \mathbb{R}^2 \rightarrow \mathbb{R}$ which is \mathcal{C}^1 , coercive, whose gradient is Lipschitz continuous on the bounded sets, and such that the (HBF) system admits an orbit $t \mapsto x(t)$ which does not converge as $t \rightarrow +\infty$. This example is an inertial version of the famous Palis–De Melo counterexample for the continuous steepest descent [76]. In this section, we examine an important situation where the convergence property is satisfied, namely when f is assumed to satisfy the property (KL), a geometric notion which is presented below.

7.1. Some basic facts concerning (KL). A function $G : \mathbb{R}^N \rightarrow \mathbb{R}$ satisfies the (KL) property if its values can be reparametrized in the neighborhood of each of its critical point, so that the resulting function become sharp. This means the existence of a continuous, concave, increasing function θ such that for all u in a slice of G

$$\|\nabla(\theta \circ G)(u)\| \geq 1.$$

The function θ captures the geometry of G around its critical point, it is called a desingularizing function; see [42], [14], [15] for further results. The tame functions satisfy the property (KL). Tameness refers to an ubiquitous geometric property of functions and sets encountered in most finite dimensional optimization problems. Sets or functions are called tame when they can be described by a finite number of basic formulas/ inequalities/Boolean operations involving standard functions such as polynomial, exponential, or max functions. Classical examples of tame objects are piecewise linear objects (with finitely many pieces), or semi-algebraic objects. The general notion covering these situations is the concept of \mathcal{o} -minimal structure; see van den Dries [82]. Tameness models nonsmoothness via the so-called stratification property of tame sets/functions. It was this property which motivated the vocable of tame topology, “la topologie modérée” according to Grothendieck. All these aspects have been well documented in a series of recent papers devoted to nonconvex nonsmooth optimization; see Ioffe [67] (“An invitation to tame optimization”), Castera–Bolte–Févotte–Pauwels [53] for an application to deep learning, and references therein. We refer to [15] for illustrations, and examples within a general optimization setting.

This property is particularly interesting in our context, because we work with an *autonomous dynamical system*, in which case the (KL) theory applies to the quasi-gradient systems. This contrasts with the accelerated gradient method of Nesterov which is based on a non-autonomous dynamic system, and for which we do not have a convergence theory based on the (KL) property. Under this property, we will obtain convergence results with convergence rates linked to the geometry of the data functions f and ϕ , via the desingularizing function.

7.2. Quasi-gradient systems. Let us first recall the main lines of the quasi-gradient approach to the inertial gradient systems as developed by Bégout–Bolte–Jendoubi in [41]. The geometric interpretation is simple: a vector field F is called quasi-gradient for a function E if it has the same singular point as E and if the angle between the field F and the gradient ∇E remains acute and bounded away from $\pi/2$. A precise definition is given below. Of course, such systems have a behavior which is very similar to those of gradient systems. We refer to Barta–Chill–Fašangová [38], [39], [56], Chergui [55], Huang [66] and the references therein for further geometrical insights on the topic.

DEFINITION 7.1. *Let Γ be a nonempty closed subset of \mathbb{R}^N , and let $F : \mathbb{R}^N \rightarrow \mathbb{R}^N$ be a locally Lipschitz continuous mapping. We say that the first-order system*

$$(7.1) \quad \dot{z}(t) + F(z(t)) = 0,$$

has a quasi-gradient structure for E on Γ , if there exist a differentiable function $E : \mathbb{R}^N \rightarrow \mathbb{R}$ and $\alpha > 0$ such that the two following conditions are satisfied:

$$\begin{aligned} (\text{angle condition}) \quad & \langle \nabla E(z), F(z) \rangle \geq \alpha \|\nabla E(z)\| \|F(z)\| \quad \text{for all } z \in \Gamma; \\ (\text{rest point equivalence}) \quad & \text{crit} E \cap \Gamma = F^{-1}(0) \cap \Gamma. \end{aligned}$$

Based on this notion, we have the following convergence properties for the bounded trajectories of (7.1). The following result is a localized version and straight adaptation of [41, Theorem 3.2].

THEOREM 7.2. *Let $F : \mathbb{R}^N \rightarrow \mathbb{R}^N$ be a locally Lipschitz continuous mapping. Let $z : [0, +\infty[\rightarrow \mathbb{R}^N$ be a bounded solution trajectory of (7.1). Take $R \geq \sup_{t \geq 0} \|z(t)\|$. Assume that F defines a quasi-gradient vector field for E_R on $\bar{B}(0, R)$, where $E_R : \mathbb{R}^N \rightarrow \mathbb{R}$ is a differentiable function. Assume further that the function E_R is (KL). Then, the following properties are satisfied:*

- (i) $z(t) \rightarrow z_\infty$ as $t \rightarrow +\infty$, where $z_\infty \in F^{-1}(0)$;
- (ii) $\dot{z} \in L^1(0, +\infty; \mathbb{R}^N)$, $\dot{z}(t) \rightarrow 0$ as $t \rightarrow +\infty$;
- (iii) $\|z(t) - z_\infty\| \leq \frac{1}{\alpha_R} \theta\left(E_R(z(t)) - E(z_\infty)\right)$

where θ is the desingularizing function for E_R at z_∞ , and α_R enters the angle condition of Definition 7.1.

7.3. Convergence of systems with closed-loop velocity control under (KL). Let us apply the above approach to the inertial system with closed-loop damping

$$(7.2) \quad \ddot{x}(t) + \nabla\phi(\dot{x}(t)) + \nabla f(x(t)) = 0,$$

by writing it as a first-order system, via its Hamiltonian formulation. We will assume that $\nabla\phi$ is locally Lipschitz continuous. Indeed, we can reduce to this situation by using a regularization procedure based on the Moreau envelope.

THEOREM 7.3. *Let $f : \mathbb{R}^N \rightarrow \mathbb{R}$ be a C^2 function whose gradient is Lipschitz continuous on the bounded sets, and such that $\inf_{\mathbb{R}^N} f > -\infty$. Let $E_\lambda : \mathbb{R}^N \times \mathbb{R}^N \rightarrow \mathbb{R}$ be defined by: for all $(x, u) \in \mathbb{R}^N \times \mathbb{R}^N$*

$$E_\lambda(x, u) := \frac{1}{2} \|u\|^2 + f(x) + \lambda \langle \nabla f(x), u \rangle.$$

Suppose that the function E_λ satisfies the (KL) property.

Let $\phi : \mathbb{R}^N \rightarrow \mathbb{R}_+$ be a damping potential (see Definition 3.1) which is differentiable, and that satisfies the following growth conditions:

- (i) (local) there exists positive constants γ, δ , and $\epsilon > 0$ such that, for all u in \mathbb{R}^N with $\|u\| \leq \epsilon$

$$\phi(u) \geq \gamma \|u\|^2 \text{ and } \|\nabla\phi(u)\| \leq \delta \|u\|.$$

- (ii) (global) there exists $p \geq 1, c > 0$, such that for all u in \mathbb{R}^N , $\phi(u) \geq c \|u\|^p$.

Let $x : [0, +\infty[\rightarrow \mathbb{R}^N$ be a bounded solution trajectory of

$$\ddot{x}(t) + \nabla\phi(\dot{x}(t)) + \nabla f(x(t)) = 0.$$

Then, the following properties are satisfied:

- (i) $x(t) \rightarrow x_\infty$ as $t \rightarrow +\infty$, where $x_\infty \in \text{crit } f$;
- (ii) $\dot{x} \in L^1(0, +\infty; \mathbb{R}^N)$, $\dot{x}(t) \rightarrow 0$ as $t \rightarrow +\infty$;
- (iii) For λ sufficiently small, and t sufficiently large

$$\|x(t) - x_\infty\| \leq \frac{1}{\alpha} \theta\left(E_\lambda(x(t), u(t)) - E_\lambda(x_\infty, 0)\right)$$

where θ is the desingularizing function for E_λ at $(x_\infty, 0)$, and α enters the corresponding angle condition.

Proof. According to the preliminary estimates established in Proposition 4.1, we have

$$\int_0^{+\infty} \phi(\dot{x}(t)) dt < +\infty \quad \text{and} \quad \sup_{t \geq 0} \|\dot{x}(t)\| < +\infty.$$

Combining the first above property with the global growth assumption on ϕ , we deduce that there exists $p \geq 1$ such that

$$\int_0^{+\infty} \|\dot{x}(t)\|^p dt < +\infty.$$

According to the constitutive equation, we have

$$\ddot{x}(t) = -\nabla\phi(\dot{x}(t)) - \nabla f(x(t)).$$

Since $x(\cdot)$ and $\dot{x}(\cdot)$ are bounded, and ∇f is locally Lipschitz continuous, we deduce that $\ddot{x}(\cdot)$ is also bounded. Classically, these properties imply that $\dot{x}(t) \rightarrow 0$ as $t \rightarrow +\infty$. Take $R \geq \sup_{t \geq 0} \|x(t)\|$. Therefore, for t sufficiently large, we have that the trajectory $t \mapsto (x(t), \dot{x}(t))$ in the phase space $\mathbb{R}^N \times \mathbb{R}^N$ belongs to the closed set $\Gamma = \bar{B}(0, R) \times \bar{B}(0, \epsilon)$.

The Hamiltonian formulation of (7.2) gives the first-order differential system

$$(7.3) \quad \dot{z}(t) + F(z(t)) = 0,$$

where $z(t) = (x(t), \dot{x}(t)) \in \mathbb{R}^N \times \mathbb{R}^N$, and $F : \mathbb{R}^N \times \mathbb{R}^N \rightarrow \mathbb{R}^N \times \mathbb{R}^N$ is defined by

$$(7.4) \quad F(x, u) = (-u, \nabla\phi(u) + \nabla f(x)).$$

Following [41], take $E_\lambda : \mathbb{R}^N \rightarrow \mathbb{R}$ defined by

$$(7.5) \quad E_\lambda(x, u) := \frac{1}{2}\|u\|^2 + f(x) + \lambda \langle \nabla f(x), u \rangle,$$

where λ is a positive parameter which will be adjusted to verify the quasi-gradient property. We have

$$\nabla E_\lambda(x, u) = \left(\nabla f(x) + \lambda \nabla^2 f(x)u, u + \lambda \nabla f(x) \right).$$

Let us analyze the angle condition with $\Gamma = \bar{B}(0, R) \times \bar{B}(0, \epsilon)$. According to the above formulation of F and ∇E_λ , we have

$$\begin{aligned} \langle \nabla E_\lambda(x, u), F(x, u) \rangle &= \left\langle \left(\nabla f(x) + \lambda \nabla^2 f(x)u, u + \lambda \nabla f(x) \right), \left(-u, \nabla\phi(u) + \nabla f(x) \right) \right\rangle \\ &= -\langle \nabla f(x) + \lambda \nabla^2 f(x)u, u \rangle + \langle u + \lambda \nabla f(x), \nabla\phi(u) + \nabla f(x) \rangle. \end{aligned}$$

After development and simplification, we get

$$\langle \nabla E_\lambda(x, u), F(x, u) \rangle = -\lambda \langle \nabla^2 f(x)u, u \rangle + \langle u, \nabla\phi(u) \rangle + \lambda \langle \nabla f(x), \nabla\phi(u) \rangle + \lambda \|\nabla f(x)\|^2.$$

According to the local Lipschitz assumption on ∇f , let

$$M := \sup_{\|x\| \leq R} \|\nabla^2 f(x)\| < +\infty.$$

Since ϕ is a damping potential, we have

$$\langle u, \nabla\phi(u) \rangle \geq \phi(u).$$

Combining the above results, we obtain

$$\begin{aligned} \langle \nabla E_\lambda(x, u), F(x, u) \rangle &\geq -\lambda M \|u\|^2 + \phi(u) + \lambda \langle \nabla f(x), \nabla\phi(u) \rangle + \lambda \|\nabla f(x)\|^2 \\ &\geq -\lambda M \|u\|^2 + \phi(u) - \frac{\lambda}{2} \|\nabla f(x)\|^2 - \frac{\lambda}{2} \|\nabla\phi(u)\|^2 + \lambda \|\nabla f(x)\|^2 \\ (7.6) \quad &\geq -\lambda M \|u\|^2 + \phi(u) - \frac{\lambda}{2} \|\nabla\phi(u)\|^2 + \frac{\lambda}{2} \|\nabla f(x)\|^2. \end{aligned}$$

At this point, we use the local growth assumption on ϕ : for all u in \mathbb{R}^N with $\|u\| \leq \epsilon$

$$(7.7) \quad \phi(u) \geq \gamma \|u\|^2 \text{ and } \|\nabla\phi(u)\| \leq \delta \|u\|.$$

By combining (7.6) with (7.7), we obtain

$$(7.8) \quad \langle \nabla E_\lambda(x, u), F(x, u) \rangle \geq \left(\gamma - \lambda M - \frac{\lambda}{2} \delta^2 \right) \|u\|^2 + \frac{\lambda}{2} \|\nabla f(x)\|^2.$$

Take λ small enough to satisfy

$$\gamma > \lambda \left(M + \frac{\delta^2}{2} \right).$$

Then,

$$(7.9) \quad \langle \nabla E_\lambda(x, u), F(x, u) \rangle \geq \alpha_0 (\|u\|^2 + \|\nabla f(x)\|^2)$$

with $\alpha_0 := \min\{\gamma - \lambda \left(M + \frac{\delta^2}{2} \right), \frac{\lambda}{2}\}$.

On the other hand,

$$\begin{aligned} \|\nabla E_\lambda(x, u)\| &\leq \sqrt{2} (1 + \lambda \max\{1, M\}) (\|u\|^2 + \|\nabla f(x)\|^2)^{\frac{1}{2}} \\ \|F(x, u)\| &\leq \sqrt{2} (1 + \delta) (\|u\|^2 + \|\nabla f(x)\|^2)^{\frac{1}{2}}. \end{aligned}$$

Therefore

$$(7.10) \quad \|\nabla E_\lambda(x, u)\| \|F(x, u)\| \leq 2 (1 + \lambda \max\{1, M\}) (1 + \delta) (\|u\|^2 + \|\nabla f(x)\|^2).$$

As a consequence, the angle condition

$$\langle \nabla E(z), F(z) \rangle \geq \alpha \|\nabla E(z)\| \|F(z)\|$$

is satisfied on Γ , by taking

$$\alpha = \frac{\min\{\gamma - \lambda \left(M + \frac{\delta^2}{2} \right), \frac{\lambda}{2}\}}{2 (1 + \lambda \max\{1, M\}) (1 + \delta)}.$$

Finally, the rest point equivalence is a consequence of the inequality (7.9). Then, apply the abstract theorem 7.2 to obtain the claims (i), (ii), (iii). \square

Remark 7.4. (i) The above result allows to consider nonlinear damping. The main restrictive assumption is that the damping potential is assumed to be nearly quadratic close to the origin. It is not necessarily quadratic close to the origin but it has to satisfy: for all u in \mathbb{R}^N with $\|u\| \leq \epsilon$

$$\phi(u) \geq \gamma\|u\|^2 \text{ and } \|\nabla\phi(u)\| \leq \delta\|u\|.$$

(ii) According to [41, Proposition 3.11], a desingularizing function of f (see [41, Definition 2.1]) is desingularizing of E_λ too, for all $\lambda \in [0, \lambda_1]$.

(iii) In Section 9.3 and Section 10.6 we will develop similar analysis for related dynamical systems which involve Hessian damping.

(iv) Following [41, Theorem 4.1 and Theorem 3.7], a key condition which induces convergence rates for the trajectories of a quasi-gradient system in the framework of the KL property is

$$(7.11) \quad \|\nabla E_\lambda(x, u)\| \leq b\|F(x, u)\| \text{ for } (x, u) \in \Gamma$$

with $b > 0$. Let us check this condition in the setting of Theorem 7.3. We have seen there that

$$\|\nabla E_\lambda(x, u)\|^2 \leq C_1(\|u\|^2 + \|\nabla f(x)\|^2),$$

with $C_1 > 0$. Further, from the Cauchy–Schwarz inequality and the properties of ϕ we derive for $\sigma > 1$:

$$\begin{aligned} \|F(x, u)\|^2 &= \|u\|^2 + \|\nabla\phi(u) + \nabla f(x)\|^2 \\ &\geq \|u\|^2 + \|\nabla\phi(u)\|^2 + \|\nabla f(x)\|^2 - 2\|\nabla\phi(u)\|\|\nabla f(x)\| \\ &\geq \|u\|^2 + \|\nabla\phi(u)\|^2 + \|\nabla f(x)\|^2 - \sigma\|\nabla\phi(u)\|^2 - \frac{1}{\sigma}\|\nabla f(x)\|^2 \\ &\geq \left(1 - \frac{1}{\sigma}\right)\|\nabla f(x)\|^2 + (1 - (\sigma - 1)\delta^2)\|u\|^2. \end{aligned}$$

From this we can chose $\sigma > 1$ such that

$$\|F(x, u)\|^2 \geq C_3(\|u\|^2 + \|\nabla f(x)\|^2),$$

with $C_3 > 0$. Condition (7.11) is fulfilled now with $b = \sqrt{C_1/C_3}$.

As in [41, Section 5], explicit convergence rates can be derived from [41, Theorem 4.1 and Theorem 3.7], based on (3.19) and Remark 3.4(c) in [41].

7.4. Application: f with polynomial growth. This concerns the question raised at the end of Section 5.1. Additionally to the hypotheses of Theorem 7.3, assume that f is convex, $\text{argmin } f \neq \emptyset$ and for each $x^* \in \text{argmin } f$, there exists $\eta > 0$ such that

$$f(x) - \inf_{\mathcal{H}} f \geq c \text{dist}(x, \text{argmin } f)^r \quad \forall x \in B(x^*, \eta),$$

with $r \geq 1$ and $c > 0$.

According to the proof of [41, Corollary 5.5], f satisfies the Łojasiewicz inequality with desingularizing function (see [41, Definition 2.1]) of the form $\varphi(s) = c's^{1-1/r}$, with $c' > 0$. According to [41, Proposition 3.11], this is a desingularizing function of E_λ too, for all $\lambda \in [0, \lambda_1]$ (with E_λ defined in Theorem 7.3). In Remark 7.4(iv) we have shown that (7.11) holds. Relying now on [41, Theorem 3.7] and [41, Remark 3.4(c)], we derive sublinear rates for $\|x(t) - x_\infty\|$ in case $r < 2$ and exponential rate in case $r = 2$.

7.5. Application: fixed damping matrix. We will recover and improve the results of Alvarez [7, Theorem 2.6], which concerns the case f convex, and the damped inertial equation

$$\ddot{x}(t) + A(\dot{x}(t)) + \nabla f(x(t)) = 0,$$

where $A : \mathcal{H} \rightarrow \mathcal{H}$ is a positive definite self-adjoint linear operator, which is possibly anisotropic (see also [45]). While the proof of Theorem 2.6 in [7] works in general Hilbert spaces, we have to restrict ourselves to finite-dimensional spaces, however, we can drop the convexity assumption on f . The following result is a direct consequence of Theorem 7.3 applied to $\phi : \mathbb{R}^N \rightarrow \mathbb{R}$, $\phi(x) = \frac{1}{2} \langle Ax, x \rangle$, where $A : \mathbb{R}^N \rightarrow \mathbb{R}^N$ is a positive definite self-adjoint linear operator. We note that in this setting, ϕ is a damping potential, it is convex continuous and attains its minimum at the origin. Moreover, the local and global growth conditions are met. Indeed, for all $u \in \mathbb{R}^N$, we have $\phi(u) \geq \frac{1}{2} \lambda_{\min} \|u\|^2$ and $\|\nabla \phi(u)\| \leq \lambda_{\max} \|u\|$, where λ_{\min} and λ_{\max} are the smallest and largest positive eigenvalues of A respectively.

THEOREM 7.5. *Let $f : \mathbb{R}^N \rightarrow \mathbb{R}$ be a C^2 function whose gradient is Lipschitz continuous on the bounded sets, and such that $\inf_{\mathbb{R}^N} f > -\infty$. Let $A : \mathbb{R}^N \rightarrow \mathbb{R}^N$ a positive definite self-adjoint linear operator. Suppose that the function E_λ is (KL) (which is true if f is (KL)) where*

$$E_\lambda(x, u) := \frac{1}{2} \|u\|^2 + f(x) + \lambda \langle \nabla f(x), u \rangle.$$

Let $x : [0, +\infty[\rightarrow \mathbb{R}^N$ be a bounded solution trajectory of

$$\ddot{x}(t) + A(\dot{x}(t)) + \nabla f(x(t)) = 0.$$

Then, the following properties are satisfied:

- (i) $x(t) \rightarrow x_\infty$ as $t \rightarrow +\infty$, where $x_\infty \in \text{crit } f$;
- (ii) $\dot{x} \in L^1(0, +\infty; \mathbb{R}^N)$, $\dot{x}(t) \rightarrow 0$ as $t \rightarrow +\infty$;
- (iii) $f(x(t)) \rightarrow f(x_\infty) \in f(\text{crit } f)$ as $t \rightarrow +\infty$.

Indeed, we can complete this result with the convergence rates which are linked to the desingularization function provided with the property (KL) for f .

8. Algorithmic results: an inertial type algorithm. Consider the following explicit temporal discretization of (ADIGE-V) with step size $h > 0$

$$\frac{1}{h^2} (x_{n+2} - 2x_{n+1} + x_n) + \nabla \phi \left(\frac{1}{h} (x_{n+1} - x_n) \right) + \nabla f(x_n) = 0.$$

This gives the algorithm

$$(8.1) \quad x_{n+2} = 2x_{n+1} - x_n - h^2 \nabla \phi \left(\frac{1}{h} (x_{n+1} - x_n) \right) - h^2 \nabla f(x_n).$$

In a parallel way to the Hamiltonian approach of (ADIGE-V) in the continuous case, let us introduce the vector field $F : \mathbb{R}^N \times \mathbb{R}^N \rightarrow \mathbb{R}^N \times \mathbb{R}^N$ defined by

$$(8.2) \quad F(x, u) = (-u, \nabla \phi(u) + \nabla f(x)).$$

Using the notation $u_n := \frac{1}{h}(x_{n+1} - x_n)$, the above scheme can be written in the condensed form

$$(8.3) \quad (x_{n+1}, u_{n+1}) - (x_n, u_n) + hF(x_n, u_n) = 0,$$

which can be seen as a discretized version of the continuous dynamic (7.3). The following convergence result is a discrete algorithmic version of Theorem 7.3. Like the latter, it is based on the quasi-gradient approach, and is valid under similar hypothesis.

THEOREM 8.1. *Let $f : \mathbb{R}^N \rightarrow \mathbb{R}$ be a \mathcal{C}^2 function whose gradient is Lipschitz continuous on the bounded sets, and such that $\inf_{\mathbb{R}^N} f > -\infty$. Let $\phi : \mathbb{R}^N \rightarrow \mathbb{R}_+$ be a damping potential (see Definition 3.1) which is differentiable. Let $(x_n)_{n \in \mathbb{N}}$ be a bounded sequence generated by the algorithm*

$$x_{n+2} = 2x_{n+1} - x_n - h^2 \nabla \phi \left(\frac{1}{h}(x_{n+1} - x_n) \right) - h^2 \nabla f(x_n).$$

We make the following assumptions on the data f , ϕ , and h :

- (assumption on f): Suppose that the function E_λ satisfies the (KL) property, where $E_\lambda : \mathbb{R}^N \times \mathbb{R}^N \rightarrow \mathbb{R}$ is defined by: for all $(x, u) \in \mathbb{R}^N \times \mathbb{R}^N$

$$E_\lambda(x, u) := \frac{1}{2} \|u\|^2 + f(x) + \lambda \langle \nabla f(x), u \rangle.$$

- (assumption on ϕ) Suppose that the function ϕ satisfies the following growth conditions:

- (i) (local) there exists positive constants γ , δ , and $\epsilon > 0$ such that, for all u in \mathbb{R}^N with $\|u\| \leq \epsilon$

$$\phi(u) \geq \gamma \|u\|^2 \text{ and } \|\nabla \phi(u)\| \leq \delta \|u\|.$$

- (ii) (global) there exists $p \geq 1$, $c > 0$, such that for all u in \mathbb{R}^N , $\phi(u) \geq c \|u\|^p$.

- (assumption on h) Suppose that the step size h is taken small enough to satisfy

$$0 < h < \min \left\{ \frac{\alpha_0}{L}, \frac{\alpha_0}{L/2 + L\delta^2} \right\},$$

where $R = \sup_n \{\|(x_n, u_n)\|\}$, L is the Lipschitz constant of ∇E_λ on the ball centered at the origin and with radius R , and α_0 enters the quasi-gradient property (7.9).

Then, the following properties are satisfied:

- (i) $x_n \rightarrow x_\infty$ as $n \rightarrow +\infty$, where $x_\infty \in \text{crit } f$;
- (ii) $\sum_{n \in \mathbb{N}} \|x_{n+1} - x_n\| < +\infty$;
- (iii) For λ small enough, and n large enough

$$\|x_n - x_\infty\| \leq \frac{1}{\alpha} \theta \left(E_\lambda(x_n, u_n) - E_\lambda(x_\infty, 0) \right),$$

where θ is the desingularizing function for E_λ at $(x_\infty, 0)$, and α enters the corresponding angle condition.

Proof. By assumption, the sequence $(x_n)_{n \in \mathbb{N}}$ is bounded. By definition of $u_n = \frac{1}{h}(x_{n+1} - x_n)$, this implies that the sequence $(u_n)_{n \in \mathbb{N}}$ is also bounded (h is fixed). Indeed, the classical energy estimate gives that $(u_n)_{n \in \mathbb{N}}$ tends to zero. Set

$$R = \sup_n \{\|(x_n, u_n)\|\}.$$

Since ∇f is Lipschitz continuous on the bounded sets, this immediately implies that the gradient of E_λ defined in (7.5) is Lipschitz on the bounded sets. According to the descent lemma, and denoting by L the Lipschitz constant of ∇E_λ on the ball centered at the origin and with radius R , we have

$$\begin{aligned} E_\lambda(x_{n+1}, u_{n+1}) &\leq E_\lambda(x_n, u_n) + \langle \nabla E_\lambda(x_n, u_n), (x_{n+1}, u_{n+1}) - (x_n, u_n) \rangle \\ &\quad + \frac{L}{2} (\|x_{n+1} - x_n\|^2 + \|u_{n+1} - u_n\|^2). \end{aligned}$$

Using the formulation of the algorithm (8.3), we get

$$E_\lambda(x_{n+1}, u_{n+1}) \leq E_\lambda(x_n, u_n) - h \langle \nabla E_\lambda(x_n, u_n), F(x_n, u_n) \rangle + \frac{L}{2} (\|x_{n+1} - x_n\|^2 + \|u_{n+1} - u_n\|^2).$$

According to the quasi-gradient property (7.9), we derive (8.4)

$$E_\lambda(x_{n+1}, u_{n+1}) \leq E_\lambda(x_n, u_n) - h\alpha_0(\|u_n\|^2 + \|\nabla f(x_n)\|^2) + \frac{L}{2} (\|x_{n+1} - x_n\|^2 + \|u_{n+1} - u_n\|^2).$$

Further, due to (8.3) and (7.4) we have $u_{n+1} - u_n = -h(\nabla\phi(u_n) + \nabla f(x_n))$. Since (u_n) tends to zero, by the local growth assumption on ϕ we have that, for n sufficiently large, say $n \geq n_1$, $\|\nabla\phi(u_n)\| \leq \delta\|u_n\|$. Therefore,

$$\|u_{n+1} - u_n\|^2 \leq 2h^2\delta^2\|u_n\|^2 + 2h^2\|\nabla f(x_n)\|^2.$$

Finally, since $u_n = \frac{1}{h}(x_{n+1} - x_n)$, we get from (8.4):

$$(8.5) \quad E_\lambda(x_{n+1}, u_{n+1}) \leq E_\lambda(x_n, u_n) - \left(\frac{\alpha_0}{h} - \frac{L}{2} - L\delta^2 \right) \|x_{n+1} - x_n\|^2 - (h\alpha_0 - Lh^2)\|\nabla f(x_n)\|^2.$$

In order to be able to telescope the above inequalities, we choose the step size h according to:

$$0 < h < \min\left\{ \frac{\alpha_0}{L}, \frac{\alpha_0}{L/2 + L\delta^2} \right\}.$$

Moreover, $E_\lambda(x, u)$ is bounded from below. Indeed,

$$E_\lambda(x, u) \geq f(x) - \frac{\lambda^2}{2} \|\nabla f(x)\|^2.$$

Further, if we suppose that ∇f is Lipschitz with constant $L_f \geq 0$, one can easily show that for λ small enough, there exists $\sigma > 0$ with the property $\frac{1}{\sigma} - \frac{L_f}{2\sigma^2} = \frac{\lambda^2}{2}$. Hence

$$E_\lambda(x, u) \geq f(x) - \left(\frac{1}{\sigma} - \frac{L_f}{2\sigma^2} \right) \|\nabla f(x)\|^2 > \inf f,$$

where we used Proposition 1 in [48] (see also Remark 3 in [70]).
By telescoping (8.5) we derive

$$(8.6) \quad \lim_{n \rightarrow +\infty} (x_{n+1} - x_n) = 0 \quad \text{and} \quad \lim_{n \rightarrow +\infty} \nabla f(x_n) = 0.$$

Moreover, there exists

$$(8.7) \quad \lim_{n \rightarrow +\infty} E_\lambda(x_n, u_n) \in \mathbb{R}.$$

Further, let us denote by $\omega((x_n)_{n \geq 0})$ the set of cluster points of the sequence $(x_n)_{n \geq 0}$, and by $\text{crit } f = \{x \in \mathbb{R}^N : \nabla f(x) = 0\}$ the set of critical points of f .

From the theory of quasi-gradient systems we have (see also (7.9))

$$\text{crit } E_\lambda = \text{crit } f \times \{0\}.$$

We notice that $\omega((x_n)_{n \geq 0}) \subseteq \text{crit } f$, thus $\omega((x_n, u_n)_{n \geq 0}) = \omega((x_n)_{n \geq 0}) \times \{0\} \subseteq \text{crit } E_\lambda$. From (8.7) one can easily conclude that E_λ is constant on $\omega((x_n, u_n))$. Indeed, for $x^* \in \omega((x_n))$, we have from above (and the definition of E_λ) that

$$(8.8) \quad \lim_{n \rightarrow +\infty} E_\lambda(x_n, u_n) = f(x^*) = E_\lambda(x^*, 0).$$

Assume now that E_λ satisfies the KL property with corresponding desingularizing function θ . We consider two cases.

I. There exists $\bar{n} \geq 0$ such that $E_\lambda(x_{\bar{n}}, u_{\bar{n}}) = E_\lambda(x^*, 0)$. From the decreasing property (8.5) we obtain that $(x_n)_{n \geq \bar{n}}$ is a constant sequence and the conclusion follows.

II. For all $n \geq 0$ we have $E_\lambda(x_n, u_n) > E_\lambda(x^*, 0)$. Since θ is concave and $\theta' > 0$, we derive from (8.5) that there exists $M > 0$ and $n' \geq 0$ such that for all $n \geq n'$ it holds

$$(8.9) \quad \begin{aligned} \Delta_{n,n+1} &:= \theta(E_\lambda(x_n, u_n) - E_\lambda(x^*, 0)) - \theta(E_\lambda(x_{n+1}, u_{n+1}) - E_\lambda(x^*, 0)) \\ &\geq \theta'(E_\lambda(x_n, u_n) - E_\lambda(x^*, 0)) \cdot (E_\lambda(x_n, u_n) - E_\lambda(x_{n+1}, u_{n+1})) \\ &\geq \theta'(E_\lambda(x_n, u_n) - E_\lambda(x^*, 0)) \cdot M \cdot (\|x_{n+1} - x_n\|^2 + \|\nabla f(x_n)\|^2) \\ &\geq \frac{M(\|x_{n+1} - x_n\|^2 + \|\nabla f(x_n)\|^2)}{\|\nabla E_\lambda(x_n, u_n)\|}, \end{aligned}$$

where the last inequality (8.9) follows from the uniformized KL property Lemma 6 in [43] applied to the nonempty compact and connected set $\Omega = \omega((x_n)_{n \geq 0}) \times \{0\}$ (notice that according to [43, Remark 5] the connectedness of this set is generic for sequences satisfying $\lim_{n \rightarrow +\infty} (x_{n+1} - x_n) = 0$). We obtain for all $n \geq 0$

$$(8.10) \quad \begin{aligned} \|x_{n+1} - x_n\| + \|\nabla f(x_n)\| &\leq \sqrt{2(\|x_{n+1} - x_n\|^2 + \|\nabla f(x_n)\|^2)} \\ &\leq \sqrt{\frac{2}{M}} \cdot \sqrt{\|\nabla E_\lambda(x_n, u_n)\| \Delta_{n,n+1}} \\ &\leq \sqrt{\frac{2}{M}} \cdot \frac{\nu \|\nabla E_\lambda(x_n, u_n)\| + (1/\nu) \Delta_{n,n+1}}{2}, \end{aligned}$$

where $\nu > 0$ will be chosen afterwards. We have seen in Section 5 that there exists $C > 0$ such that

$$\|\nabla E_\lambda(x_n, u_n)\| \leq C(\|x_{n+1} - x_n\| + \|\nabla f(x_n)\|) \quad \forall n \geq 0.$$

Combining the last two inequalities, we obtain

$$(8.11) \quad \left(1 - \sqrt{\frac{2}{M}} \cdot \frac{\nu}{2} C\right) (\|x_{n+1} - x_n\| + \|\nabla f(x_n)\|) \leq \sqrt{\frac{2}{M}} \cdot \frac{1}{\nu} \Delta_{n,n+1} \quad \forall n \geq 0.$$

We choose $\nu > 0$ such that $1 - \sqrt{\frac{2}{M}} \cdot \frac{\nu}{2} C > 0$. Since the right-hand side of (8.11) is summable, we deduce that $\sum_{n \geq 0} \|x_{n+1} - x_n\| < \infty$. Therefore, $(x_n)_{n \geq 0}$ is a Cauchy sequence in \mathbb{R}^N , and hence it converges to a critical point of f . The end of the proof is similar to the continuous case. \square

9. Closed-loop velocity control with Hessian driven damping.

9.1. Hessian damping. We propose to tackle questions similar to the previous sections, concerning the combination of closed-loop velocity control with Hessian driven damping. The following system combines closed-loop velocity control with Hessian driven damping:

$$(9.1) \quad \ddot{x}(t) + \partial\phi(\dot{x}(t)) + \beta \nabla^2 f(x(t)) \dot{x}(t) + \nabla f(x(t)) = 0.$$

This autonomous system will be our main subject of study in this section.

- The case $\phi(u) = \frac{\gamma}{2} \|u\|^2$ of a fixed viscous coefficient was first considered by Alvarez–Attouch–Bolte–Redont in [9]. In this case, (9.1) can be equivalently written as a first-order system in time and space (different from the Hamiltonian formulation), which allows to extend this system naturally to the case of a nonsmooth function f . This property has been exploited by Attouch–Maingé–Redont [29] for modeling non-elastic shocks in unilateral mechanics. To accelerate this system, several recent studies considered the case where the viscous damping is vanishing, that is

$$(9.2) \quad \ddot{x}(t) + \frac{\alpha}{t} \dot{x}(t) + \beta \nabla^2 f(x(t)) \dot{x}(t) + \nabla f(x(t)) = 0;$$

see [33], [22], [47], [53], [69], [71], [80], and Section 2.3 for the properties of this system.

- The case $\phi(u) = \frac{\gamma}{2} \|u\|^2 + r \|u\|$ which combines viscous friction with dry friction and Hessian damping has been considered by Adly–Attouch [2], [4].
- By taking $\phi(u) = \frac{r}{p} \|u\|^p$, we get

$$(9.3) \quad \ddot{x}(t) + r \|\dot{x}(t)\|^{p-2} \dot{x}(t) + \beta \nabla^2 f(x(t)) \dot{x}(t) + \nabla f(x(t)) = 0,$$

for which we will address issues similar to those of the previous theme. In addition to the fast minimization property, one can expect obtaining too the fast convergence of the gradients to zero.

9.2. Existence and uniqueness results.

Let us consider the differential inclusion

$$(9.4) \quad (\text{ADIGE-VH}) \quad \ddot{x}(t) + \partial\phi(\dot{x}(t)) + \beta \nabla^2 f(x(t)) \dot{x}(t) + \nabla f(x(t)) \ni 0,$$

which involves a damping potential ϕ (see Definition 3.1), and a geometric damping driven by the Hessian of f . The suffix V makes reference to the velocity and H to the Hessian. They both enter the damping terms. It allows to cover different situations, in particular system (9.3) corresponds to $\phi(u) = \frac{r}{p} \|u\|^p$ for $p > 1$. To prove existence and uniqueness results for the associated Cauchy problem, we make additional assumptions. We assume that f is convex, and that the Hessian mapping $x \in \mathcal{H} \mapsto \nabla^2 f(x) \in \mathcal{L}(\mathcal{H}, \mathcal{H})$ is Lipschitz continuous on the bounded sets, where $\mathcal{L}(\mathcal{H}, \mathcal{H})$ is equipped with the norm operator. Note that this property implies that ∇f is Lipschitz continuous on the bounded subsets of \mathcal{H} (apply the mean value theorem in the vectorial case). However, in the following statement, we formulate the two hypotheses for the sake of clarity.

THEOREM 9.1. *Let $f : \mathcal{H} \rightarrow \mathbb{R}$ be a convex function which is twice continuously differentiable, and such that $\inf_{\mathcal{H}} f > -\infty$. We suppose that*

- (i) ∇f is Lipschitz continuous on the bounded subsets of \mathcal{H} ;
- (ii) $\nabla^2 f$ is Lipschitz continuous on the bounded subsets of \mathcal{H} .

Let $\phi : \mathcal{H} \rightarrow \mathbb{R}$ be a convex continuous damping function. Then, for any Cauchy data $(x_0, x_1) \in \mathcal{H} \times \mathcal{H}$, there exists a unique strong global solution $x : [0, +\infty[\rightarrow \mathcal{H}$ of (ADIGE-VH) satisfying $x(0) = x_0$, and $\dot{x}(0) = x_1$.

Proof. To make the reading of the proof easier, we distinguish several steps.

Step 1: A priori estimate. Let's establish a priori energy estimates on the solutions of (9.4). After taking the scalar product of (9.4) with $\dot{x}(t)$, we get

$$\frac{d}{dt}\mathcal{E}(t) + \langle \partial\phi(\dot{x}(t)), \dot{x}(t) \rangle + \beta \langle \nabla^2 f(x(t))\dot{x}(t), \dot{x}(t) \rangle = 0,$$

where

$$\mathcal{E}(t) := f(x(t)) - \inf_{\mathcal{H}} f + \frac{1}{2}\|\dot{x}(t)\|^2$$

is the global energy. Since ϕ is a damping potential, the subdifferential inequality for convex functions, combined with $\phi(0) = 0$, gives

$$\langle \partial\phi(\dot{x}(t)), \dot{x}(t) \rangle \geq \phi(\dot{x}(t)).$$

Since f is convex, we have that $\nabla^2 f$ is positive semidefinite, which gives

$$\langle \nabla^2 f(x(t))\dot{x}(t), \dot{x}(t) \rangle \geq 0.$$

Collecting the above results, we obtain the following decay property of the energy

$$(9.5) \quad \frac{d}{dt}\mathcal{E}(t) + \phi(\dot{x}(t)) \leq 0.$$

Therefore, the energy is nonincreasing, which implies that, as long as the trajectory is defined

$$(9.6) \quad \|\dot{x}(t)\|^2 \leq 2\mathcal{E}(0).$$

Step 2: Hamiltonian formulation of (9.4). According to the Hamiltonian formulation of (9.4), it is equivalent to solve the first-order system

$$\begin{cases} \dot{x}(t) - u(t) = 0; \\ \dot{u}(t) + \partial\phi(u(t)) + \nabla f(x(t)) + \beta\nabla^2 f(x(t))u(t) \ni 0, \end{cases}$$

with the Cauchy data $x(0) = x_0$, $u(0) = x_1$. Set $Z(t) = (x(t), u(t)) \in \mathcal{H} \times \mathcal{H}$. The above system can be written equivalently as

$$\dot{Z}(t) + F(Z(t)) \ni 0, \quad Z(0) = (x_0, x_1),$$

where $F : \mathcal{H} \times \mathcal{H} \rightrightarrows \mathcal{H} \times \mathcal{H}$, $(x, u) \mapsto F(x, u)$ is defined by

$$F(x, u) = \left(0, \partial\phi(u)\right) + \left(-u, \nabla f(x) + \beta \nabla^2 f(x)u\right).$$

Hence F splits as follows $F(x, u) = \partial\Phi(x, u) + G(x, u)$, where

$$(9.7) \quad \Phi(x, u) = \phi(u) \text{ and } G(x, u) = \left(-u, \nabla f(x) + \beta \nabla^2 f(x)u\right).$$

Therefore, it is equivalent to solve the following first-order differential inclusion with Cauchy data

$$(9.8) \quad \dot{Z}(t) + \partial\Phi(Z(t)) + G(Z(t)) \ni 0, \quad Z(0) = (x_0, x_1).$$

Let us prove that the mapping $(x, u) \mapsto G(x, u)$ is Lipschitz continuous on the bounded subsets of $\mathcal{H} \times \mathcal{H}$. For any $(x, u) \in \mathcal{H} \times \mathcal{H}$, set $G(x, u) = (-u, K(x, u))$ where

$$K(x, u) := \nabla f(x) + \beta \nabla^2 f(x)u.$$

Let L_R be the Lipschitz constant of ∇f and $\nabla^2 f$ on the ball centered at the origin and with radius R , and set $M_R = \sup_{\|x\| \leq R} \|\nabla^2 f(x)\|$. Take $(x_i, u_i) \in \mathcal{H} \times \mathcal{H}$, $i = 1, 2$ with $\|(x_i, u_i)\| \leq R$. We have

$$K(x_2, u_2) - K(x_1, u_1) = \nabla f(x_2) - \nabla f(x_1) + \beta(\nabla^2 f(x_2)u_2 - \nabla^2 f(x_1)u_1).$$

According to the triangle inequality, and the local Lipschitz continuity property of ∇f and $\nabla^2 f$

$$\begin{aligned} \|K(x_2, u_2) - K(x_1, u_1)\| &\leq \|\nabla f(x_2) - \nabla f(x_1)\| + \beta\|\nabla^2 f(x_2)u_2 - \nabla^2 f(x_1)u_1\| \\ &\quad + \beta\|\nabla^2 f(x_1)u_2 - \nabla^2 f(x_1)u_1\| \\ &\leq L_R\|x_2 - x_1\| + \beta L_R\|x_2 - x_1\|\|u_2\| + \beta M_R\|u_2 - u_1\| \\ &\leq L_R(1 + R\beta)\|x_2 - x_1\| + \beta M_R\|u_2 - u_1\|. \end{aligned}$$

Therefore,

$$(9.9) \quad \|G(x_2, u_2) - G(x_1, u_1)\| \leq L_R(1 + R\beta)\|x_2 - x_1\| + (1 + \beta M_R)\|u_2 - u_1\|,$$

which gives that the mapping $(x, u) \mapsto G(x, u)$ is Lipschitz continuous on the bounded subsets of $\mathcal{H} \times \mathcal{H}$.

Step 3: Approximate dynamics. We proceed in a similar way as in Theorem 3.3 (which corresponds to the case $\beta = 0$), and consider the approximate dynamics

$$(9.10) \quad \ddot{x}_\lambda(t) + \nabla\phi_\lambda(\dot{x}_\lambda(t)) + \beta\nabla^2 f(x_\lambda(t))\dot{x}_\lambda(t) + \nabla f(x_\lambda(t)) = 0, \quad t \in [0, +\infty[$$

which uses the Moreau-Yosida approximates (ϕ_λ) of ϕ . We will prove that the filtered sequence (x_λ) converges uniformly as $\lambda \rightarrow 0$ over the bounded time intervals towards a solution of (9.4). The Hamiltonian formulation of (9.10) gives the first-order (in time) system

$$\begin{cases} \dot{x}_\lambda(t) - u_\lambda(t) = 0; \\ \dot{u}_\lambda(t) + \nabla\phi_\lambda(u_\lambda(t)) + \nabla f(x_\lambda(t)) + \beta\nabla^2 f(x_\lambda(t))u_\lambda(t) = 0, \end{cases}$$

with the Cauchy data $x_\lambda(0) = x_0$, $u_\lambda(0) = x_1$. Set $Z_\lambda(t) = (x_\lambda(t), u_\lambda(t)) \in \mathcal{H} \times \mathcal{H}$. The above system can be written equivalently as

$$\dot{Z}_\lambda(t) + F_\lambda(Z_\lambda(t)) \ni 0, \quad Z_\lambda(t_0) = (x_0, x_1),$$

where $F_\lambda : \mathcal{H} \times \mathcal{H} \rightarrow \mathcal{H} \times \mathcal{H}$, $(x, u) \mapsto F_\lambda(x, u)$ is defined by

$$F_\lambda(x, u) = \left(0, \nabla\phi_\lambda(u)\right) + \left(-u, \nabla f(x) + \beta\nabla^2 f(x)u\right).$$

Hence F_λ splits as follows $F_\lambda(x, u) = \nabla\Phi_\lambda(x, u) + G(x, u)$ where Φ and G have been defined in (9.7). Therefore, the approximate equation is equivalent to the first-order differential system with Cauchy data

$$(9.11) \quad \dot{Z}_\lambda(t) + \nabla\Phi_\lambda(Z_\lambda(t)) + G(Z_\lambda(t)) = 0, \quad Z_\lambda(0) = (x_0, x_1).$$

Let's argue with $\lambda > 0$ fixed. According to the Lipschitz continuity of $\nabla\Phi_\lambda$, and the fact that G is Lipschitz continuous on the bounded sets, we have that the sum operator $\nabla\Phi_\lambda + G$ which governs (9.11) is Lipschitz continuous on the bounded sets. As a consequence, the existence of a local solution to (9.11) follows from the classical Cauchy–Lipschitz theorem. To pass from a local solution to a global solution, we use the a priori estimate obtained in Step 1 of the proof. Note that this estimate is valid for any damping potential, in particular for ϕ_λ . According to the Cauchy data, and f minorized, this implies that, on any bounded time interval, the functions (x_λ) and (\dot{x}_λ) are bounded. According to the property (3.3) of the Yosida approximation, and the property (iii) of the damping potential ϕ , this implies that

$$\|\nabla\phi_\lambda(x_\lambda(t))\| \leq \|(\partial\phi)^0(x_\lambda(t))\|$$

is also bounded uniformly with respect to t bounded. Moreover, according to the local boundedness assumption made on the gradient and the Hessian of f , we have that $\nabla f(x_\lambda(t))$ and $\nabla^2 f(x_\lambda(t))\dot{x}_\lambda(t)$ are also bounded. According to the constitutive equation (9.10), this in turn implies that (\dot{x}_λ) is also bounded. This implies that if a maximal solution is defined on a finite time interval $[0, T]$, then the limits of $x_\lambda(t)$ and $\dot{x}_\lambda(t)$ exist, as $t \rightarrow T$. According to this property, passing from a local to a global solution is a classical argument. So for any $\lambda > 0$ we have a unique global solution of (9.10) with satisfies the Cauchy data $x_\lambda(0) = x_0$, $\dot{x}_\lambda(0) = x_1$.

Step 4: *Passing to the limit as $\lambda \rightarrow 0$.* Take $T > 0$, and $\lambda, \mu > 0$. Consider the corresponding solutions on $[0, T]$

$$\begin{aligned} \dot{Z}_\lambda(t) + \nabla\Phi_\lambda(Z_\lambda(t)) + G(Z_\lambda(t)) &= 0, & Z_\lambda(0) &= (x_0, x_1) \\ \dot{Z}_\mu(t) + \nabla\Phi_\mu(Z_\mu(t)) + G(Z_\mu(t)) &= 0, & Z_\mu(0) &= (x_0, x_1). \end{aligned}$$

Let's make the difference between the two equations, and take the scalar product by $Z_\lambda(t) - Z_\mu(t)$. We get

$$(9.12) \quad \begin{aligned} \frac{1}{2} \frac{d}{dt} \|Z_\lambda(t) - Z_\mu(t)\|^2 + \langle \nabla\Phi_\lambda(Z_\lambda(t)) - \nabla\Phi_\mu(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle \\ + \langle G(Z_\lambda(t)) - G(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle = 0. \end{aligned}$$

We now use the following ingredients:

i) According to the general properties of the Yosida approximation (see [49, Theorem 3.1]), we have

$$\langle \nabla \Phi_\lambda(Z_\lambda(t)) - \nabla \Phi_\mu(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle \geq -\frac{\lambda}{4} \|\nabla \Phi_\mu(Z_\mu(t))\|^2 - \frac{\mu}{4} \|\nabla \Phi_\lambda(Z_\lambda(t))\|^2.$$

According to the energy estimates, the sequence (Z_λ) is uniformly bounded on $[0, T]$, let

$$\|Z_\lambda(t)\| \leq C_T.$$

From these properties we immediately infer

$$\|\nabla \Phi_\lambda(Z_\lambda(t))\| \leq \sup_{\|\xi\| \leq C_T} \|(\partial\phi)^0(\xi)\| = M_T < +\infty,$$

because our assumption on ϕ gives that $(\partial\phi)^0$ is bounded on the bounded sets. Therefore

$$\langle \nabla \Phi_\lambda(Z_\lambda(t)) - \nabla \Phi_\mu(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle \geq -\frac{1}{4} M_T (\lambda + \mu).$$

ii) Since the mapping $G : \mathcal{H} \times \mathcal{H} \rightarrow \mathcal{H} \times \mathcal{H}$ is Lipschitz continuous on the bounded sets, and using again that the sequence (Z_λ) is uniformly bounded on $[0, T]$, we deduce that there exists a constant L_T such that

$$\|G(Z_\lambda(t)) - G(Z_\mu(t))\| \leq L_T \|Z_\lambda(t) - Z_\mu(t)\|.$$

Combining the above results, and using Cauchy–Schwarz inequality, we deduce from (9.12) that

$$\frac{1}{2} \frac{d}{dt} \|Z_\lambda(t) - Z_\mu(t)\|^2 \leq \frac{1}{4} M_T (\lambda + \mu) + L_T \|Z_\lambda(t) - Z_\mu(t)\|^2.$$

We now proceed with the integration of this differential inequality. According to the fact that $Z_\lambda(0) - Z_\mu(0) = 0$, elementary calculus gives

$$\|Z_\lambda(t) - Z_\mu(t)\|^2 \leq \frac{M_T}{4L_T} (\lambda + \mu) \left(e^{2L_T(t-t_0)} - 1 \right).$$

Therefore, the filtered sequence (Z_λ) is a Cauchy sequence for the uniform convergence on $[0, T]$, and hence it converges uniformly. This means the uniform convergence on $[0, T]$ of x_λ and \dot{x}_λ to x and \dot{x} respectively. Proving that x is solution of (9.4) is obtained in a similar way as in Theorem 3.3. Just rely on the classical derivation chain rule $\frac{d}{dt} (\nabla f(x_\lambda(t))) = \nabla^2 f(x_\lambda(t)) \dot{x}_\lambda(t)$ to pass to the limit on the Hessian term. \square

9.3. Convergence based on the quasi-gradient approach. Our objective is to address, from the perspective of quasi-gradient systems, the system (ADIGE-VH)

$$(9.13) \quad \ddot{x}(t) + \nabla\phi(\dot{x}(t)) + \beta \nabla^2 f(x(t)) \dot{x}(t) + \nabla f(x(t)) = 0,$$

as it was done in Section 7.3. We assume that $\mathcal{H} = \mathbb{R}^N$ is a finite-dimensional Hilbert space, and that the hypotheses of Theorem 7.3 and Theorem 9.1 hold. We follow the steps of the proof of Theorem 7.3. By using the estimates in Step 1 of the proof of Theorem 9.1, we easily derive the first part of the proof of Theorem 7.3, namely that the trajectory $t \mapsto (x(t), \dot{x}(t))$ in the phase space

$\mathbb{R}^N \times \mathbb{R}^N$ belongs to the closed bounded set $\Gamma = \bar{B}(0, R) \times \bar{B}(0, \epsilon)$. According to Step 2 in the proof of Theorem 9.1, the Hamiltonian formulation of (9.13) gives the first-order differential system

$$(9.14) \quad \dot{z}(t) + F(z(t)) = 0,$$

where $z(t) = (x(t), \dot{x}(t)) \in \mathbb{R}^N \times \mathbb{R}^N$, and $F : \mathbb{R}^N \times \mathbb{R}^N \rightarrow \mathbb{R}^N \times \mathbb{R}^N$ is defined by

$$F(x, u) = (-u, \nabla\phi(u) + \nabla f(x) + \beta\nabla^2 f(x)u).$$

Let's focus on the key point which is the angle condition (E_λ is defined as in Theorem 7.3). We have

$$\langle \nabla E_\lambda(x, u), F(x, u) \rangle = \left\langle \left(\nabla f(x) + \lambda\nabla^2 f(x)u, u + \lambda\nabla f(x) \right), \left(-u, \nabla\phi(u) + \nabla f(x) + \beta\nabla^2 f(x)u \right) \right\rangle.$$

After development and simplification, we get

$$\begin{aligned} \langle \nabla E_\lambda(x, u), F(x, u) \rangle &= -\lambda \langle \nabla^2 f(x)u, u \rangle + \langle u, \nabla\phi(u) \rangle + \lambda \langle \nabla f(x), \nabla\phi(u) \rangle + \lambda \|\nabla f(x)\|^2 \\ &\quad + \beta \langle u + \lambda\nabla f(x), \nabla^2 f(x)u \rangle \\ &\geq -\lambda \langle \nabla^2 f(x)u, u \rangle + \langle u, \nabla\phi(u) \rangle + \lambda \langle \nabla f(x), \nabla\phi(u) \rangle + \lambda \|\nabla f(x)\|^2 \\ &\quad + \lambda\beta \langle \nabla f(x), \nabla^2 f(x)u \rangle, \end{aligned}$$

where we used that $\nabla^2 f(x)$ is positive semidefinite. The only difference with respect to the next step in the proof of Theorem 7.3 is that we need to estimate the extra term $\lambda\beta \langle \nabla f(x), \nabla^2 f(x)u \rangle$. We do this by writing

$$\lambda\beta \langle \nabla f(x), \nabla^2 f(x)u \rangle \geq -\frac{\lambda}{4} \|\nabla f(x)\|^2 - \lambda\beta^2 M^2 \|u\|^2$$

and get

$$(9.15) \quad \langle \nabla E_\lambda(x, u), F(x, u) \rangle \geq \left(\gamma - \lambda M - \frac{\lambda}{2} \delta^2 - \lambda\beta^2 M^2 \right) \|u\|^2 + \frac{\lambda}{4} \|\nabla f(x)\|^2.$$

Take λ small enough to satisfy $\gamma > \lambda \left(M + \frac{\delta^2}{2} + \beta^2 M^2 \right)$. Then

$$(9.16) \quad \langle \nabla E_\lambda(x, u), F(x, u) \rangle \geq \alpha_0 (\|u\|^2 + \|\nabla f(x)\|^2),$$

with $\alpha_0 := \min\{\gamma - \lambda \left(M + \frac{\delta^2}{2} + \beta^2 M^2 \right), \frac{\lambda}{4}\}$. On the other hand, as in Theorem 7.3,

$$\begin{aligned} \|\nabla E_\lambda(x, u)\| &\leq C_1 (\|u\|^2 + \|\nabla f(x)\|^2)^{\frac{1}{2}} \\ \|F(x, u)\| &\leq C_2 (\|u\|^2 + \|\nabla f(x)\|^2)^{\frac{1}{2}}, \end{aligned}$$

where $C_2 = \sqrt{4 + 3\delta^2 + 3\beta^2 M^2}$. Therefore

$$(9.17) \quad \|\nabla E_\lambda(x, u)\| \|F(x, u)\| \leq C_1 C_2 (\|u\|^2 + \|\nabla f(x)\|^2).$$

Therefore, for $\alpha := \frac{\alpha_0}{C_1 C_2}$, the angle condition $\langle \nabla E(z), F(z) \rangle \geq \alpha \|\nabla E(z)\| \|F(z)\|$ is satisfied on Γ . let us summarize the above results.

THEOREM 9.2. *Let $f : \mathcal{H} \rightarrow \mathbb{R}$ be a convex function which is twice continuously differentiable, and such that $\inf_{\mathcal{H}} f > -\infty$. We suppose that*

- (i) ∇f is Lipschitz continuous on the bounded subsets of \mathcal{H} ;
- (ii) $\nabla^2 f$ is Lipschitz continuous on the bounded subsets of \mathcal{H} .

Let $E_\lambda : \mathbb{R}^N \times \mathbb{R}^N \rightarrow \mathbb{R}$ be defined by: for all $(x, u) \in \mathbb{R}^N \times \mathbb{R}^N$

$$E_\lambda(x, u) := \frac{1}{2}\|u\|^2 + f(x) + \lambda \langle \nabla f(x), u \rangle.$$

Suppose that the function E_λ satisfies the (KL) property.

Let $\phi : \mathbb{R}^N \rightarrow \mathbb{R}_+$ be a damping potential (see Definition 3.1) which is differentiable, and that satisfies the following growth conditions:

- (i) (local) there exists positive constants γ, δ , and $\epsilon > 0$ such that, for all u in \mathbb{R}^N with $\|u\| \leq \epsilon$

$$\phi(u) \geq \gamma\|u\|^2 \text{ and } \|\nabla\phi(u)\| \leq \delta\|u\|.$$

- (ii) (global) there exists $p \geq 1, c > 0$, such that for all u in \mathbb{R}^N , $\phi(u) \geq c\|u\|^p$.

Let $x : [0, +\infty[\rightarrow \mathbb{R}^N$ be a bounded solution trajectory of

$$\ddot{x}(t) + \nabla\phi(\dot{x}(t)) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) = 0.$$

Then, the following properties are satisfied:

- (i) $x(t) \rightarrow x_\infty$ as $t \rightarrow +\infty$, where $x_\infty \in \text{crit } f$;
- (ii) $\dot{x} \in L^1(0, +\infty; \mathbb{R}^N)$, $\dot{x}(t) \rightarrow 0$ as $t \rightarrow +\infty$;
- (iii) For λ sufficiently small, and t sufficiently large

$$\|x(t) - x_\infty\| \leq \frac{1}{\alpha} \theta \left(E_\lambda(x(t), u(t)) - E_\lambda(x_\infty, 0) \right)$$

where θ is the desingularizing function for E_λ at $(x_\infty, 0)$, and α enters the angle condition.

9.4. Numerical illustrations. We revisit the numerical examples of section 5.3 where we introduce an additional Hessian damping.

So, we take $\mathcal{H} = \mathbb{R}$, $f(x) = \frac{1}{2}|x|^2$, and $\phi(u) = \frac{1}{p}|u|^p$ with $p > 1$. Then, (ADIGE-VH) writes

$$(9.18) \quad \ddot{x}(t) + |\dot{x}(t)|^{p-2}\dot{x}(t) + \beta\dot{x}(t) + x(t) = 0.$$

For $\beta > 0$, we are in the framework of Theorem 5.1, with $\phi(u) = \frac{\beta}{2}|u|^2 + \frac{p}{2}|u|^p$. So we have convergence at an exponential rate of $x(t)$ and $\dot{x}(t)$ towards zero. This makes a big contrast with the case $\beta = 0$, for which we have convergence towards zero, but with many oscillations in the case of weak damping (p large). Note that that even for very small $\beta > 0$, we have a rapid stabilization of the trajectory towards the origin. On the other hand, taking large β is not beneficial, we can observe on Figure 9.1 that the quality of convergence is degraded in this case. Indeed, since the damping attached to $|\dot{x}(t)|^{p-2}\dot{x}(t)$ is negligible for large p with respect to the damping attached to $\beta\dot{x}(t)$, the "optimal" value of β is close to the optimal value for (HBF). So, according to Theorem 2.2, it is close to $2\sqrt{\mu}$ where μ is the coefficient of strong convexity of f (see Theorem 2.2). In our situation, this gives $\beta \sim 2$.

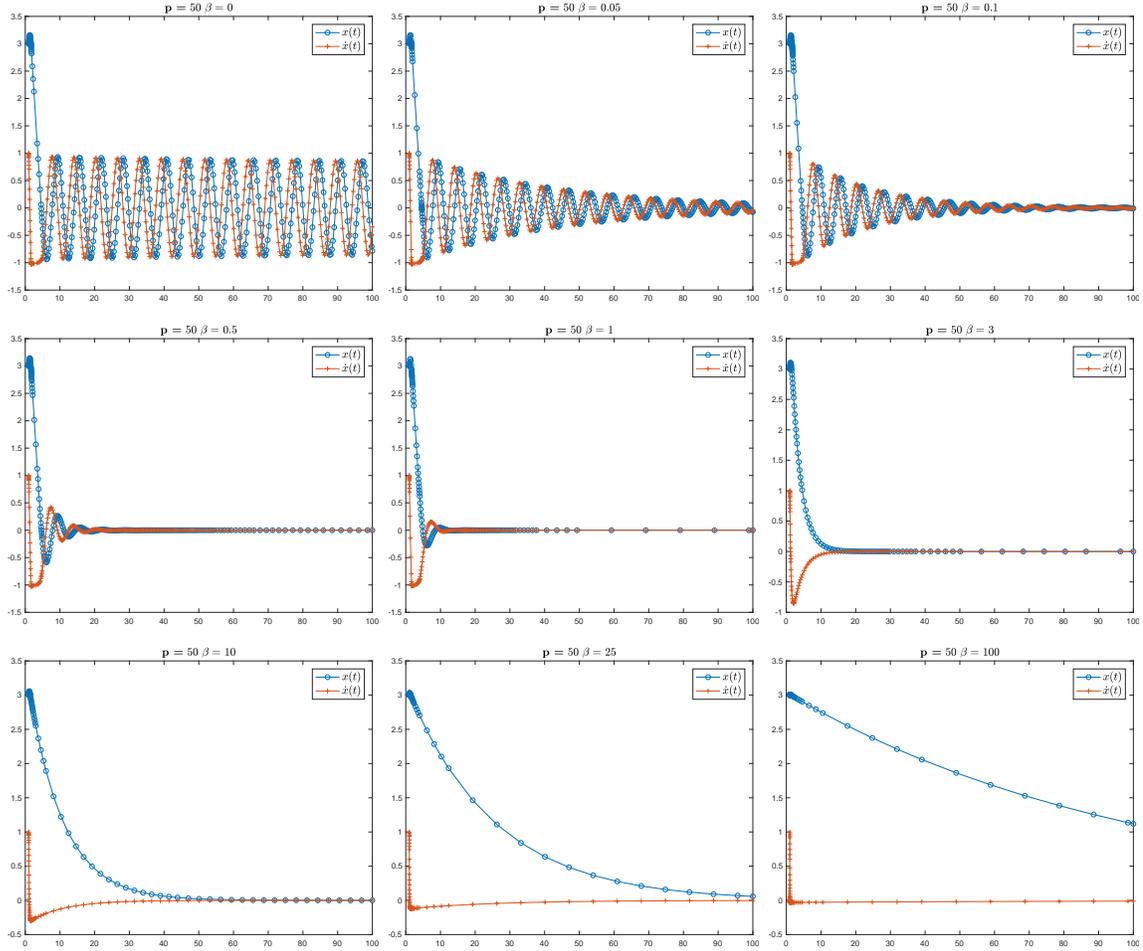


Fig. 9.1: Evolution of the trajectories $x(t)$ (blue) and $\dot{x}(t)$ (red line) of (9.18) for different values of β .

9.5. Link with the regularized Newton method. Let us specify the link between our study and Newton's method for solving $Ax \ni 0$, where A is a maximally monotone operator (for convex minimization take $A = \partial f$). To overcome the ill-posed character of the continuous Newton method, the following first-order evolution system was studied by Attouch-Svaiter [35], for a general maximally monotone operator A

$$\begin{cases} v(t) \in A(x(t)) \\ \gamma(t)\dot{x}(t) + \beta\dot{v}(t) + v(t) = 0. \end{cases}$$

The system can be considered as a continuous version of the Levenberg-Marquardt, which acts as a regularization of the Newton method. Under a fairly general assumption on the regularization parameter $\gamma(t)$, this system is well-posed and generates trajectories that converge weakly to equilib-

ria. Parallel results have been obtained for the associated proximal algorithms obtained by implicit temporal discretization, see [1], [30], [34]. Formally when A is differentiable, this system writes as

$$\gamma(t)\dot{x}(t) + \beta \frac{d}{dt} (A(x(t))) + A(x(t)) = 0.$$

When $A = \nabla f$ we obtain

$$(9.19) \quad \gamma(t)\dot{x}(t) + \beta \nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) = 0.$$

The system (ADIGE-VH) considered in the previous section can be seen as an inertial version of the above system (9.19). Most interesting, Attouch-Redont-Svaiter developed in [34] a closed-loop version of the above results. They showed the convergence of the trajectories generated by the closed-loop control system when $0 < p < 1$, where A is a general maximally monotone operator:

$$\left\{ \begin{array}{l} v(t) \in A(x(t)) \\ \|v(t)\|^p \dot{x}(t) + \dot{v}(t) + v(t) = 0 \\ x(0) = x_0, v(0) \in A(x_0), v_0 \neq 0. \end{array} \right.$$

For optimization problems, this naturally suggests to consider autonomous inertial systems where the damping coefficient is a closed-loop control of the gradient of f . A first answer to this question has been obtained by Lin-Jordan [71] who considered the autonomous system

$$(9.20) \quad \ddot{x}(t) + \gamma(t)\dot{x}(t) + \beta(t)\nabla^2 f(x(t))\dot{x}(t) + b(t)\nabla f(x(t)) = 0,$$

where γ , β and b are defined by the following formulas:

$$(9.21) \quad \left\{ \begin{array}{l} |\lambda(t)|^p \|\nabla f(x(t))\|^{p-1} = \theta \\ a(t) = \frac{1}{4} \left(\int_0^t \sqrt{\lambda(s)} ds + c \right)^2 \\ \gamma(t) = 2 \frac{\dot{a}(t)}{a(t)} - \frac{\ddot{a}(t)}{\dot{a}(t)} \\ \beta(t) = \left(\frac{\dot{a}(t)}{a(t)} \right)^2 \\ b(t) = \frac{\dot{a}(t)(\dot{a}(t) + \ddot{a}(t))}{a(t)} \end{array} \right.$$

As a specific feature, the damping coefficients are expressed with the help of $\lambda(t)$ which is equal to a power of the inverse of the norm of the gradient of f . The authors give some interesting non-trivial convergence rates for values. According to the presence of the Hessian driven damping term, they show the fast convergence towards zero of the gradient norms.

10. Closed loop damping involving the velocity and the gradient. Let's consider the following system, where the damping term $\partial\phi(\dot{x}(t) + \beta\nabla f(x(t)))$ involves both the velocity vector and the gradient of the potential function f

$$(10.1) \quad (\text{ADIGE-VGH}) \quad \ddot{x}(t) + \partial\phi(\dot{x}(t) + \beta\nabla f(x(t))) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) \ni 0.$$

The parameter $\beta \geq 0$ is attached to the geometric damping induced by the Hessian. As previously considered, ϕ is a damping potential function. The suffix V,G,H make respectively reference to the Velocity, the Gradient of f , and the Hessian of f , which enter the damping terms of the above dynamic. This model makes it possible to encompass several situations.

- When $\beta = 0$, we recover the closed loop controlled system

$$(10.2) \quad \ddot{x}(t) + \partial\phi(\dot{x}(t)) + \nabla f(x(t)) = 0,$$

studied from Section 3 to 7. So studying (10.1) can be viewed as an extension of our previous study. Still, we will see that taking $\beta > 0$ induces several favorable properties.

- When $\phi(u) = \frac{\gamma}{2}\|u\|^2$, we obtain the system

$$(10.3) \quad \ddot{x}(t) + \gamma\dot{x}(t) + \beta\nabla^2 f(x(t))\dot{x}(t) + (1 + \gamma\beta)\nabla f(x(t)) = 0,$$

studied in Section 9, and which was introduced by Alvarez-Attouch-Bolte-Redont in [9].

10.1. Existence and uniqueness results. A key property for studying (10.1) is the following equivalent formulation, different from the Hamiltonian formulation, and whose proof is immediate. Just introduce the new variable $u(t) := \dot{x}(t) + \beta\nabla f(x(t))$.

PROPOSITION 10.1. *The following are equivalent*

$$(i) \quad \ddot{x}(t) + \partial\phi(\dot{x}(t) + \beta\nabla f(x(t))) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) \ni 0.$$

$$(ii) \quad \begin{cases} \dot{x}(t) + \beta\nabla f(x(t)) - u(t) = 0 \\ \dot{u}(t) + \partial\phi(u(t)) + \nabla f(x(t)) \ni 0. \end{cases}$$

A major interest of the the formulation (ii) is that it is a first-order system in time and space (without occurrence of the Hessian). As such, it requires fewer regularity assumptions on f than in Theorem 9.1.

THEOREM 10.2. *Let $f : \mathcal{H} \rightarrow \mathbb{R}$ be a convex function which is twice continuously differentiable, and such that $\inf_{\mathcal{H}} f > -\infty$. Suppose that ∇f is Lipschitz continuous on the bounded subsets of \mathcal{H} . Let $\phi : \mathcal{H} \rightarrow \mathbb{R}$ be a convex continuous damping function. Then, for any Cauchy data $(x_0, x_1) \in \mathcal{H} \times \mathcal{H}$, there exists a unique strong global solution $x : [0, +\infty[\rightarrow \mathcal{H}$ of (ADIGE-VGH) satisfying $x(0) = x_0$, and $\dot{x}(0) = x_1$.*

Proof. The structure of the proof being similar to Theorem 9.1, we only develop the original aspects.

Step 1: *A priori estimate.* Note that (10.1) can be equivalently written as

$$(10.4) \quad \frac{d}{dt}(\dot{x}(t) + \beta\nabla f(x(t))) + \partial\phi(\dot{x}(t) + \beta\nabla f(x(t))) + \nabla f(x(t)) \ni 0.$$

After taking the scalar product of (10.4) with $\dot{x}(t) + \beta\nabla f(x(t))$, we get

$$(10.5) \quad \begin{aligned} & \frac{1}{2} \frac{d}{dt} \|\dot{x}(t) + \beta\nabla f(x(t))\|^2 + \langle \partial\phi(\dot{x}(t) + \beta\nabla f(x(t))), \dot{x}(t) + \beta\nabla f(x(t)) \rangle \\ & + \langle \nabla f(x(t)), \dot{x}(t) + \beta\nabla f(x(t)) \rangle = 0. \end{aligned}$$

Since ϕ is a damping potential, the subdifferential inequality for convex functions gives

$$\langle \partial\phi(\dot{x}(t) + \beta\nabla f(x(t))), \dot{x}(t) + \beta\nabla f(x(t)) \rangle \geq \phi(\dot{x}(t) + \beta\nabla f(x(t))).$$

Collecting the above results, we obtain

$$(10.6) \quad \frac{d}{dt} \left(\frac{1}{2} \|\dot{x}(t) + \beta\nabla f(x(t))\|^2 + f(x(t)) - \inf_{\mathcal{H}} f \right) + \phi(\dot{x}(t) + \beta\nabla f(x(t))) + \beta \|\nabla f(x(t))\|^2 \leq 0.$$

Therefore, the energy-like function

$$(10.7) \quad t \mapsto \frac{1}{2} \|\dot{x}(t) + \beta\nabla f(x(t))\|^2 + f(x(t)) - \inf_{\mathcal{H}} f \quad \text{is nonincreasing.}$$

This implies that, as long as the trajectory is defined

$$(10.8) \quad \|\dot{x}(t) + \beta\nabla f(x(t))\|^2 \leq C := \|x_1 + \beta\nabla f(x_0)\|^2 + 2(f(x_0) - \inf_{\mathcal{H}} f).$$

From this, we will obtain a bound on the trajectory. We have

$$\dot{x}(t) + \beta\nabla f(x(t)) = k(t)$$

with $\|k(t)\| \leq \sqrt{C}$. Take the scalar product of the above equation with $x(t) - x_0$.

$$\frac{1}{2} \frac{d}{dt} \|x(t) - x_0\|^2 + \beta \langle \nabla f(x(t)) - \nabla f(x_0), x(t) - x_0 \rangle + \beta \langle \nabla f(x_0), x(t) - x_0 \rangle = \langle k(t), x(t) - x_0 \rangle.$$

According to the convexity of f , and hence the monotonicity of ∇f , and by Cauchy–Schwarz inequality

$$(10.9) \quad \frac{1}{2} \frac{d}{dt} \|x(t) - x_0\|^2 \leq (\|k(t)\| + \beta \|\nabla f(x_0)\|) \|x(t) - x_0\|.$$

According to the Gronwall inequality, and $\|k(t)\| \leq \sqrt{C}$, we obtain

$$(10.10) \quad \|x(t) - x_0\| \leq t \left(\|x_1 + \beta\nabla f(x_0)\| + \sqrt{2(f(x_0) - \inf_{\mathcal{H}} f) + \beta \|\nabla f(x_0)\|^2} \right).$$

Step 2: *first-order formulation of (10.1).* According to Proposition 10.1, it is equivalent to solve the first-order system

$$\begin{cases} \dot{x}(t) + \beta\nabla f(x(t)) - u(t) = 0 \\ \dot{u}(t) + \partial\phi(u(t)) + \nabla f(x(t)) \ni 0, \end{cases}$$

with the Cauchy data $x(0) = x_0$, $u(0) = x_1$. Set $Z(t) = (x(t), u(t)) \in \mathcal{H} \times \mathcal{H}$.

The above system can be written equivalently as

$$\dot{Z}(t) + F(Z(t)) \ni 0, \quad Z(0) = (x_0, x_1),$$

where $F : \mathcal{H} \times \mathcal{H} \rightrightarrows \mathcal{H} \times \mathcal{H}$, $(x, u) \mapsto F(x, u)$ is defined by

$$F(x, u) = \left(0, \partial\phi(u) \right) + \left(\beta\nabla f(x) - u, \nabla f(x) \right).$$

Hence F splits as follows

$$F(x, u) = \partial\Phi(x, u) + G(x, u),$$

where

$$(10.11) \quad \Phi(x, u) = \phi(u) \text{ and } G(x, u) = \left(\beta \nabla f(x) - u, \nabla f(x) \right).$$

Therefore, it is equivalent to solve the following first-order differential inclusion with Cauchy data

$$(10.12) \quad \dot{Z}(t) + \partial\Phi(Z(t)) + G(Z(t)) \ni 0, \quad Z(0) = (x_0, x_1).$$

According to the local Lipschitz assumption on the gradient of f , we immediately obtain that the mapping $(x, u) \mapsto G(x, u)$ is Lipschitz continuous on the bounded subsets of $\mathcal{H} \times \mathcal{H}$.

Step 3: Approximate dynamics. We consider the approximate dynamics

$$(10.13) \quad \ddot{x}_\lambda(t) + \nabla\phi_\lambda(\dot{x}_\lambda(t) + \beta\nabla f(x_\lambda(t))) + \beta\nabla^2 f(x(t))\dot{x}_\lambda(t) + \nabla f(x_\lambda(t)) = 0, \quad t \in [0, +\infty[$$

which uses the Moreau-Yosida approximates (ϕ_λ) of ϕ . We will prove that the filtered sequence (x_λ) converges uniformly as $\lambda \rightarrow 0$ over the bounded time intervals towards a solution of (10.1). The first-order formulation of (10.13) gives the following system

$$\begin{cases} \dot{x}_\lambda(t) + \beta\nabla f(x_\lambda(t)) - u_\lambda(t) = 0; \\ \dot{u}_\lambda(t) + \nabla\phi_\lambda(u_\lambda(t)) + \nabla f(x_\lambda(t)) = 0, \end{cases}$$

with the Cauchy data $x_\lambda(0) = x_0$, $u_\lambda(0) = x_1$. Set $Z_\lambda(t) = (x_\lambda(t), u_\lambda(t)) \in \mathcal{H} \times \mathcal{H}$. The above system can be written equivalently as

$$\dot{Z}_\lambda(t) + F_\lambda(Z_\lambda(t)) \ni 0, \quad Z_\lambda(t_0) = (x_0, x_1),$$

where $F_\lambda : \mathcal{H} \times \mathcal{H} \rightarrow \mathcal{H} \times \mathcal{H}$, $(x, u) \mapsto F_\lambda(x, u)$ is defined by

$$F_\lambda(x, u) = \left(0, \nabla\phi_\lambda(u) \right) + \left(\beta\nabla f(x) - u, \nabla f(x) \right).$$

Hence F_λ splits as follows $F_\lambda(x, u) = \nabla\Phi_\lambda(x, u) + G(x, u)$ where Φ and G have been defined in (10.11). Therefore, the approximate equation is equivalent to the first-order differential system with Cauchy data

$$(10.14) \quad \dot{Z}_\lambda(t) + \nabla\Phi_\lambda(Z_\lambda(t)) + G(Z_\lambda(t)) = 0, \quad Z_\lambda(0) = (x_0, x_1).$$

Let's argue with $\lambda > 0$ fixed. According to the Lipschitz continuity of $\nabla\Phi_\lambda$, and the fact that G is Lipschitz continuous on the bounded sets, we have that the sum operator $\nabla\Phi_\lambda + G$ which governs (10.14) is Lipschitz continuous on the bounded sets. As a consequence, the existence of a local solution to (10.14) follows from the classical Cauchy–Lipschitz theorem. To pass from a local solution to a global solution, we use the a priori estimates (10.8) and (10.10) obtained in Step 1 of the proof. Note that these estimates are valid for any damping potential, in particular for ϕ_λ . Suppose that a maximal solution is defined on a finite time interval $[0, T[$. According to (10.10) we first obtain that $x_\lambda(t)$ remains bounded on $[0, T[$. Then, using (10.8) and the fact that the gradient of f is Lipschitz continuous on the bounded sets, we obtain that $\dot{x}_\lambda(t)$ is also bounded on $[0, T[$.

According to the property (3.3) of the Yosida approximation, the property (iii) of the damping potential ϕ , and (10.8), this implies that

$$\|\nabla\phi_\lambda(\dot{x}_\lambda(t) + \beta\nabla f(x_\lambda(t)))\| \leq \|(\partial\phi)^0(\dot{x}_\lambda(t) + \beta\nabla f(x_\lambda(t)))\|$$

is also bounded by on $[0, T[$. Moreover, according to the local boundedness assumption made on the gradient, and the boundedness of $x_\lambda(t)$ and $\dot{x}_\lambda(t)$, we have that $\nabla^2 f(x_\lambda(t))\dot{x}_\lambda(t)$ is also bounded. According to the constitutive equation (9.10), this in turn implies that (\ddot{x}_λ) is also bounded. This implies that the limits of $x_\lambda(t)$ and $\dot{x}_\lambda(t)$ exist, as $t \rightarrow T$. According to this property, passing from a local to a global solution is a classical argument. So, for any $\lambda > 0$ we have a unique global solution of (9.10) with satisfies the Cauchy data $x_\lambda(0) = x_0$, $\dot{x}_\lambda(0) = x_1$.

Step 4: *Passing to the limit as $\lambda \rightarrow 0$.* Take $T > 0$, and $\lambda, \mu > 0$. Consider the corresponding solutions on $[0, T]$

$$\begin{aligned} \dot{Z}_\lambda(t) + \nabla\Phi_\lambda(Z_\lambda(t)) + G(Z_\lambda(t)) &= 0, & Z_\lambda(0) &= (x_0, x_1) \\ \dot{Z}_\mu(t) + \nabla\Phi_\mu(Z_\mu(t)) + G(Z_\mu(t)) &= 0, & Z_\mu(0) &= (x_0, x_1). \end{aligned}$$

Let's make the difference between the two equations, and take the scalar product by $Z_\lambda(t) - Z_\mu(t)$. We get

$$(10.15) \quad \begin{aligned} \frac{1}{2} \frac{d}{dt} \|Z_\lambda(t) - Z_\mu(t)\|^2 + \langle \nabla\Phi_\lambda(Z_\lambda(t)) - \nabla\Phi_\mu(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle \\ + \langle G(Z_\lambda(t)) - G(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle = 0. \end{aligned}$$

We now use the following ingredients:

i) According to the general properties of the Yosida approximation (see [49, Theorem 3.1]), we have

$$\langle \nabla\Phi_\lambda(Z_\lambda(t)) - \nabla\Phi_\mu(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle \geq -\frac{\lambda}{4} \|\nabla\Phi_\mu(Z_\mu(t))\|^2 - \frac{\mu}{4} \|\nabla\Phi_\lambda(Z_\lambda(t))\|^2.$$

According to the energy estimates, the sequence (Z_λ) is uniformly bounded on $[0, T]$, let

$$\|Z_\lambda(t)\| \leq C_T.$$

From these properties we immediately infer

$$\|\nabla\Phi_\lambda(Z_\lambda(t))\| \leq \sup_{\|\xi\| \leq C_T} \|(\partial\phi)^0(\xi)\| = M_T < +\infty$$

because our assumption on ϕ gives that $(\partial\phi)^0$ is bounded on the bounded sets. Therefore

$$\langle \nabla\Phi_\lambda(Z_\lambda(t)) - \nabla\Phi_\mu(Z_\mu(t)), Z_\lambda(t) - Z_\mu(t) \rangle \geq -\frac{1}{4} M_T (\lambda + \mu).$$

ii) Since the mapping $G : \mathcal{H} \times \mathcal{H} \rightarrow \mathcal{H} \times \mathcal{H}$ is Lipschitz continuous on the bounded sets, and using again that the sequence (Z_λ) is uniformly bounded on $[0, T]$, we deduce that there exists a constant L_T such that

$$\|G(Z_\lambda(t)) - G(Z_\mu(t))\| \leq L_T \|Z_\lambda(t) - Z_\mu(t)\|.$$

Combining the above results, and using Cauchy–Schwarz inequality, we deduce from (9.12) that

$$\frac{1}{2} \frac{d}{dt} \|Z_\lambda(t) - Z_\mu(t)\|^2 \leq \frac{1}{4} M_T (\lambda + \mu) + L_T \|Z_\lambda(t) - Z_\mu(t)\|^2.$$

We now proceed with the integration of this differential inequality. According to the fact that $Z_\lambda(0) - Z_\mu(0) = 0$, elementary calculus gives

$$\|Z_\lambda(t) - Z_\mu(t)\|^2 \leq \frac{M_T}{4L_T} (\lambda + \mu) \left(e^{2L_T(t-t_0)} - 1 \right).$$

Therefore, the filtered sequence (Z_λ) is a Cauchy sequence for the uniform convergence on $[0, T]$, and hence it converges uniformly. This means the uniform convergence on $[0, T]$ of x_λ and \dot{x}_λ to x and \dot{x} respectively. Proving that x is solution of (10.1) is obtained in a similar way as in Theorem 3.3. Just rely on the property $\frac{d}{dt} (\nabla f(x_\lambda(t))) = \nabla^2 f(x_\lambda(t)) \dot{x}_\lambda(t)$ to pass to the limit on the Hessian term. \square

10.2. Convergence properties. We have the following convergence properties for the solution trajectories of the system (10.1) with closed loop damping involving both the velocity and the gradient.

THEOREM 10.3. *Let $f : \mathcal{H} \rightarrow \mathbb{R}$ be a convex function which is twice continuously differentiable, and such that $\operatorname{argmin}_{\mathcal{H}} f \neq \emptyset$. We suppose that ∇f is Lipschitz continuous on the bounded subsets of \mathcal{H} . Suppose that $\beta > 0$. Let $\phi : \mathcal{H} \rightarrow \mathbb{R}$ be a convex continuous damping function. Then, for any solution trajectory $x : [0, +\infty[\rightarrow \mathcal{H}$ of (ADIGE-VGH) we have*

- (i) *The energy-like function $t \mapsto \frac{1}{2} \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + f(x(t))$ is nonincreasing;*
- (ii) $\int_0^{+\infty} \phi(\dot{x}(t) + \beta \nabla f(x(t))) dt < +\infty$;
- (iii) $\int_0^{+\infty} \|\nabla f(x(t))\|^2 dt < +\infty$

Suppose moreover that there exists $r > 0$ such that for all $u \in \mathcal{H}$ $\phi(u) \geq r\|u\|$. Then the following properties are satisfied:

- a) *The trajectory $x(\cdot)$ converges weakly as $t \rightarrow +\infty$, and its limit belongs to $\operatorname{argmin}_{\mathcal{H}} f$.*
- b) *$\dot{x}(t)$ and $\nabla f(x(t))$ converge strongly to zero as $t \rightarrow +\infty$.*

Proof. Items (i) to (iii) are direct consequences of the estimate (10.6) established in the Step 1 of the proof of Theorem 10.2.

Let's now make the additional assumption $\phi(u) \geq r\|u\|$. According to item (ii), we obtain

$$\int_0^{+\infty} \|\dot{x}(t) + \beta \nabla f(x(t))\| dt \leq \frac{1}{r} \int_0^{+\infty} \phi(\dot{x}(t) + \beta \nabla f(x(t))) dt < +\infty.$$

Therefore, $x(\cdot)$ is solution of the non-autonomous steepest descent equation

$$\dot{x}(t) + \beta \nabla f(x(t)) = k(t)$$

with $k \in L^1(0, +\infty; \mathcal{H})$. We can apply Theorem 3.11 of [49], which gives the convergence of the trajectory to a point in $\operatorname{argmin}_{\mathcal{H}} f$. In particular, the trajectory remains bounded. According to item (i), we get that $\dot{x}(t)$ is also bounded. Returning to the constitutive equation (10.1), we deduce that the acceleration $\ddot{x}(t)$ is also bounded. This implies that $\xi(t) = \dot{x}(t) + \beta \nabla f(x(t))$ satisfies

$$\int_0^{+\infty} \|\xi(t)\| dt < +\infty \text{ and } \|\dot{\xi}(t)\| \leq M$$

for some $M > 0$. This classically implies that $\xi(t) = \dot{x}(t) + \beta \nabla f(x(t))$ tends to zero as $t \rightarrow +\infty$. According to item (iii), the same argument applied to $\nabla f(x(t))$ gives that $\nabla f(x(t))$ tends to zero as $t \rightarrow +\infty$. As a difference of the two previous quantities, we conclude that $\dot{x}(t)$ tends to zero as $t \rightarrow +\infty$. \square

Indeed, Theorem 3.11 of [49] was proved under the additional assumption that f is inf-compact. Recent progress based on Opial lemma [75] and Bruck theorem [50] allows to extend it to general convex function f , without making this additional assumption. This is made precise below.

PROPOSITION 10.4. *Let $f : \mathcal{H} \rightarrow \mathbb{R} \cup \{+\infty\}$ be a convex lower semicontinuous proper function such that $\operatorname{argmin}_{\mathcal{H}} f \neq \emptyset$, and let $k \in L^1(0, +\infty; \mathcal{H})$. Suppose that $x : [0, +\infty[\rightarrow \mathcal{H}$ is a strong global solution trajectory of*

$$\dot{x}(t) + \partial f(x(t)) \ni k(t).$$

Then, the trajectory $x(\cdot)$ converges weakly as $t \rightarrow +\infty$, and its limit belongs to $\operatorname{argmin}_{\mathcal{H}} f$.

Proof. Take $\epsilon > 0$. Since $k \in L^1(0, +\infty; \mathcal{H})$, there exists $T_\epsilon > 0$ such that $\int_{T_\epsilon}^{+\infty} \|k(t)\| dt < \epsilon$. Let's consider the solution $v : [0, +\infty[\rightarrow \mathcal{H}$ of

$$\dot{v}(t) + \nabla f(v(t)) \ni 0; \quad v(0) = x(T_\epsilon).$$

According to the semigroup of contractions property, we have, for all $t \geq T_\epsilon$

$$(10.16) \quad \|x(t) - v(t - T_\epsilon)\| \leq \|x(T_\epsilon) - v(0)\| + \int_{T_\epsilon}^t \|k(t)\| dt \leq \epsilon.$$

Take $\xi \in \mathcal{H}$. By Cauchy–Schwarz inequality, we have

$$|\langle x(t) - v(t - T_\epsilon), \xi \rangle| \leq \epsilon \|\xi\|.$$

By the triangle inequality, we deduce that, for all $t \geq T_\epsilon, t' \geq T_\epsilon$

$$|\langle x(t), \xi \rangle - \langle x(t'), \xi \rangle| \leq |\langle v(t - T_\epsilon) - v(t' - T_\epsilon), \xi \rangle| + 2\epsilon \|\xi\|.$$

According to the Bruck theorem, we know that the weak limit of $v(t)$ exists. Passing to the limsup on the above inequality we get

$$\limsup_{t, t' \rightarrow +\infty} |\langle x(t), \xi \rangle - \langle x(t'), \xi \rangle| \leq \limsup_{t, t' \rightarrow +\infty} |\langle v(t - T_\epsilon) - v(t' - T_\epsilon), \xi \rangle| + 2\epsilon \|\xi\| \leq 2\epsilon \|\xi\|.$$

This being true for any $\epsilon > 0$, we deduce that the limit of $\langle x(t), \xi \rangle$ exists, which implies that the weak limit of $x(t)$ exists as $t \rightarrow +\infty$, let x_∞ its limit. Passing to the lower limit on (10.16), according to the lower semicontinuity of the norm for the weak topology, we deduce that

$$(10.17) \quad \|x_\infty - \lim_{t \rightarrow +\infty} v(t)\| \leq \epsilon.$$

Since the weak limit of $v(t)$ belongs to $\operatorname{argmin}_{\mathcal{H}} f$, we deduce that $\operatorname{dist}(x_\infty, \operatorname{argmin}_{\mathcal{H}} f) \leq \epsilon$. This being true for any $\epsilon > 0$, and since $\operatorname{argmin}_{\mathcal{H}} f$ is closed, we finally get that $x_\infty \in \operatorname{argmin}_{\mathcal{H}} f$. \square

10.3. An approach based on Opial's lemma. Here we will prove the weak convergence of the trajectory x to a minimizer of f , based on the continuous version of the Opial Lemma [75]. As in the proof of Theorem 10.3, items (i) to (iii) hold. Assume $\phi(u) \geq r\|u\|$ for all $u \in \mathcal{H}$. According to item (ii) we obtain

$$\int_0^{+\infty} \|\dot{x}(t) + \beta \nabla f(x(t))\| dt < +\infty.$$

Equivalently, we have

$$\dot{x}(t) + \beta \nabla f(x(t)) = k(t)$$

with $k \in L^1(0, +\infty; \mathcal{H})$. Let us prove that x is bounded. Relying on step 1 of the proof of Theorem 10.2, notice that (10.9) holds for a generic $x_0 \in \mathcal{H}$. Taking an arbitrary $z \in \operatorname{argmin} f$, we derive from (10.9)

$$(10.18) \quad \frac{1}{2} \frac{d}{dt} \|x(t) - z\|^2 \leq \|k(t)\| \cdot \|x(t) - z\|.$$

Integrating we obtain

$$(10.19) \quad \frac{1}{2} \|x(T) - z\|^2 \leq \frac{1}{2} \|x_0 - z\|^2 + \int_0^T \|k(t)\| \cdot \|x(t) - z\| dt \quad \forall T \geq 0.$$

Now apply [49, Lemme A.5, pag 157] to conclude

$$\|x(T) - z\| \leq \|x_0 - z\| + \int_0^T \|k(t)\| dt \quad \forall T \geq 0.$$

Since $k \in L^1(0, +\infty; \mathcal{H})$ we obtain that x is bounded. Now we can repeat the arguments in the proof of Theorem 10.3 to conclude that $\lim_{t \rightarrow \infty} \dot{x}(t) = \lim_{t \rightarrow \infty} \nabla f(x(t)) = 0$, so we omit the proof. Let us pass forward and see how the Opial Lemma [75] can be applied.

Since x is bounded and $k \in L^1(0, +\infty; \mathcal{H})$, we get from (10.18) that $\lim_{t \rightarrow \infty} \|x(t) - z\|$ exists, hence the first condition in the Opial Lemma is fulfilled.

To check the second condition in the Opial Lemma is standard. Take $\bar{x} \in \mathcal{H}$ and $t_n \rightarrow +\infty$ such that $x(t_n)$ converges weakly to \bar{x} , as $n \rightarrow +\infty$. The convexity of f yields for all $x \in \mathcal{H}$ and all $n \in \mathbb{N}$

$$f(x) \geq f(x(t_n)) + \langle \nabla f(x(t_n)), x - x(t_n) \rangle.$$

Fixing x and taking the limit as $n \rightarrow +\infty$, and relying on the strong convergence of $\nabla f(x(t))$ to 0 and the boundedness of x , we derive

$$f(x) \geq \liminf_{n \rightarrow +\infty} f(x(t_n)) \geq f(\bar{x}),$$

where the last inequality follows from the weak lower semicontinuity of the convex function f . Since the last inequality holds for an arbitrary x , we obtain $\bar{x} \in \operatorname{argmin} f$. Therefore, the second condition in the Opial Lemma is fulfilled as well.

10.4. A finite stabilization property. As we already noticed, (ADIGE-VGH) can be equivalently written as

$$\dot{u}(t) + \partial\phi(u(t)) \ni -\nabla f(x(t))$$

where $u(t) = \dot{x}(t) + \beta \nabla f(x(t))$. After taking the scalar product of the above equation with $u(t)$, we get

$$\frac{1}{2} \frac{d}{dt} \|u(t)\|^2 + \langle \partial \phi(u(t)), u(t) \rangle = - \langle \nabla f(x(t)), u(t) \rangle.$$

When $\phi(u) \geq r\|u\|$, and by Cauchy–Schwarz inequality we get

$$\frac{1}{2} \frac{d}{dt} \|u(t)\|^2 + r\|u(t)\| \leq \|\nabla f(x(t))\| \|u(t)\|.$$

Since $\nabla f(x(t))$ converges strongly to zero as $t \rightarrow +\infty$ (that's the last point of Theorem 10.3), we get for t large enough $\|\nabla f(x(t))\| \leq \frac{1}{2}r$, and hence

$$\frac{1}{2} \frac{d}{dt} \|u(t)\|^2 + \frac{1}{2}r\|u(t)\| \leq 0.$$

This gives that $u(t) \equiv 0$ after a finite time. Let us summarize the above results in the following Proposition.

PROPOSITION 10.5. *Under the hypothesis of Theorem 10.3, and when $\phi(u) \geq r\|u\|$ for some $r > 0$, we have that after a finite time*

$$\dot{x}(t) + \beta \nabla f(x(t)) \equiv 0,$$

i.e. the trajectory follows the steepest descent dynamic.

10.5. The case f strongly convex: exponential convergence rate.

THEOREM 10.6. *Let $f : \mathcal{H} \rightarrow \mathbb{R}$ be a γ -strongly convex function (for some $\gamma > 0$) which is twice continuously differentiable, and whose gradient is Lipschitz continuous on the bounded sets. Let \bar{x} be the unique minimizer of f . Let $\phi : \mathcal{H} \rightarrow \mathbb{R}_+$ be a damping potential (see Definition 3.1) which is differentiable, and whose gradient is Lipschitz continuous on the bounded subsets of \mathcal{H} . Suppose that ϕ satisfies the following growth conditions:*

(i) (local) *there exists positive constants α , and $\epsilon > 0$ such that, for all u in \mathcal{H} with $\|u\| \leq \epsilon$*

$$\langle \nabla \phi(u), u \rangle \geq \alpha \|u\|^2.$$

(ii) (global) *there exists $p \geq 1$, $r > 0$, such that for all u in \mathcal{H} , $\phi(u) \geq r\|u\|^p$.*

Suppose $\beta > 0$. Let $x : [0, +\infty[\rightarrow \mathcal{H}$ be a solution trajectory of (ADIGE-VGH)

$$(10.20) \quad \ddot{x}(t) + \nabla \phi(\dot{x}(t) + \beta \nabla f(x(t))) + \beta \nabla^2 f(x(t)) \dot{x}(t) + \nabla f(x(t)) = 0.$$

Then, we have exponential convergence rate to zero as $t \rightarrow +\infty$ for $f(x(t)) - f(\bar{x})$, $\|x(t) - \bar{x}\|$ and $\|\dot{x}(t) + \beta \nabla f(x(t))\|$. Moreover, for $\|\dot{x}(t)\|^2$ and $\|\nabla f(x(t))\|^2$, we have exponential convergence rates to zero as $t \rightarrow +\infty$, in the following ergodic sense: there exists $C > 0$ such that

$$\frac{1}{2\beta} \frac{\int_0^t e^{\frac{1}{2\beta}s} \|\dot{x}(s)\|^2 ds}{\int_0^t e^{\frac{1}{2\beta}s} ds} + \frac{\beta}{2} \frac{\int_0^t e^{\frac{1}{2\beta}s} \|\nabla f(x(s))\|^2 ds}{\int_0^t e^{\frac{1}{2\beta}s} ds} = \mathcal{O}(e^{-Ct}).$$

Proof. Since f is strongly convex, f is a coercive function. According to the decrease property of the global energy, see (10.7) and Theorem 10.3 (i), we have that $f(x(t))$ is bounded from above, and hence the trajectory x is bounded. Item (ii) of Theorem 10.3, and the global growth assumption on ϕ give that, for some $p \geq 1$

$$\int_0^{+\infty} \|\dot{x}(t) + \beta \nabla f(x(t))\|^p dt < +\infty.$$

By a similar argument as in the proof of Theorem 10.3 (where we argued with $p = 1$) we deduce that $\lim_{t \rightarrow +\infty} \|\dot{x}(t) + \beta \nabla f(x(t))\| = 0$. Therefore, for t sufficiently large

$$\|\dot{x}(t) + \beta \nabla f(x(t))\| \leq \epsilon.$$

From (10.5) and the local property (i) we derive

$$(10.21) \quad \frac{d}{dt} \left(\frac{1}{2} \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + f(x(t)) - f(\bar{x}) \right) + \alpha \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + \beta \|\nabla f(x(t))\|^2 \leq 0.$$

Since $\dot{x}(\cdot) + \beta \nabla f(x(\cdot))$ is bounded, let $L > 0$ be the Lipschitz constant of $\nabla \phi$ on a ball that contains the vector $\dot{x}(t) + \beta \nabla f(x(t))$ for all $t \geq 0$. Since $\nabla \phi(0) = 0$ we have, for all $t \geq 0$

$$(10.22) \quad \|\nabla \phi(\dot{x}(t) + \beta \nabla f(x(t)))\| \leq L \|\dot{x}(t) + \beta \nabla f(x(t))\|.$$

Using successively (10.20), (10.22) and (5.1), we obtain

$$(10.23) \quad \begin{aligned} \frac{d}{dt} \langle x(t) - \bar{x}, \dot{x}(t) + \beta \nabla f(x(t)) \rangle &= \|\dot{x}(t)\|^2 + \beta \frac{d}{dt} (f(x(t)) - f(\bar{x})) \\ &\quad + \langle x(t) - \bar{x}, -\nabla \phi(\dot{x}(t) + \beta \nabla f(x(t))) - \nabla f(x(t)) \rangle \\ &\leq \|\dot{x}(t)\|^2 + \beta \frac{d}{dt} (f(x(t)) - f(\bar{x})) + \frac{L^2}{2\gamma} \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 \\ &\quad + \frac{\gamma}{2} \|x(t) - \bar{x}\|^2 + \langle \bar{x} - x(t), \nabla f(x(t)) \rangle \\ &\leq \|\dot{x}(t)\|^2 + \beta \frac{d}{dt} (f(x(t)) - f(\bar{x})) + \frac{L^2}{2\gamma} \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 \\ &\quad + f(\bar{x}) - f(x(t)). \end{aligned}$$

Take now $\varepsilon > 0$ (we will specify below how it should be chosen), and define

$$h_{\varepsilon, \beta}(t) := \frac{1}{2} \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + (1 - \beta\varepsilon)(f(x(t)) - f(\bar{x})) + \varepsilon \langle x(t) - \bar{x}, \dot{x}(t) + \beta \nabla f(x(t)) \rangle.$$

Multiplying (10.23) with ε and adding the result to (10.21), we derive

$$\begin{aligned} \dot{h}_{\varepsilon, \beta}(t) &\leq -\alpha \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 - \beta \|\nabla f(x(t))\|^2 + \varepsilon \|\dot{x}(t)\|^2 + \frac{\varepsilon L^2}{2\gamma} \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 \\ &\quad - \varepsilon (f(x(t)) - f(\bar{x})). \end{aligned}$$

We use the inequality

$$(10.24) \quad \varepsilon \|\dot{x}(t)\|^2 \leq 2\varepsilon \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + 2\varepsilon \beta^2 \|\nabla f(x(t))\|^2$$

and we obtain

$$(10.25) \quad \dot{h}_{\varepsilon,\beta}(t) \leq - \left(\alpha - 2\varepsilon - \frac{\varepsilon L^2}{2\gamma} \right) \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 - (\beta - 2\varepsilon\beta^2) \|\nabla f(x(t))\|^2 - \varepsilon(f(x(t)) - f(\bar{x})).$$

Choose $\varepsilon > 0$ small enough such that $C_1 := \min \left\{ \left(\alpha - 2\varepsilon - \frac{\varepsilon L^2}{2\gamma} \right), \beta - 2\varepsilon\beta^2, \varepsilon \right\} > 0$. We obtain

$$(10.26) \quad \dot{h}_{\varepsilon,\beta}(t) \leq -C_1 \left(\|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + \|\nabla f(x(t))\|^2 + f(x(t)) - f(\bar{x}) \right).$$

Further, we have

$$(10.27) \quad h_{\varepsilon,\beta}(t) = \frac{1}{2} \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + \varepsilon \beta (\langle x(t) - \bar{x}, \nabla f(x(t)) \rangle + f(\bar{x}) - f(x(t))) \\ + f(x(t)) - f(\bar{x}) + \varepsilon \langle x(t) - \bar{x}, \dot{x}(t) \rangle.$$

Since f is strongly convex, we have (see for example Theorem 2.1.10 in [74])

$$(10.28) \quad \langle x(t) - \bar{x}, \nabla f(x(t)) \rangle + f(\bar{x}) - f(x(t)) \leq \frac{1}{2\gamma} \|\nabla f(x(t))\|^2.$$

Moreover, from (5.2) and (10.24) we get

$$f(x(t)) - f(\bar{x}) + \varepsilon \langle x(t) - \bar{x}, \dot{x}(t) \rangle \leq f(x(t)) - f(\bar{x}) + \frac{\varepsilon}{2} \|x(t) - \bar{x}\|^2 + \frac{\varepsilon}{2} \|\dot{x}(t)\|^2 \\ \leq \left(1 + \frac{\varepsilon}{\gamma} \right) (f(x(t)) - f(\bar{x})) + \varepsilon \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + \varepsilon \beta^2 \|\nabla f(x(t))\|^2.$$

From this, (10.28) and (10.27) we get

$$h_{\varepsilon,\beta}(t) \leq \left(\frac{1}{2} + \varepsilon \right) \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + \left(\frac{\varepsilon\beta}{2\gamma} + \varepsilon\beta^2 \right) \|\nabla f(x(t))\|^2 + \left(1 + \frac{\varepsilon}{\gamma} \right) (f(x(t)) - f(\bar{x})) \\ \leq C_2 \left(\|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + \|\nabla f(x(t))\|^2 + f(x(t)) - f(\bar{x}) \right),$$

where $C_2 := \max \left\{ \frac{1}{2} + \varepsilon, \frac{\varepsilon\beta}{2\gamma} + \varepsilon\beta^2, 1 + \frac{\varepsilon}{\gamma} \right\} > 0$. Combining this inequality with (10.26), we obtain

$$\dot{h}_{\varepsilon,\beta}(t) + C_3 h_{\varepsilon,\beta}(t) \leq 0,$$

with $C_3 := \frac{C_1}{C_2} > 0$. Then, the Gronwall inequality classically implies

$$(10.29) \quad h_{\varepsilon,\beta}(t) \leq h_{\varepsilon,\beta}(0) e^{-C_3 t}.$$

Finally, from (5.2) and the Cauchy–Schwarz inequality we have

$$h_{\varepsilon,\beta}(t) \geq \frac{1}{2} \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + (1 - \beta\varepsilon) (f(x(t)) - f(\bar{x})) \\ - \frac{\varepsilon}{2} \|x(t) - \bar{x}\|^2 - \frac{\varepsilon}{2} \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 \\ \geq \frac{1 - \varepsilon}{2} \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + \left(1 - \beta\varepsilon - \frac{\varepsilon}{\gamma} \right) (f(x(t)) - f(\bar{x})).$$

Therefore, by taking ε small enough, we obtain

$$(10.30) \quad h_{\varepsilon,\beta}(t) \geq C_4 \left(\|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + f(x(t)) - f(\bar{x}) \right),$$

with $C_4 := \min \left\{ \frac{1-\varepsilon}{2}, 1 - \beta\varepsilon - \frac{\varepsilon}{\gamma} \right\} > 0$. Combining this inequality with (10.29) and (5.2), we obtain an exponential convergence rate to zero for $f(x(t)) - f(\bar{x})$, $\|x(t) - \bar{x}\|$ and $\|\dot{x}(t) + \beta \nabla f(x(t))\|$. Now let's go to the second part of the theorem. From (10.30) we get

$$\|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + f(x(t)) - f(\bar{x}) \leq \frac{1}{C_4} h_{\varepsilon,\beta}(t).$$

Developing the expression on the left hand side, and taking into account (10.29), we obtain

$$(10.31) \quad \frac{1}{2\beta} (\|\dot{x}(t)\|^2 + \beta^2 \|\nabla f(x(t))\|^2) + \frac{d}{dt} (f(x(t)) - f(\bar{x})) + \frac{1}{2\beta} (f(x(t)) - f(\bar{x})) \leq C_5 e^{-C_3 t},$$

with $C_5 = \frac{h_{\varepsilon,\beta}(0)}{2\beta C_4} > 0$. After multiplication of this inequality by $e^{\frac{1}{2\beta} t}$, we obtain

$$(10.32) \quad \frac{d}{dt} \left[e^{\frac{1}{2\beta} t} (f(x(t)) - f(\bar{x})) \right] + e^{\frac{1}{2\beta} t} \left(\frac{1}{2\beta} \|\dot{x}(t)\|^2 + \frac{\beta}{2} \|\nabla f(x(t))\|^2 \right) \leq C_5 e^{-C_6 t}$$

where $C_6 := C_3 - \frac{1}{2\beta} \in \mathbb{R}$. We consider now the cases:

Case I. $C_6 = 0$. Integrating (10.32) we get for $T > 0$:

$$(10.33) \quad e^{\frac{1}{2\beta} T} (f(x(T)) - f(\bar{x})) - (f(x(0)) - f(\bar{x})) + \frac{1}{2\beta} \int_0^T e^{\frac{1}{2\beta} t} \|\dot{x}(t)\|^2 dt + \frac{\beta}{2} \int_0^T e^{\frac{1}{2\beta} t} \|\nabla f(x(t))\|^2 dt \leq C_5 T,$$

hence

$$(10.34) \quad e^{\frac{1}{2\beta} T} \left(\frac{1}{2\beta} \int_0^T e^{\frac{1}{2\beta} t} \|\dot{x}(t)\|^2 dt + \frac{\beta}{2} \int_0^T e^{\frac{1}{2\beta} t} \|\nabla f(x(t))\|^2 dt \right) \leq (f(x_0) - f(\bar{x})) e^{-\frac{1}{2\beta} T} + \frac{C_5 T}{e^{\frac{1}{2\beta} T}} \\ \leq (f(x_0) - f(\bar{x})) e^{-\frac{1}{2\beta} T} + \frac{C_5}{e^{\frac{1}{4\beta} T}},$$

where the last inequality holds for sufficiently large T .

Case II. $C_6 \neq 0$. Integrating (10.32) we get for $T > 0$:

$$(10.35) \quad e^{\frac{1}{2\beta} T} (f(x(T)) - f(\bar{x})) - (f(x(0)) - f(\bar{x})) + \frac{1}{2\beta} \int_0^T e^{\frac{1}{2\beta} t} \|\dot{x}(t)\|^2 dt + \frac{\beta}{2} \int_0^T e^{\frac{1}{2\beta} t} \|\nabla f(x(t))\|^2 dt \\ \leq C_5 \left(\frac{1}{C_6} - \frac{e^{-C_6 T}}{C_6} \right).$$

Let us consider two subcases:

Case II(i). $C_6 > 0$. From (10.35) we have

$$(10.36) \quad e^{\frac{1}{2\beta} T} (f(x(T)) - f(\bar{x})) - (f(x(0)) - f(\bar{x})) + \frac{1}{2\beta} \int_0^T e^{\frac{1}{2\beta} t} \|\dot{x}(t)\|^2 dt + \frac{\beta}{2} \int_0^T e^{\frac{1}{2\beta} t} \|\nabla f(x(t))\|^2 dt < \frac{C_5}{C_6},$$

hence

$$(10.37) \quad \int_0^{+\infty} e^{\frac{1}{2\beta}t} \|\dot{x}(t)\|^2 dt < +\infty \text{ and } \int_0^{+\infty} e^{\frac{1}{2\beta}t} \|\nabla f(x(t))\|^2 dt < +\infty$$

and

$$(10.38) \quad e^{-\frac{1}{2\beta}T} \left[\frac{1}{2\beta} \int_0^T e^{\frac{1}{2\beta}t} \|\dot{x}(t)\|^2 dt + \frac{\beta}{2} \int_0^T e^{\frac{1}{2\beta}t} \|\nabla f(x(t))\|^2 dt \right] \leq \left(f(x_0) - f(\bar{x}) + \frac{C_5}{C_6} \right) e^{-\frac{1}{2\beta}T}.$$

Case II(ii). $C_6 < 0$. From (10.35) we get

$$(10.39) \quad e^{\frac{1}{2\beta}T} (f(x(T)) - f(\bar{x})) - (f(x(0)) - f(\bar{x})) + \frac{1}{2\beta} \int_0^T e^{\frac{1}{2\beta}t} \|\dot{x}(t)\|^2 dt + \frac{\beta}{2} \int_0^T e^{\frac{1}{2\beta}t} \|\nabla f(x(t))\|^2 dt < \frac{C_5}{-C_6} e^{-C_6 T}.$$

Multiplying this inequality with $e^{-\frac{1}{2\beta}T}$ we obtain

$$(10.40) \quad e^{-\frac{1}{2\beta}T} \left[\frac{1}{2\beta} \int_0^T e^{\frac{1}{2\beta}t} \|\dot{x}(t)\|^2 dt + \frac{\beta}{2} \int_0^T e^{\frac{1}{2\beta}t} \|\nabla f(x(t))\|^2 dt \right] < (f(x(0)) - f(\bar{x})) e^{-\frac{1}{2\beta}T} + \frac{C_5}{-C_6} e^{-(C_6 + \frac{1}{2\beta})T}.$$

Noticing that $C_6 + \frac{1}{2\beta} = C_3 > 0$ we get

$$(10.41) \quad e^{-\frac{1}{2\beta}T} \left[\frac{1}{2\beta} \int_0^T e^{\frac{1}{2\beta}t} \|\dot{x}(t)\|^2 dt + \frac{\beta}{2} \int_0^T e^{\frac{1}{2\beta}t} \|\nabla f(x(t))\|^2 dt \right] < (f(x(0)) - f(\bar{x})) e^{-\frac{1}{2\beta}T} + \frac{C_5}{-C_6} e^{-C_3 T}.$$

Taking into account that

$$\int_0^T e^{\frac{1}{2\beta}t} dt = 2\beta \left(e^{\frac{1}{2\beta}T} - 1 \right) \text{ and } \lim_{T \rightarrow +\infty} \frac{2\beta \left(e^{\frac{1}{2\beta}T} - 1 \right)}{e^{\frac{1}{2\beta}T}} = 2\beta,$$

the conclusion follows in each of the cases considered above. \square

Remark 10.7. Similar rates have been reported in [22, Theorem 4.2] for the heavy ball method with Hessian driven damping.

Remark 10.8. It is possible to derive similar exponential rates also for the system (9.4), however for a restrictive choice of $\beta > 0$. To see this, notice that for $\theta > 0$ we have

$$\begin{aligned} & \frac{d}{dt} \left(\frac{1}{2} \|\dot{x}(t) + \beta \nabla f(x(t))\|^2 + f(x(t)) - f(\bar{x}) \right) \\ &= -\langle \dot{x}(t), \nabla \phi(\dot{x}(t)) \rangle - \beta \langle \nabla \phi(\dot{x}(t)), \nabla f(x(t)) \rangle - \beta \|\nabla f(x(t))\|^2 \\ &\leq -\alpha \|\dot{x}(t)\|^2 + \frac{\beta\theta}{2} \|\nabla f(x(t))\|^2 + \frac{\beta L^2}{2\theta} \|\dot{x}(t)\|^2 - \beta \|\nabla f(x(t))\|^2 \\ &= -\left(\alpha - \frac{\beta L^2}{2\theta} \right) \|\dot{x}(t)\|^2 - \beta \left(1 - \frac{\theta}{2} \right) \|\nabla f(x(t))\|^2. \end{aligned}$$

10.6. Further convergence results based on the quasi-gradient approach. Let us consider the dynamical systems (ADIGE-VGH) in case ϕ is differentiable, $f : \mathbb{R}^N \rightarrow \mathbb{R}$ is a \mathcal{C}^2 function (possibly nonconvex) whose gradient is Lipschitz continuous on the bounded sets, and such that $\inf_{\mathbb{R}^N} f > -\infty$:

$$(10.42) \quad \ddot{x}(t) + \nabla\phi(\dot{x}(t) + \beta\nabla f(x(t))) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) = 0.$$

The considerations are similar to those of Section 7.3 and Theorem 7.3. According to Step 2 in the proof of Theorem 10.2, the first-order reformulation is

$$(10.43) \quad \dot{z}(t) + F(z(t)) = 0,$$

where $z(t) = (x(t), \dot{x}(t)) \in \mathbb{R}^N \times \mathbb{R}^N$, and $F : \mathbb{R}^N \times \mathbb{R}^N \rightarrow \mathbb{R}^N \times \mathbb{R}^N$ is defined by

$$F(x, u) = (\beta\nabla f(x) - u, \nabla\phi(u) + \nabla f(x)).$$

Let us check the angle condition (E_λ is defined as in Theorem 7.3). We have

$$\langle \nabla E_\lambda(x, u), F(x, u) \rangle = \left\langle \left(\nabla f(x) + \lambda\nabla^2 f(x)u, u + \lambda\nabla f(x) \right), \left(\beta\nabla f(x) - u, \nabla\phi(u) + \nabla f(x) \right) \right\rangle.$$

After development and simplification, we get

$$\begin{aligned} \langle \nabla E_\lambda(x, u), F(x, u) \rangle &= -\lambda \langle \nabla^2 f(x)u, u \rangle + \langle u, \nabla\phi(u) \rangle + \lambda \langle \nabla f(x), \nabla\phi(u) \rangle + \lambda \|\nabla f(x)\|^2 \\ &\quad + \beta \|\nabla f(x)\|^2 + \lambda\beta \langle \nabla f(x), \nabla^2 f(x)u \rangle. \end{aligned}$$

We estimate the term $\lambda\beta \langle \nabla f(x), \nabla^2 f(x)u \rangle$ by writing

$$\lambda\beta \langle \nabla f(x), \nabla^2 f(x)u \rangle \geq -\frac{\lambda}{4} \|\nabla f(x)\|^2 - \lambda\beta^2 M^2 \|u\|^2$$

and get (as in the proof of Theorem 7.3)

$$(10.44) \quad \langle \nabla E_\lambda(x, u), F(x, u) \rangle \geq \left(\gamma - \lambda M - \frac{\lambda}{2} \delta^2 - \lambda\beta^2 M^2 \right) \|u\|^2 + \left(\frac{\lambda}{4} + \beta \right) \|\nabla f(x)\|^2.$$

We also have

$$\|F(x, u)\| \leq C_2 (\|u\|^2 + \|\nabla f(x)\|^2)^{\frac{1}{2}},$$

where $C_2 = \sqrt{2(2 + \beta^2 + \delta^2)}$. The rest can be done in the lines of the proof of Theorem 7.3.

11. Conclusion, perspectives. In this article, from the point of view of optimization, we put forward some classical and new properties concerning the asymptotic convergence of autonomous damped inertial dynamics. From a control point of view, the damping terms of these dynamics can be considered as closed-loop controls of the current data: position, speed, gradient of the objective function, Hessian of the objective function, and combination of these objects. Let us list some of the main results and advantages of the autonomous approach compared to the non-autonomous approach, where the damping is given as a function of time.

11.1. PRO.

- Autonomous systems are easy to implement. It is not necessary to adjust the damping coefficient as is the case for non-autonomous systems.
- When the function to be minimized is strongly convex, there is convergence at an exponential rate, and this is valid for a large class of damping potentials.
- We were able to exploit the quasi-gradient structure of the autonomous damped dynamics and combine them with the Kurdyka-Lojasiewicz theory to obtain convergence rates for a large class of functions f , possibly non-convex. This is specific to the autonomous case because the theories mentioned above are not developed in the non-autonomous case.
- The Hessian damping naturally comes within the framework of autonomous systems. It notably improves the theoretical and numerical behavior of the trajectories, by reducing the oscillatory aspects. Its introduction into the algorithms does not change their numerical complexity (it makes appear the difference of the gradient at two consecutive steps). This makes this geometric damping very successful, several recent articles have been devoted to it.
- The closed-loop approach clearly distinguishes between the strong and weak damping effects, and the transition between them. It also shows the replacement of the theory of convergence by the notion of attractor when the damping becomes too weak.
- We have introduced a new autonomous system where the damping involves both the speed and the gradient of f , and which benefits from very good convergence properties. At the beginning of time it takes advantage of the inertial effect, then after a finite time it turns into a steepest descent dynamic, thus avoiding the oscillatory aspects. This regime change has some similarities with the restart method, and also the recent work of Poon-Liang [79] on adaptive acceleration.
- The closed-loop approach makes it possible to make the link with different fields, such as PDE and control theory, where the stabilization of oscillating systems is a central issue.
- We have developed an inertial gradient algorithm which shares the good convergence properties of the related continuous dynamics, in the case of the quasi-gradient and Kurdyka-Lojasiewicz approach. Note that the quasi-gradient approach reflects relative errors in the algorithms, and therefore gives a lot of flexibility. It is this approach that has made it possible to deal with many different algorithms in Attouch-Bolte-Svaiter [14] in the non-convex nonsmooth case. It would be interesting to develop these aspects for many other algorithms, such as proximal algorithms, proximal gradient algorithms, regularized Gauss-Seidel algorithms, and PALM.

11.2. CONS.

1. To date, we do not know in the autonomous case the equivalent of the accelerated gradient method of Nesterov and Su-Boyd-Candès damped inertial dynamic, that is to say an adjustment of the damping potential which guarantees the rate of convergence of values $1/t^2$ for any convex function. This is a current research subject, for recent progress in this direction, see Lin-Jordan [71].
2. The general approach based on the quasi-gradient and the Kurdyka-Lojasiewicz theory (as developed in Section 7) works mainly in finite dimension. The extension of the KL theory to spaces of infinite dimension is a current research subject.

11.3. Perspectives.

1. Develop closed-loop versions of the Nesterov accelerated gradient method from a theoretical and numerical point of view. Our analysis allowed us to better define the type of damping potential ϕ capable of doing this, but this remains an open question for study. Indeed, the case $p = 2$ (*i.e.* quadratic behavior of the damping potential near the origin) is the critical case separating the weak damping from the strong damping. Taking $p > 2$, with p close to 2 provides a vanishing viscosity damping coefficient, which is a specific property of the Nesterov accelerated gradient method. Our intuition is that we need to refine the power scale which is not precise enough to provide the correct setting of the vanishing damping coefficient (*i.e.* going from $p = 2$ to $p > 2$, with p even very close to 2 is too sudden a change).
2. Extend our study to the case of nonsmooth optimization possibly involving a constraint. This is an important subject, which is closely related to item 6 of this list, because a common device to deal with a constrained optimization problem is to use a gradient-projection method, which falls under fixed point methods.
3. Develop a control perspective with closed-loop damping for the restarting methods. The restarting methods take advantage of the inertial effect to accelerate the trajectories, then stop when a given criteria is deteriorate. Then restart from the current point with zero velocity, and so on. In many ways, the dynamic we developed in Section 10 follows a similar strategy. Our results are valid with general data functions f and ϕ , while the known results concerning the restart methods only concern the case where f is strongly convex. It is an important subject of study, largely to explore.
4. Obtain a closed-loop version of the Tikhonov regularization method, and make the link with the Haugazeau method. The objective is then, within the framework of the convex optimization, to obtain an autonomous dynamic whose trajectories strongly converge towards the solution of minimum norm; see Attouch–Cabot–Chbani– Riahi [20], and Boţ–Csetnek–László [47] for some recent results in the open-loop case (the Tikhonov regularization parameter tends to zero in a controlled manner, not too fast) and references therein.
5. Develop the corresponding algorithmic results. Continuous dynamics provide a valuable guide to introduce and analyze algorithms that benefit from similar convergence properties. In Theorem 8.1 we have analyzed the convergence property of an inertial algorithm with general damping potential ϕ and general (tame) function f . A similar analysis can certainly be developed on the basis of the Theorems 9.2 and 10.3 which also involve the Hessian-driven damping. Natural extensions then consist in studying structured optimization problems and the corresponding proximal-based algorithms.
6. In recent years, most of the previous themes have been extended (in the open-loop case) to the case of maximally monotone operators, see Álvarez–Attouch [8], Attouch–Maingé [28], Attouch–Peypouquet [31], Attouch–Cabot [17], Attouch–László [27], Boţ–Csetnek [45]. It would be interesting to consider the closed-loop version of these dynamics, as was done by Attouch–Redont–Svaiter [34] for first-order Newton-like evolution systems.
7. Time rescaling is a powerful tool to accelerate the inertial systems; see Attouch–Chbani–Riahi [24], Shi–Du–Jordan–Su [80] and references therein. It leads naturally to non-autonomous dynamics. It would be interesting to study autonomous closed-loop version. This means first extracting quantities which tend monotonically to $+\infty$.

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