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LONG TIME BEHAVIOUR OF SOLUTIONS TO AN EVOLUTION PDE WITH NONSTANDARD GROWTH

H. FRANKOWSKA AND G. MOSCARIELLO

Abstract. In this paper we prove time estimates for solutions to a general nonhomogeneous parabolic problem whose operator satisfies nonstandard growth conditions. We also study the asymptotic behaviour of solutons to an anisotropic problem.

Keywords. Evolution PDE, time dependence of solutions at infinity, generalized Gronwall lemma.

Mathematics Subject Classification. 35K51, 35K59, 35B40.

1. Introduction.

In this paper we investigate the time dependence of solutions for a wide class of nonuniformly parabolic equations whose model case is the following one:

$$\begin{cases} u_t - \operatorname{div} \frac{G'(|\nabla u|)}{|\nabla u|}\nabla u = 0 & \text{in } \Omega_T \\ u(x,t) = 0 & \text{on } \partial\Omega \times (0,T) \\ u(x,0) = u_0(x) & \text{in } \Omega, \end{cases} \tag{1.1}$$

where for some $N \geq 2$, Ω is a bounded domain in $\mathbb{R}^N$ with Lipschitz boundary $\partial\Omega$, T is a positive number and $\Omega_T = \Omega \times (0,T)$.

We assume that $u_0 \in L^2(\Omega)$ and $G : \mathbb{R}_+ \to \mathbb{R}_+$ is a C^1 convex function satisfying, for some $k > 0$ and $p \geq 2$ the following conditions

$$\begin{cases} (j) \qquad G(0) = 0; \qquad G(2s) \leq kG(s) \qquad \forall s > 0 \\ \\ (jj) \qquad \frac{G(s)}{s^p} \text{ is increasing on } (0,\infty) \end{cases} \tag{1.2}$$

We refer to [16, 26] for the properties of G that follow from assumptions (1.2). The functions listed below do satisfy the above assumptions for $s \geq 0$:

1) $G(s) = \frac{s^p}{p}, \quad p \geq 2;$
2) $G(s) = s^p \log(1+s), \quad p \geq 2;$
3) $G(s) = s^p L_k(s), \ p \geq 2, \ \ L_i(s) = \log(1+L_{i-1}(s)), i = 1, ..., k, \ L_0(s) = \log(1+s)$
4) $G(s) = \int_0^s g(\rho)d\rho$ where $g : \mathbb{R}_+ \to \mathbb{R}_+$ is a C^1 function satisfying

$$p - 1 \leq \frac{sg'(s)}{g(s)} \leq q - 1, \qquad \forall s > 0$$

with $2 \leq p \leq q < \infty$ such that $\lim\limits_{s\to+\infty} \dfrac{g(s)}{s} = +\infty$, see [3] where the properties of such a function G were investigated.

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Definition 1.1. *A function $u : \overline{\Omega}_T \to \mathbb{R}$ is called a weak solution to Problem* (1.1) *iff* $u \in L^\infty(0,T;L^2(\Omega)) \cap L^1(0,T;W_0^{1,1}(\Omega))$,

$$\int_0^T \int_\Omega G(|\nabla u|)dxdt < +\infty$$

and for every $\phi \in C^1(\overline{\Omega}_T)$ with $\phi(\cdot,t)|_{\partial\Omega} = 0$ the following equality holds

$$\int_\Omega u\phi(x,t)dx - \int_\Omega u_0\phi(x,0)dx + \int_0^t \int_\Omega \left[-u\phi_s + \left\langle \frac{G'(|\nabla u|)}{|\nabla u|}\nabla u, \nabla\phi \right\rangle\right] dxds = 0$$

for any $t \in [0,T]$.

If $u \in L^\infty_{loc}((0,\infty);L^2(\Omega)) \cap L^1_{loc}((0,\infty);W_0^{1,1}(\Omega))$ and the above holds true for all $T > 0$, then u is called a weak solution to Problem (1.1) *on* $(0,\infty) \times \Omega$.

In [8] the authors have shown that under (1.2) and some additional structural assumptions, there exists a unique weak solution $u \in C([0,T];L^2(\Omega))$ to problem (1.1). When G is as in example 4), the existence and the uniqueness of a solution $u \in C(\overline{\Omega}_T)$ to a Cauchy-Dirichlet problem for evolution equation in (1.1) have been investigated in [3], again under several additional hypothesis.

In the elliptic framework equation in (1.1) is the Euler equation of the energy functional

$$\int_\Omega G(|\nabla u|)dx,$$

where the convex integrand may have nonstandard growth conditions.

Note that the function G in 2) satisfies for every $\epsilon > 0$ the growth condition

$$s^p \le G(s) \le L_\epsilon(1+s)^{p+\epsilon}.$$

whenever s is sufficiently large and where $L_\epsilon > 0$ is a constant depending only on ϵ and G. Starting with the pioneering papers by P. Marcellini [19, 20], the theory in the stationary case, especially the regularity theory, has been extensively studied ([2, 10, 14, 15, 18, 24, 22, 23]). For an almost complete treatment see the survey [21] and the references therein. See also the recent paper [9].

Following the elliptic scheme, in [5] the authors recently studied variational solutions in the sense of [17] to the Cauchy Dirichlet problem (1.1). The existence of such solution has been proved in [7].

The objective of the present work is to obtain time estimates on weak solutions to (1.1) in terms of the G function. More precisely, we get the following result:

Theorem 1.1. *Assume* (1.2) *and let $u : \overline{\Omega}_T \to \mathbb{R}$ be a weak solution to problem* (1.1). *Then for any $t \in [0,T]$,*

$$\|u(\cdot,t)\|^2_{L^2(\Omega)} \le 2|\Omega|x(t)$$

where $x(\cdot)$ is the unique solution of the problem

$$\begin{cases} x'(t) = -c\,G(\sqrt{x(t)}) & a.e. \\ x(0) = \frac{1}{2|\Omega|}\|u_0\|^2_{L^2(\Omega)} \end{cases} \tag{1.3}$$

and the constant $c > 0$ depends only on p, k, $|\Omega|$.

Furthermore, if $T = \infty$, then $\lim_{t\to\infty} \|u(\cdot,t)\|^2_{L^2(\Omega)} = 0$.

Due to the nature of the problem, obtaining such estimates on solutions to (1.1) is not a trivial fact. As far as we know, Theorem 1.1 is the first result in this direction. Moreover, Theorem 1.1 implies time estimates also for evolution problems related to operators having standard growth conditions, as, for instance, for parabolic p-Laplace equation. Indeed Theorem 1.1 improves the results on the behaviour of a solution contained in [25] and in [12] (see Remark 4.2 below).

The main ingredient to achieve our result is a version of the Gronwall Lemma that has an interest by itself (see Lemma 3.1 and its two Corollaries in Section 3.). An energy balance equality for solutions to (1.1) (see Proposition 2.1) is also fundamental.

More generally, in this paper we prove time estimates, similar to those of Theorem 1.1 for solutions to a nonhomogeneous parabolic problem of the type

$$\begin{cases} u_t - \text{div } A(x,t,\nabla u) = -\text{div } f & \text{in } \Omega_T \\ u(x,t) = 0 & \text{on } \partial\Omega \times (0,T) \\ u(x,0) = u_0(x) & \text{in } \Omega, \end{cases}$$

where $A : \Omega_T \times \mathbb{R}^N \to \mathbb{R}^N$ is a Carathéodory function satisfying some nonstandard growth conditions and $f : \Omega_T \to \mathbb{R}^N$ satisfies the integrability condition (2.6) below.

Another application of our techniques is given in Section 5, where we study the asymptotic behaviour of solutions to an anisotropic problem.

The plan of the paper is as follows: In Section 2 we provide notations and preliminary results. In Section 3 we prove a generalized Gronwall Lemma. In Sections 4 and 5 we state and prove our main results.

2. Notation and preliminary results

Let us consider the parabolic problem

$$\begin{cases} u_t - \text{div } A(x,t,\nabla u) = -\text{div } f & \text{in } \Omega_T \\ u(x,t) = 0 & \text{on } \partial\Omega \times (0,T) \\ u(x,0) = u_0(x) & \text{in } \Omega, \end{cases} \tag{2.1}$$

where Ω is a Lipschitz bounded domain of $\mathbb{R}^N$, $N \geq 2$, $\Omega_T = \Omega \times (0,T)$ and $u_0 \in L^2(\Omega)$. We assume that the field $A : \Omega_T \times \mathbb{R}^N \to \mathbb{R}^N$ is a Carathéodory function such that for a.e. $(x,t) \in \Omega_T$ and for any $\xi \in \mathbb{R}^N$

$$\begin{cases} (i)\ \ \langle A(x,t,\xi),\xi\rangle \geq \nu G(|\xi|), & 0 < \nu \leq 1 \\ (ii)\ \ |A(x,t,\xi)| \leq \mu \frac{G(|\xi|)}{|\xi|} & \xi \neq 0,\ \ 1 \leq \mu \\ (iii)\ \ A(x,t,0) = 0, \end{cases} \tag{2.2}$$

where $G : \mathbb{R}_+ \to \mathbb{R}_+$ is a non zero convex function satisfying (1.2).

Let us remark that (2.2) could be imposed, equivalently, for $0 < \nu < \mu$.

If G is a C^1 non–trivial convex function satisfying (j) in (1.2), then the field $A(x,t,\xi) = \frac{G'(|\xi|)}{|\xi|}\xi$, $\xi \neq 0$, considered in the model case of Section 1 satisfies conditions (2.2). Indeed, in this case, since $G(0) = 0$. The convexity of G implies for all $\xi \in \mathbb{R}^N$, $\xi \neq 0$,

$$G(|\xi|) = \int_0^1 G'(s|\xi|)|\xi| ds \leq \left\langle \frac{G'(|\xi|)}{|\xi|}\xi, \xi \right\rangle$$

and condition (i) follows. Moreover, by the convexity of G and from the second relation in (j), we get for all $t > 0$,

$$kG(t) \geq G(2t) = \int_0^{2t} G'(s)ds \geq \int_t^{2t} G'(s)ds \geq tG'(t)$$

from which (ii) follows.

The Fenchel conjugate of G is defined by

$$\tilde{G}(r) = \sup_{s\geq 0}(rs - G(s)).$$

Then the Fenchel's inequality

$$sr \leq G(s) + \tilde{G}(r) \tag{2.3}$$

holds true for every $r, s \geq 0$. Note that $\tilde{G} : \mathbb{R} \to \mathbb{R}_+$, $\tilde{G}(0) = 0$ and $\tilde{G}$ is increasing (in this paper we say increasing for nondecreasing).

Another useful property is (see [26] Chapt. II, (1))

$$\tilde{G}\left(\frac{G(r)}{r}\right) \leq G(r) \qquad \text{for every } r > 0. \tag{2.4}$$

Observe that G is increasing and it is not difficult to realise that unless G is the zero function, (j) implies that $G(s) > 0$ for all $s > 0$.

Condition (j), which is known as the Δ_2-condition, is equivalent to the assumption that there exists $q > 1$ such that $\frac{G(t)}{t^q}$ is decreasing. On the other hand (jj) is equivalent to Δ_2 condition for the conjugate function $\tilde{G}$ of G, i.e. there exists $\tilde{k} > 0$ such that $\tilde{G}(2r) \leq \tilde{k}\tilde{G}(r)$ for every $r > 0$ (see [16] Sect. 4). Consequently, $\tilde{G}(\cdot) > 0$ on $(0, \infty)$.

Notice that Δ_2 condition yields

$$G(\lambda s) \leq \lambda^k G(s), \qquad \tilde{G}(\lambda s) \leq \lambda^{\tilde{k}}\tilde{G}(s) \qquad \forall s \geq 0,\ \lambda \geq 1. \tag{2.5}$$

For other properties of function G we refer to [16, 26].

In Problem (2.1) the datum $f : \Omega_T \to \mathbb{R}^N$ is assumed to be a measurable function such that

$$\int_0^T \int_\Omega \tilde{G}(|f(x,s)|)dxds < +\infty. \tag{2.6}$$

Definition 2.1. *A function $u : \overline{\Omega}_T \to \mathbb{R}$ is a weak solution to Problem* (2.1) *iff*

$$u \in L^\infty(0,T; L^2(\Omega)) \cap L^1(0,T; W_0^{1,1}(\Omega)) \tag{2.7}$$

$$\int_0^T \int_\Omega \langle A(x,s,\nabla u), \nabla u\rangle dxds < +\infty \tag{2.8}$$

and for every $\phi \in C^1(\overline{\Omega}_T)$ with $\phi(\cdot,t)|_{\partial\Omega} = 0$ the following equality holds

$$\int_\Omega u\phi(x,t)dx - \int_\Omega u_0\phi(x,0)dx + \int_0^t \int_\Omega [-u\phi_s + \langle A(x,s,\nabla u), \nabla\phi\rangle]dxds = \int_0^t \int_\Omega \langle f, \nabla\phi\rangle dxds \tag{2.9}$$

for any $t \in [0,T]$.

If $u \in L^\infty_{loc}((0,\infty); L^2(\Omega)) \cap L^1_{loc}((0,\infty); W^{1,1}_0(\Omega))$ *and the above holds true for all* $T > 0$, *then* u *is called a weak solution to Problem* (2.1) *on* $(0,\infty) \times \Omega$.

Notice that assumption (2.8) is equivalent to the following condition

$$\int_0^T \int_\Omega G(|\nabla u|) dx dt < +\infty. \tag{2.10}$$

Indeed, since $\tilde{G}$ is increasing, by (2.3), (*ii*), (2.4) and (2.5), for any $(x,s) \in \Omega_T$ with $\nabla u(x,s) \neq 0$ we have

$$\begin{aligned} &|\langle A(x,s,\nabla u(x,s)), \nabla u(x)\rangle| \leq |A(x,s,\nabla u(x,s))| \cdot |\nabla u(x)| \\ &\leq \tilde{G}(|A(x,s,\nabla u(x,s))|) + G(|\nabla u(x)|) \leq \tilde{G}\left(\mu \tfrac{G(|\nabla u(x,s)|)}{|\nabla u(x,s)|}\right) + G(|\nabla u(x)|) \\ &\leq \mu^{\tilde{k}} \tilde{G}\left(\tfrac{G(|\nabla u(x,s)|)}{|\nabla u(x,s)|}\right) + G(|\nabla u(x)|) \leq \mu^{\tilde{k}} G(|\nabla u(x,s)|) + G(|\nabla u(x)|). \end{aligned} \tag{2.11}$$

The equivalence between (2.8) and (2.10) follows from the first condition in (2.2) and from (2.11).

Remark 2.1. *Using conditions* (2.3), (2.6) *and* (2.10) *respectively, an argument similar to* (2.11) *proves that* $\langle f(\cdot), \nabla\phi(\cdot)\rangle$ *and* $\langle A(\cdot,\cdot,\nabla u(\cdot,\cdot)), \nabla\phi(\cdot)\rangle$ *are integrable on* Ω_T, *for every* $\phi \in C^1(\overline{\Omega}_T)$. *These considerations imply that under our assumptions, all integrals in* (2.9) *are well defined.*

Remark 2.2. *By assumptions* (1.2), *for any* $r > 0$ *we have*

$$\lambda_1 r^{q'} \leq \tilde{G}(r) \leq \lambda_2 (r^{p'} + 1)$$

for some positive constants λ_1, λ_2 *and* $pp' = p + p'$, *and* $qq' = q + q'$ *(see* [26] *Chapt. 2). Then, from* (2.9) *we deduce that* $u_t \in L^{q'}(0,T; L^{q'}(\Omega))$ *with* $1 < q' < 2$, *and so there exists* $\tilde{u} \in C([0,T]; L^{q'}(\Omega))$ *such that* $u(\cdot,t) = \tilde{u}(\cdot,t)$ *for almost all* $t \in (0,T)$. *In what folllows we identify* u *with its continuous representant* $\tilde{u}$.

Lemma 2.1. *Let* $u \in L^\infty(0,T; L^2(\Omega)) \cap L^1(0,T; W^{1,1}_0(\Omega))$, *be a solution to Problem* (2.1). *Then,* $u \in C_w([0,T]; L^2(\Omega))$.

Proof. Let $t \in [0,T]$ be fixed. Let $\{t_n\}$ be a sequence in $[0,T]$ with $t_n \to t$ as $n \to +\infty$ such that $u(\cdot,t_n) \in L^2(\Omega)$ $\forall n \in \mathbb{N}$. Since $u \in L^\infty(0,T; L^2(\Omega))$, the sequence $\{u(\cdot,t_n)\}$ is bounded in $L^2(\Omega)$. Then, there exists a subsequence $\{u(\cdot,t_{k_n})\}$ and $\xi \in L^2(\Omega)$ such that

$$u(\cdot,t_{k_n}) \to \xi \qquad \text{weakly in } L^2(\Omega) \qquad \text{as } n \to +\infty.$$

On the other hand, by Remark 2.2 $u \in C([0,T]; L^{q'}(\Omega))$. Therefore

$$u(\cdot,t_{k_n}) \to u(\cdot,t) \qquad \text{strongy in } L^{q'}(\Omega) \qquad \text{as } n \to +\infty.$$

This implies that $\xi = u(\cdot,t)$ and $u \in C_w([0,T]; L^2(\Omega))$. □

We also recall the following useful result.

Lemma 2.2. *Let $G\colon \mathbb{R}_+ \to \mathbb{R}_+$ be a function satisfying (1.2). Then there exists a convex function $F : \mathbb{R}_+ \to \mathbb{R}_+$ satisfying (j) such that*

$$F(t^p) \le G(t) \le kF(t^p). \tag{2.12}$$

where k and p are as in (1.2).

Proof. Since $G(s^{\frac{1}{p}})/s$ is increasing, it is not difficult to realize that the function

$$F(t) = \int_0^t \frac{G(s^{\frac{1}{p}})}{s} ds$$

is well defined on $(0, \infty)$ and can be extended by continuity to zero by setting $F(0) = 0$. Its first derivative is a continuous increasing function $(0, \infty)$. Hence F is convex on $\mathbb{R}_+$.

Moreover, by (jj), the integrand is increasing and therefore

$$F(t) \le \int_0^t \frac{G(t^{\frac{1}{p}})}{t} ds = G(t^{\frac{1}{p}})$$

implying the first inequality. We show next that F satisfies (j). Indeed, since G satisfies (j),

$$F(2t) = \int_0^t \frac{G((2s)^{\frac{1}{p}})}{s} ds \le \int_0^t \frac{G(2s^{\frac{1}{p}})}{s} ds \le k \int_0^t \frac{G(s^{\frac{1}{p}})}{s} ds = kF(t).$$

The second inequality in (2.12) follows from the following estimates

$$G(t^{\frac{1}{p}}) = \int_t^{2t} \frac{G(t^{\frac{1}{p}})}{t} ds \le \int_t^{2t} \frac{G(s^{\frac{1}{p}})}{s} ds \le F(2t) \le kF(t).$$

□

In this paper we do not investigate the existence of a solution to (2.1) and refer to [3, 7, 8] for some existence results. Instead, we establish some properties of weak solutions. We first show that every such solution satisfies an energy balance equality.

Proposition 2.1. *Assume* (2.2) *and* (2.6) *and let $u : \overline{\Omega}_T \to \mathbb{R}$ be a weak solution to* (2.1). *Then for any $0 < t \le T$*

$$\frac{1}{2}\|u(t)\|^2_{L^2(\Omega)} + \int_0^t \int_\Omega \langle A(x, s, \nabla u), \nabla u\rangle dxds = \frac{1}{2}\|u_0\|^2_{L^2(\Omega)} + \int_0^t \int_\Omega \langle f(x, s), \nabla u\rangle dxds.$$

where $\|u(t)\|^2_{L^2(\Omega)} := \int_\Omega |u(x, t)|^2 dx$.

Proof. Let $u : \overline{\Omega}_T \to \mathbb{R}$ be a weak solution to (2.1). We first extend solution $u(x, t)$ by the initial value $u_0(x)$ when $t < 0$. As $\partial\Omega$ is Lipschitz, we can cover $\partial\Omega$ by a finite open family $\{U_i\}_{i=1}^m$ and we can find corresponding positive numbers λ_i and vectors p_i such that for $\epsilon > 0$ the ball $B(x + \lambda_i\epsilon p_i, \epsilon) \subset \Omega$ for all $x \in U_i \cap \Omega$ (cf. [11]).

Choose an open set $U_0 \subset\subset \Omega$ such that $\{U_i\}_{i=0}^m$ forms a covering of Ω and let $\{\eta_i\}_{i=0}^m$ be a smooth partition of unity corresponding to this covering. For $x \in U_i \cap \Omega,\ i = 1, ..., m$, denote $u_i^\epsilon(x, t) = u(x + \lambda_i\epsilon p_i, t)$ and let $u_\epsilon^i(\cdot, t)$ be its mollification in $U_i \cap \Omega$. Finally we mollify $u(x, t)$ in U_0 to get u_ϵ^0.

The functions $u_\epsilon : \overline{\Omega}_T \to \mathbb{R}$ defined as

$$u_\epsilon = \sum_{i=0}^{m} \eta_i u_\epsilon^i,$$

are $C_0^\infty(\Omega)$ and converge to u in $W_0^{1,1}(\Omega)$ when $\epsilon \to 0^+$ for a.e. $t \in [0,T]$.

Consider the mollifier in time of $u_\epsilon(x,\cdot)$ defined for any natural $n \in \mathbb{N}$ such that $\frac{1}{n} < \epsilon$, as:

$$[u_\epsilon]_n (x,t) = e^{-nt} u_\epsilon(x,t_0) + n \int_0^t e^{n(s-t)} u_\epsilon(x,s) ds,$$

where $t_0 \in (0,T)$. Since $u_\epsilon \in L^\infty(0,T;L^2(\Omega)) \cap L^1(0,T;W_0^{1,1}(\Omega))$ and $u_\epsilon(x,t_0) \in L^2(\Omega)$, we get that

$$[u_\epsilon]_n \to u_\epsilon \text{ in } L^2(\Omega_T) \text{ as } n \to \infty. \tag{2.13}$$

$$[u_\epsilon]_n (x,0) = u_\epsilon(x,t_0) \quad \forall n \in \mathbb{N}.$$

$$[u_\epsilon]_n \in L^1\left([0,T]; W_0^{1,1}(\Omega)\right) \tag{2.14}$$

$$\nabla [u_\epsilon]_n (\cdot,t) = [\nabla u_\epsilon]_n (\cdot,t) \tag{2.15}$$

$$[\nabla u_\epsilon]_n \to \nabla u_\epsilon \text{ in } L^1(\Omega_T) \text{ when } n \to \infty.$$

Properties (2.13), (2.14) and (2.15) can be deduced from Lemma B2 in [5]. See also [6].

Now we define for any $\epsilon > 0$, $n \in \mathbb{N}$, $0 < h << 1$ such that $0 < \frac{1}{n} < \epsilon < h < t_0$, the Steklov average of $[u_\epsilon]_n$

$$\phi_{\epsilon,h}^n(x,t) := ([u_\epsilon]_n)_h = \frac{1}{2h} \int_{t-h}^{t+h} [u_\epsilon]_n (x,\tau) d\tau$$

where we assume that $[u_\epsilon]_n (x,t) = u_\epsilon(x,t_0)$ for any $t \le 0$.

First let us assume that u is bounded. Then the function $\phi_{\epsilon,h}^n \in C^1(\bar{\Omega}_T)$ and $\phi_{\epsilon,h}^n(\cdot,t)_{|\partial\Omega} = 0$ for every $t \in [0,T]$, and so we can use it as a test function in (2.9) to get that for any $t \in [0,T]$,

$$\int_\Omega u\phi_{\epsilon,h}^n dx \Big|_0^t + \int_0^t \int_\Omega [-u(\phi_{\epsilon,h}^n)_s + \langle A(x,s,\nabla u), \nabla \phi_{\epsilon,h}^n \rangle] dx ds = \int_0^t \int_\Omega \langle f, \nabla \phi_{\epsilon,h}^n \rangle dx ds.$$

Set

$$I_{1,n}^{\epsilon,h} := \int_\Omega u\phi_{\epsilon,h}^n(x,t) dx - \int_\Omega u_0 \phi_{\epsilon,h}^n(x,0) dx - \int_0^t \int_\Omega u(\phi_{\epsilon,h}^n)_s dx ds.$$

For $0 < h < t$, by definition of $\phi_{\epsilon,h}^n$

$$
\begin{aligned}
I_{1,n}^{\epsilon,h} \quad &= \int_\Omega u(x,t)\left(\frac{1}{2h}\int_{t-h}^{t+h}[u_\epsilon]_n(x,\tau)d\tau\right)dx \\
&- \int_\Omega u_0(x)\left(\frac{1}{2h}\int_{-h}^{h}[u_\epsilon]_n(x,\tau)d\tau\right)dx \\
&-\frac{1}{2h}\int_0^t\int_\Omega u(x,\tau)\left([u_\epsilon]_n(x,\tau+h)-[u_\epsilon]_n(x,\tau-h)\right)dxd\tau. \\
&= J_{1,n}^{\epsilon,h}+J_{2,n}^{\epsilon,h}+J_{3,n}^{\epsilon,h}.
\end{aligned}
$$

Observe that

$$
J_{2,n}^{\epsilon,h} = -\int_\Omega u_0(x)\left(\frac{1}{2h}\int_{-h}^{0}u_\epsilon(x,t_0)\,d\tau+\frac{1}{2h}\int_0^h[u_\epsilon]_n(x,\tau)\,d\tau\right)dx
$$

$$
\begin{aligned}
J_{3,n}^{\epsilon,h} \quad &= -\frac{1}{2h}\int_t^{t+h}\int_\Omega u(x,\tau-h)[u_\epsilon]_n(x,\tau)\,dx\,d\tau \\
&-\frac{1}{2h}\int_h^t\int_\Omega u(x,\tau-h)[u_\epsilon]_n(x,\tau)\,dx\,d\tau \\
&+\frac{1}{2h}\int_0^h\int_\Omega u_\epsilon(x,t_0)u(x,\tau)\,dx\,d\tau \\
&+\frac{1}{2h}\int_h^t\int_\Omega [u_\epsilon]_n(x,\tau-h)u(x,\tau)\,dx\,d\tau
\end{aligned}
$$

From (2.13), if we pass to the limit first as $n\to+\infty$ and then as $\epsilon\to 0^+$ we get

$$
\begin{aligned}
\lim_{\epsilon\to 0+}\ \lim_{n\to+\infty} I_{1,n}^{\epsilon,h} &= \int_\Omega u(x,t)\left(\tfrac{1}{2h}\int_{t-h}^{t+h}u(x,\tau)d\tau\right)dx \\
&-\tfrac{1}{2}\int_\Omega u_0(x)u(x,t_0)\,dx-\tfrac{1}{2h}\int_\Omega u_0(x)\left(\int_0^h u(x,\tau)\,d\tau\right)dx \\
&-\tfrac{1}{2h}\int_t^{t+h}\int_\Omega u(x,\tau-h)u(x,\tau)\,dx\,d\tau \\
&+\tfrac{1}{2h}\int_0^h\int_\Omega u(x,t_0)u(x,\tau)\,dx\,d\tau.
\end{aligned}
\tag{2.16}
$$

Then by Lemma 2.1, passing to the limit in (2.16) as $h\to 0^+$, we get for any $t\in[0,T]$

$$
\lim_{h\to 0}\lim_{\epsilon\to 0}\lim_{n\to+\infty} I_{1,n}^{\epsilon,h} = \frac{1}{2}\int_\Omega |u(x,t)|^2dx-\frac{1}{2}\int_\Omega |u_0(x)|^2dx.
$$

Define

$$
I_{2,n}^{\epsilon,h} := \int_0^t\int_\Omega \langle A(x,s,\nabla u),\nabla\phi_{\epsilon,h}^n\rangle dxds,\quad I_{3,n}^{\epsilon,h} := \int_0^t\int_\Omega \langle f,\nabla\phi_{\epsilon,h}^n\rangle dxds.
$$

Then,

$$
\begin{aligned}
I_{2,n}^{\epsilon,h} &= \sum_{i=0}^{m} \int_0^t \int_\Omega \langle A(x,s,\nabla u), ([u_\epsilon^i]_n)_h(x,s)\nabla\eta_i\rangle dxds \\
&+ \sum_{i=0}^{m} \int_0^t \int_\Omega \langle A(x,s,\nabla u), \eta_i(x)([\nabla u_\epsilon^i]_n)_h(x,s)\rangle dxds.
\end{aligned}
$$

Observe that

$$
\begin{aligned}
|\langle A(x,s,\nabla u), \eta_i(x)([\nabla u_\epsilon^i]_n)_h(x,s)\rangle| &\le \tilde{G}(|A(x,s,\nabla u)|) + G(\eta_i(x)|([\nabla u_\epsilon^i]_n)_h|)(x,s) \\
&\le \mu^{\tilde{k}} G(|\nabla u(x,s)|) + G(\eta_i(x)|([\nabla u_\epsilon^i]_n)_h|)(x,s), \qquad (2.17)
\end{aligned}
$$

where we used condition (ii) and inequalities (2.3), (2.4), (2.5).

Let $0 < h \le s$. By the Fubini theorem

(2.18)

$$
\begin{aligned}
|\left([\nabla u_\epsilon^i]_n\right)_h|(x,s) &\le \frac{1}{2h}\int_{s-h}^{s+h}\left(e^{-nr}|\nabla u_\epsilon^i(x,t_0)| + n\int_0^r e^{n(\tau-r)}|\nabla u_\epsilon^i(x,\tau)|d\tau\right)dr \\
&= -\frac{1}{2hn}[e^{-n(s+h)} - e^{-n(s-h)}]|\nabla u_\epsilon^i(x,t_0)| + \frac{n}{2h}\int_{s-h}^{s+h}\left(\int_\tau^{s+h} e^{n(\tau-r)}|\nabla u_\epsilon^i(x,\tau)|dr\right)d\tau \\
&= \frac{1}{2hn}[e^{-n(s-h)} - e^{-n(s+h)}]|\nabla u_\epsilon^i(x,t_0)| + \frac{1}{2h}\int_{s-h}^{s+h}\left(1 - e^{-n((s+h)-\tau)}\right)|\nabla u_\epsilon^i(x,\tau)|d\tau \\
&\le \frac{1}{2hn}|\nabla u_\epsilon^i(x,t_0)| + \frac{1}{2h}\int_{s-h}^{s+h}|\nabla u_\epsilon^i(x,\tau)|d\tau.
\end{aligned}
$$

Inequality (2.18), the monotonicity of G and the Jensen inequality imply that for $(x,s) \in \Omega \times (h,t)$

$$
\begin{aligned}
&G(\eta_i|\left([\nabla u_\epsilon^i]_n)_h|\right)(x,s) \le G(|([\nabla u_\epsilon^i]_n)_h|)(x,s) \\
&= G\left[\frac{1}{2hn}|\nabla u_\epsilon^i(x,t_0)| + \frac{1}{2h}\int_{s-h}^{s+h}|\nabla u_\epsilon^i(x,\tau)|d\tau\right] \\
&\le \frac{1}{2}G\left(\frac{1}{hn}|\nabla u_\epsilon^i(x,t_0)|\right) + \frac{1}{2}G\left(2\frac{1}{2h}\int_{s-h}^{s+h}|\nabla u_\epsilon^i(x,\tau)|d\tau\right) \qquad (2.19) \\
&\le \frac{1}{2}G\left(\frac{1}{hn}|\nabla u_\epsilon^i(x,t_0)|\right) + 2^{k-1}\frac{1}{2h}\int_{s-h}^{s+h} G\left(|\nabla u_\epsilon^i(x,\tau)|\right)d\tau.
\end{aligned}
$$

When $0 < s < h$, again by the Fubini theorem, we get

$$
\begin{aligned}
&|\left(\left[\nabla u^i_\epsilon\right]_n\right)_h|(x,s) \leq \frac{1}{2h}\int_{s-h}^{0}|\nabla u^i_\epsilon(x,t_0)|dr \\
&+\frac{1}{2h}\int_0^{s+h}\left(e^{-nr}|\nabla u^i_\epsilon(x,t_0)| + n\int_0^r e^{n(\tau-r)}|\nabla u^i_\epsilon(x,\tau)|d\tau\right)dr \\
&=\frac{h-s}{2h}|\nabla u^i_\epsilon(x,t_0)| + \frac{1}{2hn}[1-e^{-n(s+h)}]|\nabla u^i_\epsilon(x,t_0)| \\
&+\frac{n}{2h}\int_0^{s+h}\left(\int_\tau^{s+h} e^{n(\tau-r)}|\nabla u^i_\epsilon(x,\tau)|dr\right)d\tau \\
&\leq \left(\frac{1}{2}+\frac{1}{2hn}\right)|\nabla u^i_\epsilon(x,t_0)| + \frac{1}{2h}\int_0^{s+h}\left(1-e^{-n((s+h)-\tau)}\right)|\nabla u^i_\epsilon(x,\tau)|d\tau \\
&\leq \left(\frac{1}{2}+\frac{1}{2hn}\right)|\nabla u^i_\epsilon(x,t_0)| + \frac{1}{2h}\int_{s-h}^{s+h}|\nabla u^i_\epsilon(x,\tau)|d\tau
\end{aligned}
\tag{2.20}
$$

Arguing as in (2.19), from (2.20) we get, for $(x,s)\in\Omega\times(0,h)$

$$
\begin{aligned}
&G(\eta_i|\left(\left[\nabla u^i_\epsilon\right]_n)_h|\right)(x,s) \\
&\leq \frac{1}{2}G\left(\left(1+\frac{1}{hn}\right)|\nabla u^i_\epsilon(x,t_0)|\right) + 2^{k-1}\frac{1}{2h}\int_{s-h}^{s+h}G\left(|\nabla u^i_\epsilon(x,\tau)|\right)d\tau.
\end{aligned}
\tag{2.21}
$$

Moreover, by the definition of $u^i_{\epsilon,h}$, the Fubini theorem and the Jensen inequality, we know that

$$
\left[G(|\nabla u^i_\epsilon|)\right]_h(x,s) \leq \left[G(|\nabla u|)\right]_{\epsilon,h}(x,s) \tag{2.22}
$$

Combining (2.17), (2.19) and (2.22), for almost every $(x,s)\in\Omega\times(h,t)$,

$$
\begin{aligned}
&|\langle A(x,s,\nabla u(x,s)),(\left[\nabla u^i_\epsilon\right]_n)_h(x,s)\eta_i(x)\rangle| \\
&\leq \mu^{\tilde{k}}G(|\nabla u(x,s))|) + \frac{1}{2}G\left(\frac{1}{hn}|\nabla u^i_\epsilon(x,t_0)|\right) + 2^{k-1}\left[G(|\nabla u|)\right]_{\epsilon,h}(x,s)
\end{aligned}
$$

On the other hand, combining (2.17), (2.21) and (2.22), for almost every $(x,s)\in\Omega\times(0,h)$,

$$
\begin{aligned}
&|\langle A(x,s,\nabla u(x,s)),\eta_i(x)(\left[\nabla u^i_\epsilon\right]_n)_h(x,s)\rangle| \\
&\leq \mu^{\tilde{k}}G(|\nabla u(x,s))|) + \frac{1}{2}G\left(\left(1+\frac{1}{hn}\right)|\nabla u^i_\epsilon(x,t_0)|\right) + 2^{k-1}\left[G(|\nabla u|)\right]_{\epsilon,h}(x,s)
\end{aligned}
$$

Last inequalities, the continuity of G and the Lebesgue dominated convergence theorem imply that

$$\lim_{h\to 0+}\lim_{\epsilon\to 0+}\lim_{n\to\infty}\int_h^t\int_\Omega \langle A(x,s,\nabla u(x,s)),\eta_i(x)(\left[\nabla u_\epsilon^i\right]_n)_h(x,s)\rangle dxds$$
$$=\int_0^t\int_\Omega \langle A(x,s,\nabla u(x,s)),\eta_i(x)\nabla u(x,s)\rangle dxds$$

and

$$\lim_{h\to 0+}\lim_{\epsilon\to 0+}\lim_{n\to\infty}\int_0^h\int_\Omega \langle A(x,s,\nabla u(x,s)),\eta_i(x)(\left[\nabla u_\epsilon^i\right]_n)_h(x,s)\rangle dxds=0.$$

Finally, observe that

$$\sum_{i=0}^m\int_0^t\int_\Omega \langle A(x,s,\nabla u(x,s)),\nabla u(x,s)\eta_i(x)\rangle dxds=\int_0^t\int_\Omega \langle A(x,s,\nabla u(x,s)),\nabla u(x,s)\rangle dxds.$$

Consequently,

(2.23)
$$\lim_{h\to 0+}\lim_{\epsilon\to 0+}\lim_{n\to\infty}\sum_{i=0}^m\int_0^t\int_\Omega \langle A(x,s,\nabla u(x,s)),\eta_i(x)([\nabla u_\epsilon^i]_n)_h(x,s)(x)\rangle dxds$$
$$=\int_0^t\int_\Omega \langle A(x,s,\nabla u(x,s)),\nabla u(x,s)\rangle dxds.$$

On the other hand, since $\tilde G$ is increasing, by the Fenchel inequality, (ii), (2.4), (2.5), (2.13) and the boundedness of u, for some $c_0>0$, for any $i\in\{0,1,...,m\}$ and every $(x,s)\in\Omega_T$ with $\nabla u(x,s)\neq 0$ we have

$$\begin{aligned}&|\langle A(x,s,\nabla u(x,s)),([u_\epsilon^i]_n)_h(x,s)\nabla\eta_i\rangle|\le |A(x,s,\nabla u(x,s))|\cdot|([u_\epsilon^i]_n)_h(x,s)\nabla\eta_i|\\ &\le \tilde G(|A(x,s,\nabla u(x,s))|)+G(|([u_\epsilon^i]_n)_h(x,s)\nabla\eta_i|)\le \tilde G\left(\mu\tfrac{G(|\nabla u(x,s)|)}{|\nabla u(x,s)|}\right)+c_0\\ &\le \mu^{\tilde k}\tilde G\left(\tfrac{G(|\nabla u(x,s)|)}{|\nabla u(x,s)|}\right)+c_0\le \mu^{\tilde k}G(|\nabla u(x,s)|)+c_0.\end{aligned}\tag{2.24}$$

Hence (2.10) implies that $|\langle A(\cdot,\cdot,\nabla u(\cdot,\cdot)),([u_\epsilon^i]_n(\cdot,\cdot))_h\nabla\eta_i(x)\rangle|$ is integrably bounded by a function not depending on i, ϵ, n, h.

Therefore

$$\lim_{h\to 0+}\lim_{\epsilon\to 0+}\lim_{n\to\infty}\int_0^t\int_\Omega \langle A(x,s,\nabla u(x,s)),(\left[u_\epsilon^i\right]_n(x,s))_h\nabla\eta_i(x)\rangle dxds$$
$$=\int_0^t\int_\Omega \langle A(x,s,\nabla u(x,s)),u(x,s)\nabla\eta_i\rangle dxds.$$

Finally, observe that

$$\sum_{i=0}^m\int_0^t\int_\Omega \langle A(x,s,\nabla u(x,s)),u(x,s)\nabla\eta_i(x)\rangle\,dx\,ds=0.$$

Consequently,

$$\lim_{h\to 0+}\lim_{\epsilon\to 0+}\lim_{n\to\infty}\sum_{i=0}^m\int_0^t\int_\Omega \langle A(x,s,\nabla u(x,s)),(\left[u_\epsilon^i\right]_n)_h(x,s)\nabla\eta_i(x)\rangle dx\,ds=0.\tag{2.25}$$

From (2.23) and (2.25) we get

$$\lim_{h\to 0+}\lim_{\epsilon\to 0+}\lim_{n\to\infty} I^{\epsilon,h}_{2,n} = \int_0^t\int_\Omega \langle A(x,s,\nabla u),\nabla u\rangle dxds.$$

In the same way, by using (2.3), (2.5) and (2.6), we can prove that

$$\lim_{h\to 0+}\lim_{\epsilon\to 0+}\lim_{n\to\infty} I^{\epsilon,h}_{3,n} = \int_0^t\int_\Omega \langle f,\nabla u\rangle dxds.$$

If u is not bounded, define for any $j\in\mathbb{N}$, $u_j = \frac{1}{2}(|u+j|-|u-j|)$. Then,

$$u_j \in L^\infty(0,T;L^2(\Omega))\cap L^1(0,T;W^{1,1}_0(\Omega))\cap L^\infty(\Omega_T).$$

We can constuct $(u_j)_\epsilon$ as before and we can apply the above arguments by using $(\phi_j)^n_{\epsilon,h}$ as test function in (2.9) Finally we can pass to the limit as $j\to+\infty$ since $|u_j|\le|u|$ a.e. □

The uniqueness of a weak solution follows under the additional monotonicity condition

$$\langle A(x,s,\xi)-A(x,s,\eta),\xi-\eta\rangle \ge 0 \tag{2.26}$$

for a.e. $(x,s)\in\Omega_T$ and for any $\xi,\eta\in\mathbb{R}^n$.

Proposition 2.2. *Assume* (2.2), (2.6) *and* (2.26). *Let* $u:\overline{\Omega}_T\to\mathbb{R}$ *be a weak solution to* (2.1). *Then* u *is unique.*

Proof. The proof proceeds as in the homogeneous case. Let $v:\overline{\Omega}_T\to\mathbb{R}$ be a weak solution to (2.1). Apart of the approximation argument as in Proposition 2.1, we can use $w=u-v$ as test function in the definition of solutions u and v.

Then, by subtracting we get for any $0<t\le T$

$$\frac{1}{2}\|w(t)\|^2_{L^2(\Omega)} + \int_0^t\int_\Omega \langle A(x,s,\nabla u)-A(x,s,\nabla v),\nabla w\rangle dxds = 0.$$

Hence, (2.26) yields $u=v$ a.e. in Ω_T.

□

3. A Version of the Gronwall Lemma

For arbitrary non-empty closed sets in D', $D\subset\mathbb{R}^n$, define

$$d_D(D') \;:=\; \inf\{\beta>0\,|\,D'\subset D+\beta\mathcal{B}\}\ .$$

where B denotes the closed unit ball in $\mathbb{R}^n$. Consider reals $0\le t_0<T$ and a multifunction $P:[t_0,T]\rightsquigarrow\mathbb{R}^n$ with non-empty closed values. P is called *left absolutely continuous* if the following condition is satisfied: given any $\epsilon>0$ there exists $\delta>0$ such that, for any finite partition

$$t_0\le s_1<\tau_1\le s_2<\tau_2\le\ldots\le s_m<\tau_m\le T$$

satisfying $\sum_{i=1}^m(\tau_i-s_i)\;<\;\delta$, we have

$$\sum_{i=1}^m d_{P(\tau_i)}(P(s_i))\;<\;\epsilon\ .$$

For any $y \in P(t)$ the contingent derivative $DP(t,y)(1)$ is the set of all $v \in \mathbb{R}^n$ such that there exist $h_i \to 0+$ and $v_i \to v$ satisfying $y + h_i v_i \in P(t+h_i)$.

A function $\gamma : [t_0, T] \to \mathbb{R}_+$ is called upper semicontinuous from the right, if for every $t \in [t_0, T)$ we have $\limsup_{s\to t+} \gamma(s) \le \gamma(t)$. Similarly, it is called lower semicontinuous from the left, if for every $t \in (t_0, T]$ we have $\liminf_{s\to t-} \gamma(s) \ge \gamma(t)$.

Lemma 3.1. *Consider a Carathéodory function $\psi : [t_0, T] \times \mathbb{R} \to \mathbb{R}_+$ such that for every $r > 0$ there exists $k_r \in L^1(t_0, T; \mathbb{R}_+)$ satisfying for a.e. $t \in [t_0, T]$*

$$\sup_{|x|\le r} \psi(t,x) \le k_r(t).$$

Let $g \in L^1(t_0, T; \mathbb{R})$ and $\gamma : [t_0, T] \to \mathbb{R}_+$ be measurable, bounded and satisfy

$$\gamma(t_2) - \gamma(t_1) + \int_{t_1}^{t_2} \psi(t, \gamma(t))dt \le \int_{t_1}^{t_2} g(t)dt \qquad t_0 \le t_1 \le t_2 \le T. \tag{3.1}$$

Then there exists a solution $x(\cdot) \in W^{1,1}([t_0, T])$ of the Cauchy problem

$$\begin{cases} x'(t) = -\psi(t, x(t)) + g(t), & \text{a.e.} \\ x(t_0) = \gamma(t_0) \end{cases} \tag{3.2}$$

such that $\gamma(t) \le x(t)$ for all $t \in [t_0, T]$.

Furthermore, if $g \in L^1(t_0, \infty; \mathbb{R})$, ψ is defined on $[t_0, +\infty) \times \mathbb{R}$, $\gamma : [t_0, +\infty) \to \mathbb{R}_+$ is measurable and locally bounded and for every $T > t_0$ the above assumptions hold true, then there exists a solution x to (3.2) defined on $[t_0, \infty)$ such that $\gamma \le x$. In particular,

$$\limsup_{t\to\infty} \gamma(t) \le \limsup_{t\to\infty} x(t).$$

Proof. By (3.1), γ is upper semicontinuous from the right and lower semicontinuous from the left. Define the set-valued map $P : [t_0, T] \rightsquigarrow \mathbb{R}$ by

$$P(t) = \begin{cases} \gamma(t) + \mathbb{R}_+ & t \in [t_0, T) \\ \mathbb{R} & t \ge T. \end{cases}$$

By our assumptions for all $t_0 \le t_i < \tau_i < T$

$$\gamma(\tau_i) - \int_{t_i}^{\tau_i} g(t)dt \le \gamma(t_i).$$

This implies that $P(\cdot)$ is left absolutely continuous. Let $t \in (t_0, T)$ be a Lebesgue point of $\psi(\cdot, \gamma(\cdot))$ and of $g(\cdot)$. The set A of all such points is of full Lebesgue measure in $[t_0, T]$. By the assumptions of our lemma for every $t \in A$,

$$D\gamma(t) := \liminf_{\varepsilon\to 0+} \frac{\gamma(t+\varepsilon) - \gamma(t)}{\varepsilon} \le -\psi(t, \gamma(t)) + g(t).$$

This and the definition of $P(\cdot)$ imply that $-\psi(t, \gamma(t)) + g(t) \in DP(t, \gamma(t))(1)$. Since γ is upper semicontinuous from the right, it enjoys the following property: if $y > \gamma(t)$, then for every $r \in \mathbb{R}$ and all small $\varepsilon > 0$ we have $y + \varepsilon r > \gamma(t+\varepsilon)$. Therefore, $DP(t,y)(1) = \mathbb{R}$ for any $y > \gamma(t)$. Consequently for all $t \in A$ and $y \in P(t)$ we have $-\psi(t,y) + g(t) \in DP(t,y)(1)$. By [13, Theorem 4.2] and our assumptions, for some $S \in (t_0, T]$ there exists a solution of (3.2) satisfying $x(t) \in P(t)$ on $[t_0, S)$. Therefore $0 \le \gamma(t) \le x(t)$ for all $t \in [t_0, S)$. Furthermore, $x(t) \le x(t_0) + \int_{t_0}^t g(s)ds$ for all $t \in [t_0, S)$ implying

that $x(\cdot)$ is bounded. Hence, by our assumptions and since $x(\cdot)$ is a solution of (3.2), it can be extended, by continuity, on $[t_0, S]$. Furthermore, $\gamma(S) \leq x(S)$ because γ is lower semicontinuous from the left. Applying the Zorn lemma and [13, Theorem 4.2], we extend this solution $x(\cdot)$ to the time interval $[t_0, T]$.

To prove the last statement assume that $g \in L^1(t_0, \infty; \mathbb{R})$, ψ is defined on $[t_0, +\infty) \times \mathbb{R}$, γ is locally bounded on $[t_0, +\infty)$ and for every $T > t_0$ the assumptions of Lemma hold true. Then the same construction allows one to get $x(\cdot)$ defined on $[t_0, +\infty)$ satisfying for every $t > \tau \geq t_0$,

$$0 \leq \gamma(t) \leq x(t) \leq x(\tau) + \int_\tau^t g(s)ds.$$

Since $\lim_{\tau\to\infty} \int_\tau^\infty g(s)ds = 0$ the proof follows. □

Remark 3.1. *Indeed, assume that the solution $x : [t_0, T] \to \mathbb{R}_+$ of (3.2) is unique. Then for $\gamma = x$ we have*

$$\gamma(t_2) - \gamma(t_1) + \int_{t_1}^{t_2} \psi(t, \gamma(t))dt = \int_{t_1}^{t_2} g(t)dt \qquad t_0 \leq t_1 \leq t_2 \leq T.$$

That is (3.1) holds true. The estimate $\gamma \leq x$ obtained in Lemma 3.1 can not be further improved.

Corollary 3.1. *Under all the assumptions of Lemma 3.1 suppose that $\psi(t, a) = 0$ for all $a \leq 0$, $g(\cdot) \geq 0$, $\psi(t, \cdot)$ is increasing for a.e. $t \in [t_0, T]$ and that for any $R > r > 0$ there exists $k_{R,r} \in L^1(t_0, T; \mathbb{R}_+)$ satisfying for a.e. $t \in [t_0, T]$*

$$|\psi(t, x) - \psi(t, y)| \leq k_{R,r}(t)|x - y| \quad \forall\, x, y \in [r, R].$$

Then the solution $z(\cdot)$ of

$$\begin{cases} z'(t) = -\psi(t, z(t)), & \text{a.e.} \\ z(t_0) = \gamma(t_0) \end{cases} \tag{3.3}$$

is unique and well defined on $[t_0, T]$, $z(\cdot) \geq 0$ and $\gamma(t) \leq x(t) \leq z(t) + \int_{t_0}^t g(s)ds$ for all $t \in [t_0, T]$.

Proof. Clearly, if $z(\cdot)$ is a solution of (3.3), then it is decreasing on its interval of existence. Moreover, if for some $t \geq t_0$, $z(t) = 0$, then $z(s) = 0$ for any $t \leq s \leq T$. This and the Lipschitz continuity assumption imply the existence and uniqueness on the whole time interval $[t_0, T]$. Consider any solution $x(\cdot)$ of (3.2) defined on $[t_0, T]$. Then it is not difficult to realize that $x(\cdot) \geq 0$.

We claim that $z \leq x$ on $[t_0, T]$. Indeed otherwise we can find $t_0 \leq s_0 < t_1 < T$ such that $z(s_0) = x(s_0)$ and $z(s) > x(s)$ on $[s_0, t_1]$. Then $-\psi(t, z(t)) \leq -\psi(t, x(t))$ for a.e. $t \in [s_0, t_1]$. Thus $x'(s) - z'(s) \geq g(s) \geq 0$ a.e. in $[s_0, t_1]$. This implies that $x(t) \geq z(t)$ for $t \in [s_0, t_1]$ leading to a contradiction. Let $x(\cdot)$ be as in Lemma 3.1. Observe that for a.e. $t \in [t_0, T]$ we have

$$x'(t) = -\psi(t, x(t)) + g(t) \leq -\psi(t, z(t)) + g(t) = z'(t) + g(t).$$

Integrating we obtain that $x(t) \leq z(t) + \int_{t_0}^t g(s)ds$ for all $t \in [t_0, T]$. This and Lemma 3.1 yield $\gamma(t) \leq x(t) \leq z(t) + \int_{t_0}^t g(s)ds$. □

Corollary 3.2. *Consider $g \in L^1(t_0, \infty; \mathbb{R}_+)$, a Carathéodory function $\psi : [t_0, \infty) \times \mathbb{R} \to \mathbb{R}_+$, a measurable locally bounded $\gamma : [t_0, \infty) \to \mathbb{R}_+$ and let assumptions of Corollary 3.1 hold true on $[t_0, T]$ for every $T > t_0$. If for every $r > 0$ there exists $\alpha_r > 0$ such that $\inf_{x \geq r} \psi(t, x) \geq \alpha_r$ for a.e. $t \geq t_0$, then $\lim_{t\to\infty} \gamma(t) = 0$.*

Proof. Fix any sequence $t_i \to \infty$ and consider the solutions z_i to (3.3) with t_0 replaced by t_i. By Corollary 3.1 applied with t_0 replaced by t_i, we know that $z_i \geq 0$. Notice that z_i is decreasing for every i. Observe next that if for some i we have $\lim_{t\to\infty} z_i(t) =: r > 0$, then, $z_i(t) \geq r$ for every $t \geq t_i$ and

$$0 \leq z_i(t) = z_i(t_i) - \int_{t_i}^{t} \psi(s, z_i(s))ds \leq z_i(t_i) - (t - t_i)\alpha_r.$$

When t is sufficiently large, the right-hand side of the above expression becomes negative. Thus implies that $\lim_{t\to\infty} z_i(t) = 0$ for all i. This and Corollary 3.1 imply that for every i, every $\varepsilon > 0$ and all large t we have

$$\gamma(t) \leq \frac{\varepsilon}{2} + \int_{t_i}^{t} g(s)ds.$$

Taking i sufficiently large, we deduce that for every $\varepsilon > 0$ and all large t we have $\gamma(t) \leq \varepsilon$ completing the proof. □

In [25] a particular instance of $\psi(t, x)$ was considered, namely $\psi(t, x) = M|x|^{1+\nu}$, where M and ν are non negative constants. It was shown then that for any continuous $\gamma : [t_0, T] \to \mathbb{R}_+$ satisfying the inequality (3.1) we have

$$\gamma(t) \leq y(t) := \frac{\gamma(t_0)}{[1 + \nu M\gamma(t_0)^\nu (t - t_0)]^{\frac{1}{\nu}}} + \Gamma(t) \quad \forall\, t \in [t_0, T] \tag{3.4}$$

for every choice of continuous function $\Gamma : [t_0, T] \to \mathbb{R}_+$ satisfying

$$\Gamma(t_1) \leq \Gamma(t_2) + M \int_{t_1}^{t_2} \Gamma(t)^{1+\nu} dt - \int_{t_1}^{t_2} g(t)dt \qquad t_0 \leq t_1 \leq t_2 \leq T, \tag{3.5}$$

where $g \in L^1([t_0, T], \mathbb{R})$. Let Γ, M, ν be as above and x be the solution of

$$\begin{cases} x' = -M|x|^{1+\nu} + g(t), & \text{a.e. } t > t_0 \\ x(t_0) = \gamma(t_0) \end{cases} \tag{3.6}$$

We already know that $x(t) \geq \gamma(t) \geq 0$ for all $t \in [t_0, T]$.

We claim that $x(t) \leq y(t)$, which implies that our result provides a better estimate on the time behavior of γ than the one from [25]. Indeed, since $y(t_0) = \gamma(t_0) + \Gamma(t_0) \geq x(t_0)$, if for some $t_1 \in (t_0, T]$ we have $y(t_1) < x(t_1)$, then we can find $s_0 \in [t_0, t_1]$ such that $x(s_0) = y(s_0)$ and $y(s) < x(s)$ for all $s \in (s_0, t_1]$.

Note that $-y(\cdot)$ is continuous on $[t_0, T)$. By the same arguments as in the proof of Lemma 3.1, we check that $t \rightsquigarrow -\Gamma(t) + \mathbb{R}_+$ is left absolutely continuous on $[t_0, T]$ and that the map

$$t \mapsto \frac{\gamma(t_0)}{[1 + \nu M\gamma(t_0)^\nu (t - t_0)]^{\frac{1}{\nu}}}$$

is Lipschitz. Therefore $t \rightsquigarrow P(t) = -y(t) + \mathbb{R}_+$ is left absolutely continuous on $[t_0, T]$. Furthermore, for a.e. $s \in (s_0, t_1)$ we have

$$D(-\Gamma)(s) \leq M\Gamma(s)^{1+\nu} - g(s).$$

Hence, for a.e. $s \in (s_0, t_1)$,

$$\begin{aligned} D(-y)(s) &\leq \frac{M\gamma(t_0)^{1+\nu}}{[1+\nu M\gamma(t_0)^\nu(t-t_0)]^{\frac{1+\nu}{\nu}}} + M\Gamma(s)^{1+\nu} - g(s) \leq My(s)^{1+\nu} - g(s) \\ &= My(s)^{1+\nu} - Mx(s)^{1+\nu} - x'(s) \leq -x'(s). \end{aligned}$$

Therefore, for a.e. $s \in [s_0, t_1]$ and all $z \in P(s)$ we have $-x'(s) \in DP(s, z)(1)$. From [13, Theorem 4.2] applied to the single-valued map $s \mapsto \{-x'(s)\}$ we deduce that $-x(t) \in P(t)$ for all $t \in [s_0, t_1]$. This yields $x(t) \leq y(t)$ on $[s_0, t_1]$ contradicting the choice of t_1.

It was observed in [25] that if $g \geq 0$, then $\Gamma(t) = \int_{t_0}^t g(s)ds$ verifies (3.5) and that for any continuous $\gamma : [t_0, T] \to \mathbb{R}_+$ satisfying the inequality (3.1), the estimate (3.4) holds true for this choice of Γ. In particular, if $g \in L^1(t_0, \infty; \mathbb{R}_+)$, then (3.4) is verified with this Γ and t_0 replaced by any $\tau > t_0$ and $T > \tau$. This implies that $\lim_{t\to\infty} \gamma(t) = 0$.

Define $\psi(t, x) = M|x|^{1+\nu}$ for $x \geq 0$ and $\psi(t, x) = 0$ for $x < 0$ and observe that it satisfies the assumptions of Corollary 3.2. Thus the equality $\lim_{t\to\infty} \gamma(t) = 0$ is a consequence of Corollary 3.2.

4. The main results

In this section we assume that the function $G\colon \mathbb{R}_+ \to \mathbb{R}_+$ in (2.2) is continuous, convex and satisfies (1.2) for some $k > 0$ and $p \geq 2$.

We state next the main result of this paper.

Theorem 4.1. *Assume* (2.2) *and* (2.6). *Let* $u\colon \overline{\Omega}_T \to \mathbb{R}$ *be a weak solution to* (2.1). *Then, for any* $t \in [0, T]$,

$$\|u(\cdot, t)\|^2_{L^2(\Omega)} \leq 2|\Omega|x(t)$$

where $x(\cdot)$ *is the unique solution of the problem*

$$\begin{cases} \dot{x}(t) = -c\, G(\sqrt{x(t)}\,) + g(t) \\ x(0) = \frac{1}{2|\Omega|}\|u_0\|^2_{L^2(\Omega)} \end{cases} \tag{4.1}$$

for some $c > 0$ *depending only on* p, k, N, $|\Omega|$, ν *and* $g(t) = \frac{1}{|\Omega|}\|\tilde{G}(\frac{2}{\nu}|f(\cdot, t)|)\|_{L^1(\Omega)}$.

If $f : \Omega \to \mathbb{R}$ is a integrable function on Ω, we set

$$\fint_\Omega |f|dx = \frac{1}{|\Omega|}\int_\Omega |f|dx$$

where $|\Omega| > 0$ denotes the Lebesgue measure of Ω.

Proof. For $0 < t_1 < t_2 < T$, applying Proposition 2.1 at times t_1 and t_2 and subtracting the results we get

$$\frac{1}{2}\|u(t_2)\|^2_{L^2(\Omega)} - \frac{1}{2}\|u(t_1)\|^2_{L^2(\Omega)} + \int_{t_1}^{t_2}\int_\Omega \langle A(x, s, \nabla u), \nabla u\rangle dxds = \int_{t_1}^{t_2}\int_\Omega \langle f, \nabla u\rangle dxds.$$

By (i) and Fenchel's inequality (2.3) we get

(4.2)
$$\frac{1}{2}\|u(t_2)\|^2_{L^2(\Omega)} - \frac{1}{2}\|u(t_1)\|^2_{L^2(\Omega)} + \frac{\nu}{2}\int_{t_1}^{t_2}\int_\Omega G(|\nabla u|)dxds \le \int_{t_1}^{t_2}\int_\Omega \tilde{G}\left(\frac{2}{\nu}|f(x,s)|\right)dxds.$$

Applying Hölder inequality we obtain, since $p \ge 2$,

$$\left(\fint_\Omega |u|^2 dx\right)^{\frac{p}{2}} \le c\fint_\Omega |\nabla u|^p dx \tag{4.3}$$

where $c > 0$ depends only on p, N, $|\Omega|$.

Now, let F be the convex function as in the claim of Lemma 2.2. Since F is positive and $F(0) = 0$, it is nondecreasing on $\mathbb{R}_+$. From (4.3) we deduce that

$$F\left(\left(\fint_\Omega |u|^2 dx\right)^{\frac{p}{2}}\right) \le F\left(c\fint_\Omega |\nabla u|^p dx\right). \tag{4.4}$$

In what follows C denotes a positive constant depending only on N, p, k, $|\Omega|$, which may vary from line to line. From (2.5), (2.12), (4.4) and the Hölder and Jensen inequalities we deduce that for any $t \in [0, T]$,

$$G\left(\left(\frac{1}{2}\fint_\Omega |u(x,t)|^2 dx\right)^{\frac{1}{2}}\right) \le G\left(\left(\fint_\Omega |u(x,t)|^2 dx\right)^{\frac{1}{2}}\right) \le C\fint_\Omega G(|\nabla u|)dx. \tag{4.5}$$

Going back to (4.2), we deduce from (4.5) the following inequality

$$\begin{aligned} &\frac{1}{2}\fint_\Omega |u(x,t_2)|^2 dx - \frac{1}{2}\fint_\Omega |u(x,t_1)|^2 dx + \frac{\nu}{2C}\int_{t_1}^{t_2} G\left[\left(\frac{1}{2}\fint_\Omega |u(x,t)|^2 dx\right)^{1/2}\right] dt \\ &\le \frac{1}{|\Omega|}\int_{t_1}^{t_2}\int_\Omega \tilde{G}\left(\frac{2}{\nu}|f(x,s)|\right)dxds. \end{aligned} \tag{4.6}$$

Now we are in position to apply Lemma 3.1 with

$$\gamma(t) =: \frac{1}{2}\fint_\Omega |u(x,t)|^2 dx$$

and

$$\psi(t,x) =: \frac{\nu}{2C}\, G(\sqrt{x}\,).$$

By the measurable viability theorem, cf. [13], there exists a solution $x(\cdot) : \mathbb{R}_+ \to \mathbb{R}_+$ of (4.1). We claim that it is unique. Indeed, for positive reals $z \ge y$, we have $\sqrt{z} \ge \sqrt{y}$ and $\psi(t,z) \ge \psi(t,y)$. Hence $(\psi(t,z) - \psi(t,y))(z-y) \ge 0$. Similarly, it can be verified that if $z \le y$, the same inequality holds true.

Consider any solution $y : [t_0, T] \to \mathbb{R}_+$ of (4.1). Then

$$\frac{d}{dt}\frac{1}{2}|x-y|^2(t) = -(\psi(t,x(t)) - \psi(t,y(t)))(x(t)-y(t)) \le 0$$

and therefore $x = y$ on $[t_0, T]$.

Setting $\psi(t,a) = 0$ for $a < 0$, we may apply Lemma 3.1 and so to get the desired estimate.

□

Corollaries 3.1 and 3.2 yield the following result.

Theorem 4.2. *In Theorem 4.1 assume that $g \in L^1([0,\infty))$. Let $u\colon \overline{\Omega} \times [0,\infty) \to \mathbb{R}$ be a weak solution to* (2.1)*. Then, for any $t_0 \geq 0$ and $t > t_0$,*

$$\|u(t,\cdot)\|^2_{L^2(\Omega)} \leq 2|\Omega| x(t) \leq 2|\Omega| \left(z(t) + \int_{t_0}^t g(s)ds \right)$$

where $z(\cdot)$ is the unique solution of the problem

$$\begin{cases} \dot{z}(t) = -c\, G(\sqrt{z(t)}\,) \\ z(t_0) = \frac{1}{2|\Omega|} \|u(\cdot,t_0)\|_{L^2(\Omega)} \end{cases}$$

for some $c > 0$ depending only on p, k, N, $|\Omega|$, ν.

Furthermore, $\lim_{t\to\infty} \|u(t,\cdot)\|^2_{L^2(\Omega)} = 0$.

Remark 4.1. *Theorem 1.1 follows trivially from Theorem 4.1 and Theorem 4.2.*

Remark 4.2. *If $G(s) = \frac{s^p}{p}$, $p \geq 2$,, as observed in Section 3, the statement of Theorem 4.1 improves the behaviour in time of a solution to Problem* (2.1) *for a p-Laplace type operator (see* [25] *and the reference therein). For istance, when $p = 2$, Theorem 4.1 implies that for any $t \in [0,T]$*

$$\|u(t)\|^2_{L^2(\Omega)} \leq e^{-\frac{c}{2}t} \|u_0\|^2_{L^2(\Omega)} + 2 \int_0^t e^{-\frac{c}{2}(t-s)} \|f(\cdot,s)\|^2_{L^2(\Omega)} ds$$

where c is the constant in Problem (4.1)*.*

We conclude this section by an example.

Let us consider for simplicity problem (1.1) with

$$G(s) = (s^2+1)\log(1+s^2).$$

Then, if u is the weak solution of problem (1.1), we get that

$$\|u(\cdot,t)\|^2_{L^2(\Omega} \leq 2|\Omega| \left(\exp\left(\log(1 + \frac{1}{2|\Omega|} \|u_0\|_{L^2(\Omega)}) e^{-ct} \right) - 1 \right) \tag{4.7}$$

for some $c > 0$ and all $t \geq 0$.

Estimate (4.7) is new and cannot be achieved by the arguments of [25].

5. Solutions of an anisotropic problem

In this section we consider the following anisotropic problem

$$\begin{cases} u_t - \sum_{i=1}^N \frac{\partial}{\partial x_i} \left(|u_{x_i}|^{q_i-2} u_{x_i} \right) = 0 & \text{in } \Omega_T, \\ u(x,t) = 0 \qquad \text{on } \partial\Omega \times (0,T), \\ u(x,0) = u_0(x) \qquad \text{in } \Omega, \end{cases} \tag{5.1}$$

where $u_0 \in L^2(\Omega)$ and $1 < q_i$, $i \in \{1, ..., N\}$.

Here, for any $\xi \in \mathbb{R}^N$ the vector field is $A(\xi) =: \sum_{i=1}^N |\xi_i|^{q_i-2}\xi_i$.

In [8] the authors prove that there exists a unique function $u \in C([0,T], L^2(\Omega)) \cap L^2\left(0, T, W_0^{1,1}(\Omega)\right)$ with

$$\int_0^T \int_\Omega \sum_{i=1}^N |u_{x_i}|^{q_i}\, dxdt < +\infty$$

that solves (5.1). Moreover, for every $t \in (0, T]$ we have

$$\frac{1}{2}\|u(t)\|^2_{L^2(\Omega)} + \int_0^t \int_\Omega \sum_{i=1}^N |u_{x_i}|^{q_i}\, dxdt = \frac{1}{2}\|u_0\|^2_{L^2(\Omega)}. \tag{5.2}$$

The proof of (5.2) is similar to the one of Proposition 2.1. In this case the mollifier in time of test functions is not necessary and we can argue directly on the convex vector field $A(\xi)$. (See also [8]).

Set

$$\frac{1}{\bar{q}} =: \frac{1}{N}\sum_{i=1}^N \frac{1}{q_i}, \qquad q =: \max\{q_i, i = 1, ..., N\}$$

$$p =: \min\{q_i, i = 1, ..., N\}$$

and

$$B(t) =: \max\{1, \|u_{x_i}(t)\|_{L^{q_i}(\Omega),\ i=1,...,N}\}.$$

Theorem 5.1. *Assume that $2 \le q$ and $\frac{2N}{N+2} < \bar{q} < N$. Let $u : \overline{\Omega}_T \to \mathbb{R}$ be a weak solution to problem (5.1). Then for any $t \in [0, T]$,*

$$\|u(\cdot, t)\|^2_{L^2(\Omega)} \le 2x(t)$$

where $x(t)$ is the unique solution of the problem

$$\begin{cases} x'(t) = -c\, B(t)^{p-q} x(t)^{\frac{q}{2}} \\ x(0) = \frac{1}{2}\|u_0\|^2_{L^2(\Omega)}. \end{cases} \tag{5.3}$$

for some $c > 0$ depending only on $|\Omega|$, $\bar{q}$ and N.

Proof. By (5.2) for any $0 \le t_1 < t_2 \le T$

$$\frac{1}{2}\|u(t_2)\|^2_{L^2(\Omega)} - \frac{1}{2}\|u(t_1)\|^2_{L^2(\Omega)} + \int_{t_1}^{t_2} \int_\Omega \sum_{i=1}^N |u_{x_i}|^{q_i}\, dxdt = 0 \tag{5.4}$$

We claim that, for a.e. $t \in (0, T)$,

$$\left(\int_\Omega |u|^{\bar{q}^*}\, dx\right)^{\frac{q}{\bar{q}^*}} \le cB(t)^{q-p} \int_\Omega \sum_{i=1}^N |u_{x_i}|^{q_i}\, dx, \tag{5.5}$$

where c depends on $N, |\Omega|, \bar{q}$.

In fact, by the definition of $B(t)$, since $q_i \geq p$ and $B(t) \geq 1$, for any $i = 1, \dots, N$ we have

$$\|u_{x_i}\|^q_{L^{q_i}(\Omega)} \leq B(t)^{q-q_i} \|u_{x_i}\|^{q_i}_{L^{q_i}(\Omega)} \leq B(t)^{q-p} \|u_{x_i}\|^{q_i}_{L^{q_i}(\Omega)}$$

The convexity of the real function $t \mapsto t^q$ implies

$$\left(\frac{1}{N} \sum_{i=1}^N \|u_{x_i}\|_{L^{q_i}(\Omega)} \right)^q \leq \frac{1}{N} \sum_{i=1}^N \|u_{x_i}\|^q_{L^{q_i}(\Omega)}$$

Then, since $\frac{2N}{N+2} < \bar{q} < N$, we can apply the anisotropic Sobolev–type inequality [27] to obtain

$$\|u\|^q_{L^{\bar{q}^*}(\Omega)} \leq c_0^q \left(\frac{1}{N} \sum_{i=1}^N \|u_{x_i}\|_{L^{q_i}(\Omega)} \right)^q \leq c_0^q N^{-1} B(t)^{q-p} \sum_{i=1}^N \|u_{x_i}\|^{q_i}_{L^{q_i}(\Omega)}$$

where c_0 depends on $N, |\Omega|, \bar{q}$. Then, (5.5) follows.

The elliptic version of relation (5.5) has been proved in [4].

Now, from (5.5) we get

$$B(t)^{p-q} \left(\int_\Omega |u|^2 \, dx \right)^{\frac{q}{2}} \leq B(t)^{p-q} |\Omega|^{\frac{q}{2} - \frac{q}{\bar{q}^*}} \left(\int_\Omega |u|^{\bar{q}^*} \, dx \right)^{\frac{q}{\bar{q}^*}} \leq c_1 \int_\Omega \sum_{i=1}^N |u_{x_i}|^{q_i} \, dx \tag{5.6}$$

where c_1 depends on $N, |\Omega|, \bar{q}$.

Integrating (5.5) on (t_1, t_2), from the relations (5.4) and (5.6) we get

$$\frac{1}{2} \|u(t_2)\|^2_{L^2(\Omega)} - \frac{1}{2} \|u(t_1)\|^2_{L^2(\Omega)} + c \int_{t_1}^{t_2} B(t)^{p-q} \left(\frac{1}{2} \int_\Omega |u(x,t)|^2 \, dx \right)^{\frac{q}{2}} dt \leq 0 \tag{5.7}$$

for some $c > 0$ depending only on N, $|\Omega|$, $\bar{q}$, q. Now if $\gamma(t) = \frac{1}{2} \|u(t)\|^2_{L^2(\Omega)}$ we can rewrite (5.7) as

$$\gamma(t_2) - \gamma(t_1) + \int_{t_1}^{t_2} \psi(t, \gamma(t)) \, dt \leq 0$$

where

$$\psi(t,x) = cB(t)^{p-q} x^{\frac{q}{2}} \ \ \forall\, x \geq 0, \qquad \psi(t,x) = 0 \ \ \forall, x < 0.$$

We apply Corollary 3.1 to achieve the announced result.

□

Remark 5.1. *Let $\beta : [0,T] \to (1, +\infty)$ be a measurable function. The statement of Theorem 5.1 still holds if we replace the function $B(t)$, $t \in [0,T]$, with the function*

$$\tilde{B}(t) = \max\{\|u_{x_i}(t)\|_{q_i}, \beta(t) \colon \ i = 1, \dots, N\}.$$

Indeed the above proof can be applied as well.

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