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1 **The sensitivity of Southeast Pacific heat distribution to local and remote**
2 **changes in ocean properties**

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ABSTRACT

16 The Southern Ocean features ventilation pathways that transport surface wa-
17 ters into the subsurface thermocline on timescales of decades to centuries,
18 sequestering anomalies of heat and carbon away from the atmosphere and
19 thereby regulating the rate of surface warming. Despite its importance for
20 climate sensitivity, the factors that control the distribution of heat along these
21 pathways are not well understood. In this study, we use an observationally-
22 constrained, physically consistent global ocean model to examine the sensi-
23 tivity of heat distribution in the recently ventilated subsurface Pacific (RVP)
24 sector of the Southern Ocean to changes in ocean temperature and salinity.
25 First, we define the RVP using numerical passive tracer release experiments
26 that highlight the ventilation pathways. Next, we use an ensemble of adjoint
27 sensitivity experiments to quantify the sensitivity of the RVP heat content to
28 changes in ocean temperature and salinity. In terms of sensitivities to surface
29 ocean properties, we find that RVP heat content is most sensitive to anomalies
30 along the Antarctic Circumpolar Current (ACC), upstream of the subduction
31 hotspots. In terms of sensitivities to subsurface ocean properties, we find that
32 RVP heat content is most sensitive to basin-scale changes in the subtropi-
33 cal Pacific Ocean, around the same latitudes as the RVP. Despite the localized
34 nature of mode water subduction hotspots, changes in basin-scale density gra-
35 dients are an important controlling factor on heat distribution in the Southeast
36 Pacific.

37 **1. Introduction**

38 The Southern Ocean (SO) is an important region for the uptake, storage, and transport of heat and
39 carbon; SO upwelling brings up water that has not seen the surface for centuries or even millennia
40 and is thus far out of equilibrium with surface temperature and atmospheric carbon (Sabine et al.
41 2004; Sarmiento et al. 2004; Lenton and Matear 2007; Ito et al. 2010; Talley 2013; Roemmich et al.
42 2015; Armour et al. 2016). Over the period 1861-2005, the SO, defined as the ocean south of 30°S,
43 is estimated to account for $43\% \pm 3\%$ of anthropogenic carbon dioxide uptake and $75\% \pm 22\%$ of
44 heat uptake, despite only occupying 30% of global surface ocean area (Frölicher et al. 2015). The
45 SO's ability to absorb and transport this disproportionately high fraction of heat and carbon comes
46 in part from a balance between powerful overlying westerly winds, strong buoyancy fluxes, and
47 internal dynamics that maintains steeply tilted surfaces of constant density and localized pools
48 of weakly stratified water (Russell et al. 2006). These pools of weakly stratified water, referred
49 to collectively as Subantarctic Mode Water (SAMW), are refreshed by localized wintertime deep
50 convection that occurs just north of the Antarctic Circumpolar Current (ACC) (Speer et al. 2000;
51 Hanawa and Talley 2001; Lumpkin and Speer 2007; Talley 2008; Herraiz-Borreguero and Rintoul
52 2011; Speer and Forget 2013). The processes of lateral induction, eddy-induced transport, and
53 advection via the mean flow subduct the SAMW into the subsurface ocean at a rate that varies
54 across interannual and decadal timescales (Karsten and Marshall 2002; Sallée et al. 2010b; Sallée
55 and Rintoul 2011; Sallée et al. 2012; Liu and Huang 2012). Together with the denser and relatively
56 fresher Antarctic Intermediate Water (AAIW), SAMW ventilates the subtropical thermocline on
57 timescales of decades to centuries (Iudicone et al. 2007; Sallée et al. 2010a; Cerovecki et al. 2013;
58 Jones et al. 2016). Here, ventilation refers to the broad set of processes by which atmospheric
59 and surface ocean properties are able to affect subsurface ocean properties. Ventilation is driven

60 in part by the formation, subduction, and destruction of water masses and can be considered a
61 consequence of the ocean's global overturning circulation (Marshall and Speer 2012; Cerovečki
62 and Mazloff 2015).

63 SAMW and AAIW both display complex patterns of decadal variability. In the Western Pacific,
64 decadal-scale surface warming and/or freshening at high-latitude SAMW and AAIW subduction
65 sites can lead to cooling and freshening along isopycnals in the subsurface (Johnson and Orsi
66 1997). Decadal trends in SAMW temperature and salinity display complex spatial variability; in
67 recent decades, both the SAMW and AAIW have become warmer and saltier in the South Atlantic
68 but cooler and fresher in the South Atlantic (Katsumata and Fukasawa 2011). The AAIW core (its
69 salinity minimum) has warmed and shoaled overall, with small salinity trends whose sign varies
70 by region (Schmidtko and Johnson 2012). Decadal trends of SAMW and AAIW properties are
71 consistent with an amplification of the global hydrological cycle and broad-scale surface warming
72 (Helm et al. 2010; Durack and Wijffels 2010). The spatial and temporal complexity of these
73 SAMW and AAIW trends suggests that we will need a full four-dimensional description of the
74 sensitivity of ocean structure in order to make robust projections of potential future ocean states.

75 Regional differences in subduction, ventilation, and the properties of the ventilated subsurface
76 can impact (i.) the sequestration of anthropogenic heat and carbon into the subsurface thermocline
77 and (ii.) the supply of nutrients to low latitudes via the overturning circulation (Sarmiento et al.
78 2004; Sabine et al. 2004; Khatiwala et al. 2009; Ito et al. 2010; Roemmich et al. 2015). The
79 regionally specific nature of subduction and ventilation, and how those localized characteristics
80 affect the distribution of heat, carbon, and nutrients in the ocean, is an active area of oceanographic
81 study (Cerovecki et al. 2013; Jones et al. 2016). Despite these efforts, we still have relatively little
82 knowledge on how regional variations in ocean state variables (e.g. temperature, salinity) and
83 surface forcing can ultimately impact the properties of the ventilated subsurface. Advancements in

84 this area may be especially helpful for improving projections of future ocean states, as changes in
85 the Southern Ocean forcing-subduction-ventilation mechanism are expected to have a considerable
86 impact on future climate (Cessi and Otheguy 2003; Downes et al. 2009; Lovenduski and Ito 2009;
87 Morrison et al. 2011; Sallée et al. 2012).

88 In order to quantify how regional variations in ocean state variables may affect the heat distribu-
89 tion in the ventilated subsurface ocean, we perform a set of adjoint sensitivity experiments using
90 an observationally-constrained, physically consistent state estimate. First, we identify the recently
91 ventilated subsurface ocean using a combination of physical state variables (e.g. potential vor-
92 ticity, density, mixed layer depth) and numerical passive tracer distributions that track ventilation
93 pathways from the surface into the subsurface. Our method for identifying the recently ventilated
94 subsurface treats ventilation as a general process that may affect many different water masses, in
95 contrast to approaches that consider one particular water mass (e.g. SAMW). We focus our atten-
96 tion on the Eastern Pacific because it features an especially efficient export pathway of water from
97 the surface ocean into the subsurface thermocline, as measured by numerical passive tracer exper-
98 iments (Jones et al. 2016). For convenience, we refer to the recently ventilated subsurface Eastern
99 Pacific as the RVP and the heat content of the RVP as RVPh. Although there is overlap between
100 the two, we note that the RVP is more general than the region occupied by SAMW or any other
101 particular water mass. Once we have identified the RVP, we perform a set of adjoint sensitivity
102 experiments using, for consistency, the same numerical model setup that we used for the passive
103 tracer experiments. Our adjoint model produces linear, time-evolving sensitivity maps, allowing
104 us to identify locations and timescales where hypothetical anomalies could potentially impact the
105 heat distribution in the ventilated subsurface, as revealed by the heat content of the fixed-volume
106 RVP. Although we do not use this framework here, the sensitivity fields can roughly be interpreted
107 as optimal linear perturbations, i.e. perturbations with the largest possible impact on RVPh, within

108 the linear framework of the adjoint model (see a related but not identical example in Sévellec and
109 Fedorov (2015)).

110 This paper is structured as follows: in section 2, we describe the ECCOv4 model setup, the
111 design of our numerical tracer release experiments, and the design of our adjoint sensitivity ex-
112 periments. In section 3, we examine the results of the tracer experiments in order to define the
113 RVP, which is the control volume for our adjoint sensitivity experiments. In section 4, we examine
114 the time-evolving sensitivities of RVPh to changes in temperature and salinity, decomposed into
115 changes at constant density (i.e. kinematic sensitivities) and changes with varying density (i.e.
116 dynamic sensitivities). In section 5, we discuss factors that should be considered when comparing
117 our results with other studies. Finally, in section 6, we summarize our results and conclusions.

118 2. Methods

119 In this section, we describe the observationally-constrained global ocean model setup (subsec-
120 tion a), the design of our numerical passive tracer experiments used to define the RVP (subsection
121 b), and the design of our adjoint sensitivity experiments (subsection c).

122 a. Global ocean model setup

123 We use the modelling setup associated with ECCOv4 (release 2, hereafter ECCOv4-r2 or just
124 ECCOv4). ECCOv4 is a *state estimate*, meaning that it has been adjusted to minimize the mis-
125 fits between the model state and a suite of observations from various sources over the time period
126 1992-2011 (e.g. Argo float temperature and salinity profiles, ship hydrography, satellite altimetry).
127 The model’s initial conditions, surface forcing fields, and mixing parameters were iteratively ad-
128 justed in order to reduce model-data misfit. Because no artificial sources or sinks of heat were used
129 in the ocean subsurface, ECCOv4 features closed budgets of heat, salt, and volume. The model

130 setup is available for download on GitHub.com (<https://github.com/gaelforget/ECCOv4>) as
131 an instance of the open source MIT general circulation model (MITgcm, <http://mitgcm.org/>,
132 also available on GitHub). We briefly describe the relevant features of the ECCOv4 setup below;
133 a more thorough description is available in Forget et al. (2015a) and references therein.

134 ECCOv4 is a global ocean model that uses a Lat-Lon-Cap (LLC) grid. Its horizontal grid spac-
135 ing ranges from around 40-50 km at high latitudes up to roughly 110 km at the equator. It features
136 parameterized diffusion, including both diapycnal and isopycnal components, simple convective
137 adjustment, and the Gaspar-Gregoris-Lefevre (GGL) mixed layer turbulence closure scheme (Gas-
138 par et al. 1990). To represent the along-isopycnal effect of unresolved eddies, Forget et al. (2015a)
139 used a bolus transport parameterization (Gent and McWilliams 1990, hereafter GM). Although
140 the horizontal resolution of ECCOv4 is relatively coarse (roughly 1°), its mixing properties are in
141 good agreement with observations, as indicated by a comparison of simulated ECCOv4 oxygen
142 distribution and World Ocean Atlas 2013 (Forget et al. 2015b, Figure 4). This agreement is thanks
143 in part to the use of optimized, spatially-varying turbulent transport coefficients. ECCOv4 features
144 fully interactive dynamic sea ice, so that buoyancy and mass fluxes at the sea surface are calculated
145 based on the thermodynamic balance of Losch et al. (2010). Open ocean rain, evaporation, and
146 runoff simply carry (advect through the free surface) the local SST and a salinity value of zero, and
147 runoff is provided by a monthly climatology (Fekete et al. 2002). ECCOv4 calculates buoyancy,
148 radiative, and mass fluxes using the bulk formulae of Large and Yeager (2009) with 6-hourly ERA-
149 Interim re-analysis fields (Dee et al. 2011) as a “first guess” for the forcing fields. Specifically,
150 we use wind stress, 2 m air temperature, 2 m specific humidity, wind speed, downward longwave
151 radiation, and downward shortwave radiation as model inputs. These atmospheric state fields have
152 been iteratively adjusted by the state estimation process in order to minimize model-data misfits.
153 The ECCOv4-r2 setup that we use here does *not* use surface salinity restoring. For additional val-

154 idation information, see Forget et al. (2015a). For calculations relevant to mode water formation
155 and distribution in ECCOv4, see Speer and Forget (2013).

156 *b. Design of the numerical passive tracer release experiments*

157 Generally, on a selected timescale (e.g. 1 year, 10 years), we can consider the volume of the
158 ocean that has been affected by near-surface properties (e.g. the temperature and salinity charac-
159 teristics of the mode water formation regions) as having been ventilated via advection, diffusion,
160 and mixing. In terms of an operational definition, the ventilated subsurface ocean can be identified
161 using physical state variables, tracer distributions, or a combination of the two. Luyten et al. (1983)
162 showed that, in a simple, inviscid, analytical model, pathways of potential vorticity (PV) can be
163 used to separate the ventilated thermocline from the unventilated thermocline and also to separate
164 regions of weak stratification from regions of strong stratification. This result helps justify the
165 use of potential vorticity as a “dynamical tracer” that indicates thermocline ventilation pathways.
166 Alternatively, one can derive ventilation pathways using the observed distributions of tracers such
167 as tritium, helium-3 and chlorofluorocarbons, which are carried into the subsurface by both advec-
168 tion along ventilation pathways and by isopycnal and diapycnal mixing (Musgrave 1990; Speer
169 and Tziperman 1992; Primeau and Holzer 2006). In the subsurface, diapycnal mixing contributes
170 to the erosion of water masses by homogenizing their properties across density surfaces (Tross-
171 man et al. 2012). In combination with PV and tracer considerations, the ventilated subsurface is
172 generally considered to be located below the mixed layer, isolated from immediate contact with
173 the surface. In this study, we use a combination of physical state variables (e.g. PV, mixed layer
174 depth, density) and numerical passive tracer release experiments to identify the recently ventilated
175 subsurface ocean in ECCOv4-r2, described in detail below.

176 First, we use stratification (PV) and mixed layer depth (MLD) to identify the Southern Ocean re-
 177 gions that tend to ventilate the subsurface thermocline (Figure 1). We calculate $PV = f\rho_0^{-1}d\rho/dz$,
 178 where $f = 2\Omega\sin(\phi)$ is the Coriolis parameter and ρ_0 is the reference density, and we calculate
 179 mixed layer depth using a criteria based on the density change associated with a temperature vari-
 180 ation of $\Delta T = 0.8^\circ\text{C}$ (Kara et al. 2000). For each year between 1996 and 2001, we calculate the
 181 June-July-August (JJA) mean PV and both the JJA mean and annual maximum MLD, and we use
 182 these values with a set of criteria to construct a “mixed layer mask” for that year. Specifically, for
 183 a grid cell to be included in the mixed layer mask for a given year, it must:

- 184 1. be within the annual maximum mixed layer and
- 185 2. be at a latitude and longitude where one or both of the following conditions are satisfied:
 - 186 (a) the JJA mean $|PV|$ at the base of the JJA mean mixed layer must be less than $10^{-10.5}$
 187 1/ms, or
 - 188 (b) the JJA mean mixed layer must be greater than 300 m.

189 These relatively weak conditions ensure that we target likely ventilation pathways without being
 190 overly prescriptive.

191 We perform six different tracer release experiments (one for each starting year between 1996 and
 192 2001). For each experiment, we initialize passive tracer within the corresponding year’s mixed
 193 layer mask on 1 January; our method of initializing tracer above the maximum mixed layer for
 194 the release year ensures that variation associated with seasonal release timings are minimal. We
 195 integrate the tracer equations in “online” mode for 10 years, i.e. we explicitly solve the momentum,
 196 buoyancy, and tracer equations at the same time. Note that although the bulk of the tracer used to
 197 define the RVP originates in the Pacific, a smaller fraction also comes from the Indian sector. The
 198 10-year timescale of the experiments allows a large fraction of weakly stratified water to subduct

199 into the subsurface (Jones et al. 2016), and using a 10-year timescale also allows us to carry out a
200 suite of passive tracer and adjoint sensitivity runs within the 20-year observationally-constrained
201 time period of ECCOv4. We then use the annual mean tracer distribution after 10 years of model
202 integration to map out the ventilation pathways. See section 3 for more details and results.

203 *c. Design of the adjoint sensitivity experiments*

204 An adjoint model is a tool for sensitivity analysis; all other applications derive from there (Errico
205 1997). Adjoint methods can be used to construct state estimates (e.g. the ECCOv4 state estimate
206 used in this paper) via the iterative optimization of model initial conditions, model parameters,
207 and surface forcing (collectively called the “controls”), guided by *gradients* or *sensitivities* that
208 are calculated by an adjoint model. In such applications, the sensitivities indicate how one needs
209 to adjust the controls in order to reduce the *cost function*, which is a measure of model-data misfit
210 over the duration of the model run. This method enables objective, physically consistent, and
211 efficient model optimization that would otherwise require considerable ad-hoc parameter tuning
212 and/or the introduction of artificial sources or sinks of heat and salt into the subsurface.

213 More generally, adjoint methods also enable us to carry out efficient and comprehensive sensi-
214 tivity studies, such as the one presented in this manuscript. In a traditional “forward” perturbation
215 experiment, the experimenter introduces a perturbation into the model (e.g. a change in the initial
216 temperature pattern) and examines the effect on some quantity of interest (e.g. RVPh). Using this
217 method, the experimenter cannot feasibly determine if the perturbation they selected is the ideal
218 or optimal one (i.e. the one with the largest effect on RVPh) without performing a very large num-
219 ber of numerical experiments. By contrast, in an adjoint sensitivity experiment, the experimenter
220 selects a quantity of interest (e.g. annual mean RVPh) and uses an adjoint approach to calculate
221 *gradients* or *sensitivities*. The sensitivity fields indicate the perturbations with the largest possible

linear impact on RVPh. For examples of adjoint sensitivity experiments, see Verdy et al. (2014); Sévellec and Fedorov (2015); Jones et al. (2018), and many others. We refer the reader to these works for a more thorough and general description to adjoint modeling (Thacker and Long 1988; Marotzke et al. 1999b; Heimbach 2008; Mazloff et al. 2010; Griewank and Walther 2012; Verdy et al. 2014; Forget et al. 2015a; Fukumori et al. 2015; Jones et al. 2018, for example).

We perform a set of such adjoint sensitivity experiments in order to examine the sensitivity of RVPh to subsurface ocean properties. The objective function is defined as the annual- and volume-mean RVPh:

$$\tilde{J} = \frac{1}{V\Delta t} \int_V \int_{\Delta t} \rho_0 c_p \theta(\mathbf{r}, t) dt dV, \quad (1)$$

where ρ_0 is the reference density, c_p is the heat capacity of seawater, θ is the potential temperature, $\mathbf{r}$ is the position vector, t is the time, V is the control volume, and Δt is the time period of the integration. For convenience, we scale $\tilde{J}$ by the constant $\rho_0 c_p$, so $\text{RVPh} = J = \tilde{J} / \rho_0 c_p$, meaning that RVPh has units of $^{\circ}\text{C}$.

We compute an ensemble of six adjoint sensitivity experiments, with the objective function defined over the last year of each run, i.e. from 1 January to 31 December, with objective function years ranging from 2006 through 2011. Each experiment spans a minimum of 15 years, up to a maximum of 20 years. We allow the RVP to vary by ensemble member based on the results of the passive tracer experiments. Our ECCOV4 adjoint model calculates the sensitivities of these objective functions to a set of independent variables, including temperature ($\partial J / \partial T$) and salinity ($\partial J / \partial S$); the sensitivity fields vary with space (latitude, longitude, depth) and time. We use 14-day averaged sensitivity fields throughout.

3. Identifying the recently ventilated subsurface ocean

We use a six-member ensemble of passive tracer experiments with release years ranging from 1996 to 2001 to identify the RVP in ECCOv4. The ensemble mean, annual mean tracer distribution after 10 years of integration (Figure 2) highlights the ventilation pathways from the Southern Ocean mixed layer into the subsurface thermocline.

The normalized tracer distribution features several cores, i.e. local maxima with some degree of spatial coherence across latitudes and longitudes (Supplemental Animation S1). In the Atlantic sector, we find a core below the 1992-2011 maximum MLD in the subtropical latitudes (roughly between 30°S-10°S) that drifts southwards before merging with a broad pattern that stretches from the Weddell Sea, across the ACC, and into the subtropics (Figure 2(a)). The approximately zonal Atlantic distribution bears the imprint of the convergence of the Brazil current and the ACC. In the Indian sector, the tracer distribution is influenced by the Agulhas current and the broader circulation of the South Indian Gyre that is characterized by relatively shallow overturning in density space; here, we find the core of the distribution at roughly 500 m and densities lighter than the $\sigma_0 = 26.9 \text{ kg m}^{-3}$ density surface (Figure 2(b)). In the Pacific sector, we find relatively large values throughout the entire South Pacific Gyre (SPG), with two distinct cores. The core of the distribution in the western SPG is found at densities lighter than $\sigma_0 = 26.9 \text{ kg m}^{-3}$ and is affected by tracer that subducts in the Indian sector, whereas the core of the distribution in the eastern SPG largely straddles the $\sigma_0 = 26.9 \text{ kg m}^{-3}$ line and is relatively less affected by tracer that subducts in the Indian sector (Figure 2(c,d)). High tracer concentrations in the eastern region of the SPG highlights the relatively efficient Eastern Pacific export pathway that was previously identified in high-resolution simulations (Jones et al. 2016, Figure 8d).

264 In order to define the RVP, we use the annual mean tracer concentrations after 10 years of
265 transport and some additional physical and geographical criteria, selecting grid cells that satisfy
266 the following four conditions:

- 267 • be located below the maximum mixed layer over the entire ECCOv4-r2 period (i.e. 1992-
268 2011)
- 269 • tracer concentration is at least 10% of the global ocean maximum value
- 270 • be located in the Southeast Pacific, between 170°W-60°W and 60°S-20°S
- 271 • potential density is greater than or equal to $\sigma_0 = 26.9 \text{ kg m}^{-3}$

272 The $\sigma_0 = 26.9 \text{ kg m}^{-3}$ surface delineates lighter waters from waters that are at least as dense
273 as the SAMW. The resulting RVP for any given release year is a three-dimensional volume with
274 a lateral imprint that varies with depth (Figure 3 and 4). Between roughly 300-500m (Figure
275 3(a)-(c)), the areal extent of the RVP changes with depth and sits largely in the Southeast Pacific,
276 north of the regions with deep mixed layers. The structure of the RVP at these depths reflects the
277 subduction and ventilation pattern associated with the Eastern Pacific export pathway (Jones et al.
278 2016). Between roughly 500-700 m (Figure 3(d)-(e)), the RVP reaches its maximum areal extent,
279 covering a large fraction of the SPG.

280 1) THE RVP IS MORE GENERAL THAN SAMW

281 Although the RVP has considerable overlap with the density ranges associated with SAMW, it
282 is more general than SAMW alone. Our passive tracers ventilate a wide range of water masses,
283 including SAMW, AAIW, and the lightest part of the Circumpolar Deep Water (CDW). In this
284 way, our RVP approach is broader than an approach focused on a particular density range, in that
285 it allows us to consider a wider range of ventilated water mass types than SAMW alone. Because

RVPh is the heat content of a fixed volume, RVPh is a measure of the heat distribution in the target region, as RVPh may be affected by both isopycnal heave and by changes in properties along isopycnals. RVPh may have different sensitivities than the SAMW heat content, which may be affected by changes in the volume of SAMW and/or changes in the properties of water in the SAMW density range. We further explore this distinction in section 5.

Having defined the RVP in ECCOv4, we can now perform a set of adjoint sensitivity experiments using the adjoint capability of the ECCOv4 model setup.

4. Sensitivities of subsurface heat content to local and remote anomalies

In this section, we examine the sensitivity of RVPh to changes in temperature and salinity, decomposed into changes at constant density (i.e. kinematic sensitivities) and changes with varying density (i.e. dynamic sensitivities). In subsection a, we show how sensitivities to temperature and salinity perturbations can be cast in terms of kinematic and dynamic sensitivities. In subsection b, we investigate the spatial and temporal structure of the kinematic sensitivities, focusing on both the subsurface and the surface. Finally, in subsection c, we examine the spatial and temporal structure of the dynamic sensitivities, including regional patterns in the subsurface and at the surface. In this analysis, we use the time-variable -0.25 m ECCOv4 sea surface height contour as a proxy for the Subantarctic Front (SAF) that roughly divides the Southern Ocean into two regions - one dominated by gyre-like circulation and another region dominated by the circumpolar flow associated with the ACC (Kim and Orsi 2014).

a. Defining kinematic and dynamic sensitivities

Following Marotzke et al. (1999b) and Jones et al. (2018), we decompose the adjoint sensitivity fields into sensitivities to changes that propagate along isopycnals (i.e. kinematic) and sensitivities

308 to changes in density (i.e. dynamic). Writing the objective function in terms of density and
 309 temperature $J = J[\rho(T, S), T]$ allows us to write the sensitivities as follows:

$$\left(\frac{\partial J}{\partial T}\right)_S = \left(\frac{\partial J}{\partial \rho}\right)_T \left(\frac{\partial \rho}{\partial T}\right)_S + \left(\frac{\partial J}{\partial T}\right)_\rho. \quad (2)$$

310 The first term on the right-hand side of equation 2 is the “dynamic” component of the sensitivity
 311 (i.e. sensitivity to changes in density), and the second term on the right-hand side is the “kine-
 312 matic” component (i.e. dynamically-inactive sensitivities to temperature anomalies). Using the
 313 coefficient of thermal expansion α and coefficient of haline contraction β , defined as

$$\alpha \equiv -\frac{1}{\rho} \left(\frac{\partial \rho}{\partial T}\right)_S \text{ and } \beta \equiv \frac{1}{\rho} \left(\frac{\partial \rho}{\partial S}\right)_T, \quad (3)$$

314 we can write

$$\left(\frac{\partial J}{\partial S}\right)_T = \left(\frac{\partial J}{\partial \rho}\right)_T \left(\frac{\partial \rho}{\partial S}\right)_T = \beta \rho \left(\frac{\partial J}{\partial \rho}\right)_T, \quad (4)$$

315 and the dynamic sensitivity becomes:

$$F_{dyn} = \left(\frac{\partial J}{\partial \rho}\right)_T \left(\frac{\partial \rho}{\partial T}\right)_S = \frac{1}{\beta \rho} \left(\frac{\partial J}{\partial S}\right)_T \left(\frac{\partial \rho}{\partial T}\right)_S = -\frac{\alpha}{\beta} \left(\frac{\partial J}{\partial S}\right)_T. \quad (5)$$

316 We can write the kinematic sensitivity as:

$$F_{kin} = \left(\frac{\partial J}{\partial T}\right)_S + \frac{\alpha}{\beta} \left(\frac{\partial J}{\partial S}\right)_T. \quad (6)$$

317 Physically, kinematic sensitivities indicate the linear response of the objective function to the
 318 simultaneous application of a small temperature change and a compensating salinity change of
 319 the form $\Delta S = \Delta T \alpha / \beta$, such that density remains constant following the linearized equation of
 320 state for seawater, $\rho = \rho_0(1 - \alpha \Delta T + \beta \Delta S)$. We can refer to these perturbations as “density-
 321 compensated temperature anomalies”. In contrast, dynamic sensitivities indicate the linear re-
 322 sponse of the objective function to the application of a temperature change ΔT or a density-
 323 equivalent salinity change of $\Delta S = -\Delta T \alpha / \beta$ in the equation of state. The linear response is

324 given by $\Delta J = F_{dyn}\Delta T = F_{dyn}(-\Delta S\beta/\alpha)$. We use 14-day averaged, three-dimensional α/β fields
325 derived from ECCOv4-r2 potential temperatures and salinities using the TEOS-10 toolbox (Mc-
326 Dougall and Barker 2011). Next, we examine the kinematic and dynamic sensitivities of RVPh.

327 In this paper, we follow the naming and conceptual conventions established in the adjoint mod-
328 eling literature, specifically the use of kinematic and dynamic sensitivities (Marotzke et al. 1999a;
329 Jones et al. 2018, e.g.). Kinematic sensitivities are conceptually similar to sensitivities to “spice”
330 anomalies (Flament 2002), and dynamic sensitivities are conceptually similar to sensitivities to
331 “heave” anomalies (Bindoff and Mcdougall 1994). However, despite their resemblance, it would
332 be misleading to conflate kinematic sensitivities with sensitivities to spice anomalies, and it would
333 be misleading to conflate dynamic sensitivities with sensitivities to heave anomalies. Their deriva-
334 tions and resulting definitions are not rigorously identical. In addition, there is a difference in how
335 these quantities are normally used; spice and heave are typically used to decompose local potential
336 temperature anomalies into these two components, whereas kinematic and dynamic adjoint sensi-
337 tivities are more generally used to show impacts on an objective function over different spatial and
338 temporal lags.

339 At present, adjoint sensitivity fields in ECCOv4 do *not* include sensitivity to changes in mixing
340 as parameterized by the GGL mixed layer turbulence closure scheme. Also, the sensitivity fields
341 do *not* include the sensitivity to changes in sea ice. The adjoint sensitivity calculations essentially
342 treat GGL mixing and sea ice as fixed processes. We apply a mask of the form $(1 - f)$, where f is
343 the local sea ice concentration, to all surface sensitivity fields to account for the lack of sensitivity
344 to sea ice change.

345 *b. Sensitivity to changes at constant density*

346 Kinematic sensitivity fields approximately indicate potential source water pathways. More
347 specifically, one can think of the kinematic sensitivity fields as indicating the eventual response of
348 annual mean RVP heat content to a spatially uniform, density-compensated potential temperature
349 perturbation that is sustained for some time (in our case, two weeks) and allowed to propagate.
350 Regions with zero sensitivity cannot affect RVPh; such regions are not connected to the RVP via
351 advection and mixing on the associated timescales. The kinematic sensitivity fields are positive or
352 zero nearly everywhere; an increase/decrease in temperature in regions with non-zero sensitivities
353 eventually increases/decreases RVP heat content on the indicated lag timescale (Figure 5).

354 The near-surface kinematic sensitivity fields highlight the potential propagation pathways of
355 anomalies into the SAMW formation regions, where they can be subducted below the mixed layer
356 and exported into the thermocline (Figure 5, left column, and Supplemental Animation S2). In the
357 Pacific basin, the largest near-surface sensitivity fields tend to be found just north of the SAF proxy,
358 within or upstream of regions of deep mixed layers and mostly south of the RVP areal extent. In
359 the Indian basin, the largest near-surface sensitivity fields are found along the SAF proxy or just
360 south of it. The fields cross over the SAF around Campbell Plateau and South Australian Basin,
361 suggesting cross-front exchange in this region.

362 Density-compensated anomalies cannot typically propagate from directly above the RVP into
363 the RVP itself on the timescales considered here; they are not typically advected into the regions
364 of deep mixing. On timescales longer than 10 years, the near-surface sensitivities extend across the
365 entire Southern Ocean, putting a lower bound on the temperature anomaly propagation timescale
366 from the Atlantic to the RVP (Figure 5, left column). On timescales longer than roughly 5-6
367 years, we find sensitivities south of the SAF proxy in the Indian basin; here, wind stress may

368 advect anomalies across the SAF via Ekman transport. On shorter timescales, the sensitivity fields
369 are more localized around the southeast Pacific mixed layer region and are typically smaller in
370 magnitude relative to the deeper levels. Temperature anomalies can only affect RVPh if they
371 have sufficient time to propagate into the mixed layer, subduct, and get exported; only very local,
372 targeted anomalies can affect RVPh on timescales shorter than 1-2 years. The persistent local
373 maximum in the southeast Pacific indicates a relatively rapid route from the surface into the RVP,
374 which is consistent with a well-documented subduction and export pathway in that location (Sallée
375 et al. 2010b).

376 The subsurface sensitivity fields highlight the long residence time of local hypothetical pertur-
377 bations in the subsurface of the gyre, as well as the pathways of subsurface ocean advection into
378 the RVP (Figure 5, right column, and Supplemental Animation S3). On timescales longer than 10
379 years, we find sensitivities extending along the East Australian Current and, in contrast with the
380 surface sensitivities, sensitivities along the Agulhas Current and associated retroflexion as well.
381 Surface anomalies in the Agulhas Current cannot affect RVPh on the timescales considered here,
382 but anomalies in the subsurface (here, roughly 550 m) can eventually alter RVPh. We also find
383 some sensitivity at this depth extending along the Brazil current. Notably, the Agulhas current is
384 the only major Southern Hemispheric western boundary current with zero kinematic sensitivity in
385 the near-surface.

386 The global mean RMS kinematic sensitivity decreases with timescale (Figure 6(a)). It is reason-
387 ably well described ($R^2 > 0.99$) by an exponential function with a decay rate of about 10 years.
388 The decay partly reflects the fact that as we consider longer timescales, RVPh becomes sensitive
389 to a larger fraction of the ocean, so the global RMS sensitivity decreases. The decay also partly
390 reflects the tendency of air-sea exchange, represented in ECCOV4 as bulk formulae with a pre-
391 scribed set of atmospheric variables, to dampen potential temperature anomalies at the surface

and thereby reducing the sensitivities. The linear change in sensitivity between lag 0 and lag 1 yr reflects the one-year integration period over which the objective function is calculated. There is no sensitivity to anomalies after year 1, after the integration has finished. The ensemble standard deviation is small, indicating a relatively low level of interannual variability in sensitivity across our six-member ensemble.

We find non-zero sensitivities to hypothetical anomalies within the local RVP throughout the entire 16-year adjoint sensitivity experiment, indicating that RVPh is sensitive to local density-compensated temperature anomalies on timescales of at least 16 years (Figure 6(b)). The ensemble mean RMS local sensitivity time series is well-described ($R^2 > 0.99$) by a two-term exponential with decay rates of 4.4 years and 10.9 years. The non-local RMS sensitivities outside of the RVP exceed the local sensitivities on timescales longer than roughly 5 years (Figure 6(b)). RMS sensitivities outside of the MLD mask are larger than those within the MLD mask on timescales shorter than about 8 years (Figure 6(c)). This is largely due to the fact that sensitivities exist below the surface, e.g. in the RVP itself, where the overlying mixed layers are typically shallow.

Outside of the RVP, we mostly find non-zero sensitivities just north of the SAF proxy in all three ocean basins, in what may be considered the northern extent of the ACC (Figure 6(d)). The near-surface sensitivity fields display a clear seasonal cycle, with the RMS mean peaking in August in the Pacific and Indian basins, during the strong mixing phase that occurs as the mixed layer deepens before reaching its maximum depth in mid-September (Figure 6e). The sensitivity fields spread westward with lead time, reflecting the advective timescale between basins (Figure 6(e,f)). The RMS sensitivity in the West Pacific exceeds the East Pacific value on timescales longer than roughly 4.4 years.

414 *c. Sensitivity to changes in density*

415 The dynamic sensitivity fields indicate patterns of density change that can alter RVPh, poten-
416 tially by inducing changes in isopycnal tilt and thereby altering the associated transport and heat
417 convergence into the RVP. More specifically, one can think of the dynamic sensitivity fields as
418 indicating the eventual response of annual mean RVP heat content to a spatially uniform potential
419 temperature perturbation, or a density-equivalent salinity perturbation, that is sustained for some
420 time (two weeks, as with the kinematic sensitivity case) and allowed to affect the oceanic density
421 structure. Regions of zero sensitivity cannot affect RVPh; such regions are not dynamically con-
422 nected to the RVP. The dynamic sensitivity fields are both positive and negative, with the dipoles
423 highlighting regions where a change in isopycnal tilt will change the associated circulation and
424 ultimately RVPh (Figure 7).

425 The near-surface dynamic sensitivity fields are mostly negative across all lags, with local ex-
426 ceptions south of Australia and in the eastern tropical Pacific (Figure 7, left column). Negative
427 values indicate regions where a hypothetical temperature increase, which will decrease density,
428 will ultimately decrease RVP heat content. Similarly, in these regions a hypothetical temperature
429 decrease, which will increase density, will ultimately increase RVP heat content. It is important
430 to note that, due to technical limitations, our adjoint sensitivity fields *do not* represent non-linear
431 changes. They *do* reflect changes in circulation and diffusion. So the negative near-surface values
432 may represent the fact that, even under fixed mixing conditions, relatively lighter water may not
433 reach the depths of the RVP as readily as relatively denser water, where “relative” refers to the
434 background ECCOv4 state around which the adjoint sensitivities are calculated. Alternatively, the
435 negative near-surface values may represent a hypothetical decrease in RVP heat convergence in-
436 duced by a surface decrease in density, or the values may represent a hypothetical increase in RVP

437 heat convergence induced by a surface increase in density. As with the kinematic sensitivity fields,
438 on shorter timescales the dynamic sensitivities become increasingly localized to the southeastern
439 edge of the RVP, near the southern tip of South America.

440 In contrast with the near-surface kinematic sensitivity fields, the near-surface dynamic sensitiv-
441 ity fields feature a persistent negative local minimum along the western coast of South America
442 (Figure 7, left column, and Supplemental Animation S4). Considered together with the small
443 positive values in the eastern tropical Pacific, this dipole reflects a sensitivity to the across-shelf
444 pressure gradient. Variations in the across-shelf pressure gradients are associated with changes
445 in basin-scale pressure gradients and the associated basin-scale circulation, which can ultimately
446 change heat convergence and thus heat content within a selected ocean volume (Fukumori et al.
447 2015; Jones et al. 2018; Hughes et al. 2018).

448 The subsurface dynamic sensitivity fields highlight the sensitivity of annual mean RVP heat con-
449 tent to basin-wide density contrasts, equivalently expressed as tilted basin-scale isopycnal surfaces
450 (Figure 7, right column, and Supplemental Animation S5). This basin-scale structure is associated
451 with the circulation of the SPG, which is in part maintained by pressure gradients induced by wind
452 stress curl. We find a persistent sensitivity dipole stretched across the RVP, with positive values
453 to the east and negative values to the west. The positive values reach into the tropics and extend
454 westward with timescale, across both the Pacific and Indian basins. To illustrate the structure of
455 these fields, consider a hypothetical potential temperature anomaly characterized by warming to
456 the east of the RVP and cooling to the west of the RVP, imposed on the model vertical level cen-
457 tered at 553 m (Figure 7(d)). The warming to the east would decrease density at 553 m and, based
458 on the positive sign of the sensitivity field, ultimately contribute to an increase RVPh. The cooling
459 to the west would increase density at 553 m and, based on the negative sign of the sensitivity field,
460 it would also increase RVPh. The change in RVPh induced by any anomaly is the product of the

461 sensitivity field and the imposed anomaly, i.e. $\Delta J = F_{dyn}\Delta T$. In this scenario, we have imposed
462 isopycnal deepening to the east and isopycnal shoaling to the west. In the RVP, isopycnal surfaces
463 at roughly 500-600 m tend to shoal from west to east, so our imposed perturbation alters isopycnal
464 surfaces locally, ultimately changing transport and increasing heat convergence in the RVP (Jones
465 et al. 2018).

466 The subsurface sensitivity fields also feature a persistent sensitivity along the western coast of
467 South America (Figure 7, right column). In contrast to the near-surface fields, which are fairly
468 consistently negative along the shelf, the along-shelf subsurface fields are positive to the north and
469 negative to the south. The north-south sign contrast indicates a sensitivity to barotropic structure,
470 whereas the surface-subsurface sign contrast highlights an overall sensitivity to baroclinic den-
471 sity structure above the continental shelf. In addition, the subsurface fields display characteristic
472 fingerprints of baroclinic Rossby waves, as seen by westward propagating positive anomalies at
473 depth (Supplemental Animation S5). As with the surface fields, on shorter timescales the subsur-
474 face dynamic sensitivities become increasingly localized, albeit with a different structure than the
475 surface fields. The subsurface fields are localized around the edge of the RVP, indicating a sensi-
476 tivity to density contrasts across the RVP that can alter fluxes across the boundary of the RVP, and
477 the subsurface sensitivities are also localized along the western coast of South America, indicating
478 a short-timescale sensitivity to density changes along the continental shelf.

479 The globally averaged RMS dynamic sensitivity increases with timescale, reaching a maximum
480 at roughly 8.7 years (Figure 8(a)). Much of this sensitivity comes from outside the RVP, partic-
481 ularly near its boundary (Figure 8(b)). RMS sensitivity is similar inside and outside of the MLD
482 mask, indicating that when considering RVPh, we need to study dynamic changes both inside and
483 far outside of the regions of deep mixing (Figure 8(c)). The RMS sensitivity is persistently larger
484 north of the SAF proxy, although much of this is located just north of the SAF, and the RMS

sensitivity south of the SAF increases roughly linearly with timescale (Figure 8(d)). The near-surface fields also display a seasonal cycle, with similar timings as the kinematic fields (Figure 8(e)). Notably, the RMS dynamic sensitivity is persistently higher in the West Pacific than in the East Pacific, indicating the importance of dynamic basin-scale changes (Figure 8(h)). The RMS sensitivity in the Indian and Atlantic basins increases with timescale, although the values are negligibly small (less than 3% of the maximum) for short timescales (about 2.5 years for the Indian basin and 5 years for the Atlantic basin).

5. Discussion

In the canonical view of Southern Ocean meridional overturning circulation (MOC), Circumpolar Deep Water upwells south of the ACC, where it is brought to the surface and exported northwards; this water is modified by air-sea fluxes and is ultimately subducted as SAMW and AAIW (Marshall and Speer 2012). As a result of this overturning circulation, the properties of SAMW and AAIW are potentially sensitive to anomalies in near-Antarctic waters. For example, the variability of the Amundsen Sea low (ASL) can affect the interannual and interdecadal variability of SAMW in the Southeast Pacific via cross-frontal Ekman transport of Antarctic surface waters (Garabato et al. 2009; Close et al. 2013). In 2008-2010, a deep ASL enhanced the meridional wind-driven sea ice export from the eastern Ross Sea; the sea ice melt introduced a cold surface freshwater anomaly that was advected across the PF and SAF by Ekman transport and then transported into the SAMW formation region by the ACC (Cerovecki et al. 2019). In this context, our adjoint sensitivity results may appear surprising, in that they show relatively weak sensitivities to anomalies in the Ross Sea. Our results are not inconsistent with those studies, but there are several factors to consider when comparing our results with those obtained using other approaches, as discussed below.

508 *a. RVPh is a fixed-volume quantity*

509 Our control volume is fixed relative to the model grid. Although RVP captures the mean position
510 of the ventilated water, it does not move with isopycnal heave. Because it overlaps with both
511 SAMW and AAIW, it is more general than either water mass, but we must be careful not to
512 conflate SAMW/AAIW with the RVP. Isopycnal heave in the subtropical gyre can affect RVPh
513 without necessarily changing the heat content of the recently subducted SAMW. However, RVPh
514 can still tell us something about subduction. The majority of recent heat uptake by SAMW is
515 driven by a thickening of the associated isopycnal layers, driven by changes in wind forcing and
516 subduction in the deep mixed layer region Gao et al. (2018). Considering SAMW in the Pacific
517 basin, its local mean temperature is correlated with remote forcing (e.g. air-sea heat flux, wind
518 stress) with a lag that is attributed to transport, whereas its total heat content is controlled largely
519 by layer thickness, which is driven by local wintertime forcing (Meijers et al. 2019). These results
520 are consistent with the our adjoint sensitivity fields, which largely capture mechanisms that can
521 affect mean temperature within the SAMW as opposed to those that can affect total heat content
522 via changing layer thickness. It would be useful to examine the sensitivities in an adjoint model
523 that uses isopycnal surfaces as its vertical coordinate, but this is beyond the scope of this paper.

524 *b. Zonal flow dominates meridional overturning*

525 The largely zonal volume transport of the ACC is at least an order of magnitude larger than the
526 volume transport associated with the MOC (Talley 2013; Donohue et al. 2016). As such, we might
527 expect our adjoint sensitivity fields to reflect the dominance of ACC transport in the Southeast
528 Pacific. According to these fields, sensitivity to high-latitude processes is small but nonzero. To
529 examine the impacts of actual temperature and salinity anomalies on RVPh, one could examine
530 the combination of sensitivity fields together with temperature and/or salinity anomaly fields. In

531 this work, we focus on the sensitivity fields themselves, as these could potentially be combined
532 with a large variety of anomaly fields in order to address different questions (e.g. how does the
533 SAM forcing pattern affect RVPh across different timescales?)

534 *c. Cross-front exchange rates differ between data products*

535 Cross-front exchange in the SO is highly localized, occurring in relatively narrow longitude
536 bands in the lee of bathymetric features (Thompson and Sallée 2012). The propagation of anoma-
537 lies from the Ross Gyre to the subtropics occurs largely via two distinct exchange windows in the
538 Southeast Pacific and the central Pacific; on timescales of roughly 1-3 years, the central Pacific
539 pathway enables transport from Antarctic surface waters to the deep mixed layer region (Cerovecki
540 et al. 2019). The subsurface exchange pathway in the central Pacific (around the Eltanin fracture
541 zone, between roughly 110-160°W) is driven by enhanced diapycnal mixing near the relatively
542 rough topography and temporal circulation variability on timescales longer than 90 days (McAu-
543 field 2019). Because of its relatively coarse resolution, it is possible that ECCOv4-r2 may have
544 a lower cross-frontal exchange rate than that of higher-resolution data products, but a lack of ob-
545 servational data constraints makes it challenging to discriminate amongst estimates. A thorough,
546 quantitative comparison of exchange rates between ECCOv4-r2 and other data products would be
547 a useful addition to this work.

548 *d. Mixed layer depth*

549 In this work, we calculate mixed layer depth using a criteria based on the density change associ-
550 ated with a temperature variation of $\Delta T = 0.8^\circ\text{C}$ (Kara et al. 2000). Compared with other methods,
551 this threshold approach tends to be biased towards deeper values (Forget et al. 2015c, supporting
552 information). Because we discard all tracer above the maximum mixed layer, this implies that our

553 estimate of the subducted volume may be somewhat conservative. It is possible that our results
554 would be different using another mixed layer scheme. That being said, the stratification and mixed
555 layer properties in ECCOv4 compare well spatially with in-situ profiles (Forget et al. 2015c, Fig.
556 6).

557 *e. Adjoint sensitivities use fixed vertical mixing and sea ice*

558 At present, adjoint sensitivity fields in ECCOv4 do *not* include sensitivities to changes in mix-
559 ing as parameterized by the GGL mixed layer turbulence closure scheme. Also, the sensitivity
560 fields do *not* include the sensitivity to changes in sea ice. Although the forward and adjoint runs
561 do feature dynamic sea ice and vertical mixing as parameterized by GLL, the adjoint sensitivity
562 calculations essentially treat them as fixed processes. Any processes that affect RVPh primar-
563 ily through *changing* sea ice concentration and/or vertical mixing will not be represented in our
564 sensitivity fields.

565 **6. Summary and conclusions**

566 We used numerical passive tracer experiments and physical state variables (i.e. mixed layer
567 depth, potential density, potential vorticity) in an observationally-constrained state estimate to
568 identify the recently ventilated Eastern Pacific sector of the Southern Ocean (RVP), which is a
569 relatively efficient export pathway from the surface ocean into the subsurface thermocline. The
570 RVP, defined on a 10-year timescale, is located north of the -0.25 m sea surface height contour
571 (i.e. the SAF proxy) and north of the deep mixed layer regions, highlighting the ventilation path-
572 way from the tip of South America, along the eastern Pacific boundary and out across the South
573 Pacific Gyre (SPG) in the subsurface. The RVP varies with depth, reaching maximum areal extent
574 between 500-600 m, and it features some variation with tracer release year, mostly in its western

575 extent. The overall spatial pattern of the subducted tracer is consistent with estimates from an
576 eddy-permitting model (Jones et al. 2016).

577 We used the adjoint capabilities of MITgcm to calculate the sensitivities of annual mean heat
578 content within the fixed volume RVP (i.e. RVPh) to anomalies in potential temperature and salin-
579 ity. These sensitivities are four-dimensional (i.e. latitude, longitude, depth, and time) and indicate
580 the anomaly patterns that can have the largest potential impact on RVPh after the appropriate lag
581 time has elapsed (Figure 9). In terms of sensitivities to surface anomalies, RVPh is most sensitive
582 to anomalies along the SAF proxy, upstream of the subduction hotspot in the Southeast Pacific.
583 The sensitivity signal crosses the SAF proxy around Campbell Plateau and just south of Australia,
584 indicating possible wind-driven cross-frontal exchange in that region. This result is consistent
585 with estimates of the spatial pattern of meridional Ekman transport of atmospheric carbon in an
586 eddy-permitting state estimate (Ito et al. 2010). In terms of sensitivities to subsurface anomalies,
587 RVPh is most sensitive to basin-scale changes in density, as indicated by the persistent dipole in
588 subsurface dynamic sensitivity (Figure 7). The sensitivity dipole is centered around the same lati-
589 tudes as the RVP, extending upstream along the ACC and northwards into the tropical Pacific with
590 increasing timescale. This persistent dipole implies that property gradients between the SPG and
591 the ACC are important for interannual variability in RVPh.

592 By the end of the 21st century, Southern Ocean wind stress is expected to strengthen in magni-
593 tude and move poleward, resulting in a poleward shift of the gyres, and mixed layers are expected
594 to shoal (Meijers 2013). Our results suggest that the resulting subtropical, basin-scale temperature
595 and salinity changes induced by the expected shifts in wind stress and heat flux can have a consid-
596 erable impact on the heat distribution of the ventilated Southeast Pacific, across the density ranges
597 of the base of the SAMW, AAIW, and the lightest CDW. In a companion manuscript, we examine
598 the sensitivities of RVP heat content to changes in surface forcing (Jones et al. 2019).

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611 **Data availability statement:** The model used here is an instance of the ECCOv4-r2 state es-
612 timate (<https://github.com/gaelforget/ECCOv4>) as an instance of the MIT general circula-
613 tion model (MITgcm, <http://mitgcm.org/>). Specific model configuration files, input data, and
614 sample output available via Zenodo (Jones 2020).

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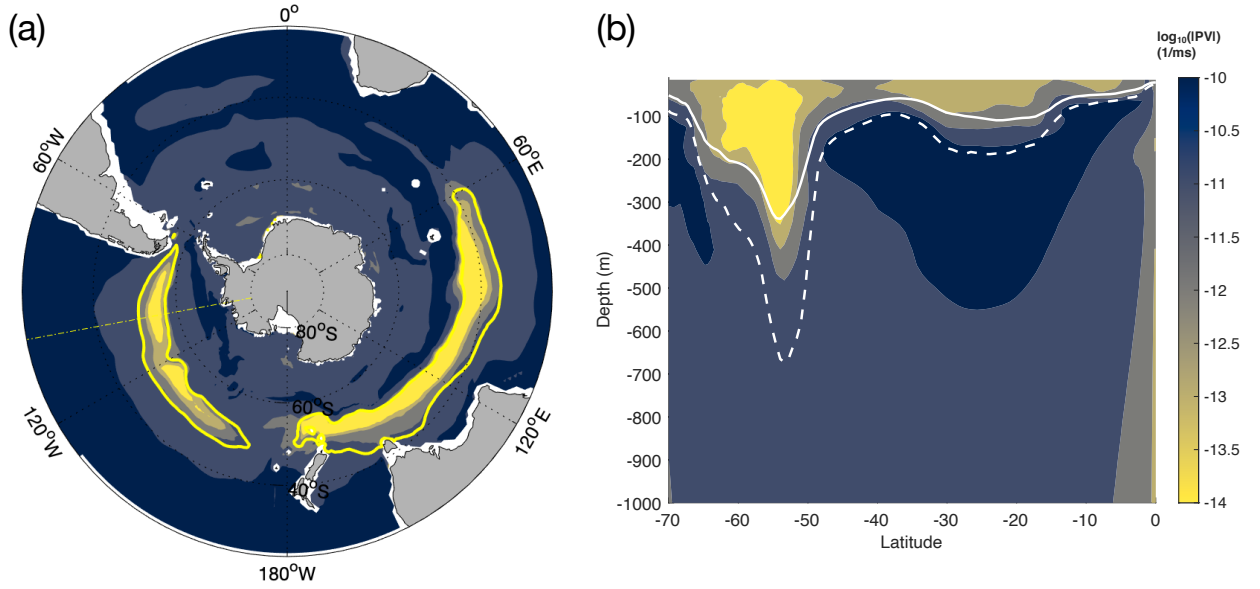


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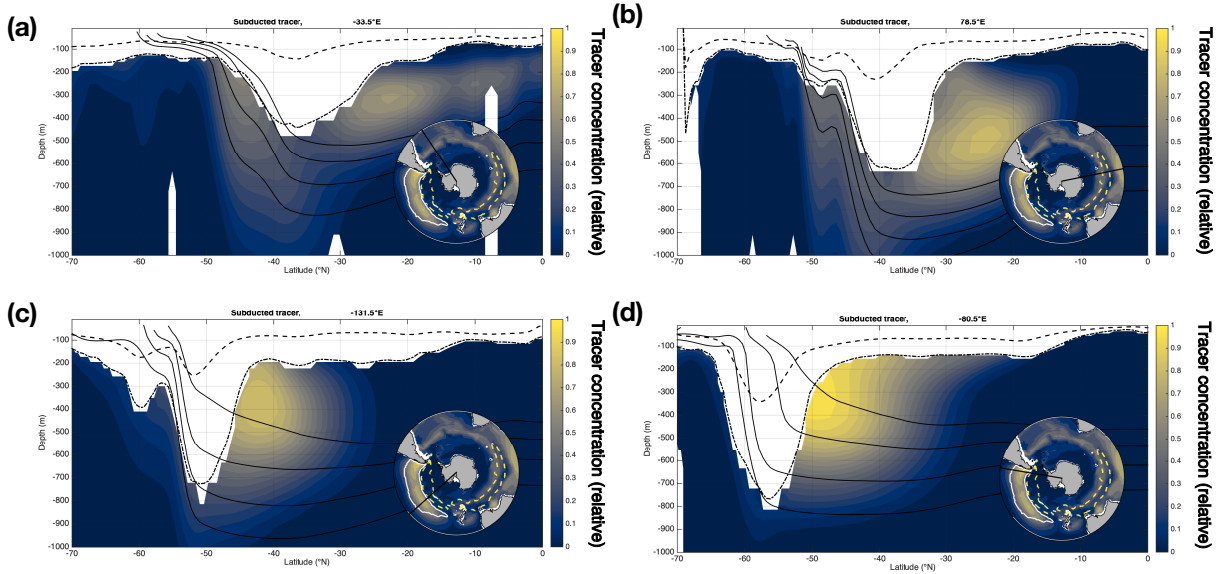


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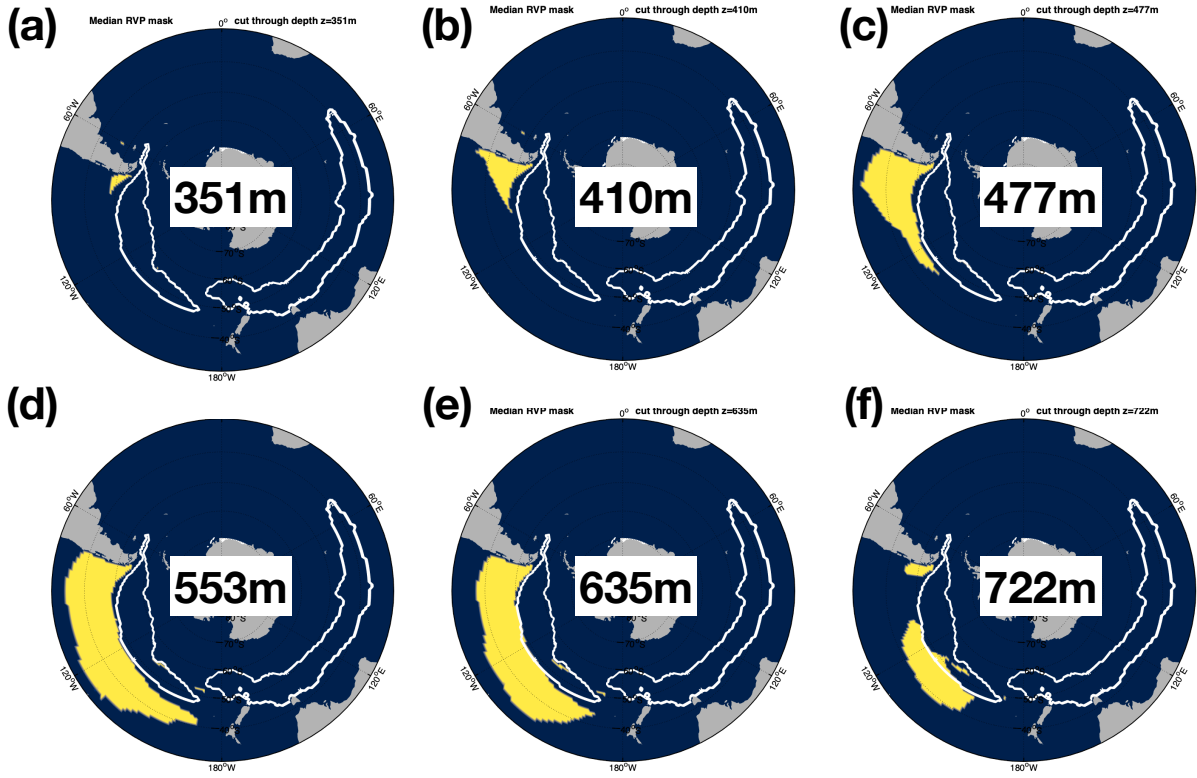


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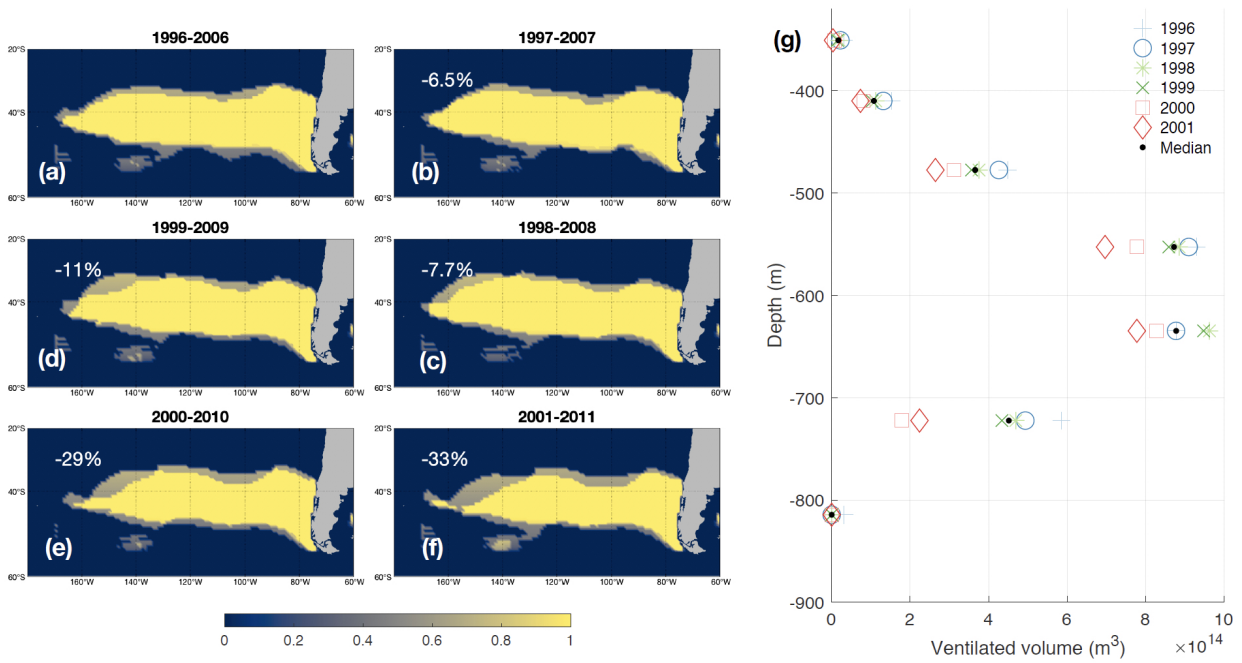


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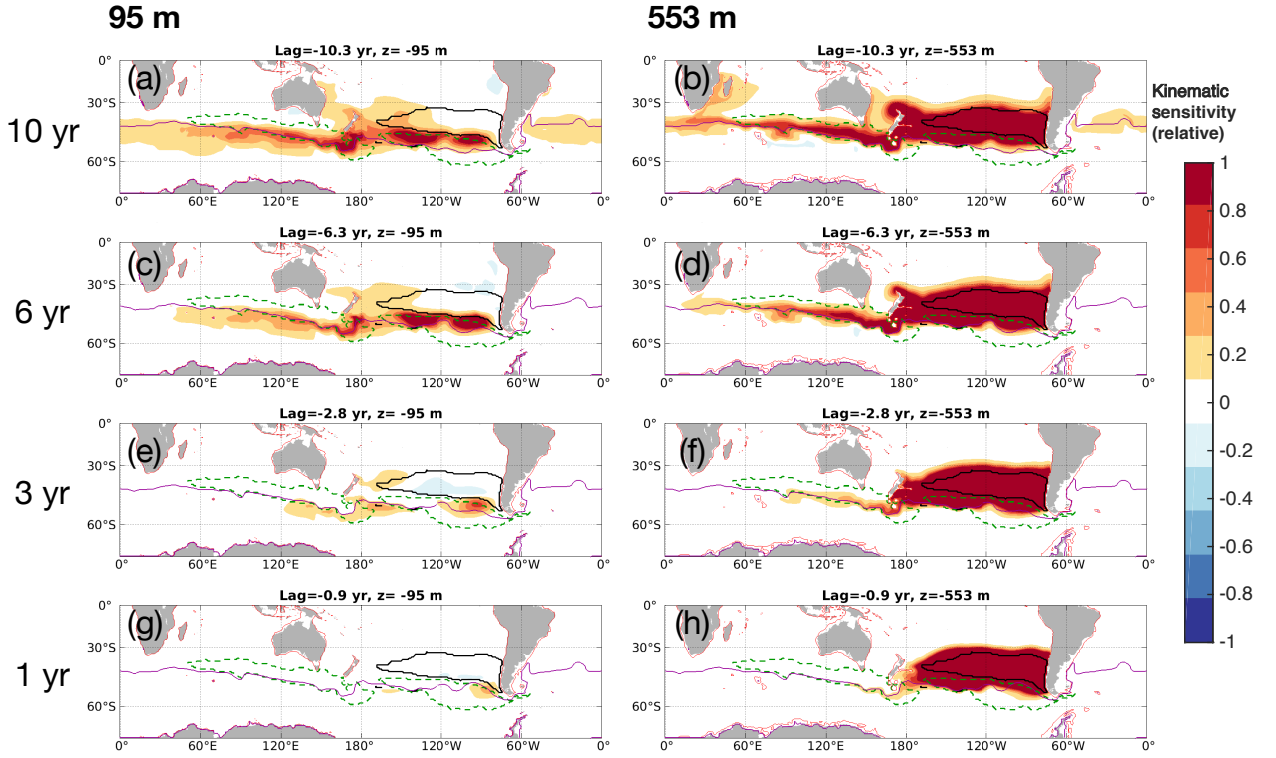


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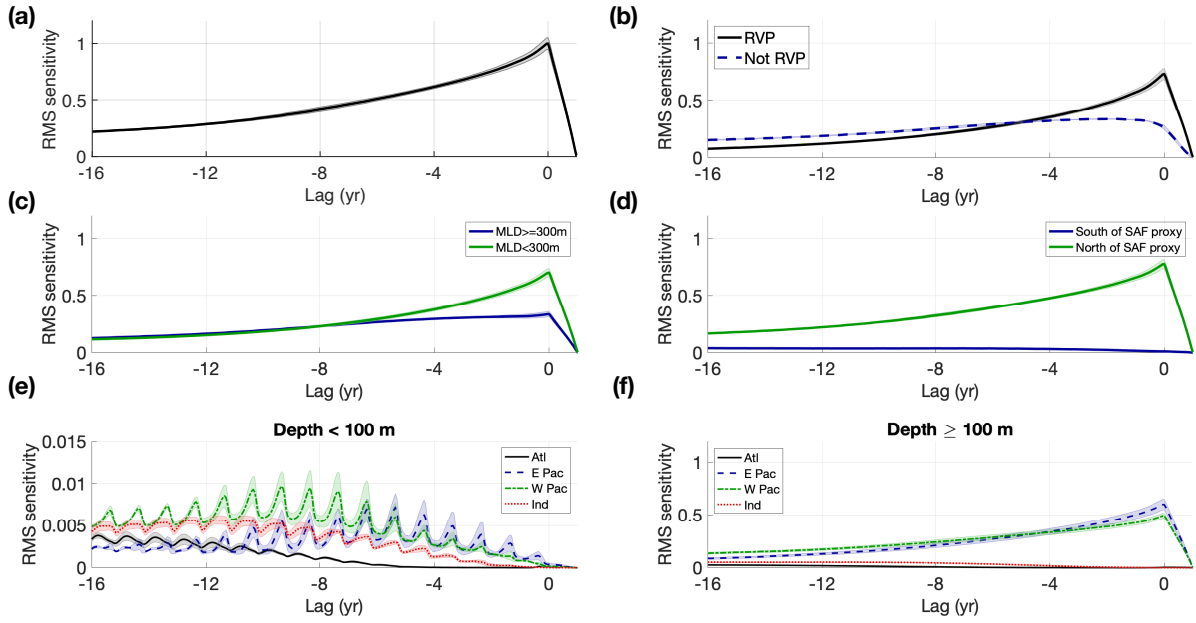


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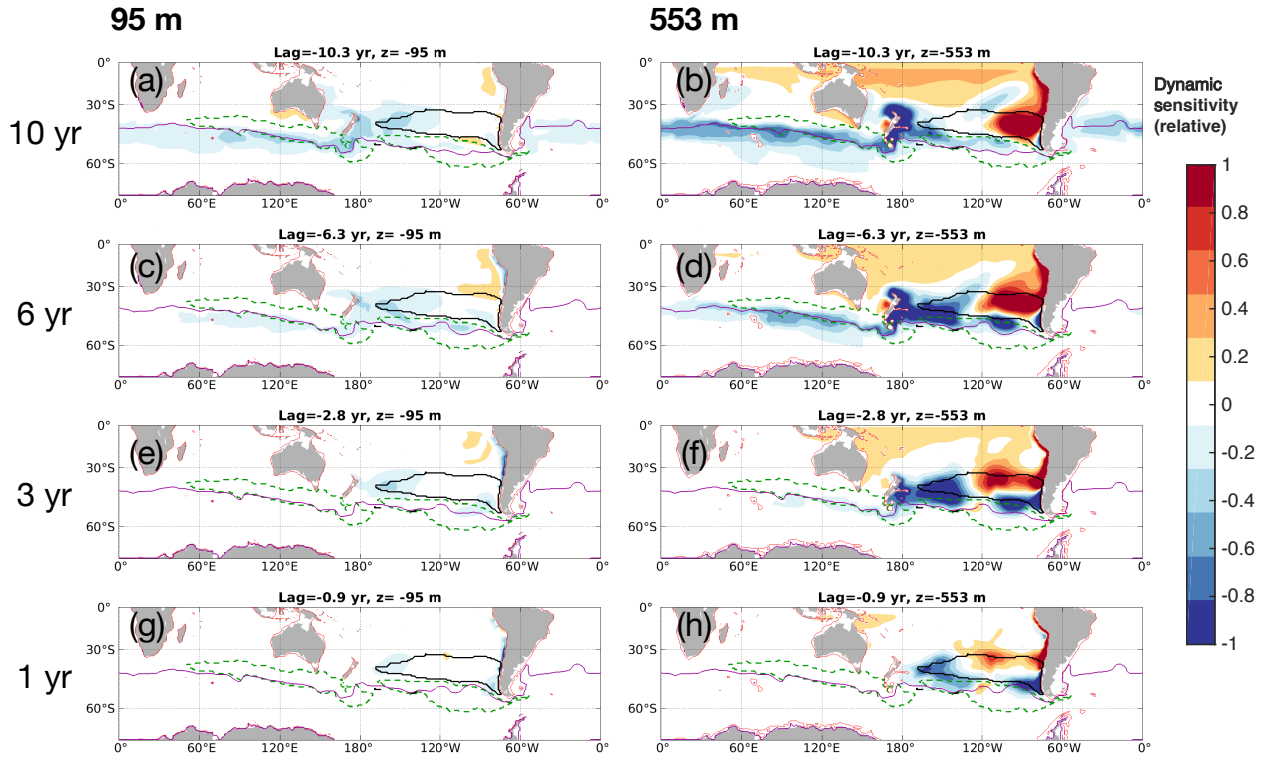


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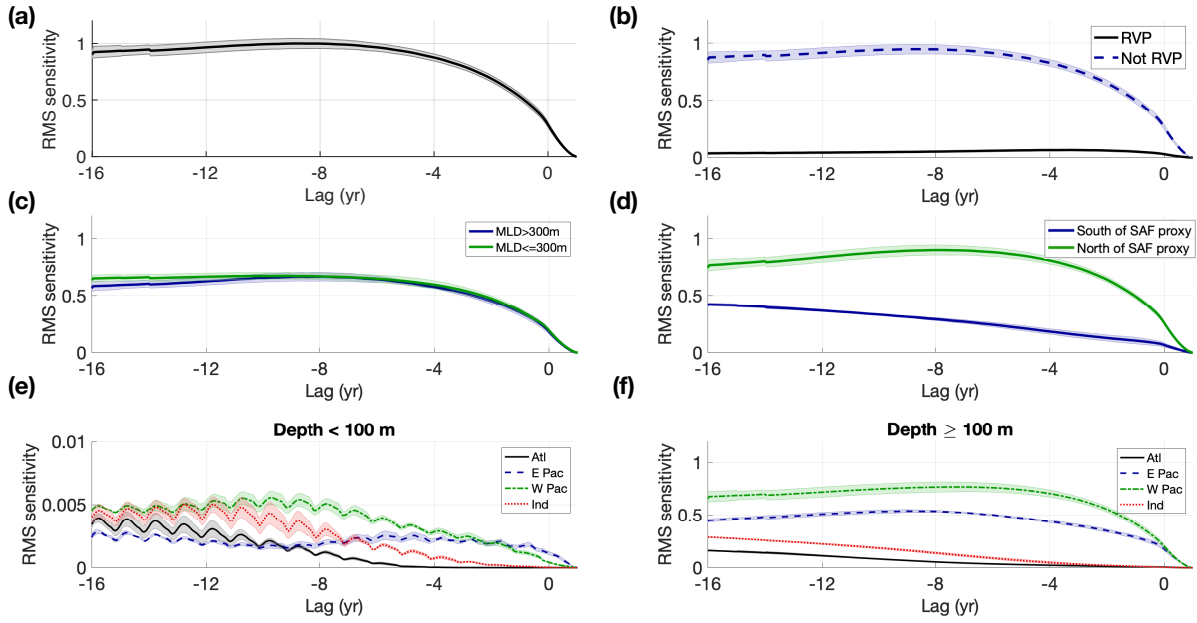


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Timescale increases from right to left.

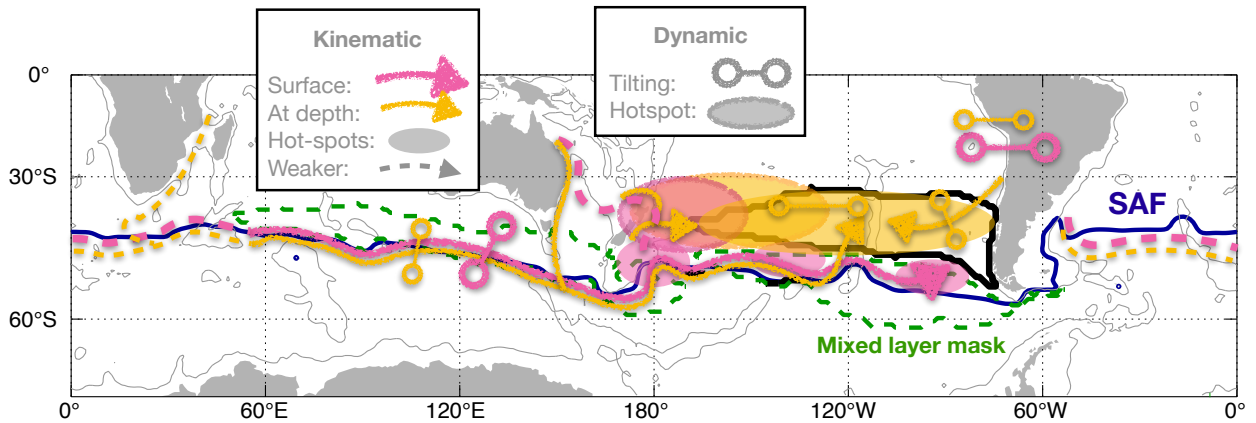


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