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► To cite this version:

Luiz Leal, Vladimir Sobes, Federico Rocchi, Danila Roubtsov, Jimmy Chow, et al.. RESONANCE EVALUATION OF GADOLINIUM ISOTOPES. International Conference on Nuclear Data for Science and Technology, ICAPP, May 2019, PEKIN, China. hal-02884352

HAL Id: hal-02884352 https://hal.science/hal-02884352v1

Submitted on 15 Dec 2020 $\,$

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RESONANCE EVALUATION OF GADOLINIUM ISOTOPES

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The objective of this paper is to present the results of an evaluation of the gadolinium isotopes with the main focus on the isotopes ¹⁵⁵Gd and ¹⁵⁷Gd. The evaluations were carried out in the resolved resonance region using the Reich-Moore formalism. The originality on the ¹⁵⁵Gd and ¹⁵⁷Gd evaluations is the addition of new high-resolution capture cross section measurements performed at the neutron time-of-flight, n_TOF facility for enriched samples and the statistical analysis of the resonance parameters. The resonance analysis was performed with the multilevel R-matrix code SAMMY together with the generalized least-squares technique based on the Bayes' theory.

KEYWORDS: R-matrix, data evaluation, experimental capture data

INTRODUCTION

Resolved resonance evaluation of the gadolinium isotopes were carried out using the multilevel R-matrix code SAMMY.[1] Capture measurements for enriched ¹⁵⁵Gd and ¹⁵⁷Gd isotopes and capture and transmission data for natural gadolinium were used in the evaluation. The natural abundances for the stable gadolinium isotopes range from 0.2 % to 24.84 %. The stable isotopes are ¹⁵²Gd (0.2%), 154 Gd (2.18%), 155 Gd (14.80%), 156 Gd (20.47%), 157 Gd (15.65%), 158 Gd (24.84%), and 160 Gd (21.86%), respectively. The isotopes with the highest thermal capture cross sections are ¹⁵⁵Gd with thermal cross section of about 65,000 barns and ¹⁵⁷Gd which the thermal cross section is of the order of 254,000 barns. Unquestionably, these two isotopes play significant role in criticality safety applications. Consequently, due to their importance a great deal of work has been devoted to the resonance evaluation of these isotopes to extract resonance parameters that described well the experimental data. In addition, detailed statistical analyses of the ¹⁵⁵Gd and ¹⁵⁷Gd resonance parameters were done on the basis of the s-wave resonances (l=0), since the penetrability for higher angular momentum $(l \ge 0)$ was not significant for the energy region where the resonance analysis was performed. The evaluations presented in this work extend the resolved resonance energy range of the ¹⁵⁵Gd and ¹⁵⁷Gd isotopes by taking advantage of the high resolution data measurements included in the evaluation. The isotopes ¹⁵⁶Gd, ¹⁵⁸Gd and ¹⁶⁰Gd were also reviewed and evaluated with the help of the natural capture cross section and transmission data. The less abundant isotopes ¹⁵²Gd and ¹⁵⁴Gd were not thoroughly evaluated but just used to complete the full set of resonance parameters in the analysis of the natural samples. Resonance parameters

for these isotopes where those listed in the Atlas of Neutron Resonances (ANR).[2] The upper energy of the resonance region for the new gadolinium evaluations is listed in Table 1. Also indicated in Table 1 is the upper resonance energy for the Joint Evaluated Fission and Fusion (JEFF) project, JEFF3.3 library.

Isotope	JEFF3.3	New Gd
	(eV)	evaluation
		(eV)
152 Gd	2658	2658
154 Gd	2760	2760
155 Gd	181.8	500
¹⁵⁶ Gd	1580	2250
157 Gd	215	500
¹⁵⁸ Gd	6037.6	10000
160 Gd	2883.7	10000

 Table 1 Resonance region upper energy

EXPERIMENTAL DATA

Transmission and capture data measurements performed at the Gaerttner time-of-flight (TOF) linac facility located at Rensselaer Polytechnic Institute (RPI) [3,4] and data taken at the n_TOF facility [5] were included in the evaluation. The measurements were done at the temperature of 293.6 K. Three transmission data and capture yield measurements for natural gadolinium done at RPI were included in the evaluation. In addition, two capture measurements for enriched ¹⁵⁵Gd and ¹⁵⁷Gd from RPI were also included in the analysis. From n_TOF, four measurements of different thicknesses for enriched ¹⁵⁵Gd and ¹⁵⁷Gd and ¹⁵⁷Gd were carefully analyzed. The descriptions of the measured data used in the evaluation are indicated in Table 2. Thermal values and uncertainties indicated in the ANR

were also considered in the evaluation.

METHOD OF EVALUATION

3.1 Fitting procedure

A Reich-Moore [6] resonance analysis and evaluation of the experimental data displayed in Table 2 were carried out with the computer code SAMMY. The experimental data were entered sequentially taking into account the temperature effects, data resolution, normalization, background effects, etc. At each step of the evaluation process an updated resonance parameter and covariance is generated, which are, subsequently feedback into the analysis of the next experimental data in the sequence. This process is repeated several times until a good fit of the experimental data is achieved.

Table 2 Experimental	transmission and	l capture data
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Data Set	Enrichment	Energy Range	Flight	Density
	(%)	(eV)	Path	(at/b)
			(m)	
		Natural Gadoliniu	m	
Transmission	-	0.2 - 300.0	25.585	7.806 x 10 ⁻⁴
Transmission (RPI)	-	0.3 - 500.0	25.597	1.566 x 10 ⁻³
Transmission (RPI)	-	0.3 - 1000.0	25.597	1.566 x 10 ⁻³
Capture (RPI)	-	0.2 - 1000.0	25.585	7.806 x 10 ⁻⁴
		¹⁵⁵ Gd		
Capture (RPI)	91.74	0.2 - 1000.0	25.567	3.083 x 10 ⁻⁴
Capture (n_TOF)	91.74	0.025 - 50.0	183.90	1.236 x 10 ⁻⁶
Capture (n_TOF)	91.74	1.0 - 1000.0	183.90	1.244 x 10 ⁻⁴
¹⁵⁷ Gd				
Capture (RPI)	90.96	0.2 - 1000.0	25.569	5.820 x 10 ⁻⁴
Capture (n_TOF)	88.32	0.025 - 50.0	183.90	5.753 x 10 ⁻⁶
Capture (n_TOF)	88.32	1.0 - 1000.0	183.90	2.340 x 10 ⁻⁴

An important figure-of-merit that indicates a good fitting of the experimental data is the χ^2 provided at the end of each SAMMY run. The resulting set of resonance parameters is expected to reproduce well the experimental data within the quoted experimental uncertainty. SAMMY also generates resonance parameter covariance. The latter is not part of the work presented in this paper. It should be stressed that while it looks a simple matter, the process of

fitting the experimental data is, actually a very demanding task. It requires to the evaluator a good perception of the theory as well a decent understanding of the experimental data used in the evaluation. One may derived a set of parameters that fit the data but it can be totally unphysical. Further tests, such as statistical test must be employed to guarantee that the resonance parameters fall close to the underlining physics. The latter issue will be discussed in this paper.

3.2 SAMMY fitting results

Comparisons of the SAMMY R matrix fit to the data of the n_TOF capture yield data for the ¹⁵⁵Gd and ¹⁵⁷Gd, in the energy region 0.02 eV to 10 eV, are shown in Figures 1. The top curve represents the capture yield for ¹⁵⁵Gd and ¹⁵⁷Gd comparison is the lower curve.



Fig. 1 Comparisons of SAMMY calculations with the resonance parameter (solid line) of the capture yield of the n_TOF data in the energy region 0.02 eV to 10 eV.

Comparison of the SAMMY fitting of the n_TOF and RPI capture yields for 155 Gd in the energy 50 eV to 250 eV are shown Figure 2.



Fig. 2 Comparisons of SAMMY calculations with the resonance parameter (solid line) of the capture yield of the n_TOF and RPI data in the energy region 50 eV to 250 eV for 155 Gd.



Fig. 3 Comparisons of SAMMY calculations with the resonance parameter (solid line) of the capture yield of the n_TOF and RPI data in the energy region 50 eV to 250 keV for 157 Gd.

Comparisons of the capture yield for the n_TOF and RPI data in the energy range of 250 eV to 500 eV for 155 Gd and 157 Gd are shown in Figure 4. Results for the fitting of the transmission data for natural gadolinium are shown in Figure 5 in the energy range from 3 eV to 300 eV. The contribution of each gadolinium isotopes is taken in account in the fitting of the transmission data.



Fig. 4 Comparisons of SAMMY calculations with the resonance parameter (solid line) of the n_{TOF} and RPI capture yield in the energy range 250 eV to 500 eV.



Fig. 5 Comparisons of SAMMY calculations with the resonance parameter (solid line) of the RPI transmission data for

natural gadolinium in the energy range 3 eV to 300 eV.

Capture and scattering cross section at thermal (0.0253 eV), capture resonance integral and capture Westcott factor calculated from the resonance parameters obtained in the evaluation for ¹⁵⁵Gd, ¹⁵⁶Gd, ¹⁵⁷Gd, ¹⁵⁸Gd, and ¹⁶⁰Gd are listed in Table 3.

 Table 3 Thermal values, resonance integral and Westcott factors

	Gadolinium Isotopes				
Quantities	¹⁵⁵ Gd	¹⁵⁶ Gd	¹⁵⁷ Gd	¹⁵⁸ Gd	¹⁶⁰ Gd
$\sigma_{s(\mathrm{barns})}$	61.58	4.88	1018.58	5.28	10.31
$\sigma_{\gamma(\mathrm{barns})}$	61857.42	1.93	254570.90	2.25	1.52
$I_{\gamma (\text{barns})}$	1567	99.40	814.98	78.27	10.32
g_{γ}	0.83326	1.00037	0.85204	1.00046	0.99994

The impact of the values indicated in Table 3 in benchmark calculations are under investigation.

4. Resonance Parameters Statistics

Statistical analysis of the *s*-wave resonance parameters derived in this work for ¹⁵⁵Gd and ¹⁵⁷Gd were carried out using the SAMDIST [7] component of the SAMMY code. The level spacing distribution and the Dyson-Mehta Δ_3 -statistics [8] were carefully examined.

4.1 Level spacing distribution

The isotopes ¹⁵⁵Gd and ¹⁵⁷Gd have identical target spin and parity of $3/2^{-}$ and consequently the compound nucleus total angular momentum J^{π} and parity, for *s*-wave resonances, are 1⁻ and 2⁻, respectively. Assuming a level spacing density proportional to 2J+1, the ratio of the average spacing of the two states will be, approximately, $D_1^{-}/D_2^{-} \sim 5/3$ which indicates that resonances in the spin state 1⁻ is less frequent as compared to spin state 2⁻. The average level spacing values for each spin and the mixed spin for ¹⁵⁵Gd and ¹⁵⁷Gd in the energy range from thermal to 500 eV are shown in Table 4.

Table 4 Average level spacing

$D(J^{\pi})$	¹⁵⁵ Gd	¹⁵⁷ Gd
	(eV)	(eV)
D(1 ⁻)	5.00±0.32	15.20±1.75
D(2)	2.46±0.12	5.99±0.43
D	1.64±0.35	4.31±1.80

Comparisons of the spacing distributions to the Wigner distributions are shown in Figures 6 and 7 for ¹⁵⁵Gd and ¹⁵⁷Gd, respectively.



Fig. 6 Wigner distribution of the resonances spacing in the energy 0 eV to 500 eV for 155 Gd.



Fig. 7 Wigner distribution of the resonances spacing in the energy 0 eV to 500 eV for 157 Gd.

The cumulative number of energy levels for ¹⁵⁵Gd and ¹⁵⁷Gd in the energy range 0 to 500 eV are shown in Figure 8 and Figure 9, respectively.



Fig. 8 Cumulative number of energy levels in the energy 0 eV to 500 eV for 155 Gd.



Fig. 9 Cumulative number of energy levels in the energy 0 eV to 500 eV for 157 Gd.

The results of the Δ_3 -statistics tests for ¹⁵⁵Gd and ^{157}Gd for the two J^{π} are displayed in Table 5. The Δ_3 -statistics provides a good insight on the long range correlation between the energy resonance levels, resonance spin assignment, missing levels, etc. It can be noted from Table 5 that in the energy range from 0 to 500 eV the $\Delta_3\text{-statistics}$ results are in better shape for ^{157}Gd as compared to the results for ¹⁵⁵Gd. This observation does not come as a surprise since the resonance spacing for 155 Gd is smaller than that for ¹⁵⁷Gd, that is, more energy levels in ¹⁵⁵Gd increases the chance of missing levels. Another observation is that the distribution of levels may not follow strictly a Gaussian orthogonal distribution from which the Wigner distribution is derived but rather a unitary orthogonal distribution. Further tests are under way on this issue.

Table 5 Δ_3 -statistics Test

	¹⁵⁵ Gd		¹⁵⁷ Gd	
\mathbf{J}^{π}	Theory	Experimental	Theory	Experimental
1-	0.459±0.109	0.744	0.344±0.109	0.519
2-	0.531±0.109	0.904	0.441±0.109	0.586

4. CONCLUDING REMARKS

This paper presents a new resonance evaluation of the gadolinium isotopes ¹⁵⁵Gd, ¹⁵⁶Gd, ¹⁵⁷Gd, ¹⁵⁸Gd, and ¹⁶⁰Gd. Compared to existing evaluations, the resonance energies ranges have been increased on the base of high resolution transmission and capture cross section data. Natural transmission and capture cross section data and capture data for enriched samples of ¹⁵⁵Gd and ¹⁵⁷Gd have been used in the analysis evaluation. New capture cross section values at thermal have been proposed. The impact of the new

evaluation in benchmark calculation is presently under study. Not described in this work nevertheless resonance parameter covariances for the evaluated isotopes have also been derived.

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