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6 **1 Effects of pore-scale precipitation on permeability and flow**

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32 Crystal growth rates and velocities are highly variable

33 Crystal shape, density and growth rate depends on mineral substrate

34 Different pore space geometry changes leads to different permeability changes

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36 **Keywords**

37 Calcite precipitation, growth rate, growth velocity, X-ray micro-tomography, permeability, porosity

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39 **Abstract**

40 The effects of calcite precipitation on porous media permeability and flow were evaluated with a
41 combined experimental and modeling approach. X-ray microtomography images of two columns packed
42 with glass beads and calcite (spar crystals) or aragonite (Bahamas ooids) injected with a supersaturated
43 solution ($\log \Omega = 1.42$) were processed in order to calculate rates of calcite precipitation with a spatial
44 resolution of 4.46 μm . Identification and localization of the newly precipitated crystals on the 3D images
45 was performed and results used to calculate the crystal growth rates and velocities. The effects of
46 carbonate precipitation was also evaluated in terms of the integrated precipitation rate over the length of
47 the column, crystal shape, surface area and pore roughness changes. While growth was epitaxial on calcite
48 spar, calcite rhombohedra formed on glass beads and clusters of polyhedrons formed on aragonite ooids.
49 Near the column inlet, calcite precipitation occurred preferentially on carbonate grains compared to glass
50 beads, with almost 100% of calcite spar surface area covered by new crystals versus 92% in the case of
51 aragonite and 11% in the case of glass beads. Although the experimental chemistry and flow boundary
52 conditions in the two columns were similar, their porosity-permeability evolution was different because
53 the nucleation and subsequent crystal growth on the two substrates (i.e., calcite spar and aragonite ooids)
54 was very different. The impact of mineral precipitation on pore-scale flow and permeability was evaluated
55 using a pore-scale Stokes solver that accounted for the changes in pore geometry. For similar magnitude
56 reductions in porosity, the decrease in permeability was highest within the sample that experienced the
57 greatest increase in pore roughness. Various porous media models were generated to show the impact of
58 different crystal growth patterns and pore roughness changes on flow and permeability-porosity
59 relationship. Under constant flow rate boundary conditions, precipitation resulted in an increase in both
60 the average and maximum velocities. Increases in pore roughness led to a more heterogeneous flow field,
61 principally through the effects on the fastest and slowest velocities within the domain.

62 **Keywords**

63 Calcite precipitation, growth velocity, growth rate, porosity, permeability, X-ray microtomography.

64 **1. Introduction**

65 Mineral precipitation plays a critical role in many geological processes, including sediment diagenesis,
66 hydrothermal circulation and alteration, biomineralization, and CO₂ sequestration [1-6]. Precipitation of
67 minerals like calcite can cause significant reduction in permeability and a reorganization of the flow field

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4 68 in reservoirs by altering the shape, size and connectivity of the pores, the roughness of their surfaces, and
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6 69 by partial or complete obstruction of flow in pore throats. Developing a full understanding of the
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8 70 feedbacks between geochemical reactions and flow and transport characteristics at the reservoir scale
9
10 71 requires investigations of precipitation mechanisms at the pore-scale.

11 72 Three rate-determining mechanisms have been identified to describe calcite precipitation from aqueous
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13 73 solution [7], (i) diffusive transport and/or adsorption of lattice ion, (ii) surface spiral growth, and (iii)
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15 74 surface nucleation involving polynuclear growth. These mechanisms involve different dependences of the
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17 75 precipitation rate, r_{ppt} ($\text{mol}\cdot\text{m}^{-2}\cdot\text{s}^{-1}$), on the supersaturation ($1-\Omega$):

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20 76
$$r_{ppt} = k \left[\exp\left(\frac{m \Delta G}{R^* T}\right) - 1 \right]^n = k (\Omega^m - 1)^n \quad (1)$$

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24 77 with k the rate constant ($\text{mol}\cdot\text{m}^{-2}\cdot\text{s}^{-1}$), ΔG the Gibbs free energy change of the overall reaction ($\text{J}\cdot\text{mol}^{-1}$),
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26 78 Ω the saturation index, R^* the gas constant ($\text{J}\cdot\text{K}^{-1}\cdot\text{mol}^{-1}$), T the absolute temperature (K), and n and m the
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28 79 semi-empirical constants that depend on the kinetic behavior involved in the chemical reaction. Here r_{ppt}
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30 80 is the rate normalized against the reactive surface area, S_r (m^2). A value of $n = 1$ (linear rate law) has been
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32 81 attributed to crystallization limited by adsorption of lattice ions [7], and in some cases to a multi-source
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34 82 spiral growth mechanism [8]. A second order equation ($n = 2$) may be used to describe growth at single
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36 83 screw dislocation by the spiral mechanism [8], while higher order dependences ($n > 2$) can be applied to
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38 84 growth both at screw and edge dislocation [9] or growth by 2D nucleation [8].

38 85 Porosity-permeability evolution in rocks results from the interplay between reaction kinetics and mass
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40 86 transport. Two parameters, i.e., the Péclet (Pe) and Damköhler (Da) numbers, defined locally as:
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42 87 $\text{Pe} = u L^* / D_m$ and $\text{Da} = k_r L^{*2} / D_m$, where u is the fluid velocity ($\text{m}\cdot\text{s}^{-1}$), D_m is molecular diffusion ($\text{m}^2\cdot\text{s}^{-1}$),
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44 88 k_r is a first order kinetic constant (s^{-1}), and L^* is a characteristic length (m), e.g., the pore size [10, 11],
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46 89 are commonly used to describe the physical evolution of the pore space to capture the possible feedback
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48 90 between the flow regime and the geochemical alteration, e.g., [12-15]. Although precipitation of calcite in
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50 91 natural systems is often assumed to be kinetically-controlled, i.e., independent of mass transport rates [16],
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52 92 precipitation rates can vary along flow paths due to changes in the fluid saturation index, temperature and
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54 93 pressure, and the presence of bacterial communities, inhibitors or (in)organic impurities [17-31]. In
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56 94 addition, several studies have reported that both crystal location, morphology and growth rate are closely
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58 95 related to the mineral surface properties [17, 32-36], which make mineralogical heterogeneities of the pore
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60 96 walls an important factor for understanding calcite nucleation and growth in porous media.

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97 The distribution of new precipitates leads to a decrease of porosity with resulting changes in the pore-size
98 distribution, and in some cases, in pore roughness. These effects in turn can modify flow and permeability.
99 Description of permeability evolution in the presence of pore-scale heterogeneities is still challenging, as
100 micro-heterogeneities are often ignored in flow and reactive transport codes at the macroscopic (or Darcy)
101 scale.

102 Most permeability reduction models consider at least an exponential or power-law dependence of
103 permeability to porosity, and sometimes integrate the effects of tortuosity, shape factor, specific surface
104 area or pore shape [37-39]. Using a simple model of cylindrical pores, Ghezzehei [40] established
105 differences in permeability of up to three orders of magnitude between uniform and non-uniform
106 precipitation models. The uniform model corresponds to a coating of the cylindrical pore with precipitate
107 of uniform thickness, while in the non-uniform model the pore contains several rectangular precipitate
108 crystals of various sizes separated by equal distance. The author proposed a modified formulation of the
109 Kozeny-Carman model [37] to correct for pore-scale distribution of the precipitates. Although this model
110 has been extensively modified to improve the estimation of the permeability when porosity changes, it
111 seems that modifications at micro-scale are too diverse to be captured by only a small number of macro-
112 scale parameters such as porosity and/or tortuosity of cylindrical pores.

113 Advances in non-invasive and non-destructive techniques like X-ray microtomography (XMT) have
114 allowed for characterization of flow or geometry changes over time at the pore-scale during dynamic
115 processes or experiments, see e.g., [14, 41] for a review. In particular, characterization of pore geometry
116 and pore-space distribution is possible [42-44]. The dynamics of the fluid-rock interface can also be
117 followed through time [45, 46]. Even if it is difficult to define pore roughness to the extent that it depends
118 on the resolution and scale of observation, the changes in pore or surface roughness can be examined [47].
119 The 3D geometry derived in this way can be used for discrete calculations of porous media properties such
120 as permeability and transport, e.g., see [48] for a review.

121 The present study is focused on the effects of calcite precipitation on changes in pore geometry,
122 permeability and flow at the pore-scale. These effects are investigated both experimentally and
123 numerically through an idealized representation of porous rock columns representing very simple analogs
124 of a calcareous sandstone. Glass beads and calcite spar crystals or aragonite ooids were selected to test the
125 effects of the substrate composition on precipitation under similar chemistry and flow conditions. The
126 characterization of the pore structure, which is crucial for determination of flow in porous media, was
127 carried out before and after precipitation experiments using X-ray microtomography (XMT) at a
128 resolution of 4.46 μm . The first part of the paper is devoted to the identification, localization and
129 characterization of the newly precipitated crystals and to the determination of crystal growth velocity and

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4 130 growth rate from XMT images. The discrete 3D geometry is then used as an input for permeability
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6 131 calculations, and several models of precipitation were generated and examined to test the effects of
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8 132 precipitation-induced micro-scale changes on permeability and flow properties.
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11 133 **2. Experimental and modeling procedure**

14 134 Experiments were designed to study the effects of precipitation and substrate composition in columns
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16 135 packed with glass beads and calcite crystals or aragonite ooids. X-ray microtomography (XMT) allowed
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18 136 for the investigation of the crystallization processes in porous media by monitoring changes in porosity,
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20 137 pore geometry, and surface roughness. Because of the high initial porosity and permeability of the
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22 138 columns, experimental determination of permeability based on the pressure difference between the inlet
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24 139 and outlet was not possible. Instead, permeability was computed directly from the 3D discrete geometry of
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26 140 the samples after image segmentation using a Stokes solver.

27 141 **2.1. Precipitation experiments**

29 142 Two samples (REAC-1 and REAC-3) consisting of the contents of cylindrical plug-flow column
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31 143 reactors of 6.5 mm diameter and approximately 12 mm length were prepared by packing a mixture of
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33 144 glass beads and crushed calcite spar or aragonite ooids (Table 1, see also [36] for more details about
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35 145 sample REAC-1). The glass beads used here were acid-washed soda-lime glass in the range 425-600 μm
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37 146 (Sigma G8772). Calcite that began as centimeter-size Iceland spar crystals was crushed in an agate mortar
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39 147 and then sieved in the range 355-500 μm . Aragonite ooids, which come from oolitic sand sampled on the
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41 148 seabed in Bahamas, were also sieved in the range 355-500 μm . All the materials were washed with
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43 149 deionized water, cleaned ultrasonically, and dried at 50°C prior to column assembly.

44 150 Each reactor was injected with a 50/50 mixture of 0.002 M CaCl_2 and 0.004 M NaHCO_3 at a flow rate
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46 151 of $0.5 \text{ cm}^3 \cdot \text{h}^{-1}$ ($Q = 1.38 \cdot 10^{-10} \text{ m}^3 \cdot \text{s}^{-1}$) using a multi-channel syringe pump (Figure 1), inducing a calcite
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48 152 saturation index ($\log \Omega$) of 1.42 (see [36] for calculation details). The reactors were saturated with
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50 153 deionized water under vacuum prior to experiments and injected with an acidic solution (deionized water
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52 154 + HCl, pH 4.0) for 5 h at a flow rate of $20 \text{ ml} \cdot \text{h}^{-1}$ ($Q = 5.56 \cdot 10^{-9} \text{ m}^3 \cdot \text{s}^{-1}$) to further clean the mineral
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54 155 surfaces and eliminate fine particles.

55 156 The experiments were conducted at room temperature (22°C) over either 30 days (REAC-1) or 15 days
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57 157 (REAC-3). The second experiment (REAC-3) was stopped at 15 days after clogging of the column inlet.
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59 158 The pH of both the inlet and outlet solutions was measured using an Orion pH electrode. In addition,
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61 159 major and minor elements were measured in solution by inductively coupled plasma mass spectroscopy
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(ICP-MS), after acidification with nitric acid (pH < 2). After the experiments were finished, materials of the inlet and exit portions of each column were collected for scanning electron microscopy (SEM) observations. Raman spectroscopy was used to determine the identity of the calcium carbonate precipitates. Only calcite was found as a precipitate on both carbonate or glass bead surfaces.

The overall calcite precipitation rate (mol·s⁻¹) through the columns is calculated from the chemical analysis of the column effluent (outlet) as:

$$r_{ppt} = \frac{1}{\nu_{cal}} \frac{\Delta V_{cal}}{\Delta t} = Q \times \overline{\Delta Ca} \quad (2)$$

with V_{cal} the volume of precipitated calcite, Q the volumetric flow rate (m³·s⁻¹), ν_{cal} the molar volume of calcite (m³·mol⁻¹), and $\overline{\Delta Ca}$ the average variation in calcium concentration (mol·m⁻³) between the inlet and the outlet of the column during the experiment.

Table 1. Composition of the plug-flow columns.

Sample name	REAC-1	REAC-3
	Precipitation on calcite spar and glass beads	Precipitation on aragonite ooids and glass beads
Mass (g)	0.6300	0.6699
Carbonate (% w)	25*	26**
Glass beads (% w)	75	74

* calcite spar; ** aragonite

Figure 1. Schematic of the plug-flow column experimental setup

2.2. Image acquisition and processing

The geometry of the two packed columns was characterized with X-ray microtomography (XMT) at Beamline 8.3.2 at the Advanced Light Source of Lawrence Berkeley National Laboratory. The optical system used in the experiment, i.e., a Mitutoyo ×2 lens combined with a Cooke PCO4000 CCD camera, provided a spatial resolution of 4.46 μm (voxel size).

Two XMT datasets were collected for each column, both before (initial time, t_0) and after (final time, t_1) each experiment. The X-ray source was diffracted through a multilayer, yielding a focused monochromatic and parallel beam with energy of 28 keV. Exposure time for each radiograph was 700 ms. 3D image reconstruction was performed from 1441 radiographs taken over 180° using the software package Octopus [49] and image processing carried out using Avizo® software on data sets of $1600 \times 1600 \times 3000$ voxels. Image processing procedure, including noise reduction, segmentation and registration is described in [36].

Several volumes of interest (VOI) were extracted from the original datasets for quantification. Column-scale volumes of $1000 \times 1000 \times 1000$ voxels and pore-scale sub-volumes of either $300 \times 300 \times 300$ voxels or $350 \times 350 \times 350$ voxels were extracted near the column inlets, where precipitation was highly visible. The VOI were named according to the volume size, column name, and stage of experiment (Table 2).

Table 2. Description of the volumes of interest (VOI) extracted from original datasets for quantification and flow computation.

Volume of interest (VOI)	Sample	Stage	Size (voxels)
Vol-REAC-1-initial	REAC-1 Precipitation on calcite spar and glass beads	initial	$1000 \times 1000 \times 1000$
Vol-REAC-1-final		final	
SubVol-REAC-1-initial		initial	$300 \times 300 \times 300$
SubVol-REAC-1-final		final	
Vol-REAC-3-initial	REAC-3 Precipitation on aragonite oids and glass beads	initial	$1000 \times 1000 \times 1000$
Vol-REAC-3-final		final	
SubVol-REAC-3-initial		initial	$350 \times 350 \times 350$
SubVol-REAC-3-final		final	

After segmentation of the XMT images, the total porosity of the sample ϕ is the ratio of the pore space to the total volume. The local porosity $\phi(z)$ can also be calculated for each slice along the flow axis. It is also possible to calculate the rate of calcite precipitation from the difference in the number of solid voxels (or pixels) observed at the end (t_1) and at the beginning (t_0) of the experiment. The overall precipitation rate ($\text{mol}\cdot\text{s}^{-1}$) obtained from XMT is expressed as:

$$r'_{ppt} = \frac{1}{v_{cal}} \frac{\Delta V'_{cal}}{\Delta t} = \frac{\sum_{VOI} N_{cal}(t_1) - N_{cal}(t_0) \times V_{pix}}{t_1 - t_0} \quad (3)$$

with V'_{cal} the volume of precipitated calcite calculated from XMT, N_{cal} the number of calcite voxels in the image, and V_{pix} the volume of a voxel ($V_{pix} = 4.46 \times 4.46 \times 4.46 \mu\text{m}^3$).

2.2.1. Mineral separation and quantification

Every grain (i.e., either calcite spar or aragonite ooid and glass bead) was also individually mapped in the individual volumes of interest. From the binary images, the 3D Euclidean distance map based on a 3-4-5 chamfer distance transform [50, 51] was computed in order to identify the centroids of the different grains. Following this, a fast 3D watershed algorithm [52] was applied to invert the distance map in order to find the separation lines around the different centroids that correspond to the contact areas between grains. Once separated, every grain and contact area was labeled. For REAC-1, it was possible to discriminate between spar crystal and glass beads based on the difference in sphericity of the grains using a threshold value of 0.83. The sphericity, ψ , is defined as: $\psi = \pi^{1/3} (6V_{grain})^{2/3} / S_{grain}$, with V_{grain} the volume and S_{grain} the surface area of the grains, respectively. The procedure is illustrated in Figure 2. For REAC-3, the morphological descriptors were less efficient for this sample since glass beads and aragonite ooids have very similar morphologies. As a result, some of the grains had to be labeled manually from visual inspection, particularly the truncated ones near the VOI edges.

Once labeled, the surface area S_{mat} (m^2) of the different materials (i.e., calcite spar, glass beads or aragonite ooids) could be calculated as the number of pore-material pixels $N_{mat-fluid}$ multiplied by their surface-area ($S_{pix} = 4.46 \times 4.46 \mu\text{m}^2$). The geometric surface-area normalized to the volume of material, S'_{mat} ($\text{m}^2 \cdot \text{m}^{-3}$), is defined by:

$$S'_{mat} = \sum_{V_t} N_{mat-fluid} \times S_{pix} / V_{mat} \quad (4)$$

with V_{mat} the material volume. The ratio of final surface area of the different materials to their initial surface area, S/S_0 , with S_0 (m^2) the surface area before experiment., provides a good indicator of the changes in surface roughness during the experiment, i.e.,:

$$S/S_0 = \sum_{mat} S_{mat} / S_{0mat} \quad (5)$$

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4 228 **Figure 2. Glass beads and calcite spar crystals separation and labeling. (a) Grayscale image (size:**
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6 229 **4.46×4.46 mm), Vol-REAC-1-initial (calcite spar: high grayscale, glass bead: intermediate**
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8 230 **grayscale, porosity: low grayscale). (b) Segmentation and separation of every grain (in blue), i.e.,**
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10 231 **calcite spar and glass beads, in the image. (c) Shape-based identification and labeling of individual**
11 232 **glass beads (color scale). Note that mislabeling of truncated spheres (those with a sphericity < 0.83)**
12 233 **occurred near the edges of the image.**
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17 235 **2.2.2. Precipitation localization and growth velocity**

20 236 Image difference between final and initial pore-scale sub-volume was also computed for REAC-3 in
21 237 order to identify which of the new crystals had grown on aragonite ooids versus the glass beads. It was not
22 238 possible to apply the procedure to REAC-1 due to slight displacement of some glass beads. Where a slight
23 239 displacement occurred, volume subtraction failed to map the target volume to the reference one. Newly
24 240 precipitated crystals were labeled individually after cleaning with an erosion-dilatation filter. The growth
25 241 velocity ($\mu\text{m}\cdot\text{d}^{-1}$) of individual crystals was calculated based on the change in crystal length normal to the
26 242 substrate surface, as:
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$$32 \quad v_{\text{growth}} = \frac{d \mathbf{I}_{\text{fs}} \cdot \mathbf{n}}{dt} \quad (6)$$

35 244 with \mathbf{I}_{fs} the vector position of the fluid-solid interface and \mathbf{n} the normal to the substrate surface. Crystal
36 245 length normal to the substrate surface was determined by combining a 3D Euclidean distance map of the
37 246 newly formed crystals compared to the initial material. The map is based on a 3-4-5 chamfer distance
38 247 transform [50, 51], and labels each pixel of the newly formed crystals with the distance to the nearest
39 248 boundary pixel [53], starting from the initial position of the fluid-solid interface. Local maximum of the
40 249 distance map for every newly formed crystal gives their maximal length. In addition, the average crystal
41 250 growth rate ($\text{mol}\cdot\text{s}^{-1}$) can be calculated from the volume of precipitated crystals:
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$$48 \quad r_{\text{growth}} = \frac{1}{V_{\text{cal}}} \int_{V_{\text{cryst}}} \frac{dV_{\text{cryst}}}{dt} \quad (7)$$

51 252 with V_{cryst} the crystal volume (m^3). The procedure is illustrated in Figure 3.
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57 254 **Figure 3. Crystal growth velocity and growth rate distribution within pore-scale volume SubV-**
58 255 **REAC-3. (a) Identification between newly precipitated crystals on glass beads and on aragonite**
59 256 **ooids from image difference. (b) Distance map of the new crystals normal to their substrate; the**
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maxima are dotted in blue. The color scale denotes the distance of every pixel of the newly formed crystal from the fluid-rock interface they are attached to. (c) Growth velocity distribution ($\mu\text{m}\cdot\text{d}^{-1}$) based on maximum crystal length. (d) Crystal growth rate ($\text{mol}\cdot\text{s}^{-1}$) based on crystal volume. Note that although the procedure is illustrated here in 2D, it was applied in 3D in this study.

2.3. Flow model

The complete permeability tensors \bar{k} were obtained by solving the closure problem associated with the volume averaging of the Stokes' equations within the fluid phase [54], assuming that the inertial forces were negligible:

$$\begin{cases} \nabla \mathbf{u} = 0 \\ \mu \nabla^2 \mathbf{u} - \nabla P = 0 \end{cases} \quad (8)$$

with $\mathbf{u} = [u_x, u_y, u_z]$ the fluid velocity vector, μ the viscosity and P the pressure. A solver implemented in Avizo[®] under the XLabHydro package was used. Details about the solver can be found in [55]. Spatial discretization is carried out following the classical finite volume method [56]. The discretized system is solved numerically on a staggered grid with a cubic grid spacing of $4.46 \mu\text{m}$ defined directly from the XMT images. Time integration is carried out using a fully explicit scheme for which a stability criterion of 10^{-4} was chosen so as to obtain a symmetrical tensor with a precision of 3 digits. Determination of the permeability tensor components in the x - and y -directions (i.e., perpendicular to the flow direction) directly on the columns was not possible due to their elongated cylindrical shape, as the use of periodic boundary conditions requires that the inflow and outflow boundaries match. The calculations were performed instead on cubic volumes extracted from the columns, i.e., the column-scale volumes of interest of REAC-1, and also on the different pore-scale volumes of REAC-1 and REAC-3 (Table 2).

The velocity field was also resolved in the fluid flow direction in the pore-scale volumes using a flow solver of the Stokes' equations. The advantage of this solver over the calculation of the permeability tensor is that it is possible to calculate the flow properties in the conditions of experiments, e.g., the velocity magnitude (at the voxel center): $u = \|\mathbf{u}\| = \sqrt{\mathbf{u} \cdot \mathbf{u}}$, and the average velocity magnitude: \bar{u} . Boundary conditions are constant volumetric flux ($q = Q \times s / S$, with Q the experiment flow rate, S the column sectional area, s the considered volume sectional area) at the inlet, constant pressure at the outlet (10^5 Pa , i.e., atmospheric pressure) and no-flow for the other faces. Inflow and outflow zones are added to accommodate the complex porous shape at the inlet and outlet of the volumes and to increase numerical stability. Periodicity is forced for fluid velocity at the inlet and outlet. The pressure is solved at the center

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4 287 of the grid and the velocity is solved at the boundary of the grid using an artificial compressibility scheme
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6 288 [57]. After pressure and flow fields are resolved, pressure drop (ΔP) between inlet and outlet is used to
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8 289 calculate permeability k_D (m^2) from Darcy's law [58]. Darcy's seepage velocity is calculated as:
9 290 $u_p = u_D / \phi$, with u_D Darcy's average velocity ($u_D = q / s$).

13 291 3. Results

16 292 3.1. Calcite nucleation and growth

18 293 Precipitation occurred preferentially near the inlet of the columns where the supersaturation with
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20 294 respect to calcite was the highest. In this region, the newly precipitated crystals are both larger and more
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22 295 numerous. The shape, number, and size of the crystals, however, are closely linked to the substrate for
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24 296 precipitation (Figure 4) and the saturation index. Nucleation is favored on calcite compared to aragonite
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26 297 and glass beads. Crystal growth on calcite spar is epitaxial, i.e., crystallographically oriented on the single
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28 298 spar surfaces; as a result, the newly formed crystals cover the spar surface as a homogenous layer (Figure
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30 300 5a). In contrast, single individual crystals are observed on both aragonite and glass beads, but with
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32 301 differing shapes, orientations, densities and sizes (Figure 4 and Figure 5). The surface of glass beads
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34 302 display sparsely distributed rhombohedral crystals (Figure 5b), while the surface of aragonite ooids is
covered by many polyhedral crystals that are aggregated (Figure 5c,d).

35 303 For sample SubVol-REAC-1, the volume of new crystals precipitated on calcite spar was 4.7 times
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37 304 higher than on glass beads. The new crystals cover almost 100 % of the surface area of calcite spar, except
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39 305 near the contact areas with glass beads (Figure 6). For sample SubVol-REAC-3, the volume of new
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41 306 crystals precipitated on aragonite ooids was 3.7 times higher than on glass beads. In the case of SubVol-
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43 307 REAC-3, the new crystals cover only 11% of the surface area of glass beads, while 92% of the surface
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45 308 area of aragonite ooids was covered. A higher density of nucleation sites on calcite spar or aragonite
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47 309 compared with glass beads likely explains the greater number of new crystals precipitated on carbonate
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49 310 phases.

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52 312 **Figure 4. XMT observation of pore-scale volumes extracted near columns inlets before (initial)**
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54 313 **and after (final) experiment. (a) SubVol-REAC-1 ($1.34 \times 1.34 \times 1.34 \text{ mm}^3$). (b) SubVol-REAC-3**
55 314 **($1.56 \times 1.56 \times 1.56 \text{ mm}^3$). Precipitation is visible on the final volumes.**

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4 317 **Figure 5. Scanning electron microscope (SEM) observations of the surface of (a) calcite spar**
5 **crystal, (b) glass bead, and (c) and (d) aragonite ooid after experiment. Note that the areas where no**
6 318 **precipitation occurred in (c) correspond to contact areas with other grains.**
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14 322 **Figure 6. 3D rendering of precipitation on a calcite spar crystal, sample REAC-1. The yellow**
15 **volume was obtained by image difference and corresponds to the precipitation on the outer surface**
16 323 **of spar crystals (transparent) during the experiment. Precipitation is not detectable near the contact**
17 324 **with glass beads (arrows).**
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24 327 **3.2. Precipitation and growth rate**

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26 328 The chemical composition of the inlet and outlet solutions show a decrease of calcium concentration over
27 the length of the column ($\Delta Ca > 0$) as a result of calcite precipitation. At the column scale, the overall
28 329 precipitation rate is more than two times higher for the experiment with calcite spar and glass beads
29 the length of the column ($\Delta Ca > 0$) as a result of calcite precipitation. At the column scale, the overall
30 330 precipitation rate is more than two times higher for the experiment with calcite spar and glass beads
31 ($r_{ppt, REAC-1} = 5.5 \cdot 10^{-11} \text{ mol}\cdot\text{s}^{-1}$ for $\overline{\Delta Ca} = 4.0 \cdot 10^{-4} \text{ M}$) compared to aragonite ooids and glass beads
32 331 ($r_{ppt, REAC-3} = 2.9 \cdot 10^{-11} \text{ mol}\cdot\text{s}^{-1}$ for $\overline{\Delta Ca} = 2.0 \cdot 10^{-4} \text{ M}$). By comparison, rates measured from XMT are 5.2
33 $10^{-11} \text{ mol}\cdot\text{s}^{-1}$ for REAC-1, and $1.8 \cdot 10^{-11} \text{ mol}\cdot\text{s}^{-1}$ for REAC-3. As glass bead and carbonate proportions are
34 332 almost equal for the two columns, different rates indicate that the rate of precipitation on the calcite spar
35 crystals is higher than on the aragonite ooids. Differences in the rate can be explained by the differences in
36 333 reactivity of calcite surfaces compared to those of aragonite. In fact, newly precipitated calcite was
37 observed with SEM on calcite spar over the entire length of the column, as there is a nucleation threshold
38 334 for aragonite (which is nevertheless lower than for glass beads on which new crystals were observed to
39 335 nucleate only within the first 1.8 mm of the columns).
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48 340 At the pore-scale, the crystals grown on glass beads were often larger than those grown on the nearby
49 aragonite ooids even though the saturation index and the flow conditions were nearly identical. In
50 341 addition, the crystals formed on the two substrates show very different growth velocities and average
51 342 growth rates (Figure 7). Statistical details calculated on pore-scale Subvol-REAC-3 are given in Table 3.
52 343 The growth velocity of new crystals on glass beads is higher on average but also more disperse. In
53 344 contrast, the average growth rate of individual crystals is lower on aragonite ooids. This is largely because
54 many of the new crystals labeled on aragonite are actually aggregated and form clusters, the volume of
55 345 which is higher than for individual crystals.
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Table 3. Statistics of new calcite crystals precipitated in pore-scale volume SubVol-REAC-3

SubVol-REAC-3	$\bar{v}_{\text{growth}} \pm \sigma$ ($\mu\text{m}\cdot\text{d}^{-1}$)	$\bar{r}_{\text{growth}} \pm \sigma$ ($\times 10^{-15}$ mol $\cdot\text{s}^{-1}$)
New crystals on aragonite	3.9 ± 2.8	1.7 ± 4.8
New crystals on glass beads	4.6 ± 5.0	1.1 ± 2.2

Figure 7. Growth velocity distribution of new calcite crystals on glass beads and aragonite ooids for SubVol-REAC-3. (a) Pore-scale volume with newly precipitated calcite crystals on glass beads (in red) and aragonite (in orange). (b) Growth velocity distribution of new crystals on glass beads. (c) Growth velocity distribution of new crystals on aragonite ooids; note that crystals often form clusters. (d) Growth velocity ($\mu\text{m}\cdot\text{d}^{-1}$) distribution of new crystals on glass beads (left) and aragonite (right). (e) Growth rate (mol $\cdot\text{s}^{-1}$) distribution of new crystals on glass beads (left) and aragonite (right).

3.3. Changes in surface area and roughness

Changes in surface roughness were studied through changes of the surface area from pore-scale volumes. Because of the strong dependence of precipitation on the mineral substrate, the roughness and surface area changes are very different between the two columns. For SubVol-REAC-1 (i.e., experiment using calcite spar and glass beads), the roughness of the calcite spar crystals changes at the micro-scale due to 2D heterogeneous and spiral growth, but these changes occur at a very small scale (Figure 5a, see also [36] for discussion) and are not visible at the resolution of XMT imaging, thus the crystals appear to remain flat. The total surface area increases by only 5% due to calcite precipitation. We calculated an increase of the surface area of calcite spar as 7%, while the increase on glass beads was 3%. In contrast, the growth of individual crystals at the surface of aragonite and glass beads induces complex precipitation patterns associated with an increase in the roughness of the pores. For SubVol-REAC-3 (i.e., experiment with aragonite ooids and glass beads), the entire surface area increases by 20% as a result of calcite precipitation. An increase of the surface area of aragonite of 41% was calculated, while the increase was only 7% for the glass beads.

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3.4. Porosity, permeability and flow changes

Estimated porosity and permeability changes are presented for the different volumes of interest in Table 4. Both porosity and permeability changes differ between the column-scale and the pore-scale volumes, as precipitation was not uniform over the column length. As mentioned above, precipitation primarily occurred near the inlet of the columns where the supersaturation was the highest, with the result that porosity reduction is higher near the column inlet (Figure 8). For the pore-scale volume SubVolREAC-1, permeability in the flow direction, k_{zz} , decreases from $96 \cdot 10^{-12}$ to $60 \cdot 10^{-12} \text{ m}^2$ for a porosity reduction of 3.9% (from 36.2 to 32.3%). The variations are quite comparable for the sub-volume SubVolREAC-3, with permeability in the flow direction decreasing from $115 \cdot 10^{-12}$ to $62 \cdot 10^{-12} \text{ m}^2$ for a porosity reduction of 3.5% (from 36.3 to 32.8%). The average fluid velocity magnitude in the flow direction increases from 19.5 to 22.5 $\mu\text{m}\cdot\text{s}^{-1}$ for SubVolREAC-1, and from 14.3 to 16.3 $\mu\text{m}\cdot\text{s}^{-1}$ for SubVolREAC-3.

Changes in permeability are closely linked to the closing of pore throats associated with porosity reduction, but also to the increase in roughness at the surface of the aragonite and glass beads. Distribution of the velocity magnitude u along the main flow direction shows similar trends between SubVolREAC-1 and SubVolREAC-3. The distribution of the fluid velocity magnitude normalized to the Darcy's seepage velocity, \bar{u}/u_p , is presented in Figure 9 as log-log and semi-log plots. The magnitude of the principal velocity peak is centered on approximately $\bar{u}/u_p = 1.36$. Precipitation results in both REAC-1 and REAC-3 in: (1) an increase of the frequency of locations of the slowest velocities (underlined with an arrow in Figure 9a), (2), an increase of the highest velocities (underlined with an arrow in Figure 9b), and (3) a decrease of the magnitude of the principal velocity peak resulting from an increase of the standard deviation of the fluid velocity distribution (Table 5). However, for REAC-3 (i.e., aragonite and glass beads) the volume of the pore space within which the velocity decreases is far greater. Although fluid velocity does not seem to be disturbed in the vicinity of the smooth grains (i.e., calcite spar and glass beads), the velocity decrease in the vicinity of rough grains (i.e., aragonite) is more pronounced (Figure 9c,d).

Table 4. Porosity and permeability changes within the volumes of interest (VOI)

VOI	Vol-REAC-1		Vol-REAC-3		Subvol-REAC-1		Subvol-REAC-3	
Size (voxel)	1000 × 1000 × 1000		1000 × 1000 × 1000		300 × 300 × 300		350 × 350 × 350	
	initial	final	initial	final	initial	final	initial	final

ϕ (%)	35.9	34.6	35.9	34.8	36.3	32.3	36.3	32.8
k_{zz} ($\times 10^{-12}$ m ²)	198	170	-	-	112	69	137	77
k_D ($\times 10^{-12}$ m ²)	-	-	-	-	115	79	111	67
$\bar{u} \pm \sigma$ ($\mu\text{m}\cdot\text{s}^{-1}$)	-	-	-	-	19.5 \pm 14.6	22.3 \pm 16.2	14.3 \pm 10.3	16.3 \pm 11.9
u_P ($\mu\text{m}\cdot\text{s}^{-1}$)	-	-	-	-	11.5	12.8	11.6	13.0

Figure 8. Porosity changes along the flow direction during the experiments for the column-scale volumes of interest Vol-REAC-1 (calcite and glass) and Vol-REAC-3 (aragonite and glass).

Figure 9. Fluid velocity distribution in pore-scale volumes SubVol-REAC-1 (calcite and glass) and SubVol-REAC-3 (aragonite and glass) before and after experiment, respectively. (a) and (b) Histograms of normalized velocity magnitude in log-log and semi-log plots, respectively. (c) and (d) example of 2D cross-section of the fluid velocity along the principal flow direction.

4. Pore evolution model

Imaging of the 3D geometry of the samples has shown that the density, shape and size of the newly precipitated crystals depended strongly on the pre-existing substrate surface properties. Although the experimental chemistry and flow boundary conditions in samples REAC-1 and REAC-3 were similar, their porosity-permeability evolution is very different because (1) the initial geometry was different, and (2) the geometry modifications resulting from precipitation on calcite spar compared with aragonite ooids were quite different. To overcome the differences in the initial geometry of the columns (or samples), we tested for the influence of precipitation patterns on permeability evolution of pore-scale volumes that are directly comparable. To do that, a random growth model was used to generate different cases (or scenarios) of precipitation, from highly uniform to highly non-uniform, within the same starting porous

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4 422 medium. Then, permeability and flow changes were computed in the new generated volumes at two stages
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6 423 of porosity reduction.

9 424 **4.1. Description of the model and scenarios of precipitation**

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11 425 Volume SubVol-REAC-1-initial was used as initial porous medium for all the precipitation cases
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13 426 studied. Four phases were identified in this starting volume: the fluid phase, the contact surfaces between
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15 427 grains, and two different solid phases, labeled as solid 1 and solid 2. Solid 1 corresponds to the original
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17 428 glass beads and solid 2 corresponds to the calcite spar crystals.

18 429 The random growth model is used to nucleate and grow new crystals in the fluid phase from one or
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20 430 several fluid-solid interfaces, i.e., at the surface of solid 1 and/or solid 2. Homogeneous nucleation (i.e.,
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22 431 directly from solution) was not allowed (or considered).

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24 432 The new crystals are added through a number of discrete steps. First, a nucleation stage consists of
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26 433 attaching the center of gravity of nuclei to the solid surface. The nuclei, which are defined as cubes of 27
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28 434 voxels (i.e., $13.38 \times 13.38 \times 13.38 \mu\text{m}^3$), were attached onto the mineral surface following a uniform
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30 435 probability model, i.e., every voxel at the mineral surface has the same probability to host a nucleus. Note
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32 436 that the term nucleus is used to refer to a small crystal which has started to grow at the surface, and does
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34 437 not strictly correspond to the critical nucleus discussed in classical nucleation theory. Second, a growth
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36 438 stage follows the nucleation stage. Cubic crystals of 1, 8 or 27 voxels were added to the nuclei. They can
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38 439 be attached at the surface of the nuclei only to promote uniform precipitation, or added in a cumulative
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40 440 way to any crystals of the domain to promote non-uniform precipitation. It is also possible to set up the
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42 441 growth from both the nuclei and the original surface of one or several solids in order to induce a mix of
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44 442 growth and nucleation at the same time. The growth is repeated step by step until a desired fraction of the
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46 443 pore space has been filled.

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48 444 In this study, four different scenarios (Sc1 to Sc4) of precipitation were tested: (1) uniform growth at
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50 445 the surface of both solid 1 and solid 2, (2) no crystal growth at the surface of solid 1 and uniform growth
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52 446 at the surface of solid 2, (3) no crystal growth at the surface of solid 1 and non-uniform nucleation and
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54 447 growth at the surface of solid 2 using a low density of nucleus (filling 0.1% of the pore volume) at the
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56 448 surface of solid 2, and (4) no crystal growth at the surface of solid 1 and non-uniform nucleation and
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58 449 growth at the surface of solid 2, using a high density of nucleus at the surface of solid 2 (filling 0.5% of
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60 450 the pore volume). The different scenarios are illustrated in Figure 10. Scenario 1 (Sc1) is the simplest and
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62 451 corresponds to uniform precipitation over essentially the entire surface of the porous medium. This
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64 452 scenario could occur if the column were composed entirely of calcite spar. In this scenario, the nuclei
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66 453 cover 31% of the surface of solids 1 and 2 in contact with fluid. Growth of cubes of 1 voxel was imposed

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4 454 only at the surface of both the nuclei and the minerals to obtain a relatively smooth surface. Scenario 2 (Sc
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6 455 2) is close to what took place in REAC-1, with uniform growth of calcite onto spar crystals and no growth
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8 456 on glass beads, which are considered to be entirely non-reactive in this case (no nucleation is possible).
9 457 The nuclei cover 67% of the surface of solid 2 in contact with fluid. Successive growth of cubes of 8 and 1
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11 458 voxels was imposed only at the surface of both the nuclei and solid 2 to obtain a very smooth surface at
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13 459 the end. Scenario 3 (Sc 3) is close to what occurred in REAC-3, with complex, non-uniform growth of
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15 460 calcite crystals on aragonite beads, with the exception that glass beads are also considered in the
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17 461 simulations as totally non-reactive in contrast to what was observed during the experiment (some growth
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19 462 on glass beads did occur). The nuclei cover 4% of the surface of solid 2 in contact with fluid. Growth of
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21 463 cubes of 27 voxels is cumulative at the surface of the nuclei only. Scenario 4 (Sc 4) is similar to Sc 3, but
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23 464 with a higher density of nuclei at the surface of solid 2 (18%).

23 465 After the nucleation stage, two stages (stage 1 and 2) of growth were imposed in the pore evolution
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25 466 modeling. The porosity was reduced by approximately 3.7 % after stage 1 (corresponding to the average
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27 467 decrease of porosity in pore-scale volumes SubVol-REAC-1 and SubVol-REAC-3), and by approximately
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29 468 7.4 % after stage 2 so as to produce a porosity decrease two times that of the experiments.

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35 471 **Figure 10. Scenarios of precipitation Sc1 to Sc4, shown as: (a) cross-sections at stage 1**
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37 472 **(corresponding to a porosity reduction of ~3.7%), (b) cross-sections at stage 2 (corresponding to a**
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39 473 **porosity reduction of ~7.4%), and (c) 3D volumes at Stage 2. New precipitate is shown in white.**

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43 475 **4.2. Flow simulation results**

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45 476 For each scenario of precipitation (Sc1 to Sc4) generated by the pore evolution model, permeability
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47 477 and flow changes were computed as the porosity decreased due to calcite precipitation. The results are
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49 478 presented in Table 5 and can be compared with the results obtained for the samples REAC-1 and REAC-3.
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51 479 The changes in permeability are shown as a function of porosity (Figure 11a) and roughness changes
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53 480 (S/S_0) (Figure 11b). There is a cumulative effect of porosity decrease and roughness increase
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55 481 contributing to the decrease in permeability. Initially, permeability decreases as the porosity decreases. At
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57 482 the end of stage 1, the permeability reduction is between 72 % and 54 %, depending on the precipitation
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59 483 scenario. By comparison, the permeability changes of SubVolREAC-1 and SubVolREAC-3 are between
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61 484 these values (Figure 11a). Pore roughness (Figure 11b) is also a parameter contributing to the permeability

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4 485 decrease by increasing the pressure gradients at the fluid-mineral interface: the highest permeability
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6 486 reductions are for Sc3 and Sc4, for which the increases in pore roughness are highest.
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12 489 **Figure 11. (a) Permeability versus porosity (normalized to initial permeability k_{D0} and porosity**
13 **ϕ_0 , respectively) for the different pore-scale volumes SubVol-REAC-1 and SubVol-REAC-3 and**
14 490 **scenarios Sc1 to Sc4. (b) Permeability versus roughness changes.**
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20 493 Under constant flow rate boundary conditions, the consequence of mineral precipitation is an increase
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22 494 of both the average (\bar{u}) and maximum (u_{\max}) velocities within the porous media (Figure 12). Average
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24 495 fluid velocity and standard deviation of the distribution of velocities also depend on the surface roughness,
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26 496 and is greatest for scenarios for which the increase in roughness is the highest, i.e., Sc3 and Sc4 (Table 5).
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28 497 The fluid velocity distribution is broader when roughness is highest (Figure 13), which confirms the trend
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30 498 observed with the pore-scale volumes SubV-REAC-1 and SubV-REAC-3. More precisely, the principal
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32 499 differences are observed for the slowest ($u/u_p < 0.1$) and highest ($u/u_p \sim 10$) velocities. The flow field
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34 500 is also more affected in the vicinity of the fluid-rock interface for Sc3 and Sc4 (Figure 14), and becomes
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36 501 more heterogeneous with the development of areas with low flow velocities (i.e., “immobile zones”).
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38 502 Although the average velocity increases close to the interface due to the increases in roughness, a larger
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40 503 number of slow fluid flow volumes of velocity ~ 0 were noted.
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44 506 **Figure 12. (a) Average fluid velocity versus maximum velocity in the different sub-volumes**
45 **(SubVolREAC-1, SubVolREAC-3 and scenarios Sc1 to Sc4 at stage 1 and 2 of growth). (b) Average**
46 507 **fluid velocity in the vicinity of the fluid-rock interface, i.e., at a distance of 9 μm (i.e., 2 voxels)**
47 **maximum from the solid phase (for SubVolREAC-1, and the two scenarios with the lowest and**
48 508 **highest roughness S/S_0 , i.e., Sc2 and Sc4, respectively).**
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Figure 13. Normalized fluid velocity distribution for the different scenarios Sc1 to S4 at stage 2 (i.e., porosity reduction of ~7.4%), with comparison to the pore-scale volume SubV-REAC-1-initial. The histograms are provided in log-log (a) and (b) semi-log plots.

Figure 14. 2D cross-section of the flow velocities (a) perpendicular and (b) parallel to the flow main direction for the different scenarios Sc1 to Sc4, stage 1. White arrows underline areas where flow becomes more heterogeneous compared to Sc1.

Table 5. Permeability and flow properties for the different scenarios of precipitation. Porosity at stage 1 is about the average porosity of the two sub-volumes SubV-REAC-1-final and SubV-REAC-3-final.

Scenario	initial	Sc1		Sc2		Sc3		Sc4	
		stage 1	stage 2	stage 1	stage 2	stage 1	stage 2	stage 1	stage 2
$k_D (\times 10^{-12} \mu\text{m}^2)$	115*	73	55	82	63	61	38	64	41
$\bar{u} \pm \sigma (\mu\text{m}\cdot\text{s}^{-1})$	19.5 ±	22.2 ±	25.3 ±	22.2 ±	25.7 ±	22.8 ±	26.8 ±	22.7 ±	26.3 ±
	14.6	17.3	19.1	16.7	18.8	18.2	22.2	18.3	21.9
$u_P (\mu\text{m}\cdot\text{s}^{-1})$	11.5	12.9	14.5	12.9	14.5	12.8	14.3	12.8	14.3

* Identical initial permeability as SubVol-REAC-1 for every scenario.

5. Discussion

5.1. Calcite nucleation and growth

The heterogeneous statistical distributions, shapes, growth rates and growth velocities of the new crystals within the columns are best explained by a combination of heterogeneous nucleation and subsequent crystal growth. According to nucleation theory, the energy barrier for nucleation, Δg_n , depends on the

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4 534 degree of fluid supersaturation with respect to the nucleating mineral, $(1-\Omega)$, and on the interfacial free
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6 535 energy, α , e.g., [17]: $\Delta g_n \propto \alpha^3 / (1-\Omega)^2$. The lack of measurable precipitation with XMT at a distance of
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8 536 about 1.8 mm from the column inlet on glass beads and about 4.7 mm on aragonite, together with the fact
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10 537 that precipitation was not detected with SEM on glass beads and on many of the aragonite ooids in the
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12 538 distal portion of the columns (close to the exit) suggest that a high degree of supersaturation is required to
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14 539 lower the energy barrier for nucleation on the non-calcite-based material. In contrast, calcite spar crystals
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16 540 are covered uniformly by a thin layer of new precipitate. Lin and Singer [35] observed in short-term
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18 541 experiments that a saturation index ($\log \Omega$) of 0.72 was insufficient to allow nucleation of calcite on
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20 542 quartz or dolomite seed materials, although heterogeneous nucleation did take place on calcite. Liolou et
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22 543 al. [34] also observed no precipitation on quartz for solutions with a saturation index ($\log \Omega$) below 0.73.
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24 544 In our experiments, an evaluation using geochemical model JChess [59] based on an estimation of the
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26 545 solution composition that took into account the mass of calcite precipitated along the flow direction
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28 546 indicates a critical saturation index ($\log \Omega_{\text{crit}}$) for nucleation on glass beads of ~ 1.26 in the experiments. In
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30 547 contrast, a saturation index ($\log \Omega$) of 1.20 at the REAC-1 column outlet was sufficient to allow
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32 548 nucleation on calcite spar, but made it more difficult or impossible on aragonite ooids. It seems that the
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34 549 surface area that allowed initiation of calcite precipitation is correlated with the energy barrier for
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36 550 nucleation which might explain why only sparse calcite crystals are observed on glass beads, while 92%
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38 551 of aragonite surface of SubVol-REAC-3 and almost 100% of calcite spar surface of SubVol-REAC-1
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40 552 were covered by newly precipitated crystals.

37 553 Stockmann et al. [33] observed different induction times for calcite nucleation during mixed-flow
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39 554 precipitation experiments at 25°C in the presence of different solid substrates. For calcite and other silicate
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41 555 substrates, precipitation was measurable within half a day, while it took up to 5 days for precipitation to be
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43 556 detectable on basaltic glass. We have no indication of the time required for nucleation on glass beads in
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45 557 the columns, but the difference in size of neighboring crystals (e.g., Figure 5b) suggests that the induction
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47 558 time for nucleation may contribute to the heterogeneity of crystal size distributions.

48 559 Crystal shape is also highly dependent on the growth substrate. Indeed, the interfacial free energy is
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50 560 reduced when the atomic structure of the substrate surface closely matches a particular plane of the
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52 561 nucleating phase (minimizing the lattice strain) and the substrate presents a set of chemical functionalities
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54 562 that promotes strong bonding to the nucleus [17]. The degree of epitaxy, which can be compared to the
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56 563 degree of matching between the substrate and the overgrowth crystals, directly constrains crystal growth,
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58 564 e.g., [60]. Starting from isolated nuclei, calcite grows as randomly oriented rhombohedra on glass beads,
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60 565 whereas growth occurs layer-by-layer after two-dimensional surface nucleation on calcite spar [36]. The
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62 566 non-equivalent crystal lattice of aragonite (orthorhombic system) makes it difficult for calcite to grow on
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4 567 its surface compared to calcite spar (trigonal system). Consequently, three-dimensional island growth is
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6 568 observed. The polyhedral shape of crystals on aragonite is therefore quite surprising, but might be the
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8 569 result of a release of Sr ($1.1 \cdot 10^{-6}$ M on average at the column outlet) by the ooids (which Sr/Ca ratio is 1.4
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10 570 10^{-2}) during the experiment. This interpretation is perhaps supported by the observation that the shapes of
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12 571 crystals growing on glass beads are not affected. Sr is known to affect calcite surface morphology during
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14 572 growth by scalloping the obtuse and calcite acute step edges [61, 62], although no effect on step velocities
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16 573 has been measured at such low concentrations. However, a concentration gradient may exist at the surface
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18 574 of aragonite ooids, leading very locally to higher concentrations of Sr.

18 575 Precipitation rates as a function of crystal growth velocities and growth rates are also highly variable.
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20 576 Precipitation rate at the column-scale is highly dependent on the saturation index, with the highest rates
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22 577 occurring near the column inlets. Reactive transport modeling of REAC-1 [36] has shown that a good fit
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24 578 of experimental data was obtained using a nearly linear rate expression of eq (1), with $k = 1.79 \cdot 10^{-9}$
25
26 579 $\text{mol}\cdot\text{m}^{-2}\cdot\text{s}^{-1}$ and $n = 1.12$, suggesting that spiral growth is the main mechanism of precipitation. However, a
27
28 580 reasonable agreement was also obtained using a mixed-control reaction combining multi-source spiral
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30 581 growth with a higher order rate expression (as suggested by [63]).

30 582 For individual crystals, growth velocities and rates depend on saturation index and crystal orientation,
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32 583 shape and growth competition. The distance used to calculate the growth velocity is the maximal length
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34 584 normal to the substrate surface, so that the velocity depends largely on the crystal shape and orientation.
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36 585 For example, orientation of the longest crystal diagonal perpendicular to the substrate surface will result in
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38 586 highest velocities. In addition, the different facets of a crystal generally exhibit different growth rates, e.g.,
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40 587 [8, 64, 65]. Finally, growth competition, which is highly visible on aragonite surfaces (Figure 5d), also
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42 588 involves variation of growth rate over time [3, 66].

43 589 **5.2. Porosity, permeability and flow**

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45 590 These interpretations concerning calcite nucleation and growth indicate that non-uniform precipitates
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47 591 can develop readily at the pore-scale even if the reactions are kinetically-controlled. Development of
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49 592 heterogeneity is expected in rocks containing multiple minerals for which surface properties for nucleation
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51 593 and growth differ. The pores that experience the development of pore-scale heterogeneity as a result of
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53 594 precipitation exhibit a higher reduction in permeability compared to those that experienced uniform
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55 595 precipitation. Reduction in permeability is linked to the decrease of porosity [38] through the reduction in
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57 596 the pore-throat size. However, similar reductions in porosity can lead to differing reductions in
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59 597 permeability. These differences arise from the development of pore roughness that is linked to the location
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61 598 and size of the newly formed crystals. More importantly, roughness increase causes the decrease in
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599 average fluid velocity in the vicinity of the mineral surface (despite the fact that the maximum velocity is
600 increased). Although it has not been quantified in this study, this can have significant implications for the
601 hydrodynamics and the transport of reactants and products near the fluid rock interface. For example, [67]
602 explained the origin of early breakthrough and long tailing plume behavior in various carbonate rock
603 samples by the shape of the probability density function of normalized velocity, which exhibits different
604 characteristics in terms of magnitude of the main velocity peak and spreading between low and high
605 velocities. Transport in samples with a relative narrow spread of velocities is characterized by a small
606 peak (of probability of molecular displacement) representing stagnant portion of the pore space and a
607 dominant secondary peak of solute moving at approximately the average flow velocity, while transport in
608 samples with a wider velocity distribution is characterized by a significant peak of stagnant solute and an
609 elongated tail of moving fluid (see [67] for more details).

610 **6. Conclusion**

611 We performed a set of experiments to obtain detailed measurements of calcite precipitation in two
612 porous columns, one consisting of a mix of calcite spar and glass beads, another consisting of aragonite
613 ooids and glass beads. XMT has provided direct information about the porous geometry changes and
614 localization of the precipitates. Despite the similar chemical compositions of the inlet solution and flow
615 conditions, the shape, size and distribution of the new calcite crystals have shown to be very different
616 depending on the material substrate. The results indicate that both calcite nucleation and growth stages are
617 important in the development of non-uniform precipitation patterns. As a consequence, the pore geometry
618 evolution was very different in the two experiments and led to differing evolution of pore roughness. The
619 modeling analysis carried out here to supplement the XMT characterization highlighted the significance of
620 the pore-scale distribution of the new precipitates on the evolution of permeability. Permeability changes
621 are linked to the porosity reduction as well as to the changes of the pore surface roughness. Roughness
622 increase in the vicinity of grains increases the heterogeneity of the flow field and affects the transport of
623 reactants and products near the mineral surface.

624 The results also demonstrate that the spatial distribution of different minerals should be taken into
625 account for the calculation of permeability changes in porous rocks, as both localization of precipitation
626 and precipitation patterns depend on it. Further investigation is needed to explore systems that experience
627 higher porosity decreases, i.e., down to the point where the percolation threshold is reached, although it
628 can be difficult to obtain experimentally a constant decrease of the porosity along the porous media.
629 Increasing the number of XMT data sets during an experiment could also confirm that the growth rates
630 and velocities are not constant with space and time.

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631 This study does not consider the impact of hydrodynamics changes close to the fluid-mineral interface
632 on transport, and possible feedbacks between chemical reaction and transport, which could be significant.
633 One consequence of the growth of crystals is the generation of stagnant or low velocity zones in the
634 vicinity of rough crystals, while the reduction of pore throats implies an increase of the fluid velocity
635 through them. Consequently, changes of Damköhler and Péclet numbers with space and time [14] might
636 reduce or increase the probability of nucleating and growing new crystals. Such an analysis will be only
637 possible with a full pore-scale transient reactive transport model.

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Figure 1

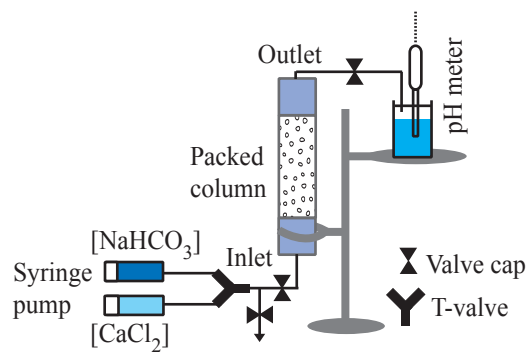


Figure 1. Schematic of the plug-flow column experimental setup

Figure2

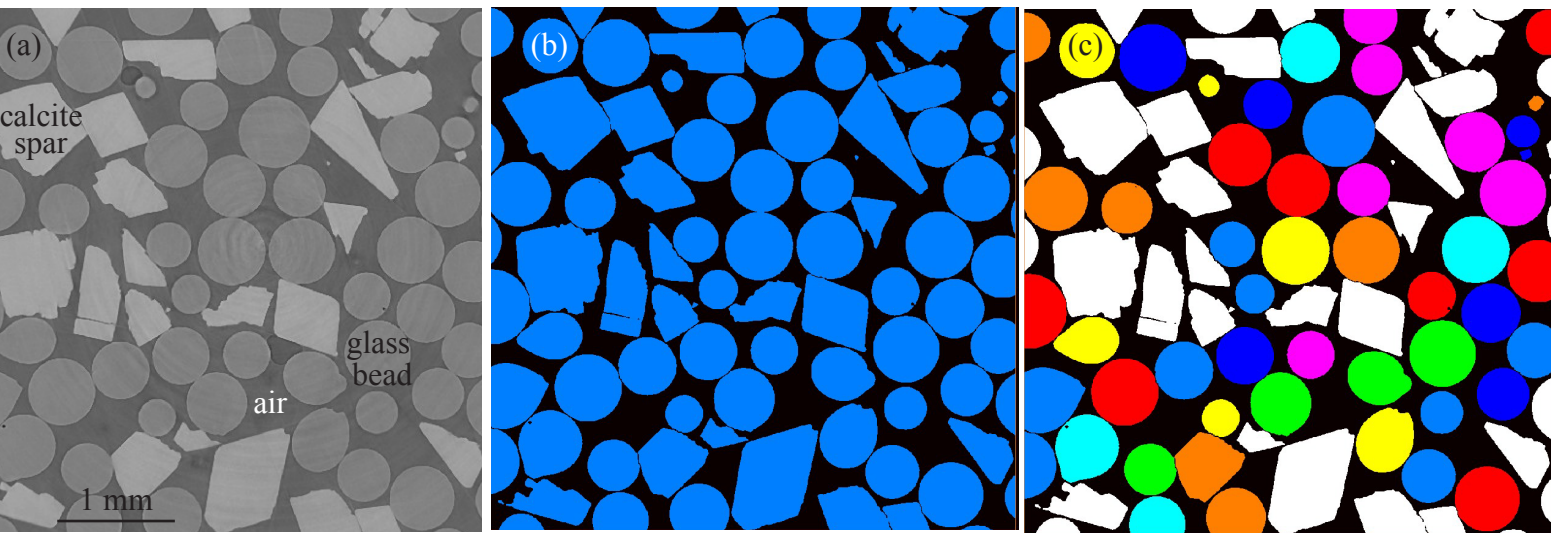
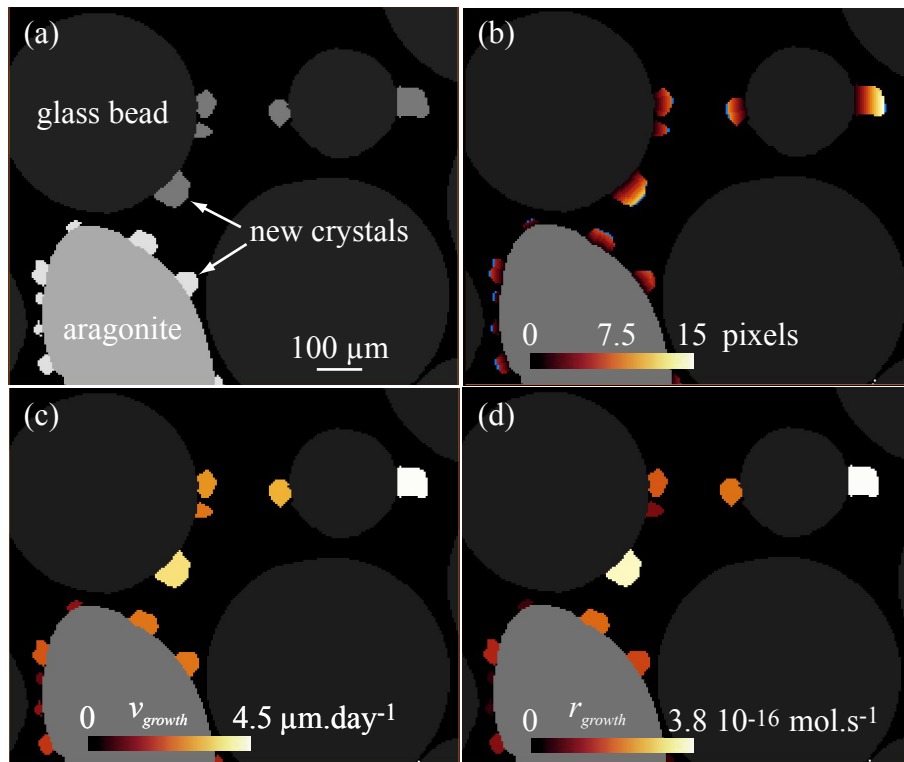


Figure3



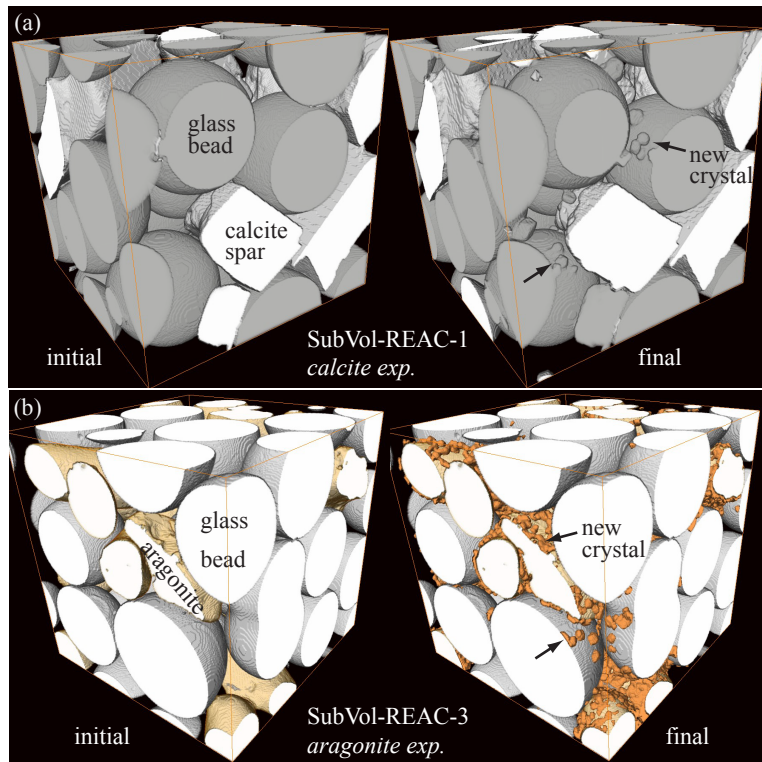


Figure 5

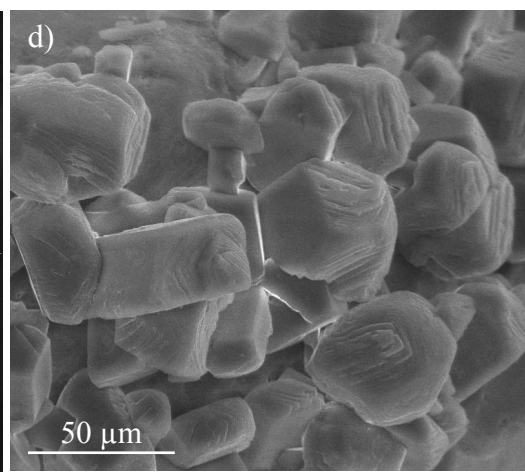
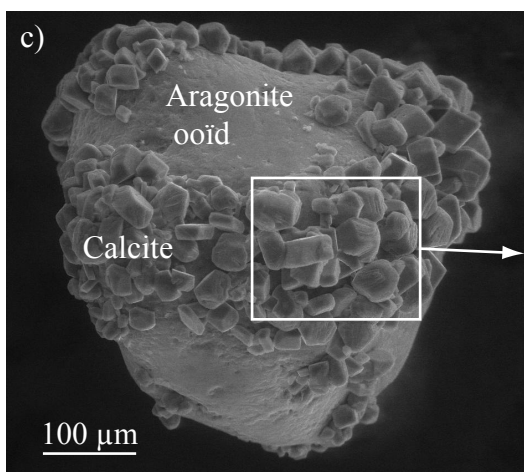
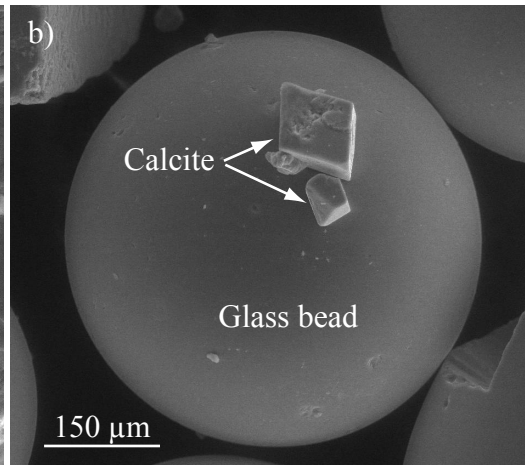
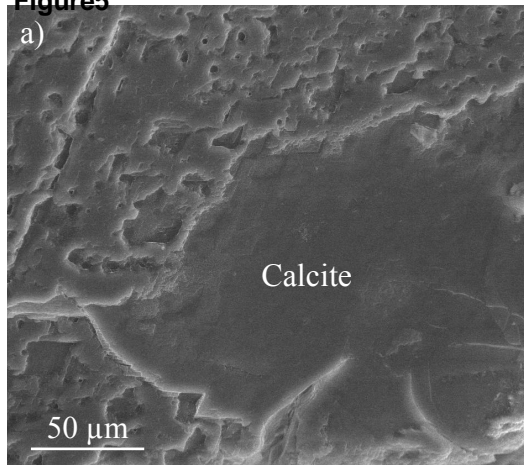


Figure6

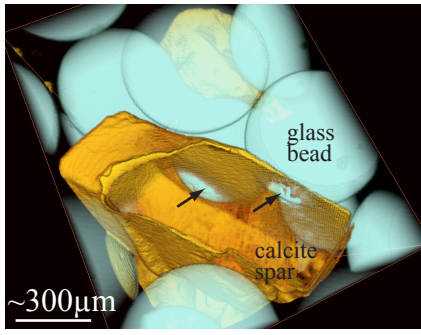


Figure 7

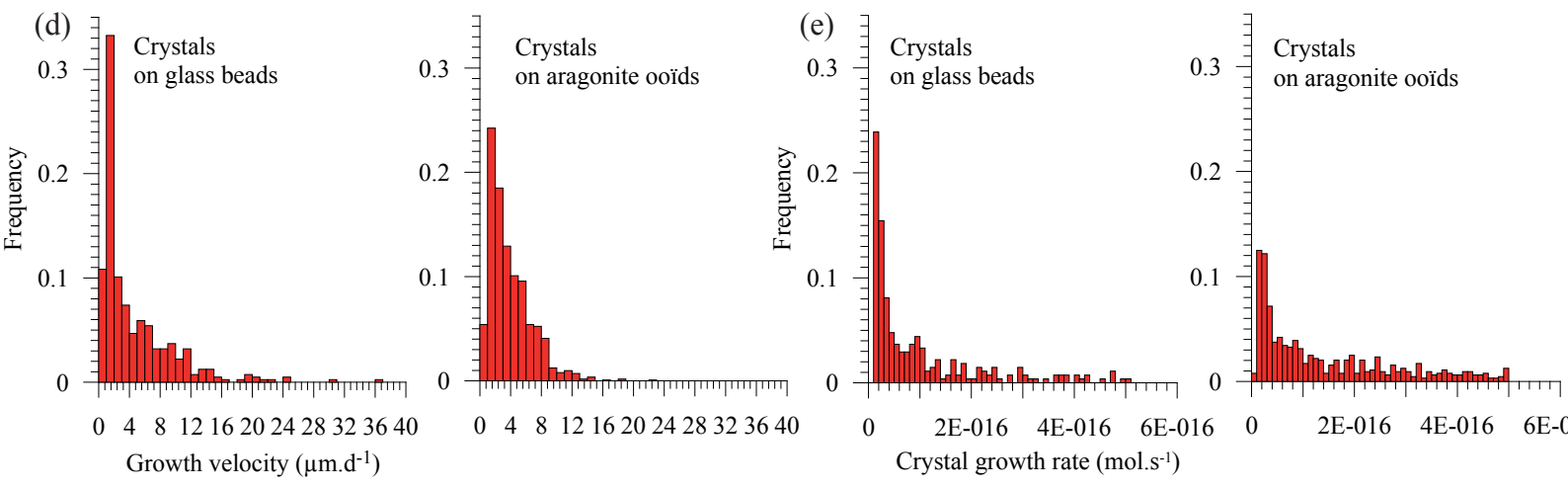
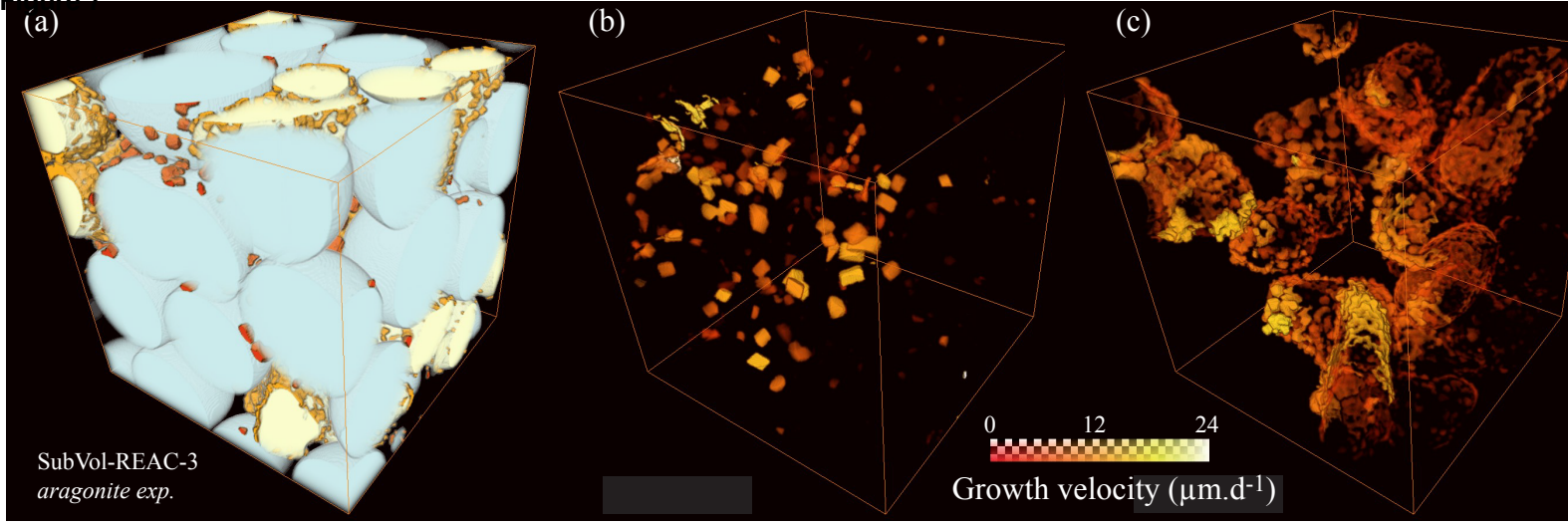


Figure 8

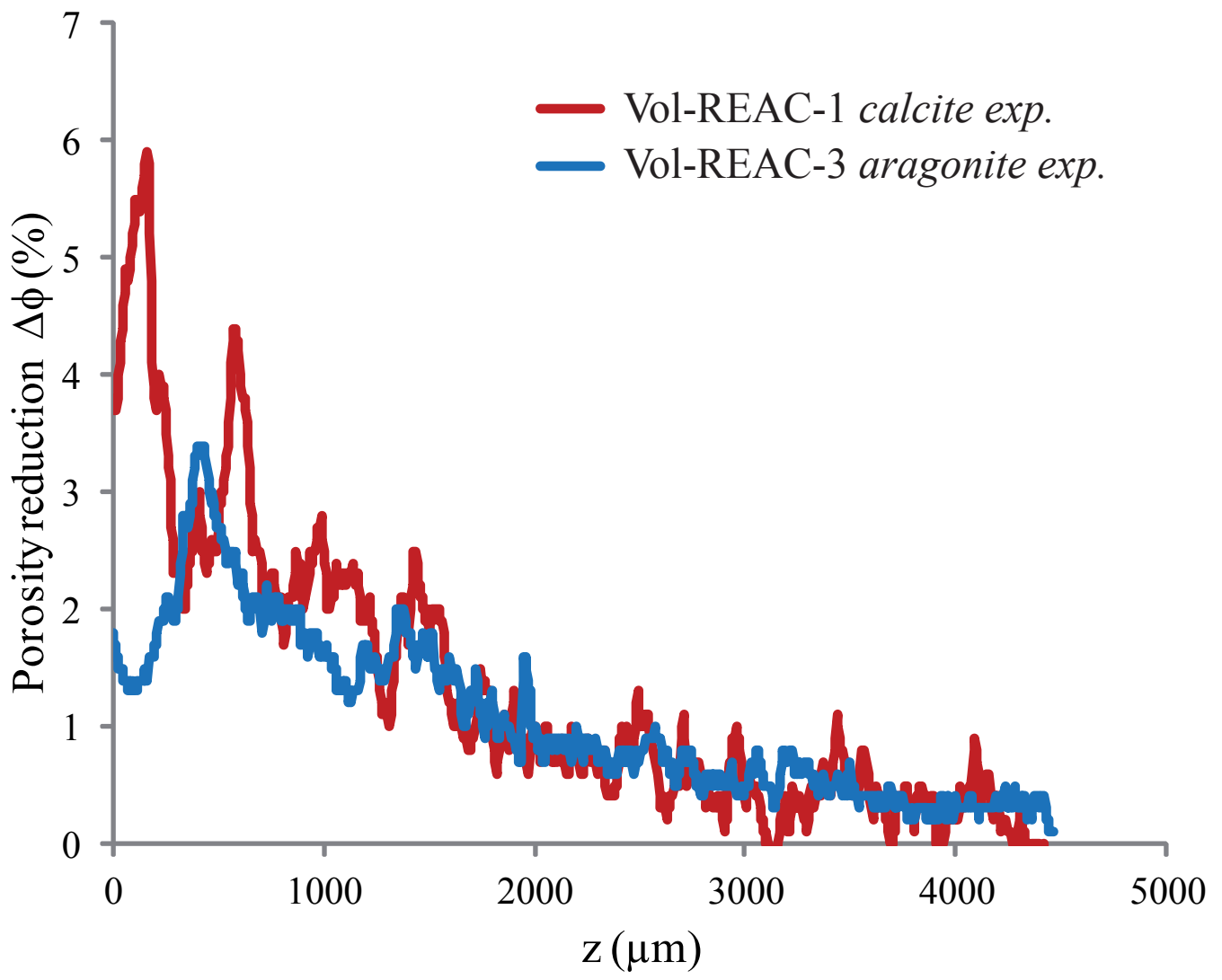


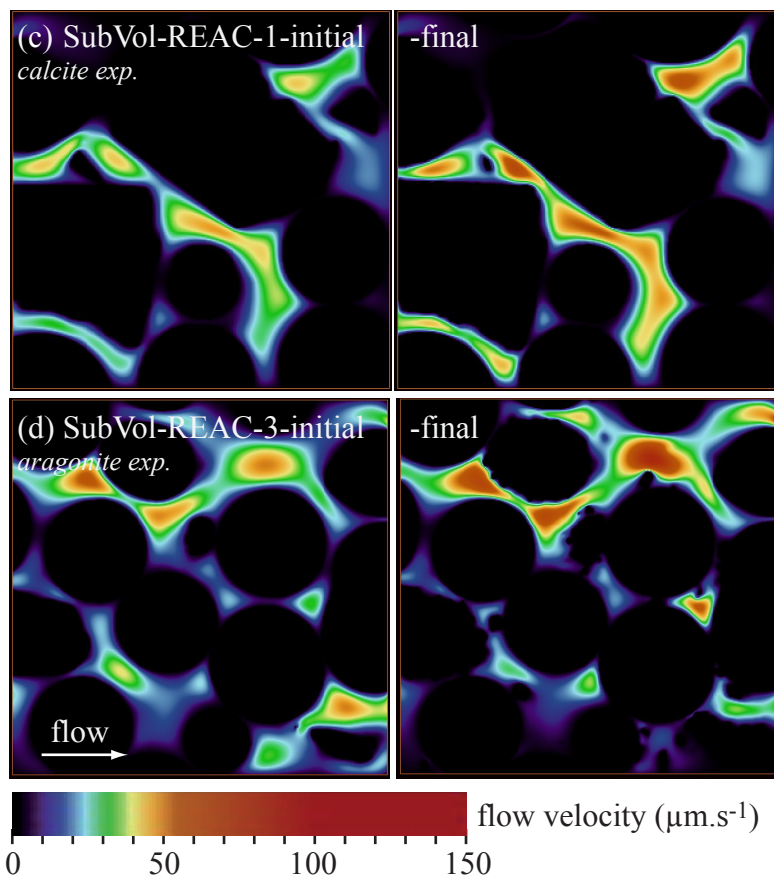
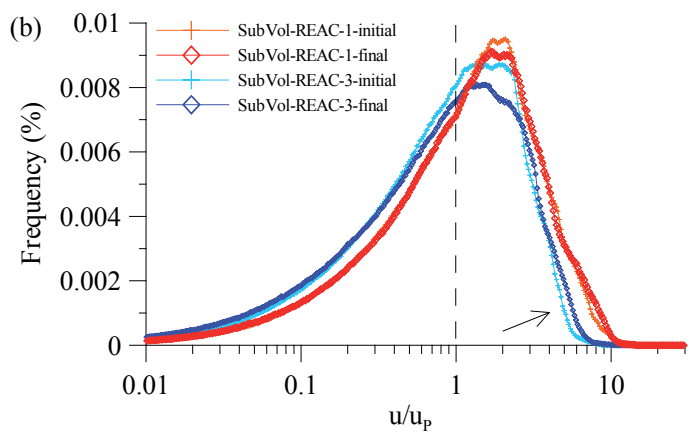
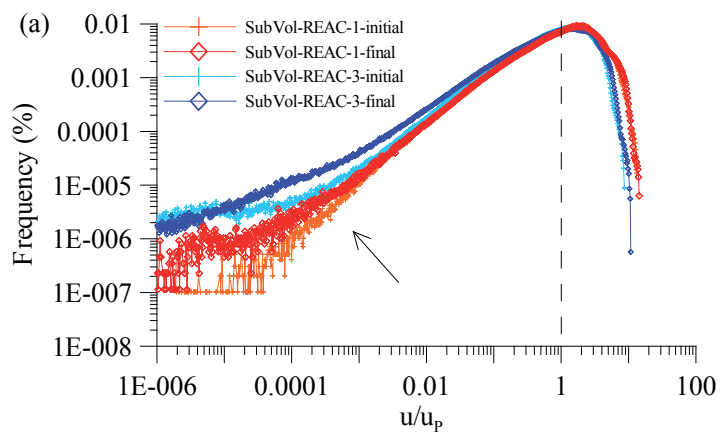
Figure 9

Figure10

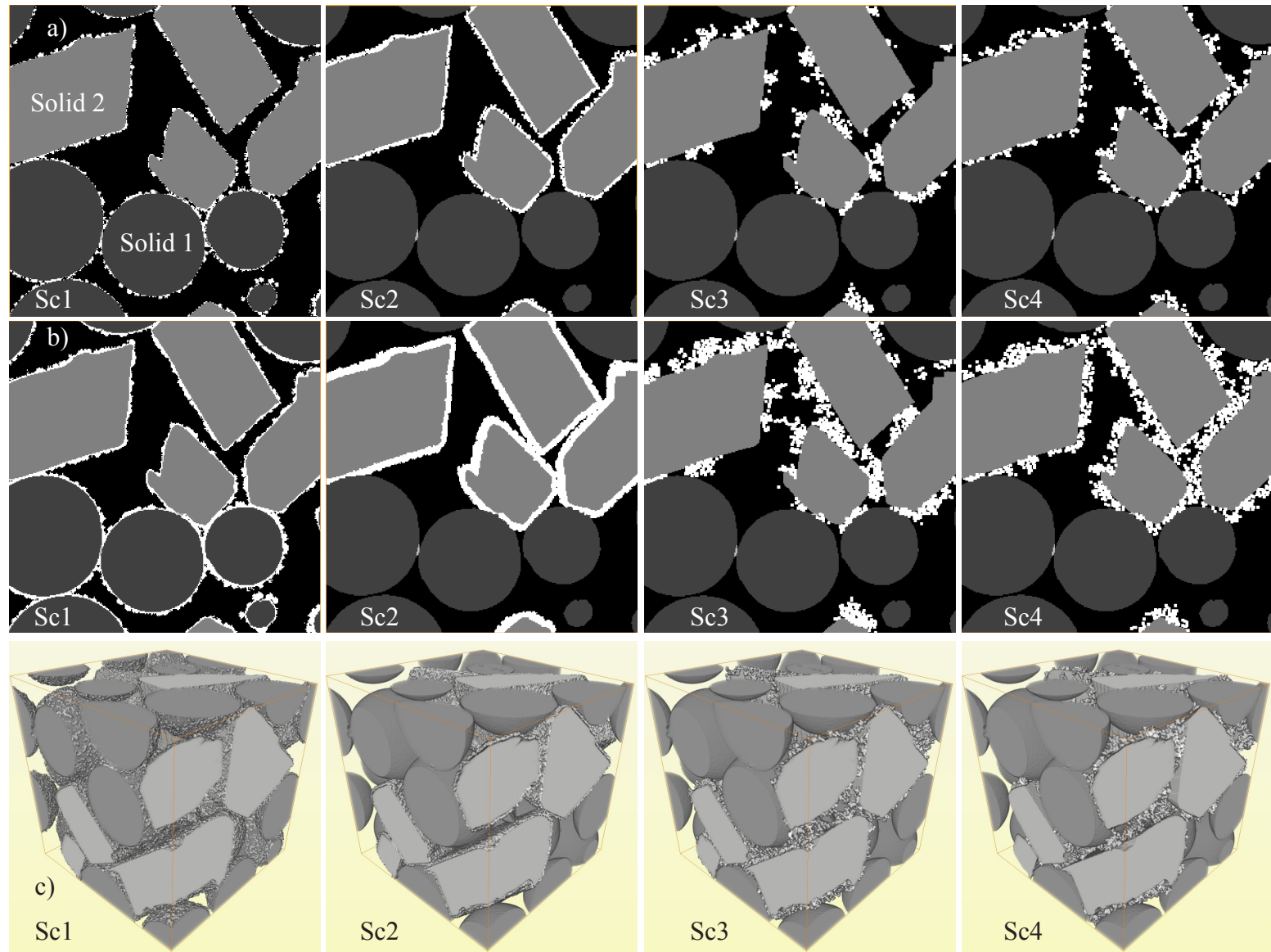


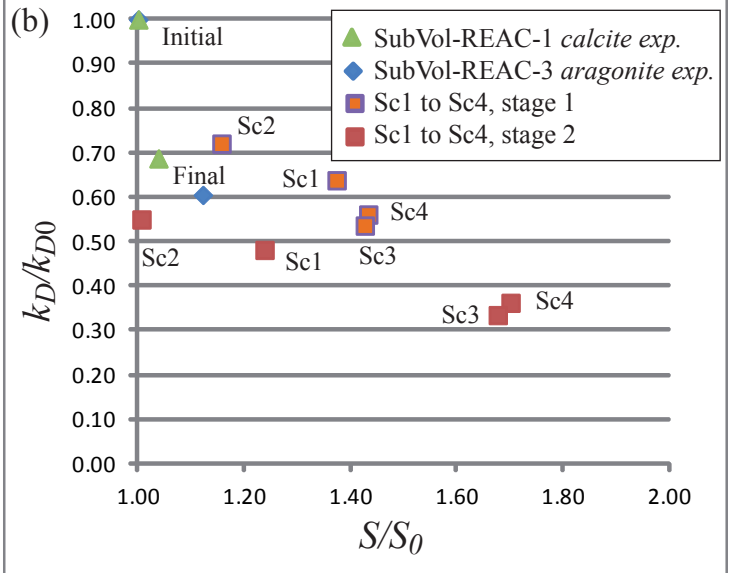
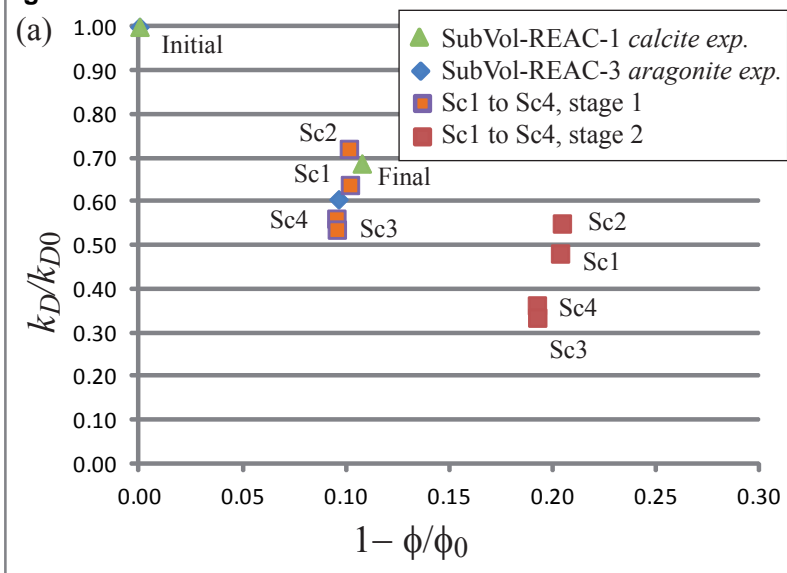
Figure 11

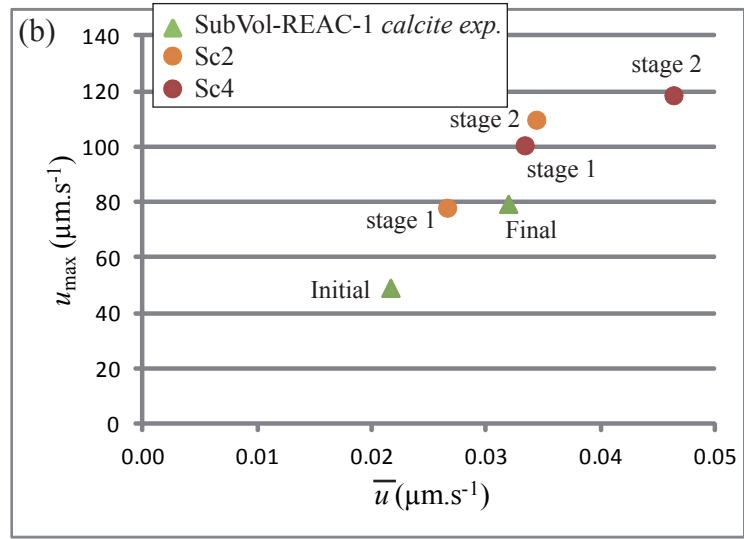
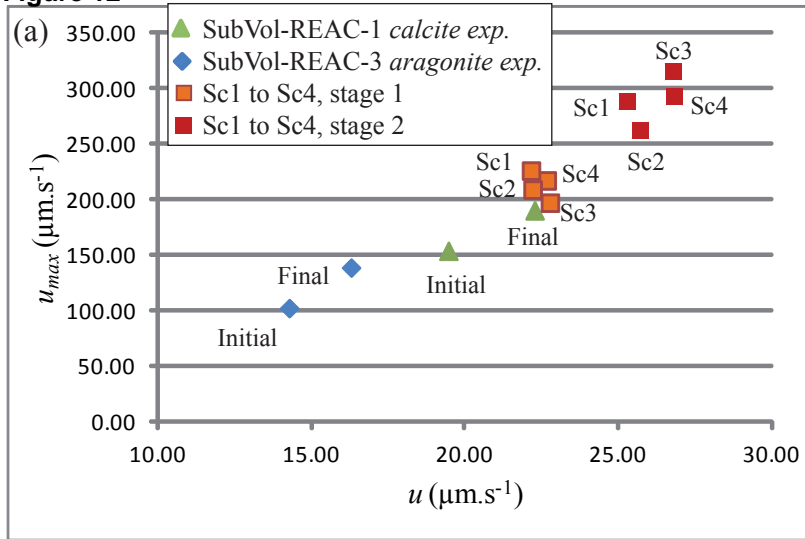
Figure 12

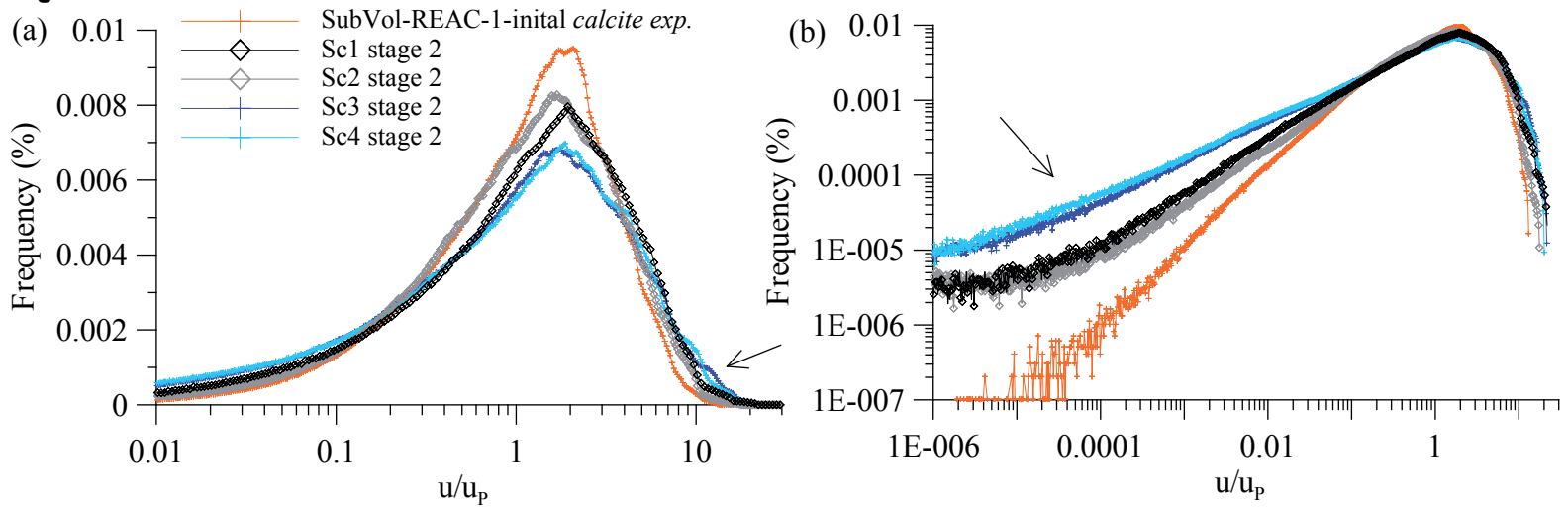
Figure 13

Figure14

