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On the long-cycle variability of the Algol OGLE-LMC-DPV-065 and its stellar, orbital, and disc parameters

R. E. Mennickent,^{1★} M. Cabezas,^{1,2} G. Djurašević,^{3,4} T. Rivinius,⁵ P. Hadrava,²
 R. Poleski,^{6,7} I. Soszyński,⁷ L. Celedón,¹ N. Astudillo-Defru^{1,8},^{1,8} A. Raj,⁹
 J. G. Fernández-Trincado,^{1,10,11} L. Schmidtobreick^{1,5},^{1,5} C. Tappert,¹² V. Neustroev^{1,13}
 and I. Porritt¹⁴

¹Universidad de Concepción, Departamento de Astronomía, Casilla 160-C, Concepción, Chile

²Astronomical Institute of the Academy of Sciences of the Czech Republic, Boční II 1401/1, Prague, 141 00, Czech Republic

³Astronomical Observatory, Volgina 7, 11060 Belgrade 38, Serbia

⁴Isaac Newton Institute of Chile, Yugoslavia Branch

⁵ESO - European Organization for Astronomical Research in the Southern Hemisphere, Casilla 19001, Santiago 19, Chile

⁶Department of Astronomy, Ohio State University, 140 W. 18th Ave., Columbus, OH 43210, USA

⁷Warsaw University Observatory, Al. Ujazdowskie 4, PL-00-478 Warszawa, Poland

⁸Departamento de Matemática y Física Aplicadas, Universidad Católica de la Santísima Concepción, Alonso de Rivera 2850, Concepción, Chile

⁹Indian Institute of Astrophysics, II Block Koramangala, Bangalore 560034, India

¹⁰Institut Utinam, CNRS UMR6213, Univ. Bourgogne Franche-Comté, OSU THETA, Observatoire de Besançon, BP 1615, F-25010 Besançon Cedex, France

¹¹Instituto de Astronomía y Ciencias Planetarias, Universidad de Atacama, Copayapu 485, Copiapó, Chile

¹²Instituto de Física y Astronomía, Universidad de Valparaíso, Valparaíso, Chile

¹³Astronomy Research Unit, PO Box 3000, FL-90014 University of Oulu, Oulu, Finland

¹⁴Turitea Observatory, Palmerston North, New Zealand

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ABSTRACT

OGLE-LMC-DPV-065 is an interacting binary whose double-hump long photometric cycle remains hitherto unexplained. We analyze photometric time series available in archive data sets spanning 124 yr and present the analysis of new high-resolution spectra. A refined orbital period is found of $10^d 031\,6267 \pm 0^d 000\,0056$ without any evidence of variability. In spite of this constancy, small but significant changes in timings of the secondary eclipse are detected. We show that the long period continuously decreases from 350 to 218 d during 13 yr, then remains almost constant for about 10 yr. Our study of radial velocities indicates a circular orbit for the binary and yields a mass ratio of 0.203 ± 0.001 . From the analysis of the orbital light curve, we find that the system contains 13.8 and 2.81 M_{\odot} stars of radii 8.8 and 12.6 R_{\odot} and absolute bolometric magnitudes -6.4 and -3.0 , respectively. The orbit semimajor axis is 49.9 R_{\odot} and the stellar temperatures are 25 460 K and 9825 K. We find evidence for an optically and geometrically thick disc around the hotter star. According to our model, the disc has a radius of 25 R_{\odot} , central and outer vertical thickness of 1.6 R_{\odot} and 3.5 R_{\odot} , and temperature of 9380 K at its outer edge. Two shock regions located at roughly opposite parts of the outer disc rim can explain the light-curve asymmetries. The system is a member of the double periodic variables and its relatively high-mass and long photometric cycle make it similar in some aspects to β Lyrae.

Key words: stars: binaries: eclipsing, close, spectroscopic – stars: activity, circumstellar matter – fundamental parameters.

1 INTRODUCTION

Stellar magnetic dynamos are relatively common in solar-type stars, and magnetic activity in binaries containing GK dwarfs is well documented in the RSCVn systems (Hall 1989). The situation in

* E-mail: rmennick@astro-udec.cl

Algol-type variables is less clear. Algols are semidetached binaries with intermediate-mass components, where the less-massive star, dubbed secondary or donor, is more evolved than the more massive star, named gainer or primary. Some authors have proposed that orbital period changes observed in some close binaries might be driven by magnetic cycles through the Applegate (1992) mechanism; the angular momentum of the star and the binary is redistributed during the magnetic cycle producing the observed orbital period changes (Lanza, Rodonò & Rosner 1998; Lanza & Rodonò 1999). Further studies indicate that the presence of a dynamo may modulate the mass transfer rate in Algol systems, leading to a characteristic impact of the dynamo cycle on the system luminosity (Bolton 1989; Meintjes 2004). In this context the existence of a group of hot Algols showing a long photometric cycle lasting on average about 33 times the orbital period might be relevant, since this variability has been recently interpreted in terms of a magnetic dynamo (Schleicher & Mennickent 2017). If this hypothesis turns to be correct, one may deduce that the stellar dynamo is also active in the hot, rapidly rotating (orbitally synchronized) A-type giants in some semidetached Algols. In fact, for the Algol binary V393 Sco indirect evidence for magnetism in the secondary star has been deduced from the presence of chromospheric emission lines (Mennickent, Schleicher & San Martín-Perez 2018). These authors note that the spin-up of the donor during mass-transfer stage increases its dynamo number, likely enhancing the probability of occurrence of a magnetic dynamo at the semidetached stage.

The aforementioned group of hot Algols showing long photometric cycles in addition to their orbital variability is named double periodic variables (DPVs; Mennickent et al. 2003; Poleski et al. 2010; Pawlak et al. 2013; Mennickent, Otero & Kołaczowski 2016; Mennickent 2017). DPVs are semidetached binaries typically consisting of an A/F/G giant star filling its Roche lobe and transferring mass onto a B-type primary surrounded by a circumprimary disc. Among Galactic DPVs, one famous example is β Lyrae (Guinan 1989; Harmanec et al. 1996; Harmanec 2002).

Few extragalactic DPVs have been studied at some detail. Among them, the case of OGLE-LMC-DPV-065 (OGLE05200407-6936391; $RA_{2000} = 05^h20^m04^s.07$, $Dec_{2000} = -69^{\circ}36'39''.1$) is notable, since it is one of the brightest DPVs in the LMC ($V = 14.74$, $B - V = -0.07$), and shows a remarkable change in the long period from 350 to 210 d in 15 yr that clearly stands out among the rest of the DPVs. In addition, the system is eclipsing, with a 1.4 mag deep main eclipse and a comparatively large amplitude of the long cycle of 0.3 mag, in the I band. To date, there is no other DPV with such a remarkable change in the long cycle. The orbital period has been reported as $P_o = 10^d031\,645 \pm 0.000\,033$ (Poleski et al. 2010).

The above credentials make OGLE-LMC-DPV-065 an ideal target for a deeper study. If the long cycle is related to changes in the mass transfer driven by a magnetic dynamo, it might show up in spectroscopic and photometric signatures during the long cycle. For this reason we conducted a long-term spectroscopic monitoring of this target with UVES at the VLT (Section 2.3). For the sake of order and clarity we have divided our work into two parts. In this first paper we analyze the available photometric time series making use of archive data, present our new high-resolution spectroscopic observations, calculate the system and orbital parameters, and provide a solution for the stellar radius, mass, luminosity, and surface gravity along with a characterization of the accretion disc. In a second forthcoming paper, we will provide an analysis of the spectroscopic changes during the long cycle and a study of the evolutionary stage of the binary. We notice that a short

Table 1. Summary of photometric observations considered in this paper. The number of measurements, starting and ending times for the series, average magnitude, and standard deviation (in magnitudes) are given. Single-point uncertainties in the I band and V band for OGLE data are between 4 and 6 mmag. The zero point for HJD is 2 400 000. See comment on the average magnitudes in the text.

Source	N	HJD _{start}	HJD _{end}	Mag	Std.	Band
DASCH	460	12 697.8482	34 399.4995	14.996	0.219	B
OGLE-II	915	50 455.6744	51 873.7744	14.898	0.218	I
OGLE-III	504	52 123.9345	54 953.5268	14.907	0.246	I
OGLE-IV	73	55 326.4931	57 710.7482	14.901	0.302	I
CTIO	97	56 964.7927	57 327.7354	14.901	0.317	I
OGLE-II	95	50 467.7237	51 631.5633	14.908	0.244	V
OGLE-III	90	52 990.6851	54 948.4703	14.929	0.245	V
Turitea	106	56 342.9193	56 467.8381	14.918	0.368	y
ASAS-SN	664	56 789.4535	57 974.8870	14.918	0.116	V
CTIO	95	56 964.7954	57 327.7381	14.918	0.330	V

and preliminary spectroscopic study of this object based on the data presented in this paper has been presented in a recent conference (Cabezas et al. 2019).

2 OBSERVATIONS AND METHODS

2.1 Photometric observations

We included OGLE-II (Szymanski 2005)¹ and OGLE-III/IV data.² The OGLE-IV project is described by Udalski, Szymański & Szymański (2015). Poleski et al. (2010) published the OGLE-II and OGLE-III I -band³ and V -band⁴ data of this star. We also considered 460 B magnitudes from the Digitalized Harvard plates (DASCH project)⁵ covering 59.4 yr, since 1893 August to 1953 January. In addition, we obtained new photometry with the CTIO 1.3m telescope operated by the SMARTS consortium in service mode between 2014 November and 2015 October, with the ANDICAM camera and filters V and I . Another data set was collected by Ian Porritt in Turitea Observatory, New Zealand, with the 0.36-m Meade telescope and a yellow filter. These new data were reduced in the usual way, removing bias and performing flat field corrections in the images and calculating differential magnitudes with respect to constant comparison stars. Finally, 664 V -band magnitudes were included from the ASAS-SN catalogue. The photometric observations analyzed in this paper amount to 3099 data points, cover 124 yr, and are summarized in Table 1.

2.2 Light-curve disentangling

We separated the light curve into long- and short-period components. For that we used an algorithm especially designed to disentangle multiperiodic light curves through the analysis of their Fourier component amplitudes. The method is described in Mennickent et al. (2012) and a short summary is given here. The

¹<http://ogledb.astrow.edu.pl/~ogle/photdb/>

²OGLE-III/IV data kindly provided by the OGLE team.

³<ftp://ftp.astrow.edu.pl/ogle/ogle3/OIII-CVS/lmc/dpv/phot/I/OGLE-LMC-DPV-065.dat>

⁴<ftp://ftp.astrow.edu.pl/ogle/ogle3/OIII-CVS/lmc/dpv/phot/V/OGLE-LMC-DPV-065.dat>

⁵<http://dasch.rc.fas.harvard.edu/project.php>

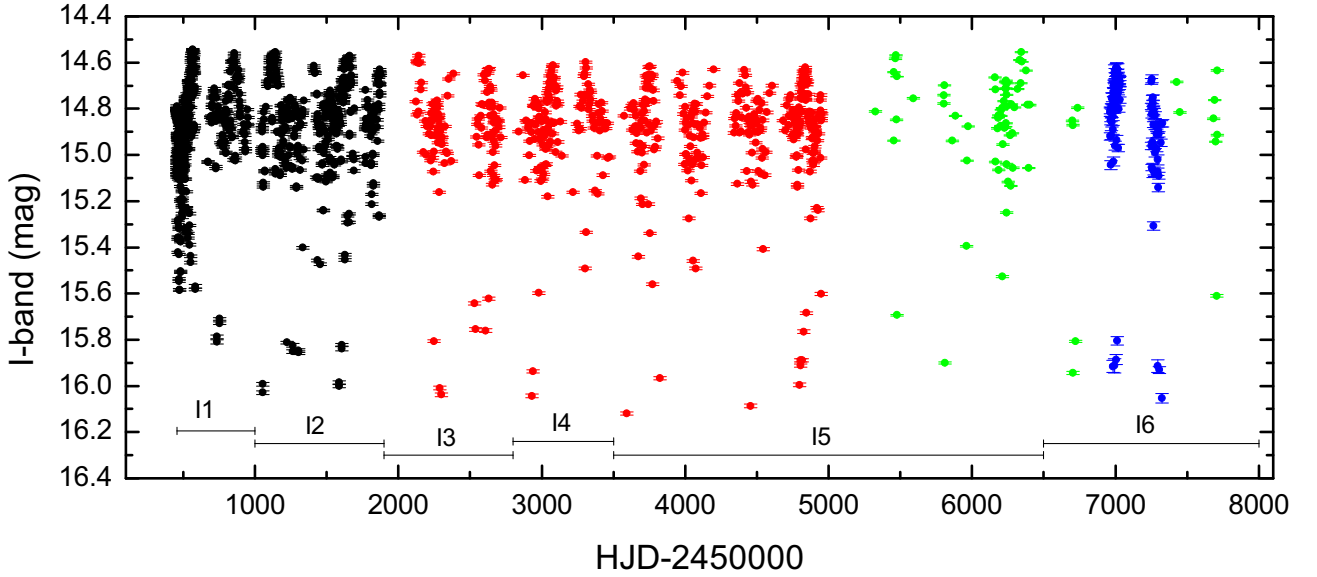


Figure 1. Light curve for OGLE II (black), III (red), and IV (green) data sets along with CTIO (blue) data. The upper envelope is produced by the long-period variability. The data points fainter than about 15.2 mag are taken during primary eclipse. The data intervals given in Table 4 are also shown.

Table 2. Summary of the search for the orbital period P_o .

P_o	Method	note
$10^d031\ 6450 \pm 0^d000\ 0330$	-	Poleski et al. (2010)
$10^d031\ 7800 \pm 0^d000\ 0780$	Eclipse timings	<i>I</i> band
$10^d031\ 6259 \pm 0.000\ 0230$	Period04	<i>I</i> band
$10^d031\ 6173 \pm 0.000\ 0128$	Period04	<i>I</i> residuals
$10^d031\ 6257 \pm 0.000\ 0066$	Period04	<i>B</i> and <i>I</i> bands
$10^d031\ 6267 \pm 0^d000\ 0056$		Weighted mean

main frequency f_1 is found with a period searching algorithm; this is usually the orbital frequency. A least-square fit is then applied to the light curve with a fitting function consisting of a sum of sine functions representing the main frequency and their additional significant harmonics. Afterwards the residuals are inspected for a new periodicity f_2 . This new periodicity (the long cycle in the case of DPVs) and their harmonics are included in the new fitting procedure. Finally, we obtain the light curve represented by a sum of Fourier components of frequency f_1 and f_2 and their harmonics. Data residuals with respect to the second and first theoretical light curves are the photometric series representing the orbital and long cycle, respectively.

2.3 Spectroscopic observations

We were granted 25 h for spectroscopic observations of the target with the ESO Ultraviolet and Visual Echelle Spectrograph UVES at the Kueyen telescope in the Paranal Observatory in service mode. This is a cross-dispersed Echelle spectrograph designed to operate with high efficiency from the atmospheric cut-off at 300 nm to the long-wavelength limit of the CCD detectors (about 1100 nm). The light beam is split into two arms, UV-Blue and Visual-Red, within the instrument. The two arms can be operated separately or in parallel with a dichroic beam splitter. The instrument provides

accurate calibration of the wavelength scale down to an accuracy of at least 50 m s^{-1} .

With the aim of covering both the orbital and long cycle 27 spectra were secured between 2013 October 1 and 2015 February 1, with the dichroic#2 mode in the ranges 3760–4985, 5700–7520, and 7665–9445 Å. The slit widths of 0'9 at the blue channel and 0'8 at the red channels provided a resolving power of 50 000 and 55 000, respectively. The object was observed at typical airmass 1.4 and with 3000 s exposure time per single exposure. A typical S/N of 65 was achieved at 480 nm. A summary of the spectroscopic observations is given in Table 7.

3 DATA ANALYSIS

3.1 The orbital period

Zero points can be an issue when combining photometry obtained at different sites with different detectors, filters and sky conditions. We have shifted OGLE IV, ASAS-SN, Turitea, DASCH, and CTIO magnitudes to the OGLE-II and OGLE-III averages before performing the analysis described in this paper. In Table 1 we provide only the original magnitude average for OGLE-II and OGLE-III data and also for the DASCH *B*-band magnitude. The light curve in the *I*-band is shown in Fig. 1.

We conducted a search for the orbital period using standard methodologies: eclipse times were measured interactively in the light curve with the computer cursor and a straight line fit was performed with the measured (epoch, time) pairs; the resulting slope gave the orbital period and their error. We also used the `period04` program, which calculates errors based on a Monte Carlo technique (Lenz & Breger 2005). The *I*-band residuals were obtained after removing the long-term cycle (see the next section). The periodicities found in different data sets are given in Table 2. We can see that the data are consistent with a constant orbital period; we find the following ephemerides for the main eclipse:

$$HJD_{\min} = 2\ 450\ 453.025 + (10^d031\ 6267 \pm 0^d000\ 0056) E. \quad (1)$$

Table 3. Observed and calculated times of long-cycle maxima. N is the cycle number. HJD are referred to 2 450 000. O stands for HJD_O and C for HJD_C . Errors are given when available. We have used the linear ephemerides with zero point 2617.89 and period 240 d. An estimate of the local period P_l is given subtracting the observed timing from the previous one and dividing by the number of elapsed cycles.

N	HJD_O	HJD_C	$[O - C]$ (d)	P_l (d)	band/source
-8	571.68 ± 4.49	697.89	-126.21	-	I OGLE
-7	859.75 ± 7.05	937.89	-78.14	288.06	I OGLE
-6	1125.58 ± 9.48	1177.89	-52.31	265.83	I OGLE
-4	1645.23 ± 8.04	1657.89	-12.66	259.83	I OGLE
-2	2131.23 ± 10.49	2137.89	-6.66	243.00	I OGLE
0	2617.89 ± 15.47	2617.89	0.00	243.33	I OGLE
2	3077.87 ± 7.04	3097.89	-20.02	229.99	I OGLE
3	3311.28 ± 4.98	3337.89	-26.61	233.41	I OGLE
5	3747.12 ± 7.00	3817.89	-70.77	217.92	V CTIO
5	3749.41 ± 6.93	3817.89	-68.48	219.06	I OGLE
6	3973.60 ± 2.52	4057.89	-84.29	224.19	I OGLE
8	4403.01 ± 10.98	4537.89	-134.88	214.71	I OGLE
8	4408.67 ± 7.00	4537.89	-129.22	217.54	V CTIO
10	4835.37 ± 9.55	5017.89	-182.52	213.35	I OGLE
13	5477.80	5737.89	-260.09	214.14	I OGLE
17	6343.85	6697.89	-354.04	216.51	y Turitea
17	6347.25	6697.89	-350.64	217.36	I OGLE
17	6353.84 ± 8.00	6697.89	-344.05	219.01	I CTIO
20	7011.15 ± 4.00	7417.89	-406.74	219.10	I OGLE
20	7011.20 ± 2.00	7417.89	-406.69	219.12	I CTIO
20	7011.23 ± 10.00	7417.89	-406.66	219.13	V ASAS-SN
22	7448.75	7897.89	-449.14	218.76	I OGLE
22	7453.28 ± 10.00	7897.89	-444.61	221.03	V ASAS-SN

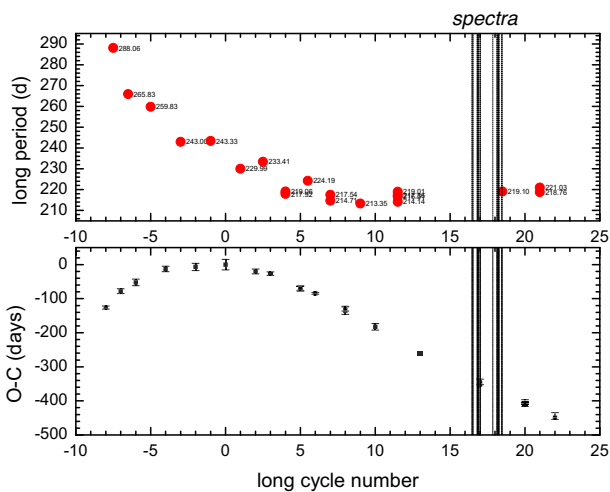


Figure 2. (up): Long period versus long cycle number. (down): Observed minus calculated times of the long cycle maximum versus long cycle number. Data are based on a test period of 240 d and are given in Table 3. Vertical dashed lines indicate spectroscopic observation times.

Table 4. Data intervals used for the long-period analysis. HJD are referred to 2 450 000. n is the number of I -band data points. Times for long-cycle maxima are given.

label	n	HJD range	P_l (d)	$T_0(max)$
I1	538	455 : 1000	282	855.66
I2	377	1000 : 1900	258	1128.82
I3	109	1900 : 2800	240	2618.78
I4	125	2800 : 3500	231	3316.75
I5	330	3500 : 6500	216	3753.70
I6	104	6500 : 7500	219	7009.62

3.2 The orbital and long cycle light curves

Maxima of the long cycle were measured directly from the light curves and compared with the ephemerides for a 240-d test period, as reported in Table 3. The observed minus calculated ($O - C$) diagram, constructed with the observed times of maxima (HJD_O) and the predicted ones (HJD_C), shows that the long cycle decreased at the beginning of the observations and then remained more or less constant during about 14 cycles (Fig. 2). We notice that considering the MACHO data analyzed by Mennickent et al. (2005, HJD: 2448900-2451500, not included in this paper), which is previous to the OGLE data reported here, the long period has decreased from about 350 to 218 d continuously during about 13 yr, before entering in a phase of almost constant period, that lasted for slightly more than 10 yr.

At every epoch we defined a local long cycle period P_l , subtracting the observed maximum timing from the previous one and dividing by the number of elapsed times, as given in Table 3. After inspection of Fig. 2 we choose six data intervals characterized by a more or less constant long cycle and a large number of observations (Table 4). This procedure allowed us to apply the disentangling to every data block considering the variability of the long cycle. The resulting disentangled light curves are shown in Fig. 3; they reveal that the long cycle is double-humped and that its shape remains relatively constant. In addition, the orbital light curve shows a small but significant variability (Figs 4 and 5): (i) on the 5th interval between HJD 2455 804 and 2456 405 the system is brighter at quadratures, and produces larger scatter in the long cycle light curve, (ii) on the first interval the main eclipse seems to be shallower, (iii) significant variability is observed during secondary eclipse; the secondary eclipse seems to occur earlier in interval 1 than in interval 2, and (iv) the shapes of the eclipses vary minimally during

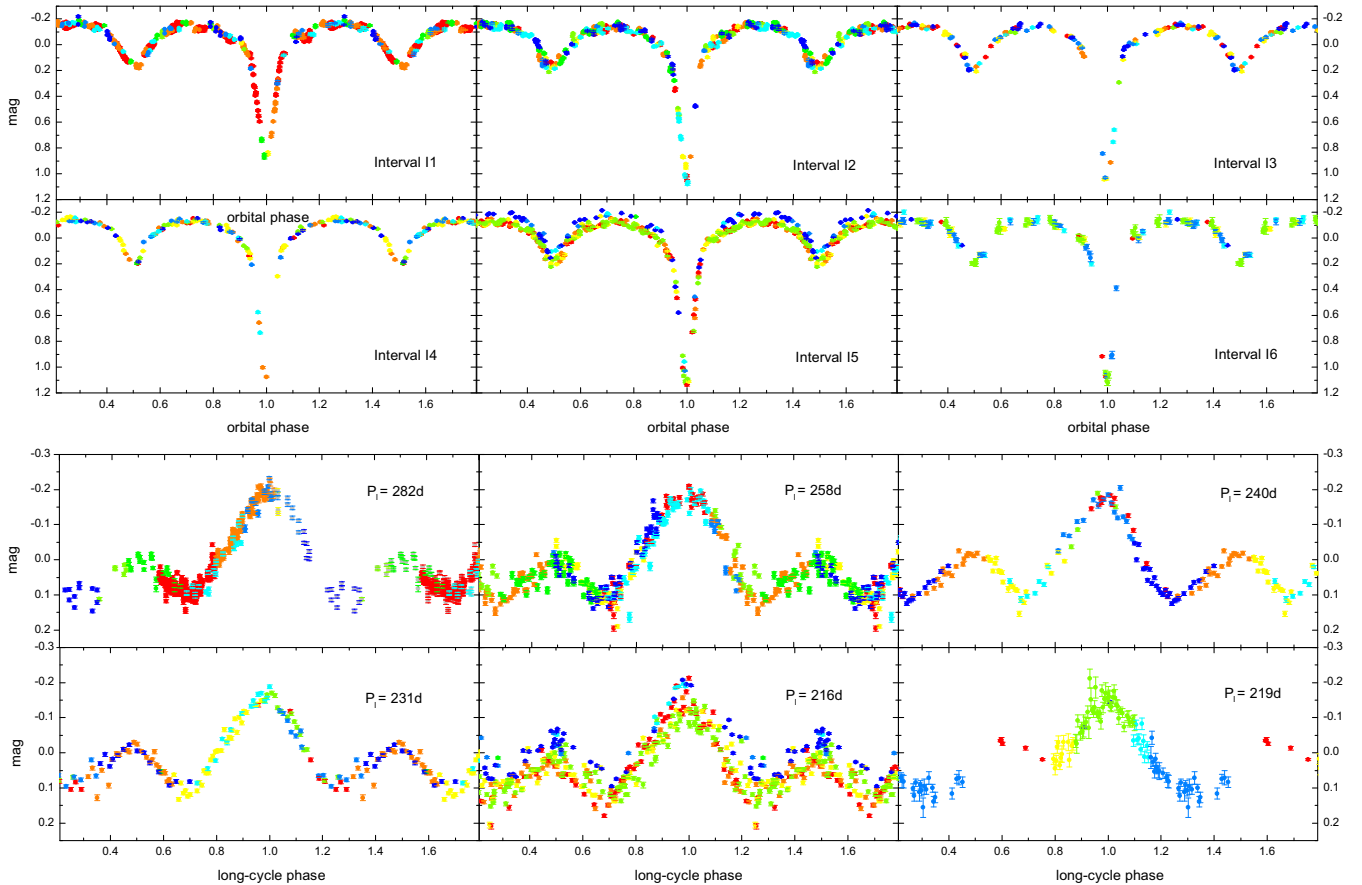


Figure 3. (up): Orbital phase curves at the intervals 1–6 defined in Table 4. (Down): Long cycle phase curves with different periods. Intervals 1–6 defined in Table 4 are illustrated from the top-left to down-right panels. In both panels the magnitude is differential *I*-band and the color is used to indicate time strings of nearby data points.

the maximum, the minimum, and the secondary maximum of the long cycle, perhaps the egress of the main eclipse around phases 0.1–0.2 is shallower during the low stage. The changes in timing of minima during the secondary eclipse might indicate changes in the photo-center of the eclipsed or eclipsing source, or changes in circumstellar matter or the donor hemisphere facing the gainer. An unseen/undetected body that dynamically affects the photo-center is another possibility.

We did the same exercise with the *V* band but we had to use a smaller number of intervals due to the smaller number of observations in this band. The intervals are documented in Table 5. The long cycle usually has a smaller amplitude than in the *I*-band and the orbital light curve shows subtle variability. These changes are better visualized in the combined light curve (Fig. 5).

3.3 Spectra components and orbital/system parameters

In order to obtain the radial velocities and orbital parameters, we used the `korel` code (Hadrava 1995, 1997) based on the method of Fourier disentangling, yielding directly the orbital parameters together with the decomposed spectra of the multiple stellar system under study. In addition, we also used the code `fotel` (Hadrava 1990) to estimate the errors of the orbital parameters.

We notice that the system can be classified as SB2, i.e. both stellar components are detected in the spectrum, in particular in helium and hydrogen lines. The detected components correspond to an early B-

type (primary or gainer) and an early A-type (secondary or donor). The method of spectra disentangling does not use any template or another information about the laboratory wavelengths of the spectral lines, therefore the systemic velocity is set to zero. For this reason we adopted an average of systemic velocities calculated by Gaussian adjustments for different spectral lines of each component. For the gainer we obtain $\gamma_{\text{pri}} = 275.1 \pm 2.3 \text{ km s}^{-1}$ and for the donor $\gamma_{\text{sec}} = 279.8 \pm 2.8 \text{ km s}^{-1}$. The lines used in this calculation are shown in Table 6.

We notice that S II/III measurements systematically differ from other lines, suggesting a different formation place. For this reason they are not included in the above calculation. The disentangled spectra are shown in Figs 6 and 7.

We performed the calculation of radial velocity in seven regions of every spectrum. These regions were chosen because they include several narrow, unblended, and well-identified metallic lines. All our spectra were prepared with a routine written in `iraf6` and the sampling auxiliary code `prekor` (Hadrava 2004) was used. To diminish the numerical errors of the disentangling, we sampled each spectral region in the maximum number of bins allowed by the code, viz. 4096. This results in the average resolution 0.726 km s^{-1} per bin, which is higher than the original resolution on the spectrograph detector.

Radial velocities obtained with `korel` for the cases of circular orbit are given in Table 7 and their best fit is shown in Fig. 8. The operation of the `KOREL` code is described in Hadrava (2004).

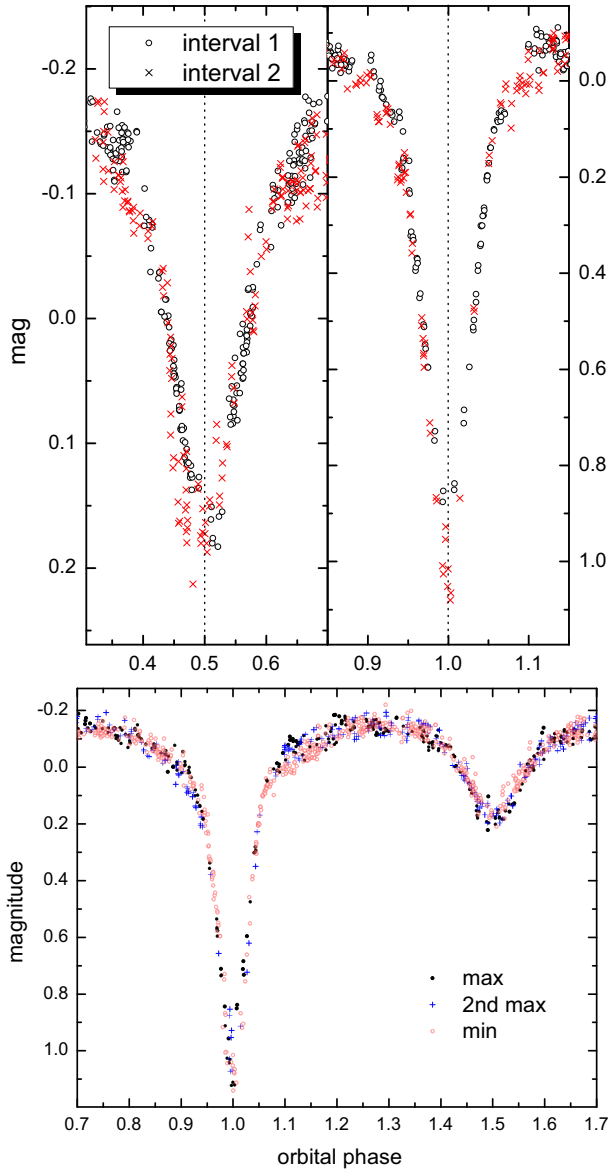


Figure 4. Up: A zoom into the secondary (left) and primary (right) eclipses for data intervals I1 and I2, as defined in Table 4. Vertical axis shows the differential *I*-band magnitude. Down: Eclipses during main and secondary long cycle maxima ($0.9 < \Phi_l \leq 1.1$ and $0.4 < \Phi_l \leq 0.6$, respectively) and minima ($0.2 < \Phi_l \leq 0.4$ and $0.6 < \Phi_l \leq 0.8$). Vertical axis shows the differential *I*-band magnitude.

The radial velocity for each components is given by

$$v_j(t, p) = \sum_0 K(\cos(\omega + v) + e \cos \omega), \quad (2)$$

where the sum is realized on the orbits that influence the movement of the star. The true anomaly v is calculated according to

$$v = 2 \arctan \left(\sqrt{\frac{1+e}{1-e}} \tan \frac{E}{2} \right), \quad (3)$$

where E is obtained from the solution of Kepler's equation.

The orbital parameters obtained by disentangling the seven spectral regions are summarized in Table 8. Solution I, which we accept for our modeling of photometry, has been obtained using an independent disentangling of each region separately and then cal-

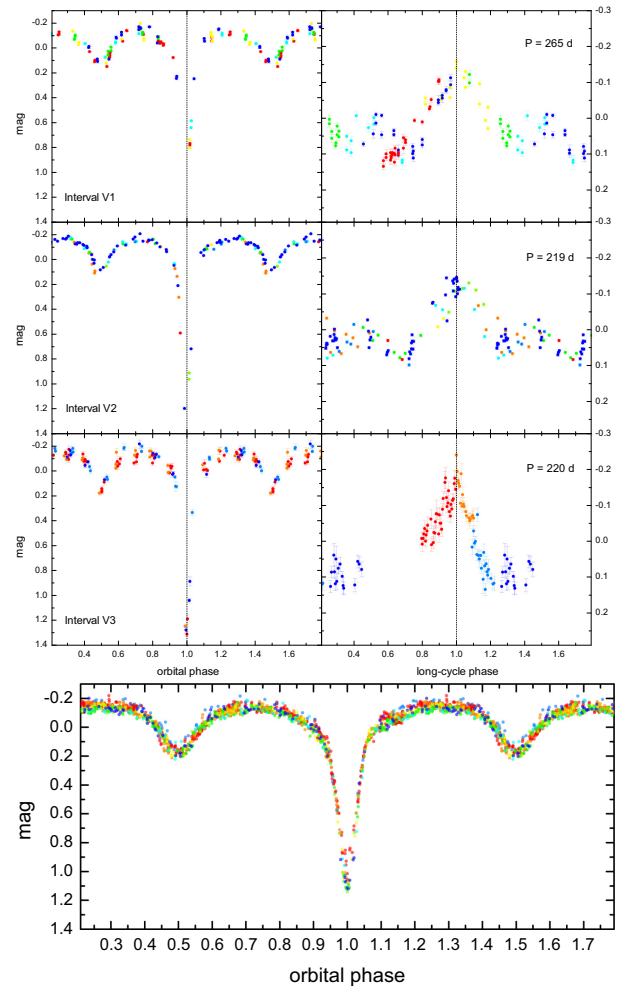


Figure 5. Up: Orbital (left) and long cycle (right) *V*-band differential light curves at the intervals defined in Table 5. The color is used to indicate time strings of nearby data points. Down: Combined data *V*-band light curve. Colors indicate time strings of subsequent points.

Table 5. Data intervals used for disentangling the *V*-band light curve. HJD are referred to 2 450 000. n is the number of data points. Times for long cycle maxima are given.

Label	n	HJD range	P_l (d)	$T_0(max)$
V1	95	468 : 1632	265	580.73
V2	90	2991 : 4949	219	2858.65
V3	95	6965 : 7328	220	6789.63

culating mean solutions and standard deviations of each parameter. Solution II is the simultaneous ('multiregion') disentangling of all the regions together. The errors of the parameters were obtained using the Bayesian estimate, i.e. from the moments of the Bayesian probability distribution (Hadrava 2016). We have also solved the radial-velocity curve using the `fote1` code with the input radial velocities obtained from the disentangling. The resulting values of parameters were within the error bars of Solution I, but their errors were for about one order underestimated, so we skipped this solution. Finally, the multiregion Solution III is to verify that a possible eccentricity of the orbit can be neglected.

Table 6. RV zero points derived from different lines.

Spectral line ion λ (Å)	γ_{pri} (km s ⁻¹)	γ_{sec} (km s ⁻¹)
SiII 4128.054	-	281.136
SiII 4130.894	-	278.431
HeI 4143.76	279.219	-
SII 4153.068	290.429	-
NII 4227.74	276.360	-
FeII 4233.172	-	280.476
NII 4236.91	277.24	-
NII 4241.78	272.029	-
CrII 4242.364	-	275.769
ScII 4246.822	-	275.291
SIII 4253.589	289.218	-
OII 4414.884	275.352	-
OII 4416.974	274.255	-
HeI 4437.551	276.146	-
TiII 4443.794	-	277.682
NII 4447.04	273.909	-
TiII 4533.960	-	281.188
SII 4552.410	289.434	-
SiIII 4567.840	276.484	-
SiIII 4574.757	274.978	-
TiII 4549.617	-	284.985
CrII 4558.650	-	277.747
TiII 4563.757	-	281.192
TiII 4571.968	-	282.472
OII 4590.971	274.638	-
OII 4596.174	273.442	-
NII 4607.153	272.240	-
SiII 4621.418	271.289	-
FeII 4629.336	-	278.677
OII 4638.854	276.561	-
OII 4641.811	276.864	-
OII 4649.138	276.155	-
OII 4699.21	270.561	-
OII 4705.355	276.678	-
HeI 4713.143	278.230	-
FeII 4731.453	-	280.792
Mean (no SiII/III)	275.132 ± 2.325	279.763 ± 2.832
Mean (all)	277.030 ± 5.464	-

Once disentangled the donor spectrum, we compared it with a grid of solar-metallicity synthetic models constructed with `spectrum6` and search for the synthetic spectrum minimizing residuals. We find the best fit with a stellar spectrum of $T_{\text{eff},2} = 9825 \pm 75$ K, $v_{\text{rot},2} \sin i = 53 \pm 3$ km s⁻¹ and $\log g = 3.2 \pm 0.2$. Comparisons of the donor disentangled spectrum with the best-fitting model are shown in Fig. 6. Similarly, from the region 4120–4199 Å, we obtained a model with $T_{\text{eff},1} = 22\,000$ K and $v_{\text{rot},1} \sin i = 70.6$ km s⁻¹ for the gainer.

4 MODELS FOR THE SYSTEM

4.1 Model for an optically thick disc around the gainer

Part of the phenomenology of DPVs has been associated with the presence of an optically thick disc around the gainer, probably fed by a Roche-lobe filling donor star (e.g. Garcés et al. 2018). Consistently, we model the orbital light curve of OGLE-LMC-DPV-065 considering the stellar fluxes of the two stars, the contribution

of an accretion disc around the primary and eventually the light contribution of hot/bright spots located in the outer disc rim. The basic elements of the binary system model with a plane-parallel disc and the corresponding light curve synthesis procedure are described by Djurašević (1992a,b, 1996). The code has been successfully applied to several close binaries including the well-studied binary system β Lyrae (e.g. Djurašević et al. 2010; Djurašević et al. 2012; Mennickent et al. 2012; Garrido et al. 2013; Mennickent & Djurašević 2013).

We assume that the disc is optically and geometrically thick and that its outer edge is approximated by a cylindrical surface. The vertical thickness of the disc can change linearly with radial distance, allowing different disc's conical shapes: plane-parallel, concave, or convex. The geometrical parameters of the disc are its radius (\mathcal{R}_d), its vertical thickness at the outer edge (d_e), and the vertical thickness at the inner boundary (d_c). The cylindrical edge of the disc is characterized by its temperature, T_d , and the conical surface of the disc by a radial temperature profile inspired in the temperature distribution proposed by Zola (1991):

$$T(r) = T_d + (T_h - T_d) \left[1 - \left(\frac{r - \mathcal{R}_h}{\mathcal{R}_d - \mathcal{R}_h} \right) \right]^{a_T} \quad (4)$$

We further assume that the disc is in physical contact and thermal equilibrium with the gainer, so its inner radius and corresponding temperature are equal to the radius and temperature of the star (\mathcal{R}_h , T_h), respectively. The temperature of the disc at the edge (T_d) and the temperature exponent (a_T), as well as the radii of the gainer (\mathcal{R}_h) and of the disc (\mathcal{R}_d) are free parameters, determined by solving the inverse problem (see Section 4.2).

Motivated by previous research on DPVs (Mennickent et al. 2016), our model includes active regions on the edge of the disc. These active regions are usually revealed in Doppler tomography maps of discs in Algol binaries (Richards 2004). These regions have higher local temperatures than the disc, and produce a non-uniform distribution of radiation. We consider two active regions: a hot spot (h_s) and a bright spot (b_s), characterized by their temperatures T_{h_s, b_s} , angular dimensions (radius) θ_{h_s, b_s} , and longitudes λ_{h_s, b_s} . The longitude λ is measured from the line joining the gainer and donor centres in the direction opposite to the orbital motion in the orbital plane. These parameters are also determined by solving the inverse problem. We also consider a possible departure of symmetry of light emerging from the hot spot due, for instance, to the impact of the gas stream in the disc. This deviation is described by the angle θ_{rad} between the line perpendicular to the local disc edge surface and the direction of the hot spot maximum radiation in the orbital plane. The second spot in the model (here named bright spot) simulates the spiral structure of an accretion disc, observed in hydrodynamical calculations (Heemskerk 1994). The tidal forces exerted by the donor star produce a spiral shock, observed as one or two extended spiral arms in the outer disc regions. This bright spot can also be interpreted as a region where the disc significantly deviates from the circular shape.

Two potential limitations of the code need to be briefly mentioned: the lack of treatment of the donor irradiation by the hot spot and the lack of inclusion of a possible not eclipsed *additional* third light, considering that the long-cycle light was already removed with the process of disentangling. However, the very good fit obtained (based on χ^2 minimization) suggests that these additional light sources, if present, are much fainter than the stars and the disc/spots. In addition, while the donor irradiation by the hot spot is not included, the much larger effect of the donor irradiation by the gainer is implemented in our code.

⁶<http://www.appstate.edu/grayro/spectrum/spectrum.html>

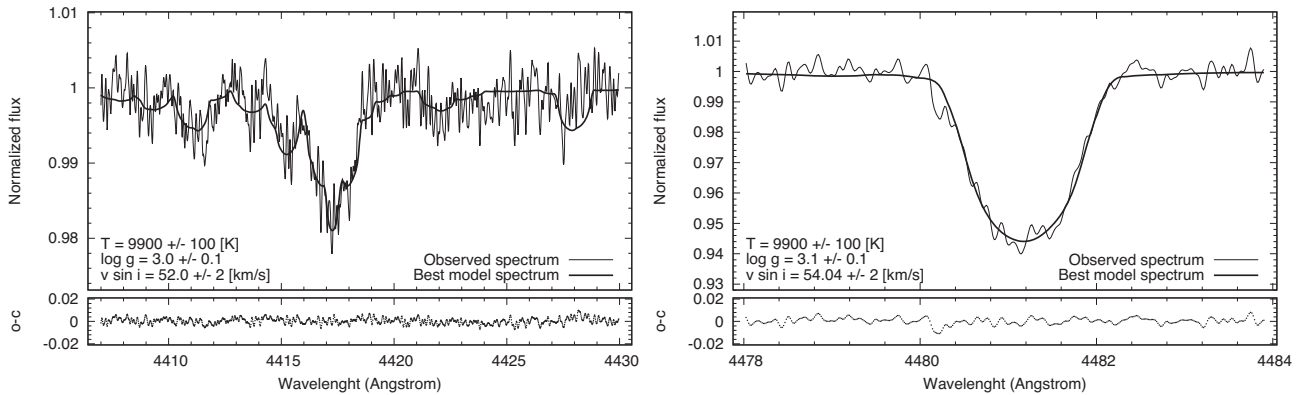


Figure 6. Disentangled donor spectrum and the best-fitting model at two different spectral regions.

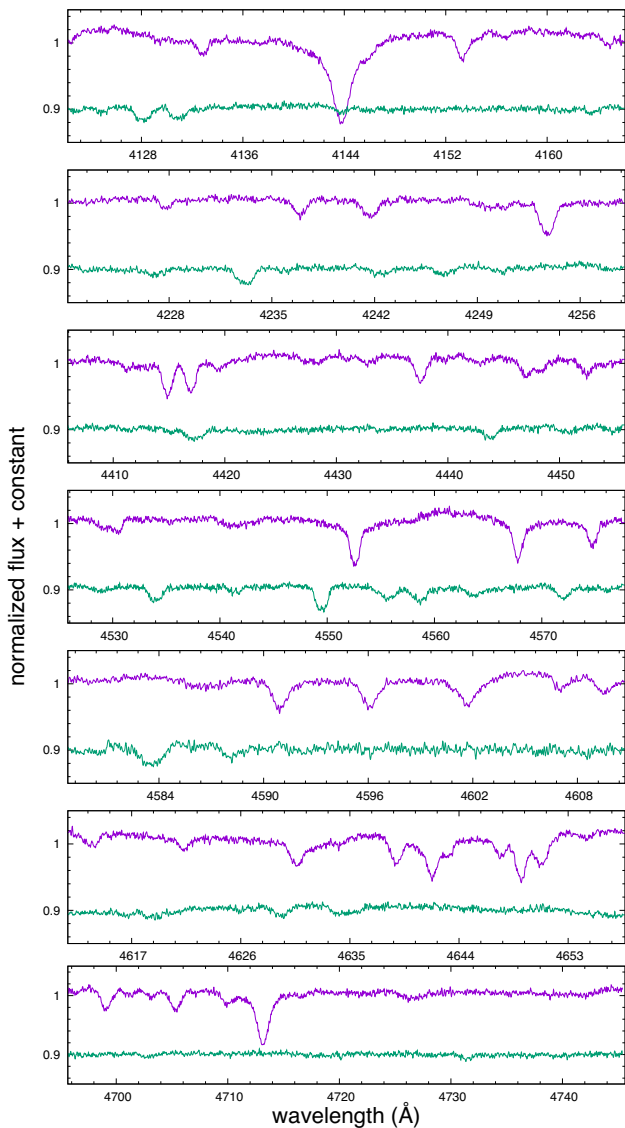


Figure 7. The component spectra in our seven selected regions disentangled using *korel*. For each panel, the upper spectrum corresponds to the primary component and the lower spectrum to the secondary component. The wavelength scale is shifted for the mean γ -velocity to correspond to the laboratory wavelengths of the lines.

4.2 The fit to the orbital light curve

The fit to the orbital light curve was performed using the inverse-problem-solving method based on the simplex algorithm, and the model for the binary system described in the previous section. The inverse-problem method is the process of finding the set of parameters that will optimally fit the synthetic light curve to the observed data. We used the Nelder–Mead simplex algorithm (Press et al. 1992) with the optimization described by Dennis & Torczon (1991). For details, see Djurašević (1992b).

Based on results of the previous sections, we fixed the spectroscopic mass ratio to $q = 0.203$ and the donor temperature to $T_2 = 9825$ K. In addition, we set the gravity darkening coefficient and the albedo of the gainer and the donor to $\beta_{h,c} = 0.25$ and $A_{h,c} = 1.0$, following von Zeipel’s law for radiative shells and complete re-radiation (von Zeipel 1924). The limb darkening for the components was calculated as described by Djurašević et al. (2010).

We assume that the donor is rotating synchronously, i.e. the non-synchronous rotation coefficient, defined as the ratio between the actual and the Keplerian angular velocity is $f_c = 1.0$. This is justified since it is assumed that the donor has filled its Roche lobe (i.e. the filling factor of the donor was set to $F_c = 1.0$), then the synchronization is expected as a consequence of the tidal forces.

The case for the gainer is different, since the accreted material from the disc is expected to transfer angular momentum increasing the rotational speed of the gainer up to the critical velocity as soon as even a small fraction of the mass has been transferred (Packet 1981; de Mink, Pols & Glebbeek 2007; Deschamps et al. 2013). For this reason we assumed critical rotation for the gainer, and estimated a non-synchronous rotation factor $f_h = 8.9$ in the critical rotation regime.

The best fit along with the $O - C$ residuals, individual donor, disc, and gainer flux contributions and the view of the optimal model at orbital phases 0.25, 0.50, and 0.75 are shown in Fig. 9. We note that the residuals show no dependence on orbital or long-cycle phases, except a larger random scatter around main eclipse. Parameters are given in Table 9 and the sensitivity of χ^2 with some parameters is illustrated in Fig. 10. We find that at quadrature and I band, the gainer contributes 27 per cent more flux than the donor and the disc only 48 per cent of the donor to the total flux.

We find that the system contains a $13.8 M_\odot$ star and a $2.81 M_\odot$ star with absolute magnitude $M_{bol} = -6.4$ and -3.0 , respectively, separated by $49.9 R_\odot$. The stellar temperatures are $T_h = 25\,460$ K and T_c (fixed) = 9825 K. The best-fitting model contains an

Table 7. RVs for the primary and secondary components from `kore1` solutions in a circular orbit. The radial velocity is the average from each spectral region and we considered the standard deviation as error. Orbital phases Φ_o are given for the orbital ephemerides given by equation (1) and long cycle phases Φ_l for a long period of 219 d and $T_0(\text{max}) = 57\,009.62$.

Date-ut	HJD -2450000.0	ϕ_o	ϕ_l	RV_g km s ⁻¹	RV_d km s ⁻¹
2013-10-02	6567.7689	0.547	0.982	12.082 ± 0.789	-59.521 ± 3.382
2013-10-04	6569.8218	0.751	0.992	43.182 ± 1.020	-209.288 ± 5.953
2013-10-06	6571.8038	0.949	0.001	14.536 ± 0.830	-67.676 ± 2.214
2013-10-07	6572.8075	0.049	0.005	-13.402 ± 0.618	61.623 ± 3.562
2013-10-19	6584.7162	0.236	0.060	-43.399 ± 1.163	207.660 ± 7.547
2013-10-22	6587.8479	0.548	0.074	12.647 ± 0.685	-60.792 ± 3.658
2013-12-22	6648.5833	0.603	0.351	25.914 ± 0.940	-124.089 ± 5.008
2013-12-24	6650.7561	0.819	0.361	39.975 ± 0.886	-190.796 ± 4.926
2013-12-31	6657.5987	0.501	0.393	0.115 ± 0.649	0.015 ± 2.346
2014-01-04	6661.6647	0.907	0.411	24.468 ± 0.615	-117.341 ± 2.902
2014-01-18	6675.6735	0.303	0.475	-41.331 ± 0.952	198.184 ± 6.020
2014-01-19	6676.6249	0.398	0.480	-26.363 ± 0.566	126.834 ± 2.904
2014-02-11	6699.5976	0.688	0.584	40.177 ± 0.934	-193.185 ± 5.941
2014-02-15	6703.5569	0.083	0.603	-21.199 ± 0.864	103.465 ± 3.704
2014-02-16	6704.5526	0.182	0.607	-39.697 ± 0.839	189.726 ± 6.078
2014-09-01	6901.8820	0.853	0.508	34.533 ± 0.659	-169.289 ± 4.214
2014-11-03	6964.8370	0.128	0.796	-31.257 ± 0.812	149.933 ± 5.166
2014-11-20	6981.7855	0.818	0.873	39.944 ± 0.924	-191.460 ± 5.220
2014-11-22	6983.7327	0.012	0.882	-3.961 ± 1.684	13.814 ± 2.907
2014-11-25	6986.7551	0.313	0.896	-40.347 ± 0.907	193.721 ± 5.611
2014-11-26	6987.7523	0.413	0.900	-23.026 ± 0.504	111.558 ± 3.302
2014-11-27	6988.7666	0.514	0.905	3.562 ± 0.391	-16.488 ± 3.285
2014-12-08	6999.6148	0.595	0.954	24.437 ± 1.068	-116.546 ± 5.316
2014-12-09	7000.7687	0.710	0.960	42.190 ± 1.088	-202.936 ± 5.340
2014-12-14	7005.6695	0.199	0.982	-41.140 ± 0.994	198.112 ± 6.064
2015-01-20	7042.5709	0.877	0.151	30.758 ± 0.632	-147.288 ± 3.797
2015-02-01	7054.5922	0.075	0.205	-19.228 ± 0.538	94.275 ± 3.972

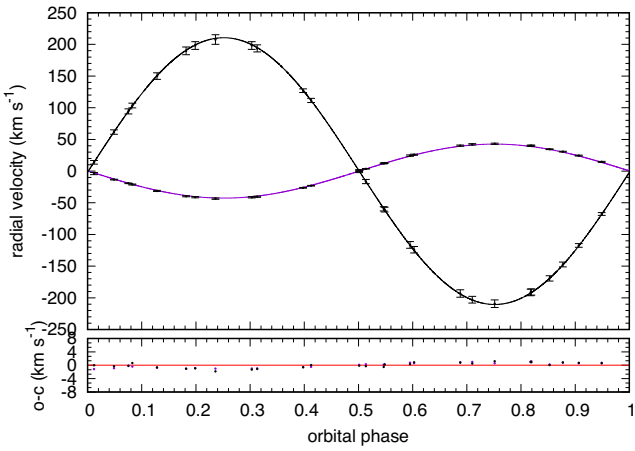


Figure 8. Fit of theoretical RVs to the average velocities from Table 7.

Table 8. Orbital parameters obtained in different solutions.

Parameter	I	II	III
P (d)		10.031 6267 (fixed)	
τ^*	92.31 ± 0.02	92.305 ± 0.004	94.79 ± 0.33
K_1 (km s ⁻¹)	42.60 ± 0.97	42.44 ± 0.33	42.45 ± 0.32
K_2 (km s ⁻¹)	210.5 ± 6.4	214.1 ± 1.8	213.5 ± 1.7
q	0.203 ± 0.008	0.198 ± 0.002	0.199 ± 0.002
ω (deg)	90	90	178.7 ± 11.6
e	0	0	0.021 ± 0.006

optically and geometrically thick disc around the hotter, more massive gainer star. With a radius of $\mathcal{R}_d \approx 25 R_\odot$, the disc is 2.8 times larger than the central star ($\mathcal{R}_h \approx 8.8 R_\odot$). The disc has a convex shape, with central thickness $d_c \approx 1.6 R_\odot$ and edge thickness $d_e \approx 3.5 R_\odot$. The temperature of the disc decreases from $T_d = 9380$ K at its edge, to $T_h = 25460$ K at the inner radius.

We notice that the hot spot has 20:5 angular radius and covers 12 per cent of the disc outer rim, and it is situated at longitude $\lambda_{\text{hs}} = 312^\circ.4$, roughly between the components of the system, at the place where the gas stream falls on to the disc (Lubow & Shu 1975). The temperature of the hot spot is approximately 18 per cent higher than the disc edge temperature, i.e. $T_{\text{hs}} = 11068$ K. Although including the hot spot region into the model improves the fit, it cannot explain the light-curve asymmetry completely. By introducing one additional bright spot, larger than the hot spot and located on the disc edge at $\lambda_{\text{bs}} = 114^\circ.9$, the fit becomes much better. This bright spot is 26 per cent hotter than the disc at its edge, i.e. $T_{\text{bs}} = 11819$ K, and has an angular radius of $27^\circ.9$, covering 16 per cent of the disc outer rim.

5 DISCUSSION

Only a few DPVs have been studied spectroscopically in detail and therefore few of them possess relatively well-determined orbital and stellar parameters; nine Galactic DPVs and the LMC DPV OGLE05155332-6925581 are documented by Mennickent et al. (2016) and recently stellar and orbital parameters were provided for V 495 Cen by Rosales Guzmán et al. (2018). Our study of OGLE-

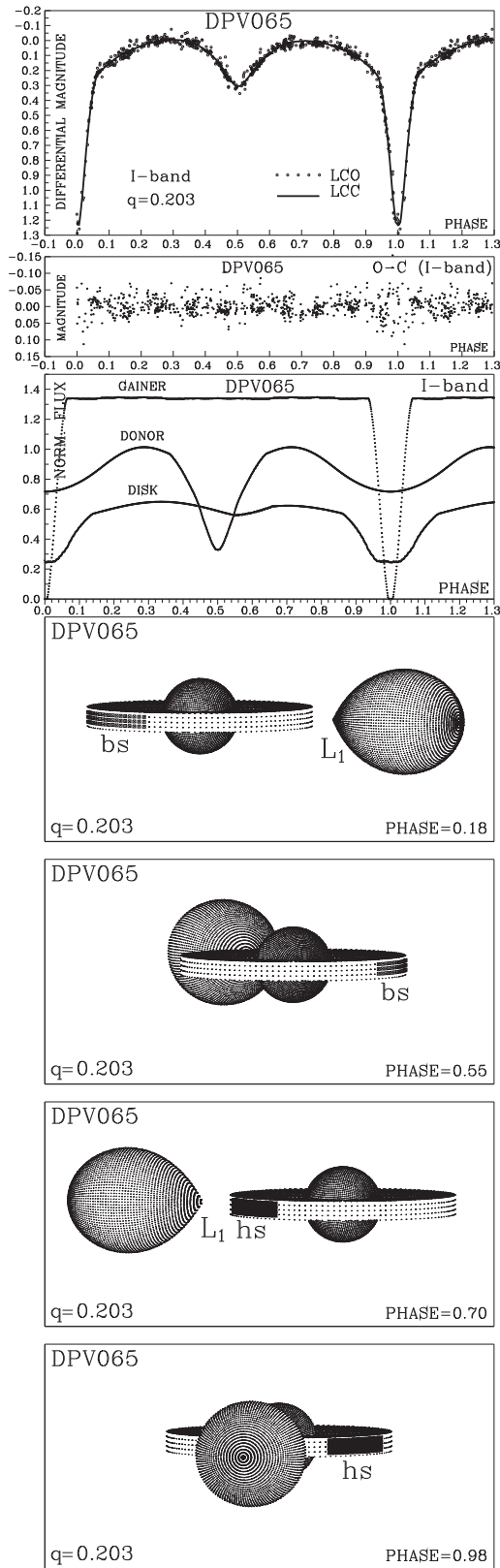


Figure 9. Observed (LCO), synthetic (LCC) light-curves and the final O–C residuals between the observed and synthetic light curves of OGLE-LMC-DPV-065; fluxes of donor, gainer and of the accretion disc, normalized to the donor flux at phase 0.25; the views of the model at orbital phases 0.18, 0.55, 0.70, and 0.98, obtained with parameters estimated by the light-curve analysis. This is for a gainer in critical non-synchronous rotation regime.

Table 9. Results of the analysis of DPV065 light curves obtained by solving the inverse problem for the Roche model with a large accretion disc partially obscuring the more-massive (hotter) gainer in critical non-synchronous rotation regime.

Quantity	Quantity		
n	718		
$\Sigma(O - C)^2$	0.4842		
σ_{rms}	0.0260		
i ($^\circ$)	86.7 ± 0.4	$\mathcal{M}_h [M_\odot]$	13.8 ± 0.3
F_d	0.948 ± 0.03	$\mathcal{M}_c [M_\odot]$	2.81 ± 0.2
T_d (K)	9380 ± 400	$\mathcal{R}_h [R_\odot]$	8.8 ± 0.3
d_c (a_{orb})	0.071 ± 0.005	$\mathcal{R}_c [R_\odot]$	12.6 ± 0.3
d_c (a_{orb})	0.032 ± 0.008	$\log g_h$	3.68 ± 0.05
a_T	4.6 ± 0.5	$\log g_c$	2.69 ± 0.06
F_h	1.000	M_{bol}^h	-6.39 ± 0.2
T_h (K)	25460 ± 1500	M_{bol}^c	-3.01 ± 0.1
T_c (K)	9825	$a_{\text{orb}} [R_\odot]$	49.91 ± 0.2
$A_{\text{hs}} = T_{\text{hs}}/T_d$	1.18 ± 0.07	$\mathcal{R}_d [R_\odot]$	25.0 ± 0.4
θ_{hs} ($^\circ$)	20.5 ± 0.6	$d_e [R_\odot]$	3.5 ± 0.2
λ_{hs} ($^\circ$)	312.4 ± 8.0	$d_c [R_\odot]$	1.6 ± 0.1
θ_{rad} ($^\circ$)	-21.5 ± 2.2		
$A_{\text{bs1}} = T_{\text{bs}}/T_d$	1.26 ± 0.07		
θ_{bs} ($^\circ$)	27.9 ± 4.0		
λ_{bs} ($^\circ$)	114.9 ± 10.0		
f_h	8.9 ± 0.3		
Ω_h	7.067 ± 0.02		
Ω_c	2.240 ± 0.02		

Note: Fixed parameters: $q = \mathcal{M}_c/\mathcal{M}_h = 0.203$ – mass ratio of the components; $T_c = 9825$ K – temperature of the less-massive (cooler) donor; $F_c = 1.0$ – filling factor for the critical Roche lobe of the donor; $f_h = 8.9$; $f_c = 1.00$ – non-synchronous rotation coefficients of the gainer and the donor, respectively, $\beta_h = 0.25$; $\beta_c = 0.25$ – gravity-darkening coefficients of the gainer and the donor, $A_h = 1.0$; $A_c = 1.0$; $A_d = 1.0$ – albedo coefficients of the gainer, the donor, and disc.

Quantities: n – number of observations; $\Sigma(O - C)^2$ – final sum of squares of residuals between observed (LCO) and synthetic (LCC) light-curves; σ_{rms} – root mean square of the residuals; i – orbit inclination (in arc degrees); $F_d = R_d/R_{yc}$ – disc dimension factor (ratio of the disc radius to the critical Roche lobe radius along y -axis), T_d – disc-edge temperature, d_e , d_c – disc thicknesses (at the edge and at the center of the disc, respectively) in the units of the distance between the components, a_T – disc temperature distribution coefficient, $F_h = R_h/R_{zc}$ – filling factor for the critical Roche lobe of the hotter, more-massive gainer (ratio of the stellar polar radius to the critical Roche lobe radius along z -axis for a star in critical rotation regime), T_h – temperature of the more-massive (hotter) gainer; $A_{\text{hs, bs}} = T_{\text{hs, bs}}/T_d$ – hot and bright spots’ temperature coefficients; $\theta_{\text{hs, bs}}$ and $\lambda_{\text{hs, bs}}$ – spots’ angular dimensions and longitudes (in arc degrees), θ_{rad} – angle between the line perpendicular to the local disc edge surface and the direction of the hot-spot maximum radiation, f_h – non-synchronous rotation coefficients of the gainer in critical rotation regime, $\Omega_{h, c}$ – dimensionless surface potentials of the hotter gainer and cooler donor, $\mathcal{M}_{h, c} [M_\odot]$, $\mathcal{R}_{h, c} [R_\odot]$ – stellar masses and mean radii of stars in solar units, $\log g_{h, c}$ – logarithm (base 10) of the system components effective gravity; $M_{\text{bol}}^{h, c}$ – absolute stellar bolometric magnitudes; $a_{\text{orb}} [R_\odot]$, $\mathcal{R}_d [R_\odot]$, $d_e [R_\odot]$, $d_c [R_\odot]$ – orbital semimajor axis, disc radius, and disc thicknesses at its edge and center, respectively, given in the solar radius units.

LMC-DPV-065 presented in this paper is the second spectroscopic study of an LMC DPV.

In Fig. 11 we compare OGLE-LMC-DPV-065 data with those of other DPVs and classical Algols, these later taken as reference. It is clear that DPVs are hotter and more massive than ordinary Algols, a fact already noticed in previous studies. In addition, it is clear that OGLE-LMC-DPV-065 is a comparatively massive and hot DPV, in many aspects similar to β Lyrae. In Table 10 we provide a compar-

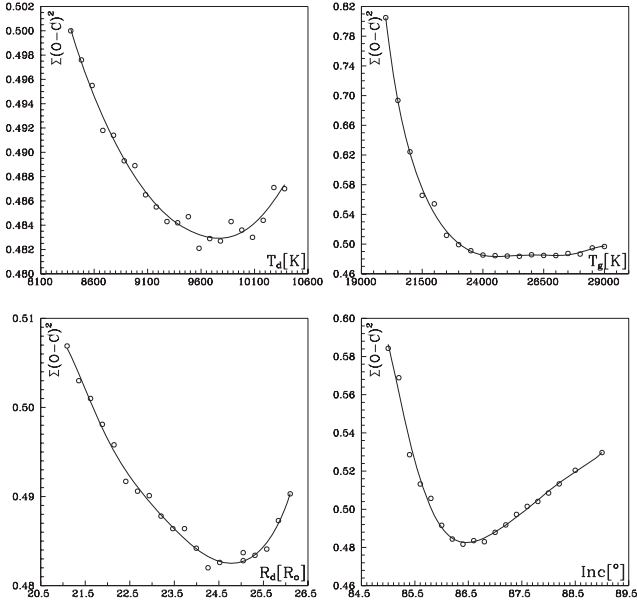


Figure 10. Plots showing the dependence of χ^2 .

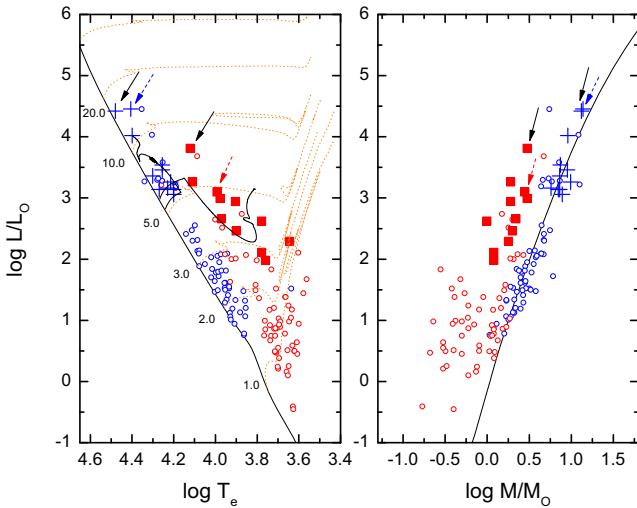


Figure 11. Comparison of physical data for semidetached Algols from Dervişoğlu, Tout & Ibañoğlu (2010, primaries open blue circles and secondaries open red circles) and DPVs (Mennickent et al. 2016, primaries blue crosses and secondaries red squares). The zero-age main sequence for $Z = 0.02$ is plotted with a solid black line and evolutionary tracks for single stars with initial masses (in solar masses) labelled at the track footprints are also shown (Pols et al. 1998). The positions of the gainer and the donor of OGLE-LMC-DPV-065 are indicated by color arrows, whereas β Lyrae is labelled with black arrows. The best evolutionary tracks for the primary and secondary of the DPV HD 170582 are also plotted by two solid lines in the left-hand panel (Mennickent et al. 2016).

ison between these systems based on the results of Mennickent & Djurašević (2013), although see also the recent research on the β Lyrae disc by Mourard et al. (2018) for a complementary approach confirming the existence of a hot spot and obtaining roughly the same disc size but from an interferometric study.

The similarity is especially significant in inclination angle, stellar masses, surface gravities, and time-scale of the long-cycle length. Both systems are found in a mass transfer stage, harbor a comparatively hot accretion disc and massive B + A type stars for the DPV standard (Fig. 9). The radial extension of the disc is also

Table 10. Comparison between the OGLE-LMC-DPV-065 data obtained in this paper and those of β Lyrae obtained by Mennickent & Djurašević (2013) and references therein.

Quantity	OGLE-LMC-DPV-065	beta Lyrae
P_o (d)	10.031 6267	12.95 (variable)
P_l (d)	350-210 (variable)	282.4
P_l/P_o	35-21	22
i ($^\circ$)	86.7	86.1
T_d (K)	9380	8200
d_e (a_{orb})	0.071	0.192
d_c (a_{orb})	0.032	0.01
a_T	4.6	3.8
T_h (K)	25460	30000
T_c (K)	9825	13300
$A_{hs} = T_{hs}/T_d$	1.18	1.21
λ_{hs} ($^\circ$)	312.4	324.6
$A_{bs} = T_{bs}/T_d$	1.26	1.12
λ_{bs} ($^\circ$)	114.9	107.3
\mathcal{M}_h (\mathcal{M}_\odot)	13.8	13.16
\mathcal{M}_c (\mathcal{M}_\odot)	2.81	2.97
\mathcal{R}_h (R_\odot)	8.8	6.0
\mathcal{R}_c (R_\odot)	12.6	15.2
$\log g_h$	3.68	4.0
$\log g_c$	2.69	2.5
M_{bol}^h	-6.39	-6.3
M_{bol}^c	-3.01	-4.7
a_{orb} (R_\odot)	49.91	58.5
\mathcal{R}_d (R_\odot)	25.0	28.3
d_e (R_\odot)	3.5	11.2
d_c (R_\odot)	1.6	0.6

similar along with the location of the hot and bright spots. As a jet has been detected in β Lyrae (Harmanec et al. 1996; Ak et al. 2007; Lomax & Hoffman 2011), it is then possible that the same structure exists in OGLE-LMC-DPV-065 and could be related to the long-cycle through a magnetic dynamo as suggested by Schleicher & Mennickent (2017). On the other hand, an important difference is the remarkable long-cycle change observed in OGLE-LMC-DPV-065, which is not observed in β Lyrae. The large amplitude of the long-cycle in OGLE-LMC-DPV-065 is also remarkable. In comparison, the long-cycle in β Lyrae is of low amplitude and relatively constant in period. Orbital period changes can be explained in terms of conservative mass transfer in a binary system. Hence it is possible that both systems are in different stages of the mass transfer episode. A much larger mass transfer in β Lyrae might explain why this binary shows a variable orbital period, whereas OGLE-LMC-DPV-065, eventually with a smaller mass transfer rate, does not. In addition, β Lyrae has a larger and brighter secondary star, which might also play a role in the observed differences between both systems. These issues will be investigated in a forthcoming paper.

If a magnetic dynamo is the cause for the long-cycle, then these two systems with similar parameters but different long cycle light curve morphology, constitute constraints to be satisfied by any competent detailed physical model of the long variability. Our next study will explore this point, establishing the evolutionary stage of OGLE-LMC-DPV-065 and analyzing the spectroscopic changes during the long cycle. We will also present numerical calculations aimed to test the hypothesis of variable mass transfer driven by a magnetic dynamo as proposed by Schleicher & Mennickent (2017).

6 CONCLUSIONS

We have analyzed the variability of the eclipsing Algol OGLE-LMC-DPV-065 considering new and published photometric data

spanning 124 yr. The orbital and long-cycle light curves have been disentangled and characterized. We also presented the first spectroscopic study of this binary system obtaining the mass ratio and temperature of the cooler stellar component. These quantities served as fixed input parameters in our model of the light curve that was done following an inverse-problem methodology. The best solution shows a reasonable fit to the light curve providing additional parameters for the binary, the stellar components, and the circumprimary accretion disc. The main results of our research can be summarized as follows:

- (i) We find a refined orbital period of $10^d031\ 6267 \pm 0^d000\ 0056$ without any evidence of variability.
- (ii) Small but significant changes in timings of the secondary eclipse are detected. They might be caused by circumstellar material.
- (iii) The long cycle is characterized by a double hump light curve at *I* and *V* bands, of amplitude about 0.3 and 0.2 mag, respectively, whose general shape is more or less constant, with only minor variability.
- (iv) We find that after a continuous decrease of the long-period during about 13 yr, from 350 to 218 d, it remained almost constant by about 10 yr.
- (v) The study of radial velocities indicates a binary in a circular orbit with a mass ratio of 0.203 ± 0.001 .
- (vi) We find that the system consists of a pair of stars of 13.8 and $2.81 M_{\odot}$ of radii 8.8 and $12.6 R_{\odot}$ and absolute bolometric magnitudes -6.4 and -3.0 , respectively.
- (vii) We find stellar temperatures of 25 460 K and 9825 K for the gainer and the donor, respectively.
- (viii) We find an orbital semimajor axis of $49.9 R_{\odot}$.
- (ix) We find evidence of an accretion disc with a radius of $25 R_{\odot}$, central thickness $1.6 R_{\odot}$, and edge thickness $3.5 R_{\odot}$. The temperature of the disc decreases from 25 460 K at the inner radius to 9380 K at its outer edge.
- (x) As happens in other DPVs, two hot shock regions located at roughly opposite parts of the outer disc rim can explain the light-curve asymmetries.
- (xi) OGLE-LMC-DPV-065 resembles in some aspects to the well-studied binary β Lyrae. However, its orbital period does not change; this could indicate a smaller mass transfer rate.

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