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Proton conduction study of a new selenate-based hybrid compound

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ABSTRACT

Ongoing studies aimed at developing a novel proton conducting solids, resulted in the new organic inorganic compound $[(C_2H_5)_4N]$ HSeO₄ (H₂SeO₄)₂. The crystals were prepared by slow evaporation technique at room temperature. Crystal structure, thermal and electrical measurements were provided to characterize this compound. Single crystal X ray study showed that the new phase crystallizes belongs to the Cc space group of the monoclinic system with lattice parameters: a = 19.705 (1) Å; b = 7.5420 (3) Å; c = 19.199 (1) Å, $\beta = 115^{\circ}$. The structure consists of infinite two dimensional inter linkers via strong hydrogen bonds (O H···O) giving birth to trimers. Two phase transition were disclosed by differential thermal analysis and confirmed by electrical measurements and dielectric measurements. The complex impedance of the compound was investigated in the temperature range 290 373 K and in the frequency range 1 13.10³ KHz. Besides, an electrical equivalent circuit was proposed to explain the impedance results. Electric and dielectric analyses reveal the presence of a relaxation phenomenon and highlight the good protonic conduction of this material.

Protons migration can be ascribed to the reorientation of protons carriers through hydrogen bonds, and then it occurs by the Grotthuss mechanism.

1. Introduction

During the past decade, proton conducting compounds have garnered huge attention as candidate for electrolytes in batteries, fuel cells, and sensors [1-4]. Different classes of compounds gained increasing interest as proton conductors such as ceramic oxides, carbon nanotubes, and organic polymers. Particularly, organic inorganic hybrid materials have been recently considered as one of the best promising candidates in term of applications regarding protons conducting. Their high degree of crystallinity, the design [5], the structural diversities and the ability to combine physical properties derived from the co existence of organic and inorganic elements appear as key factors [6–10].

In fact, across the systematic research that aim to better un derstand these features, a significant number of organic inorganic materials based on selenate have been already reported:

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 $[C_3H_7N_6^+]_4(SeO_4^{2-})_2.3H_2O$ [11], $[C_6H_5CH_2C_5H_4NH]$.HSeO₄ [12], $C(NH_2).3HSeO_4$ [13–15], $(NH_4)_4H_2(SeO_4)_3$ [16,17], $[N(C_2H_5)_4]$.HSeO₄ [18]. In these materials, protons migration is ensured through hydrogen bonds involving and the coordinating atom, which is oxygen. The more extensive and stronger the cluster is, the faster the migration will be. Thus, a stronger hydrogen bond reduces the O···O distance and facilitates the close approach intended for transfer. Generally, the distances of the O···O interactions full in the range of 2.4–3 Å. Such phenomenon gives rise directly to the two plausible main processes of proton conduction transfer known as Grotthus and vehicle mechanisms [19,20].

In the present study, focus is laid on the XRD and physical properties of new hybrid compound tetraethylammonium mono hydrogenselenate bis selenic acid $[(C_2H_5)_4N]$.HSeO₄. $(H_2SeO_4)_2$ (hereafter abbreviated TESe).

The ultimate aim of this work is to reveal the phase transition from differential thermal analysis, conductivity analysis through complex impedance spectroscopy and dielectric measurements and give more information about the relaxation phenomenon. Thus, this makes it feasible to point the proton transport and transfer phenomena pathway, thereby leading to understand the

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proton conduction mechanism.

2. Experimental details

2.1. Synthesis

2.1.1. Materials

Reagents and Solvents were commercially available from Sigma Aldrich and used without further purification: selenic acid H₂SeO₄ (purity 98%; Sigma Aldrich) and tetraethylammonium hydroxide [(CH₃CH₂)₄N] OH (purity 98%; Sigma Aldrich).

2.1.2. Preparation of TESe

Crystal of TESe with the formula $[(C_2H_5)_4N]$ HSeO₄ (H₂SeO₄)₂ was prepared by slow evaporation technique, at room temperature, from an aqueous solution containing selenic acid and tetraethy lammonium hydroxide in the stochiometric ratio 3:1. This mixture was stirred well till it became clear. Over a period of several days, colorless parallelepipedic crystals with high quality suitable for X ray single crystal experiments were obtained. The reaction scheme is given in the following equation (Scheme 1):

 $[(C_2H_5)_4N]^+.$ OH^ + 3H_2SeO_4 $[(C_2H_5)_4N]^+.$ HSeO_4^ _. (H_2SeO_4)_2+ H_2O

The density of the synthesized compound was measured at room temperature by pycnometry by flotation in carbon tetra chloride (CCl₄). The mean value of the measured density, D_m 1.709, is in good accordance with that obtained from the crystal structure determination (see below).

2.2. X ray data collection

In order to define the crystal structure, X ray diffraction in tensity data were collected within the limits of $2,1^{\circ} \le \theta \le 27,1^{\circ}$ at room temperature and then investigated using Bruker Kappa APEX II diffractometer using Mo K α radiation source.

5609 reflections were recorded and among them 5472 checked the condition of intensity $I \ge 3\sigma$. Then, the structure of the title compound was solved within the WINGX software [21] in the monoclinic system with the non centro symmetric space group Cc. Heavy atoms such us selenic were fixed with Patterson methods using the SHELXS 86 program [22]. SHELXL 79 program [23] was used to facilitate the identification of carbon, oxygen and nitrogen atoms. The final cycle of refinement led to reliability factors $R_1 = 0.071$ and $wR_2 = 0.2187$. Furthermore, the drawings of the structure was made using Diamond [24].

Crystallographic and structure refinement details for the title compound are depicted in Table 1. Atomic coordinates, U_{eq}/U_{iso} , inter atomic distances and angles are listed in Table S1 and Table S2.

2.3. Thermal analysis

Thermogravimetric analysis (TGA) and differential thermal analysis (DTA) measurements for TESe were performed by heating 3.809 mg of powdered sample from 293 K to 523 K. The TGA analysis was conducted in air at a heating rate of 10 K min⁻¹.

2.4. Electrical measurements

The electrical studies were carried out in frequency ranges be tween 1 and 13.10³ KHz by means of two silver electrode configu rations. In fact, the polycrystalline sample was pressed into 1.1 mm thick pellets with 8 mm diameter. These measurements were recorded in the temperature range from 290 to 373 K using a HEWLETT PACKARD 4192 ALF component analyzer.

3. Results and discussion

3.1. Structure description

The TESe hybrid compound crystallizes in the monoclinic crystal system with the Cc space group and lattice parameters: a 19.705 (1) Å; b 7.5420 (3) Å; c 19.199 (1) Å, β 115° and Z 4. Fig. 1 illustrate the perspective view of the sample in a three dimensional framework which reveals the organic and inorganic groups architecture. The unit formula is constituted of one independent [(C₂H₅) N]+cation, one [HSeO₄]⁻ and two molecules of selenic acid H₂SeO₄ (Fig. 2). The projection of the atomic arrangement of Fig. 3 shows a projection of the structure along the *b* axis. This latter illustrates a layered arrangement perpendicular to *a* axis. The planes of the tetrahedron [HSeO₄] and H₂SeO₄ alternate with planes of organic groups.

The $[(C_2H_5) N]$ +cation processes a tetrahedral geometry in the form of a broken cross. The N–C and C–C distances vary from 1.534 (19) to 1.64 (2) Å and from 1.566 (13) to 1.668 (15) respectively. The C–N–C angle ranges between 95.5° (10) and 118.4° (12). The mid planes of this unit are located at x 0 and x 1/2. These average values are comparable to the previously reported date ones for other similar hybrid compound [18,25].

On the other hand, the structure consists of strong two dimensional character based on sheets of $HSeO_4^-$ and H_2SeO_4 tetrahedra giving rise to an infinite 2D chains built up through $[(HSeO_4)(H_2SeO_4)]_n^{n-}$ layers perpendicular to the *a* axis. The mid planes of these layers are located at x 1/4 and x 3/4 (Fig. 4). Within these inorganic groups, the Se–O distances in the [HSe(1) O₄] anions vary from 1.776 (1) Å to 1.824 (2) Å.

In contrast, the Se–O distances in the $H_2Se(2)O4$ and $H_2Se(3)O_4$ are included between 1.688 (2) and 1.784 (1) respectively (Table 2). In fact, the longest Se–O distance [d Se–O] 1.824 (2) Å] is ascribed to the presence of an acidic hydrogen on the atom of the SeO₄ tetrahedron. The O-Se-O angles are in the range from 104.7 $(6)^{\circ}$ to 115.8 $(6)^{\circ}$ in HSe(1)O₄ anions and 101.7 $(6)^{\circ}$ 115.7 $(3)^{\circ}$ and 102.4 (6)° 114.7 (7)° in H₂Se(2)O₄ and H₂Se(3)O₄, respectively. These values clearly justify the distortion from the ideal tetrahedral geometry. The calculated mean values of the distortion indices are listed in Table 3. For the two HSe(1) O_{4}^{-} and H₂Se (2)O₄ tetrahedral, a strong distortion of the Se-O distances from those of O-O. These values confirm that the molecule H_2 Se (3) O_4 is a slightly deformed tetrahedron. Bond distances and angles are comparable to those reported in [N (CH₃)₄].HSeO₄ [26], (C₃H₇N₆⁺)₄ (SeO₄²⁻)₂.3H₂O [11] and $C_6H_5CH_2C_5H_4NH^+$. $HSeO_4^-$ [12], $(C_6H_{14}N_2)$ [MII(H₂O)₆](SeO₄)₂ [27].

The structure packing is stabilized through intermolecular hydrogen bonding interactions occurring between the inorganic moieties, giving rise to layers perpendicular to the a axis. Indeed, the O-H...O bonds range from 2.480 (16) to 2.784 (17) Å. These interactions play a major role on the cohesion of the crystal

$$[(C_2H_5)_4N]^+$$
. OH + 3 H₂SeO₄ \longrightarrow $[(C_2H_5)_4N]^+$. HSeO₄² (H₂SeO₄)₂+ H₂O

Table 1

Crystallographic and structure refinement details for the title compound for TESe.

Formula	$[(C_2H_5)_4N]$ HSeO ₄ (H ₂ SeO ₄) ₂
Color/Shape	Colorless/Parallelipipedic
Mr (g mol ¹)	564.17
Crystal system	Monoclinic
Space group	Cc
Density	1.717
Crystal size (mm)	$0.40 \times 0.20 \text{ x } 0.10$
Unit-cell volume (A° ³)	2181.9 (2)
Temperature (K)	293 (2)
Diffractometer	Kappa CCD
a, b, c (Å)	19.705(1), 7.542(3), 16.199(1)
α, β, γ (°)	90, 115, 90
Z	4
Radiation type	Mo K_{α}
θ range for data collection (°)	$2.1 \le \theta \le 27.1$
No. of measured, independent and observed $[I > 2\sigma(I)]$ reflections	5066, 4383, 3366
Range of h, k, l	-25/259/119/19
F (000)	1112
Rint	0.029
$(\sin \theta \lambda) \max (\dot{A}^{-1})$	0.650
$R [F^2 > 2\sigma(F^2)], wR (F^2), S$	0.071, 0.217, 1.03
$\Delta \rho$ max, $\Delta \rho$ min (e Å ⁻³)	1.45, -2.40
w	$1/[\sigma^2(Fo^2) + (0.1662P)^2 + 2.0887P]$ where $P = (F_o^2 + 2F_c^2)/3$



Fig. 1. Perspective view of TESe in a tree-dimensional framework.

structure and the construction of infinite bi dimensional frame works. Distances and angles describing hydrogen bonds are sum marized in Table 4.

3.2. Thermal analysis

In order to obtain information about the thermal decomposition of the synthesized compound, thermogravimetric measurements (TGA) and differential thermal analysis (DTA) were carried out in the temperature range from 293 k to 523 K. The results are reported in Figs. 5 and 6.

The TGA curve of the thermal decomposition reveals that the sample is devoid of water molecules since no mass loss was detected before the decomposition at 420 K. The DTA curve shows

three peaks: the first endothermic peak reaches 335 K, the second transition at 350 K may be attributed to order disorder phase transition and the last exothermic peak observed at 393 K corre sponds to the decomposition of the material. These results suggest interesting electrical properties.

3.3. Electrical properties

3.3.1. Impedance analysis and equivalent circuit

The variation of the Nyquist plot (*Z*" versus *Z*') of the TESe obtained from 223 K to 353 K is represented in Fig. 7. These dia grams exhibit some dispersion instead being a semi circle centered on the real axis, which indicates a Cole Cole empirical behavior [28]. The value of the parameter which characterizes the deviation



Fig. 2. Asymmetric unit of the TESe cristal.



Fig. 3. Atomic arrangement of TESe along the b axis.

from the Debye law is of the order of 0.19. It is observed that temperature is directly proportional to the semi circles radius. Therefore, the increase of the temperature is followed by a decrease of the radius of semi circles and then a decrease in resistance. These observations are an indication of the existence of an activated conduction mechanism [28].

The series of Nyquist representations were modeled by means of Z view software, fitted to an equivalent circuit composed of two elements connected in parallel: resistance R and fractal capacity CPE. In this case, the impedance of the fractal capacitance is. $Z_{(CPE)} = \frac{1}{Q(i\omega)^{\alpha}}$ Where, Q represents the value of capacitance and α indicates the fractal exponent. (Z') and (Z'') were calculated from the following expressions:

Z'
$$\frac{R^2 A_0 \omega^{\alpha} \cos\left(\frac{\alpha \pi}{2}\right) + R \infty}{\left(1 + R A_0 \ \omega^{\alpha} \cos\left(\frac{\alpha \pi}{2}\right)\right)^2 + \left(R A_0 \ \left(\omega\right)^{\alpha} \sin\left(\frac{\alpha \pi}{2}\right)\right)^2}$$
(1)

$$Z'' \quad \frac{R^2 A0\omega^{\alpha} \sin\left(\frac{\alpha \pi}{2}\right) \infty}{\left(1 + RA_0 \ \omega^{\alpha} \cos\left(\frac{\alpha \pi}{2}\right)\right)^2 + \left(RA_0 \ \left(\omega\right)^{\alpha} \sin\left(\frac{\alpha \pi}{2}\right)\right)^2} \tag{2}$$

The variation of the real part (Z') according to the frequency at different temperatures is depicted in Fig. 8. It is observed that the amplitude of Z' is higher in low frequency region at lower tem peratures. Afterwards, Z' decreases progressively with increasing both temperature and frequency. This behavior stands for an



Fig. 4. Two dimensional sheet of the inorganic clusters where groups are linked through hydrogen bonds.

increase in the DC conductivity of the compound. All the curves consolidate at all temperatures and at higher frequencies, hence Z' becomes frequency independent. These observations confirm a possible release of the space charge, subsequently, the lowering of barrier properties [29].

The variation of the imaginary part Z" as a function of frequency at different temperatures is reported in Fig. 9. These curves show that Z" moves to high frequency with increasing temperature until it reaches a maximum (Z"_{max}), then decreases progressively. In fact, the broadening of the peaks and the shift of (Z"_{max}) to higher temperature confirm the existence of relaxation phenomena in the system as well as spreading of the relaxation time [30]. Further more, when frequency decreases, all curves merge. Such behavior indicates a possible release of the space charge mainly H⁺ protons. At high frequency, the space charges have less time to relax, therefore the polarization of the space charge reduces with increasing frequency leading to the fusion.

Fig. 10 exhibits the variation of the real and the imaginary parts of the impedance (Z' and Z'') according the frequency at 323 K. It is clear that Z'' rises with increasing frequency while Z' decreases. This trend remains until Z'' reaches a maximum and intersects with Z'. The maximum value corresponds to a particular frequency p (relaxation frequency). Afterwards, Z' and Z'' decline with increasing frequency until they coincide and reach zero.

Plots reported in Figs. 7–10 indicate a good agreement between experimental and simulated data denoting that the proposed equivalent circuit describes fairly well the electrical properties of the compound.

The dc conductivity was calculated from the following relationship:

$$\sigma(dc) = \frac{e}{R * S}$$

where, S and e represent the area and the thickness of the pellets, respectively.

The variation of specific conductivity (Ln (σ_{dc} T)) versus 10³/T is given in Fig. 11. As it is observed, this curve indicates a rise in the

conductivity of the material with an increase of the temperature. Such behavior indicates that the conduction in the compound is a thermally activated transport mechanism governed by the Arrhe nius law: $\sigma_{dc}T = \sigma_0 \exp(E_a/kT)$, where σ_0 is the pre exponential factor, Ea is the activation energy, k is the Boltzmann constant and T is the absolute temperature. Three regions were identified with changing of slope at $T_1 = 329$ and $T_2 = 338$.

An Arrhenius type law is shown and confirms the two phase transition already observed by DTA measurements at the same temperature. The phase I (T < T₁) was characterized by an increasing conductivity ($\sigma_{dc} > 10^{-4} \ \Omega^{-1} \text{cm}^{-1}$), which obeys the Arrhenius law: $\sigma T = \exp\left(\frac{-Ea}{kT}\right)$. At these temperatures, the slope of log (σ T) versus 10³/T gives an activation energy ΔE_1 0.91 eV. Afterwards, a decrease in the activation energy with increasing temperature from ΔE_1 0.91 eV (T < T₁) to ΔE_2 0.43 eV (T1 < T2) was observed. This anomaly may be attributed to the disorder caused from the reorientation motion of inorganic groups. The phase II ($T_2 > T$) was characterized by an increasing conduc tivity when the temperature increases ($\sigma_{dc} > 10^{-3} \ \Omega^{-1} cm^{-1}$). The Arrhenius law is obeyed and the activation energy increases to ΔE_2 0.61 eV. These high values of conductivity allow us to classify this salt among the super protonic conductors. The experimental values of the conductivity fitted according to the Arrhenius low are depicted in Table 5.

3.3.2. Modulus analysis

The formalism of the electric modulus: M* is recognized as:

- M^* j ω COZ*
- M^* M' + iM''

where, M' and M" are the real and imaginary parts of the complex modulus, obtained according to the following expression:

Table 2

Selected interatomic distances and angles (Å, °) for TESe sample.

	8	(,,)	
Se (1)-01	1.824 (11)	C4 N C1	109.4 (11)
Se (1)-02	1.703 (9)	C7 N C1	95.5 (10)
Se (1)-03	1.549 (14)	C2 N C8	118.4 (12)
Se (1)-04	1.766 (8)	C7 C1 H1A	112.7
Se (2)-05	1.719 (10)	N C1 H1A	112.7
Se (2)-06	1.695 (9)	C7 C1 H1B	112.7
Se (2)-07	1.730 (13)	N C1 H1B	112.7
Se (2)-08	1.784 (11)	H1A-C1-H1B	110.1
Se (3)-09	1.723 (13)	N C2 H2A	107.7
Se (3)-O10	1.688 (16)	C8 C2 H2A	107.7
Se (3)-011	1.688 (16)	H2A-C2-H2B	107.1
Se (3)-012	1.774 (7)	N C2 H2B	107.7
C1 N	1.64 (2)	N C3 H3A	109.5
C2 N	1.563 (14)	C5 C3 H3A	109.5
C3 N	1.488 (14)	N C3 H3B	109.5
C4 N	1.534 (19)	C5 C3 H3B	109.5
C1 C7	1.615 (17)	НЗА-СЗ-НЗВ	108.1
C2 C8	1.668 (15)	C3 C5 H5A	109.5
C3 C5	1.566 (13)	C3 C5 H5B	109.5
C1 H1A	0.9700	H5A-C5-H5B	109.5
C1 H1B	0.9700	C3 C5 H5C	109.5
C2 H2A	0.9700	H5A-C5-H5C	109.5
C2 H2B	0.9700	C2 C8 H8A	109.5
C3 H3A	0.9700	C2 C8 H8B	109.5
C3 H3B	0.9700	H8A-C8-H8B	109.5
C5 H5A	0.9700	C2 C8 H8C	109.5
C5 H5B	0.9600	H8A-C8-H8C	109.5
C5 H5C	0.9600	H8B C8 H8C	109.5
C8 H8A	0.9600	N C4 H4A	107.4
C8 H8B	0.9600	C6 C4 H4B	107.4
C2 N C1	101.4 (12)	N C4 H4B	107.0
C3 N C5	110.8 (11)	H4A-C4-H4B	109.5
C3 N C2	110.1 (12)	C4 C6 H6A	109.5
C3 N C4	111.9 (12)	C4 C6 H6B	109.5
C4 N C2	109.4 (11)	H6A-C6-H6B	109.5
C3 N C1	110.7 (12)	C4 C6 H6C	109.5
C1 C7 H7B	109.5	C1 C7 H7C	109.5
H7A-C7-H7B	109.5	H7A-C7-H7C	109.5
H7B C7 H7C	109.5	06 Se2 05	112.8 (5)
03 Se (1)-01	106.6 (6)	05 Se2 07	107.2 (6)
02 Se (1)-04	112.3 (5)	05 Se2 08	115.7 (3)
O3 Se (1)-O2	115.8 (6)	06 Se2 07	107.1 (6)
04 Se (1)-01	109.4 (5)	06 Se2 08	101.7 (6)
03 Se (1)-04	104.7 (6)	010 Se3 011	114.7 (7)
010 Se3 09	109.3 (7)	010 Se3 012	102.4 (6)
09 Se3 012	113.4 (5)	011 Se3 012	108.3 (5)
011 Se3 09	108.7 (6)	02 Se (1)-01	107.8 (5)

Table 3	
Calculated mean values of the distortion indices.	

	ID (OSeO)	ID (OSeO)	ID (00)
[HSe(1)O ₄]	0.0127	0.0012	0.00008
H ₂ Se(2)O ₄	0.028	0.0002	0.00006
H ₂ Se(3)O ₄	0.0001	0.007	0.0083

Table 4
Geometrical characteristics of hydrogen bonds (Å, °).



Fig. 5. Thermogravimetric analysis plot of TESe.



Fig. 6. DTA curves of $[(C_2H_5)_4N]$ HSeO₄ $(H_2SeO_4)_2$.

M ωC0*Z*

M["] ωC0*Z*['],

where, C₀ presents the vacuum capacitance of the cell.

This formalism has proved to be a forcefulparameter to obtain further information about the relaxation mechanism as well as charge transport processes in the system and optimize the con centration of the charge carriers, mainly H⁺ protons.

In order to verify the contribution of these phenomena, we examined the variation of M' versus (log f) at several temperatures (Fig. 12). At all measured temperatures, curves display a low value of the real part M' at low frequencies. Such observation can be related to a lack of force that governs the mobility of the charge carriers within the influence of the induced electric field. This

D–H A	d (D–H) (Á)	d (H A) (Á)	d (D A) (Á)	$\angle D-H = A(^{\circ})$
O1 H1 O5 (i)	0.82	1.90	2.480 (16)	172
O8 H8 O11 (ii)	0.82	1.86	2.668 (15)	172
O7 H7 O9(iii)	0.82	2.00	2.788 (19)	172
O9 H9 O3	0.82	1.97	2.784 (17)	170
O12 H12 O2	0.82	1.95	2.759 (13)	162

Symmetry codes: (i) *x*, *y*-1, *z*; (ii) *x*, -*y*+1, *z*-1/2; (iii) *x*, -*y*, *z*-1/2.



Fig. 7. Nyquist plot (-Z" versus Z') of TESe from 290 K to 373 K.



Fig. 8. Fluctuation of the real part (Z') according to the frequency at different temperatures for TESe.



Fig. 9. Plot of the imaginary part (Z') according to the frequency at different temperatures.

behavior promotes the long range mobility of H⁺ protons. Furthermore, the rises of the value of M' until it reaches a constant value $\left(M'\infty \quad \frac{1}{\epsilon'\infty}\right)$ may be assigned to a conduction mechanism ascribed to the short range mobility of H⁺ protons.

The 3D and 2D presentations of the imaginary part of modulus normalized (M''/M_{max}) as a function of frequency at different temperatures is depicted in Fig. 13. In fact, these curves show a well defined asymmetric peak which shifts to higher Frequencies,



Fig. 10. Variation of the real and the imaginary parts of impedance (Z' and Z'') according the frequency at 323 K.



Fig. 11. Variation of specific conductivity (Ln ($\sigma_{dc}T$)) versus 10³/T for TESe compound.

Table 5The experimental values of the conductivity for TESe sample.

T (K)	$\sigma_{dc} (10^{-4}\Omega^{-1} \text{cm}^{-1})$	Ea (ev)
290	3.09	
308	4.12	
313	4.57	0.91
318	7.79	
323	6.51	
328	6.54	
338	6.95	
338 T (K)	6.95 $\sigma_{dc}(10^{-3}\Omega^{-1}cm^{-1})$	Ea(ev)
338 T (K) 343	6.95 σ _{dc} (10 ⁻³ Ω ⁻¹ cm ⁻¹) 1.13	Ea(ev)
338 T (K) 343 348	6.95 σ _{dc} (10 ³ Ω ¹ cm ¹) 1.13 1.29	Ea(ev)
338 T (K) 343 348 353	6.95 σ _{dc} (10 ³ Ω ¹ cm ¹) 1.13 1.29 1.37	Ea(ev) 0.61
338 T (K) 343 348 353 358	6.95 σ _{dc} (10 ³ Ω ¹ cm ¹) 1.13 1.29 1.37 1.45	Ea(ev) 0.61
338 T (K) 343 348 353 358 363	6.95 σ _{dc} (10 ³ Ω ¹ cm ¹) 1.13 1.29 1.37 1.45 1.51	Ea(ev) 0.61

as the temperature increases. These peaks (M" $_{max}$) can be associ ated to characteristic relaxation rate and exhibit asymmetric fre quency dependence. At low frequency region, the H⁺ protons are mobile over long distances. At high frequencies, H⁺ protons are confined to their potential wells, being mobile over short distances. However, the region of transition from a long distance to short



Fig. 12. Variation of M' according to (log f) at several temperatures.



Fig. 13. Frequency dependence of the imaginary part of modulus normalized (M''/M_{max}) at various temperatures.

distance mobility is close to the asymmetric peak [31]. Moreover, it is revealed that the maximum of the asymmetric peak (M''_{max}) of modulus curves correspond to the bulk of relaxation frequency fp ($fp = 2\pi\tau_{max}$). On the other hand, this frequency is defined by the relation: $\tau_{max}\omega_{max}$ 1, where τ_{max} presents the most probable proton relaxation time [32].

The electric relaxation process is well described, in general, by the KWW decay function (numerical Laplace transform of the Kohlrausch– ψ illiams_b- ψ atts) presented as.

 $\varphi(t) \exp \left| \left(\frac{t}{\tau_0} \right) \right|$ [33]. The exponent β presents the Kohl rausch parameter and characterizes the degree of non Debye behavior (1.14 decades). This parameter ranging between [0–1] is found to be 0.75.

In Fig. 14, the variation of Ln ($\sigma_{dc}T$) and log (fp) with the inverse of the temperature. The conductivity Ln ($\sigma_{dc}T$) and the relaxation



Fig. 14. Variation of Ln $(\sigma_{dc}T)$ and log $({\it fp})$ according to the inverse of temperature for TESe.

frequency $\log (fp)$ are depicted by two quasi-parallel line segments and they present a discontinuity at 329 and 338 K. These results confirm the two phase transitions observed in DTA analysis and conductivity measurements.

3.3.3. Ac electrical conductivity

The electric conductivity presents key approach in term of dis playing information about electrical dynamic including conduc tivity, capacitance and loss factor in the sample. It is also useful in identifying the nature of the conduction.

The Ac conductivity is derived from impedance complex data according to the following formula:

$$\sigma_{ac} \quad \left(\frac{e}{s}\right) \left(\frac{Z'}{Z'^2 + Z''^2}\right)$$

The variation of the ac conductivity according to the frequency is shown in Fig. 15. It is noteworthy that, the ac conductivity is almost constant at low frequency. Then, it exhibits dispersion at high fre quency. Consequently, the rises in frequency ameliorate the jumping frequency of charge carriers, leading to a rise of the con duction process and then an increase of the conductivity. In addi tion, this feature denotes the presence of a thermally activated process generated from an increase pertaining to the energy of the charge carriers. At lower frequencies, the applied electric field compels the charge carriers to move over large distance [34]. Thus, the mean movements of the charge carriers are reduced with the increasing frequency.



Fig. 15. Variation of the ac conductivity according to the frequency for different temperatures.

3.4. Dielectric studies

Dielectric relaxation is an alternative approach to reveal reliable information about mechanisms of conduction, the origin of the dielectric losses, and to characterize phase transition detected in the ATD curve at 396 K. The dielectric relaxation is described by a non Debye model which can be expressed as:

$$\varepsilon^* \quad \varepsilon' + i\varepsilon''$$

where ε_{∞} is the permittivity at high frequency value of ε_{1} , τ_{1} is the relaxation time, ε_{S} presents the static permittivity, ω is the angular frequency and α is the tilting angle of the circular arc from the real axis in the complex permittivity plane [35].

Real and imaginary parts of the electric permittivity were calculated according to the following relations:

$$\varepsilon' \quad \frac{Z''}{\omega C_0 \left(Z'^2 + Z''^2\right)}$$
$$\varepsilon'' \quad \frac{Z'}{\omega C_0 \left(Z'^2 + Z''^2\right)}$$

Figs. 16 and 17 presents the temperature dependence of (ε') and (ε'') (real and imaginary parts) of the electric permittivity for TESe. It can be noted that, at low temperature, the fluctuation of ε' and ε'' with temperature is almost constant. This behavior can be explained by limited reorientation movements of the inorganic groups ([HSeO₄] and H₂SeO₄) which are unable to orient them selves to the direction of the applied electric field. Hence, they acquire a feeble contribution to the polarization [36]. Afterwards, an asymmetric loss peak was detected at 333 K, this anomaly is may be due to a disorder in the lattice generated from the reorientation motion of inorganic chains which reduces the Van Der Waals in teractions between the ions that leave to another unit cell let behind a vacancy [37,38]. This anomaly arises when the hopping frequency of localized charge carriers becomes approximately even to that of the applied electric field [39]. At high temperature, all curves of (ε') and (ε'') increase until they reach a large maximum approximately at 359 K, featuring the phase transition, and then decrease at all measured frequencies. Over the phase transition, the units have enough thermal energy to adapt more easily to the change of the external electric field. This behavior can be elucidated by a weakness of the Van Der Waals interaction between ions as a result of the disorder and the inorganic groups reorientation motion.

The ratio of the loss factor ε''_r divided by the actual permittivity ε'_r is defined as the tangent losses (tan δ), this parameter is described as:

$$\tan \delta = \frac{\varepsilon'' r(\omega)}{\varepsilon' r(\omega)}$$

Fig. 18 shows the fluctuation of the dissipation factor $(\tan \delta)$ as a function of temperature for different frequencies. At low temper ature, the variation of $\tan \delta$ is almost constant until it reaches an asymmetric loss peak at 333 K, confirming the limited reorientation motion and then reorientation motions of these groups. At high temperature, curve rises within the whole frequency range, pre sents a wide maximum and then decreases at 363 K. Consequently, this behavior can be explained by the interaction weakness be tween ions induced by the disorder and the reorientation move ment of the inorganic groups.

Fig. 19 presents the variations of the loss factor (ε''_r) depending to the frequency at different temperatures. As observed, ε''_r is high at low frequency then it decreases with rising frequency. At low frequency region, the dielectric loss increases clearly with tem perature. Then, it indicates a dispersive behavior owing to the po larization of space charge, explained through the spreading and accumulation of the proton close to the electrode. As frequency rises, the fluctuation of the applied field becomes too fast and the dipoles will be unable to attend it. After, ε''_r approaches approxi mately to values close to zero [40].

The fluctuation of the dissipation factor $(\tan \delta)$ versus frequency is depicted in Fig. 20. These plots show a similar behavior to that of the variation of the loss factors (ε''_r) which, decreases progressively with increasing both temperature and frequency. Such behavior confirms the presence of the dielectric relaxation phenomenon in the sample.

The dielectric study shows clearly two phase transitions observed in the title compound. The first can be assigned to a



Fig. 16. Temperature dependence of the electric permittivity in TESe.



Fig. 17. Variation of the dielectric loss according to the temperature for different frequencies.



Fig. 18. Fluctuation of the dissipation factor (tan δ) as a function of temperature for different frequencies.

disorder generated from the reorientation motion of the inorganic groups and the second may be attributed to order disorder phase transition stressed by the diffusion of the H⁺ protons.

3.5. Investigation of the proton conduction mechanism

The diffusion of protons through the chemical species is described within the framework of the vehicle and Grotthuss mechanisms [19,20]. In the first mechanism, proton conductivity is achieved through the migration of the vehicle or proton solvent. Clearly, the vehicle diffusion rate is highly related to the overall proton conductivity. In contrast, in the Grotthuss mechanism (known also as proton jumping) proton solvent or vehicle are not necessary and proton conductivity is achieved by transfer of pro tons from one site to another via the formation and breakage of hydrogen bonds.

In the case of TESe, the SeO₄ tetrahedral are donor and acceptor of charges. The phase transition observed is closely related to the rotation or liberation of the anionic groups when heating. Thus, the movement of the charge carriers is closely related to the tetrahedral dynamic. The increase in SeO₄ groups libration with the rise of temperature allow to lengthen short inter atomic distances, the hydrogen bonds breaking, H⁺ protons movement and the anionic groups reorientation. Consequently, the process of charge transport follows the Grotthuss mechanism **[20]**.



Fig. 19. The variations of the loss factor ε''_r depending on the frequency at different temperatures for the title compound.



Fig. 20. Fluctuation of the dissipation factor (tan $\delta)$ versus frequency for some temperature.

4. Conclusion

In summary, a novel selenate hybrid compound based on Tet raethylammonium with the general formula [(C₂H₅)₄N] HSeO₄ $(H_2SeO_4)_2$ was prepared and crystallized at room temperature by the slow evaporation technique.

Crystallographic studies using Single crystal X ray diffraction revealed that TESe belongs to the monoclinic symmetry with the Cc space group. The asymmetric unit contains one independent $[(C_2H_5) N^+]$ cation, one $[HSeO_4]^{2-}$ and two neutral selenic acid H₂SeO₄. The structure consists of a layered organization of organic and inorganic sheet along the b axis and the inorganic groups are interconnected through strong O-H···O hydrogen bond yielding infinite 2D planes.

Thermal measurements were carried out revealing the appear ance of two phase transitions which is subsequently confirmed by the electrical and dielectric analysis. Complex impedance analyses over temperature and frequency allow us to determine the equiv alent circuit and electric measurements reveal that this material is a super protonic conductor. Dielectric analysis indicates that the title compound presents a dielectric relaxation phenomenon.

Finally, the transport process results from the jumping of H⁺ protons between SeO₄ groups. Proton conduction can be ascribed to the transfer through hydrogen bonds and movements of charges carriers, and then it arises by the Grotthuss mechanism. Therefore, these results appear to be an interesting route to develop a new proton conducting compounds for specific applications in electro lytes and batteries.

Declaration of competing InterestCOI

The authors declare that they have no known competing financialinterestsor personal relationships that could have appeared to influence the work reported in this paper.

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Appendix A. Supplementary data

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