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Gilles Daigre, Patrick Gougeon, Philippe Gall, Odile Merdrignac-Conanec, Rabih Al Rahal Al Orabi, et al.. Unravelling the Beneficial Influence of Ag insertion on the Thermoelectric Properties of the Cluster Compound K2Mo15Se19. ACS Applied Energy Materials, 2020, 3 (3), pp.2846-2855. 10.1021/ac-saem.9b02488 . hal-02569139

HAL Id: hal-02569139 https://hal.science/hal-02569139

Submitted on 27 Nov 2020

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Unravelling the beneficial influence of Ag insertion on the thermoelectric properties of the cluster compound K₂Mo₁₅Se₁₉

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ABSTRACT

Mo-based cluster compounds are an emerging class of potential candidates for thermoelectric applications due to several key crystallographic characteristics that drive their ability to transport heat close to that of amorphous systems. Here, we report on a detailed investigation of the crystal structure and high-temperature transport properties of the cluster compound K₂Mo₁₅Se₁₉ and of its Ag-filled variant Ag₃K₂Mo₁₅Se₁₉. Single-crystal X-ray diffraction confirm that both compounds crystallize with a hexagonal crystal structure, successfully solved in the space group $R\bar{3}c$, built by octahedral Mo₆Se₈ and bioctahedral Mo₉Se₁₁ clusters with K and Ag cations filling the large inter-cluster voids. In agreement with electronic band structure calculations, the electron transfer that occurs from the filling cations to the cluster subnetwork provides a simple guiding rule to predict the evolution of the transport properties upon filling

with Ag. The metallic state observed in K₂Mo₁₅Se₁₉ is turned into a more heavily-doped semiconducting character in the presence of Ag resulting in enhanced thermopower values in Ag₃K₂Mo₁₅Se₁₉. The beneficial influence of Ag on the thermoelectric properties is also reflected by lowered lattice thermal conductivity values that reach 0.45 W m⁻¹ K⁻¹ at 750 K. These results indicate that inserting an additional element in the ternaries M_2 Mo₁₅Se₁₉ is an interesting route to improving their thermoelectric properties.

Keywords: Thermoelectric; molybdenum clusters; semiconductors; thermal conductivity; single crystal

INTRODUCTION

In 1971, R. Chevrel, M. Sergent and J. Prigent discovered the "Chevrel phases" of generic chemical formula M_x Mo₆ X_8 (M = rare-earth, transition metals, Na, Ca, Sr, Ba, Sn, Pb; X = S, Se, Te), the crystal structure of which is built up by octahedral Mo₆ clusters.¹ These compounds exhibit many interesting physical properties such as high-field superconductivity in PbMo₆S8 or the coexistence of superconductivity and ferromagnetic order in HoMo₆Se8.² The inherent metallic character of the electrical conduction in these compounds originates from the nature of the top of the valence bands, the presence of inter-cluster interactions and the flexibility of the electron count, that is, of the charge transfer between the Mo₆ clusters and the *M* cations.

Following these pioneering works, condensed molybdenum-based clusters Mo₉ and Mo₁₂ were obtained by increasing the metal-to-chalcogen ratio.³⁻⁵ The condensation process is realized by sharing Mo₃ opposite triangular faces of the Mo₆ octahedra whose ultimate stage leads to the uniaxial infinite chain of face-faced octahedra $|Mo_{6/2}|_{\infty}$ observed in the quasi-1D $M_2Mo_6X_6$ compounds (M = Na, Rb, Cs, Tl, In and X = S, Se, Te).⁶ Of note, the nearly 1D character of the crystal structure of these compounds gives rise to exotic superconductivity and possibly to novel types of fermionic excitations due to topologically non-trivial electronic band structures.⁷⁻¹⁰ So far, the intermediate Mo₉, Mo₁₂, Mo₁₅, Mo₁₈, Mo₂₁, Mo₂₄, Mo₃₀ and Mo₃₆ clusters have been observed in various compounds.^{3-5,11-13}

However, the thermoelectric properties of some of these ternary and quaternary molybdenum chalcogenides have only been assessed recently.¹⁴⁻²³ Their inherent structural complexity, arising from a high number of atoms per unit cell and the presence of cations in the inter-cluster voids, is the key ingredient leading to very low lattice thermal conductivity κ_{ph} . Combined with electrical properties that can gradually evolve from metallic to semiconducting-like by inserting additional cations, this property can potentially yield high dimensionless

thermoelectric figures of merit $ZT = \alpha^2 T / \rho(\kappa_e + \kappa_{ph})$ which determine the efficiency of a thermoelectric material at an absolute temperature T.¹⁴⁻²³ The main challenge then lies in achieving concomitantly high thermopower α , low electrical resistivity ρ , and low electronic thermal conductivity κ_e .

First studies performed on the Chevrel phases revealed their potential for hightemperature thermoelectric applications with a maximum *ZT* of 0.6 at 1150 K for Cu₂FeMo₆Ses, rivaling those obtained in the state-of-the-art Si-Ge solid solutions.^{24,25} Later on, initial investigations of the thermoelectric properties of the Ag_xMo₉Se₁₁ ($3.4 \le x \le 3.8$) compounds revealed equivalent *ZT* values at 800 K (*ZT* of 0.65 for x = 3.8) that notably arise from their very low lattice thermal conductivity approaching the amorphous limit at 800 K (0.5 W m⁻¹ K⁻¹).^{14,15,21} The physical origin of this glass-like thermal transport is related to the large voids separating the Mo₉Seⁱ₁₁Se^a₆ bioctahedral clusters in which Ag⁺ cations occupy four independent crystallographic sites and show large anisotropic thermal displacement parameters. The polymorph compounds Ag₂Tl₂Mo₉Se₁₁ and Ag₃RbMo₉Se₁₁ that crystallize in the hexagonal space groups $R\bar{3}c$ and $P6_3/m$, respectively, also combine very low κ_{ph} values and semiconducting behavior yielding peak *ZT* values of 0.6 and 0.4 at 800 K, respectively.^{17,22}

The M_2 Mo₁₅Se₁₉ (M = K, Ba, In, Tl, Rb) compounds are another family investigated for their thermoelectric properties.^{16,19,23} Their crystal structure contains both Mo₆Seⁱ₈Se^a₆ and Mo₉Seⁱ₁₁Se^a₆ cluster units in equal proportion that delineates large voids where the M cations reside. In agreement with electronic band structure calculations,²⁶ metallic electron counts (MEC) indicate that these materials are metallic, assuming a monovalent character of the Mcations. Driving these metallic properties towards semiconducting-like properties can be realized in filled variants by inserting additional cations such as Ag. This strategy has been successfully applied for M = In and Tl resulting in higher thermoelectric performances than in Ag-free compounds with peak ZT values of 0.45 and 0.20 at 1100 and 800K, respectively.^{16,23} Following these encouraging results, we explore herein the transport properties of the isostructural compounds $K_2Mo_{15}Se_{19}$ and $Ag_3K_2Mo_{15}Se_{19}$ to further elucidate the influence of the nature of the *M* cations on the thermoelectric properties. The experimental results are compared with those reported on the In and Tl compounds outlined above and complemented by calculations of the electronic density of states.

RESULTS AND DISCUSSION

The two compounds K₂Mo₁₅Se₁₉ and Ag₃K₂Mo₁₅Se₁₉ are isostructural to In₂Mo₁₅Se₁₉ and Ag₃In₂Mo₁₅Se₁₉, respectively, that both crystallize in the space group $R\bar{3}c$ (Tables 1 and 2). Their crystal structures are based on the two cluster units Mo₆Se₈Se₆ and Mo₉Se₁₁Se₆, the threedimensional arrangement of which is shown in Figure 1. These two units share ligands via the Mo1-Se₂ and Se₁-Mo₂ bonds to ensure the three-dimensional cohesion of the crystal structure (Figure 2). Compared to the K₂Mo₁₅Se₁₉ compound, the insertion of Ag atoms results in distorted Se sites forming trigonal bipyramids located between two consecutive K sites and in a modification of the K environment (Figures 3a and 3b). A clear shift of the K atoms towards the Se₅ atoms is indeed observed with a significant decrease in the distance K-Se₅ from 3.6760(15) Å in K₂Mo₁₅Se₁₉ to 3.260(2) Å in Ag_{2.2}K₂Mo₁₅Se₁₉. In contrast, the Se₃ and Se₄ atoms move away from K with K-Se₃ and K-Se₄ distances of 4.062(1) Å and 3.674(2) Å in Ag_{2.2}K₂Mo₁₅Se₁₉ instead of 3.7891(7) Å and 3.4256(15) Å in K₂Mo₁₅Se₁₉.

Ag insertion leads to variations in the Mo-Mo distances in both clusters due to the increase in the cationic charge transfer. In the Mo₆ cluster (Table 3), a shortening of the inter-triangular distance Mo1-Mo1 from 2.7675(3) Å in K₂Mo₁₅Se₁₉ to 2.7146(4) Å in Ag_{2.2}K₂Mo₁₅Se₁₉ is observed upon increasing the cationic charge. In contrast, the intra-triangular Mo1-Mo1 distances remain practically unchanged. In the Mo₉ cluster, in addition to the decrease in the inter-triangular Mo2-Mo3 distances (2.7138(2) Å to 2.6873(2) Å), the intra-angular Mo3-Mo3 distance increases from 2.6821(3) Å in K₂Mo₁₅Se₁₉ to 2.7158(5) Å in Ag_{2.2}K₂Mo₁₅Se₁₉. This modification is accompanied by a shortening of the intra-triangular Mo2-Mo2 distances from 2.6682(3) Å to 2.6424(4) Å. All these variations in the interatomic distances are consistent with the theoretical calculations (see below) carried out on the host structure K₂Mo₁₅Se₁₉.

The PXRD patterns of both polycrystalline specimens, shown in Figure 4, were successfully indexed by the structure model derived from single-crystal X-ray diffraction data. These data indicate the absence of secondary phases within experimental uncertainty. Rietveld refinements against the diffraction data for Ag₃K₂Mo₁₅Se₁₉ showed that the occupancy factor of the Ag atom is 0.493(9), in very good agreement with the nominal chemical composition that corresponds to an occupancy factor of 0.5. The crystallographic parameters determined from these refinements are given in Table S1 in Supporting Information. The atomic coordinates, equivalent isotropic displacement parameters and selected interatomic distances are listed in Table S2 and S3, respectively.

The good chemical homogeneity of the two samples is further confirmed by elemental Xray mapping (Figures S4 and S5 in Supporting Information) indicating that Ag and K are homogeneously distributed. In the Ag₃K₂Mo₁₅Se₁₉ compound, these experiments further indicate the presence of very small Ag-rich or K-rich areas, suggesting that local, weak compositional fluctuations in the Ag/K ratio exist in this sample. The possibility that the overall cation content slightly varies spatially is in agreement with the single-crystal X-ray diffraction data evidencing an under-stoichiometric Ag concentration in the single-crystals studied. Nevertheless, as shown by these images, these Ag- or K-rich areas are very small and are thus not expected to significantly impact the measured transport properties. Fracture images collected on both samples evidence some differences in the microstructure of both samples (Figure S6 in Supporting Information). While the Ag₃K₂Mo₁₅Se₁₉ compound is composed of micron-sized particles homogeneously distributed, the Ag-free compound is rather characterized by a coexistence of similar large grains with smaller ones, suggesting a broader size distribution than in Ag₃K₂Mo₁₅Se₁₉. The corresponding elemental X-ray mappings (Figures S7 and S8 in Supporting Information) show that this difference in the microstructure does not affect the homogeneous distribution of Ag and K in both polycrystalline samples.

The metallic electron count (MEC), that is, the number of electrons available for the formation of Mo-Mo bonds, optimizes metal-metal bonding in the Mo₉Se₁₁ and Mo₆Se₈ clusters when equal to 36 and 24, respectively.¹⁶ Assuming the absence of significant interaction between these clusters, the optimum MEC for the Mo15Se19 subnetwork is thus equal to 60. The electronic DOS and COHP curves calculated for different Mo-Mo bonds confirm this analysis. The Fermi level is located almost at the top of the valence bands, predominantly centered on the Mo atoms. The small contribution of K states to the occupied bands confirms the ionic interaction with the transfer of 4s electrons from K to the Mo₁₅Se₁₉ subnetwork (Figure 5). The COHP curves calculated for the Mo-Mo bonds in the Mo6 and Mo9 clusters show that the vacant bands at the top of the valence bands have mostly a weakly bonding or weakly anti-bonding character (Mo3-Mo3). Hence, these vacant bands can be considered as non-bonding as a first approximation. The COHP for the Mo1-Mo2 inter-cluster, shown in Figure 6, confirms the weak interaction between these patterns. A band gap of about 0.5 eV separates the valence bands from the anti-wear strips. Within a rigid-band model (that is, independent of the electron count), an increase in the MEC up to 60 should therefore be possible without significantly altering the architecture of K₂Mo₁₅Se₁₉.

For Ag₃K₂Mo₁₅Se₁₉, each inserted Ag atom transfers its valence electrons to the Mo-Se framework which increases the overall MEC of the Mo₁₅Se₁₉ clusters to 57. The variations in the Mo-Mo distances observed in single-crystal X-ray diffraction data (Table 3) are in agreement with those expected in view of the COHP curves, considering the occupation of the

bands located just above the Fermi level in K₂Mo₁₅Se₁₉. Indeed, the observed shortening of the distances Mo1-Mo1, Mo2-Mo2 and Mo2-Mo3 is consistent with the Mo-Mo bonding character that binds bands above the Fermi level in K₂Mo₁₅Se₁₉ (Figures 7 and 8). In contrast, the antibonding character of the Mo3-Mo3 and the intra-triangular Mo1-Mo1 bonds explains the observed increase in these distances upon inserting Ag.

According to the above-mentioned considerations of the MEC in these two compounds, the insertion of Ag into K₂Mo₁₅Se₁₉ should drive this compound towards a more heavily-doped semiconducting state. This expectation should therefore be synonymous with an increase in the thermopower α in Ag₃K₂Mo₁₅Se₁₉ with respect to K₂Mo₁₅Se₁₉. A comparison of the temperature dependences of α , shown in Figure 9a, confirms this simple picture with a significant increase in the α values in Ag₃K₂Mo₁₅Se₁₉ over the entire temperature range. However, a concomitant increase in ρ that may be expected due to a decrease in the carrier concentration is not observed (Figure 9b). In contrast, $Ag_3K_2Mo_{15}Se_{19}$ shows lower ρ values compared to K₂Mo₁₅Se₁₉, resulting in higher power factor compared to the Ag-free compound (Figure S9 in supporting Information). This trend is not due to differences in the relative density of the two polycrystalline samples since both consolidated pellets exhibit similar geometric densities. This difference may be related to the different microstructures of the two samples evidenced by fracture images. The broader distribution of grain size suggested by these images in K₂Mo₁₅Se₁₉ may result in lower carrier mobilities resulting in increased ρ values. Alternatively to microstructure effects, a decrease in the electron-phonon coupling strength induced by the addition of Ag could also explain higher hole mobilities in Ag₃K₂Mo₁₅Se₁₉ compared to K₂Mo₁₅Se₁₉, as suggested recently in the Cs₂Mo₁₂Se₁₄ cluster compound upon Cu insertion.¹⁸ A detailed analysis of this possibility based on Hall effect measurements is however complicated in the present compounds due to the high hole concentrations giving rise to small Hall signals difficult to resolve experimentally, should it be measured at low or high temperatures. Indeed, as documented for several other Mo-based cluster compounds, the Hall signal is often composed of electron-like and hole-like contributions which tend to cancel out.¹⁴⁻

The total thermal conductivity κ , shown in Figure 9c, exhibits a similar linear temperature dependence in both compounds, varying upon warming from 1.0 and 1.2 W m⁻¹ K⁻¹ at 300 K to 1.6 and 2.0 W m⁻¹ K⁻¹ at 800 K in the Ag₃K₂Mo₁₅Se₁₉ and K₂Mo₁₅Se₁₉ compounds, respectively. The electronic thermal conductivity κ_e has been estimated using the Wiedemann-Franz law $\kappa_e = LT/\rho$ where *L* is the Lorenz number (Figure S10 in Supporting Information). In order to take into account possible variations in *L* in this temperature range, the temperature dependence of *L* has been calculated assuming a single-parabolic band model with acoustic phonon scattering. The subtraction of κ_e from κ leaves the lattice contribution κ_{ph} , shown in Figure 9c as a function of temperature for both compounds. The insertion of Ag leads to a substantial reduction in the heat transport from 1.05 W m⁻¹ K⁻¹ in K₂Mo₁₅Se₁₉ to 0.45 W m⁻¹ K⁻¹ in Ag₃K₂Mo₁₅Se₁₉ at 800 K. This strong decrease in κ_{ph} is a result of an increased complexity in the unit cell and of the probable presence of additional low-energy optical modes related to the large thermal motion of the Ag atoms. The addition of Ag helps to bring the heat transport down to its theoretical minimum estimated by the model developed by Cahill and Pohl where the high-temperature limit of the minimum thermal conductivity κ_{min} is expressed as²⁷

$$\kappa_{min} = \frac{1}{2} \left(\frac{\pi}{6}\right)^{1/3} k_B V^{-2/3} (2v_T + v_L) \tag{1}$$

In this expression, k_B is the Boltzmann constant, V is the atomic volume and v_T and v_L are the transverse and longitudinal sound velocities. In the present case, room-temperature measurements of v_T and v_L by a pulse-echo method yields $\kappa_{min} \approx 0.45$ W m⁻¹ K⁻¹ for both compounds.

The temperature dependence of the dimensionless thermoelectric figure of merit ZT, shown in Figure 9d, evidences the overall beneficial influence of Ag on the thermoelectric

properties of K₂Mo₁₅Se₁₉. The concomitant increase in the power factor α^2/ρ and decrease in κ_{ph} driven by the insertion of Ag results in a nearly eightfold increase in the *ZT* values with a peak value of 0.18 at 800 K. This maximum value is similar to that achieved in the Tl-containing analogue (0.19 at 800 K) but lower than that obtained in the In-containing compound (0.26 at 800 K).^{16,23}

Although all these compounds seem similar according to their respective MEC, these results show that the nature of the transport properties are sensitive to the nature of the M cation. The fact that these different elements are not electronically equivalent in these compounds may have several origins. First, the MEC count relies on the assumption that the M cations exhibit a +1 valence state. In the Ag_xMo₉Se₁₁ compounds, combined high-resolution fluorescencedetection X-ray absorption spectroscopy (HERFD-XAS) and resonant inelastic X-ray scattering (RIXS) measurements have evidenced the presence of sub-monovalent Ag atoms.²⁸ The presence of a similar property in the $Ag_3M_2Mo_{15}Se_{19}$ compounds would explain the differences observed in their transport properties. In addition, the actual charge of the Ag and M atoms in these structures may likely vary depending on the nature and atomic size of the M cations and on the actual concentrations of both Ag and M. Alternatively, these differences may be due to slight but systematic variations in the valence bands as a function of the M cations. In particular, these modifications might change the balance between the electron- and hole-like contributions usually observed in the Hall signal which, in turn, would alter their impact on the thermopower values and hence, on the ZT values. Because both mechanisms could be at play in these compounds, further experimental investigations with HERFD-XAS and RIXS would be valuable to systematically probe the valence state of the Ag and M cations.

CONCLUSION

The novel reduced molybdenum selenides K2M015Se19 and Ag3K2M015Se19 were synthesized by solid-state reaction at 1300°C in sealed Mo crucibles. As in In₂Mo₁₅Se₁₉, the crystal structure contains an equal proportion of octahedral Mo₆ and bioctahedral Mo₉ clusters, forming Mo₆Seⁱ₈Se^a₆ and Mo₉Seⁱ₁₁Se^a₆ cluster units, separated from each other by large voids. These voids are filled up by eleven-coordinated K atoms in K₂Mo₁₅Se₁₉ and by five-coordinated Ag atoms in Ag₃K₂Mo₁₅Se₁₉. In agreement with simple metallic electron count and electronic band structure calculations, K₂Mo₁₅Se₁₉ is a *p*-type metal with moderate thermopower values. Upon inserting Ag, a clear evolution of these properties towards a heavily-doped semiconducting state is observed. In addition to influencing the electronic properties, the presence of Ag further lowers the phonon transport close to the amorphous limit. Very low-temperature thermodynamic measurements and spectroscopic probes will be of interest to establish in detail the microscopic origin of this behavior. The comparison of the transport properties measured herein with those reported for the In and Tl-containing analogues suggests that the nature of M cation is an important parameter to tune the thermoelectric performances. Future investigations of the transport properties of these compounds with other M atoms will be helpful to better understand its role and to develop structure-properties relationships in this family of cluster compounds.

Supporting Information.

Experimental and computational details, details of the crystallographic data collection and powder X-ray structural analysis, final atomic coordinates and equivalent isotropic displacement parameters, selected interatomic distances for polycrystalline Ag₃K₂Mo₁₅Se₁₉ determined from profile matching against the PXRD data, temperature dependence of the thermal diffusivity for both samples, temperature dependence of the thermal conductivity κ measured perpendicular (low temperature) and parallel (high temperature) to the pressing direction for K₂Mo₁₅Se₁₉ and Ag₃K₂Mo₁₅Se₁₉, temperature dependence of the specific heat C_p for K₂Mo₁₅Se₁₉ and Ag₃K₂Mo₁₅Se₁₉, elemental X-ray mapping for K₂Mo₁₅Se₁₉ and Ag₃K₂Mo₁₅Se₁₉, elemental X-ray mapping for K₂Mo₁₅Se₁₉, elemental fracture images and their corresponding elemental X-ray mapping for K₂Mo₁₅Se₁₉ and Ag₃K₂Mo₁₅Se₁₉, temperature dependence of the power factor, temperature dependence of the electronic thermal conductivity.

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Notes

The authors declare no competing financial interest.

Conflicts of interest

There are no conflicts to declare.

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Cluster Compounds Ag_xMo₉Se₁₁ (*x* = 3.4, 3.9). *ACS Appl. Energy Mater.* **2018**, *1*, 4032–4039.

Tables

Table 1. Single-crystal data collection and structure refinements of $K_2Mo_{15}Se_{19}$ and

$Ag_{2.2}K_2M$	015Se19 at	300 K.
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Empirical chemical formula	K2M015Se19	Ag2.2K2M015Se19	
Molar mass (g.mol ⁻¹)	3017.54	3251.76	
Symmetry	Trigonal		
Space group	R3c		
<i>a</i> (Å)	9.7229(10)	9.96150(10)	
<i>c</i> (Å)	58.2467(8)	57.5198(9)	
V (Å ³)	4768.63(10)	4943.08(10)	
Ζ	6		
d (g.cm ⁻³)	6.305	6.558	
Crystal color and habit	black, irregular	black, irregular	
Crystal size (mm ³)	$0.142 \times 0.094 \times 0.071$	$0.205 \times 0.117 \times 0.108$	
Radiation	Μο Κα	Μο Κα	
θ range (°)	3.70-44.56	3.69-39.996	
Absorption coefficient (mm ⁻¹)	27.705	27.992	
Limiting indices	-19≤ <i>h</i> ≤17, -18≤ <i>k</i> ≤19, -114≤ <i>l</i> ≤95	-18≤h≤17, -18≤k≤16, -103≤l≤103	
Reflections collected/unique	57775 / 4288	45416 / 3410	
R (int)	0.0776	0.0807	
Absorption correction	analytical	analytical analytical	
Max./min. transmission	0.2811 / 0.1219	0.1689 / 0.0376	
Data/restraints/parameters	4288 / 0 / 57	3410 / 0 / 67	
<i>R</i> indices $[I > 2\sigma(I)]$	R1 = 0.0310, wR2 = 0.0520	R1 = 0.0296, wR2 = 0.0485	
Extinction coefficient	0.000053(3)	0.000138(4)	
Largest diff. peak and hole $(e^{A^{-3}})$	1.795 and -2.192	2.584 and -2.170	
Goodness-of-fit on F^2	1.058	1.140	

Atom	x	у	Z	$\overline{U}_{ m eq}$	s.o.f
Mo1	0.64652(2)	-0.51856(2)	0.35303(2)	0.00774(4)	1
Mo2	0.84743(2)	-0.16372(2)	0.28925(2)	0.00774(4)	1
Mo3	0.84074(3)	0.0000	0.2500	0.00727(5)	1
Se1	0.62840(3)	-0.37968(3)	0.31684(2)	0.00935(4)	1
Se2	0.65235(3)	-0.30839(3)	0.38094(2)	0.01032(5)	1
Se3	0.69207(3)	-0.30793(3)	0.2500	0.01105(6)	1
Se4	0.6667	-0.6667	0.38756(2)	0.01267(8)	1
Se5	1.0000	0.0000	0.32386(2)	0.01081(7)	1
K	1.0000	0.0000	0.38697(2)	0.0314(3)	1
		•			·]
Atom	x	У	Ζ.	$U_{ m eq}$	s.o.f
Atom Mo1	x 0.65070(3)	y -0.51996(3)	z 0.35272(2)	U _{eq} 0.00546(4)	s.o.f
Atom Mo1 Mo2	x 0.65070(3) 0.85270(3)	y -0.51996(3) -0.15840(3)	z 0.35272(2) 0.28920(2)	U _{eq} 0.00546(4) 0.00581(5)	s.o.f 1 1
Atom Mo1 Mo2 Mo3	x 0.65070(3) 0.85270(3) 0.84259(3)	y -0.51996(3) -0.15840(3) 0.0000	z 0.35272(2) 0.28920(2) 0.2500	Ueq 0.00546(4) 0.00581(5) 0.00480(5)	s.o.f 1 1 1
Atom Mo1 Mo2 Mo3 Se1	x 0.65070(3) 0.85270(3) 0.84259(3) 0.63359(3)	y -0.51996(3) -0.15840(3) 0.0000 -0.38357(3)	z 0.35272(2) 0.28920(2) 0.2500 0.31594(2)	Ueq 0.00546(4) 0.00581(5) 0.00480(5) 0.00793(6)	s.o.f 1 1 1 1 1
Atom Mo1 Mo2 Mo3 Se1 Se2	x 0.65070(3) 0.85270(3) 0.84259(3) 0.63359(3) 0.64789(3)	y -0.51996(3) -0.15840(3) 0.0000 -0.38357(3) -0.30520(4)	z 0.35272(2) 0.28920(2) 0.2500 0.31594(2) 0.37892(2)	Ueq 0.00546(4) 0.00581(5) 0.00480(5) 0.00793(6) 0.00943(6)	s.o.f 1 1 1 1 1 1 1
Atom Mo1 Mo2 Mo3 Se1 Se2 Se3	x 0.65070(3) 0.85270(3) 0.84259(3) 0.63359(3) 0.64789(3) 0.70133(4)	y -0.51996(3) -0.15840(3) 0.0000 -0.38357(3) -0.30520(4) -0.29867(4)	z 0.35272(2) 0.28920(2) 0.2500 0.31594(2) 0.37892(2) 0.2500	Ueq 0.00546(4) 0.00581(5) 0.00480(5) 0.00793(6) 0.00943(6) 0.00858(7)	s.o.f 1 1 1 1 1 1 1 1 1
Atom Mo1 Mo2 Mo3 Se1 Se2 Se3 Se4	x 0.65070(3) 0.85270(3) 0.84259(3) 0.63359(3) 0.64789(3) 0.70133(4) 0.6667	y -0.51996(3) -0.15840(3) 0.0000 -0.38357(3) -0.30520(4) -0.29867(4) -0.6667	z 0.35272(2) 0.28920(2) 0.2500 0.31594(2) 0.37892(2) 0.2500 0.38847(2)	Ueq 0.00546(4) 0.00581(5) 0.00480(5) 0.00793(6) 0.00943(6) 0.00858(7) 0.01002(9)	s.o.f 1 1 1 1 1 1 1 1 1 1 1 1 1
Atom Mo1 Mo2 Mo3 Se1 Se2 Se3 Se4 Se5	x 0.65070(3) 0.85270(3) 0.84259(3) 0.63359(3) 0.64789(3) 0.70133(4) 0.6667 1.0000	y -0.51996(3) -0.15840(3) 0.0000 -0.38357(3) -0.30520(4) -0.29867(4) -0.6667 0.0000	$\begin{array}{c}z\\0.35272(2)\\0.28920(2)\\0.2500\\0.31594(2)\\0.37892(2)\\0.2500\\0.38847(2)\\0.32432(2)\end{array}$	Ueq 0.00546(4) 0.00581(5) 0.00480(5) 0.00793(6) 0.00943(6) 0.00858(7) 0.01002(9) 0.01029(9)	s.o.f 1 1 1 1 1 1 1 1 1 1 1 1 1
Atom Mo1 Mo2 Mo3 Se1 Se2 Se3 Se4 Se5 K	x 0.65070(3) 0.85270(3) 0.84259(3) 0.63359(3) 0.64789(3) 0.70133(4) 0.6667 1.0000 1.0000	y -0.51996(3) -0.15840(3) 0.0000 -0.38357(3) -0.30520(4) -0.29867(4) -0.6667 0.0000 0.0000	$\begin{array}{c}z\\0.35272(2)\\0.28920(2)\\0.2500\\0.31594(2)\\0.37892(2)\\0.2500\\0.38847(2)\\0.32432(2)\\0.38101(4)\end{array}$	U_{eq} 0.00546(4) 0.00581(5) 0.00480(5) 0.00793(6) 0.00943(6) 0.00858(7) 0.01002(9) 0.01002(9) 0.0400(4)	s.o.f 1 1 1 1 1 1 1 1 1 1 1 1 1

Table 2. Fractional atomic coordinates, equivalent isotropic thermal displacement parameters(in Å²) and site occupancy factors for single-crystalline K2M015Se19 and Ag2.2K2M015Se19.

Table 3. Selected interatomic distances (in Å) for single-crystalline K₂Mo₁₅Se₁₉ and Ag_{2.2}K₂Mo₁₅Se₁₉.

	$K_2Mo_{15}Se_{19}$	$Ag_{2.2}K_2Mo_{15}Se_{19}$
Mo6Se8 unit		
Mo1-Mo1 (×2) intratriangle	2.6799(3)	2.6797(4)
Mo1-Mo1 (×2) intertriangle	2.7675(3)	2.7146(4)
Mo1-Se4	2.5374(4)	2.5733(4)
Mo1-Se2	2.5898(3)	2.6282(4)
Mo1-Se1	2.5968(3)	2.5902(4)
Mo1-Se1	2.6092(4)	2.6030(3)
Mo ₉ Se ₁₁ unit		
Mo2-Mo2 (×2) intratriangle	2.6682(3)	2.6424(4)
Mo3-Mo3 (×2) intratriangle	2.6821(3)	2.7158(5)
Mo2-Mo3 intertriangle	2.7138(2)	2.6873(2)
Mo1-Mo2 intercluster	3.459(2)	3.579(3)
Mo2-Se5	2.5370(4)	2.5317(4)
Mo2-Se2	2.5818(3)	2.5747(4)
Mo2-Se2	2.6105(3)	2.6300(4)
Mo2-Se1	2.6563(3)	2.6957(4)
Mo2-Se3	2.7123(3)	2.6837(3)
Mo3-Se2	2.5345(3)	2.6086(3)
Mo3-Se3	2.5933(3)	2.5779(3)
K-Se1	3.7959(8)	3.7050(10)
K-Se2	3.2255(3)	3.3010(3)
K-Se3	3.7891(7)	4.0618(11)
K-Se4	3.4256(15)	3.673(2)
K-Se5	3.6760(15)	3.260(2)
Ag-Se2		2.5829(14)
Ag-Se2		2.7194(13)
Ag-Se3		2.5505(8)
Ag-Se4		3.0135(10)

Figure Captions

Figure 1. Views of the crystal structures of (a) K₂Mo₁₅Se₁₉ and (b) Ag_{2.2}K₂Mo₁₅Se₁₉ projected along the [010] direction. Thermal displacement ellipsoids were drawn at the 97% probability level.

Figure 2. Inter-unit linkage of the Mo₉Se₁₁Se₆ and Mo₆Se₈Se₆ cluster units. Thermal displacement ellipsoids were drawn at the 97% probability level.

Figure 3. a) Se environment of the K^+ cations in $K_2Mo_{15}Se_{19}$. b) Se environment of the Ag^+ and K^+ cations in $Ag_{2.2}K_2Mo_{15}Se_{19}$.

Figure 4. Profile matching refinement of the powder X-ray diffraction patterns of K₂Mo₁₅Se₁₉ (top) and Ag₃K₂Mo₁₅Se₁₉ (bottom) ($\lambda = 1.5406$ Å). The black circles, red line and bottom black line correspond to the experimental, calculated and difference profile, respectively. All the observed reflections are consistent with the crystal structure determined from refinements against the single-crystal X-ray diffraction data. The vertical ticks mark the calculated positions of the reflections.

Figure 5. Total and partial densities of states as a function of the energy in K₂Mo₁₅Se₁₉.

Figure 6. COHP curves calculated for the Mo-Mo intra-bonds in the Mo_6 cluster for $K_2Mo_{15}Se_{19}$.

Figure 7. COHP curves calculated for the Mo-Mo intra-bonds in the Mo9 cluster for K₂Mo₁₅Se₁₉.

Figure 8. COHP curves calculated for the inter-cluster bonds Mo1-Mo2 for K₂Mo₁₅Se₁₉.

Figure 9. Temperature dependence of the a) electrical resistivity, b) thermopower, c) total (filled symbols) and lattice thermal conductivity (open symbols) and d) *ZT* of the K₂Mo₁₅Se₁₉ and Ag₃K₂Mo₁₅Se₁₉ compounds. The color-coded symbols are similar for all panels. In panel c), the solid horizontal black line corresponds to the estimated minimum thermal conductivity.



Figure 1



Figure 2



Figure 3





Figure 4



Figure 5



Figure 6



Figure 7



Figure 8



Figure 9

Table of Contents Graphic



This work reports a detailed study of the influence of Ag insertion on the thermoelectric properties of the Mo-based cluster compound $K_2Mo_{15}Se_{19}$ that yields enhanced thermoelectric properties at high temperatures.