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Pentagonal Bipyramidal Ln(III) Complexes Containing an Axial Phosphine Oxide Ligand: Field-induced Single-ion Magnetism Behavior of the Dy(III) Analogues

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Abstract

A series of neutral homologous complexes $[(L)Ln(Cy_3PO)Cl]$ {where $Ln = Gd (1)$, Tb (2), Dy (3) and Er (5)} and $[(L)Dy(Ph_3PO)Cl]$ (4) $[H_2L = 2,6$ -diacetylpyridine *bis*benzoylhydrazone] have been isolated. In these complexes, the central lanthanide ion possesses a pentagonal bipyramidal (PBP) geometry with an overall pseudo D_{5h} symmetry. The coordination environment around the lanthanide ion comprises of three nitrogen and two oxygen donors in an equatorial plane. The axial positions are taken up by a phosphine oxide (O donor) and a chloride ion. Among these compounds, the Dy(III) (**3** and **4**) analogues were found to be field-induced single-ion magnets.

Introduction

There has been a significant renaissance in the chemistry of the *rare earth* elements because of their applications in catalysis¹, photophysical properties² and in magnetic materials³. In recent years lanthanide- 4 and some actinide complexes⁵ are finding increasing utility as molecular magnets (single-molecule- and single-ion magnets; SMMs and SIMs). These molecular systems, once magnetized, retain their magnetization even after the removal of the external magnetic field and are characterized by a slow reversal of magnetization below certain temperatures. ⁶ This is because in SMMs, on application of a magnetic field, a doublewell potential comprising of the various m_j or m_s states with an energy barrier (U_{eff}) is generated which prevents the reversal of magnetization below certain temperatures (the blocking temperature, $T_{\rm B}$).⁷ Various relaxation mechanisms including quantum tunnelling are prevalent to allow the magnetization to be lost. $8⁸$

The evidence of SMM behaviour in a molecular Tb(III) complex, where the Tb(III) is sandwiched by two phthalocyanine ligands, was first reported by Ishikawa and co-workers.⁹ One of the intrinsic features of the lanthanide elements is that the 4f electrons are deeply buried inside the [Xe] core and are considerably shielded by the 5s and 5p electrons. This results in an almost unquenched orbital angular momentum (*L*) which couples with the spin angular momentum (*S*) giving rise to the total angular momentum, *J*. 4d, 10 Unlike transition metal ions, the magnitude of spin-orbit coupling in the case of 4f metal ions is comparatively much larger than the crystal field and which splits the ground $^{2S+1}L_1$ term into different *J* multiplets. Although the crystal field effects are small, it has a significant impact on removing the degeneracy of the $(2J + 1)$ m_J microstates corresponding to each of the *J* multiplets. Since the dynamics of magnetization relies on the relative energies of the ground *J* manifold, therefore, a suitable crystal field renders the requirement of large splitting between the energy levels giving rise to high energy barrier for magnetization reversal.¹¹

Soon after this discovery, many mononuclear Ln(III) complexes were reported to be SMMs with high energy barriers (U_{eff}) and high blocking temperatures (T_B).¹² Among various types of lanthanide complexes, the mononuclear complexes are of considerable interest as they provide a very good understanding of the influence of the ligand field on the observed magnetic properties. The recent reports on mononuclear $Dy(III)$ complexes $[Dy(Cp^{ttt})₂]⁺$ $(Cp^{ttt} = C_5H_2^tBu_3-1,2,4)$ and $[Dy(Cp^{iPr5})(Cp^*)]^+$ $(Cp^{iPr5} =$ penta-iso-propylcyclopentadienyl, Cp^* = pentamethylcyclopentadienyl) revealing the highest magnetization blocking temperatures of 60 K and 80 K respectively has further spurred activity in this area of mononuclear $Ln(III)$ complexes.¹³

Although ligand fields are much smaller compared to spin-orbit coupling among lanthanide complexes, ironically the former happens to be the most decisive in controlling the performance of SMMs/SIMs. The spatial distribution of the electrons in the different 4f orbitals leads to inherent anisotropic shapes in the Ln(III) ions (except Gd, Eu, La, and Lu). Based on an electrostatic argument, Rinehart and Long have proposed a qualitative model that assists in the designing of $SMMs/SIMs.$ ¹⁴ According to this model, an axial ligand field stabilizes the oblate-shaped Ln(III) ions while a prolate-shaped Ln(III) ion requires an equatorial ligand field because such a ligand field minimizes the electrostatic repulsion between the ligands and the metal center and maximizes the molecular magnetic anisotropy. Using this clue a large number of monometallic Ln(III) complexes were prepared with interesting magnetic properties.¹⁵ Among them, the *pseudo*-linear pentagonal bipyramidal complexes in the DyO7, DyClO6, DyXN4O2 ($X = Cl$ and Br), and DyN5O2 coordination environment with strong axial ligand field and weak equatorial ligand field stand out as the most effective systems for the observation of high energy barriers of magnetization reversal.¹⁶

We have been utilizing various types of multidentate ligands for the synthesis of mononuclear Ln(III) complexes. In our previous work, we have synthesized mononuclear pentagonal bipyramidal Ln(III) complexes by employing a pentadentate chelating ligand that provides a rigid equatorial plane.¹⁷ The axial sites in these complexes were occupied by the chloride ions which are considerably weak field ligands compared to N and O donors present in the ligand backbone. We have thoroughly studied the magnetic properties of the Dy(III), Tb(III) and the diluted Dy(III) (in an isostructural Y(III) host) complexes which reveal the molecular origin of slow magnetic relaxation in the Dy(III) derivative with an energy barrier of magnetization reversal of 70 K.¹⁷ To understand the role of the axial ligands in this system we have now prepared a series of neutral mononuclear PBP complexes, $[(L)Ln^{III}(R_3PO)Cl]$ $[(\text{Ln} = \text{Gd} (1), \text{Tb} (2), \text{Dy} (3), \text{Er} (5); \text{R} = \text{cyclohexyl})$ and $(\text{Ln} = \text{Dy} (4); \text{R} = \text{phenyl})]$ where one phosphine oxide ligand replaces one of the two chloride ligands in the axial sites. Herein, we report the synthesis, structural characterization and magnetic properties of **1**-**5**. The difference in the magnetic properties has been analyzed by *ab initio* CASSCF/RASSI-SO/SINGLE_ANISO calculation which provides the insight to design the potential SIM.

Experimental Section

Materials and methods. All the reagents and solvents used for the syntheses were used as received from commercial sources. The organic ligand 2,6-diacetylpyridine *bis*benzoylhydrazone $(H₂L)$ was synthesized following a reported procedure.¹⁸ Fourier transform infrared (FT-IR) spectroscopy was performed with a Bruker FT-IR spectrometer. Elemental analyses were performed with a Perkin–Elmer 2400 series II instrument. Powder X-ray diffraction study was performed on finely ground polycrystalline material with Bruker D8 Advance Powder X-ray diffractometer.

Magnetic Measurements. Magnetic measurements for all the samples were carried out with a Quantum Design MPMS 5S SQUID magnetometer in the temperature range 2−300 K. The measurements were performed on polycrystalline samples. The crystalline powders of the complexes were mixed with grease (except for Gd derivative) and put in gelatin capsules. The temperature dependences of the magnetization were measured in an applied field of 1kOe, and the isothermal field dependence of the magnetizations were collected up to 5 T. The molar susceptibility (χ_M) was corrected for sample holder and for the diamagnetic contribution of all the atoms by using Pascal's tables. AC susceptibility data have been collected in zero field and with applied fields in the frequency range 1–1500 Hz.

X-ray crystallographic studies. The single-crystal X-ray diffraction data of **1-5** were collected on a Rigaku Xtal LAB X-ray Diffractometer system equipped with a CCD area detector and operated at 30 W power (50 kV, 0.6 mA) to generate MoK α radiation (λ = 0.71073 Å) at 120(2) K. Data were integrated using CrysAlis^{Pro} software with a narrow frame algorithm. Data were subsequently corrected for absorption by the program SCALE3 ABSPACK scaling algorithm.^{13a} All the structures were solved by the direct methods in SHELXTL¹⁹ and refined by the full-matrix least-squares method on F^2 (SHELXL-2014)²⁰ using the Olex-2 software.²¹ All the non-hydrogen atoms were refined with anisotropic thermal parameters. All the hydrogen atoms were included in idealized positions, and a riding model was used. All the mean plane analyses and crystallographic figures have been generated using the DIAMOND software (version $3.2k$).²² The crystal data and refinement parameters for **1**–**5** are summarized in Table 1. More details on the crystallographic data are given in the X-ray crystallographic files in the CIF format.

Synthesis

General procedure. The following general protocol was employed for the synthesis of complexes **1**-**5**.

The organic ligand, $H₂L$ (1 eq.) was suspended in 30 mL of EtOH and cyclohexyl/phenyl phosphine oxide (1 eq.) was added to it. To this white cloudy solution, the respective LnCl₃^{\cdot 6H₂O (1 eq.) salts were added which results in a yellow solution. The reaction mixture} was then heated under reflux conditions for 1 h and allowed to cool to room temperature. To this solution 2 eq. of NEt₃ was added and the solution further stirred at room temperature for 10 minutes. The solvent was evaporated to dryness and the resulting yellow precipitate was washed with diethyl ether. The dried yellow precipitate was then dissolved in 10 mL of EtOH and filtered. The filtrate was kept under vapor diffusion with diethyl ether to afford needleshaped crystals suitable for X-ray crystallography after one week. The stoichiometry of the reactants involved in each reaction, yield of the products, and their characterization data are provided below:

 $[(L)Gd(Cy_3PO)Cl]$ (1). H₂L (0.040 g, 0.100 mmol), GdCl₃·6H₂O (0.037 g, 0.100 mmol), Cy₃PO (0.030 g, 0.100 mmol), and Et₃N (28 μ L, 0.200 mmol) were used. Yield: 0.053 g, 60% (based on Gd). M.P.: >250 ℃. IR (KBr ν/cm-1): 3439(br), 3062(w), 2929(s), 2852(m), 1632(w), 1587(m), 1552(m), 1503(s), 1446(m), 1411(m), 1371(s), 1324(m), 1297(m), 1258(w), 1197(w), 1169(m), 1148(m), 1103(s), 1069(w), 1040(s), 987(w), 895(m), 854(w), 809(m), 744(m), 716(s), 679(s), 650(w), 534(m). Anal. Calcd for $C_{41}H_{52}Cl_1N_5O_3P_1Gd_1$ (886.57): C, 55.55; H, 5.91; N, 7.90. Found: C, 55.21; H, 6.36; N, 7.61.

 $[(L)Tb(Cy_3PO)Cl]$ (2). H₂L (0.040 g, 0.100 mmol), TbCl₃·6H₂O (0.037 g, 0.100 mmol), Cy₃PO (0.030 g, 0.100 mmol), and Et₃N (28 μ L, 0.200 mmol) were used. Yield: 0.059 g, 67% (based on Tb). M.P.: >250 ℃. IR (KBr ν/cm-1): 3441(br), 3064(w), 2927(s), 2854(m), 1634(w), 1587(m), 1552(m), 1505(s), 1446(m), 1409(m), 1368(s), 1326(m), 1299(m), 1256(w), 1197(w), 1169(m), 1148(m), 1105(s), 1067(w), 1040(s), 987(w), 897(m), 856(w), 809(m), 744(m), 714(s), 679(s), 652(w), 532(m). Anal. Calcd for $C_{41}H_{52}Cl_1N_5O_3P_1Tb_1$ (888.25): C, 55.44; H, 5.90; N, 7.88. Found: C, 55.02; H, 6.56; N, 7.75.

 $[(L)Dy(Cy_3PO)Cl]$ (3). H₂L (0.040 g, 0.100 mmol), DyCl₃·6H₂O (0.038 g, 0.100 mmol), Cy₃PO (0.030 g, 0.100 mmol), and Et₃N (28 μ L, 0.200 mmol) were used. Yield: 0.061 g, 69% (based on Dy). M.P.: >250 ℃. IR (KBr ν/cm-1): 3443(br), 3064(w), 2929(s), 2852(m), 1630(w), 1587(m), 1554(m), 1505(s), 1446(m), 1411(m), 1368(s), 1326(m), 1299(m), 1258(w), 1197(w), 1171(m), 1150(m), 1105(s), 1067(w), 1042(s), 989(w), 897(m), 854(w), 809(m), 744(m), 714(s), 679(s), 650(w), 532(m). Anal. Calcd for $C_{41}H_{52}Cl_1N_5O_3P_1D_{11}$ (891.82): C, 55.22; H, 5.88; N, 7.85. Found: C, 54.85; H, 6.39; N, 7.71.

 $[(L)Y_{0.90}Dy_{0.10}(Cy_3PO)CI]$ (3'): Anal. Calcd for $C_{41}H_{52}Cl_1N_5O_3P_1Y_{0.90}Dy_{0.10}$: C, 60.19; H, 6.41; N, 8.56. Found: C, 59.91; H, 6.52; N, 8.39.

 $[(L)Dy(Ph_3PO)Cl]$ (4). H₂L (0.040 g, 0.100 mmol), DyCl₃·6H₂O (0.038 g, 0.100 mmol), Ph₃PO (0.028 g, 0.100 mmol), and Et₃N (28 *μ*L, 0.200 mmol) were used. Yield: 0.064 g, 72% (based on Dy). M.P.: >250 °C. IR (KBr v/cm^{-1}): 3429(br), 3054(w), 2919(s), 1632(w), 1587(m), 1552(m), 1499(m), 1438(m), 1409(), 1366(s), 1325(m), 1297(m), 1258(w), 1160(w), 1122(m), 1093(m), 1067(s), 1044(w), 989(w), 897(m), 809(m), 744(m), 714(s), 691(), 650(2), 540(m). Anal. Calcd for C41H34Cl1N5O3P1Dy1 (873.68): C, 56.37; H, 3.92; N, 8.02. Found: C, 55.98; H, 3.61; N, 7.89

 $[(L)E_r(Cy_3PO)Cl]$ (5). H₂L (0.040 g, 0.100 mmol), ErCl₃·6H₂O (0.038 g, 0.100 mmol), Cy₃PO (0.030 g, 0.100 mmol), and Et₃N (28 μ L, 0.200 mmol) were used. Yield: 0.065 g, 73% (based on Er). M.P.: >250 IR (KBr v/cm⁻¹): 3447(br), 3068(w), 2929(s), 2852(m), 1636(w), 1587(m), 1554(m), 1505(s), 1446(m), 1413(m), 1366(s), 1326(m), 1299(m), 1260(w), 1199(w), 1169(m), 1152(m), 1107(s), 1067(w), 1044(s), 989(w), 897(m), 854(w), 809(m), 746(m), 714(s), 679(s), 650(w), 534(m). Anal. Calcd for $C_{41}H_{52}Cl_1N_5O_3P_1Er_1$ (896.58): C, 54.93; H, 5.85; N, 7.78. Found: C, 54.65; H, 6.33; N, 7.53.

7

Computational Details

Post-Hartree-Fock *ab initio* calculations were carried out on the X-ray crystal structures of all complexes using the CASSCF+RASSI-SO+SINGLE_ANISO approach as implemented in MOLCAS 8.2 programme package.²³ The relativistic effects of the lanthanide ions have been incorporated using the DKH Hamiltonian.²⁴ The basis set of all the atoms (Table S8 in ESI) has been taken from the ANO-RCC library implemented in MOLCAS 8.2 suite. First, we have performed CASSCF calculations by considering 8, 9 and 11 electrons in the seven 4f orbitals of Tb(III), Dy(III) and Er(III) ions in their respective complexes. Using this active space, we have computed 7 septet, 140 quintet and 195 triplet states for Tb(III), 21 sextet states for Dy(III), 35 quartet and 112 doublet states for Er(III) ion. These spin free states of each complex were mixed by RASSI-SO to calculate the spin-orbit energy levels. Finally, the g tensors and mechanism of magnetization relaxations were estimated using the SINGLE_ANISO which interfaced with the RASSI-SO energies. In a nutshell, CASSCF/RASSI-SO/SINGLE_ANISO method was used to compute the magnetic anisotropy, energy of the spin free and spin-orbit states and the magnetic relaxation dynamics.

Results and Discussion

Synthetic aspects. The PBP geometry is regarded as one of the most promising coordination geometries around the Ln(III) center that can bring axiality in the ground state of Ln(III) ions provided the axial sites are occupied by comparatively strong donor ligands. We have previously reported the synthesis of mononuclear pentagonal bipyramidal Ln(III) complexes using a pentadentate chelating ligand which effectively provides a rigid equatorial plane. In these complexes, the two axial sites were occupied by chloride ions which can be regarded as weak field ligands compared to the N and O donor atoms of the pentadentate chelating ligand. Theoretical studies accompanied by experimental evidence show that oblate shaped Ln(III) ions show high energy barriers of magnetization in the PBP geometry when the axial sites are occupied by relatively strong donor ligands compared to the equatorial sites. Keep this in mind we have chosen tri-alkyl/aryl phosphine oxides to replace the chloride ions in the axial sites. Accordingly, when we treated the ligand $H₂L$ with lanthanide chlorides in the presence of phosphine oxides followed by addition of base we obtained neutral mononuclear $[(L)Ln^{III}(R_3PO)Cl]$ (Ln = Dy Tb, Gd, Er when R = cyclohexyl; Ln = Dy when R = phenyl) complexes (Scheme 1).

Scheme 1. Reaction scheme for the synthesis of **1**–**5**.

X-ray Crystallography

The complexes 1, 3 and 5 crystallize in the monoclinic crystal system with $P2_1/c$ (for **3**) and *P*21/n (for **1** and **5**) space groups whereas the complexes **2** and **4** crystallize in the triclinic crystal system with *P*-1 space group. Crystallographic data and refinement parameters of all the complexes are given in Table S1. The overall molecular structures of the complexes **1-5** are essentially identical. The molecular structure of complex **3** is shown in Figure 1, while those of **1**, **2**, **4** and **5** are given in the Supporting Information (Figures S1-S4). In view of the structural similarities present in the complexes we discuss below the molecular structures of complexes **3** and **4**.

Figure 1. Molecular structure of complex **3**. (The H atoms are removed for clarity)

The complexes are formed by the coordination action of ligand *pyridyl* N atom, two *imino* N atoms, and two *carboxy* O atoms in the equatorial positions. One of the two axial sites is occupied by one chloride anion in both the two complexes. The remaining axial site is occupied by one $Cy₃PO$ ligand in the case of 3 and $Ph₃PO$ ligand in the case of **4**. The ligand upon chelation with the Ln(III) ions generates four *five*-membered rings revealing its excellent ability to stabilize the Ln(III) ions in its pentagonal coordination environment. The equatorial Dy–O/N bond distances are in the range of 2.259(2)–2.462(2) Å for **3** and 2.282(2)–2.456(2) Å for **4**. The Dy−Oaxial bond distances are 2.237(2) for **3** and 2.275(2) for **4**. The Dy−Cl bond distances are 2.625(8) Å for **3** and 2.622 (7) Å for **4**. Interestingly, the Dy–O_{axial} bond distance in both **3** and **4** are shorter compared to the Dy−Oequatorial distances (Table S2 and Table S4) indicating the strong-field nature of the phosphine oxide ligand in comparison to the equatorial oxygen donors. The Ophos−Dy−Cl bond angles are 169.62(5)° for **3** and $174.07(5)$ ^o for **4**. The immediate coordination environment the Ln(III) ions are analyzed with Continuous-Shape Measures using the SHAPE program.²⁵ It reveals a distorted pentagonal bipyramid geometry around the Dy(III) ions with *D*5h (*pseudo*) CF symmetry (Table S3, see ESI). The pentagonal bipyramidal geometry of the Dy(III) ion in complex **3** is shown in Figure 2 (*left*). The shortest intermolecular Ln^{$...$}Ln distance in **1**, **3** and **4** is 8.47 Å, 8.56 Å, and 8.80 Å respectively as revealed in the solid-state packing diagram (see Figure 2 (*right*) for complex **3** and Figures S5 and S6 for **1** and **4**). The solid state phase purity of the complex **3'** was confirmed by powder X-ray diffraction studies (Figure S7 see ESI). The selected bond lengths and bond angles of complexes **1**, **2**, **4** and **5** are summarized in Table S4 (see ESI).

Figure 2. (*left*) Coordination polyhedra of Dy(III) and (*right*) solid-state packing diagram of complex **3**.

Magnetic Properties

The temperature dependence of the molar magnetic susceptibility (χ_M) for 1-3 and 5 are plotted in Figure 3 and the field dependence of the magnetization for these compounds can be found in Figure S9. The corresponding behavior for **4** is summarized in Figure S8. The values of the product of the molar magnetic susceptibility with temperature, $\chi_M T$ (in cm³ mol⁻¹ K) found at 300 K are 7.85 (Gd), 11.79 (Tb), 14.14 and 14.10 (Dy), and 11.18 (Er), in good accordance with the expected values expected for the isolated ions (i.e. 7.88, 11.81, 14.17, and 11.48 cm³ mol⁻¹ K, respectively). For the Tb, Dy, and Er derivatives the $\chi_M T$ slowly decreases as T is lowered in agreement with the anticipated crystal field effect. The absence of any contribution from intermolecular exchange interactions is confirmed by the perfect Curie behavior down to 2 K for Gd. For the Tb and Dy derivatives, the field-dependent

magnetization at 2-5 K show a fast rise at lower field regions and remain almost unchanged above 15 kOe (at 2 K) to reach 4.8 μ B (2), 5.14 (3), and 5.09 (4) at high field (5 T). To probe the dynamics of the magnetization relaxation, AC susceptibility behavior was investigated without and with applied static fields. No out-of-phase component (*χ*^{*n*}_M) were found for the Tb and Er derivatives down to 2 K (see Figure S10 in ESI). However, the Dy complexes **3** and 4 exhibited a χ''_M signal but no maximum was observed above 2 K (see Figure S10 in ESI).

Figure 3. Temperature dependence of the product of the molar magnetic susceptibility with temperature, $\chi_M T$, under 1 kOe applied magnetic field for complexes 1-3 and 5. Solid red lines represent the *ab initio* calculated data. The computed data for **5** was reduced by 2% to meet the experimentally observed.

Such a behaviour was suggesting relaxation driven by QTM which was reduced by applying a static field. For **3**, the optimum applied field was estimated at $H_{DC} = 1$ kOe (See Figure S11). AC susceptibility study carried out in this field within the frequency domain 1–1500 Hz yielded well-defined maxima for χ''_M between 4 and 18 K (Figure 4). Relaxation times (*τ*)

have been assessed by fitting $\chi''_M = f(\text{Frq})$ for different temperatures with an extended Debye model, best fit parameters are gathered in Table S5. The very small *a* parameter is indicative for a narrow distribution width for the relaxation time over the whole temperature domain, suggesting that mainly one relaxation process is operative. The temperature dependence of the relaxation time, plotted in log scale in Figure 4, develops to a more linear variation above 10 K, which is the behavior anticipated for a thermally activated process (Orbach). Deviation from linearity for lower *T* indicates that other processes also come into play. Analysis of the behavior over the whole *T* range was obtained by summing the contributions of the Orbach, Raman, and direct processes $(\tau = \tau_0 \exp(U_{\text{eff}}/k_B T) + 1/(CT^n) + 1/(AT))$. The latter were required to reproduce the lower *T* behavior. Best fit gave a thermal energy barrier for magnetization reversal, $U_{\text{eff}}/k_B = 204 \pm 3 \text{ K}$ with $\tau_0 = (6 \pm 1) \times 10^{-9} \text{ s}$, $C = 0.015 \text{ K}^{-1} \text{s}^{-1}$, $n = 4.5$, and $A = -1.30$ s^{-1} .

Figure 4. (*a,b*) AC susceptibility behaviors for **3**; (*c*) temperature dependence of the relaxation times for **3**, and (*d*) for diluted compound **3**'; the solid red lines are the best fits (see text).

The dilution of **3** in a matrix of isomorphous Y complex, **3'** hereafter, gave very similar results (see Figure S12 in ESI). For this sample however, a maximum was observed in the $\chi''_M = f(T)$ behavior in zero field but a QTM contribution was discernible at low temperature. This was suppressed upon applying a small DC field of 750 Oe. The temperature dependence of the relation times between 2 and 17.5 K (Figure 4 and Table S6) parallels that obtained for the pure Dy derivative **3**, and contributions of the Orbach, Raman, and direct processes had to be taken into account to reach a good modeling. Best fit to the experimental data for **3**' yielded $U_{\text{eff}}/k_B = 208 \pm 5 \text{ K}$, $\tau_0 = (6 \pm 2) \times 10^{-9} \text{ s}$, $C = 4.3 \times 10^{-3} \text{ K}^{-1} \text{s}^{-1}$, $n = 5.0$, and $A = -0.30 \text{ s}^{-1}$. The comparison of the behavior for **3** and its diluted form **3**' shows that the observed behavior is clearly of molecular origin.

The replacement of $OPCy₃$ for $OPPh₃$ in the apical position of Dy appeared to have no significant incidence on the magnetic behaviours; the AC susceptibility features for **4** (Figure S13 and Table S7 in ESI) are very similar to that obtained for **3**. To reproduce the temperature dependence of the relaxation time for **4** required to consider contributions from Orbach, Raman and direct processes, best fit to the experimental behavior gave $U_{\text{eff}}/k_B = 241$ \pm 7 K, $\tau_0 = (2.3 \pm 0.9) \times 10^{-10}$ s, C = $(5.97 \pm 0.03) \times 10^{-3}$ K⁻¹ s⁻¹, n = 5.1, and A = 0.2 \pm 0.7 s⁻¹.

It is satisfying to see that the energy barrier for magnetization reversal, U_{eff}/k_B , for the Dy(III) complexes reported herein is significantly increased with respect to the homologue complex with Cl[–] ligands in the apical positions. This could be attributed to the stronger axial field due to the phosphine oxide ligand. A detailed comparison of the bond and magnetic properties of the present complexes with the reported pentagonal bipyramidal Dy(III) complexes having other equatorial ligands (including monodentate ones) and axial groups of various crystal field strengths are given in Table 1.

SI. No.	Complex	L_{ax} -Ln- L_{ax} (°)	Average axial Ln- $O/Ln-X$ distance (\check{A})	Average equatorial Ln- $O/Ln-N/Ln-X$ distance (\check{A})	$U_{\rm eff}$ (K) (H_{dc})	$\tau_{0}(s)$	Ref.
$\mathbf{1}$	$[Dy(O'Bu)2(py)5][BPh4]$	155.80	2.112	2.557	1815 $(0$ Oe $)$	1.17×10^{-12}	\lceil ^{16e}]
\overline{c}	[Dy(bbpen)Br]	155.79	2.163	2.586/2.851	1025 (0 0e)	4.21×10^{-12}	\lceil ^{16c}]
$\overline{3}$	$[Dy(^tBuO)Cl(THF)_5]$ $[BPh_4]$ 2THF	178.28	2.043/2.662	2.410	950 (0 0e)	3×10^{-12}	[26]
$\overline{4}$	[Dy(bbpen)Cl]	154.24	2.166	2.584/2.682	708 (0 0e)	9.46×10^{-11}	$\left[16c\right]$
5	$[L^{1}_{2}Dy(H_{2}O)_{5}][I]_{3}(L^{1})_{2}$ H_2O	177.9	2.205	2.363	651 (0 0e)	5.63×10^{-12}	\lceil ^{16b}]
6	$[Dy(L^2)2(H2O)5]2·Br6·$ $2L^2.2H2O$	177.82	2.221	2.354	556 (0 0e)	9.33×10^{-12}	$\left[\begin{matrix} 27 \\ 1 \end{matrix} \right]$
$\boldsymbol{7}$	$[Dy(Cy_3PO)_2(H_2O)_5]Br_3.$ $2(Cy_3PO)$ $2H_2O$ $2EtOH$	179.04	2.200	2.352	543 (0 0e)	2.0×10^{-11}	$[$ ²⁸ $]$

Table 1. Selected bond and magnetic parameters of PBP Dy(III) complexes

 L^1 = ('BuPO(NHⁱPr)₂; L²: hexamethylphosphoric triamide; H₂bbpen = *N,N*^{*r*}-bis(2-hydroxybenzyl)-*N,N^{<i>r*}-bis(2-hydroxybenzyl)-*N,N^r*-bis(2-hydroxybenzyl)-*N,N^r*-bis(2-hydroxybenzyl)-*N,Nr*⁻bis(2-hydrox methylpyridyl)ethylenediamine); NCN: $[2,6-(2,6-C_6H_3R_2N_5CH)_2C_6H_3]$; Bpen: N , N '-bis(2methylenepyridinyl)ethylenediamine; H_4L^3 : 2,6-diacetylpyridine bis-salicylhydrazone; H_2L^4 : : 2,6 diacetylpyridine bis-benzoylhydrazone

Computational study

In order to rationalize the experimentally observed magnetization relaxation, we have performed *ab initio* CASSCF/RASSI+SO/SINGLE_ANISO calculations²³ on the X-ray structures of complexes 2-5 and the previously reported complex 6^{17} in which both the axial positions were occupied by the Cl⁻ ligands. The X-ray analysis revealed that complexes 2 and **5** consist of two molecules in the asymmetric unit and we have performed our calculations on the one molecule for each of the complexes.

Figure 5. Orientation of the anisotropy axis of Dy(III) ion in complex **3** (a) and **4** (b). (c) and (d) represent the plausible relaxation mechanism for **3** and **4** respectively. The red arrows show the QTM and TA-QTM via ground and higher excited KD respectively. The sky dotted arrows show the Orbach process for the relaxation. The green arrows show the possible mechanism of magnetic relaxation. The blue characters imply the m_J composition of the KDs.

Figure 6. (a) The beta electron density of complex **3**. (b) The beta electron density of complex **4**. Colour code: Dy-yellow, Cl-green, P-violet, O-red, N-blue, C-grey, H-white.

The computed energy of the eight low lying KDs generated from ${}^{6}H_{15/2}$ state spans from zero to 903 K in complexes 3 and 4 and 942 K in complex 6 (Table 2). The g tensors; $g_x \approx g_y \approx 0$ and $g_z \approx 20$ for complexes 3, 4 and 6 demonstrate the *Ising* behavior of Dy(III) ion in these complexes. The g_{zz} axis of ground KD is found to be lying on the equatorial plane due to the strong crystal field generated from the equatorial ligand (Figure 5 and S14). We have plotted the beta electron density of Dy in complexes **3**, **4** and **6** using the procedure introduced by Ruiz and coworkers.³³ Since the beta electron density is found to orient along the axial Cl– Dy–O bond (Cl–Dy–Cl bond in 6), the g_{zz} anisotropy axis is oriented along the perpendicular direction of it to minimize the electrostatic repulsion (Figure 6). The large thermally assisted quantum tunneling (TA-QTM) value of 0.41, 0.64 and 0.24 μ_B in the first excited KDs suggest the relaxation will occur ideally through the first excited state KD (Figure 5 and S14). This suggests U_{cal} value of 222.6, 222.9 and 211.0 K for complexes 3, 4 and 6 respectively. These values are consistent with the estimate obtained from the experiments except for **6** where U_{cal} values are overestimated (for 3 208 K, 4 241 K and 6 70 K). We have also analyzed the composition of the computed wave function to get more insight about the extent of mixing of the m_J levels. The ground KD is found to be consist of m_J = \pm 15/2 (\approx 97%) with the negligible ($\approx 0.02\%$) mixing with m_J = \pm 11/2 states. The strong mixing in the first excited KD leads to very large TA-QTM which forces the complex to relax via this state.

The *ab initio* computed energy barrier for the magnetization reversal does not affect significantly when one Cl⁻ ion in the axial position (complex 6) is replaced by the neutral Cy3PO or Ph3PO ligand (complexes **3** and **4** respectively). The energy spectrum depends on the ligand field generated from the equatorial ligands which is found to be same in all the complexes. We have analyzed the LoProp charges to get more insight into the axiality of the three complexes. The LoProp charge of the axial oxygen atoms is found to be larger than the equatorial nitrogen and oxygen atoms in **3** and **4** (Table S9 and Figure S15). More importantly, the ratio of the av. axial/equatorial LoProp charges is found to be almost same in the complexes 3 and 4 which is consistent with similar U_{cal} values estimated in these complexes. The LoProp charge of the axial chlorine atoms and equatorial oxygen atoms is found to be very close in **6** which imply the less axial nature of this complex compared to **3** and **4**. However the av. axial/equatorial LoProp charge >1.7 implies significant axiality in this complex.

The crystal field parameter of the three complexes has been estimated using the Steven formalism ($\widehat{H}_{CF} = \sum \sum_{k=-q}^{+q} B_k^q \widetilde{O}_k^q$ $+q \partial P^q \tilde{\Omega}^q$ $_{k=-q}^{+q} B_k^q \tilde{O}_k^q$, where B_k^q a_k^q and \tilde{O}_k^q $\frac{q}{k}$ is computed crystal field parameter and Stevens operator respectively) as implemented in SINGLE_ANISO suite to get more insight on the relaxation mechanism. The axial crystal field parameter B_2^0 is found to be larger in 4 compare to **3** implies the stronger axiality in **4** which agrees well with the Ueff value (Table S10). However the B_2^2 value in complex 6 is found to be similar implies lower axiality here compare to **3** and **4**.

Furthermore, we have modeled the complex **3a**, **4a** and **6a** (where the equatorial ligand has been removed) from complexes **3**, **4** and **6** respectively to achieve the axial limit in each of these complexes. The anisotropy axis is found to be oriented along the Cl–Dy–O bond due to the absence of equatorial ligand field (Figure S16). The beta electron density is found to be oblate in nature and the g_{zz} axis is oriented along the axial Cl–Dy–O bond (Cl–Dy–Cl bond in **6**) to minimize the electrostatic repulsion (Figure 7). The *ab initio* calculations on these models reveal that the energy splitting of the eight KDs generated from ${}^{6}H_{15/2}$ state raises to 2864, 2787 and 2616 K for **3a**, **4a** and **6a** respectively (Table S11). Absence of equatorial ligand field leads to higher LoProp charge in the Dy center and metal coordinated oxygen atoms revealing the axial limit that can be achieved in these set of complexes (Table S12). It is also reflected in the computed crystal field parameter in which larger B_2^0 value is found compare to original complexes (Table S13). The magnetization relaxation is found to occur via fourth excited KD due to the strong transverse anisotropy which leads to U_{cal} value of 2540, 2471 and 2248 K in **3a**, **4a** and **6a** respectively (Figure S17).

Figure 7. (a) The beta electron (spin-down) density of model **3a**. (b) The beta electron density of model **4a**. (c) The beta electron density of model **6a**. Colour code: Dy-yellow, Clgreen, P-violet, O-red, N-blue, C-grey, H-white.

In the case of complex **2**, a large tunnel splitting (Δ_{tun}) of 2.26 cm⁻¹ is found between the two pseudo ground KDs which triggers the magnetization relaxation via the ground state (Table S14). The ground state anisotropy axis is found to be lying on the equatorial plane due to the strong equatorial ligand field (Figure S18). In case of complex **5** the computed large QTM value of 0.22 μ ^B between the ground state KD also suggests the relaxation via ground state which is supported by no out-of-phase signals in ac susceptibility (Figure S18 and Table S15).

Summary

Mononuclear 4f complexes $(Gd^{III}, Dy^{III}, Er^{III})$ and Tb^{III}) have been assembled using a multifunctional chelating ligand. The ligand enforces a pentagonal equatorial architecture around the metal ion. The two axial sites are taken up by a chloride and a phosphine oxide ligand. Thus, the overall coordination geometry around the lanthanide metal ion is distorted pentagonal bipyramidal. Magnteic studies on these complexes reveal that the Dy ^{III} analogues show a slow relaxation of magnetization under the presence of small DC fields. A large tunneling gap between the ground as well as consequent excited non-Kramers doublets of Tb^{III} and a very high quantum tunneling between the ground state Kramers doublets in the E^{III} complex prevents these complexes from revealing a single-ion magnet behavior. The strategy adopted in the present instance of using a rigid ligand that enforces a pentagonal equatorial plane around the lanthanide ion is clearly effective in enabling the designed assembly of single-ion magnets. An improvement in this strategy would be to effect modifications in the ligand design to decrease the effective crystal field in the equatorial plane and increase it in the axial sites. We are currently examining these possibilities.

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Supporting Information

CCDC 1989266 (**1**), 1989267(**2**), 1989268(**3**), 1989270(**4**), 1989271(**5**) contain crystallographic supplementary data for this paper. These data can be obtained free of charge from The Cambridge Crystallographic Data Centre via www.ccdc.cam.ac.uk/data_request/cif.

X-ray crystallographic table of complexes **1**–**5**; molecular structures of complexes **1**, **2**, **4** and **5**: crystal packing diagrams of **1** and **4**. table of bond distances and angles; table of continuous shape measures calculations; powder X-ray diffraction pattern of **3**'. details of magnetic properties of **1**−**5**; details of computaional calculations of complexes **1**−**5** (PDF).

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28

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Supporting Information

Table S1. Crystallographic data and refinement parameters of **1**–**5**.

Figure S1. Molecular structure of complex **1**. (H-atoms are omitted for clarity)

Figure S2. Molecular structure of complex **2**. (H-atoms are omitted for clarity)

Figure S3. Molecular structure of complex **4**. (H-atoms are omitted for clarity)

Figure S4. Molecular structure of complex **5**. (H-atoms are omitted for clarity)

Figure S5. Solid-state packing diagram of complex **1** viewed along the crystallographic *c* axis. (H-atoms are omitted for clarity)

Figure S6. Solid-state packing diagram of complex **4** viewed along the crystallographic *c* axis (H-atoms are omitted for clarity).

Bond distances (Å)		Bond angles (°)		
			O2-Dy1-Cl1 93.19(5) O2-Dy1-O1	100.23(7)
$Dy1 - C11$	2.625(8)		O2-Dy1-N3 130.25(7) O2-Dy1-N2	165.17(7)
Dy1-01	2.272(2)		O2-Dy1-N4 65.59(7) O3-Dy1-Cl1	169.62(5)
Dy1-02	2.259(2)		O3-Dy1-O1 92.19(7) O3-Dy1-O2	90.35(7)
$Dy1-O3$	2.237(2)		O3-Dy1-N3 84.83(7) O3-Dy1-N2	91.41(7)
$Dy1-N3$	2.462(2)		O3-Dy1-N4 81.19(7) N3-Dy1-Cl1	85.44(6)
$Dy1-N2$	2.464(2)		N3-Dy1-N2 64.57(7) N2-Dy1-Cl1	87.67(6)
$Dy1-N4$	2.448(2)		O1-Dy1-Cl1 96.78(6) O1-Dy1-N3	129.37(7)
			O1-Dy1-N2 65.00(7) O1-Dy1-N4	164.10(7)
			N4-Dy1-Cl1 91.39(6) N4-Dy1-N3	64.74(7)

Table 2. Bond distance and bond angle parameters of complex **3**.

Table S3. Continuous shape analysis

	Structure [']						
Complex	$HP-7$	HPY-7	PBPY-7	COC-7	CTPR-7	JPBPY-7	JETPY-7
1	31.572	19.350	2.233	7.665	6.207	6.298	18.718
$\overline{2}$	32.740	20.143	1.732	7.872	6.409	5.652	20.810
3	32.944	22.380	1.446	8.229	6.678	5.549	21.175
4	32.742	21.540	1.505	7.661	6.338	5.725	21.267
5	32.781	20.809	1.353	7.774	6.403	5.232	21.450

¹HP-7: Heptagon (D_{7h}); HPY-7: Hexagonal pyramid (C_{6v}); PBPY-7: Pentagonal bipyramid (D_{5h}) ; COC-7: Capped octahedron (C_{3v}) ; CTPR-7: Capped trigonal prism (C_{2v}) ; JPBPY-7: Johnson pentagonal bipyramid J13 (D_{5h}); JETPY-7: Johnson elongated triangular pyramid J7 (C_{3v})

Table S4. Crystallographic details of **1**, **2**, **4** and **5**.

Figure S7. The experimental and simulated powder XRD pattern of complex **3'**.

Figure S8. Magnetic behavior for **4**. *(a)* Temperature dependence of $\chi_M T$ and *(b)* field dependence of the magnetization (inset), solid lines represent the *ab initio* computed data using MOLCAS 8.2 program.

Figure S9. Field dependence of the magnetization (A-C) for complexes **2**-**3** and **5** respectively, solid lines are the ab initio computed data using MOLCAS 8.2 program.

Figure S140. temperature dependence of the AC susceptibility in zero field and in an applied field of 1 kOe or 750 Oe for **2** (a), **3** (b), **3'** (c), **4** (c), and **5** (e).

Figure S10. Field dependence of the AC susceptibility *(left),* and of the relaxation time, ^τ *(right),* for 3 at 8 K. The larger τ is reached for 1 kOe which was used as applied DC field in the AC studies for this complex.

T(K)	$\chi_{\scriptscriptstyle \rm T}$	$\chi_{\rm S}$	$\tau(s)$	α
4.5	5.71	3.29	0.11675	0.12
5.0	5.61	3.39	0.065181	0.12
5.5	5.52	3.47	0.038369	0.11
6.0	5.45	3.55	0.024021	0.10
6.5	5.38	3.61	0.015912	0.095
7.0	5.33	3.67	0.011003	0.088
7.5	5.28	3.72	0.007867	0.082
8.0	5.24	3.76	0.005804	0.078
8.5	5.20	3.80	0.004379	0.075
9.0	5.16	3.84	0.003369	0.073
9.5	5.13	3.87	0.002638	0.070
10.0	5.10	3.90	0.002096	0.068
10.5	5.07	3.93	0.001685	0.065
11.0	5.05	3.95	0.001365	0.065
11.5	5.02	3.97	0.001128	0.064
12.0	5.00	4.00	0.000905	0.063
12.5	4.98	4.01	0.00075741	0.060
13.0	4.97	4.03	0.0006327	0.059
13.5	4.95	4.05	0.00053185	0.057
14.0	4.94	4.06	0.00044954	0.056
14.5	4.92	4.08	0.00037815	0.056
15.0	4.91	4.09	0.00031952	0.055
15.5	3.90	3.10	0.00027045	0.055
16.0	3.88	3.11	0.00022847	0.052
16.5	3.87	3.13	0.00019351	0.050
17.0	3.86	3.14	0.0001633	0.048
17.5	4.35	3.65	0.00013939	0.042
18.0	3.84	3.16	0.00011626	0.041

Table S5. Best fit parameters of the Debye model to the $\chi''_M = f$ (frq) behaviors for **3**. χ_T stands for the isothermal susceptibility, χ_s for the adiabatic susceptibility, τ for the relaxation time, and α accounts for the distribution width of the relaxation time.

Figure S12. Magnetic behavior for **3'**. The composition in Dy^{III} for the sample was estimated by adjusting the magnetization reached for **3'** under 5 t at 2 K (blue plot in (a)) to that of **3** at same temperature (red plot).

Table S4. Best fit parameters of the Debye model to the $\chi''_M = f(\text{frq})$ behaviors for 3'. χ_T stands for the isothermal susceptibility, χ s for the adiabatic susceptibility, τ for the relaxation time, and α accounts for the distribution width of the relaxation time.

T(K)	$\chi_{\scriptscriptstyle \rm T}$	χ_{s}	τ (s)	α
4.0	5.825	3.175	0.166	0.16
4.5	5.6816	3.3184	0.080	0.11
5.0	5.5944	3.4056	0.04511	0.062
5.5	5.5049	3.4951	0.02776	0.064
6.0	5.4324	3.5677	0.0178	0.0558
6.5	5.3645	3.6355	0.01177	0.0573
7.0	5.3137	3.6863	0.00817	0.052
7.5	5.2669	3.7331	0.00577	0.048
8.0	5.2233	3.7767	0.00417	0.044
8.5	5.1878	3.8122	0.00306	0.034
9.0	5.1514	3.8486	0.00230	0.035
9.5	5.119	3.881	0.001747	0.037
10.0	5.0922	3.9078	0.001341	0.027
10.5	5.0658	3.9342	0.001049	0.032
11.0	5.0435	3.9565	0.000825	0.026
11.5	5.0224	3.9776	0.000656	0.023
12.0	5.001	3.999	0.000517	0.018
12.5	4.9822	4.0178	0.000421	0.023
13.0	4.9661	4.0339	0.0003419	0.020
13.5	4.9513	4.0487	0.0002823	0.011
14.0	4.9367	4.0633	0.000233	0.008
14.5	4.9235	4.0765	0.0001895	0.013
15.0	4.9127	4.0873	0.000156	0.01
15.5	4.9004	4.0996	0.000128	6e-5
16.0	4.89	4.11	0.0001045	0.002
16.5	4.8769	4.1231	8.57e-5	0.0005

Figure S12. AC susceptibiliy behavior for **4**.

Table S7. Best fit parameters of the Debye model to the $\chi''_M = f(\text{frq})$ behaviors for **4**. χ_T stands for the isothermal susceptibility, χ_s for the adiabatic susceptibility, τ for the relaxation time, and α accounts for the distribution width of the relaxation time.

T(K)	$\chi_{\scriptscriptstyle \rm T}$	$\chi_{\rm s}$	τ (s)	α
4.0	5.825	3.175	0.166	0.16
4.5	5.6816	3.3184	0.080	0.11
5.0	5.5944	3.4056	0.04511	0.062
5.5	5.5049	3.4951	0.02776	0.064
6.0	5.4324	3.5677	0.0178	0.0558
6.5	5.3645	3.6355	0.01177	0.0573
7.0	5.3137	3.6863	0.00817	0.052
7.5	5.2669	3.7331	0.00577	0.048
8.0	5.2233	3.7767	0.00417	0.044
8.5	5.1878	3.8122	0.00306	0.034
9.0	5.1514	3.8486	0.00230	0.035
9.5	5.119	3.881	0.001747	0.037
10.0	5.0922	3.9078	0.001341	0.027
10.5	5.0658	3.9342	0.001049	0.032
11.0	5.0435	3.9565	0.000825	0.026
11.5	5.0224	3.9776	0.000656	0.023
12.0	5.001	3.999	0.000517	0.018
12.5	4.9822	4.0178	0.000421	0.023
13.0	4.9661	4.0339	0.0003419	0.020
13.5	4.9513	4.0487	0.0002823	0.011
14.0	4.9367	4.0633	0.000233	0.008
14.5	4.9235	4.0765	0.0001895	0.013
15.0	4.9127	4.0873	0.000156	0.01
15.5	4.9004	4.0996	0.000128	6e-5
16.0	4.89	4.11	0.0001045	0.002
16.5	4.8769	4.1231	8.57e-5	0.0005

Table S8. Basis set used for all the elements in our calculation.

Figure S14. Computed g_{zz} orientation in complex 6 (a) and possible magnetization relaxation pathways (b). The red arrows show the QTM and TA-QTM via ground and higher excited KD respectively. The sky dotted arrows show the Orbach process for the relaxation. The green arrows show the possible mechanism of magnetic relaxation. The blue characters imply the m_J composition of the KDs.

Figure S15. Atom number for Loprop charge of complex (A) **3** (B) **4** (C) **6**. Colour code: Dygreenish yellow, Cl-green, P-violet, O-red, N-blue, C-grey. Hydrogens are removed for clarity.

Atom	3	4	6
Dy	2.5006	2.4993	2.4656
O6/O5/Cl1	-1.1346	-1.1107	-0.8799
Cl ₂	-0.8781	-0.8753	-0.8713
04/04/05	-0.8314	-0.8228	-0.8236
05/06/031	-0.8242	-0.8283	-0.8239
N7/N7/N4	-0.3882	-0.2262	-0.3726
N8/N8/N7	-0.2568	-0.2542	-0.2660
N10/N9/N33	-0.2497	-0.3908	-0.2662
Av. axial/av. Equatorial	1.9730	1.9684	1.7153

Table S9. LoProp charge of the metal center and its coordinated atom of complex **3**, **4** and **6**.

Table S10. The *ab initio* computed crystal field parameter of complex **3**, **4** and **6**.

k	Q		B_k^q	
		3	4	6
	-2	1.20E-01	$-1.63E-02$	$-1.52E-02$
	-1	$-5.59E-02$	$-2.09E-02$	1.34E-02
2	0	$-2.18E + 00$	$-2.32E+00$	$-2.80E+00$
	1	$-1.13E-01$	$-7.01E - 01$	4.89E-03
	2	$-2.14E-01$	2.34E-01	1.77E+00
	-4	1.70E-02	$-2.97E-03$	1.26E-03
	-3	3.73E-04	$-5.85E-04$	4.58E-04
	-2	$-1.22E - 02$	2.32E-03	$-8.19E - 05$
	-1	$-3.69E-04$	$-1.81E-04$	1.91E-04
4	0	$-1.57E-03$	$-1.41E-03$	$-8.93E - 04$
	1	$-7.93E - 04$	2.33E-03	$-2.42E - 07$
	$\overline{2}$	4.28E-02	4.22E-02	4.04E-02
	3	6.25E-03	$-2.79E-03$	$-7.15E - 05$
	4	$-3.02E-02$	$-3.35E-02$	$-2.75E-02$
	-6	$-4.13E-05$	2.50E-05	1.59E-05
	-5	$-8.33E-06$	1.29E-05	$-1.89E-06$
	-4	$-2.58E-05$	2.40E-05	8.27E-06
	-3	4.59E-05	1.16E-05	$-4.35E-06$
	-2	7.59E-05	2.63E-05	2.53E-05
	$^{\rm -1}$	3.82E-05	1.42E-05	$-7.95E - 06$
6	0	$-5.50E-05$	$-5.30E - 05$	$-4.64E-05$
	1	9.50E-06	2.02E-05	$-4.09E - 07$
	2	$-4.03E-04$	$-4.02E-04$	$-4.20E - 04$
	3	$-7.82E - 05$	$-3.79E-05$	3.72E-07
	4	6.56E-06	7.68E-06	5.92E-05
	5	3.19E-05	6.33E-05	$-3.51E-07$
	6	$-1.06E-05$	3.32E-05	$-4.81E-05$

Table S11. The *ab initio* computed energy and the associated g-tensors of the low lying KDs generated from ⁶H15/2 state of complexes **3a**, **4a** and **6a**.

Table S12. LoProp charge of the metal center and its coordinated atom of complex **3a**, **4a** and **6a**.

k	Q		B_k^q	
		3a	4a	6a
	-2	9.10E-02	2.95E-02	3.18E-07
	-1	$-9.54E-02$	6.71E-02	$-2.31E-08$
$\overline{2}$	0	$-1.21E + 01$	$-1.18E + 01$	$-1.10E + 01$
	1	5.94E-02	$-1.57E-02$	9.50E-02
	$\overline{2}$	2.24E-01	1.81E-01	$-4.22E - 01$
	-4	$-2.18E - 04$	$-8.87E - 06$	$-3.74E-09$
	-3	$-1.11E-03$	6.57E-04	5.79E-10
	-2	1.22E-04	$-1.63E-05$	$-8.69E-10$
	-1	1.82E-03	$-5.31E-04$	1.48E-10
4	0	$-6.49E - 03$	$-6.26E - 03$	$-4.60E-03$
	1	$-3.44E - 04$	4.15E-05	$-4.11E-04$
	$\overline{2}$	$-4.40E - 04$	$-7.07E - 04$	1.10E-03
	3	1.62E-04	2.90E-04	$-5.07E-04$
	4	$-1.44E-03$	$-1.42E-03$	2.49E-03
	-6	7.25E-06	4.93E-07	$2.17E-10$
	-5	$-1.04E-05$	4.44E-06	$-2.51E-10$
	-4	$-3.17E-06$	$-1.78E-07$	$-6.33E-11$
	-3	$-6.28E - 06$	4.31E-06	3.35E-11
	-2	$-3.51E-06$	1.24E-07	3.57E-12
	-1	$-1.95E-05$	1.86E-06	4.14E-13
6	0	4.16E-05	4.05E-05	2.85E-05
	1	$-2.99E-06$	1.34E-06	$-4.29E - 06$
	$\overline{2}$	$-1.31E-06$	4.71E-06	$-5.29E-06$
	3	$-3.64E-06$	$-1.01E-06$	$-3.47E - 05$
	4	$-2.27E-05$	$-2.39E - 05$	4.47E-05
	5	1.68E-05	$-2.43E-06$	1.39E-04
	6	3.38E-05	3.37E-05	$-1.02E-04$

Table S13. The *ab initio* computed crystal field parameter of complex **3a**, **4a** and **6a**.

Table S14. The *ab initio* computed energy with the associated g-tensors and tunneling splitting (Δ_{tun}) of the low lying energy levels of complexes 2.

Complex 5						
Energy (K)	g_{x}	g _v	g_{z}			
0.0	0.089	1.396	14.176			
20.7	1.608	3.471	12.510			
50.3	0.220	4.998	11.662			
185.5	3.176	4.926	9.333			
320.7	0.882	2.076	12.297			
403.1	1.227	3.722	14.512			
610.4	0.026	0.672	15.353			
632.3	0.229	0.671	15.370			

Table S15. The *ab initio* computed energy (cm⁻¹) with the associated g-tensors of the eight low lying KDs of complexes **5**.

Figure S16. (a) Anisotropy axis of **3a.** (b) Anisotropy axis of **4a**. (c) Anisotropy axis of **6a**. Colour code: Dy-greenish yellow, Cl-green, P-orange, O-red, N-blue, C-grey, H-white.

Figure S17. Mechanism of magnetic relaxation of (a) **3a** (b) **4a** (c) **6a**. The red arrows show the QTM and TA-QTM via ground and higher excited KD respectively. The sky dotted arrows show the Orbach process for the relaxation. The green arrows show the mechanism of magnetic relaxation. The blue characters imply the m_J composition of the KDs.

Figure S18. Computed g_{zz} orientation in complex 2 and 5 $(a \& b)$ and possible magnetization relaxation pathways (c & d) respectively. Colour code: Er-Pink, Tb-Dark violet, Cl-green, Porange, O-red, N-blue, C-grey, H-white. The red dotted arrows show the QTM and TA-QTM via ground and higher excited KD respectively. The sky dotted arrows show the Orbach process for the relaxation. The green arrows show the mechanism of magnetic relaxation. The blue characters imply the m_J composition of the KDs.