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# Finite time stabilization of continuous inertial dynamics combining dry friction with Hessian-driven damping

Samir ADLY\* and Hedy ATTOUCH†

April 29, 2020

*This paper is dedicated to Professor Umberto Mosco on the occasion of his 80th birthday.*

**ABSTRACT.** In a Hilbert space  $\mathcal{H}$ , we study the stabilization in finite time of the trajectories generated by a continuous (in time  $t$ ) damped inertial dynamic system. The potential function  $f : \mathcal{H} \rightarrow \mathbb{R}$  to be minimized is supposed to be differentiable, not necessarily convex. It enters the dynamic via its gradient. The damping results from the joint action of dry friction, viscous friction, and a geometric damping driven by the Hessian of  $f$ . The dry friction damping function  $\phi : \mathcal{H} \rightarrow \mathbb{R}_+$ , which is convex with a sharp minimum at the origin (typically  $\phi(x) = r\|x\|$  with  $r > 0$ ), enters the dynamic via its subdifferential. It acts as a soft threshold operator on the velocities, and is at the origin of the stabilization property in finite time. The Hessian driven damping, which enters the dynamics in the form  $\nabla^2 f(x(t))\dot{x}(t)$ , permits to control and attenuate the oscillations which occur naturally with the inertial effect. We give two different proofs, in a finite dimensional setting, of the existence of strong solutions of this second-order differential inclusion. One is based on a fixed point argument and the use of Leray-Schauder theorem, the other one is based on the Yosida approximation technique and the Mosco convergence. Then, we study the convergence properties of the trajectories as  $t \rightarrow +\infty$ , and show their stabilization property in finite time. The convergence results tolerate the presence of perturbations, errors, under the sole assumption of their asymptotic convergence to zero. Then, we extend the study to the case of a nonsmooth convex function  $f$ .

*Mathematics Subject Classifications:* 37N40, 34A60, 34G25, 49K24, 70F40.

*Key words and phrases:* damped inertial dynamics; differential inclusion; dry friction; Hessian-driven damping; finite time stabilization;

## 1 Introduction and preliminary results

Throughout the paper  $\mathcal{H}$  is a real Hilbert space, with the scalar product  $\langle \cdot, \cdot \rangle$  and the associated norm  $\|\cdot\|$ , and  $f : \mathcal{H} \rightarrow \mathbb{R}$  is a  $\mathcal{C}^1$  function (not necessarily convex) whose gradient is Lipschitz continuous. In the introduction, when we consider continuous dynamics in which the Hessian intervenes, to simplify the presentation, we can assume that  $f$  is a  $\mathcal{C}^2$  function.

### 1.1 Presentation of the dynamic

We will analyze the finite time stabilization of the following second-order differential inclusion

$$(IGDH) \quad \ddot{x}(t) + \gamma\dot{x}(t) + \partial\phi(\dot{x}(t)) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) \ni 0, \quad t \in [t_0, +\infty[ \quad (1.1)$$

where (IGDH) stands shortly for Inertial Gradient system with Dry friction and Hessian-driven damping.

Let us briefly describe the three types of damping that occur in (IGDH), and which provide stabilization:

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\*Laboratoire XLIM, Université de Limoges, 87060 Limoges, France. E-mail: samir.adly@unilim.fr

†IMAG, Université Montpellier, CNRS, 34095 Montpellier CEDEX 5, France. E-mail: hedy.attouch@umontpellier.fr

- $\gamma > 0$  is a positive viscous damping coefficient.
- The dry friction damping function  $\phi : \mathcal{H} \rightarrow \mathbb{R}_+$  is convex with a sharp minimum at the origin, typically  $\phi(x) = r\|x\|$  with  $r > 0$ . It is at the origin of the finite time stabilization. It enters the dynamic via its subdifferential, which is defined by

$$\partial\phi(x) = \{p \in \mathcal{H} : \langle p, y - x \rangle \leq \phi(y) - \phi(x), \forall y \in \mathcal{H}\}.$$

- The geometrical damping driven by the Hessian of  $f$  enters the dynamics in the form  $\nabla^2 f(x(t))\dot{x}(t)$ . It allows to attenuate and control the oscillations which naturally occur with inertial systems.  $\beta > 0$  is the corresponding damping coefficient.

Let us mention that, from an abstract point of view, the viscous friction  $\gamma\dot{x}(t)$  could enter the dynamic (IGDH) into two ways:

- since  $\gamma\dot{x}(t) + \partial\phi(\dot{x}(t)) = \partial(\frac{\gamma}{2}\|\cdot\|^2 + \phi)(\dot{x}(t))$ , it is possible to introduce a new dry friction function  $\psi = \frac{\gamma}{2}\|\cdot\|^2 + \phi$ .
- since  $\gamma\dot{x}(t) + \beta\nabla^2 f(x(t))\dot{x}(t) = \frac{d}{dt}\left(\nabla\left(\frac{\gamma}{2}\|\cdot\|^2 + \beta f\right)\right)(x(t))$ , it is possible to combine the viscous friction with the damping driven by the Hessian.

As we have in mind the optimization algorithms, and in order to place our results in the context of the literature on the subject, we will keep the three frictions separated in the dynamic (IGDH).

### 1.1.1 Link with optimization

Our main motivation for studying the continuous dynamic (IGDH) comes from optimization. Its temporal discretization provides algorithms which share its good convergence properties. As a key property, we use that  $\nabla^2 f(x(t))\dot{x}(t)$  is the time derivative of  $\nabla f(x(t))$ . This gives, in discretized form, first-order algorithms, which we describe below. Given a constant time step  $h > 0$ , we consider the following temporal discretization of (IGDH)

$$\begin{aligned} & \frac{1}{h^2}(x_{k+1} - 2x_k + x_{k-1}) + \frac{\gamma}{h}(x_{k+1} - x_k) + \partial\phi\left(\frac{1}{h}(x_{k+1} - x_k)\right) \\ & + \frac{\beta}{h}(\nabla f(x_k) - \nabla f(x_{k-1})) + \nabla f(x_k) \ni 0. \end{aligned} \quad (1.2)$$

In (1.2), the terms  $\frac{1}{h^2}(x_{k+1} - 2x_k + x_{k-1})$  and  $\frac{1}{h}(x_{k+1} - x_k)$  respectively represent the discrete acceleration, and the discrete speed. According to  $\nabla^2 f(x(t))\dot{x}(t) = \frac{d}{dt}\nabla f(x(t))$ , the correcting term  $\frac{1}{h}(\nabla f(x_k) - \nabla f(x_{k-1}))$  is directly linked to the temporal discretization of the Hessian-driven damping term. By solving (1.2) with respect to  $x_{k+1}$ , we obtain the following first-order algorithm where the dry friction enters via the proximal operator of  $\phi$ , and the function to minimize  $f$  enters via its gradient:

$$(\text{IPAHDD}) \begin{cases} z_k = \frac{1}{h(1+h\gamma)}(x_k - x_{k-1}) - \frac{\beta}{1+h\gamma}(\nabla f(x_k) - \nabla f(x_{k-1})) - \frac{h}{1+h\gamma}\nabla f(x_k) \\ x_{k+1} = x_k + h \operatorname{prox}_{\frac{h}{1+h\gamma}\phi}(z_k), \end{cases}$$

(IPAHDD) stands shortly for Inertial Proximal-gradient Algorithm with Hessian-Damping and Dry friction. Note that the temporal discretization (1.2) of (IGDH) is implicit with respect to the nonsmooth function  $\phi$ , and explicit with respect to the smooth function  $f$ . In the above formula,  $\operatorname{prox}_{\phi}$  denotes the proximal mapping associated with the convex function  $\phi$ . Recall that, for any  $x \in \mathcal{H}$ , for any  $\lambda > 0$

$$\operatorname{prox}_{\lambda\phi}(x) := \operatorname{argmin}_{\xi \in \mathcal{H}} \left\{ \lambda\phi(\xi) + \frac{1}{2}\|x - \xi\|^2 \right\}.$$

The proximal mapping can be equivalently formulated as the resolvent of the maximally monotone operator  $\partial\phi : \mathcal{H} \rightarrow 2^{\mathcal{H}}$ , that is, for any  $\lambda > 0$

$$\text{prox}_{\lambda\phi} = (I + \lambda\partial\phi)^{-1}.$$

In [3], it is proved the finite convergence of the sequences  $(x_k)$  generated by the (IPAHDD) algorithm. Indeed, for continuous optimization algorithms, convergence in a finite number of steps, with the estimation of this number, is a very favorable property. The limit  $x_\infty$  of the sequence  $(x_k)$  satisfies

$$\nabla f(x_\infty) + \partial\phi(0) \ni 0.$$

Thus,  $x_\infty$  is an “approximate” critical point of  $f$ . In practice, we choose  $\phi$  with  $\partial\phi(0)$  ”small”. By taking  $\phi(x) = r\|x\|$ , this means taking a small  $r > 0$ . This amounts to solving the optimization problem  $\min_{\mathcal{H}} f$  with the variational principle of Ekeland, instead of the Fermat rule.

(IPAHDD) is a splitting algorithm in the sense that the two constitutive potential functions  $f$  and  $\phi$  are treated separately, via the gradient and the proximal mapping respectively. Since the proximal mapping of  $\phi$  can be easily calculated in most practical situations, this makes (IPAHDD) a first-order algorithm that can be used to deal with problems of large size.

The continuous dynamic (IGDH) is the basis of these algorithmic developments. Its mechanical interpretation naturally suggests Lyapunov functions. This therefore leads us to study (IGDH) in depth, and explore some of its extensions. This is the main purpose of this article.

### 1.1.2 Link with mechanics, damped shocks

Another motivation for the study of (IGDH) comes from mechanics, and the modeling of damped shocks. In [18], Attouch-Maingé-Redont consider the second-order differential system with Hessian-driven damping

$$\ddot{x}(t) + \gamma\dot{x}(t) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) + \nabla g(x(t)) = 0, \quad (1.3)$$

where  $g : \mathcal{H} \rightarrow \mathbb{R}$  is a smooth real-valued function. An interesting property of this system is that, after introduction of an auxiliary variable  $y$ , it can be equivalently written as a first-order system involving only the time derivatives  $\dot{x}(t)$ ,  $\dot{y}(t)$  and the gradient terms  $\nabla f(x(t))$ ,  $\nabla g(x(t))$ . More precisely, the system (1.3) is equivalent to the following first order differential equation

$$\begin{cases} \dot{x}(t) + \beta\nabla f(x(t)) + ax(t) + by(t) = 0, \\ \dot{y}(t) - \beta\nabla g(x(t)) + ax(t) + by(t) = 0, \end{cases} \quad (1.4)$$

where  $a$  and  $b$  are real numbers such that:  $a + b = \gamma$  and  $\beta b = 1$ . Note that (1.4) is different from the classical Hamiltonian formulation, which would still involve the Hessian of  $f$ . In contrast, the formulation (1.4) uses only first-order information from the function  $f$  (no occurrence of the Hessian of  $f$ ). Replacing  $\nabla f$  by  $\partial f$  in (1.4) allows us to extend the analysis to the case of a convex lower semicontinuous function  $f : \mathcal{H} \rightarrow \mathbb{R} \cup \{+\infty\}$ , and so to introduce constraints in the model. When  $f = \delta_K$  is the indicator function of a closed convex set  $K \subset \mathcal{H}$ , the subdifferential operator  $\partial f$  takes account of the contact forces, while  $\nabla g$  takes account of the driving forces. In this setting, by playing with the geometrical damping parameter  $\beta$ , one can describe nonelastic shock laws with restitution coefficient (for more details we refer to [18] and references therein). Introducing the dry friction into this dynamic would allow to obtain finite type stabilization of damped oscillating systems with obstacle constraint.

## 1.2 Some historical facts

Let us explain the role and the importance of each of the three damping terms which enter into the continuous dynamics (IGDH). Our main contribution in this article is to show how to combine them to obtain fast stabilization properties.

### 1.2.1 Viscous friction

The use of inertial dynamics to accelerate the gradient method in optimization was first considered by B. Polyak in [35]. Based on the inertial system with a fixed viscous damping coefficient  $\gamma > 0$

$$(HBF) \quad \ddot{x}(t) + \gamma \dot{x}(t) + \nabla f(x(t)) = 0,$$

he introduced the Heavy Ball with Friction method. This system was further developed by Attouch-Goudou-Redont [17] as a tool to explore the local minima of  $f$ . For a strongly convex function  $f$ , and  $\gamma$  judiciously chosen, (HBF) provides convergence at exponential rate of  $f(x(t))$  to  $\min_{\mathcal{H}} f$ . For a general convex function  $f$ , the asymptotic convergence rate of (HBF) is  $\mathcal{O}(\frac{1}{t})$  (in the worst case). This is however not better than the steepest descent. A decisive step to obtain a faster asymptotic convergence was taken by Su-Boyd-Candès [39] with the introduction of an Asymptotic Vanishing Damping coefficient  $\gamma(t) = \frac{\alpha}{t}$ , that is

$$(AVD)_{\alpha} \quad \ddot{x}(t) + \frac{\alpha}{t} \dot{x}(t) + \nabla f(x(t)) = 0.$$

As a specific feature, the viscous damping coefficient  $\frac{\alpha}{t}$  vanishes (tends to zero) as time  $t$  goes to infinity, hence the terminology. For a general convex function  $f$ , it provides a continuous version of the accelerated gradient method of Nesterov. For  $\alpha \geq 3$ , each trajectory  $x(\cdot)$  of  $(AVD)_{\alpha}$  satisfies the asymptotic convergence rate of the values  $f(x(t)) - \inf_{\mathcal{H}} f = \mathcal{O}(1/t^2)$  as  $t \rightarrow +\infty$ . The convergence properties of the dynamic  $(AVD)_{\alpha}$  have been the subject of many recent studies, see [7, 10, 11, 12, 14, 16, 19, 22, 23, 32, 39]. The case  $\alpha = 3$ , which corresponds to Nesterov's historical algorithm, is critical. In the case  $\alpha = 3$ , the question of the convergence of the trajectories remains an open problem (except in one dimension where convergence holds [16]). For  $\alpha > 3$ , it has been shown by Attouch-Chbani-Peypouquet-Redont [14] that each trajectory converges weakly to a minimizer. The corresponding algorithmic result has been obtained by Chambolle-Dossal [30]. For  $\alpha > 3$ , it is shown in [19] and [32] that the asymptotic convergence rate of the values is actually  $o(1/t^2)$ . The subcritical case  $\alpha \leq 3$  has been examined by Apidopoulos-Aujol-Dossal [7] and Attouch-Chbani-Riahi [16], with the convergence rate of the objective values  $\mathcal{O}\left(t^{-\frac{2\alpha}{3}}\right)$ . These rates are optimal, that is, they can be reached, or approached arbitrarily close.

### 1.2.2 Dry friction

Some first results concerning the finite convergence property under the action of dry friction have been obtained by Adly-Attouch-Cabot [4] for the continuous dynamics

$$\ddot{x}(t) + \partial\phi(\dot{x}(t)) + \nabla f(x(t)) \ni 0, \quad t \in [t_0, +\infty[. \quad (1.5)$$

Assuming that the potential friction function  $\phi$  has a sharp minimum at the origin (dry friction), they showed that, generically with respect to the initial data, the solution trajectories converge in finite time to equilibria. Similar results for the corresponding proximal-based algorithms have been obtained by Baji-Cabot [24] and Adly-Attouch [3].

Let's make precise the tools that will be useful for the mathematical analysis of the set-valued term  $\partial\phi(\dot{x}(t))$  in (1.5) which models dry friction. We say that the friction potential function  $\phi$  has the Dry Friction property (denoted by (DF)) if it satisfies the following properties

$$(DF) \quad \begin{cases} \phi : \mathcal{H} \rightarrow \mathbb{R} \text{ is convex continuous;} \\ \min_{\xi \in \mathcal{H}} \phi(\xi) = \phi(0) = 0; \\ 0 \in \text{int}(\partial\phi(0)). \end{cases}$$

The particular case  $\phi = r\|\cdot\|$ , with  $r > 0$ , models dry friction (also called Coulomb friction) in mechanics. The key assumption  $0 \in \text{int}(\partial\phi(0))$  expresses that  $\phi$  has a sharp minimum at the origin. This is specified

in the following elementary lemma, see [1, Lemma 4.1 page 83], where, in item (iv),  $\phi^*$  is the Fenchel conjugate of  $\phi$ .

**Lemma 1.1** *Let  $\phi : \mathcal{H} \rightarrow \mathbb{R}$  be a convex continuous function such that  $\min_{\xi \in \mathcal{H}} \phi(\xi) = \phi(0) = 0$ . Then, the following formulations of the dry friction are equivalent:*

- (i)  $0 \in \text{int}(\partial\phi(0))$ ;
- (ii) *there exists some  $r > 0$  such that  $B(0, r) \subset \partial\phi(0)$ ;*
- (iii) *there exists some  $r > 0$  such that, for all  $\xi \in \mathcal{H}$ ,  $\phi(\xi) \geq r\|\xi\|$ .*
- (iv) *there exists some  $r > 0$  such that,  $\|g\| \leq r \implies \partial\phi^*(g) \ni 0$ .*

The positive parameter  $r$  will play a crucial role in our analysis. To enlighten its role, we will say that the friction potential function  $\phi$  satisfies the property  $(DF)_r$  if  $\phi$  satisfies the Dry Friction property (DF) with  $B(0, r) \subset \partial\phi(0)$ . The property (iv) above expresses that, when the force  $g$  exerted on the system is less than a threshold  $r > 0$ , then the system stabilizes, *i.e.* the velocity  $v = 0 \in \partial\phi^*(g)$ . This contrasts with the viscous damping that can asymptotically produce many small oscillations.

### 1.2.3 Hessian-driven damping

The inertial system

$$(\text{DIN})_{\gamma, \beta} \quad \ddot{x}(t) + \gamma\dot{x}(t) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) = 0,$$

was introduced in [6]. In the same spirit as (HBF), the dynamic  $(\text{DIN})_{\gamma, \beta}$  contains a *fixed* positive friction coefficient  $\gamma > 0$ . The introduction of the Hessian-driven damping allows to damp the transversal oscillations that might arise with (HBF), as observed in [6] in the case of the Rosenbrock function. The need to take a geometric damping adapted to  $f$  had already been observed by Alvarez [5] who considered the inertial system

$$\ddot{x}(t) + \Gamma\dot{x}(t) + \nabla f(x(t)) = 0,$$

where  $\Gamma : \mathcal{H} \rightarrow \mathcal{H}$  is a linear positive anisotropic operator. But still this damping operator is fixed. For a general convex function, the Hessian-driven damping in  $(\text{DIN})_{\gamma, \beta}$  performs a similar operation in a closed-loop adaptive way. The terminology (DIN) stands shortly for Dynamical Inertial Newton. It refers to the natural link between this dynamic and the continuous Newton method, see Attouch-Svaiter [21]. Recent studies have been devoted to the study of the inertial dynamic

$$\ddot{x}(t) + \frac{\alpha}{t}\dot{x}(t) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) = 0,$$

which combines asymptotic vanishing damping with Hessian-driven damping. The corresponding algorithms involve a correcting term in the Nesterov accelerated gradient method which reduces the oscillatory aspects, see Attouch-Peypouquet-Redont [20], Attouch-Chbani-Fadili-Riahi [13], Shi-Du-Jordan-Su [38].

### 1.3 Mosco convergence and Attouch theorem

In 1969, U. Mosco [33] introduced a fundamental concept of convergence for sequence of closed convex sets in reflexive Banach spaces, that can be transposed to functions via their epigraphs. This concept, called the Mosco convergence, is widely applicable in convex optimization, stability of variational inequalities and optimal control. Precisely, a sequence  $(\varphi_n) \in \Gamma_0(\mathcal{H})$  of convex proper and lower semicontinuous functions, is declared to be Mosco-convergent to  $\varphi \in \Gamma_0(\mathcal{H})$  provided, (i) and (ii) are satisfied:

- (i) for all  $x \in \text{dom}\varphi$ , there exists a sequence  $(x_n)$  strongly convergent to  $x$  s.t.  $\lim_{n \rightarrow +\infty} \varphi_n(x_n) = \varphi(x)$
- (ii) whenever  $(x_n)$  converges weakly to  $x$ , we have  $\liminf_{n \rightarrow +\infty} \varphi_n(x_n) \geq \varphi(x)$ .

An other important result was given by U. Mosco in [34] where the bicontinuity of the Fenchel conjugate is established: for any sequence of convex functions  $(\varphi_n)$  we have the equivalence:

$$\varphi_n \text{ is Mosco-convergent to } \varphi \iff (\varphi_n^*) \text{ is Mosco-convergent to } \varphi^*.$$

This notion was used in 1977 by H. Attouch (see Theorem 3.66 in [8] page 373) to establish that a sequence  $(\varphi_n) \subset \Gamma_0(\mathcal{H})$  converges in the sense of Mosco to  $\varphi \in \Gamma_0(\mathcal{H})$  if and only if  $\partial\varphi_n$  converges in the sense of Painlevé-Kuratowski to  $\partial\varphi$ , provided some normalization condition is satisfied. Here the notation  $\partial\varphi$  denotes the convex subdifferential of  $\varphi \in \Gamma_0(\mathcal{H})$ , and the operators are identified with their graphs.

## 1.4 Contents

The paper is organized as follows. In sections 2 and 3, we give two different proofs of the existence of solutions to (IGDH): they are based respectively on the Leray-Schauder fixed point theorem, and the Yosida approximation. In section 4, we present our main results, which concern the convergence in finite time of the trajectories generated by the inertial dynamic (IGDH). In section 5, we examine the effect of the introduction of perturbations, errors in the dynamic (IGDH). In section 6, based on the variational properties of Moreau's envelope, we extend these results to the case where  $f : \mathcal{H} \rightarrow \mathbb{R} \cup \{+\infty\}$  is a convex lower semicontinuous and proper function such that  $\inf f > -\infty$ . We complete the paper with some perspectives.

## 2 Existence and Uniqueness for the Cauchy problem

### 2.1 The notion of strong solution of (IGDH)

Unless specified, we make the following assumptions on the potential functions  $f$  and  $\phi$ :

- $f : \mathcal{H} \rightarrow \mathbb{R}$  is a differentiable function whose gradient is Lipschitz continuous;
- $\phi : \mathcal{H} \rightarrow \mathbb{R}$  is a convex continuous function which satisfies (DF) as well as the following property: there exists a constant  $c > 0$  such that for all  $x \in \mathcal{H}$

$$\|\partial\phi(x)\| \leq c(1 + \|x\|).$$

where the above formula means: for any  $\xi \in \partial\phi(x)$ , we have  $\|\xi\| \leq c(1 + \|x\|)$ .

We call  $t_0 \in \mathbb{R}$  the origin of time. Since (IGDH) is an autonomous system, we can take an arbitrary real number for  $t_0$ .

**Definition 2.1** *We say that  $x : [t_0, +\infty[ \rightarrow \mathcal{H}$  is a strong global solution of (IGDH) if the following properties 1 and 2 are satisfied:*

1. For all  $T > t_0$ ,  $x \in \mathbb{H}^2(t_0, T; \mathcal{H})$ . Equivalently,  $x \in \mathbb{L}^2(t_0, T; \mathcal{H})$ ,  $\dot{x} \in \mathbb{L}^2(t_0, T; \mathcal{H})$  and  $\ddot{x} \in \mathbb{L}^2(t_0, T; \mathcal{H})$ .
2. The following inclusion is satisfied: for almost all  $t \geq t_0$

$$\ddot{x}(t) + \gamma\dot{x}(t) + \partial\phi(\dot{x}(t)) + \beta \frac{d}{dt} \left( \nabla f(x(t)) \right) + \nabla f(x(t)) \ni 0. \quad (2.1)$$

Let us comment the above definition. First note that, since  $\nabla f$  is Lipschitz continuous (say with constant  $L$ ), we have for all  $s, t \geq t_0$

$$\|\nabla f(x(t)) - \nabla f(x(s))\| \leq L\|x(t) - x(s)\|.$$

Since  $x \in \mathbb{H}^1(t_0, T; \mathcal{H})$ ,  $x$  is absolutely continuous. According to the above inequality, this implies that  $t \mapsto \nabla f(x(t))$  is also absolutely continuous. Therefore,  $t \mapsto \nabla f(x(t))$  is almost everywhere differentiable. Using again the above inequality, this gives

$$\left\| \frac{d}{dt} \left( \nabla f(x(t)) \right) \right\| \leq L \|\dot{x}(t)\|.$$

According to  $\dot{x} \in \mathbb{L}^2(t_0, T; \mathcal{H})$  this gives

$$\frac{d}{dt} \left( \nabla f(x(\cdot)) \right) \in \mathbb{L}^2(t_0, T; \mathcal{H})$$

where  $\frac{d}{dt} \left( \nabla f(x(t)) \right)$  can be taken indifferently in the distribution or pointwise sense (the two notions coincide for an absolutely continuous function). Without ambiguity, we will write shortly

$$\frac{d}{dt} \left( \nabla f(x(t)) \right) = \nabla^2 f(x(t)) \dot{x}(t)$$

which is justified when  $f$  is twice differentiable.

In addition, we have

$$\|\partial\phi(\dot{x}(t))\| \leq c(1 + \|\dot{x}(t)\|),$$

where we briefly write  $\partial\phi(\dot{x}(t))$  to designate a measurable selection  $\xi(t)$  with  $\xi(t) \in \partial\phi(\dot{x}(t))$  for almost every  $t \geq t_0$ . According to  $x \in \mathbb{H}^1(t_0, T; \mathcal{H})$  this gives

$$\partial\phi(x(\cdot)) \in \mathbb{L}^2(t_0, T; \mathcal{H}).$$

Therefore, for all  $T > t_0$ , all the constitutive elements of (IGDH) belong to  $\mathbb{L}^2(t_0, T; \mathcal{H})$ . Finally note that since for all  $T > t_0$ ,  $x \in \mathbb{H}^2(t_0, T; \mathcal{H})$ , we have that  $x$  and  $\dot{x}$  are continuous functions of  $t$ , which will allow to give a sense to the Cauchy data.

## 2.2 The case without the Hessian driven damping term, and with right handside

As a preliminary result, we consider (IGDH) in the case  $\beta = 0$ , *i.e.* without the Hessian driven damping term, but with a second member  $e(\cdot)$ . This is a preliminary step to study (IGDH) in the general case, using a fixed point argument. More precisely, we have the following equivalence, which is immediate to obtain, and which is just the Hamiltonian formulation of (IGDH).

**Proposition 2.1** *The following are equivalent: (i)  $\iff$  (ii)*

$$(i) \quad \ddot{x}(t) + \gamma(t)\dot{x}(t) + \partial\phi(\dot{x}(t)) + \nabla f(x(t)) \ni e(t).$$

$$(ii) \quad \begin{cases} \dot{x}(t) = u(t); \\ \dot{u}(t) \in -\nabla f(x(t)) - \gamma(t)u(t) + e(t) - \partial\phi(u(t)) \end{cases}$$



**Proof.** (i)  $\implies$  (ii). Set  $u(t) := \dot{x}(t)$ . We have  $\dot{u}(t) = \ddot{x}(t)$ , which, by the constitutive equation, gives

$$\dot{u}(t) + \gamma(t)\dot{x}(t) + \partial\phi(\dot{x}(t)) + \nabla f(x(t)) \ni e(t). \quad (2.2)$$

This gives the first-order system (ii).

(ii)  $\implies$  (i). By differentiating the first equation of (ii) we obtain  $\ddot{x}(t) = \dot{u}(t)$ . According to the second equation of (ii) we obtain (i). ■

We can now state the following existence and uniqueness result.

**Theorem 2.1** *Let  $f : \mathcal{H} \rightarrow \mathbb{R}$  be a differentiable function whose gradient is Lipschitz continuous, and let  $\phi : \mathcal{H} \rightarrow \mathbb{R} \cup \{+\infty\}$  be a convex lower semicontinuous proper function. Suppose that, for any  $T > t_0$ , the function  $\gamma : [t_0, +\infty[ \rightarrow \mathbb{R}_+$  belongs to  $L^1(t_0, T; \mathbb{R})$ , and  $e : [t_0, +\infty[ \rightarrow \mathcal{H}$  belongs to  $L^2(t_0, T; \mathcal{H})$ . Then, for any Cauchy data  $(x_0, \dot{x}_0) \in \mathcal{H} \times \mathcal{H}$ , there exists a unique strong global solution of*

$$\ddot{x}(t) + \gamma(t)\dot{x}(t) + \partial\phi(\dot{x}(t)) + \nabla f(x(t)) \ni e(t) \quad (2.3)$$

satisfying  $x(t_0) = x_0$ , and  $\dot{x}(t_0) = \dot{x}_0$ .

**Proof.** According to Proposition 2.1, it is equivalent to solve the first-order system (ii) with the Cauchy data  $x(t_0) = x_0$ ,  $u(t_0) = \dot{x}_0$ . Set

$$Z(t) = (x(t), u(t)) \in \mathcal{H} \times \mathcal{H}.$$

The system (ii) can be written equivalently as

$$\dot{Z}(t) + F(t, Z(t)) \ni 0, \quad Z(t_0) = (x_0, \dot{x}_0),$$

where  $F : [t_0, +\infty[ \times \mathcal{H} \times \mathcal{H} \rightrightarrows \mathcal{H} \times \mathcal{H}$ ,  $(t, x, u) \mapsto F(t, x, u)$  is defined by

$$F(t, x, u) = \left(0, \partial\phi(u)\right) + \left(-u, \nabla f(x) + \gamma(t)u - e(t)\right).$$

Hence  $F$  splits as follows

$$F(t, x, u) = F_1(x, u) + F_2(t, x, u)$$

where

$$F_1(x, u) = \partial\Phi(x, u)$$

is the subdifferential of the convex lower semicontinuous function (which depends only on  $u$ )

$$\Phi(x, u) = \phi(u)$$

and

$$F_2(t, x, u) = \left(-u, \nabla f(x) + \gamma(t)u - e(t)\right).$$

Therefore, the second-order dynamic (2.3) is equivalent to the following first-order differential inclusion

$$\dot{Z}(t) + \partial\Phi(Z(t)) + F_2(t, Z(t)) \ni 0, \quad Z(t_0) = (x_0, \dot{x}_0), \quad (2.4)$$

Since  $\nabla f$  is Lipschitz continuous, so is  $F_2$  with respect to  $Z = (x, u)$ . According to [28, Proposition 3.12] which considers the evolution equation governed by the Lipschitz perturbation (which can be time dependent as here) of a convex subdifferential, we deduce the existence and uniqueness of the solution to the Cauchy problem formulated as (ii), which gives the claim. ■

### 2.3 The general case of (IGDH)

In this section, we assume that  $\mathcal{H}$  is a finite dimensional Hilbert space.

**Theorem 2.2** *Let  $f : \mathcal{H} \rightarrow \mathbb{R}$  be a differentiable function whose gradient is  $L$ -Lipschitz continuous. Suppose that the dry potential friction  $\phi : \mathcal{H} \rightarrow \mathbb{R}$  is a convex continuous function that satisfies  $(DF)_r$ , and  $\|\partial\phi(x)\| \leq c(1 + \|x\|)$ . Suppose that the damping parameters satisfy*

$$\gamma > \beta L.$$

*Then, for any Cauchy data  $(x_0, \dot{x}_0) \in \mathcal{H} \times \mathcal{H}$ , there exists a strong global solution  $x : [t_0, +\infty[ \rightarrow \mathcal{H}$  of*

$$(IGDH) \quad \ddot{x}(t) + \gamma\dot{x}(t) + \partial\phi(\dot{x}(t)) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) \ni 0 \quad (2.5)$$

*satisfying the Cauchy data  $x(t_0) = x_0$  and  $\dot{x}(t_0) = \dot{x}_0$ .*

**Proof.** Since no geometrical assumption is made on the function  $f$  (it is not assumed to be convex), we will treat the Hessian driven damping term as a perturbation. According to the growth assumptions made on this term, we will apply a fixed point argument. To this end, let us write the equation (IGDH) in the following form

$$\ddot{x}(t) + \gamma\dot{x}(t) + \partial\phi(\dot{x}(t)) + \nabla f(x(t)) \ni -\beta\nabla^2 f(x(t))\dot{x}(t). \quad (2.6)$$

Then, to solve it, we will apply the Leray-Schauder fixed point theorem in the space  $\mathbb{H}^1(t_0, T; \mathcal{H})$ , with  $T \geq t_0$ . Recall the definition of the Sobolev space

$$\mathbb{H}^1(t_0, T; \mathcal{H}) = \{x \in \mathbb{L}^2(t_0, T; \mathcal{H}) \text{ such that } \dot{x} \in \mathbb{L}^2(t_0, T; \mathcal{H})\}$$

where  $\dot{x}$  is taken in the distribution sense.  $\mathbb{H}^1(t_0, T; \mathcal{H})$  is equipped with the norm

$$\|x\|_{\mathbb{H}^1(t_0, T; \mathcal{H})} = \left( \|x\|_{\mathbb{L}^2(t_0, T; \mathcal{H})}^2 + \|\dot{x}\|_{\mathbb{L}^2(t_0, T; \mathcal{H})}^2 \right)^{\frac{1}{2}},$$

which makes it a Hilbert space, see [28, Appendix] for a detailed presentation of these spaces.

We can formulate equation (2.6) as

$$\mathcal{T}(x) = x$$

where  $\mathcal{T} : \mathbb{H}^1(t_0, T; \mathcal{H}) \rightarrow \mathbb{H}^1(t_0, T; \mathcal{H})$  is the mapping which associates to  $x \in \mathbb{H}^1(t_0, T; \mathcal{H})$  the unique solution  $z = \mathcal{T}x$  of the evolution equation

$$\ddot{z}(t) + \gamma\dot{z}(t) + \partial\phi(\dot{z}(t)) + \nabla f(z(t)) \ni -\beta\nabla^2 f(x(t))\dot{x}(t), \quad (2.7)$$

which satisfies the Cauchy data  $z(t_0) = x_0$ , and  $\dot{z}(t_0) = \dot{x}_0$ . Let us first verify that this equation is well posed. For  $x \in \mathbb{H}^1(t_0, T; \mathcal{H})$  and according to the  $L$ -Lipschitz continuity of  $\nabla f$ , we have

$$\|\nabla^2 f(x(t))\dot{x}(t)\| \leq L\|\dot{x}(t)\|$$

which belongs to  $\mathbb{L}^2(t_0, T; \mathcal{H})$ . So, we can apply Theorem 2.1 which gives the existence and uniqueness of the solution of the solution  $z$  of (2.7) with the given Cauchy data.

Let us show that  $\mathcal{T} : \mathbb{H}^1(t_0, T; \mathcal{H}) \rightarrow \mathbb{H}^1(t_0, T; \mathcal{H})$  is a compact operator. Suppose that  $\|x\|_{\mathbb{H}^1(t_0, T; \mathcal{H})} \leq C$ , where  $C$  is a positive number. Equivalently  $\|x\|_{\mathbb{L}^2(t_0, T; \mathcal{H})} \leq C$  and  $\|\dot{x}\|_{\mathbb{L}^2(t_0, T; \mathcal{H})} \leq C$ . By definition of  $z = \mathcal{T}(x)$ , we have that (2.7) is verified. Take the scalar product of (2.7) with  $\dot{z}(t)$ . We obtain

$$\langle \ddot{z}(t), \dot{z}(t) \rangle + \gamma\|\dot{z}(t)\|^2 + \langle \partial\phi(\dot{z}(t)), \dot{z}(t) \rangle + \langle \nabla f(z(t)), \dot{z}(t) \rangle = \langle -\beta\nabla^2 f(x(t))\dot{x}(t), \dot{z}(t) \rangle. \quad (2.8)$$

Since  $\phi : \mathcal{H} \rightarrow \mathbb{R}$  is a convex continuous function that satisfies  $(DF)_r$ , we obtain

$$\frac{1}{2} \frac{d}{dt} \|\dot{z}(t)\|^2 + \gamma \|\dot{z}(t)\|^2 + r \|\dot{z}(t)\| + \frac{d}{dt} f(z(t)) \leq \langle -\beta \nabla^2 f(x(t)) \dot{x}(t), \dot{z}(t) \rangle.$$

According to the  $L$ -Lipschitz continuity of  $\nabla f$ , and the Cauchy-Schwarz inequality, we have

$$|\langle \nabla^2 f(x(t)) \dot{x}(t), \dot{z}(t) \rangle| \leq L \|\dot{x}(t)\| \|\dot{z}(t)\| \leq \frac{L}{2} (\|\dot{x}(t)\|^2 + \|\dot{z}(t)\|^2).$$

Collecting the above results, we obtain

$$\frac{d}{dt} \left( \frac{1}{2} \|\dot{z}(t)\|^2 + f(z(t)) \right) + \left( \gamma - \frac{\beta L}{2} \right) \|\dot{z}(t)\|^2 + r \|\dot{z}(t)\| \leq \frac{\beta L}{2} \|\dot{x}(t)\|^2. \quad (2.9)$$

After integration, we obtain

$$\left( \gamma - \frac{\beta L}{2} \right) \int_{t_0}^T \|\dot{z}(t)\|^2 dt \leq \frac{1}{2} \|\dot{x}_0\|^2 + f(x_0) - \inf_{\mathcal{H}} f + \frac{\beta L}{2} \int_{t_0}^T \|\dot{x}(t)\|^2 dt.$$

According to the hypothesis  $\gamma > \beta L$ , and since  $\|\dot{x}\|_{\mathbb{L}^2(t_0, T; \mathcal{H})} \leq C$ , we deduce that

$$\sup \left\{ \|\mathcal{T}(x)\|_{\mathbb{H}^1(t_0, T; \mathcal{H})} : \|\dot{x}\|_{\mathbb{L}^2(t_0, T; \mathcal{H})} \leq C \right\} < +\infty. \quad (2.10)$$

Returning to equation (2.7), and according to the growth condition on  $\partial\phi$  we deduce that

$$\sup \left\{ \|\mathcal{T}(x)\|_{\mathbb{H}^2(t_0, T; \mathcal{H})} : \|\dot{x}\|_{\mathbb{L}^2(t_0, T; \mathcal{H})} \leq C \right\} < +\infty. \quad (2.11)$$

Since  $\mathcal{H}$  is a finite dimensional space, we can use the compact embedding of  $\mathbb{H}^2(t_0, T; \mathcal{H})$  into  $\mathbb{H}^1(t_0, T; \mathcal{H})$ , and so conclude that  $\mathcal{T}$  is a compact mapping from  $\mathbb{H}^1(t_0, T; \mathcal{H})$  into itself. To conclude to the existence of a fixed point to  $\mathcal{T}$  we will use the Schaefer theorem (see [37]), which, besides the fact that  $\mathcal{T}$  is a compact mapping, requires the following property: the set

$$K := \left\{ x \in \mathbb{H}^1(t_0, T; \mathcal{H}) : \exists \lambda \in ]0, 1[ : x = \lambda \mathcal{T}(x) \right\} \text{ is bounded.}$$

So suppose that  $x = \lambda \mathcal{T}(x)$  for some  $\lambda \in ]0, 1[$ . This is equivalent to  $\mathcal{T}(x) = \frac{1}{\lambda} x$ , that is

$$\frac{1}{\lambda} \ddot{x}(t) + \frac{1}{\lambda} \gamma \dot{x}(t) + \partial\phi\left(\frac{1}{\lambda} \dot{x}(t)\right) + \nabla f\left(\frac{1}{\lambda} x(t)\right) + \beta \nabla^2 f(x(t)) \dot{x}(t) \ni 0. \quad (2.12)$$

After multiplication by  $\lambda$ , we get

$$\ddot{x}(t) + \gamma \dot{x}(t) + \lambda \partial\phi\left(\frac{1}{\lambda} \dot{x}(t)\right) + \lambda \nabla f\left(\frac{1}{\lambda} x(t)\right) + \beta \lambda \nabla^2 f(x(t)) \dot{x}(t) \ni 0. \quad (2.13)$$

Take the scalar product of (2.13) with  $\dot{x}(t)$ . A similar computation as above gives

$$\frac{d}{dt} \left( \frac{1}{2} \|\dot{x}(t)\|^2 + \lambda^2 f\left(\frac{1}{\lambda} x(t)\right) \right) + (\gamma - \beta L) \|\dot{x}(t)\|^2 + r \|\dot{x}(t)\| \leq 0. \quad (2.14)$$

After integration on  $[t_0, T]$  we get

$$(\gamma - \beta L) \int_{t_0}^T \|\dot{x}(t)\|^2 dt \leq \frac{1}{2} \|\dot{x}_0\|^2 + \lambda^2 f\left(\frac{1}{\lambda} x_0\right) - \inf_{\mathcal{H}} f.$$

Since  $\nabla f$  is  $L$ -Lipschitz continuous, the classical gradient descent lemma gives

$$f(x) \leq f(0) + \langle \nabla f(0), x \rangle + \frac{L}{2} \|x\|^2.$$

Hence,

$$\sup_{\lambda \in ]0, 1[} \lambda^2 f\left(\frac{1}{\lambda} x_0\right) \leq \sup_{\lambda \in ]0, 1[} \left( \lambda^2 f(0) + \lambda \langle \nabla f(0), x_0 \rangle + \frac{L}{2} \|x_0\|^2 \right) < +\infty.$$

Therefore the set  $K$  is bounded in  $\mathbb{H}^1(t_0, T; \mathcal{H})$ , which completes the proof of the existence of a fixed point to  $\mathcal{T}$ , thanks to the Schaefer theorem. ■

### 3 Another approach to existence using Mosco convergence

The major difficulty in (IGDH) is the presence of the term  $\partial\phi(\dot{x}(t))$ , which is attached to dry friction, and which involves the nonsmooth operator  $\partial\phi$ . A natural idea is to regularize this operator, and thus obtain a classical evolution equation. To this end, we assume that  $\mathcal{H}$  is a finite dimensional Hilbert space, and we will use the Moreau-Yosida regularization (this technique was used in [4] but without the Hessian term). Let us recall some basic facts concerning this regularization procedure. For any  $\lambda > 0$ , the Moreau envelope of  $\phi$  of index  $\lambda$  is the function  $\phi_\lambda : \mathcal{H} \rightarrow \mathbb{R}$  defined by: for all  $x \in \mathcal{H}$ ,

$$\phi_\lambda(x) = \min_{\xi \in \mathcal{H}} \left\{ \phi(\xi) + \frac{1}{2\lambda} \|x - \xi\|^2 \right\}.$$

The function  $\phi_\lambda$  is convex, of class  $\mathcal{C}^{1,1}$ , and satisfies  $\inf_{\mathcal{H}} \phi_\lambda = \inf_{\mathcal{H}} \phi$ ,  $\operatorname{argmin}_{\mathcal{H}} \phi_\lambda = \operatorname{argmin}_{\mathcal{H}} \phi$ . One can consult [9, section 17.2.1], [25], [28] for an in-depth study of the properties of the Moreau envelope in a Hilbert framework. In our context, since  $\phi : \mathcal{H} \rightarrow \mathbb{R}$  is a convex function that satisfies  $(DF)_r$ , we will have that  $\phi_\lambda$  is convex with  $\phi_\lambda(0) = \inf_{\mathcal{H}} \phi_\lambda = 0$ . This implies that, for all  $x \in \mathcal{H}$

$$\langle \nabla \phi_\lambda(x), x \rangle \geq 0. \quad (3.1)$$

This inequality will be very useful later. So, for each  $\lambda > 0$ , we consider the regularized evolution equation

$$(IGDH)_\lambda \quad \ddot{x}(t) + \gamma \dot{x}(t) + \nabla \phi_\lambda(\dot{x}(t)) + \beta \nabla^2 f(x(t)) \dot{x}(t) + \nabla f(x(t)) = 0, \quad t \in [t_0, +\infty[ \quad (3.2)$$

As a key property, this equation can be written in an equivalent way as a first-order system in time and space. The following writing is different from the Hamiltonian formulation used in the previous section. More precisely, in our context, we have the following equivalence, which follows from elementary differential calculus.

**Proposition 3.1** *The following are equivalent: (i)  $\iff$  (ii)*

$$(i) \quad \ddot{x}(t) + \gamma \dot{x}(t) + \nabla \phi_\lambda(\dot{x}(t)) + \beta \nabla^2 f(x(t)) \dot{x}(t) + \nabla f(x(t)) = 0.$$

$$(ii) \quad \begin{cases} \dot{x}(t) = u(t) - \beta \nabla f(x(t)); \\ \dot{u}(t) = -(\gamma I + \nabla \phi_\lambda)(u(t) - \beta \nabla f(x(t))) - \nabla f(x(t)). \end{cases}$$

**Proof.** (i)  $\implies$  (ii). Set

$$u(t) := \dot{x}(t) + \beta \nabla f(x(t)). \quad (3.3)$$

According to the classical derivation chain rule, we have

$$\dot{u}(t) = \ddot{x}(t) + \beta \nabla^2 f(x(t)) \dot{x}(t).$$

By using (i), we obtain

$$\dot{u}(t) + \gamma \dot{x}(t) + \nabla \phi_\lambda(\dot{x}(t)) + \nabla f(x(t)) = 0. \quad (3.4)$$

Putting together (3.3) and (3.4) we obtain the first-order system in time and space, as given in (ii).

(ii)  $\implies$  (i). Suppose that  $\nabla f$  is differentiable (that is  $f$  is twice differentiable). By differentiating the first equation of (ii) we obtain

$$\ddot{x}(t) = \dot{u}(t) - \beta \nabla^2 f(x(t)) \dot{x}(t).$$

According to the second equation of (ii), we deduce that

$$\begin{aligned} \ddot{x}(t) &= -(\gamma I + \nabla \phi_\lambda)(u(t) - \beta \nabla f(x(t))) - \nabla f(x(t)) - \beta \nabla^2 f(x(t)) \dot{x}(t) \\ &= -(\gamma I + \nabla \phi_\lambda)(\dot{x}(t)) - \nabla f(x(t)) - \beta \nabla^2 f(x(t)) \dot{x}(t) \end{aligned}$$

which gives (i). ■

This approach leads to a slightly more general existence statement than in Theorem 2.2.

**Theorem 3.1** *Let  $f : \mathcal{H} \rightarrow \mathbb{R}$  be a  $C^1$  function whose gradient is  $L$ -Lipschitz continuous, and let  $\phi : \mathcal{H} \rightarrow \mathbb{R}$  be a convex continuous function that satisfies  $(DF)_r$  and the following growth condition*

$$\|(\partial\phi)^o(x)\| \leq c(1 + \|x\|),$$

where  $(\partial\phi)^o(x)$  denotes the element of minimal norm of the set  $\partial\phi(x)$  and  $c > 0$ . Suppose that

$$\gamma > \beta L.$$

Then, for any Cauchy data  $(x_0, \dot{x}_0) \in \mathcal{H} \times \mathcal{H}$ , there exists a strong global solution of

$$(IGDH) \quad \ddot{x}(t) + \gamma\dot{x}(t) + \partial\phi(\dot{x}(t)) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) \ni 0. \quad (3.5)$$

satisfying  $x(t_0) = x_0$  and  $\dot{x}(t_0) = \dot{x}_0$ .

**Proof.** For each  $\lambda > 0$ ,  $(IGDH)_\lambda$  is relevant of the classical Cauchy-Lipschitz theorem. Indeed, according to Proposition 3.1, we can formulate it equivalently as follows. Set

$$Z(t) = (x(t), u(t)) \in \mathcal{H} \times \mathcal{H}.$$

The system (ii) of Proposition 3.1 can be written equivalently as

$$\dot{Z}(t) + F_\lambda(Z(t)) = 0, \quad Z(t_0) = (x_0, \dot{x}_0)$$

where

$$F_\lambda(x, u) = \left( \beta\nabla f(x) - u, (\gamma I + \nabla\phi_\lambda)(u - \beta\nabla f(x)) + \nabla f(x) \right).$$

The Lipschitz continuity property is stable by sum and composition. Since  $\nabla f$  and  $\nabla\phi_\lambda$  are Lipschitz continuous, the components of  $F_\lambda$  are Lipschitz continuous with respect to  $(x, u)$ , as well as  $F_\lambda$ . Therefore, according to the Cauchy-Lipschitz theorem, there exists a unique global classical solution  $(x_\lambda, u_\lambda)$  of the above system, which gives that  $x_\lambda$  is solution of

$$(IGDH)_\lambda \quad \ddot{x}_\lambda(t) + \gamma\dot{x}_\lambda(t) + \nabla\phi_\lambda(\dot{x}_\lambda(t)) + \beta\nabla^2 f(x_\lambda(t))\dot{x}_\lambda(t) + \nabla f(x_\lambda(t)) = 0, \quad t \in [t_0, +\infty[ \quad (3.6)$$

Let us first establish energy estimates on the sequence  $(x_\lambda)$ . By taking the scalar product of  $(IGDH)_\lambda$  with  $\dot{x}_\lambda(t)$ , we obtain

$$\langle \ddot{x}_\lambda(t), \dot{x}_\lambda(t) \rangle + \gamma \|\dot{x}_\lambda(t)\|^2 + \langle \nabla\phi_\lambda(\dot{x}_\lambda(t)), \dot{x}_\lambda(t) \rangle + \beta \langle \nabla^2 f(x_\lambda(t))\dot{x}_\lambda(t), \dot{x}_\lambda(t) \rangle + \langle \nabla f(x_\lambda(t)), \dot{x}_\lambda(t) \rangle = 0.$$

According to the classical derivation chain rule and inequality (3.1), we obtain

$$\frac{1}{2} \frac{d}{dt} \|\dot{x}_\lambda(t)\|^2 + \gamma \|\dot{x}_\lambda(t)\|^2 + \beta \langle \nabla^2 f(x_\lambda(t))\dot{x}_\lambda(t), \dot{x}_\lambda(t) \rangle + \frac{d}{dt} (f(x_\lambda(t))) = 0.$$

According to the  $L$ -Lipschitz continuity of  $\nabla f$ , and the Cauchy-Schwarz inequality, we have

$$|\langle \nabla^2 f(x_\lambda(t))\dot{x}_\lambda(t), \dot{x}_\lambda(t) \rangle| \leq L \|\dot{x}_\lambda(t)\|^2.$$

Therefore

$$\frac{d}{dt} \left( \frac{1}{2} \|\dot{x}_\lambda(t)\|^2 + f(x_\lambda(t)) \right) + (\gamma - \beta L) \|\dot{x}_\lambda(t)\|^2 \leq 0. \quad (3.7)$$

According to the hypothesis  $\gamma > \beta L$ , we deduce that the global energy

$$E_\lambda(t) = \frac{1}{2} \|\dot{x}_\lambda(t)\|^2 + f(x_\lambda(t))$$

is non increasing. From this, and by using that  $f$  is minorized we deduce that

$$\sup_{t \geq t_0, \lambda > 0} \|\dot{x}_\lambda(t)\| < +\infty. \quad (3.8)$$

According to the mean value theorem, and  $x_\lambda(t_0) = x_0$  fixed, we deduce that, for each  $T \geq t_0$

$$\sup_{t \in [t_0, T], \lambda > 0} \|x_\lambda(t)\| < +\infty. \quad (3.9)$$

Since  $\nabla f$  is Lipschitz continuous, this immediately implies that, for each  $T \geq t_0$

$$\sup_{t \in [t_0, T], \lambda > 0} \|\nabla f(x_\lambda(t))\| < +\infty. \quad (3.10)$$

In addition, by integrating (3.7), we deduce that

$$\sup_{\lambda > 0} \int_{t_0}^t \|\dot{x}_\lambda(t)\|^2 dt < +\infty. \quad (3.11)$$

As a classical property of the Yosida approximation of a maximally monotone operator  $A$ , we have

$$\|A_\lambda x\| \leq \|A^o(x)\|$$

where  $A^o(x)$  is the element of minimal norm of  $A(x)$ , see Brézis [28, Proposition 2.6]. According to the assumption  $\|(\partial\phi)^o(x)\| \leq c(1 + \|x\|)$ , we obtain

$$\|\nabla\phi_\lambda(x)\| \leq c(1 + \|x\|).$$

Therefore, and according to (3.8), for each  $T \geq t_0$

$$\sup_{t \in [t_0, T], \lambda > 0} \|\nabla\phi_\lambda(\dot{x}_\lambda(t))\| < +\infty. \quad (3.12)$$

Finally, for each  $T \geq t_0$

$$\sup_{t \in [t_0, T], \lambda > 0} \|\nabla^2 f(x_\lambda(t))\dot{x}_\lambda(t)\| \leq L, \quad \sup_{t \in [t_0, T], \lambda > 0} \|\dot{x}_\lambda(t)\| < +\infty. \quad (3.13)$$

Let us now return to  $(\text{IGDH})_\lambda$  that we write as follows:

$$\ddot{x}_\lambda(t) = -\gamma\dot{x}_\lambda(t) - \nabla\phi_\lambda(\dot{x}_\lambda(t)) - \beta\nabla^2 f(x_\lambda(t))\dot{x}_\lambda(t) - \nabla f(x_\lambda(t)). \quad (3.14)$$

Using the above estimates, we deduce that, for each  $T \geq t_0$

$$\sup_{\lambda > 0} \int_{t_0}^T \|\ddot{x}_\lambda(t)\|^2 dt < +\infty. \quad (3.15)$$

Therefore, for each  $T \geq t_0$ , the sequence  $(x_\lambda)$  is bounded in  $\mathbb{H}^2(t_0, T; \mathcal{H})$ . Since  $\mathcal{H}$  has been supposed to be a finite dimensional Hilbert space, this implies that the sequence  $(x_\lambda)$  is relatively compact in  $\mathbb{H}^1(t_0, T; \mathcal{H})$  for all  $T \geq t_0$ . After extraction of a subsequence, still denoted  $(x_\lambda)$ , we will have the existence of  $x$  such that

1.  $x \in \mathbb{H}^2(t_0, T; \mathcal{H})$  for all  $T \geq t_0$ .
2.  $x_\lambda \rightarrow x$  uniformly on  $[t_0, T]$  for all  $T \geq t_0$ .
3.  $\nabla f(x_\lambda) \rightarrow \nabla f(x)$  uniformly on  $[t_0, T]$  for all  $T \geq t_0$ .
4.  $\dot{x}_\lambda \rightarrow \dot{x}$  uniformly on  $[t_0, T]$  for all  $T \geq t_0$ .
5.  $\ddot{x}_\lambda \rightarrow \ddot{x}$  weakly in  $\mathbb{L}^2(t_0, T; \mathcal{H})$  for all  $T \geq t_0$ .
6.  $\nabla^2 f(x_\lambda)\dot{x}_\lambda \rightarrow \nabla^2 f(x)\dot{x}$  weakly in  $\mathbb{L}^2(t_0, T; \mathcal{H})$  for all  $T \geq t_0$ .

This last property comes from

$$\nabla^2 f(x_\lambda)\dot{x}_\lambda = \frac{d}{dt}\nabla f(x_\lambda)$$

and the continuity of the derivation in the distribution sense.

To pass on the limit on  $(\text{IGDH})_\lambda$ , we write it as follows

$$\nabla\phi_\lambda(\dot{x}_\lambda(t)) = \xi_\lambda(t) \tag{3.16}$$

where

$$\xi_\lambda(t) = -\ddot{x}_\lambda(t) - \gamma\dot{x}_\lambda(t) - \beta\nabla^2 f(x_\lambda(t))\dot{x}_\lambda(t) - \nabla f(x_\lambda(t)).$$

We now rely on the variational convergence properties of the Yosida approximation. Since  $\phi_\lambda$  converges increasingly to  $\phi$ , the sequence of integral functionals

$$\Phi^\lambda(\xi) := \int_{t_0}^T \phi_\lambda(\xi(t))dt$$

converges increasingly to

$$\Phi(\xi) = \int_{t_0}^T \phi(\xi(t))dt.$$

Therefore it converges in the Mosco sense in  $\mathbb{L}^2(t_0, T; \mathcal{H})$ , see Mosco [33]. According to the theorem which makes the link between the Mosco convergence of a sequence of convex lower semicontinuous functions and the graph convergence of their subdifferentials, see Attouch [8, Theorem 3.66], we have that

$$\partial\Phi^\lambda \rightarrow \partial\Phi$$

with respect to the topology strong  $-\mathbb{L}^2(t_0, T; \mathcal{H}) \times$  weak  $-\mathbb{L}^2(t_0, T; \mathcal{H})$ . We have

$$\xi_\lambda = \nabla\Phi^\lambda(\dot{x}_\lambda).$$

Since

$$\dot{x}_\lambda \rightarrow \dot{x} \text{ strongly in } \mathbb{L}^2(t_0, T; \mathcal{H})$$

and  $\xi_\lambda$  converges weakly in  $\mathbb{L}^2(t_0, T; \mathcal{H})$  to  $\xi$  given by

$$\xi(t) = -\ddot{x}(t) - \gamma\dot{x}(t) - \beta\nabla^2 f(x(t))\dot{x}(t) - \nabla f(x(t)), \tag{3.17}$$

we deduce that  $\xi \in \partial\Phi(\dot{x})$ , that is

$$\xi(t) \in \partial\phi(\dot{x}(t)).$$

According to the formulation (3.17) of  $\xi$ , we finally obtain that  $x$  is a solution of (IGDH). ■

## 4 Finite time convergence of the trajectories

Let's analyze the asymptotic behavior as  $t \rightarrow +\infty$ , and the finite convergence property, of the solution trajectories of the second-order differential inclusion

$$(IGDH) \quad \ddot{x}(t) + \gamma \dot{x}(t) + \partial\phi(\dot{x}(t)) + \beta \nabla^2 f(x(t)) \dot{x}(t) + \nabla f(x(t)) \ni 0, \quad t \in [t_0, +\infty[.$$

**Theorem 4.1** *Let  $f : \mathcal{H} \rightarrow \mathbb{R}$  be a  $C^1$  function whose gradient is  $L$ -Lipschitz continuous, and let  $\phi : \mathcal{H} \rightarrow \mathbb{R}$  be a convex continuous function that satisfies  $(DF)_r$  and which is bounded on the bounded sets. Suppose that*

$$\gamma > \beta L.$$

*Then, for any solution trajectory  $x(\cdot)$  of (IGDH) we have:*

- (i)  $\|\dot{x}\| \in \mathbb{L}^1(t_0, +\infty; \mathbb{R})$ , and therefore the strong limit  $x_\infty := \lim_{t \rightarrow +\infty} x(t)$  exists.
- (ii) The limit point  $x_\infty$  is an equilibrium point of (IGDH), i.e.  $-\nabla f(x_\infty) \in \partial\phi(0)$ .
- (iii) If  $-\nabla f(x_\infty) \notin \text{boundary}(\partial\phi(0))$ , then there exists  $t_1 \geq 0$  such that  $x(t) = x_\infty$  for every  $t \geq t_1$ .

**Proof.** (i) Take the scalar product of (IGDH) with  $\dot{x}(t)$ . We obtain

$$\langle \ddot{x}(t), \dot{x}(t) \rangle + \gamma \|\dot{x}(t)\|^2 + \langle \partial\phi(\dot{x}(t)), \dot{x}(t) \rangle + \beta \langle \nabla^2 f(x(t)) \dot{x}(t), \dot{x}(t) \rangle + \langle \nabla f(x(t)), \dot{x}(t) \rangle = 0, \quad (4.1)$$

which gives

$$\frac{1}{2} \frac{d}{dt} \|\dot{x}(t)\|^2 + \gamma \|\dot{x}(t)\|^2 + \langle \partial\phi(\dot{x}(t)), \dot{x}(t) \rangle + \beta \langle \nabla^2 f(x(t)) \dot{x}(t), \dot{x}(t) \rangle + \frac{d}{dt} (f(x(t)) - \inf_{\mathcal{H}} f) = 0.$$

According to the  $L$ -Lipschitz continuity of  $\nabla f$ , and the Cauchy-Schwarz inequality, we have

$$|\langle \nabla^2 f(x(t)) \dot{x}(t), \dot{x}(t) \rangle| \leq L \|\dot{x}(t)\|^2.$$

According to the assumption  $(DF)_r$  on  $\phi$  and Lemma 1.1,

$$\langle \partial\phi(\dot{x}(t)), \dot{x}(t) \rangle \geq \phi(\dot{x}(t)) \geq r \|\dot{x}(t)\|.$$

Collecting the above results, we obtain

$$\frac{d}{dt} \left( \frac{1}{2} \|\dot{x}(t)\|^2 + f(x(t)) - \inf_{\mathcal{H}} f \right) + (\gamma - \beta L) \|\dot{x}(t)\|^2 + r \|\dot{x}(t)\| \leq 0. \quad (4.2)$$

According to the hypothesis  $\gamma > \beta L$ , we deduce that the global energy

$$E(t) = \frac{1}{2} \|\dot{x}(t)\|^2 + f(x(t)) - \inf_{\mathcal{H}} f$$

is non increasing. Moreover, by integrating (4.2) we obtain

$$\int_{t_0}^{\infty} \|\dot{x}(t)\|^2 dt < +\infty \quad \text{and} \quad \int_{t_0}^{\infty} \|\dot{x}(t)\| dt < +\infty. \quad (4.3)$$

This last property expresses that the trajectory has a finite length, and hence  $\lim_{t \rightarrow +\infty} x(t) := x_\infty$  exists.

(ii) Since  $E(\cdot)$  is non increasing, we have that the velocity  $\dot{x}(t)$  remains bounded. Since  $\phi$  is bounded on the bounded sets, so is  $\partial\phi$ . Therefore, from equation (IGDH) we deduce that the acceleration  $\ddot{x}(t)$  remains bounded. This combined with  $\int_{t_0}^{\infty} \|\dot{x}(t)\| dt < +\infty$  implies that the velocity  $\dot{x}(t)$  converges



strongly to zero, as  $t \rightarrow +\infty$ . Let us now pass to the limit on (IGDH). Set  $u(t) = \dot{x}(t)$ . Let us write (IGDH) equivalently as

$$\dot{u}(t) + (\gamma I + \partial\phi)(u(t)) = h(t)$$

with  $h(t) := -\beta \nabla^2 f(x(t))\dot{x}(t) - \nabla f(x(t))$ . The operator  $A = \gamma I + \partial\phi$  is strongly monotone because  $\gamma > 0$ . According to the above results, we have that  $h(t)$  converges strongly to  $-\nabla f(x_\infty)$ . We now apply the Theorem 3.9 of Brezis [28], which tells us that the strong limit of  $u(t)$ , that's zero, satisfies  $A(0) \ni -\nabla f(x_\infty)$ . Equivalently

$$\partial\phi(0) \ni -\nabla f(x_\infty).$$

(iii) The assumption  $-\nabla f(x_\infty) \in \text{int}(\partial\phi(0))$  implies the existence of  $\varepsilon > 0$  such that

$$-\nabla f(x_\infty) + B(0, 2\varepsilon) \subset \partial\phi(0).$$

On the other hand, since  $\lim_{t \rightarrow +\infty} \nabla f(x(t)) = \nabla f(x_\infty)$ , there exists  $t_1 \geq t_0$  such that for all  $t \geq t_1$

$$\nabla f(x(t)) \in \nabla f(x_\infty) + B(0, \varepsilon).$$

Hence,

$$-\nabla f(x(t)) + B(0, \varepsilon) \subset -\nabla f(x_\infty) + B(0, 2\varepsilon) \subset \partial\phi(0).$$

Equivalently, for every  $t \geq t_1$  and for every  $u \in B(0, 1)$ , we have:

$$-\nabla f(x(t)) + \varepsilon u \in \partial\phi(0).$$

Let's write the corresponding subdifferential inequality at the origin (recall that  $\phi(0) = 0$ ). For every  $t \geq t_1$

$$\forall u \in B(0, 1), \quad \phi(\dot{x}(t)) \geq \langle -\nabla f(x(t)) + \varepsilon u, \dot{x}(t) \rangle.$$

Taking the supremum over  $u \in B(0, 1)$ , we obtain that, for every  $t \geq t_1$ ,

$$\phi(\dot{x}(t)) + \langle \nabla f(x(t)), \dot{x}(t) \rangle \geq \varepsilon \|\dot{x}(t)\|. \quad (4.4)$$

Let's return to (4.1). According to the above results, we obtain

$$\frac{d}{dt} \frac{1}{2} \|\dot{x}(t)\|^2 + (\gamma - \beta L) \|\dot{x}(t)\|^2 + \varepsilon \|\dot{x}(t)\| \leq 0. \quad (4.5)$$

a) Neglecting the nonnegative term  $\varepsilon \|\dot{x}(t)\|$  we obtain

$$\frac{d}{dt} \frac{1}{2} \|\dot{x}(t)\|^2 + (\gamma - \beta L) \|\dot{x}(t)\|^2 \leq 0, \quad (4.6)$$

whose integration gives

$$\|\dot{x}(t)\| \leq \|\dot{x}(t_0)\| e^{-(\gamma - \beta L)t}.$$

b) Neglecting the nonnegative term  $(\gamma - \beta L) \|\dot{x}(t)\|^2$  we obtain

$$\frac{d}{dt} \|\dot{x}(t)\|^2 + 2\varepsilon \|\dot{x}(t)\| \leq 0, \quad (4.7)$$

Set  $v(t) = \|\dot{x}(t)\|^2$ . We have  $\dot{v}(t) + 2\varepsilon \sqrt{v(t)} \leq 0$ . As long as  $v(t) > 0$  we will have  $\frac{d}{dt} \sqrt{v(t)} \leq -\varepsilon$ . This forces  $v(t)$  to be equal to zero after some finite time. ■

**Remark 4.1** With the condition  $-\nabla f(x_\infty) \notin \text{boundary}(\partial\phi(0))$ , the finite time convergence of the trajectory to a stationary point of the dynamic (IGDH) is ensured, i.e. there exists  $t_1 \geq 0$  such that  $x(t) = x_\infty$  for every  $t \geq t_1$ . In addition, an estimate of the final time can be given. Thanks to a detailed analysis of the end of the previous demonstration, we can show that

$$t_1 \leq \tau_0 + \frac{2\|\dot{x}(\tau_0)\|}{\text{dist}\left(-\nabla f(x_\infty), \text{boundary}(\partial\phi(0))\right)},$$

where  $\tau_0$  is the first time instant such that

$$\nabla f(x(t)) \in \nabla f(x_\infty) + B(0, \varepsilon), \text{ for all } t \geq t_0, \text{ with } \varepsilon = \frac{1}{2} \text{dist}\left(-\nabla f(x_\infty), \text{boundary}(\partial\phi(0))\right).$$

**Soft thresholding of the velocities** As a model situation for dry friction, take  $\phi : \mathcal{H} \rightarrow \mathbb{R}$  given by  $\phi(x) = r\|x\|$ , with  $r > 0$ . We have

$$\partial\phi(x) = \begin{cases} r \frac{x}{\|x\|} & \text{if } x \neq 0; \\ B(0, r) & \text{if } x = 0. \end{cases} \quad (4.8)$$

A direct application of Theorem 4.1 gives the following result:

**Corollary 4.1** Let  $f : \mathcal{H} \rightarrow \mathbb{R}$  be a  $C^1$  function whose gradient is  $L$ -Lipschitz continuous. Assume that the potential friction function  $\phi$  is given by  $\phi(x) = r\|x\|$ . Suppose that

$$\gamma > \beta L.$$

Then, for any solution trajectory  $x(\cdot)$  of (IGDH) we have:

(i)  $\|\dot{x}\| \in \mathbb{L}^1(t_0, +\infty; \mathbb{R})$ , and therefore the strong limit  $x_\infty := \lim_{t \rightarrow +\infty} x(t)$  exists.

(ii) The limit point  $x_\infty$  satisfies

$$\|\nabla f(x_\infty)\| \leq r.$$

(iii) If  $\|\nabla f(x_\infty)\| < r$ , then there exists  $t_1 \geq 0$  such that  $x(t) = x_\infty$  for every  $t \geq t_1$ .

Clearly, taking  $r$  small is the interesting situation for optimization.

## 5 Errors, perturbations

Let's analyze the asymptotic behavior as  $t \rightarrow +\infty$ , and the finite convergence property, of the solution trajectories of the second-order differential inclusion

$$\text{(IGDH)-pert} \quad \ddot{x}(t) + \gamma\dot{x}(t) + \partial\phi(\dot{x}(t)) + \beta\nabla^2 f(x(t))\dot{x}(t) + \nabla f(x(t)) \ni e(t)$$

which comes from the introduction of a second member  $e : [t_0, +\infty[ \rightarrow \mathcal{H}$  in (IGDH). Depending on the context,  $e$  can be interpreted as an external action, a control term, or coming from perturbations, errors.

**Theorem 5.1** Let  $f : \mathcal{H} \rightarrow \mathbb{R}$  be a  $C^1$  function whose gradient is  $L$ -Lipschitz continuous, and let  $\phi : \mathcal{H} \rightarrow \mathbb{R}$  be a convex continuous function that satisfies  $(\text{DF})_r$  and which is bounded on the bounded sets. Suppose that

$$\gamma > \beta L.$$

Suppose that the external action  $e : [t_0, +\infty[ \rightarrow \mathcal{H}$  satisfies

$$\lim_{t \rightarrow +\infty} \|e(t)\| = 0.$$

Then, for any solution trajectory  $x(\cdot)$  of (IGDH)-pert we have:

- (i)  $\|\dot{x}\| \in \mathbb{L}^1(t_0, +\infty; \mathbb{R})$ , and therefore the strong limit  $x_\infty := \lim_{t \rightarrow +\infty} x(t)$  exists.
- (ii) The limit point  $x_\infty$  satisfies:  $-\nabla f(x_\infty) \in \partial\phi(0)$ .
- (iii) If  $-\nabla f(x_\infty) \notin \text{boundary}(\partial\phi(0))$ , then there exists  $t_1 \geq 0$  such that  $x(t) = x_\infty$  for every  $t \geq t_1$ .

**Proof.** (i) Take the scalar product of (IGDH)-pert with  $\dot{x}(t)$ . We obtain

$$\langle \ddot{x}(t), \dot{x}(t) \rangle + \gamma \|\dot{x}(t)\|^2 + \langle \partial\phi(\dot{x}(t)), \dot{x}(t) \rangle + \beta \langle \nabla^2 f(x(t))\dot{x}(t), \dot{x}(t) \rangle + \langle \nabla f(x(t)), \dot{x}(t) \rangle = \langle e(t), \dot{x}(t) \rangle \quad (5.1)$$

which gives

$$\frac{1}{2} \frac{d}{dt} \|\dot{x}(t)\|^2 + \gamma \|\dot{x}(t)\|^2 + \langle \partial\phi(\dot{x}(t)), \dot{x}(t) \rangle + \beta \langle \nabla^2 f(x(t))\dot{x}(t), \dot{x}(t) \rangle + \frac{d}{dt} f(x(t)) = \langle e(t), \dot{x}(t) \rangle.$$

According to the L-Lipschitz continuity of  $\nabla f$ , and the Cauchy-Schwarz inequality, we have

$$|\langle \nabla^2 f(x(t))\dot{x}(t), \dot{x}(t) \rangle| \leq L \|\dot{x}(t)\|^2.$$

According to the assumption  $(DF)_r$  on  $\phi$  and Lemma 1.1,

$$\langle \partial\phi(\dot{x}(t)), \dot{x}(t) \rangle \geq \phi(\dot{x}(t)) \geq r \|\dot{x}(t)\|.$$

By Cauchy-Schwarz inequality

$$|\langle e(t), \dot{x}(t) \rangle| \leq \|e(t)\| \|\dot{x}(t)\|.$$

Collecting the above results, we obtain

$$\frac{d}{dt} \left( \frac{1}{2} \|\dot{x}(t)\|^2 + f(x(t)) - \inf_{\mathcal{H}} f \right) + (\gamma - \beta L) \|\dot{x}(t)\|^2 + (r - \|e(t)\|) \|\dot{x}(t)\| \leq 0. \quad (5.2)$$

According to the hypothesis  $\gamma > \beta L$ , and  $\lim_{t \rightarrow +\infty} \|e(t)\| = 0$ , we deduce that there exists  $t_1 \geq t_0$  such that, for  $t \geq t_1$ , the global energy

$$E(t) = \frac{1}{2} \|\dot{x}(t)\|^2 + f(x(t)) - \inf_{\mathcal{H}} f$$

is non increasing. Moreover, by integrating (5.2) we obtain

$$\int_{t_0}^{\infty} \|\dot{x}(t)\|^2 dt < +\infty \quad \text{and} \quad \int_{t_0}^{\infty} \|\dot{x}(t)\| dt < +\infty. \quad (5.3)$$

This last property expresses that the trajectory has a finite length, and hence  $\lim_{t \rightarrow +\infty} x(t) := x_\infty$  exists.

(ii) Since  $E(\cdot)$  is non increasing, we have that the velocity  $\dot{x}(t)$  remains bounded. Since  $\phi$  is bounded on the bounded sets, so is  $\partial\phi$ . Therefore, from equation (IGDH)-pert we deduce that the acceleration  $\ddot{x}(t)$  remains bounded. This combined with  $\int_{t_0}^{\infty} \|\dot{x}(t)\| dt < +\infty$  implies that the velocity  $\dot{x}(t)$  converges strongly to zero, as  $t \rightarrow +\infty$ . Let us now pass to the limit on (IGDH)-pert. Set  $u(t) = \dot{x}(t)$ . Let us write (IGDH)-pert equivalently as

$$\dot{u}(t) + (\gamma I + \partial\phi)(u(t)) = h(t)$$

with  $h(t) := -\beta \nabla^2 f(x(t)) \dot{x}(t) - \nabla f(x(t)) + e(t)$ . The operator  $A = \gamma I + \partial\phi$  is strongly monotone because  $\gamma > 0$ . According to the above results, we have that  $h(t)$  converges strongly to  $-\nabla f(x_\infty)$ . We now apply the Theorem 3.9 of Brezis [28], which tells us that the strong limit of  $u(t)$ , that's zero, satisfies  $A(0) \ni -\nabla f(x_\infty)$ . Equivalently

$$\partial\phi(0) \ni -\nabla f(x_\infty).$$

(iii) The assumption  $-\nabla f(x_\infty) \in \text{int}(\partial\phi(0))$  implies the existence of  $\varepsilon > 0$  such that

$$-\nabla f(x_\infty) + B(0, 2\varepsilon) \subset \partial\phi(0).$$

On the other hand, since  $\lim_{t \rightarrow +\infty} \nabla f(x(t)) = \nabla f(x_\infty)$ , there exists  $t_1 \geq t_0$  such that for all  $t \geq t_1$

$$\nabla f(x(t)) \in \nabla f(x_\infty) + B(0, \varepsilon).$$

Hence,

$$-\nabla f(x(t)) + B(0, \varepsilon) \subset -\nabla f(x_\infty) + B(0, 2\varepsilon) \subset \partial\phi(0).$$

Equivalently, for every  $t \geq t_1$  and for every  $u \in B(0, 1)$ , we have:

$$-\nabla f(x(t)) + \varepsilon u \in \partial\phi(0).$$

Let's write the corresponding subdifferential inequality at the origin (recall that  $\phi(0) = 0$ ). For every  $t \geq t_1$

$$\forall u \in B(0, 1), \quad \phi(\dot{x}(t)) \geq \langle -\nabla f(x(t)) + \varepsilon u, \dot{x}(t) \rangle.$$

Taking the supremum over  $u \in B(0, 1)$ , we obtain that, for every  $t \geq t_1$ ,

$$\phi(\dot{x}(t)) + \langle \nabla f(x(t)), \dot{x}(t) \rangle \geq \varepsilon \|\dot{x}(t)\|. \quad (5.4)$$

Let's return to 5.1. According to the above results, we obtain

$$\frac{d}{dt} \frac{1}{2} \|\dot{x}(t)\|^2 + (\gamma - \beta L) \|\dot{x}(t)\|^2 + (\varepsilon - \|e(t)\|) \|\dot{x}(t)\| \leq 0. \quad (5.5)$$

a) Neglecting the nonnegative term  $(\varepsilon - \|e(t)\|) \|\dot{x}(t)\|$  (recall that  $\varepsilon > 0$  and  $\|e(t)\| \rightarrow 0$ ) we obtain

$$\frac{d}{dt} \frac{1}{2} \|\dot{x}(t)\|^2 + (\gamma - \beta L) \|\dot{x}(t)\|^2 \leq 0, \quad (5.6)$$

whose integration gives

$$\|\dot{x}(t)\| \leq \|\dot{x}(t_0)\| e^{-(\gamma - \beta L)t}.$$

b) Neglecting the nonnegative term  $(\gamma - \beta L) \|\dot{x}(t)\|^2$  we obtain that for  $t$  large enough

$$\frac{d}{dt} \|\dot{x}(t)\|^2 + \varepsilon \|\dot{x}(t)\| \leq 0, \quad (5.7)$$

Set  $v(t) = \|\dot{x}(t)\|^2$ . We have  $\dot{v}(t) + \varepsilon \sqrt{v(t)} \leq 0$ . As long as  $v(t) > 0$  we will have  $\frac{d}{dt} \sqrt{v(t)} \leq -\frac{\varepsilon}{2}$ . This forces  $v(t)$  to be equal to zero after some finite time. ■

## 6 Nonsmooth case

In this section, we assume that  $f : \mathcal{H} \rightarrow \mathbb{R} \cup \{+\infty\}$  is a convex lower semicontinuous and proper function such that  $\inf f > -\infty$ . The preceding sections deal with a differentiable function  $f$ , without convexity assumption on  $f$ . Now, when considering nonsmooth functions, we assume the convexity of  $f$ . This allows us to use the regularity properties of the Moreau envelope in the convex case. Indeed, to reduce to the previous situation, where  $f : \mathcal{H} \rightarrow \mathbb{R}$  is a  $\mathcal{C}^1$  function whose gradient is Lipschitz continuous, the idea is to replace  $f$  by its Moreau envelope  $f_\lambda : \mathcal{H} \rightarrow \mathbb{R}$  which is defined by: for all  $x \in \mathcal{H}$ ,

$$f_\lambda(x) = \min_{\xi \in \mathcal{H}} \left\{ f(\xi) + \frac{1}{2\lambda} \|x - \xi\|^2 \right\}.$$

The properties of Moreau envelopes useful for our study were recalled in the section 3. Since the infimal value and the set of minimizers are preserved by taking the Moreau envelope, the idea is to replace  $f$  by  $f_\lambda$  in the previous study, and take advantage of the fact that  $f_\lambda$  is continuously differentiable. The dynamic (IGDH) becomes formally

$$\text{(IGDH)-regularized} \quad \ddot{x}(t) + \gamma \dot{x}(t) + \partial\phi(\dot{x}(t)) + \beta \nabla^2 f_\lambda(x(t)) \dot{x}(t) + \nabla f_\lambda(x(t)) \ni 0, \quad t \in [t_0, +\infty[$$

In fact  $f_\lambda$  is only a  $\mathcal{C}^{1,1}$  function. The term  $\nabla^2 f_\lambda(x(t)) \dot{x}(t)$  has to be understood as the derivative in the distribution sense of the Lipschitz continuous function  $t \mapsto \nabla f_\lambda(x(t))$ . Based on the properties of the Moreau envelope, a direct adaptation of Theorem 2.2 gives the following existence result for the regularized dynamic (IGDH)-regularized .

**Theorem 6.1** *Let  $f : \mathcal{H} \rightarrow \mathbb{R} \cup \{+\infty\}$  be a convex lower semicontinuous and proper function such that  $\inf_{\mathcal{H}} f > -\infty$ . Let  $\phi : \mathcal{H} \rightarrow \mathbb{R}$  be a convex continuous function that satisfies (DF)<sub>r</sub> and  $\|(\partial\phi)^o(x)\| \leq c(1 + \|x\|)$ . Take  $\lambda > 0$  and consider the corresponding differential equation (IGDH)-regularized .*

*Suppose that*

$$\gamma\lambda > \beta.$$

*Then, for any Cauchy data  $(x_0, \dot{x}_0) \in \mathcal{H} \times \mathcal{H}$ , there exists a strong global solution of the differential equation (IGDH)-regularized satisfying  $x(t_0) = x_0$ , and  $\dot{x}(t_0) = \dot{x}_0$ .*

**Proof.** We have that  $\nabla f_\lambda$  is L-Lipschitz continuous with  $L = \frac{1}{\lambda}$ . So, the condition  $\gamma > \beta L$  of Theorems 2.2 and 3.1 becomes  $\gamma\lambda > \beta$ . Then, just apply the conclusion of these theorems. ■

In the following we denote by  $x_\lambda$  a solution of (IGDH)-regularized . Note that  $\lambda$  is fixed, it is not intended to go to zero.

**Theorem 6.2** *Under the assumptions of Theorem 6.1, let  $x_\lambda$  be a solution of (IGDH)-regularized . Then, the following properties are satisfied:*

- (i)  $\|\dot{x}_\lambda\| \in \mathbb{L}^1(t_0, +\infty; \mathbb{R})$ , and therefore the strong limit  $x_{\lambda, \infty} := \lim_{t \rightarrow +\infty} x_\lambda(t)$  exists.
- (ii) The limit point  $x_{\lambda, \infty}$  is an equilibrium point of (IGDH)-regularized , i.e.  $-\nabla f_\lambda(x_{\lambda, \infty}) \in \partial\phi(0)$ .
- (iii) If  $-\nabla f_\lambda(x_{\lambda, \infty}) \notin \text{boundary}(\partial\phi(0))$ , then there exists  $t_1 \geq 0$  such that  $x_\lambda(t) = x_{\lambda, \infty} \forall t \geq t_1$ .
- (iv) Set  $p(t) := \text{prox}_{\lambda f}(x_\lambda(t))$ . We have that  $p(\cdot)$  has also a finite length, and  $p(t)$  converges strongly to  $p_\infty := \text{prox}_{\lambda f}(x_{\lambda, \infty})$ . In addition

$$\partial f(p_\infty) + \partial\phi(0) \ni 0.$$

*In addition, when  $x_\lambda$  is finitely convergent, so is  $p$ .*

**Proof.** The proof consists in replacing  $f$  by  $f_\lambda$  in Theorem 4.1, and in using that  $\nabla f_\lambda$  is  $\frac{1}{\lambda}$ -Lipschitz continuous. Since the proximal mapping is nonexpansive, we immediately deduce that for any  $s, t \geq t_0$

$$\|p(t) - p(s)\| \leq \|x_\lambda(t) - x_\lambda(s)\|. \quad (6.1)$$

This implies that  $p$  is Lipschitz continuous, hence almost everywhere differentiable and, for a.e.  $t \geq t_0$

$$\|\dot{p}(t)\| \leq \|\dot{x}_\lambda(t)\|.$$

Since  $\|\dot{x}_\lambda\| \in \mathbb{L}^1(t_0, +\infty; \mathbb{R})$ , we also have  $\|\dot{p}\| \in \mathbb{L}^1(t_0, +\infty; \mathbb{R})$ . Therefore, the trajectory  $p(\cdot)$  has a finite length. It converges strongly towards  $p_\infty = \text{prox}_{\lambda f} x_{\lambda, \infty}$ . Using the relation

$$\nabla f_\lambda(x_{\lambda, \infty}) \in \partial f(\text{prox}_{\lambda f} x_{\lambda, \infty}) = \partial f(p_\infty),$$

we obtain the approximate optimality property:

$$\partial f(p_\infty) + \partial \phi(0) \ni 0.$$

According to (6.1), when  $x_\lambda$  is finitely convergent, so is  $p$ . ■

## 7 Perspectives

As a major property, the inertial system (IGDH) provides strong global solutions which converge in finite time. Proof of the existence of such solutions requires restrictive assumptions about the data. It would be interesting to obtain such existence results in infinite dimensional spaces, thus opening the door to PDE's and nonlinear wave equations. Also, the study of the uniqueness of the solution trajectory for the Cauchy problem must be studied. This would be important to study damped shocks and their finite time stabilization (a natural property in mechanics). From the point of view of optimization, the main challenge is to be obtain a precise estimate of the finite time stabilization. This is based on the study of the condition  $-\nabla f(x_\infty) \notin \text{boundary}(\partial \phi(0))$ , which is a non-trivial topic. Thus, this paper opens the door to several interesting problems arising in various domains.

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