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Statistical study of non-adiabatic energization and transport in Kelvin-Helmholtz vortices at Mercury

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Abstract. Non-adiabatic energization and transport of planetary ions in the vicinity of Kelvin-Helmholtz (KH) vortices at Mercury are statistically investigated using a magnetohydrodynamic model combined with a test particle tracing technique. The measurements from the MErcury Surface, Space ENvironment, GEOchemistry, and Ranging spacecraft show the signature of KH waves in the magnetospheric flanks, especially on the duskside. We examine here the energization and transport of heavy ions of planetary origin such as Na^+ associated with the growth of the KH instability on both the dawn and dusk sides of the magnetosphere, and for both northward and southward interplanetary magnetic field configurations. We find that in all cases picked up planetary ions in the magnetosphere can be accelerated non-adiabatically with a level of acceleration larger on the dawnside. The transport of ions is controlled by the convection electric field in the magnetosheath. The planetary ions

26 picked up in the duskside magnetosphere stagnate inside the magnetosphere where the KH
27 waves have been mainly observed at Mercury and, hence, are not expected to escape. The
28 present study indicates that these energized ions may likely stay inside the magnetosphere and
29 sputter the planetary surface on the duskside preferentially, being the source of secondary
30 neutrals from the planetary surface.

1. Introduction

The Mercury Surface, Space ENvironment, GEochemistry, and Ranging (MESSENGER) spacecraft, the first orbiter around Mercury, was in orbit from March 11, 2011 until April 30, 2015. The many data collected provided in particular key information on the structure of the magnetosphere as well as on the dynamics of the plasma environment around Mercury. A number of plasma processes similar to the ones identified at Earth have been reported: magnetic reconnection and Flux Transfer Events (FTE) at the magnetopause [Slavin et al., 2009, 2010b, 2014; DiBraccio et al., 2013; Imber et al., 2014], substorm-like dipolarization events in the magnetotail [e.g., Sundberg et al., 2012a, Imber et al., 2017; Dewey et al., 2017, 2018], and Kelvin-Helmholtz (KH) vortices in the magnetospheric flanks [Boardsen et al., 2010; Sundberg et al., 2011, 2012b; Liljeblad et al., 2014, 2016; Gershman et al., 2015]. Moreover, the MESSENGER data have shown evidence of heavy ions of planetary origin in several regions of Mercury's magnetosphere, those resulting from the ionization of exospheric neutrals by solar UV [Zurbuchen et al., 2011; Schriver et al., 2011a, b; Raines et al., 2011, 2013; DiBraccio et al., 2013; Gershman et al., 2015]. These planetary ions can contribute to the mass loading and structural stability of Mercury's environment, and can also drive original plasma processes. Mercury, like Earth, has its own intrinsic magnetic field. Since this magnetic field is weaker than the magnetic field of Earth, the size of Mercury's magnetosphere is relatively small. As a result, the plasma processes at Mercury should occur faster, e.g., the Dungey plasma cycle at Mercury lasts for a few minutes only while it lasts for a few hours at Earth [Siscoe et al., 1975; Slavin et al., 2010a]. In addition, the heavy ions of planetary origin have a scale comparable to the scale of plasma processes due to their large gyration motion. Under these circumstances, the ion motion and the electromagnetic field variations within the magnetosphere are strongly coupled and thus the plasma processes at Mercury are expected to be affected by these coupling. The interactions between the planetary heavy ions and the variable electromagnetic fields at

Mercury have been investigated by various authors. In particular, Delcourt et al. (2003) investigated how magnetic field variations affect essentially the ion motion. They showed that a dipolarization in the magnetotail can accelerate ions non-adiabatically by violating the first adiabatic invariant (the magnetic moment μ). Aizawa et al. (2018) further investigated how sudden electric field variations - E burst – occurring during the development of the KH instability can also lead to non-adiabatic acceleration of ions. They show how these accelerated ions can effectively sputter the planetary surface and produce secondary neutrals and ions that populate both the exosphere and magnetosphere of Mercury. In addition, Gingell et al. (2015) revealed that the presence of heavy planetary ions can modify the growth of the KH instability, which was previously reported by MESSENGER observations [Gershman et al., 2015]. The KH instability is a macroscopic instability that occurs at several planets' magnetopause and known to be an important process for mixing plasma of different origin under a particular interplanetary magnetic field (IMF) configuration (e.g., northward for Earth and Mercury). Indeed, when the instability fully developed, it forms large rolled-up KH vortices where magnetic reconnection can occur.

In the present study, the non-adiabatic acceleration and related transport in both the dawn and dusk flanks of the magnetosphere of Mercury during the development of the KH instability are investigated in detail. Although Aizawa et al. (2018) already revealed that heavy ions can be accelerated by this process, the configuration they used was restricted to the dawn side of the magnetosphere. Since MESSENGER observations show that a strong dawn-dusk asymmetry on the KH occurrence at Mercury exists [Liljeblad et al., 2014], it appears necessary to extend the study of Aizawa et al. (2018) to the dusk side of the magnetosphere. The present study focuses on the transport of ions due to the field variations and investigates in more details the role of the convection electric field. In order to explore the dynamics of heavy ions, we performed the same numerical method that Aizawa et al. (2018) conducted, launching single

particle trajectory computations in the KH field obtained by MHD simulations. In addition, previous study considered only sodium ions whereas here we include other ions such as solar wind plasma, oxygen and potassium planetary ions, and discuss from a statistical analysis the dependence of the acceleration and transport of ions with their Larmor radius. In Section 2, we introduce our modeling approach and initial settings used. We detail the results of our simulations in Section 3 and discuss their interpretation in Section 4. Section 5 summarizes our conclusions.

2. Field Configuration and Overview of ion behavior

In order to investigate the dynamics of heavy ions of planetary origin under realistic conditions for Mercury, we adopt the same approach than the one used in Aizawa et al. (2018). We trace the motion of particles in the KH field configurations obtained from MHD simulations. A detailed description of the numerical scheme and conditions used can be found in Aizawa et al. (2018). In the present study, quantitative and systematic calculations were performed in configurations representative for both the dawn and dusk sides of Mercury's magnetosphere. All parameters are carefully chosen from observational data when available. In addition, various ion species, not only Na^+ but also O^+ , K^+ (heavy ions of planetary origin), H^+ , H_2^+ , (lighter ions of planetary origin), and H^+ , He^{2+} (solar wind plasmas) are included in the present study, as characterized by Zurbuchen et al. (2008). All ions are injected as picked up ions at given times. The size of a simulation box in the x-direction, corresponding to the KH wavelength, is set to exactly $1.5 R_M$ where R_M is Mercury radius, and the size in the y-direction is taken as $12 R_M$ with a spatial resolution of 32.2 km (114 x 909 grids). The thickness of the magnetopause (boundary layer) is ~ 515 km (referred to $2a$ hereafter) as estimated in Gershman et al. (2015). Ions are traced in the frame of the average motion of the KH wave, the periodical boundary condition being set in the x-direction. All our parameter settings are summarized in Table 1.

The variations of all physical quantities across the magnetopause follow hyperbolic tangent profiles.

Although MESSENGER observations did not enable us to derive the magnetospheric convection speed, we consider 300 km/s for the anti-sunward flow speed in the magnetosheath and 50 km/s for the sunward flow speed in the magnetosphere as used in previous numerical studies [e.g., Seki et al., 2013; Fujimoto et al., 2007; Jia et al., 2015; Wang et al., 2010]. The magnitude of the B_z magnetic field component does not vary across the magnetopause and was constrained by MESSENGER data [e.g., Wang et al., 2010]. We set three different injection times for ions so that we could investigate the characteristics of the ions dynamics during different stages of the KH development. Illustrations of our simulation conditions are shown in Figure 1. Although the figure shows only a scale of $\pm 1.0 R_M$ in the y -direction, actual calculation domain is wide enough so that the gyration motions of all ions are closed within the domain. In addition, although it appears visually that the development of the KH instability seems to be different between the dawn and dusk side configurations studied here, we confirm that the growth of the KH instability is almost the same. This behavior may be due to the initial perturbation that we used this time, i.e., white noise. Although the MHD approximation may not be valid especially at the dawnside magnetopause of Mercury according to previous studies [e.g., Paral and Rankin, 2013], we consider in the present study both the dawn and dusk sides of the magnetosphere in order to globally understand the effect of a large convection electric field on the ion dynamics and transport. In addition to a configuration with a northward IMF, we also consider a configuration with a southward IMF so that we could explore the physical mechanisms of energization and transport of ions more thoroughly.

An overview of the properties of Na^+ ions for the northward IMF Case is shown in Figure 2. Here 10,000 ions were injected in three different regions as picked up ions. Blue dots represent individual ions in the magnetosheath, while magenta and yellow dots show individual

ions in the boundary layer (magnetopause) and magnetosphere region, respectively. The spatial scale is normalized here by the radius of Mercury. Because the magnetopause (boundary layer) is the transition layer for the physical quantities, only blue (magnetosheath) and yellow (magnetosphere) dots are analyzed in this study. Initial ion positions were determined by a given function and its distribution is shown in one of the right panels of Figure 2 at $t = 32.2$ [s]. Because of the different directions of the large-scale convection electric field between the dawn and dusk sides (see left panel of Figure 2), the ion distribution during the later stage of the development of the KH instability is significantly different between the dawn and dusk sides. In particular, a large ion density gap can be observed near the magnetopause on the dawnside, while such a gap is not seen on the duskside. Moreover, ions of magnetospheric origin (yellow dots) can enter the magnetosheath region at $t = 96.6$ [s] due to the development of the KH vortices on the dawnside, whereas those on the duskside could not cross the magnetopause. Still on the dawnside, once ions penetrate inside the magnetosheath region, they seem to move freely from the filament structure of the KH vortex. On the contrary, planetary ions picked up in the magnetosheath (blue dots) did not penetrate into the magnetosphere. The situation is opposite for the duskside configuration. Ions in the magnetosheath (blue dots) can enter inside the magnetosphere region, whereas ions in the magnetosphere (yellow dots) penetrate less inside the magnetosheath region compared to those in the case of the dawnside configuration.

3. Results

3.1 From magnetosphere to magnetosheath

Since all ions were injected as picked up ions, their initial energy depends on the energy calculated from the ExB drift speed in each region. As such, the analysis and discussion should be done separately. In this section, we focus on the ions picked up in the magnetosphere region which are presented by yellow dots in Figure 2. Figure 3 shows the properties of Na^+ ions for

157 the case of northward IMF at given times. Although all colored dots represent Na^+ ions picked
 158 up in the magnetosphere, different colors correspond to the different distances from the
 159 magnetopause. Each color has an initial thickness of that of the magnetopause, $2a$. For instance,
 160 the region within $\pm 0.1 R_M$ of the center of the boundary layer is defined as a magnetopause.
 161 Magenta dots indicate ions initially picked up in the region $0.1 - 0.3 R_M$ from the magnetopause,
 162 green dots represent the ions in the region of $0.3 - 0.5 R_M$ from the magnetopause, cyan
 163 represents ions in the region of $0.5 - 0.7 R_M$ from the magnetopause, and yellow dots show the
 164 ions in the region from 0.7 to $0.9 R_M$ from the magnetopause. The Na^+ ions were injected with
 165 a thermal energy of 1 eV at $t = 32.2 \text{ [s]}$, corresponding to the time when the KH instability
 166 begins to develop. For simplicity, we chose 90° of pitch angle for all ions. At $t = 128.8$
 167 $[\text{s}]$, magenta and green dots on the dawnside seem to be distributed randomly and the
 168 corresponding ions are transported to the magnetosheath region. On the other hand, no ions are
 169 observed in the magnetosheath region on the duskside. This is because of the orientation of
 170 convection electric field. At $t = 193.2 \text{ [s]}$, ions on the dawnside are distributed throughout the
 171 simulation box, showing that they are well mixed. In contrast, only a few ions are observed
 172 in the magnetosheath region on the duskside, while they are well-mixed only inside the
 173 magnetosphere. The three bottom panels of Figure 3 shows, from top to bottom, the growth of
 174 the KH instability, the number of energized ions, and the number of transport ions as a function
 175 of time. The 2nd and 3rd panels indicate the occurrence rate, obtained by counting the number
 176 of energized/transported ions and taking the ratio to the total number of ions in each region. For
 177 instance, to obtain the data in the second panel, when a magenta ion is accelerated, and its
 178 magnetic moment becomes $\mu/\mu_0 > 10$, it is counted as an energized ion. Next, the total number
 179 of these ions is divided by the total number of ions initially injected in the magenta region. A
 180 similar counting is applied for the third panel, where net transport means that an ion crosses the
 181 magnetopause and is located up to a distance of $+2a$ from the magnetopause. For example, we

count an ion picked up in the $0.3 - 0.5 R_M$ region, colored by green for the case of the dawnside, as it crosses the magnetopause and is located above $-0.3 R_M$ in Figure 3. These rates were calculated at each computational time interval $t_{com} = 10$ (or $t = 32.2$ [s] in real time). Magenta, dark, and orange profiles in the 2nd and 3rd panels in Figure 3(b) correspond to the color dots represented in magenta and yellow, which are respectively, the closest region to and the farthest region from the magnetopause, respectively. Profiles for the properties of ions behavior colored by green and cyan are omitted here for readability.

Looking at Figure 3(b), we first find that non-adiabatic acceleration occurs on both the dawn and dusk sides, but ions are more significantly accelerated on the dawnside. The solid and dashed lines refer to the dawn and dusk configurations in Figure 3. For example, 85% of the ions are accelerated in the closest region to the magnetopause at dawn, while only 35% are seen at dusk. The starting time for acceleration closely matches the time of KH growth, suggesting that the formation of KH vortices efficiently accelerate ions. We point out here that ion transport here means that ions are observed in another region after crossing the magnetopause. On the dawnside, we observe ions having moved in another region only under the action of their gyration motion. We also observe ions with a large Larmor radius near the magnetopause after having crossed the magnetopause during their one of their cyclotron motion. This latter reason can lead to an oscillatory motion, which is illustrated in the bottom last panel in Figure 3(b). The fact that there is no dashed line in the 3rd panel of Figure 3(b) indicates that no transport occurs from the magnetosphere to the magnetosheath side on the duskside. We note the absence of ion transport even though the ions are energized. In the present analysis, we count all energized ions with $\mu/\mu_0 > 10$, thus some of them maybe even more energized. Ions with $\mu/\mu_0 \sim 10$ may not cross the boundary but ions with $\mu/\mu_0 \sim 100$ may do so because of their large Larmor radius and may in addition show the oscillation signature in the transport panel. To investigate this effect in more detail, we obtain the magnetic moment changes in the

energization and transport of ions picked up ion in the closest region to the magnetopause (indicated by A and A' in Figure 3). Figure 4 shows the breakdown of energization and transport of each range of the magnetic moment for the case of the closest region, as a function of time. First of all, the major difference between the dusk and dawn sides in the figure is seen in the energization profile. The left panel shows that energized ions have their magnetic moment changing from 10 to 50 whereas profiles in the right panel show more dispersed magnetic moment changes. On the duskside, 78 % of energized ions represented by the magenta profile have their magnetic moment changing from 10 to 25, while the blue profile shows their magnetic moment changing from 25 to 50. Figure 4 clearly shows that ions on the dawnside are effectively accelerated, and the profile in the 3rd panel shows that the lower magnetic moment change is also associated with some transport. The level of energization does not correspond to that of transport. The question that emerges here is why do we have effective acceleration and transport on the dawnside, and fewer ions are significantly energized with no transport on the duskside? To address this question, we selected model trajectories to understand the dawn-dusk difference of the energization and transport of planetary ions.

Figure 5 shows two model trajectories of picked up sodium ions in the magnetosphere for both (a) dusk and (b) dawn configurations. In this figure, we selected an ion with $\mu/\mu_0 < 10$ at dusk as an example of limited energization and an ion with $\mu/\mu_0 > 100$ at dawn as an example of efficient energization. From Top to bottom, we show the ion position in the y-direction with arrows showing the electric field orientation, the electric field intensity that the ion experiences, the magnetic moment which is normalized to the initial value, and energy. The electric field orientation shown on the 1st panel corresponds to the direction of the Lorentz force that ion experiences. For example, in Figure 5(a) at $t = 50$ [s], green arrows point towards the magnetosheath side indicating that the ion is accelerated towards that side. In contrast, at $t = 250$ [s] in the same figure, green arrows point in the opposite direction of the electric field. At

this time, the ions are essentially accelerated towards the inside of the magnetosphere. Two selected ions were initially injected at the same distance from the magnetopause. In figure 5(b), the magnetic moment changes can be seen around the linear growth time (indicated by red dashed line), and μ -jump occurred at $t \sim 115$ [s]. On the other hand, figure 5(a) does not show such clear magnetic moment changes. By also checking snapshots of the properties of an ion and the electric field that ion experiences, acceleration steps can be explained. At dawn, when the ions is just picked up in the magnetosphere, it is subject to normal gyration motion at first. The ion may experience a small electric field variation with the development of KH instability and thus the ion starts the polarization drift which gradually accelerates the ion (equivalent to Larmor radius increases). This is an adiabatic acceleration. When the ion enters the initial boundary layer region, here at $t \sim 115$ [s], the ion suddenly experiences a large electric field towards the magnetosheath where the μ -jump occurred. After that, the ion keeps its gyration motion in the magnetosheath region. A model trajectory on the duskside presented in Figure 5(a) shows that although E burst-like signatures can be seen on the 2nd panel, the ion does not gain a large amount of energy. Noting the direction of the electric field (green arrows at the 1st panel), the E burst occurs when the ion moves against the electric field, and thus deceleration occurs. Some ions are able to gain large energy such as the case of $\mu/\mu_0 \sim 30$ for instance, due to the stronger electric field inside the rotating vortex (rotating electric field hereinafter), which directs to the magnetosphere side, in the later stage of the KH development. Investigating other model trajectories allow us to understand the transport of ions, which is controlled by the electric field configuration in the magnetosheath under northward IMF.

The magnetic reconnection generally occurs during the case of southward IMF, meaning that the configuration of the magnetopause will be more complicated, thus the KH instability less likely occurs. However, this case presents a good example of understanding the role of the large convection electric field. In this whole study, we perform the ideal MHD simulation with

the reversal magnetic field (from the magnetosheath to the magnetosphere, varies from $-B_z$ to $+B_z$) across the magnetopause with the same magnitude in the entire simulation domain. It should be stressed here is that there is no magnetic reconnection so that we could have the KH growth. In the dawnside configuration, we have a large electric field towards the magnetopause in the magnetosheath region, while another small electric field points away from the magnetopause in the magnetosphere. Conversely, in the duskside configuration, we have a large electric field pointing away from the magnetopause in the magnetosheath region, and a small electric field directed towards the magnetopause in the magnetosphere. As a consequence of the southward IMF, ions picked up in the just adjacent region to the magnetopause, the energized ratio is much higher than that of the case of northward IMF for all three heavy ion species (O^+ , Na^+ , and K^+). The energized ratio increases for both dawn and dusk sides in the case of southward IMF compared to the case of northward IMF. The net electric field changes across the magnetopause, i.e. the electric fields in the case of northward IMF cancel each other whereas they add themselves up in the case of southward IMF, resulting in an effective energization in the case of southward IMF. We also investigated the breakdown of energization in the same manner as for Northward Case. The breakdown contains the lower and higher energized ions, and the amount of transport corresponding to the energization. In contrast, we found that for highly energized ions, $\mu/\mu_0 > 100$, the energization increases in time due to the rotating electric field component. On the contrary, although the magnetosheath electric field is directed toward the magnetopause at dawn, the ions are transported to the magnetosheath which is inconsistent with our results for the case of northward IMF. The dawnside configuration under the southward IMF provides the large rotating electric field inside the KH vortex pointing to the magnetosheath region, leading significant ion energization and transport in the mixing region of magnetosheath side. Under a southward IMF, signatures of energization and transport of ions picked up in the magnetosphere region can be seen on both the dawn and dusk sides. In

both cases. ions are transported and there is no stagnation inside the magnetopause. Overall, ion transport starts at a later stage in the KH linear growth, being equivalent in both cases. Even though the large convection electric field affects the ion transport, they cannot cross the magnetopause layer without the development of the KH instability. However, in reality in the magnetospheric flanks exist additional components of the magnetic field, B_x and B_y , and this may lead to the effective transport of magnetospheric ions.

3.2 From magnetosheath to magnetosphere

Here we study the properties of ions picked up in the magnetosheath region. A noticeable difference is the electric field which ions initially experience: ions picked up in the magnetosheath imply that the ion might not experience a large E burst because of the other electric field intensity affecting the ion is small compared to the electric field in the magnetosheath. The results for the northward IMF case are shown in Figure 6. We can first notice that non-adiabatic acceleration does not occurs on both the dawn and dusk sides of the magnetosphere (see 2nd panel of Figure 6(b)). Ions cannot be transported to the magnetosphere on the dawnside, while moving outward from the magnetopause following the large electric field in the magnetosheath. On the other hand, some ions can cross the magnetopause and move inside the magnetosphere region on the duskside where they are well mixed even although they are not energized. This no energization feature can be understood due to the ions initially experiencing a large electric field and thus they do not experience any stronger electric field during their gyration motions. Ion transport on the duskside starts at the time corresponding to the linear KH growth time and its depth of penetration seems to be the same order of the fluid structure represented in black and gray color, suggesting that the formation of the KH vortex plays a role in the transport of ions even though they initially have a large Larmor radius.

Mercury's exosphere extends to a few R_M [McClintok et al., 2008] and so many ions of exospheric origin can be picked up in the magnetosheath. Since these ions can have a large Larmor radius (approximately of the same order or more than the magnetopause thickness), these ions could penetrate inside the magnetosphere region. However, according to our results, without the x and y components of the magnetic field, it is difficult to cross on the dawnside. In addition, even if the ions were picked up in the duskside configuration, fewer ions were transported to the magnetosphere region compared to the results in the previous section.

Under southward IMF, a small number of ions can be energized and some transport on both the dawn and dusk sides can be seen. However, ion transport on the dawnside suddenly starts when ions are injected at the beginning of the KH growth, not because of the development of the KH instability but because of the effect of the large electric field itself. This leads to the ion transport without KH instability. In this study, because the effect of the field variations due to the development of the KH instability has been discussed, such an ion transport without the KH growth should not be taken into consideration. Additional experiments were performed to see ion transport without the KH growth and similar transport could be seen in the dawnside configuration. In addition, we found that the ion transport without the KH growth shows no large time differences even when we take into account the distance from the magnetopause. If we take into account of the ion transport without the KH growth, we can conclude that small (negligible) energization on both sides and ion transport on the duskside occur due to the development of the KH instability. Ion energization and transport occur significantly on both sides because of the same directed electric field across the magnetopause.

3.3 Properties of lighter ion species

Compared to the properties of heavy ions of planetary origin, lighter ion species that were injected with initial energies of 0.047 eV (= 540 K) have different properties. Figure 7

shows the properties of lighter ions initially picked up in the magnetosphere for northward IMF. These lighter species have a Larmor radius small enough compared to the magnetopause thickness so that they can be treated as fluid, Although some ion energization can be seen after the linear KH growth, no transport occurs. Our results indicate that the fine structure due to the development of the KH instability causes acceleration. Although the scale of fine structures cannot be characterized, it is proven by checking different species. Acceleration of hydrogen ions is observed in the region closest to the magnetopause (in magenta) at $t \sim 120$ [s], whereas it is observed at $t \sim 140$ [s] for protons with a smaller Larmor radius. Ion transport does not occur either at dawn or dusk: it is influenced only within the expanded magnetopause (or called mixing layer) due to the development of the KH instability. On the other hand, ions picked up in the magnetosheath when the IMF is northward are not energized. Small transport on the duskside can be observed similar to heavy ions. As we discussed previously, because the rotating components of the electric field work effectively on the energization and transport of ions. When the IMF is southward we found that the energization starts earlier than when the IMF is northward and that significant transport can be observed at both dawn and dusk.

Lighter ions picked up in the magnetosheath region are also observed to be energized and transported both at dawn and dusk. The main difference compared to the properties of heavy ions of planetary origin is the fact that lighter ions picked up in the magnetosheath region are significantly energized and transported when the IMF is southward. Indeed, because of their lighter masses, it is easier for them to be trapped inside KH vortices and be affected by small variations of the rotating electric field variations inside the structure.

4. Discussion

The summary tables of particle energization and transport of planetary ions for both northward and southward IMF are shown in Table 2. Generally, the level of energization is

superior when we have southward IMF, and always in the dawnside. In particular, the energization and transport of heavy ions of planetary origin picked up in the magnetosphere are illustrated in Figure 8. When the IMF is northward, we found that no acceleration of heavy planetary ions picked up in the magnetosheath is observed and that ion transport from the magnetosphere to the magnetosheath occurs only at dawn, while a small amount of heavy ion transport in the opposite direction occurs at dusk because of the electric field orientation. Because of their smaller Larmor radius, lighter ions behave magnetohydrodynamically. Acceleration of lighter planetary ions picked up in the magnetosphere occurs at both dawn and dusk but no transport towards the magnetosheath is observed.

On the other hand, when the IMF is southward, the energization and transport of both heavy and lighter ions initially picked up in the magnetosphere to the magnetosheath occurs at both dawnside and dusks, whereas the energization of heavy ions initially picked up in the magnetosheath appears negligible and there are only very little ions transported to the magnetosphere. However, in this case, the heavy ion transport from the magnetosheath to the magnetosphere at dawn can occur at dawn without the development of the KH instability. Contrary to heavy ions, lighter planetary ions picked up in the magnetosheath can be accelerated and transported due to the initial direction of the large electric field across the magnetopause. Since they are light, they have a small Larmor radius and thus, they stay inside the KH vortices. Those lighter ions are affected by small electric variations within the KH vortex, leading to efficient acceleration and transport. In the previously described cases, we have shown only two species, sodium and hydrogen ions. We now discuss the dependence on Larmor radius for the ion acceleration and transport. The different properties of heavy ions compared to lighter ions indicates that there is a relationship between the size of the KH vortex (equivalent to the magnetopause thickness) and the Larmor radius of the ion. Figure 9 shows the dependence of Larmor radius for the energization of ions picked up in the magnetosphere region, especially in

the region closest to the magnetopause (in magenta in previous figures), for the northward IMF. The x- and y-axes are scaled by ion Larmor radius normalized to the initial half magnetopause thickness and the time where the energization reaches 30% normalized to the linear growth time ($t \sim 90$ [s]). The red dashed line represents $x = 1$ and $y = 1$. The circle and triangle symbols are used for the duskside and dawnside cases, respectively. We performed additional calculations for helium, carbon, and iron ions. According to this figure, ions are less energized at dusk compared to dawn. The acceleration of lighter species take place during a late stage of the KH growth, whereas ions with a larger Larmor radius, such as Fe^+ and K^+ , are accelerated just after formation of the KH vortex. Though the oxygen ions seem to be accelerated effectively, the Larmor radius dependence on the ion energization seems to vary gradually (see Figure 9). A similar feature is observed for southward IMF. In this study, when the ions move farther than the distance of $0.3 R_M$ from the magnetopause we assume they are transported to the other region. However, the lighter planetary species follow the background fluid structure (MHD-like) and background structure does not expand more than $0.3 R_M$ this time, they are not counted. To first order, the number of planetary ions transported increases with increasing mass. For example, for northward IMF 42% of carbon ions, doubly-charged iron ions and potassium ions, 44% of oxygen ions, and 41% of sodium ions are transported from the magnetosphere to the magnetosheath at dawn, whereas in contrast none of the protons and hydrogen ions, and only 10% of helium ions are. For the southward IMF case, the corresponding values are 25% for protons, 31% for hydrogen ions, 30% for oxygen and sodium ions, and 32% for potassium ions. We may say that some threshold for MHD-like behavior can be set between helium and hydrogen which can be seen in Figure 9.

After having discussed ion acceleration and transport within KH vortices, we now discuss the possibility of solar wind plasma entry inside the magnetosphere. Two different ion species, H^+ and He^{++} are included in our modeled solar wind. These ions were injected with

406 initial energies of 10 eV on the ExB frame at the simulation start time. Although H^+ can
 407 originate both from the solar wind and from planetary ions, He^{++} cannot be produced in the
 408 Hermean environment, and thus the detection of He^{++} can be used as a proxy for the entry of
 409 solar wind plasmas inside the magnetopause. Previous results suggested that lighter ions cannot
 410 be accelerated under northward IMF and that the small amount of solar wind penetration inside
 411 the magnetosphere is expected because those lighter ions follow the background KH structure.
 412 In addition, because we assume a proton dominant background for the MHD simulation, the 1st
 413 adiabatic invariant of H^+ violation leads to the violation of the MHD approximation. Previous
 414 studies on the KH instability did not discuss the limitation of the MHD approximation on the
 415 nonlinear stage of the KH instability development because of its difficulty. However, when the
 416 1st adiabatic invariant of H^+ is not conserved, it is a clear signature of the MHD approximation
 417 violation. Here, we discuss our results for both northward and southward IMF. In both cases,
 418 we consider protons picked up in the magnetosheath region closest to the magnetopause. Their
 419 Larmor radius normalized to the magnetopause thickness is 0.018, whereas the Larmor radius
 420 of He^{++} is ~ 0.026 . In northward IMF case, similar to the behavior of lighter ions of planetary
 421 origin, solar wind plasma can also be energized within KH waves. The energized ion ratio
 422 becomes almost the same and corresponds to 20% of injected ions in the closest region to the
 423 magnetopause. Conversely, no ion transport was found on either side. Considering ion injection
 424 time (here $t = 0$ [s]) and that acceleration happens at the same time on both the dawn and dusk
 425 sides, fine structures created due to the KH vortex breaking accelerate H^+ ions. Similar
 426 behavior is observed for He^{++} , where the acceleration starts at $t \sim 122$ [s]. The time delay of
 427 the energization start for protons is reasonably understood because they have twice smaller
 428 different Larmor radius compared to that of He^{++} . Negligible ion transport has been observed
 429 at $t \sim 130$ [s] for protons, and their transport starts before the acceleration starts. This transport
 430 is understood as MHD-like transport, as the proton motion seems to follow the motion of fluid

background proton. In addition, the Larmor radius of He^{++} is large enough not to follow the fluid background. On the other hand, in Southward Case, both He^{++} and H^+ show qualitatively similar results. Ion energization and transport can be seen on both sides. The ion energization starts at $t \sim 110$ [s] for H^+ and $t \sim 108$ [s] for He^{++} . Moreover, ion transport starts at $t \sim 120$ [s] for H^+ and $t \sim 120$ [s] for He^{++} in contrast to Northward Case. Ion acceleration occurs both on the dawn and dusk sides with almost the same ratio, with ion transport being more effective on the duskside. In addition, ion transport in Southward Case starts at the same time as energization starts. This suggests that ions can be transported due to acceleration from the rotational components of the electric field inside the KH vortex in the same manner as for ions of planetary origin. Those results are summarized in Table 3. At same time, the 1st adiabatic moment of a proton at $t = 145$ [s], for Northward Case, indicates that the MHD approximation is not valid anymore. Though our injected proton had different characteristics as the background proton had, it implies that we need to consider the ion kinetics at a later stage of the KH non-linear phase. In this study, we do not focus on this, however this aspect could be studied in the future. Is it possible to have a solar wind plasma entry through a KH vortex? Our results suggest that the ions can be transported magnetohydrodynamically, meaning that the size of the solar wind plasma entry region is determined by the vortex size. Especially, even if we do not have any magnetic reconnections, solar wind plasmas are subject to enter into the magnetosphere region during southwards IMF due to the convection electric field.

Combining all these results, we can show that the ions on the dawnside are less dense than the ions on the duskside under a northward IMF. Indeed, on the duskside, ions picked up in the magnetosphere cannot cross the magnetopause, and those in the magnetosheath can penetrate inside the magnetosphere. In order to be accelerated by variations of the electric field, the ions should be inside the magnetopause where the initial electric field is smaller than that

in the other region. Moreover, the energized and transported ratio in both sides is leveling off in Figure 3 for instance, 85% (35%) of the energized ratio at dawn (dusk) in the closest region. This seems to indicate that heavy ions are not affected by small-scale structures (so-called fine structures) due to the breaking of KH vortices in the non-linear growth (saturation) phase of the instability but we note that there fine structures have an effect on the acceleration of lighter ions that have smaller Larmor radius.

5. Conclusions

Particle trajectory computations have been performed in a time-varying MHD model of KH vortices in order to understand the properties of ions within KH vortices. First of all, we revealed that the large-scale convection electric field plays an important role for the ion acceleration and transport. The results from both Northward and Southward cases showed that the ion transport is well controlled by the electric field orientation in the magnetosheath having a stronger electric field. Ions picked up in the magnetosphere are energized on both the dawn and dusk sides under both northward and southward IMF. In particular, we found when the IMF is southward that ions are efficiently accelerated within KH vortices due to the developed electric field orientation inside the vortex. Regarding lighter ions, including the solar wind plasmas, their Larmor radius is small enough compared to the magnetopause thickness and thus small enough not to be accelerated during the linear growth phase of the KH instability, resulting in the MHD-like behavior. In their non-linear stage, KH vortices start to break and make fine structures, whose scale becomes comparable to the Larmor radius. Ions can be energized with these fine structures. Non-adiabatic energization of solar wind protons goes beyond the limitation of MHD validation. In the later stage of the nonlinear phase, the ion kinetics should be taken into account for physics. In this study, although we did not consider any parallel component of the magnetic field to the wavenumber of KH and the existence of

ions already moving non-adiabatically, some general features should be detected by instruments aboard a spacecraft. First, energized ions in the magnetosphere region on both the dawn and dusk sides with the development of the KH instability are an important feature. If we have a wide range of plasma ion analyzer from a few eV to keV, ion acceleration should be observed and thus, the energy distribution would be changed between the KH and non-KH events. The FIPS instrument onboard MESSENGER had an energy range from 100 eV/e to 13 keV/e and, therefore, was too limited to identify the acceleration process for picked up ions. The BepiColombo mission now en route to Mercury includes a full set of more sophisticated charged particle instruments. In particular, the MSA instrument onboard the MIO spacecraft has a wider energy range that can cover the picked up ion acceleration theoretically found in this study, and with a wider FOV will be used for a better understanding of this non-adiabatic energization and transport within KH vortices at Mercury.

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http://rosina1.irap.omp.eu/documents/PUBLICATIONS/AIZAWA/AIZAWA2020_PSS/.

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Table 1. Summary of Parameter settings

region	Magnetosheath (MSH)	Magnetosphere (MSP)
Background H^+ number density	100 ions/cc	10 ions/cc
Flow velocity (MSO)	- 300 km/s	+ 50 km/s
Injected Planetary ion species (m/q)	H^+ (1), H_2^+ (2), O^+ (16), Na^+ (23), K^+ (39)	
Solar wind plasma in ExB frame (m/q)	H^+ (1), He^{++} (2)	
Initial particle energy	H^+ (1), H_2^+ (2) : 0.047 eV (540K) in Mercury frame	
	O^+ (16), Na^+ (23), K^+ (39) : 1 eV in Mercury frame	
	H^+ (1), He^{++} (2) : 10 eV in ExB frame	
Injection time	32.2 s, 80.5 s, 128.8 s	
Magnetic field (Bz)	+/- 48.7 nT	+ 48.7 nT

Case

	IMF (Bz)	Magnetopause thickness	KH wavelength
Northward IMF	+ 48.7 nT	515 km	1.5 R_M
Southward IMF	- 48.7 nT	515 km	1.5 R_M

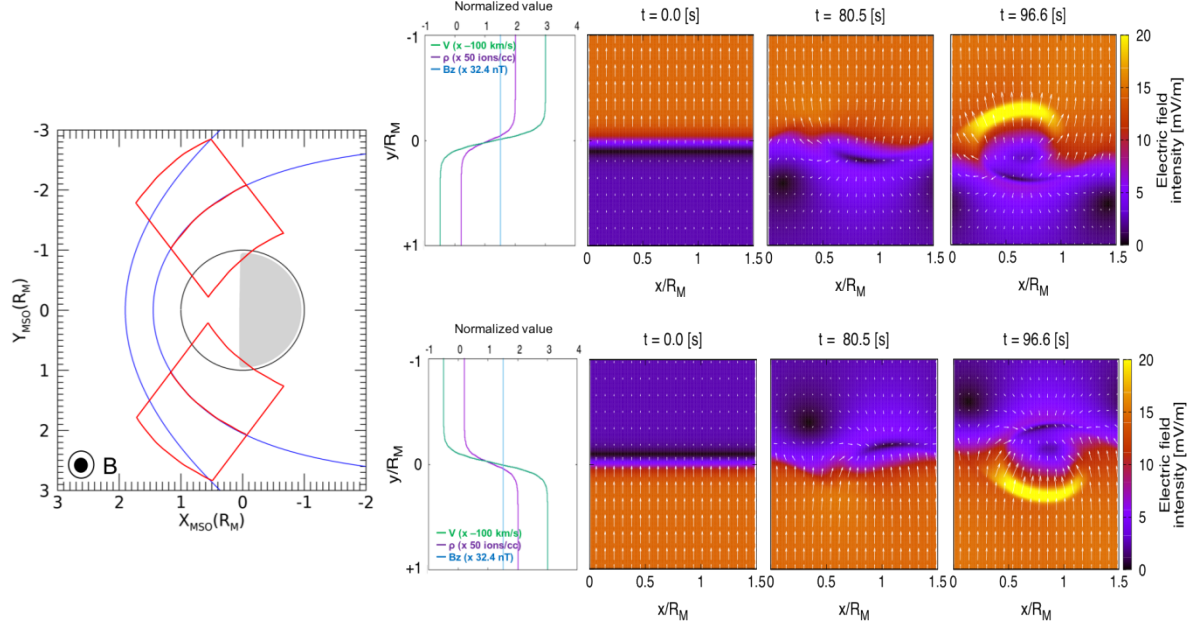


Fig 1. Illustration of simulation conditions for the case of Northward IMF.

The left panel shows the location of the bow shock (BS) and magnetopause (MP) from models by Slavin et al. (2009) for BS and Shue et al. (1997) for MP in MSO coordinates, respectively. The parameters for those paraboloid are taken from Winslow et al. (2013). The two red boxes delineate the area of the KH occurrence on both flanks of the magnetosphere of Mercury and, therefore, our simulation domain. Initial profiles of normalized physical quantities and color-coded electric field intensity obtained from MHD modeling are shown on the right panel. White arrows indicate the electric field orientation. The upper and lower color-coded panels correspond to electric field maps at distinct simulation times during the evolution of the KH instability on the dawn and dusk sides, respectively. Here we assumed that the solar wind is flowing tailward with a velocity of 300 km/s and that the magnetospheric convection flows from the nightside to the dayside with a velocity of 50 km/s. Initial variations of physical quantities across the magnetopause follow hyperbolic tangent profiles. The thickness of the

magnetopause is 515 km and the wavelength of the KH instability is set to $1.5 R_M$ km, which is consistent with MESSENGER observations.

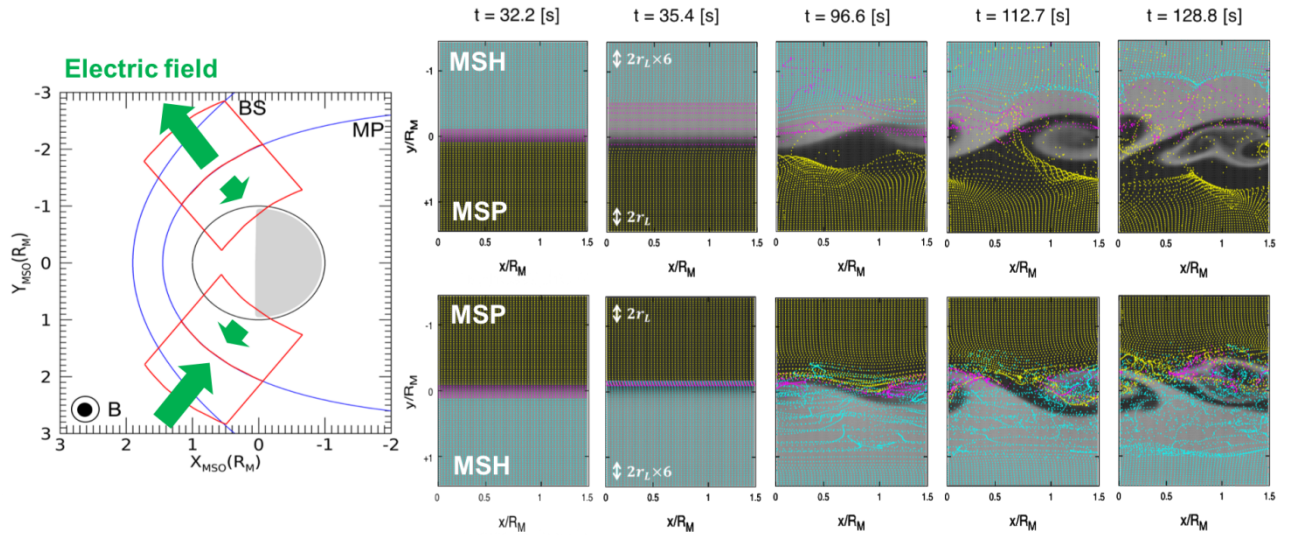


Fig 2. Properties of sodium ions at distinct simulation times during the development of the KH instability for the case of Northward IMF. Background grey scale color shows the proton number density (light (dark) colors indicate higher (lower) density). Yellow, magenta and cyan dots represent sodium ions picked up in the magnetosphere, magnetopause (boundary layer), and magnetosheath regions, respectively. The upper and lower panels show the dawn and dusk side configurations and the white arrows in the panels at $t = 35.4$ [s] indicate the Larmor radii of the sodium ions in each region. The electric field orientation and intensity is given in the left

panel. Each ion is injected with an initial thermal energy of 1 eV as a picked up ion at $t = 32.2$ [s], which corresponds to the beginning of the KH growth.

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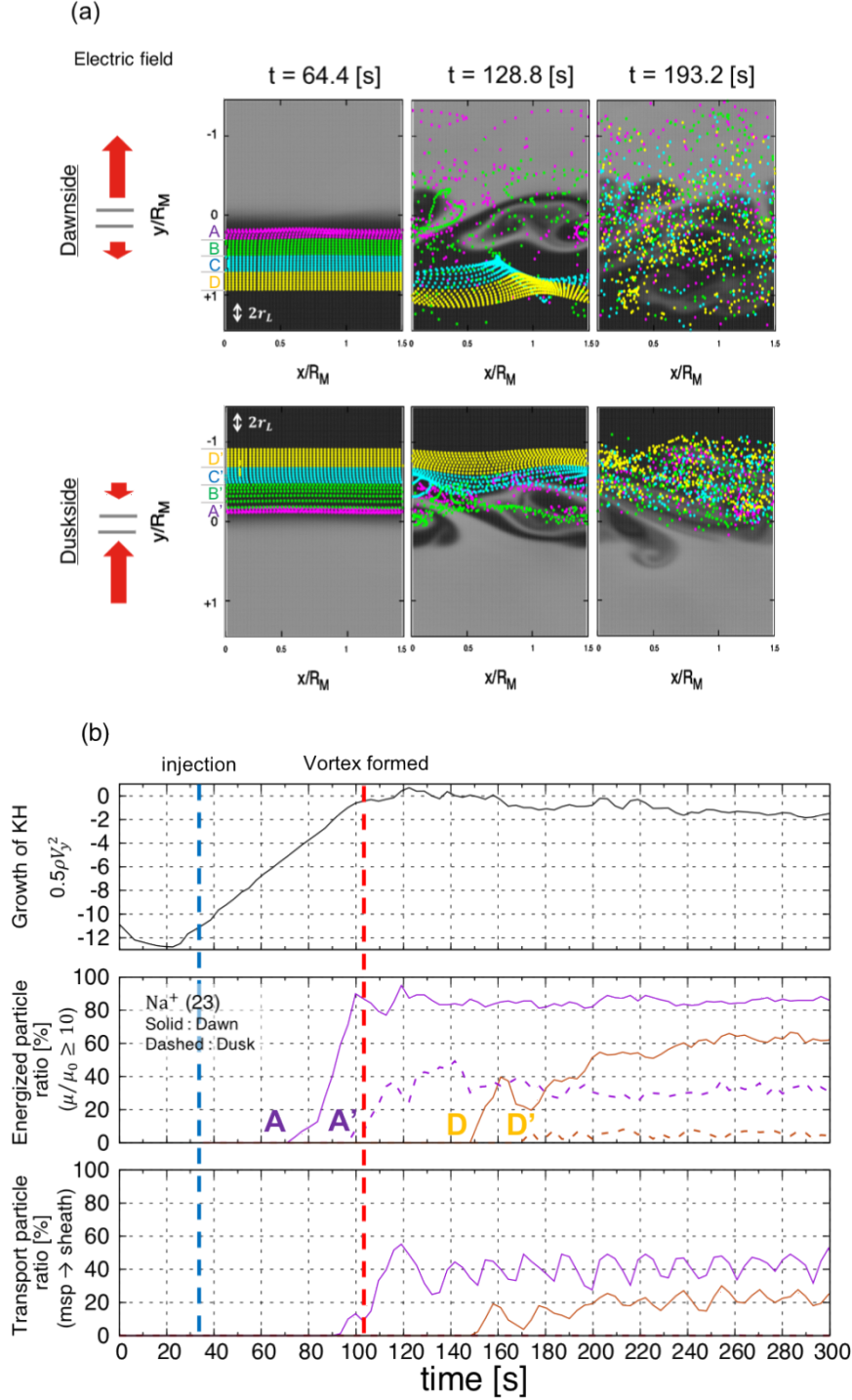


Fig 3. Energization and transport of sodium ions picked up in the magnetosphere. (a) Colored panels in the upper part of the figure represent snapshots of the properties of sodium ions at a given time in the dawn

(top) and dusk (bottom) configuration. The region in magenta (indicated by A and A') is the closest region to the magnetopause whereas the region in yellow (indicated by D and D') is the farthest region from the magnetopause. Each region has the same thickness as the magnetopause, i.e. $2a$. The white arrows indicate the Larmor radii of sodium ions at the time of injection. The initial direction of the electric field for each field configuration is given on the left. (b) The top

to bottom panels show the time series of the peak vertical kinetic energy as a proxy for the KH growth, the ratio of energized ions ($\mu/\mu_0 > 10$) to the total number of ions in each region, and the ratio of transported ions that entered the magnetosheath region ($y > -0.3 R_M$ at dawn and $y > +0.3 R_M$ at dusk) to the total number of ions in each region. The magenta and yellow line profiles refer to ions originating from the regions A, A' and D, D' in Fig. 3a, respectively. Solid and dashed lines in the 2nd and 3rd panels are used for the dawn and dusk sides, respectively. The vertical red dashed line shows the linear growth time of the KH. All ions are injected at $t = 32.2$ [s] (the vertical blue dashed line) with an initial thermal energy of 1 eV.

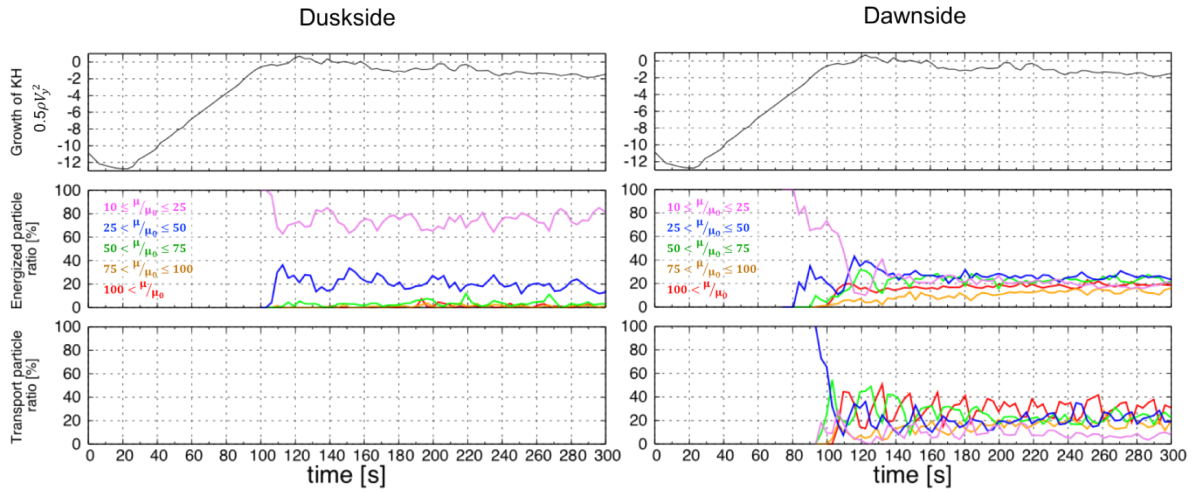


Figure 4. Breakdown of ion energization and transport for ions initially picked up in the magnetosphere in the case of Northward IMF. The left (right) panel shows the results on the duskside (dawnside). Each colored line corresponds to a change in the magnetic moment of the ions. The red line shows ions energized with $\mu/\mu_0 > 100$, the orange line is for ions with

magnetic moment in the range of $75 < \mu/\mu_0 \leq 100$, the green line is for ions with magnetic moment in the range of $50 < \mu/\mu_0 \leq 75$, the blue line is for ions with magnetic moment in the range of $25 < \mu/\mu_0 \leq 50$, and the magenta line is for ions with magnetic moment in the range of $10 \leq \mu/\mu_0 \leq 25$.

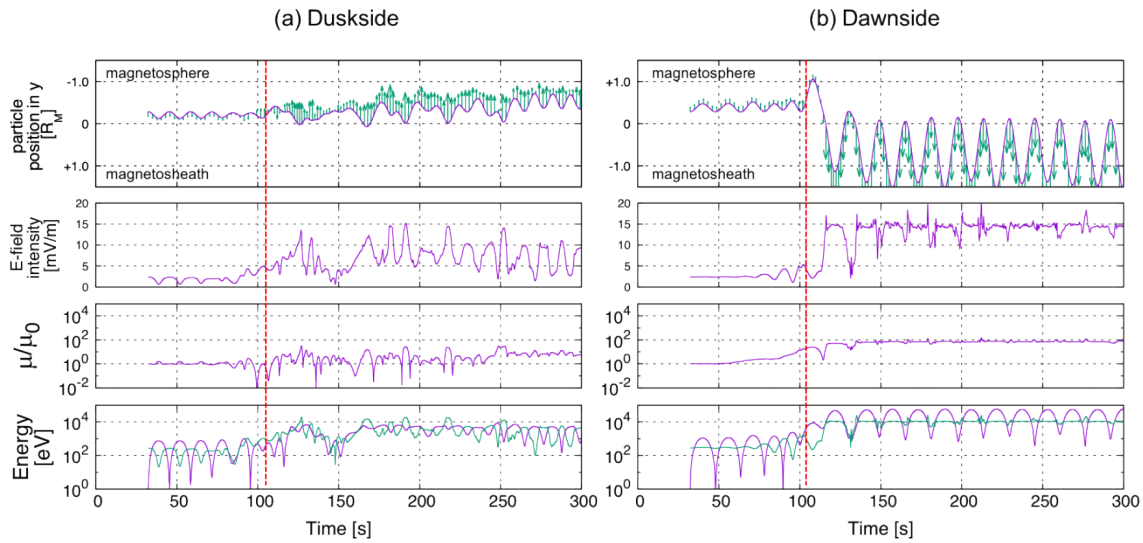
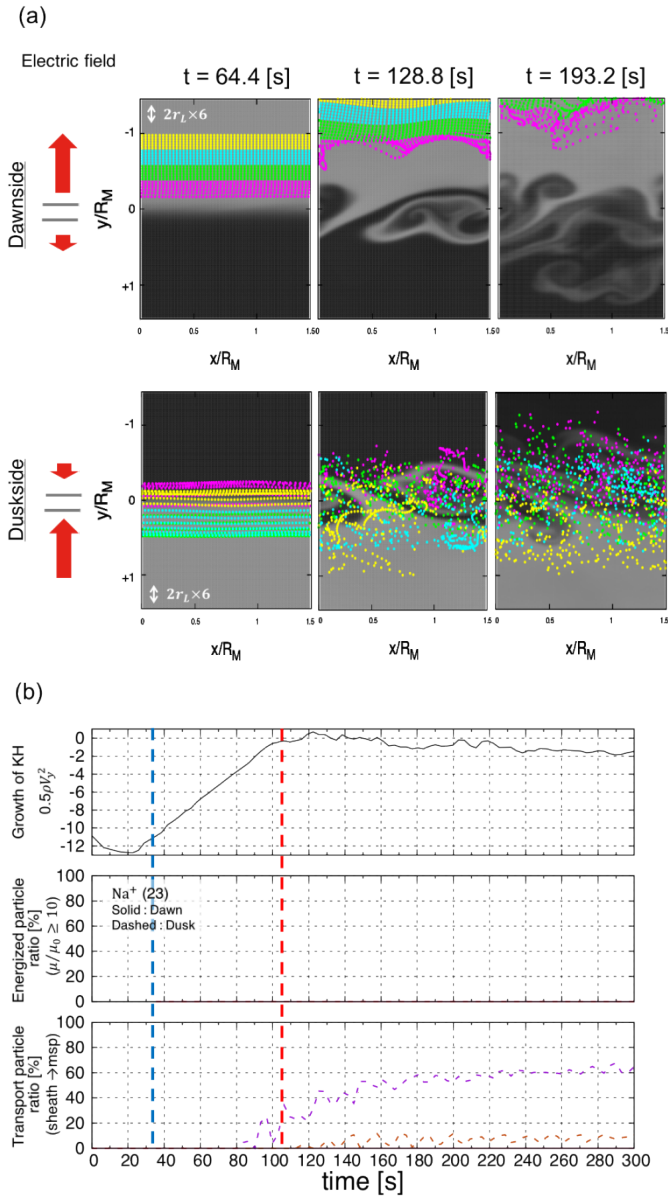


Figure 5. Model trajectories at (a) dusk and (b) dawn. The red dashed line represents the linear growth time of the KH instability. From top to bottom, each panel shows the ion position in the y-direction with the electric field orientation (green arrows), the electric field intensity, the magnetic moment (normalized to the initial value), and the energy as a function of time.

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803 Figure 6. Energization and transport of sodium ions picked up in the magnetosheath. Same
804 format as the one used in Figure 3 except that transported ions now move in the opposite
805 direction.

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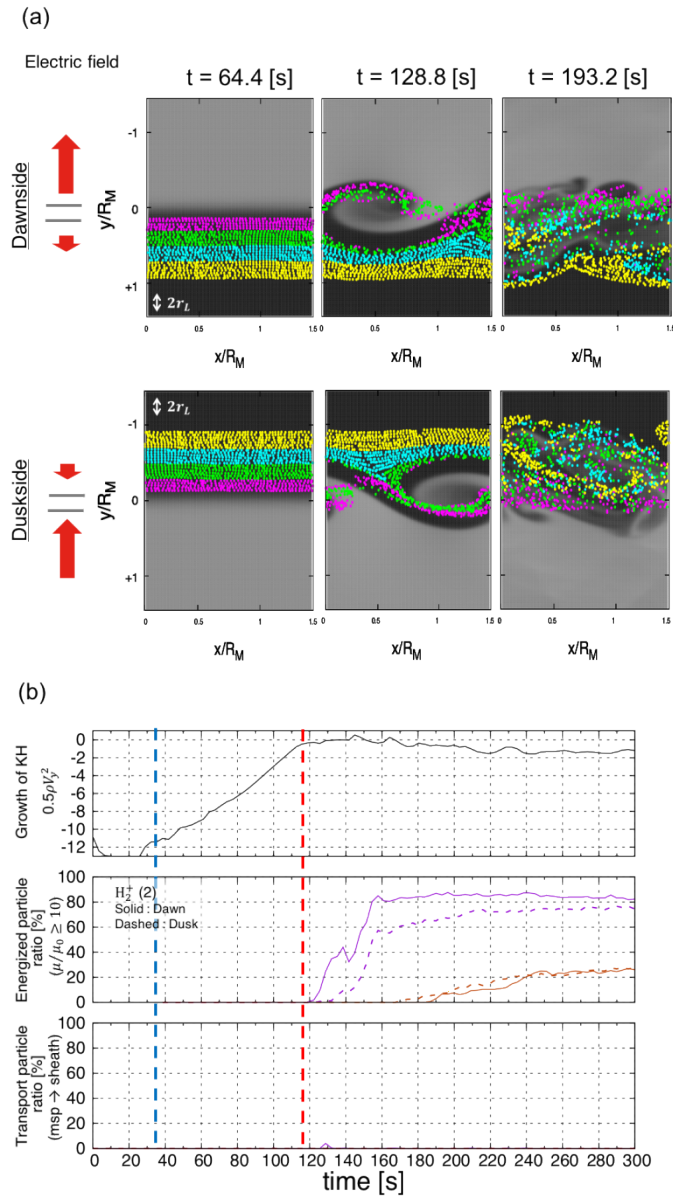


Figure 7. Energization and transport of hydrogen ions picked up in the magnetosphere. Same format as the one used in Figure 3.

(a) Northward IMF Case

IMF +Bz	Picked up in		dawn	dusk	comment
energization	MSH	h	×	×	
		l	×	×	MHD-like
	MSP	h	⊙	○	
		l	⊙	○	MHD-like
transport	MSH → MSP	h	×	○*	* transport can occur due to convection electric field
		l	×	△*	
	MSP → MSH	h	○	×	
		l	×	×	

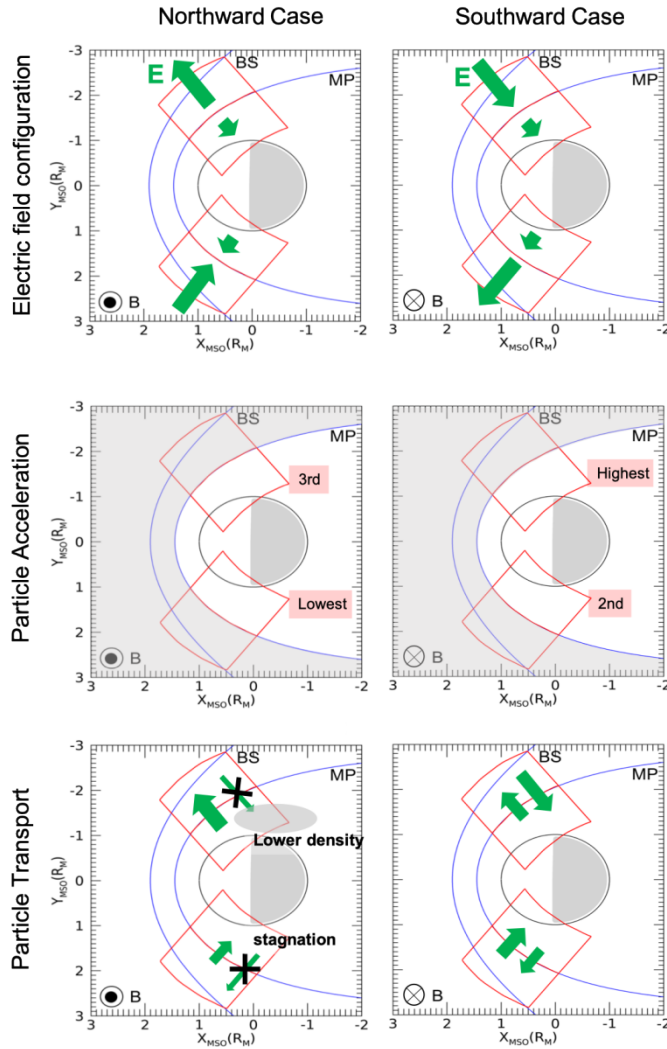
(b) Southward IMF Case

IMF -Bz	Picked up in		dawn	dusk	comment
energization	MSH	h	△	△	negligible
		l	○	○	
	MSP	h	⊙	⊙	
		l	⊙	⊙	
transport	MSH → MSP	h	○*	○	* transport can occur due to convection electric field
		l	○	○	
	MSP → MSH	h	○	○	
		l	○	○	

Table 2. Summary tables for the energization and transport of planetary ions (h for heavy ions, l for light ions) in the closest region to the magnetopause. (a) and (b) show the northward and southward IMF case, respectively. Each symbol, double circle, circle, and triangle indicates the level of energization and transport. A double circle is used when the level of energization (transport) is higher than 80%, whereas a triangle is used when those are lower than 20%. A circle represents its level between 20 – 80 %.

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848 Figure 8. Illustration of general features of ion energization and transport of heavy ions of
 849 planetary origin (O^+ , Na^+ , and K^+ in this study) for north and southward IMF cases. From
 850 top to bottom, each panel shows the initial large electric field configuration, ion acceleration,
 851 and ion transport in the x-y plane. The grey hashed area in the ion acceleration panel shows no
 852 energization.

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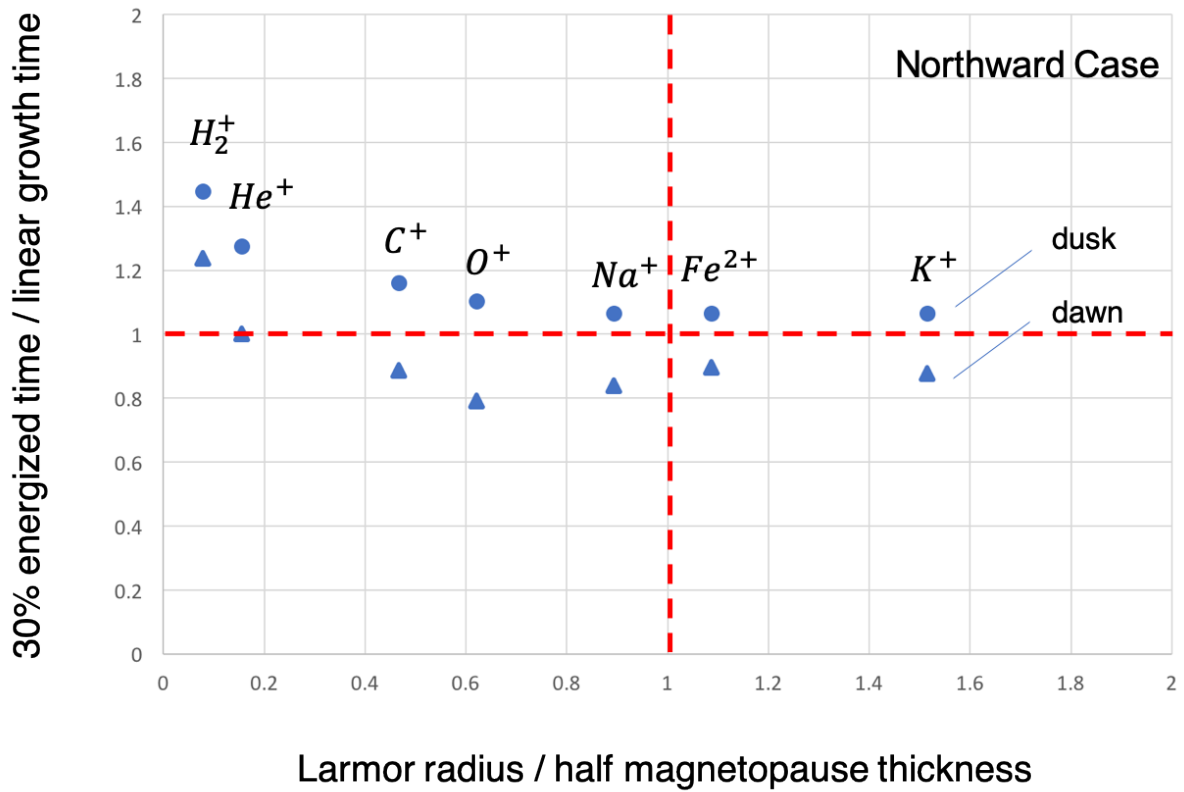


Figure 9. Larmor radius dependence on the ion energization of ions picked up in the magnetosphere, which is the region closest to the magnetopause under the northward IMF. The area below the horizontal red dashed line corresponds to the linear growth stage and the area above is nonlinear growth phase of the development of the KH instability.

	Region	IMF Bz	dawn	dusk	comment
Energization	MSH	+	⊙	⊙	

		-	⊙	⊙	
Transport	MSH → MSP	+	×	×	MHD-like
		-	○	○	MHD-like

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880 Table 3. The summary tables of energization and transport of solar wind.

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Declaration of interests

☒ The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

☐ The authors declare the following financial interests/personal relationships which may be considered as potential competing interests:

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