L^p -Theory for the Stokes and Navier-Stokes equations with different boundary conditions

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XIV Workshop on PDE, Petropolis, September 25, 2015

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I. Introduction and motivation

Let us consider in a bounded cylindrical domain $\Omega \times (0, T)$ the following Navier-Stokes equations

$$\begin{cases} \frac{\partial \boldsymbol{u}}{\partial t} - \Delta \boldsymbol{u} + \boldsymbol{u} \cdot \nabla \boldsymbol{u} + \nabla \boldsymbol{\pi} = \boldsymbol{f}, & \operatorname{div} \boldsymbol{u} = \boldsymbol{0} & \operatorname{in} & \Omega \times (0, T), \\ \boldsymbol{u}(0) = \boldsymbol{u}_0 & \operatorname{in} & \Omega, \end{cases}$$

where the unknowns \boldsymbol{u} and π stand respectively for the velocity field and the pressure of a fluid occupying a domain Ω . Given data are the external force \boldsymbol{f} and the initial velocity \boldsymbol{u}_0 . To study Navier-Stokes equations it is necessary to add appropriate boundary conditions.

- The Navier-Stokes equations are more studied, since the famous works of Jean Leray (1933, 1934), with a Dirichlet boundary condition which is not always realistic since it does not reflect the behavior of the fluid on or near the boundary.
- In the applications and in particular in the electromagnetism problems, it is possible to find problems where it is necessary to consider others boundary conditions. This boundary conditions (BC) are also used to simulate flows near rough walls, such as in aerodynamics, in weather forecasts and in hemodynamics, as well as perforated walls. BC involving the pressure, such as for example in cases of pipes, hydraulic gears using pomps, containers, etc ...

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H. Navier has suggested in 1824 a type of boundary conditions based on a proportionality between the tangential components of the normal dynamic tensor and the velocity

$$\boldsymbol{u} \cdot \boldsymbol{n} = 0, \qquad 2\nu \left[\mathbb{D} \boldsymbol{u} \cdot \boldsymbol{n} \right]_{\boldsymbol{\tau}} + \alpha \, \boldsymbol{u}_{\boldsymbol{\tau}} = 0 \qquad \text{on } \Gamma \times (0, T) \quad (1)$$

where ν is the viscosity and $\alpha \ge 0$ is the coefficient of friction and

$$\mathbb{D}\boldsymbol{u} = \frac{1}{2}(\nabla \boldsymbol{u} + \nabla \boldsymbol{u}^T)$$

denotes the deformation tensor associated to the velocity field u. Observe that if α tends to infinity, we get formally

$$\boldsymbol{u} = \boldsymbol{0}$$
 on Γ .

The Navier boundary conditions are often used to simulate flows near rough walls as well as perforated walls. Such slip boundary conditions are used in the Large Eddy Simulations (LES) of turbulent flows. Using the following relation

$$[2\mathbf{D}(\boldsymbol{v})\boldsymbol{n}]_{\boldsymbol{\tau}} = -\mathbf{curl}\,\boldsymbol{v}\times\boldsymbol{n} - 2\boldsymbol{\Lambda}\boldsymbol{v} \quad \text{on}\,\boldsymbol{\Gamma},$$

where Λ is an operator of order 0 defined by

$$oldsymbol{\Lambda} oldsymbol{v} = \sum_{k=1}^2 ig(oldsymbol{v}_{oldsymbol{ au}} \cdot rac{\partialoldsymbol{n}}{\partial s_k} ig) oldsymbol{ au}_k,$$

one can observe that in the case of a flat boundary and when $\alpha = 0$ the conditions (1) may be replaced by

$$\boldsymbol{u} \cdot \boldsymbol{n} = 0, \qquad \operatorname{\mathbf{curl}} \boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{0} \text{ on } \qquad \Gamma \times (0, T).$$
 (2)

We call them Navier-type boundary conditions.

We are interested here, in a first part, by the stationary case for the Stokes equations:

$$-\Delta \boldsymbol{u} + \nabla \boldsymbol{\pi} = \boldsymbol{f} \qquad \text{in } \Omega,$$
$$\operatorname{div} \boldsymbol{u} = 0 \qquad \text{in } \Omega,$$

with one of the following boundary conditions:

$$\boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{g} \times \boldsymbol{n} \text{ and } \boldsymbol{\pi} = \pi_0 \text{ on } \boldsymbol{\Gamma},$$
 (3)

$$\boldsymbol{u} \cdot \boldsymbol{n} = g \text{ and } \operatorname{\mathbf{curl}} \boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{h} \times \boldsymbol{n} \text{ on } \Gamma,$$
 (4)

$$\boldsymbol{u} \cdot \boldsymbol{n} = g \quad \text{and} \quad 2 \left[\mathbf{D}(\boldsymbol{u}) \boldsymbol{n} \right]_{\boldsymbol{\tau}} + \alpha \boldsymbol{u}_{\boldsymbol{\tau}} = \boldsymbol{h} \quad \text{on } \Gamma.$$
 (5)

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We will study, in a second part, the case of stationary Navier-Stokes equations:

Find
$$\boldsymbol{u}, \pi, \alpha_1, \ldots, \alpha_I$$
, with $\alpha_i \in \mathbb{R}$

$$\begin{cases} -\Delta \boldsymbol{u} + \boldsymbol{u} \cdot \nabla \boldsymbol{u} + \nabla \boldsymbol{\pi} = \boldsymbol{f} \quad \text{and} \quad \nabla \cdot \boldsymbol{u} = 0 \quad \text{in } \Omega, \\ \boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{g} \times \boldsymbol{n} & \text{on } \Gamma, \\ \boldsymbol{\pi} = \pi_0 \text{ on } \Gamma_0 \text{ and } \boldsymbol{\pi} = \pi_0 + \alpha_i & \text{on } \Gamma_i, \ i = 1, \dots, I, \end{cases}$$

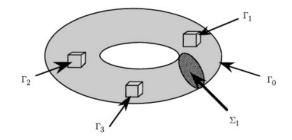
where we suppose that Ω is an open set possibly multiply connected sufficiently regular with a boundary Γ possibly non connected. We denote

$$\Gamma = \bigcup_{i=0}^{I} \Gamma_{i}, \qquad \Sigma = \bigcup_{j=1}^{J} \Sigma_{j}$$

with Γ_i the connected components of Γ and Σ_j a finite number of cuts such that

$$\Omega^{\circ} = \Omega \setminus \bigcup_{j=1}^{J} \Sigma_j$$

is simply connected.



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Some works concerning the stationary Stokes and Navier-Stokes equations with non Dirichlet boundary conditions:

- Scadilov-Solonnikov, 1973
- Conca-Murat-Pironneau, 1994
- Neustupa-Penel, 2007
- Maruzic, 2007
- Beirao da Veiga, 2009
- Berselli, 2009

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- Mitrea-Monniaux, 2009
- Amrouche-Seloula, 2011-2013
- Amrouche-Rejaiba, 2014

Considering for example the case of the Stokes equations with the homogeneous boundary conditions

$$(\mathcal{S}_T^0) \quad \begin{cases} -\Delta \, \boldsymbol{u} + \nabla \, \boldsymbol{\pi} = \boldsymbol{f} & \text{and} & \text{div} \, \boldsymbol{u} = 0 & \text{in } \Omega, \\ \boldsymbol{u} \cdot \boldsymbol{n} = 0, & \text{and} & \textbf{curl} \, \boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{0} & \text{on } \Gamma. \end{cases}$$

Because these BC, we write

$-\Delta u = \operatorname{curl} \operatorname{curl} u - \nabla \operatorname{div} u$

For the variational formulation, we will consider the following spaces:

$$V = \{ v \in L^{2}(\Omega), \text{ curl } v \in L^{2}(\Omega), \text{ div } v = 0, v \cdot n = 0 \text{ on } \Gamma \},$$
$$H_{0}^{2}(\text{div}, \Omega) = \{ v \in L^{2}(\Omega), \text{ div } v \in L^{2}(\Omega), v \cdot n = 0 \text{ on } \Gamma \}.$$

The Stokes problem (\mathcal{S}_T^0) is then equivalent to

$$\begin{cases} \text{Find} \ \boldsymbol{u} \in \boldsymbol{V} \ \text{such that} \\ \forall \ \boldsymbol{v} \in \boldsymbol{V}, \quad \int_{\Omega} \operatorname{\mathbf{curl}} \boldsymbol{u} \cdot \operatorname{\mathbf{curl}} \boldsymbol{v} \, \mathrm{d} \boldsymbol{x} = \langle \boldsymbol{f}, \ \boldsymbol{v} \rangle_{\Omega}. \end{cases}$$

We note that $[\boldsymbol{H}_0^2(\operatorname{div}, \Omega)]' \hookrightarrow \boldsymbol{H}^{-1}(\Omega).$

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Questions:

• The bilinear form is it coercive to apply the Lax-Milgram lemma ?

We will see that if Ω is simply connected, we have:

 $\forall \boldsymbol{v} \in \boldsymbol{V}, \quad \|\boldsymbol{v}\|_{\boldsymbol{H}^{1}(\Omega)} \leq C \|\operatorname{curl} \boldsymbol{v}\|_{\boldsymbol{L}^{2}(\Omega)}.$

- What happens if Ω is not simply connected ?
- Can we find generalized solution in $\boldsymbol{W}^{1,p}(\Omega)$ with 1 ?
- Can we find strong solution in $\boldsymbol{W}^{2,p}(\Omega)$ with 1 ?
- Can we find very weak solution in $L^p(\Omega)$ with 1 ?

We need the following spaces, for 1 :

$$\begin{split} \boldsymbol{H}^{p}(\boldsymbol{\mathrm{curl}},\Omega) &= \{\boldsymbol{v} \in \boldsymbol{L}^{p}(\Omega); \, \boldsymbol{\mathrm{curl}} \, \boldsymbol{v} \in \boldsymbol{L}^{p}(\Omega)\}, \ \boldsymbol{H}^{p}(\mathrm{div},\Omega) = \{\boldsymbol{v} \in \boldsymbol{L}^{p}(\Omega); \, \mathrm{div} \, \boldsymbol{v} \in L^{p}(\Omega)\}\\ \boldsymbol{X}^{p}(\Omega) &= \boldsymbol{H}^{p}(\boldsymbol{\mathrm{curl}},\Omega) \cap \boldsymbol{H}^{p}(\mathrm{div},\Omega), \end{split}$$

and their subspaces:

$$\begin{split} \boldsymbol{H}_{0}^{p}(\boldsymbol{\operatorname{curl}},\Omega) &= \left\{ \boldsymbol{v} \in \boldsymbol{H}^{p}(\boldsymbol{\operatorname{curl}},\Omega); \ \boldsymbol{v} \times \boldsymbol{n} = \boldsymbol{0} \ \mathrm{sur} \ \Gamma \right\}, \\ \boldsymbol{H}_{0}^{p}(\mathrm{div},\Omega) &= \left\{ \boldsymbol{v} \in \boldsymbol{H}^{p}(\mathrm{div},\Omega); \ \boldsymbol{v} \cdot \boldsymbol{n} = 0 \ \mathrm{sur} \ \Gamma \right\}, \\ \boldsymbol{X}_{N}^{p}(\Omega) &= \left\{ \boldsymbol{v} \in \boldsymbol{X}^{p}(\Omega); \ \boldsymbol{v} \times \boldsymbol{n} = \boldsymbol{0} \ \mathrm{sur} \ \Gamma \right\}, \quad \boldsymbol{X}_{T}^{p}(\Omega) = \left\{ \boldsymbol{v} \in \boldsymbol{X}^{p}(\Omega); \boldsymbol{v} \cdot \boldsymbol{n} = 0 \ \mathrm{sur} \ \Gamma \right\} \end{split}$$

II. Sobolev's inequalities for vector fields

Lemma (Sobolev's inequalities)

i) Any function $v \in W^{1,p}(\Omega)$ with $v \times n = 0$ on Γ satisfies:

$$\|\nabla \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} \leq C\Big(\|\operatorname{div} \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\operatorname{\mathbf{curl}} \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} + \sum_{i=1}^{I} |\langle \boldsymbol{v} \cdot \boldsymbol{n}, 1 \rangle_{\Gamma_{i}}|\Big).$$
(6)

ii) Any function $v \in W^{1,p}(\Omega)$ with $v \cdot n = 0$ on Γ satisfies:

$$\|\nabla \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} \leq C\Big(\|\operatorname{div} \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\operatorname{\mathbf{curl}} \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} + \sum_{j=1}^{J} |\langle \boldsymbol{v} \cdot \boldsymbol{n}, 1 \rangle_{\Sigma_{j}}|\Big).$$
(7)

W. Von Wahl (1992) $(I = 0, i.e \Gamma \text{ is connected and } J = 0, i.e \Omega \text{ is simply connected}).$

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Idea of the proof of (6):

• We consider the following double layer potential:

$$T \lambda(\boldsymbol{x}) = -\frac{1}{2\pi} \int_{\Gamma} \lambda(\boldsymbol{\xi}) \frac{\partial}{\partial \boldsymbol{n}} |\boldsymbol{x} - \boldsymbol{\xi}|^{-1} \, \mathrm{d}\sigma_{\boldsymbol{\xi}},$$

 $T: L^{p}(\Gamma) \longrightarrow W^{1,p}(\Gamma)$ is continuous and then compact from $L^{p}(\Gamma)$ into $L^{p}(\Gamma)$. The Fredholm alternative implies that the dimension of Ker(Id + T) is finite(equal to I) and thanks to the theorem of the open map we have for all $\boldsymbol{v} \in \boldsymbol{W}^{1,p}(\Omega)$:

$$\|\boldsymbol{v}\cdot\boldsymbol{n}\|_{L^{p}(\Gamma)} \leq C\big(\|(Id+T)(\boldsymbol{v}\cdot\boldsymbol{n})\|_{L^{p}(\Gamma)} + \sum_{i=1}^{l}|\langle \boldsymbol{v}\cdot\boldsymbol{n},1\rangle_{\Gamma_{i}}|\big)$$
(8)

• We write the following integral representation: for any $v \in W^{1,p}(\Omega)$ with $v \times n = 0$ on Γ :

$$(Id+T)(\boldsymbol{v}\cdot\boldsymbol{n}) = -\frac{1}{2\pi} \left(\operatorname{\mathbf{grad}} \int_{\Omega} \frac{1}{|\boldsymbol{x}-\boldsymbol{y}|} \operatorname{div}_{\boldsymbol{y}} \boldsymbol{v}(\boldsymbol{y}) \, \mathrm{d}\boldsymbol{y} \right) \cdot \boldsymbol{n} - \frac{1}{2\pi} \left(\operatorname{\mathbf{rot}} \int_{\Omega} \frac{1}{|\boldsymbol{x}-\boldsymbol{y}|} \operatorname{\mathbf{rot}}_{\boldsymbol{y}} \boldsymbol{v}(\boldsymbol{y}) \, \mathrm{d}\boldsymbol{y} \right) \cdot \boldsymbol{n}$$
(9)

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• Using the traces theorem, we prove that:

$$\begin{split} \|(Id+T)(\boldsymbol{v}\cdot\boldsymbol{n})\|_{L^p(\Gamma)} &\leq \quad C\Big(\Big\|\mathbf{grad}\int_{\Omega}\frac{1}{|\boldsymbol{x}-\boldsymbol{y}|}\operatorname{div}_{\boldsymbol{y}}\boldsymbol{v}(\boldsymbol{y})\,\mathrm{d}\boldsymbol{y}\Big\|_{\boldsymbol{W}^{1,p}(\Omega)} \\ &+ \quad \Big\|\operatorname{rot}\int_{\Omega}\frac{1}{|\boldsymbol{x}-\boldsymbol{y}|}\operatorname{rot}_{\boldsymbol{y}}\boldsymbol{v}(\boldsymbol{y})\,\mathrm{d}\boldsymbol{y}\Big\|_{\boldsymbol{W}^{1,p}(\Omega)}\Big). \end{split}$$

• Using the Calderón-Zygmund inequalities, we get:

 $\|(Id+T)(\boldsymbol{v}\cdot\boldsymbol{n})\|_{L^p(\Gamma)} \leq C \left(\|\operatorname{div}\boldsymbol{v}\|_{L^p(\Omega)} + \|\operatorname{rot}\boldsymbol{v}\|_{L^p(\Omega)}\right).$

• From (8), we deduce:

$$\|\boldsymbol{v}\cdot\boldsymbol{n}\|_{L^{p}(\Gamma)} \leq C\left(\|\operatorname{div}\boldsymbol{v}\|_{L^{p}(\Omega)} + \|\operatorname{rot}\boldsymbol{v}\|_{L^{p}(\Omega)} + \sum_{i=1}^{l} |\langle\boldsymbol{v}\cdot\boldsymbol{n},1\rangle_{\Gamma_{i}}|\right).$$
(10)

• Moreover, from (9) we have:

$$\begin{split} \|\boldsymbol{v}\cdot\boldsymbol{n}\|_{W^{1-\frac{1}{p},p}(\Gamma)} &\leq \left(\|T(\boldsymbol{v}\cdot\boldsymbol{n})\|_{W^{1-\frac{1}{p},p}(\Gamma)} + \left\|\mathbf{grad}\int_{\Omega} \frac{1}{|\boldsymbol{x}-\boldsymbol{y}|}\operatorname{div}_{\boldsymbol{y}}\boldsymbol{v}(\boldsymbol{y})\,\mathrm{d}\boldsymbol{y}\right\|_{W^{1-\frac{1}{p},p}(\Gamma)} \\ &+ \left\|\operatorname{rot}\int_{\Omega} \frac{1}{|\boldsymbol{x}-\boldsymbol{y}|}\operatorname{rot}_{\boldsymbol{y}}\boldsymbol{v}(\boldsymbol{y})\,\mathrm{d}\boldsymbol{y}\right\|_{W^{1-\frac{1}{p},p}(\Gamma)}\right), \end{split}$$

and by the traces theorem:

$$\begin{split} \|\boldsymbol{v}\cdot\boldsymbol{n}\|_{W^{1-\frac{1}{p},p}(\Gamma)} &\leq \quad C\Big(\|\boldsymbol{v}\cdot\boldsymbol{n}\|_{L^{p}(\Gamma)} + \left\|\operatorname{grad}\int_{\Omega}\frac{1}{|\boldsymbol{x}-\boldsymbol{y}|}\operatorname{div}_{\boldsymbol{y}}\boldsymbol{v}(\boldsymbol{y})\,\mathrm{d}\boldsymbol{y}\right\|_{W^{1,p}(\Omega)} \\ &+ \quad \left\|\operatorname{rot}\int_{\Omega}\frac{1}{|\boldsymbol{x}-\boldsymbol{y}|}\operatorname{rot}_{\boldsymbol{y}}\boldsymbol{v}(\boldsymbol{y})\,\mathrm{d}\boldsymbol{y}\right\|_{W^{1,p}(\Omega)}\Big). \end{split}$$

• Using again the Calderón-Zygmund inequalities and (10), we obtain:

$$\|\boldsymbol{v}\cdot\boldsymbol{n}\|_{W^{1-\frac{1}{p},p}(\Gamma)} \leq C\left(\|\operatorname{div}\boldsymbol{v}\|_{L^{p}(\Omega)} + \|\operatorname{rot}\boldsymbol{v}\|_{L^{p}(\Omega)} + \sum_{i=1}^{I} |\langle \boldsymbol{v}\cdot\boldsymbol{n}, 1\rangle_{\Gamma_{i}}|\right).$$
(11)

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• As $v \in W^{1-\frac{1}{p},p}(\Gamma)$, by the traces theorem, there exists $u \in W^{1,p}(\Omega)$ such that:

$$\boldsymbol{v} = \boldsymbol{u} \quad ext{sur } \Gamma \quad ext{ and } \quad \| \boldsymbol{u} \|_{\boldsymbol{W}^{1,p}(\Omega)} \leq C \| \boldsymbol{v} \|_{\boldsymbol{W}^{1-\frac{1}{p},p}(\Gamma)}.$$

Because $\boldsymbol{v} \times \boldsymbol{n} = 0$ on Γ , we have $\boldsymbol{v}|_{\Gamma} = (\boldsymbol{v} \cdot \boldsymbol{n})\boldsymbol{n}$. Using then (11), we get

$$\begin{aligned} \|\boldsymbol{u}\|_{\boldsymbol{W}^{1,p}(\Omega)} &\leq C \|\boldsymbol{v}\cdot\boldsymbol{n}\|_{\boldsymbol{W}^{1-\frac{1}{p},p}(\Gamma)} \\ &\leq C \Big(\|\operatorname{div}\boldsymbol{v}\|_{L^{p}(\Omega)} + \|\operatorname{rot}\boldsymbol{v}\|_{L^{p}(\Omega)} + \sum_{i=1}^{I} |\langle \boldsymbol{v}\cdot\boldsymbol{n},1\rangle_{\Gamma_{i}}| \Big). \end{aligned}$$
(12)

As u − v ∈ W₀^{1,p}(Ω), we know that for any function w belonging to W₀^{1,p}(Ω), we have the following integral representation integrale:

$$\boldsymbol{w} = -\mathbf{grad} \ \frac{1}{4\pi} \int_{\Omega} \frac{1}{|\boldsymbol{x} - \boldsymbol{y}|} \operatorname{div}_{\boldsymbol{y}} \boldsymbol{w}(\boldsymbol{y}) \, \mathrm{d}\boldsymbol{y} + \mathbf{rot} \ \frac{1}{4\pi} \int_{\Omega} \frac{1}{|\boldsymbol{x} - \boldsymbol{y}|} \operatorname{rot}_{\boldsymbol{y}} \boldsymbol{w}(\boldsymbol{y}) \, \mathrm{d}\boldsymbol{y}.$$

Thanks to the Calderón-Zygmund inequalities, we have

$$\|\nabla \boldsymbol{w}\|_{\boldsymbol{L}^{p}(\Omega)} \leq C \big(\|\operatorname{div} \boldsymbol{w}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\operatorname{rot} \boldsymbol{w}\|_{\boldsymbol{L}^{p}(\Omega)} \big).$$
(13)

• Applying (13) to $\boldsymbol{v} - \boldsymbol{u} \in \boldsymbol{W}_0^{1,p}(\Omega)$, we obtain:

$$\|\nabla \left(\boldsymbol{v} - \boldsymbol{u}\right)\|_{\boldsymbol{L}^{p}(\Omega)} \leq C \Big(\|\mathrm{div}\,\boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\mathrm{div}\,\boldsymbol{u}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\mathrm{rot}\,\boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\mathrm{rot}\,\boldsymbol{u}\|_{\boldsymbol{L}^{p}(\Omega)}\Big)$$

Finally, we deduce the required estimate by using (12), $\Box \rightarrow \langle \neg \rangle \rightarrow \langle \neg \rangle \rightarrow \langle \neg \rangle$

Idea of the proof of (7):

We use the same ideas with the following linear integral operator

$$R \boldsymbol{\lambda}(\boldsymbol{x}) = \frac{1}{2\pi} \int_{\Gamma} \operatorname{\mathbf{curl}} \left(\frac{\boldsymbol{\lambda}(\boldsymbol{\xi})}{|\boldsymbol{x} - \boldsymbol{\xi}|} \right) \times \boldsymbol{n} \, \mathrm{d}\sigma_{\boldsymbol{\xi}}.$$

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Because the embeddings $\boldsymbol{X}_{N}^{p}(\Omega) \hookrightarrow \boldsymbol{L}^{p}(\Omega)$ and $\boldsymbol{X}_{T}^{p}(\Omega) \hookrightarrow \boldsymbol{L}^{p}(\Omega)$ are compact, by using Peetre-Tartar Theorem and the previous inequality, we deduce the following Poincaré's type inequalities:

Corollary (Poincaré's Inequalities)

i) We have
$$\boldsymbol{v} \in \boldsymbol{X}_N^p(\Omega) \hookrightarrow \boldsymbol{W}^{1,p}(\Omega)$$
 and for any $\boldsymbol{v} \in \boldsymbol{X}_N^p(\Omega)$

$$\|\boldsymbol{v}\|_{\boldsymbol{W}^{1,p}(\Omega)} \leq C\Big(\|\operatorname{div}\boldsymbol{v}\|_{L^{p}(\Omega)} + \|\operatorname{curl}\boldsymbol{v}\|_{L^{p}(\Omega)} + \sum_{i=1}^{I} |\langle \boldsymbol{v} \cdot \boldsymbol{n}, 1 \rangle_{\Gamma_{i}}|\Big).$$

ii) We have
$$\boldsymbol{v} \in \boldsymbol{X}_T^p(\Omega) \hookrightarrow \boldsymbol{W}^{1,p}(\Omega)$$
 and for any $\boldsymbol{v} \in \boldsymbol{X}_T^p(\Omega)$

$$\|\boldsymbol{v}\|_{\boldsymbol{W}^{1,p}(\Omega)} \leq C\Big(\|\operatorname{div}\boldsymbol{v}\|_{L^{p}(\Omega)} + \|\operatorname{\mathbf{curl}}\boldsymbol{v}\|_{L^{p}(\Omega)} + \sum_{j=1}^{J} |\langle \boldsymbol{v} \cdot \boldsymbol{n}, 1 \rangle_{\Sigma_{j}}|\Big).$$

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Regularity properties

The previous results can be extended to the case of inhomogeneous boundary conditions :

• Let Ω of class $\mathcal{C}^{1,1}$ and $\boldsymbol{v} \in \boldsymbol{X}^p(\Omega)$ with

$$\boldsymbol{v} \cdot \boldsymbol{n} \in W^{1-\frac{1}{p},p}(\Gamma) \quad \text{or} \quad \boldsymbol{v} \times \boldsymbol{n} \in \boldsymbol{W}^{1-\frac{1}{p},p}(\Gamma).$$

Then

$$\boldsymbol{v} \in \boldsymbol{W}^{1,p}(\Omega).$$

Moreover

$$\|\boldsymbol{v}\|_{\boldsymbol{W}^{1,p}(\Omega)} \leq C \left(\|\boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\mathbf{curl}\,\boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\mathrm{div}\,\boldsymbol{v}\|_{L^{p}(\Omega)} + \|\boldsymbol{v}\cdot\boldsymbol{n}\|_{W^{1-\frac{1}{p},p}(\Gamma)}\right)$$

or

 $\|\boldsymbol{v}\|_{\boldsymbol{W}^{1,p}(\Omega)} \leq C \big(\|\boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\mathbf{curl}\,\boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\mathrm{div}\,\boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\boldsymbol{v}\times\boldsymbol{n}\|_{\boldsymbol{W}^{1-\frac{1}{p},p}(\Omega)} + \|\boldsymbol{v}\times\boldsymbol{v}\|_{\boldsymbol{W}^{1-\frac{1}{p},p}(\Omega)} + \|\boldsymbol{v}\|_{\boldsymbol{W}^{1-\frac{1}{p},p}(\Omega)} +$

• Let Ω of class $\mathcal{C}^{2,1}$ and $\boldsymbol{v} \in \boldsymbol{L}^p(\Omega), \quad \operatorname{\mathbf{curl}} \boldsymbol{v} \in \boldsymbol{W}^{1,p}(\Omega), \quad \operatorname{div} \boldsymbol{v} \in W^{1,p}(\Omega)$ with

$$oldsymbol{v}\cdotoldsymbol{n}\in W^{2-rac{1}{p},p}(\Gamma) \quad ext{or} \quad oldsymbol{v} imesoldsymbol{n}\inoldsymbol{W}^{2-rac{1}{p},p}(\Gamma).$$

Then

$$\boldsymbol{v} \in \boldsymbol{W}^{2,p}(\Omega),$$

with the corresponding estimates.

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III. *L*^{*p*}**-Theory for Vector Potentials**

To study the Stokes problems, we need some results concerning the vector potentials and the inf-sup conditions.

Theorem (Tangential vector potentials)

$$\begin{split} \boldsymbol{u} \in \boldsymbol{H}^{\,p}(\text{div},\,\Omega) \text{ satisfies: div } \boldsymbol{u} = 0 \quad \text{in }\Omega, \quad \langle \boldsymbol{u} \cdot \boldsymbol{n},\, 1 \rangle_{\Gamma_i} = 0, \quad 0 \leq i \leq I, \\ \text{iff there exists a vector potential } \boldsymbol{\psi} \text{ in } \boldsymbol{X}^{\,p}(\Omega) \text{ such that} \end{split}$$

$$\boldsymbol{u} = \operatorname{\mathbf{curl}} \boldsymbol{\psi} \quad \text{and} \quad \operatorname{div} \boldsymbol{\psi} = 0 \text{ in } \Omega, \tag{14}$$

$$\boldsymbol{\psi} \cdot \boldsymbol{n} = 0 \qquad \text{on } \boldsymbol{\Gamma},$$
 (15)

$$\langle \boldsymbol{\psi} \cdot \boldsymbol{n}, 1 \rangle_{\Sigma_j} = 0, \ 1 \le j \le J.$$
(16)

This function is unique and moreover, we have

$$\|\boldsymbol{\psi}\|_{\boldsymbol{W}^{1,p}(\Omega)} \leq C \|\boldsymbol{u}\|_{\boldsymbol{L}^{p}(\Omega)}.$$

For p = 2, see : C. Amrouche, C. Bernardi, M. Dauge, V. Girault [1998].

Remark: If

 $\boldsymbol{u} \in \boldsymbol{L}^p(\Omega)$ and div $\boldsymbol{u} = 0$,

the condition

$$\langle \boldsymbol{u} \cdot \boldsymbol{n}, 1 \rangle_{\Gamma_i} = 0, \quad 0 \le i \le I,$$

is necessary and sufficient for the existence of vector potential ψ satisfying (14) and (15). While the condition (16)

$$\langle \boldsymbol{\psi} \cdot \boldsymbol{n}, \quad 1 \rangle_{\Sigma_j} = 0, \ 1 \leq j \leq J$$

ensures the uniqueness of ψ . When this last condition is not satisfied, the vector potential ψ exits up an additive vector function of the the following corresponding kernel:

$$oldsymbol{K}_T^p(\Omega) = \{oldsymbol{w} \in oldsymbol{X}_T^p(\Omega); \ \mathbf{curl} \, oldsymbol{w} = oldsymbol{0}, \ \mathrm{div} \, oldsymbol{w} = 0 \ \mathrm{in} \ \Omega\}.$$

To characterize this kernel, let us introduce the following space

$$\Theta = \left\{ r \in H^1(\Omega^\circ); \ [r]_j = \text{constant}, \ 1 \le j \le J \right\}$$

For each $1 \leq j \leq J$, there exists a unique solution $q_j^T \in H^1(\Omega^\circ)$, up to an additive constant, satisfying

$$orall r \in \Theta, \quad \int_{\Omega^\circ}
abla q_j^T \cdot
abla r = [r]_j$$

where $[\chi]_j$ is the jump of χ through Σ_j .

In fact, we have

$$\begin{cases} -\Delta q_j^T = 0 & \text{in } \Omega^\circ, \\ \partial_n \, q_j^T = 0 & \text{on } \Gamma, \\ \left[\begin{array}{c} q_j^T \end{array} \right]_k = \text{constant} & \text{and} \quad \left[\begin{array}{c} \partial_n \, q_j^T \end{array} \right]_k = 0, \ 1 \le k \le J, \\ \left\langle \begin{array}{c} \partial_n \, q_j^T, \, 1 \right\rangle_{\Sigma_k} = \delta_{j\,k}, \ 1 \le k \le J. \end{cases} \end{cases}$$

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Let us consider any extension to $L^2(\Omega)$ of ∇q_j^T which is denoted by $\widetilde{\mathbf{grad}} q_j^T$. It is easy to verify that

div
$$\widetilde{\mathbf{grad}} q_j^T = 0$$
, $\mathbf{curl} \, \widetilde{\mathbf{grad}} \, q_j^T = \mathbf{0}$

and thanks to the regularity results of vector fields, we deduce first that

$$\widetilde{\operatorname{grad}} \, q_j^T \in \boldsymbol{H}^1(\Omega),$$

and then by boostrap argument that

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$$\widetilde{\operatorname{\mathbf{grad}}} q_j^T \in W^{1,p}(\Omega),$$

Moreover, $\boldsymbol{K}_{T}^{p}(\Omega)$ is spanned by

$$\widetilde{\operatorname{\mathbf{grad}}} q_j^T, \qquad j = 1, \dots, J.$$

Remark that $\mathbf{K}_T^p(\Omega) = \{0\}$ if J = 0 (i.e Ω is simply connected).

Theorem (Inf-Sup condition in Banach spaces)

Let X and M be two reflexive Banach spaces and X' and M' their dual spaces. Let a be the continuous bilinear form defined on $X \times M$, let $A \in \mathcal{L}(X; M')$ and $A' \in \mathcal{L}(M; X')$ be the operators defined by

$$\forall v \in X, \; \forall w \in M, \; a(v,w) = < Av, w > = < v, A'w >$$

and V = Ker A. The following statements are equivalent: i) There exist $\beta > 0$ such that

$$\inf_{\substack{w \in M \\ w \neq 0}} \sup_{\substack{v \in X \\ v \neq 0}} \frac{a(v, w)}{\|v\|_X \|w\|_M} \ge \beta.$$
(17)

ii) The operator $A: X/V \mapsto M'$ is an isomophism and $1/\beta$ is the continuity constant of A^{-1} . iii) The operator $A': M \mapsto X' \perp V$ is an isomophism and $1/\beta$ is the continuity constant of $(A')^{-1}$.

Remark

As consequence, if the Inf-Sup condition (17) is satisfied, then we have the following properties:

i) Because $A': M \mapsto X' \perp V$ is an isomophism, then for any $f \in X'$, satisfying the compatibility condition

 $\forall v \in V, < f, v > = 0,$

there exists a unique $w \in M$ such that

$$\forall v \in X, \ a(v,w) = \langle f, v \rangle \text{ and } \|w\|_M \le \frac{1}{\beta} \|f\|_{X'}.$$
 (18)

ii) Because $A : X/V \mapsto M'$ is an isomophism, then for any $g \in M', \exists v \in X$, unique up an additive element of V, such that:

$$\forall w \in M, \ a(v,w) = \langle g, w \rangle \quad \text{and} \quad \|v\|_{X/V} \le \frac{1}{\beta} \|g\|_{M'}$$

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We define the space

 $\boldsymbol{V}^p_T(\Omega) = \{ \boldsymbol{w} \in \boldsymbol{X}^p_T(\Omega); \, \text{div} \, \boldsymbol{w} = 0 \text{ in } \Omega \text{ and } \langle \boldsymbol{w} \cdot \boldsymbol{n}, \, 1 \rangle_{\Sigma_j} = 0, \, 1 \leq j \leq J \},$

which is a Banach space for the norm $\|\cdot\|_{\boldsymbol{X}^p(\Omega)}$.

Using some classical Helmholtz decomposition and the previous tangential vector potential, we can deduce:

Lemma (Inf Sup Condition)

The following Inf-Sup Condition holds: there exists a constant $\beta > 0$, such that

$$\inf_{\substack{\boldsymbol{\varphi} \in \boldsymbol{V}_T^{p'}(\Omega) \\ \boldsymbol{\varphi} \neq \boldsymbol{0}}} \sup_{\substack{\boldsymbol{\xi} \in \boldsymbol{V}_T^p(\Omega) \\ \boldsymbol{\xi} \neq \boldsymbol{0}}} \frac{\int_{\Omega} \operatorname{curl} \boldsymbol{\xi} \cdot \operatorname{curl} \boldsymbol{\varphi} \, \mathrm{d} \boldsymbol{x}}{\|\boldsymbol{\xi}\|_{\boldsymbol{X}_T^p(\Omega)} \|\boldsymbol{\varphi}\|_{\boldsymbol{X}_T^{p'}(\Omega)}} \ge \beta. \tag{19}$$

Theorem (Normal vector potentials)

 $\boldsymbol{u} \in \boldsymbol{L}^{p}(\Omega)$ satisfies:

$$\boldsymbol{u} \cdot \boldsymbol{n} = 0$$
 on Γ and div $\boldsymbol{u} = 0$ in Ω , $\langle \boldsymbol{u} \cdot \boldsymbol{n}, 1 \rangle_{\Sigma_{i}} = 0$, $0 \leq j \leq J$,

iff there exists a vector potential $\boldsymbol{\psi}$ in $\boldsymbol{X}^{p}(\Omega)$ such that

$$\boldsymbol{u} = \operatorname{\mathbf{curl}} \boldsymbol{\psi} \quad \text{and} \quad \operatorname{div} \boldsymbol{\psi} = 0 \text{ in } \Omega,$$
 (20)

$$\psi \times n = 0$$
 on Γ , (21)

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$$\langle \boldsymbol{\psi} \cdot \boldsymbol{n}, 1 \rangle_{\Gamma_i} = 0, \ 1 \le i \le I.$$

This function is unique and moreover, we have:

 $\|\boldsymbol{\psi}\|_{\boldsymbol{W}^{1,p}(\Omega)} \leq C \|\boldsymbol{u}\|_{\boldsymbol{L}^{p}(\Omega)}.$

Remark : For

 $\boldsymbol{u} \in \boldsymbol{L}^{p}(\Omega),$ with div $\boldsymbol{u} = 0$ in Ω and $\boldsymbol{u} \cdot \boldsymbol{n} = 0$ on Γ , the condition

$$\langle \boldsymbol{u} \cdot \boldsymbol{n}, 1 \rangle_{\Sigma_j} = 0, \quad \text{for } 1 \le j \le J,$$

is necessary and sufficient for the existence of the vector potential $\boldsymbol{\psi}$ satisfying (20) and (21). The condition

$$\langle \boldsymbol{\psi} \cdot \boldsymbol{n}, 1 \rangle_{\Gamma_i} = 0, \quad \text{for } 0 \le i \le I$$

ensures the uniqueness of ψ .

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Idea of the proof

• We use the Inf-Sup condition (19) to solve the problem: Find $\boldsymbol{\xi} \in \boldsymbol{W}^{1,p}(\Omega)$ such that:

$$(\mathcal{P}) \quad \begin{cases} -\Delta \boldsymbol{\xi} = 0 \quad \text{and } \operatorname{div} \boldsymbol{\xi} = 0 & \text{in } \Omega, \\ \boldsymbol{\xi} \cdot \boldsymbol{n} = 0, \ \mathbf{curl} \boldsymbol{\xi} \times \boldsymbol{n} = \boldsymbol{\psi}_0 \times \boldsymbol{n} & \text{on } \Gamma, \\ \langle \boldsymbol{\xi} \cdot \boldsymbol{n}, 1 \rangle_{\Sigma_j} = 0, \end{cases}$$

where $\boldsymbol{\psi}_0$ is the tangential vector potential.

• The required vector potential is given by:

$$oldsymbol{\psi} = oldsymbol{\psi}_0 - \mathbf{curl}\,oldsymbol{\xi} - \sum_{i=1}^I \langle (oldsymbol{\psi}_0 - \mathbf{curl}\,oldsymbol{\xi}) \cdot oldsymbol{n}, \, 1
angle_{\Gamma_i} \, \mathbf{grad}\, q_i^N$$

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We define now the space

 $\boldsymbol{V}_{N}^{p}(\Omega) = \{\boldsymbol{w} \in \boldsymbol{X}_{N}^{p}(\Omega); \text{ div } \boldsymbol{w} = 0 \text{ in } \Omega \text{ and } \langle \boldsymbol{w} \cdot \boldsymbol{n}, 1 \rangle_{\Gamma_{i}} = 0, \ 1 \leq i \leq I\},\$

which is a Banach space for the norm $\|\cdot\|_{\boldsymbol{X}^p(\Omega)}$.

Using a Helmholtz decomposition and the previous normal vector potential, we can deduce:

Lemma (Inf Sup Condition)

The following Inf-Sup Condition holds: there exists a constant $\beta > 0$, such that

$$\inf_{\substack{\boldsymbol{\varphi} \in \boldsymbol{V}_{N}^{p'}(\Omega) \\ \boldsymbol{\varphi} \neq \boldsymbol{0}}} \sup_{\substack{\boldsymbol{\xi} \in \boldsymbol{V}_{N}^{p}(\Omega) \\ \boldsymbol{\xi} \neq \boldsymbol{0}}} \frac{\int_{\Omega} \operatorname{curl} \boldsymbol{\xi} \cdot \operatorname{curl} \boldsymbol{\varphi} \, \mathrm{d} \boldsymbol{x}}{\|\boldsymbol{\xi}\|_{\boldsymbol{X}_{N}^{p}(\Omega)} \|\boldsymbol{\varphi}\|_{\boldsymbol{X}_{N}^{p'}(\Omega)}} \ge \beta. \tag{23}$$

IV. Stokes stationary problem with normal boundary conditions

$$(\mathcal{S}_T) \quad \begin{cases} -\Delta \, \boldsymbol{u} + \nabla \, \boldsymbol{\pi} = \boldsymbol{f} & \text{in } \Omega, \\ \text{div} \, \boldsymbol{u} = 0 & \text{in } \Omega, \\ \boldsymbol{u} \cdot \boldsymbol{n} = \boldsymbol{g}, \, \, \operatorname{\mathbf{curl}} \boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{h} \times \boldsymbol{n} & \text{on } \Gamma, \\ \langle \boldsymbol{u} \cdot \boldsymbol{n}, \, 1 \rangle_{\Sigma_j} = 0, & 1 \leq j \leq J. \end{cases}$$

• We solve the pressure π :

 $\Delta \pi = \operatorname{div} f \text{ in } \Omega, \quad \frac{\partial \pi}{\partial n} = f \cdot n + \operatorname{div}_{\Gamma}(h \times n) \text{ on } \Gamma$

• We set $F = f - \nabla \pi$. Then, u is a solution of the following elliptic problem:

$$(E_T) \quad \begin{cases} -\Delta \, \boldsymbol{u} = \boldsymbol{F} & \text{in } \Omega, \\ \operatorname{div} \boldsymbol{u} = \boldsymbol{0} & \operatorname{in} \Omega, \\ \boldsymbol{u} \cdot \boldsymbol{n} = \boldsymbol{g}, \, \operatorname{\mathbf{curl}} \boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{h} \times \boldsymbol{n} & \text{on } \Gamma, \\ \langle \boldsymbol{u} \cdot \boldsymbol{n}, 1 \rangle_{\Sigma_j} = \boldsymbol{0}, & 1 \leq j \leq J \end{cases}$$

• To solve the Stokes problem (\mathcal{S}_T) , we are reduced to solve the problem (E_T) .

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IV. Stokes stationary problem with normal boundary conditions

$$(\mathcal{S}_T) \quad \begin{cases} -\Delta \, \boldsymbol{u} + \nabla \, \boldsymbol{\pi} = \boldsymbol{f} & \text{in } \Omega, \\ \operatorname{div} \, \boldsymbol{u} = 0 & \text{in } \Omega, \\ \boldsymbol{u} \cdot \boldsymbol{n} = \boldsymbol{g}, \, \operatorname{\mathbf{curl}} \boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{h} \times \boldsymbol{n} & \text{on } \Gamma, \\ \langle \boldsymbol{u} \cdot \boldsymbol{n}, \, 1 \rangle_{\Sigma_j} = 0, & 1 \leq j \leq J. \end{cases}$$

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• To solve the Stokes problem (\mathcal{S}_T) , we are reduced to solve the problem (\mathcal{E}_T) .

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IV. Stokes stationary problem with normal boundary conditions

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IV. Stokes stationary problem with normal boundary conditions

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• To solve the Stokes problem (S_T) , we are reduced to solve the problem (E_T) .

• Contrarly to the Stokes problem with Dirichlet boundary condition, it appears that when

div $\boldsymbol{f} = 0$ in Ω and $\boldsymbol{f} \cdot \boldsymbol{n} = \operatorname{div}_{\Gamma}(\boldsymbol{h} \times \boldsymbol{n}) = 0$ on Γ ,

the pressure π can be constant.

• In the case of the Stokes problem with Dirichlet, even if

div
$$\boldsymbol{f} = 0$$
 in Ω and $\boldsymbol{f} \cdot \boldsymbol{n} = 0$ on Γ ,

we have $\Delta \boldsymbol{u} \cdot \boldsymbol{n} \neq 0$ on Γ and then π is non constant.

Study of the elliptic problem :

$$\int -\Delta \boldsymbol{u} = \boldsymbol{f} \qquad \text{in } \Omega,$$

$$\int \operatorname{div} \boldsymbol{u} = 0 \qquad \qquad \text{in } \Omega,$$

$$(E_T)$$
 $\begin{cases} \mathbf{u} \cdot \mathbf{u} = g, \quad \mathbf{curl} \, \mathbf{u} \times \mathbf{n} = \mathbf{h} \times \mathbf{n} \\ \langle \mathbf{u} \cdot \mathbf{n}, 1 \rangle_{\Sigma_j} = 0, \quad 1 \le j \le J. \end{cases}$

Proposition (Weak and Strong solutions of (E_T))

i) Let $\boldsymbol{f} \in \boldsymbol{L}^{p}(\Omega)$ with div $\boldsymbol{f} = 0$ in Ω , $g \in W^{1-\frac{1}{p},p}(\Gamma)$ and $\boldsymbol{h} \in \boldsymbol{W}^{-\frac{1}{p},p}(\Gamma)$ satisfying the following compatibility condition:

$$egin{aligned} &orall oldsymbol{v} \in oldsymbol{K}_T^{\ p'}(\Omega), \ \int_\Omega oldsymbol{f} \cdot oldsymbol{v} \,\mathrm{d}oldsymbol{x} + \langleoldsymbol{h} imes oldsymbol{n}, oldsymbol{v} \ oldsymbol{W}^{-rac{1}{p},p}(\Gamma) imes oldsymbol{W}^{rac{1}{p},p'}(\Gamma) &= 0, \ oldsymbol{f} \cdot oldsymbol{n} + \operatorname{div}_\Gamma \left(oldsymbol{h} imes oldsymbol{n}\right) &= 0 \ \operatorname{on} \Gamma, \ \int_\Gamma g \,\mathrm{d}oldsymbol{\sigma} &= 0. \end{aligned}$$

Then, the problem (E_T) has a unique solution $u \in W^{1,p}(\Omega)$ satisfying the estimate:

$$\| \boldsymbol{u} \|_{\boldsymbol{W}^{1,p}(\Omega)} \leq C \Big(\| \boldsymbol{f} \|_{L^{p}(\Omega)} + \| \boldsymbol{g} \|_{W^{1-1/p,p}(\Gamma)} + \| \boldsymbol{h} \times \boldsymbol{n} \|_{\boldsymbol{W}^{-1/p,p}(\Gamma)} \Big).$$

ii) If $g \in W^{2-1/p,p}(\Gamma)$ and $h \in W^{1-1/p,p}(\Gamma)$, then the solution u belongs to $W^{2,p}(\Omega)$.

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Theorem (Weak and strong solution of (\mathcal{S}_T))

i) Let f, g, h with:

$$\boldsymbol{f} \in (\boldsymbol{H}_0^{p'}(\operatorname{div}, \Omega))', \ \boldsymbol{g} \in W^{1-\frac{1}{p}, p}(\Gamma), \ \boldsymbol{h} \in \boldsymbol{W}^{-\frac{1}{p}, p}(\Gamma),$$

verifying the compatibility conditions: for any \boldsymbol{v} in $\boldsymbol{K}_{T}^{p'}(\Omega)$:

$$\begin{split} \langle \boldsymbol{f}, \, \boldsymbol{v} \rangle_{(\boldsymbol{H}_0^{\, p'}(\operatorname{div}, \Omega))' \times \boldsymbol{H}_0^{\, p'}(\operatorname{div}, \Omega)} &+ \langle \boldsymbol{h} \times \boldsymbol{n}, \, \boldsymbol{v} \rangle_{\boldsymbol{W}^{-\frac{1}{p}, p}(\Gamma) \times \boldsymbol{W}^{\frac{1}{p}, p'}(\Gamma)} = 0. \\ & \int_{\Gamma} g \, \mathrm{d}\boldsymbol{\sigma} &= 0. \end{split}$$

Then, the Stokes problem (\mathcal{S}_T) has a unique solution $(\boldsymbol{u}, \pi) \in \boldsymbol{W}^{1,p}(\Omega) \times L^p(\Omega)/\mathbb{R}$ satisfying the estimate:

$$\|\boldsymbol{u}\|_{\boldsymbol{W}^{1,p}(\Omega)}+\|\boldsymbol{\pi}\|_{L^{p}(\Omega)}\leq \Big(\|\boldsymbol{f}\|_{(\boldsymbol{H}_{0}^{p'}(\operatorname{div},\Omega))'}+\|\boldsymbol{g}\|_{W^{1-\frac{1}{p},p}(\Gamma)}+\|\boldsymbol{h}\times\boldsymbol{n}\|_{\boldsymbol{W}^{-\frac{1}{p},p}(\Gamma)}\Big).$$

ii) If $f \in L^{p}(\Omega)$, $g \in W^{2-1/p,p}(\Gamma)$, $h \in W^{1-1/p,p}(\Gamma)$ then the solution (u, π) of the problem (S_{T}) given previously belongs to $W^{2,p}(\Omega) \times W^{1,p}(\Omega)$.

Remark: We can consider the case of the problem (\mathcal{S}_T) with div $\boldsymbol{u} = \chi$ in Ω , where $\chi \in \boldsymbol{L}^p(\Omega)$ satisfying $\int_{\Gamma} g \, \mathrm{d}\boldsymbol{\sigma} = \langle \chi, 1 \rangle_{\Gamma}$.

Very weak solution for (\mathcal{S}_T)

Let f, χ, g , and h with

$$f \in (T^{p'}(\Omega))', \ \chi \in L^{p}(\Omega), \ g \in W^{-1/p,p}(\Gamma), \ h \in W^{-1-1/p,p}(\Gamma),$$

with $T^{p'}(\Omega) = \left\{ \varphi \in H_0^{p'}(\operatorname{div}, \Omega); \operatorname{div} \varphi \in W_0^{1,p'}(\Omega) \right\}$ and satisfying the compatibility conditions:

$$\begin{split} \forall \, \boldsymbol{\varphi} \in \boldsymbol{K}_{T}^{\, p'}(\Omega), \quad \langle \boldsymbol{f}, \, \boldsymbol{\varphi} \rangle_{(T^{\, p'}(\Omega))' \times T^{\, p'}(\Omega)} + \langle \boldsymbol{h} \times \boldsymbol{n}, \, \boldsymbol{\varphi} \rangle_{\Gamma} = 0. \\ \int_{\Omega} \chi \, \mathrm{d} \, \boldsymbol{x} = \langle g, \, 1 \rangle_{\Gamma}. \end{split}$$

Then, the Stokes problem (\mathcal{S}_T) has exactly one solution $u \in L^p(\Omega)$ and $\pi \in W^{-1,p}(\Omega)/\mathbb{R}$. Moreover, there exists a constant C > 0 depending only on p and Ω such that:

$$\begin{split} \|\boldsymbol{u}\|_{L^{p}(\Omega)} + \|\pi\|_{W^{-1,p}(\Omega)/\mathbb{R}} &\leq C \Big(\|\boldsymbol{f}\|_{(T^{p'}(\Omega))'} + \|\chi\|_{L^{p}(\Omega)} + \|g\|_{W^{-1/p,p}(\Gamma)} + \\ &+ \|\boldsymbol{h} \times \boldsymbol{n}\|_{W^{-1-1/p,p}(\Gamma)} \Big). \end{split}$$

Helmholtz Decomposition for vector fields in $L^p(\Omega)$

For any vector field $\boldsymbol{v} \in \boldsymbol{L}^p(\Omega)$, we have the first following decomposition:

 $\boldsymbol{v} = \boldsymbol{z} + \nabla \, \boldsymbol{\chi} + \operatorname{\mathbf{curl}} \boldsymbol{u},$

- $\boldsymbol{z} \in \boldsymbol{K}_N^p(\Omega)$ is unique,
- $\chi \in W_0^{1,p}(\Omega)$ is unique,
- $\boldsymbol{u} \in \boldsymbol{W}^{1,p}(\Omega)$ is the unique solution, up to an additive element of the kernel $\boldsymbol{K}_T^p(\Omega)$, of the problem :

$$\begin{cases} -\Delta \boldsymbol{u} = \operatorname{\mathbf{curl}} \boldsymbol{v} \quad \text{and} \quad \operatorname{div} \boldsymbol{u} = 0 \quad \text{in } \Omega, \\ \boldsymbol{u} \cdot \boldsymbol{n} = 0, \quad (\operatorname{\mathbf{curl}} \boldsymbol{u} - \boldsymbol{v}) \times \boldsymbol{n} = 0 \quad \text{on } \Gamma. \end{cases}$$

Helmholtz Decomposition for vector fields in $L^p(\Omega)$

For any vector field $\boldsymbol{v} \in \boldsymbol{L}^{p}(\Omega)$, we have the second following decomposition:

$$\boldsymbol{v} = \boldsymbol{z} + \nabla \, \boldsymbol{\chi} + \operatorname{\mathbf{curl}} \boldsymbol{u},$$

- $\boldsymbol{z} \in \boldsymbol{K}_T^p(\Omega)$ is unique,
- $\chi \in W^{1,p}(\Omega)$ is unique up an additive constant,
- $\boldsymbol{u} \in \boldsymbol{W}^{1,p}(\Omega)$ is the unique solution, up to an additive element of the kernel $\boldsymbol{K}_{N}^{p}(\Omega)$, of the problem :

$$\begin{cases} -\Delta \boldsymbol{u} = \operatorname{\mathbf{curl}} \boldsymbol{v} & \text{and} & \operatorname{div} \boldsymbol{u} = 0 & \text{in } \Omega, \\ \boldsymbol{u} \times \boldsymbol{n} = 0, & \text{on } \Gamma. \end{cases}$$

V. Stokes Equations with Pressure Boundary Condition

$$(\mathcal{S}_N) \quad \begin{cases} -\Delta \, \boldsymbol{u} + \nabla \, \boldsymbol{\pi} = \boldsymbol{f} & \text{in } \Omega, \\ \text{div} \, \boldsymbol{u} = 0 & \text{in } \Omega, \\ \boldsymbol{u} \times \, \boldsymbol{n} = \boldsymbol{g} \times \, \boldsymbol{n}, \quad \boldsymbol{\pi} = \pi_0 & \text{on } \Gamma, \\ \langle \boldsymbol{u} \cdot \, \boldsymbol{n}, \, 1 \rangle_{\Gamma_i} = 0, \quad 1 \leq i \leq I. \end{cases}$$

• The pressure can be found independently of the velocity as a solution of the Dirichlet problem:

$$\Delta \pi = \operatorname{div} \boldsymbol{f}$$
 in Ω , $\pi = \pi_0$ on Γ

• The velocity is a solution of the following system:

$$\int -\Delta u = F$$
 in Ω ,

$$(E_N) \begin{cases} \operatorname{div} \boldsymbol{u} = 0 & \text{in } \Omega, \end{cases}$$

$$\begin{pmatrix} \boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{g} \times \boldsymbol{n} & \text{on } \Gamma, \\ \langle \boldsymbol{u} \cdot \boldsymbol{n}, 1 \rangle_{\Gamma_i} = 0, \text{ for all } 1 \leq i \leq I, \end{cases}$$

where $\boldsymbol{F} = \boldsymbol{f} - \nabla \pi$

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Remarks:

- The condition div $\mathbf{F} = 0$ in Ω is necessary to solve (E_N) .
- The condition div *u* = 0 in Ω ⇐⇒ div *u* = 0 on Γ on the one hand. On the other hand, since

$$\operatorname{div} \boldsymbol{u} = \operatorname{div}_{\Gamma} \boldsymbol{u}_{\boldsymbol{\tau}} + K \boldsymbol{u} \cdot \boldsymbol{n} + \frac{\partial \boldsymbol{u}}{\partial \boldsymbol{n}} \cdot \boldsymbol{n} \quad \operatorname{sur} \, \Gamma,$$

where K denotes the mean curvature of Γ , the condition div $\boldsymbol{u} = 0$ on Γ is itself equivalent, if $\boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{0}$ on Γ , to the Fourier-Robin condition:

$$K\boldsymbol{u}\cdot\boldsymbol{n}+\frac{\partial \boldsymbol{u}}{\partial \boldsymbol{n}}\cdot\boldsymbol{n}=0 \text{ on } \Gamma.$$

That means that the problem (E_N) is equivalent to the following:

$$\begin{cases} -\Delta \boldsymbol{u} = \boldsymbol{F} & \text{in } \Omega, \\ \boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{0} & \text{on } \Gamma \\ K\boldsymbol{u} \cdot \boldsymbol{n} + \frac{\partial \boldsymbol{u}}{\partial \boldsymbol{n}} \cdot \boldsymbol{n} = 0 & \text{on } \Gamma. \end{cases}$$

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Using the Inf-Sup condition, we deduce the following theorem

Theorem (Weak and Strong solutions for (\mathcal{S}_N))

i) Let f, g, π_0 with

$$\boldsymbol{f} \in (\boldsymbol{H}_0^{p'}(\operatorname{\mathbf{curl}}, \Omega))', \ \boldsymbol{g} \in \boldsymbol{W}^{1-1/p,p}(\Gamma), \ \pi_0 \in W^{1-1/p,p}(\Gamma),$$

satisfying the compatibility condition:

$$\forall \boldsymbol{v} \in \boldsymbol{K}_{N}^{p'}(\Omega), \qquad \langle \boldsymbol{f}, \, \boldsymbol{v} \, \rangle_{\Omega} - \int_{\Gamma} \pi_{0} \, \boldsymbol{v} \cdot \boldsymbol{n} \, \mathrm{d}\boldsymbol{\sigma} = 0, \tag{24}$$

with $\langle \cdot, \cdot \rangle_{\Omega} = \langle \cdot, \cdot \rangle_{[H_0^{p'}(\operatorname{curl},\Omega)]' \times H_0^{p'}(\operatorname{curl},\Omega)]}$.

Then, the Stokes problem (\mathcal{S}_N) has a unique solution $(u, \pi) \in W^{1,p}(\Omega) \times W^{1,p}(\Omega)$ and satisfies the estimate: $\| u \|_{W^{1,p}(\Omega)} + \| \pi \|_{W^{1,p}(\Omega)} \leq C (\| f \|_{(H_0^{p'}(\operatorname{curl}, \Omega))'} + \| g \times n \|_{W^{1-1/p,p}(\Gamma)} + \| \pi \|_{W^{1-1/p,p}(\Gamma)}).$

ii) If $f \in L^p(\Omega)$ and $g \in W^{2-1/p,p}(\Gamma)$, then $u \in W^{2,p}(\Omega)$.

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Question:

What happens if the previous compatibility condition is not satisfied?

Variant of the system (\mathcal{S}_N) :

Find $(\boldsymbol{u}, \pi, \boldsymbol{c})$ such that:

$$(\mathcal{S}'_N) \quad \begin{cases} -\Delta \, \boldsymbol{u} + \nabla \, \boldsymbol{\pi} = \boldsymbol{f} \text{ and div } \boldsymbol{u} = 0 & \text{in } \Omega, \\ \boldsymbol{u} \times \, \boldsymbol{n} = \boldsymbol{g} \times \, \boldsymbol{n} & \text{on } \Gamma, \\ \boldsymbol{\pi} = \pi_0 \text{ on } \Gamma_0 \text{ and } \boldsymbol{\pi} = \pi_0 + \boldsymbol{c_i} & \text{on } \Gamma_i, \ 1 \le i \le I \\ \langle \boldsymbol{u} \cdot \boldsymbol{n}, \ 1 \rangle_{\Gamma_i} = 0, \ 1 \le i \le I, \end{cases}$$

where $\boldsymbol{c} = (c_i)_{1 \leq i \leq I}$.

Theorem (Weak and Strong solutions for (\mathcal{S}'_N))

Let f, g and π_0 such that:

$$\boldsymbol{f} \in [\boldsymbol{H}_0^{p'}(\operatorname{\mathbf{curl}}, \Omega)]', \quad \boldsymbol{g} \in \boldsymbol{W}^{1-1/p, p}(\Gamma), \quad \pi_0 \in W^{1-1/p, p}(\Gamma).$$

Then, the problem (S'_N) has a unique solution $u \in W^{1,p}(\Omega)$, $\pi \in W^{1,p}(\Omega)$ and constants c_1, \ldots, c_I satisfying the estimate:

$$\|u\|_{W^{1,p}(\Omega)} + \|\pi\|_{W^{1,p}(\Omega)} \le C \big(\|f\|_{[H_0^{p'}(\operatorname{curl},\Omega)]'} + \|g\|_{W^{1-1/p,p}} + \|\pi_0\|_{W^{1-1/p,p}} \big),$$

and where c_1, \ldots, c_I are given by

$$c_i = \langle \boldsymbol{f}, \, \nabla \, \boldsymbol{q}_i^N \rangle_{\Omega} - \langle \pi_0, \, \nabla \, \boldsymbol{q}_i^N \cdot \boldsymbol{n} \rangle_{\Gamma}.$$
(25)

In particular, if $f \in L^p(\Omega)$ and $g \in W^{2-1/p,p}(\Gamma)$, then $u \in W^{2,p}(\Omega)$.

Remark :

• Observe that if we suppose that the compatibility condition (24) is verified, we have that: $c_i = 0$ for all i = 1, ..., I. Then, we have reduced to solve the problem (S'_N) without the constant c_i and (S'_N) is anything other then (S_N) .

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The assumption on f in the previous theorem can be weakened by considering the space defined for all $1 < r, p < \infty$:

 $\boldsymbol{H}_{0}^{r,p}(\boldsymbol{\mathrm{curl}},\,\Omega) = \{ \boldsymbol{\varphi} \in \boldsymbol{L}^{r}(\Omega); \ \boldsymbol{\mathrm{curl}}\, \boldsymbol{\varphi} \in \boldsymbol{L}^{p}(\Omega), \ \boldsymbol{\varphi} \times \boldsymbol{n} = \boldsymbol{0} \text{ on } \Gamma \}.$

which is a Banach space for the norm

$$\|\varphi\|_{\boldsymbol{H}_{0}^{r,p}(\operatorname{\mathbf{curl}},\Omega)}=\|\varphi\|_{\boldsymbol{L}^{r}(\Omega)}+\|\operatorname{\mathbf{curl}}\varphi\|_{\boldsymbol{L}^{p}(\Omega)}.$$

We can prove that

$$\mathcal{D}(\Omega)$$
 is dense in $H_0^{r', p'}(\mathbf{curl}, \Omega)$

and its dual space can be characterized as:

$$[{oldsymbol H}_0^{r',\,p'}({f curl},\,\Omega)]'=\{{oldsymbol F}+{f curl}\,{oldsymbol \psi};\;\;{oldsymbol F}\in{oldsymbol L}^r(\Omega),\;{oldsymbol \psi}\in{oldsymbol L}^p(\Omega)\}.$$

Theorem (Second Version for Weak solutions for (\mathcal{S}'_N))

Let f, g and π_0 such that

$$\boldsymbol{f} \in [\boldsymbol{H}_{0}^{r',p'}(\mathbf{curl},\,\Omega)]', \quad \boldsymbol{g} \times \boldsymbol{n} \in \boldsymbol{W}^{1-1/p,p}(\Gamma), \quad \pi_{0} \in W^{1-1/r,r}(\Gamma),$$

with

$$r \leq p \quad ext{and} \quad rac{1}{r} \leq rac{1}{p} + rac{1}{3}$$

Then, the problem (\mathcal{S}'_N) has a unique solution

$$\boldsymbol{u} \in \boldsymbol{W}^{1,p}(\Omega), \quad \pi \in W^{1,r}(\Omega)$$

and constants c_1, \ldots, c_I satisfying the estimate:

$$\begin{aligned} \|\boldsymbol{u}\|_{\boldsymbol{W}^{1,p}(\Omega)} &+ \|\pi\|_{W^{1,r}(\Omega)} \leq C \big(\|\boldsymbol{f}\|_{[\boldsymbol{H}_{0}^{r',p'}(\mathbf{curl},\Omega)]'} \\ &+ \|\boldsymbol{g} \times \boldsymbol{n}\|_{\boldsymbol{W}^{1-1/p,p}(\Gamma)} + \|\pi_{0}\|_{W^{1-1/r,r}(\Gamma)} \big), \end{aligned}$$

and c_1, \ldots, c_I are given by (25), where we replace the duality brackets on Ω by

$$\langle \cdot, \cdot \rangle_{\Omega} = \langle \cdot, \cdot \rangle_{[H_0^{r',p'}(\operatorname{\mathbf{curl}},\Omega)]' \times H_0^{r',p'}(\operatorname{\mathbf{curl}},\Omega)}.$$

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Theorem (Very weak solutions for (\mathcal{S}_N))

Let $\boldsymbol{f}, \, \boldsymbol{g}, \, \text{and} \, \pi_0$ with

$$oldsymbol{f} \in [oldsymbol{H}_0^{p'}(\mathbf{curl},\,\Omega)]', \ oldsymbol{g} \in oldsymbol{W}^{-1/p,p}(\Gamma), \ \pi_0 \in W^{-1/p,p}(\Gamma),$$

and satisfying the compatibility conditions (24). Then, the Stokes problem (\mathcal{S}_N) has exactly one solution

$$\boldsymbol{u} \in \boldsymbol{L}^p(\Omega) \quad \text{and} \quad \pi \in L^p(\Omega)/\mathbb{R}.$$

Moreover, there exists a constant C > 0 depending only on p and Ω such that:

$$\begin{split} \|\boldsymbol{u}\|_{\boldsymbol{L}^{p}(\Omega)} + \|\pi\|_{L^{p}(\Omega)/\mathbb{R}} &\leq C \Big(\|\boldsymbol{f}\|_{[\boldsymbol{H}_{0}^{p'}(\mathbf{curl},\Omega)]'} + \|\boldsymbol{g}\|_{\boldsymbol{W}^{-1/p,p}(\Gamma)} + \\ &+ \|\pi_{0}\|_{\boldsymbol{W}^{-1/p,p}(\Gamma)} \Big). \end{split}$$

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Remark. To study the case of Navier boundary conditions:

$$\boldsymbol{u} \cdot \boldsymbol{n} = 0$$
 and $[\mathbf{D}(\boldsymbol{u})\boldsymbol{n}]_{\boldsymbol{\tau}} = \boldsymbol{h},$

it suffices to observe that

$$[2\mathbf{D}(\boldsymbol{v})\boldsymbol{n}]_{\boldsymbol{\tau}} = -\mathbf{curl}\,\boldsymbol{v}\times\boldsymbol{n} - 2\mathbf{\Lambda}\boldsymbol{v} \quad \text{on}\,\Gamma,$$

where Λ is an operator of order 0 defined by

$$oldsymbol{\Lambda}oldsymbol{w} = \sum_{k=1}^2ig(oldsymbol{w}_{oldsymbol{ au}}\cdotrac{\partialoldsymbol{n}}{\partial s_k}ig)oldsymbol{ au}_k,$$

and to use the results obtained with the Navier type boundary conditions.

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VI. Oseen and Navier-Stokes Problem with Pressure Boundary Condition

We are interested to study the following problem: Find u, q and $\alpha \in \mathbb{R}^{I}$ satisfying:

$$(\mathcal{NS}) \qquad \begin{cases} -\Delta \, \boldsymbol{u} + \boldsymbol{u} \cdot \nabla \, \boldsymbol{u} + \nabla \, \boldsymbol{q} = \boldsymbol{f} & \text{and} & \text{div} \, \boldsymbol{u} = \chi & \text{in } \Omega, \\ \boldsymbol{u} \times \, \boldsymbol{n} = \boldsymbol{g} & \text{on } \Gamma, \\ \boldsymbol{q} = q_0 & \text{on } \Gamma_0 & \text{and} \, \boldsymbol{q} = q_0 + \alpha_i & \text{on } \Gamma_i, \ i = 1, \dots, I, \\ \int_{\Gamma_i} \boldsymbol{u} \cdot \boldsymbol{n} \, \mathrm{d}\boldsymbol{\sigma} = 0, \ i = 1, \dots, I, \end{cases}$$

- Note that α is a supplementary unknown Stokes which depends in fact on u
- If we take $\chi = 0$ and g = 0, unlike the Navier-Stokes problem with Dirichlet boundary conditions de Dirichlet, the property: $\int_{\Omega} (\boldsymbol{u} \cdot \nabla \boldsymbol{u}) \cdot \boldsymbol{u} \, d\boldsymbol{x} = 0$ does not hold.
- But, we have

$$oldsymbol{u} \cdot
abla oldsymbol{u} = \mathbf{curl}\,oldsymbol{u} imes oldsymbol{u} + rac{1}{2}\,
abla |oldsymbol{u}|^2$$

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We rewrite then (\mathcal{NS}) under the following form:

$$(\mathcal{NS}_N) \qquad \begin{cases} -\Delta \, \boldsymbol{u} + \operatorname{curl} \boldsymbol{u} \times \boldsymbol{u} + \nabla \, \boldsymbol{\pi} = \boldsymbol{f} & \text{in } \Omega, \\ \operatorname{div} \boldsymbol{u} = \chi & \operatorname{in} \Omega, \\ \boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{g} & \text{on } \Gamma, \\ \boldsymbol{\pi} = \pi_0 \, \operatorname{sur} \, \Gamma_0 \, \operatorname{et} \, \boldsymbol{\pi} = \pi_0 + \alpha_i & \text{on } \Gamma_i, \, i = 1, \dots, I, \\ \int_{\Gamma_i} \boldsymbol{u} \cdot \boldsymbol{n} \, \mathrm{d} \boldsymbol{\sigma} = 0, \, i = 1, \dots, I, \end{cases}$$

Remarks.

- We can search directely weak solutions $\boldsymbol{u} \in \boldsymbol{H}^1(\Omega)$ and $\pi \in L^2(\Omega)$ of the system (\mathcal{NS}_N) by using a fixed point method.
- We can then obtain solutions $\boldsymbol{u} \in \boldsymbol{W}^{1,p}(\Omega)$ for p > 2 thanks to the Stokes problem theory.
- The case p < 2 to study the (NS_N) system is more complicated.
- For this reason, we will study the Oseen problem (OS_N) .

Remarks.

- We can search directely weak solutions $\boldsymbol{u} \in \boldsymbol{H}^1(\Omega)$ and $\pi \in L^2(\Omega)$ of the system (\mathcal{NS}_N) by using a fixed point method.
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- The case p < 2 to study the (NS_N) system is more complicated.
- For this reason, we will study the Oseen problem (\mathcal{OS}_N) .

$$(\mathcal{OS}_N) \qquad \begin{cases} -\Delta \, \boldsymbol{u} + \operatorname{curl} \boldsymbol{a} \times \boldsymbol{u} + \nabla \, \boldsymbol{\pi} = \boldsymbol{f} & \text{in } \Omega, \\ \operatorname{div} \, \boldsymbol{u} = 0 & \text{in } \Omega, \\ \boldsymbol{u} \times \, \boldsymbol{n} = \boldsymbol{0} & \text{on } \Gamma, \\ \boldsymbol{\pi} = \pi_0 + c_i & \text{on } \Gamma_i, \ \boldsymbol{0} = 1, \dots, I, \\ \int_{\Gamma_i} \boldsymbol{u} \cdot \boldsymbol{n} \, \mathrm{d}\boldsymbol{\sigma} = 0, \ i = 1, \dots, I, \end{cases}$$

where we have take $\chi = 0$ and g = 0. We suppose also that

 $\operatorname{\mathbf{curl}} \boldsymbol{a} \in \boldsymbol{L}^{3/2}(\Omega)$

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We introduce the following Hilbert space:

$$oldsymbol{V}_N = \left\{ oldsymbol{v} \in H^1(\Omega); ext{ div } oldsymbol{v} = 0 ext{ in } \Omega, \ oldsymbol{v} imes oldsymbol{n} = oldsymbol{0} ext{ on } \Gamma$$

and $\int_{\Gamma_i} oldsymbol{v} \cdot oldsymbol{n} = 0, \ 1 \leq i \leq I
ight\}$

and recall that

$$oldsymbol{v}\mapsto \Big(\int_{\Omega}|{f curl}\,oldsymbol{v}|^2\Big)^{1/2}$$

is a norm on V_N equivalent to the full norm of $H^1(\Omega)$.

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Before establishing the result of existence of a weak solution for the problem (\mathcal{OS}_N) , we will see in what functional space it is reasonable to take π_0 and to find the pressure π appearing in (\mathcal{OS}_N) , knowing that we are first interesting to velocity fields in

 $\boldsymbol{u} \in \boldsymbol{H}^{1}(\Omega)$ with $\boldsymbol{f} \in \boldsymbol{L}^{6/5}(\Omega)$.

With a such vector \boldsymbol{u} , we have

$$\operatorname{curl} \boldsymbol{a} \times \boldsymbol{u} \in \boldsymbol{L}^{6/5}(\Omega) \hookrightarrow \boldsymbol{H}^{-1}(\Omega).$$

Since $\Delta u \in H^{-1}(\Omega)$, we deduce from the first equation in (\mathcal{OS}_N) that $\nabla \pi \in H^{-1}(\Omega)$. Then the pressure π belongs to $L^2(\Omega)$. Furthermore,

$$-\Delta \pi = \operatorname{div} \boldsymbol{f} - \operatorname{div} \left(\operatorname{\mathbf{curl}} \boldsymbol{a} \times \boldsymbol{u}\right) \quad \text{in } \Omega,$$

so that $\Delta \pi \in W^{-1,6/5}(\Omega)$ and the trace of π on Γ belongs to $H^{-1/2}(\Gamma)$ so that we must assume that

 $\pi_0 \in H^{-1/2}(\Gamma).$

Theorem

Let

 $f \in L^{6/5}(\Omega), \quad \pi_0 \in H^{-1/2}(\Gamma) \text{ and } a \in \mathcal{D}'(\Omega) \text{ such that } \operatorname{\mathbf{curl}} a \in L^{3/2}(\Omega).$ Then, the problem:

Find $(\boldsymbol{u}, \pi, \boldsymbol{c}) \in \boldsymbol{V}_N \times L^2(\Omega) \times \mathbb{R}^{I+1}$ satisfying (\mathcal{OS}_N) with $\langle \pi, 1 \rangle_{\Gamma} = 0$ (26) is equivalent to the problem: Find $\boldsymbol{u} \in \boldsymbol{V}_N$ such that

$$\forall v \in V_N, \quad \int_{\Omega} \operatorname{\mathbf{curl}} u \cdot \operatorname{\mathbf{curl}} v \, \mathrm{d} x + \int_{\Omega} (\operatorname{\mathbf{curl}} u \times u) \cdot v = \int_{\Omega} f \cdot v \, \mathrm{d} x - \langle \pi_0, v \cdot n \rangle_{\Gamma}$$
(27)

and find constants c_0, \ldots, c_I satisfying $\sum_{i=0}^{I} c_i \max \Gamma_i + \langle \pi_0, 1 \rangle_{\Gamma} = 0$ and such that for any $i = 1, \ldots, I$:

$$c_i - c_0 = \int_{\Omega} \boldsymbol{f} \cdot \nabla q_i^N \, \mathrm{d}\boldsymbol{x} - \int_{\Omega} (\operatorname{\mathbf{curl}} \boldsymbol{a} \times \boldsymbol{u}) \cdot \nabla q_i^N \, \mathrm{d}\boldsymbol{x} - \langle \pi_0, \, \nabla \, q_i^N \cdot \boldsymbol{n} \rangle_{\Gamma}.$$
(28)

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Using the Lax Milgram theorem and some regularity result of the Laplacian, we prove the following theorem.

Theorem

Let $\mathbf{f} \in \mathbf{L}^{6/5}(\Omega)$, $\operatorname{curl} \mathbf{a} \in \mathbf{L}^{3/2}(\Omega)$ and $\pi_0 \in H^{-1/2}(\Gamma)$, then the problem (56) has a unique solution $(\mathbf{u}, \pi, \mathbf{c}) \in \mathbf{H}^1(\Omega) \times L^2(\Omega) \times \mathbb{R}^{I+1}$ with $\langle \pi, 1 \rangle_{\Gamma} = 0$ and we have the following estimates:

$$\|\boldsymbol{u}\|_{\boldsymbol{H}^{1}(\Omega)} \leq C \big(\|\boldsymbol{f}\|_{\boldsymbol{L}^{6/5}(\Omega)} + \|\pi_{0}\|_{H^{-1/2}(\Gamma)} \big),$$
(29)

 $\|\pi\|_{L^{2}(\Omega)} \leq C \left(1 + \|\mathbf{curl}\,a\|_{L^{3/2}(\Omega)}\right) \left(\|f\|_{L^{6/5}(\Omega)} + \|\pi_{0}\|_{H^{-1/2}(\Gamma)}\right),$ (30)

where $\mathbf{c} = (c_0, \ldots, c_I)$. Moreover, if $\pi_0 \in W^{1/6, 6/5}(\Gamma)$ and Ω is $\mathcal{C}^{2,1}$, then $\mathbf{u} \in \mathbf{W}^{2, 6/5}(\Omega)$ and $\pi \in W^{1, 6/5}(\Omega)$.

Remark

Even if the pressure π change in $\pi - c_0$, the system (OS_N) is equivalent to the following type-Oseen problem:

 $\begin{cases} -\Delta \boldsymbol{u} + \operatorname{curl} \boldsymbol{a} \times \boldsymbol{u} + \nabla \pi = \boldsymbol{f} \quad \text{and} \quad \operatorname{div} \boldsymbol{u} = 0 & \text{in } \Omega, \\ \boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{0} & \text{on } \Gamma, \\ \pi = \pi_0 \quad \text{on } \Gamma_0, \quad \text{and} \quad \pi = \pi_0 + \alpha_i, \ i = 1, \dots, I, \quad \text{on } \Gamma_i, \\ \int_{\Gamma_i} \boldsymbol{u} \cdot \boldsymbol{n} \, \mathrm{d}\boldsymbol{\sigma} = 0, \ i = 1, \dots, I, \end{cases}$

where the unknowns constants satisfy for any i = 1, ..., I:

$$\alpha_i = \int_{\Omega} \boldsymbol{f} \cdot \nabla \, q_i^N \, \mathrm{d}\boldsymbol{x} - \int_{\Omega} (\operatorname{\mathbf{curl}} \boldsymbol{a} \times \boldsymbol{u}) \cdot \nabla \, q_i^N \, \mathrm{d}\boldsymbol{x} - \langle \pi_0, \, \nabla \, q_i^N \cdot \boldsymbol{n} \rangle_{\Gamma}.$$

But, it is clear that the new pressure π does not satisfy the condition $\langle \pi, 1 \rangle_{\Gamma} = 0$.

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Remark

If we suppose that

 $oldsymbol{f} \in [oldsymbol{H}_0^{6,2}(\mathbf{curl},\,\Omega)]', \quad \mathbf{curl} \ oldsymbol{a} \in oldsymbol{L}^{3/2}(\Omega) \quad ext{and} \quad \pi_0 \in H^{-1/2}(\Gamma),$

then the problems (26) and (27)-(28) are again equivalent, with the difference that we use here the duality brackets between $[\mathbf{H}_0^{6,2}(\mathbf{curl}, \Omega)]'$ and $\mathbf{H}_0^{6,2}(\mathbf{curl}, \Omega)$ in place of the integral on Ω in the right hand side of (27) and the density of $\mathcal{D}_{\sigma}(\overline{\Omega}) \times \mathcal{D}(\overline{\Omega})$ in the space

$$\boldsymbol{\mathcal{M}} = \left\{ (\boldsymbol{u}, \pi) \in \boldsymbol{H}_{\sigma}^{1}(\Omega) \times L^{2}(\Omega); \ -\Delta \, \boldsymbol{u} + \nabla \, \pi \in [\boldsymbol{H}_{0}^{6,2}(\boldsymbol{\mathrm{curl}}, \, \Omega)]' \right\}$$

It is easy now to extend the previous theorem to the case where

- $\boldsymbol{f} \in [\boldsymbol{H}_0^{6,2}(\operatorname{\mathbf{curl}}, \Omega)]',$
- the divergence operator does not vanish
- the boundary conditions are nonhomogeneous.

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Theorem

Let $f \in [H_0^{6,2}(\operatorname{curl}, \Omega)]'$, $\operatorname{curl} a \in L^{3/2}(\Omega)$, $\chi \in W^{1,6/5}(\Omega)$, $\pi_0 \in H^{-1/2}(\Gamma)$ and $g \in H^{1/2}(\Gamma)$. Then the problem

$$\begin{cases}
-\Delta \boldsymbol{u} + \operatorname{curl} \boldsymbol{a} \times \boldsymbol{u} + \nabla \pi = \boldsymbol{f} \quad \text{and} \quad \operatorname{div} \boldsymbol{u} = \chi \quad \text{in } \Omega, \\
\boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{g} \times \boldsymbol{n} \quad \text{on } \Gamma, \\
\pi = \pi_0 \quad \text{on } \Gamma_0, \quad \text{and} \quad \pi = \pi_0 + \alpha_i, \quad i = 1, \dots, I \quad \text{on } \Gamma_i, \\
\int_{\Gamma_i} \boldsymbol{u} \cdot \boldsymbol{n} \, \mathrm{d} \boldsymbol{\sigma} = 0, \quad i = 1, \dots, I,
\end{cases}$$
(31)

has a unique solution $(\boldsymbol{u}, \pi, \boldsymbol{\alpha}) \in \boldsymbol{H}^1(\Omega) \times L^2(\Omega) \times \mathbb{R}^I$ verifying the estimate:

$$\begin{split} \|\boldsymbol{u}\|_{\boldsymbol{H}^{1}(\Omega)} &\leq \quad C\Big(\|\boldsymbol{f}\|_{[\boldsymbol{H}_{0}^{6,2}(\mathbf{curl},\,\Omega)]'} + \|\pi_{0}\|_{\boldsymbol{H}^{-1/2}(\Gamma)} + \big(1 + \|\mathbf{curl}\,\boldsymbol{a}\|_{L^{3/2}(\Omega)}\big) \times \\ &\times \quad \big(\|\chi\|_{W^{1,6/5}(\Omega)} + \|\boldsymbol{g}\|_{\boldsymbol{H}^{1/2}(\Gamma)}\big)\Big), \end{split}$$

$$\begin{split} \|\pi\|_{L^{2}(\Omega)} &\leq C \big(1 + \|\mathbf{curl}\, \boldsymbol{a}\|_{L^{3/2}(\Omega)} \big) \Big(\|\boldsymbol{f}\|_{[\boldsymbol{H}_{0}^{6,2}(\mathbf{curl},\Omega)]'} + \|\pi_{0}\|_{H^{-1/2}(\Gamma)} + \\ &+ \big(1 + \|\mathbf{curl}\, \boldsymbol{a}\|_{L^{3/2}(\Omega)} \big) \times \big(\|\chi\|_{W^{1,6/5}(\Omega)} + \|\boldsymbol{g}\|_{H^{1/2}(\Gamma)} \big) \Big), \end{split}$$

where $\boldsymbol{\alpha} = (\alpha_1, \dots, \alpha_I)$. Moreover, if $\boldsymbol{f} \in \boldsymbol{L}^{6/5}(\Omega)$, $\pi_0 \in W^{1/6,6/5}(\Gamma)$, $\boldsymbol{g} \in W^{7/6,6/5}(\Gamma)$ and Ω is $\mathcal{C}^{2,1}$, then $\boldsymbol{u} \in W^{2,6/5}(\Omega)$ and $\pi \in W^{1,6/5}(\Omega)$. C. AMROUCHE Chérif AMROUCHE, University of PAU In the rest of this talk, we suppose that

 Ω is $\mathcal{C}^{2,1}$

and we are interested in the study of strong solutions $\boldsymbol{u} \in W^{2,p}(\Omega)$ for the system (\mathcal{OS}_N) when $p \geq 6/5$.

When $p < \frac{3}{2}$, because the embedding $W^{2,p}(\Omega) \hookrightarrow W^{1,p*}(\Omega)$, the term **curl** $\boldsymbol{a} \times \boldsymbol{u} \in L^p(\Omega)$ and we can use the regularity results on the Stokes problem. But this is not more the case when $p \geq \frac{3}{2}$ and that **curl** \boldsymbol{a} belongs only to $\boldsymbol{L}^{3/2}(\Omega)$.

We give in the following theorem the good conditions to ensure the existence of strong solutions.

Theorem

Let $p \geq 6/5$,

$$\boldsymbol{f} \in \boldsymbol{L}^p(\Omega), \ \pi_0 \in W^{1-1/p,p}(\Gamma), \ \operatorname{\mathbf{curl}} \boldsymbol{a} \in \boldsymbol{L}^s(\Omega)$$

with

$$s = \frac{3}{2}$$
 if $p < \frac{3}{2}$, $s = p$ if $p > \frac{3}{2}$, $s = \frac{3}{2} + \varepsilon$ if $p = \frac{3}{2}$, (32)

for $\varepsilon > 0$ arbitrary. Then the solution (\boldsymbol{u}, π) given by the previous theorem belongs to $\boldsymbol{W}^{2,p}(\Omega) \times W^{1,p}(\Omega)$ and satisfies the estimate:

 $\|\boldsymbol{u}\|_{\boldsymbol{W}^{2,p}(\Omega)} + \|\pi\|_{W^{1,p}(\Omega)} \le C \big(1 + \|\mathbf{curl}\,\boldsymbol{a}\|_{L^{s}(\Omega)}\big) \big(\|\boldsymbol{f}\|_{L^{p}(\Omega)} + \|\pi_{0}\|_{W^{1-1/p,p}(\Gamma)}\big).$

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In the following theorem, we prove the existence of Generalized Solutions with p > 2.

Theorem

Let

$$p>2, \quad f\in [\boldsymbol{H}_0^{r',p'}(\operatorname{\mathbf{curl}},\,\Omega)]', \quad \chi\in W^{1,r}(\Omega) \quad ext{and} \quad g\in \, \boldsymbol{W}^{1-1/p,p}(\Gamma).$$

We suppose that $\pi_0 \in W^{1-1/r,r}(\Gamma)$ and $\operatorname{curl} a \in L^s(\Omega)$ with $\frac{1}{r} = \frac{1}{p} + \frac{1}{3}$ and s satisfies:

$$s = \frac{3}{2}$$
 if $2 , $s = \frac{3}{2} + \varepsilon$ if $p = 3$ and $s = r$ if $p > 3$,$

for some arbitrary $\varepsilon > 0$. Then the problem (31) has a unique solution $(\boldsymbol{u}, \pi, \boldsymbol{\alpha}) \in \boldsymbol{W}^{1,p}(\Omega) \times \boldsymbol{W}^{1,r}(\Omega) \times \mathbb{R}^{I}$ satisfying the estimate

$$\|\boldsymbol{u}\|_{\boldsymbol{W}^{1,p}(\Omega)} + \|\pi\|_{W^{1,r}(\Omega)} \leq C \left(1 + \|\mathbf{curl}\,\boldsymbol{a}\|_{\boldsymbol{L}^{s}(\Omega)}\right)^{2} \left(\|\boldsymbol{f}\|_{\boldsymbol{H}_{0}^{r',p'}(\mathbf{curl},\Omega)]'} + \|\boldsymbol{g}\|_{\boldsymbol{W}^{1-1/p,p}(\Gamma)} + \|\pi_{0}\|_{W^{1-1/r,r}(\Gamma)} + \|\chi\|_{W^{1,r}(\Omega)}\right)$$
(33)

where $\boldsymbol{\alpha} = (\alpha_1, \ldots, \alpha_I).$

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• Generalized Solutions (p < 2):

Using a duality argument, we obtain the following result :

Theorem

We suppose that p < 2. Soit $\boldsymbol{f} \in [\boldsymbol{H}_0^{r',p'}(\operatorname{curl}, \Omega)]'$, curl $\boldsymbol{a} \in \boldsymbol{L}^s(\Omega)$ and $\pi_0 \in W^{1-1/r,r}(\Gamma)$ with

$$r = 1 + \epsilon'$$
 if $p < \frac{3}{2}$, $r = \frac{9 + 6\epsilon}{9 + 2\epsilon}$ if $p = \frac{3}{2}$ and $r = \frac{3p}{3+p}$ if $\frac{3}{2} , (34)$

$$s = (1+\epsilon')\frac{3p}{4p-3-\epsilon'(3-p)} \text{ if } p < \frac{3}{2}, \ s = \frac{3}{2}+\epsilon \text{ if } p = \frac{3}{2} \text{ and } s = \frac{3}{2} \text{ if } \frac{3}{2} < p < 2$$
(35)

where $\epsilon, \epsilon' > 0$ are arbitrary. Problem (\mathcal{OS}_N) has a unique solution $(u, \pi, \alpha) \in W^{1,p}(\Omega) \times W^{1,r}(\Omega) \times \mathbb{R}^I$ satisfying the estimate:

$$\|\boldsymbol{u}\|_{\boldsymbol{W}^{1,p}(\Omega)} \leq C(1 + \|\mathbf{curl}\,\boldsymbol{a}\|_{\boldsymbol{L}^{s}(\Omega)})^{2} (\|\boldsymbol{f}\|_{[\boldsymbol{H}_{0}^{r',p'}(\mathbf{curl},\Omega)]'} + \|\pi_{0}\|_{W^{1-1/r,r}(\Omega)}),$$

$$\|\pi\|_{W^{1,r}(\Omega)} \leq C(1 + \|\mathbf{curl}\, a\|_{L^{s}(\Omega)})^{3}(\|f\|_{[H_{0}^{r',p'}(\mathbf{curl},\Omega)]'} + \|\pi_{0}\|_{W^{1-1/r,r}(\Omega)})$$

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The Navier-Stokes problem (\mathcal{NS}_N)

$$\begin{cases} -\Delta \, \boldsymbol{u} + \operatorname{\mathbf{curl}} \, \boldsymbol{u} \times \boldsymbol{u} + \nabla \, \boldsymbol{\pi} = \boldsymbol{f} & \text{in } \Omega, \\ \operatorname{div} \, \boldsymbol{u} = \chi & \text{in } \Omega, \end{cases}$$

$$(\mathcal{NS}_N)$$

$$\boldsymbol{u} \times \boldsymbol{n} = \boldsymbol{g}$$
 on Γ ,

$$\pi = \pi_0 \text{ on } \Gamma_i \text{ and } \pi = \pi_0 + c_i \quad \text{ on } \Gamma_i,$$

 $\int_{\Gamma_i} \boldsymbol{u} \cdot \boldsymbol{n} \, \mathrm{d}\boldsymbol{\sigma} = 0, \ i = 1, \dots, I,$

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In the search of a proof of the existence of generalized solution for the Navier-Stokes equations (\mathcal{NS}_N) , we consider the case of small enough data. Let

$$\boldsymbol{f} \in [\boldsymbol{H}_0^{r',p'}(\operatorname{\mathbf{curl}},\,\Omega)]', \quad \chi \in W^{1,r}(\Omega), \quad \boldsymbol{g} \in \boldsymbol{W}^{1-1/p,p}(\Gamma), \quad \pi_0 \in W^{1-1/r,r}(\Gamma)$$

with

$$p > \frac{3}{2}$$
 and $r = \frac{3p}{3+p}$.

i) There exists a constant $\alpha_1 > 0$ such that, if

$$\|f\|_{[H_0^{r',p'}(\operatorname{curl},\Omega)]'} + \|\chi\|_{W^{1,r}(\Omega)} + \|g\|_{W^{1-1/p,p}(\Gamma)} + \|\pi_0\|_{W^{1-1/r,r}(\Gamma)} \le \alpha_1,$$

then, there exists a solution

$$(\boldsymbol{u}, \pi, \boldsymbol{c}) \in \boldsymbol{W}^{1,p}(\Omega) \times W^{1,r}(\Omega) \times \mathbb{R}^{I}$$

to problem (NS $_{N})$ verifying the estimate

$$\|u\|_{W^{1,p}(\Omega)} \le C \big(\|f\|_{[H_0^{r',p'}(\operatorname{curl},\Omega)]'} + \|\chi\|_{W^{1,r}(\Omega)} + \|g\|_{W^{1-1/p,p}(\Gamma)} + \|\pi_0\|_{W^{1-1/r,r}(\Gamma)} \big),$$

with
$$c_i = \langle \mathbf{f}, \nabla q_i^N \rangle_{\Omega_{r', p'}} + \int_{\Gamma} (\chi - \pi_0) \nabla q_i^N \cdot \mathbf{n} - \int_{\Omega} (\operatorname{\mathbf{curl}} \mathbf{u} \times \mathbf{u}) \cdot \nabla q_i^N.$$

ii) Moreover, there exists a constant $\alpha_2 \in]0, \alpha_1]$ such that this solution is unique, if

$$\|f\|_{[H_0^{r',p'}(\operatorname{curl},\Omega)]'} + \|\chi\|_{W^{1,r}(\Omega)} + \|g\|_{W^{1-1/p,p}(\Gamma)} + \|\pi_0\|_{W^{1-1/r,r}(\Gamma)} \le \alpha_2.$$

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