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Temporal optical besselon waves

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Abstract: We analyse the temporal properties of the optical pulse wave that is obtained by applying a set of spectral $\pi/2$ phase shifts to continuous-wave light that is phase-modulated by a temporal sinusoidal wave. We develop an analytical model to describe this new optical waveform that we name ‘besselon’. We also discuss the reduction of sidelobes in the wave intensity profile by means of an additional spectral π phase shift, and show that the resulting pulses can be efficiently time-interleaved. The various predicted properties of the besselon are confirmed by experiments demonstrating the generation of low-duty cycle, high-quality pulses at repetition rates up to 28 GHz.

Keywords: phase modulation ; high-repetition rates sources ; optical waveforms.

I. Introduction

Ultrafast optics has provided extremely efficient means to generate various pulse waveforms with durations of a few picoseconds and very high-repetition rates. In principle, arbitrarily complex optical waveforms can be synthesised at high repetition rates by careful phase-intensity spectral shaping of frequency comb sources [1]. It remains however of high importance to have simple experimental methods for optical waveform generation as well as a clear mathematical description of the generated pulse profiles. Sinusoidal intensity profiles can be efficiently generated by use of standard high-bandwidth modulators. Gaussian and hyperbolic secant pulses are routinely delivered from fibre lasers. Within the toolbox of optical signal processing, triangular, parabolic and rectangular shapes can be achieved by linear [2] or nonlinear [3, 4] sculpturing. In the context of linear shaping or nonlinear fibre propagation, the properties of other pulse waveforms have also been the subject of recent discussion, including solitons over finite background [4, 5], superregular breathers [6], Riemann waves [7], flaticon waves [8], Airy pulses [9], and Hermite-Gaussian structures [10]. Some of these waves feature very strong oscillations in their temporal profiles [8-10].

An attractive method to generate high-repetition-rate and stable pulse trains relies on the spectral processing of periodically phase-modulated continuous-wave light [11]. Within this approach, the use of a quadratic spectral phase profile has indeed enabled the generation of ultra-short pulses at high repetition rates [12, 13] as well as of flat-top profiles [14]. We have recently introduced the line-by-line application of $\pi/2$ spectral phase shifts to a periodically phase-modulated continuous wave as an enhanced technique to obtain high-quality ultra-short pulse trains [15]. This pulse generation method can sustain multi-wavelength modulation and temporal multiplexing [16], as well as operation with a dual-tone signal [17]. Furthermore, compared to the original approach, the extinction ratio of the pulses and the suppression of sidelobes in their temporal intensity profiles are remarkably enhanced [15]. However, to date, no clear analytical description and explanation of the generated pulses have been reported. The discussion has essentially relied on numerical simulations with the amplitude of the initial phase modulation restricted to 2 radians.

In this paper, we fill this gap by providing an insight into the new type of an optical wave structure generated with this method, which we name ‘besselon’. We unveil its properties for a broad range of phase-modulation amplitudes, and we derive simple analytical guidelines to predict

the main pulse features. An approximate model is also introduced to help explain the evolution of the wave profile. Further, we theoretically discuss the possibility of doubling the pulse repetition rate, and identify optimum operating conditions to realise this. Our analytical predictions are validated by the experimental demonstration of the generation of besselon pulse patterns at repetition rates of 14 GHz and 28 GHz.

II. Theoretical model for besselon waves

A Pulse generation and properties

We consider a continuous optical wave with amplitude ψ_0 and carrier angular frequency ω_c , $\Psi(t) = \psi_0 \psi(t) e^{i\omega_c t}$ whose phase is temporally modulated by a sinusoidal wave:

$$\psi(t) = e^{i A_m \cos(\omega_m t)}$$

where A_m is the amplitude of the phase modulation and ω_m is its frequency. As a result of this phase modulation, the spectrum of frequency components of the wave envelope will consist of a series of discrete lines evenly spaced by ω_m , which can be obtained from the Jacobi-Anger expansion [18-20] of $\psi(t)$:

$$\psi(t) = \sum_{n=-\infty}^{\infty} i^n J_n(A_m) e^{i n \omega_m t}$$

Therefore, the n^{th} spectral component will have an intensity proportional to $J_n^2(A_m)$, where $J_n(x)$ is the Bessel function of the first kind and order n . An important characteristic of this spectrum is the existence of a $\pi/2$ phase shift between successive frequency components. The application of a quadratic spectral phase profile - such as produced by propagation through a purely dispersive element - to $\psi(t)$, with a curvature chosen so as to optimise the peak power of the resulting pulses, will enable partial compensation of the initial sinusoidal phase. This will in turn lead to the emergence of temporally localised, periodic pulse structures with a period $T_0 = 2\pi / \omega_m$ [12, 13], as shown in Fig. 1(b). We can see in Fig. 1(b) that increasing the amplitude A_m of the initial modulation causes shortening of the central part of the pulses accompanied by a concomitant increase of the pulse peak power. However, since the initial spectral phase is not perfectly

cancelled, the presence of a residual background impairs the resulting pulse intensity profiles [14]. On the other hand, the recent advances in linear pulse shaping technology make it now possible to manipulate optical frequency combs line by line. With our pulse-shaping approach, imprinting a $\pi/2$ phase shift to each individual spectral component of the phase-modulated continuous wave enables the synthesis of a new optical pulse field, which we call ‘besselon’ because of the important role played by Bessel functions in the description of the waveform features. The besselon wave envelope, ψ_B , can be represented as:

$$\psi_B(t) = \sum_{n=-\infty}^{\infty} J_n(A_m) e^{i n \omega_m t} = J_0(A_m) + 2 \sum_{n=1}^{\infty} J_n(A_m) \cos(n \omega_m t) \quad (1)$$

In Ref. [15], we have emphasised that this pulse structure exhibits an excellent extinction ratio (defined as the ratio of the power at $t = T_0/2$ to the power at $t = 0$) combined with the absence of spurious pedestals for values of the initial phase-modulation amplitude A_m around 1.1 rad. Figure 1(a) shows the evolution of the wave intensity profile $|\psi_B|^2$ with A_m across a wide range of modulation amplitude values, and features distinctly different from that of the pulse pattern resulting from spectral manipulation with a quadratic spectral phase. The besselon pulses do not experience a decrease of their temporal duration with increasing modulation amplitude, and their extinction ratio may also vary significantly with varying amplitude. The central part of the pulses also evolves in a different manner: sidelobes develop and tend to merge with the central part that displays an increasingly large number of strong oscillations.

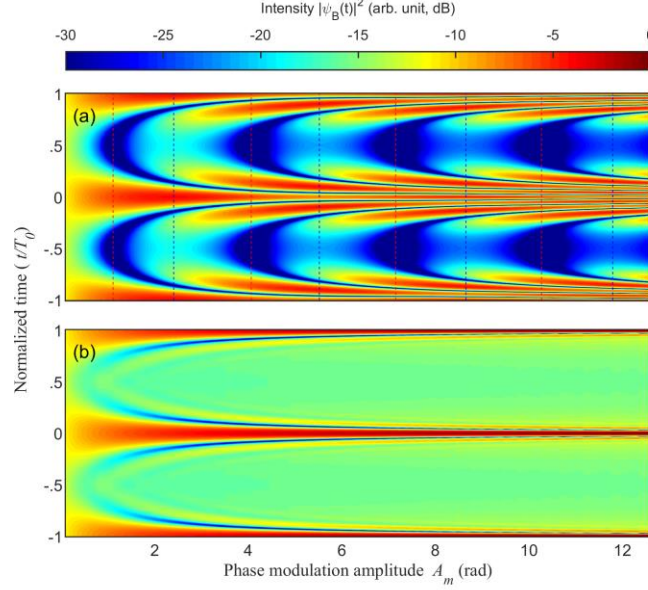


Figure 1: Evolution of the wave intensity profile with the amplitude of the initial sinusoidal phase modulation. **(a)** Besselon intensity profile $|\psi_B|^2$ as given by Eq. (1), and **(b)** profile synthesised by use of an optimal quadratic spectral phase profile as obtained from numerical simulations. The wave profiles are displayed using a logarithmic scale color map. The red vertical dashed lines in panel (a) represent the values of A_m satisfying the condition for maximal extinction ratio, $I_{J_0}(A_m) = 1$. The blue vertical dashed lines represent the zeros $A_{m,p}$ of $J_0(A_m)$, corresponding to the locations of the maxima of the peak power of the Besselon.

To gain better insight into the Besselon pulse wave, we have plotted on Fig. 2 the evolution of the pulse peak power at $t = 0$, the minimum power at $t = T_0/2$ and the extinction ratio of the pulse train with the phase-modulation amplitude. The Fourier-transform limited nature of the Besselon wave for $A_m \leq 2.40$ rad brings about a continuous increase of the peak power with A_m , which reaches higher values than the peak power of the pulse wave obtained after a dispersive element. For further increase of the modulation amplitude, however, the peak power of the Besselon wave oscillates and remains below the value of 6.10 obtained at $A_m = 2.40$ rad. From Eq. (1), one may derive analytically the value of the maximum of the field (i.e. at $t = 0$) :

$$\psi_B(t = 0) = J_0(A_m) + 2 \sum_{n=1}^{\infty} J_n(A_m)$$

and using the relations [18, 21]

$$2 \sum_{n=1}^{\infty} J_{2n}(A_m) = 1 - J_0(A_m)$$

$$2 \sum_{n=0}^{\infty} J_{2n+1}(A_m) = \int_0^{A_m} J_0(x) dx$$

we can obtain the peak amplitude of the beselson:

$$\psi_B(t=0) = 1 + I_{J_0}(A_m) \quad (2)$$

with I_{J_0} being plotted on Fig. 2(d) (red line) and defined as

$$I_{J_0}(A_m) = \int_0^{A_m} J_0(x) dx \quad (3)$$

We can see in Fig. 2(a) that the peak-power predictions from Eq. (2) are in perfect agreement with the results of numerical simulations. The amplitude values $A_{m,p}$ of the phase modulation at which the peak power of the beselson is maximal are given by $dI_{J_0}/dA_m|_{A_{m,p}} = 0$, that is $J_0(A_{m,p}) = 0$. In other words, the maxima of the peak power occur at the zeros of the Bessel function of zero order. For these amplitudes ($A_{m,p} = 2.4048, 5.5201, 8.6537, \dots$), the central frequency component of the spectrum vanishes and the pulse train becomes carrier-suppressed. Using the following asymptotic form of the zeroth-order Bessel function for large arguments [18] (Fig. 2(d), black circles)

$$J_0(x) \approx \sqrt{\frac{2}{\pi x}} \cos\left(x - \frac{\pi}{4}\right) \quad (4)$$

we can obtain an approximate expression for the location of the p^{th} maximum:

$$A_{m,p} \approx \frac{3\pi}{4} + (p-1)\pi, \quad p \geq 1 \quad (5)$$

Equation (5) indicates that the power at $t=0$ of the beselson experiences periodic growth and decay with a period of π rad. We also note that as $I_{J_0}(x)$ asymptotically tends to 1 for large values of x [21], the amplitude of the oscillations decreases with increasing values of A_m .

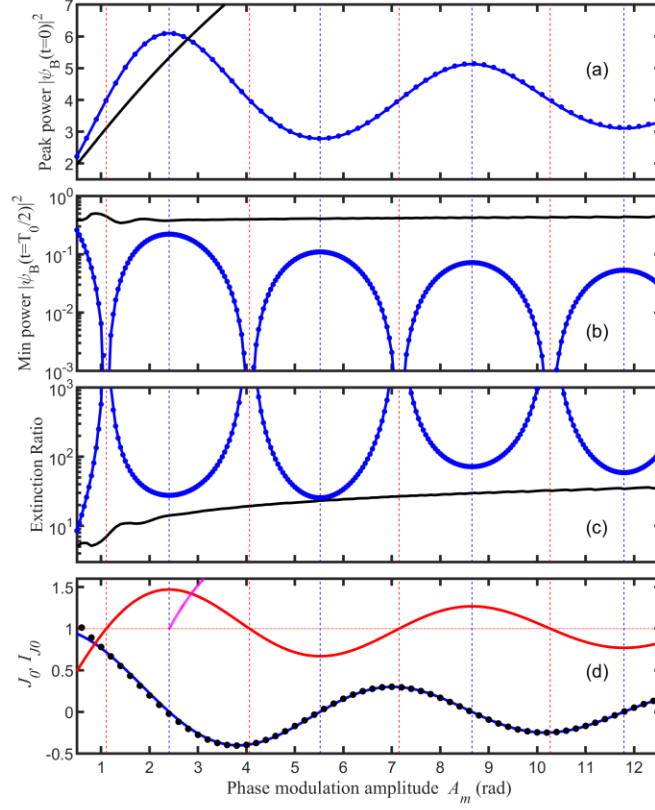


Figure 2: (a-c) Evolution of the **(a)** peak power, **(b)** minimum power, and **(c)** extinction ratio of the beselson pulse wave with the amplitude of the initial sinusoidal phase modulation. The predictions provided by (2), (6) and (7) (blue circles) are compared with the results of numerical simulations (blue curves). Also shown are the numerical simulation results for the pulse wave synthesised by use of an optimal quadratic spectral phase profile (black curves). **(d)** Bessel function of the first kind and order 0, $J_0(A_m)$, (blue curve) and its integral $I_{J_0}(A_m)$ as defined by Eq. (3) (red curve). The filled black circles represent the asymptotic form of $J_0(A_m)$ for large argument (Eq. (4)). The red vertical dashed lines represent the values of A_m satisfying the condition for maximal extinction ratio, $I_{J_0}(A_m) = 1$. The blue vertical dashed lines represent the zeros $A_{m,p}$ of $J_0(A_m)$, corresponding to the locations of the maxima of the peak power. The purple curve in panel (d) represents the function $1 - 2J_0(A_m)$ (Eq. (19) in section D).

The minimum power of the beselson pulse wave (Fig. 2(b)) also displays strong variations with varying amplitude of the initial phase modulation. This contrasts with the pulse wave obtained after a dispersive element, where the intensity level between consecutive pulses remains almost constant for A_m above 2 rad. Similarly to the temporal maximum of the beselson, we can obtain the minimum amplitude of the beselson from Eq. (1):

$$\psi_B\left(t = \frac{T_0}{2}\right) = 1 - I_{J_0}(A_m), \quad (6)$$

which yields the following expression for the extinction ratio of the pulses:

$$ER(A_m) = \frac{1 + I_{J0}(A_m)}{1 - I_{J0}(A_m)} \quad (7)$$

We can therefore deduce that for amplitudes of the sinusoidal modulation satisfying the equation $I_{J0}(A_m) = 1$, the background level between pulses is zero hence the extinction ratio is maximal. This condition is consistent with the condition derived in [15] based on numerical simulation data. Similarly to the peak and minimum powers, the extinction ratio of the pulses varies strongly with A_m , while remaining larger than the extinction ratio of the pulse wave synthesised by use of a quadratic spectral phase profile across the whole A_m variation range. Note that, for $I_{J0}(A_m) = 1$, the maximum amplitude of wave is simply 2, leading to an intensity of 4.

B Details of the pulse shape

We now try to better describe the shape of the besselon at the points of maximal extinction ratio, i.e. when $I_{J0}(A_m) = 1$. Temporal profiles ψ_B obtained for $A_m = 1.1086, 4.0628, 7.15$ and 10.2695 rad are plotted on Fig. 3 (panels a). From Eq. (1), we can first note that the besselon is a waveform without any complex part, contrary to the waveform achieved after quadratic phase compensation. If the besselon does not present any negative part for $A_m < A_{m,l}$, this is not the case for higher values of A_m and we can clearly see negative and positive parts that alternate, stressing that two successive oscillations of the wave are π -shifted. We also note that for an amplitude of oscillation $A_{m,p}$, we obtain $2p+1$ peaks in the temporal intensity profile.

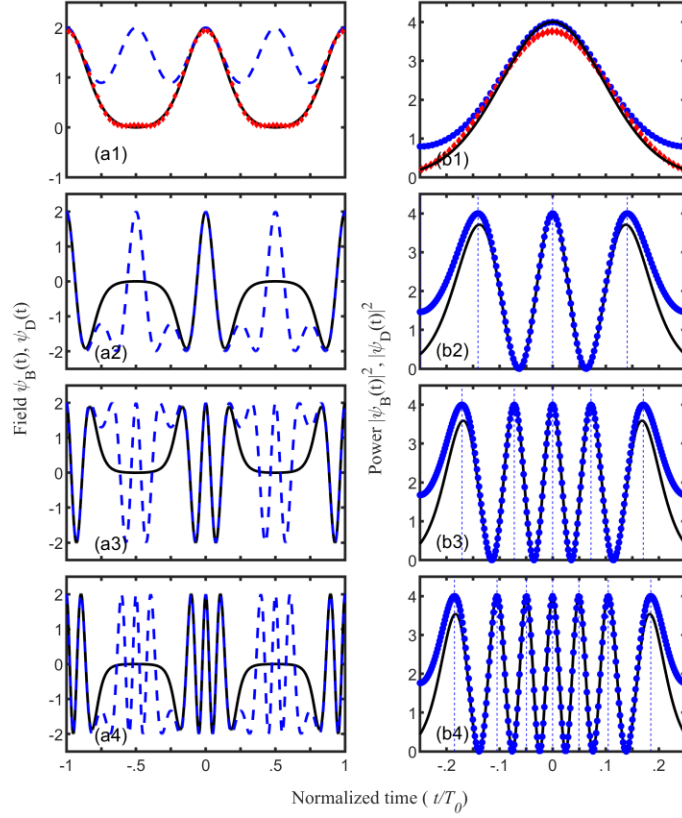


Figure 3: Temporal (a) amplitude and (b) power of the Bessel function for various values of A_m leading to an optimum extinction ratio (the condition $I_{J_0}(A_m) = 1$ is fulfilled). Results obtained for $A_m = 1.1086, 4.0628, 7.15$ and 10.2695 are plotted on panel 1, 2, 3 and 4 respectively. The profile of the Bessel function ψ_B (solid black line, Eq. (1)) is compared with the profile of $2\psi_D$ (dashed blue line, Eq. (8)). Approximation of the Bessel function proposed by Eq. (12) is plotted with red diamonds. The vertical dashed line in panels (b) indicate the position of the peak as predicted by Eq. (9).

In order to better understand the origin and the properties of these peaks, it could be interesting to express the Bessel function as:

$$\psi_B(t) = J_0(A_m) + 2 \sum_{n=1}^{\infty} J_{2n}(A_m) \cos(2n \omega_m t) + 2 \sum_{n=0}^{\infty} J_{2n+1}(A_m) \cos((2n+1) \omega_m t)$$

and to take into account that [18, 21]:

$$J_0(A_m) + 2 \sum_{n=1}^{\infty} J_{2n}(A_m) \cos(2n \omega_m t) = \cos(A_m \sin(\omega_m t)) = \psi_D(t) \quad (8)$$

It is interesting to superimpose the profiles of $2 \psi_D$ and ψ_B (see Fig. 3, blue dashed line and solid black lines respectively) and to note that the central part of the beselson is remarkably well described by $2 \psi_D$. One may also note that the repetition rate of ψ_D is twice the repetition rate of the beselson and that $2 \psi_D(t) = \psi_B(t) + \psi_B(t-T_0/2)$. Note that ψ_D can be generated directly from a Mach-Zehnder device driven by a sinusoidal electrical wave. From $2 \psi_D$, it is possible to deduce several interesting features of the central part of the beselson. First the power of the various peaks equals 4 (amplitude of 2). Moreover, the temporal location of the maxima t_M of the temporal intensity profiles are:

$$t_M = \text{asin}\left(\frac{n \pi}{A_m}\right) / \omega_m \quad (9),$$

with $n \leq (N-1)/2$, N being the number of peaks that is given by $\lfloor 2A_m / \pi \rfloor + 1$. We can also derive the full width at half maximum of the beselson:

$$t_B = \frac{2}{\omega_m} \text{asin}\left(\frac{\pi}{A_m} \left(\frac{N}{2} - \frac{1}{4}\right)\right) \quad (10),$$

as well as the temporal width of the central part:

$$t_C = \frac{2}{\omega_m} \text{asin}\left(\frac{\pi}{4A_m}\right) \quad (11).$$

Results of those analytical predictions are reported on Fig. 4 where we can make out an excellent agreement with the numerical simulations. The difference between the waveform resulting from a quadratic spectral phase compensation and the beselson is also highlighted.

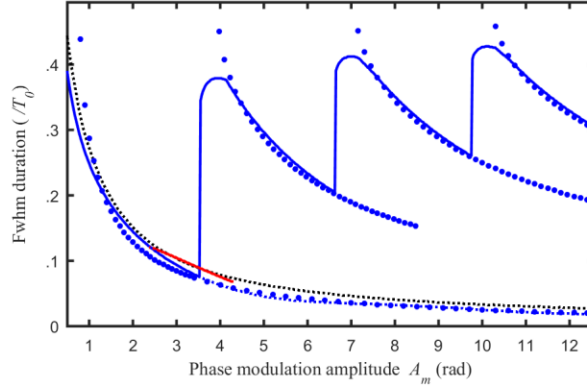


Figure 4: Evolution of the full-width at half maximum duration of the Bessel wave. Numerical simulations (blue solid line) are compared with the analytical predictions of Eq. (10) and (11) (blue circles). The fwhm duration of the pulse obtained after optimal quadratic spectral phase compensation is plotted with black dotted line, whereas the duration of the central part of the Bessel is plotted with dotted blue line. Results obtained for ψ_B' are plotted with red line.

C Simplified expression for low values of A_m

In this section, we now focus on a simplified model that efficiently describes the central part of the pulse for low values of A_m . Indeed, from Fig. 3(a1) and (b1) it appears that $2\psi_D$ is not accurate enough for low values of A_m (typically below 1.5 rad). In this context, we have shown in previous works that for $A_m = 1.1$ rad, the generated waveform was very close from a Gaussian profile [15]. We have also recently discussed the similarity between the Bessel profile obtained for low A_m and the typical profile of an Akhmediev breather [22] that may exist in nonlinear fibre optics [23]. Here, we propose to interpret the waveform from another viewpoint. As shown in Fig. 5(a), we can consider that the spectrum obtained for low A_m is the result of the coherent superposition of three waves: a continuous wave ψ_1 with an amplitude $J_0 - 2J_2$ and two sinusoidally partially modulated wave (a continuous background with an amplitude J_1 that is modulated with an amplitude $2J_2$ at a frequency ω_m) that are frequency shifted by $-\omega_m$ and ω_m (ψ_2 and ψ_3 respectively). In the temporal domain, this leads to

$$\begin{aligned} \psi_B(t) \simeq & [J_0 - 2J_2] \\ & + \left[(J_1 + 2J_2 \cos(\omega_m t)) e^{i\omega_m t} \right], \\ & + \left[(J_1 + 2J_2 \cos(\omega_m t)) e^{-i\omega_m t} \right] \end{aligned}$$

which can be rewritten as

$$\psi_B(t) \simeq [J_0 - 2J_2] + 2[J_1 + 2J_2 \cos(\omega_m t)] \cos(\omega_m t) \quad (12),$$

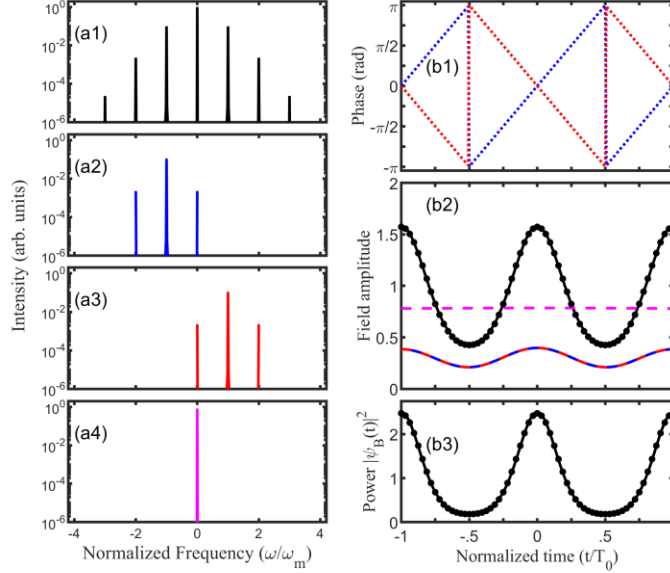


Figure 5: Comparison of the Bessel function properties obtained for $A_m = .6$ rad with the predictions of a simplified 3-wave model. **(a)** Optical spectrum of ψ_B , ψ_2 , ψ_3 and ψ_1 (panels 1,2,3,4 respectively). **(b)** Temporal properties: chirp profile, field and intensity profiles (panel 1, 2, 3, respectively). The results for waves ψ_1 , ψ_2 and ψ_3 are plotted with purple, blue and red colors respectively. Results ψ_B are plotted with solid black line and compared with the approximate model (black circles, Eq. (12)).

The amplitude and phase profiles of the three waves are plotted in Fig. 3(b1-b2) which helps to understand how the waveform is formed. For $t = 0$, the three waves are in phase, they interfere constructively and the amplitude is maximum so that

$$\psi_B(t=0) \simeq J_0(A_m) + 2J_1(A_m) + 2J_2(A_m) \quad (13),$$

On the contrary, for $t = T_0/2$, we note that the modulated waves are both π -shifted with respect to the continuous wave. They will interfere destructively and the amplitude of the Bessel function is minimum as follows:

$$\psi_B(t=T_0/2) \simeq J_0(A_m) - 2J_1(A_m) + 2J_2(A_m) \quad (14),$$

The Bessel function therefore reaches a null value at $T_0/2$ when

$$J_0(A_m) = 2(J_1(A_m) - J_2(A_m)) \quad (15),$$

that gives an optimum value of $A_m = 1.187$ rad, which is consistent with the value found previously (1.109 rad). Finally, we can note that for $t = T_0/4$, the two sinusoidal waves present a phase difference of π . Consequently, they cancel each other and $\psi_B(t = T_0/4) \approx J_0(A_m) - 2J_2(A_m)$. Therefore, for low values of A_m , $\psi_B(t = T_0/4)$ continuously decreases indicating that the pulse gets temporally shorter and shorter. We have plotted the intensity profiles resulting from the approximate Eq. (12) on Fig. 5(b3) for a value of $A_m = .6$ rad. The approximation is nearly indiscernible compared to the exact prediction from Eq. (1). The agreement is also rather good for $A_m = 1.10$ rad as can be seen on Fig. 3(a1-b1). In Fig. 6(a-b), we compare the peak-power and extinction ratios predicted by this approximate model (purple circles) and once again, we find an excellent agreement as long as the level of the component J_3 is low enough.

D Reduction of the lateral temporal sidelobes

As can be seen in Fig. 1(a) and in the magnified view provided in Fig. 6(c1), the evolution is marked for $A_m > A_{m,1}$ by the development of noticeable sidelobes in the pulse. Given the non-negligible energy fraction contained in those sidelobes, this leads to a decrease of the peak power as stressed in panel (a) of Fig. 6. Recall that the value $A_{m,1}$ corresponds to the amplitude of the initial phase modulation for which the central spectral component decreases down to zero. It then becomes negative up to to $A_m = A_{m,2}$. Consequently, from the simplified model we developed in the previous section, we can understand that at $t = 0$ and between $A_{m,1}$ and $A_{m,2}$, the continuous background becomes π -shifted with respect to the sinusoidally modulated waves. Therefore, the interference changes from constructive interference (below $A_m < A_{m,1}$) to destructive interference. The resulting structure has a reduced peak power and is not Fourier-transform limited anymore. In order to overcome this effect, a solution is to imprint an additional π shift on the central component for $A_m > A_{m,1}$. In this case, the modified Bessel pulse ψ'_B is provided by $\psi'_B(t) = \psi_B(t)$ for $A_m < A_{m,1}$ and:

$$\psi_B'(t) = -J_0(A_m) + 2 \sum_{n=1}^{\infty} J_n(A_m) \cos(n \omega_m t) = \psi_B(t) - 2 J_0(A_m) \quad (16)$$

elsewhere. The corresponding peak amplitude becomes for $A_m > A_{m,0}$:

$$\psi_B'(t=0) = 1 + I_{J_0}(A_m) - 2J_0(A_m) \quad (17)$$

The maximum amplitude is therefore increased by $2 |J_0(A_m)|$ with respect to the case without π phase shift. The condition on A_m that now leads to a maximum amplitude is that $J_0(A_m) = -2 J_1(A_m)$, leading to a value of the maximum for $A'_{m,1} = 3.40$ rad (cyan vertical line in Fig. 6a) for a peak-power value that is 8.93, representing a 46% increase compared to the results obtained for ψ_B . The extinction ratio becomes:

$$ER'(A_m) = \frac{1 + I_{J_0}(A_m) - 2J_0(A_m)}{1 - I_{J_0}(A_m) - 2J_0(A_m)}, \quad (18)$$

which tends to infinity when:

$$I_{J_0}(A_m) = 1 - 2J_0(A_m) \quad (19)$$

which leads to a condition $A_m = 2.86$ rad as can be found from Fig. 1(d). This value differs quite significantly than the one found for the unmodified beselson. Note that an ER above 10^2 is obtained for A_m in the range of 2.6-3.15 rad. The analytical results from Eq. (17), (18) are in perfect agreement with the numerical results as can be seen on Fig. 6(a-b). Comparison of the temporal profiles obtained without or with an additional π phase shift on the central component is provided in Fig. 6(c). It clearly highlights the decrease of the sidelobes of the pulse that are at -13 dB of the peak power for $A_m = 2.86$ rad (while in the case of ψ_B the decrease is only -7.22 dB for the same A_m). If there is an improvement in terms of peak-power, we do not observe such an improvement in terms of temporal duration (see Fig. 4). For $A_m = 2.86$ rad, pulse with a fwhm temporal duration of $0.1 T_0$ is obtained, which represents a significant reduction when compared to the results achieved for $A_m = 1.1086$. Finally, note that for $A_m > 3.83$, the first lateral sidebands will experience a π phase shift as J_1 changes of sign. Therefore, insertion of additional phase shifts could be required to further improve the peak-power.

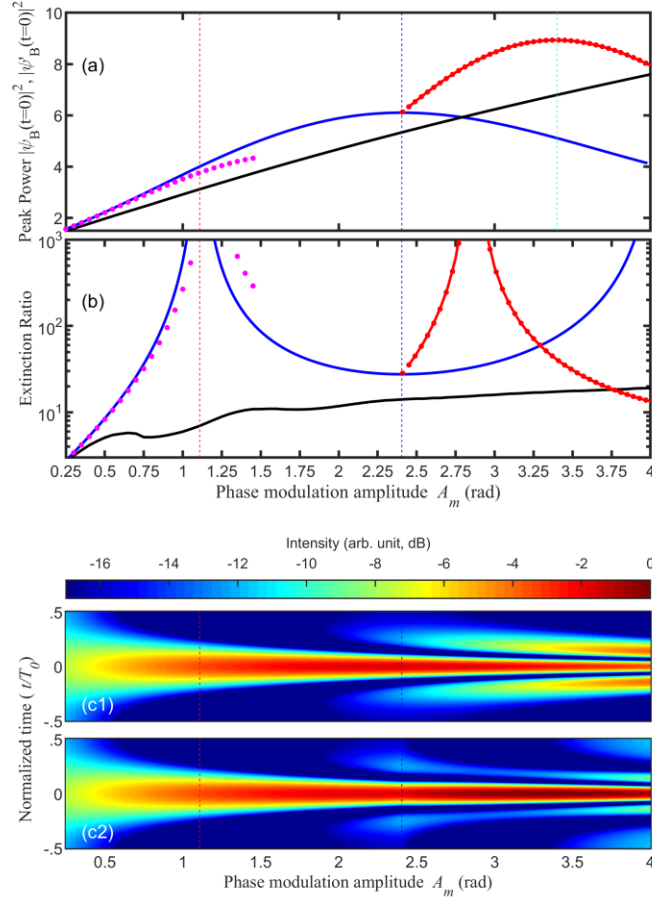


Figure 6: (a-b) Evolution of the properties of the temporal intensity profile according to the initial amplitude of the sinusoidal phase modulation. **(a)** Peak power. **(b)** Extinction ratio. Results obtained after optimum quadratic phase compensation (black solid lines) are compared with the properties of the Besselon (blue lines, Eq. (2) and (7)) and of the modified Besselon (red lines, analytical predictions solid line, analytical predictions – Eq. (17) and (18), red circles). Results from the approximate model (derived from Eq. (12)) are plotted with purple circles. **(c)** Intensity profiles obtained for the Besselon $|\psi_B|^2$ (Eq. (1)) and modified Besselon $|\psi'_B|^2$ (Eq. (16)).

E Time multiplexing

In this last section, we consider the possibility to temporally interleave the Besselon pulses ψ'_B in order to double the repetition rate. Many solutions may exist such as the use of the fractional Talbot effect [24, 25], temporal delay lines [26], or spectrum manipulation such as intensity

modification [27] or phase only shaping [28]. Here, in order to double the repetition rate, we suppress all odd spectral components. In this context, the resulting field is restricted to:

$$\begin{cases} \psi'_{2B}(t) = J_0(A_m) + 2 \sum_{n=1}^{\infty} J_{2n}(A_m) \cos(2n \omega_m t) & \text{for } A_m \leq A_{m,1} \\ \psi'_{2B}(t) = -J_0(A_m) + 2 \sum_{n=1}^{\infty} J_{2n}(A_m) \cos(2n \omega_m t) & \text{for } A_m > A_{m,1} \end{cases}$$

That, given Eq. (8), can be rigorously expressed as :

$$\begin{cases} \psi'_{2B}(t) = \psi_D(t) & \text{for } A_m \leq A_{m,1} \\ \psi'_{2B}(t) = \psi_D(t) - 2 J_0(A_m) & \text{for } A_m > A_{m,1} \end{cases} \quad (20)$$

Results of the repetition doubled pulse train are shown on Fig. 7(a). One can make out that the optimal condition in terms of extinction ratio is obtained for $A_m = 1.57$ and $A_m = 3.78$ and not for $A_m = 1.10$ and $A_m = 2.86$ rad as one could have initially anticipated from sections IIA,D (Eq. (7) and Eq. (19)). Indeed, the value at $t = T_0/4$ is given by:

$$\begin{cases} \psi'_{2B}(t = T_0/4) = \cos(A_m) & \text{for } A_m \leq A_{m,1} \\ \psi'_{2B}(t = T_0/4) = \cos(A_m) - 2 J_0(A_m) & \text{for } A_m > A_{m,1} \end{cases}$$

and vanishes when

$$\begin{cases} A_m = \pi/2 & \text{for } A_m \leq A_{m,1} \\ 2 J_0(A_m) = \cos(A_m) & \text{for } A_m > A_{m,1} \end{cases} \quad (21)$$

In Fig. 7(b), we compare the profiles obtained before repetition rate doubling (panel 1) with the results achieved after processing (panel 2). We can make out that suppressing the odd components enables to get rid of the residual component observed between two pulses [26]. Nice pulse shapes are achieved especially for $A_m = 3.78$ rad (optimum operating point for $A_m > A_{m,1}$) where the spurious sidelobes are decreased down to a level of nearly -20dB the peak power. Temporal duration of $0.092 T_0$ leads to a duty cycle of 0.18. The cost of the approach is mainly in the drop of the peak power, which decreases by a factor more than two.

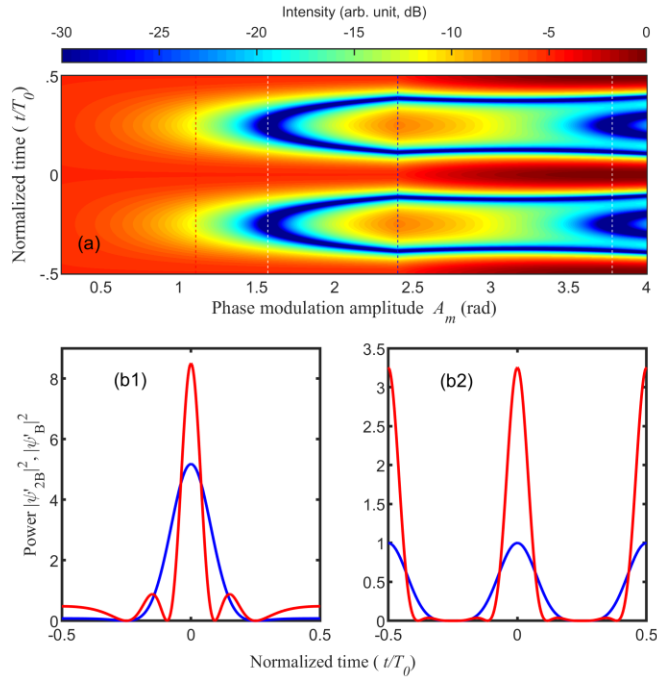


Figure 7 : (a) Evolution of the intensity profile obtained for modified besselon $|\psi'_{2B}|^2$ at the doubled repetition rate (Eq. **Erreur ! Source du renvoi introuvable.**). (b) Detail of the temporal intensity profile obtained before and after repetition rate doubling $|\psi'_B|^2$ and $|\psi'_{2B}|^2$ (panels 1 and 2 respectively), for $A_m = \pi/2$ and 2.78 rad respectively (blue and red curves).

III. Experimental validation

A. Experimental setup

The experimental setup is sketched in Fig. 8(a) and is based on devices that are commercially available and typical of the telecommunication industry. A continuous wave laser at 1550 nm is first temporally phase modulated using a Lithium Niobate electro-optic device driven by a sinusoidal electrical signal at a frequency $f_m = 14$ GHz. The use of a low V_π modulator enables us to investigate modulation amplitudes A_m up to 3.75 rad. Note that amplitudes of modulations exceeding 10 rad have already been demonstrated in the past using a resonant microwave modulator [29, 30] or cross-phase modulation in a highly non-linear fiber [8].

A linear spectral shaper (Finisar Waveshaper) based on liquid crystal on silicon technology [31] is then used to apply the suitable spectral phase profiles made of simple discrete spectral phase

shifts of $\pi/2$ applied between two successive components for the experimental demonstration of the waves ψ_B and ψ'_B (see panels 8(b1) and 8(b2) respectively). For the doubling of the repetition rates, we have inserted π phase shifts at frequencies $(2p+1) f_m$, with $p \geq 0$ for the wave ψ'_{2B} ($A_m \leq A_{m,l}$) and with $p \geq 1$ for ($A_m > A_{m,l}$) (see panels 8(b3) and 8(b4), red curves). Indeed, the discrete spectral π -phase shifts create a notch filter [32] that dramatically attenuates the spectral harmonics to be suppressed. Attenuation curves recorded using an amplified spontaneous emission source (ASE) as an input signal are shown with black solid lines and confirm that we can simultaneously achieve the spectral amplitude and phase targets using a simple phase-only spectral shaping. Let us mention that for operation at a fixed wavelength and a fixed repetition rate, the linear shaping stage can also be fully realized by a cascaded uniform fibre Bragg grating [33]. Note also that in order to increase the repetition rate by suppression one spectral component of two, other approaches can be implemented such as the use of a birefringent fibre [34] or Fabry-Perot interferometers [35].

The resulting signal is directly recorded by means of a high-speed optical sampling oscilloscope (1 ps resolution) and with a high-resolution optical spectrum analyzer. A low noise erbium-doped amplifier has been inserted in order to reach the optimum power level on the temporal detection device.

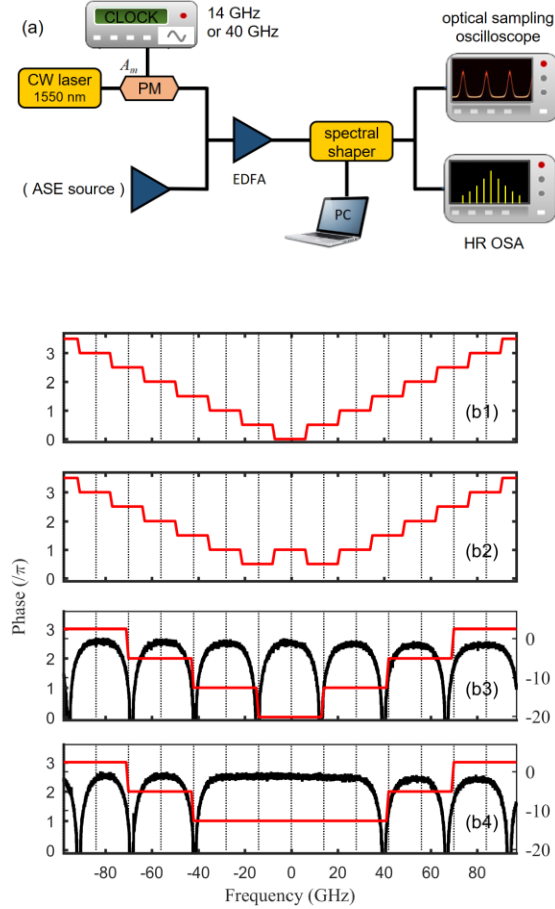


Figure 8: (a) Experimental setup : CW : Continuous Wave ; PM : Phase Modulator ; RF : Radio-Frequency ; ASE : Amplified Spontaneous Emission ; EDFA : Erbium Doped Fiber Amplifier ; HR OSA : High Resolution Optical Spectrum Analyzer. (b) Spectral phase profile (red line) applied on the programmable phase shaper for the generation of ψ_B , for ψ'_B and for the simultaneous phase shaping and repetition rate doubling of ψ_B and ψ'_B (panels 1,2, 3 and 4 respectively). The optical attenuation induced by the phase attenuation is plotted with black line.

B. Pulse waveform

A set of experimental temporal and spectral power profiles recorded for different levels of initial phase modulation A_m is summarized on Fig. 9. The agreement between the experimental results and the theoretical predictions is excellent, both in the temporal and spectral domains. We can recognize all the various features that have been discussed in the preceding sections. For low values of A_m (typically below 0.8), the wave is only partially modulated and a continuous background remains (panel a). For values around 1.1 rad (panel b), the extinction ratio becomes very high and a Gaussian-like waveform is obtained. For increasing values of A_m (2.47 rad, panel c) the

background impairs significantly the pulse train. Significant sidelobes then develop (panel d1). Those pedestals can be suppressed by inserting an additional π phase shift on the central spectral component (panel d3).

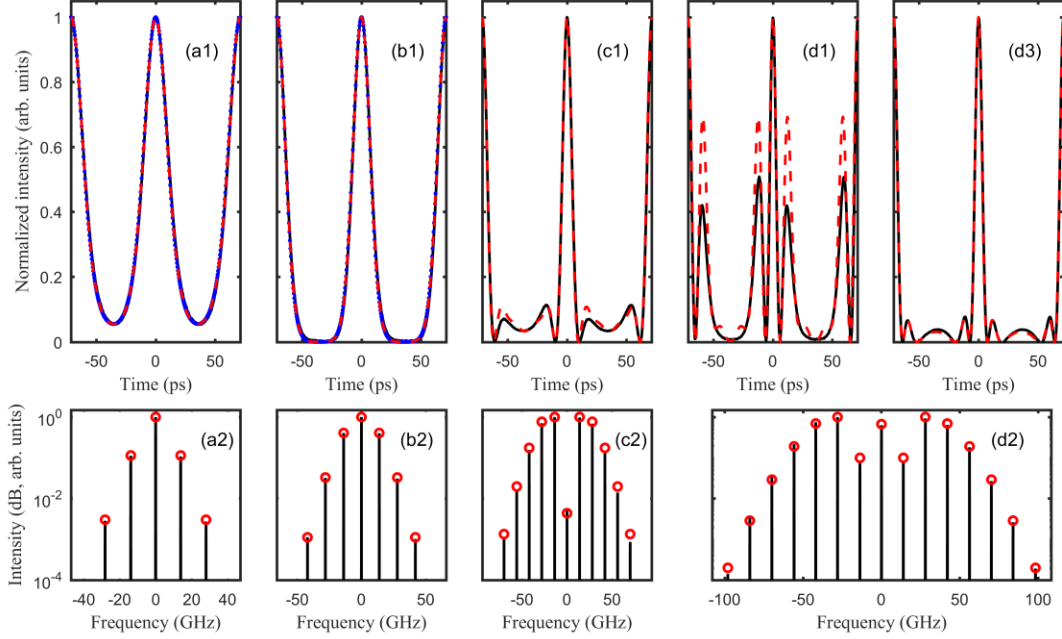


Figure 9: Temporal and spectral intensity profiles (panels 1 and 2, respectively) ψ_B obtained for levels of amplitude modulation A_m of 0.64 rad, 1.08 rad, 2.47 rad and 3.48 rad (panels a, b, c and d respectively). For 3.48 rad, the panel (d3) corresponds to the modified Besselon ψ'_B . Experimental results (black lines) are compared with the theoretical predictions (red dashed lines and red circles, Eq. (1) or (16) for panel (d3)) as well as the predictions of the approximate model (blue dashed line, Eq. (12)). All the profiles are normalized to their maximum.

The results of a more systematic experimental study are reported in Fig. 10 for both structure type ψ_B and ψ'_B . These results are to be compared with the results obtained in Fig. 6(c) and we note once again a remarkable agreement, both in terms of central shape and also in terms of the evolution of the sidelobes. We have also put the details of the time profiles recorded at the compression point corresponding to the best extinction ratio (2.83 rad), which further stresses the beneficial impact of the π phase shift on the central spectral component.

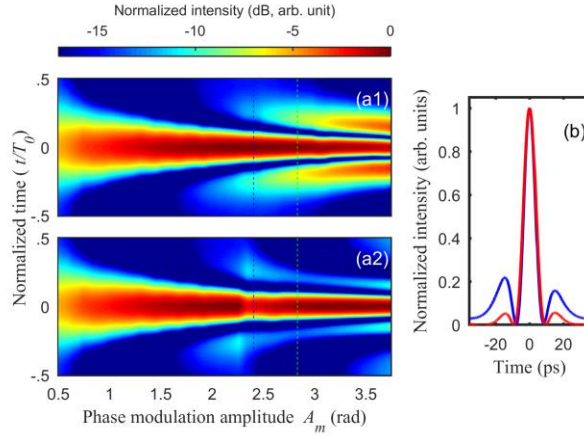


Figure 10: (a) Experimental evolution of the temporal intensity profile according to the initial amplitude of the sinusoidal phase modulation. Results for ψ_B and ψ'_B are plotted on panels (a1) and (a2) respectively. White dotted line corresponds to $A_m = 2.83$ rad whereas the black dotted line is for $A_m = A_{m,1}$. (b) Details of the intensity profile for $A_m = 2.83$ rad

The excellent agreement between theory and experiment is also visible in Fig. 11 where we have summarized the evolution of the peak-power and the fwhm duration. All the various features we previously predicted are reproduced, thus confirming our modelling.

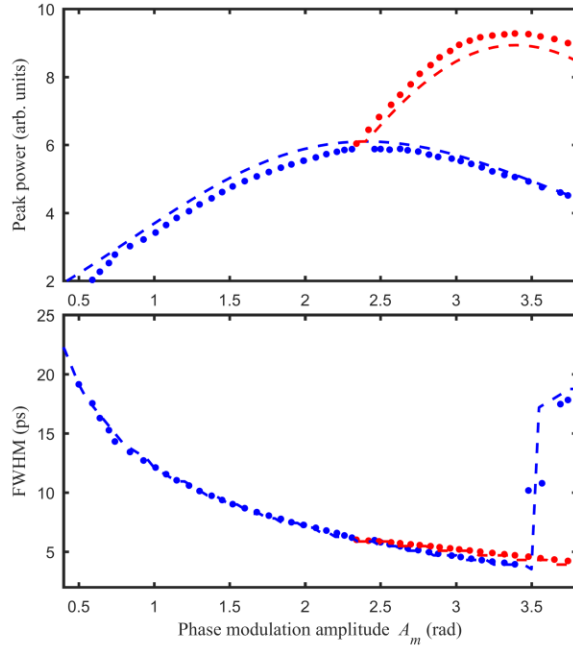


Figure 11: Evolution of (a) the peak-power and (b) the fwhm duration of the pulse train according to the amplitude of modulation A_m . The experimental results (circles) are compared with the results of Eq. (2) and (17) or with numerical simulations (dashed lines). Results obtained with the besselon ψ_B and the modified besselon ψ'_B are plotted with blue and red colors, respectively.

C. Doubling of the repetition rate

We finally experimentally investigate the doubling of the repetition rate. Resulting pulse trains obtained at 28 GHz are reported in Fig. 12 for $A_m = \pi/2$ and $A_m = 3.72$ rad. The optical spectra are displayed in panels (2) where we can note the significant decrease of the odd spectral components. The π phase shifts applied on those harmonics induces a reduction by more than 20 dB of their power. The odd spectral components are not fully suppressed but their very low level no longer influences the overall stability of the resulting train as it is confirmed by the eye diagrams. In both cases, the experimental results are in excellent agreement with the theoretical predictions from Eq. (20). The extinction ratio is excellent. The level of the pedestals is very low and for $A_m = 3.72$ rad, a fwhm duration of 6.9 ps is achieved, leading to a duty-cycle of 0.19.

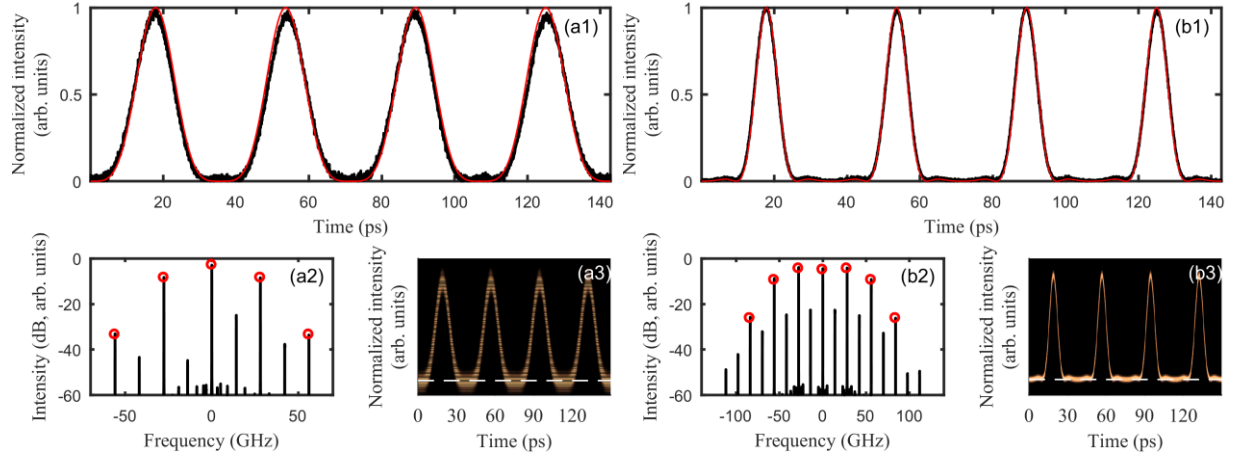


Figure 12: Pulse trains generated at a repetition rate of 28 GHz obtained for (a) $A_m = \pi/2$ rad and (b) $A_m = 3.72$ rad. (panel 1) Temporal power profile. (panel 2) Optical power spectrum. (panel 3) Temporal profile recorded in persistent mode. Experimental results (black line) are compared with theoretical results (red line, Eq. (20); circles). The horizontal white dashed lines indicate the zero-level.

IV. Conclusions

In conclusion, we have theoretically introduced a new waveform we have named *besselon* for which we have provided accurate and simple analytical predictions of the main features. We have stressed the significant differences between the *besselon* and the waveform that arises from quadratic phase compensation of a sinusoidally phase modulated wave. For a low level of initial phase modulation, an approximate model has been proposed to interpret the *besselon* shape as the result of an interference between three waves. For high levels of initial phase modulation, the single pulse structure then splits into several peaks. By inserting an additional π -phase shift on the spectral component, the peak power can be further increased and the temporal duration reduced, making it attractive for high-repetition sources as well as for temporal multiplexing. Duty cycles as low as 0.1 without residual background are demonstrated without involving another complex stage of fiber-based nonlinear compression [36, 37]. By carefully choosing the level of initial modulation, we have shown that efficient repetition-rate doubling can be achieved with nice temporal profiles free from residual background. Compared to existing solutions that have relied on the use of an additional intensity modulator tightly synchronized with the phase modulation [38, 39] or on the use of a nonlinear optical loop mirror [40], our proposed method only implies a single phase modulator. Our approach could also be of interest for noise-free amplification [41] as well as for optical sampling [42]. In the present contribution, we have mainly focused on the generation of the pulse waveform. Further studies will explore the propagation in linear media (where it will be subject to Talbot self-imaging [24]) or in nonlinear media such as anomalous dispersive fibers where the spectrum with many components could be of extreme interest to seed new nonlinear regimes.

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References

- [1] S.T. Cundiff, A.M. Weiner, Optical arbitrary waveform generation, *Nat. Photonics*, 4 (2010) 760-766.
- [2] F. Parmigiani, P. Petropoulos, T.T. Ng, M. Ibsen, D.J. Richardson, *Applications of Superstructured Fibre Bragg Gratings in All-Optical Signal Processing*, Ieee, New York, 2009.
- [3] C. Finot, I. Gukov, K. Hammani, S. Boscolo, Nonlinear sculpturing of optical pulses with normally dispersive fiber-based devices, *Opt. Fiber Technol.*, 45 (2018) 306-312.
- [4] S. Boscolo, C. Finot, *Shaping Light in Nonlinear Optical Fibers*, 2017.
- [5] K. Hammani, B. Kibler, C. Finot, P. Morin, J. Fatome, J.M. Dudley, G. Millot, Peregrine soliton generation and breakup in standard telecommunications fiber, *Opt. Lett.*, 36 (2011) 112-114.
- [6] B. Kibler, A. Chabchoub, A. Gelash, N. Akhmediev, V.E. Zakharov, Superregular Breathers in Optics and Hydrodynamics: Omnipresent Modulation Instability beyond Simple Periodicity, *Phys. Rev. X*, 5 (2015) 041026.
- [7] B. Wetzell, D. Bongiovanni, M. Kues, Y. Hu, Z. Chen, S. Trillo, J.M. Dudley, S. Wabnitz, R. Morandotti, Experimental Generation of Riemann Waves in Optics: A Route to Shock Wave Control, *Phys. Rev. Lett.*, 117 (2016) 073902.
- [8] B. Varlot, S. Wabnitz, J. Fatome, G. Millot, C. Finot, Experimental generation of optical flaticon pulses, *Opt. Lett.*, 38 (2013) 3899-3902.
- [9] R. Driben, Y. Hu, Z. Chen, B.A. Malomed, R. Morandotti, Inversion and tight focusing of Airy pulses under the action of third-order dispersion, *Opt. Lett.*, 38 (2013) 2499-2501.
- [10] R. Slavík, Y. Park, M. Kulishov, R. Morandotti, J. Azaña, Ultrafast all-optical differentiators, *Opt. Express*, 14 (2006) 10699-10707.

- [11] Z. Jiang, D.E. Leaird, A.M. Weiner, Optical processing based on spectral line-by-line pulse shaping on a phase-modulated CW laser, *IEEE J. Quantum Electron.*, 42 (2006) 657-665.
- [12] T. Kobayashi, H. Yao, K. Amano, Y. Fukushima, A. Morimoto, T. Sueta, Optical pulse compression using high-frequency electrooptic phase modulation, *IEEE J. Quantum Electron.*, 24 (1988) 382-387.
- [13] T. Komukai, Y. Yamamoto, S. Kawanishi, Optical pulse generator using phase modulator and linearly chirped fiber Bragg gratings, *IEEE Photon. Technol. Lett.*, 17 (2005) 1746-1748.
- [14] V. Torres-Company, J. Lancis, P. Andrés, Unified approach to describe optical pulse generation by propagation of periodically phase-modulated CW laser light, *Opt. Express*, 14 (2006) 3171-3180.
- [15] U. Andral, J. Fatome, B. Kibler, C. Finot, Triangular spectral phase tailoring for the generation of high-quality picosecond pulse trains, *Opt. Lett.*, 44 (2019) 4913.
- [16] U. Andral, C. Finot, Multiwavelength high-repetition rate source, *Laser Phys.*, 30 (2020) 016203.
- [17] U. Andral, C. Finot, High-repetition-rate source delivering optical pulse trains with a controllable level of amplitude and temporal jitters, submitted, [arXiv:1910.11739](https://arxiv.org/abs/1910.11739).
- [18] M. Abramowitz, I.A. Stegun, *Handbook of mathematical functions: with formulas, graphs, and mathematical tables*, Courier Corporation 1964.
- [19] C. Finot, F. Chaussard, S. Boscolo, Impact of a temporal sinusoidal phase modulation on the optical spectrum, *Eur. J. Phys.*, 39 (2018) 055303.
- [20] K. Hammani, J. Fatome, C. Finot, Applications of sinusoidal phase modulation in temporal optics to highlight some properties of the Fourier transform, *Eur. J. Phys.*, 40 (2019) 055301.
- [21] J. Spanier, K.B. Oldham, *An atlas of functions*, Hemisphere Publishing Corporation 1987.

- [22] U. Andral, B. Kibler, J.M. Dudley, C. Finot, Akhmediev breather signatures from dispersive propagation of a periodically phase-modulated continuous wave, *Wave Motion*, 95 (2020) 1025-1045.
- [23] N.N. Akhmediev, V.I. Korneev, Modulation instability and periodic-solutions of the nonlinear Schrödinger equation, *Theor. Math. Phys.*, 69 (1986) 1089-1093.
- [24] L. Romero Cortés, R. Maram, H. Guillet de Chatellus, J. Azaña, Arbitrary Energy-Preserving Control of Optical Pulse Trains and Frequency Combs through Generalized Talbot Effects, *Laser Photonics Rev.*, 13 (2019) 1900176.
- [25] J. Fatome, S. Pitois, C. Fortier, B. Kibler, C. Finot, G. Millot, C. Courde, M. Lintz, E. Samain, Multiple four-wave mixing in optical fibers : 1.5-3.4-THz femtosecond pulse sources and real-time monitoring of a 20-GHz picosecond source, *Opt. Commun.*, 283 (2010) 2425-2429.
- [26] J. Fatome, B. Kibler, C. Finot, High-quality optical pulse train generator based on solitons on finite background, *Opt. Lett.*, 38 (2013) 1663-1665.
- [27] P. Petropoulos, M. Ibsen, M.N. Zervas, D.J. Richardson, Generation of a 40-GHz pulse stream by pulse multiplication with a sampled fiber Bragg grating, *Opt. Lett.*, 25 (2000) 521-523.
- [28] J. Caraquitená, Z. Jiang, D.E. Leaird, A.M. Weiner, Tunable pulse repetition-rate multiplication using phase-only line-by-line pulse shaping, *Opt. Lett.*, 32 (2007) 716-718.
- [29] A.A. Godil, B.A. Auld, D.M. Bloom, Picosecond time-lenses, *IEEE J. Quantum Electron.*, 30 (1994) 827-837.
- [30] T. Khayim, M. Yamauchi, D.-S. Kim, T. Kobayashi, Femtosecond optical pulse generation from a CW laser using an electrooptic phase modulator featuring lens modulation, *IEEE J. Quantum Electron.*, 35 (1999) 1412-1418.

- [31] M.A.F. Roelens, S. Frisken, J. Bolger, D. Abakounov, G. Baxter, S. Poole, B.J. Eggleton, Dispersion trimming in a reconfigurable wavelength selective switch, *J. Lightw. Technol.*, 26 (2008) 73-78.
- [32] N.Q. Ngo, Y. Song, On the interrelations between an optical differentiator and an optical Hilbert transformer, *Opt. Lett.*, 36 (2011) 915-917.
- [33] N.K. Berger, B. Levit, B. Fischer, Reshaping periodic light pulses using cascaded uniform fiber Bragg gratings, *J. Lightw. Technol.*, 24 (2006) 2746-2751.
- [34] C. Yu, Y. Wang, Z. Pan, T. Luo, S. Kumar, B. Zhang, A.E. Willner, Carrier-suppressed 160 GHz pulse-train generation using a 40 GHz phase modulator with polarization-maintaining fiber, *Opt. Lett.*, 34 (2009) 1657-1659.
- [35] T. Sizer, Increase in laser repetition rate by spectral selection, *IEEE J. Quantum Electron.*, 25 (1989) 97-103.
- [36] T. Inoue, S. Namiki, Pulse compression techniques using highly nonlinear fibers, *Laser Photonics Rev.*, 2 (2008) 83-99.
- [37] C. Finot, J. Fatome, S. Pitois, G. Millot, All-Fibered High-Quality Low Duty-Cycle 20-GHz and 40-GHz Picosecond Pulse Sources, *IEEE Photon. Technol. Lett.*, 19 (2007) 1711-1713.
- [38] Y. Ji, Y. Li, J. Wu, F.Z. Zhang, K. Xu, W. Li, X.B. Hong, J.T. Lin, A Phase Stable Short Pulses Generator Using an EAM and Phase Modulators for Application in 160-GBaud DQPSK Systems, *IEEE Photon. Technol. Lett.*, 24 (2012) 64-66.
- [39] T. Otsuji, M. Yaita, T. Nagatsuma, E. Sano, 10-80-Gb/s highly extinctive electrooptic pulse pattern generation, *IEEE J. Sel. Top. Quantum Electron.*, 2 (1996) 643-649.
- [40] S. Yang, X. Bao, Generating a high-extinction-ratio pulse from a phase-modulated optical signal with a dispersion-imbalanced nonlinear loop mirror, *Opt. Lett.*, 31 (2006) 1032-1034.

[41] R. Maram, J. Van Howe, M.J. Li, J. Azaña, Noiseless intensity amplification of repetitive signals by coherent addition using the temporal Talbot effect, *Nat Commun*, 5 (2014).

[42] J. Nuno, C. Finot, J. Fatome, Linear Sampling and Magnification Technique Based on Phase Modulators and Dispersive Elements: the Temporal Lenticular Lens, *Opt. Fiber Technol.*, 36 (2017) 125-129.