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# A MIXED FINITE ELEMENT DISCRETIZATION OF DYNAMICAL OPTIMAL TRANSPORT 

ANDREA NATALE AND GABRIELE TODESCHI


#### Abstract

In this paper we introduce a new class of finite element discretizations of the quadratic optimal transport problem based on its dynamical formulation. These generalize to the finite element setting the finite difference scheme proposed by Papadakis et al. [SIAM J Imaging Sci, 7(1):212-238, 2014]. We solve the discrete problem using a proximal-splitting approach and we show how to modify this in the presence of regularization terms which are relevant for imaging applications.


## 1. Introduction

Optimal transport provides a convenient framework for density interpolation as a convex optimization problem. Its most remarkable feature is its sensitivity to horizontal displacement, which generally allows one to retrieve translations when interpolating between two densities. This property has motivated the application of optimal transport to many imaging problems, especially in physical sciences and fluid dynamics. A typical example comes from satellite image interpolation in oceanography. In this case, one is interested in reconstructing the evolution of a quantity of interest such as Sea Surface Temperature (SST) or Sea Surface Height (SSH) between two given observations. As highlighted in [17], for this type of applications one needs to include appropriate regularization terms to avoid the appearance of unphysical phenomena such as mass concentration in the reconstructed density evolution.

In this paper we propose a finite element approach to solve the dynamical formulation of optimal transport which is numerically stable on unstructured meshes and that can be easily modified to include different type of regularizations which are relevant for image interpolation problems. For some choices of finite element spaces, using the framework introduced in [19], we can prove convergence of our discrete solutions to the ones of the continuous problem.

The dynamical formulation of optimal transport inspired the first numerical methods for this problem. This reads as follows: given two probability measures $\rho_{0}, \rho_{1} \in \mathcal{P}(D)$ on a compact domain $D \subset \mathbb{R}^{d}$, find the curve $t \in[0,1] \mapsto \rho(t, \cdot) \in$ $\mathcal{P}(D)$

$$
\begin{equation*}
\inf _{\rho, v}\left\{\int_{0}^{1} \int_{D} \frac{|v(t, \cdot)|^{2}}{2} \mathrm{~d} \rho(t, \cdot) \mathrm{d} t ; \partial_{t} \rho+\operatorname{div}_{x}(\rho v)=0, \rho(0)=\rho_{0}, \rho(1)=\rho_{1}\right\} \tag{1.1}
\end{equation*}
$$

where $v:[0,1] \times D \rightarrow \mathbb{R}^{d}$ is a time-dependent velocity field on $D$ tangent to the boundary $\partial D$, and $|\cdot|$ denotes the Euclidean norm. In other words, problem (1.1) selects the curve of minimal kinetic energy with fixed endpoints $\rho_{0}$ and $\rho_{1}$.

Benamou and Brenier [5] realized that introducing the momentum $m:=\rho v$, problem (1.1) can be recast into a convex optimization problem in the variables $(\rho, m)$, with a linear constraint, since the continuity equation becomes

$$
\begin{equation*}
\partial_{t} \rho+\operatorname{div}_{x} m=0 . \tag{1.2}
\end{equation*}
$$

[^0]If we define $\sigma:=(\rho, m)$, regarded as a measure on $[0,1] \times D$, this constraint is equivalent to $\operatorname{div} \sigma=0$, where now div denotes the divergence operator on the space-time domain $[0,1] \times D$. Introducing the dual variable $q=(a, b)$ wher $a \in$ $L^{\infty}([0,1] \times D)$ and $b \in L^{\infty}\left([0,1] \times D ; \mathbb{R}^{d}\right)$, the kinetic energy minimized in (1.1) can be written in the form

$$
\begin{equation*}
\sup _{q}\left\{\int_{0}^{1} \int_{D} q \cdot \mathrm{~d} \sigma ; a+\frac{|b|^{2}}{2} \leq 0\right\} . \tag{1.3}
\end{equation*}
$$

Combining this expression with (1.1) we obtain a saddle point problem in the variables $(q, \sigma)$ with a nonlinear constraint on $q$ and a linear one on $\sigma$. The numerical method proposed in [5] involves discretizing $q$ and $\sigma$ by their values on a regular grid, and expressing the constraint on $\sigma$ via a Lagrange multiplier; then the method can be solved by an augmented Lagrangian ADMM approach, optimizing separately in $\sigma$ and $q$ and then performing a gradient ascent step on the Lagrange multiplier. Disregarding the discretization in space-time, the convergence of the method has been studied in $[15,18]$. The same approach was used to discretize different problems related to optimal transport (e.g., gradient flows [8], mean field games [7], unbalanced optimal transport [13]) using a finite element discretization in spacetime: in these cases the numerical method is obtained by discretizing the several steps of the augmented Lagrangian approach rather than as a discrete optimization algorithm.

In [22], using a staggered finite difference approach on a regular grid, the authors considered the discrete optimization problem arising from the discretization of (1.1), and they considered different proximal splitting algorithms to solve it. The computational bottleneck for these methods as well as for the original augmented Lagrangian approach is the solution of a Laplace equation which appears when computing the proximal operator associated with the constraint $\operatorname{div} \sigma=0$. This however can be avoided by exploiting the Helmholtz decomposition of vector fields, as recently showed in [16].

In [20] a numerical scheme was proposed using tools from finite element and finite volume methods, where one explicitly constructs a duality structure for the discrete variables. Later Lavenant [19] proved convergence of this scheme, unconditionally with respect to the time/space step size, to the solutions of the optimal transport problem, proposing a general framework for convergence of discretizations of problem (1.1) between two arbitrary probability measures. This filled a critical gap for the analysis of discrete dynamical transport models, since previously convergence results were only known in the Hilbert space setting or conditional to the relative time/space step sizes.

In this paper we leverage this result to propose a mixed finite element discretization of (1.1) which generalizes to unstructured meshes the scheme proposed in [22] and that can also be proved to converge using the framework of [19]. The paper is structured as follows. We establish the notation in section 2. In section 3 we give the precise formulation of problem (1.1) and describe the proximal splitting algorithm for the continuous problem in the Hilbert space setting. In section 4 we describe the finite element discretization and the major tools we use in our method. In section 5 we describe the discrete algorithm for different choices of finite element spaces and state the convergence result. In section 7 we describe how to introduce regularization terms in the formulation. Finally in section 8 we present some numerical results.

## 2. Notation

Throughout the paper we will denote by $D \subset \mathbb{R}^{d}$ a compact domain with Lipschitz boundary and by $\Omega:=[0,1] \times D$ the space-time domain. We use $\langle\cdot, \cdot\rangle$ to denote the $L^{2}$ inner product on either $D$ or $\Omega$. For differential operator such as $\nabla$ or div, we use the subscript $x$ to highlight that these are defined on $D$ rather than $\Omega$, but we will drop this subscript when this is clear from the context.

We use the standard notation for Sobolev spaces on $D$ or $\Omega$. In particular, $L^{p}\left(D ; \mathbb{R}^{d}\right)$ denotes the space of functions $f: D \rightarrow \mathbb{R}^{d}$ whose Euclidean norm $|f|$ is in $L^{p}(D)$. We use a similar notation for functions taking values on a subset $K \subset \mathbb{R}^{d}$, or defined on $\Omega$. We denote by $H($ div; $D)$ the space of vector fields $f: D \rightarrow \mathbb{R}^{d}$ in $L^{2}\left(D ; \mathbb{R}^{d}\right)$ whose divergence is in $L^{2}(D)$. Similarly, $H(\operatorname{div} ; \Omega)$ the space of vector fields $f: \Omega \rightarrow \mathbb{R}^{d+1}$ in $L^{2}\left(\Omega ; \mathbb{R}^{d+1}\right)$ whose divergence is in $L^{2}(\Omega)$.

Finally, we denote by $\mathcal{M}(D)$ the set of finite signed measures on $D$, by $\mathcal{M}_{+}(D) \subset$ $\mathcal{M}(D)$ the convex subset of positive measures, and by $\mathcal{P}(D) \subset \mathcal{M}_{+}(D)$ the set of positive measures of total mass equal to one. We use a similar notation for the spaces of measures on $\Omega$.

## 3. DYnamical formulation of optimal transport

The Benamou-Brenier formula allows one to formulate the optimal transport problem (1.1) as a saddle point problem on the space of measures $\sigma:=(\rho, m) \in$ $\mathcal{M}(\Omega) \times \mathcal{M}(\Omega)^{d}$. This can be written as follows

$$
\begin{equation*}
\inf _{\sigma \in \mathcal{C}} \mathcal{A}(\sigma), \quad \mathcal{A}(\sigma):=\sup _{q \in L^{\infty}(\Omega ; K)}\langle q, \sigma\rangle \tag{3.1}
\end{equation*}
$$

where $\mathcal{C}$ is the set of measures $\sigma \in \mathcal{M}(\Omega)^{d+1}$ satisfying $\operatorname{div} \sigma=0$ in distributional sense with boundary conditions

$$
\sigma \cdot n_{\partial \Omega}=\mathcal{X}, \quad \mathcal{X}:= \begin{cases}\rho_{0} & \text { on }\{0\} \times D  \tag{3.2}\\ \rho_{1} & \text { on }\{1\} \times D \\ 0 & \text { otherwise }\end{cases}
$$

with $\rho_{0}, \rho_{1} \in \mathcal{P}(D)$, and where

$$
\begin{equation*}
K:=\left\{(a, b) \in \mathbb{R} \times \mathbb{R}^{d} ; a+\frac{|b|^{2}}{2} \leq 0\right\} . \tag{3.3}
\end{equation*}
$$

It will be convenient to treat time and space as separate variables. In particular we will also use the action defined by

$$
\begin{equation*}
A(\rho, m):=\sup _{(a, b) \in L^{\infty}(D ; K)}\langle\rho, a\rangle+\langle m, b\rangle, \tag{3.4}
\end{equation*}
$$

for any $(\rho, m) \in \mathcal{M}(D)^{d+1}$. Then, $A(\rho, m)$ is finite if and only if $m$ has a density with respect to $\rho$ and in that case $A(\rho, m)=\int_{D} B(\rho, m)$, where $B: \mathbb{R} \times \mathbb{R}^{d} \rightarrow$ $[0,+\infty]$ is the function given by

$$
B(a, b):= \begin{cases}\frac{|b|^{2}}{2 a} & \text { if } a>0  \tag{3.5}\\ 0 & \text { if } a=0, b=0 \\ +\infty & \text { if } a=0, b \neq 0 \text { or } a<0\end{cases}
$$

Due to the definition of the function (3.5), a saddle point of problem (3.1) must satisfy $\rho \geq 0$.
3.1. Hilbert space setting and proximal splitting. Before discussing the discretization of problem (3.1), we review its reformulation on Hilbert spaces, and discuss the convergence of the proximal splitting algorithm.
Proposition 3.1 (Guittet [15]; Hug et al. [18]). Suppose $\rho_{0}, \rho_{1} \in L^{2}(D)$. Then problem (3.1) is equivalent to

$$
\begin{equation*}
\inf _{\sigma \in \mathcal{C}} \sup _{q \in L^{2}(\Omega ; K)}\langle q, \sigma\rangle \tag{3.6}
\end{equation*}
$$

where $\mathcal{C}$ is the set of functions $\sigma \in H(\operatorname{div} ; \Omega)$ satisfying $\operatorname{div} \sigma=0$ in weak sense with boundary conditions given by (3.2). Moreover, there exists a saddle point $\left(\sigma^{*}, q^{*}\right) \in \mathcal{C} \times L^{2}(\Omega ; K)$ solving problem (3.7).

The equivalence of problem (3.7) to (3.1) can be easily deduced by a regularization argument on $\sigma$ and then applying Lusin's theorem as in Proposition 5.18 in [25]. The proof for the existence of a saddle point problem is more delicate and can be found in [18].

In order to apply a proximal splitting algorithm to solve problem (3.7), we first write it in the form

$$
\begin{equation*}
\inf _{\sigma \in L^{2}\left(\Omega ; \mathbb{R}^{d+1}\right)} \sup _{q \in L^{2}\left(\Omega ; \mathbb{R}^{d+1}\right)}\langle q, \sigma\rangle+\iota_{\mathcal{C}}(\sigma)-\iota_{\mathcal{K}}(q) \tag{3.7}
\end{equation*}
$$

where $\iota$ denotes the convex indicator function and

$$
\begin{equation*}
\mathcal{K}:=L^{2}(\Omega ; K)=\left\{q \in L^{2}\left(\Omega ; \mathbb{R}^{d+1}\right) ; q \in K \text { a.e. }\right\} . \tag{3.8}
\end{equation*}
$$

Note in particular that $\mathcal{C}$ and $\mathcal{K}$ are closed convex sets of $L^{2}$.
We apply to (3.7) the primal-dual projection algorithm proposed in [23]. In particular, given $\tau_{1}, \tau_{2}>0$ and an admissible $\left(\sigma^{0}, q^{0}\right) \in \mathcal{C} \times \mathcal{K}$, we define the sequence $\left\{\left(\sigma^{k}, q^{k}\right)\right\}_{k}$ by the two-step algorithm:

$$
\begin{array}{ll}
\text { Step 1: } & \sigma^{k+1}=P_{\mathcal{C}}\left(\sigma^{k}-\tau_{1} q^{k}\right) \\
\text { Step 2 : } & q^{k+1}=P_{\mathcal{K}}\left(q^{k}+\tau_{2}\left(2 \sigma^{k+1}-\sigma^{k}\right)\right) \tag{3.9b}
\end{array}
$$

where $P_{\mathcal{C}}$ and $P_{\mathcal{K}}$ are the $L^{2}$ projections on the closed convex sets $\mathcal{C}$ and $\mathcal{K}$, respectively. The projection onto $\mathcal{C}$ amounts to computing the Helmholtz decomposition of $\sigma^{k}-\tau_{1} q^{k}$ and selecting the divergence-free part, whereas the projection onto $\mathcal{K}$ is a pointwise projection on a representative of $q^{k}+\tau_{2}\left(2 \sigma^{k+1}-\sigma^{k}\right)$.

The proof of convergence in [23] holds also in our setting. More precisely, the following convergence theorem holds.
Theorem 3.2 (Pock et al. [23]). If $\tau_{1} \tau_{2}<1$ then $\left(\sigma^{k}, q^{k}\right) \rightarrow\left(\sigma^{*}, q^{*}\right) \in \mathcal{C} \times \mathcal{K}$ which solves (3.7).

Discretizing problem (3.7), and consequently the proximal splitting algorithm (3.9), with finite elements requires choosing finite-dimensional spaces for $\sigma$ and $q$ so that the steps in (3.9) are well-posed and solvable. However, satisfying these requirements is not enough to guarantee convergence of the discrete solutions to the ones of the infinite dimensional problem. Hereafter we will identify a class of finite element spaces for which the theory developed in [19] applies, which allows us to deduce convergence in the setting of problem (3.1), i.e. even when $\rho_{0}$ and $\rho_{1}$ are arbitrary probability measures.

## 4. Mixed finite element setting

4.1. Finite element spaces on $D$. For simplicity, from now on we will assume that $D$ is a polygonal domain in $\mathbb{R}^{d}$. We consider a triangulation of $D$ which we denote $\mathcal{T}_{h}$, i.e. a decomposition of $D$ in either simplicial or quadrilateral elements,
where $h$ is the maximum diameter of the elements in $\mathcal{T}_{h}$. We assume that there exists a constant $C_{\text {mesh }}$ such that

$$
\begin{equation*}
|h|^{d} \leq C_{m e s h}|T|, \quad \forall T \in \mathcal{T}_{h} . \tag{4.1}
\end{equation*}
$$

This implies that the mesh is quasiuniform, meaning that the ratio of any two element diameters is uniformly bounded by a constant depending only on $C_{m e s h}$, and shape-regular, that is, for each element $T \in \mathcal{T}_{h}$, the ratio of its diameter and the diameter of the largest inscribed ball is uniformly bounded by a constant depending only on $C_{\text {mesh }}$ (see, e.g., [2]).

For any $T \in \mathcal{T}_{h}$, we denote by $\mathcal{P}_{k}(T)$ the space of polynomials of degree up to $k$ on $T$. If $T$ is a quadrilateral element, i.e., in general, if $T$ is obtained by an affine transformation $\phi: I^{d} \rightarrow T$ where $I$ is the unit interval, then we define $\mathcal{P}_{k_{1}, \ldots k_{d}}\left(I^{d}\right):=\mathcal{P}_{k_{1}}(I) \otimes \ldots \otimes \mathcal{P}_{k_{d}}(I)$ and $\mathcal{P}_{k_{1}, \ldots k_{d}}(T):=\mathcal{P}_{k_{1}, \ldots k_{d}}\left(I^{d}\right) \circ \phi^{-1}$.

We now define the finite elment spaces $Q_{h}$ and $V_{h}$ which will serve to construct approximations of the denisty $\rho$ and the momentum $m$, respectively. We set

$$
\begin{gather*}
Q_{h}:=\left\{\phi \in L^{2}(D) ;\left.\phi\right|_{T} \in \mathcal{P}_{0}(T), \forall T \in \mathcal{T}_{h}\right\}  \tag{4.2}\\
V_{h}:=\left\{v \in H(\operatorname{div} ; D) ;\left.v\right|_{T} \in V_{h}(T), \forall T \in \mathcal{T}_{h}\right\} . \tag{4.3}
\end{gather*}
$$

where $V_{h}(T)$ is the so-called shape function space. We distinguish two cases:
(1) for simplicial elements (triangles or tetrahedrons), we take $V_{h}(T)$ to be either
$\mathcal{R} \mathcal{T}_{0}(T):=\left\{v=v_{0}+v_{1} \hat{x} ; v_{0} \in\left(\mathcal{P}_{0}(T)\right)^{d}, v_{1} \in \mathcal{P}_{0}(T)\right\} \subset\left(\mathcal{P}_{1}(T)\right)^{d}$,
where $\hat{x}=\left(x_{1}, \ldots, x_{d}\right) \in\left(\mathcal{P}_{1}(T)\right)^{d}$, which generates the lowest order Raviart-Thomas space; or $\mathcal{B D} \mathcal{M}_{1}(T)=\left(\mathcal{P}_{1}(T)\right)^{d}$, which generates the lowest order Brezzi-Douglas-Marini $H$ (div)-conforming space;
(2) for quadrilateral elements, we set $T=\phi\left(I^{d}\right)$, where $I$ is an interval and $\phi$ an affine transformation, and we take $V_{h}(T)$ to be the tensor product space which generates the lowest order Raviart-Thomas space on quadrilateral elements. This is defined as follows:

$$
\mathcal{R} \mathcal{T}_{[0]}(T):= \begin{cases}\mathcal{P}_{1,0}(T) e_{1}+\mathcal{P}_{0,1}(T) e_{2} & \text { if } d=2 \\ \mathcal{P}_{1,0,0}(T) e_{1}+\mathcal{P}_{0,1,0}(T) e_{2}+\mathcal{P}_{0,0,1}(T) e_{3} & \text { if } d=3\end{cases}
$$

where $\left\{e_{i}\right\}_{i}$ is the basis for $\mathbb{R}^{d}$ aligned with the edges of $T$.
In practice, the space $V_{h}$ is chosen among the standard lowest order $H$ (div)conforming spaces. In fact, the property of being piece-wise linear will be crucial in the following, namely to prove the convergence result (5.3) (see, in particular, proposition A. 2 in the appendix). A graphical representation of the degrees of freedom associated with these spaces is shown in figure 1.

Importantly, with the choices mentioned above, there exist bounded projections $\Pi_{Q_{h}}: L^{2}(D) \rightarrow Q_{h}$ and $\Pi_{V_{h}}: H(\operatorname{div} ; D) \rightarrow V_{h}$ that commute with the divergence operator [ 2,10 ]. In other words, the following diagram commutes


As a consequence, the divergence operator is surjective onto $Q_{h}$ when restricted on $V_{h}$, i.e. $\operatorname{div} V_{h}=Q_{h}$. Finally, we let $Q_{h}^{+} \subset Q_{h}$ the convex subset of non-negative piecewise constant functions.


Figure 1. Degrees of freedom for different choices of shape function space $V_{h}(T)$

Remark 4.1. The precise definition of the projection operator can be found in section 5.4 in [2] or in [10]. Here, we will only need the explicit definition of $\Pi_{Q_{h}}$, which is given by

$$
\begin{equation*}
\left.\Pi_{Q_{h}} \rho\right|_{T}=\frac{1}{|T|} \int_{T} \rho \tag{4.5}
\end{equation*}
$$

for any $T \in \mathcal{T}_{h}$. Therefore, $\Pi_{Q_{h}}$ is well-defined on $\mathcal{M}(D)$ and its resctriction on $\mathcal{M}_{+}(D)$ is surjective onto $Q_{h}^{+}$.
4.2. Finite element spaces on $\Omega$. We now introduce finite element spaces on the space-time domain $[0,1] \times D$. We first define a decomposition $\mathcal{T}_{h, \tau}$, obtained by a tensor product construction. In other words, we assume that $\mathcal{T}_{h, \tau}$ is obtained by tensor product of a triangulation $\mathcal{T}_{h}$ of $D$ and a decocomposition of $[0,1]$ of maximum size $\tau$, so that any element $S \in \mathcal{T}_{h, \tau}$ is of the form $S=\left[t_{0}, t_{1}\right] \otimes T$ where $T \in \mathcal{T}_{h}$.

We now define the finite element spaces $F_{h, \tau}$ and $Z_{h, \tau}$ on the space-time domain. The space $Z_{h, \tau}$ will be constructed using the standard tensor product construction based on the spaces $Q_{h}$ and $V_{h}$ defined on $D$, and continuous $\mathcal{P}_{1}$ and discontinuous $\mathcal{P}_{0}$ spaces on $[0,1]$. In our discretization, the space-time vector field $(\rho, m)$ will be an element of $Z_{h, \tau}$ whereas $F_{h, \tau}$ will be the space of discrete Lagrange multipliers associated to the continuity equation, which is just the constraint that the spacetime divergence of $(\rho, m)$ is zero.

More precisely, we define

$$
\begin{gather*}
F_{h, \tau}:=\left\{\phi \in L^{2}(\Omega) ;\left.\phi\right|_{S} \in \mathcal{P}_{0}(S), \forall S \in \mathcal{T}_{h, \tau}\right\},  \tag{4.6}\\
Z_{h, \tau}:=\left\{v \in H(\operatorname{div} ; \Omega) ;\left.v\right|_{S} \in Z_{h, \tau}(S), \forall S \in \mathcal{T}_{h, \tau}\right\} \tag{4.7}
\end{gather*}
$$

For $S=\left[t_{0}, t_{1}\right] \otimes T$, the shape function space $Z_{h, \tau}(S)$ is built by defining a shape function space for the density, in the space-time domain, which is given by

$$
\begin{equation*}
Q_{h, \tau}(S):=\mathcal{P}_{1}\left(\left[t_{0}, t_{1}\right]\right) \otimes Q_{h}(T) \tag{4.8}
\end{equation*}
$$

(i.e. the density is piecewise linear in time), and a shape function space for the momentum, in the space-time domain, which is given by

$$
\begin{equation*}
V_{h, \tau}(S):=\mathcal{P}_{0}\left(\left[t_{0}, t_{1}\right]\right) \otimes V_{h}(T) \tag{4.9}
\end{equation*}
$$

(i.e. the momentum is piecewise constant in time). Then, we set

$$
\begin{equation*}
Z_{h, \tau}(S):=V_{h, \tau}(S) \oplus\left(Q_{h, \tau}(S) \hat{t}\right) \tag{4.10}
\end{equation*}
$$

where $\hat{t}$ is the unit vector oriented in the time direction. The spaces $F_{h, \tau}$ and $Z_{h, \tau}$ inherit from $Q_{h}$ and $V_{h}$ the commuting diagram property mentioned above.

In particular, there exist bounded projections $\Pi_{F_{h, \tau}}: L^{2}(\Omega) \rightarrow F_{h, \tau}$ and $\Pi_{Z_{h, \tau}}$ : $H($ div $; \Omega) \rightarrow Z_{h, \tau}$ for which the following diagram commutes

where the divergence is the one associated with the space-time domain $\Omega$. Then, as before, the divergence operator is surjective onto $F_{h, \tau}$ when restricted on $Z_{h, \tau}$, i.e. div $Z_{h, \tau}=F_{h, \tau}$. Note that the precise definition for the projection operators on tensor product meshes can be found in [1].
4.3. Discrete projection on the divergence-free subspace. Denote by $\mathcal{B}$ the kernel of the divergence operator on $H(\operatorname{div} ; \Omega)$. Given $\xi \in L^{2}(\Omega)$ we define the projection $P_{\mathcal{B}}(\xi)$ to be the divergence-free vector field $\sigma$ minimizing the $L^{2}$ distance from $\xi$. This can be obtained solving the following problem for $(\sigma, \phi) \in H(\operatorname{div} ; \Omega) \times$ $L^{2}(\Omega)$

$$
\begin{cases}\langle\sigma, v\rangle+\langle\phi, \operatorname{div} v\rangle=\langle\xi, v\rangle, & \forall v \in H(\operatorname{div} ; \Omega)  \tag{4.12}\\ \langle\operatorname{div} \sigma, \psi\rangle=0, & \forall \psi \in L^{2}(\Omega)\end{cases}
$$

The system is well-posed thanks to the surjectivity of the divergence operator onto $L^{2}(\Omega)$. Let $\mathcal{B}_{h, \tau}$ be the kernel of the divergence operator restricted on $Z_{h, \tau}$. We define the projection $P_{\mathcal{B}_{h, \tau}}(\xi)$ to be the divergence-free vector field $\sigma_{h, \tau} \in Z_{h, \tau}$ minimizing the $L^{2}$ distance from $\xi$. This can be obtained solving the following problem for $\left(\sigma_{h, \tau}, \phi_{h, \tau}\right) \in Z_{h, \tau} \times F_{h, \tau}$

$$
\begin{cases}\left\langle\sigma_{h, \tau}, v\right\rangle+\left\langle\phi_{h, \tau}, \operatorname{div} v\right\rangle=\left\langle\xi, v_{h}\right\rangle, & \forall v \in Z_{h, \tau}  \tag{4.13}\\ \left\langle\operatorname{div} \sigma_{h, \tau}, \psi\right\rangle=0, & \forall \psi \in F_{h, \tau}\end{cases}
$$

The commuting diagram (4.11) implies well-posedness of the discrete system. In particular, it implies the following inf-sup condition: there exists a constant $\beta>0$ independent of $h$ and $\tau$ such that

$$
\begin{equation*}
\inf _{\phi \in F_{h, \tau}} \sup _{\sigma \in Z_{h, \tau}} \frac{\langle\phi, \operatorname{div} \sigma\rangle}{\|\sigma\|_{H(\text { div })}\|\phi\|_{L^{2}}}>\beta \tag{4.14}
\end{equation*}
$$

see for example proposition 5.4.2 in [10]. Then, problem (4.13) is well-posed, i.e. it has a unique solution ( $\sigma_{h, \tau}, \phi_{h, \tau}$ ) which verifies $\sigma_{h, \tau} \in \mathcal{B}$ and

$$
\begin{gather*}
\left\|\sigma_{h, \tau}\right\|_{L^{2}} \leq C_{1}\|\xi\|_{L^{2}}  \tag{4.15}\\
\left\|\phi_{h, \tau}\right\|_{L^{2}} \leq C_{2}\|\xi\|_{L^{2}}  \tag{4.16}\\
\left\|\sigma_{h, \tau}-\sigma\right\|_{L^{2}}+\left\|\phi_{h, \tau}-\phi\right\|_{L^{2}} \leq C_{3}\left\|\xi_{h, \tau}-\xi\right\|_{L^{2}} \tag{4.17}
\end{gather*}
$$

where $C_{1}, C_{2}, C_{3}>0$ are constants independent of $h$ and $\tau, \xi_{h}$ is the $L^{2}$ projection of $\xi$ onto $Z_{h, \tau}$ and ( $\sigma, \phi$ ) is the unique solution of problem (4.12) (e.g., these results can be derived as particular cases of theorems 4.3.2, 5.2.1 and 5.2.5 in [10]).

In the following we will need to compute the discrete version of the projection on $\mathcal{C}$. In particular we define

$$
\begin{equation*}
\mathcal{C}_{h, \tau}:=\left\{\sigma \in \mathcal{B}_{h, \tau}, \sigma \cdot n_{\partial \Omega}=\mathcal{X}_{h, \tau}\right\} \tag{4.18}
\end{equation*}
$$

where, since $\Pi_{Q_{h}}$ can be defined on $\mathcal{M}(D)$ (see equation (4.5)), we set

$$
\mathcal{X}_{h, \tau}:= \begin{cases}\Pi_{Q_{h}} \rho_{0} & \text { on }\{0\} \times D  \tag{4.19}\\ \Pi_{Q_{h}} \rho_{1} & \text { on }\{1\} \times D \\ 0 & \text { otherwise }\end{cases}
$$

The well-posedness arguments above for the projection onto $\mathcal{B}$ and $\mathcal{B}_{h, \tau}$ hold also for the projections onto $\mathcal{C}$ and $\mathcal{C}_{h, \tau}$ up to adding boundary conditions to the spaces $H(\operatorname{div} ; \Omega)$ and $Z_{h, \tau}$, and replacing $L^{2}(\Omega)$ and $F_{h, \tau}$ by $L^{2}(\Omega) / \mathbb{R}$ and $F_{h, \tau} / \mathbb{R}$, respectively.

## 5. Discrete dynamical formulation and convergence

In this section we formulate the discrete problem and state a convergence result obtained by applying the theory developed in [19]. For this, we need to introduce a space for the discrete dual variable $q$. We adopt the same notation as for the spaces defined in section 4 . In particular, we set for $r \in\{0,1\}$,

$$
\begin{equation*}
X_{h}^{r}:=\left\{\phi \in L^{2}(D) ;\left.\phi\right|_{T} \in X_{h}^{r}(T), \forall T \in \mathcal{T}_{h}\right\} . \tag{5.1}
\end{equation*}
$$

The superscript $r$ denotes the polynomial order of the shape function space $X_{h}^{r}(T)$. We distinguish two cases:
(1) for simplicial elements (triangles or tetrahedrons), we take $X_{h}^{r}(T):=\mathcal{P}_{r}(T)$.
(2) for quadrilateral elements, we set $T=\phi\left(I^{d}\right)$, where $I$ is an interval and $\phi$ an affine transformation, and we take $X_{h}^{r}(T):=\mathcal{P}_{r}(I)^{d}$.
The associated space-time space is defined by

$$
\begin{equation*}
X_{h, \tau}^{r}:=\left\{\phi \in L^{2}(\Omega) ;\left.\phi\right|_{S} \in X_{h, \tau}^{r}(S), \forall S \in \mathcal{T}_{h, \tau}\right\} \tag{5.2}
\end{equation*}
$$

with $X_{h, \tau}^{r}(S)=\mathcal{P}_{0}\left(\left[t_{0}, t_{1}\right]\right) \otimes X_{h}^{r}(T)$. In order to simplify the notation, we will omit the superscript $r$ when not relevant to the discussion.
Remark 5.1. The choice $r \in\{0,1\}$ is dictated by computational feasibility of the algorithm. In fact, for these spaces, we can compute explicitly the projection on $\mathcal{K} \cap X_{h, \tau}$ as it will be explained in the next section. On the other hand, we restrict ourselves to piecewise constant functions in time since this is crucial for the convergence of the algorithm, as shown in [19].

The discrete action (at fixed time) is defined as follows

$$
\begin{equation*}
A_{h}(\rho, m)=\sup _{(a, b) \in\left(X_{h}\right)^{d+1}}\{\langle\rho, a\rangle+\langle m, b\rangle ;(a, b) \in K \text { a.e. }\} \tag{5.3}
\end{equation*}
$$

for any $(\rho, m) \in Q_{h} \times V_{h}$. By construction, $A_{h}: Q_{h} \times V_{h} \rightarrow[0,+\infty]$ is a proper convex function -1 -positively homogeneous in its first variable and 2-positively homogeneous in its second variable. Moreover, it is non-increasing in its first argument, i.e. $A_{h}\left(\rho_{1}+\rho_{2}, m\right) \leq A_{h}\left(\rho_{1}, m\right)$ for any $\rho_{1}, \rho_{2} \in Q_{h}^{+}$and $m \in V_{h}$. In fact, suppose that $A_{h}\left(\rho_{1}+\rho_{2}, m\right)<+\infty$. Then there exists $\left(a^{*}, b^{*}\right) \in K$ such that $\left\langle\rho_{1}+\rho_{2}, a^{*}\right\rangle+\left\langle m, b^{*}\right\rangle=A_{h}\left(\rho_{1}+\rho_{2}, m\right)$; in particular $a^{*} \leq 0$. Then

$$
\begin{equation*}
A_{h}\left(\rho_{1}, m\right) \geq A_{h}\left(\rho_{1}+\rho_{2}, m\right)-\left\langle\rho_{2}, a^{*}\right\rangle \geq A_{h}\left(\rho_{1}+\rho_{2}, m\right) \tag{5.4}
\end{equation*}
$$

and by a similar reasoning we obtain that if $A\left(\rho_{1}+\rho_{2}, m\right)=+\infty$ then we also have $A\left(\rho_{1}, m\right)=+\infty$.

The space-time discretization of problem (3.1) is given by

$$
\begin{equation*}
\inf _{\sigma \in \mathcal{C}_{h, \tau}} \mathcal{A}_{h, \tau}(\sigma), \quad \mathcal{A}_{h, \tau}(\sigma):=\sup _{q \in\left(X_{h, \tau}\right)^{d+1} \cap \mathcal{K}}\langle q, \sigma\rangle . \tag{5.5}
\end{equation*}
$$

Suppose that the time discretization is given by a decomposition of the interval $[0,1]$ in $N$ elements, i.e. fixing the points $0=t_{0}<t_{1}<\ldots<t_{N+1}=1$. Given $\sigma=(\rho, m) \in V_{h}$, we can identify the density $\rho$ with the collection $\left\{\rho_{i}\right\}_{i=1}^{N+1}$ with $\rho_{i} \in Q_{h}$, and the momentum $m$ with the collection $\left\{m_{i}\right\}_{i=1}^{N}$ with $m_{i} \in V_{h}$. Since $q$ is piecewise constant in time, we have the following equivalent formulation

$$
\begin{equation*}
\mathcal{A}_{h, \tau}(\sigma)=\sum_{i=1}^{N} A_{h}\left(\frac{\rho_{i}+\rho_{i+1}}{2}, m_{i}\right)\left|t_{i+1}-t_{i}\right| \tag{5.6}
\end{equation*}
$$

Note that in order to obtain (5.6) from (5.5), we relied on the particular choice of finite element spaces for density (piecewise linear in time), momentum (piecewise constant in time) and the corresponding dual variables (piecewise constant in time).
Remark 5.2 (Positivity constraint). A disadvantage of having the action evaluated on the mean density (in time) is that the positivity constraint $\rho \geq 0$ is then only enforced on $\frac{\rho_{i}+\rho_{i+1}}{2}$, rather than on each $\rho_{i}$ separately. For certain choices of spaces this may lead to oscillations as shown in the numerical tests in section 8.

The objects introduced until now define a finite dimensional model of optimal transport in the sense of Definition 2.5 in [19]. The framework developed therein can be used to deduce a convergence result for our scheme.

Theorem 5.3. Let $\rho_{0}, \rho_{1} \in \mathcal{P}(D)$ be given and $\left\{\mathcal{T}_{h, \tau}\right\}_{h, \tau \geq 0}$ a family of tensorproduct decomposition of $\Omega$ such that the time discretization is uniform, i.e. $t_{i+1}-$ $t_{i}=\tau$ for all $i=1, \ldots, N$, and the space discretization $\mathcal{T}_{h}$ satisfies equation (4.1). Let $\sigma_{h, \tau}$ be the solution of problem (5.5) associated with $\mathcal{T}_{h, \tau}$ and for $r=1$. Then, as $h, \tau \rightarrow 0$, up to extraction of a subsequence, $\sigma_{h, \tau}$ converges weakly to $\sigma \in \mathcal{M}(\Omega)^{d+1}$ a minimizer of problem (3.1).

The proof is essentially an adaptation of the one presented in [19] and is postponed to the appendix.
Remark 5.4. Suppose that $D=[0,1]^{d}$ and that $\mathcal{T}_{h, \tau}$ is a uniform quadrilateral discretization of $\Omega=[0,1]^{d+1}$. Then for $r=0$, the discrete problem (5.5) coincides with the discretization proposed in [22]. Theorem 5.3 shows that modifying this method with $r=1$, one can prove convergence to the solution of the continuous problem (3.1).

## 6. The proximal splitting algorithm

We now describe in detail the discrete version of the proximal splitting algorithm introduced in section 3.1. Note that this is a proximal splitting algorithm applied to the discrete optimization problem rather than the discretization of the optimality conditions of the algorithm in the continuous setting.

First of all, we define

$$
\begin{equation*}
\mathcal{K}_{h, \tau}^{r}:=\mathcal{K} \cap\left(X_{h, \tau}^{r}\right)^{d+1}:=\left\{q \in\left(X_{h, \tau}^{r}\right)^{d+1} ; q \in \text { Ka.e. }\right\} \tag{6.1}
\end{equation*}
$$

We write the discrete problem as follows:

$$
\begin{equation*}
\inf _{\sigma \in L^{2}\left(\Omega ; \mathbb{R}^{d+1}\right)} \sup _{q \in L^{2}\left(\Omega ; \mathbb{R}^{d+1}\right)}\langle q, \sigma\rangle+\iota_{\mathcal{C}_{h, \tau}}(\sigma)-\iota_{\mathcal{K}_{h, \tau}}(q) \tag{6.2}
\end{equation*}
$$

where $\mathcal{C}_{h, \tau}$ is defined in (4.18). Then, applying the proximal splitting algorithm to problem (6.2) gives the following algorithm: given $\tau_{1}, \tau_{2}>0$ and an admissible $\left(\sigma^{0}, q^{0}\right) \in \mathcal{C}_{h, \tau} \times \mathcal{K}_{h, \tau}$, we define the sequence $\left\{\left(\sigma^{k}, q^{k}\right)\right\}_{k}$ by the two-step algorithm:

$$
\begin{array}{ll}
\text { Step 1 : } & \sigma^{k+1}=P_{\mathcal{C}_{h, \tau}}\left(\sigma^{k}-\tau_{1} q^{k}\right) \\
\text { Step 2 : } & q^{k+1}=P_{\mathcal{K}_{h, \tau}}\left(q^{k}+\tau_{2}\left(2 \sigma^{k+1}-\sigma^{k}\right)\right) \tag{6.3b}
\end{array}
$$

The convergence result in theorem 3.2 clearly holds also in the discrete setting and gives convergence of the algorithm to a discrete saddle point ( $\sigma_{h, \tau}, F_{h, \tau}$ ), assuming that it exists, with the condition $\tau_{1} \tau_{2}<1$. The two steps in the algorithm can be computed as follows.

Step 1. As discussed in section 4.3 the projection $P_{\mathcal{C}_{h, \tau}}$ can be computed modifying the system given by (4.13) by adding the Neumann boundary conditions associated with the function (4.19).

Step 2. Since $P_{\mathcal{K}_{h, \tau}}$ is an $L^{2}$ projection, it is easy to check that $P_{\mathcal{K}_{h, \tau}}=P_{\mathcal{K}_{h, \tau}} \circ$ $P_{\left(X_{h, \tau}^{r}\right)^{d+1}}$, where $P_{\left(X_{h, \tau}^{r}\right)^{d+1}}$ denotes the $L^{2}$ projection onto $\left(X_{h, \tau}^{r}\right)^{d+1}$. This means that we only need to be able to compute $P_{\mathcal{K}_{h, \tau}}$ when applied to an element of $X_{h, \tau}^{r}$. In addition, since $X_{h, \tau}^{r}$ is discontinuous across elements, we can compute the projection element by element, and since functions in $X_{h, \tau}^{r}(S)$ are constant along the time direction, we can also eliminate the time variable in the projection. More precisely, let $\xi \in\left(X_{h, \tau}^{r}(S)\right)^{d+1}$ and $S=\left[t_{0}, t_{1}\right] \times T$, we can identify $\xi$ with an element of $\left(X_{h}^{r}(T)\right)^{d+1}$. Then,

$$
\begin{equation*}
\xi_{\mathcal{K}}:=\left.P_{\mathcal{K}_{h, \tau}}\right|_{S} \xi=\operatorname{argmin}\left\{\|\xi-q\|_{T}^{2} ; q \in\left(X_{h}^{r}(T)\right)^{d+1} \cap K\right\} . \tag{6.4}
\end{equation*}
$$

We distinguish two cases:
(1) if $r=0$, the projection (6.4) is just the projection of a vector $\xi \in \mathbb{R}^{d+1}$ onto the convex set $K$;
(2) if $r=1, \xi_{\mathcal{K}}$ is fully determined by its value on the vertices $\left\{v_{i}\right\}_{i}$ of $T$, and by linearity of the shape functions and convexity of the constraint, we obtain

$$
\begin{equation*}
\xi_{\mathcal{K}}\left(v_{i}\right)=\operatorname{argmin}\left\{\left\|\xi\left(v_{i}\right)-q\right\|_{T}^{2} ; q \in K\right\} . \tag{6.5}
\end{equation*}
$$

In both cases we only need to compute the projection of vectors in $\mathbb{R}^{d+1}$ onto $K$, which can be done explicitly by finding the largest root of a third order polynomial, as described in [22] for example.

Remark 6.1. As for the finite difference discretization studied in [22], different optimization techniques could be applied to solve problem (5.5). In particular, it should be noted that the ADMM approach orginally proposed by Benamou and Brenier [5] could also be applied. This would lead to a very similar algorithm to (6.3), but it would require the introduction of an additional variable which avoids coupling of the degrees of freedom in the optimization step with respect to $q$. In other words, this is needed in order to be able to perform the projection on $K$ for each degree of freedom separately. More details on this issue can be found in [22] for the discretization studied therein, and they hold also in the finite element setting.

## 7. Regularization

The optimal transport problem does not involve any regularizing effect on the interpolation between two measures. In fact, one can even expect a loss of regularity in some cases, namely if one is interpolating between two smooth densities on a smooth but non-convex domain. Such a a loss of regularity (which is often unphysical when the density represents a physical quantity) can be avoided introducing additional regularization terms in the formulation. In this section we describe how to do so, and how these modifications translate at the algorithmic level.

We consider the Hilbert space setting discribed in section 3.1 and we study problems in the form

$$
\begin{equation*}
\inf _{\sigma \in \mathcal{C}} \mathcal{A}(\sigma)+\alpha \mathcal{R}(\sigma) \tag{7.1}
\end{equation*}
$$

where $\mathcal{R}: L^{2}(\Omega) \rightarrow \mathbb{R}$ is a convex, proper and l.s.c. functional, and $\alpha>0$. For this type of problem, we can still apply the proximal splitting algorithm (3.9) replacing
the projection onto $\mathcal{C}$ by $\operatorname{prox}_{\tau_{1} \mathcal{F}}$ the proximal operator of $\mathcal{F}:=\iota_{\mathcal{C}}+\alpha \mathcal{R}$, defined by

$$
\begin{equation*}
\operatorname{prox}_{\tau_{1} \mathcal{F}}(\xi)=\underset{\eta \in L^{2}\left(\Omega ; \mathbb{R}^{d+1}\right)}{\operatorname{argmin}} \frac{\|\xi-\eta\|^{2}}{2 \tau_{1}}+\mathcal{F}(\eta) . \tag{7.2}
\end{equation*}
$$

This leads to the so-called PDGH algorithm, which for $\tau_{1} \tau_{2}<1$ can be seen just as a proximal point method applied to a monotone operator [11], and therefore we still have convergence in the Hilbert space setting. As mentioned in [19] convergence of the discrete problem with mesh refinement is more delicate and will not be discussed here.
7.1. Mixed $L^{2}$-Wasserstein distance. Define for any $\sigma=(\rho, m) \in L^{2}(\Omega) \times$ $L^{2}\left(\Omega ; \mathbb{R}^{d}\right)$

$$
\mathcal{R}(\sigma):= \begin{cases}\frac{1}{2}\left\|\partial_{t} \rho\right\|_{L^{2}(\Omega)}^{2} & \text { if } \partial_{t} \rho \in L^{2}(\Omega)  \tag{7.3}\\ +\infty & \text { otherwise }\end{cases}
$$

With this functional, problem (7.1) yields an interpolation between the Wasserstein distance and the $L^{2}$ distance. It was originally considered in [6], where a conjugate gradient method was proposed to compute the minimizers. Let $V:=$ $H^{1}\left([0,1] ; L^{2}(D)\right) \times L^{2}([0,1] ; H(\operatorname{div} ; D))$ and let $\stackrel{\circ}{V}$ be the same space with homogenous boundary conditions on the fluxes. For any $\xi \in L^{2}(\Omega)^{d+1}, \sigma=\operatorname{prox}_{\tau_{1} \mathcal{F}}(\xi)$ is obtained by solving the following system for $(\sigma, \phi) \in V \times L^{2}(\Omega) / \mathbb{R}$

$$
\begin{cases}\langle\sigma, v\rangle+\alpha \tau_{1}\left\langle\partial_{t} \rho, \partial_{t} v_{t}\right\rangle+\langle\phi, \operatorname{div} v\rangle=\langle\xi, v\rangle, & \forall v \in \stackrel{\circ}{V}  \tag{7.4}\\ \langle\operatorname{div} \sigma, \psi\rangle=0, & \forall \psi \in L^{2}(\Omega) / \mathbb{R} \\ \sigma \cdot n_{\partial \Omega}=\mathcal{X}, & \end{cases}
$$

where $v_{t}=v \cdot \hat{t}$ is the component of $v$ in the time direction. Well-posedness can be obtained by standard methods for saddle point problems [10] and it translates directly into well-posedness of the discrete system obtained by replacing $V$ with $Z_{h, \tau}, L^{2}(\Omega)$ with $F_{h, \tau}$, and $\mathcal{X}$ with $\mathcal{X}_{h, \tau}$.
7.2. $H^{1}$ regularization. Define for any $\sigma=(\rho, m) \in L^{2}(\Omega) \times L^{2}\left(\Omega ; \mathbb{R}^{d}\right)$

$$
\mathcal{R}(\sigma):= \begin{cases}\frac{1}{2}\left\|\nabla_{x} \rho\right\|_{L^{2}(\Omega)}^{2} & \text { if } \rho \in L^{2}\left([0,1] ; H^{1}(D)\right)  \tag{7.5}\\ +\infty & \text { otherwise }\end{cases}
$$

In this case we set $V:=H(\operatorname{div} ; \Omega), W:=L^{2}\left([0,1] ; H\left(\operatorname{div}_{x} ; D\right)\right)$ and let $\stackrel{\circ}{V}$ and $\stackrel{\circ}{W}$ be the same spaces with homogenous boundary conditions on the fluxes. Then, for any $\xi \in L^{2}(\Omega)^{d+1}, \sigma=\operatorname{prox}_{\tau_{1} \mathcal{F}}(\xi)$ is obtained by solving the following system for $(\sigma, \eta, \phi) \in V \times \stackrel{\circ}{W} \times L^{2}(\Omega) / \mathbb{R}$

$$
\begin{cases}\langle\sigma, v\rangle-\alpha \tau_{1}\left\langle\operatorname{div}_{x} \eta, v_{t}\right\rangle+\langle\phi, \operatorname{div} v\rangle=\langle\xi, v\rangle, & \forall v \in \stackrel{\circ}{V}  \tag{7.6}\\ \left\langle\rho, \operatorname{div}_{x} w\right\rangle+\langle\eta, w\rangle=0, & \forall w \in \stackrel{\circ}{W} \\ \langle\operatorname{div} \sigma, \psi\rangle=0, & \forall \psi \in L^{2}(\Omega) / \mathbb{R} \\ \sigma \cdot n_{\partial \Omega}=\mathcal{X}, & \end{cases}
$$

where $v_{t}=v \cdot \hat{t}$ is the component of $v$ in the time direction. As before, well-posedness can be obtained by standard methods for saddle point problems [10].

We introduce the space $W_{h, \tau} \subset L^{2}\left([0,1] ; H\left(\operatorname{div}_{x} ; D\right)\right)$ whose shape functions on $S=\left[t_{0}, t_{1}\right] \otimes T$ are given by

$$
\begin{equation*}
W_{h, \tau}(S):=\mathcal{P}_{1}\left(\left[t_{0}, t_{1}\right]\right) \otimes V_{h}(T) \tag{7.7}
\end{equation*}
$$

We denote by ${ }^{\circ}{ }_{h, \tau}$ the same space with the boundary conditions $\eta \cdot n_{\partial \Omega}=0$ on $[0,1] \times \partial D$. Denote by $\nabla_{x}^{h}: L^{2}(\Omega) \rightarrow{ }^{\circ}{ }_{h, \tau}$ the adjoint of $-\operatorname{div}_{x}$ defined by

$$
\begin{equation*}
\left\langle\nabla_{x}^{h} \phi, \eta\right\rangle=-\left\langle\phi, \operatorname{div}_{x} \eta\right\rangle, \quad \forall(\phi, \eta) \in L^{2}(\Omega) \times \stackrel{\circ}{W}_{h, \tau} \tag{7.8}
\end{equation*}
$$

We define a discrete version of (7.5) as follows:

$$
\begin{equation*}
\mathcal{R}_{h, \tau}(\sigma):=\frac{1}{2}\left\|\nabla_{x}^{h} \rho\right\|_{L^{2}(\Omega)}^{2} \tag{7.9}
\end{equation*}
$$

Let $\mathcal{F}_{h, \tau}:=\iota_{\mathcal{C}_{h, \tau}}+\alpha \mathcal{R}_{h, \tau}$. Then for any $\xi \in L^{2}(\Omega)^{d+1}, \sigma=\operatorname{prox}_{\tau_{1} \mathcal{F}_{h, \tau}}(\xi)$ is obtained by solving the following system for $(\sigma, \eta, \phi) \in \stackrel{\circ}{V}_{h, \tau} \times \stackrel{\circ}{W}_{h, \tau} \times F_{h, \tau} / \mathbb{R}$ :

$$
\begin{cases}\langle\sigma, v\rangle-\alpha \tau_{1}\left\langle\operatorname{div}_{x} \eta, v_{t}\right\rangle+\langle\phi, \operatorname{div} v\rangle=\langle\xi, v\rangle, & \forall v \in \stackrel{\circ}{V}_{h, \tau},  \tag{7.10}\\ \left\langle\rho, \operatorname{div}_{x} w\right\rangle+\langle\eta, w\rangle=0, & \forall w \in \stackrel{\rightharpoonup}{W}_{h, \tau}, \\ \langle\operatorname{div} \sigma, \psi\rangle=0, & \forall \psi \in F_{h, \tau} / \mathbb{R} \\ \sigma \cdot n_{\partial \Omega}=\mathcal{X}_{h, \tau}\end{cases}
$$

## 8. Numerical Results

In this section we describe two numerical tests that demonstrate the behaviour of the proposed discretization both qualitatively and in terms of convergence of the algorithm. For both tests the time discretization is fixed with a time step $\tau:=\left|t_{i+1}-t_{i}\right|=1 / 30$, but we will use different meshes and finite element spaces for the discretization in space. The results shown hereafter have been obtained using the finite element software Firedrake [24] (see [21, 9], for the tensor product constructions) and the linear solver for the Poisson equation is based on PETSc [3, 4].
8.1. Gaussian interpolation. We set $D=[0,1]^{2}$ and

$$
\begin{equation*}
\rho_{0}=\exp \left(-\frac{\left|x-x_{0}\right|^{2}}{2 s^{2}}\right), \quad \rho_{1}=\exp \left(-\frac{\left|x-x_{1}\right|^{2}}{2 s^{2}}\right) \tag{8.1}
\end{equation*}
$$

where $s=0.1, x_{0}=(0.35,0.35)$ and $x_{1}=(0.65,65)$. The optimal transport interpolation between the density in equation (8.1) when the domain is $\mathbb{R}^{2}$ is simply given by

$$
\begin{equation*}
\rho_{t}=\exp \left(-\frac{\left|x-(1-t) x_{0}-t x_{1}\right|^{2}}{2 s^{2}}\right) \tag{8.2}
\end{equation*}
$$

This is still a good approximation on our bounded domain and it allows us to assess qualitatively the different discretizations we have proposed. In figure 2, it is shown the interpolation at time $t=0.5$ for different choices of spaces $V_{h}$ and $X_{h}$. In general, for the same space $V_{h}$, the discretizations corresponding to $X_{h}^{1}$ are slightly more diffusive than those corresponding to $X_{h}^{0}$, but they also lead to smaller oscillations on the positivity. On the other hand for the couple of spaces $V_{h}=\mathcal{B D} \mathcal{M}_{1}, X_{h}^{0}$, we observe some oscillations appearing in the interpolation (see also remark 5.2). In figure (3), the different schemes are compared in terms of convergence of the proximal splitting algorithm. The cases corresponding to $X_{h}^{1}$ appear to converge faster than $X_{h}^{0}$ (when used jointly with the spaces $\mathcal{B D} \mathcal{M}_{1}$ and $\left.\mathcal{R} \mathcal{T}_{[0]}\right)$.
8.2. Non-convex domain. Here we consider a non-convex domain $D$ and boundary conditions still given by (8.1), but with $s=0.1, x_{0}=(0.5,0.1)$ and $x_{1}=$ ( $0.5,0.9$ ). The mesh $\mathcal{T}_{h}$ and the boundary conditions are illustrated in figure 4. Note that we have refined the mesh along the diagonal where we expect the mass to concentrate. In figure 5,6 and 7 we show the density evolution up to time $t=0.5$ (the other half of the time evolution being symmetric in space given the
boundary conditions and the domain shape) for the non-regularized case, the $H^{1}$ regularization and the $L^{2}$ regularization, respectively. For both regularizations the density profile appears to be smoothened, but only the $H^{1}$ regularization avoids concetration at the the corners. The proximal operator of the projection on the continuity equation is more expensive computationally for the $H^{1}$ regularization than for the other two cases, since we have to solve a larger mixed system at each iteration. However, for both regularizations, the proximal splitting algorithm itself converges much faster than the non-regularized case, as it can be seen in figure 8 .


Figure 2. Comparison between Optimal Transport interpolation of two Gaussians for different spaces.

## Appendix A. Proof of theorem 5.3

In Definition 2.9 of [19], we find a number of requirements that are sufficient to prove convergence of a finite dimensional optimal transport model to its continuous counterpart in the sense of theorem 5.3. Here, we prove that such conditions hold for the finite element discretization proposed in this work. In the following, we assume $r=1$, and $X_{h}$ stands for $X_{h}^{1}$. We also denote by $\mathcal{I}$ is the standard nodal interpolant onto $X_{h}$, defined element by element.

First of all, we introduce some notation and list some technical results [12]. Denote by $P_{X_{h}}$ and $P_{V_{h}}$ the $L^{2}$ projections onto $X_{h}$ and $V_{h}$, respectively. Then,

$$
\begin{equation*}
\left\|P_{X_{h}} \varphi\right\|_{L^{p}} \leq C\|\varphi\|_{L^{p}}, \quad \forall \varphi \in L^{p}, 1 \leq p \leq \infty \tag{A.1}
\end{equation*}
$$

and moreover $\forall T \in \mathcal{T}_{h}$

$$
\begin{equation*}
\left\|\varphi-P_{X_{h}} \varphi\right\|_{L^{p}(T)} \leq C h_{T}\|\nabla \varphi\|_{L^{p}(T)}, \quad \forall \varphi \in W^{1, p}(T), 1 \leq p \leq \infty \tag{A.2}
\end{equation*}
$$

where, with an abuse of notation, we have used $P_{X_{h}}$ to denote the $L^{2}$ projection onto $X_{h}(T)$. These imply the following lemma.

Lemma A.1. Given the regularity assumption in (4.1) on $\mathcal{T}_{h}$, we have

$$
\begin{equation*}
\left\|\mathcal{I}\left|P_{V_{h}} b\right|^{2}\right\|_{L^{\infty}} \leq C\|b\|_{L^{\infty}}^{2} \tag{A.3}
\end{equation*}
$$

for any $b \in L^{\infty}(D)$, and

$$
\begin{equation*}
\left\|\mathcal{I}\left|P_{V_{h}} b\right|^{2}-|b|^{2}\right\|_{L^{\infty}} \leq C h|b|_{W^{1, \infty}}\|b\|_{L^{\infty}} \tag{A.4}
\end{equation*}
$$

for any $b \in W^{1, \infty}(D)$.
Proof. For the first inequality, using standard inverse inequalities, we have

$$
\begin{aligned}
\left\|\mathcal{I}\left|P_{V_{h}} b\right|^{2}\right\|_{L^{\infty}} & \leq\left\|\left|P_{V_{h}} b\right|^{2}\right\|_{L^{\infty}} \\
& \leq C h^{-d}\left\|\left|P_{V_{h}} b\right|^{2}\right\|_{L^{1}} \\
& =C h^{-d}\left\|P_{V_{h}} b\right\|_{L^{2}}^{2} \\
& \leq C h^{-d}\left\|P_{\left(X_{h}\right)^{d}} b\right\|_{L^{2}}^{2} \\
& \leq C\left\|P_{\left(X_{h}\right)^{d}} b\right\|_{L^{\infty}}^{2} \\
& \leq C\|b\|_{L^{\infty}}^{2}
\end{aligned}
$$

For the second inequality, we observe that

$$
\begin{equation*}
\left\|\mathcal{I}\left|P_{V_{h}} b\right|^{2}-|b|^{2}\right\|_{L^{\infty}} \leq\left\|\mathcal{I}\left|P_{V_{h}} b\right|^{2}-\mathcal{I}|b|^{2}\right\|_{L^{\infty}}+\left\|\mathcal{I}|b|^{2}-|b|^{2}\right\|_{L^{\infty}} \tag{A.5}
\end{equation*}
$$

The second term of the right-hand side is easy to control. For the first term, we have

Again, the second term is easy to control. For the first tem, using the same reasoning as above,

$$
\begin{aligned}
\left\|\left|P_{V_{h}} b\right|^{2}-\left|P_{\left(X_{h}\right)^{d}} b\right|^{2}\right\|_{L^{\infty}} & \leq C h^{-d}\left\|\left|P_{V_{h}} b\right|^{2}-\left|P_{\left(X_{h}\right)^{d}} b\right|^{2}\right\|_{L^{1}} \\
& \leq C h^{-d} \sum_{i=1}^{d}\left\|\left(P_{V_{h}} b\right)_{i}^{2}-\left(P_{X_{h}} b_{i}\right)^{2}\right\|_{L^{1}} \\
& \leq C h^{-d} \sum_{i=1}^{d}\left\|\left(P_{V_{h}} b\right)_{i}-P_{X_{h}} b_{i}\right\|_{L^{1}}\|b\|_{L^{\infty}} \\
& \leq C h^{-\frac{d}{2}}\left\|P_{V_{h}} b-P_{X_{h}} b\right\|_{L^{2}}\|b\|_{L^{\infty}} \\
& \leq C h\left\|\nabla P_{X_{h}} b\right\|_{L^{\infty}}\|b\|_{L^{\infty}} \leq C h\|\nabla b\|_{L^{\infty}}\|b\|_{L^{\infty}}
\end{aligned}
$$

We will need also some approximation properties for the projection operators $\Pi_{Q_{h}}$ and $\Pi_{V_{h}}$. Specifically, we will assume that these verify for any $\phi \in H^{1}(D)$ and $\eta \in H^{1}(D)^{d}$

$$
\begin{equation*}
\left\|\Pi_{Q_{h}} \phi-\phi\right\|_{L^{2}(D)} \leq C h\|\phi\|_{H^{1}(D)}, \quad\left\|\Pi_{V_{h}} \eta-\eta\right\|_{L^{2}(D)^{d}} \leq C h\|\eta\|_{H^{1}(D)^{d}} \tag{A.6}
\end{equation*}
$$

Given the mesh regularity assumption (4.1), this is a standard property for $\Pi_{Q_{h}}$ as defined in equation (4.5), whereas for $\Pi_{V_{h}}$ (since this has not been defined explicitly) we refer to the construction discussed in [2] based on a Clément interpolant, for example.

Proposition A. 2 below contains the properties needed for convergence: it can be seen as a specific instance of Definition 2.9 of [19]. Note that a few of the properties listed therein are omitted here because they are either unnecessary or true by construction in our setting. Note also that the sampling operators used in [19] are replaced here with the bounded projections $\Pi_{Q_{h}}$ and $\Pi_{V_{h}}$, where $\Pi_{Q_{h}}$ can be naturally extended to $\mathcal{M}(D)$ (see equation (4.5)) and $\Pi_{V_{h}}$ is considered to be defined on a domain different than $\mathcal{M}(D)^{d}$. Moreover the reconstruction operators are simply the injection operators from $Q_{h}$ and $V_{h}$ to $\mathcal{M}(D)$ and $\mathcal{M}(D)^{d}$, respectively. Finally, we define for any $(\rho, b) \in \mathcal{M}(D) \times C\left(D ; \mathbb{R}^{d}\right)$

$$
\begin{equation*}
A^{*}(\rho, b):=\int_{D} \frac{|b|^{2}}{2} \rho \tag{A.7}
\end{equation*}
$$

so that if $(\rho, m) \in \mathcal{M}_{+}(D) \times \mathcal{M}(D)^{d}$ then

$$
\begin{equation*}
A(\rho, m)=\sup _{b \in C\left(D ; \mathbb{R}^{d}\right)}\langle m, b\rangle-A^{*}(\rho, b) ; \tag{A.8}
\end{equation*}
$$

and for any $(\rho, b) \in Q_{h} \times V_{h}$,

$$
\begin{equation*}
A_{h}^{*}(\rho, b):=\sup _{b \in V_{h}}\langle m, b\rangle-A_{h}(\rho, m) \tag{A.9}
\end{equation*}
$$

Proposition A.2. The following properties hold:
(1) For any $\rho \in \mathcal{M}_{+}(D), \Pi_{Q_{h}} \rho \rightarrow \rho$ as $h \rightarrow 0$ weakly in $\mathcal{M}(D)$.
(2) Let $B \subset\left(C^{1}(D)\right)^{d}$ a bounded subset. Then there exists a constant $\varepsilon_{h}$ tending to 0 as $h \rightarrow 0$ such that for any $b \in B$ and $\rho \in Q_{h}$

$$
\begin{equation*}
A_{h}^{*}\left(\rho, P_{V_{h}} b\right) \leq A^{*}(\rho, b)+\epsilon_{h}\|\rho\| \tag{A.10}
\end{equation*}
$$

where $P_{V_{h}}$ denotes the $L^{2}$ projection onto $V_{h}$. Moreover there exists a constant $C \geq 1$ such that for any $b \in C(D)^{d}$, there holds

$$
\begin{equation*}
A_{h}^{*}\left(\rho, P_{V_{h}} b\right) \leq \frac{C}{2}\|\rho\|\|b\|_{L^{\infty}}^{2} \tag{A.11}
\end{equation*}
$$

(3) Let $B \subset C^{1}(D)$ a bounded subset such that for all $\rho \in B$ there holds $\rho>$ $C>0$, and let $B^{\prime} \subset\left(C^{1}(D)\right)^{d}$ a bounded subset. There exists a constant $\varepsilon_{h}$ tending to 0 as $h \rightarrow 0$ such that, given $(\rho, m) \in \mathcal{M}(D)^{d+1}$ such that $\rho$ has density in $B$ and $m$ in $B^{\prime}$, then

$$
\begin{equation*}
A_{h}\left(\Pi_{Q_{h}} \rho, \Pi_{V_{h}} m\right) \leq A(\rho, m)+\varepsilon_{h} \tag{A.12}
\end{equation*}
$$

(4) There exists $\varepsilon_{h}$ tending to 0 as $h \rightarrow 0$ and a continuous function $\omega$ satisfying $\omega(0)=0$ such that: for any $x, y \in D$ and $h>0$ there exists $\rho \in Q_{h}^{+}$and $m_{1}, m_{2} \in V_{h}$ satisfying

$$
\left\{\begin{array}{l}
\operatorname{div} m_{1}=\rho-\Pi_{Q_{h}}\left(\delta_{x}\right)  \tag{A.13}\\
\operatorname{div} m_{2}=\rho-\Pi_{Q_{h}}\left(\delta_{y}\right)
\end{array} \quad \text { and } \quad A_{h}\left(\rho, m_{i}\right) \leq \omega(|x-y|)+\varepsilon_{h}, \forall i \in\{1,2\} .\right.
$$

Proof. The first point is immediate from the definition of $\Pi_{Q_{h}}$ in equation (4.5). For (2), we observe that

$$
\begin{equation*}
\left.A_{h}(\rho, m)=\sup _{b \in X_{h}}\langle m, b\rangle-\left.\frac{1}{2}\langle\rho, \mathcal{I}| b\right|^{2}\right\rangle \tag{A.14}
\end{equation*}
$$

where we recall that $\mathcal{I}$ is the standard element-wise nodal interpolant onto $X_{h}$. In fact, for any $b \in\left(X_{h}\right)^{d}$, we have $b^{2} \leq \mathcal{I}|b|^{2}$, and therefore when $\rho \geq 0$ we can "saturate" the constraint setting $a=-\mathcal{I}|b|^{2} / 2$. On the other hand if $\rho<0$ on some element both sides of the equality are $+\infty$. For $(\rho, b, m) \in Q_{h} \times V_{h} \times V_{h}$ define

$$
\begin{equation*}
\left.A_{\mathcal{I}, h}^{*}(\rho, b):=\left.\frac{1}{2}\langle\rho, \mathcal{I}| b\right|^{2}\right\rangle, \quad \bar{A}_{\mathcal{I}, h}(\rho, m):=\sup _{b \in V_{h}}\langle m, b\rangle-A_{\mathcal{I}, h}^{*}(\rho, b) . \tag{A.15}
\end{equation*}
$$

Then, since when $\rho<0$ on some element $A_{h}^{*}(\rho, b)=-\infty$,

$$
\begin{equation*}
A_{h}(\rho, m) \geq \bar{A}_{\mathcal{I}, h}(\rho, m), \quad A_{h}^{*}(\rho, b) \leq \bar{A}_{\mathcal{I}, h}^{*}(\rho, b) \leq A_{\mathcal{I}, h}^{*}(\rho, b) \tag{A.16}
\end{equation*}
$$

and we can prove (2) for $A_{\mathcal{I}, h}^{*}$. In particular, we have

$$
\begin{equation*}
A_{\mathcal{I}, h}^{*}\left(\rho, P_{V_{h}} b\right) \leq A^{*}(\rho, b)+\frac{1}{2}\left\|\mathcal{I}\left|P_{V_{h}} b\right|^{2}-|b|^{2}\right\|_{L^{\infty}}\|\rho\|, \tag{A.17}
\end{equation*}
$$

and we obtain the result applying lemma A.1. Using again lemma A.1, we easily obtain the second bound as well.

For point (3), observe first that $A_{h}\left(\Pi_{Q_{h}} \rho, \Pi_{V_{h}} m\right) \leq A\left(\Pi_{Q_{h}} \rho, \Pi_{V_{h}} m\right)$ by definition. Then given the assumption on $\rho$ and $m$ we can simply write

$$
\begin{align*}
A_{h}\left(\Pi_{Q_{h}} \rho, \Pi_{V_{h}} m\right)-A(\rho, m) & \leq \int_{D} \frac{\left|\Pi_{V_{h}} m\right|^{2}}{2 \Pi_{Q_{h}} \rho}-\frac{|m|^{2}}{2 \rho} \\
& \leq \frac{1}{2} \int_{D}\left|\frac{\left|\Pi_{V_{h}} m\right|^{2}-|m|^{2}}{\Pi_{Q_{h}} \rho}\right|+\left|\frac{|m|^{2}}{\Pi_{Q_{h}} \rho}-\frac{|m|^{2}}{\rho}\right|  \tag{A.18}\\
& \leq C\left(\left\|\Pi_{Q_{h}} \rho-\rho\right\|_{L^{2}}+\left\|\left|\Pi_{V_{h}} m\right|^{2}-|m|^{2}\right\|_{L^{1}}\right)
\end{align*}
$$

where the constant $C$ depends on the uniform lower bound on $\rho$ and on the $L^{\infty}$ norm of $|m|$. We conclude using Cauchy-Schwarz inequality on the second term and then equation (A.6).

For the last point, we will establish a connection between our scheme and the one proposed by Gladbach, Kopfer and Mass [14] and then use propoperty (A.13) for this scheme which was proved in [19]. We will consider the case of a simplicial mesh, the quadrilateral case being completely analogous upon substituting $\mathcal{R} \mathcal{T}_{0}$ with $\mathcal{R} \mathcal{T}_{[0]}$ in the discussion below.

First, we introduce some notation. For each $T \in \mathcal{T}_{h}$, let $\mathcal{T}_{h, T}$ be the set of neighbouring elements $L \in \mathcal{T}_{h}$ such that $f_{T, L}:=\bar{T} \cap \bar{L} \neq \emptyset$, which we assume to be oriented. Define by $\mathcal{F}_{h}$ the set of $(d-1)$-dimensional facets in the triangulation. Let $T, L \in \mathcal{T}_{h}$ be neighbouring elements, we denote by $\varphi_{T, L} \in \mathcal{R} \mathcal{T}_{0}$ the canonical basis function associated with the oriented facet $f_{T, L}$. Then, any $m \in \mathcal{R} \mathcal{T}_{0}$ can be written as

$$
\begin{equation*}
m=\sum_{f_{T, L} \in \mathcal{F}_{h}} m_{T, L} \varphi_{T, L} \tag{A.19}
\end{equation*}
$$

where $m_{T, L}$ is the flux of $m$ on the oriented facet $f_{T, L}$. In other words we can identify functions in $(\rho, m) \in Q_{h} \times V_{h}$ with their finite volume representation $\left\{\rho_{T}, m_{T, L}\right\}_{T, L}$. Then, we can interpret the action for the finite volume scheme [14], which we denote by $A_{h}^{F V}(\rho, m)$, as a function on $Q_{h} \times V_{h}$. This is given by the following expression

$$
\begin{equation*}
A_{h}^{F V}(\rho, m):=\sum_{f_{T, L} \in \mathcal{F}_{h}} \frac{m_{T, L}^{2}}{2 \theta\left(\rho_{T}, \rho_{L}\right)}\left|f_{T, L}\right|\left|x_{T}-x_{L}\right| \tag{A.20}
\end{equation*}
$$

where $\theta: \mathbb{R}^{+} \times \mathbb{R}^{+} \rightarrow \mathbb{R}^{+}$is an appropriate function (see [14]) which we take to be the harmonic mean.

Now, in order to construct $\rho \in Q_{h}^{+}$and $m_{1}, m_{2} \in V_{h}$ satisfying (A.13), we use the same construction as in [19] for the finite volume scheme, and interpolate this to the spaces $\mathcal{R} \mathcal{T}_{0} \subseteq V_{h}$ and $Q_{h}^{+}$to obtain $\rho, m_{1}$ and $m_{2}$ satisfying

$$
\left\{\begin{array}{l}
\operatorname{div} m_{1}=\rho-\Pi_{Q_{h}}\left(\delta_{x}\right)  \tag{A.21}\\
\operatorname{div} m_{2}=\rho-\Pi_{Q_{h}}\left(\delta_{y}\right)
\end{array}\right.
$$

In particular the support of $\rho, m_{1}$ and $m_{2}$ is a chain of neighbouring elements $T_{1}, \ldots, T_{N}$. To prove the bound on the action, we observe that $A_{h}\left(\rho, m_{i}\right) \leq$ $A\left(\rho, m_{i}\right)$. Then, we only need to bound $A\left(\rho, m_{i}\right)$ by the action of the finite-volume scheme $A_{h}^{F V}\left(\rho, m_{i}\right)$, since $A_{h}^{F V}$ satisfies the desired inequality thanks to the regularity assumption (4.1) on the mesh [19].

By the regularity assumption on the triangulation, we can assume

$$
\begin{equation*}
\int_{T \cup L}\left|\varphi_{T, L}\right|^{2} \mathrm{~d} x \leq C\left|f_{T, L}\right|\left|x_{T}-x_{L}\right| \tag{A.22}
\end{equation*}
$$

uniformly. Then, by explicit calculations we obtain $A\left(\rho, m_{i}\right) \leq C A_{h}^{F V}\left(\rho, m_{i}\right)$ and we are done.

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$V_{h}=\mathcal{R} \mathcal{T}_{0}$


$$
V_{h}=\mathcal{B D} \mathcal{M}_{1}
$$



$$
V_{h}=\mathcal{R} \mathcal{T}_{[0]}
$$

Figure 3. Convergence of the proximal splitting algorithm measured by $\left\|\sigma_{n+1}-\sigma_{n}\right\|_{L^{2}(\Omega)}$ for the optimal transport problem between two Gaussians for different spaces $V_{h}$ and $X_{h}^{r}$.


Figure 4. Initial and final density (top) and mesh (bottom) for the non-convex domain test.


Figure 5. Density evolution for the non-convex domain test without regularization, $V_{h}=\mathcal{R} \mathcal{T}_{0}, X_{h}^{1}$.

$t=0$

$t=0.2$

$t=0.4$

$t=0.1$


$$
t=0.3
$$


$t=0.5$

Figure 6. Density evolution for the non-convex domain test with the $H^{1}$ regularization $\alpha=0.002, V_{h}=\mathcal{R} \mathcal{T}_{0}, X_{h}^{1}$.


Figure 7. Density evolution for the non-convex domain test with the $L^{2}$ regularization $\alpha=0.002, V_{h}=\mathcal{R} \mathcal{T}_{0}, X_{h}^{1}$.


Figure 8. Convergence of the proximal splitting algorithm measured by $\left\|\sigma_{n+1}-\sigma_{n}\right\|_{L^{2}(\Omega)}$ for non-convex domain test without regularization (a); with the $H^{1}$ regularization and $\alpha=0.002$ (b); with the $L^{2}$ regularization and $\alpha=0.002$ (c).


[^0]:    Date: March 9, 2020.

