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► **To cite this version:**

H. Kagan, A. Alexopoulos, M. Artuso, F. Bachmair, L. Bäni, et al.. Diamond detector technology, status and perspectives. Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment, 2019, 924, pp.297-300. 10.1016/j.nima.2018.06.009 . hal-02498265

**HAL Id: hal-02498265**

**<https://hal.science/hal-02498265v1>**

Submitted on 19 Mar 2020

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**DOI:**

[10.1016/j.nima.2018.06.009](https://doi.org/10.1016/j.nima.2018.06.009)

**Document Version**

Accepted author manuscript

[Link to publication record in Manchester Research Explorer](#)

**Citation for published version (APA):**

Kagan, H., Alexopoulos, A., Artuso, M., Bachmair, F., Bäni, L., Bartosik, M., ... Zavrtanik, M. (2018). Diamond detector technology, status and perspectives. *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment*. <https://doi.org/10.1016/j.nima.2018.06.009>

**Published in:**

Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment

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# Diamond detector technology, status and perspectives

H. Kagan<sup>m</sup>, A. Alexopoulos<sup>c</sup>, M. Artuso<sup>l</sup>, F. Bachmair<sup>x</sup>, L. Bäni<sup>x</sup>, M. Bartosik<sup>c</sup>, J. Beacham<sup>m</sup>,  
H. Beck<sup>w</sup>, V. Bellini<sup>b</sup>, V. Belyaev<sup>l</sup>, B. Bentele<sup>s</sup>, P. Bergonzo<sup>k</sup>, A. Bes<sup>aa</sup>, J-M. Brom<sup>g</sup>,  
M. Bruzzi<sup>d</sup>, G. Chiodini<sup>z</sup>, D. Chren<sup>r</sup>, V. Cindro<sup>j</sup>, G. Claus<sup>g</sup>, J. Collot<sup>aa</sup>, J. Cumalat<sup>s</sup>,  
A. Dabrowski<sup>c</sup>, R. D'Alessandro<sup>d</sup>, D. Dauvergne<sup>aa</sup>, W. de Boer<sup>j</sup>, S. Dick<sup>m</sup>, C. Dorfer<sup>x</sup>,  
M. Dunser<sup>c</sup>, V. Eremin<sup>f</sup>, G. Forcolin<sup>v</sup>, J. Forneris<sup>o</sup>, L. Gallin-Martel<sup>aa</sup>, M-L. Gallin-Martel<sup>aa</sup>,  
K.K. Gan<sup>m</sup>, M. Gastal<sup>c</sup>, C. Giroletti<sup>q</sup>, M. Goffe<sup>g</sup>, J. Goldstein<sup>q</sup>, A. Golubev<sup>h</sup>, A. Gorišek<sup>i</sup>,  
E. Grigoriev<sup>h</sup>, J. Grosse-Knetter<sup>w</sup>, A. Grummer<sup>u</sup>, B. Gui<sup>m</sup>, M. Guthoff<sup>c</sup>, I. Haughton<sup>v</sup>,  
B. Hiti<sup>i</sup>, D. Hits<sup>x</sup>, M. Hoferkamp<sup>u</sup>, T. Hofmann<sup>c</sup>, J. Hosslet<sup>g</sup>, J-Y. Hostachy<sup>aa</sup>, F. Hügging<sup>a</sup>,  
C. Hutton<sup>q</sup>, J. Janssen<sup>a</sup>, K. Kanxheri<sup>ab</sup>, G. Kasieczka<sup>x</sup>, R. Kass<sup>m</sup>, F. Kassel<sup>j</sup>, M. Kis<sup>c</sup>,  
G. Kramberger<sup>i</sup>, S. Kuleshov<sup>h</sup>, A. Lacoste<sup>aa</sup>, S. Lagomarsino<sup>d</sup>, A. Lo Giudice<sup>o</sup>, E. Lukosi<sup>v</sup>,  
C. Maazouzi<sup>g</sup>, I. Mandic<sup>i</sup>, C. Mathieu<sup>g</sup>, M. Menichelli<sup>ab</sup>, M. Mikuž<sup>i</sup>, A. Morozzi<sup>ab</sup>,  
J. Moss<sup>ac</sup>, R. Mountain<sup>l</sup>, S. Murphy<sup>v</sup>, M. Muškinja<sup>i</sup>, A. Oh<sup>v</sup>, P. Olivero<sup>o</sup>, D. Passeri<sup>ab</sup>,  
H. Pernegger<sup>c</sup>, R. Perrino<sup>z</sup>, F. Picollo<sup>o</sup>, M. Pomorski<sup>k</sup>, R. Potenza<sup>b</sup>, A. Quadt<sup>w</sup>, A. Re<sup>o</sup>,  
M. Reichmann<sup>x</sup>, G. Riley<sup>v</sup>, S. Roe<sup>c</sup>, D. Sanz<sup>x</sup>, M. Scaringella<sup>d</sup>, D. Schaefer<sup>c</sup>, C.J. Schmidt<sup>e</sup>,  
D.S. Smith<sup>m</sup>, S. Schnetzer<sup>n</sup>, S. Sciortino<sup>d</sup>, A. Scorzoni<sup>ab</sup>, S. Seidel<sup>u</sup>, L. Servoli<sup>ab</sup>, B. Sopko<sup>r</sup>,  
V. Sopko<sup>r</sup>, S. Spagnolo<sup>z</sup>, S. Spanier<sup>v</sup>, K. Stenson<sup>s</sup>, R. Stone<sup>n</sup>, C. Sutura<sup>b</sup>, A. Taylor<sup>u</sup>,  
B. Tannenwald<sup>m</sup>, M. Traeger<sup>e</sup>, D. Tromson<sup>k</sup>, W. Trischuk<sup>p</sup>, C. Tuve<sup>b</sup>, J. Velthuis<sup>q</sup>,  
N. Venturi<sup>c</sup>, E. Vittone<sup>o</sup>, S. Wagner<sup>s</sup>, R. Wallny<sup>x</sup>, J.C. Wang<sup>t</sup>, J. Weingarten<sup>w</sup>, C. Weiss<sup>c</sup>,  
T. Wengler<sup>c</sup>, N. Wermes<sup>a</sup>, M. Yamouni<sup>aa</sup>, M. Zavrtanik<sup>i</sup>

<sup>a</sup>Universität Bonn, Bonn, Germany

<sup>b</sup>INFN/University of Catania, Catania, Italy

<sup>c</sup>CERN, Geneva, Switzerland

<sup>d</sup>INFN/University of Florence, Florence, Italy

<sup>e</sup>GSI, Darmstadt, Germany

<sup>f</sup>Ioffe Institute, St. Petersburg, Russia

<sup>g</sup>IPHC, Strasbourg, France

<sup>h</sup>ITEP, Moscow, Russia

<sup>i</sup>Jozef Stefan Institute, Ljubljana, Slovenia

<sup>j</sup>Universität Karlsruhe, Karlsruhe, Germany

<sup>k</sup>CEA-LIST Technologies Avancees, Saclay, France

<sup>l</sup>MEPHI Institute, Moscow, Russia

<sup>m</sup>The Ohio State University, Columbus, OH, USA

<sup>n</sup>Rutgers University, Piscataway, NJ, USA

<sup>o</sup>University of Torino, Torino, Italy

<sup>p</sup>University of Toronto, Toronto, ON, Canada

<sup>q</sup>University of Bristol, Bristol, UK

<sup>r</sup>Czech Technical Univ., Prague, Czech Republic

<sup>s</sup>University of Colorado, Boulder, CO, USA

<sup>t</sup>Syracuse University, Syracuse, NY, USA

<sup>u</sup>University of New Mexico, Albuquerque, NM, USA

<sup>v</sup>University of Manchester, Manchester, UK

<sup>w</sup>Universität Goettingen, Goettingen, Germany

<sup>x</sup>ETH Zürich, Zürich, Switzerland

<sup>y</sup>University of Tennessee, Knoxville, TN, USA

<sup>z</sup>INFN-Lecce, Lecce, Italy

<sup>aa</sup>LPSC-Grenoble, Grenoble, France

<sup>ab</sup>INFN-Perugia, Perugia, Italy

<sup>ac</sup>California State University, Sacramento, CA, USA

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50 **Abstract**

51 Detectors based on Chemical Vapor Deposition (CVD) diamond have been used extensively  
52 and successfully in beam conditions/beam loss monitors as the innermost detectors in the highest  
53 radiation areas of Large Hadron Collider (LHC) experiments. The startup of the LHC in 2015  
54 brought a new milestone where the first polycrystalline CVD (pCVD) diamond pixel modules  
55 were installed in an LHC experiment and successfully began operation. The RD42 collaboration  
56 at CERN is leading the effort to develop polycrystalline CVD diamond as a material for tracking  
57 detectors operating in extreme radiation environments. The status of the RD42 project with  
58 emphasis on recent beam test results is presented.

59 *Keywords:* Chemical Vapor Deposition, pCVD diamond, diamond detectors, 3D diamond  
60 detectors, radiation tolerant detectors

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61 **1. Introduction**

62 The RD42 collaboration [1, 2] at CERN is leading the effort to develop radiation tolerant  
63 devices based on pCVD diamond as a material for tracking detectors operating in harsh radiation  
64 environments. Diamond has properties which make it suitable for such detector applications.  
65 During the last few years the RD42 group has succeeded in producing and measuring a number  
66 of devices to address specific issues related to use at the HL-LHC [3, 4]. This paper presents  
67 the status of the RD42 project with emphasis on recent beam test results. In particular, results  
68 are presented on the status of the first diamond pixel detector based on pCVD material, on the  
69 independence of signal size on incident particle rate in pCVD diamond detectors over a range of  
70 particle fluxes up to 20 MHz/cm<sup>2</sup> and on the 3D diamond detectors fabricated in pCVD diamond.

71 **2. Status of the ATLAS Diamond Beam Monitor**

72 The startup of the LHC in 2015 brought a new milestone for diamond detector development  
73 where the first planar diamond pixel modules based on pCVD diamond were installed in an  
74 LHC experiment, the ATLAS experiment [5], and successfully began operation. The ATLAS  
75 Diamond Beam Monitor (DBM) [6, 7] was designed to measure the instantaneous luminosity,  
76 the background rates and the beam spot position. A single DBM module consists of an 18 mm ×  
77 21 mm pCVD diamond 500 μm thick instrumented with a FE-I4 pixel chip [8]. The 26,880 pixels  
78 are arranged in 80 columns on 250 μm pitch and 336 rows on 50 μm pitch resulting in an active  
79 area of 16.8 mm × 20.0 mm. This fine granularity provides high precision particle tracking. The  
80 deposited charge from a particle is measured in the FE-I4 by Time-over-Threshold.

81 The ATLAS DBM uses diamonds with a charge collection distance (the average distance an  
82 electron-hole pair move apart under the influence of the applied electric field) of 200-220 μm at  
83 an applied bias voltage of 500 V. Three telescopes each with 3 diamond DBM modules (plus 1  
84 telescope with silicon sensors) mounted as a three layer tracking device were installed inside the  
85 pixel detector services on each side of the ATLAS interaction point at 90 cm < |z| < 111 cm,

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*Email address:* harris.kagan@cern.ch (H. Kagan)

86  $3.2 < |\eta| < 3.5$  and at a radial distance from 5 cm to 7 cm from the center of the beam pipe. The  
 87 modules have an inclination of  $10^\circ$  with respect to the ATLAS solenoid magnetic field direction  
 88 to suppress erratic dark currents [9] in the diamonds. The ATLAS DBM data-acquisition system  
 89 is shared with the ATLAS IBL [10]. After initial installment, data were collected in the July 2015  
 90 run. These data have been analyzed and the first results of the ATLAS DBM tracking capabilities  
 91 are shown in Fig. 1. A clear separation between background particles from unpaired bunches  
 92 (open circles) and collision particles from colliding bunches (filled circles) is observed. After  
 93 two electrical incidents in 2015 with consequent loss of several silicon and diamond modules,  
 94 the DBM has now been re-commissioned and is again in the operation phase.

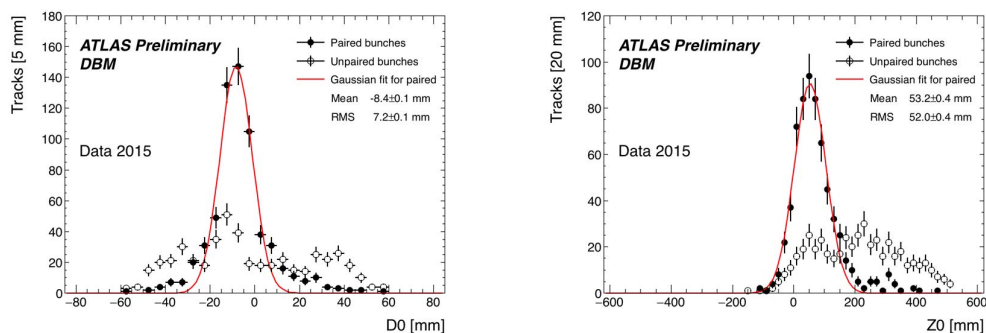


Figure 1: Radial distance (left plot) and longitudinal distance (right plot) of the closest approach of the projected particle tracks to the interaction point as recorded by a single DBM telescope with preliminary alignment.

### 95 3. Rate Studies in pCVD diamond

96 In order to study the dependence of signal size on incident particle rate, RD42 performed a  
 97 series of beam tests in the  $\pi M1$  beam line of the High Intensity Proton Accelerator (HIPA) at  
 98 Paul Scherrer Institute (PSI) [11]. This beam line is able to deliver  $260 \text{ MeV}/c \pi^+$  fluxes from a  
 99 rate of  $\sim 5 \text{ kHz}/\text{cm}^2$  to a rate  $\sim 20 \text{ MHz}/\text{cm}^2$  in bunches spaced 19.8 ns apart.

100 Sensors using pCVD material [12] were tested in a tracking telescope [13] based on  $100 \mu\text{m}$   
 101  $\times 150 \mu\text{m}$  silicon pixel sensors read out by the PSI46v2 pixel chip [14]. The diamond signals  
 102 were amplified with custom-built front-end electronics with a peaking time of  $\sim 6 \text{ ns}$ , return-  
 103 to-baseline in  $\sim 16 \text{ ns}$  and  $550e$  noise with 2 pf input capacitance. The amplified signals were  
 104 recorded with a DRS4 evaluation board [15] operating at 2 GS/s. The entire system was triggered  
 105 with a scintillator which determined the timing of the beam particles with a precision of  $\sim 0.7 \text{ ns}$ .

106 A series of cuts were applied to the data including: removing 60 s of triggers at the beginning  
 107 of each run, removing triggers from heavily ionizing particles with saturated waveforms (mostly  
 108 protons), removing calibration triggers, removing triggers in the wrong beam bucket, removing  
 109 triggers with no tracks in the telescope and removing triggers with large angle tracks in the  
 110 telescope. After applying this procedure all telescope tracks which project into the diamond  
 111 fiducial region have a pulse height well separated from the pedestal distribution in the diamond  
 112 i.e. the diamond is 100% efficient at all rates. The same procedure was applied to all particle  
 113 flux points and the resulting mean pulse height (in arbitrary units) versus rate is shown in Fig. 2  
 114 for both positive and negative bias voltage. The uncertainty on the data points in the plot include  
 115 both statistical and systematic sources. The systematic uncertainty was determined by assuming  
 116 any deviations in pulse height for rates below  $80 \text{ kHz}/\text{cm}^2$  were due to systematic effects. Thus

117 the spread in the data points at a given rate indicates the reproducibility of the data. Fig. 2  
 118 indicates the mean pulse height in pCVD diamond detectors irradiated up to  $5 \times 10^{14}$  n/cm<sup>2</sup> does  
 119 not depend strongly on rate up to rates of 20 MHz/cm<sup>2</sup>.

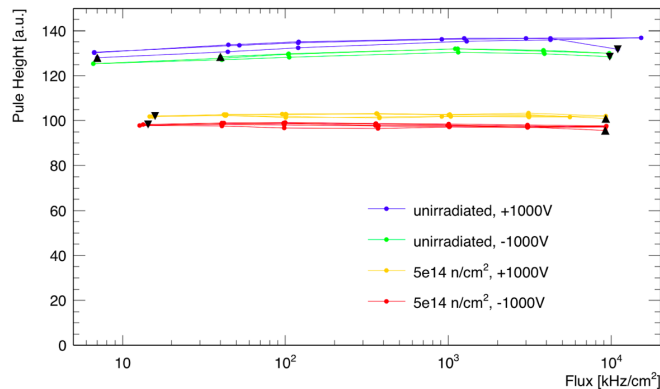


Figure 2: The average pulse height versus rate for an un-irradiated and irradiated pCVD diamond pad detector at positive and negative bias. The beam line parameters were adjusted to set the different particle rates. The data was taken by scanning up and down in rate multiple times. The pulse height units are arbitrary since the un-irradiated and irradiated detectors used different readout electronics. The resulting electronics gain corrections and the relative gain correction for positive versus negative signals in the electronics is still being determined and has not been applied.

#### 120 4. 3D Diamond Pixel Detectors

121 3D sensors with electrodes in the bulk of the sensor material were first proposed in 1997 [16]  
 122 in order to reduce the drift distance of the charged carriers to much less than the sensor thickness.  
 123 In order to achieve this goal a series of alternating + and – electrodes perpendicular to the read  
 124 out face were created in the bulk detector material. This idea is particularly beneficial in detectors  
 125 with a limited mean free path such as trap dominated sensor materials like heavily irradiated  
 126 silicon and pCVD diamond where the observed signal size is related to the mean free path divided  
 127 by the drift distance. Under these circumstances one gains radiation tolerance (larger signals) by  
 128 keeping the drift distance less than the mean free path. With this geometrical structure charge  
 129 carriers drift inside the bulk parallel to the surface over a typical drift distance of 25-100  $\mu$ m  
 130 instead of perpendicular to the surface over a distance of 250-500  $\mu$ m.

131 In 2015 RD42 published results of a 3D device fabricated in single-crystal CVD diamond [17]  
 132 showing that the 3D structure works in diamond. In 2016 RD42 fabricated the first 3D device in  
 133 pCVD diamond [18]. The electrodes in the bulk of the pCVD diamond 3D device were fabricated  
 134 with lasers as described in [17]. The bias electrodes were placed at the corners and the readout  
 135 electrodes were placed in the middle of the cells. This pCVD device was shown to collect more  
 136 than 75% of the deposited charge which translates in more than a factor of two more charge than  
 137 a planar diamond strip detector fabricated on the same pCVD diamond.

138 In 2017 RD42 successfully constructed the first pCVD diamond 3D pixel detector with 50 $\mu$ m  
 139  $\times$  50  $\mu$ m cells. This pixel device is designed to be read out with the RD53 pixel readout chip [19]  
 140 which is not yet available. In order to read this device out with an existing pixel readout chip a  
 141 small number of cells were ganged together to match the pitch of the pixel readout chip. RD42  
 142 is proceeding to make 3D diamond pixel devices compatible with both the CMS pixel readout

143 chip ( $3 \times 2$  ganging) and the ATLAS pixel readout chip ( $1 \times 5$  ganging). The first  $50 \mu\text{m} \times 50$   
 144  $\mu\text{m}$  pCVD diamond 3D pixel device which was bump-bonded used the CMS pixel readout chip.  
 145 This first diamond 3D pixel device was tested during the Aug 2017 beam test at PSI at a  
 146 single voltage and with rates from  $7 \text{ kHz/cm}^2$  to  $7 \text{ MHz/cm}^2$ . During the initial lab test it was  
 147 discovered that the bump bonding had a small issue on one edge. We decided to take data with the  
 148 device rather than try to repair this small bump bonding issue. Fig. 3 (left) shows the preliminary  
 149 efficiency as a function of  $xy$  position for every cell in the device with a  $1500e$  pixel threshold.  
 150 The red box marks the fiducial region used to measure the hit efficiency. The blue circle indicates  
 151 the position of the one non-working pixel cell in the central region of the device. Fig. 3 (right)  
 152 shows the hit efficiency in the fiducial region with the  $1500e$  pixel threshold as a function of time  
 153 during an up-down scan of incident particle rates from  $7 \text{ kHz/cm}^2$  to  $7 \text{ MHz/cm}^2$  and back to  $7$   
 154  $\text{kHz/cm}^2$ . The overall measured efficiency is  $99.2\%$  and no change in efficiency as a function of  
 155 rate is observed. The corresponding efficiency for a planar silicon CMS pixel detector in this test  
 156 was  $99.7\%$  with no change in efficiency as a function of rate. The slight loss of efficiency ( $0.5\%$ ),  
 157 assuming it holds through the completion of the analysis, is most likely due to charge loss in the  
 158 column electrodes. If this explanation is correct, then this effect can be easily remedied by tilting  
 159 the detector at a small angle with respect to the incident beam.

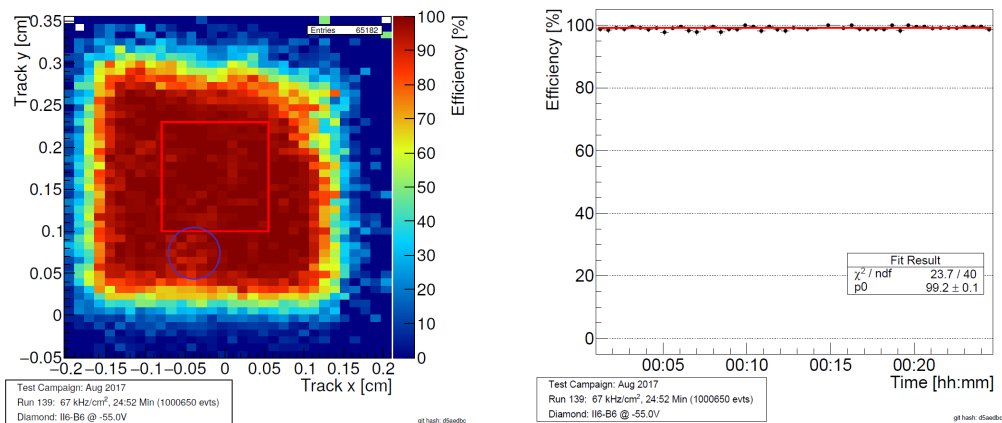


Figure 3: The hit efficiency of the first  $50 \mu\text{m} \times 50 \mu\text{m}$  cell pCVD 3D pixel detector with  $3 \times 2$  ganged cells read out with CMS pixel electronics with a  $1500e$  threshold. The left plot shows the efficiency of each ganged cell as a function of  $xy$  position in the device. The right plot shows the average efficiency in the fiducial region (the red box in the left plot) as a function of time during the run at  $67 \text{ kHz/cm}^2$ .

## 160 5. Conclusions

161 The recent progress in the design, fabrication and testing of polycrystalline CVD diamond  
 162 detectors was presented. The following milestones have been achieved: successful operation of  
 163 the first pCVD diamond planar pixel detector in the ATLAS experiment at the LHC; demonstra-  
 164 tion that the average signal pulse height of pCVD diamond detectors irradiated up to  $5 \times 10^{14}$   
 165  $\text{n/cm}^2$  is independent of the particle flux up to  $\sim 20 \text{ MHz/cm}^2$ ; successful fabrication and opera-  
 166 tion of the first pCVD diamond 3D pixel detector with  $50 \mu\text{m} \times 50 \mu\text{m}$  pixels read out with CMS  
 167 pixel electronics where the efficiency for a MIP was  $>99\%$  and the average charge collected in  
 168 the device was  $>90\%$  of the deposited charge. In the future RD42 plans to study the pulse height  
 169 dependence of CVD diamond sensors with pad and pixel electrodes with radiation doses up to

170  $10^{17}$  n/cm<sup>2</sup> and continue the development of 3D diamond detectors with the production of a 50  
171  $\mu\text{m} \times 50 \mu\text{m}$  cell pCVD diamond 3D pixel detector compatible with ATLAS readout electronics.

## 172 6. Acknowledgments

173 The RD42 Collaboration gratefully acknowledges the staff at CERN for test beam time and  
174 their help in setting up the beam conditions. We would especially like to thank Henric Wilkins,  
175 the test beam coordinator, for his assistance in making our tests a success. We also thank the  
176 beam line staff at the PSI High Intensity Proton Accelerator especially Konrad Deiters, Manuel  
177 Schwarz and Davide Reggiani for their assistance in carrying out the diamond detector tests. We  
178 extend our gratitude to Prof. Lin Li and David Whitehead of the University of Manchester Laser  
179 Processing Center for their assistance in the production of 3D diamond devices. The research  
180 leading to these results received funding from the European Union's Horizon 2020 research and  
181 innovation program under grant agreement No. 654168. This work was also partially supported  
182 by the Swiss National Science Foundation grant #20FL20\_154216, ETH grant 51 15-1, Royal  
183 Society Grant UF120106 and the U.S. Department of Energy through grant DE-SC0010061.

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