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Weak solutions for the Oseen system in 2D and when the given velocity is not sufficiently regular

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Abstract

The aim of this work is twofold: proving the existence of solution $(\boldsymbol{u}, \pi) \in \mathbf{H}^1(\Omega) \times \mathbf{L}^2(\Omega)$ in bounded domains of \mathbb{R}^2 and the whole plane for the Oseen problem (O) for solenoidal vector fields \boldsymbol{v} in $\mathbf{L}^2(\Omega)$, and analyzing the same problem in bounded domains of \mathbb{R}^n for n = 2, 3 when $h = 0, \boldsymbol{g} = \boldsymbol{0}$ and the solenoidal field \boldsymbol{v} belongs to $\mathbf{L}^s(\Omega)$ for s < n.

Keywords: Oseen equations, stationary solutions, irregular data, Hardy spaces 2010 MSC: 35Q35, 25J15, 76D03, 76D07

1. Introduction

This work is dedicated to the study of some existence aspect related to the Oseen problem in bounded domain $\Omega \subset \mathbb{R}^n$, n = 2, 3:

(O) $-\Delta \boldsymbol{u} + \boldsymbol{v} \cdot \nabla \boldsymbol{u} + \nabla \pi = \boldsymbol{f}, \quad \nabla \cdot \boldsymbol{u} = h \text{ in } \Omega, \quad \boldsymbol{u} = \boldsymbol{g} \text{ on } \Gamma.$

In the 3-dimensional case, the existence of weak solutions $(\boldsymbol{u}, \pi) \in \mathbf{H}^{1}(\Omega) \times L^{2}(\Omega)$, regular solution in $\mathbf{H}^{2}(\Omega) \times H^{1}(\Omega)$ and $\mathbf{W}^{1,p}(\Omega) \times L^{p}(\Omega)$ (and intermediate Sobolev spaces) together with the analysis of the existence of very weak solutions in $\mathbf{L}^{p}(\Omega) \times W^{-1,p}(\Omega)$ have been analyzed by the authors in [1], assuming \boldsymbol{v} a solenoidal field belonging to $\mathbf{L}^{s}(\Omega)$ for $s \geq 3$ (from now on, we will denote this solenoidal space by $\mathbf{L}^{s}_{\sigma}(\Omega)$). However, the existence of solution for the 2-dimensional Oseen system has not been attacked in [1] because the "logical" assumption of considering the solenoidal field $\boldsymbol{v} \in \mathbf{L}^{2}(\Omega)$ (in order to obtain weak solutions for (O)) poses some difficulties in the treatment of the convective term $(\boldsymbol{v} \cdot \nabla) \boldsymbol{u}$: On the one hand, it is not clear if the bilinear form associated is coercive and continuous. Some related results can be found in [2] for the scalar case (instead of considering a vector field solution \boldsymbol{u} , one considers a scalar unknown θ) and for $\boldsymbol{g} = \boldsymbol{0}$. On the other hand, when $\Omega = \mathbb{R}^{2}$ an additional awkwardness appears because

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even if we can prove $\nabla u \in \mathbf{L}^2(\mathbb{R}^2)$, it is not evident that $u \in \mathbf{L}^p(\mathbb{R}^2)$ (for any p). Giving a successful answer to both previous problems is our first aim.

Our second aim is to give a first answer to the question of the existence of solution for the Oseen problem (O) when \boldsymbol{v} only belongs to $\mathbf{L}_{\sigma}^{s}(\Omega)$, with s < n and n = 2, 3.

2. Solutions for the Oseen problem in the 2-dimensional case

The existence of weak solutions in $\mathbf{H}^1(\Omega)$ for Problem (O) in 2-dimensional domains is not known when a solenoidal field \boldsymbol{v} that only belongs to $\mathbf{L}^2(\Omega)$ is considered. In this case, the term $(\boldsymbol{v} \cdot \nabla)\boldsymbol{u}$ belongs only to $\mathbf{L}^1(\Omega)$. It is then not clear neither if the bilinear form associated to the Problem (O), with h = 0 and $\boldsymbol{g} = \boldsymbol{0}$:

$$a(\boldsymbol{u}, \boldsymbol{w}) = \int_{\Omega} \nabla \boldsymbol{u} \cdot \nabla \boldsymbol{w} \, d\boldsymbol{x} + \int_{\Omega} (\boldsymbol{v} \cdot \nabla) \boldsymbol{u} \cdot \boldsymbol{w} \, d\boldsymbol{x}$$

is coercive on the space $\mathbf{V}(\Omega) = \{ \boldsymbol{w} \in \mathbf{H}_0^1(\Omega); \text{ div } \boldsymbol{w} = 0 \text{ in } \Omega \}$ nor if it is continuous on $\mathbf{V}(\Omega) \times \mathbf{V}(\Omega)$. In order to overcome this difficulty, we use the Hardy space $\mathcal{H}^1(\mathbb{R}^2)$. One equivalent definition of such a space (in the *n*-dimensional case) is ([3]):

$$\mathcal{H}^{1}(\mathbb{R}^{n}) = \{ f \in L^{1}(\mathbb{R}^{n}), R_{j}f \in L^{1}(\mathbb{R}^{n}), 1 \leq j \leq n \} \text{ where } R_{j} = \frac{\partial}{\partial x_{j}}(-\Delta)^{-1/2}.$$

A partial study of the BMO spaces (Bounded Mean Oscillation) will be also necessary taking into account the duality between \mathcal{H}^1 and the BMO (see [4]). Moreover, the VMO-space (Vanishing Mean Oscillator) is a subspace of the BMO: a function f in $BMO(\mathbf{R}^n)$ is said to be in $VMO(\mathbf{R}^n)$ if

$$\lim_{r \to 0} \sup_{\boldsymbol{x}_0 \in \mathbb{R}^n} \frac{1}{r^n} \int_{B(\boldsymbol{x}_0, r)} |f - \overline{f}| \, \mathrm{d}\boldsymbol{x} = 0, \quad \text{where } \overline{f} = \frac{1}{|B(\boldsymbol{x}_0, r)|} \int_{B(\boldsymbol{x}_0, r)} f.$$

It is also crucial the fact that $H^1(\mathbb{R}^2) \hookrightarrow VMO(\mathbb{R}^2)$ (see [5]).

With these ingredients, we will prove one of the two main results of this work, namely Theorem 2.2 in bounded domains and Theorem 2.5 if $\Omega = \mathbb{R}^2$. In order to prove them, we use the following result:

Lemma 2.1. Assume $\boldsymbol{v} \in \mathbf{L}^2_{\sigma}(\Omega)$ and $\boldsymbol{y} \in H^1_0(\Omega)$. Then $(\boldsymbol{v} \cdot \nabla) \boldsymbol{y} \in H^{-1}(\Omega)$ and

$$\|(\boldsymbol{v}\cdot\nabla)\boldsymbol{y}\|_{H^{-1}(\Omega)} \le C \,\|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)} \|\nabla\boldsymbol{y}\|_{\mathbf{L}^{2}(\Omega)}.$$
(1)

Moreover, we have that

15

20

$$\langle \boldsymbol{v} \cdot \nabla \boldsymbol{z}, \boldsymbol{z} \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_0^1(\Omega)} = 0 \quad \text{for all } \boldsymbol{z} \in \mathbf{H}_0^1(\Omega).$$
 (2)

PROOF. Indeed, considering $\boldsymbol{w} \in \mathbf{L}^2(\mathbb{R}^2)$ the extension of \boldsymbol{v} to \mathbb{R}^2 given by: $\boldsymbol{w} = \boldsymbol{v}$ in Ω , and $\boldsymbol{w} = \nabla \theta$ in $\Omega' = \mathbb{R}^2 \setminus \overline{\Omega}$ where θ is the solution of the following problem:

$$\Delta \theta = 0 \text{ in } \Omega', \quad \frac{\partial \theta}{\partial \boldsymbol{n}} = -\boldsymbol{v} \cdot \boldsymbol{n} \text{ on } \Gamma,$$

2

with $\nabla \theta \in \mathbf{L}^2(\Omega')$ that satisfies

$$\|\nabla \theta\|_{\mathbf{L}^{2}(\Omega')} \leq C \|\boldsymbol{v} \cdot \boldsymbol{n}\|_{H^{-1/2}(\Gamma)} \leq C \|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)}$$

because $\nabla \cdot \boldsymbol{v} = 0$ in Ω . Moreover, $\nabla \cdot \boldsymbol{v} = 0$ in Ω implies $\langle \boldsymbol{v} \cdot \boldsymbol{n}, 1 \rangle_{\Gamma} = 0$ and the existence of θ is ensured by Theorem 3.1 [6]. Observe that $\nabla \cdot \boldsymbol{w} = 0$ in \mathbb{R}^2 because for $\varphi \in \mathcal{D}(\mathbb{R}^2)$:

$$\langle \nabla \cdot \boldsymbol{w}, \varphi \rangle = -\int_{\Omega} \boldsymbol{v} \cdot \nabla \varphi \, d\boldsymbol{x} - \int_{\Omega'} \nabla \theta \cdot \nabla \varphi \, d\boldsymbol{x} = \langle \boldsymbol{v} \cdot \boldsymbol{n}, \varphi \rangle_{\Gamma} - \langle \boldsymbol{v} \cdot \boldsymbol{n}, \varphi \rangle_{\Gamma} = 0,$$

and $\|\boldsymbol{w}\|_{\mathbf{L}^2(\mathbb{R}^2)} \leq C \|\boldsymbol{v}\|_{\mathbf{L}^2(\Omega)}$. On the other hand, we consider \tilde{y} the extension by zero of y that satisfies $\tilde{y} \in H^1(\mathbb{R}^2)$. Using Theorem II.2 point 2) or Theorem II.1 point 2) in [3], we can deduce that $\boldsymbol{w} \cdot \nabla \tilde{y} \in \mathcal{H}^1(\mathbb{R}^2)$ and the bound

$$\|\boldsymbol{w}\cdot\nabla\widetilde{\boldsymbol{y}}\|_{\mathcal{H}^{1}(\mathbb{R}^{2})} \leq C\|\boldsymbol{w}\|_{\mathbf{L}^{2}(\mathbb{R}^{2})}\|\nabla\widetilde{\boldsymbol{y}}\|_{\mathbf{L}^{2}(\mathbb{R}^{2})} \leq C\|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)}\|\nabla\boldsymbol{y}\|_{\mathbf{L}^{2}(\Omega)}.$$

Now, we have to prove that $\boldsymbol{v} \cdot \nabla y \in H^{-1}(\Omega)$ and $\langle \boldsymbol{v} \cdot \nabla y, y \rangle_{H^{-1}(\Omega) \times H^1_0(\Omega)} = 0$. Indeed, for $\varphi \in \mathcal{D}(\Omega)$

$$\begin{split} \int_{\Omega} \varphi \boldsymbol{v} \cdot \nabla y \, d\boldsymbol{x} \bigg| &= \left| \int_{\mathbb{R}^2} \widetilde{\varphi} \boldsymbol{w} \cdot \nabla \widetilde{y} \, d\boldsymbol{x} \right| &\leq \| \boldsymbol{w} \cdot \nabla \widetilde{y} \|_{\mathcal{H}^1(\mathbb{R}^2)} \| \widetilde{\varphi} \|_{BMO(\mathbb{R}^2)} \\ &\leq C \| \boldsymbol{v} \|_{\mathbf{L}^2(\Omega)} \| \nabla y \|_{\mathbf{L}^2(\Omega)} \| \widetilde{\varphi} \|_{H^1(\mathbb{R}^2)} \\ &\leq C \| \boldsymbol{v} \|_{\mathbf{L}^2(\Omega)} \| \nabla y \|_{\mathbf{L}^2(\Omega)} \| \varphi \|_{H^1(\Omega)} \end{split}$$

because $H^1(\mathbb{R}^2) \hookrightarrow VMO(\mathbb{R}^2) \hookrightarrow BMO(\mathbb{R}^2)$. In that way, as $\mathcal{D}(\Omega)$ is dense in $H^1_0(\Omega)$, we can deduce that $\boldsymbol{v} \cdot \nabla \boldsymbol{y} \in H^{-1}(\Omega)$ and estimate (1).

For the proof of (2), let us consider $\boldsymbol{z}_k \in \boldsymbol{\mathcal{D}}(\Omega)$ be such that $\boldsymbol{z}_k \to \boldsymbol{z}$ in $\mathbf{H}_0^1(\Omega)$. Then,

$$egin{aligned} &\langle m{v} \cdot
abla m{z}, m{z}
angle_{\mathbf{H}^{-1}(\Omega) imes \mathbf{H}_0^1(\Omega)} - \langle m{v} \cdot
abla m{z}_k, m{z}_k
angle_{\mathbf{H}^{-1}(\Omega) imes \mathbf{H}_0^1(\Omega)} | \ &\leq |\langle m{v} \cdot
abla (m{z} - m{z}_k), m{z}
angle_{\mathbf{H}^{-1}(\Omega) imes \mathbf{H}_0^1(\Omega)}| + |\langle m{v} \cdot
abla m{z}_k, (m{z}_k - m{z})
angle_{\mathbf{H}^{-1}(\Omega) imes \mathbf{H}_0^1(\Omega)}| \ \end{aligned}$$

Using (1) and the convergence of z_k to z in $\mathbf{H}_0^1(\Omega)$, both duality terms on the righthand-side of the previous inequality tend to 0 when $k \to +\infty$.

Finally, from
$$\langle \boldsymbol{v} \cdot \nabla \boldsymbol{z}_k, \boldsymbol{z}_k \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}^1_{\mathbf{c}}(\Omega)} = 0$$
, we can deduce (2).

Theorem 2.2 (Existence of weak solution for (O)). Let Ω be a Lipschitz bounded domain in \mathbb{R}^2 . Let

$$\boldsymbol{f}\in \mathbf{H}^{-1}(\Omega), \quad \boldsymbol{v}\in \mathbf{L}^2_\sigma(\Omega), \quad h\in L^2(\Omega) \quad \text{and} \quad \boldsymbol{g}\in \mathbf{H}^{1/2}(\Gamma)$$

satisfy the compatibility condition

$$\int_{\Omega} h(\boldsymbol{x}) \, d\boldsymbol{x} = \int_{\partial \Omega} \boldsymbol{g} \cdot \boldsymbol{n} \, d\sigma. \tag{3}$$

Then, the problem (O) has a unique solution $(\mathbf{u}, \pi) \in \mathbf{H}^1(\Omega) \times L^2(\Omega)/\mathbb{R}$. Moreover, there exist some constants $C_1 > 0$ and $C_2 > 0$ such that:

$$\|\boldsymbol{u}\|_{\mathbf{H}^{1}(\Omega)} \leq C_{1} \Big(\|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)} + \big(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)}\big) \big(\|\boldsymbol{h}\|_{L^{2}(\Omega)} + \|\boldsymbol{g}\|_{\mathbf{H}^{1/2}(\Gamma)} \big) \Big),$$
(4)

$$\|\pi\|_{L^{2}(\Omega)/\mathbb{R}} \leq C_{2} \Big(\|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)} + \big(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)}\big) \big(\|h\|_{L^{2}(\Omega)} + \|\boldsymbol{g}\|_{\mathbf{H}^{1/2}(\Gamma)} \big) \Big),$$
(5)

30 where $C_1 = C(\Omega)$ and $C_2 = C_1 (1 + \|\boldsymbol{v}\|_{\mathbf{L}^2(\Omega)}).$

PROOF. Although some parts of this proof are identical to the proof made in [1], we include the whole argument here for completeness. In order to prove the existence of solution, first (using Lemma 3.3 in [7], for instance) we lift the boundary and the divergence data. Then, there exists $\boldsymbol{u}_0 \in \mathbf{H}^1(\Omega)$ such that $\nabla \cdot \boldsymbol{u}_0 = h$ in Ω , $\boldsymbol{u}_0 = \boldsymbol{g}$ on Γ and:

$$\|\boldsymbol{u}_{0}\|_{\mathbf{H}^{1}(\Omega)} \leq C \left(\|h\|_{L^{2}(\Omega)} + \|\boldsymbol{g}\|_{\mathbf{H}^{1/2}(\Gamma)} \right).$$
(6)

Therefore, it remains to find $(\boldsymbol{z}, \pi) = (\boldsymbol{u} - \boldsymbol{u}_0, \pi)$ in $\mathbf{H}_0^1(\Omega) \times L^2(\Omega)$ such that:

$$-\Delta \boldsymbol{z} + \boldsymbol{v} \cdot \nabla \boldsymbol{z} + \nabla \boldsymbol{\pi} = \mathbf{F} \text{ and } \nabla \cdot \boldsymbol{z} = 0 \text{ in } \Omega, \quad \boldsymbol{z} = \mathbf{0} \text{ on } \Gamma.$$
(7)

being $\boldsymbol{F} = \boldsymbol{f} + \Delta \boldsymbol{u}_0 - (\boldsymbol{v} \cdot \nabla) \boldsymbol{u}_0$. From Lemma 2.1, we deduce that $(\boldsymbol{v} \cdot \nabla) \boldsymbol{u}_0 \in \mathbf{H}^{-1}(\Omega)$, then $\boldsymbol{F} \in \mathbf{H}^{-1}(\Omega)$. Since the space $\mathcal{D}_{\sigma}(\Omega) = \{\boldsymbol{\varphi} \in \mathcal{D}(\Omega); \nabla \cdot \boldsymbol{\varphi} = 0\}$ is dense in the space $\mathbf{V}(\Omega)$, the previous problem is equivalent to:

Find
$$\boldsymbol{z} \in \mathbf{V}(\Omega)$$
 such that: $\forall \boldsymbol{\varphi} \in \mathbf{V}(\Omega)$
$$\int_{\Omega} \nabla \boldsymbol{z} \cdot \nabla \boldsymbol{\varphi} \, d\boldsymbol{x} + \langle (\boldsymbol{v} \cdot \nabla) \boldsymbol{z}, \boldsymbol{\varphi} \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_{0}^{1}(\Omega)} = \langle \boldsymbol{F}, \boldsymbol{\varphi} \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_{0}^{1}(\Omega)}.$$

Now, using (2) by Lax-Milgram's Theorem, if we assume that $\mathbf{F} \in \mathbf{H}^{-1}(\Omega)$, then we can deduce the existence of a unique $\mathbf{z} \in \mathbf{H}_0^1(\Omega)$ solution of (7) verifying:

$$\begin{aligned} \|\boldsymbol{z}\|_{\mathbf{H}^{1}(\Omega)} &\leq C \|\boldsymbol{F}\|_{\mathbf{H}^{-1}(\Omega)} \\ &\leq C \Big(\|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)} + (1 + \|\boldsymbol{v}\|_{\mathbf{L}^{2}(\Omega)}) \left(\|\boldsymbol{h}\|_{L^{2}(\Omega)} + \|\boldsymbol{g}\|_{\mathbf{H}^{1/2}(\Gamma)}\right) \Big), \end{aligned}$$
(8)

which added to estimate (6) makes (4). We can recover the pressure π thanks to the De Rham's Lemma (see Lemma 6 in [1] and Corollary III.5.1 in [8]). Now, $-\Delta \boldsymbol{z} + \boldsymbol{v} \cdot \nabla \boldsymbol{z} - \boldsymbol{F} \in \mathbf{H}^{-1}(\Omega)$ and:

$$\forall \boldsymbol{\varphi} \in \mathbf{V}(\Omega), \qquad \langle -\Delta \boldsymbol{z} + \boldsymbol{v} \cdot \nabla \boldsymbol{z} - \boldsymbol{F}, \boldsymbol{\varphi} \rangle_{\mathbf{H}^{-1}(\Omega) \times \mathbf{H}_0^1(\Omega)} = 0.$$

Thanks to De Rham's Lemma, there exists a unique $\pi \in L^2(\Omega)/\mathbb{R}$ such that

$$-\Delta \boldsymbol{z} + \boldsymbol{v} \cdot
abla \boldsymbol{z} +
abla \boldsymbol{\pi} = \mathbf{F}$$

with $\|\pi\|_{L^2(\Omega)/\mathbb{R}} \leq C \|\nabla\pi\|_{\mathbf{H}^{-1}(\Omega)}$. Finally, estimate (5) follows from the previous equation and estimate (8) for \mathbf{z} .

With the same procedure than in [1], we can prove strong and weak- $W^{1,p}(\Omega)$ regularity for (O) in the 2-dimensional bounded case. These results can be stated as follows:

Theorem 2.3 (Existence of strong solution for (O)). Let p > 1,

$$\boldsymbol{f} \in \boldsymbol{L}^p(\Omega), \quad h \in W^{1,p}(\Omega), \quad \boldsymbol{v} \in \mathbf{L}^s_{\sigma}(\Omega) \quad \text{and} \quad \boldsymbol{g} \in \mathbf{W}^{2-1/p,p}(\Gamma)$$

satisfying the compatibility condition (3) with s = 2 if p < 2, s = p if p > 2 and $s = 2 + \varepsilon$ ($\varepsilon > 0$) if p = 2. Then, the unique solution of (O) given by Theorem 2.2 (\mathbf{u}, π) belongs to $\mathbf{W}^{2,p}(\Omega) \times W^{1,p}(\Omega)$, and there exists a constant C > 0 such that:

$$\begin{aligned} \|\boldsymbol{u}\|_{\mathbf{W}^{2,p}(\Omega)} &+ \|\boldsymbol{\pi}\|_{W^{1,p}(\Omega)/\mathbb{R}} \leq C \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{s}(\Omega)}\right) \\ &\times \left(\|\boldsymbol{f}\|_{\mathbf{L}^{p}(\Omega)} + \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{s}(\Omega)}\right) \left(\|\boldsymbol{h}\|_{W^{1,p}(\Omega)} + \|\boldsymbol{g}\|_{\mathbf{W}^{2-1/p,p}(\Gamma)}\right)\right) \\ & 4 \end{aligned}$$

Theorem 2.4. Let

$$p > 1, \quad \boldsymbol{f} \in \mathbf{W}^{-1,p}(\Omega), \quad h \in L^p(\Omega), \quad \boldsymbol{v} \in \mathbf{L}^3_{\sigma}(\Omega) \quad \text{and} \quad \boldsymbol{g} \in \mathbf{W}^{1-1/p,p}(\Gamma)$$

satisfying the compatibility condition (3). Then, the problem (O) has a unique solution $(\boldsymbol{u}, \pi) \in \mathbf{W}^{1,p}(\Omega) \times L^p(\Omega)/\mathbb{R}$, and there exists a constant C > 0 such that:

$$\begin{aligned} \|\boldsymbol{u}\|_{\mathbf{W}^{1,p}(\Omega)} &+ \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{3}(\Omega)}\right)^{\gamma} \|\boldsymbol{\pi}\|_{L^{p}(\Omega)/\mathbb{R}} \leq C \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{3}(\Omega)}\right) \\ &\times \left(\|\boldsymbol{f}\|_{\mathbf{W}^{-1,p}(\Omega)} + \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{3}(\Omega)}\right) \left(\|\boldsymbol{h}\|_{L^{p}(\Omega)} + \|\boldsymbol{g}\|_{\mathbf{W}^{1-1/p,p}(\Gamma)}\right)\right) \end{aligned}$$

 $\text{ so with } \gamma = 0 \text{ if } p \geq 2 \text{ and } \gamma = -1 \text{ if } p < 2.$

If we treat the case of $\Omega = \mathbb{R}^2$, we have to introduce the Sobolev spaces:

$$\begin{split} W_0^{1,2}(\mathbb{R}^2) &= \left\{ \varphi \in \mathcal{D}'(\mathbb{R}^2); \, \frac{\varphi}{w_1} \in L^2(\mathbb{R}^2), \, \nabla \varphi \in L^2(\mathbb{R}^2) \right\}, \\ W_0^{2,2}(\mathbb{R}^2) &= \left\{ \varphi \in \mathcal{D}'(\mathbb{R}^2); \, \frac{\varphi}{w_2} \in L^2(\mathbb{R}^2), \, \frac{\nabla \varphi}{w_1} \in L^2(\mathbb{R}^2, \, \nabla^2 \varphi \in L^2(\mathbb{R}^2) \right\}, \end{split}$$

where $w_1 = (1+|\boldsymbol{x}|) \ln(2+|\boldsymbol{x}|)$ and $w_2 = (1+|\boldsymbol{x}|)^2 \ln(2+|\boldsymbol{x}|)$ (see Definition (7.1),p. 593 in [9]). We denote by $W_0^{-1,2}(\mathbb{R}^2)$ the dual space of $W_0^{1,2}(\mathbb{R}^2)$. Recall ([5]) that the space $W_0^{1,2}(\mathbb{R}^2)$ is densely embedded in $VMO(\mathbb{R}^2)$, and therefore $\mathcal{H}^1(\mathbb{R}^2) = [VMO(\mathbb{R}^2)]' \hookrightarrow W_0^{-1,2}(\mathbb{R}^2)$.

Theorem 2.5 (Case $\Omega = \mathbb{R}^2$). *i*) Let

$$f = \operatorname{div} \mathbb{F}$$
 with $\mathbb{F} \in \mathbf{L}^2(\mathbb{R}^2)$ and $\mathbf{h} \in \mathbf{L}^2(\mathbb{R}^2)$.

Then, the problem (O) has a unique solution (\mathbf{u}, π) satisfying $\mathbf{u} \in \mathbf{W}_0^{1,2}(\mathbb{R}^2)$ and $\pi \in L^2(\mathbb{R}^2)$, where π is unique and \mathbf{u} is unique up to an additive constant vector field. ii) Moreover, if

$$\boldsymbol{f} \in \mathcal{H}^1(\mathbb{R}^2)$$
 and $\nabla h \in \mathcal{H}^1(\mathbb{R}^2),$

then

$$\nabla^2 \boldsymbol{u} \in \mathcal{H}^1(\mathbb{R}^2), \quad \nabla \pi \in \mathcal{H}^1(\mathbb{R}^2), \quad \nabla \boldsymbol{u} \in \boldsymbol{L}^{2,1}(\mathbb{R}^2) \quad and \quad \boldsymbol{u} \in L^{\infty}(\mathbb{R}^2), \tag{9}$$

being $L^{2,1}(\mathbb{R}^2)$ is the Lorentz space of all measurable functions f satisfying

$$\int_0^\infty t^{-1/2} f^*(t) \, dt < +\infty,$$

where the rearrangement function f^* is defined by $f^*(t) = \sup\{s \in (0,\infty); \mu(\{x \in \mathbb{R}^2; |f(x)| > s\}) > t\}$, for μ the Lebesgue measure on \mathbb{R}^2 .

PROOF. i) Existence: Let $\chi \in W_0^{2,2}(\mathbb{R}^2)$ be the unique solution, up to a polynomial function of degree one, of $\Delta \chi = h$ in \mathbb{R}^2 (see Theorem 9.6 in [9]). Then, we take $\boldsymbol{u}_h = \nabla \chi \in \boldsymbol{W}_0^{1,2}(\mathbb{R}^2)$. Problem (O) is then written as:

$$-\Delta \boldsymbol{z} + \boldsymbol{v} \cdot \nabla \boldsymbol{z} + \nabla \pi = \mathbf{k}, \quad \nabla \cdot \boldsymbol{z} = 0 \quad \text{in } \mathbb{R}^2,$$

with $\mathbf{k} = \mathbf{f} + \Delta \mathbf{u}_h - \mathbf{v} \cdot \nabla \mathbf{u}_h$. Because of $(\mathbf{v} \cdot \nabla) \mathbf{u}_h \in \mathcal{H}^1(\mathbb{R}^2) \hookrightarrow W_0^{-1,2}(\mathbb{R}^2)$, by using Lax-Milgram's Lemma (as in the bounded case) we can deduce the existence of a solution $\mathbf{z} \in \mathbf{W}_0^{1,2}(\mathbb{R}^2)$ with $\nabla \cdot \mathbf{z} = 0$, unique up to a constant vector of \mathbb{R}^2 . Lax-Milgram's

⁴⁵ Lemma hypotheses are satisfied because, on the one hand, we know that the quotient norm $\|\boldsymbol{z}\|_{\boldsymbol{W}_0^{1,2}(\mathbb{R}^2)/\mathbb{R}^2}$ is equivalent to that one defined as $\|\nabla \boldsymbol{z}\|_{\boldsymbol{L}^2(\mathbb{R}^2)}$, and, on the other hand, $(\boldsymbol{v} \cdot \nabla) \boldsymbol{z} \in \mathcal{H}^1(\mathbb{R}^2)$ for any $\boldsymbol{z} \in \boldsymbol{W}_0^{1,2}(\mathbb{R}^2)$. The pressure can be recovered by using Theorem 1 in [10].

ii) Regularity: Assume that $\mathbf{f} \in \mathcal{H}^1(\mathbb{R}^2)$ and $\nabla h \in \mathcal{H}^1(\mathbb{R}^2)$ (which, in particular, implies that $h \in L^{2,1}(\mathbb{R}^2)$). Therefore,

$$-\Delta oldsymbol{u}+
abla \pi=oldsymbol{f}-oldsymbol{v}\cdot
abla u\in\mathcal{H}^1(\mathbb{R}^2) \quad ext{and} \quad
abla \cdotoldsymbol{u}=h.$$

By using Theorem 3.14 in [10], one deduces (9).

⁵⁰ 3. The Oseen problem in bounded domains for a less regular v

The aim of this section is the analysis of the existence of solutions of (O) in a bounded domain (n = 2 or 3) when $\boldsymbol{v} \in \mathbf{L}^s_{\sigma}(\Omega)$ for s < n. We analyze the case for $\boldsymbol{f} \in \mathbf{H}^{-1}(\Omega)$, h = 0 and $\boldsymbol{g} = \boldsymbol{0}$. Observe that the term $(\boldsymbol{v} \cdot \nabla)\boldsymbol{u}$ can also be written as $\nabla \cdot (\boldsymbol{u} \otimes \boldsymbol{v})$. The proof of Theorem 3.2 (n = 2) applies directly from that one of Theorem 3.1 (n = 3).

Theorem 3.1. Let $\Omega \subset \mathbb{R}^3$ a Lipschitz bounded domain,

$$\boldsymbol{f} \in \mathbf{H}^{-1}(\Omega), \quad h = 0, \quad \boldsymbol{g} = \boldsymbol{0} \quad \text{and} \quad \boldsymbol{v} \in \boldsymbol{L}_{\sigma}^{6/5+\alpha}(\Omega)$$

for any $0 < \alpha \leq 9/5$. Then, there exists a solution of (O) such that $(\boldsymbol{u}, \pi) \in \mathbf{H}_0^1(\Omega) \times L^{q(\alpha)}(\Omega)/\mathbb{R}$ for $q(\alpha) = (6(6+5\alpha))/(36+5\alpha)$ with the estimate:

$$\|\boldsymbol{u}\|_{\mathbf{H}^{1}(\Omega)} + \|\boldsymbol{\pi}\|_{L^{q(\alpha)}(\Omega)/\mathbb{R}} \leq C \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{6/5+\alpha}(\Omega)}\right) \|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)}$$
(10)

PROOF. We approximate \boldsymbol{v} by $\boldsymbol{v}_{\lambda} \in \boldsymbol{\mathcal{D}}_{\sigma}(\overline{\Omega})$ in the $\mathbf{L}^{6/5+\alpha}(\Omega)$ -norm and look for the solution of the problem:

$$(O_{\lambda}) \quad -\Delta \boldsymbol{u}_{\lambda} + \nabla \cdot (\boldsymbol{u}_{\lambda} \otimes \boldsymbol{v}_{\lambda}) + \nabla \pi_{\lambda} = \boldsymbol{f} \quad \text{and} \quad \nabla \cdot \boldsymbol{u}_{\lambda} = 0 \quad \text{in } \Omega, \quad \boldsymbol{u}_{\lambda} = \boldsymbol{0} \quad \text{on } \Gamma$$

Taking u_{λ} as test function in (O_{λ}) , we get the estimate:

$$\|\boldsymbol{u}_{\lambda}\|_{\mathbf{H}_{0}^{1}(\Omega)} \leq C(\Omega) \|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)}.$$
(11)

By De Rham Theorem, there exists $\pi_{\lambda} \in L^2(\Omega)$ (unique up to a constant) such that:

$$abla \pi_{\lambda} = \boldsymbol{f} + \Delta \boldsymbol{u}_{\lambda} - \nabla \cdot (\boldsymbol{u}_{\lambda} \otimes \boldsymbol{v}_{\lambda}).$$

Moreover, $\boldsymbol{v}_{\lambda} \otimes \boldsymbol{u}_{\lambda}$ belongs to a bounded set of $\mathbb{L}^{q(\alpha)}(\Omega)$ with $q(\alpha) = (6(6+5\alpha))/(36+5\alpha)$ and which implies that $\nabla \cdot (\boldsymbol{v}_{\lambda} \otimes \boldsymbol{u}_{\lambda})$ belongs to a bounded set of $\mathbf{W}^{-1,q(\alpha)}(\Omega)$. Note that if $0 < \alpha \leq 9/5$ then $1 < q(\alpha) \leq 2$. Using (11),

$$\begin{aligned} \|\nabla \pi_{\lambda}\|_{W^{-1,q(\alpha)}(\Omega)} &\leq C_{1}\left(1+C(\Omega)\right)\|f\|_{\mathbf{H}^{-1}(\Omega)}+\|\boldsymbol{u}_{\lambda}\otimes\boldsymbol{v}_{\lambda}\|_{\mathbb{L}^{q(\alpha)}(\Omega)} \\ &\leq C_{1}\left(1+C(\Omega)\right)\|f\|_{\mathbf{H}^{-1}(\Omega)}+C_{2}\|\boldsymbol{v}_{\lambda}\|_{\mathbf{L}^{6/5+\alpha}(\Omega)}\|\boldsymbol{u}_{\lambda}\|_{\mathbf{H}^{1}_{0}(\Omega)} \\ &\leq C(\Omega)\left(1+\|\boldsymbol{v}\|_{\mathbf{L}^{6/5+\alpha}(\Omega)}\right)\|f\|_{\mathbf{H}^{-1}(\Omega)} \end{aligned}$$
(12)

where C_1 and C_2 are the constant of the Sobolev embeddings $\mathbf{H}^{-1}(\Omega) \hookrightarrow \mathbf{W}^{-1,q(\alpha)}(\Omega)$ and $\mathbf{H}_0^1(\Omega) \hookrightarrow \mathbf{L}^6(\Omega)$, respectively. Therefore, from (12) we obtain:

$$\inf_{K \in \mathbb{R}} \|\pi_{\lambda} + K\|_{L^{q(\alpha)}(\Omega)} \le C(\Omega) \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{6/5+\alpha}(\Omega)}\right) \|f\|_{\mathbf{H}^{-1}(\Omega)}$$

Now, it is necessary to take the limit when $\lambda \to 0$: We can extract a subsequence of (u_{λ}) and $(\pi_{\lambda} + C_{\lambda})$ (that will be called in the same way that the original one) such that:

$$\boldsymbol{u}_{\lambda} \rightharpoonup \boldsymbol{u} \quad \text{in } \mathbf{H}_{0}^{1}(\Omega), \qquad \pi_{\lambda} + C_{\lambda} \rightharpoonup \pi \quad \text{in } L^{q(\alpha)}(\Omega),$$

⁵⁵ where (\boldsymbol{u}, π) is solution of (O) and satisfies (10).

Theorem 3.2. Let $\Omega \subset \mathbb{R}^2$ a Lipschitz bounded domain, $\mathbf{f} \in \mathbf{H}^{-1}(\Omega)$, h = 0, $\mathbf{g} = \mathbf{0}$ and $\mathbf{v} \in \mathbf{L}^{1+\alpha}_{\sigma}(\Omega)$ with $0 < \alpha \leq 1$. Then, there exists a solution of (O) such that $(\mathbf{u}, \pi) \in \mathbf{H}^{1}_{0}(\Omega) \times L^{q(\beta)}(\Omega)/\mathbb{R}$ for $q(\beta) = 1 + \beta$, for any $0 < \beta < \alpha$, with the estimate:

$$\|\boldsymbol{u}\|_{\mathbf{H}^{1}(\Omega)} + \|\pi\|_{L^{q(\beta)}(\Omega)/\mathbb{R}} \leq C \left(1 + \|\boldsymbol{v}\|_{\mathbf{L}^{1+\alpha}(\Omega)}\right) \|\boldsymbol{f}\|_{\mathbf{H}^{-1}(\Omega)}$$

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