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# Lenses on very curved zones of a singular line field of $\mathbb{C}^{2}$ or of a singular plane field of $\mathbb{C}^{3}$ 

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#### Abstract

We renormalize, using suitable lenses, small domains of a singular holomorphic line field of $\mathbb{C}^{2}$ or plane field of $\mathbb{C}^{3}$ where the curvature of a plane-field is concentrated. At a proper scale the field is almost invariant by translations. When the field is integrable, the leaves are locally almost translates of a surface that we will call profile. When the singular rays of the tangent cone (a generalization to a plane-field of the tangent cone of a singular surface is defined) are isolated, we obtain more precise results. We also generalize a result of Merle ([Me]) concerning the contact order of generic polar curves with the singular level $f=0$ when $\omega=d f$. On the way we obtain some classical results (Lê's carousels) on the knot $K=\left(\{f=0\} \cap B_{\varepsilon}(0,0,0)\right)$ in dimension 2 an a maybe less classical ones in dimension 3.


key words
complex polynomial, complex one-form, isolated singularity, polar curve, profile

## domains

[math.AG], [math.IT], [math.GT], [math.DG]

## 1 Introduction



Figure 1: $\left\{x^{3}-y^{2}=0\right\}$ and $\left\{x^{3}-y^{2}=\lambda\right\}$, a real picture.
In 1968 J. Milnor [Mil], published a book where he shows that levels of a complex polynomial $f: \mathbb{C}^{n} \rightarrow \mathbb{C}$ present some limit topology near an isolated singular point. Ness [Ne], Langevin [La1] studied the total curvature of the intersection of a level of $f$ with a small ball centered at the singular point. Then, studying polynomials of two variables, Teissier [Tei3],

Garcia-Barroso [GarBar] and Garcia-Barroso and Teissier [GarBar-Te] got more and more precise results about where curvature is situated on a given level $\{f=\lambda\}, \lambda$ close to 0 . The first results, in dimension 2 , about the pattern of the very curved zones of levels $\{f=\lambda\}, \lambda$ close to 0 , or of the leaves of a singular foliation defined by an integrable algebraic 1 -form, were obtained by the author and J.C Sifre (see [La-Si]).

The results given here are also closely related to Merle's ([Me]) about bouquets (see Definition 3.2.2) of branches of polar curves in dimension 2.

The goal of this text is, given a 1 -form in $\mathbb{C}^{3}$ with an isolated singular point at the origin, to renormalize small domains where the curvature is concentrated in order to find profiles, that is limit non-trivial plane field or foliation (see Definition 3.2.3). The choice of the dimension $n=3$ of the ambient space is first intended to avoid heavy notations. It is also the last dimension where the author has still some geometrical intuition. When $\omega=d f$, renormalizing simultaneously $d f$ and $f$, we prove that the profiles are in this case obtained translating the graph of a polynomial. We also observe the contact order of generic polar curves and the singular level $f=0$, generalizing a result of Merle ([Me]). We obtain also a decomposition of the link $K=\{f=0\} \cap S_{\varepsilon}$ in pieces which are either $\mathbb{S}^{1}$ fiber bundles or fiber bundles over $\mathbb{S}^{1}$. The pieces are glued along boundary tori.
The figures implying the use of a computer were made by J-C. Sifre and are taken from [La-Si] or [La2].

## 2 Preliminaries

### 2.1 Gauss map and polar curves (in dimension 2)

Given a 1-form $\omega$ defined on $\mathbb{R}^{2}$, we can define, at a point $m$ where $\omega \neq 0$, a Gauss map $G_{\omega}(m)=\operatorname{ker} \omega(m)$. It has values in $\mathbb{R} \mathbb{P}^{1}$. The line field ker $\omega$ defines, where $\omega \neq 0$, a foliation $\mathcal{F}_{\omega}$. The Gauss map $G_{\omega}$ depends only on the foliation.

We can now define the polar sets of a foliation $\mathcal{F}_{\omega}$ of $\mathbb{R}^{2}$ by curves. They are inverse images of the points of $\mathbb{R} \mathbb{P}^{1}$ by the Gauss map $G_{\omega}$. These polar sets are in general curves, maybe with singularities. That is why we will abusively use the term polar curves.

Definition 2.1.1. - The polar curve $\Gamma_{[a, b]}$ is the closure of the set (in general a curve, maybe with singularities) of points of the plane where the line tangent at a point $m$ to the leaf $L_{m}$ of the foliation which contains the point $m$ is parallel to a direction $[a, b]$.

- When the foliation is defined by a 1 -form $\omega$, it is the closure of the set $\{\omega(a, b)=0\}$

This definition holds as well when the foliation $\mathcal{F}$ is a foliation of $\mathbb{C}^{2}$ by complex curves. The second part of this definition holds also when the 1 -form in $\mathbb{C}^{n}$ or $\mathbb{R}^{n}$ defines only a hyperplane field.

### 2.2 Examples in dimension 2

## Polar curves near an isolated singularity, the seminal example in dimension 2

Let $\ell$ be the line generated by the vector $(a, b)$. The equation of the polar curve $\Gamma_{\ell}$ is then $d f(a, b)=0$. Here it writes $3 a x^{2}-2 b y=0$. The generic $(b \neq 0)$ polar curves form a family of parabolas tangent to the $x$-axis and the $x$-axis.


Figure 2: Polar curves of $x^{3}-y^{2}=\lambda$ using a suitable lens, the line-field tangent to $\mathcal{F}$ after enlargement

Notice that, when one observes a family of tangent generic polar curves $\Gamma_{[a, b]}$, they look like lines parallel to the common tangent direction after enlarging enough a small enough neighborhood of a point on the common tangent at the origin close to the origin (see Figure 2).

By definition of polar curves, the direction tangent to $\mathcal{F}$ at the points of a given segment of polar curve are parallel. If the direction $[a, b]$ is not tangent to $\Gamma_{[a, b]}$ at the origin, the integration of the line field will give pieces looking like parallel graphs (maybe with asymptotes parallel to the $x$-axis, the saddle-node of equation $\omega=y d x-x^{2} d y$ is already an example of this phenomenon in $\mathbb{R}^{2}$ or in $\mathbb{C}^{2}$.

Our tool providing a profile is a Newton lens.
Definition 2.2.1. A Newton lens is a pair

- an analytic curve $\gamma(t)$ ending at the origin and tangent at the origin to a privileged direction, say the $x$-axis,
- an enlargement rate $1 / \rho=1 / \rho(x(\gamma(t))), \lim _{|x(t)| \rightarrow 0} \rho(x(t))=0$

The analytic curve will first be the $x$-axis, then suitably chosen curves tangent to the $x$-axis at the singular point.

The definition will stay unchanged in $\mathbb{C}^{3}$.

### 2.3 Newton-lens cloud and Newton-lens polygon

The Newton cloud of a polynomial $f(x, y)=\sum a_{i, j} x^{i} y^{j}$ is the set $\{i, j\}$ of points of $\mathbb{N} \times \mathbb{N}$ such that $a_{i, j} \neq 0$ (see Figure 3).

The Newton polygon of $f$ is the boundary of the convex hull of the union of the upper quadrants of $\mathbb{N} \times \mathbb{N}$ of vertices the points of the Newton cloud of $f$ (see Figure 3).

Let $\Phi_{1}$ be the change of variables

$$
\begin{aligned}
& x=x_{0}+\rho_{1} X_{1} \\
& y=\rho_{1} Y_{1}
\end{aligned}
$$

Then, we can see this change of variables as a moving lens settled to observe a neighborhood of the point $\left(x_{0}, 0\right)$ if $\rho_{1}$ is a function of $x_{0} \rightarrow 0$.


Figure 3: Newton cloud (black dots) and Newton polygon (in red) of $f(x, y)=x^{3}-y^{2}+x^{2} y^{2}+x^{2} y$


Figure 4: First Newton-lens cloud (black dots and green stars) and first Newton-lens polygon (in red) of $f(x, y)=x^{3}-y^{2}+x^{2} y^{2}+x^{2} y$

Definition 2.3.1. The first Newton-lens cloud of $f$ is the Newton cloud associated to the polynomial $\Phi_{1}^{*}(f)$ of the variables $x_{0}$ and $\rho_{1}$.

The first Newton-lens polygon of $f$ is the boundary of the convex hull of the union of the upper quadrants of $\mathbb{N} \times \mathbb{N}$ of vertices the points of the first Newton cloud of $f$ (see Figures 3 and 4).

We will first consider the exponents of the polynomial $\Phi_{1}^{*}(f)$ considered as a polynomial of the two variables $x_{0}$ and $\rho_{1}$.

Then the Newton-lens polygon allows us to chose the value $\rho_{1}=x_{0}^{r_{1}}$. Let us denote by $\widetilde{\Phi_{1}}$
the change of variables

$$
\begin{aligned}
& x=x_{0}+x_{0}^{r_{1}} X_{1} \\
& y=x_{0}^{r_{1}} Y_{1}
\end{aligned}
$$

Then we will iterate the construction. After replacing $\rho_{i-1}$ by a rational power of $x_{0}$, obtaining a change of variables $\widetilde{\Phi_{i-1}}$ we consider the change of variables $\Phi_{i}$

$$
\begin{aligned}
X_{i-1} & =\rho_{i} X_{i} \\
Y_{i-1} & =y_{i}+\rho_{i} Y_{i}
\end{aligned}
$$

We consider the exponents in $x_{0}$ and $\rho_{i}$ of the polynomial $\Phi_{i}^{*}{\widetilde{\Phi_{i-1}}}^{*} \cdots{\widetilde{\Phi_{1}}}^{*}(f)$. The Newton-lens clouds are now in $\mathbb{Q}^{+} \times \mathbb{N}$.

Then we will apply a similar construction to a polynomial one-form $\omega=A d x+B d y$ plotting, after changes of variables, all the exponents of the monomials of $A$ and of $B$.

We also apply this constructions to a polynomial $f(x, y, z)$ and a polynomial one-form $\omega=$ $A(x, y, z) d x+B(x, y, z) d y+C(x, y, z) d z$. In this case, as it would be with more variables, the Newton-lens clouds are still planar (contained in $\mathbb{Q}^{+} \times \mathbb{N}$ ) after the first step.

## 3 1-forms in dimension 3

The advantage of our method is to provide, in particular when $\omega$ is integrable, a precise description of the limit shapes of the leaves which appear near the origin. When $\omega=d f$, we unfortunately loose the global structure of the levels of $f, f=\lambda, \lambda \rightarrow 0$.

Let us now consider a 1 -form $\omega=A d x+B d y+C d z$.
Let $\Sigma$ be the singular set of $\omega$ that is where $\omega$ is zero. The planes $\{\operatorname{ker}(\omega)\}$ define a plane field $\mathcal{P}$ of $\mathbb{C}^{3} \backslash \Sigma$.

In this text, we assume that the origin $O$ is an isolated singular point of $\omega$.

### 3.1 The tangent cone

Inspired by Euler's formula valid when $\omega=d f$, we define the tangent cone of a 1 -form $\omega$. Let low $(\omega)$ be the homogeneous form selecting globally the lowest degree terms of the polynomials $A, B$ and $C$ coefficients of $\omega=A d x+B d y+C d z$; if this lowest degree is $k$, then valA $\geq k$; let $A_{k}(x, y)=\operatorname{low}(A)$ if the valuation of $A$ is $k$, zero if the valuation of $A$ is larger than $k$; we use the same convention for $B$ and $C$.

Definition 3.1.1. The equation of the tangent cone of $\omega=A(x, y, z) d x+B(x, y, z) d y+C(x, y, z) d z$ at the origin is

$$
\operatorname{low}(\omega(x, y, z))(x, y, z))=A_{k}(x, y, z) x+B_{k}(x, y, z) y+C_{k}(x, y, z) z=0
$$

Recall that, given in $\mathbb{C}^{2}$ the 1-form $\omega=A(x, y) d x+B(x, y) d y$, the polynomials $x A(x, y)$ and $y B(x, y)$ appear in [Ca-Li-Sa]. Notice that, when $\omega$ as an homogeneous 1 -form, the plane field $\operatorname{ker} \omega$ is constant along rays. In particular, along rays of the tangent cone, one has $\omega(x, y, z)=0$, that is the plane $\operatorname{ker} \omega(x, y, z)$ contains the ray $\lambda \cdot(x, y, z), \lambda \in \mathbb{C}$.

Notice also that, when $\omega=d f, \operatorname{low}(\omega(x, y, z))(x, y, z)=k \cdot \operatorname{low}(f)$, where $k$ is the degree of $\operatorname{low}(f)$ (Euler's equality).

## Example: the tangent cone of a linear forms

Let us suppose that $\omega=A d x+B d y+C d z$,

$$
\begin{aligned}
A & =a_{1} x+b_{1} y+c_{1} z \\
B & =a_{2} x+b_{2} y+c_{2} z \\
C & =a_{3} x+b_{3} y+c_{3} z
\end{aligned}
$$

Let us denote by $\mathcal{M}$ the matrix $\left(\begin{array}{lll}a_{1} & b_{1} & c_{1} \\ a_{2} & b_{2} & c_{2} \\ a_{3} & b_{3} & c_{3}\end{array}\right)$ we can write $\omega=\left(\begin{array}{lll}d x & d y & d z\end{array}\right) \cdot \mathcal{M} \cdot\left(\begin{array}{l}x \\ y \\ z\end{array}\right)$.
The equation of the tangent cone is then $\left(\begin{array}{lll}x & y & z\end{array}\right) \cdot \mathcal{M} \cdot\left(\begin{array}{l}x \\ y \\ z\end{array}\right)$.
It is a quadratic equation which depends only on symmetrization $\frac{1}{2}\left({ }^{t} \mathcal{M}+\mathcal{M}\right)$.
When the matrix $\mathcal{M}$ is antisymmetric, say $\mathcal{M}=\left(\begin{array}{ccc}0 & -a_{2} & -a_{3} \\ a_{2} & 0 & -b_{3} \\ a_{3} & b_{3} & 0\end{array}\right)$.
The form $\omega=\left(\begin{array}{lll}d x & d y & d z\end{array}\right) \cdot \mathcal{M} \cdot\left(\begin{array}{l}x \\ y \\ z\end{array}\right)$ is an example where the tangent cone is the whole space as

$$
\mathfrak{t}\left(\left(\begin{array}{ccc}
d x & d y & d z
\end{array}\right) \mathcal{M} \cdot\left(\begin{array}{l}
x \\
y \\
z
\end{array}\right)\right)\left(\begin{array}{l}
x \\
y \\
z
\end{array}\right)=(-1)^{3}\left(\left(\begin{array}{lll}
d x & d y & d z
\end{array}\right) \mathcal{M} \cdot\left(\begin{array}{l}
x \\
y \\
z
\end{array}\right)\right)\left(\begin{array}{l}
x \\
y \\
z
\end{array}\right) .
$$

Notice that all rays are separatrices (see Definition 7.5.1).
In fact it is the general form of a linear 1-form such that the tangent cone is the whole space, as in that case $\left(\mathcal{M}+{ }^{\mathbf{t}} \mathcal{M}\right)$ should be zero.

Remark 3.1.2. A parenthesis in $\mathbb{R}^{3}$ When the matrix $\mathcal{M}$ is of rank 3, the integrability condition on the whole $\mathbb{R}^{3}, \omega \wedge d \omega=0 \forall(x, y, z)$, writes $\mathcal{M}-{ }^{\mathbf{t}} \mathcal{M}=0$, in other words, the matrix is a symmetrical one.

Remark When the image of the Gauss map is exactly a projective line $\mathcal{L}$, the foliation defined by ker $\omega$ is a family of planes rotating around an axis.

When $\omega$ is linear, we can represent the Gauss map using the matrix of a map $\mathcal{M}: \mathbb{C}^{3} \rightarrow \mathbb{C} 3$ as $\operatorname{ker} \omega=\left(\mathcal{M} \cdot\left(\begin{array}{l}x \\ y \\ z\end{array}\right)\right)^{\perp}$.

When the Gauss map is of rank 1, and the linear map representing it of rank 2, the kernel of the matrix $\mathcal{M}$ is a line $L$. All the planes $\operatorname{ker} \omega$ contain the line $L$. These planes form the projective line $\mathcal{L}$. The form $\omega$ defines a pencil of planes (see Figure ??).
Remark Notice that, even if the tangent cone is a plane, the map $\gamma_{\omega_{\text {low }}}$ maybe locally surjective and guarantees, when $\omega$ is homogeneous, that all the polar curves are rays. This is the case of $\omega=\omega_{\text {low }}=x d x+z d y-y d z$ (the tangent cone has equation $x^{2}=0$, it is singular, see below Definition 3.1.3).

### 3.1.1 Singular rays of the tangent cone in dimension 3

Definition 3.1.3. A ray $\{\lambda \cdot(x, y, z), \lambda \in \mathbb{C}\}$ of this cone is singular for $\omega$ if $A_{k}(x, y, z)=$ $B_{k}(x, y, z)=C_{k}(x, y, z)=0$, that is if $\operatorname{low}(\omega)(x, y, z) \equiv 0$.

Definition 3.1.4. The equation of the tangent cone degenerates when all the lines of the tangent cone are singular (see Definition 3.1.3).

This is the case when $\omega=d f$ and if $f=g^{2}$.
The equation of the tangent cone provides a projective curve $\mathcal{C}_{\omega} \subset \mathbb{C P}^{2}$.
Proposition 3.1.5. When $\omega=d f$, a line of the tangent cone is singular for $\omega$ if it is a singular point of $C_{\omega}$.

Proof: Let us suppose that the $x$-axis is a singular ray of the tangent cone of $\omega=d f$, and that $f_{\text {low }}$ is homogeneous of degree $(k+1)$. It means that $\omega_{\text {low }}(1,0,0)=d\left(f_{\text {low }}\right)(1,0,0)=0$. The section of the tangent cone by the plane $x=1$ has the equation $F(y, z)=A_{k}(1, y, z)+y B_{k}(1, y, z)+$ $z C_{k}(1, y, z)=0$. The gradient of the function $F$ writes

$$
\begin{aligned}
& \frac{\partial A(1, y, z)}{\partial y}+B(1, y, z)+y \frac{\partial B}{\partial y}+z \frac{\partial C}{\partial y} \\
& \frac{\partial A(1, y, z)}{\partial z}+y \frac{\partial B}{\partial z}+C(1, y, z)+z \frac{\partial C}{\partial z}
\end{aligned}
$$

The point $(1,0,0)$ is a singular point of the curve of equation $F=0$ if

$$
\frac{\partial A(1, y, z)}{\partial y}(1,0,0)+B(1,0,0)=0=\frac{\partial A(1, y, z)}{\partial z}(1,0,0)+C(1,0,0)
$$

As the $x$-axis is singular we know already that $B(1,0,0)=C(1,0,0)=0$.
The term of $f$ which may contribute to $\frac{\partial A(1, y, z)}{\partial y}(1,0,0)$ and $\frac{A(1, y, z)}{\partial z}(1,0,0)$ are respectively of the form $x^{k}(a y)$ and $x^{k}(b z)$. But if $a \neq 0$ then $\frac{\partial f}{\partial y}(1,0,0)=B(1,0,0) \neq 0$ Similarly, if $b \neq 0$, $\frac{\partial f}{\partial z}(1,0,0)=C(1,0,0) \neq 0$. Therefore $a=b=0$ and the point $(1,0,0)$ is a critical point of $F=0$ ).

In our search of profiles, the pertinent objects are the rays singular for $\omega$.
Example of the type $\omega=d f: x^{4}+y^{4}-x \cdot y^{2}+z^{3}=0$.
The tangent cone has the equation $z^{3}-x y^{2}=0$. As $\operatorname{low}(\omega)=-y^{2} d x-2 x y d y+3 z^{2} d z$ we see that the only degenerate ray is the $x$-axis.

The reader will find below a profile associated to this example.

### 3.2 Gauss map, polar surfaces and polar curves in dimension 3

Given a 1 -form $\omega$ defined on $\mathbb{C}^{3}$, we can define, at a point $m$ where $\omega \neq 0$, a Gauss map $G_{\omega}(m)=\operatorname{ker} \omega(m)$. It has values in $\mathbb{C P}^{2}$. The map $G_{\omega}$ is defined on the complement of the singular locus of $\omega$. It is therefore defined on a small enough neighborhood of the origin but for the origin itself.

## Definition 3.2.1.

- The polar set $V_{\ell}=V_{v}$ associated to a line $\ell$ generated by a vector $v$ is the closure of the set of points where $\omega(v)=0$. It is in general a surface.
- The polar set $\Gamma_{h}$ associated to a plane $h$ is the closure of the set of points where $k e r \omega=h$. In general it is a curve.

We will use only the terms "polar surface", "polar curve" and deal with degenerate cases only if unavoidable.

In other terms the polar curve $\Gamma_{h}$ is the inverse image of $h$ by the Gauss map $G_{\omega}$. A polar curve $\Gamma_{h}, h$ a plane of equation $A x+B y+C z=0$, is the intersection of two polar surfaces $V_{a_{i}, b_{i}, c_{i}}$ such that the vectors $\left(a_{1}, b_{1}, c_{1}\right)$ and $\left(a_{2}, b_{2}, c_{2}\right)$ span the plane $h$.

An example of very degenerate case is the form $\omega=x d y-y d x$. It defines the foliation by planes containing the $z$-axis. The polar sets corresponding to planes transverse to the $z$-axis are empty and they coincide with the plane when it contains the $z$-axis.

## Example

Suppose that $\omega=\omega_{\text {low }}$ is homogeneous. Then the plane field $\mathcal{P}=\{\operatorname{ker} \omega\}$ is invariant by homotheties of center the origin.

Given a plane $h$, in general the set of points where $\operatorname{ker} \omega=h$ is a union of rays (lines containing the origin).

In this case, the polar curves are union of rays where the Gauss map $G_{\omega}:\left(\mathbb{C}^{3} \backslash O\right) \rightarrow$ $\mathbb{C P}^{2} ;[x, y, z] \mapsto[A, B, C]$ is locally surjective.

In dimension 3, let us suppose that the plane $h$ is transverse to the $x$-axis. Then $h$ is spanned by the vectors $(a, b, 0)$ and $(c, 0, d), b \neq 0, d \neq 0$ and the polar curve $\Gamma_{h}$ is the intersection of the surface $S_{a, b, 0}$ of equation $a A+b B=0$ and the surface $S_{c, 0, d}$ of equation $c A+d C=0$.

Definition 3.2.2. Given a subset $G \subset \mathbb{C P}^{2}$ of dimension 1 , a neighborhood $N_{G}$ of $G$, a family of branches of polar curves $\Gamma_{h}, h \notin N_{G}$ tangent to a line $\ell$ at the origin form a bouquet.

A typical case, pertinent when the 1 -form $\omega$ is the differential of a function $f$ is $G_{\ell}=$ $\{h$ such that $\ell \subset h\}, \ell$ a line of $\mathbb{C P}^{2}$.

Definition 3.2.3. When a bouquet of polar curves is tangent to an isolated ray of the tangent cone of $\omega$, say the $x$-axis, enlarging neighborhoods of a sequence of points of the $x$-axis converging to the origin may give rise to a limit plane field invariant by translations parallel to the $x$-axis that we will call profile. When $\omega$ is integrable we also call profile a typical leaf of the enlarged foliation, which is invariant by translations parallel to the $x$-axis.

More systematically, we will use Newton lenses (see Definition 2.2.1) to find profiles.
The choice of the dimension $n=3$ of the ambient space is intended to avoid heavy notations. The methods and most of the results extend straightforward to dimension $n$.

The induction leading to the construction of the profiles is similar to the Newton-Puiseux induction leading to a parametrization of an analytic curve of $\mathbb{C}^{2}$, but it relies on a different choice of variables and change of variables.

## 4 The Newton-lens algorithm, first step

We consider a complex polynomial differential form

$$
\omega=A(x, y, z) d x+B(x, y, z) d y+C(x, y, z) d z
$$

singular at the origin, that is such that $A(0,0,0)=B(0,0,0)=C(0,0,0)=0$. We will suppose also that the origin is an isolated singular point of $\omega$.

The Newton-lens algorithms provide Newton lenses, that is compute the suitable enlargement rates $\rho_{i}\left(x_{0}\right)$ along successive curves $\gamma_{i}$ tangent to a singular ray of the tangent cone, that we will suppose to be the $x$-axis.

As in the dimension 2 examples (see Figure 2), the effect of a moving lens with adapted stronger and stronger strength when approaching the origin will be to straighten the polar curves tangent to the $x$-axis keeping the directions of the planes $\operatorname{ker}(\omega)$.

### 4.1 The "quiet" regions

Let us first rule out the case where the $x$-axis is not in the tangent cone or is in the tangent cone but is not singular for $\omega$. In other terms the three coefficients of $\operatorname{low}(\omega)(x, 0,0)=A_{k}(x, 0,0) d x+$ $B_{k}(x, 0,0) d y+C_{k}(x, 0,0)$ are not simultaneously zero.

Proposition 4.1.1. When the $x$-axis is not in the tangent cone or is in the tangent cone but is not singular for $\omega$, the limit of the enlargement of a ball of size $\left|x_{0}\right|^{r}, r>0$ shows a plane field which is a family of parallel planes.

Proof. The differential form $\operatorname{low}(\omega)=A_{k}(x, y, z) d x+B_{k}(x, y, z) d y+C_{k}(x, y, z) d z$, with $A_{k}, B_{k}$ and $C_{k}$ homogeneous polynomials null or of degree $k$, verifies :

$$
A_{k}\left(x_{0}, 0,0\right) \neq 0 \quad \text { or } \quad B_{k}\left(x_{0}, 0,0\right) \neq 0 \quad \text { or } \quad C_{k}\left(x_{0}, 0,0\right) \neq 0 .
$$

For any exponent $r>1$

$$
\begin{aligned}
& \operatorname{low}(\omega)\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}, x_{0}^{r} Z_{1}\right) \\
& \quad=x_{0}^{k}\left(A_{k}(1,0,0) d x+B_{k}(1,0,0) d y+C_{k}(1,0,0) d z\right)+o\left(x_{0}^{k}\right)
\end{aligned}
$$

Notice that $d x=x_{0}^{r} d X_{1}, d y=x_{0}^{r} d Y_{1}, d z=x_{0}^{r} d Z_{1}$. Factorizing in $\operatorname{low}(\omega)\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}, x_{0}^{r} Z_{1}\right)$ the maximal possible power of $x_{0}$, we get the differential form $\omega_{1}$ which defines the profile (see Equation 1)

$$
\omega_{1}\left(X_{1}, Y_{1}, Z_{1}\right)=A_{k}(1,0,0) d X_{1}+B_{k}(1,0,0) d Y_{1}+C_{k}(1,0,0) d Z_{1} .
$$

The form $\omega_{1}$ is constant (and non-zero), therefore the plane field $\operatorname{ker} \omega_{1}$ is just a family of parallel planes.

In this case, after enlargement and when $x_{0} \rightarrow 0$, when $\omega$ is integrable, the leaves of the foliation look more and more like parallel planes.

A simple example is the singularity $\omega=d f ; f(x, y, z)=x y+y z+z x=0$. The leaves of the foliation are the levels $f(x, y, z)=\lambda$; they are surfaces, which, along any line through the origin $O=(0,0,0)$, give, near a point $\left(t_{0} x_{0}, t_{0} y_{0}, t_{0} z_{0}\right)$, at any scale $x_{0}^{r}, r>1$, a limit foliation by parallel planes. More generally, when $\omega$ is homogeneous and integrable, even if the $x$-axis is a singular ray of the tangent cone, near a point $\left(x_{0}, 0,0\right)$, after enlargement with ratio $1 /\left(x_{0}\right)^{r}, r>1$, the leaves of the foliation $\mathcal{F}$ look like parallel planes.

### 4.2 The first change of variables

The focus of the first lens is again the point $\left(x_{0}, 0,0\right)$ of the-axis, where $\left|x_{0}\right|$ is small. The first enlargement is of strength $\left(1 / \rho_{1}\right) \gg 1$. We will use values of $\rho_{1}$ of the form $\rho_{1}=x_{0}^{r_{1}}, r_{1}>$ $1, r_{1}$ rational. In our construction we can choose any determination of $x_{0}^{r}$. A choice of a complex enlargement ratio may introduce a rotation of the picture, but does not change the profile we want to observe. Now we suppose that the $x$-axis is a singular ray.

For any exponent $r>1$, the differential form $\omega$ writes, after enlargement using a lens centered at the point $\left(x_{0}, 0,0\right)\left(\left|x_{0}\right|\right.$ small $)$ of the $x$-axis (i.e. change of variables $\left.\Phi_{1}\right)$, as above :

$$
\left\{\begin{align*}
\Phi_{1}^{*}(\omega)\left(X_{1}, Y_{1}, Z_{1}\right) & =A\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}, x_{0}^{r} Z_{1}\right) d X_{1}+B\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}, x_{0}^{r} Z_{1}\right) d Y_{1}  \tag{1}\\
& +C\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}, x_{0}^{r} Z_{1}\right) d Z_{1} \\
& =x_{0}^{\theta} \widetilde{\omega_{1}}=x_{0}^{\theta} \omega_{1}\left(X_{1}, Y_{1}, Z_{1}\right)+o\left(x_{0}^{\theta}\right), \quad \text { with } \\
\omega_{1}\left(X_{1}, Y_{1}, Z_{1}\right) & =A_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d X_{1}+B_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d Y_{1}+C_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d Z_{1}
\end{align*}\right.
$$

where $x_{0}^{\theta}$ is the highest power of $x_{0}$ which is a factor of $\Phi_{1}^{*}(\omega)\left(X_{1}, Y_{1}, Z_{1}\right) ; x_{0}^{\theta} \omega_{1}\left(X_{1}, Y_{1}, Z_{1}\right)$ is the sum of terms of lowest degree in $x_{0}$ of $\Phi_{1}^{*}(\omega)$ and $\theta$ the valuation (in $\left.x_{0}\right)$ of $\Phi_{1}^{*}(\omega)$. This defines the differential form

$$
\omega_{1}\left(X_{1}, Y_{1}, Z_{1}\right)=A_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d X_{1}+B_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d Y_{1}+C_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d Z_{1}
$$

If $r$ is not strictly superior to 1 , when $x_{0}$ tends to 0 , the origin $(0,0,0)$ will stay at a finite distance from the new origin, the eventual renormalized plane field contains the origin and will stay singular.

When $\left[A_{1}\left(X_{1}, Y_{1}, Z_{1}\right), B_{1}\left(X_{1}, Y_{1}, Z_{1}\right), C_{1}\left(X_{1}, Y_{1}, Z_{1}\right)\right]$ define a point of $\mathbb{C P}^{2}$ independent of ( $X_{1}, Y_{1}, Z_{1}$ ), the planes of the plane-field $\operatorname{ker} \omega$ defined by
$\omega_{1}=A_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d X_{1}+B_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d Y_{1}+C_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d Z_{1}$ are parallel planes.
From now on, we suppose that the $x$-axis is an isolated singular ray of the tangent cone.
We shall give necessary conditions on $r_{1}$ to obtain a profile where the planes $\operatorname{ker} \omega_{1}$ are not all vertical.

The exponent $r_{1}$ of $\rho_{1}=x_{0}^{r_{1}}$ is determined using a vertex or a side of the first Newton-lens polygon of $\omega$.

Recall that the Newton-lens cloud of the polynomial

$$
\mathcal{A}\left(x_{0}, \rho_{1}\right)=A\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}, \rho Z_{1}\right)
$$

is the set of pairs $(i, j)$ of exponents of a non zero monomial of $\mathcal{A}$ (in the variables $\left(x_{0}, \rho_{1}\right)$ ). We obtain the same way Newton-lens clouds for

$$
\mathcal{B}\left(x_{0}, \rho_{1}\right)=B\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}, \rho_{1} Z_{1}\right) \text { and } C\left(x_{0}, \rho_{1}\right)=C\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}, \rho_{1} Z_{1}\right)
$$

Definition 4.2.1. The first Newton-lens cloud of $\omega$ is the union of the Newton-lens clouds of the three polynomials $\mathcal{A}, \mathcal{B}$ and $\mathcal{C}$. The first Newton-lens polygon of $\omega$ is the lower convex hull of the union of the upper quadrants of vertices the points of the first Newton-lens cloud of $\omega$.

An example in dimension 2, $f(x, y)=a y^{3}+b x^{2} y^{2}-c x^{4} y+x^{5}, \omega=d f=\left(2 b x y^{2}-4 c x^{3} y+\right.$ $\left.5 x^{4}\right) d x+\left(3 a y^{2}+2 b x^{2} y-c x^{4}\right) d y$ We get, setting $\rho_{1} \widetilde{\omega_{1}}\left(x_{0}, \rho_{1}\right)=\omega\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}\right)$ (we forget


Figure 5: A cubic profile.
about the $\rho_{1}$-factor coming from the differentials).
$\widetilde{\omega_{1}}\left(x_{0}, \rho_{1}\right)=\left[2 b\left(x_{0}+\rho_{1} X_{1}\right) \rho_{1}^{2} Y_{1}^{2}-4 c\left(x_{0}+\rho_{1} X_{1}\right)^{3} \rho_{1} Y_{1}+5\left(x_{0}+\rho_{1} X_{1}\right)^{4}\right] \rho_{1} d X_{1}+\left[3 a \rho_{1}^{2} Y_{1}^{2}+\right.$ $\left.2 b\left(x_{0}+\rho_{1} X_{1}\right)^{2} \rho_{1} Y_{1}-c\left(x_{0}+\rho_{1} X_{1}\right)^{4}\right] \rho_{1} d Y_{1}$.


Figure 6: Newton-lens polygon for a cubic profile.

Taking, as the Newton-lens polygon suggests, $\rho_{1}=x_{0}^{2}$, the lowest power of $x_{0}$ is $x_{0}^{4}$. Factorizing $x_{0}^{4}$ provides the differential equation $10 d X_{1}=3 a Y_{1}^{2}+2 b Y_{1}$ and therefore a cubic profile $X_{1}=\frac{a}{10} Y_{1}^{3}+\frac{b}{10} Y_{1}^{2}$.

## Examples in dimension 3

- Example $1 f(x, y, z)=x^{3}-y^{2}-z^{2}$ We get $d f=3 x^{2} d x-2 y d y-2 z d z$ and $\operatorname{low}(d f)=2 y d y-2 z d z$. The tangent cone is the cylinder of equation $y^{2}+z^{2}=0$; the $x$-axis is degenerate. It is the only singular line of the tangent cone.

Its Newton-lens polygon has, as the Newton-lens polygon of $f(x, y)=x^{3}-y^{2}$, just one side of slope $-1 / 2$, therefore take $\rho_{0}=x_{0}^{2}$.

Therefore we chose $M \cdot x_{0}^{2}, M$ as large as we want, as radius of the ball $B=B\left(x_{0}, R=M \cdot x_{0}^{2}\right)$ where we will look at the enlarged plane-field.

Performing the change of variable

$$
\begin{aligned}
& \left(x-x_{0}\right)=x_{0}^{2} \cdot X_{1} \\
& y=x_{0}^{2} \cdot Y_{1}
\end{aligned}
$$

we see that the level $C_{x_{0}^{2}}=\left\{f=x_{0}^{2}\right\}$ have in the ball $B$ a shape similar to the solution of the differential equation

$$
\begin{aligned}
& a \simeq \partial X_{1} / \partial Y_{1} \simeq(2 / 3) Y_{1} \\
& b \simeq \partial X_{1} / \partial Z_{1} \simeq(2 / 3) Z_{1}
\end{aligned}
$$

which admits the solution $X_{1}=-(1 / 3)\left(Y_{1}^{2}+Z_{1}^{2}\right)$

- Example 2, $f(x, y, z)=x^{4}+y^{4}-x y^{2}+z^{3}$ (see Figure ??).

$$
\begin{aligned}
& \omega=d f=\left(4 x^{3}-y^{2}\right) d x+\left(4 y^{3}-2 x y\right) d y+3 z^{2} d z \\
& \omega_{\text {low }}=-y^{2} d x-2 x y d y+3 z^{2} d z
\end{aligned}
$$

The tangent cone at the origin has the equation $z^{3}-x y^{2}=0$,
$\omega_{\text {low }}=-y^{2} d x-2 x y d y+3 z^{2} d z$. The unique singular ray has equations $z=y=0$, so it is the $x$-axis.

After the change of variables

$$
x=x_{0}+\rho_{1} X_{1}, y=\rho_{1} Y_{1}, z=\rho_{1} Z_{1}
$$

we get, setting $\rho_{1} \widetilde{\omega_{1}}\left(x_{0}, \rho_{1}\right)=\omega\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}, \rho_{1} Z_{1}\right)$ (we forget about the $\rho_{1}$-factor coming from the differentials).

$$
\begin{aligned}
\widetilde{\omega_{1}} & =\left[4\left(x_{0}+\rho_{1} X_{1}\right)^{3}-\rho_{1}^{2} Y_{1}^{2}\right] d X_{1}+\left[4 \rho_{1}^{3} Y_{1}^{3}-2 \rho_{1}\left(x_{0}+\rho_{1} X_{1}\right) Y_{1}\right] d Y_{1} \\
& +3 \rho_{1}^{2} Z_{1}^{2} d Z_{1}
\end{aligned}
$$

In our construction, we will choose $\rho_{1}=x_{0}^{r_{1}}, r_{1}>1, r_{1}$ rational (we can choose any determination of $x_{0}^{r_{1}}$ ) using the slope of a side of the Newton-lens polygon.


Figure 7: Newton-lens cloud and polygon of $\widetilde{\omega_{1}}=\left[4\left(x_{0}+\rho_{1} X_{1}\right)^{3}-\rho_{1}^{2} Y_{1}^{2}\right] d X_{1}+\left[4 \rho_{1}^{3} Y_{1}^{3}-2 \rho_{1}\left(x_{0}+\right.\right.$ $\left.\left.\rho_{1} X_{1}\right) Y_{1}\right] d Y_{1}+3 \rho_{1}^{2} Z_{1}^{2} d Z_{1}$

Here (see Figure 7) we take $\rho_{1}=x_{0}^{2}$, then $\omega_{1}=4 d X_{1}-2 Y_{1} d Y_{1}$, so

$$
\begin{aligned}
& \frac{d X_{1}}{d Y_{1}}=\frac{1}{2} Y_{1} \text { and } \\
& X_{1}=\frac{Y_{1}^{2}}{4}+c
\end{aligned}
$$

If we want an estimation in terms of the level $\lambda$ of $f=\lambda$, we see that the intersection point $x_{0}$ of $O x$ and the level $f=\lambda$ satisfies $\left|x_{0}\right|=|\lambda|^{1 / 4}$, therefore $\left|\rho_{1}\right|=|\lambda|^{1 / 2}$.

Notice that we obtained interesting profiles when the bottom-left point of the Newton-lens polygon, end of a side we considered, was coming from the " $x_{0}^{k} d X_{1}$ " term of $\omega_{1}$ itself "coming" from the a monomial of $f$ which is a power of $x$ only. This computation takes care of half of what we can observe on Figure ??; the conical shape should be attributed to the homogeneous part.

### 4.3 Why the variables $X_{1}$ does not appear after the first step

We now proceed to understand the effect of zooming near a point $\left(x_{0}, 0,0\right)$ close to the origin in the $x$-axis. Let us first consider slopes $-1<(-1 / r)<0$ which are not slope of a side of the first Newton-lens polygon.

We perform the change of variables

$$
\left\{\begin{array}{l}
x=x_{0}+x_{0}^{r} X_{1}  \tag{2}\\
y=x_{0}^{r} Y_{1} \\
z=x_{0}^{r} Z_{1}
\end{array}\right.
$$

Then

$$
\left\{\begin{align*}
& \Phi_{1}^{*}(\omega)\left(X_{1}, Y_{1}, Z_{1}\right)=A\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}, x_{0}^{r} Z_{1}\right) d X_{1}  \tag{3}\\
&+B\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}, x_{0}^{r} Z_{1}\right) d Y_{1} \\
&+C\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}, x_{0}^{r} Z_{1}\right) d Z_{1} \\
&=x_{0}^{\theta} \widetilde{\omega_{1}}\left(X_{1}, Y_{1}, Z_{1}\right) \quad \text { with } \\
& \tilde{\omega}_{1}\left(X_{1}, Y_{1}, Z_{1}\right) \widetilde{A_{1}}\left(X_{1}, Y_{1}, Z_{1}\right)+\widetilde{B_{1}}\left(X_{1}, Y_{1}, Z_{1}\right)+\widetilde{C_{1}}\left(X_{1}, Y_{1}, Z_{1}\right) \\
&=x_{0}^{\theta} \omega_{1}\left(X_{1}, Y_{1}, Z_{1}\right)+o\left(x_{0}^{\theta}\right), \quad \text { with } \\
& \omega_{1}\left(X_{1}, Y_{1}, Z_{1}\right)=A_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d X_{1}+B_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d Y_{1} \\
&+C_{1}\left(X_{1}, Y_{1}, Z_{1}\right)
\end{align*}\right.
$$

Proposition 4.3.1. The polynomials $A_{1}\left(X_{1}, Y_{1}, Z_{1}\right), B_{1}\left(X_{1}, Y_{1}, Z_{1}\right)$ and $C_{1}\left(X_{1}, Y_{1}, Z_{1}\right)$ (see relations (3) above) do not contain the variable $X_{1}$.



Figure 8: A tail coming from a point of the Newton-lens polygon of $\omega_{1}$, left: from a point on a side, right: from a vertex

Proof. We prove separately that $A_{1}\left(X_{1}, Y_{1}, Z_{1}\right), B_{1}\left(X_{1}, Y_{1}, Z_{1}\right), C_{1}\left(X_{1}, Y_{1}, Z_{1}\right)$ do not contain $X_{1}$.
Writing $A(x, y, z)=\sum a_{p, q, \ell} x^{p} y^{q} z^{\ell}$, a non zero monomial $a_{p, q, \ell} x^{p} y^{q} z^{\ell}$ gives a polynomial $a_{i, j, \ell}\left(x_{0}+\rho_{1} X_{1}\right)^{p} \rho_{1}^{q} Y_{1}^{q} \rho_{1}^{\ell} Z_{1}^{\ell}$ in $A\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}, \rho Z_{1}\right)$.

One point comes from the monomial $x_{0}^{p} \rho_{1}^{q} Y_{1}^{q} \rho_{1}^{\ell} Z_{1}^{\ell}$.
The other points introduced by the terms of $a_{p, q, \ell}\left(x_{0}+\rho_{1} X_{1}\right)^{i} \rho_{1}^{j} Y_{1}^{j} \rho_{1}^{k} Z_{1}^{k}$ containing $X_{1}$ form a tail of slope -1 on the left of ( $p, q, \ell$ ) (in black in Figure 8).

Then a line in $\mathbb{R}^{2}$ of slope $-1 / r, r>1$ below the Newton-lens cloud of $\mathcal{A}\left(x_{0}, \rho_{1}\right)$ cannot contain a point in this tail.

We proceed in the same way with $\mathcal{B}\left(x_{0}, \rho_{1}\right)=B\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}, \rho_{1} Z_{1}\right)$ and $\mathcal{C}\left(x_{0}, \rho_{1}\right)=$ $C\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}, \rho_{1} Z_{1}\right)$.

A support line of the Newton-lens polygon, contains either a single pair $(i, j)$, or an edge of the Newton-lens polygon of $\omega$. All the corresponding monomials $a_{i, j, k} x_{0}^{i} \rho_{1}^{j} Y_{1}^{j} \rho_{1}^{k} Z_{1}^{k}, b_{i, j, k} x_{0}^{i} \rho_{1}^{j} Y_{1}^{j} \rho_{1}^{k} Z_{1}^{k}$ or $c_{i, j, k} x_{0}^{i} \rho_{1}^{j} Y_{1}^{j} \rho_{1}^{k} Z_{1}^{k}$ give terms which do not contain $X_{1}$. They are followed by a tail of slope -1 the points of which cannot belong to any side of slope $-1<r<0$ of the Newton polygon.


Figure 9: Support lines of slopes between $-1 / a_{1}$ and $-1 / b_{1}$ used to build a basic safety annulus

### 4.4 Vertices of the first Newton-lens polygon

First let $\rho=x_{0}^{r_{1}}, r_{1}>1$ an enlargement rate such that $-1 / r_{1}$ is not a slope of a side of the first Newton-lens polygon. Performing the change of variables (Formula 1) we get:

$$
\left\{\begin{align*}
\Phi_{1}^{*}(\omega)\left(X_{1}, Y_{1}, Z_{1}\right) & =A\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}, x_{0}^{r} Z_{1}\right) d X_{1}  \tag{4}\\
& +B\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}, x_{0}^{r} Z_{1}\right) d Y_{1} \\
& +C\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}, x_{0}^{r} Z_{1}\right) d Z_{1} \\
& =x_{0}^{\theta} \omega_{1}\left(X_{1}, Y_{1}, Z_{1}\right)+o\left(x_{0}^{\theta}\right), \quad \text { with } \\
\omega_{1}\left(X_{1}, Y_{1}, Z_{1}\right) & =A_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d X_{1}+B_{1}\left(X_{1}, Y_{1}, Z_{1}\right) d Y_{1} \\
& +C_{1}\left(X_{1}, Y_{1}, Z_{1}\right)
\end{align*}\right.
$$

where $A_{1}, B_{1}$ and $C_{1}$ are homogeneous polynomials in the variables $Y_{1}$ and $Z_{1}$ of the same degree which depend only on the vertex of the Newton-lens polygon "first" touched by an affine line of slope $-1 / r$. Indeed, the degrees in $\rho$ of the monomials of $A\left(x_{0}, \rho Y_{1}, \rho Z_{1}\right), B\left(x_{0}+\rho X_{1}, \rho Y_{1}\right)$, $C\left(x_{0}+\rho X_{1}, \rho Y_{1}, \rho Z_{1}\right)$ are the same as the degrees of the homogeneous polynomials in $Y_{1}, Z_{1}$.

Then the plane field $\operatorname{ker} \omega_{1}$ is constant on planes projecting on lines containing the origin of the $\left(Y_{1}, Z_{1}\right)$-plane. It defines a Gauss map $G_{\omega_{1}}\left(X_{1}, Y_{1}, Z_{1}\right) \mapsto \operatorname{ker} \omega_{1}\left(X_{1}, Y_{1}, Z_{1}\right)$ In this case image of $G_{\omega_{1}}$ is of dimension 1 in $\mathbb{C P}^{2}$.

Definition 4.4.1. - Given $x_{0} \neq 0 \in \mathbb{C}$, a positive constants $1<a$, the zone $Z_{a}\left(x_{0}\right) \subset\left\{x=x_{0}\right\} \subset$ $\mathbb{C}^{3}$ is the set of points $|(y, z)|<x_{0}^{a}$. The zone $Z_{O x, a}$ is the union $\bigcup_{x_{0} \in \mathbb{C} \backslash 0} Z_{a}\left(x_{0}\right)$

- Given $x_{0} \neq 0 \in \mathbb{C}$, positive constants $1<a<b$, the zone $Z_{a, b}\left(x_{0}\right) \subset\left\{x=x_{0}\right\} \subset \mathbb{C}^{3}$ is the set of points $x_{0}^{b}<|(y, z)|<x_{0}^{a}$. The zone $Z_{O x, a, b}$ is the union $\cup_{x_{0} \in \mathbb{C} \backslash 0} Z_{a, b}\left(x_{0}\right)$

The zone $Z_{O x, a}$ looks like a thinned cone.
The zone $Z_{O x, a, b}$ looks like a thinned cone on an annulus.
Both admit a tangent line at the origin: the $x$-axis.
Observe that any vertex, but the upper left one, of the Newton-lens polygon is the boundary of two sides of slopes $-\left(1 /\left(r_{\text {low }}\right)\right)>-\left(1 /\left(r_{\text {sup }}\right)\right)$. Lemma 4.4.3 describes the limit, when $\left|x_{0}\right| \rightarrow 0$ of the plane-field $\operatorname{ker} \omega$ in the region $Z_{O x, r_{\text {sup }}-\delta, r_{\text {low }}+\delta}, \delta$ is supposed to be very small compared to $r_{\text {sup }}$ and $r_{l o w}$. We will call the region $Z_{O x, r_{\text {sup }}-\delta, r_{l o w}+\delta}, \delta$, a safety funnel; this funnel has a very thick side and a very thin hole.

Let us concentrate on the upper vertex of the union of sides of slope strictly between -1 and 0 (see Figure 9). The support lines touching the Newton-lens polygon at this vertex define a
zone $Z_{0 x, a_{1}, b_{1}}$. The form $\omega_{1}$ has coefficients $A, B, C$ which are homogeneous polynomials in $X_{1}$ and $Z_{1}$ and define a Gauss map of image $G$ of dimension 1. Therefore the image of the zone $Z_{0 x, a_{1}, b_{1}}$ by the Gauss map of the form $\omega$ is contained in a neighborhood of $G$. Getting closer to the origin makes the neighborhood of $G$ as thin as desired.

Remark 4.4.2. Given a neighborhood $V_{0}$ of $G \subset \mathbb{C P}^{2}$, any polar curve $\Gamma_{h}, h \notin V_{0}$ having a point in $Z_{O x, b_{1}}$ is trapped in $Z_{O x, b_{1}}$ when approaching the origin. In particular it is tangent to $O x$ at the origin. Therefore it will look closer and closer, in any $C^{k}$-topology, to a line parallel to $O x$ when $x_{0}$ goes to zero.

As polar curves are analytic curves, the limits of the enlargement of such a polar curve trough a Newton-lens will therefore always be a line parallel to the $x$-axis.

The definition of polar curves implies already that the plane field $\operatorname{ker} \omega_{1}$, that we will construct using the Newton-lens induction if it ends at the first step do not depend on the variable $X_{1}$. More precisely

Lemma 4.4.3. - Through a lens of magnifying ratio $x_{0}^{-r},\left(r_{\text {sup }}>r>r_{\text {low }}\right)$ centered on the $x$-axis at a point close enough from the origin, where $-1 / r_{\text {low }}$ and $-1 / r_{\text {sup }}$ are the slopes of the adjacent sides to a vertex on the first Newton-lens polygon ( $r_{\text {sup }}=1$, if one adjacent side is of slope $(-1)$ or if one extremity belongs to the $\rho_{1}$-axis (we do not consider the case $r_{\text {low }}=\infty$ corresponding to a point on the $x_{0}$-axis) we observe a region around the $x$-axis where, except along the $x$-axis, the plane-field looks, through the lens, invariant by translations parallel to the $x$-axis. The image of the Gauss map of the limit plane-field is of dimension 1.

Proof: A vertex of the first Newton-lens polygon corresponds to terms of the form $x_{0}^{p} \rho_{1}^{q}$. The coefficients of the terms of the form $x_{0}^{p} \rho_{1}^{q}$ are homogeneous polynomials a priori in the new variables $X_{1}, Y_{1}$ and $Z_{1}$. The coefficients of the 1-form $\omega_{1}$ depend in fact only in the variables $Y_{1}$ and $Z_{1}$ (see Proposition 4.3). In other terms the variable $X_{1}$ does not appear (see Figure 8, as the picture is the same for the Newton-lens polygon of a polynomial $f$ or of the three coefficients $(A, B, C)$ of a 1 -form $\omega$ together. This proves again that the limit plane-field is invariant by translation in the $X_{1}$ direction.

### 4.5 First step, $\omega=d f$

Lemma 4.5.1. When $\omega=d f$, choosing an enlarging rate given by a support line of the first Newton-lens polygon containing only one vertex, the foliation looks through the lens more and more, when $x_{0}$ tends to zero, like the product of homogeneous foliation of the $\left(Y_{1}, Z_{1}\right)$-plane and the $x$-axis.

Remark 4.5.2. The separatrices (see definition 7.5.1) of the homogeneous foliation defined in the $Y_{1}, Z_{1}$-plane are lines containing the origin. Their products with the $x$-axis are pieces of the limit of the enlargement of $f=0$ contained in the domain enlarged by the Newton lens.

Proof: Proof of the lemma When $\omega=d f$, a vertex of the first Newton-lens polygon not situated on the $x_{0}$-axis cannot correspond to a term $a_{p, q, \ell} x^{p-1} y^{q} z^{\ell}$ of $\partial f / \partial x$ as $\partial x^{p} y^{q} z^{\ell} / \partial y$ and/or $\partial x^{p} y^{q} z^{\ell} / \partial z$ would create a term below the support line containing the vertex. Then, the plane field defined by the form $\omega_{1}$ is the product of a line field defined in the $\left(Y_{1}, Z_{1}\right)$ plane by the $x$-axis (a priori, the planes of $\operatorname{ker} \omega_{1}$ need not to be parallel to the $x$-axis when $\omega \neq d f$.).

Corollary 4.5.3. When $\omega=d f$ there is, associated to each vertex of the Newton-lens polynomial, for any $\delta<\left(r_{\text {sup }}-r_{\text {low }}\right) / 2$, a "security funnel" centered on the $x$-axis, crossing planes transverse to the $x$-axis in annuli of inner radius $x_{0}^{r_{\text {low }}+\delta}$ and outer radius $x_{0}^{r_{\text {sup }}-\delta}$ where the foliation defined by $\omega$ looks through a lens of strength $x_{0}^{r_{\text {sup }}-\delta}$ like a product of a transverse line-field and the $x$-axis.

Remark 4.5.4. When $\rho_{1}=x_{0}^{r}, r>1$ is an enlargement rate such that $-1 / r$ is not a slope of a side of the first Newton-lens polygon, the limit Gauss map seen through a Newton lens is constant on the planes projecting on rays of the $\left(Y_{1}, Z_{1}\right)$-plane, and therefore is of rank less or equal to 1 . Therefore a necessary condition in order to get a map

$$
\left(X_{1}, Y_{1}, Z_{1}\right) \rightarrow\left[A_{1}\left(X_{1}, Y_{1}, Z_{1}\right), B_{1}\left(X_{1}, Y_{1}, Z_{1}\right), C_{1}\left(X_{1}, Y_{1}, Z_{1}\right)\right] \in \mathbb{C P}^{2}
$$

of rank 2 is that $\rho_{1}=x_{0}^{r_{1}}$ where $-1 / r_{1}$ is the slope of an edge of the first Newton-lens polygon of $\omega$.

### 4.6 First step, $\omega=d f$, polynomial profile

From now on, we suppose that $r_{1}=p / q$, where $s=-q / p$ is a slope of an edge $\Delta$ of the Newton-lens polygon. It gives, after enlargement, a differential form (see Equation 4)

$$
\omega_{1}=A_{1}\left(Y_{1}, Z_{1}\right) d X_{1}+B_{1}\left(Y_{1}, Z_{1}\right) d Y_{1}+C_{1}\left(Y_{1}, Z_{1}\right) d Z_{1} .
$$

Notice that at least one of the coefficients $A(x, y, z), B(x, y, z), C(x, y, z)$ of $d x, d y, d z$ in $\omega$ should contain a monomial of the form $x^{n}$, otherwise the whole $x$-axis would be a set of singular points.

The coefficients $A_{1}, B_{1}, C_{1}$ of $d X_{1}, d Y_{1}, d Z_{1}$ in $\omega_{1}$ need not a priori to contain a constant term. Let us suppose that a monomial $a_{n, 0,0} x^{n}$ of $A(x, y, z)$ generates a monomial $a_{n, 0,0} x_{0}^{n}$ of $A\left(x_{0}+\rho X_{1}, \rho_{1} Y_{1}, \rho_{1} Z_{1}\right)$ such that the Newton-lens polygon of $\omega$ contains the point $(n, 0,0)$. This point is the extremity of an edge of the Newton-lens polygon of slope $-\left(1 / r_{1}\right), r_{1}>1$. Then $A_{1}$ contains a constant term. Therefore we get

$$
d X_{1}=\frac{B_{1}}{A_{1}} d Y_{1}+\frac{C_{1}}{A_{1}} d Z_{1}
$$

A solution of the system of partial differential equations

$$
\begin{equation*}
\frac{\partial X_{1}}{\partial Y_{1}}=-\frac{B_{1}\left(Y_{1}, Z_{1}\right)}{A_{1}\left(Y_{1}, Z_{1}\right)} \frac{\partial X_{1}}{\partial Z_{1}}=-\frac{C_{1}\left(Y_{1}, Z_{1}\right)}{A_{1}\left(Y_{1}, Z_{1}\right)} \tag{5}
\end{equation*}
$$

is a profile (see Definition 3.2.3).
Otherwise the induction should be continued.
When $\omega$ if integrable, if $A_{1}\left(Y_{1}, Z_{1}\right)$ is non zero, then the partial differential equations 5 provide a foliation, as its solution is a limit of leaves of foliations. In particular this is the case when $\omega=d f$.

Theorem 4.6.1. When $\omega=d f, f(x, y, z)$ a polynomial with an isolated singularity at the origin such that the $x$-axis is a singular ray of the tangent cone of $\omega=d f$, when the first step of the induction provides a profile, lenses of strength $x_{0}^{-\left(r_{1}\right)}$ centered at $\left(x_{0}, 0,0\right)$ provide, when $x_{0} \rightarrow 0$, a polynomial profile, that is the leaves of $\mathcal{F}_{\text {df }}$ look as parallel graphs of the form $X_{1}=P_{1}\left(Y_{1}, Z_{1}\right)+c$.

We will use the notation $d_{y, z} f=\frac{\partial f}{\partial y} d y+\frac{\partial f}{\partial z} d z$.
On Figure 10, left, the red bullets would correspond to a term of $f$ providing a term of $d_{y, z} f$ on a side or slope $-(1 / r), r>1$ of the Newton-lens polygon, and a term of $d_{y z} f$ under the side of the Newton-lens polygon.

On Figure 10, right, the red bullets correspond to a term of $f$ providing a term of $d_{y, z} f$ on the $x_{0}$-axis of the Newton-lens polygon. The following lemma repeat with words the idea conveyed


Figure 10: Impossible and possible contributions when $\omega=d f$
by Figure 10 and proves Theorem 4.6.1.
Lemma 4.6.2. If $\omega=d f$, and if the induction ends at the first step, then one of the coefficients $B_{1}\left(Y_{1}, Z_{1}\right)$ of $d Y_{1}$ or $C_{1}\left(Y_{1}, Z_{1}\right)$ of $d Z_{1}$ in $\omega_{1}$ is not zero and $A_{1}$ is reduced to a $x_{0}^{n}$ term.

Proof. A monomial $a x^{p} y^{q} z^{\ell}$ of $f$ provides a monomial pax $x^{p-1} y^{q} z^{\ell}$ of $\frac{\partial f}{\partial x}$ and monomials $q a x^{p} y^{q-1} z^{\ell}$ and $\ell a x^{p} y^{q} z^{\ell-1}$ of $d_{y z} f$, except if $q=\ell=0$. Then all the points of the edge $\Delta$ of the Newton-lens polygon but for the right extremity if it is on the $O x$-axis come from a monomial of $\frac{\partial}{\partial y}$ or $\frac{\partial}{\partial z}$. Indeed, all the points corresponding to monomials of $\frac{\partial f}{\partial x}$ are strictly above the edge $\Delta$ of slope in ] - 1,0 ], except maybe the last point on the right if it is on the $x_{0}$ axis.

Let us now prove a result analogous to the 2-dimensional one proved in Rouille's thesis [Rou2].

Proposition 4.6.3. The graph of any polynomial in the variables $\left(Y_{1}, Z_{1}\right)$ can be a profile obtained at the first step of Newton-lens induction.

Proof: It suffices to consider the polynomial

$$
\begin{array}{r}
f(x, y, z)=a_{n}(y, z)+a_{n-1}(y, z) x^{2}+\cdots+a_{n-k}(y, z) x^{2 k} \\
+\cdots a_{1}(y, z) x^{2(n-1)}+x^{2 n-1} .
\end{array}
$$



Figure 11: The contributions of a term of $f$.
where the $a_{i}$ are homogeneous polynomials of degree $i$ in the variables $(y, z)$. We again use the notation $d_{y, z}$ which follows Theorem 4.6.1.

Then

$$
\begin{aligned}
d f & =d_{y, z} a_{n}(y, z)+d_{y, z} a_{n-1}(y, z) \cdot x^{2}+\cdots+d_{y, z} a_{n-k}(y, z) \cdot x^{2 k} \\
& +\cdots+d_{y, z} a_{1}(y, z) x^{2(n-1)} \\
& +2 x a_{n-1}(x, y) d x+\cdots+2 k x^{2 k-1} a_{n-k}(y, z) d x \\
& +\cdots+(2 n-1) x^{2 n-2} .
\end{aligned}
$$

The Newton-lens polygon is therefore analogous to Rouillé's example (see Figure 11).

### 4.7 Merle-type results (1)

Proposition 4.7.1. In $\mathbb{C}^{3}$, when $\omega=d f$ and when only one step is needed to obtain a profile, a bouquet of generic polar curves containing the profile stays apart from the singular level $f=0$.

Remark 4.7.2. Of course, Proposition 4.7 .1 is valid in $\mathbb{C}^{2}$. In that case, it writes in terms of contact order: the contact order of two curves of a bouquet is larger than the contact order of a curve of the bouquet and the curve $f=0$.

Proof: For that, let us compare the Newton-lens polygons of $f\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}, \rho_{1} Z_{1}\right)$ and of $d f\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}, \rho_{1} Z_{1}\right)$.

The slope of the bottom segment of the Newton-lens polygon of $f(3 / 2$ in Figure 12) is (in absolute value) bigger than ( $1 / 2$ in Figure 12) the slope of the bottom side of the Newton-lens polygon of $d f$.

This means that a renormalization ball centered at $\left(x_{0}, 0,0\right)$ catching the profile (radius $\simeq x_{0}^{2}$ ) is much smaller that a renormalization ball centered at $\left(x_{0}, 0,0\right)$ catching the singular level $f=0$ (radius $\simeq x_{0}^{3 / 2}$ ). In other terms, Merle's bouquet of polar curves corresponding to the tangent


The two sides have different slopes

Figure 12: Comparison of Newton-lens polygons of $f$ (purple) and $d f$ (yellow); the tails are not represented.


Figure 13: Comparison of $\{f=0\}$ and a bouquet of polar curves.
directions at points of balls whose limits will provide the profile is "far" from the singular level $f=0$.

When $\omega=d f$, the slope of the line joining terms coming from a term $x^{p} y^{q} z^{\ell}$ of $\partial f / \partial x$ as $\partial x^{p} y^{q} z^{\ell} / \partial y$ and/or $\partial x^{p} y^{q} z^{\ell} / \partial z$ would create a term below the support line containing the vertex. In other words, the bottom side of the Newton-lens polygon of $d f$ is the only one which is not parallel to a side of the Newton-lens polygon of $f$. The only way to get a side of $\mathcal{N}(d f)$ non parallel to a side of $\mathcal{N}(f)$ is when the vertex on the $x$-axis of $\mathcal{N}(d f)$ comes from $\frac{\partial}{\partial x}(f)$. When $\omega=d f$ the Gauss image $G$ of $Z_{O x, r_{a}, r_{b}}$ is a neighborhood of the set of $h$ containing the $x$-axis which is a projective line in $\mathbb{C P}^{2}$. In that case, the bouquets of strands of polar curves $\Gamma_{h}, h$ not too close from the line $D_{x}$ of $\mathbb{C P}^{2}$ of planes containing $O x$, look more and more like parallel lines, the leaves of $\mathcal{F}_{\omega_{\ell}}$ are graphs of polynomials.

## 5 Further steps

### 5.1 Example with more than one step

We give now one example, in $\mathbb{C}^{2}$, of a profile obtained with two steps of induction.
Example 5.1.1. This example is certainly the shortest possible (three terms) :

$$
f(x, y)=y^{2}+x^{2} y+x^{4} .
$$

Its first Newton-lens polygon (figure (??) on the left) has only one side. The differential form $\omega_{1}=A_{1}\left(Y_{1}\right) d X_{1}+B_{1}\left(Y_{1}\right) d Y_{1}$ is $\omega_{1}=\left(2 Y_{1}+1\right) d Y_{1}$.

A second step is needed with the root $c=-1 / 2$. The second Newton-lens polygon is shown
in Figure (??) on the right. It gives the differential equation :

$$
\frac{d X_{2}}{d Y_{2}}=-\frac{2 Y_{2}}{3}
$$

The integral curves of the profile are parabolas.

### 5.2 The general construction

We will detail the second step of the induction. We need to prove that the terms $A_{2}, B_{2}, C_{2}$ of $\omega_{2}=A_{2} d X_{2}+B_{2} d Y_{2}+C_{2} d Z_{2}$ do not depend on the variable $X_{2}$. The further steps, when needed, are analogous.

Let us now suppose that $A_{1}\left(Y_{1}, Z_{1}\right)=0$. Let $\rho_{1}=x_{0}^{r_{1}}$.
After dividing $\omega\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}, \rho_{1} Z_{1}\right)$ by the maximal power of $x_{0}$, we get a one-form in the variables $X_{1}, Y_{1}, Z_{1}$
$\tilde{\omega_{1}}=\tilde{A_{1}} d X_{1}+\tilde{B_{1}} d Y_{1}+\tilde{C_{1}} d Z_{1}$
and, considering only the term $\omega_{1}$ of $\tilde{\omega}_{1}$ which does not contain a factor $x_{0}^{s}$, as $A_{1}\left(Y_{1}, Z_{1}\right)=0$,
$\omega_{1}=B_{1} d Y_{1}+C_{1} d Z_{1}$
In the ( $Y_{1}, Z_{1}$ )-plane, $\omega_{1}$ defines a one-dimensional foliation with singular points the roots $\left(y_{i}^{1}, z_{i}^{1}\right)$ of $\left(B_{1}=C_{1}=0\right)$

After choosing a root $\left(y_{i}^{1}, z_{i}^{1}\right)$ of $B_{1}+C_{1}=0$, let us perform the second change of variables

$$
\begin{aligned}
& X_{1}=\rho_{2} X_{2} \\
& Y_{1}=y_{i}^{1}+\rho_{2} Y_{2} \\
& Z_{1}=z_{i}^{1}+\rho_{2} Z_{2}
\end{aligned}
$$

We get a Newton-lens cloud and a Newton-lens polygon, now with vertices in $\mathbb{Q}^{+} \times \mathbb{N}$, plotting the coefficient of form $\tilde{\omega_{1}}\left(\rho_{2} X_{2}, y_{i}^{1}+\rho_{2} Y_{2}, z_{i}^{1}+\rho_{2} Z_{2}\right)$; the horizontal axis corresponds to exponents of $x_{0}$ and the vertical axis to exponents of $\rho_{2}$.

Let $\rho_{2}=x_{0}^{r_{2}}, r_{2}$ any positive rational number. In particular we will soon chose it using the slope of a side $\Delta$ of the second Newton-lens polygon.

We get

$$
\begin{cases}\Phi_{2}^{*}\left(\widetilde{\omega_{1}}\right)\left(X_{2}, Y_{2}, Z_{2}\right) & =x_{0}^{r_{2}} \widetilde{\omega_{2}}=x_{0}^{r_{2}} \omega_{2}\left(X_{2}, Y_{2}, Z_{2}\right)+o\left(x_{0}^{r_{2}}\right), \quad \text { with }  \tag{6}\\ \omega_{2}\left(X_{2}, Y_{2}, Z_{2}\right) & =A_{2}\left(X_{2}, Y_{2}, Z_{2}\right) d X_{2}+B_{2}\left(X_{2}, Y_{2}, Z_{2}\right) d Y_{2} \\ & +C_{2}\left(X_{2}, Y_{2}, Z_{2}\right) d Z_{2}\end{cases}
$$

Notice that $\widetilde{\omega_{2}}$ keeps terms of different orders in $x_{0}$. When dealing with the equation $f=0$ we will use similarly the notations $\widetilde{f_{2}}$.

Let us first chose a vertex $v$ of the Newton-lens polygon. It corresponds to a term of the form $x_{0}^{s} Q_{A, 2}, x_{0}^{s} Q_{B, 2}$ or $x_{0}^{s} Q_{C, 2}$ of $\widetilde{A}, \widetilde{B}$ or $\widetilde{C}$, where the polynomials $Q_{A, 2}, Q_{B, 2}$ or $Q_{C, 2}$ are homogeneous of the same degree.

The plane-field is then invariant on rays. Therefore the rays are contained in the polar set at this scale.

The roots $\left(y_{i, 1} z_{i, 1}\right)$ belong to the plane $x=x_{0}$
Moving the point $x_{0}$ we can follow the roots $\left(y_{i, 1} z_{i, 1}\right)$. At the scale of the previous change of variable, and even more at the scale of the present one, they describe vertical lines.

Therefore the Gauss map associated to $\omega 2$ cannot contain an open set. Its image is therefore of dimension 1 . The polar loci are then generically surfaces union of rays. The existence of two
different cones is also incompatible with a small change of $x_{0}$. The polar surfaces should then be planes containing the axis $0 x$. This means that the direction $[A, B, C]$ depends only on its value on the $Y_{2}, Z_{2}$-plane where it is constant on rays.

We chose now an exponent $r$ corresponding to the slope of a support line of the Newton-lens polygon at $v$.

Choosing a side of the Newton-lens polygon ending on the $x$-axis if it exists, and if $A_{2}(x, y, z)$ contributes to this term, will provide a limit plane field. As before, we will prove that the limit plane-field is invariant by translations. Otherwise, we should continue the induction.

In both case we need to prove, first that the limit plane-field is invariant by translations along lines parallel to the $X_{2}$-axis, then that the polynomials $A_{2}, B_{2}$ and $C_{2}$ do not depend on the variable $X_{2}$.

Theorem 5.2.1. The limit plane field is invariant by translations and therefore, when $\omega$ is integrable, the leaves of the foliation it defines are graphs, maybe with some vertical asymptotes, from the $\left(Y_{2}, Z_{2}\right)$-plane to the $X_{2}$-axis. We call profile one of these graphs.

The reader can find in [La-Si] a less geometric proof in $\mathbb{C}^{2}$ of Theorem 5.2.1.
Remark 5.2.2. In [GarBar-Te] Garcia- Barroso and Teissier compute the lowest exponent of the Puiseux serie of a strand of generic polar curve involving the slope of the line direction defining the polar curve in terms of the value $\lambda$ defining the level $f=\lambda$. It gives both the size and the localization in terms of $\lambda$ of a domain of our profile meeting the level $f=\lambda$.

Proof: of theorem 5.2.1 Corollary 4.5.3 implies that there exists two values of $r_{a}>r_{b}>1$ close to 1 which determine a zone $Z_{O x, r_{a}, r_{b}}$ (see Definition 4.4.1)where the planes $h=\operatorname{ker} \omega$ define a Gauss map of very thin image $G$. Therefore a polar curve $\Gamma_{h}, h$ not too close from $G$, which starts out of $Z_{O x,\left(r_{b}-\delta\right)}$ (see also Definition 4.4.1) cannot cross $Z_{O x, r_{a}}$. The polar curve $\Gamma_{h}$, if it starts at a point of $Z_{O x, r_{a}}, r_{a}>r_{c}>r_{b}$ has to stay in $Z_{O x, r_{b}}$ and therefore is tangent to the axis $O x$ at the origin. Then the strands of polar curves $\Gamma_{h}, h$ not too close from $G$, form a bouquet, and look through a lens of strength $x_{0}^{r}, r_{a}>r>r_{b}$, more and more like parallel lines when $\left|x_{0}\right| \rightarrow 0$. Therefore the plane field defined by $\operatorname{ker} \omega_{2}$ is invariant by translation along the $X_{2}$-axis. Then, when $\omega$ is integrable, so is $\omega_{2}$ and the leaves of $\omega_{2}$ are graphs (we accept vertical asymptotes) of functions not depending on the variable $X_{2}$.

Remark 5.2.3. Remark 4.4.2 shows that the consideration of the larger zone $Z_{O x, a_{1}, b_{1}}$ is already sufficient to trap the polar branches.

In order to continue the induction we need a more "algebraic" result.
Theorem 5.2.4. The polynomials $A_{2}\left(X_{2}, Y_{2}, Z_{2}\right), B_{2}\left(X_{2}, Y_{2}, Z_{2}\right)$ and $C_{2}\left(X_{2}, Y_{2}, Z_{2}\right)$ do not contain the variable $X_{2}$.

Proof: We know that the plane field $\operatorname{ker} \omega_{2}$ is invariant by translations parallel to the $X_{2}$-axis, therefore the point $\left[A_{2}\left(X_{2}, Y_{2}, Z_{2}\right), B_{2}\left(X_{2}, Y_{2}, Z_{2}\right), C_{2}\left(X_{2}, Y_{2}, Z_{2}\right)\right]$ does not depend of $X_{2}$ but $A_{2}$, $B_{2}, C_{2}$ may a priori depend on $X_{2}$.

Let us consider some (fixed) value $Y_{2,0}$ and $Z_{2,0}$ of $Y_{2}$ and $Z_{2}$. The polynomial $A_{2}\left(X_{2}, Y_{2,0}, Z_{2,0}\right)$, if it is not constant, should have some root $X_{2,0}$. If one of the polynomials $B_{2}\left(X_{2,0}, Y_{2,0}, Z_{2,0}\right)$ and $C_{2}\left(X_{2,0}, Y_{2,0}, Z_{2,0}\right)$ is not zero, it is also not zero in a neighborhood of ( $X_{2,0}, Y_{2,0}, Z_{2,0}$ ). Then the plane $\operatorname{ker} \omega_{2}$ has a position which is different in the neighborhood of the point $X_{2} \in$
the line $\left\{Y_{2}=Y_{2,0}, Z_{2}=Z_{2,0}\right\}$ and at the point $X_{2,0} \in$ the line $\left\{Y_{2}=Y_{2,0}, Z_{2}=Z_{2,0}\right\}$. Therefore the polynomial $A_{2}\left(X_{2}, Y_{2,0}, Z_{2,0}\right)$ should be constant. Similarly $B_{2}$ and $C_{2}$ should be constant on the line $\left\{Y_{2}=Y_{2,0}, Z_{2}=Z_{2,0}\right\}$, unless the three polynomials $A_{2}, B_{2}, C_{2}$ are simultaneously zero on the line $\left\{Y_{2}=Y_{2,0}, Z_{2}=Z_{2,0}\right\}$.

Notice that one or two of the constant values of $A_{2}\left(X_{2}, Y_{2,0}, Z_{2,0}\right), B_{2}\left(X_{2,0}, Y_{2,0}, Z_{2,0}\right)$ or $C_{2}\left(X_{2,0}, Y_{2,0}, Z_{2,0}\right)$ maybe zero.

Let us now suppose that the three values $A_{2,0}\left(X_{2,0}, Y_{2,0}, Z_{2,0}\right), B_{2,0}\left(X_{2,0}, Y_{2,0}, Z_{2,0}\right)$ and $C_{2,0}\left(X_{2,0}, Y_{2,0}, Z_{2,0}\right)$ are zero.

A priori the set $\Sigma_{2}=\left\{A_{2}=B_{2}=C_{2}=0\right\}$ is an algebraic set.
If a zero $m_{2} \in \Sigma_{2}$ of the map

$$
\left(A_{2}, B_{2}, C_{2}\right): \mathbb{C}^{3} \rightarrow \mathbb{C}^{3}
$$

is isolated, it is of non-zero index. Therefore, for $x_{0}$ small enough, the map $\left(\widetilde{A_{2}}, \widetilde{B_{2}}, \widetilde{C_{2}}\right): \mathbb{C}^{3} \rightarrow$ $\mathbb{C}^{3}$ (see Equation 3 and the beginning of Subsection 5.2 for a definition of $\tilde{A}, \tilde{B}$ and $\tilde{C}$ ) has also a zero in a neighborhood of ( $X_{2,0}, Y_{2,0}, Z_{2,0}$ ). This contradicts the fact that the origin in an isolated singular point of $\omega$.

Otherwise the algebraic set $\Sigma_{2}$ may contain curves and surfaces. The curves cannot be only lines parallel to the $X_{2}$-axis as some common zeros of $\left(A_{2}, B_{2}\right.$ and $\left.C_{2}\right)$ are isolated on the line parallel to the $X_{2}$-axis trough it.

Let us now consider a regular point $m_{2}$ of an arc $\alpha_{2}$ contained in $\Sigma_{2}$, and a transverse holomorphic disc $D_{2}$ to the arc at the point. It should contain an isolated zero of the map $\left(B_{2}, C_{2}\right): D_{2} \rightarrow \mathbb{C}^{2}$. For the same reason as above the map $\left(\widetilde{B_{2}}, \widetilde{C_{2}}\right): D_{2} \rightarrow \mathbb{C}^{2}$ should have a zero close to $m_{2}$. Moving the point $m_{2}$ on a neighborhood of $m_{2}$ on $\alpha_{2}$ we get a piece of the polar curve $\Gamma_{h}, h=\left(O Y_{2}, O Z_{2}\right)$. It should have $O x$ as limit tangent at the origin. Therefore it should look more and more parallel through the second lens to the $O X_{2}$ axis.

Let us now consider a regular point $m_{2}$ of a surface $S_{2}$ contained in $\Sigma_{2}$ where the tangent plane does not contain $O X_{2}$. The polar curve trough $m_{2}$ cannot have a tangent having a limit in $T_{m_{2}} S_{2}$ providing a contradiction.

Remark 5.2.5. Using a Newton-lens providing a profile, we do not loose the polar branches going to the limit $x_{0} \rightarrow 0$. When a branch of the Newton-lens induction ends at a vertex, the situation is different. Nevertheless the limit 1-form is still independent of the last $X$-variable say $X_{p}$. The conclusion of Remark 4.4.2 are still valid. Therefore, in any $C^{k}$ topology the dependence of the form $\omega_{p}$ on the variable $X_{p}$, considering derivatives of order up to $k$, is going to zero. The limit form, which is polynomial, therefore does not depend on the variable $X_{p}$. This is true although the Gauss map of the limit form $\omega_{p}$ is of rank strictly less that 2 as the limit polar loci are planes containing the $X_{p}$-axis.

If needed, the following steps of the induction provide forms $\omega_{p}=A_{p} d X_{p}+B_{p} d Y_{p}+C_{p} d Z_{p}$ where the polynomial $A_{p}, B_{p}$ and $C_{p}$ do not depend on the variable $X_{p}$.

The induction stops when $A_{p} \neq 0$.
Remark 5.2.6. The profiles, up to linear isomorphism, are analytic invariants of the isolated singularity. In dimension 2, Zariski defined analytic invariant of an irreducible curve of equation $f(x, y)=0$ using relevant coefficients of the Puiseux expansion of the curve (see [Za]). The dimension of Zariski's module is larger that the dimension provided by the profiles. It would be interesting to understand which information the profiles give about Zariski's invariants.

## 6 What does Newton lenses tell about the level $f=0$ ?

### 6.1 The singular level $\{f=0\} \subset \mathbb{C}^{2}$

Remark 6.1.1. We know that $f$ is reduced, as a multiple factor will introduce a curve of singular points. The induction should be stopped when the induction concerning df shows a profile. We may as well suppose that $f$ is irreducible, and study the possible component one by one.

The tangent cone of the singular level $\{f=0\} \subset \mathbb{C}^{2}$ is where the homogeneous polynomial sum of the terms of lowest degree of $f$ is zero; it is a finite union of lines. The interesting lines for us are the lines of multiplicity at least 2 .

The change of variables $\Phi_{1}$ is defined by

$$
\begin{aligned}
& x=x_{0}+\rho_{1} X_{1} \\
& y=\rho_{1} Y_{1}
\end{aligned}
$$

Let us chose one side of slope $-1<1 / r_{1}<0$ of the Newton-lens polygon of $f$ (case 1 ) or a vertex bounded by two sides of slope $1 / r 0,1$ and $1 / r 0,2-1<1 / r_{0,1}<1 / r_{0,2}<0$ (case 2).

Then, taking $\rho_{1}=x_{0}^{r}, r=r_{1}$ in the first case, $1 / r_{0,1}<1 / r<1 / r_{0,2}$ in the second case, consider

$$
\widetilde{\Phi}_{1}^{*}(f)\left(X_{1}, Y_{1}\right)=f\left(x_{0}+x_{0}^{r} X_{1}, x_{0}^{r} Y_{1}\right)
$$

Proposition 6.1.2. Choosing $\rho=x_{0}^{r}$ the sum of the terms corresponding to points of this side of the Newton-lens polygon provides, after factorizing the highest possible power of $x_{0}$, an equation $f_{1}\left(Y_{1}\right)$ (which does not depend of the variable $X_{1}$ ). Choosing a slope corresponding to a support line of a vertex of the first Newton-lens polygon, we get $f_{1}\left(Y_{1}\right)=Y_{1}^{p}$.

Proof. Writing $f(x, y)=\sum a_{p, q} x^{p} y^{q}$, a non zero monomial $a_{p, q} x^{p} y^{q}$ gives a polynomial $a_{i, j}\left(x_{0}+\right.$ $\left.\rho_{1} X_{1}\right)^{p} \rho_{1}^{q} Y_{1}^{q}$ in $f\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}\right)$.

One point comes from the monomial $x_{0}^{p} \rho_{1}^{q} Y_{1}^{q}$.
The other points introduced by the terms of $a_{p, q}\left(x_{0}+\rho_{1} X_{1}\right)^{i} \rho_{1}^{j} Y_{1}^{j}$ containing $X_{1}$ form a tail of slope -1 on the left of ( $p, q, \ell$ ) (in black in Figure 8, here we apply the same reasoning to $f$ ).

Two possibilities can occur:
-1) the equation $f_{1}\left(Y_{1}\right)=0$ has only simple roots.
-2) some roots are multiple, more precisely, some factors of $f_{1}\left(Y_{1}\right)$ have an exponent strictly greater than $1 ; f_{1}=f_{1,1}^{a_{1}} \cdot f_{1,2}^{a_{2}} \cdots$.
Example with only simple roots $f(x, y)=x^{3}-y^{2}$.
$\Phi_{1}^{*}(f)=\left(x_{0}+\rho_{1} X_{1}\right)^{3}-\left(\rho_{1} Y_{1}\right)^{2}$; Choosing $\rho_{1}=x_{0}^{3 / 2}$, we get $f_{1}=1-Y_{1}^{2}$.
Remembering that $y=x_{0}^{3 / 2} Y_{1}$ and following a circle of radius $\left|x_{0}\right|$ centered at the origin in the $x$-plane, we recover a Lê carrousel (see [Le]), and a trefoil knot in $S_{\left|x_{0}\right|}^{1} \times(y$-axis) (the $y$-axis is a complex line, that is of real dimension 2 ).

### 6.2 Multiple roots of $f_{1}\left(Y_{1}\right)=0$

Example with multiple roots $f(x, y)=y^{4}-2 x^{3} y^{2}-4 x^{5} y+x^{6}-x^{7}$ (example taken from [Bri])
From the construction of the Newton-lens polygon, we get again $\rho_{1}=x_{0}^{3 / 2}$ and $f_{1}=\left(Y_{1}^{2}-1\right)^{2}$; the two roots $Y_{1}= \pm 1$ are double roots.

Next step provides $\rho_{2}=x_{0}^{1 / 4}$ and $f_{2}\left(Y_{2}\right)=Y_{2}^{2}-1$, equation which have simple roots $Y_{2}= \pm 1$.

When $f_{1}=0$ ) has multiple roots, we need to proceed and perform another change of variables $\Phi_{2}$, with $c_{1}$ a roots of one of the multiples factors $f_{1, i}^{a_{i}}$.

$$
\begin{aligned}
& X_{1}=\rho_{2} X_{2} \\
& Y_{1}=c_{1}+\rho_{2} Y_{2}
\end{aligned}
$$

We get a polynomial in the variables $x_{0}, \rho_{2}$ from the expression $\Phi_{2}^{*}{\widetilde{\Phi_{1}}}^{*}(f)$. Choosing a side of the Newton-lens polygon indicates the enlarging rate $\left(1 / \rho_{2}\right)$. If the polynomial obtained from the sum of the terms of $\Phi_{2}^{*}{\widetilde{\Phi_{1}}}^{*}(f)$ corresponding to point of the side of the second Newton-lens polygon, after factorizing the maximal power of $x_{0}$, has only simple roots the induction stops.

The form of the two successive changes of variables in the $y$-plane

$$
\begin{aligned}
& y=\rho_{1} Y_{1} \\
& Y_{1}=c_{1}+\rho_{2} Y_{2}
\end{aligned}
$$

provides, following $x_{0}=\left|x_{0}\right| e^{i \theta}$ in the $x$-plane, a two-step Lê carousel in the $y$-plane. It provides also an iterated torus knot, component of $\{f=0\} \cap S_{\left|x_{0}\right|}^{1} \times y$ - plane. In general, the number of steps of the induction provides the number of iterations of the torus $\operatorname{knot}\left\{(f=0) \cap \mathbb{S}_{\varepsilon}^{3}\right\}$.

### 6.3 In $\mathbb{C}^{3},\{f=0\}$

The use of Newton lenses leads, when the tangent cone of $\{f=0\}$ has only isolated singular rays, to a decomposition of the link $K=\{f=0\} \cap \mathbb{S}_{\varepsilon}^{5}$ in components which are either $\mathbb{S}^{1}$-fiber bundles over an algebraic curves deprived of a finite number of discs or fiber bundles over $\mathbb{S}^{1}$, glued along tori.

The tangent cone of $\{f=0\}$ defines a curve $C_{\text {tang }}$ of $\mathbb{C P}^{2}$ with isolated singularities $\left\{\sigma_{i}\right\}$; deleting small balls $B_{\sigma_{i}, \varepsilon}^{4}$ of radius $\varepsilon$ centered at the singular points $\left\{\sigma_{i}\right\}$, we get a regular complex curve $C_{\text {tang }, \varepsilon}$. Its inverse image by the Hopf map $\mathbb{S}^{5} \rightarrow \mathbb{C P}^{2}$ is a $\mathbb{S}^{1}$ fiber bundle $A$ bounded by tori product of a Hopf fiber and a component of one of the (usual) links $\kappa_{i}=\mathbb{S}_{\sigma_{i, \varepsilon}}^{3} \cap C_{\text {tang }}$.

Let us now suppose as usual that the $x$-axis is an isolated singular ray of the tangent cone of $\{f=0\}$. Consider the change of variables

$$
\begin{align*}
& x=x_{0}+\rho_{1} X_{1} \\
& y=\rho_{1} Y_{1}  \tag{7}\\
& z=\rho_{1} Z_{1}
\end{align*}
$$

Then

$$
\Phi_{1}^{*}(f)\left(X_{1}, Y_{1}, Z_{1}\right)=f\left(x_{0}+\rho_{1} X_{1}, \rho_{1} Y_{1}, \rho_{1} Z_{1}\right)
$$

Lemm 6.3.1. Factorizing the maximal power of $x_{0}$ from the sum of the terms corresponding to a side of the Newton-lens polygon of slope $-1 / r_{1}$ and choosing $\rho_{1}=x_{0}^{r_{1}}$, we get a polynomial $f_{1}\left(Y_{1}, Z_{1}\right)$ which does not depend on the variable $X_{1}$.

The proof is the same as the proof of Proposition 4.3 .1 (see Figure 8).
Remark 6.3.2. The vertices of the Newton-lens polygon also give rise to an homogeneous polynomial $f_{1}$ of the two variables $Y_{1}$ and $Z_{1}$ only. The equation $f_{1}=0$ represents in this case a finite number of lines.

Corollary 6.3.3. The knot $K$ is, when $f$ is irreducible, when the singular rays of the tangent cone of $f$ are isolated, and when one step is sufficient to provide a smooth non-multiple transverse profile, obtained from a fiber space A over $C_{\text {tang, }}$ gluing along the boundary tori of $A$ a fiber space over $\mathbb{S}^{1}$ (a circle in a complex line parallel to the $x$-axis) with fiber an algebraic curve deprived of a finite number of discs (a bounded part of an algebraic curve of $\mathbb{C}^{2}$ ).

## 7 Further steps, polynomial profiles

### 7.1 The second step for the level $f=0$

Recall that the $x$-axis is singular for $\omega$.
Chose first a regular point $\left(c_{1}, c_{2}\right)$ of the algebraic curve of equation $f_{1}\left(Y_{1}, Z_{1}\right)=0$.
Let us prove the
Proposition 7.1.1. When the point $\left(c_{1}, c_{2}\right)$ is regular, the function $f_{2}\left(X_{2}, Y_{2}, Z_{2}\right)$ is cylindrical, that is, depends only on the variables $Y_{2}$ and $Z_{2}$.

Proposition 5.2.4 shows that the assertion is true for $d f$. In fact, the weaker statement: the plane field $\operatorname{ker} \omega_{2}$ is invariant by translation parallel to the $X_{2}$-axis is sufficient to prove Proposition 7.1.1. We know also that the level $f_{2}=0$ need to be of the form $A_{2} \times x$-axis, as it is the limit of approximations of the level $f=0$ of $f$.
Proof: As we look at analytic curves tangent to the $x$-axis at the origin, after the second enlargement, a small neighborhood of a point will look more and more like a piece of the $x$-axis. Theorem 5.2.4 shows that the coefficients of the forms $\omega_{2}$ do not depend on the variable $X_{2}$. The places where ker $\omega$ differs significantly from a direction containing $O x$ is a finite union of zones union of branches of polar curves of bouquets. A bouquet is visible through a Newton lens constructed at the last step of the Newton-lens induction, when the slope of the bottom right segment of the Newton-lens polygon of $f$ is (negative and ) smaller than the bottom right segment of the Newton-lens $d f$ (see Figure 12). Then the last enlargement obtained by an induction adapted to $f$ does not catch the nearest bouquet and the previous ones do not see them as their limit is a line parallel to the $O x$-axis above a root of the equations ( $B_{1}=C_{1}=0$ ) (see the end of Subsection 5.2). Notice that we already know that $f_{2}=0$ can contain only lines parallel to the $x$-axis and that $(d f)_{2}$ does not depend on $X_{2}$. Therefore $f_{2}$ cannot depend on the variable $X_{2}$.

When the slope $-1<s<0$ of the support line is such that the support point is a vertex bounded by sides of slope $-1 / r_{2,1}$ and $-1 / r_{2,2}$ of a second Newton-lens polygon, the function $f_{2}$ is homogeneous in the two variables $Y_{2}$ and $Z_{2}$. In the plane $x=x_{0}$, the equation $f_{2}=0$ has then as solutions a finite number of lines through the origin.
Lemma 7.1.2. - When the enlarging ratio $\rho_{2}=x_{0}^{r_{2,1}}$, the regular points of $f_{2}\left(Y_{1}, Z_{1}\right)$ provide a transverse Lê carousels and therefore fibered pieces of a link $K$, defined by the action of $S^{1}$ obtained following a circle $S_{x_{0}}^{1}=\left\{x_{o} e^{i \theta}\right\}$.

- a vertex bounded by sides of slope $-1<s<0$ provides 3-dimensional regions of the form (annulus) $\times \mathbb{S}^{1}$.

The arguments necessary to continue the induction, if required, are similar.
The hypothesis "isolated singularity" rules out the possibility of multiple components, therefore the induction stops when $f_{p}\left(Y_{p}, Z_{p}\right)$ has no multiple points. Simultaneously, following the same path of inductive steps, $p$ is the step where a profile of $\omega=d f$ appears (see Diagram 7.2).

Lemma 7.1.2 implies Theorem 7.1.3
Theorem 7.1.3. Let $f$ be an algebraic hypersurface with an isolated singular point at the origin. Suppose also that its tangent cone has only isolated singular rays. Let $S_{\varepsilon}^{5}$ be the sphere centered at the origin of radius $\varepsilon$. When $\varepsilon$ is small enough, the link $K=\mathbb{S}_{\varepsilon}^{5} \cap\{f=0\}$ is obtained from the $\mathbb{S}^{1}$-fiber bundle $A$ defined at the beginning of the subsection by gluing a sequence of fiber bundles over $\mathbb{S}^{1}$ of fiber the intersection of a complex curve deprived of neighborhoods of its singular points with a large 4-ball of $\mathbb{C}^{2}$, and at the end of each sequence, a fiber bundle over $S^{1}$ obtained following a circle $S_{x_{0}}^{1}=\left\{x_{o} e^{i \theta}\right\}$, and of fiber the intersection of an algebraic curve of $\mathbb{C}^{2}$ with no singularity and a large 4-ball .

### 7.2 Polynomial profile when $\omega=d f$

Theorem 7.2.1. When $\omega=d f, f$ a polynomial with an isolated singularity at the origin and such that $O x$ is an isolated singular ray of the tangent cone, then profiles are graphs of polynomials.

Theorem 7.2.1 is already proved in $\mathbb{C}^{2}$ (see [La-Si]).
Remark 7.2.2. Theorem 7.2.1 is closely related to the position of the bouquets of generic polar curves with respect to the singular level $f=0$, generalizing the 2-dimensional result of Merle (see [Me].

Given a form $\omega$, the profile may be given by differential equations with coefficient rational functions. The saddle-node of equation $\omega=y d x-x^{2} d y$ is already an example of this phenomenon in $\mathbb{C}^{2}$ (see figure ??).

The diagram 7.2 is a scheme of the proof of Theorem7.2.1. The notation for Newton polygons is $\mathcal{N}$.

not a double root of $f_{1}=0$ double root of $f_{1}=0$

separatrices are kept away $\underset{\mathcal{N}}{\mathcal{N}}(d f)$ is obtained from $\mathcal{N}(f)$ by a vertical translation, but, maybe, for a term on $O x$
a term on $O x$ is the end of segments of different slopes of $\mathcal{N}(f)$ and $\mathcal{N}(d f) \Longrightarrow$ polar curves are kept away from $\vec{f}=0$
$A_{2}=0$
-
$\bullet$

### 7.3 A proof of Theorem 7.2.1

Lemma 7.3.1. The change of variables of the Newton-lens induction and differentiation commute when the points chosen in the $\left(Y_{p}, Z_{p}\right)$-plane used to continue the induction are the same for $f_{p}$ and $A_{p}$. In other terms

$$
d_{X_{p}, Y_{p}, Z_{p}} \Phi_{p} \circ \varphi_{p-1} \circ \Phi_{1} f=\Phi_{p} \circ \Phi_{p-1} \circ \Phi_{1} d_{x, y, z} f
$$

This fact is sufficient to compare efficiently the Newton-lens polygons of $f$ and $d f$ when the induction needs more than one step. We need now at each extra step to compare the multiple roots of the equation $f_{p}\left(Y_{p}, Z_{p}\right)=0$ and the roots of the system of equations $B_{p}\left(Y_{p}, Z_{p}\right)=0$ and $C_{p}\left(Y_{p}, Z_{p}\right)=0$.

The two-dimensional case will guide us. We had noticed in Section 6, observing Newtonlens polygon at different stages that the double roots of $f_{p}\left(Y_{p}\right)=0$ are roots of $B_{p}(Y)=0$.

Lemma 7.3.2. Again, let us suppose that $\omega=d f$.

- When $A_{p}=0$, the Newton-lens cloud of $\Phi_{p} \circ \Phi_{p-1} \circ \Phi_{1} d f$ can be obtained from the Newtonlens cloud of $\Phi_{p} \circ \Phi_{p-1} \circ \Phi_{1} f$ using a vertical translation of vector $(0,-1)$ as the Newton-lens cloud of df can be obtained from the Newton-cloud of $f$ when $A_{1}=0$.
- When $A_{p} \neq 0$ its only contribution to the Newton-lens polygon of $\widetilde{d f_{p}}$ correspond to a point on the $x$-axis. This point is then the extremity of a segment of (negative) slope larger than the slope of the bottom segment of the Newton-lens polygon of $\widetilde{f_{p}}$.

When $\omega=d f$, if the polynomial $A_{1}\left(X_{2}, Y_{2}, Z_{2}\right)$ is non-zero, it is reduced to a constant term. The proof is the same as in first step we considered parallel side of the first Newtonlens polygons of $f$ and $d f$. Moreover the double roots of $f_{1}$ are the roots of $B_{1}=C_{1}=0$. Therefore $\left(\Phi_{2} \circ \Phi_{1}\right)^{*} d f=d\left(\Phi_{2} \circ \Phi_{1}\right)^{*} f$. The only term of $A_{2}$ should then come from the term of $\left(\Phi_{2} \circ \Phi_{1}\right) * \frac{\partial f}{\partial x} f$ on the $x_{0}$-axis. We get a profile from the differential equations

$$
\begin{aligned}
& \frac{\partial X_{2}}{\partial Y_{2}}=B_{2}\left(Y_{2}, Z_{2}\right) / A_{2} \\
& \frac{\partial X_{2}}{\partial Z_{2}}=C_{2}\left(Y_{2}, Z_{2}\right) / A_{2}
\end{aligned}
$$

when the polynomial $A_{2}$ is a non-zero constant. If $A_{2}=0$, we should continue the induction.

### 7.4 Merle-type results (2)

Theorem 7.4.1. When $\omega=d f$, the bouquets of polar curves tangent to the $x$-axis are separated from the singular level $f=0$ by a security funnel.

Proof: The computation above shows that the small ball of radius $\left|x_{0}\right|^{r_{1}+\ldots+r_{p}^{d f}}$ containing points of the polar curves of the bouquet is much smaller that the ball of radius $\left|x_{0}\right|^{r_{1}+\ldots+r_{p}^{f}}$ containing points of $f=0$ nearby. Then the bouquet of polar branches contained in the zone $Z_{O x, r_{p}^{d f}}$ is separated by a zone $Z_{O x, r_{p}^{f}, r_{p}^{d f}}$ from the component of $f=0$ corresponding to the path leading to the profile corresponding to the bouquet of polar branches we a considering.

The zone $Z_{O x, r_{p}^{f}, r_{p}^{d f}}$ (see Remark 4.4.2) guarantees also the separation of the bouquet of polar branches from components of $f=0$ obtained by other paths, in particular if the component of $f=0$ we want to avoid is not tangent to the $x$-axis at the origin.


Figure 14: A bouquet of polar curves is separated from $f=0$ by a security funnel

When, at the end of a branch of the Newton-lens induction tree, we get a profile, we simultaneously get a bouquet of branches of polar curves corresponding to directions in $\mathbb{C P}^{2}$ avoiding a neighborhood of the line of planes containing the $x$-axis. The diagram (see Figure 14) shows that bouquets of branches of polar curves corresponding to different branches of the Newtonlens induction tree are separated and separated from $f=0$ by annular zones. This provide a "pedestrian" proof of Merle's result [Me].

Remark 7.4.2. - The domains of these bouquets are good candidates to provide Lipschitz equivalence of the intrinsic and extrinsic distances on a level $f=\lambda$ (see [Neu-Pi2] and [Pi-Te]).

- In $\mathbb{C}^{2}$ we can repeat Remark 4.7.2, giving an interpretation of Theorem 7.4.1 in terms of contact orders.


### 7.5 Separatrices

Definition 7.5.1. A separatrix of $\omega$ is an analytic arc $S$ ending at the origin such that

$$
\dot{S}(x(t), y(t), z(t)) \in \operatorname{ker} \omega \quad \forall t .
$$

In particular, when $\omega=d f$, and when the origin is the only singular point and is contained in the level $f=0$, the analytic curves contained in the level $f=0$ ending at the origin are the separatrices.

In the saddle-node example, the $x$-axis is a separatrix. Notice that in this case the separatrix stays in the bouquet of polar curves $\Gamma_{h}, h$ "far" from $D_{x}$.

In dimension 2, Rouillé (see [Rou3]) considered the role of separatrices in a particular case (quasi functions).

Definition 7.5.2. The separatrices of a 1 -form $\omega$ are kept away from the bouquets of polar curves $\Gamma_{h}$, $h$ not too close from the curve $D_{x} \subset \mathbb{C P}^{2}$ of planes containing the $x$-axis, if, in a small enough neighborhood of the singular point, none of these bouquets contains a separatrix.

Therefore, when $\omega$ is integrable, the "graph of polynomial" condition will be guaranteed if the separatrices tangent to the $x$-axis at the origin stay away from the very curved zones, or using Merle's ideas, are kept away from all the bouquets of polar curves $\Gamma_{h}, h$ "far from" $D_{x}$.

## 8 Non-isolated singular rays of the tangent cone, an example

- $f(x, y, z)=x^{4}-y^{3}+z^{2}($ see Figure 15$)$.
$\omega=d f=4 x^{3} d x-3 y^{2} d y+2 z d z$
The equation of the tangent cone is $z^{2}=0$, all the rays of this plane are degenerate.


Figure 15: The surface of equation $x^{4}-y^{3}+z^{2}=0$ and a level $x^{4}-y^{3}+z^{2}=\lambda$,
$f(x, y, z)=\lambda$. Let us compute $x_{0}$ and $\rho_{1}$ in terms of $\lambda$, using in fact [GarBar-Te] track.
$x_{0}^{4}=\lambda,\left|x_{0}\right|=|\lambda|^{1 / 4},\left|\rho_{1}\right|=|\lambda|^{3 / 4}$
Choosing the singular ray $O y$, that is using the change of variables
$x=x_{0}+\rho_{1} X_{1}, y=y_{0}+\rho_{1} Y_{1}, z=\rho_{1} Z_{1}$, we get
$\omega_{1}=4 \rho_{1}^{3} X_{1}^{3} d x-3\left(y_{0}+\rho_{1} Y_{1}\right)^{2} d y+2 \rho_{1} Z_{1} d z$.
The Newton-lens polygon provides
$\omega_{1}=-3 d Y_{1}+2 Z_{1} d Z_{1}$
Therefore

$$
\frac{d Y_{1}}{d Z_{1}}=(2 / 3) Z_{1}, \text { so } Y_{1}=\frac{Z_{1}^{2}}{3}+c
$$

$f(x, y, z)=\lambda$. Let us compute $y_{0}$ and $\rho_{1}$ in terms of $\lambda$.
$-y_{0}^{3}=\lambda,\left|y_{0}\right|=|\lambda|^{1 / 3},\left|\rho_{1}\right|=|\lambda|^{2 / 3}$
In order to understand the "bag" shape of the levels $f=\lambda$, let us look at the critical points of the orthogonal projection of the levels on the tangent plane $z=0$. These critical points have the


Figure 16: $f(x, y, z)=x^{4}-y^{3}+z^{2}$, Newton-lens polygon when $x_{0} \in O x$ and when $y_{0} \in O y$
equation $d f(0,0,1)=0$, that is $z=0$. These critical points are all in the horizontal plane. The form $\omega_{x, y}=4 x^{3} d x-3 y^{2} d y$ defines a foliation, in fact the levels of the function $f_{x, y}=x^{4}-y^{3}$, which are the projection of the critical locus of the restriction of the orthogonal projection on the horizontal $(z=0)$ plane of the levels $f=\lambda$. We can study the 2-dimensional function


Figure 17: $f(x, y, z)=x^{4}-y^{3}+z^{2}$, real picture (sketch) of the critical locus of the projection on the $(x, y)$-plane of polar curves and of two real levels
$f_{x, y}=x^{4}-y^{3}$. Its differential is $\omega_{x, y}=4 x^{3} d x-3 y^{2} d y$. The Newton-lens polygons of $f_{x, y}$ and $d f_{x, y}$ indicate that the choice $\rho_{1}=x_{0}^{3 / 2}$ will provide a profile.

We get the differential equation

$$
\frac{d X_{1}}{d Y_{1}}=\frac{3}{4} Y_{1}
$$

providing the solutions

$$
X_{1}=\frac{1}{4} Y_{1}^{3}+c
$$

This explains the shape of the sides of the "bag". The flat bottom of the bag comes from the flat piece of the level through the point $\left(x_{0}, 0,0\right)$, of equation $f_{x, y}=x_{0}^{4}$, which is at distance of the order $\left|x_{0}\right|^{4 / 3}$ of the origin.

Understanding the polar sets $V_{\ell}, \ell=H^{\perp}$, for planes $H$ containing a non-isolated ray of the tangent cone of $f$, and the image by the orthogonal projection on $H$ of the critical loci of the restriction to the levels $f=\lambda$ of the orthogonal projection on $\ell$ may help to understand better
the geometry of the neighborhood of an isolated singularity when the tangent cone have nonisolated singular rays. The symmetry of the example $f(x, y, z)=x^{4}-y^{3}+z^{2}$ with respect to the "horizontal plane" makes the study deceptively easy. The general case implies the study of the restriction of differential form $\omega$ to the surface $V_{\ell}$, which makes sense even if $\omega$ is not of the form $d f$.

This step will be generically unavoidable when we will study 1-forms in $\mathbb{C}^{4}$ or $\mathbb{C}^{n}, n \geq 4$.

## References

[Bri] E. Brieskorn and H. Knörrer. Planes algebraic curves, Modern Birkhäuser classics, Birkhäuser, (1986).
[Ful] W. Fulton. Algebraic Curves, an Introduction to Algebraic Geometry (1969), offered by his author at http://www.math.lsa.umich.edu/ wfulton/ under the title CurveBook.
[Ca-Li-Sa] C. Camacho, A. Lins Neto, P. Sad. Topological invariants and equidesingularisation for holomorphic vector fields, J. differential geometry 20 (1984) pp. 143-174.
[GarBar] E. Garcia Barroso. Sur les courbes polaires d'une courbe plane réduite, Proc. London Math. Soc; (3) 81 (2000) pp. 1-28.
[GarBar-Te] E. Garcia Barroso and B. Teissier. Concentration multi-échelles de courbure dans des fibres de Milnor, Commentarii Mathematici Helvetici 74 (1999) pp.398-418.
[La1] R. Langevin. Courbure et singularités complexes, Comment. Math. Helvetici 54 (1979) pp. 6-16.
[La2] R. Langevin. Integral geometry from Buffon to geometers of today. Cours spécialisés, Société Mathématique de France (2016).
[La3] R. Langevin. Focalisation de la courbure près d'une singularité, Preprint IMPA (1988).
[La4] R. Langevin. Geometry of algebraic curves, hypersurfaces, singularities and foliations Preprint, Juillet 1991 (after a congress in Oran, 1989).
[La-Si] R. Langevin and J-C. Sifre. Lenses on very curved zones of a singular foliation of $\mathbb{C}^{2}$, Topology and its Applications 234 (2018) pp. 397-414.
[Le] Lê Dung Trang. The geometry of the monodromy, Studies in Mathematics 8 (1978) pp. 157-173.
[Me] M. Merle. Invariants polaires des courbes planes, Inventiones Mathematicae 41 (1977) pp. 508-524.
[Mil] J. Milnor. Singular Points of Complex Hypersurfaces, Annals of Mathematics Studies 61 Princeton University Press (1968).
[Ne] L. Ness. Curvature of algebraic plane curves. I, Compositio Mathematicae 35 issue 1 (1977) pp. 57-63.
[Ne2] L. Ness. Curvature on holomorphic plane curves. II, Compositio Mathematicae 35 issue 2 (1977) pp. 129-137.
[Neu-Pi1] W.D. Neumann and A. Pichon. Lipschitz Geometry does not determine embedded topological type, arXiv 1506.0384.
[Neu-Pi2] W.D. Neumann and A. Pichon. Lipschitz Geometry of complex curves, Journal of Singularities (2014) pp. 225-234.
[Pi-Te] A. Pichon and B. Teissier. In preparation.
[PePei] M. Pe Peireira. Tête à tête twists Notes São Carlos workshop on singularities (2018).
[Rol] Rolfsen. Knots and links, Publish or perish, University of british Columbia (1976).
[Rou1] P. Rouillé. Courbes polaires et courbure, Thèse, Dijon, France (1996).
[Rou2] P. Rouillé. Sur les polaires de certaines 1-formes, [On the polars of certain oneforms] C. R. Acad. Sci. Paris Sér. I Math. 326, $n^{\circ} 6$ (1998) pp. 677-680.
[Rou3] P. Rouillé. Théorème de Merle : cas des 1-formes de type courbe généralisées, Boletim da sociedade brasileira de matemática vol. $30 n^{o} 3$ (1999) pp. 293-314.
[Tei1] B. Teissier. Variétés polaires. I. Invariants polaires des singularités d'hypersurfaces, Inventiones Math. $40 n^{o} 3$ (1977) pp. 267-292.
[Te2] B. Teissier. Variétés polaires. II, La Rabida conference (1981) Springer Lecture Notes $N^{o} 961$.
[Tei3] B. Teissier. Courbes polaires relatives et courbure d'hypersurfaces de niveau, Preprint (1990).
[Za] O. Zarisky. Le problème des modules pour les branches planes, lectures given at the mathematical center of the École Polytechnique (France), Hermann (1986). English translation : The Moduli Problem for Plane Branches, University Lecture Series Volume 39, American Mathematical Society, Providence (2006).

