



HAL
open science

Revisiting asphaltenes instability predictions by probing destabilization using a fully immersed quartz crystal resonator

Mohamed Saidoun, Thierry Palermo, Nicolas Passade Poubat, Jean-Philippe Gingras, Carrier Hervé, Jean-Luc Daridon

► To cite this version:

Mohamed Saidoun, Thierry Palermo, Nicolas Passade Poubat, Jean-Philippe Gingras, Carrier Hervé, et al.. Revisiting asphaltenes instability predictions by probing destabilization using a fully immersed quartz crystal resonator. *Fuel*, 2019. hal-02480610

HAL Id: hal-02480610

<https://hal.science/hal-02480610>

Submitted on 22 Oct 2021

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons Attribution - NonCommercial 4.0 International License

Revisiting Asphaltene Instability Predictions by Probing Destabilization using a Fully Immersed Quartz Crystal Resonator

Mohamed Saidoun^{a,c}, Thierry Palermo^{b,*}, Nicolas Passade-Boupat^a, Jean-Philippe Gingras^b, Hervé Carrier^c, Jean-Luc Daridon^{c,**}

^aTOTAL SA, PERL Pôle Economique 2, BP 47 - RD 817 - 64170 LACQ, France

^bTOTAL SA, CSTJF Avenue Larribau - 64018 PAU Cedex, France

^cLaboratoire des Fluides Complexes et leurs Réservoirs-IPRA, UMR5150, CNRS/TOTAL/Univ Pau & Pays Adour, 64000 PAU, France

Abstract

The method of Asphaltene Instability Trend (ASIST) was derived from London dispersions first principles and Flory-Huggins theory to model the microscopical appearance of unstable flocs using limited number of refractive index (n_D) measures on stock-tank oils (STO), series of liquid precipitants (n-alkanes) and routine PVT data [1].

In order to eliminate the tuning of the aging time of solutions in the method, the experimental protocol was modified by monitoring the signal of a fully immersed Quartz Crystal Resonator (QCR) to detect the instantaneous destabilization of asphaltene in the surrounding media. Isothermal titrations were performed under atmospheric pressure while mixing. The test matrix was designed to vary the chain length of the n-alkanes titrant and the temperature. Subsequent refractive indices at detected conditions were used to calculate solubility parameters (δ) and to predict the instantaneous detection of unstable asphaltene during the expansion of gas dissolved systems using the exact same immersed apparatus.

Predictions were then challenged against experimental observations of artificial live oil systems created by recombining dead oils with gas mixtures at pressures up to 1000bar. Isothermal Constant Mass Expansion (CME) experiments provided measured volumes as a function of pressure which served for modeling, along with the QCR sensor signal to probe the first aggregation and deposition of unstable asphaltene for each system at various conditions. Prediction results were compared to experimental detections for 2 different crude oils, temperatures ranging from 45°C to 120°C and gas dissolved concentrations ranging from 0 to 60 %mol.

Keywords: Asphaltene instability; Pressure; Prediction; Crude oil; Quartz Crystal Microbalance

*using Asphaltene Instability Trend (ASIST)[1]

*Thierry Palermo

**Jean-Luc Daridon

Email addresses: thierry.palermo@total.com (Thierry Palermo), jean-luc.daridon@univ-pau.fr (Jean-Luc Daridon)

1. Introduction

Spooky stories of oil & gas industry have been centering asphaltenes in many fundamental research interests since decades. Indeed, anticipation of unstable asphaltenes fouling in upstream oil production and in refining systems is a key step in early project phase. Deposit build-up have been observed in many field cases forming localized flow restrictions in the production facilities. These macro-molecules are classified as soluble in aromatic solvents and insoluble in light n-alkanes, which are produced along with the complex mixtures defined as crude oils. Resulting from changes of pressure, temperature and subsequently volumetric composition of their surrounding media during transport, their destabilization is a precursor to deposition. Previous fundamental investigations involved mapping of two-dimensional solubility parameters of model oil and crude oil systems [2]. Based on results of a large number of crude oils, authors stated that asphaltene stability was predominantly governed by Van Der Waals (VdW) interactions. A crude oil system contains very few permanent dipoles, for this reason generalizing VdW interactions to solely London dispersion interactions is a fair approximation. Consequently, predictive techniques of appearance of asphaltene instability were designed using laboratory measures of solvent parameters affecting dispersive forces such as densities [3, 4, 5, 6], molar volumes [7, 8, 9], volume ratios [10] or refractive indices [11, 12, 13, 1, 12]. Vargas and Chapman [14] related solubility parameters to density and refractive indices of crude oil in order to predict solubility parameters of oil with gas dissolved. Summaries of existing modelization tools of the asphaltene precipitation were recently published by Subramanian et al. [15] and Wiehe [16].

Among, these models, the Asphaltene Instability Trend (ASIST) was developed from concepts mentioned above by New Mexico Institute of Mining and Technology in the frame of DeepStar consortium [11, 12]. The experimental approach uses microscopical detection of unstable asphaltenes with various liquid n- alkanes additions, limited number of refractive indices (n_D) measurements on stock-tank oils (STO) and routine PVT data to define a linear threshold of stability. The trend is extrapolated to predict the appearance of the least stable asphaltene due to light alkanes expansion during live crude oil transport in wells and pipelines. Comparisons of ASIST predictions for pressure depletion experiments performed in crude oils with dissolved gas systems have shown a very good agreement between the measurements and the model expectations [17, 18, 19]. Authors claim the existence of an optimal aging time of mixtures (ranging from 1 to 24h) after addition of liquid precipitants. Indeed the aging time was used as a tuning parameter for extrapolation to low molar volumes of precipitants in order to fit predictions to observations of flocs upon depressurization. This practice suggests that despite the linear thresholds of ASIST, the aggregation kinetics would vary on a single linear asphaltene instability trend.

However, we should note that experimental detections of particles during depressurization were carried out by near-infrared (NIR) light transmittance or by high pressure filtration which are different compared to the microscopy (used at atmospheric pressure) and can generate uncertainties subsequently causing the need to tune the aging time of ambient pressure measures. The microscopy observations of unstable asphaltenes are subject to human visual interpretation; but more importantly, as illustrated by Maqbool et al. [20], the appearance of clear particles of about 2 to 3 μm is necessary to confirm the presence of unstable asphaltenes. In pressurized apparatus, the source of laser-light wavelength in NIR solid detection systems (SDS) is around 850nm while the pore size of high pressure filtration systems used was 500nm. All experimental probing techniques used to validate ASIST have a resolution comprised in the order of 0.5 to 2 μm . However, since the high pressure apparatus were slightly more sensitive than microscopy, the tuning of the aging to longer times for microscopy detections to fit NIR or high pressure filtration data was expected. Therefore, we suspect that the tuning of the aging time is an artificial way to correct the difference in sensitivity of the techniques used.

47 On the other hand, some other authors have extensively studied the kinetics of asphaltenes destabilization
 48 and showed how long the aggregation process can take to reach micrometer sized flocs and even longer
 49 times to reach equilibrium[20, 21, 22]. They showed that the time to detect micrometer sized particles
 50 scales inversely exponentially to the concentration of precipitant added to the oil in time scales ranging from
 51 minutes to thousands of hours. In the same range of investigations; Haji-Akbari et al. [23, 24] eventually
 52 found a proportionality between the difference of solubility parameters (δ_{asph} the solubility parameter of
 53 asphaltenes, $\delta_{solution}$ the solubility parameter of the mixture solution) and a natural logarithm function of the
 54 total concentration of unstable asphaltenes called $C_1(0)$, the microscopy detection time of particles $t_{detection}$
 55 and the viscosity μ of the media.

$$\ln(t_{detection} \sqrt{C_1(0)}/\mu) \propto \frac{1}{\delta_{asph} - \delta_{solution}} \quad (1)$$

56 This so-called unified model described by equation 1 reminds us that despite the very long time that
 57 asphaltenes can take to reach a micrometer detectable size, the concentration of unstable asphaltenes can be
 58 significantly high in a relatively viscous media. The relevance of small unstable aggregates to the deposi-
 59 tion was revealed by Hoepfner and Fogler [25] [26] and Vilas Bôas Fávero et al. [27]. Indeed, significant
 60 accumulations of deposit were experimentally observed while flowing mixtures of n-heptane and crude oil
 61 for which unstable flocs had not reach a microscope detectable size ($\sim 0.5-2\mu m$). The size of depositing par-
 62 ticles was back-calculated from their mass diffusivities to diameters between 10 and 100nm, 5 to 100 times
 63 smaller than the detected flocs in high pressure monitoring techniques mentioned earlier. These findings
 64 suggest that the size of unstable asphaltenes matters to potential deposition mechanisms, smaller aggregates
 65 will cause a faster deposition rate under diffusion-limited regime. One should note that the concentration of
 66 unstable particles is another significant parameter to such a deposition mechanism.

67
 68 Therefore a better resolution of the detection of unstable asphaltenes was necessary and methods re-
 69 ferred as direct detections cited above were judged not sensitive enough according to Tavakkoli et al. [28].
 70 Consequently, an indirect method of detection was developed by these authors for ambient pressure desta-
 71 bilization experiments. The methodology was tested and validated using a model oil, which conveniently
 72 has a low viscosity and enables centrifuging out relatively small particles thanks to the good separation ef-
 73 ficiency. However, higher viscosities or high pressure recombined oil fluids still would require large flocs
 74 of asphaltenes to form for their reliable detection using the latter technique. At elevated pressures, the
 75 wide-spread SDS using NIR light transmittance, the high pressure filtration mentioned earlier or even high
 76 pressure microscopy are commonly used [29, 30] but their sensitivity is much larger than the diameters of
 77 diffusing particles calculated by Vilas Bôas Fávero et al. [27].

78 The purpose of this study is to validate ASIST without tuning the aging time and to confirm the contribution
 79 of small particles (10-100nm) to deposition upon depressurization. Consequently we revisited the instru-
 80 mentation used to detect asphaltene destabilization in ASIST and strictly applied the same data processing
 81 method. The following advantages were combined:

- 82 • instantaneously sensitive to orders of 10nm sized unstable aggregates [31, 32, 33]
- 83 • applicable to atmospheric pressure titration of crude oils with liquid n-alkanes
- 84 • applicable to depressurization of light alkanes (C1 to C4) recombined crude oils.

85 For that purpose, an acoustic sensor working in thickness shear mode was considered in this work. Indeed,
 86 Quartz Crystal Microbalance with Dissipation monitoring (QCM-D) is a well-known technique widely used
 87 to probe liquid adsorption of monolayer deposits or colloidal material on its surface thanks to piezo-electrical

88 properties of the quartz material. Its sensitivity has been advantageous to quantitatively study asphaltene
 89 destabilization [34, 35] and deposition in flow cells under atmospheric pressure[36, 37, 38, 39]. Moreover,
 90 it has been shown that QCM-D can be used in full immersion in live crude oil as a resonator (QCR) to
 91 probe fluid phase transitions and asphaltene destabilization [33, 40]. This work involves the use of the same
 92 experimental methods to define phase diagrams including asphaltene instability envelope for 2 recombined
 93 crude oils. The results obtained from high pressure measurements were then compared to predictions of the
 94 revisited version of ASIST using an immersed quartz sensor during n-alkane liquid titrations.

95 2. Materials & Method

96 2.1. Experimental revisit of ASIST

97 All liquid titrants (e.g., toluene and n-alkanes) were 98+% purity supplied by Sigma Aldrich. The 2 dead
 98 crude oil samples later referred to as crude oil A and crude oil B were provided by TOTAL from 2 different
 99 fields located offshore West Africa, both were free from contamination of water, solid particles, production
 100 additives or drilling fluid, their STO properties are summarized in Table 1.

101 **Table 1:** STO properties of crude A and B

Crude oil	Density	Viscosity	Refractive Index	Asphaltenes content (by IP143)
-	kg/m^3	$mPa.s$	-	wt%
crude A	898	64	1.5148	11
crude B	861.4	15.2	1.4902	4.5

102
 103
 104 All the following solubility parameters of liquid solutions were calculated by the linear proportionality to
 105 the Lorentz-Lorenz refractive index function proposed by Wang and Buckley [12]. In the following sections,
 106 equation 2 is applied to temperatures ranging from 20 to 120 °C ; we make the assumption that the predom-
 107 inant parameter affected by the temperature is the refractive index of the liquid solutions. According to van
 108 Laar- Lorenz definition of enthalpy, the other parameter affected by the temperature is the size of molecules.
 109 By using equation 2 at multiple temperatures we make the assumption that the distance between molecules
 110 (captured by the refractive index) is predominantly affecting the cohesive energy density and the variation in
 111 size of molecules is neglected.

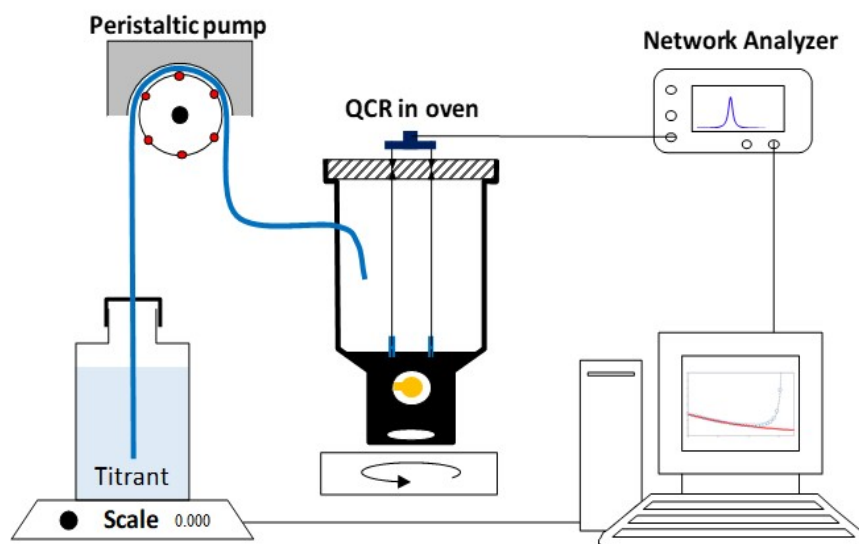
$$\delta = 52.042 \frac{(n_D^2 - 1)}{(n_D^2 + 2)} + 2.904 \quad (2)$$

112 The refractive indices (n_D) of crude oils and mixtures were measured at the wavelength of yellow Sodium
 113 D line (589.3nm) using a benchtop refractometer with a temperature control. The mentioned wavelength is
 114 usually chosen because it corresponds to the peak of absorption of organic matter and to a high enough
 115 frequency ($\sim 10^{15}Hz$) to assume that the subsequent total polarizability of the molecules is well represented
 116 by the electronic polarizability [41] (see equation 4 in the following section).

117 *QCR Atmospheric pressure titration experiments.*

118 The use of a 10nm sensitive instrument was decided to liberate ourselves from potential experimental and
 119 interpretation uncertainties thanks to the use of the same exact sensor during liquid titrations and during
 120 recombined oil depressurizations. The detection is instantaneous and is monitored in live while the titrations
 121 or depressurizations were carried out. As introduced, it was also meant to extend ASIST validations to a
 122 much finer detection with regards to small particles contribution to deposition. Holder mounted 3MHz AT-
 123 cut quartz crystal resonators are used through the entire results presented, each sensor was bought with

124 surfaces polished and electrodes made of Gold-Titanium. Resonance frequency and dissipation shifts (Δf
125 & $\Delta\Gamma$) are both calculated by subtracting the monitored signal to a reference signal; this reference was
126 calibrated prior to each experiment in the air for consistency to account for a potential aging of sensors.
127 During the first step of each experiment, as illustrated by Figure 1, the quartz crystal sensor was fully
128 immersed into a crude oil in a sealed custom glass vessel at atmospheric pressure. The mass of crude oil
129 for each experiment was kept constant. The temperature was controlled in an oven ($\pm 0.5^\circ C$) and monitored
130 by an inserted thermocouple probe. Titration experiments were achieved by pumping a pure titrant into the
131 vessel at controlled flow rate while the mixture was continuously stirred at rates enabling an instantaneous
132 mixing, ensured by visual inspection of the presence of a vortex. The volumetric rate of addition of the
133 solvent was steadily increased in order to keep a nearly linear rate of change in the volume fraction of the
134 titrant.



135
136 **Figure 1:** Atmospheric pressure experimental titration QCR set-up

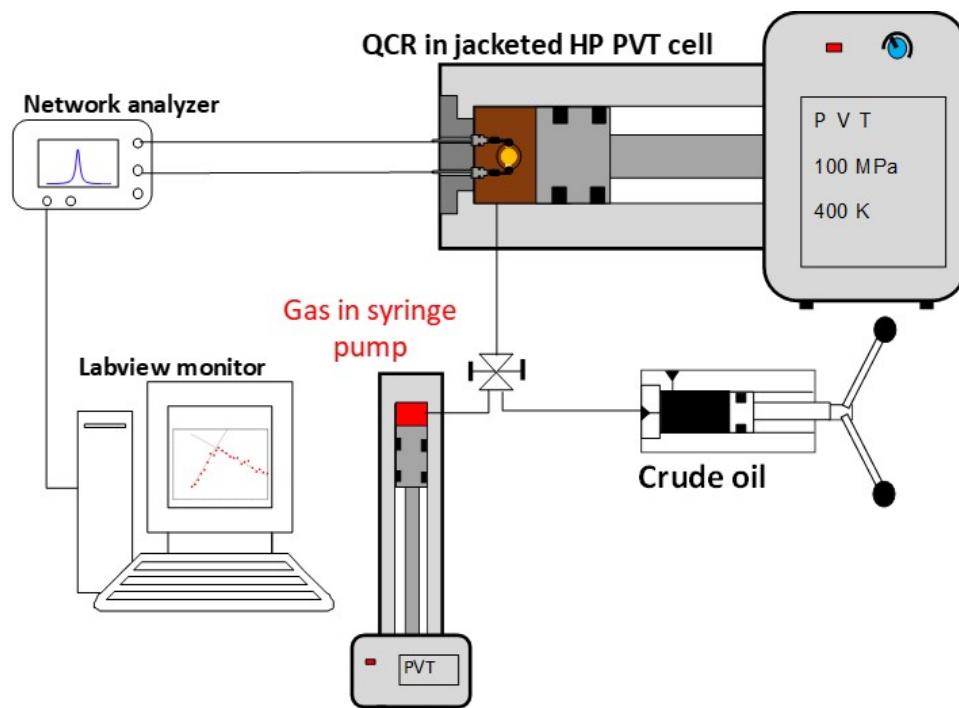
137 While titrating the oil the electrical signal transmitted to the sensor electrodes was controlled and re-
138 sponse was monitored through a network analyzer Agilent E5071C connected by coaxial cables to a monitor
139 with LabVIEW internally automated program. The record of quartz frequency and dissipation signals was
140 achieved by scans of several odd number of overtones (usually 1, 3, 5 and 7) from the first immersion and
141 titration was started once the noise signal had reached a stable plateau ($\pm 10Hz$). As already well established
142 by Cassiède et al. [42][43] and by Daridon et al. [35] the immersed AT-cut quartz crystals frequency of
143 resonance is sensitive to the pressure, to the mass loaded on its surface (deposit of unstable asphaltenes) and
144 to the contact liquid density viscosity product within a penetration depth ($\sim 100 - 300nm$) at each overtone
145 of the nominal frequency. The dissipation of the signal is mainly affected by the ability of the surrounding
146 media to absorb the shear waves induced by the resonance. Assuming a rigid deposit on the surface of the
147 sensor, the dissipation is mainly affected by the density viscosity product of the mixture. The shift in reso-
148 nance frequency signal was therefore chosen in order to interpret conditions of first appearance of deposited
149 unstable asphaltenes. Treatment of the recorded raw data were performed manually and coherency of de-
150 tection was further guaranteed by analyzing data of each overtone recorded. Details and validations of the
151 detection technique used in this study are further described elsewhere [35].

152

153 2.2. Pressurized systems experimental methods

154 *Sample preparation.*

155 The high pressure experimental set up is schemed on Figure 2. It consists of a stainless steel custom jacketed
156 and continuously stirred PVT cell with an integrated quartz crystal resonator. The equipment can undergo
157 pressures measured through a transducer up to 1000bar and temperatures measured by an inserted thermo-
158 couple probe from ambient to 120°C. Pressure, volume and temperature calibration of the equipment was
159 achieved using pure toluene prior to each experiment. Referring to previous work of Cassiède et al. [42][43]
160 and Daridon et al. [44][45], this calibration was not only necessary to know the exact parameters of the
161 piston cell body but also to account for the hydrostatic pressure effect on the immersed quartz sensor later
162 referred as Δf_p .



163

164 **Figure 2:** High pressure QCR experimental set-up

165 As opposed to injections of precipitant at constant pressure [8, 46], recombined oil experiments were
166 based on the volume expansion of a constant molar composition by isothermal depressurization of a gas
167 dissolved fluid. For such experiments, it is crucial to reduce experimental uncertainties; one of them being
168 the control of volumes and compositions. The parameters were precisely fixed by keeping simple gas com-
169 positions with well documented thermodynamic properties of the pure methane or of the binary mixture of
170 methane and carbon dioxide. A controlled mass of crude oil was first loaded to the PVT cell containing
171 the sensor followed by isobar and isotherm transfer of the desired volume of the gas used as a precipitant,
172 initially calculated using REFPROP software. Gases were supplied by Linde France SA, one bottle of pure
173 methane and one containing a blend of methane and carbon dioxide (80:20) molar composition.

174 *Pressure scanning.*

175 Daridon and Carrier [33], Cardoso et al. [45] proposed a phase transition detection protocol along with a
 176 detailed treatment of recorded data to calculate the thickness of the deposited layer. These authors showed
 177 that the appearance of unstable asphaltenes and the growth of a deposit layer on the quartz crystal's surface
 178 was well correlated to the deviation from linearity of the resonance frequency shift, Δf , during pressure
 179 depletion experiments. Consequently, in our interpretation the frequency shift is used as a criterion to sim-
 180 plify data treatment since the detection point was the only point of interest to this study. The pressure was
 181 decreased at a rate of 3 bars per minute, a realistic depressurization rate that a fluid would experience in an
 182 oil & gas production well.

183

184 *High pressure experiments.*

185 High pressure experiments were performed at conditions summarized in Table 2 and the solubility parameter
 186 of the recombined oil solution mixtures were calculated as described in the prediction methods section. For
 187 clarity, we will only discuss results for one set of conditions investigated for the same crude oil (e. i. crude oil
 188 A recombined with 57.4%mol of CH_4 at 80°C). Results at other conditions are discussed and summarized
 189 later in subsection 4.2 and in the appendix.

Table 2: Summary of the high pressure experimental range of conditions

Crude oil	Precipitant	Temperature °C	Dissolved precipitant concentration mol%
-	-	-	-
crude A	CH_4	45 - 80	8 - 60
crude B	$CH_4:CO_2(80:20)\%mol$	45 - 120	55 - 60

191 3. Theory & Calculation

192 3.1. Calculation of the refractive index of recombined oils

193 The refractive indices (n_D) at elevated pressures with gas dissolved (single liquid phase) were calculated
 194 as proposed by Buckley et al. [11], assuming a volume fraction ϕ_i averaged Lorenz-Lorentz function of 2
 195 pseudo-components in the mixture: Stock Tank Oil (STO) and its associated gas dissolved at the respective
 196 conditions of study.

$$\left(\frac{n_D^2 - 1}{n_D^2 + 2}\right)_{solution}^{P,T} = \phi_{oil} \left(\frac{n_D^2 - 1}{n_D^2 + 2}\right)_{STO}^{P,T} + \phi_{gas} \left(\frac{n_D^2 - 1}{n_D^2 + 2}\right)_{gas}^{P,T} \quad (3)$$

197 The use of Clausius-Mossotti and Lorenz-Lorentz equations lets us write the equation 4 [41].

$$\left(\frac{n_D^2 - 1}{n_D^2 + 2}\right)_i^{P,T} = \frac{\alpha_i \rho_i^{P,T} N_0}{3M_i \epsilon_0} \quad (4)$$

198 Where $\rho_i^{P,T}$ is density of component i at the pressure and temperature of test, α_i and M_i are respectively
 199 the electronic polarizability and the molecular weight of the component i, ϵ_0 is the permittivity of free space
 200 and N_0 is the Avogadro Number. By using a convenient reference condition of pressure and temperature,
 201 e. g. ambient conditions thereafter denoted by the exponent 0, Lorenz-Lorentz equation 3 combined to
 202 equation 4 leads us to the following relation which allows to calculate the refractive index function of the
 203 mixture at each pressure and temperature.

$$\left(\frac{n_D^2 - 1}{n_D^2 + 2}\right)_{solution}^{P,T} = \frac{V_{STO}^0}{V_{solution}^{P,T}} \left(\frac{n_D^2 - 1}{n_D^2 + 2}\right)_{STO}^0 + \frac{N_{gas}}{V_{solution}^{P,T}} Rm_{gas}^0 \quad (5)$$

204 Where V_{STO}^0 and $V_{solution}^{P,T}$ are respective volumes of oil at reference conditions and volume of the solution
 205 mixture at the pressure and temperature of test, N_{gas} and Rm_{gas}^0 are respectively the molar quantity and the
 206 molar refractivity of the gas dissolved.

207 We should note that this relation assumes molecular polarizabilities of components to be independent of
 208 pressure and temperature in our range of conditions (1 to 1000 *bar* and 20 to 120°C). Indeed, the density
 209 of the mixture is the only pressure and temperature dependent parameter in the equation 4. In our case the
 210 density is a measured value for single liquid phase conditions, i. e. at pressures greater than the saturation
 211 pressure. One can note that at under saturated pressures, the apparition of a gas phase no longer enables us
 212 to precisely measure the volume of the liquid phase and the molar quantity of gas dissolved accurately. This
 213 reinforces the choice of constant mass expansion instead of compression or constant pressure experiments
 214 to destabilize the asphaltenes.

215 3.2. Calculation of the molar volume of gas dissolved

216 Linear extrapolation of ASIST method was empirically found on the plot of solubility parameter at
 217 detection against molar volumes of the precipitants. Although mixtures of crude oil and liquid solvents
 218 are certainly not completely ideal mixtures, it is still a fair approximation to say that the molar volumes of
 219 the pure n-heptane (later referred as n-C7), n-undecane (n-C11) and n-pentadecane (n-C15) have the same
 220 values as when mixed to crude oil at atmospheric pressure. Therefore, for these liquid precipitants, their
 221 molar volumes were retrieved from the NIST database.

222 Unlike these components, $v_{gas}^{P,T}$ the molar volume of light compressible alkanes can be significantly differ-
 223 ent when dissolved in crude oil at elevated pressures compared to their pure properties reported in databases.
 224 Therefore, we back-calculated these values from experimental measurement. An internally TOTAL devel-
 225 oped PVT simulator was used to calculate the density of the crude oil at pressures and temperatures of tests.
 226 We tuned binary parameters of the Peng-Robinson cubic equations of state (P-R EoS) in order to match
 227 density measurements from an Anton Paar DMA 4500 M density meter at several temperatures. Simula-
 228 tions were then performed to replicate results obtained during tests on drill stem tests (DST) downhole fluid
 229 samples to check the reliability of the parameters tuned. PVT simulations were then computed to generate
 230 $\rho_{STO}^{P,T}$, the density of the degassed crude oil at pressures and temperature conditions of tests.

231 In our case, we recombined crude oils with well-known molar compositions, again in the following
 232 calculations (equation 6) we consider two pseudo-components system composed of crude oil and the asso-
 233 ciated gas. The molar volume of the gas dissolved at specific pressure and temperature condition can then
 234 be expressed as follows.

$$v_{gas}^{P,T} = \frac{V_{gas}^{P,T}}{N_{gas}} = \frac{V_{solution}^{P,T} - \frac{m_{STO}}{\rho_{STO}^{P,T}}}{N_{gas}} \quad (6)$$

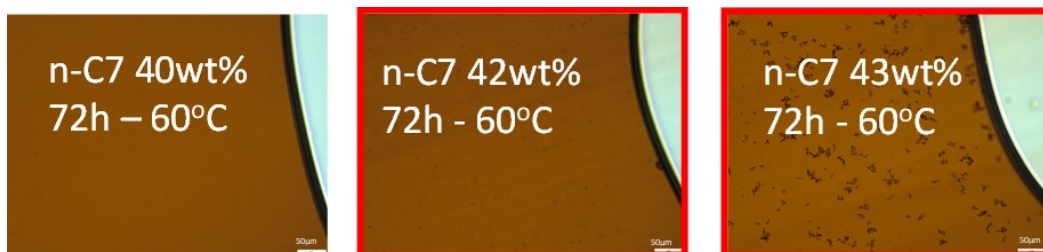
235 We should note that $V_{solution}^{P,T}$, N_{gas} and m_{STO} the mass of crude oil are all measured parameters in our
 236 study.

237 4. Results & Discussions

238 *Conventional ASIST.*

239 This empirical method first consists of preparing mixtures of crude oil and liquid n-alkane usually referred

240 as precipitants. The variable parameters are the temperature, the nature of the precipitant (chain length of
 241 n- alkane) and its concentration. Passed arbitrary aging times, the mixtures are inspected by microscopy to
 242 determine the lowest concentration of precipitant at which particles can be observed as illustrated in Figure 3.
 243 In our case, the aging time of 72h was chosen only for practical reasons on the microscopy detections. These
 244 results are later compared to the revisited ASIST by immersed quartz detection of unstable asphaltenes.



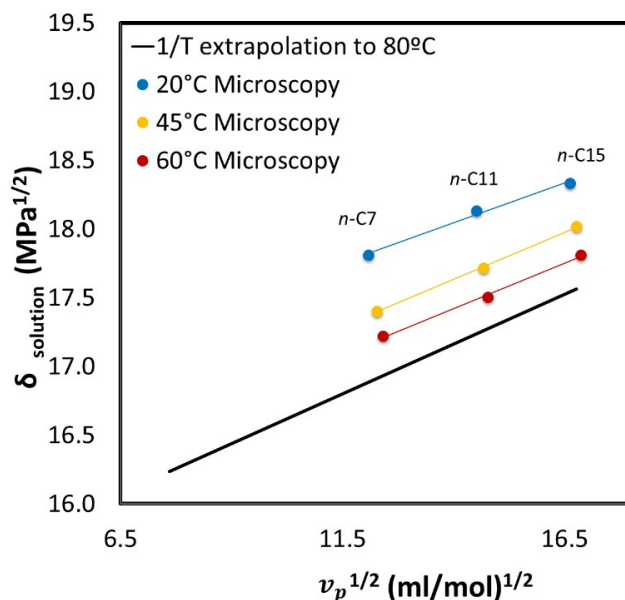
245

246

Figure 3: Illustration of microscopy detection of unstable asphaltenes - crude oil A

247 The solution with the lowest concentrations of precipitant showing flocs of asphaltenes were reproduced
 248 to measure refractive indices and the corresponding solubility parameter are computed using equation 2.
 249 For each temperature, the solubility parameters of threshold solutions are plotted against $v_p^{1/2}$, the square
 250 root of precipitant's molar volume (see Figure 4). Upon addition of precipitants at a respective temperature,
 251 $\delta_{solution}$ decreases; a solution having a higher solubility parameter compared to the instability linear threshold
 252 (or trend) is defined as stable and ones having a smaller solubility parameter contain unstable asphaltenes
 253 because they correspond to mixtures with higher concentrations of the precipitant.

254 The detection threshold can then be extrapolated to both parameters of interest; the temperature and lower
 255 molar volumes of the precipitant (e. g. gas dissolved). Linear extrapolation of the inverse of temperature
 256 was used because it was observed to be in fair agreement with experimental data in the range investigated.
 257 Results of the applied conventional ASIST method are plotted in Figure 4.

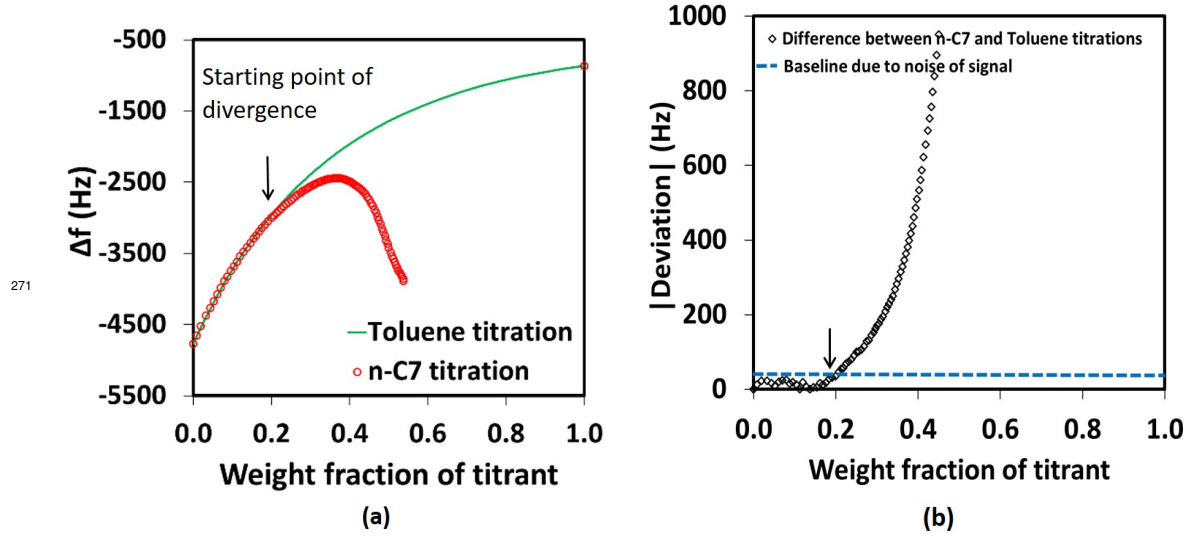


258

259 **Figure 4:** ASIST applied to crude oil A by microscopy detection of unstable asphaltenes(aging time of 72h)

260 The experimental concept of an optimal aging time (here arbitrarily chosen to be 72h) discussed by
 261 several authors brings questionable uncertainties to the method as discussed in the introduction[17, 18, 19].
 262 Indeed, the kinetics of asphaltenes' aggregation depends on the precipitant causing the instability [23], how-
 263 ever the definition of the instability trends of ASIST suggests that the aging time should be kept constant in
 264 the empirical inputs along the entire trend. Therefore, the extrapolation of thresholds to lower molar volume
 265 of precipitant should involve similar aging time to any precipitant. In the present study, we suspected that
 266 the necessary tuning of the aging time could be due to the difference sensitivity of instrumentations used to
 267 probe the presence of flocs in solutions at atmospheric pressure and at gas dissolved pressurized conditions.

268 *QCR atmospheric titration results.* Crude oil A was titrated by a good solvent, toluene, and a precipitant,
 269 n-C7 at the exact same conditions of stirring, injection and temperature in 2 distinct experiments. The Figure
 270 5(a) shows the results of the frequency shift evolution for both experiments.



272 **Figure 5:** (a) Δf signal recorded for titration experiments of crude oil B at $20^\circ C$ – (b) difference between Δf signals of
 n-C7 and toluene titrations

273 By considering that no mass was accumulated on the surfaces during the titration of a good solvent,
 274 according to equations established by Kanazawa and Gordon [47], the change in frequency shift for a quartz
 275 oscillator is proportional to the change of the density viscosity product of the contacting liquid for homoge-
 276 neous systems with no adsorption or deposition later referred as mass loading effects:

$$\Delta f_{n,\mu} = -\sqrt{n} \frac{C_m}{\sqrt{\pi f_0}} \sqrt{\rho_{liq} \mu_{liq}} \quad (7)$$

$$C_m = \frac{2f_0^2}{\sqrt{\rho_q \mu_q}} \quad (8)$$

277 Where C_m is the Sauerbrey constant defined by equation 8, f_0 is the fundamental resonance frequency of
 278 the quartz crystal used, ρ_q and μ_q are the density and the shear modulus of the quartz material and n is the
 279 overtone investigated [48].

280 Therefore, thanks to the similar density viscosity products of toluene and n-C7, toluene titration results
 281 provides insights of the dilution effect of n-C7. In other words, if n- C7 was a good solvent to all the
 282 compounds of the oil, its titration would result in the same data as the ones of toluene. As illustrated by
 283 the data of Figure 5(a) and 5(b), for precipitants the measured resonance frequency diverges from dilution
 284 effect upon increase of their fraction. At the interpreted detection of a divergence, the change in $\sqrt{\rho_{liq}\mu_{liq}}$ of
 285 the surrounding media and the surface mass load are two competing effects. Indeed the response is first due
 286 to the presence of particles in the vicinity of the sensor; it could be interpreted by a change in viscosity and
 287 density of the surrounding media. When the titration of precipitant is carried further, the change in slope
 288 with an opposite sign effect indicates accumulation of multiple layers deposit on the surface of the sensor as
 289 opposed to a monolayer adsorption. This effect becomes predominant compared to the presence of particles
 290 in the surrounding bulk and the dilution effects. The behavior is explained by a combination of Sauerbrey
 291 and Kanazawa equations resulting in equation 9; indeed an increase of the theoretical mass deposited per unit
 292 of area $\rho_{deposit}h_{deposit}$ makes the shift in frequency absolute value decrease while the reduction of $\sqrt{\rho_{liq}\mu_{liq}}$
 293 due to dilution induces an increase of the absolute value $|\Delta f|$ [47, 48].

$$\Delta f_{n,load} = -n2C_m\rho_{deposit}h_{deposit} - \sqrt{n}\frac{C_m}{\sqrt{\pi f_0}}\sqrt{\rho_{liq}\mu_{liq}} \quad (9)$$

294 The difference between both signals recorded; titration by toluene and titration by n-C7, was then com-
 295 puted in order to locate the lowest heptane concentration at which the signal significantly deviates from a
 296 dilution curve due to the presence of unstable asphaltenes to the vicinity of the sensor before deposit ac-
 297 cumulates. The Figure 5b illustrates how the detection was geometrically determined by intersection on a
 298 narrowed region of weight fraction at the departure of the curve. After interpretation, the mixtures at which
 299 divergence was detected were replicated in order to measure their respective refractive indices. Results are
 300 summarized in Table 3 in the appendix .

301 The effect of the aging time of solutions was studied using the methodology described by Maqbool et al.
 302 [20]. Figure 6 shows how the detection time curve of unstable asphaltenes by microscopy compares to the
 303 instantaneous detection of the QCR signal while titrating. At the QCR detection conditions, we would have
 304 to age the solution $10^5 h$ in order to detect micrometer sized particles.

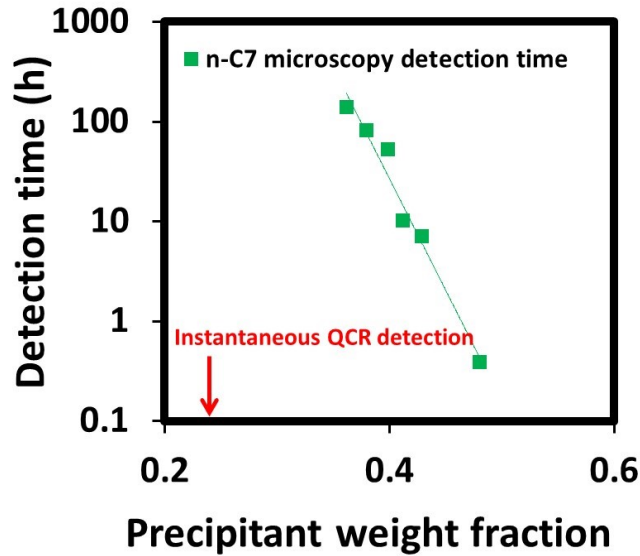
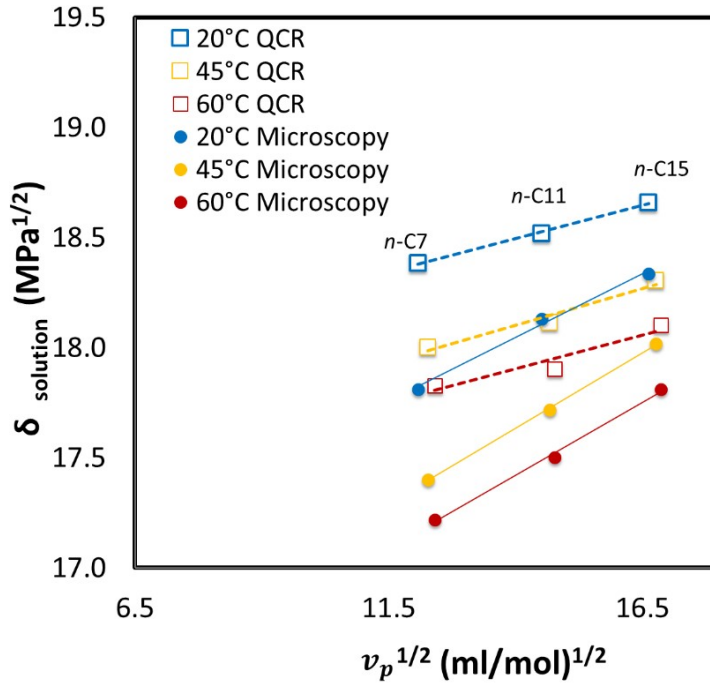


Figure 6: Comparison the detection time curve by microscopy and QCR instantaneous detection of unstable asphaltenes for crude oil A at 60°C

In a similar manner to microscopy data showed in Figure 4 and for comparison to microscopy detected trends, Figure 7 incorporates the data for crude oil A using QCR first detection of unstable asphaltenes in the vicinity of the sensor. The extrapolation to higher temperature was again achieved by an inversely linear proportionality and the range of investigated temperatures was verified to be a good approximation of a linearized inverse dependence.

As expected, larger concentrations of precipitants are required to observe flocculation through an optical microscope even aged for 72h, this is put in perspective by inferior solubility parameters required by microscopy compared to QCR detection of unstable asphaltenes at a given temperature with a given n-alkane. We can note that the results do not only show a change in the sensitivity (magnitude of $\delta_{solution}$) but also a significant change in slope of the trends against $v_p^{1/2}$. In other words, for this particular case it seems that the titration of a system containing nanometer sized unstable asphaltenes with different precipitants requires different solubility parameter change to reach micro-aggregates. This result suggests that the nature of the precipitant has an effect on the kinetics of aggregation between nanometer and micrometer sized unstable asphaltenes. The experimental data are summarized in Table 4 of the appendix.



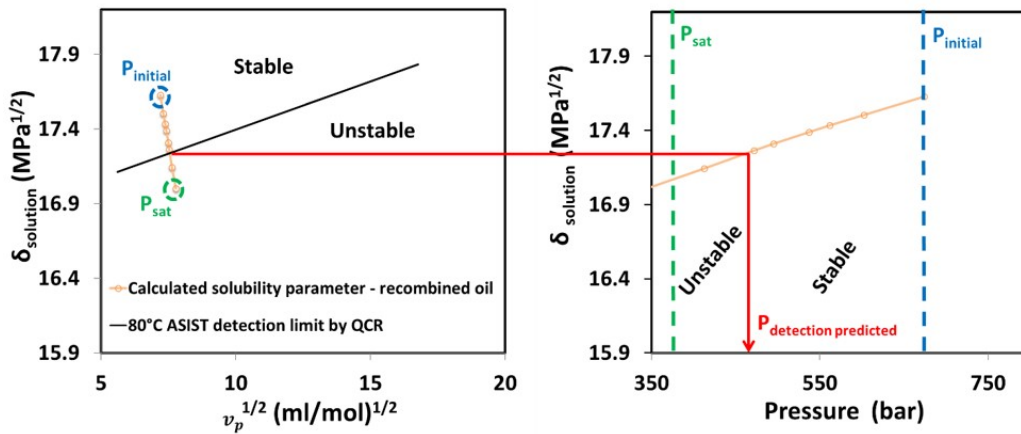
321

Figure 7: Comparison of ASIST using microscopy detections of unstable asphaltenes and using QCR detection of first deposit for crude oil A

322

323 *4.1. Predictions of ASIST by QCR*

324 In Figure 8, the graphical representations are showed in parallel to illustrate the data process to predict
 325 the pressure at which detection should be expected.



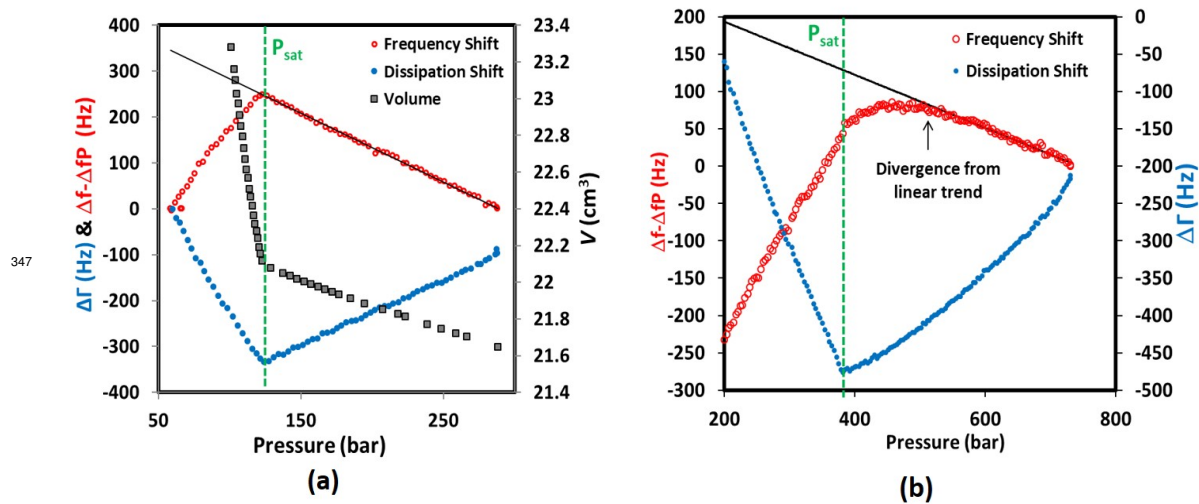
326

327 **Figure 8:** Prediction of the deposit detection pressure for Crude oil A recombined at 80°C with 57.4%mol of CH₄

328 For the condition showed, the ASIST predicted pressure of appearance of unstable asphaltenes to the
 329 vicinity of the sensor is 465bar, the experimentally observed detection is provided in the following section
 330 along with results of different conditions.

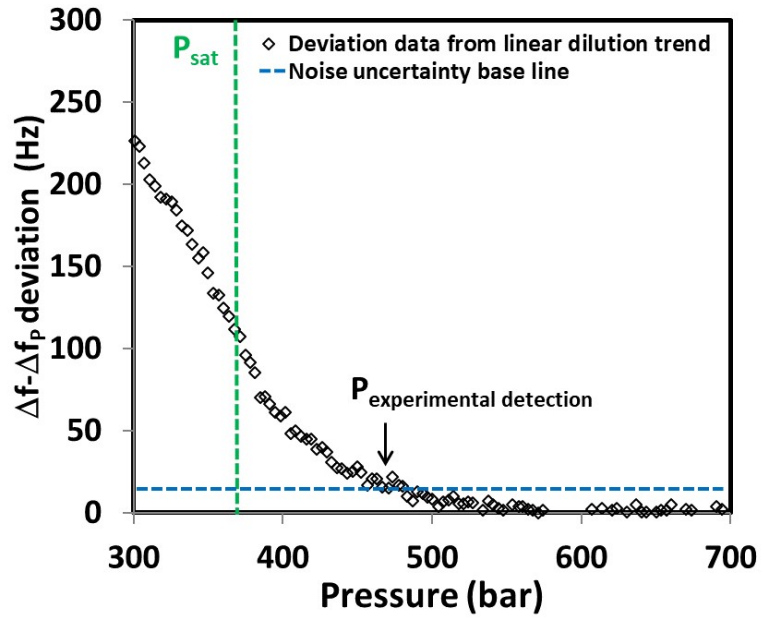
331 4.2. Constant Mass Expansions experimental results

332 Constant mass expansion (CME) scanning results are showed in Figure 9 for two different compositions
 333 in gas content, both experiments were achieved using crude oil A at a controlled temperature of 80°C.
 334 Figure 9(a) shows results for a system containing 29.95%mol of methane, frequency and dissipation shifts
 335 show as expected a sharp change in slope due to the appearance of the gas phase at the saturation pressure,
 336 confirmed by the behavior of volume of the system. Both frequency and dissipation shifts experience a
 337 monotonous linear trend before reaching the bubble pressure. Analogically to the atmospheric pressure
 338 titrations, this is due to a dilution effect of the expanding compressible components (mainly methane in
 339 this case), which induces a reduction of the $\sqrt{\rho_{liq}\mu_{liq}}$. In this case, the system passes the pressure scanning
 340 without destabilizing asphaltenes. Unlike to the system containing 29.95%mol, the recombined crude oil
 341 with 57.4%mol of methane experienced a diverging frequency signal from linearity at pressures above the
 342 saturation pressure (which can still clearly be determined by the sharp changes of dissipation or volume
 343 not showed in Figure 9(b) for clarity). Such behavior of the resonance frequency was demonstrated to be a
 344 signature of the appearance of unstable asphaltene by other authors[33]. We should note that the explained
 345 combination of Sauerbrey and Kanazawa in equation 9 for atmospheric pressure titrations also applies to
 346 pressure depletion [33, 42, 43].



347
 348 **Figure 9:** Depressurization experimental data recorded for crude oil A at 80 °C with CH₄ composition of (a) 29.95%mol and (b) 57.4%mol

349 The difference between the actual recorded data and the averaged linear behavior of high enough pressure
 350 data is then graphically analyzed in Figure 10 to observe the deviation and determine the experimental
 351 pressure of detection.



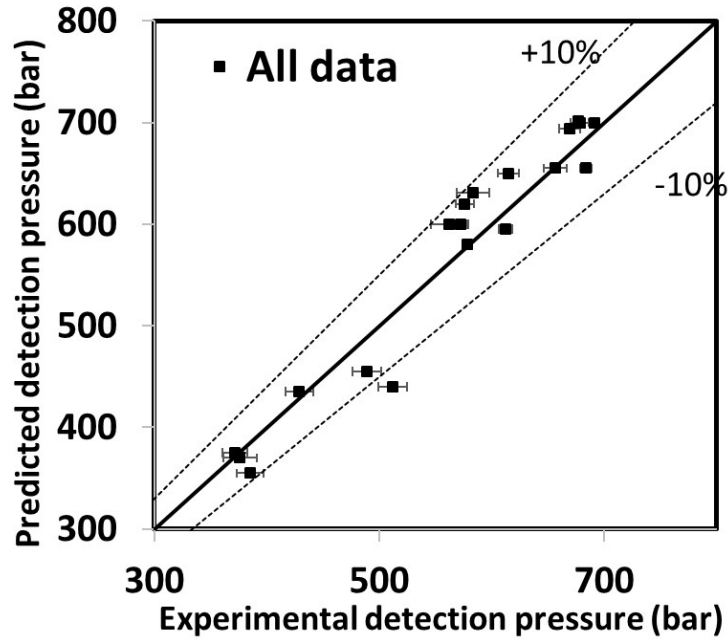
352

Figure 10: Deviation from linear signal due to asphaltenes instability during expansion of CH_4 (57.4%mol) for crude oil A at 80°C

353

354 4.3. Comparison of predictions to experimental observations

355 In this section, the overall experimental data collected are compared to predictions using the modified
 356 ASIST method. All the conditions are summarized by Table 3 of the Appendix.



357

358

Figure 11: Revisited ASIST predictions vs experimental pressure of first deposit detected

359

360

361

362

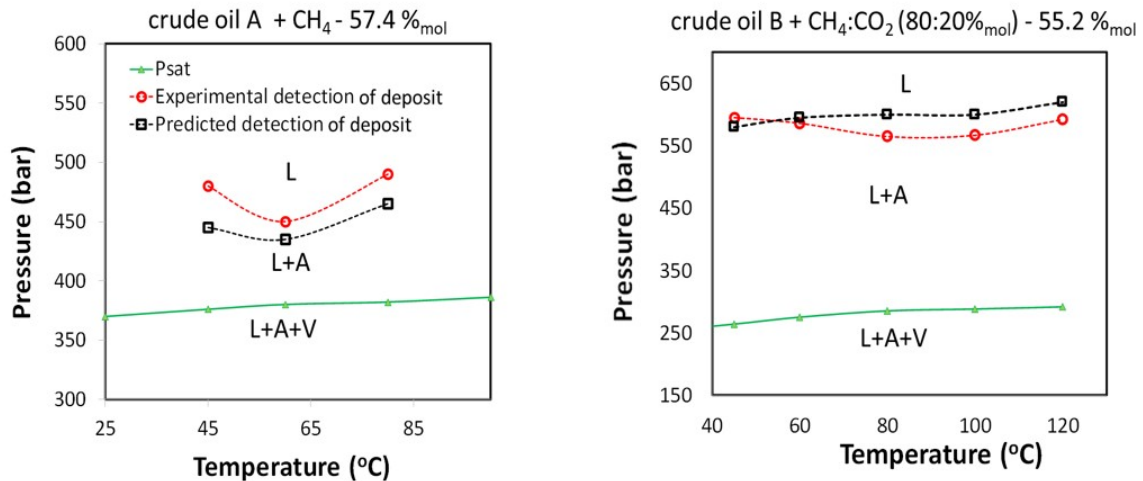
363

364

365

As it can be noticed in the Figure 11, despite assumptions, potential experimental errors and theoretically unclear empirical trends, results of predicted pressures fall within $\sim 10\%$ of error compared to experimental observations. One should note additional validation results are absent from Figure 11; no detections were predicted neither observed for mixtures containing less than $55\%_{mol}$ of gas dissolved. The pressure versus temperature phase diagrams on Figure 12 show a different aspect of ASIST predictions; despite their slight offset, the curvatures follow the same trend as experimental observations, which confirms that the capture of underlying physics on the effects of temperature and pressure.

366



367

368 **Figure 12:** Pressure vs Temperature diagrams for recombined crude oils A and B. L stands for liquid, V for vapor and
A for unstable asphaltenes

369 The conventional ASIST by microscopy and an aging time of 72h was applied to the same conditions,
370 however, no destabilization was predicted to the depressurization experiments. This was an expected result
371 due to the difference of resolution of the techniques.

372 5. Conclusions

373 In this study, ASIST predictive concepts were successfully implemented and revisited by using data gath-
374 ered from a novel and highly sensitive method of detection of the first instantaneously unstable asphaltenes:
375 a fully immersed Quartz Crystal Resonator (QCR). Destabilization was first caused by the volume addition
376 of several n-alkanes in order to extrapolate to light alkanes and carbon dioxide dissolved conditions. Results
377 of the predicted pressure at which instability would be detected using the same experimental apparatus were
378 in excellent agreement to the actual observed detected pressures, evidencing that the tuning of the aging
379 time is certainly an artificial matching of detections using different apparatus. This is even reinforced by the
380 observation of a change in slope of trends depending on the size of particles, suggesting different kinetics of
381 aggregation depending on the nature of the precipitant. In fact, using several apparatus to validate the method
382 would add a variable. This founding do not only bring an enhanced understanding of the brutal asphaltenes
383 destabilization during the production and transport of oil & gas, but the modified practice proposed in this
384 publication provides prediction of different conditions at which the unstable asphaltenes are small particles
385 that could not be detected instantaneously by microscopy but still could contribute significantly to a diffusive
386 deposition mechanism.

387
388 Conveniently and controversially, this study shows that principles of ASIST can be applied and adapted
389 to more or less sensitive detection methods and can give significantly different results, some of the detection
390 techniques may be too conservative to the prediction of industrial risks while others may not be enough
391 sensitive. It is also important to note that slight uncertainties in the measurements of compositions, refractive
392 indices or in calculations of $\delta_{solution}^{P,T}$ and $v_{gas}^{P,T}$ at gas dissolved conditions can induce major errors in the
393 prediction of the detection pressure. Indeed, the input parameters were randomly varied in order to observe
394 the numerical sensitivity of the outputs. Consequently, in this study all the measurements and data processing
395 were repeated 2 to 3 times ensuring the precision required by the method.

396 Although the empirical predictive tool showed excellent agreements with experimental observations, we
397 should note that the principles of ASIST bring remarkable theoretical questions on how it relates to the
398 extensively used Flory-Huggins regular solution theory and to other foundings. Despite the polydisperse
399 nature of asphaltenes, the least-stable ones have hypothetically (according to the regular solution theory)
400 a constant solubility parameter δ_{asph} for a given crude oil regardless of the precipitant used to destabilized
401 them. However, some authors recently showed how the polydispersity can improve the modeling of unstable
402 asphaltenes' properties [49, 50], agreeing with the extensive experimental research on the morphological
403 structure (e. g. size, fractal dimension) of asphaltenes during precipitation[51, 52, 53, 54] which is expected
404 to be of relevance to their transport during the deposition mechanism(s)[55, 56, 27]. In relation to the
405 aforementioned, the predictive quantification of the destabilization and deposition of asphaltenes in gas-
406 dissolved environment remains poorly documented. Due to the complexity of multiple variables and to the
407 simultaneous phenomena hapenning during the flow of a live-oil, the discussed aspects in this conclusion
408 have never been related to the asphaltenes destabilized during the expansion of light components. More
409 comparative efforts in this direction are necessary.

410 6. Acknowledgements

411 The authors thank TOTAL SA for the financial support and the permission to publish this study.

412 References

- 413 [1] J. S. Buckley, J. Wang, J. L. Creek, Solubility of the least-soluble asphaltenes, *Asphaltenes, Heavy*
414 *Oils, and Petroleomics* (2007) 401–437.
- 415 [2] I. A. Wiehe, Two-dimensional solubility parameter mapping of heavy oils, *Fuel Science and Technol-*
416 *ogy International* 14 (1996) 289–312.
- 417 [3] N. Passade-Boupat, H. Zhou, M. Rondon-Gonzalez, Asphaltene Precipitation From Crude Oils : How
418 To Predict It And To Anticipate Treatment?, *SPE Middle East Oil and Gas Show and Conference*
419 (2010).
- 420 [4] D. Powers, H. Sadeghi, H. Yarranton, F. Van Den Berg, Regular solution based approach to modeling
421 asphaltene precipitation from native and reacted oils: Part 1, molecular weight, density, and solubility
422 parameter distributions of asphaltenes, *Fuel* 178 (2016) 218–233.
- 423 [5] H. Yarranton, D. Powers, J. Okafor, F. van den Berg, Regular solution based approach to modeling
424 asphaltene precipitation from native and reacted oils: Part 2, molecular weight, density, and solubility
425 parameter of saturates, aromatics, and resins, *Fuel* 215 (2018) 766–777.
- 426 [6] D. Barrera, D. Ortiz, H. Yarranton, Molecular weight and density distributions of asphaltenes from
427 crude oils, *Energy & fuels* 27 (2013) 2474–2487.
- 428 [7] J. Wang, J. L. Creek, J. S. Buckley, Screening for potential asphaltene problems, *SPE Annual Technical*
429 *Conference and Exhibition* (2006).
- 430 [8] K. Kraiwattanawong, H. S. Fogler, S. G. Gharfeh, P. Singh, W. H. Thomason, S. Chavadej, Thermo-
431 dynamic solubility models to predict asphaltene instability in live crude oils, *Energy & fuels* 21 (2007)
432 1248–1255.
- 433 [9] H. Alboudwarej, K. Akbarzadeh, J. Beck, W. Y. Svrcek, H. W. Yarranton, Regular solution model for
434 asphaltene precipitation from bitumens and solvents, *AIChE Journal* 49 (2003) 2948–2956.
- 435 [10] I. A. Wiehe, H. W. Yarranton, K. Akbarzadeh, P. M. Rahimi, A. Tecler, The paradox of asphalt-
436 ene precipitation with normal paraffins, *Energy & Fuels* 19 (2005) 1261–1267.
- 437 [11] J. Buckley, G. Hirasaki, Y. Liu, S. Von Drasek, J. Wang, B. Gill, Asphaltene precipitation and solvent
438 properties of crude oils, *Petroleum Science and Technology* 16 (1998) 251–285.
- 439 [12] J. Wang, J. Buckley, A two-component solubility model of the onset of asphaltene flocculation in crude
440 oils, *Energy & Fuels* 15 (2001) 1004–1012.
- 441 [13] J. Wang, J. S. Buckley, Asphaltene stability in crude oil and aromatic solvents the influence of oil
442 composition, *Energy & fuels* 17 (2003) 1445–1451.
- 443 [14] F. M. Vargas, W. G. Chapman, Application of the one-third rule in hydrocarbon and crude oil systems,
444 *Fluid Phase Equilibria* 290 (2010) 103–108.

- 445 [15] S. Subramanian, S. Simon, J. Sjoblom, Asphaltene precipitation models: A review, *Journal of Disper-*
446 *sion Science and Technology* (2015).
- 447 [16] I. A. Wiehe, Asphaltene solubility and fluid compatibility, *Energy & Fuels* 26 (2012) 4004–4016.
- 448 [17] J. L. Creek, J. Wang, J. S. Buckley, et al., Verification of asphaltene-instability-trend (asist) predictions
449 for low-molecular-weight alkanes, *SPE Production & Operations* 24 (2009) 360–368.
- 450 [18] D. L. Gonzalez, F. M. Vargas, E. Mahmoodaghdam, F. Lim, N. Joshi, Asphaltene stability prediction
451 based on dead oil properties: Experimental evaluation, *Energy & fuels* 26 (2012) 6218–6227.
- 452 [19] S. Dolati, H. Zarei, R. Kharrat, Asphaltene Instability Trends to Predict Asphaltene Precipitation Onset
453 Pressure: Constrained for Light and Heavy Crude Oils, *Journal of Dispersion Science and Technology*
454 36 (2015) 103–110.
- 455 [20] T. Maqbool, P. Srikiratiwong, H. S. Fogler, Effect of temperature on the precipitation kinetics of
456 asphaltenes, *Energy & Fuels* 25 (2011) 694–700.
- 457 [21] T. Maqbool, S. Raha, M. P. Hoepfner, H. S. Fogler, Modeling the aggregation of asphaltene nanoag-
458 gregates in crude oil- precipitant systems, *Energy & Fuels* 25 (2011) 1585–1596.
- 459 [22] C. V. B. Fávero, T. Maqbool, M. Hoepfner, N. Haji-Akbari, H. S. Fogler, Revisiting the flocculation
460 kinetics of destabilized asphaltenes, *Advances in colloid and interface science* 244 (2017) 267–280.
- 461 [23] N. Haji-Akbari, P. Masirisuk, M. P. Hoepfner, H. S. Fogler, A unified model for aggregation of as-
462 phaltenes, *Energy & Fuels* 27 (2013) 2497–2505.
- 463 [24] N. Haji-Akbari, P. Teeraphakul, A. T. Balgoa, H. S. Fogler, Effect of n-alkane precipitants on aggre-
464 gation kinetics of asphaltenes, *Energy & Fuels* 29 (2015) 2190–2196.
- 465 [25] M. P. Hoepfner, H. S. Fogler, Multiscale scattering investigations of asphaltene cluster breakup,
466 nanoaggregate dissociation, and molecular ordering, *Langmuir* 29 (2013) 15423–15432.
- 467 [26] M. P. Hoepfner, V. Limsakoune, V. Chuenmeechao, T. Maqbool, H. S. Fogler, A fundamental study of
468 asphaltene deposition, *Energy & fuels* 27 (2013) 725–735.
- 469 [27] C. Vilas Bôas Fávero, A. Hanpan, P. Phichphimok, K. Binabdullah, H. S. Fogler, Mechanistic investi-
470 gation of asphaltene deposition, *Energy & Fuels* 30 (2016) 8915–8921.
- 471 [28] M. Tavakkoli, M. R. Grimes, X. Liu, C. K. Garcia, S. C. Correa, Q. J. Cox, F. M. Vargas, Indirect
472 method: a novel technique for experimental determination of asphaltene precipitation, *Energy & Fuels*
473 29 (2015) 2890–2900.
- 474 [29] A. Khaleel, M. Abutaqiya, M. Tavakkoli, A. A. Melendez-Alvarez, F. M. Vargas, et al., On the pre-
475 diction, prevention and remediation of asphaltene deposition, in: *Abu Dhabi International Petroleum*
476 *Exhibition and Conference*, Society of Petroleum Engineers, 2015, pp. 1–23.
- 477 [30] M. Haghshenas, S. Balashanmugam, D. Gonzalez, M. Pietrobon, et al., Prediction of asphaltene onset
478 pressure from dead oil stability, in: *Offshore Technology Conference*, Offshore Technology Confer-
479 ence, 2016, pp. 1–20.

- 480 [31] A. Arnau, Y. Montagut, J. V. García, Y. Jiménez, A different point of view on the sensitivity of quartz
481 crystal microbalance sensors, *Measurement Science and Technology* 20 (2009) 124004.
- 482 [32] X. Huang, Q. Bai, J. Hu, D. Hou, A practical model of quartz crystal microbalance in actual applica-
483 tions, *Sensors* 17 (2017) 1785.
- 484 [33] J.-L. Daridon, H. Carrier, Measurement of phase changes in live crude oil using an acoustic wave
485 sensor: Asphaltene instability envelope, *Energy & Fuels* 31 (2017) 9255–9267.
- 486 [34] L. Goual, G. Horváth-Szabó, J. H. Masliyah, Z. Xu, Adsorption of bituminous components at oil/water
487 interfaces investigated by quartz crystal microbalance: Implications to the stability of water-in-oil
488 emulsions, *Langmuir* 21 (2005) 8278–8289.
- 489 [35] J. L. Daridon, M. Cassiede, D. Nasri, J. Pauly, H. Carrier, Probing asphaltene flocculation by a quartz
490 crystal resonator, *Energy & Fuels* 27 (2013) 4639–4647.
- 491 [36] S. Subramanian, S. Simon, B. Gao, J. Sjöblom, Asphaltene fractionation based on adsorption onto
492 calcium carbonate: Part 1: Characterization of sub-fractions and QCM-measurements, *Colloids and
493 Surfaces A: Physicochemical and Engineering Aspects* 495 (2016) 136–148.
- 494 [37] M. Tavakkoli, S. R. Panuganti, F. M. Vargas, V. Taghikhani, M. R. Pishvaie, W. G. Chapman, As-
495 phaltene deposition in different depositing environments: part 1. model oil, *Energy & fuels* 28 (2013)
496 1617–1628.
- 497 [38] M. Tavakkoli, S. R. Panuganti, V. Taghikhani, M. R. Pishvaie, W. G. Chapman, Asphaltene deposition
498 in different depositing environments: Part 2. real oil, *Energy & Fuels* 28 (2014) 3594–3603.
- 499 [39] S. Campen, B. Smith, J. S. Wong, Deposition of asphaltene from destabilized dispersions in heptane-
500 toluene, *Energy & Fuels* (2018).
- 501 [40] J. L. Daridon, E. Orlandi, H. Carrier, Measurement of bubble point pressure in crude oils using an
502 acoustic wave sensor, *Fluid Phase Equilibria* 427 (2016) 152–160.
- 503 [41] J. N. Israelachvili, *Intermolecular and Surface Forces*, third edition ed., Elsevier Academic Press, 2011.
- 504 [42] M. Cassiède, J.-L. Daridon, J. Paillol, J. Pauly, Impedance analysis for characterizing the influence
505 of hydrostatic pressure on piezoelectric quartz crystal sensors, *Journal of Applied Physics* 108 (2010)
506 034505.
- 507 [43] M. Cassiède, J.-L. Daridon, J. Paillol, J. Pauly, Characterization of the behaviour of a quartz crystal
508 resonator fully immersed in a newtonian liquid by impedance analysis, *Sensors and Actuators A:
509 Physical* 167 (2011) 317–326.
- 510 [44] J.-L. Daridon, M. Cassiède, J. Paillol, J. Pauly, Viscosity measurements of liquids under pressure by
511 using the quartz crystal resonators, *Review of Scientific Instruments* 82 (2011) 095114.
- 512 [45] F. Cardoso, H. Carrier, J.-L. Daridon, J. Pauly, P. Rosa, Co₂ and temperature effects on the asphaltene
513 phase envelope as determined by a quartz crystal resonator, *Energy & Fuels* 28 (2014) 6780–6787.
- 514 [46] A. Hirschberg, L. deJong, A. Schipper, J. Meijer, Influence of temperature and pressure on asphaltene
515 flocculation, *Society of Petroleum Engineers Journal* 24 (1984).

- 516 [47] K. K. Kanazawa, J. G. Gordon, Frequency of a quartz microbalance in contact with liquid, *Analytical*
517 *Chemistry* 57 (1985) 1770–1771.
- 518 [48] G. Sauerbrey, Verwendung von Schwingquarzen zur Wagungdiinner Schichten und zur Mikrowagung,
519 *Zeitschrift fur Physik* 155 (1959) 206–222.
- 520 [49] G. Javanbakht, M. Sedghi, W. R. Welch, L. Goual, M. P. Hoepfner, Molecular polydispersity improves
521 prediction of asphaltene aggregation, *Journal of Molecular Liquids* (2018).
- 522 [50] M. Tavakkoli, S. R. Panuganti, V. Taghikhani, M. R. Pishvaie, W. G. Chapman, Understanding the
523 polydisperse behavior of asphaltenes during precipitation, *Fuel* 117 (2014) 206–217.
- 524 [51] L. Barré, S. Simon, T. Palermo, Solution properties of asphaltenes, *Langmuir* 24 (2008) 3709–3717.
- 525 [52] J. Eyssautier, P. Levitz, D. Espinat, J. Jestin, J. Gummel, I. Grillo, L. Barré, Insight into asphaltene
526 nanoaggregate structure inferred by small angle neutron and x-ray scattering, *The Journal of Physical*
527 *Chemistry B* 115 (2011) 6827–6837.
- 528 [53] J. Eyssautier, D. Frot, L. Barré, Structure and dynamic properties of colloidal asphaltene aggregates,
529 *Langmuir* 28 (2012).
- 530 [54] Y. Yang, W. Chaisoontornyotin, M. P. Hoepfner, The Structure of Asphaltenes During Precipitation
531 Investigated by Ultra-Small-Angle X-ray Scattering, *Langmuir* (2018) [acs.langmuir.8b01873](https://doi.org/10.1021/acs.langmuir.8b01873).
- 532 [55] W. Chaisoontornyotin, N. Haji-Akbari, H. S. Fogler, M. P. Hoepfner, Combined Asphaltene Aggrega-
533 tion and Deposition Investigation, *Energy and Fuels* 30 (2016) 1979–1986.
- 534 [56] Q. Guan, Y. Yap, A. Goharzadeh, J. Chai, F. Vargas, W. Chapman, M. Zhang, Integrated one-
535 dimensional modeling of asphaltene deposition in wellbores/pipelines, in: *Modeling, Simulation, and*
536 *Applied Optimization (ICMSAO)*, 2017 7th International Conference on, IEEE, 2017, pp. 1–6.

Table 3: : Experimental QCR detection results for atmospheric titration of liquid precipitants

Crude oil	Precipitant	Temperature °C	Precipitant concentration weight%	Refractive Index
-	-	-	-	-
		20	24	1.4846
	n-C7	45	23	1.4708
		60	24	1.4644
crude A	n-C11	20	26	1.4896
		45	26	1.4748
		60	26	1.4672
	n-C15	20	22	1.4948
		45	23	1.4818
		60	23	1.4744
		20	20	1.4710
	n-C7	45	20	1.4580
		60	20	1.4527
crude B	n-C11	20	24	1.4765
		45	24	1.4615
		60	24	1.4558
	n-C15	20	20	1.4793
		45	21	1.4677
		60	21	1.4633

Table 4: Experimental QCR gas dissolved detection results compared to predictions using revisited ASIST

Crude oil	Precipitant	Temperature oC	Dissolved precipitant concentration mol %	Experimental detection pressure bar	ASIST predicted pressure bar	
crude A	C1	45	8.27	No detection	No prediction	
			20.01			
			29.95			
			39.84			
			50.02			
			55.19			
		60	57.41	376	370	
			8.27	512	440	
			20.01	No detection	No prediction	
			29.95			
			39.84			
			50.02			
		55.19	371.5			375
		57.41	429			435
		80	8.27	No detection	No prediction	
			20.01			
			29.95			
			39.84			
50.02						
55.19	385		355			
crude B	C1:CO2(80:20) % mol	45	57.41	489	455	
			59.90	583.5	631	
			55.22	578.5	580	
		60	58.12	684	655	
			55.22	612.5	595	
			58.12	677	701	
		80	60.01	691	700	
			55.22	572.5	600	
			58.12	656.5	655	
			60.01	669.5	694	
			55.22	562.5	600	
			58.12	615	650	
		100	60.01	679	700	
			58.12	615	650	
			60.01	679	700	
		120	55.22	576	620	