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A three-dimensional Dionne model for multipactor simulations

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The multipactor phenomenon is characterized by a very fast growth of the electronic population in Radio-Frequency (RF) devices under vacuum. As this effect limits the transmissible RF power and can harm RF systems, it has been widely studied during the last decades. Due to the high cost of experimental tests, simulation tools are heavily used to predict the threshold of multipactor growth. However, their reliability is limited for complex configurations, e.g. when dielectrics or magnetic fields are present. A crucial element of these multipactor simulations is the secondary-emission model. Dionne's model is able to model both metals and dielectrics secondary emission but is one-dimension only. As the three-dimensional aspect is essential for complex configurations, the Dionne model is extended to three-dimensions. Measurements of the total electron emission yield have been carried out at the ONERA and shows a good agreement for low-impact energy and low-impact angle electrons, which is relevant in multipactor simulations.

I. INTRODUCTION

The multipactor effect is an electronic avalanche that may appear in vacuum RF systems when two conditions are met: (i) electrons enter in resonance with the RF field and (ii) their energy is sufficiently high to extract additional electrons from the system walls. In practice, multipactor appears when a RF electric field threshold E_{\max} is reached. This phenomenon has been observed in telecommunications satellites¹ and in RF plasma heating systems of experimental fusion reactors (tokamaks)², but also in particle accelerators³. The multipactor creates an electron cloud, leading to a disturbance or even a reflection to the source of RF signal. A breakdown can also be created, leading to a rapid heating of the surfaces, an alteration of the system material's properties or an out-gassing triggering the destructive corona effect⁴.

As multipactor threshold measurements are expensive, alternative methods have been developed to estimate it: theoretical^{5,6} and statistical⁷ studies, Monte-Carlo⁸ or Particle-in-Cell (PIC) simulations^{9,10}. A PIC code has been developed in Refs. 11–13 in order to understand the apparition of this phenomenon in complex RF devices from a physical point of view. This approach requires an accurate modelling of the physics at stake, and especially of the electron emission phenomenon.

Let us consider a sample impacted by a flux of electrons, called Primary Electrons (PEs). A second flux of electrons may be emitted by the sample; it is made up of Secondary Electrons (SEs, which are electrons removed from the material), of Inelastically and Elastically

Backscattered Electrons (IBEs and EBEs, which are PEs rediffused by the sample respectively with and without energy loss). The Secondary Electron Emission Yield (δ , SEEY) is defined as the ratio between the SEs and PEs fluxes. It is higher than one if there are more SEs than PEs. The Elastically Backscattered Electron Emission Yield (η_e , EBEEY) and Inelastically Backscattered Electron Emission Yield (η_i , IBEEY) are similarly defined. Finally, the Total Electron Emission Yield (σ , TEEY) is the ratio between all emitted electrons and PEs fluxes. $\sigma > 1$ is a necessary condition for the electron population to grow and thus for the multipactor to appear. We have the relation:

$$\sigma = \delta + \eta_e + \eta_i \quad (1)$$

A full modelling of the three emitted electrons populations require the knowledge of their emission yields, as well as their energies and angular distributions. All these quantities rely on the energy and direction of the PEs flux¹⁴, on the surface morphology of the sample¹⁵, on the presence of contaminants¹⁶ or external magnetic fields¹⁷. In addition, dielectric materials, which are widely used in spacecraft applications, can hold a net electric charge. It influences electron emission through two processes: firstly, the created surface potential may have an effect on electrons trajectories in the vacuum. Secondly, the excess of holes or electrons inside the material modifies the transport of electrons when they are in the sample. SEs are the population most influenced as they are the least energetic.

As SEs are the predominant population in waveguides, the SEEY mostly conditions the multipactor apparition. SEEY models exist^{18–20}, but, to our knowledge, the Dionne model²¹ is the only one able to take into account the influence of dielectric charge on the transport of SEs inside the material²² while being sufficiently simple for PIC simulations, which involve large numerical resources. Another asset of Dionne's model is that it gives δ , while some other models only give σ ^{23–25}. Thus, Dionne's

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model allows SEs to be discriminated from backscattered electrons. EBEs have been shown to be of great importance in multipactor apparition²⁶, especially when magnetic fields are present²⁷. EBEs were also shown to be significant in Hall plasma thrusters energy balance²⁸.

The main limitation of Dionne’s model is that it is one-dimension only and cannot take into account for the incidence angle of PEs; this restriction becomes a limitation when simulating complex 3D geometries such as circulators. Circulators are used to avoid RF power reflection to the source, that would cause its degradation. The central part of circulators is composed of ferrites that may have a dielectric behavior; they also create a static magnetic field parallel to the RF electric field, reason why a three-dimension simulation becomes mandatory. A preliminary study has been published¹³. We present here the complete study providing detailed explanations, comparison of the model on materials with various surface treatments and the introduction of the equivalent incidence angle notion.

In this paper, we develop a three-dimension version of Dionne’s model for such configurations. In section II, we make a simple simulation to evaluate the validity range relatable for multipactor study. In section III, we develop the secondary emission model. Finally, we benchmark it with in-house electron emission yield measurements in section IV.

II. MODEL FRAMEWORK

We define here the validity range of the proposed SEEY model in terms of PEs’ impact energy and incidence angle. Preliminary multipactor simulations have been made with the simulation code POTOMAC (Physical simulatiOn TOol for Multipactor in Advanced Configurations)^{11–13}. POTOMAC is a PIC code that computes the evolution of the electron population in a domain with a given RF electric field. Each time a collision occurs, δ , η_e and η_i are calculated and the appropriate emission mechanism is chosen. Here, the multipactor threshold was computed for a WR75 rectangular waveguide with silver coated walls. The RF signal has a frequency of 11 GHz and propagates in the $TE_{1,0}$ mode, which is the fundamental mode for this geometry. Dimensions are 19.05 mm \times 9.525 mm. The emission yield, angular emission distribution and energy emission distribution for the three populations are listed in Tab. I. All data calculated with the Monte-Carlo code SLAB²⁹ were tabulated in order to reduce the computational time.

We recorded the impact energy of all PEs at the multipactor threshold and represented it as an histogram in Fig. 1. Seed Electrons are the electrons injected in the waveguide during the first periods of the simulation. These results first show that SEs are predominant in POTOMAC simulations, and thus must be modelled with great attention. Secondly, 99% of PEs impact the walls with an energy lower than 230 eV. In order to take into

	Emission Yield	Angular distrib.	Energy distrib.
SE	Dionne (this paper)	Cosine law ³⁰	Chung <i>et al.</i> ³¹
EBE	SLAB ²⁹	SLAB ²⁹	Monoenergetic
IBE	Constant	Cosine law ³⁰	Uniform

TABLE I: Full electron emission model used with POTOMAC for the preliminary simulation

account for the various configurations, the model should be valid up to 300 eV.

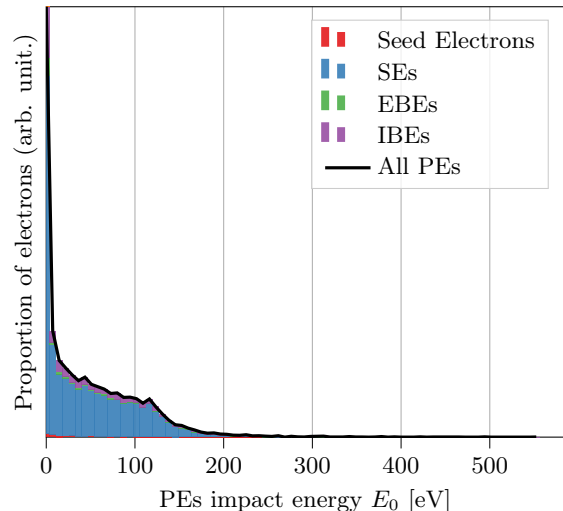


FIG. 1: Histogram of the PEs impact energy

Fig. 2 is a histogram of the SEs angular flux when they impact a wall $\theta_0/\sin\theta_0$, θ_0 being defined relative to the wall normal (see Fig. 3). It shows us that the majority of SEs impact surfaces almost perpendicularly; as we want our model to be usable in the broadest variety of scenarios possible, it should be valid for angles ranging from 0° to 40° .

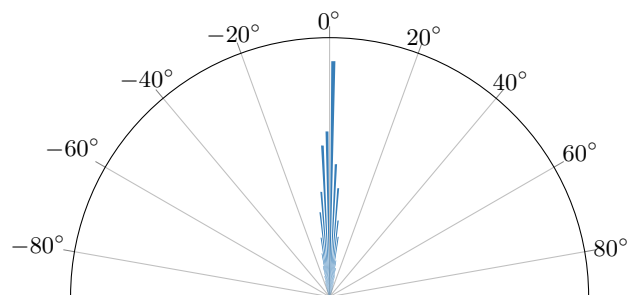


FIG. 2: Histogram of the flux of impacting SEs according to their collision angle θ_0

This basic configuration was chosen for the sake of simplicity, and for that it is representative. Using POTOMAC, we saw that in waveguides with smaller gaps,

electrons tended to impact walls with a smaller energy. The distribution of impact angle was slightly more spread around the central value of 0° . On the contrary, in waveguides with greater gaps, the impacting energy of electrons rose. This is not a problem, given that multipactor generally occurs where the gap is small and the electric field high. We also found that adding magnetic fields or dielectrics with a low surface charge did not change the order of magnitude of PEs' impact energy and incidence angle. In presence of highly charged dielectrics, some electrons can impact walls with energies up to a few keV; however, both the mean and the median of the E_0 distribution are shifted towards lower energies. This E_0 and θ_0 study should be led again for more complex geometries. We do not see any reason that a significant fraction of electrons impacted surfaces with energies higher than 300 eV. In the contrary, the θ_0 distribution may be severely distorted towards extreme angles. This could be the case in iris for example.

Most RF systems use copper or silver, which have comparable emission properties. The nature of the material do not influence much E_0 nor θ_0 . However, very poor emitting materials would drastically increase the electric field threshold \mathbf{E}_{\max} , which would in consequence increase E_0 .

The model we propose can be used in other fields than multipactor, as long as electrons impact surfaces with a low energy and a low incidence angle.

III. SEEY MODEL

A. Dionne in 1D

Numerous secondary emission models already exist; for example, Monte-Carlo methods are precise and can take into account all the physical processes involved, but are very slow, which is prohibitive for PIC simulations. On the other side of the spectrum, very fast and simple models exist, such as Vaughan^{23,24}. This model is 3D and can take into account the materials rugosity. However, it has the inconvenience of returning a null TEEY for very low energy PEs, which constitutes an important proportion of impacting electrons (see Fig. 1) and is not consistent with measurements³². Plus, this model cannot discriminate SEs from IBEs and EBEs.

A modified version of Vaughan has been proposed³³, which manually forces a TEEY equal to unity for low-energy electrons; it also fits a parameter $E_{\text{threshold}}$ to enforce the model to match $\sigma(E_{c1})$. E_{c1} is the first cross-over energy, defined as the first energy to verify $\sigma(E_{c1}) = 1$. This point has been shown of being of primordial importance for multipactor simulations³⁴. The modified Vaughan model is widely used by the multipactor community.

Dionne's model is a fast, simple and physical model for the SEEY. It makes the assumption that secondary emission process can be divided into three stpdf:

1. a PE enter a sample and emit SEs on its trajectory inside the material;
2. SEs are transported from their emission point to the surface;
3. SEs escape the sample through the surface.

All terms in the Dionne model have a physical meaning. Thus, it is for example possible to analytically reflect the influence thin coatings have on the SEEY and hence on the multipactor^{12,35}. However, Dionne's model does not take into account for the incidence angle of the PE, unlike Vaughan.

For the sake of clarity, Dionne's model is explained here. Under the hypothesis of a constant energy loss for the PE and no scattering for the SEs, it may be written as:

$$\delta = G \cdot T \cdot S \quad (2)$$

Where the generation term G represents the mean number of SEs produced by the PE, the transmission term T represents the probability for the produced SEs to reach the surface and the escape term S is a sample dependent constant representing the probability for the SEs to cross the surface. This terms can be expressed as:

$$G = \frac{E_0^n}{\Phi \cdot R(E_0)} \quad (3)$$

$$T = d \cdot \left(1 - e^{-R(E_0)/d}\right) \quad (4)$$

Where n is a material dependent constant, E_0 is the energy of the PE and Φ the material effective work function³⁶ defined relative to the vacuum level in eV. For a metal, $\Phi = W_f$ with W_f the work function and for a dielectric, $\Phi = \chi + E_g$ with χ the electron affinity and E_g the gap energy. R is the range of the PE (see Fig. 3) and d the diffusion length of the PE in the material in nm. Using the continuous slowing down approximation, with A a surface-dependent constant, the range can be expressed as^{21,37}:

$$R(E_0) = \frac{E_0^n}{A \cdot n} \quad (5)$$

Φ can be measured or calculated³⁸. It was suggested that d was equal to the inelastic mean-free path³⁹; however, very low energy electrons trajectories are mostly driven by elastic collisions¹⁸, making this assumption unsuitable for our study. d also can be calculated with the density and the effective absorption coefficient⁴⁰. Thermionic emission experiments allow us to measure S ⁴⁰. A and n can be extracted from Monte-Carlo simulations or measurements³⁷.

In this study, Φ will be taken from literature while A , n , d and S will be fitted, for that they drastically vary from one sample to another – especially when samples are contaminated.

B. Dionne in 3D

At a given energy and if the considered sample is perfectly flat, PEs penetrate less and less deep while θ_0 increases (see Fig. 3). Thus, SEs are emitted closer and closer to the surface and are more likely to be emitted in the vacuum, increasing the SEEY.

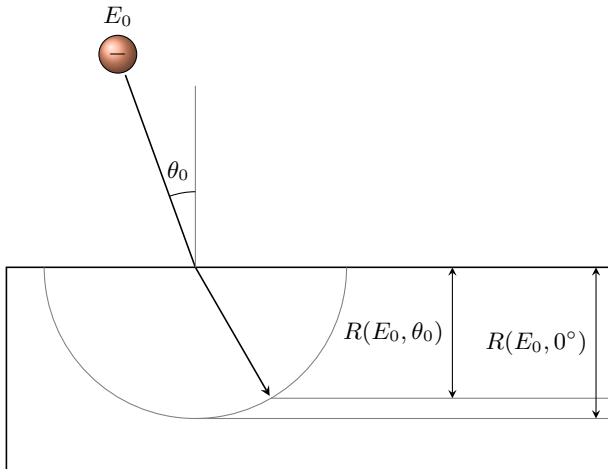


FIG. 3: Macroscopic view of a PE impacting a sample with an energy E_0 and an angle θ_0

The Dionne model was not designed to take the incidence angle θ_0 into account. As it is physical, it is however possible to add an angle dependency. Modifying θ_0 does not influence the distance travelled by the PE but only its range (see Fig. 3). The Snell's law for electron reads⁴¹:

$$\sqrt{E_0} \sin \theta_0 = \sqrt{E_1} \sin \theta_1 \quad (6)$$

where E_1 and θ_1 are respectively the energy and angle of the electron after it has been refracted in the sample. The range thus can be expressed as:

$$R(E_0, \theta_0) = R(E_0, 0^\circ) \sqrt{1 - \frac{E_0}{E_0 + \Phi} \cdot \sin^2 \theta_{\text{equ}}} \quad (7)$$

θ_{equ} is the equivalent incidence angle. For a perfectly flat sample, θ_{equ} can be replaced by θ_0 . For any other material, θ_{equ} is a function of θ_0 and of the sample roughness profile. It firstly takes into account the fact that the macroscopic angle θ_0 is different from the microscopic impact angle, due to the surface roughness (oblique incidence effect). Other phenomena are taken into account: if the sample presents a small roughness, low-energy high incidence angle electrons that could not be emitted on a flat surface can escape an inclined plane, increasing the SEEY at low energy¹⁵. For high roughness materials, SEs are more likely to be recollected by the sample, which causes a reduction of the SEEY (shading effect)¹⁵. This effect is more pronounced in presence of magnetic fields¹⁷. Recollected SEs can also emit other electrons (multigeneration).

Zhang *et al.*⁴² showed with Monte-Carlo simulations that it was possible to replace θ_0 with θ_{equ} in TEEY calculations to take all these effects into account, for random rough surfaces with Gaussian height distribution. Their results suggested that a first order approximation could be:

$$\theta_{\text{equ}}(\theta_0) = k_a \theta_0 + k_b \quad (8)$$

With k_a and k_b roughness profile dependent constants. For a perfectly flat material, $\theta_{\text{equ}} = \theta_0$ which implies $k_a = 1$ and $k_b = 0^\circ$.

IV. MODEL VERIFICATION

A. Experimental procedure

The TEEY measurements were performed in an Ultra High Vacuum (UHV) facility designed for the study of electron emission. A dry turbo-molecular pump associated with an oil-free primary pump allows the system to be maintained at a vacuum level down to 3×10^{-9} mbar. The sample holder allows the variation of the electron incidence angle from 0° (normal incidence angle) to 80° . An ELG-2022B electron beam from Kimball Instrument was used. The electron beam was pulsed during TEEY measurements to limit the surface conditioning effect⁴³. The facility is equipped with a Faraday cup, a hemispherical electron energy analyser and a X-ray MgK α source. The sample holder and the collector can be independently biased to choose the desired potential. The sample holder is ordinarily negatively biased in order to prevent the low energy emitted electron to come back to the sample surface and prevent the generation of tertiary electrons.

TEEY measurements were made on several samples [Associated dataset available at <http://dx.doi.org/10.5281/zenodo.3648704>] (Ref. 44). In this paper, we decided to compare the developed model to copper only, which is often used in RF components. We did so to facilitate the comparison between samples and to focus on their emission properties evolution with surface treatment.

The first sample, called hereafter Cu #1, was heated *in situ* and under UHV at 200°C for 2 h. A second set of measurements was made after 3 h of *in situ* erosion under 1 keV Ar^+ flux at normal incidence. XPS study showed that after the heating, the sample was contaminated, mainly with oxygen and carbon; after the erosion, it was almost pure.

A second set of measurements was made on another copper sample, Cu #2. Measurements were made on the sample without any treatment (as received), after 2 h of heating at 200°C *in situ* and under UHV, and after 1 h of *in situ* erosion under 1 keV Ar^+ flux at normal incidence.

B. Fitting process

The total backscattered electron emission yield η is defined as:

$$\eta = \eta_e + \eta_i \quad (9)$$

We computed η with the Monte-Carlo software CASINO v3.3.0.4⁴⁵. ELSEPA model was used for the total and partial cross-section calculation⁴⁶, the direction cosines were calculated with Lowney *et al.*⁴⁷, the stopping power and ionisation potential with Joy *et al.*⁴⁸. The pseudo-random number generator was a Lagged Fibonacci generator. Results showed a very weak dependence of η on θ_0 .

We computed η_e with the SLAB model²⁹. Due to the lack of literature on the subject, we considered that η_i was a constant, independent from E_0 and θ_0 . As our values of η and η_e are accurate for $E_0 = 500$ eV and $\theta_0 = 0^\circ$, we set η_i so that Eq. (9) holds at 500 eV and 0° . The Dionne model was then fitted on measurements with E_0 in 0–300 eV and θ_0 in $0^\circ - 40^\circ$.

C. Evaluation of the model's capacity to fit σ

The Dionne model was compared to modified Vaughan model as defined in Ref. 33. For almost all materials, Vaughan rugosity parameters k_E and k_θ should be equal to unity. They may be lower for intentionally roughened materials, down to 0 for textured carbon. They may be up to 2 for extremely smooth surfaces. The modified Vaughan parameter $E_{\text{threshold}}$ and Vaughan rugosity parameters were also fitted on 0–300 eV and $0^\circ - 40^\circ$, k_E and k_θ being constrained between 0.5 and 2.

Fil *et al.* studied the sensitivity to surfaces' TEEY curves of the multipactor threshold obtained by simulation⁴⁹. They realised more than four hundred multipactor simulations with the software Spark3D⁵⁰. The simulated geometry was a waveguide of width 9.525 mm and a gap of 0.1 mm, working at 12 GHz. They showed that, for this system, the calculated multipactor threshold was particularly sensitive to the value of E_{c1} and to the profile of $\sigma(E_{c1} < E_0 < E_{0,\text{max}})$. $E_{0,\text{max}}$ is the energy for which the TEEY is maximum and is generally between 200 eV and 1 keV. Very few electrons impact surfaces with energies this high, and multipactor apparition is probably most influenced by the profile of σ between E_{c1} and E_{c1} plus a few tenths of eV. POTOMAC simulations showed that the E_0 and θ_0 validity range defined in Section II were still relevant in this geometry.

These requirements were shown to be valid in one particular geometry only, and did not address incidence angle problematics. In order to compare Dionne's and Vaughan's models in a way that is not relevant only for multipactor in small gap waveguides, we decided to evaluate their performance on more criteria. For θ_0 in $0^\circ - 40^\circ$, we calculated the standard deviation between

the measured and calculated TEEY for $E_0 < E_{c1}$ and for $E_0 > E_{c1}$; the standard deviation on the whole energy range was also computed. Finally, we calculated the standard deviation between the measured and calculated first cross-over energies. Although this approach does not quantify the contribution of backscattered electrons to multipactor, it is a convenient way to compare the different models to the measures.

D. Sample Cu #1 results

The measured and fitted TEEY for the sample Cu #1 after heating are represented in Fig. 4. The standard deviations between the models and the measurements are presented in Tab. II. Results show an overall good agreement between both models and measurements, on the validity range we defined. For low energies under E_{c1} and for all θ_0 , Dionne's model underestimates the TEEY and Vaughan's model fits better. On the contrary, Dionne's model show better agreement than Vaughan's for energies above E_{c1} – excluding $\theta_0 = 60^\circ$. On the whole energy range, Vaughan's model shows a smaller standard deviation from the measurements than Dionne's model at $\theta_0 = 0^\circ$, but it rapidly grows with θ_0 and Dionne fits better for 40° and 60° . Dionne's error is more independent from the incidence angle than Vaughan's. As the modified Vaughan model was designed to fit $E_{c1}(0^\circ)$, its deviation from the measurements at this point is very low. It increases with θ_0 and, at 60° , Dionne fits better than Vaughan.

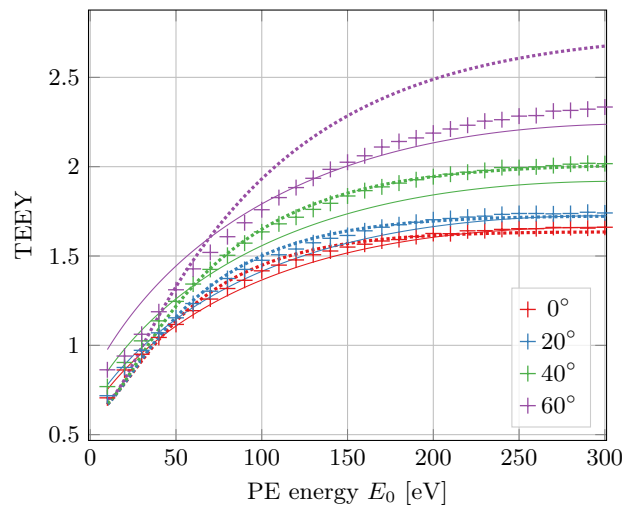


FIG. 4: Measured and fitted TEEY on Cu #1 (heated; + measurements, Dionne fit, — Vaughan fit) [Associated dataset available at <http://dx.doi.org/10.5281/zenodo.3648704>] (Ref. 44)

The results for eroded Cu #1 are represented in Fig. 5, and standard deviations in Tab. III. Excluding $\theta_0 = 20^\circ$, results are similar to heated Cu #1. Vaughan fits better σ

		$\sigma(E_0 < E_{c1})$	E_{c1} [eV]	$\sigma(E_0 > E_{c1})$	$\sigma(0-300\text{ eV})$
Dionne	0°	0.054	1.9	0.025	0.029
	20°	0.066	3.1	0.020	0.028
	40°	0.101	5.7	0.023	0.034
Vaughan	0°	0.027	0.0	0.026	0.026
	20°	0.037	1.6	0.038	0.038
	40°	0.080	6.1	0.086	0.086

TABLE II: Heated Cu #1 standard deviation for the different regions of interest

under E_{c1} , and Dionne fits better after E_{c1} . On the whole energy range, the error of Dionne is more constant than Vaughan's. Again, the Vaughan model perfectly matches $E_{c1}(0^\circ)$. Here, $\sigma(E_0, 0^\circ) > \sigma(E_0, 20^\circ)$, which cannot be reproduced by Dionne nor Vaughan. Plus, the TEEY drastically increases at 60° and is thus underestimated by the models. Vaughan results are similar to Dionne.

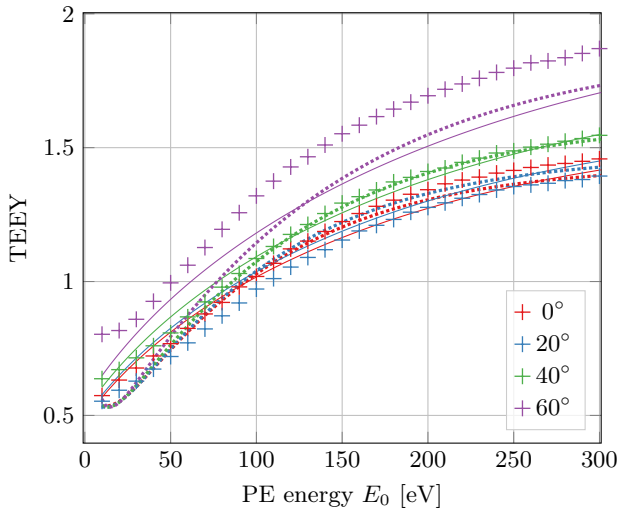


FIG. 5: Measured and fitted TEEY on Cu #1 (eroded; + measurements, Dionne fit, — Vaughan fit)

[Associated dataset available at <http://dx.doi.org/10.5281/zenodo.3648704>] (Ref. 44)

		$\sigma(E_0 < E_{c1})$	E_{c1} [eV]	$\sigma(E_0 > E_{c1})$	$\sigma(0-300\text{ eV})$
Dionne	0°	0.042	1.3	0.040	0.041
	20°	0.047	15.3	0.054	0.052
	40°	0.075	2.5	0.010	0.040
Vaughan	0°	0.026	0.0	0.057	0.050
	20°	0.082	16.6	0.038	0.057
	40°	0.036	4.9	0.026	0.029

TABLE III: Eroded Cu #1 standard deviation for the different regions of interest

The obtained fitting parameters for Cu #1 are summarized in Tab. IV.

		Heated	Eroded
Dionne	A	6.22	5.43
	n	1.00	1.00
	d [nm]	7.76	17.42
	S	0.13	0.06
	k_a	0.92	0.74
	k_b [°]	2.26	0.00
Vaughan	Φ [eV]	4.65	—
	η_i	0.25	—
	σ_{\max}	1.66	1.53
	$E_{0, \max}$ [eV]	300	550
	$E_{\text{threshold}}$ [eV]	-21.81	-28.93
	k_E	0.50	1.72
	k_θ	2.00	2.00

TABLE IV: Evolution of Cu #1 fitting parameters

E. Sample Cu #2 results

The results for the sample Cu #2 without any surface treatment are represented in Fig. 6 and the standard errors in Tab. V. Both Dionne and Vaughan models work very well. Dionne fits better σ on the whole energy range and for all angles – except for $\sigma(E_0 > E_{c1}, \theta_0 = 40^\circ)$. Dionne even fits 60° which is out of the validity range. Meanwhile, Vaughan fits better E_{c1} .

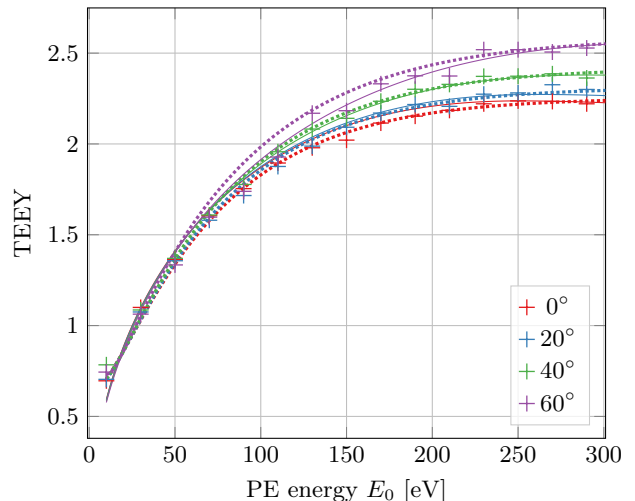


FIG. 6: Measured and fitted TEEY on Cu #2 (as received; + measurements, Dionne fit, — Vaughan fit)

[Associated dataset available at <http://dx.doi.org/10.5281/zenodo.3648704>] (Ref. 44)

Heated Cu #2 results are represented in Fig. 7, standard errors in Tab. VI. For this material, the

		$\sigma(E_0 < E_{c1})$	E_{c1} [eV]	$\sigma(E_0 > E_{c1})$	$\sigma(0-300\text{ eV})$
Dionne	0°	0.007	3.2	0.026	0.025
	20°	0.002	2.1	0.031	0.030
	40°	0.080	3.6	0.023	0.030
Vaughan	0°	0.102	0.0	0.038	0.045
	20°	0.113	0.8	0.043	0.051
	40°	0.200	1.3	0.020	0.055

TABLE V: As received Cu #2 standard deviation for the different regions of interest

Dionne model fits much better the measurements on all the regions of interest, as Vaughan underestimates the TEEY. However, Vaughan matches better $E_{c1}(0^\circ)$. Here, $\sigma(60^\circ) > \sigma(40^\circ)$, leading to a wild overestimation of the TEEY at 60° for both models.

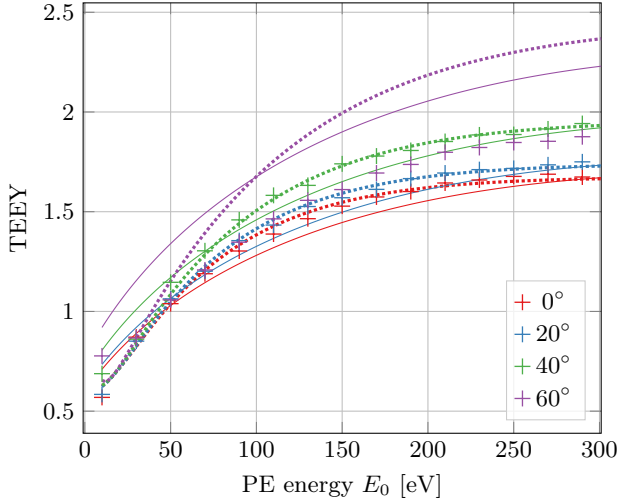


FIG. 7: Measured and fitted TEEY on Cu #2 (heated; + measurements, Dionne fit, — Vaughan fit) [Associated dataset available at <http://dx.doi.org/10.5281/zenodo.3648704>] (Ref. 44)

		$\sigma(E_0 < E_{c1})$	E_{c1} [eV]	$\sigma(E_0 > E_{c1})$	$\sigma(0-300\text{ eV})$
Dionne	0°	0.052	0.9	0.022	0.028
	20°	0.033	1.0	0.017	0.020
	40°	0.044	2.9	0.023	0.027
Vaughan	0°	0.101	0.1	0.060	0.067
	20°	0.117	3.7	0.056	0.068
	40°	0.134	11.0	0.051	0.068

TABLE VI: Heated Cu #2 standard deviation for the different regions of interest

Finally, the measured and fitted TEEY of eroded Cu #2 are represented in Fig. 8 and the standard de-

viations in Tab. VII. According to Fil *et al.* criteria⁴⁹ and considering only $\theta_0 = 0^\circ$, Vaughan model is more appropriate for multipactor simulations. However, for θ_0 in $0^\circ - 40^\circ$, the standard deviation of Dionne is lower for $\sigma(0-300\text{ eV})$ and $\sigma(E_0 < E_{c1})$. For $\theta_0 \geq 20^\circ$, Dionne also fits better $\sigma(E_0 > E_{c1})$. The measured TEEY at 60° is lower than the TEEY at 40° , leading to a wide overestimation of the measurements by both models.

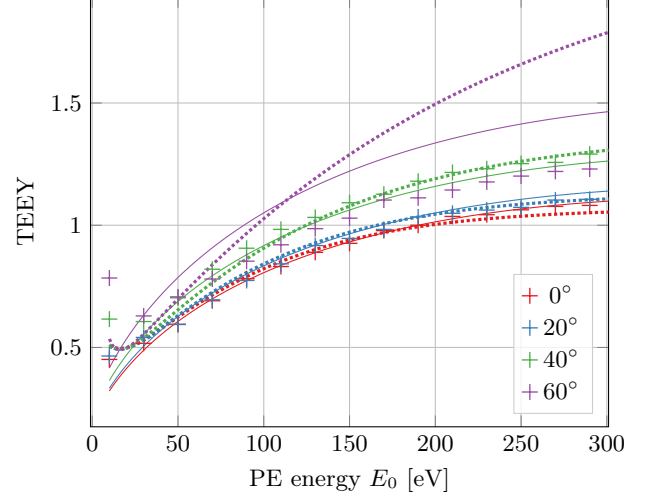


FIG. 8: Measured and fitted TEEY on Cu #2 (eroded; + measurements, Dionne fit, — Vaughan fit) [Associated dataset available at <http://dx.doi.org/10.5281/zenodo.3648704>] (Ref. 44)

		$\sigma(E_0 < E_{c1})$	E_{c1} [eV]	$\sigma(E_0 > E_{c1})$	$\sigma(0-300\text{ eV})$
Dionne	0°	0.024	6.4	0.024	0.025
	20°	0.029	10.9	0.008	0.023
	40°	0.059	6.7	0.012	0.039
Vaughan	0°	0.042	0.1	0.007	0.035
	20°	0.050	8.4	0.023	0.041
	40°	0.107	9.2	0.030	0.072

TABLE VII: Eroded Cu #2 standard deviation for the different regions of interest

All fitting parameters are listed in Tab. VIII.

V. DISCUSSION

The proposed model works very well on the defined validity range (i.e., E_0 in $0-300\text{ eV}$ and θ_0 in $0^\circ - 40^\circ$).

According to Fil *et al.*⁴⁹, TEEY models must fit precisely $\sigma(E_0 > E_{c1})$ to ensure a precise prediction of the multipactor threshold. This conclusion is valid only in small gaps structures. The Dionne model address this requirement on more materials and incidence angles than

		As received	Heated	Eroded
Dionne	A	6.67	4.98	8.50
	n	1.00	1.00	1.00
	d [nm]	9.47	12.40	9.84
	S	0.14	0.10	0.04
	k_a	0.40	0.83	1.13
	k_b [°]	8.04	1.12	0.00
	Φ [eV]	4.65	–	–
	η_i	0.25	–	–
Vaughan	σ_{\max}	2.24	1.70	1.12
	$E_{0,\max}$ [eV]	250	390	390
	$E_{\text{threshold}}$ [eV]	1.11	–25.09	–6.51
	k_E	2.00	0.50	0.50
	k_θ	0.82	2.00	2.00

TABLE VIII: Evolution of Cu #2 fitting parameters

the Vaughan model. However, the second most important requirement is the precision around E_{c1} , which is better answered by Vaughan’s model. Still, as Dionne’s model can discriminate secondary from backscattered electrons, we think that it is more suitable for multipactor simulations.

On a more general point of view, the standard deviation of $\sigma(0-300\text{ eV})$ is generally lower for Dionne than for Vaughan. For Cu #2, it is always the case. For Cu #1, this difference is less significant. Dionne standard deviation of $\sigma(E_0 < E_{c1})$ was lower for Cu #2, but higher for Cu #1.

It is to be noted that Vaughan’s fit was much better than Dionne for energies higher than a few hundreds of eV, and that its implementation is easier. Vaughan rugosity parameters are often inconsistent, for example for heated Cu #2, where $k_E = 0.50$ suggests that the sample is very rough, while $k_\theta = 2.00$ tells the exact opposite.

Results however show some discrepancy in the very low energy range ($\lesssim 20\text{ eV}$), especially for non-normal incidences and eroded materials. We propose several explanations; firstly, the range expression (Eq. (5)) was obtained with the hypothesis that electrons lost their energy in the material in a continuous way (continuous slowing down approximation), which is inexact at low energies. In reality, the range may be constant below a few tenths of eV^{51,52}. Secondly, we set a constant IBEEY as we lack literature on the subject, which is a very important hypothesis. Thirst, the potential of the sample holder (-9 V for these experiments) can deflect significantly very low-energy electrons at non-normal incidence. A method that polarizes the collector instead of the sample has been developed⁵³. Finally, the magnetic field of the Earth may have an influence for very low energies⁵³.

In Fig. 9, we represented θ_{equ} as a function of θ_0 (Eq. (8)) for the different materials, with the values of k_a and k_b listed in Tab. IV and VIII. We also plotted Zhang *et al.* results, obtained with Monte-Carlo simulations; the three black lines represent the evolution of the equivalent angle for three copper samples with different random ru-

gosity profiles⁴². As we can see from the shape of Zhang’s curves, the choice of modelling the θ_{equ} law with a first-order polynomial seems to be justified. However, the materials studied in this paper are all softer than Zhang’s.

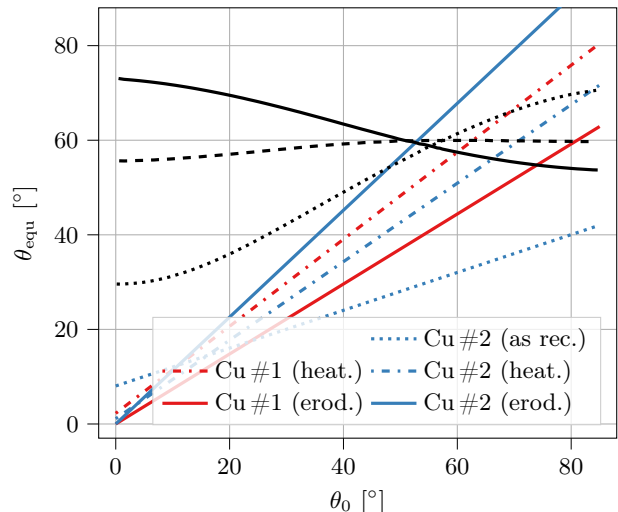


FIG. 9: θ_{equ} laws for different materials; black lines represent Zhang *et al.* simulation results on different random surfaces with a Gaussian height profile⁴²

All trends are not faithfully reproduced when increasing θ_0 leads to a reduction of σ (i.e. Cu #1 eroded, Cu #2 heated, Cu #2 eroded). In these cases, the variation of σ with θ_0 is non monotonous, which cannot be reproduced by Eq. (8). Using higher-order polynomials of the form $\theta_{\text{equ}} = \sum k_i \theta_0^i$ may solve the issue. However, the values of the coefficients k_i have to be chosen carefully so that, for the defined range of θ_0 , $\theta_{\text{equ}} \in [0, \pi]$. We noted that, even for polynomials of order two respecting this condition, θ_{equ} could take very unlikely values for $\theta_0 > 40^\circ$.

New emission yield measurements campaigns associated with rugosity profile studies would be required to investigate further the relation between θ_0 and θ_{equ} . Also, the hypothesis of a random rough surface with a Gaussian height distribution may be not representative of the real surface profile.

For the two materials, the surface treatments reduce contamination and thus the emission yield. As SEs have a low energy, the only ones able to escape are those produced in the first nanometers below the surface. Thus, they are very influenced by the presence of contaminants¹⁶. For both samples, the reduction of the contamination layer reduces the parameter k_b , which is 0 for eroded coppers. Our interpretation is that, for this materials, the contamination layer at the surface of the samples has a rougher surface state than the bulk material underneath; this is particularly noticeable on Cu #2, where the difference of rugosity between the sample without any surface treatment and after the heating is tremendous. It is to be noted that, even for clean ma-

materials with $k_a \approx 1$ and $k_b \approx 0$, the introduction of θ_{equ} increased the accuracy of our model. We observed that the simpler law $\theta_{\text{equ}} = k\theta_0$ ¹³ was precise enough for all materials, including Cu #2 as received; though, it cannot model the Zhang *et al.* materials.

The validity of the present model has not been tested yet on dielectrics. As the dielectric material charge would deflect electrons, the angle between the electron canon and the sample θ_{canon} would be different from the electrons incidence angle θ_0 . Proper measurements require additional work and are not in the scope of the present study.

VI. CONCLUSIONS

In this paper, we developed a simple extension of Dionne’s model to three-dimension. We used the refraction law for electrons to compute the range of electrons for various incidence angles. We used the concept of equivalent angle θ_{equ} to model the differences between macroscopic and microscopic incidence angle, as well as shading and multigeneration effects.

The present model was compared to in-house measurements on two copper samples with several surface treatments. The fit was made on the low-energy range (0–300 eV) and low-angle range (0°–40°), which is relevant for multipactor simulations. As a whole, the model fit experimental data in these ranges. The fit is however less good for high incidence angles and low energies, as well as for some materials where the evolution of σ with θ_0 is not monotonous.

The introduction of θ_{equ} was indispensable for the most contaminated material (Cu #2 as received), but also greatly increased the accuracy of the fit for purer materials. The relation between θ_0 and θ_{equ} should be explored, with Monte-Carlo simulations and emission yield measurements associated with rugosity profile studies. As magnetic fields were shown to influence electron emission from rough samples¹⁷, studying the link between θ_{equ} and magnetic fields would be another step towards extensive multipactor studies in presence of magnetic fields. Fits on silver led to similar results but were not included in this publication.

A comparison was made with modified Vaughan model, widely used in the multipactor community. At very low energy, the Dionne model showed to be more precise. At higher energies up to 300 eV, Dionne’s model gave similar or better results than Vaughan.

Several points could enhance the model: firstly, it has been proven that SLAB model was too simplistic⁵⁴ and thus could be replaced by OKGM model even if it generally underestimates η_e ^{55,56}. A simple and physical model would be of great value, as it would avoid the expensive computational cost of tabulating large arrays. Plus, it would allow a deeper understanding of the physics at stake.

Secondly, the modeling of η_i is also simplistic – we lack

a simple physical model for the IBEEY. Finally, the used range function is not valid on the whole energy range, R being almost constant below a few tenths of eV^{51,52}.

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