



GHz-rate pulse pairs generated by a frequency-shifting loop containing an electro-optic modulator

Hongzhi Yang, Marc Brunel, Marc Vallet, Haiyang Zhang, Changming Zhao

► To cite this version:

Hongzhi Yang, Marc Brunel, Marc Vallet, Haiyang Zhang, Changming Zhao. GHz-rate pulse pairs generated by a frequency-shifting loop containing an electro-optic modulator. Journée du Club Optique et Micro-ondes (JCOM 2019), Jun 2019, Brest, France. hal-02470354

HAL Id: hal-02470354

<https://hal.science/hal-02470354>

Submitted on 12 Feb 2020

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Pulse doublets generated by a frequency-shifted loop containing an electro-optic amplitude modulator

H. YANG,^{1,2} M. Brunel,¹ M. Vallet,¹ H. ZHANG,² C. ZHAO²

¹Univ. Rennes, CNRS, Institut FOTON - UMR 6082, 35000 Rennes, France

²Beijing Institute of Technology, School of Optics and Photonics, China

Introduction

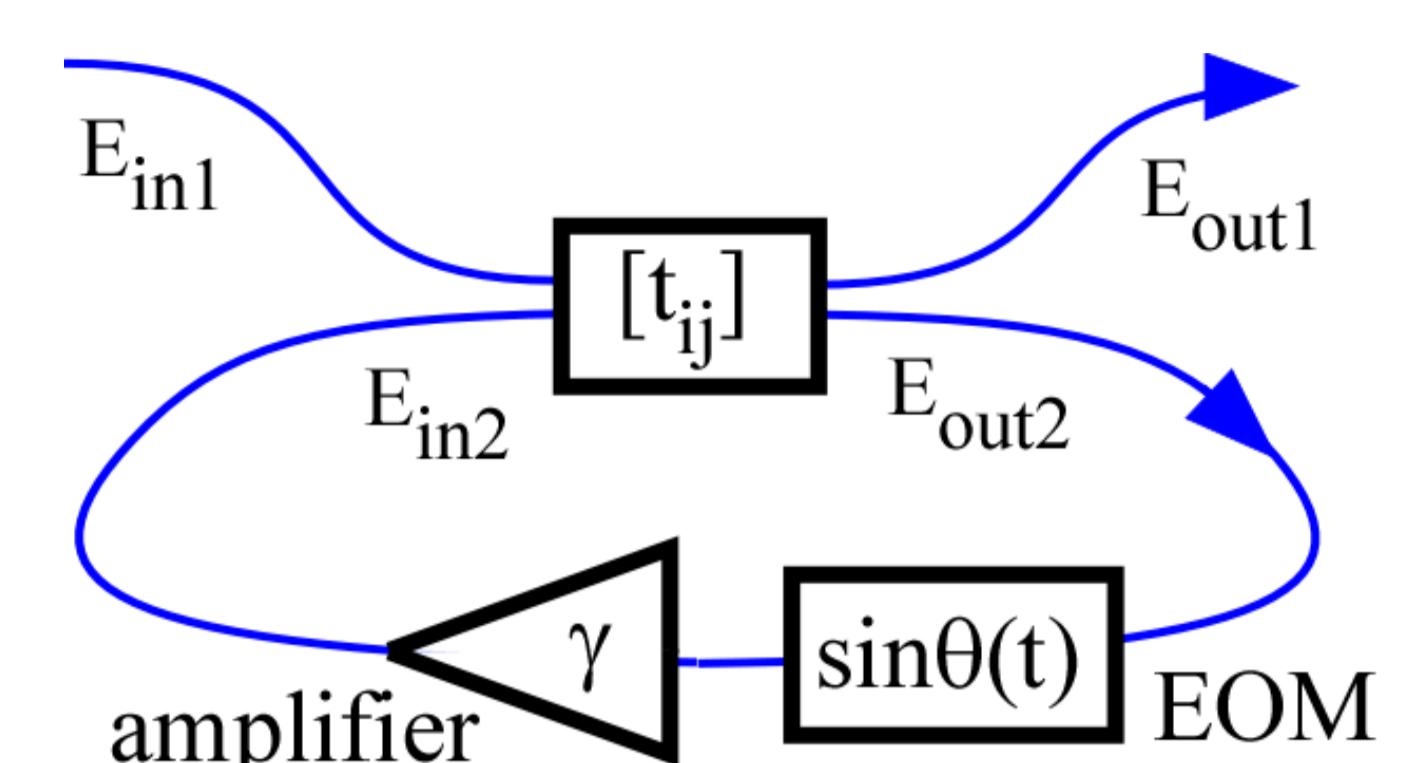
Frequency-shifted feedback loops, both active or passive, are promising solutions for pulse generation with high repetition rate. FSFL usually contains an acousto-optic frequency shifter AOFS for single-side band modulation [1-3].

This leads to poor tunability and low modulation frequency f_m .

Conversely: the use of in-loop electro-optic amplitude modulators EOM permits high modulation, wide bandwidth and integration [4-5].

However, due to dual-side band modulation, EOM lead to new regimes

→ Here, theoretical and experimental study of FSFL with EOM



Theoretical model E_{in1} cw + modulation at f_m

$$E_{out1}(t) = \left(t_{11} + t_{12}t_{21} \sum_{p=1}^N \left(t_{22}^p \gamma^p \prod_{q=1}^p \sin(\theta(t - q\tau)) \right) \right) E_{in1}$$

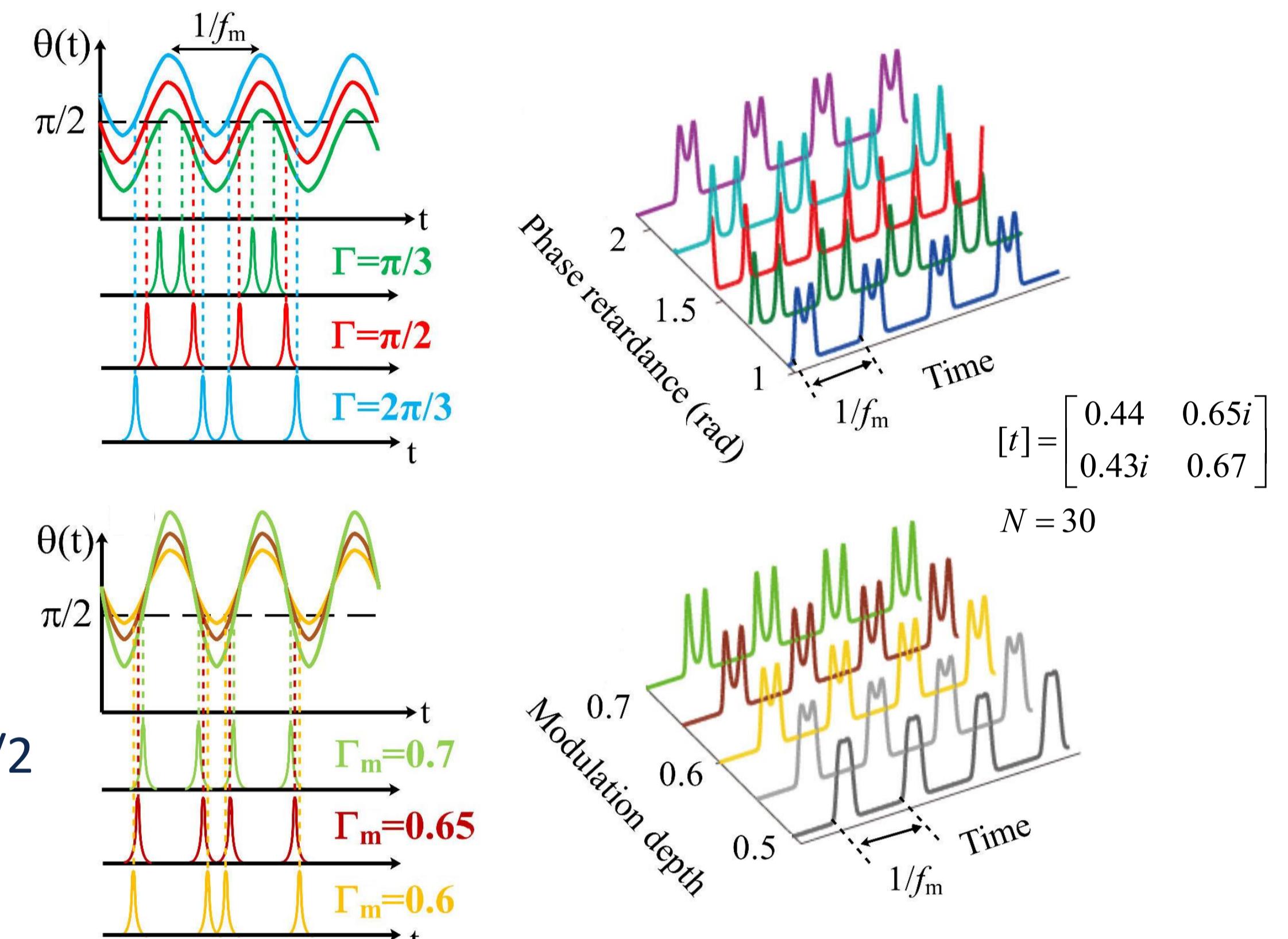
Γ phase retardance
 Γ_m modulation depth

$$\theta(t) = \Gamma + \Gamma_m \sin(2\pi f_m t)$$

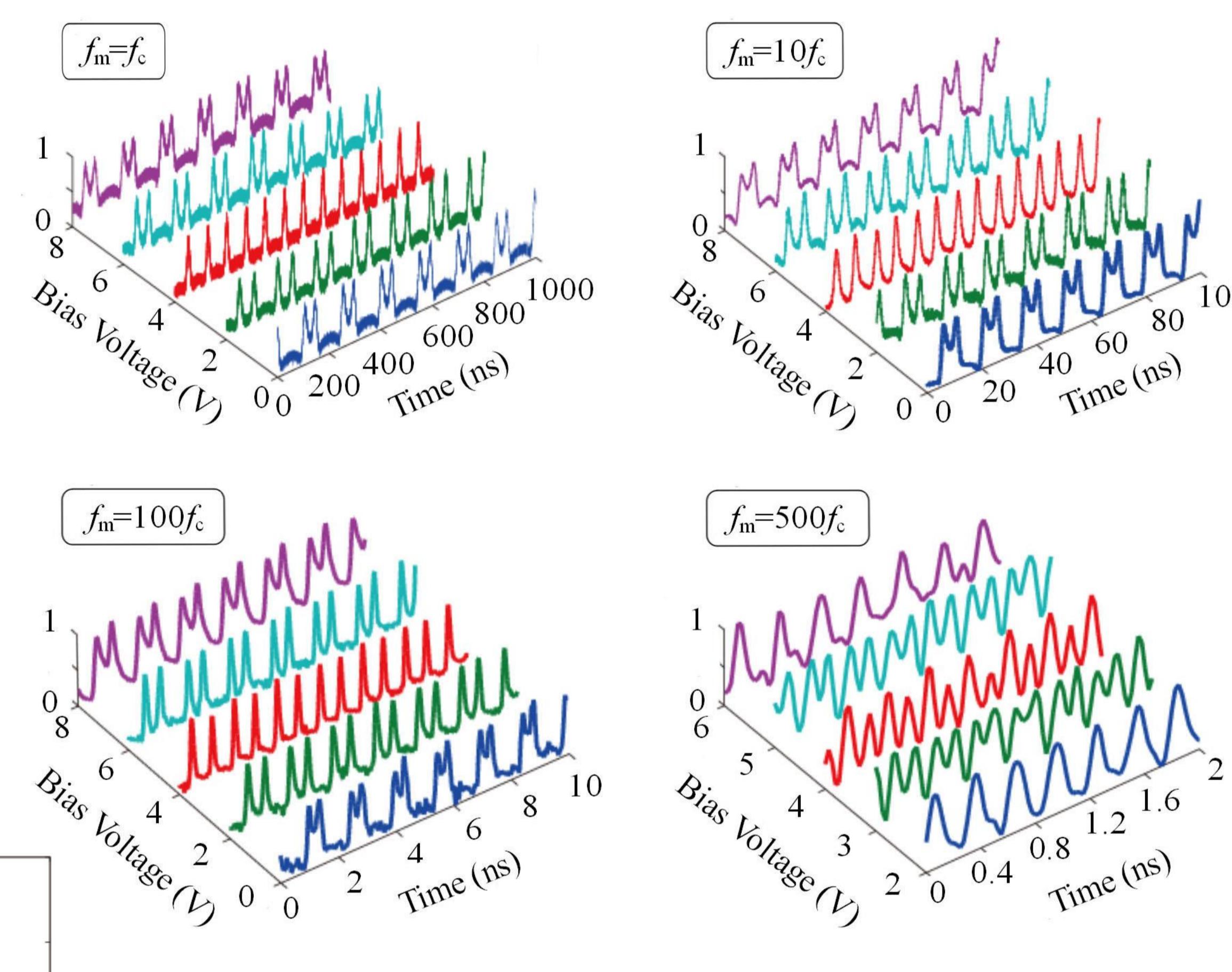
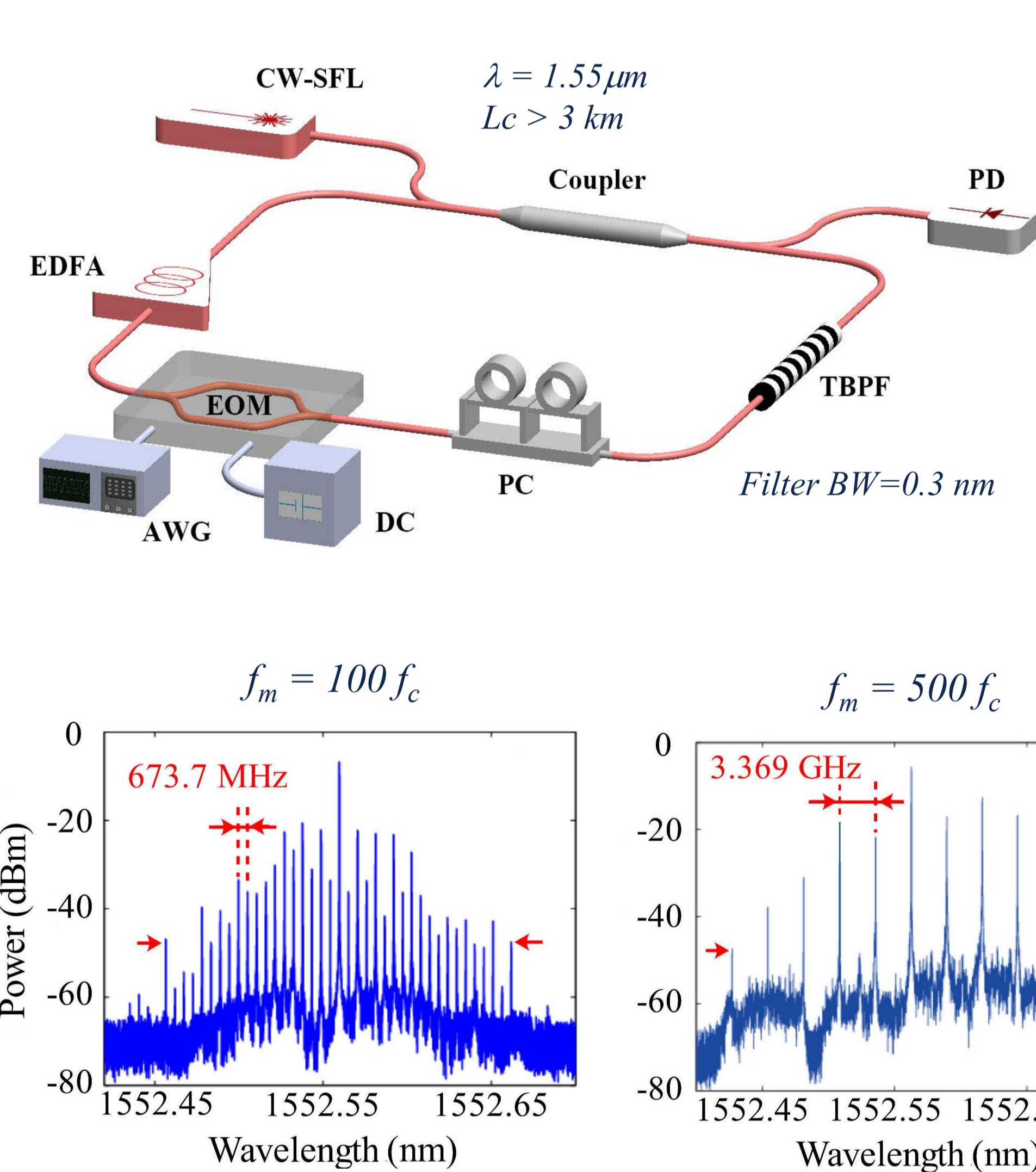
For integer Talbot condition [1-2]: $f_m = n/\tau$

$$\text{If } N \gg 1 \quad E_{out1}(t) = \left(\frac{t_{11} - \gamma(t_{11}t_{22} - t_{12}t_{21}) \sin \theta(t)}{1 - \gamma t_{22} \sin \theta(t)} \right) E_{in1}$$

→ Pulses for $\theta(t) \approx \pi/2$

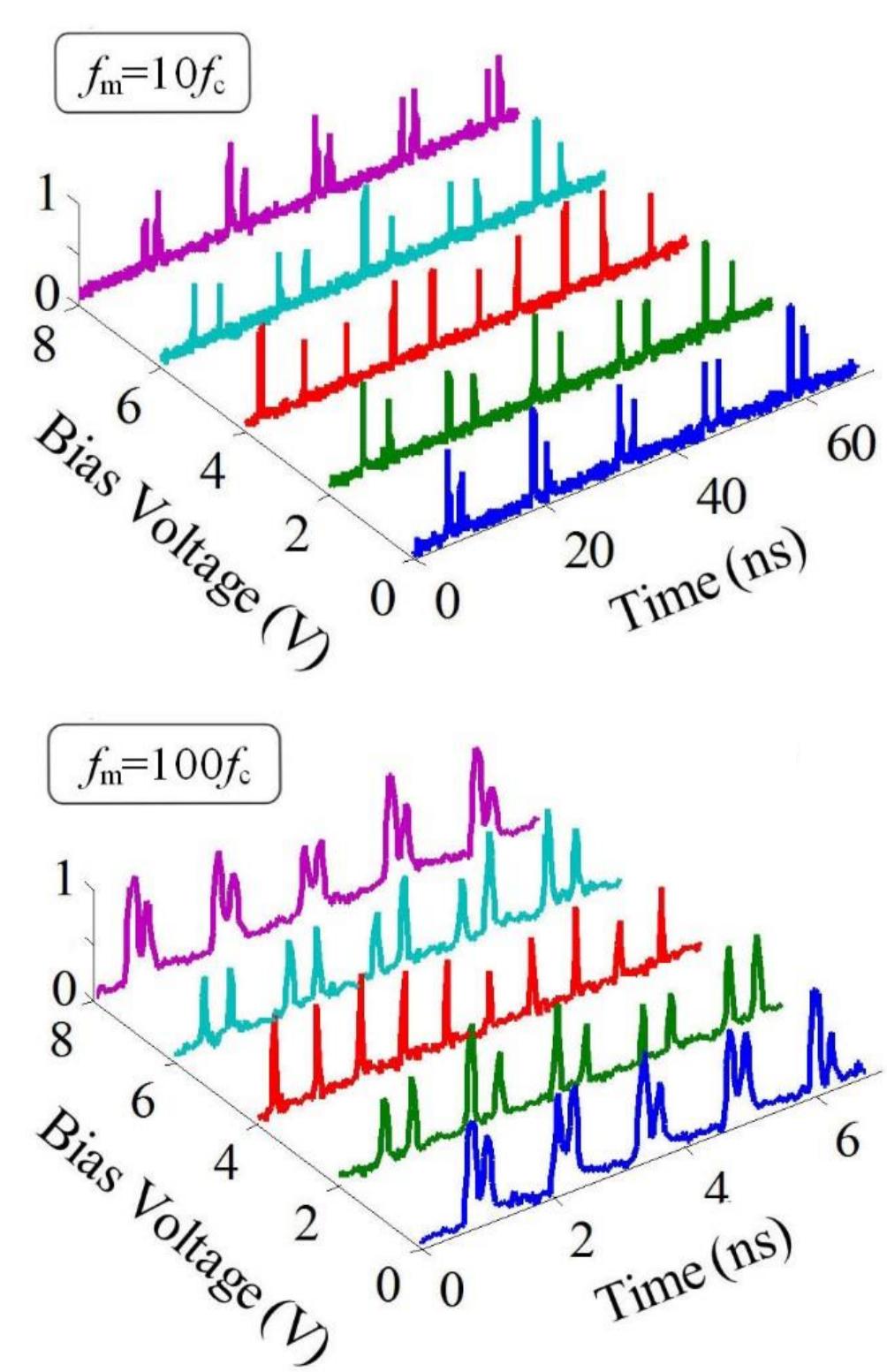
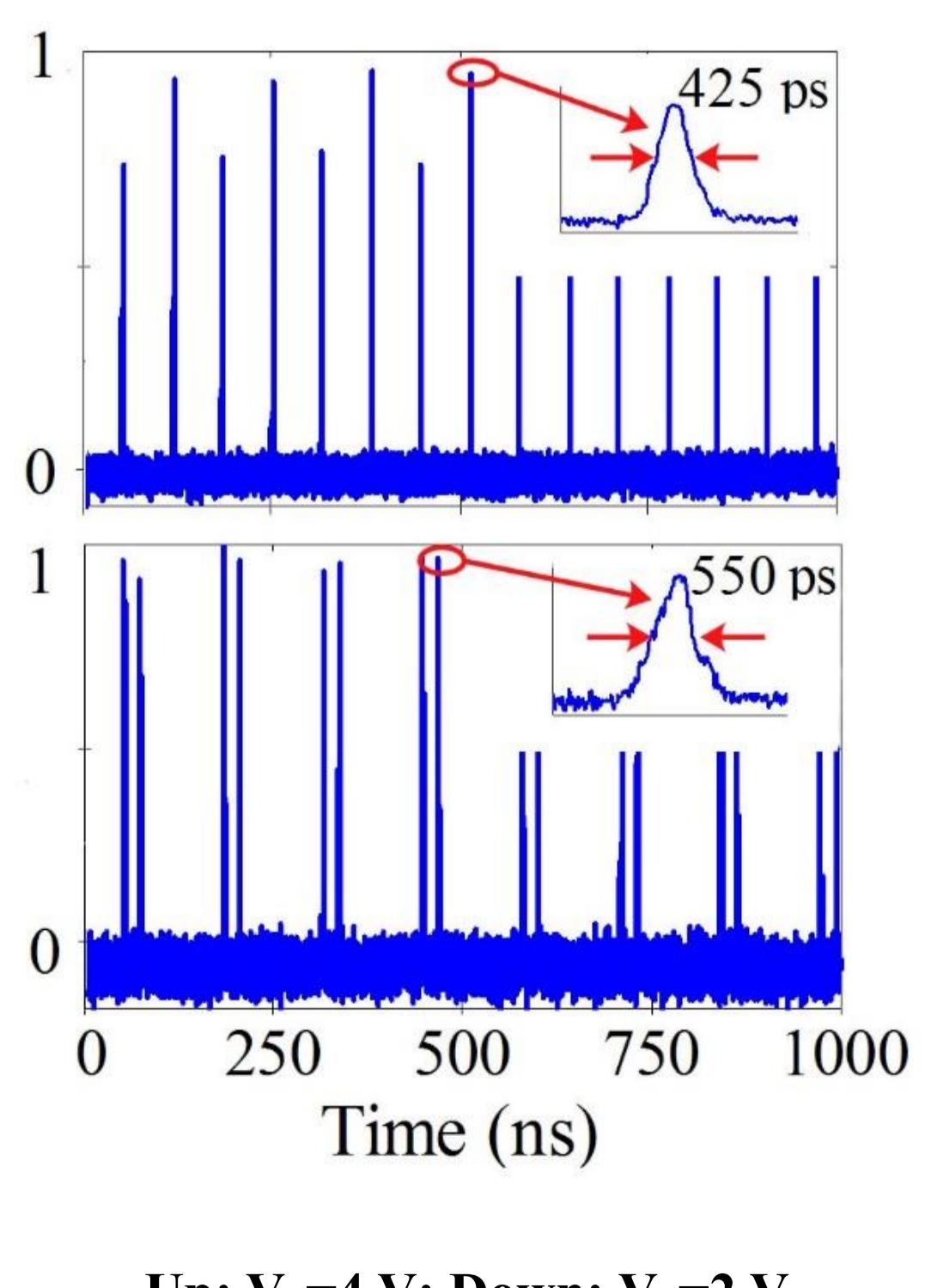


Experimental results



(H. Yang, Optics Express, submitted)

Mode-locked Double-pulse regime



Conclusion

- Observation of original double-pulse regime
- Theoretical model in agreement with experiments
- Repetition-rate tunability from MHz to GHz
- frequency comb width up to 40 GHz
- waveform generation: rectangle, sawtooth,...

Forthcoming

- Phase modulation vs intensity modulation
- Frequency-to-time mapping (as in [6] with AOFS)
- Repetition-rate tunability from MHz to GHz
- Fractional Talbot condition $f_m = p/q f_c$: towards high repletion rates

References

- [1] H.G. de Chatellus *et al*, Phys. Rev. A 2013
[2] C. Schnébelin *et al*, CLEO Techn. Digest 2017
[3] H. Yang *et al*, IEEE Photon. J. 2017
[4] L. Wang *et al*, CLEO Techn. Digest 2017
[5] F. Tian, IEEE J. Lightwave technol. 2011
[6] H.G. de Chatellus *et al*, Optica 2016