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# $L^{q}$-weak solutions to the time-dependent Oseen system: decay estimates. 

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#### Abstract

This article deals with $L^{q}$-weak solutions to the 3D time-dependent Oseen system. This type of solution is defined in terms of the velocity only. It is shown that the velocity may be represented by a sum of integrals none of which involves the pressure and without a surface integral of the spatial gradient of the velocity. On the basis of this representation formula, an estimate of the spatial decay of the velocity and its spatial gradient is derived. No boundary conditions need to be imposed for these results.


AMS subject classifications. 35Q30, 65N30, 76D05.
Key words. Oseen system, weak solutions, integral representation, spatial decay.

## 1 Introduction

In this work, we deal with the time-dependent Oseen system

$$
\begin{equation*}
u_{t}-\Delta_{x} u+\tau \partial x_{1} u+\nabla_{x} \pi=f, \quad \operatorname{div}_{x} u=0 \quad \text { in } \bar{\Omega}^{c} \times\left(0, T_{0}\right), \tag{1.1}
\end{equation*}
$$

with the exterior domain $\bar{\Omega}^{c}$ defined by $\bar{\Omega}^{c}:=\mathbb{R}^{3} \backslash \bar{\Omega}$, where $\Omega$ is an open, bounded set in $\mathbb{R}^{3}$ with Lipschitz boundary and connected complement. System (1.1) arises as a linearization of the time-dependent Navier-Stokes system with Oseen term,

$$
\begin{equation*}
u_{t}-\Delta_{x} u+\tau \partial x_{1} u+\left(u \cdot \nabla_{x}\right) u+\nabla_{x} \pi=f, \quad \operatorname{div}_{x} u=0 \quad \text { in } \bar{\Omega}^{c} \times\left(0, T_{0}\right) . \tag{1.2}
\end{equation*}
$$

The latter system is the usual model of the flow of a viscous incompressible fluid around a rigid body moving with constant velocity and without rotation. The parameter $\tau \in(0, \infty)$ corresponds to the Reynolds number of the fluid, and the function $f: \bar{\Omega}^{c} \times\left(0, T_{0}\right) \mapsto \mathbb{R}^{3}$ represents a volume force acting on the fluid. The unknowns in (1.1) are the velocity $u: \bar{\Omega}^{c} \times\left(0, T_{0}\right) \mapsto \mathbb{R}^{3}$ and the pressure $\pi: \bar{\Omega}^{c} \times\left(0, T_{0}\right) \mapsto \mathbb{R}$, hence $T_{0} \in(0, \infty]$ is the life-span of the solution. We impose the initial condition $u(0)=U_{0}$ in $\bar{\Omega}^{c}$, where the initial velocity $U_{0}: \bar{\Omega}^{c} \mapsto \mathbb{R}^{3}$ is given as well.

It was shown in [19] (see [19, Theorem 6.1] and the remark following it) that if $u$ is the velocity part of a regular solution to (1.1) as specified in [19, Corollary 5.2 and Theorem 6.1], and if $|f(x, t)|$ and $|u(x, 0)|$ decrease sufficiently fast for $|x| \rightarrow \infty$, then the estimate

$$
\begin{equation*}
\left|\partial_{x}^{\alpha} u(x, t)\right| \leq \mathfrak{C}\left[(|x| \nu(x))^{-(3+|\alpha|) / 2+1 /\left(2 \max \left\{\gamma_{1}^{\prime}, \gamma_{2}^{\prime}, \gamma_{3}^{\prime}\right\}\right)}+|x|^{-\gamma-|\alpha|}\right] \tag{1.3}
\end{equation*}
$$

holds for $x \in B_{R_{0}}^{c}:=\mathbb{R}^{3} \backslash B_{R_{0}}, t \in\left(0, T_{0}\right), \alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha|:=\alpha_{1}+\alpha_{2}+\alpha_{3} \leq 1$. The parameter $R_{0}$ is some fixed positive real, sufficiently large so that $\bar{\Omega} \subset B_{R_{0}}$. The function $\nu: \mathbb{R}^{3} \mapsto[1, \infty)$ is defined by $\nu(x):=1+|x|-x_{1} \quad$ for $x \in \mathbb{R}^{3}$. Its presence in (1.3) should
be interpreted as a mathematical manifestation of the wake extending downstream behind the rigid body. The condition $|\alpha| \leq 1$ expresses the fact that inequality (1.3) yields an upper bound for the velocity $u$ itself as well as for its spatial gradient $\nabla_{x} u$. The number $\gamma$ on the right-hand side of (1.3) equals 3 if the zero flux condition

$$
\begin{equation*}
\int_{\partial \Omega} u(t) \cdot n^{(\Omega)} d x=0 \quad\left(t \in\left(0, T_{0}\right)\right) \tag{1.4}
\end{equation*}
$$

is fulfilled, with $n^{(\Omega)}$ denoting the outward unit normal to $\Omega$. Otherwise $\gamma$ takes the value 2. The parameters $\gamma_{1}, \gamma_{2}, \gamma_{3} \in[0, \infty], q \in(1, \infty)$ in (1.3) are introduced via the assumptions $u\left|Z_{R_{0}, T_{0}} \in L^{\gamma_{1}}\left(0, T_{0}, L^{q}\left(\Omega_{R_{0}}\right)^{3}\right), \nabla_{x} u\right| Z_{R_{0}, T_{0}} \in L^{\gamma_{2}}\left(0, T_{0}, L^{q}\left(\Omega_{R_{0}}\right)^{3}\right)$ and $f \mid Z_{R_{0}, T_{0}} \in L^{\gamma_{3}}\left(0, T_{0}, L^{q}\left(\Omega_{R_{0}}\right)^{3}\right)$, where $\Omega_{R_{0}}:=B_{R_{0}} \backslash \bar{\Omega}$ and $Z_{R_{0}, T_{0}}:=\Omega_{R_{0}} \times(0, \infty)$. The appearance of these parameters $\gamma_{j}$ in (1.3) means that the spatial decay of $u$ and $\nabla_{x} u$ depends on $L^{p}$-integrability in time of $u\left|Z_{R_{0}, T_{0}}, \nabla_{x} u\right| Z_{R_{0}, T_{0}}$ and $f \mid Z_{R_{0}, T_{0}}$. Another important feature of the theory in [19] is that no boundary conditions are imposed on $u$ and $\pi$.

The work at hand aims to extend the decay estimate (1.3) to $L^{q}$-weak solutions of (1.1). Solutions of this type verify a weak form (equation (5.1)) of (1.1) involving only the velocity $u$, which is supposed to be continuous as a function of $t \in\left[0, T_{0}\right)$ with values in $L^{q_{0}}\left(\bar{\Omega}^{c}\right)^{3}$, and satisfy the relation $\nabla_{x} u \in L_{l o c}^{1}\left(\left[0, T_{0}\right), L^{q_{1}}\left(\bar{\Omega}^{9}\right)\right.$, for some numbers $q_{0}, q_{1} \in(1, \infty)$. In addition, $u$ and $f$ are required to fulfill the conditions involving the parameters $\gamma_{j}$ as stated above. Actually our assumptions are somewhat more general (see the beginning of Section 5 and Theorem 5.2), but the preceding conditions correspond to what we essentially have in mind when we use the term " $L^{q}$-weak solutions". Note that in particular nothing is supposed on the pressure. This point is the main difficulty of the theory presented in the work at hand. The greatest part of our proof of (1.3) consists in establishing the representation formula (5.24) for $L^{q}$-weak solutions (Section 5). This formula should be considered as one of the main results of this work. In particular it is an essential element of the theory developed in [20] on the nonlinear problem (1.6); see the remarks further below. Inequality (1.3), for its part, is the starting point of the theory in [21], where we study the spatial asymptotics of mild solutions to (1.1) under homogeneous Dirichlet boundary conditions.

Let us compare the results of the work at hand with what is already available in literature. Spatial decay of $L^{2}$-weak solutions to (1.1) with Dirichlet boundary conditions was studied in [11] and [14], where these solutions were shown to satisfy the estimate

$$
\begin{equation*}
\left|\partial_{x}^{\alpha} u(x, t)\right| \leq \mathfrak{C}(|x| \nu(x))^{-1-|\alpha| / 2} \tag{1.5}
\end{equation*}
$$

for $x, t, \alpha$ as in (1.3), if $f$ and $U_{0}$ decrease sufficiently fast. Inequality (1.3) applied in the situation considered in [11] and [14] yields the stronger decay bound $(|x| \nu(x))^{-5 / 4-|\alpha| / 2}$ even though we do no longer impose a specific boundary condition. Again this result is valid under the usual condition that the data $f$ and $U_{0}$ decrease in an appropriate way These points are discussed in more detail in Section 6. In [15], we considered the nonlinear stability problem

$$
\begin{align*}
& \partial_{t} u-\Delta_{x} u+\tau \partial_{x_{1}} u+\tau\left(U \cdot \nabla_{x}\right) u+\tau(u \cdot \nabla) U+\tau\left(u \cdot \nabla_{x}\right) u+\nabla_{x} \pi=f  \tag{1.6}\\
& \operatorname{div}_{x} u=0 \quad \text { in } \bar{\Omega}^{c} \times\left(0, T_{0}\right)
\end{align*}
$$

where $U: \bar{\Omega}^{c} \mapsto \mathbb{R}^{3}$ is the velocity part of a solution to the stationary Navier-Stokes system, and is considered as given. It was shown in that reference that inequality (1.5) remains valid for $L^{2}$-strong solutions to (1.6) satisfying Dirichlet boundary conditions, again if the asymptotic behaviour of $f$ and $U_{0}$ allows this rate of decay. In the successor paper [20] mentioned above, we will show that (1.3) with $p=2$ is valid for these solution, so that the decay bound $\mathfrak{C}(|x| \nu(x))^{-1-|\alpha| / 2}$ from [15] may be replaced by $\mathfrak{C}(|x| \nu(x))^{-5 / 4-|\alpha| / 2}$, under the usual caveat on $f$ and $U_{0}$. No specific boundary conditions will be imposed in [20].
A remark is in order with respect to our proof of (1.3) for $L^{q}$-weak solutions. As already indicated, our main effort consists in showing that the integral representation (5.24) of the velocity is valid. This representation is a slightly modified version of equation [19, (5.7)] proved in [19] for regular solutions. Since all the integrals appearing in (5.24) were already estimated in [19], we obtain (1.3) almost immediately once this representation is available (Theorem 5.2). A major tool in the proof of (5.24) is Friedrich's mollifier for functions with values in Banach spaces, smoothing weak solutions with respect to the time variable. We will consider these smoothed solutions as weak solutions of the stationary Oseen system, with the time derivative subsumed into the right-hand side. In this way we will be able to use the regularity theory of this latter system in order to construct an associated pressure. Once this result is available, the mollified version of the $L^{q}$-weak solution we started out with will turn out to be a sufficiently regular solution of (1.1) so that we may apply formula $[19,(5.7)]$ (Theorem 5.1). By letting certain parameters tend to zero in this formula, we will then obtain that (5.24) holds for the original weak solution; see Lemma 5.3 and Corollary 5.1. The mollifying precedure requires that $T_{0}=\infty$. However, once (5.24) has been proved in the case $T_{0}=\infty$, it is not difficult to handle the case $T_{0}<\infty$ as well (Corollary 5.2).
Let us mention some references more distantly related to the work at hand. Knightly [33] considered even the case that the velocity of the rigid body changes with time. However, his results are valid only under various smallness assumptions. Mizumachi [37] proved (1.5) with $\alpha=0, f=0$ for a certain class of solutions to the nonlinear system (1.2). Takahashi [43] deals with spatial decay of solutions to the Navier-Stokes system with Oseen term in the case $\Omega=\emptyset$ under a smallness condition. In [3], [4], solutions to (1.1) and (1.2) are estimated in weighted $L^{p}$-norms, with the weights adapted to the wake in the flow field downstream to the rigid body. Reference [18] by the present author combines decay estimates in time and in space, for solutions of (1.1) and (1.6), as a continuation of [14] (Oseen system (1.1)) and [15] (problem (1.6)), under the same assumptions and with the same methods as in these articles. Various technical aspects of the theory in [11], [14], [15] and [18] are dealt with in predecessor papers [7] - [10], [12], [13]. Questions of existence, regularity and stability related to (1.1), (1.2) or (1.6) are addressed in [25], [26], [27], [30], [31], [34], [35], [36], [40], [42].

## 2 Notation. Some auxiliary results.

The parameters $T_{0} \in(0, \infty]$ and $\tau \in(0, \infty)$ introduced in Section 1 are kept fixed throughout, as is the open, bounded set $\Omega \subset \mathbb{R}^{3}$ with Lipschitz boundary.

The symbol $\left|\mid\right.$ denotes the Euclidean norm of $\mathbb{R}^{n}$ for any $n \in \mathbb{N}$, as well as the length $\alpha_{1}+\alpha_{2}+\alpha_{3}$ of a multi-index $\alpha \in \mathbb{N}_{0}^{3}$. For $R \in(0, \infty), x \in \mathbb{R}^{3}$, put $B_{R}(x):=\left\{y \in \mathbb{R}^{3}\right.$ : $|x-y|<R\}$. In the case $x=0$, we write $B_{R}$ instead of $B_{R}(0)$.
Recall that in Section 1, we introduced the function $\nu: \mathbb{R}^{3} \mapsto[1, \infty)$ by setting $\nu(x):=$ $1+|x|-x_{1}$ for $x \in \mathbb{R}^{3}$.
We fix numbers $S_{0}, R_{0} \in(0, \infty)$ with $S_{0}<R_{0}$ and $\bar{\Omega} \subset B_{S_{0}}$, as well as a function $\varphi_{0} \in$ $C_{0}^{\infty}\left(B_{\left(R_{0}+S_{0}\right) / 2}\right)$ with $0 \leq \varphi_{0} \leq 1$ and $\varphi_{0} \mid B_{S_{0}+\left(R_{0}-S_{0}\right) / 4}=1$. We put $R_{1}:=\left(R_{0}+S_{0}\right) / 2$.
For $n \in \mathbb{N}, I \subset \mathbb{R}^{n}$, let $\chi_{I}$ stand for the characteristic function of $I$ in $\mathbb{R}^{n}$. If $A \subset \mathbb{R}^{3}$, we denote by $A^{c}$ the complement $\mathbb{R}^{3} \backslash A$ of $A$ in $\mathbb{R}^{3}$. Put $e_{l}:=\left(\delta_{j l}\right)_{1 \leq j \leq 3}$ for $1 \leq l \leq 3$ (unit vector in $\mathbb{R}^{3}$ ). If $A$ is an open bounded set in $\mathbb{R}^{3}$ with Lipschitz boundary, we write $n^{(A)}$ for the outward unit normal to $A$. If $R, S \in(0, \infty)$ with $S<R$, we write $A_{R, S}$ for the annular domain $B_{R} \backslash \overline{B_{S}}$.
Let $p \in[1, \infty)$ and $n \in \mathbb{N}$. For any open set $A \subset \mathbb{R}^{3}$, the norm of the Lebesgue space $L^{p}(A)$ is denoted by $\left\|\|_{p}\right.$, and the usual norm of the Sobolev space $W^{m, p}(A)$ of order $m$ and exponent $p$ is designated by $\left\|\|_{m, p}\right.$. Again for an open set $A \subset \mathbb{R}^{3}$, we define $C_{0, \sigma}^{\infty}(A):=\left\{V \in C_{0}^{\infty}(A)^{3}: \operatorname{div} V=0\right\}$, and write $L_{l o c}^{p}(A)$ and $W_{l o c}^{m, p}(A)$ for the set of all functions $V$ from $A$ into $\mathbb{R}$ such that $V \mid K \in L^{p}(K)$ and $V \mid K \in W^{1, p}(K)$, respectively, for any open, bounded set $K \subset \mathbb{R}^{3}$ with $\bar{K} \subset A$. We put $\nabla V:=\left(\partial_{k} V_{j}\right)_{1 \leq j, k \leq 3}$ for $V \in W_{l o c}^{1,1}(A)^{3}$.
Let $\mathcal{V}$ a normed space, and let the norm of $\mathcal{V}$ be denoted by $\|\|$. Take $n \in \mathbb{N}$. Then we will use the same notation $\left\|\|\right.$ for the norm on $\mathcal{V}^{n}$ defined by $\|\left(f_{1}, \ldots, f_{n}\right) \|:=\left(\sum_{j=1}^{n}\left\|f_{j}\right\|^{2}\right)^{1 / 2}$ for $\left(f_{1}, \ldots, f_{n}\right) \in \mathcal{V}^{n}$. The space $\mathcal{V}^{3 \times 3}$, as concerns its norm, is identified with $\mathcal{V}^{9}$. If $p \in(1, \infty), n \in\{1,3\}$ and $A \subset \mathbb{R}^{3}$ open, the dual space of $W_{0}^{1, p^{\prime}}(A)^{n}$ will be denoted by $W_{0}^{-1, p}(A)^{n}$ (although in the case $n=3$ this notation is not coherent with the usual custom of letting $\mathfrak{A}^{n}$ stand for the Cartesion product of a given set $\left.\mathfrak{A}\right)$.
We additionally introduce some weighted Sobolev spaces. To this end, for $x \in \mathbb{R}^{3}$ we put

$$
\begin{aligned}
& \omega_{p}^{(j)}(x):=(1+|x|)^{-j} \text { for } j \in\{1,2\}, \quad p \in(1, \infty) \backslash\{3 / 2,3\}, \\
& \omega_{3 / 2}^{(1)}(x):=(1+|x|)^{-1}, \omega_{3 / 2}^{(2)}(x):=(1+|x|)^{-2}(\ln (2+|x|))^{-1}, \\
& \omega_{3}^{(j)}(x):=(1+|x|)^{-j}(\ln (2+|x|))^{-1} \text { for } j \in\{1,2\} .
\end{aligned}
$$

Then we set

$$
\begin{aligned}
& W_{2}^{2, p}\left(\mathbb{R}^{3}\right):=\left\{V \in W_{l o c}^{2,1}\left(\mathbb{R}^{3}\right): \omega_{p}^{(2)} V, \omega_{p}^{(1)} \partial_{l} V, \partial_{l} \partial_{k} V \in L^{p}\left(\mathbb{R}^{3}\right) \text { for } 1 \leq k, l \leq 3\right\}, \\
& W_{1}^{1, p}\left(\mathbb{R}^{3}\right):=\left\{V \in W_{\text {loc }}^{1,1}\left(\mathbb{R}^{3}\right): \omega_{p}^{(1)} V, \partial_{l} V \in L^{p}\left(\mathbb{R}^{3}\right) \text { for } 1 \leq l \leq 3\right\} \quad(p \in(1, \infty)) .
\end{aligned}
$$

We will not work with a norm of $W_{2}^{2, p}\left(\mathbb{R}^{3}\right)$. However, the norm of $W_{1}^{1, p}\left(\mathbb{R}^{3}\right)$ defined by $\|V\|_{W_{1}^{1, p}\left(\mathbb{R}^{3}\right)}:=\left(\left\|\omega_{p}^{(1)} V\right\|_{p}^{p}+\|\nabla V\|_{p}^{p}\right)^{1 / p}$ will be relevant.
Let $p \in[1, \infty], \mathfrak{B}$ a Banach space and $J \subset \mathbb{R}$ an interval. Then the norm of the space $L^{p}(J, \mathfrak{B})$ is denoted by $\left\|\|_{L^{p}(J, \mathcal{B})}\right.$. Let $a, b \in \mathbb{R} \cup\{\infty\}$ with $a<b$, take $m \in \mathbb{N}$ and $q \in[1, \infty)$. Then we write $L^{p}(a, b, \mathcal{B})$ and $W^{1, q}(a, b, \mathcal{B})$ instead of $L^{p}((a, b), \mathcal{B})$ and $W^{1, q}((a, b), \mathcal{B})$, respectively. We use the expression $L_{l o c}^{p}([a, b), \mathcal{B})$ for the space of all functions $v:(a, b) \mapsto$
$\mathcal{B}$ such that $v \mid(a, T) \in L^{p}(a, T, \mathcal{B})$ for any $T \in(a, b)$. This space is to be distinguished from the space $L_{l o c}^{p}(a, b, \mathcal{B})$, defined in the usual way. Let $T \in(0, \infty], A \subset \mathbb{R}^{3}$ open, and $n \in \mathbb{N}$. Then we will write $\left\|\left\|\|_{q, p ; T} \text { and }\right\|\right\|_{q, p ; \mathbb{R}}$ instead of $\left\|\|_{L^{p}\left(0, T, L^{q}(A)^{n}\right)}\right.$ and $\| \|_{L^{p}\left(\mathbb{R}, L^{q}(A)^{n}\right)}$, respectively.
If $v \in W^{1,1}(a, b, \mathcal{B})$, then, possibly after a modification on a subset of $[a, b)$ with measure zero, the function $v$ belongs to $C^{0}([a, b), \mathcal{B})([44$, Lemma 3.1.1]). If the latter relation is already valid, we write $v \in W^{1,1}(a, b, \mathcal{B}) \cap C^{0}([a, b), \mathcal{B})$.
Of course, a function $v \in L^{p}\left(J, L^{q}(A)^{n}\right)$ may be considered also as a function on $A \times J$, although there is a minor issue with respect to measurability on $A \times J$, settled in [12, Lemma 2.1] and [19, Lemma 2.3]. We will write $v(t)(x)$ or $v(x, t)$, depending on whether we consider $v$ as a function on $J$ with values in $L^{q}(A)^{n}$, or as a function on $A \times J$. For an interval $J \subset \mathbb{R}$ and a function $v: J \mapsto W_{l o c}^{1,1}(A)^{3}$, the notation $\nabla_{x} v$ stands for the gradient of $v$ with respect to $x \in A$, in the sense that

$$
\nabla_{x} v: J \mapsto L_{l o c}^{1}(A)^{3 \times 3}, \nabla_{x} v(t)(x):=\left(\partial x_{k}\left(v_{j}(t)\right)(x)\right)_{1 \leq j, k \leq 3} \text { for } t \in J, x \in A
$$

(spatial gradient of $v$ ). Similar conventions are to be valid with respect to the expressions $\Delta_{x} v, \operatorname{div}_{x} v$ and $\partial x_{j} v$.
Concerning Bochner integrals, if $J \subset \mathbb{R}$ is open, $\mathcal{B}$ a Banach space and $w: J \mapsto \mathcal{B}$ an integrable function, it is convenient sometimes to write $\mathcal{B}-\int_{J} w(t) d t$ instead of $\int_{J} w(t) d t$ for the corresponding $\mathcal{B}$-valued Bochner integral. For the definition of the Bochner integral, we refer to [46, p. 132-133], or to [32, p. 78-80.].
We define the Fourier transform $\hat{f}$ of $f \in L^{1}\left(\mathbb{R}^{3}\right)$ by $\hat{f}(\xi):=(2 \pi)^{-3 / 2} \int_{\mathbb{R}^{n}} e^{-i \xi \cdot z} f(z) d z$ for $\xi \in \mathbb{R}^{n}$. An analogous definition is to hold for functions belonging to $L^{2}\left(\mathbb{R}^{3}\right)$.
We write $C$ for numerical constants and $C\left(\gamma_{1}, \ldots, \gamma_{n}\right)$ for constants depending exclusively on paremeters $\gamma_{1}, \ldots, \gamma_{n} \in[0, \infty)$ for some $n \in \mathbb{N}$. However, such a precise bookkeeping will be possible only at some places. Mostly we will use the symbol $\mathfrak{C}$ for constants whose dependence on parameters must be traced from context. Sometimes we write $\mathfrak{C}\left(\gamma_{1}, \ldots, \gamma_{n}\right)$ in order to indicate that the constants in question is influenced by the quantities $\gamma_{1}, \ldots, \gamma_{n}$. But in such cases, this constant depends on other parameters as well.
The following simple version of Young's inequality for integrals will be used frequently. Stated her for the convenience of the reader, we will refer to it as "Young's inequality".
Lemma 2.1 ([1, Corollary 2.25]) Let $n \in \mathbb{N}$ and $q \in[1, \infty]$. Then

$$
\left(\int_{\mathbb{R}^{n}}\left|\int_{\mathbb{R}^{n}} U(x-y) V(y) d y\right|^{q} d x\right)^{1 / q} \leq C\|U\|_{1}\|V\|_{q} \quad \text { for } U \in L^{1}\left(\mathbb{R}^{n}\right), V \in L^{q}\left(\mathbb{R}^{n}\right)
$$

We will use Minkowski's inequality for integrals, which we restate, too.
Theorem 2.1 ([1, Theorem 2.9]) Let $m, n \in \mathbb{N}, p \in[1, \infty), F: \mathbb{R}^{n} \times \mathbb{R}^{m} \mapsto \mathbb{R} a$ measurable function. Then

$$
\left(\int_{\mathbb{R}^{n}}\left(\int_{\mathbb{R}^{m}}|F(x, y)| d y\right)^{p} d x\right)^{1 / p} \leq \int_{\mathbb{R}^{m}}\left(\int_{\mathbb{R}^{n}}|F(x, y)|^{p} d x\right)^{1 / p} d y
$$

The next theorem deals with solenoidal $W_{0}^{1, q}$-functions.

Theorem 2.2 ([29, Theorem III.4.2, III.6.1]) Let $n \in \mathbb{N}, q, r_{1}, \ldots, r_{n} \in(1, \infty), A \subset$ $\mathbb{R}^{3}$ open, bounded and with Lipschitz boundary. Let $V \in W_{0}^{1, q}\left(\bar{A}^{c}\right)^{3} \cap L^{r_{j}}\left(\bar{A}^{c}\right)^{3}$ for $1 \leq j \leq n$ with div $V=0$. Then there is a sequence $\left(\vartheta_{n}\right)$ in $C_{0, \sigma}^{\infty}\left(\bar{A}^{c}\right)$ such that $\left\|V-\vartheta_{n}\right\|_{1, q} \rightarrow 0$ and $\left\|V-\vartheta_{n}\right\|_{r_{j}} \rightarrow 0(n \rightarrow \infty)$ for $1 \leq j \leq n$.
The ensuing theorem presents a result on $L^{p}$-integrability of functions defined in an exterior domain and possessing an $L^{q}$-integrable gradient.

Theorem 2.3 Let $A \subset \mathbb{R}^{3}$ be open, bounded and with Lipschitz boundary. Let $q \in(1,3)$ and $V \in W_{l o c}^{1,1}\left(\bar{A}^{c}\right)$ with $\nabla V \in L^{q}\left(\bar{A}^{c}\right)^{3}$. Suppose there is some $\kappa \in(1, \infty)$ with $V \in L^{\kappa}\left(\bar{A}^{c}\right)$. Then $V \in L^{3 q /(3-q)}\left(\bar{A}^{c}\right)$ and $\|V\|_{3 q /(3-q)} \leq \mathfrak{C}\|V\|_{q}$.
Proof: This theorem may be deduced from [29, Theorem II.6.1]; see [17, Theorem 2.4] and its proof.

We mention some results about Bochner's integral. Our basic tool in this context is the following theorem (compatibility of bounded operators and Bochner integrals).

Theorem 2.4 Let $B_{1}, B_{2}$ be Banach spaces, $A: B_{1} \mapsto B_{2}$ a linear and bounded operator, $n \in \mathbb{N}, J \subset \mathbb{R}^{n}$ an open set and $f: J \mapsto B_{1}$ a Bochner integrable mapping. Then $A \circ f: J \mapsto B_{2}$ is Bochner integrable, too, and $A\left(B_{1}-\int_{J} f d x\right)=B_{2}-\int_{J} A \circ f d x$.
Proof: See [46, p. 134, Corollary 2], [32, Theorem 3.7.12 and the remark on p. 84].
Next we indicate a compatibility result for Bochner integrals with values in $L^{p}$-spaces.
Lemma 2.2 Let $J \subset \mathbb{R}$ be an interval, $n \in \mathbb{N}, B \subset \mathbb{R}^{n}$ and $A \subset B$ open sets, $q_{1}, q_{2} \in$ $[1, \infty)$ and $f: J \mapsto L^{q_{1}}(B)^{3}$ a Bochner integrable mapping with $f(t) \mid A \in L^{q_{2}}(A)^{3}$ for $t \in J$ and $f \mid A: J \mapsto L^{q_{2}}(A)^{3}$ Bochner integrable as well. Then $\left(L^{q_{1}}(B)^{3}-\int_{J} f(s) d s\right) \mid A=$ $L^{q_{2}}(A)^{3}-\int_{J} f(s) \mid A d s$.

Proof: According to [19, Lemma 2.3], for a. e. $x \in A$, the integral on the left-hand side of the equation at the end of the lemma taken at the point $x$ equals the integral on the right-hand side taken at that point.
Theorem 2.5 Let $B$ be a Banach space, $n \in \mathbb{N}, p \in[1, \infty)$ and $J \subset \mathbb{R}^{n}$ measurable. Then the set of integrable functions from $J$ into $B$ only taking a finite number of values ("simple functions") is dense in $L^{p}(J, B)$.
Proof: See [24, Section 8.18 .1 and Exercise 8.29]. On the basis of Lebesgue's theorem, the proof can be done by first approximating $f \in L^{p}(J, B)$ by functions with bounded support, and then by functions with bounded support and such that the set of their values is bounded with respect to the norm of $B$. Functions of the latter kind belong to $L^{1}(J, B)$ and thus, by the definition of Bochner's integral, may be approximated in $L^{1}(J, B)$ by simple functions, which implies approximation in $L^{p}(J, B)$ in this situation, due to Lebesgue's theorem.
Corollary 2.1 Let $B$ be a Banach space, $A$ a dense subset of $B, p \in[1, \infty)$ and $J \subset \mathbb{R}^{n}$ open. Then the set $\left\{\sum_{j=1}^{k} \varphi_{j} a_{j}: k \in \mathbb{N}, \varphi_{j} \in C_{0}^{\infty}(J), a_{j} \in A\right.$ for $\left.1 \leq j \leq n\right\}$ is dense in $L^{p}(J, B)$. In particular the set of continuous functions $f: J \mapsto B$ with supp $(f)$ compact is dense in $L^{p}(J, B)$.

Proof: Use Theorem 2.5 and the density of $C_{0}^{\infty}(J)$ in $L^{p}(J)$.

In order to define Friedrich's mollifier for functions with values in Banach spaces, we fix a function $\rho \in C_{0}^{\infty}((-1,1))$ with $\rho \geq 0$ and $\int_{\mathbb{R}} \rho(s) d s=1$, and put $\rho_{\delta}(r):=\delta^{-1} \rho\left(\delta^{-1} r\right)$ for $\delta \in(0, \infty), r \in \mathbb{R}$. If B is a Banach space and $f \in L_{l o c}^{1}(\mathbb{R}, B)$, define $f_{\delta}(t):=$ $B-\int_{\mathbb{R}} \rho_{\delta}(t-s) f(s) d s$ for $t \in \mathbb{R}, \delta \in(0, \infty)$.
Key properties of Friedrich's mollifier of functions with values in $\mathbb{R}$ carry over to functions with values in Banach spaces. Properties of this type needed in the work at hand are collected in the ensuing Theorem 2.6.

Theorem 2.6 Let $B$ be a Banach space and $f \in L_{l o c}^{1}(\mathbb{R}, B)$. Then $f_{\delta} \in C^{\infty}(\mathbb{R}, B)$ and $f_{\delta}^{(n)}(t)=\int_{\mathbb{R}} \rho_{\delta}^{(n)}(t-s) f(s) d s(n \in \mathbb{N}, t \in \mathbb{R}, \delta \in(0, \infty))$. Moreover, if $f \in W_{l o c}^{1,1}(\mathbb{R}, B)$, then $\left(f_{\delta}\right)^{\prime}=\left(f^{\prime}\right)_{\delta}$.
Let $g \in L^{p}(\mathbb{R}, B)$. Then $\left\|g_{\delta}\right\|_{L^{p}(\mathbb{R}, B)} \leq\|g\|_{L^{p}(\mathbb{R}, B)}$ for $\delta \in(0, \infty)$ and $\left\|g_{\delta}-g\right\|_{L^{p}(\mathbb{R}, B)} \rightarrow$ $0(\delta \downarrow 0)$.
Let $h \in C^{0}(\mathbb{R}, B)$ and $t \in \mathbb{R}$. Then $\left\|\left(h_{\delta}-h\right)(t)\right\| \rightarrow 0(\delta \downarrow 0)$, where $\|\|$ denotes the norm of $B$.
Proof: The relation $f_{\delta} \in C^{\infty}(\mathbb{R}, B)$ and the equation for $f_{\delta}^{(n)}(t)$ follow from the assumption $\rho_{\delta} \in C_{0}^{\infty}((-\delta, \delta))$. If $f \in W_{l o c}^{1,1}(\mathbb{R}, B)$ and $t \in \mathbb{R}$, the equation $\left(f_{\delta}\right)^{\prime}(t)=\left(f^{\prime}\right)_{\delta}(t)$ holds due to the above equation for $f_{\delta}^{(n)}(t)$ with $n=1$, and since the function $s \mapsto$ $\rho_{\delta}(t-s)(s \in \mathbb{R})$ belongs to $C_{0}^{\infty}(\mathbb{R})$. The inequality $\left\|g_{\delta}\right\|_{L^{p}(\mathbb{R}, B)} \leq\|g\|_{L^{p}(\mathbb{R}, B)}$ for $\delta>0$ is an immediate consequence of Young's inequality and the choice of $\rho_{\delta}$. We further note that for $\epsilon \in(0, \infty)$, Corollary 2.1 yields existence of a function $g^{(\epsilon)} \in C^{0}(\mathbb{R}, B)$ with compact support such that $\left\|g-g^{(\epsilon)}\right\|_{L^{p}(\mathbb{R}, B)} \leq \epsilon / 2$. With this result available, the relation $\left\|g_{\delta}-g\right\|_{L^{p}(\mathbb{R}, B)} \rightarrow 0(\delta \downarrow 0)$ follows by the same arguments as in the case $B=\mathbb{R}$; see $[1, \mathrm{p}$. 37-38] for example. The same reference yields the last claim of the theorem.

## 3 Some results on the Poisson equation and the stationary Oseen system.

In the ensuing theorem, we state some properties of the Newton potential. The proof of this theorem is well known (Use Hardy-Littlewood-Sobolev's inequality, Calderon-Zygmund's inequality and density arguments.), although we cannot give a reference.
Theorem 3.1 Put $\mathfrak{N}(z):=(4 \pi|z|)^{-1}$ for $z \in \mathbb{R}^{3} \backslash\{0\}$ (fundamental solution of the Poisson equation). Let $q \in(1,3 / 2), F \in L^{q}\left(\mathbb{R}^{3}\right)$. Then $\int_{\mathbb{R}^{3}} \mathfrak{N}(x-y)|F(y)| d y<\infty$ for a. e. $x \in \mathbb{R}^{3}$. Put $(\mathfrak{N} * F)(x):=\int_{\mathbb{R}^{3}} \mathfrak{N}(x-y) F(y) d y$ for $x \in \mathbb{R}^{3}$ ("Newton potential"). Then $\mathfrak{N} * F \in W_{l o c}^{2, q}\left(\mathbb{R}^{3}\right), \Delta(\mathfrak{N} * F)=-F$ and $\nabla(\mathfrak{N} * F) \in L^{(1 / q-1 / 3)^{-1}}\left(\mathbb{R}^{3}\right)$. If $F \in W^{1, q}\left(\mathbb{R}^{3}\right)$, then $\mathfrak{N} * F \in W_{l o c}^{3, q}\left(\mathbb{R}^{3}\right)$ and $\left\|\partial_{k} \partial_{l} \partial_{m}(\mathfrak{N} * F)\right\|_{q}+\left\|\partial_{l} \partial_{m}(\mathfrak{N} * F)\right\|_{q} \leq C(q)\|F\|_{1, q}$ for $1 \leq k, l, m \leq 3$.
In the next theorem, we introduce a pressure $\Pi$ associated with the velocity part $U$ of a weak solution to the Oseen system $(\lambda=0)$ or the Oseen resolvent system $(\lambda \neq 0)$ in $\mathbb{R}^{3}$. The case $\lambda \neq 0$ is included in view of an application in [20].
Theorem 3.2 Let $A \subset \mathbb{R}^{3}$ be open, $\lambda \in \mathbb{C}, q \in(1, \infty), U \in W_{\text {loc }}^{1, q}(A)^{3}$ with

$$
\begin{equation*}
\int_{A}\left(\nabla U \cdot \nabla \vartheta+\left(\tau \partial_{1} U+\lambda U-F\right) \cdot \vartheta\right) d x=0 \text { for } \vartheta \in C_{0, \sigma}^{\infty}(A), \quad \operatorname{div} U=0 \tag{3.1}
\end{equation*}
$$

Then there is a function $\Pi \in L_{\text {loc }}^{q}(A)$ with

$$
\begin{equation*}
\int_{A}\left(\nabla U \cdot \nabla \vartheta+\left(\tau \partial_{1} U+\lambda U-F\right) \cdot \vartheta-\Pi \operatorname{div} \vartheta\right) d x=0\left(\vartheta \in C_{0}^{\infty}(A)^{3}\right), \operatorname{div} U=0 \tag{3.2}
\end{equation*}
$$

If $B \subset \mathbb{R}^{3}$ is open, bounded, with Lipschitz boundary, $\bar{B} \subset A$ and $\int_{B} \Pi=0$, then

$$
\begin{equation*}
\|\Pi\|_{q} \leq \mathfrak{C}\left(\|\nabla U\|_{q}+|\lambda|\|U\|_{q}+\|F\|_{q}\right) \tag{3.3}
\end{equation*}
$$

Proof: Let $B \subset \mathbb{R}^{3}$ be open and bounded, with $\bar{B} \subset A$. Since $F \in L_{l o c}^{q}(A)^{3}$, we have $F \mid B \in L^{q}(B)^{3}$, hence $F \mid B \in W_{0}^{-1, q}(B)$ by Poincaré's inequality. Obviously $-\tau \partial_{1} U \mid B \in$ $W_{0}^{-1, q}(B)^{3}$. Thus, by [27, Lemma IV.1.1], there is a function $\Pi \in L_{l o c}^{q}(A)$ such that (3.2) holds. The last statement of the theorem including (3.3) also follows from [27, Lemma IV.1.1] and Poincaré's inequality.

Theorem 3.3 Let $A \subset \mathbb{R}^{3}$ be open, $\lambda \in \mathbb{C}, q, s \in(1, \infty), F \in L_{l o c}^{q}(A)^{3}, U \in W_{l o c}^{1,1}(A)^{3}$ with $\nabla U \in L_{l o c}^{s}(A)^{9}$ such that (3.1) holds. Then $U \in W_{l o c}^{2, q}(A)^{3}$. Suppose in addition that $\Pi \in L_{l o c}^{s}(A)$ is such that the pair $(U, \Pi)$ satisfies (3.2). Then $\Pi \in W_{l o c}^{1, q}(A)$ and

$$
\begin{equation*}
-\Delta U+\tau \partial_{1} U+\lambda U+\nabla \Pi=F, \quad \operatorname{div} U=0 \tag{3.4}
\end{equation*}
$$

Proof: The theorem is a consequence of interior regularity of solutions to the Stokes system; see [17, Theorem 3.2] and its proof.
Theorem 3.4 Let $A \subset \mathbb{R}^{3}$ be open, $\lambda \in \mathbb{C}, F \in C^{\infty}(A)^{3}, q \in(1, \infty), U \in W_{\text {loc }}^{1,1}(A)^{3}$ with $\nabla U \in L_{l o c}^{q}(A)^{9}$, and $\Pi \in L_{\text {loc }}^{q}(A)$. Suppose that (3.2) holds. Then $U \in C^{\infty}(A)^{3}, \Pi \in$ $C^{\infty}(A)$, and (3.4) (Oseen system if $\lambda=0$ ) is valid.

Proof: [17, Corollary 3.2]. The assumption $v \in L^{q}(A)^{3}$ in that reference should read $v \in L_{l o c}^{q}(A)^{3}$.

We will need a solution theory for the Oseen system in weighted Sobolev spaces of functions defined in the whole space $\mathbb{R}^{3}$. This theory is useful for us because it holds in an $L^{q_{-}}$ framework with any $q \in(1, \infty)$ being admitted.
Theorem 3.5 Let $q \in(1, \infty)$. Then for any $F \in L^{q}\left(\mathbb{R}^{3}\right)^{3}$, there is a pair of functions $(U(F), \Pi(F))$ such that $U(F) \in W_{2}^{2, q}\left(\mathbb{R}^{3}\right)^{3}, \Pi(F) \in W_{1}^{1, q}\left(\mathbb{R}^{3}\right)$, the pair $(U(F), \Pi(F))$ solves (3.4) with $\lambda=0, \quad A=\mathbb{R}^{3}$, and $\left\|\partial_{l} \partial_{k} U(F)\right\|_{q}+\|\Pi\|_{W_{1}^{1, q}\left(\mathbb{R}^{3}\right)}+\left\|\partial_{1} U(F)\right\|_{q} \leq$ $C(q, \tau)\|F\|_{q}$ for $1 \leq k, l \leq 3$.
Proof: The theorem reproduces some of the statements of [2, Theorem 3.3].
In [17], we proved uniqueness of the velocity part $U$ of a weak solution to the Oseen system (equation (3.4) with $\lambda=0$ ) or to the Oseen resolvent problem (equation (3.4) with $\lambda \in \mathbb{C} \backslash\{0\}, \Re \lambda \geq 0)$ in the whole space $\mathbb{R}^{3}$, under the assumptions that $|\lambda| \leq(\tau / 2)^{2}$ and $U \mid B_{R}^{c} \in \sum_{j=1}^{3} L^{r_{j}}\left(B_{R}^{c}\right)^{3}, \quad \nabla V \in \sum_{j=1}^{3} L^{q_{j}}\left(\mathbb{R}^{3}\right)^{9}$ for some $r_{j}, q_{j} \in(1, \infty)(j \in\{1,2,3\})$ and some $R \in(0, \infty)$ ([17, Theorem 5.1]). In the following theorem, we generalize this result in the sense that we admit functions $U$ growing polynomially for $|x| \rightarrow \infty$. However, such weak solutions need not vanish, but they turn out to be polynomials. The case $\lambda \neq 0$ will not be needed in what follows, but it is included because it can be handled without additional effort.

Theorem 3.6 Let $U \in W_{\text {loc }}^{1,1}\left(\mathbb{R}^{3}\right)^{3}$ satisfy (3.1) with $A=\mathbb{R}^{3}, F=0$ and $\lambda \in \mathbb{C}$ with $\Re \lambda \geq$ $0,|\lambda| \leq(\tau / 2)^{2}$. Suppose there are numbers $R \in(0, \infty), r \in[0, \infty), k_{0}, m_{0} \in \mathbb{N}, \varrho_{k}, q_{m} \in$ $(1, \infty)$ for $1 \leq k \leq k_{0}, 1 \leq m \leq m_{0}$, and functions $U^{(k)} \in L_{l o c}^{\varrho_{k}}\left(\mathbb{R}^{3}\right)^{3}, V^{(m)} \in L_{l o c}^{q_{m}}\left(\mathbb{R}^{3}\right)^{3 \times 3}$ for $k, m$ as before such that $U=\sum_{k=1}^{k_{0}} U^{(k)}, \partial_{\mu} U_{\sigma}=\sum_{q_{m}}^{m_{0}} V_{\mu, \sigma}^{(m)}$ for $1 \leq \mu, \sigma \leq 3$, and $\int_{B_{R}^{c}}\left(\left|U^{(k)}(x)\right||x|^{-r}\right)^{\varrho_{k}} d x<\infty, \int_{B_{R}^{c}}^{k=1}\left(\left|V^{(m)}(x)\right||x|^{-r}\right)^{q_{m}} d x<\infty$ for $1 \leq k \leq k_{0}, 1 \leq$ $m \leq m_{0}$. Then $U$ is a polynomial.
Proof: The theorem follows by the standard theory of topological vector spaces ([38, Section $7.3,7.11]$ ), as used in the proof of [17, Theorem 5.1]. The only point that should still be checked in more detail is whether the operators $T, S: \mathcal{S}\left(\mathbb{R}^{3}\right)^{3} \mapsto \mathbb{R}$ defined by $T(\phi):=\int_{\mathbb{R}^{3}}\left(\nabla U \cdot \nabla \phi+\tau \partial_{1} U \cdot \phi\right) d x$ and $S(\phi):=\int_{\mathbb{R}^{3}} U \cdot \phi d x$ for $\phi \in \mathcal{S}\left(\mathbb{R}^{3}\right)^{3}$, respectively, are tempered distributions, where $\mathcal{S}\left(\mathbb{R}^{3}\right)$ denotes the set of rapidly decreasing functions in $\mathbb{R}^{3}$, equipped in the usual way with a topology. So let us show that $T$ is a tempered distribution. To that end, put $\bar{R}:=\max \{1, R\}, \bar{r}:=\min \{n \in \mathbb{N}: n \geq r\}, p_{\alpha, \beta}(\phi):=$ $\sup \left\{\left|x^{\alpha} \partial^{\beta} \phi(x)\right|: x \in \mathbb{R}^{3}\right\}, G(\phi):=\sup \left\{|x|^{r+3} \sum_{\gamma \in \mathbb{N}_{0}^{3},|\gamma| \leq 1}\left|\partial^{\gamma} \phi(x)\right|: x \in B \frac{c}{\bar{R}}\right\}$ for $\phi \in \mathcal{S}\left(\mathbb{R}^{3}\right)^{3}, \alpha, \beta \in \mathbb{N}_{0}^{3}$. Then we find for $\phi \in \mathcal{S}\left(\mathbb{R}^{3}\right)^{3}$ that

$$
G(\phi) \leq \mathfrak{C} \sup \left\{\sum_{\gamma \in \mathbb{N}_{0}^{3},|\gamma| \leq 1} \sum_{m=1}^{3}\left|x_{m}^{\bar{r}+3} \partial^{\gamma} \phi(x)\right|: x \in B \frac{c}{R}\right\}
$$

and thus $G(\phi) \leq \mathfrak{C} \sum_{\gamma \in \mathbb{N}_{0}^{3},|\gamma| \leq 1} \sum_{m=1}^{3} p_{(\bar{r}+3) e_{m}, \gamma}(\phi)$. On the other hand,

$$
\begin{aligned}
& \int_{B_{\frac{c}{R}}}|x|^{-r-3} \sum_{\mu=1}^{3}\left|\partial_{\mu} U(x)\right| d x \leq C \sum_{m=1}^{m_{0}} \sum_{\mu, \sigma=1}^{3} \int_{B_{\bar{R}}^{c}}|x|^{-r-3}\left|V_{\mu \sigma}^{(m)}(x)\right| d x \\
& \leq C \sum_{m=1}^{m_{0}} \sum_{\mu, \sigma=1}^{3}\left(\int_{B_{\bar{R}}^{c}}\left|V_{\mu \sigma}^{(m)}(x)\right|^{q_{m}}|x|^{-r q_{m}} d x\right)^{1 / q_{m}}\left(\int_{B_{\bar{R}}^{c}}|x|^{-3 q_{m}^{\prime}} d x\right)^{1 / q_{m}^{\prime}}
\end{aligned}
$$

Since $\int_{B_{\bar{c}}}|x|^{-3 q_{m}^{\prime}} d x<\infty$ for $1 \leq m \leq m_{0}$, our assumptions on the functions $V^{(m)}$ imply that the right-hand side of the preceding estimate is finite. Thus we may conclude that the integral $\int_{B \frac{c}{R}}|x|^{-r-3} \sum_{\mu=1}^{3}\left|\partial_{\mu} U(x)\right| d x$ is finite. As a consequence

$$
\left|\int_{B_{\frac{c}{R}}^{c}}\left(\nabla U \cdot \nabla \phi+\tau \partial_{1} U \cdot \phi\right) d x\right| \leq \mathfrak{C} G(\phi) \int_{B_{\bar{R}}^{c}}|x|^{-r-3} \sum_{\mu=1}^{3}\left|\partial_{\mu} U(x)\right| d x
$$

with the right-hand side being bounded, in turn, by $\mathfrak{C} \sum_{\alpha \in \mathbb{N}_{0}^{3},|\alpha| \leq 1} \sum_{m=1}^{3} p_{\bar{r}+3 e_{m}, \alpha}(\phi)$, for $\phi \in \mathcal{S}\left(\mathbb{R}^{3}\right)^{3}$. Since $\left|\int_{B_{\bar{R}}}\left(\nabla U \cdot \nabla \phi+\tau \partial_{1} U \cdot \phi\right) d x\right| \leq \mathfrak{C}\left\|U \mid B_{\bar{R}}\right\|_{1,1} p_{0,0}(\phi)$, we have thus found that $|T(\phi)|$ is bounded by $\mathfrak{C}\left(p_{0,0}(\phi)+\sum_{\alpha \in \mathbb{N}_{0}^{3},|\alpha| \leq 1} \sum_{m=1}^{3} p_{\bar{r}+3 e_{m}, \alpha}(\phi)\right)$ for $\phi \in \mathcal{S}\left(\mathbb{R}^{3}\right)^{3}$. This shows our claim for $T$. A similar but simpler reasoning is valid for $S$.

In Section 5, when we exploit the preceding theorem, the next two lemmas will be useful.
Lemma 3.1 Let $m_{0} \in \mathbb{N}, p_{1}, \ldots, p_{m_{0}} \in(1, \infty), \widetilde{p}:=\max \left\{p_{m}: 1 \leq m \leq m_{0}\right\}, \epsilon_{0} \in$ $[0,1 / \widetilde{p}), R_{1} \in(0, \infty), W^{(1)}, \ldots, W^{\left(m_{0}\right)} \in L_{l o c}^{1}\left(\mathbb{R}^{3}\right)^{3}$ such that $W:=\sum_{m=1}^{m_{0}} W^{(m)}$ is a polynomial and $\int_{B_{R_{1}}^{c}}\left(\left|W^{(m)}(x)\right||x|^{-2-\epsilon_{0}}\right)^{p_{m}} d x<\infty$ for $1 \leq m \leq m_{0}$. Then the degree of $W$ is at most 1.

Proof: Abbreviate $\mathfrak{A}(x):=\sum_{m=1}^{m_{0}}\left(\left|W^{(m)}(x)\right||x|^{-2-\epsilon_{0}}\right)^{p_{m}}$ for $x \in \mathbb{R}^{3}$. Suppose there is $\widetilde{R} \in\left[R_{1}, \infty\right)$ such that $\int_{\partial B_{R}} \mathfrak{A}(x) d o_{x} \geq 1 / R$ for $R \in[\widetilde{R}, \infty)$. Then $\int_{B_{\widetilde{R}^{c}}} \mathfrak{A}(x) d x=$ $\int_{\widetilde{R}}^{\infty} \int_{\partial B_{R}} \mathfrak{A}(x) d o_{x} d r \geq \int_{\widetilde{R}}^{\infty} r^{-1} d r=\infty$. Since $\int_{B_{R_{1}}^{c}} \mathfrak{A}(x) d x<\infty$ by our assumptions, we have arrived at a contradiction. Thus there is a sequence $\left(R_{n}\right)$ in $\left[\max \left\{R_{1}, 1\right\}, \infty\right)$ with $\int_{\partial B_{R}} \mathfrak{A}(x) d o_{x} \leq R_{n}^{-1}$ for $n \in \mathbb{N}$ and $R_{n} \rightarrow \infty$. It follows that $\mathfrak{K}_{n}:=\int_{\partial B_{1}} \mathfrak{A}\left(R_{n} x\right) d o_{x} \leq$ $R_{n}^{-3}$ for $n \in \mathbb{N}$. But $\left(R_{n}^{-2-\epsilon_{0}}\right)^{p_{m}}=\left(R_{n}^{-2-\epsilon_{0}+1 / \widetilde{p}}\right)^{p_{m}} R_{n}^{-p_{m} / \widetilde{p}} \geq\left(R_{n}^{-2-\epsilon_{0}+1 / \widetilde{p}}\right)^{p_{m}} R_{n}^{-1}$ for $n \in \mathbb{N}, 1 \leq m \leq m_{0}$, where the last inequality holds because $R_{n} \geq 1$ and $p_{m} \leq \widetilde{p}$. Setting $\widetilde{\mathfrak{K}}_{n}:=\int_{\partial B_{1}} \sum_{m=1}^{m_{0}}\left(\left|W^{(m)}\left(R_{n} y\right)\right| R_{n}^{-2-\epsilon_{0}+1 / \widetilde{p}}\right)^{p_{m}} d o_{y}$ for $n \in \mathbb{N}$, we get $\widetilde{\mathfrak{K}}_{n} \leq \mathfrak{K}_{n} R_{n} \leq$ $R_{n}^{-2}(n \in \mathbb{N})$ by the preceding inequalities. On the other hand, for $m \in\left\{1, \ldots, m_{0}\right\}, n \in \mathbb{N}$, we find with Hölder's inequality that $R_{n}^{-2-\epsilon_{0}+1 / \widetilde{p}} \int_{\partial B_{1}}\left|W^{(m)}\left(R_{n} y\right)\right| d o_{y} \leq(4 \pi)^{1 / p_{m}^{\prime}} \widetilde{\mathfrak{K}}_{n}^{1 / p_{m}}$. It follows that $R_{n}^{-2-\epsilon_{0}+1 / \widetilde{p}} \int_{\partial B_{1}} \sum_{m=1}^{m_{0}}\left|W^{(m)}\left(R_{n} y\right)\right| d o_{y} \rightarrow 0(n \rightarrow \infty)$, so we have found that $R_{n}^{-2-\epsilon_{0}+1 / \widetilde{p}} \int_{\partial B_{1}}\left|W\left(R_{n} y\right)\right| d o_{y} \rightarrow 0(n \rightarrow \infty)$.
Now suppose that the degree of $W$ is larger than 1 . Then there is $k_{0} \in \mathbb{N}$ and for any $\alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha| \leq k_{0}$ a number $a_{\alpha} \in \mathbb{R}$ such that $k_{0} \geq 2, W(x)=\sum_{\alpha \in \mathbb{N}_{0}^{3},|\alpha| \leq k_{0}} a_{\alpha} x^{\alpha}$ for $x \in \mathbb{R}^{3}$, and $a_{\alpha_{0}} \neq 0$ for some $\alpha_{0} \in \mathbb{N}_{0}^{3}$ with $\left|\alpha_{0}\right|=k_{0}$. Put $P(x):=\sum_{\alpha \in \mathbb{N}_{0}^{3},|\alpha|=k_{0}} a_{\alpha} x^{\alpha}\left(x \in \mathbb{R}^{3}\right)$. We distinguish two cases. In the first, we suppose that $\int_{\partial B_{1}}|P(y)| d o_{y}>0$. Since for $S \in(0, \infty)$

$$
\int_{\partial B_{1}}|W(S y)| d o_{y} \geq S^{k_{0}} \int_{\partial B_{1}}|P(y)| d o_{y}-\sum_{l=0}^{k_{0}-1} S^{l} \sum_{\alpha \in \mathbb{N}_{0}^{3},|\alpha|=l}\left|a_{\alpha}\right| \int_{\partial B_{1}}\left|y^{\alpha}\right| d o_{y}
$$

we may conclude there is $\widetilde{R} \in\left[R_{1}, \infty\right)$ such that $\int_{\partial B_{1}}|W(S y)| d o_{y} \geq S^{k_{0}} \int_{\partial B_{1}}|P(y)| d o_{y} / 2$ for $S \in[\widetilde{R}, \infty)$. Thus $R_{n}^{-2-\epsilon_{0}+1 / \widetilde{p}} \int_{\partial B_{1}}\left|W\left(R_{n} y\right)\right| d o_{y} \geq R_{n}^{k_{0}-2-\epsilon_{0}+1 / \widetilde{p}} \int_{\partial B_{1}}|P(y)| d o_{y} / 2$ for $n \in \mathbb{N}$ with $R_{n} \geq \widetilde{R}$. On the other hand, since we assumed $k_{0} \geq 2$ and $\epsilon_{0} \in[0,1 / \widetilde{p})$, and because $R_{n} \rightarrow \infty$ and $\int_{\partial B_{1}}|P(y)| d o_{y}>0$, we get $R_{n}^{k_{0}-2-\epsilon_{0}+1 / \widetilde{p}} \int_{\partial B_{1}}|P(y)| d o_{y} / 2 \rightarrow$ $\infty(n \rightarrow \infty)$. We may thus conclude that $R_{n}^{-2-\epsilon_{0}+1 / \widetilde{p}} \int_{\partial B_{1}}\left|W\left(R_{n} y\right)\right| d o_{y} \rightarrow \infty$ for $n$ tending to infinity, which is a contradiction to what was shown above.
In the second case, we assume that $\int_{\partial B_{1}}|P(y)| d o_{y}=0$. Then $P(y)=0$ for a. e. $y \in \partial B_{1}$, and hence for any $y$ on $\partial B_{1}$ by continuity. Since $P$ is homogeneous, it follows that $P=0$ in $\mathbb{R}^{3}$. For $x_{2}, x_{3} \in \mathbb{R}$, the function $P\left(\cdot, x_{2}, x_{3}\right)$ is a polynomial in one variable, so we may conclude that all coefficients of this polynomial vanish. But this polynomial may be written as $\sum_{l=0}^{k_{0}} P_{l}\left(x_{2}, x_{3}\right) r^{l}$ for $r \in \mathbb{R}$, with $P_{l}\left(x_{2}, x_{3}\right):=\sum_{m=0}^{k_{0}-l} a_{\left(l, m, k_{0}-l-m\right)} x_{2}^{m} x_{3}^{k_{0}-l-m}$ for $x_{2}, x_{3} \in \mathbb{R}, 0 \leq l \leq k_{0}$. The numbers $P_{l}\left(x_{2}, x_{3}\right)$ are the coefficients of $P\left(\cdot, x_{2}, x_{3}\right)$ and must therefore vanish $\left(x_{2}, x_{3} \in \mathbb{R}\right)$. But each function $P_{l}$ is a polynomial, too. So, in the next step, the same sort of reasoning may be applied to these polynomials, implying that their coefficients must vanish as well. In the end we get that $a_{\alpha}=0$ for $\alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha|=k_{0}$, which is a contradiction to the fact that $a_{\alpha_{0}} \neq 0$ for some $\alpha_{0} \in \mathbb{N}_{0}^{3}$ with $|\alpha|=k_{0}$. Thus we arrive at a contradiction in any case. This proves that the degree of $W$ cannot exceed 1.
Lemma 3.2 Let $m_{0} \in \mathbb{N}, p_{m} \in(1, \infty)$, $V^{(m)} \in L^{p_{m}}\left(\mathbb{R}^{3}\right)^{3}$ for $1 \leq m \leq m_{0}$. Suppose that $V:=\sum_{m=1}^{m_{0}} V^{(m)}$ is constant. Then $V=0$.

Proof: By an argument already used in the proof of [19, Lemma 5.2] and Lemma 3.1, we may choose a sequence $\left(R_{n}\right)$ in $(0, \infty)$ with $\int_{\partial B_{R_{n}}} \sum_{m=1}^{m_{0}}\left|V^{(m)}(x)\right|^{p_{m}} d o_{x} \leq R_{n}^{-1}$ for $n \in \mathbb{N}$ and with $R_{n} \rightarrow \infty$. Put $\mathfrak{A}_{n}:=\int_{\partial B_{1}} \sum_{m=1}^{m_{0}}\left|V^{(m)}\left(R_{n} y\right)\right|^{p_{m}} d o_{y}$ for $n \in \mathbb{N}$, so $\mathfrak{A}_{n} \leq R_{n}^{-3}$. But $\int_{\partial B_{1}}\left|V^{(m)}\left(R_{n} y\right)\right| d o_{y} \leq(4 \pi)^{1 / p_{m}^{\prime}} \mathfrak{A}_{n}^{1 / p_{m}}$ for $n \in \mathbb{N}, 1 \leq m \leq m_{0}$ by Hölder's inequality, hence $\int_{\partial B_{1}} \sum_{m=1}^{m_{0}}\left|V^{(m)}\left(R_{n} y\right)\right| d o_{y} \rightarrow 0(n \rightarrow \infty)$. Therefore $\int_{\partial B_{1}}\left|V\left(R_{n} y\right)\right| d o_{y} \rightarrow 0$ for $n$ tending to infinity. Since $V$ is constant, this means that $V=0$.

## 4 Some fundamental solutions and potential functions.

We recall that the fundamental solution $\mathfrak{N}$ of the Poisson equation ("Newton kernel") and a convolution ("Newton potential") with this fundamental solution was introduced in Theorem 3.1. We define the usual heat kernel in 3D by setting
$\mathfrak{H}(z, t):=(4 \pi t)^{-3 / 2} e^{-|z|^{2} /(4 t)}$ for $z \in \mathbb{R}^{3}, t \in(0, \infty), \quad \mathfrak{H}(z, 0):=0$ for $z \in \mathbb{R}^{3} \backslash\{0\}$.
Thus, in our context, $\mathfrak{H}$ is defined on $\mathfrak{B}:=\left(\mathbb{R}^{3} \times(0, \infty)\right) \cup\left(\left(\mathbb{R}^{3} \backslash\{0\}\right) \times\{0\}\right)$.
Theorem 4.1 The relations $\mathfrak{H} \in C^{\infty}(\mathfrak{B})$ and $\int_{\mathbb{R}^{3}} \mathfrak{H}(z, t) d t=1$ for $t \in(0, \infty)$ hold. If $\alpha \in \mathbb{N}_{0}^{3}, \sigma \in \mathbb{N}_{0}$, the inequality $\left|\partial_{z}^{\alpha} \partial_{t}^{\sigma} \mathfrak{H}(z, t)\right| \leq C(\alpha, \sigma)\left(|z|^{2}+t\right)^{-(3+|\alpha|+2 \sigma) / 2}$ is valid for $z \in \mathbb{R}^{3}, t \in(0, \infty)$.
Proof: See [41] for the preceding estimate.
The estimate in Theorem 4.1 in the case $|\alpha|=2, \sigma=0$ allows to define the velocity part $\Gamma$ of a fundamental solution to the time-dependent Stokes system,

$$
\Gamma_{j k}(z, t):=\mathfrak{H}(z, t) \delta_{j k}+\int_{t}^{\infty} \partial z_{j} \partial z_{k} \mathfrak{H}(z, s) d s \quad \text { for }(z, t) \in \mathfrak{B}, j, k \in\{1,2,3\},
$$

as well as the velocity part $\Lambda$ of a fundamental solution to the time-dependent Oseen system (1.1),

$$
\Lambda_{j k}(z, t):=\Gamma_{j k}\left(z-\tau t e_{1}, t\right) \quad \text { for }(z, t) \in \mathfrak{B}, j, k \in\{1,2,3\} .
$$

We will need the following properties of $\Lambda$.
Lemma 4.1 The relations $\Lambda \in C^{\infty}(\mathfrak{B})^{3 \times 3}$ and $\sum_{k=1}^{3} \partial z_{k} \Lambda_{j k}(z, t)=0$ are valid for $1 \leq$ $j \leq 3, z \in \mathbb{R}^{3}, t \in(0, \infty)$.
Let $K>0$. Then $\left|\partial_{z}^{\alpha} \partial_{t}^{\sigma} \Lambda(z, t)\right| \leq C(K, \tau) \sum_{\mu=1}^{2}(|z| \nu(z)+t)^{-(3+|\alpha|+\mu \sigma) / 2}$ for $z, t, \alpha, \sigma$ as before, with the additional assumption $|z| \geq K$.
The estimate $\int_{\mathbb{R}^{3}}\left|\partial_{x}^{\alpha} \partial_{t}^{\sigma} \Lambda(x, t)\right|^{q} d x \leq C(\tau, q) \sum_{\mu=1}^{2} t^{-(3+|\alpha|+\mu \sigma) q / 2+3 / 2}$ is valid for $q \in$ $(1, \infty), t \in(0, \infty), \alpha \in \mathbb{N}_{0}^{3}, \sigma \in\{0,1\}$ with $|\alpha|+\sigma \leq 1$. If $|\alpha|+\sigma=1$, the case $q=1$ is also admitted.
Proof: See [19, Lemma 3.3, Corollary 3.3, 3.4].
We introduce the first of our potential functions.

Lemma 4.2 ([19, Corollary 3.5]) Let $A \subset \mathbb{R}^{3}$ be measurable, $q \in[1, \infty)$, $V \in L^{q}(A)^{3}$, and let $\widetilde{V}$ denote the zero extension of $V$ to $\mathbb{R}^{3}$. Then $\int_{\mathbb{R}^{3}}\left|\partial_{x}^{\alpha} \Lambda(x-y, t) \widetilde{V}(y)\right| d y<\infty$ for $\alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha| \leq 1, x \in \mathbb{R}^{3}, t \in(0, \infty)$. Define the function $\mathfrak{I}^{(\tau)}(V): \mathbb{R}^{3} \times(0, \infty) \mapsto \mathbb{R}^{3}$ by $\mathfrak{I}^{(\tau)}(V)(x, t):=\int_{\mathbb{R}^{3}} \Lambda(x-y, t) \cdot \widetilde{V}(y) d y \quad$ for $x \in \mathbb{R}^{3}, t \in(0, \infty)$.
The derivative $\partial x_{l} \mathfrak{I}^{(\tau)}(V)(x, t)$ exists and equals $\int_{\mathbb{R}^{3}} \partial x_{l} \Lambda(x-y, t) \cdot \tilde{V}(y) d y$ for $x, t$ as above, and for $l \in\{1,2,3\}$. The functions $\mathfrak{I}^{(\tau)}(V)$ and $\partial x_{l} \mathfrak{I}^{(\tau)}(V)$ are continuous in $\mathbb{R}^{3} \times(0, \infty)$. If $q>1$, then $\left\|\mathfrak{I}^{(\tau)}(V)(t)\right\|_{q} \leq C(q, \tau)\|V\|_{q}$ for $t \in(0, \infty)$.
If $m_{0} \in \mathbb{N}, p_{l} \in[1, \infty), V^{(l)} \in L^{p_{l}}(A)^{3}$ for $1 \leq l \leq m_{0}$, put $\mathfrak{I}^{(\tau)}\left(\sum_{l=1}^{m_{0}} V^{(l)}\right):=$ $\sum_{l=1}^{m_{0}} \mathfrak{I}^{(\tau)}\left(V^{(l)}\right)$.
We will need a variant of $\mathfrak{I}^{(\tau)}(V)$.
Lemma 4.3 Let $q \in(1, \infty)$, $A \subset \mathbb{R}^{3}$ be measurable, $V \in L^{q}(A)^{3}$. Write $\widetilde{V}$ for the zero extension of $V$ to $\mathbb{R}^{3}$. Then $\int_{\mathbb{R}^{3}}\left|\partial_{x}^{\alpha} \mathfrak{H}\left(x-y-\tau t e_{1}, t\right) \widetilde{V}(y)\right| d y<\infty$ for $x \in \mathbb{R}^{3}$, $t \in$ $(0, \infty), \alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha| \leq 1$.
Define $\mathcal{H}^{(\tau)}(V)(x, t):=\int_{\mathbb{R}^{3}} \mathfrak{H}\left(x-y-\tau t e_{1}, t\right) \widetilde{V}(y) d y, \quad \mathcal{H}^{(\tau)}(V)(x, 0):=\widetilde{V}(x, 0)$ for $x \in \mathbb{R}^{3}, t \in(0, \infty)$. Then $\mathcal{H}^{(\tau)}(V)$ belongs to $C^{0}\left([0, \infty), L^{2}\left(\mathbb{R}^{3}\right)^{3}\right)$, and the estimate $\left\|\mathcal{H}^{(\tau)}(V)(t)\right\|_{2} \leq C(\tau)\|V\|_{2}$. holds. Moreover, the derivative $\partial x_{l} \mathcal{H}^{(\tau)}(V)(x, t)$ exists and equals $\int_{\mathbb{R}^{3}} \partial_{x}^{\alpha} \mathfrak{H}\left(x-y-\tau t e_{1}, t\right) \widetilde{V}(y) d y$, for $x \in \mathbb{R}^{3}, t \in(0, \infty), 1 \leq l \leq 3$. In addition $\nabla_{x} \mathcal{H}^{(\tau)}(V) \in L^{2}\left(0, \infty, L^{2}\left(\mathbb{R}^{3}\right)^{3}\right)$.
Proof: All the claims of the lemma except the continuity of $\mathcal{H}^{(\tau)}(V)$ and the integrability of $\left|\nabla_{x} \mathcal{H}^{(\tau)}(V)\right|^{2}$ follow by the same arguments as used in [19, proof of Corollary 3.5] with respect to $\mathfrak{I}^{(\tau)}(V)$. The continuity of $\mathcal{H}^{(\tau)}(V)$ as a mapping from $[0, \infty)$ to $L^{2}\left(\mathbb{R}^{3}\right)^{3}$ at $t=0$ holds according to [19, Theorem 3.3]. Continuity at $t>0$ may be shown by the same reasoning as in [19, proof of Corollary 3.6] As for $L^{2}$-integrability of the spatial gradient, we use that the Fourier transform $[\mathfrak{H}(\cdot, t)]^{\wedge}$ of the function $x \mapsto \mathfrak{H}(x, t) \quad\left(x \in \mathbb{R}^{3}\right)$ (heat kernel as a function of the space variables) is given by $[\mathfrak{H}(\cdot, t)]^{\wedge}(\xi)=(2 \pi)^{-3 / 2} e^{-|\xi|^{2} t}$ for $\xi \in \mathbb{R}^{3}, t \in(0, \infty)$. Therefore by Parseval's formula $\left\|\partial x_{l} \mathcal{H}^{(\tau)}(V)\right\|_{2}^{2}=\int_{0}^{\infty} \int_{\mathbb{R}^{3}}\left|(2 \pi)^{-3 / 2} \xi_{l} e^{\left(i \tau \xi_{1}-|\xi|^{2}\right) t} \widehat{V}(\xi)\right|^{2} d \xi d t$. By applying Hölder's inequality with respect to the variable $\xi$, and then integrating with respect to $t$, we get $\left\|\partial x_{l} \mathcal{H}^{(\tau)}(V)\right\|_{2}^{2} \leq \mathfrak{C}\|\widehat{V}\|_{2}^{2}$, so Parseval's formula yields $\left\|\partial x_{l} \mathcal{H}^{(\tau)}(V)\right\|_{2}^{2} \leq \mathfrak{C}\|V\|_{2}^{2}$. This completes the proof of Lemma 4.3

We turn to the definition of another potential function.
Lemma 4.4 Let $T_{0} \in(0, \infty], A \subset \mathbb{R}^{3}$ measurable, $q \in[1, \infty)$ and $f$ a function from $L_{l o c}^{1}\left(\left[0, T_{0}\right), L^{q}(A)^{3}\right)$. Let $\widetilde{f}$ denote the zero extension of $f$ to $\mathbb{R}^{3} \times(0, \infty)$. Then the integral $\int_{\mathbb{R}^{3}}\left|\partial_{x}^{\alpha} \Lambda(x-y, t-\sigma) \cdot \tilde{f}(y, \sigma)\right| d y$ is finite for any $x \in \mathbb{R}^{3}, t \in(0, \infty), \sigma \in$ $(0, t), \alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha| \leq 1$. Moreover, for a. e. $t \in(0, \infty)$ and for $\alpha$ as before, the integral $\int_{0}^{t} \int_{\mathbb{R}^{3}}\left|\partial_{x}^{\alpha} \Lambda(x-y, t-\sigma) \cdot \widetilde{f}(y, \sigma)\right| d y d \sigma$ is finite for a. e. $x \in \mathbb{R}^{3}$. Thus we may define

$$
\mathfrak{R}^{(\tau)}(f)(x, t):=\int_{0}^{t} \int_{\mathbb{R}^{3}} \Lambda(x-y, t-\sigma) \cdot \widetilde{f}(y, \sigma) d y d \sigma
$$

for such $t$ and $x$. The relation $\mathfrak{R}^{(\tau)}(f)(t) \in W_{l o c}^{1,1}\left(\mathbb{R}^{3}\right)^{3}$ holds for a. e. $t \in(0, \infty)$, and for such $t \partial x_{l} \mathfrak{R}^{(\tau)}(f)(t)(x)=\int_{0}^{t} \int_{\mathbb{R}^{3}} \partial x_{l} \Lambda(x-y, t-\sigma) \cdot \widetilde{f}(y, \sigma) d y d \sigma$ for $l \in\{1,2,3\}$ and $a$.
e. $x \in \mathbb{R}^{3}$.

Moreover the integral $\int_{0}^{t}\left|\int_{\mathbb{R}^{3}} \Lambda(x-y, t-s) \cdot \widetilde{f}(y, s) d y\right| d s$ is finite for any $t \in(0, \infty)$ and for $a$. e. $x \in \mathbb{R}^{3}$. Thus the function $\mathfrak{R}^{(\tau)}(f)$ is well defined even for any $t \in(0, \infty)$ (instead of only for a. e. $t \in(0, \infty)$ ) and for a. e. $x \in \mathbb{R}^{3}$. In addition the inequality $\left\|\mathfrak{R}^{(\tau)}(f)(t)\right\|_{q} \leq C(q)\left\|\widetilde{f} \mid \mathbb{R}^{3} \times(0, t)\right\|_{q, 1 ; t}$ holds for $t>0$.
Let $m_{0} \in \mathbb{N}, p_{l} \in(1, \infty)$ and $f^{(l)} \in L_{l o c}^{1}\left(\left[0, T_{0}\right), L^{p_{l}}(A)^{3}\right)$ for $1 \leq l \leq m_{0}$. Then define $\mathfrak{R}^{(\tau)}\left(\sum_{j=1}^{m_{0}} f^{(l)}\right):=\sum_{j=1}^{m_{0}} \mathfrak{R}\left(f^{(l)}\right)$.
Proof: [19, Lemma 3.8, Corollary 3.7].
Lemma 4.5 Let $q \in(1, \infty)$ and $f \in L^{2}\left(0, \infty, L^{q}\left(\mathbb{R}^{3}\right)^{3}\right)$. Then the function $\mathfrak{R}^{(\tau)}(f)$ belongs to the space $C^{0}\left([0, \infty), L^{q}\left(\mathbb{R}^{3}\right)^{3}\right)$.
Let $B \subset \mathbb{R}^{3}$ be measurable and bounded. Suppose that $q<3$. Then $\mathfrak{R}^{(\tau)}(f) \mid B \times(0, \infty) \in$ $L^{\infty}\left(0, \infty, L^{q}(B)^{3}\right)$ and $\nabla_{x} \mathfrak{R}^{(\tau)}(f) \mid B \times(0, \infty) \in L^{2}\left(0, \infty, L^{q}(B)^{3}\right)$.

Proof: Put $\mathfrak{B}(x, t, h):=\int_{t}^{t+h} \mathfrak{I}^{(\tau)}(f(s))(x, t+h-s) d s$ for $t \in[0, \infty), h \in(0, \infty), x \in$ $\mathbb{R}^{3}$. Then $\|\mathfrak{B}(\cdot, t, h)\|_{q} \leq \int_{t}^{t+h}\left\|\mathfrak{I}^{(\tau)}(f(s))(\cdot, t+h-s)\right\|_{q} d s \leq \mathfrak{C} \int_{t}^{t+h}\|f(s)\|_{q} d s \leq$ $\mathfrak{C} h^{1 / 2}\|f\|_{q, 2 ; \infty}$ by Theorem 2.1 and Lemma 4.2. We may conclude that $\|\mathfrak{B}(\cdot, t, h)\| q \rightarrow$ $0(h \downarrow 0)$. It follows in particular that $\mathfrak{R}^{(\tau)}(f)$ as a mapping from $[0, \infty)$ in $L^{q}\left(\mathbb{R}^{3}\right)^{3}$ is continuous in $t=0$. Let $t \in(0, \infty)$. Then $\mathfrak{R}^{(\tau)}(f)(x, t+h)-\mathfrak{R}^{(\tau)}(f)(x, t)=\mathfrak{B}(x, t, h)+$ $\mathcal{C}(x, t, h)$, with $\mathcal{C}(x, t, h):=\int_{0}^{t}\left[\mathfrak{I}^{(\tau)}(f(s))(x, t+h-s)-\mathfrak{I}^{(\tau)}(f(s))(x, t-s)\right] d s$, for $x \in$ $\mathbb{R}^{3}, h>0$. Define $\mathfrak{A}(s, h):=\left\|\mathfrak{I}^{(\tau)}(f(s))(t+h-s)-\mathfrak{I}^{(\tau)}(f(s))(t-s)\right\|_{q}$ for $s \in(0, t), h>0$. Then Lemma 4.2 yields that $\mathfrak{A}(s, h) \leq \sum_{r \in\{0, h\}}\left\|\mathfrak{I}^{(\tau)}(f(s))(t+r-s)\right\|_{q} \leq \mathfrak{C}\|f(s)\|_{q}$ for $s, h$ as before. The function $s \mapsto\|f(s)\|_{q}(s \in(0, t))$ is integrable. Moreover, for $s \in(0, t), h>0$, we get by the definition of $\mathfrak{I}^{(\tau)}(f(s))$, Young's inequality and the last inequality in Lemma 4.1 that

$$
\begin{aligned}
& \mathfrak{A}(s, h) \leq\left(\int_{\mathbb{R}^{3}}|\Lambda(z, t+h-s)-\Lambda(z, t-s)| d z\right)\|f(s)\|_{q} \\
& =\left(\int_{\mathbb{R}^{3}}\left|\int_{0}^{1} \partial_{r} \Lambda(z, r)_{\mid r=t+\vartheta h-s} d \vartheta\right| d z\right)\|f(s)\|_{q} h \leq \mathfrak{C}^{\mathfrak{C}} \sum_{\mu=1}^{2}(t-s)^{-\mu / 2}\|f(s)\|_{q} h .
\end{aligned}
$$

Therefore $\mathfrak{A}(s, h) \rightarrow 0$ for $h \downarrow 0$. At this point it follows with Lebesgue's theorem that $\|\mathcal{C}(\cdot, t, h)\|_{q} \rightarrow 0(h \downarrow 0)$. Since we have already shown that $\|\mathfrak{B}(\cdot, t, h)\|_{q} \rightarrow 0(h \downarrow 0)$, we may conclude that $\left\|\mathfrak{R}^{(\tau)}(f)(t+h)-\mathfrak{R}^{(\tau)}(f)(t)\right\|_{q} \rightarrow 0(h \downarrow 0)$. Altogether it follows that $\mathfrak{R}^{(\tau)}(f)$ as a mapping from $[0, \infty)$ into $L^{q}\left(\bar{\Omega}^{c}\right)^{3}$ is continuous.
In view of the second part of the lemma, we obtain with Lemma 4.4 that $\partial_{x}^{\alpha} \mathfrak{R}^{(\tau)}(f)(x, t)=$ $\mathfrak{A}_{1}^{(\alpha)}+\mathfrak{A}_{2}^{(\alpha)}$, with $\mathfrak{A}_{1}^{(\alpha)}(x, t):=\int_{0}^{t} \chi_{(0,1)}(t-s) \int_{\mathbb{R}^{3}} \partial_{x}^{\alpha} \Lambda(x-y, t-s) f(y, s) d y d s$ for $x \in$ $\mathbb{R}^{3}, t \in(0, \infty), \alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha| \leq 1$, and with $\mathfrak{A}_{2}^{(\alpha)}(x, t)$ defined in the same way as $\mathfrak{A}_{1}^{(\alpha)}(x, t)$, except that the term $\chi_{(0,1)}(t-s)$ is replaced by $\chi_{[1, \infty)}(t-s)$. Then for $t \in(0, \infty)$, we have $\left\|\mathfrak{A}_{1}^{(0)}(t)\right\|_{q} \leq \int_{0}^{t} \chi_{(0,1)}(t-s)\left\|\mathfrak{I}^{(\tau)}(f(s))(t-s)\right\|_{q} d s$, so we may conclude with Lemma 4.2 that $\left\|\mathfrak{A}_{1}^{(0)}(t)\right\|_{q} \leq \mathfrak{C} \int_{0}^{t} \chi_{(0,1)}(t-s)\|f(s)\|_{q} d s \leq \mathfrak{C}\|f\|_{q, 2 ; \infty}$, where the last estimate follows with Hölder's inequality. Hence we get $\left\|\mathfrak{A}_{1}^{(0)}\right\|_{q, \infty ; \infty} \leq\|f\|_{q, 2 ; \infty}$. Moreover,
by Hölder's inequality and the last inequality in Lemma 4.1 , for $x \in \mathbb{R}^{3}, t>0$,

$$
\begin{aligned}
& \left|\mathfrak{A}_{2}^{(0)}(x, t)\right| \leq \int_{0}^{t} \chi_{[1, \infty)}(t-s)\left(\int_{\mathbb{R}^{3}}|\Lambda(x-y, t-s)|^{q^{\prime}} d y\right)^{1 / q^{\prime}}\|f(s)\|_{q} d s \\
& \leq \mathfrak{C} \int_{0}^{t} \chi_{[1, \infty)}(t-s)(t-s)^{-3 /(2 q)}\|f(s)\|_{q} d s
\end{aligned}
$$

Suppose that $q<3$. Then $3 /(2 q)>1 / 2$, so by Hölder's inequality $\left|\mathfrak{A}_{2}^{(0)}(x, t)\right| \leq \mathfrak{C}\|f\|_{q, 2 ; \infty}$. It follows that $\left\|\mathfrak{A}_{2}^{(0)}\right\|_{\infty, \infty ; \infty} \leq \mathfrak{C}\|f\|_{q, 2 ; \infty}$. This estimate and the one given above for $\left\|\mathfrak{A}_{1}^{(0)}\right\|_{q, \infty ; \infty}$ imply $\left\|\mathfrak{R}^{(\tau)}(f) \mid B \times(0, \infty)\right\|_{q, \infty ; \infty} \leq \mathfrak{C}\|f\|_{q, 2 ; \infty}$. Now let $l \in\{1,2,3\}$. Young's inequality and the last estimate in Lemma 4.1 yield

$$
\begin{aligned}
& \left\|\mathfrak{A}_{1}^{\left(e_{l}\right)}(t)\right\|_{q} \leq \int_{0}^{t} \chi_{(0,1)}(t-s)\left(\int_{\mathbb{R}^{3}}\left|\partial z_{l} \Lambda(z, t-s)\right| d z\right)\|f(s)\|_{q} d s \\
& \leq \mathfrak{C} \int_{0}^{t} \chi_{(0,1)}(t-s)(t-s)^{-1 / 2}\|f(s)\|_{q} d s
\end{aligned}
$$

for $t \in(0, \infty)$, and thus $\left\|\mathfrak{A}_{1}^{\left(e_{l}\right)}\right\|_{q, 2 ; \infty} \leq \int_{0}^{\infty} \chi_{(0,1)}(r) r^{-1 / 2} d r\|f\|_{q, 2 ; \infty} \leq \mathfrak{C}\|f\|_{q, 2 ; \infty}$ again by Young's inequality. Applying Hölder's inequality and the last estimate in Lemma 4.1 as in the estimate of $\left|\mathfrak{A}_{2}^{(0)}(x, t)\right|$, we obtain for $x \in \mathbb{R}^{3}, t>0$ that $\left|\mathfrak{A}_{2}^{\left(e_{l}\right)}(x, t)\right| \leq$ $\mathfrak{C} \int_{0}^{t} \chi_{[1, \infty)}(t-s)(t-s)^{-3 /(2 q)-1 / 2}\|f(s)\|_{q} d s$. But $\left\|\mathfrak{A}_{2}^{\left(e_{l}\right)}(t)\left|B\left\|_{q} \leq \mathfrak{C}\right\| \mathfrak{A}_{2}^{\left(e_{l}\right)}(t)\right| B\right\|_{\infty}$. Thus, combining the two preceding inequalities, and recalling the assumption $q<3$, hence $3 /(2 q)+1 / 2>1$, we may conclude by another application of Young's inequality that $\left\|\mathfrak{A}_{2}^{\left(e_{l}\right)} \mid B \times(0, \infty)\right\|_{q, 2 ; \infty} \leq \mathfrak{C} \int_{0}^{\infty} \chi_{[1, \infty)}(r) r^{-3 /(2 q)-1 / 2} d r\|f\|_{q, 2 ; \infty} \leq \mathfrak{C}\|f\|_{q, 2 ; \infty}$. Combining this estimate of $\left\|\mathfrak{A}_{2}^{\left(e_{l}\right)} \mid B \times(0, \infty)\right\|_{q, 2 ; \infty}$ with the one of $\left\|\mathfrak{A}_{1}^{\left(e_{l}\right)}\right\|_{q, 2 ; \infty}$ given above, we obtain that $\left\|\partial x_{l} \mathfrak{R}^{(\tau)}(f) \mid B \times(0, \infty)\right\|_{q, 2 ; \infty} \leq \mathfrak{C}\|f\|_{q, 2 ; \infty}$. This completes the proof of Lemma 4.5.

The next lemma deals with still another potential function, this one defined on the surface of an open bounded set.
Lemma 4.6 Let $q \in[1, \infty], T_{0} \in(0, \infty], A \subset \mathbb{R}^{3}$ open and bounded, with Lipschitz boundary, $\phi \in L_{l o c}^{1}\left(\left[0, T_{0}\right), L^{q}(\partial A)^{3}\right)$, $\widetilde{\phi}$ the zero extension of $\phi$ to $\partial A \times(0, \infty)$. For $t \in(0, \infty), x \in \mathbb{R}^{3} \backslash \partial A, \alpha \in \mathbb{N}_{0}^{3}$, the term $\left|\partial_{x}^{\alpha} \Lambda(x-y, t-s) \cdot \widetilde{\phi}(y, s)\right|$ is integrable as a function of $(y, s) \in \partial A \times(0, t)$. Define $\mathfrak{V}^{(\tau)}(\phi):=\mathfrak{V}^{(\tau, A)}(\phi):\left(\mathbb{R}^{3} \backslash \partial A\right) \times(0, \infty) \mapsto \mathbb{R}^{3}$ by

$$
\mathfrak{V}^{(\tau)}(\phi)(x, t):=\int_{0}^{t} \int_{\partial A} \Lambda(x-y, t-s) \cdot \widetilde{\phi}(y, s) d o_{y} d s \quad \text { for } x \in \mathbb{R}^{3} \backslash \partial A, t \in(0, \infty)
$$

Then, for any $t \in(0, \infty)$, the integral $\int_{0}^{t} \int_{\partial A} \Lambda(x-y, t-s) \cdot \widetilde{\phi}(y, s) d o_{y} d s$ as a function of $x \in$ $\mathbb{R}^{3} \backslash A$ belongs to $C^{\infty}\left(\mathbb{R}^{3} \backslash A\right)^{3}$, and $\partial_{x}^{\alpha} \mathfrak{V}^{(\tau)}(\phi)(x, t)=\int_{0}^{t} \int_{\partial A} \partial_{x}^{\alpha} \Lambda(x-y, t-s) \cdot \widetilde{\phi}(y, s) d o_{y} d s$ for $\alpha \in \mathbb{N}_{0}^{3}, x \in \mathbb{R}^{3} \backslash A$.
Proof: The function $\Lambda$ is $C^{\infty}$ on $\mathbb{R}^{3} \times(0, \infty)$ (Lemma 4.1), so the lemma follows from Lebesgue's theorem.
Lemma 4.7 Let $\phi \in L^{2}(\partial \Omega \times(0, \infty))^{3}$ and $S \in(0, \infty)$ with $\bar{\Omega} \subset B_{S}$. Then $\mathfrak{V}^{(\tau)}(\phi) \mid B_{S}^{c} \times$ $[0, \infty) \in C^{0}\left([0, \infty), L^{4}\left(B_{S}^{c}\right)^{3}\right)$.

Proof: Since the function $\Lambda$ and its derivatives are continuous on $\left(\mathbb{R}^{3} \backslash\{0\}\right) \times[0, \infty)$ (Lemma 4.1), and because $\operatorname{dist}\left(\partial \Omega, B_{S}^{c}\right)>0$ and $|x-y| \geq|x| / 2$ for $y \in \partial \Omega, x \in B_{2 S}^{c}$, we may conclude from the first estimate in Lemma 4.1 that

$$
\begin{align*}
& \left|\partial_{r}^{\sigma} \Lambda(x-y, r)\right| \leq \mathfrak{C} \sum_{\mu=1}^{2}\left(\chi_{A_{2 S, S}}(x)+\chi_{B_{2 S}^{c}}(x)|x|^{-3 / 2-\mu \sigma / 2}\right)  \tag{4.5}\\
& \leq \mathfrak{C}\left(\chi_{A_{2 S, S}}(x)+\chi_{B_{2 S}^{c}}(x)|x|^{-3 / 2-\sigma / 2}\right)
\end{align*}
$$

for $x \in B_{S}^{c}, y \in \partial \Omega, r \in[0, \infty), \sigma \in\{0,1\}$. As a consequence,

$$
\begin{align*}
& |\Lambda(x-y, t+h-s)-\Lambda(x-y, t-s)|=\left|\int_{0}^{1} \partial_{r}^{\sigma} \Lambda(x-y, r)_{\mid r=t+\vartheta h-s} d \vartheta\right| h  \tag{4.6}\\
& \leq \mathfrak{C}\left(\chi_{A_{2 S, S}}(x)+\chi_{B_{2 S}^{c}}(x)|x|^{-2}\right) h \text { for } x \in B_{S}^{c}, y \in \partial \Omega, h, t \in(0, \infty), s \in(0, t) .
\end{align*}
$$

Put $\mathcal{B}(x, t, h):=\int_{t}^{t+h} \int_{\partial \Omega} \Lambda(x-y, t+h-s) \phi(y, s) d o_{y} d s$ for $x \in B_{S}^{c}, t \in[0, \infty), h \in$ $(0, \infty)$, and $\mathcal{C}(x, t, h):=\int_{0}^{t} \int_{\partial \Omega}(\Lambda(x-y, t+h-s)-\Lambda(x-y, t-s)) \phi(y, s) d o_{y} d s$ for $x, h$ as before and $t \in(0, \infty)$. Then by Theorem 2.1 and (4.5),

$$
\|\mathcal{B}(\cdot, t, h)\|_{4} \leq \mathfrak{C} \int_{t}^{t+h} \int_{\partial \Omega}|\phi(y, s)| d o_{y} d s\left(\int_{B_{S}^{c}}\left[\chi_{A_{2 S, S}}(x)+\chi_{B_{2 S}^{c}}(x)|x|^{-3 / 2}\right]^{4} d x\right)^{1 / 4}
$$

hence $\|\mathcal{B}(\cdot, t, h)\|_{4} \leq \mathfrak{C} h^{1 / 2}\|\phi\|_{2}$ for $t \in[0, \infty), h \in(0, \infty)$. Thus $\|\mathcal{B}(\cdot, t, h)\|_{4} \rightarrow 0$ when $h \downarrow 0$. This means in particular that $\mathfrak{V}^{(\tau)}(\phi) \mid B_{S}^{c} \times[0, \infty)$ as a mapping from $[0, \infty)$ into $L^{4}\left(B_{S}^{c}\right)^{3}$ is continuous in $t=0$. Let $t \in(0, \infty)$. Theorem 2.1 and inequality (4.6) imply that $\|\mathcal{C}(\cdot, t, h)\|_{4}$ may be estimated in the same way as $\|\mathcal{B}(\cdot, t, h)\|_{4}$, but with the integral over $(t, t+h)$ replaced by an integral over $(0, t)$, and with an additional factor $h$. It follows that $\|\mathcal{C}(\cdot, t, h)\|_{4} \leq \mathfrak{C} t^{1 / 2}\|\phi\|_{2} h$. Therefore $\left\|\mathfrak{V}^{(\tau)}(\phi)(t+h)-\mathfrak{V}^{(\tau)}(\phi)(t)\right\|_{4}=$ $\|\mathcal{B}(\cdot, t, h)+\mathcal{C}(\cdot, t, h)\|_{4} \leq \mathfrak{C}\left(1+t^{1 / 2}\right)\|\phi\|_{2}\left(h+h^{1 / 2}\right)$ for $h>0$. Thus $\mathfrak{V}^{(\tau)}(\phi) \mid B_{S}^{c} \times[0, \infty)$ is continuous on $(0, \infty)$. Since continuity in $t=0$ is already proved, the lemma follows.

In the rest of this section, we recapitulate results from [19] that will be needed in what follows.

We introduce another kernel function, for the definition of which we refer to [19]. This kernel is a truncated version of $\Lambda$. Its relevant properties of this kernel are collected in the ensuing theorem. Recall that the numbers $S_{0}, R_{0}, R_{1}$ and the function $\varphi_{0}$ are introduced in Section 1.

Theorem 4.2 There is a function $\mathfrak{G}:=\mathfrak{G}_{R_{0}, S_{0}, \varphi_{0}}: B_{R_{0}}^{c} \times B_{R_{1}} \times[0, \infty) \mapsto \mathbb{R}^{3 \times 3}$ with the following properties.
Let $x \in B_{R_{0}}^{c}, r \in[0, \infty)$. Then $\mathfrak{G}(x, \cdot, r) \in C_{0}^{\infty}\left(B_{R_{1}}\right)^{3 \times 3}, \sum_{k=1}^{3} \partial y_{k} \mathfrak{G}_{j k}(x, y, r)=0$ for $1 \leq j \leq 3, y \in B_{R_{1}}$, and $\mathfrak{G}(x, y, r)=\Lambda(x-y, r)$ for $y \in B_{S_{0}+\left(R_{0}-S_{0}\right) / 4}$.
Let $x \in B_{R_{0}}^{c}, q \in(1, \infty)$. Then the mapping $r \mapsto \mathfrak{G}(x, \cdot, r)(r \in[0, \infty))$ belongs to $C^{1}\left([0, \infty), W^{1, q}\left(B_{R_{1}}\right)^{3 \times 3}\right)$. Thus a function $G^{\prime} \in C^{0}\left([0, \infty), W^{1, q}\left(B_{R_{1}}\right)^{3 \times 3}\right)$ may be defined by the condition $\left\|(\mathfrak{G}(x, \cdot, r+h)-\mathfrak{G}(x, \cdot, r)) / h-G^{\prime}(r)\right\|_{1, q} \rightarrow 0(h \rightarrow 0)$ for $r \in[0, \infty)$. We write $\partial_{r} \mathfrak{G}(x, y, r)$ instead of $G^{\prime}(r)(y)\left(r \in[0, \infty), y \in B_{R_{1}}\right)$.

Let $r \in[0, \infty), q \in(1, \infty)$.
Let $\sigma \in\{0,1\}$, and define $L(x): B_{R_{1}} \mapsto \mathbb{R}^{3 \times 3}$ by $L(x)(y):=\partial_{r}^{\sigma} \mathfrak{G}(x, y, r)$ for $x \in B_{R_{0}}^{c}, y \in$ $B_{R_{1}}$. Then $L(x) \in C_{0}^{\infty}\left(B_{R_{1}}\right)^{3 \times 3} \cap W^{1, q}\left(B_{R_{1}}\right)^{3 \times 3}$ for $x \in B_{R_{0}}^{c}$, and $L$ as a mapping from $B_{R_{0}}^{c}$ into $W^{1, q}\left(B_{R_{1}}\right)^{3 \times 3}$ is partially differentiable on ${\overline{B_{R_{0}}}}^{c}$. Thus we may define $D_{m} L:{\overline{B_{R_{0}}}}^{c} \mapsto W^{1, q}\left(B_{R_{1}}\right)^{3 \times 3}$ by the condition $\left\|\left(L\left(x+h e_{m}\right)-L(x)\right) / h-D_{m} L(x)\right\|_{1, q} \rightarrow$ $0(h \rightarrow 0)$, for $m \in\{1,2,3\}, x \in{\overline{B_{R_{0}}}}^{c}$. Instead of $D_{m} L(x)(y)$, we write $\partial x_{m} \partial_{r}^{\sigma} \mathfrak{G}(x, y, r)$. Let $l \in\{1,2,3\}$ and define $\widetilde{L}(x): B_{R_{1}} \mapsto \mathbb{R}^{3 \times 3}$ by $\widetilde{L}(x)(y):=\partial y_{l} \mathfrak{G}(\underset{\sim}{x}, y, r)$ for $x \in$ $B_{R_{0}}^{c}, y \in B_{R_{1}}$. Then $\widetilde{L}(x) \in C_{0}^{\infty}\left(B_{R_{1}}\right)^{3 \times 3} \cap L^{q}\left(B_{R_{1}}\right)^{3 \times 3}$ for $x \in B_{R_{0}}^{c}$, and $\widetilde{L}$ considered as
 define $D_{m} \widetilde{L}:{\overline{B_{R_{0}}}}^{c} \mapsto L^{q}\left(B_{R_{1}}\right)^{3 \times 3}$ by the condition $\left\|\left(\widetilde{L}\left(x+h e_{m}\right)-\widetilde{L}(x)\right) / h-D_{m} \widetilde{L}(x)\right\|_{q} \rightarrow$ $0(h \rightarrow 0)\left(m \in\{1,2,3\}, x \in{\overline{B_{R_{0}}}}^{c}\right)$. Instead of $D_{m} \widetilde{L}(x)(y)$, we write $\partial x_{m} \partial y_{l} \mathfrak{G}(x, y, r)$.
Let $q \in(1, \infty), p \in[1, \infty]$. Then

$$
\begin{equation*}
\int_{B_{R_{1}}}\left|\partial_{x}^{\alpha} \partial_{t}^{\sigma} \partial_{y}^{\beta} \mathfrak{G}(x, y, t) \cdot V(y)\right| d y \leq \mathfrak{C}(|x| \nu(x))^{-(3+|\alpha|+\sigma) / 2}\|V\|_{q} \tag{4.7}
\end{equation*}
$$

for $V \in L^{q}\left(B_{R_{1}}\right)^{3}, t \in(0, \infty), x \in B_{R_{0}}^{c}, \alpha, \beta \in \mathbb{N}_{0}^{3}, \sigma \in\{0,1\}$ with $|\alpha| \leq 1,|\beta|+\sigma \leq 1$,
$\int_{0}^{t} \int_{B_{R_{1}}}\left|\partial_{x}^{\alpha} \partial_{t}^{\sigma} \partial_{y}^{\beta} \mathfrak{G}(x, y, t-s) \cdot v(y, s)\right| d y d s \leq \mathfrak{C}(|x| \nu(x))^{-(3+|\alpha|+\sigma) / 2+1 /\left(2 p^{\prime}\right)}\|v\|_{q, p ; t}$
for $t, x, \alpha, \beta, \sigma$ as in (4.7), and $v \in L^{p}\left(0, t, L^{q}\left(B_{R_{1}}\right)^{3}\right)$.
Proof: [19, (3.13), Lemma 3.11, 3.12, 3.13, Theorem 4.2].
We note a consequence of the preceding theorem.
Corollary 4.1 ([19, Corollary 4.2]) Let $\beta \in \mathbb{N}_{0}^{3}, \sigma \in\{0,1\}$ with $|\beta|+\sigma \leq 1$. Let $q \in(1, \infty)$, and let the function $v$ belong to $L_{l o c}^{1}\left([0, \infty), L^{q}\left(B_{R_{1}}\right)^{3}\right)$ and the function $V$ to $L^{q}\left(B_{R_{1}}\right)^{3}$. Define

$$
F(x, t):=\int_{0}^{t} \int_{B_{R_{1}}} \partial_{s}^{\sigma} \partial_{y}^{\beta} \mathfrak{G}(x, y, t-s) \cdot v(y, s) d y d s, \quad H(x, t):=\int_{B_{R_{1}}} \mathfrak{G}(x, y, t) \cdot V(y) d y
$$

for $x \in{\overline{B_{R_{0}}}}^{c}, t \in[0, \infty)$. Take a number $l \in\{1,2,3\}$. Then the derivatives $\partial x_{l} F(x, t)$ and $\partial x_{l} H(x, t)$ exist pointwise, and they equal $\int_{0}^{t} \int_{B_{R_{1}}} \partial x_{l} \partial_{s}^{\sigma} \partial_{y}^{\beta} \mathfrak{G}(x, y, t-s) \cdot v(y, s) d y d s$ and $\int_{B_{R_{1}}} \partial x_{l} \mathfrak{G}(x, y, t) \cdot V(y) d y$, respectively, for $x \in{\overline{B_{R_{0}}}}^{c}, t \in[0, \infty)$
It will be convenient to subsume a number of terms in a single operator, which we define here, and whose definition makes sense due to the preceding Corollary 4.1. The parameters $T_{0}, S_{0}, R_{0}, R_{1}$ and the set $\Omega$ appearing in the following were fixed at the beginning of Section 1.

Let $A \subset B_{S_{0}}$ be open and bounded, with Lipschitz boundary. Put $A_{R}:=B_{R} \backslash \bar{A}, Z_{R, T}:=$ $A_{R} \times(0, T)$ for $R \in\left[S_{0}, \infty\right), T \in(0, \infty]$. Let $\mathfrak{A} \subset \mathbb{R}^{3} \times \mathbb{R}$ be such that $Z_{R_{1}, T_{0}} \subset \mathfrak{A}$. Let $q \in(1, \infty)$ and let $v: \mathfrak{A} \mapsto \mathbb{R}^{3}$ with $v\left|Z_{R_{1}, T_{0}} \in C^{0}\left(\left[0, T_{0}\right), L^{q}\left(A_{R_{1}}\right)^{3}\right), v(s)\right| A_{R_{1}} \in$ $W_{l o c}^{1,1}\left(A_{R_{1}}\right)^{3}$ for $s \in\left(0, T_{0}\right)$ and $\nabla_{x} v \mid Z_{R_{1}, T_{0}} \in L_{l o c}^{1}\left(\left[0, T_{0}\right), L^{q}\left(A_{R_{1}}\right)^{9}\right)$. Then, for $t \in\left(0, T_{0}\right)$
and $x \in B_{R_{0}}^{c}$, we define

$$
\begin{align*}
& \mathfrak{K}_{R_{0}, S_{0}, \varphi_{0}, A, T_{0}}(v)(x, t):=\int_{0}^{t} \int_{\Omega_{R_{1}}}\left(\sum_{l=1}^{3} \partial y_{l} \mathfrak{G}(x, y, t-s) \cdot \partial_{y_{l}} v(y, s)\right.  \tag{4.9}\\
& \left.-\tau \partial y_{1} \mathfrak{G}(x, y, t-s) \cdot v(y, s)-\partial_{s} \mathfrak{G}(x, y, t-s) \cdot v(y, s)\right) d y d s+\int_{\Omega_{R_{1}}} \mathfrak{G}(x, y, 0) \cdot v(y, t) d y
\end{align*}
$$

Next we reproduce some decay estimates proved in [19], beginning with a decay estimate of $\mathfrak{K}_{R_{0}, S_{0}, \varphi_{0}, A, T_{0}}(v)$. We use the same notation as in (4.9).

Corollary 4.2 ([19, Corollary 4.3]) Let $p_{1}, p_{2} \in[1, \infty]$ and take $A, \mathfrak{A}, q$ as in (4.9). Then, forv $: \mathfrak{A} \mapsto \mathbb{R}^{3}$ such that $v\left|Z_{R_{1}, T_{0}} \in C^{0}\left(\left[0, T_{0}\right), L^{q}\left(A_{R_{1}}\right)^{3}\right), v(s)\right| A_{R_{1}} \in W_{\text {loc }}^{1,1}\left(A_{R_{1}}\right)^{3}$ $\left(s \in\left(0, T_{0}\right)\right)$ and $\nabla_{x} v \mid Z_{R_{1}, T_{0}} \in L^{p_{2}}\left(0, T_{0}, L^{q}\left(A_{R_{1}}\right)^{9}\right)$, and for $x \in{\overline{B_{R_{0}}}}^{c}, t \in\left(0, T_{0}\right), \alpha \in$ $\mathbb{N}_{0}^{3}$ with $|\alpha| \leq 1$, the term $\left|\partial_{x}^{\alpha} \mathfrak{K}_{R_{0}, S_{0}, \varphi_{0}, A, T_{0}}(v)(x, t)\right|$ is bounded by

$$
\mathfrak{C}\left(\left\|v\left|Z_{R_{1}, t}\left\|_{q, p_{1} ; t}+\right\| \nabla_{x} v\right| Z_{R_{1}, t}\right\|_{q, p_{2} ; t}+\left\|v(t) \mid \Omega_{R_{1}}\right\|_{q}\right) \max _{j \in\{1,2\}}(|x| \nu(x))^{-(3+|\alpha|) / 2+1 /\left(2 p_{j}^{\prime}\right)} .
$$

Lemma 4.8 ([19, Lemma 4.3]) Let $A, \mathfrak{A}, q$ be given as in (4.9), and let $p_{1}, p_{2} \in[1, \infty]$. Then, for $v: \mathfrak{A} \mapsto \mathbb{R}^{3}$ with $v\left|Z_{R_{1}, T_{0}} \in L^{p_{1}}\left(0, T_{0}, L^{q}\left(A_{R_{1}}\right)^{3}\right), v(s)\right| A_{R_{1}} \in W_{l o c}^{1,1}\left(A_{R_{1}}\right)^{3}$ $\left(s \in\left(0, T_{0}\right)\right)$ and $\nabla_{x} v \mid Z_{R_{1}, T_{0}} \in L^{p_{2}}\left(0, T_{0}, L^{q}\left(A_{R_{1}}\right)^{9}\right)$, as well as for $x \in B_{R_{0}}^{c}, t \in$ $\left(0, T_{0}\right), \alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha| \leq 2, l \in\{1,2,3\}$, the term $\left|\partial_{x}^{\alpha} \mathfrak{V}^{(\tau)}\left(n_{l}^{(A)} v\right)(x, t)\right|$ is bounded by

$$
\mathfrak{C}\left(\left\|v\left|Z_{R_{1}, t}\left\|_{q, p_{1} ; t}+\right\| \nabla_{x} v\right| Z_{R_{1}, t}\right\|_{q, p_{2} ; t}\right) \max _{j \in\{1,2\}}(|x| \nu(x))^{-(3+|\alpha|) / 2+1 /\left(2 p_{j}^{\prime}\right)},
$$

where $\left(n_{l}^{(A)} v\right)(y, s):=n_{l}^{(A)}(y) v(y, s)$ for $y \in \partial A, s \in\left(0, T_{0}\right)$.
Lemma 4.9 ([19, Lemma 4.4]) Recall that the Newton kernel $\mathfrak{N}$ was introduced in Theorem 3.1. Let $A \subset B_{S_{0}}$ open and bounded, with Lipschitz boundary. Let $q \in(1, \infty)$. Then the estimate $\left|\int_{\partial A}\left(\partial^{\alpha} \nabla \mathfrak{N}\right)(x-y)\left(n^{(A)} \cdot V\right)(y) d o_{y}\right| \leq \mathfrak{C}|x|^{-2-|\alpha|}\|V\|_{q}$ holds for $V \in L^{q}\left(A_{R_{1}}\right)^{3} \cap W^{1,1}\left(A_{R_{1}}\right)^{3}$ with div $V=0$, and for $t \in(0, \infty), x \in B_{R_{0}}^{c}, \alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha| \leq 1$. If $\int_{\partial A} n^{(A)} \cdot V d o_{y}=0$, the factor $|x|^{-2-|\alpha|}$ may be replaced by $|x|^{-3-|\alpha|}$.

Finally we restate the representation formula proved in [19] for regular solutions of (1.1). This formula will play a crucial role in the ensuing sections.

Theorem 4.3 ([19, Corollary 5.2]) Let $A \subset B_{S_{0}}$ be open and with Lipschitz boundary. Let $k_{0} \in \mathbb{N}, \varrho_{k} \in(1, \infty)$, and let $u^{(k)}$ belong to $C^{0}\left(\left[0, T_{0}\right), L^{\varrho_{k}}\left(\bar{A}^{c}\right)^{3}\right)$ and to $W_{\text {loc }}^{1,1}\left(0, T_{0}, L^{\varrho_{k}}\left(\bar{A}^{c}\right)^{3}\right)$, for $1 \leq k \leq k_{0}$. Put $u=\sum_{k=1}^{k_{0}} u^{(k)}$. (This means in particular that $u \mid A_{R} \times[0, \infty) \in C^{0}\left([0, \infty), L^{q}\left(A_{R}\right)^{3}\right)$ for $R \in\left[S_{0}, \infty\right), q \in(1, \infty)$ with $q \leq \varrho_{k}$ for $1 \leq$ $k \leq k_{0}$.) Let $\pi:\left(0, T_{0}\right) \mapsto W_{l o c}^{1,1}\left(\bar{A}^{c}\right), n_{0} \in \mathbb{N}, p_{j} \in(1, \infty)$ and $f^{(j)} \in L_{l o c}^{1}\left(0, T_{0}, L^{p_{j}}\left(\bar{A}^{c}\right)^{3}\right)$ for $1 \leq j \leq n_{0}$. Suppose that $u(s) \in W_{\text {loc }}^{2,1}\left(\bar{A}^{c}\right)^{3}$ for a. e. $t \in\left(0, T_{0}\right)$, and

$$
\begin{equation*}
u^{\prime}(s)-\Delta_{x} u(s)+\tau \partial x_{1} u(s)+\nabla_{x} \pi(s)=f(s), \operatorname{div}_{x} u(s)=0 \text { for a. e. } s \in\left(0, T_{0}\right), \tag{4.10}
\end{equation*}
$$

with $f:=\sum_{j=1}^{n_{0}} f^{(j)}, u^{\prime}:=\sum_{k=1}^{k_{0}}\left(u^{(k)}\right)^{\prime}$.
In addition suppose there is $q_{1} \in(1, \infty)$ such that $\nabla_{y} u(s) \mid B_{R_{0}}^{c} \in L^{q_{1}}\left(B_{R_{0}}^{c}\right)^{9}$ for a. e. $s \in\left(0, T_{0}\right)$ and $\nabla_{y} u \mid A_{R_{0}} \times\left(0, T_{0}\right) \in L_{\text {loc }}^{1}\left(\left[0, T_{0}\right), L^{q_{1}}\left(A_{R_{0}}\right)^{9}\right)$. Assume that $u \mid A_{R_{0}} \times\left(0, T_{0}\right) \in$
$L_{l o c}^{1}\left(0, T_{0}, W^{2,1}\left(A_{R_{0}}\right)^{3}\right), \pi \mid A_{R_{0}} \times\left(0, T_{0}\right) \in L_{l o c}^{1}\left(0, T_{0}, W^{1,1}\left(A_{R_{0}}\right)\right)$, and there are numbers $a_{0} \in \mathbb{N}, \gamma_{j} \in(1, \infty)$ as well as functions $\pi^{(j)}:\left(0, T_{0}\right) \mapsto L_{\text {loc }}^{1}\left({\overline{B_{R_{0}}}}^{c}\right)$ for $1 \leq j \leq a_{0}$ such that

$$
\pi(s) \mid{\overline{B_{R_{0}}}}^{c}=\sum_{j=1}^{a_{0}} \pi^{(j)}(s), \quad \sum_{j=1}^{a_{0}} \int_{\overline{B_{R_{0}}}}\left(\left|\pi^{(j)}(y, s)\right|[(1+|y|) \ln (2+|y|)]^{-1}\right)^{\gamma_{j}} d y<\infty
$$

for a. e. $s \in\left(0, T_{0}\right)$. Then there is a subset $\mathfrak{T}_{T_{0}}$ of $\left(0, T_{0}\right)$ of measure zero such that for $t \in\left(0, T_{0}\right) \backslash \mathfrak{T}_{T_{0}}$, the equation

$$
\begin{align*}
& u(x, t)=\mathfrak{R}^{(\tau)}(f)(x, t)+\mathfrak{I}^{(\tau)}(u(0))(x, t)-\sum_{l=1}^{3} \partial x_{l} \mathfrak{V}^{(\tau)}\left(n_{l}^{(A)} u\right)(x, t)  \tag{4.11}\\
& \quad-\int_{\partial A}(\nabla \mathfrak{N})(x-y)\left(n^{(A)}(y) \cdot u(y, t)\right) d o_{y}+\mathfrak{K}_{R_{0}, S_{0}, \varphi_{0}, A, T_{0}}(u)(x, t) \\
& \quad-\int_{A_{R_{1}}} \mathfrak{G}_{R_{0}, S_{0}, \varphi_{0}}(x, y, t) \cdot u(y, 0) d y-\int_{0}^{t} \int_{A_{R_{1}}} \mathfrak{G}_{R_{0}, S_{0}, \varphi_{0}}(x, y, t-s) \cdot f(y, s) d y d s
\end{align*}
$$

holds for $x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t}$, where $N_{t}$ is some subset of ${\overline{B_{R_{0}}}}^{c}$ with measure zero. The function $\mathfrak{K}_{R_{0}, S_{0}, \varphi_{0}, A, T_{0}}(u)$ is defined in (4.17), and the function $\mathfrak{G}_{R_{0}, S_{0}, \varphi_{0}}$ in Theorem 4.3.
Next we restate a tool which will help to extend the integral representation (4.11) to solutions of (4.10) that are more general than the ones in the preceding theorem.
Lemma 4.10 ([19, Lemma 5.3]) Let $A \subset B_{S_{0}}$ be open and bounded, with Lipschitz boundary. Let $k_{0} \in \mathbb{N}, \varrho_{k} \in(1, \infty)$, $u^{(k)} \in C^{0}\left(\left[0, T_{0}\right), L^{\varrho_{k}}\left(\bar{A}^{c}\right)^{3}\right)$ for $1 \leq k \leq k_{0}$, and put $u=\sum_{k=1}^{k_{0}} u^{(k)}$. In particular $u \mid A_{R} \times[0, \infty) \in C^{0}\left([0, \infty), L^{q}\left(A_{R}\right)^{3}\right)$ for $R \in$ $\left[S_{0}, \infty\right), q \in\left(1, \min \left\{\varrho_{k}: 1 \leq k \leq k_{0}\right\}\right]$. Further suppose that $u(s) \in W_{l o c}^{1,1}\left(\bar{A}^{c}\right)^{3}$ for a. $e . \quad s \in\left(0, T_{0}\right)$, div $v_{x} u=0$, and $\nabla_{x} u \mid A_{R_{1}} \times\left(0, T_{0}\right) \in L_{l o c}^{1}\left(\left[0, T_{0}\right), L^{q_{1}}\left(A_{R_{1}}\right)^{9}\right)$ for some $q_{1} \in(1, \infty)$. Furthermore let $n_{0} \in \mathbb{N}, p_{j} \in(1, \infty)$, $f^{(j)} \in L_{\text {loc }}^{1}\left(\left[0, T_{0}\right), L^{p_{j}}\left(\bar{A}^{c}\right)^{3}\right)$ for $1 \leq j \leq n_{0}$, and put $f:=\sum_{j=1}^{n_{0}} f^{(j)}$.
Fix some function $\zeta \in C^{\infty}(\mathbb{R})$ with $\zeta|(-\infty, 1]=0, \zeta|[2, \infty)=1,0 \leq \zeta \leq 1$ and $\zeta^{\prime} \geq 0$. For $\epsilon \in(0, \infty), r \in \mathbb{R}$, define $\zeta_{\epsilon}(r):=\zeta(r / \epsilon)$. Put $u_{\epsilon}^{(k)}(s):=\zeta_{\epsilon}(s) u^{(k)}(s)$, $u_{\epsilon}(s):=$ $\zeta_{\epsilon}(s) u(s), f_{\epsilon}(s):=\zeta_{\epsilon}(s) f(s)$ and $g_{\epsilon}(s):=f_{\epsilon}(s)+\zeta_{\epsilon}^{\prime}(s) u(s)$ for $s \in\left(0, T_{0}\right), \epsilon \in(0, \infty), 1 \leq$ $k \leq k_{0}$.
Let $t \in\left(0, T_{0}\right)$. Suppose there is some $\epsilon_{0} \in(0, \infty)$ such that for $\epsilon \in\left(0, \epsilon_{0}\right]$, equation (4.11) holds with $u$, $f$ replaced by $u_{\epsilon}$ and $g_{\epsilon}$, respectively, if $x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t, \epsilon}$ for some subset $N_{t, \epsilon}$ of ${\overline{B_{R_{0}}}}^{c}$ of measure zero. (This means in particular that the second from last term in (4.11) vanishes.) Then there is some zero-measure set $N_{t} \subset{\overline{B_{R_{0}}}}^{c}$ such that equation (4.11) remains valid for $u$ and $f$ if $x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t}$.

## 5 Weak solutions of the Oseen system: representation formula and pointwise spatial decay.

We begin by introducing a $L^{q}$-weak solution to (1.1), with assumptions chosen in such a way that it will be possible to prove the representation formula (4.11) for this solution.

In order to deduce a decay estimate from this formula, we will require somewhat stronger conditions (see Theorem 5.2). We again point out that $T_{0}, S_{0}, R_{0}, R_{1}, \varphi_{0}$ and $\Omega$ are introduced at the beginning of Section 2.
Fix parameters $n_{0}, m_{0} \in \mathbb{N}, \widetilde{p}, q_{0}, q_{1}, p_{1}, \ldots, p_{n_{0}}, \varrho_{1}, \ldots, \varrho_{m_{0}} \in(1, \infty)$, as well as functions $u:\left(0, T_{0}\right) \mapsto W_{l o c}^{1,1}\left(\bar{\Omega}^{c}\right)^{3}, f^{(j)} \in L_{l o c}^{1}\left(\left[0, T_{0}\right), L^{p_{j}}\left(\bar{\Omega}^{c}\right)^{3}\right)$ for $1 \leq j \leq n_{0}, G^{(l)} \in$ $C^{0}\left(\left[0, T_{0}\right), L^{\varrho_{l}}\left(\overline{B_{S_{0}}}\right)^{3}\right)$ for $1 \leq l \leq m_{0}, U_{0} \in L^{\tilde{p}}\left(\bar{\Omega}^{c}\right)^{3}$ with the following properties:
$u \mid \Omega_{S_{0}} \times\left(0, T_{0}\right) \in L_{l o c}^{1}\left(\left[0, T_{0}\right), L^{q_{0}}\left(\Omega_{S_{0}}\right)^{3}\right), \operatorname{div}_{x} u(t)=0$ and $u(t) \mid{\overline{B_{S_{0}}}}^{c}=\sum_{l=1}^{m_{0}} G^{(l)}(t)$ for $t \in\left(0, T_{0}\right), \nabla_{x} u \in L_{l o c}^{1}\left(\left[0, T_{0}\right), L^{q_{1}}\left(\bar{\Omega}^{c}\right)^{3}\right)$,

$$
\begin{align*}
& \int_{0}^{T_{0}} \int_{\bar{\Omega}^{c}}\left(-\gamma^{\prime}(t) u(t) \cdot \vartheta+\gamma(t)\left[\nabla_{x} u(t) \cdot \nabla \vartheta+\tau \partial x_{1} u(t) \cdot \vartheta-f(t) \cdot \vartheta\right]\right) d x d t  \tag{5.1}\\
& \quad-\gamma(0) \int_{\bar{\Omega}^{c}} U_{0} \cdot \vartheta d x=0 \quad \text { for } \gamma \in C_{0}^{\infty}\left(\left[0, T_{0}\right)\right), \vartheta \in C_{0, \sigma}^{\infty}\left(\bar{\Omega}^{c}\right), \quad \text { with } f:=\sum_{j=1}^{n_{0}} f^{(j)} .
\end{align*}
$$

There will be no claim that the preceding assumptions imply $u(0)=U_{0}$. But as it will turn out, they are sufficient to derive a representation formula. The case $m_{0}>1$ or $n_{0}>1$ and the distinction between the functions $G^{(l)}$ on ${\overline{B_{S_{0}}}}^{c} \times(0, \infty)$ with $u(t) \mid{\overline{B_{S_{0}}}}^{c}=\sum_{l=1}^{m_{0}} G^{(l)}(t)$ on the one hand and $u \mid \Omega_{S_{0}} \times\left(0, T_{0}\right)$ on the other one are introduced in order to take account of some technical difficulties arising when the theory in this work is applied to a nonlinear problem in [20], to the proof of Theorem 6.2 below, and in order to avoid any assumptions on $\partial \Omega$ stronger than Lipschitz regularity. Any reader who wants to avoid these technicalities may consider the case that $\partial \Omega$ is smooth, $u \in C^{0}\left(\left[0, T_{0}\right), L^{q_{0}}\left(\bar{\Omega}^{c}\right)^{3}\right), m_{0}=$ $n_{0}=1$ and $p_{1}=q_{1}$. All the main difficulties of our proofs will then still be present.
Until further notice, we suppose that $T_{0}=\infty$. It is only at the end of the present section that we will turn to the case $T_{0}<\infty$ (Corollary 5.2).
In the ensuing lemma, we cut off the functions $u, f^{(j)}$ and $G^{(l)}$ near the instant $t=0$, and then present the version of (5.1) satisfied by the extension of these modified functions to the whole real axis. In the lemma after that, we apply Friedrich's mollifier to these extensions. Both lemmas mainly serve to introduce notation and collect obvious facts which constitute the basis for the rest of this section.
Lemma 5.1 Fix some $q \in(1, \infty)$ with

$$
q \leq \min \left(\left\{q_{0}, q_{1}, 5 / 4\right\} \cup\left\{\varrho_{l}: 1 \leq l \leq m_{0}\right\} \cup\left\{p_{j}: 1 \leq l \leq n_{0}\right\}\right) .
$$

Then the function $u \mid \Omega_{R} \times(0, \infty)$ belongs to $L_{\text {loc }}^{1}\left([0, \infty), L^{q}\left(\Omega_{R}\right)^{3}\right)$ and $u \mid A_{R, S_{0}} \times[0, \infty)$ to $C^{0}\left([0, \infty), L^{q}\left(A_{R, S_{0}}\right)^{3}\right)$, for $R \in\left[S_{0}, \infty\right)$.
Choose functions $\zeta_{\epsilon}$ for $\epsilon \in(0, \infty)$ as in Lemma 4.10 For $s, \epsilon \in(0, \infty)$, define $u_{\epsilon}(s):=$ $\zeta_{\epsilon}(s) u(s), f_{\epsilon}^{(j)}(s):=\zeta_{\epsilon}(s) f^{(j)}(s)\left(1 \leq j \leq n_{0}\right)$ and $G_{\epsilon}^{(l)}(s):=\zeta_{\epsilon}(s) G^{(l)}(s)\left(1 \leq l \leq m_{0}\right)$,
$p_{n_{0}+1}:=q_{0}, f_{\epsilon}^{\left(n_{0}+1\right)}:=\zeta_{\epsilon}^{\prime}(s) \chi_{\Omega_{S_{0}}} u(s), \quad p_{n_{0}+1+l}:=\varrho_{l}, f_{\epsilon}^{\left(n_{0}+1+l\right)}(s):=\zeta_{\epsilon}^{\prime}(s) \widetilde{G}^{(l)}(s)$ for $1 \leq l \leq m_{0}$, where $\widetilde{G}^{(l)}$ denotes the zero extension of $G^{(l)}:{\overline{B_{S_{0}}}}^{c} \times[0, \infty) \mapsto \mathbb{R}^{3}$ to $\bar{\Omega}^{c} \times[0, \infty)$. Then $f_{\epsilon}^{\left(n_{0}+k\right)} \in L_{l o c}^{1}\left([0, \infty), L^{p_{n}+k}\left(\bar{\Omega}^{c}\right)^{3}\right)$ for $1 \leq k \leq m_{0}+1$, and $\zeta_{\epsilon}^{\prime}(s) u(s)=\sum_{k=1}^{m_{0}+1} f_{\epsilon}^{\left(n_{0}+k\right)}(s)(s \in(0, \infty))$.

The functions $u_{\epsilon}(t), f_{\epsilon}^{(j)}(t), G_{\epsilon}^{(l)}(t), \widetilde{G}_{\epsilon}^{(l)}(t)(t \in(0, \infty))$ are extended by zero to the real axis $\mathbb{R}\left(1 \leq j \leq n_{0}+m_{0}+1,1 \leq l \leq m_{0}\right)$, without change of notation. Then $f_{\epsilon}^{(j)} \in L_{\text {loc }}^{1}\left(\mathbb{R}, L^{p_{j}}\left(\bar{\Omega}^{c}\right)^{3}\right)$ for $j \in\left\{1, \ldots, n_{0}+m_{0}+1\right\}, G_{\epsilon}^{(l)} \in C^{0}\left(\mathbb{R}, L^{\varrho_{l}}\left({\overline{B_{S_{0}}}}^{c}\right)^{3}\right)$ and $\widetilde{G}_{\epsilon}^{(l)} \in C^{0}\left(\mathbb{R}, L^{\varrho_{l}}\left(\bar{\Omega}^{c}\right)^{3}\right)$ for $1 \leq l \leq m_{0}$. In addition $u_{\epsilon} \mid \Omega_{R} \times \mathbb{R} \in L_{l o c}^{1}\left(\mathbb{R}, L^{q}\left(\Omega_{R}\right)^{3}\right)$ and $u_{\epsilon} \mid A_{R, S_{0}} \times \mathbb{R} \in C^{0}\left(\mathbb{R}, L^{q}\left(A_{R, S_{0}}\right)^{3}\right)$ for $R \in\left[S_{0}, \infty\right)$, $u_{\epsilon}=\chi_{\Omega_{S_{0}}} u_{\epsilon}+\sum_{j=1}^{m_{0}} \widetilde{G}_{\epsilon}^{(l)}$ and $u_{\epsilon} \mid{\overline{B_{S_{0}}}}^{c} \times \mathbb{R}=\sum_{l=1}^{m_{0}} G_{\epsilon}^{(l)}, u(s) \in W_{l o c}^{1,1}\left(\bar{\Omega}^{c}\right)^{3}$ for $s \in \mathbb{R}, \operatorname{div}_{x} u_{\epsilon}=0, \nabla_{x} u_{\epsilon} \in$ $L_{l o c}^{1}\left(\mathbb{R}, L^{q_{1}}\left(\bar{\Omega}^{c}\right)^{9}\right)$,

$$
\begin{align*}
& \int_{\mathbb{R}} \int_{\bar{\Omega}^{c}}\left(-\gamma^{\prime}(t) u_{\epsilon}(t) \cdot \vartheta+\gamma(t)\left[\nabla_{x} u_{\epsilon}(t) \cdot \nabla \vartheta+\tau \partial x_{1} u_{\epsilon}(t) \cdot \vartheta\right.\right.  \tag{5.2}\\
& \left.\left.\quad-\sum_{j=1}^{n_{0}+m_{0}+1} f_{\epsilon}^{(j)}(t) \cdot \vartheta\right]\right) d x d t=0 \quad \text { for } \gamma \in C_{0}^{\infty}(\mathbb{R}), \vartheta \in C_{0, \sigma}^{\infty}\left(\bar{\Omega}^{c}\right) .
\end{align*}
$$

We remark that the functions $f_{\epsilon}^{\left(n_{0}+1\right)}, \ldots, f_{\epsilon}^{\left(n_{0}+m_{0}+1\right)}$ are perturbation terms which arise in (5.2) due to the cut-off function $\zeta_{\epsilon}$ and the decomposition of $u(s)$ into the sum $\chi_{\Omega_{S_{0}}} u(s)+$ $\sum_{l=1}^{m_{0}} \widetilde{G}^{(l)}$. In the ensuing lemma, we use Friedrich's mollifier with respect to the time variable, as defined in the passage preceding Theorem 2.6. This mollifier is applied to functions with domain $\mathbb{R}$ and values in Banach spaces, in this context spaces $L^{r}(A)^{3}$, for certain $r \in(1, \infty)$ and certain open subsets $A$ of $\mathbb{R}^{3}$. The functions with index $\epsilon$ are those introduced in Lemma 5.1.
Concerning $u_{\epsilon}$, there is no Banach space $B$ which we deem useful and which is such that $u_{\epsilon} \in L_{l o c}^{1}(\mathbb{R}, B)$. So we cannot directly apply our definition of Friedrich's mollifier to $u_{\epsilon}$. Instead we use the equation $u_{\epsilon}=\chi_{\Omega_{S_{0}}} u_{\epsilon}+\sum_{j=1}^{m_{0}} \widetilde{G}_{\epsilon}^{(l)}$ in order to define a more regular version of $u_{\epsilon}$, denoted by $u_{\epsilon, \delta}$. Note that due to the relation $\nabla_{x} u_{\epsilon} \in L_{l o c}^{1}\left(\mathbb{R}, L^{q_{1}}\left(\bar{\Omega}^{c}\right)^{3}\right)$, no such problem arises with $\nabla_{x} u_{\epsilon}$.
Lemma 5.2 Let $\epsilon, \delta \in(0, \infty)$. Put $f_{\epsilon, \delta}^{(j)}:=\left(f_{\epsilon}^{(j)}\right)_{\delta}\left(1 \leq j \leq n_{0}+m_{0}+1\right), G_{\epsilon, \delta}^{(l)}:=$ $\left(G_{\epsilon}^{(l)}\right)_{\delta}, \quad \widetilde{G}_{\epsilon, \delta}^{(l)}:=\left(\widetilde{G}_{\epsilon}^{(l)}\right)_{\delta}\left(1 \leq l \leq m_{0}\right)$. Then, for $j, l$ as before, the relations $f_{\epsilon, \delta}^{(j)} \in$ $C^{\infty}\left(\mathbb{R}, L^{p_{j}}\left(\bar{\Omega}^{c}\right)^{3}\right), G_{\epsilon, \delta}^{(l)} \in C^{\infty}\left(\mathbb{R}, L^{\varrho_{l}}\left({\overline{B_{S_{0}}}}^{c}\right)^{3}\right)$ and $\widetilde{G}_{\epsilon, \delta}^{(l)} \in C^{\infty}\left(\mathbb{R}, L^{\varrho_{l}}\left(\bar{\Omega}^{c}\right)^{3}\right)$ hold. Define $u_{\epsilon, \delta}:=\left(\chi_{\Omega_{S_{0}}} u_{\epsilon}\right)_{\delta}+\sum_{l=1}^{m_{0}} \widetilde{G}_{\epsilon, \delta}^{(l)}$. Then $u_{\epsilon, \delta} \mid \Omega_{R} \times \mathbb{R}=\left(u_{\epsilon} \mid \Omega_{R} \times \mathbb{R}\right)_{\delta} \in C^{\infty}\left(\mathbb{R}, L^{q}\left(\Omega_{R}\right)^{3}\right)$ for $R \in\left[S_{0}, \infty\right)$, and $u_{\epsilon, \delta} \mid{\overline{B_{S_{0}}}}^{c} \times \mathbb{R}=\sum_{l=1}^{m_{0}} G_{\epsilon, \delta}^{(l)}$. Moreover $u_{\epsilon, \delta}(t) \in W_{l o c}^{1,1}\left(\bar{\Omega}^{c}\right)^{3}$ for $t \in \mathbb{R}, \partial x_{k} u_{\epsilon, \delta}=\left(\partial x_{k} u_{\epsilon}\right)_{\delta} \in C^{\infty}\left(\mathbb{R}, L^{q_{1}}\left(\bar{\Omega}^{c}\right)^{3}\right)$ for $1 \leq k \leq 3$, and div $u_{x, \delta}=0$.
Define $s_{0}:=n_{0}+2 m_{0}+2, f_{\epsilon, \delta}^{\left(n_{0}+m_{0}+2\right)}:=-\left(\chi \Omega_{S_{0}} u_{\epsilon}\right)_{\delta}^{\prime}, p_{n_{0}+m_{0}+2}:=q, f_{\epsilon, \delta}^{\left(n_{0}+m_{0}+2+l\right)}:=$ $-\left(\widetilde{G}_{\epsilon, \delta}^{(l)}\right)^{\prime}$ and $p_{n_{0}+m_{0}+2+l}:=\varrho_{l}$ for $1 \leq l \leq m_{0}$. Then $f_{\epsilon, \delta}^{(j)} \in C^{\infty}\left(\mathbb{R}, L^{p_{j}}\left(\bar{\Omega}^{c}\right)^{3}\right)$ for $n_{0}+m_{0}+2 \leq j \leq s_{0}$.
For $t \in \mathbb{R}$, the function $u_{\epsilon, \delta}(t)$ satisfies (3.1) with $A, U, F$ replaced by $\bar{\Omega}^{c}, u_{\epsilon, \delta}(t)$ and $\sum_{j=1}^{s_{0}} f_{\epsilon, \delta}^{(j)}(t)$, respectively, and with $\lambda=0$.
The last statement of the lemma is the crucial one. It means that $u_{\epsilon, \delta}(t)$ is a weak solution in $\bar{\Omega}^{c}$ of the stationary Oseen system, with right-hand side $F$ as indicated.
Proof of Lemma 5.2: The regularity properties of $f_{\epsilon, \delta}^{(j)}, G_{\epsilon, \delta}^{(l)}, \widetilde{G}_{\epsilon, \delta}^{(l)}$ and $\left(\chi \Omega_{R} u_{\epsilon}\right)_{\delta}$ follow
with Theorem 2.6 from the properties of $f_{\epsilon}^{(j)}, G_{\epsilon}^{(l)}, \widetilde{G}_{\epsilon}^{(l)}$ and $\chi_{\Omega_{R}} u_{\epsilon}$, respectively, as stated in Lemma $5.1\left(1 \leq j \leq s_{0}, 1 \leq l \leq m_{0}\right.$ and $R \in\left[S_{0}, \infty\right)$ ). Lemma 2.2 and the equation $u_{\epsilon} \mid{\overline{B_{S_{0}}}}^{c} \times \mathbb{R}=\sum_{l=1}^{m_{0}} G_{\epsilon}^{(l)}(t)$ (Lemma 5.1) yield that $\left(u_{\epsilon} \mid \Omega_{R} \times \mathbb{R}\right)_{\delta}=u_{\epsilon, \delta} \mid \Omega_{R} \times \mathbb{R}$ for $R \in\left[S_{0}, \infty\right)$ and $u_{\epsilon, \delta} \mid \overline{{S_{0}}^{c}} \times \mathbb{R}=\sum_{l=1}^{m_{0}} G_{\epsilon, \delta}^{(l)}$.
In order to check existence of weak derivatives of $u_{\epsilon, \delta}$ with respect to the space variables, let $k \in\{1,2,3\}, t \in \mathbb{R}, \psi \in C_{0}^{\infty}\left(\mathbb{R}^{3}\right)^{3}$ and $R \in\left[S_{0}, \infty\right)$ with $\operatorname{supp}(\psi) \subset \Omega_{R}$. Then $\int_{\Omega^{c}} \partial_{k} \psi \cdot u_{\epsilon, \delta}(t) d x=\int_{\Omega_{R}} \partial_{k} \psi \cdot\left(u_{\epsilon} \mid \Omega_{R} \times \mathbb{R}\right)_{\delta}(t) d x$ by what was shown above. The operator $V \mapsto \int_{\Omega_{R}} \partial_{k} \psi \cdot V d x\left(V \in L^{q}\left(\Omega_{R}\right)^{3}\right)$ is linear and bounded, so Theorem 2.4 yields that $\int_{\Omega_{R}} \partial_{k} \psi \cdot\left(u_{\epsilon} \mid \Omega_{R} \times \mathbb{R}\right)_{\delta}(t) d x=\int_{\mathbb{R}} \rho_{\delta}(t-s) \int_{\Omega_{R}} \partial_{k} \psi \cdot u_{\epsilon}(s) d x d s$, with $\rho_{\delta}$ defined in the passage preceding Theorem 2.6. Since $\int_{\Omega_{R}} \partial_{k} \psi \cdot u_{\epsilon}(s) d x=-\int_{\Omega_{R}} \psi \cdot \partial x_{k} u_{\epsilon}(s) d x$ for a. e. $s \in \mathbb{R}$, and by a similar reasoning as before, we further get that the right-hand side of the second from last equation equals $-\int_{\bar{\Omega}^{c}} \psi \cdot\left(\partial x_{k} u_{\epsilon}\right)_{\delta}(t) d x$. Altogether we arrive at the equation $\int_{\bar{\Omega}^{c}} \partial_{k} \psi \cdot u_{\epsilon, \delta}(t) d x=-\int_{\bar{\Omega}^{c}} \psi \cdot\left(\partial x_{k} u_{\epsilon}\right)_{\delta}(t) d x$. Thus we have shown that $u_{\epsilon, \delta}(t) \in W_{l o c}^{1,1}\left(\bar{\Omega}^{c}\right)^{3}$ for $t \in \mathbb{R}$ and $\partial x_{k} u_{\epsilon, \delta}=\left(\partial x_{k} u_{\epsilon}\right)_{\delta}$ for $1 \leq k \leq 3$. It follows with Theorem 2.6 and Lemmma 5.1 that $\partial x_{k} u_{\epsilon, \delta} \in C^{\infty}\left(\mathbb{R}, L^{q_{1}}\left(\bar{\Omega}^{c}\right)^{3}\right)$. The same references yield $f_{\epsilon, \delta}^{(j)} \in C^{\infty}\left(\mathbb{R}, L^{p_{j}}\left(\bar{\Omega}^{c}\right)^{3}\right)$ for $n_{0}+m_{0}+2 \leq j \leq s_{0}$.
In order to determine the PDE satisfied by $u_{\epsilon, \delta}$, take $t \in \mathbb{R}$ and $\vartheta \in C_{0, \sigma}^{\infty}\left(\bar{\Omega}^{c}\right)$. Our starting point is equation (5.2) with $\gamma(s)=\rho_{\delta}(t-s)$ for $s \in \mathbb{R}$. (We recall again that $\rho_{\delta}$ is introduced in the passage preceding Theorem 2.6.) For $\sigma \in\{0,1\}, k \in\{1,2,3\}, p \in\left\{q_{1}, q\right\} \cup\left\{p_{j}\right.$ : $\left.1 \leq j \leq n_{0}+m_{0}+1\right\} \cup\left\{\varrho_{l}: 1 \leq l \leq m_{0}\right\}$, the functional $V \mapsto \int_{\bar{\Omega}^{c}} V \cdot \partial_{k}^{\sigma} \vartheta d x\left(V \in L^{p}\left(\bar{\Omega}^{c}\right)^{3}\right)$ is linear and bounded. Therefore by Theorem 2.4 and because $\left(\partial x_{k} u_{\epsilon}\right)_{\delta}=\partial x_{k} u_{\epsilon, \delta}$,

$$
\begin{align*}
& \int_{\mathbb{R}} \rho_{\delta}(t-s) \int_{\bar{\Omega}^{c}} \partial x_{k} u_{\epsilon}(s) \cdot \partial_{k} \vartheta d x d s=\int_{\bar{\Omega}^{c}} \partial x_{k} u_{\epsilon, \delta}(t) \cdot \partial_{k} \vartheta d x,  \tag{5.3}\\
& \int_{\mathbb{R}} \rho_{\delta}(t-s) \int_{\bar{\Omega}^{c}} \partial x_{1} u_{\epsilon}(s) \cdot \vartheta d x d s=\int_{\bar{\Omega}^{c}} \partial x_{1} u_{\epsilon, \delta}(t) \cdot \vartheta d x,  \tag{5.4}\\
& \int_{\mathbb{R}} \rho_{\delta}(t-s) \int_{\bar{\Omega}^{c}} f_{\epsilon}^{(j)}(s) \cdot \vartheta d x d s=\int_{\bar{\Omega}^{c}} f_{\epsilon, \delta}^{(j)}(t) \cdot \vartheta d x, \tag{5.5}
\end{align*}
$$

for $k \in\{1,2,3\}, 1 \leq j \leq n_{0}+m_{0}+1$. Put $K^{(1)}:=\chi_{\Omega_{S_{0}}} \cdot u_{\epsilon}, K^{(l)}:=\widetilde{G}_{\epsilon}^{(l)}$ for $2 \leq l \leq m_{0}+1$. Then, using Theorem 2.4, the equation $u_{\epsilon}=\sum_{l=1}^{m_{0}+1} K^{(l)}$ (Lemma 5.1) and Theorem 2.6, we obtain

$$
\begin{align*}
& \int_{\mathbb{R}} \int_{\bar{\Omega}^{c}} \partial_{s}\left(\rho_{\delta}(t-s)\right) u_{\epsilon}(s) \cdot \vartheta d x d s=-\int_{\bar{\Omega}^{c}}\left(\int_{\mathbb{R}} \rho_{\delta}^{\prime}(t-s) u_{\epsilon}(s) d s\right) \cdot \vartheta d x  \tag{5.6}\\
& -\int_{\bar{\Omega}^{c}}\left(\int_{\mathbb{R}} \rho_{\delta}^{\prime}(t-s) \sum_{l=1}^{m_{0}+1} K^{(l)}(s) d s\right) \cdot \vartheta d x=\sum_{j=m_{0}+n_{0}+2}^{s_{0}} \int_{\bar{\Omega}^{c}} f_{\epsilon, \delta}^{(j)}(t) \cdot \vartheta d x .
\end{align*}
$$

We may deduce from (5.2) and (5.3) - (5.6) that $u_{\epsilon, \delta}(t)$ satisfies the variational problem in (3.1) with $A, U, F$ replaced as indicated in the lemma, and with $\lambda=0$. Again by the equation $\partial x_{k} u_{\epsilon, \delta}=\left(\partial x_{k} u_{\epsilon}\right)_{\delta}$, and because $\operatorname{div}_{x} u_{\epsilon}=0$ (Lemma 5.2), we obtain $\operatorname{div}_{x} u_{\epsilon, \delta}=0$. Thus $u_{\epsilon, \delta}$ satisfies (3.1) in full.

Now we are in a position to deal with the main difficulty of this section, that is, to show that $u_{\epsilon, \delta}$ satisfies the representation formula (4.11). The notation from Lemma 5.1 and 5.2 will be used without further notice.

Theorem 5.1 Let $\epsilon, \delta \in(0, \infty)$. Then, for any $t \in(0, \infty)$, there is a set $N_{t} \subset{\overline{B_{R_{0}}}}^{c}$ of measure zero such that equation (4.11) holds for $x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t}$, with $A=B_{S_{0}}$ and $u$, $f$ replaced by $u_{\epsilon, \delta} \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty)$ and $\sum_{j=1}^{n_{0}+m_{0}+1} f_{\epsilon, \delta}^{(j)} \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty)$, respectively.
Proof: Since $\bar{\Omega} \subset B_{S_{0}}$ by our assumptions at the beginning of this section, we may fix $S_{1} \in\left(0, S_{0}\right)$ with $\bar{\Omega} \subset S_{S_{1}}$. Let $t \in \mathbb{R}$. For brevity, we set $U:=u_{\epsilon, \delta}(t), F^{(j)}:=f_{\epsilon, \delta}^{(j)}(t)$ for $1 \leq j \leq s_{0}$. Due to the choice of $q$ in Lemma 5.1 and by the properties of $u_{\epsilon, \delta}$ and $f_{\epsilon, \delta}^{(j)}$ (Lemma 5.2), we have $U \in W_{l o c}^{1, q}\left(\bar{\Omega}^{c}\right)^{3}, \nabla U \in L^{q_{1}}\left(\bar{\Omega}^{c}\right)^{9}, F^{(j)} \in L_{l o c}^{q}\left(\bar{\Omega}^{c}\right)^{3} \cap L^{p_{j}}\left(\bar{\Omega}^{c}\right)^{3}$ for $1 \leq j \leq s_{0}$, and the function $U$ satisfies (3.1) with $\lambda=0$ and with $A, F$ replaced by $\bar{\Omega}^{c}$ and $\sum_{j=1}^{s_{0}} F^{(j)}$, respectively. Thus, by Theorem 3.2, there is $\Pi \in L_{l o c}^{p}\left(\bar{\Omega}^{c}\right)^{3}$ such that (3.2) holds with $A=\bar{\Omega}^{c}, \lambda=0$ and $F=\sum_{j=1}^{s_{0}} F^{(j)}$. Since the latter equation remains valid when we replace $\Pi$ by $\Pi+c$, for any number $c \in \mathbb{R}$, we may assume in addition that $\int_{A_{S_{0}, S_{1}}} \Pi d x=0$. Equation (3.2) is valid in particular for $\vartheta \in C_{0}^{\infty}\left(A_{S_{0}, S_{1}}\right)^{3}$. Therefore, according to Theorem 3.2, this function $\Pi$ satisfies the estimate

$$
\begin{equation*}
\left\|\Pi \mid A_{S_{0}, S_{1}}\right\|_{q} \leq C\left(q, S_{0}, S_{1}\right)\left(\left\|\nabla U\left|A_{S_{0}, S_{1}}\left\|_{q}+\right\| \sum_{j=1}^{s_{0}} F^{(j)}\right| A_{S_{0}, S_{1}}\right\|_{q}\right) \tag{5.7}
\end{equation*}
$$

Due to (3.2) and the regularity properties indicated above for $U, F^{(j)}$ and $\Pi$, Theorem 3.3 now yields that $U \in W_{l o c}^{2, q}\left(\bar{\Omega}^{c}\right)^{3}, \Pi \in W_{l o c}^{1, q}\left(\bar{\Omega}^{c}\right)$ and the pair $(U, \Pi)$ solves $(3.4)$ in $\bar{\Omega}^{c}$ with $\lambda=0$ and $F=\sum_{j=1}^{s_{0}} F^{(j)}$. We may choose a function $\widetilde{\varphi} \in C^{\infty}\left(\mathbb{R}^{3}\right)$ with $0 \leq \widetilde{\varphi} \leq$ 1, $\widetilde{\varphi}\left|B_{S_{1}+\left(S_{0}-S_{1}\right) / 4}=0, \widetilde{\varphi}\right| B_{S_{1}+\left(S_{0}-S_{1}\right) / 2}^{c}=1$. Since $\bar{\Omega} \subset B_{S_{1}}$ and $\operatorname{supp}\left(\partial^{\alpha} \widetilde{\varphi}\right) \subset A_{S_{0}, S_{1}}$ for $\alpha \in \mathbb{N}_{0}^{3}$ with $\alpha \neq 0$, and by the regularity properties of $U$ and $\Pi$ just derived, the function

$$
\begin{equation*}
\widetilde{F}:=-\Delta \widetilde{\varphi} U-2\left(\nabla \widetilde{\varphi} \cdot \nabla U_{j}\right)_{1 \leq j \leq 3}+\tau \partial_{1} \widetilde{\varphi} U+\Pi \nabla \widetilde{\varphi} \tag{5.8}
\end{equation*}
$$

belongs to $W^{1, q}\left(\mathbb{R}^{3}\right)$, as does $\nabla \widetilde{\varphi} \cdot U$. In addition $\operatorname{supp}(\widetilde{F}) \cup \operatorname{supp}(\nabla \widetilde{\varphi} \cdot U) \subset A_{S_{0}, S_{1}}$. Put $\widetilde{W}:=\mathfrak{N} *(\nabla \widetilde{\varphi} \cdot U)$; see Theorem 3.1. Since $q<3 / 2$ by the choice of $q$ in Lemma 5.1, we know by Theorem 3.1 that $\widetilde{W}$ is well defined and

$$
\begin{align*}
& \widetilde{W} \in W_{l o c}^{3, q}\left(\mathbb{R}^{3}\right), \nabla \widetilde{W} \in L^{(1 / q-1 / 3)^{-1}}\left(\mathbb{R}^{3}\right)^{3}, \partial_{l} \partial_{m} \widetilde{W} \in W^{1, q}\left(\mathbb{R}^{3}\right)^{9}  \tag{5.9}\\
& \Delta \widetilde{W}=-\nabla \widetilde{\varphi} \cdot U, \quad\left\|\partial_{k} \partial_{l} \partial_{m} \widetilde{W}\right\|_{q}+\left\|\partial_{k} \partial_{l} \widetilde{W}\right\|_{q} \leq C(q)\|\nabla \widetilde{\varphi} \cdot U\|_{1, q} \quad(1 \leq k, l, m \leq 3)
\end{align*}
$$

We put

$$
\begin{align*}
& \widetilde{U}:=\widetilde{\varphi} U+\nabla \widetilde{W}, \widetilde{\Pi}:=\widetilde{\varphi} \Pi, \widetilde{F}^{(j)}:=\widetilde{\varphi} F^{(j)} \text { for } 1 \leq j \leq s_{0}  \tag{5.10}\\
& p_{s_{0}+1}:=q, \widetilde{F}^{\left(s_{0}+1\right)}:=\widetilde{F}-\Delta \nabla \widetilde{W}+\tau \partial_{1} \nabla \widetilde{W}
\end{align*}
$$

Note that $\operatorname{div} \nabla \widetilde{W}=\Delta \widetilde{W}=-\nabla \widetilde{\varphi} \cdot U$, hence $\operatorname{div} \widetilde{U}=0$. As a consequence of these observations and definitions, we have $\widetilde{U} \in W_{l o c}^{2, q}\left(\mathbb{R}^{3}\right)^{3}, \widetilde{\Pi} \in W_{l o c}^{1, q}\left(\mathbb{R}^{3}\right), \widetilde{F}^{(j)} \in L^{p_{j}}\left(\mathbb{R}^{3}\right)^{3}$ for $1 \leq j \leq s_{0}+1, \widetilde{\varphi} U\left|B_{S_{0}}^{c}=U\right| B_{S_{0}}^{c}, \widetilde{\Pi}\left|B_{S_{0}}^{c}=\Pi\right| B_{S_{0}}^{c}$, and the pair $(\widetilde{U}, \widetilde{\Pi})$ solves $(3.4)$ in $\mathbb{R}^{3}$ with $\lambda=0$ (stationary Oseen system) and with $F=\sum_{j=1}^{s_{0}+1} \widetilde{F}^{(j)}$. Since $\widetilde{F}^{(j)} \in L^{p_{j}}\left(\mathbb{R}^{3}\right)^{3}$ for $1 \leq j \leq s_{0}+1$, Theorem 3.5 yields functions $W^{(j)} \in W_{2}^{2, p_{j}}\left(\mathbb{R}^{3}\right)^{3}, \Pi^{(j)} \in W_{1}^{1, p_{j}}\left(\mathbb{R}^{3}\right)$ such that $\partial_{1} W^{(j)} \in L^{p_{j}}\left(\mathbb{R}^{3}\right)^{3}$, equation (3.4) is satisfied in $\mathbb{R}^{3}$ with $U, \Pi, F$ replaced by $W^{(j)}, \Pi^{(j)}, \widetilde{F}^{(j)}$, respectively, and the inequality

$$
\begin{equation*}
\left\|\partial_{l} \partial_{m} W^{(j)}\right\|_{p_{j}}+\left\|\Pi^{(j)}\right\|_{W_{1}^{1, p_{j}}\left(\mathbb{R}^{3}\right)} \leq C\left(\tau, p_{j}\right)\left\|\widetilde{F}^{(j)}\right\|_{p_{j}}\left(1 \leq l, m \leq 3,1 \leq j \leq s_{0}+1\right) \tag{5.11}
\end{equation*}
$$

holds. Then $\bar{U}:=\widetilde{U}-\sum_{j=1}^{s_{0}+1} W^{(j)} \in W_{l o c}^{2, q}\left(\mathbb{R}^{3}\right)^{3}, \bar{\Pi}:=\widetilde{\Pi}-\sum_{j=1}^{s_{0}+1} \Pi^{(j)} \in W_{l o c}^{1, q}\left(\mathbb{R}^{3}\right)$, and $-\Delta \bar{U}+\tau \partial_{1} \bar{U}+\nabla \bar{\Pi}=0, \operatorname{div} \bar{U}=0$ in $\mathbb{R}^{3}$. It follows with Theorem 3.4 that $\bar{\Pi}$ is a $C^{\infty}$ function. Let us specify in which way the function $\bar{U}$ satisfies the assumptions of Theorem 3.6. Put $q_{j}(0):=q_{j}(1):=p_{j}$ for $j \in\left\{1, \ldots, s_{0}+1\right\}$. Then the definition of the space $W_{2}^{2, p_{j}}\left(\mathbb{R}^{3}\right)$ implies that $\int_{B_{R_{0}}^{c}}\left(\left|\partial^{\alpha} W^{(j)}(x)\right|\left[(1+|x|)^{2-|\alpha|} \ln (2+|x|)\right]^{-1}\right)^{q_{j}(|\alpha|)} d x<\infty$ for $\alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha| \leq 1$ and $1 \leq j \leq s_{0}+1$. Put

$$
\begin{equation*}
\widetilde{q}:=\max \left(\left\{p_{j}: 1 \leq j \leq n_{0}\right\} \cup\left\{(1 / q-1 / 3)^{-1}, q_{0}\right\} \cup \max \left(\left\{\varrho_{l}: 1 \leq l \leq m_{0}\right\}\right)\right. \tag{5.12}
\end{equation*}
$$

and $\epsilon_{0}:=1 /(2 \widetilde{q})$. By the choice of $p_{n_{0}+1}, \ldots, p_{s_{0}+1}$ (Lemma 5.1, 5.2, (5.10)), this means in particular that $\widetilde{q} \geq p_{j}$ for $1 \leq j \leq s_{0}+1$, and $\widetilde{q} \geq q$. Obviously $1+|x| \leq C\left(R_{0}\right)|x|$ and $\ln (2+|x|) \leq C\left(R_{0}, \epsilon_{0}\right)|x|^{\epsilon_{0}}$ for $x \in B_{R_{0}}^{c}$. Thus we may conclude for $1 \leq j \leq s_{0}+1$ that

$$
\begin{equation*}
\int_{B_{R_{0}}^{c}}\left(\left|\partial^{\alpha} W^{(j)}(x)\right||x|^{-2-\epsilon_{0}}\right)^{q_{j}(|\alpha|)} d x<\infty \quad \text { for } \alpha \in \mathbb{N}_{0}^{3} \text { with }|\alpha| \leq 1 \tag{5.13}
\end{equation*}
$$

Next we put $W^{\left(s_{0}+2\right)}:=-\nabla \widetilde{W}, q_{s_{0}+2}(0):=(1 / q-1 / 3)^{-1}, q_{s_{0}+2}(1):=q$. Then, due to (5.9), the relation in (5.13) is valid for $j=s_{0}+2$ as well. Since $\nabla U \in L^{q_{1}}\left(\bar{\Omega}^{c}\right)^{3}$ as mentioned above, $\widetilde{\varphi} \mid B_{S_{0}}^{c}=1$ and $R_{0}>S_{0}$, we get $\int_{B_{R_{0}}^{c}}\left(|\nabla(\widetilde{\varphi} U)(x)||x|^{-2-\epsilon_{0} \mid}\right)^{q_{1}} d x<\infty$. We further note that by the definition of $\widetilde{U}$ (see (5.10)), $\bar{U}$ and $W^{\left(s_{0}+2\right)}$, the equation $\nabla \bar{U}=$ $\nabla(\widetilde{\varphi} U)-\sum_{j=1}^{s_{0}+2} \nabla W^{(j)}$ holds. Thus, if in Theorem 3.6 we replace $U, m_{0},\left(V^{(m)}\right)_{1 \leq m \leq m_{0}}$ by $\bar{U}, s_{0}+3,\left(\left(\partial_{\mu}\left[\widetilde{\varphi} U_{\sigma}\right]\right)_{1 \leq \mu, \sigma \leq 3},\left(\partial_{\mu} W_{\sigma}^{(1)}\right)_{1 \leq \mu, \sigma \leq 3}, \ldots,\left(\partial_{\mu} W_{\sigma}^{\left(s_{0}+2\right)}\right)_{1 \leq \mu, \sigma \leq 3}\right)$, respectively, the assumptions on $\nabla U$ are satisfied with $r=2+\epsilon_{0}$. Concerning those on $U$, we recall that by the definition of $U$ and $u_{\epsilon, \delta}$, we have $\widetilde{\varphi} U=-\sum_{j=s_{0}+3}^{s_{0}+3+m_{0}} W^{(j)}$, with $W^{\left(s_{0}+3\right)}:=$ $-\widetilde{\varphi}\left(\chi_{\Omega_{S_{0}}} u_{\epsilon}\right)_{\delta}(t), W^{\left(s_{0}+3+l\right)}:=-\widetilde{\varphi} \widetilde{G}_{\epsilon, \delta}^{(l)}(t)$ for $1 \leq l \leq m_{0}$. Lemma 5.2 yields $W^{(j)} \in$ $L^{p_{j}(0)}\left(\mathbb{R}^{3}\right)^{3}$ for $s_{0}+3 \leq j \leq s_{0}+3+m_{0}$, with $p_{s_{0}+3}(0):=q, p_{s_{0}+3+l}(0):=\varrho_{l}$ for $1 \leq l \leq m_{0}$. It follows the relation in (5.13) holds for $s_{0}+3 \leq j \leq s_{0}+3+m_{0}$ with $\alpha=0$. We may conclude with the definition of $\widetilde{U}$ and $W^{\left(s_{0}+2\right)}$ that $\widetilde{U}=-\sum_{j=s_{0}+2}^{s_{0}+3+m_{0}} W^{(j)}$. Now the definition of $\bar{U}$ yields that $\bar{U}=-\sum_{j=1}^{s_{0}+m_{0}+3} W^{(j)}$. Therefore the assumptions on $U$ in Theorem 3.6 are satisfied with $\bar{U}, s_{0}+m_{0}+3,\left(-W^{(k)}\right)_{1 \leq k \leq s_{0}+m_{0}+3}$ in the role of $U, k_{0},\left(U^{(k)}\right)_{1 \leq k \leq k_{0}}$, respectively, and with $r=2+\epsilon_{0}$. At this point Theorem 3.6 yields that $\bar{U}$ is a polynomial. By the definitions in (5.12) and the choice of the exponents $q_{j}(0)$, we have $\widetilde{q}=\max \left\{q_{j}(0): 1 \leq j \leq s_{0}+3+m_{0}\right\}$ and $\epsilon_{0} \in(0,1 / \widetilde{q})$. Thus it follows from (5.13) with $\alpha=0,1 \leq j \leq s_{0}+m_{0}+3$ and from Lemma 3.1 that the degree of $\bar{U}$ cannot exceed 1. As a first consequence of this result, we get $\partial^{\alpha} \bar{U}=0$ for $\alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha|=2$. Since $\bar{U}=\widetilde{\varphi} U-\sum_{j=1}^{s_{0}+2} W^{(j)}, \widetilde{\varphi} \mid B_{S_{0}}^{c}=1$ and $W^{\left(s_{0}+2\right)}=-\nabla \widetilde{W}$, we thus find

$$
\begin{equation*}
\partial^{\alpha} U\left|B_{S_{0}}^{c}=-\partial^{\alpha}(\nabla \widetilde{W})+\sum_{j=1}^{s_{0}+1} \partial^{\alpha} W^{(j)}\right| B_{S_{0}}^{c} \quad \text { for } \alpha \in \mathbb{N}_{0}^{3} \text { with }|\alpha|=2 \tag{5.14}
\end{equation*}
$$

It further follows there is $c_{1} \in \mathbb{R}$ with $\partial_{1} \bar{U}=c_{1}$. On the other hand, $\widetilde{\varphi} \mid B_{S_{0}}^{c}=1$ and $S_{0}<R_{0}$, so $\partial_{1} U \mid B_{R_{0}}^{c} \in L^{q_{1}}\left(B_{R_{0}}\right)^{3}$. Moreover $\partial_{1} W^{\left(s_{0}+2\right)}=\partial_{1} \nabla \widetilde{W} \in L^{q}\left(\mathbb{R}^{3}\right)^{3}$ (see (5.9)), and $\partial_{1} W^{(j)} \in L^{p_{j}}\left(\mathbb{R}^{3}\right)^{3}$ for $1 \leq j \leq s_{0}+1$ because $W^{(j)} \in W_{2}^{2, p_{j}}\left(\mathbb{R}^{3}\right)^{3}$. It follows from Lemma 3.2 and the equation $\bar{U}=\widetilde{\varphi} U-\sum_{j=1}^{s_{0}+2} W^{(j)}$ that $c_{1}$ vanishes. Recalling that
$-\Delta \bar{U}+\tau \partial_{1} \bar{U}+\nabla \bar{\Pi}=0$, we thus obtain $\nabla \bar{\Pi}=0$. Since $\bar{\Pi} \in C^{\infty}\left(\mathbb{R}^{3}\right)$, as shown above, and because $\bar{\Omega}^{c}$ is connected, we may conclude there is $c \in \mathbb{R}$ with $\bar{\Pi}=c$. But $\bar{\Pi}=$ $\widetilde{\varphi} \Pi-\sum_{j=1}^{s_{0}+1} \Pi^{(j)}$ by the definition of $\bar{\Pi}$, so we arrive at the equation

$$
\begin{equation*}
\Pi-c\left|{\overline{B_{S_{0}}}}^{c}=\sum_{j=1}^{s_{0}+1} \Pi^{(j)}\right|{\overline{B_{S_{0}}}}^{c} . \tag{5.15}
\end{equation*}
$$

Hölder's inequality yields the estimate $\left\|\Pi^{(j)}\left|A_{R_{0}, S_{0}}\left\|_{1} \leq \mathfrak{C}\right\| \Pi^{(j)}\right| A_{R_{0}, S_{0}}\right\|_{p_{j}}$, and it is obvious that $\left\|\Pi^{(j)} \mid A_{R_{0}, S_{0}}\right\|_{p_{j}} \leq \mathfrak{C}\left[\int_{A_{R_{0}, S_{0}}}\left(\left|\Pi^{(j)}(x)\right|[(1+|x|) \ln (2+|x|)]^{-1}\right)^{p_{j}} d x\right]^{1 / p_{j}}$, so we arrive at the inequality $\left\|\Pi^{(j)} \mid A_{R_{0}, S_{0}}\right\|_{1} \leq \mathfrak{C}\left\|\Pi^{(j)}\right\|_{W_{1}^{1, p_{j}}\left(\mathbb{R}^{3}\right)}$, for $1 \leq j \leq s_{0}+1$. Obviously $\left\|\nabla \Pi^{(j)}\left|A_{R_{0}, S_{0}}\left\|_{1} \leq \mathfrak{C}\right\| \nabla \Pi^{(j)}\right| A_{R_{0}, S_{0}}\right\|_{p_{j}} \leq \mathfrak{C}\left\|\nabla \Pi^{(j)}\right\|_{p_{j}} \leq \mathfrak{C}\left\|\Pi^{(j)}\right\|_{W_{1}^{1, p_{j}}\left(\mathbb{R}^{3}\right)}$. The constants in the estimates following (5.15) only depend on $S_{0}, R_{0}$ and $p_{j}$, with the relevant index $j$. These estimates and (5.11) yield $\left\|\Pi^{(j)} \mid A_{R_{0}, S_{0}}\right\|_{1,1} \leq C\left(S_{0}, R_{0}, p_{j}, \tau\right)\left\|\widetilde{F}^{(j)}\right\|_{p_{j}}$ for $1 \leq j \leq s_{0}+1$. Due to (5.15), we thus get

$$
\begin{equation*}
\left\|\Pi-c \mid A_{R_{0}, S_{0}}\right\|_{1,1} \leq C\left(S_{0}, R_{0}, p_{1}, \ldots, p_{s_{0}+1}, \tau\right) \sum_{j=1}^{s_{0}+1}\left\|\widetilde{F}^{(j)}\right\|_{p_{j}} \tag{5.16}
\end{equation*}
$$

Once more starting with Hölder's inequality, then using (5.9) and the relations $q \leq q_{1}$ (Lemma 5.1) and $\partial_{l} U=\left(\partial x_{l} u_{\epsilon}\right)_{\delta}(t)$ (Lemma 5.2), we get

$$
\begin{align*}
& \left\|\partial_{k} \partial_{l} \partial_{m} \widetilde{W}\left|A_{R_{0}, S_{0}}\left\|_{1} \leq \mathfrak{C}\right\| \partial_{k} \partial_{l} \partial_{m} \widetilde{W}\right| A_{R_{0}, S_{0}}\right\|_{q} \leq \mathfrak{C}\left\|\partial_{k} \partial_{l} \partial_{m} \widetilde{W}\right\|_{q} \leq \mathfrak{C}\|\nabla \widetilde{\varphi} \cdot U\|_{1, q}  \tag{5.1}\\
& \leq \mathfrak{C}\left\|U \mid A_{S_{0}, S_{1}}\right\|_{1, q} \leq \mathfrak{C}\left(\left\|U \mid \Omega_{S_{0}}\right\|_{q}+\|\nabla U\|_{q_{1}}\right) \leq \mathfrak{C}\left(\left\|U \mid \Omega_{S_{0}}\right\|_{q}+\sum_{l=1}^{3}\left\|\left(\partial x_{l} u_{\epsilon}\right)_{\delta}(t)\right\|_{q_{1}}\right) .
\end{align*}
$$

Obviously with (5.11), $\left\|\partial^{\alpha} W^{(j)}\left|A_{R_{0}, S_{0}}\left\|_{1} \leq \mathfrak{C}\right\| \partial^{\alpha} W^{(j)}\right| A_{R_{0}, S_{0}}\right\|_{p_{j}} \leq \mathfrak{C}\left\|\widetilde{F}^{(j)}\right\|_{p_{j}}$ for $1 \leq j \leq$ $s_{0}+1, \alpha \in \mathbb{N}_{0}^{3},|\alpha|=2$. Now we may conclude with (5.14) that

$$
\begin{equation*}
\left\|\partial^{\alpha} U \mid A_{R_{0}, S_{0}}\right\|_{1} \leq \mathfrak{C}\left(\left\|U \mid \Omega_{S_{0}}\right\|_{q}+\sum_{l=1}^{3}\left\|\left(\partial x_{l} u_{\epsilon}\right)_{\delta}(t)\right\|_{q_{1}}+\sum_{j=1}^{s_{0}+1}\left\|\widetilde{F}^{(j)}\right\|_{p_{j}}\right) \tag{5.18}
\end{equation*}
$$

for $\alpha$ as before. As in the last two inequalities in (5.17), we obtain that $\left\|U \mid A_{R_{0}, S_{0}}\right\|_{1,1}$ is bounded by the right-hand side of (5.17). The constants in the inequalities from (5.16) onwards only depend on $\tau, S_{0}, R_{0}, q, q_{1}$ or on the respective exponent $p_{j}$. This observation, (5.18) and (5.16) yield

$$
\begin{align*}
& \left\|U\left|A_{R_{0}, S_{0}}\left\|_{2,1}+\right\| \Pi-c\right| A_{R_{0}, S_{0}}\right\|_{1,1}  \tag{5.19}\\
& \leq C\left(\tau, S_{0}, R_{0}, q, q_{1}, p_{1}, \ldots, p_{s_{0}+1}\right)\left(\left\|U \mid \Omega_{S_{0}}\right\|_{q}+\sum_{l=1}^{3}\left\|\left(\partial x_{l} u_{\epsilon}\right) \delta(t)\right\|_{q_{1}}+\sum_{j=1}^{s_{0}+1}\left\|\widetilde{F}^{(j)}\right\|_{p_{j}}\right) .
\end{align*}
$$

Let us estimate $\left\|\widetilde{F}^{\left(s_{0}+1\right)}\right\|_{p_{s_{0}+1}}$. (See (5.10) for the definition of $\widetilde{F}^{\left(s_{0}+1\right)}$ and $p_{s_{0}+1}$.) By (5.9) and the last four inequalities in (5.17), we get

$$
\begin{equation*}
\|\Delta \nabla \widetilde{W}\|_{q}+\left\|\partial_{1} \nabla \widetilde{W}\right\|_{q} \leq C(q)\|\nabla \widetilde{\varphi} \cdot U\|_{1, q} \leq \mathfrak{C}\left(\left\|U \mid \Omega_{S_{0}}\right\|_{q}+\sum_{l=1}^{3}\left\|\left(\partial x_{l} u_{\epsilon}\right)_{\delta}(t)\right\|_{q_{1}}\right) . \tag{5.20}
\end{equation*}
$$

With $\widetilde{F}$ defined in (5.8), we further find that $\|\widetilde{F}\|_{q} \leq \mathfrak{C}\left(\left\|U\left|A_{S_{0}, S_{1}}\left\|_{1, q}+\right\| \Pi\right| A_{S_{0}, S_{1}}\right\|_{q}\right)$. The quantity $\left\|\Pi \mid A_{S_{0}, S_{1}}\right\|_{q}$ was evaluated in (5.7). In view of the right-hand side of this
inequality, and by an estimate of $\left\|U \mid A_{S_{0}, S_{1}}\right\|_{1, q}$ as in the last two inequalities in (5.17), we may conclude that $\|\widetilde{F}\|_{q} \leq \mathfrak{C}\left(\left\|U\left|\Omega_{S_{0}}\left\|_{q}+\sum_{l=1}^{3}\right\|\left(\partial x_{l} u_{\epsilon}\right)_{\delta}(t)\left\|_{q_{1}}+\sum_{j=1}^{s_{0}}\right\| F^{(j)}\right| A_{S_{0}, S_{1}}\right\|_{q}\right)$. It follows with the preceding inequality, (5.20), (5.10) and the choice of $q$ in Lemma 5.1 that

$$
\begin{equation*}
\left\|\widetilde{F}^{\left(s_{0}+1\right)}\right\|_{p_{s_{0}+1}}=\left\|\widetilde{F}^{\left(s_{0}+1\right)}\right\|_{q} \leq \mathfrak{C}\left(\left\|U \mid \Omega_{S_{0}}\right\|_{q}+\sum_{l=1}^{3}\left\|\left(\partial x_{l} u_{\epsilon}\right)_{\delta}(t)\right\|_{q_{1}}+\sum_{j=1}^{s_{0}}\left\|F^{(j)}\right\|_{p_{j}}\right) \tag{5.21}
\end{equation*}
$$

Since $0 \leq \widetilde{\varphi} \leq 1$, we deduce from the definition of $\widetilde{F}^{(j)}$ in (5.10) that $\left\|\widetilde{F}^{(j)}\right\|_{p_{j}} \leq\left\|F^{(j)}\right\|_{p_{j}}$ for $1 \leq j \leq s_{0}$. Returning to the notation $u_{\epsilon, \delta}(t)$ and $f_{\epsilon, \delta}^{(j)}(t)$ introduced at the beginning of this proof to replace $U$ and $F^{(j)}$, respectively $\left(1 \leq j \leq s_{0}\right)$, and putting $\pi_{\epsilon, \delta}(t):=\Pi-c$, we thus obtain from (5.21) and (5.19) that

$$
\begin{align*}
& \left\|u_{\epsilon, \delta}(t)\left|A_{R_{0}, S_{0}}\left\|_{2,1}+\right\| \pi_{\epsilon, \delta}(t)\right| A_{R_{0}, S_{0}}\right\|_{1,1}  \tag{5.22}\\
& \leq \mathfrak{C}\left(\left\|u_{\epsilon, \delta}(t) \mid \Omega_{R_{0}}\right\|_{q}+\sum_{l=1}^{3}\left\|\left(\partial x_{l} u_{\epsilon}\right)_{\delta}(t)\right\|_{q_{1}}+\sum_{j=1}^{s_{0}}\left\|f_{\epsilon, \delta}^{(j)}(t)\right\|_{p_{j}}\right)
\end{align*}
$$

In addition, equation (5.15) takes the form $\pi_{\epsilon, \delta}(t)\left|{\overline{B_{S_{0}}}}^{c}=\sum_{j=1}^{s_{0}+1} \Pi^{(j)}\right|{\overline{B_{S_{0}}}}^{c}$. Since $\Pi^{(j)} \in$
 $\infty\left(1 \leq j \leq s_{0}+1\right)$. By what was shown for $U$ and $\Pi$ at the beginning of this proof, the relations $u_{\epsilon, \delta}(t) \in W_{l o c}^{2, q}\left(\bar{\Omega}^{c}\right)^{3}$ and $\pi_{\epsilon, \delta}(t) \in W_{l o c}^{1, q}\left(\bar{\Omega}^{c}\right)$ hold, and the pair $\left(u_{\epsilon, \delta}(t), \pi_{\epsilon, \delta}(t)\right)$ solves (3.4) in $\bar{\Omega}^{c}$, with $\lambda=0$ (stationary Oseen system) and $F=\sum_{j=1}^{s_{0}} f_{\epsilon, \delta}^{(j)}(t)$.
Recall that $G_{\epsilon, \delta}^{(l)} \in C^{\infty}\left(\mathbb{R}, L^{\varrho_{l}}\left({\overline{B_{S_{0}}}}^{c}\right)^{3}\right)$ for $1 \leq l \leq m_{0}, u_{\epsilon, \delta}(t) \mid{\overline{B_{S_{0}}}}^{c}=\sum_{j=1}^{m_{0}} G_{\epsilon, \delta}^{(l)}(t)$ and $f_{\epsilon, \delta}^{(j)} \in C^{\infty}\left(\mathbb{R}, L^{p_{j}}\left(\bar{\Omega}^{c}\right)^{3}\right)$ for $1 \leq j \leq n_{0}+m_{0}+1$ (Lemma 5.2). Since $f_{\epsilon, \delta}^{\left(n_{0}+m_{0}+2\right)}(t) \mid{\overline{B_{S_{0}}}}^{c}=$ 0 and $f_{\epsilon, \delta}^{\left(n_{0}+m_{0}+2+l\right)}(t) \mid{\overline{B_{S_{0}}}}^{c}=\left(G_{\epsilon, \delta}^{(l)}\right)^{\prime}(t)$ for $1 \leq l \leq m_{0}$ (Lemma 2.2, Theorem 2.6), we obtain that $\left(u_{\epsilon, \delta} \mid{\overline{B_{S_{0}}}}^{c} \times \mathbb{R}\right)^{\prime}(t)=\sum_{j=n_{0}+m_{0}+2}^{n_{0}+2 m_{0}+2} f_{\epsilon, \delta}^{(j)}(t) \mid{\overline{B_{S_{0}}}}^{c}$. At this point the relation $s_{0}=n_{0}+2 m_{0}+2$ and the fact that the pair $\left(u_{\epsilon, \delta}(t), \pi_{\epsilon, \delta}(t)\right)$ solves (3.4) in $\bar{\Omega}^{c}$ with $\lambda=0$ and $F=\sum_{j=1}^{s_{0}} f_{\epsilon, \delta}^{(j)}(t)$ and with $t$ arbitrary in $\mathbb{R}$, allow us to conclude that the pair $\left(u_{\epsilon, \delta}\left|{\overline{B_{S_{0}}}}^{c} \times(0, \infty), \pi_{\epsilon, \delta}\right|{\overline{B_{S_{0}}}}^{c} \times(0, \infty)\right)$ solves the time-dependent Oseen system (4.10) with $A=B_{S_{0}}$ and $f=\sum_{j=1}^{n_{0}+m_{0}+1} f_{\epsilon, \delta}^{(j)} \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty)$, without any exceptional values of $t$. In particular the equation $\operatorname{div}_{x} u(t)=0$ holds for any $t \in(0, \infty)$, without exceptional values.
Lemma 5.2 states in particular that the function $\nabla_{x} u_{\epsilon, \delta}$ belongs to $C^{\infty}\left(\mathbb{R}, L^{q_{1}}\left(\bar{\Omega}^{c}\right)^{9}\right)$ and $u_{\epsilon, \delta} \mid \Omega_{R_{0}} \times \mathbb{R}$ to $C^{\infty}\left(\mathbb{R}, L^{q}\left(\Omega_{R_{0}}\right)^{3}\right)$. Thus the right-hand side of (5.22) is integrable with respect to $t \in J$ for any bounded interval $J \subset \mathbb{R}$. Therefore we may conclude from (5.22) that $u_{\epsilon, \delta} \mid A_{R_{0}, S_{0}} \times(0, \infty) \in L_{l o c}^{1}\left(\mathbb{R}, W^{2,1}\left(A_{R_{0}, S_{0}}\right)^{3}\right)$ and $\pi_{\epsilon, \delta} \mid A_{R_{0}, S_{0}} \times \mathbb{R} \in$ $L_{l o c}^{1}\left(\mathbb{R}, W^{1,1}\left(A_{R_{0}, S_{0}}\right)\right)$.
Thus all assumptions of Theorem 4.3 are verified with obvious replacements, in particular with $B_{S_{0}}$ in the role of $A$. The theorem now follows from Theorem 4.3.

Next we show that (4.11) remains valid when $\delta$ tends to zero.
Lemma 5.3 Let $\epsilon, t \in(0, \infty)$. Then there is a set $N_{t, \epsilon} \subset{\overline{B_{R_{0}}}}^{c}$ of measure zero such that
equation (4.11) holds for $x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t, \epsilon}$, with $A=B_{S_{0}}$ and $u$, $f$ replaced by $u_{\epsilon} \mid{\overline{B_{S_{0}}}}^{c} \times$ $(0, \infty)$ and $\sum_{j=1}^{n_{0}+m_{0}+1} f_{\epsilon}^{(j)} \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty)$, respectively.
Proof: Abbreviate $Z_{r}:={\overline{B_{S_{0}}}}^{c} \times(0, r)$ for $r \in(0, \infty], \mathfrak{G}:=\mathfrak{G}_{R_{0}, S_{0}, \varphi_{0}}$, with $\mathfrak{G}_{R_{0}, S_{0}, \varphi_{0}}$ defined in Theorem 4.2. Recall that $R_{1}:=\left(S_{0}+R_{0}\right) / 2$, as defined at the beginning of this section. Let $l \in\left\{1, \ldots, m_{0}\right\}, k \in\{1,2,3\}, j \in\left\{1, \ldots, n_{0}+m_{0}+1\right\}, v \in$ $\left\{G_{\epsilon}^{(l)}, \partial x_{k} u_{\epsilon}, f_{\epsilon}^{(j)}\right\}$. Put $p:=\varrho_{l}, A:={\overline{B_{S_{0}}}}^{c}$ if $v=G_{\epsilon}^{(l)}, p:=q_{1}, A:=\bar{\Omega}^{c}$ if $v=\partial x_{k} u_{\epsilon}$, and $p:=p_{j}, A:=\bar{\Omega}^{c}$ if $v=f_{\epsilon}^{(j)}$. Due to the choice of $p$ and $A$, and by Lemma 5.1, the relation $v \in L_{l o c}^{1}\left(\mathbb{R}, L^{p}(A)^{3}\right)$. holds. As a consequence $\chi_{(-1, t+1)} v \in L^{1}\left(\mathbb{R}, L^{p}(A)^{3}\right)$, and thus by Theorem 2.6, $\left\|\chi_{(-1, t+1)} v-\left(\chi_{(-1, t+1)} v\right)_{\delta}\right\|_{p, 1 ; \mathbb{R}} \rightarrow 0(\delta \downarrow 0)$. But for $\delta \in(0,1], s \in(0, t)$, we have $\left(\chi_{(-1, t+1)} v\right)_{\delta}(s)=v_{\delta}(s)$, so $\left\|v_{\delta}-v \mid A \times(0, t)\right\|_{p, 1 ; t} \rightarrow 0(\delta \downarrow 0)$. Thus we have shown that
$\left\|G_{\epsilon, \delta}^{(l)}-G_{\epsilon}^{(l)}\left|{\overline{B_{S_{0}}}}^{c} \times(0, t)\left\|_{\varrho_{l}, 1 ; t} \rightarrow 0,\right\|\left(\partial x_{k} u_{\epsilon}\right)_{\delta}-\partial x_{k} u_{\epsilon}\right| \bar{\Omega}^{c} \times(0, t)\right\|_{q_{1}, 1 ; t} \rightarrow 0$,
$\left\|f_{\epsilon, \delta}^{(j)}-f_{\epsilon}^{(j)} \mid \bar{\Omega}^{c} \times(0, t)\right\|_{p_{j}, 1 ; t} \rightarrow 0(\delta \downarrow 0)$ for $1 \leq l \leq m_{0}, 1 \leq k \leq 3,1 \leq j \leq n_{0}+m_{0}+1$.
In particular, since $\partial x_{k} u_{\epsilon, \delta}=\left(\partial x_{k} u_{\epsilon}\right)_{\delta}$ according to Lemma 5.2, we deduce from (5.23) that $\left\|\partial x_{k} u_{\epsilon, \delta}-\partial x_{k} u_{\epsilon} \mid \bar{\Omega}^{c} \times(0, t)\right\|_{q_{1}, 1 ; t} \rightarrow 0(\delta \downarrow 0)$ for $1 \leq k \leq 3$. Recall that $u_{\epsilon} \mid{\overline{B_{S_{0}}}}^{c} \times$ $\mathbb{R}=\sum_{l=1}^{m_{0}} G_{\epsilon}^{(l)}$ (Lemma 5.1), and $u_{\epsilon, \delta} \mid{\overline{B_{S_{0}}}}^{c} \times \mathbb{R}=\sum_{l=1}^{m_{0}} G_{\epsilon, \delta}^{(l)}$ (Lemma 5.2). Using the parameter $q$ introduced in Lemma 5.1, and observing that $q \leq \varrho_{l}$ for $1 \leq l \leq m_{0}$, we may thus further conclude from (5.23) that $\left\|u_{\epsilon, \delta}-u_{\epsilon} \mid A_{R, S_{0}} \times(0, t)\right\|_{q, 1 ; t} \rightarrow 0(\delta \downarrow 0)$ for any $R \in\left[S_{0}, \infty\right)$. For $l \in\left\{1, \ldots, m_{0}\right\}$, since $G_{\epsilon}^{(l)}(t)=0$ if $t \in(-\infty, \epsilon]$, we have $G_{\epsilon, \delta}^{(l)}(t)=0$ for $t \in(-\infty, \epsilon / 2], \delta \in(0, \epsilon / 2]$, hence $\left(u_{\epsilon, \delta}-u_{\epsilon}\right)(y, 0)=0$ for $y \in{\overline{B_{S_{0}}}}^{c}$ and thus $\mathfrak{I}^{(\tau)}\left(\left[u_{\epsilon, \delta}-u_{\epsilon} \mid{\overline{B_{S_{0}}}}^{c} \times[0, \infty)\right](0)\right)=0$, for $\delta \in(0, \epsilon / 2]$. Since $G^{(l)} \in C^{0}\left(\mathbb{R}, L^{\varrho}\left({\overline{B_{S_{0}}}}^{c}\right)^{3}\right)$, the last statement in Theorem 2.6 yields that $\left\|\left(G_{\epsilon, \delta}^{(l)}-G_{\epsilon}^{(l)}\right)(t)\right\|_{\varrho_{l}} \rightarrow 0(\delta \downarrow 0)$ for $1 \leq$ $l \leq m_{0}$. Again because $q \leq \varrho_{l}$, it follows that $\left\|\left(G_{\epsilon, \delta}^{(l)}-G_{\epsilon}^{(l)}\right)(t) \mid A_{R, S_{0}}\right\|_{q} \rightarrow 0(\delta \downarrow 0)$ for $1 \leq l \leq m_{0}, R \in\left[S_{0}, \infty\right)$, and thus $\left\|\left(u_{\epsilon, \delta}-u_{\epsilon} \mid A_{R, S_{0}} \times(0, \infty)\right)(t)\right\|_{q} \rightarrow 0(\delta \downarrow 0)$ for the same range of $R$.
Let $x \in B_{R_{0}}^{c}$. By Theorem 4.2, the function $y \mapsto \mathfrak{G}(x, y, 0)\left(y \in B_{R_{1}}\right)$ belongs to $L^{q^{\prime}}\left(B_{R_{1}}\right)^{3}$, so we may conclude that $\int_{A_{R_{1}, S_{0}}} \mathfrak{G}(x, y, 0) \cdot\left(u_{\epsilon, \delta}-u_{\epsilon}\right)(y, t) d y \rightarrow 0$ for $\delta \downarrow 0$. Since $\left\|u_{\epsilon, \delta}-u_{\epsilon} \mid \Omega_{R} \times(0, t)\right\|_{q, 1 ; t} \rightarrow 0$ and $\left\|\partial y_{k} u_{\epsilon, \delta}-\partial y_{k} u_{\epsilon} \mid Z_{t}\right\|_{q_{1}, 1 ; t} \rightarrow 0$ for $\delta \downarrow 0$ if $1 \leq k \leq 3, R \in\left[S_{0}, \infty\right)$, as shown above, inequality (4.8) in Theorem 4.2 yields that $\int_{0}^{t} \int_{A_{R_{1}, S_{0}}} \partial y_{k}^{\sigma} \partial_{s}^{1-\sigma} \mathfrak{G}(x, y, t-s) \cdot\left(\partial y_{k}^{\mu} u_{\epsilon, \delta}-\partial y_{k}^{\mu} u_{\epsilon}\right)(y, s) d y d s \rightarrow 0(\delta \downarrow 0)$ for $\sigma, \mu \in$ $\{0,1\}, k \in\{1,2,3\}$. Altogether we arrive at the relation $\mathfrak{K}\left(u_{\epsilon, \delta}-u_{\epsilon}\right)(x, t) \rightarrow 0(\delta \downarrow 0)$, with $\mathfrak{K}\left(u_{\epsilon, \delta}-u_{\epsilon}\right)=\mathfrak{K}_{R_{0}, S_{0}, \varphi_{0}, B_{S_{0}}, T_{0}}\left(u_{\epsilon, \delta}-u_{\epsilon}\right)$ defined in (4.9). Recall that $\left(u_{\epsilon, \delta}-u_{\epsilon}\right)(y, 0)=$ 0 for $\delta \in(0, \epsilon / 2], y \in{\overline{B_{S_{0}}}}^{c}$, so $\int_{A_{R_{1}, S_{0}}} \mathfrak{G}(x, y, t) \cdot\left(u_{\epsilon, \delta}-u_{\epsilon}\right)(y, 0) d y \rightarrow 0(\delta \downarrow 0)$. Since $\left\|f_{\epsilon, \delta}^{(j)}-f_{\epsilon}^{(j)}\right\|_{p_{j}, 1 ; t} \rightarrow 0(\delta \downarrow 0)$ (see above), we may apply (4.8) again, to obtain that $\int_{0}^{t} \int_{A_{R_{1}, S_{0}}} \mathfrak{G}(x, y, t-s) \cdot\left(f_{\epsilon, \delta}^{(j)}-f_{\epsilon}^{(j)}\right)(y, s) d y d s \rightarrow 0(\delta \downarrow 0)\left(j \in\left\{1, \ldots, n_{0}+m_{0}+1\right\}\right)$. Using the function $\varphi_{0} \in C_{0}^{\infty}\left(B_{R_{1}}\right)$ fixed at the beginning of this section, we set $\mathfrak{E}(y):=$ $\left(-\left(\partial_{j} \mathfrak{N}\right)(x-y) \partial_{k} \varphi_{0}(y)+\left(\partial_{k} \partial_{j} \mathfrak{N}\right)(x-y) \varphi_{0}(y)\right)_{1 \leq j, k \leq 3}$ for $y \in \overline{B_{R_{1}}}$. Since $x \in B_{R_{0}}^{c}$, this function $\mathfrak{E}$ is well defined and belongs to $C^{1}\left(\overline{B_{R_{1}}}\right)^{3 \times 3}$. Hence, by the Divergence theorem
and because $\operatorname{div}_{x}\left(u_{\epsilon, \delta}-u_{\epsilon}\right)=0$,

$$
\int_{\partial S_{0}}(\nabla \mathfrak{N})(x-y) \cdot\left(S_{0}^{-1} y \cdot\left(u_{\epsilon, \delta}-u_{\epsilon}\right)(y, t)\right) d o_{y}=\int_{A_{R_{1}, S_{0}}} \mathfrak{E}(y) \cdot\left(u_{\epsilon, \delta}-u_{\epsilon}\right)(y, t) d y .
$$

Again referring to the relation $\left\|\left(u_{\epsilon, \delta}-u_{\epsilon} \mid A_{R_{1}, S_{0}} \times(0, \infty)\right)(t)\right\|_{q} \rightarrow 0(\delta \downarrow 0)$, it follows that $\int_{\partial S_{0}}(\nabla \mathfrak{N})(x-y) \cdot\left(S_{0}^{-1} y \cdot\left(u_{\epsilon, \delta}-u_{\epsilon}\right)(y, t)\right) d o_{y} \rightarrow 0(\delta \downarrow 0)$. By Lemma 4.8, we find

$$
\begin{aligned}
& K_{\delta}:=\left|\sum_{l=1}^{3} \partial x_{l} \mathfrak{V}^{\left(\tau, B_{S_{0}}\right)}\left(n_{l}^{\left(B_{S_{0}}\right)}\left(u_{\epsilon, \delta}-u_{\epsilon}\right)\right)(x, t)\right| \\
& \leq \mathfrak{C}\left(\left\|u_{\epsilon, \delta}-u_{\epsilon}\left|A_{R_{1}, S_{0}} \times(0, t)\left\|_{q, 1 ; t}+\right\| \nabla_{y}\left(u_{\epsilon, \delta}-u_{\epsilon}\right)\right| A_{R_{1}, S_{0}} \times(0, t)\right\|_{q, 1 ; t}\right) .
\end{aligned}
$$

But $\left\|\nabla_{y}\left(u_{\epsilon, \delta}-u_{\epsilon}\right) \mid A_{R_{1}, S_{0}} \times(0, t)\right\|_{q, 1 ; t}$ is bounded by $\mathfrak{C}\left\|\nabla_{y}\left(u_{\epsilon, \delta}-u_{\epsilon}\right) \mid Z_{t}\right\|_{q_{1}, 1 ; t}$ for any $\delta \in(0, \infty)$, so $K_{\delta} \rightarrow 0(\delta \downarrow 0)$ by what we have proved above for the convergence of $u_{\epsilon, \delta}-u_{\epsilon}$ and $\nabla_{x}\left(u_{\epsilon, \delta}-u_{\epsilon}\right)$.
Up to this point, $x$ was arbitrary but fixed in $B_{R_{0}}^{c}$. Since $\left\|f_{\epsilon, \delta}^{(j)}-f_{\epsilon}^{(j)} \mid \bar{\Omega}^{c} \times(0, t)\right\|_{p_{j}, 1 ; t} \rightarrow 0$ for $\delta \downarrow 0$ by (5.23), Lemma 4.4 yields $\left\|\mathfrak{R}^{(\tau)}\left(f_{\epsilon, \delta}^{(j)}-f_{\epsilon}^{(j)}\right)(t)\right\|_{p_{j}} \rightarrow 0(\delta \downarrow 0)$, for $1 \leq$ $j \leq n_{0}+m_{0}+1$. Recalling that $\left\|\left(u_{\epsilon, \delta}-u_{\epsilon} \mid A_{R, S_{0}} \times(0, \infty)\right)(t)\right\|_{q} \rightarrow 0(\delta \downarrow 0)$ for $R \in$ $\left[S_{0}, \infty\right)$, we see there is a zero-measure set $M_{t, \epsilon} \subset \mathbb{R}^{3}$ and a sequence $\left(\delta_{n}\right)$ in $(0, \infty)$ with $\delta_{n} \rightarrow 0$ such that $\mathfrak{R}^{(\tau)}\left(f_{\epsilon, \delta_{n}}^{(j)}-f_{\epsilon}^{(j)}\right)(x, t) \rightarrow 0(n \rightarrow \infty)$ and $\left(u_{\epsilon, \delta}-u_{\epsilon}\right)(x, t) \rightarrow 0$ for
 $N_{t, n} \subset{\overline{B_{R_{0}}}}^{c}$ of measure zero such that equation (4.11) holds for $x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t, n}$, with $A=B_{S_{0}}$, with $u_{\epsilon, \delta_{n}}\left|{\overline{B_{S_{0}}}}^{c} \times(0, \infty), \pi_{\epsilon, \delta_{n}}\right|{\overline{B_{S_{0}}}}^{c} \times(0, \infty)$ in the role of $u$ and $\pi$, respectively, with $n_{0}$ replaced by $n_{0}+m_{0}+1$, and the functions $f^{(j)}\left(1 \leq j \leq n_{0}\right)$ by $f_{\epsilon, \delta_{n}}^{(j)} \mid{\overline{B_{S_{0}}}}^{c} \times$ $(0, \infty)\left(1 \leq j \leq n_{0}+m_{0}+1\right)$. (The function $\pi$ only appears in the assumptions of equation (4.11), not in the equation itself.) Letting $n$ tend to zero in that equation, we may conclude by the preceding convergence results that (4.11) is satisfied by $u_{\epsilon}, \pi_{\epsilon}, f_{\epsilon}^{(j)}$ as well, in the way stated in the lemma.

Finally we let $\epsilon$ tend to zero in (4.11).
Corollary 5.1 Define ${ }^{\left(S_{0}\right)}(y):=S_{0}^{-1} y$ for $y \in \partial B_{S_{0}}$. Let $t \in(0, \infty)$. Then there is a zero-measure set $N_{t} \subset{\overline{B_{R_{0}}}}^{c}$ such that the equation

$$
\begin{align*}
& u(x, t)=\mathfrak{R}^{(\tau)}(f)(x, t)+\mathfrak{I}^{(\tau)}\left(U_{0} \mid{\overline{B_{S_{0}}}}^{c}\right)(x, t)  \tag{5.24}\\
& \left.\quad-\sum_{l=1}^{3} \partial x_{l} \mathfrak{V}^{\left(\tau, B_{S_{0}}\right)}\left(n_{l}^{\left(S_{0}\right)} u\right)(x, t)-\int_{\partial B_{S_{0}}}(\nabla \mathfrak{N})(x-y)\left(n^{\left(S_{0}\right)}\right)(y) \cdot u(y, t)\right) d o_{y} \\
& \quad+\mathfrak{K}(u)(x, t)-\int_{A_{R_{1}, S_{0}}} \mathfrak{G}(x, y, t) \cdot U_{0}(y) d y-\int_{0}^{t} \int_{A_{R_{1}, S_{0}}} \mathfrak{G}(x, y, t-s) \cdot f(y, s) d y d s
\end{align*}
$$

holds for $x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t}$, with $f=\sum_{j=1}^{n_{0}} f^{(j)} \mid{\overline{S_{S_{0}}}}^{c} \times(0, \infty)$, where $\mathfrak{G}=\mathfrak{G}_{R_{0}, S_{0}, \varphi_{0}}$ is introduced in Theorem 4.2, $\mathfrak{K}(u)=\mathfrak{K}_{R_{0}, S_{0}, \varphi_{0}, B_{S_{0}}, T_{0}}(u)$ in (4.9), and $R_{1}$ at the beginning of this section.

Proof: By defintions in Lemma 5.1, the equation $\zeta_{\epsilon}^{\prime} u(s)\left|{\overline{B_{S_{0}}}}^{c}=\sum_{j=n_{0}+1}^{n_{0}+m_{0}+1} f_{\epsilon}^{(j)}(s)\right|{\overline{B_{S_{0}}}}^{c}$ holds for $s \in(0, \infty)$. (Actually $f_{\epsilon}^{\left(n_{0}+1\right)} \mid{\overline{B_{S_{0}}}}^{c} \times \mathbb{R}=0$, but this does not matter here.) Therefore Lemma 5.3 implies that for any $\epsilon \in(0, \infty)$, there is a set $N_{t, \epsilon} \subset{\overline{B_{R_{0}}}}^{c}$ of measure zero such that equation (4.11) holds for $x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t, \epsilon}$ with $A=B_{S_{0}}$ and with $u, f$ replaced by $u_{\epsilon} \mid \overline{B_{S_{0}}} \times(0, \infty)$ and $g_{\epsilon}$, respectively, where the function $g_{\epsilon}$ is defined by $g_{\epsilon}(s):=\sum_{j=1}^{n_{0}} f_{\epsilon}^{(j)}(s)+\zeta_{\epsilon}^{\prime}(s) u(s) \mid{\overline{B_{S_{0}}}}^{c}$ for $s \in(0, \infty)$. We recall that $G^{(l)} \in$ $C^{0}\left([0, \infty), L^{\varrho} l\left({\overline{B_{S_{0}}}}^{c}\right)^{3}\right)$ for $1 \leq l \leq m_{0}$ and $u(s) \mid{\overline{B_{S_{0}}}}^{c}=\sum_{l=1}^{m_{0}} G^{(l)}(s)$ for $s \in(0, \infty)$. Thus we see that all assumptions of Lemma 4.10 are valid with $A, u, f^{(j)}\left(1 \leq j \leq n_{0}\right)$ replaced by $B_{S_{0}}, u \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty)$ and $f^{(j)} \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty)\left(1 \leq j \leq n_{0}\right)$, respectively, with $m_{0}$ in the role of $k_{0}$ and $G^{(l)}\left(1 \leq l \leq m_{0}\right)$ in that of $u^{(l)}\left(1 \leq l \leq k_{0}\right)$. Lemma 4.10 implies that there is a measurable set $N_{t} \subset{\overline{B_{R_{0}}}}^{c}$ with properties as stated in the corollary, provided the equation

$$
\begin{equation*}
\mathfrak{I}^{(\tau)}\left(u(0)-U_{0} \mid{\overline{B_{S_{0}}}}^{c}\right)(x, t)+\int_{A_{R_{1}, S_{0}}} \mathfrak{G}(x, y, t) \cdot\left(U_{0}(y)-u(y, 0)\right) d y=0 \tag{5.25}
\end{equation*}
$$

holds. (Note that in view of our assumptions at the beginning of this section, we do not claim that $u(0)=U_{0}$.) In order to show (5.25), fix some $T \in(0, \infty)$. Let $\vartheta \in$ $C_{0, \sigma}^{\infty}\left({\overline{B_{S_{0}}}}^{c}\right)$. In a first step, we prove that $\int_{{\overline{S_{0}}}^{c}} U_{0} \cdot \vartheta d x=\int{\overline{B_{S_{0}}}}^{c} u(0) \cdot \vartheta d x$. Since $\nabla_{x} u$ belongs to $L_{l o c}^{1}\left([0, \infty), L^{q_{1}}\left(\bar{\Omega}^{c}\right)^{9}\right)$ and $f^{(j)}$ to $L_{l o c}^{1}\left([0, \infty), L^{p_{j}}\left(\bar{\Omega}^{c}\right)^{3}\right)$ for $1 \leq j \leq n_{0}$, the function $K(s):=\int{\overline{B_{S_{0}}}}^{c}\left(\nabla_{x} u(s) \cdot \nabla \vartheta+\tau \partial x_{1} u(s) \cdot \vartheta-f(s) \cdot \vartheta\right) d x \quad(s \in(0, \infty))$ belongs to $L_{l o c}^{1}([0, \infty))$, in particular $K \mid(0, T) \in L^{1}((0, T))$. On the other hand, since $u \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty)=\sum_{l=1}^{m_{0}} G^{(l)}$ and $G^{(l)} \in C^{0}\left([0, \infty), L^{\varrho_{l}}\left({\overline{B_{S_{0}}}}^{c}\right)^{3}\right)\left(1 \leq l \leq m_{0}\right)$, the function $H(s):=\int{\overline{B_{S_{0}}}} c u(s) \cdot \vartheta d x$ with $s \in[0, \infty)$ belongs to $C^{0}([0, \infty))^{3}$. Moreover, due to equation (5.1), we get

$$
\begin{equation*}
\int_{0}^{\infty}\left(-\gamma^{\prime}(s) H(s)+\gamma(s) K(s)\right) d s=\gamma(0) \int_{{\overline{S_{S_{0}}}}^{c}} U_{0} \cdot \vartheta d x \tag{5.26}
\end{equation*}
$$

for $\gamma \in C_{0}^{\infty}([0, \infty))$. This equation for $\gamma \in C_{0}^{\infty}((0, T))$ and the relations $H \mid[0, T] \in$ $C^{0}([0, T])$ and $K \mid(0, T) \in L^{1}((0, T))$ yield that $H \mid(0, T) \in W^{1,1}((0, T))$ and $H^{\prime}(s)=$ $-K(s)$ for a. e. $s \in(0, T)$. Let $\gamma \in C_{0}^{\infty}([0, T))$. It follows that $\gamma H \in C^{0}([0, T]) \cap$ $W^{1,1}((0, T))$ and $(\gamma H)^{\prime}=-\gamma K+\gamma^{\prime} H$, hence $\int_{0}^{T}\left(\gamma K-\gamma^{\prime} H\right) d s=\gamma(0) H(0)$. Comparing this equation with (5.26) and recalling that $H(0)=\int{\overline{B_{S_{0}}}}^{c} u(0) \cdot \vartheta d x$, we arrive at the equation $\int{\overline{B_{S_{0}}}} c U_{0} \cdot \vartheta d x=\int{\overline{B_{S_{0}}}} c u(0) \cdot \vartheta d x$ we wanted to show. Now let $t \in(0, \infty)$ and $x \in B_{R_{0}}^{c}$, and put $Y(y):=\Lambda(x-y, t)-\mathfrak{G}(x, y, t)$ for $y \in B_{R_{1}}$. Recall that $\mathfrak{G}(x, \cdot, t) \in$ $C_{0}^{\infty}\left(B_{R_{1}}\right)^{3 \times 3}, \mathfrak{G}(x, y, t)=\Lambda(x-y, t)$ for $y \in B_{S_{0}+\left(R_{0}-S_{0}\right) / 4}$ (Theorem 4.2), and $\Lambda \mid \mathbb{R}^{3} \times$ $(0, \infty) \in C^{\infty}\left(\mathbb{R}^{3} \times(0, \infty)\right)\left(\right.$ Lemma 4.1). Thus we get $\operatorname{supp}(Y) \subset B_{S_{0}+\left(R_{0}-S_{0}\right) / 4}^{c}$ and $Y \in$ $C^{\infty}\left(\mathbb{R}^{3}\right)^{3 \times 3}$. Moreover $\sum_{k=1}^{3} \partial y_{k} Y_{j k}(y)=0$ for $y \in \mathbb{R}^{3}, 1 \leq j \leq 3$ and $Y \in W^{1, r}\left(\mathbb{R}^{3}\right)^{3 \times 3}$ for $r \in(1, \infty)$ also by Lemma 4.1. Therefore, for $1 \leq \mu \leq 3$, Theorem 2.2 provides a sequence $\left(\vartheta_{n}^{(\mu)}\right)_{n \geq 1}$ in $C_{0, \sigma}^{\infty}\left({\overline{B_{S_{0}}}}^{c}\right)$ such that in particular $\left\|\left(Y_{\mu k}\right)_{1 \leq k \leq 3}-\vartheta_{n}^{(\mu)}\right\|_{r} \rightarrow 0(n \rightarrow \infty)$ for $r=\widetilde{p}$ and $r \in\left\{\varrho_{m}: 1 \leq m \leq m_{0}\right\}$, with $\widetilde{p}$ introduced at the beginning of this section in the assumption $U_{0} \in L^{\widetilde{p}}\left(\bar{\Omega}^{c}\right)^{3}$. But $\int{\overline{B_{S_{0}}}}^{c} \vartheta_{n}^{(\mu)} \cdot U_{0} d x=\int{\overline{B_{S_{0}}}}^{c} \vartheta_{n}^{(\mu)} \cdot u(0) d x(1 \leq \mu \leq$ $3, n \in \mathbb{N}$ ) by what has already been proved, and $u(0) \mid{\overline{B_{S_{0}}}}^{c}=\sum_{l=1}^{m_{0}} G^{(l)}(0), G^{(l)}(0) \in$
$L^{\varrho_{l}}\left({\overline{B_{S_{0}}}}^{c}\right)^{3}\left(1 \leq l \leq m_{0}\right)$ by our assumptions. Hence, by letting $n$ tend to infinity, we obtain $\int_{\overline{B_{S_{0}}}} c Y \cdot U_{0} d y=\int_{\overline{S_{S_{0}}}} c Y \cdot u(0) d y$. Due to the definition of $Y$, this implies (5.25). So the proof of Corollary 5.1 is completed.

Up to now, we considered the case $T_{0}=\infty$ in our assumptions at the beginning of this section. But in the present context, the transition from this case to the case $T_{0}<\infty$ is easy to perform:
Corollary 5.2 Suppose that $T_{0}<\infty$. Then, for $t \in\left(0, T_{0}\right)$, there is a zero-measure set $N_{t} \subset{\overline{B_{R_{0}}}}^{c}$ such that equation (5.24) holds for $x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t}$ with $f=\sum_{j=1}^{n_{0}} f^{(j)} \mid{\overline{B_{S_{0}}}}^{c} \times$ $\left(0, T_{0}\right)$.
Proof: Let $T^{\prime} \in\left(0, T_{0}\right)$, and choose a function $\zeta \in C^{\infty}(\mathbb{R})$ such that $\zeta$ equals 1 on the interval $\left(-\infty, T^{\prime}+\left(T_{0}-T^{\prime}\right) / 4\right)$ and vanishes on $\left(T^{\prime}+\left(T_{0}-T^{\prime}\right) / 2, \infty\right)$. Define $\widetilde{u}(s):=\zeta(s) u(s), \widetilde{f}^{(j)}(s):=\zeta(s) f^{(j)}(s), \widetilde{G}^{(l)}(s):=\zeta(s) G^{(l)}(s)$ for $1 \leq j \leq n_{0}, 1 \leq$ $l \leq k_{0}, s \in\left(0, T_{0}\right)$. On $\left[T_{0}, \infty\right)$, the value zero is assigned to these functions. Then all assumptions listed at the beginning of this section except equation (5.1) are valid with $\infty$ in the role of $T_{0}$, and $\widetilde{u}, \widetilde{f}^{(j)}, \widetilde{G}^{(l)}$ in that of $u, f^{(j)}$ and $G^{(l)}$, respectively $\left(1 \leq j \leq n_{0}, 1 \leq\right.$ $\left.l \leq m_{0}\right)$. Define $\widetilde{f}^{\left.\left(n_{0}+1\right)\right)}(s):=\zeta^{\prime}(s) \widetilde{u}(s)$ for $s \in(0, \infty)$. Let $\gamma \in C_{0}^{\infty}([0, \infty))$. Then $\gamma \zeta \in C_{0}^{\infty}\left(\left[0, T_{0}\right)\right)$, so we may replace $\gamma$ by $\zeta \gamma$ in (5.1). It follows that (5.1) is valid with $T_{0}$ replaced by $\infty, u$ by $\widetilde{u}, n_{0}$ by $n_{0}+1$, and $f^{(j)}\left(1 \leq j \leq n_{0}\right)$ by $\widetilde{f}^{(j)}\left(1 \leq j \leq n_{0}+1\right)$. Let $t \in(0, \infty)$. Then Corollary 5.1 implies there is a set $N_{t} \subset{\overline{B_{R_{0}}}}^{c}$ of measure zero such that (5.24) holds for $x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t}$ with $u$ replaced by $\widetilde{u}$, and with $f=\sum_{j=1}^{n_{0}+1} \widetilde{f}(j) \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty)$. Now suppose that $t \in\left(0, T^{\prime}\right)$. Then, for $s \in(0, t]$, the functions $\widetilde{u}(s)$ and $u(s)$ coincide, as do $\widetilde{f}^{(j)}(s)$ and $f^{(j)}(s)$ for $1 \leq j \leq n_{0}$, whereas $\widetilde{f}^{\left(n_{0}+1\right)}(s)=0$. By the definitions in Lemma 4.4, we get that $\mathfrak{R}^{(\tau)}\left(\sum_{j=1}^{n_{0}+1} \widetilde{f}^{(j)} \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty)\right)(x, t)=\mathfrak{R}\left(\sum_{j=1}^{n_{0}} f^{(j)} \mid{\overline{B_{S_{0}}}}^{c} \times\left(0, T_{0}\right)\right)(x, t)$ for $x \in \mathbb{R}^{3}$. Therefore we may conclude that (5.24) is valid for $x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t}$ with $f=$ $\sum_{j=1}^{n_{0}} f^{(j)} \mid{\overline{B_{S_{0}}}}^{c} \times\left(0, T_{0}\right)$. Since $T^{\prime}$ was taken arbitrarily in $\left(0, T_{0}\right)$, the corollary is proved.

In order to exploit our representation formula (5.24) for the purpose of decay estimates, we need somewhat stronger assumptions on $u$. The ensuing theorem gives the details.
Theorem 5.2 Abbreviate $\mathfrak{A}:=A_{R_{1}, S_{0}} \times\left(0, T_{0}\right), f:=\sum_{j=1}^{n_{0}} f^{(j)} \mid{\overline{B_{S_{0}}}}^{c} \times\left(0, T_{0}\right)$. In addition to the assumptions at the beginning of this section, suppose that the function $u \mid \mathfrak{A}$ belongs to $L^{\infty}\left(0, T_{0}, L^{q}\left(A_{R_{1}, S_{0}}\right)^{3}\right)$ and to $L^{\gamma_{1}}\left(0, T_{0}, L^{q}\left(A_{R_{0}, S_{0}}\right)^{3}\right)$, the function $\nabla_{x} u \mid \mathfrak{A}$ to $L^{\gamma_{2}}\left(0, T_{0}, L^{q}\left(A_{R_{1}, S_{0}}\right)^{9}\right)$, and $f \mid \mathfrak{A}$ is in the space $L^{\gamma_{3}}\left(0, T_{0}, L^{q}\left(A_{R_{1}, S_{0}}\right)^{3}\right)$, for certain parameters $\gamma_{1}, \gamma_{2}, \gamma_{3} \in[1, \infty]$, with $q$ chosen in Lemma 5.1. Then there is a zero-measure set $\mathfrak{S}_{T_{0}} \subset\left(0, T_{0}\right)$ such that

$$
\begin{align*}
& \left|\partial_{x}^{\alpha} u(x, t)-\partial_{x}^{\alpha} \mathfrak{R}^{(\tau)}(f)(x, t)-\partial_{x}^{\alpha} \mathfrak{I}^{(\tau)}\left(U_{0} \mid B_{S_{0}}^{c}\right)(x, t)\right|  \tag{5.27}\\
& \leq \mathfrak{C}\left(\left\|u\left|\mathfrak{A}\left\|_{q, \infty ; T_{0}}+\right\| u\right| \mathfrak{A}\right\|_{q, \gamma_{1} ; T_{0}}+\left\|\nabla_{x} u\left|\mathfrak{A}\left\|_{q, \gamma_{2} ; T_{0}}+\right\| f\right| \mathfrak{A}\right\|_{q, \gamma_{3} ; T_{0}}+\left\|U_{0}\right\|_{\tilde{p}}\right) \\
& \left.\quad\left((|x| \nu(x))^{-(3+|\alpha|) / 2+1 /\left(2 \min \left\{\gamma_{1}^{\prime}, \gamma_{2}^{\prime}, \gamma_{3}^{\prime}\right\}\right)}+|x|^{-\gamma-|\alpha|}\right)\right)
\end{align*}
$$

for $t \in\left(0, T_{0}\right) \backslash \mathfrak{S}_{T_{0}}, x \in{\overline{B_{R_{0}}}}^{c} \backslash N_{t}$ with some set $N_{t} \subset{\overline{B_{R_{0}}}}^{c}$ of measure zero, and for $\alpha \in \mathbb{N}_{0}^{3}$ with $|\alpha| \leq 1$, where $\gamma=3$ if $\int_{\partial S_{0}} u(t) \cdot n^{(\Omega)} d o_{x}=0$ for $t \in\left(0, T_{0}\right)$, and $\gamma=2$ else. If $\gamma=3$, the last line in (5.27) may be replaced by $(|x| \nu(x))^{-(3+|\alpha|) / 2+1 /\left(2 \min \left\{\gamma_{1}^{\prime}, \gamma_{2}^{\prime}, \gamma_{3}^{\prime}\right\}\right)}$.

Proof: Since $q \leq \min \left\{\varrho_{l}: 1 \leq l \leq m_{0}\right\}, u \mid{\overline{B_{S_{0}}}}^{c} \times\left(0, T_{0}\right)=\sum_{l=1}^{m_{0}} G^{(l)}$ and $G^{(l)} \in$ $C^{0}\left(\left[0, T_{0}\right), L^{\varrho_{l}}\left({\overline{B_{S_{0}}}}^{c}\right)^{3}\right)$ for $1 \leq l \leq m_{0}$, we have $u \mid \mathfrak{A} \in C^{0}\left(\left[0, T_{0}\right), L^{q}\left(A_{R_{1}, S_{0}}\right)^{3}\right)$. As a consequence, we get $\left\|u(t)\left|A_{R_{1}, S_{0}}\left\|_{q} \leq\right\| u\right| \mathfrak{A}\right\|_{q, \infty ; T_{0}}$ for $t \in\left[0, T_{0}\right)$. Thus inequality (5.27) follows from Corollary 5.1 and 5.2 , (4.7), (4.8), Corollary 4.2, Lemma 4.8 and 4.9. Obviously $|x| \geq C(R) \nu(x)$ for $x \in B_{R}^{c}, R \in(0, \infty)$, so in the case $\gamma=3$, inequality (5.27) holds without the term $|x|^{-\gamma-|\alpha|}$ in the last line.

## 6 Some comments and applications

In view of inequality (5.27), we remark that estimates of the terms $\left|\partial_{x}^{\alpha} \mathfrak{\Re}^{(\tau)}(f)(x, t)\right|$ and $\left|\partial_{x}^{\alpha} \mathfrak{I}^{(\tau)}\left(U_{0} \mid{\overline{B_{S_{0}}}}^{c}\right)(x, t)\right|$ may be found elsewhere. They are independent of the theory presented here, only depending on the decay properties of $U_{0}$ and $f$. An optimal decay bound, that is, a bound also valid for $\left|\partial_{x}^{\alpha} \Lambda(x, t)\right|$ if $|x|$ is large (see Lemma 4.1), is given by $\mathfrak{C}(|x| \nu(x))^{-(3+|\alpha|) / 2}$. This bound is obtained if, for example, $f$ and $U_{0}$ have compact support ([19, Lemma 4.1, 4.2]). For conditions on $f$ and $U_{0}$ leading to the decay bound $\mathfrak{C}(|x| \nu(x))^{-(2+|\alpha|) / 2}$, we refer to [14, Theorem 3.1] and [13, Theorem 1.1], respectively. These indications explain why inequality (1.3) follows from (5.27) under suitable assumptions on $U_{0}$ and $f$.
It should further be mentioned that the $\operatorname{sum} \partial_{x}^{\alpha} \mathfrak{R}^{(\tau)}(f)+\mathfrak{I}^{(\tau)}\left(U_{0}\right)$ constitutes a solution to (1.1) in the whole space $\mathbb{R}^{3}(\Omega=\emptyset)$, with initial data $U_{0}$ if $U_{0}$ is solenoidal ([19, (3.3), Lemma 3.7, Theorem 3.3]). So left-hand side of inequality (6.11) may be interpreted as the perturbation generated in the fluid by the presence of the rigid body. But it is precisely this perturbation which is of interest here. Thus inequality (6.11) may be considered as a decay estimate of that aspect of the flow which is relevant in the present context.
Let us compare the theory developed here with that in [14] and its predecessor paper [11]. The former paper extends the theory derived in the latter one, so we only consider the former one. Our general approach in the work at hand consists in taking existence of a solution to (1.1) for granted. We require conditions on the velocity $u$ and on the right-hand side $f$ in (1.1) as well as on the initial data $U_{0}$ in order to establish an integral representation of $u$ (equation (5.24)), which, in turn, serves to derive a decay estimate of $u$ (inequality (5.27)), under some additional assumptions on $u$ and $f$. In contrast to that, the approach in [14] leads simultaneously to an existence result for the velocity part $u$ of a solution to (1.1) and to an integral representation of $u$, and in a supplementary step to the estimate

$$
\begin{equation*}
\left|\partial_{x}^{\alpha} u(x, t)-\partial_{x}^{\alpha} \mathfrak{R}^{(\tau)}(f)(x, t)-\partial_{x}^{\alpha} \mathfrak{I}^{(\tau)}\left(U_{0}\right)(x, t)\right| \leq \mathfrak{C}(|x| \nu(x))^{-1-|\alpha| / 2} \tag{6.1}
\end{equation*}
$$

for $x, t, \alpha$ as in (5.27) ([14, Corollary 2.28, Lemma 3.2, (3.7)]). No assumptions on the velocity $u$ are needed because its existence is a result, not an assumption. The method of proof in [14] consists in solving an integral equation on $\partial \Omega \times\left(0, T_{0}\right)$, as proposed by Shen [39] for the time-dependent Stokes system in the case $f=0, U_{0}=0$. The regularity of the solutions constructed in [39] corresponds to the regularity of the boundary data $b: \partial \Omega \times\left(0, T_{0}\right) \mapsto \mathbb{R}^{3}$. Thus the assumptions on $b$ in $[39]$ may be considered as natural. However, when in [14] this method is extended to the Oseen system with nonvanishing $f$
or $U_{0}$, the assumptions on these latter functions are more restrictive than the ones in the work at hand. And if our theory here is applied to the solutions constructed in [14], under the same assumptions on the data as in [14], we obtain decay rates which are stronger than those in this reference. In fact, we are able to replace the factor $(|x| \nu(x))^{-1-|\alpha| / 2}$ in (6.1) by $(|x| \nu(x))^{-5 / 4-|\alpha| / 2}$. The details are given in the ensuing Theorem $6.1\left(f=0, U_{0}=0\right)$ and $6.2(b=0)$.
Theorem 6.1 For $\varphi \in C_{0}^{\infty}\left(\mathbb{R}^{4}\right)^{3}$ with $\varphi \mid \mathbb{R}^{3} \times(-\infty, 0]=0$, define the fractional derivative $\partial_{t}^{1 / 2} \varphi$ by $\left(\partial_{t}^{1 / 2} \varphi\right)(x, t):=\pi^{-1 / 2} \partial_{t}\left(\int_{0}^{t}(t-s)^{-1 / 2} \varphi(x, s) d s\right)$ for $t \in(0, \infty), x \in \mathbb{R}^{3}$.
Put $S_{T_{0}}:=\partial \Omega \times\left(0, T_{0}\right)$. Let $b \in L^{2}\left(S_{T_{0}}\right)^{3}$ with $\int_{\partial \Omega} b(t) \cdot n^{(\Omega)} d o_{x}=0$ for $t \in\left(0, T_{0}\right)$. Suppose there is a sequence $\left(\varphi_{n}\right)$ in $C_{0}^{\infty}\left(\mathbb{R}^{4}\right)^{3}$ such that $\varphi_{n} \mid \mathbb{R}^{3} \times(-\infty, 0]=0$ for $n \in \mathbb{N},\left\|b-\varphi_{n}\right\|_{2} \rightarrow$ $0(n \rightarrow \infty)$ and $\int_{0}^{T_{0}}\left\|\left(\varphi_{n}-\varphi_{m}\right)(t)\left|\partial \Omega\left\|_{H^{1}(\partial \Omega)^{3}}^{2} d t \rightarrow 0, \int_{0}^{T_{0}}\right\| \partial_{t}^{1 / 2}\left(\varphi_{n}-\varphi_{m}\right)(t)\right| \partial \Omega\right\|_{2}^{2} d t \rightarrow$ $0, \int_{0}^{T_{0}}\left\|\partial_{t}\left(\varphi_{n}-\varphi_{m}\right)(t) \cdot n^{(\Omega)}\right\|_{H^{1}(\partial \Omega)^{\prime}}^{2} d t \rightarrow 0$ for $m, n \rightarrow \infty$. Here the space $H^{1}(\partial \Omega)$ is to be defined in the usual way, and the symbol $\left\|\|_{H^{1}(\partial \Omega)^{3}}\right.$ denotes the usual norm of $H^{1}(\partial \Omega)^{3}$ with respect to some local coordinates (see [28, Section III.6.7] for example). The symbol $\left\|\|_{H^{1}(\partial \Omega)^{\prime}}\right.$ stands for the canonical norm of the dual space of $H^{1}(\partial \Omega)$.

Then there is a unique function $u \in L_{\text {loc }}^{2}\left(\left[0, T_{0}\right), H^{1}\left(\bar{\Omega}^{c}\right)^{3}\right)$ such that div ${ }_{x} u(t)=0$ and $u(t) \mid \partial \Omega=b(t)$ for $t \in\left(0, T_{0}\right)$, and such that equation (5.1) is satisfied with $f=0$ and $U_{0}=0$. Moreover $u \in L^{\infty}\left(0, T_{0}, L^{2}\left(\bar{\Omega}^{c}\right)^{3}\right)$ and $\nabla_{x} u \in L^{2}\left(0, T_{0}, L^{2}\left(\bar{\Omega}^{c}\right)^{9}\right)$. In addition $\left|\partial_{x}^{\alpha} u(x, t)\right| \leq \mathfrak{C}(|x| \nu(x))^{-5 / 4-|\alpha| / 2}$ for $t, x, \alpha$ as in (5.27).
Proof: The uniqueness statement follows as in the Stokes case; see [12, Theorem 3.7] and its proof. For the existence result we refer to [14, Theorem 2.26], which yields a function $u \in L_{l o c}^{2}\left(\left[0, T_{0}\right), H^{1}\left(\bar{\Omega}^{c}\right)^{3}\right)$ such that $\operatorname{div} u(t)=0, u(t) \mid \partial \Omega=b(t)$ for $t \in\left(0, T_{0}\right)$ and equation (5.1) is valid with $f=0, U_{0}=0$. Due to [14, (2.13)] and [9, Theorem 2.3], this function $u$ additionally belongs to $L^{\infty}\left(0, T_{0}, L^{2}\left(\bar{\Omega}^{c}\right)^{3}\right)$, and $\nabla_{x} u \in L^{2}\left(0, T_{0}, L^{2}\left(\bar{\Omega}^{c}\right)^{9}\right)$. This leaves us to consider the estimate stated at the end of this theorem. Theorem 2.3 implies that $u \in L^{2}\left(0, T_{0}, L^{6}(\bar{\Omega})^{3}\right)$. In addition, by $[14,(2.13)]$ and Lemma 4.7 we have $u \mid{\overline{B_{S_{0}}}}^{c} \times\left(0, T_{0}\right) \in C^{0}\left(\left[0, T_{0}\right), L^{4}\left({\overline{B_{S_{0}}}}^{c}\right)^{3}\right)$. Therefore the estimate in question follows from Theorem 5.2 with $U_{0}=0, n_{0}=1, p_{1}=2, f^{(1)}=0, m_{0}=1, \varrho_{1}=4, G^{(1)}=$ $u \mid{\overline{B_{S_{0}}}}^{c} \times\left(0, T_{0}\right), q_{0}=q_{1}=\widetilde{p}=\gamma_{j}=2$ for $j \in\{1,2,3\}$.
Theorem 6.2 Let $\kappa \in(0,1 / 2]$ and define $H_{\sigma}^{1 / 2+\kappa}\left(\bar{\Omega}^{c}\right)$ as the closure of $C_{0, \sigma}^{\infty}\left(\bar{\Omega}^{c}\right)$ with respect to the norm of the fractional-order Sobolev space $W^{1 / 2+\kappa, 2}\left(\bar{\Omega}^{c}\right)^{3}$. Suppose that $U_{0} \in H_{\sigma}^{1 / 2+\kappa}\left(\bar{\Omega}^{c}\right)$. Let $\widetilde{q} \in[1,3 / 2)$ and let $f$ belong to the spaces $L^{2}\left(0, \infty, L^{3 / 2}\left(\bar{\Omega}^{c}\right)^{3}\right)$ and $L^{2}\left(0, \infty, L^{\widetilde{q}}\left(\bar{\Omega}^{c}\right)^{3}\right)$. Then there is a unique function $u \in L_{l o c}^{2}\left([0, \infty), W^{1,2}\left(\bar{\Omega}^{c}\right)^{3}\right)$ such that $u(t) \mid \partial \Omega=0$ and $\operatorname{div} u(t)=0$ for $t \in(0, \infty)$, and such that $u$ satisfies (5.1). In addition

$$
\begin{align*}
& \left|\partial_{x}^{\alpha} u(x, t)-\partial_{x}^{\alpha} \mathfrak{R}^{(\tau)}\left(f \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty)\right)(x, t)-\partial_{x}^{\alpha} \mathfrak{I}^{(\tau)}\left(U_{0} \mid{\overline{B_{S_{0}}}}^{c}\right)(x, t)\right|  \tag{6.2}\\
& \leq \mathfrak{C}(|x| \nu(x))^{-5 / 4-|\alpha| / 2} \quad \text { for } x, t, \alpha \text { as in (5.27). }
\end{align*}
$$

Proof: Concerning uniqueness, as in the proof of Theorem 6.1, we refer to [12, Theorem 3.7]. The assumptions on $U_{0}$ and $f$ were chosen in such a way that the existence result in Theorem 6.2 holds according to $[14$, Theorem 2.26$]$ with $T_{0}=\infty, d=0, c=U_{0}, h=f$.

The solution provided by this reference is given by

$$
\begin{equation*}
u=\left(\mathfrak{R}^{(\tau)}(f)+\mathcal{H}^{(\tau)}\left(U_{0}\right)+\mathfrak{V}^{(\tau)}(\varphi)\right) \mid \bar{\Omega}^{c} \times(0, \infty) \tag{6.3}
\end{equation*}
$$

(see [14, (2.13)]), for a certain function $\varphi \in L^{2}(\partial \Omega \times(0, \infty))^{3}$ specified in [14, Theorem 2.26], and with $\mathcal{H}^{(\tau)}\left(U_{0}\right)$ defined in Lemma 4.3. We use this representation formula in order to check whether $u$ satisfies the assumptions imposed at the beginning of Section 5 and in Theorem 5.2. To this end, we observe that Lemma 4.3 yields $\mathcal{H}^{(\tau)}\left(U_{0}\right) \in$ $C^{0}\left([0, \infty), L^{2}\left(\mathbb{R}^{3}\right)^{3}\right) \cap L^{\infty}\left(0, \infty, L^{2}\left(\mathbb{R}^{3}\right)^{3}\right)$ and $\nabla_{x} \mathcal{H}^{(\tau)}\left(U_{0}\right) \in L^{2}\left(0, \infty, L^{2}\left(\mathbb{R}^{3}\right)^{9}\right)$. Moreover, by Lemma 4.5 , we have $\mathfrak{R}^{(\tau)}(f) \in C^{0}\left([0, \infty), L^{3 / 2}\left(\mathbb{R}^{3}\right)^{3}\right), \mathfrak{R}^{(\tau)}(f) \mid B_{R_{0}} \times(0, \infty) \in$ $L^{\infty}\left(0, \infty, L^{3 / 2}\left(B_{R_{0}}\right)^{3}\right)$ and $\nabla_{x} \mathfrak{R}^{(\tau)}(f) \mid B_{R_{0}} \times(0, \infty) \in L^{2}\left(0, \infty, L^{3 / 2}\left(B_{R_{0}}\right)^{9}\right)$. According to [10, Corollary 2.17] with $s=2, q=3 / 2$, the function $\nabla_{x} \mathfrak{R}^{(\tau)}(f)$ belongs to $L_{\text {loc }}^{2}\left([0, \infty), L^{2}\left(\mathbb{R}^{3}\right)^{9}\right)$. Moreover $\left[9\right.$, Theorem 2.3] yields that $\mathfrak{V}^{(\tau)}(\varphi) \mid \bar{\Omega}^{c} \times(0, \infty)$ belongs to $\in L^{\infty}\left(0, \infty, L^{2}\left(\bar{\Omega}^{c}\right)^{3}\right)$, and $\nabla_{x} \mathfrak{V}^{(\tau)}(\varphi) \mid \bar{\Omega}^{c} \times(0, \infty)$ to $L^{2}\left(0, \infty, L^{2}\left(\bar{\Omega}^{c}\right)^{9}\right)$. By Lemma 4.7 we have $\mathfrak{V}^{(\tau)}(\varphi) \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty) \in C^{0}\left([0, \infty), L^{4}\left({\overline{B_{S_{0}}}}^{c}{ }^{3}\right)\right.$. At this point Theorem 2.3 implies that $\mathcal{H}^{(\tau)}\left(U_{0}\right) \mid \bar{\Omega}^{c} \times(0, \infty)$ and $\mathfrak{V}^{(\tau)}(\varphi) \mid \bar{\Omega}^{c} \times(0, \infty)$ belong to $L^{2}\left(0, \infty, L^{6}\left(\bar{\Omega}^{c}\right)^{3}\right)$. Since $\nabla_{x} \mathfrak{R}^{(\tau)}(f) \mid B_{R_{0}} \times(0, \infty) \in L^{2}\left(0, \infty, L^{3 / 2}\left(B_{R_{0}}\right)^{9}\right)$, as mentioned above, and because of (6.3), we thus get that $\nabla_{x} u \mid \Omega_{R_{0}} \times(0, \infty) \in L^{2}\left(0, \infty, L^{3 / 2}\left(\Omega_{R_{0}}\right)^{9}\right)$. But $u(t) \mid \partial \Omega=$ $0(t \in(0, \infty))$, hence by Poincaré's lemma, $u \mid \Omega_{R_{0}} \times(0, \infty) \in L^{2}\left(0, \infty, L^{3 / 2}\left(B_{R_{0}}\right)^{3}\right)$. Again recalling (6.3), we thus see that $u$ satisfies the assumptions imposed at the beginning of Section 5 and in Theorem 5.2 with $T_{0}=\infty, \widetilde{p}=q_{1}=2, q_{0}=3 / 2, n_{0}=1, p_{1}=$ $3 / 2, \gamma_{3}=2, f^{(1)}=f, m_{0}=3, \varrho_{1}=2, G^{(1)}=\mathcal{H}^{(\tau)}\left(U_{0}\right) \mid{\overline{B_{S_{0}}}}^{c} \times(0, \infty), \varrho_{2}=3 / 2, G^{(2)}=$ $\mathfrak{R}^{(\tau)}(f)\left|\overline{B_{S_{0}}}{ }^{c} \times(0, \infty), \varrho_{3}=4, G^{(3)}=\mathfrak{V}{ }^{(\tau)}(\varphi)\right|{\overline{B_{S_{0}}}}^{c} \times(0, \infty), \gamma_{1}=\gamma_{2}=2, q=5 / 4$. Now Theorem 5.2 yields (6.2).

As another application of Theorem 5.2, we can handle standard $L^{2}$-weak solutions to (1.1) with side conditions $u(t) \mid \partial \Omega=0$ for $t \in\left(0, T_{0}\right)$ and $u(0)=U_{0}$, under assumptions that are only slightly more restrictive than the conditions needed for existence $\left(f \in L^{2}\left(0, T_{0}, L_{\sigma}^{2}\left(\bar{\Omega}^{c}\right)\right)\right.$ instead of $f \in L^{2}\left(0, T_{0}, \mathcal{V}^{\prime}\right)$, with $\mathcal{V}^{\prime}$ defined below). The theory in [13] does not cover this situation at all. The next theorem gives the details on existence and on our decay result for this type of solution.
Theorem 6.3 Put $\mathcal{V}:=\left\{V \in W_{0}^{1,2}\left(\bar{\Omega}^{c}\right)^{3}: \operatorname{div} V=0\right\}$, equip $\mathcal{V}$ with the norm of $W^{1,2}\left(\bar{\Omega}^{c}\right)^{3}$, and let $\mathcal{V}^{\prime}$ denote the dual space to $\mathcal{V}$. Write $L_{\sigma}^{2}\left(\bar{\Omega}^{c}\right)$ for the closure of $C_{0, \sigma}^{\infty}\left(\bar{\Omega}^{c}\right)$ with respect to the norm of $L^{2}\left(\bar{\Omega}^{c}\right)^{3}$. Let $U_{0} \in L_{\sigma}^{2}\left(\bar{\Omega}^{c}\right)$ and $f \in L^{2}\left(0, T_{0}, \mathcal{V}^{\prime}\right)$. Then there is a unique function $u:\left(0, T_{0}\right) \mapsto \mathcal{V}$ such that $u \in L^{\infty}\left(0, T_{0}, L^{2}\left(\bar{\Omega}^{c}\right)^{3}\right), \nabla v \in$ $L^{2}\left(0, T_{0}, L^{2}\left(\bar{\Omega}^{c}\right)^{9}\right)$, and equation (5.1) is fulfilled with the modification that the term $f(t) \cdot \vartheta$ is dropped, and instead the term $f(t)(\vartheta)$ is added outside the integral over $\bar{\Omega}^{c}$. This function $u$ is in $C^{0}\left(\left[0, T_{0}\right), L^{2}\left(\bar{\Omega}^{c}\right)^{3}\right)$.
If $f \in L^{2}\left(0, T_{0}, L^{2}\left(\bar{\Omega}^{c}\right)^{3}\right)$, then inequality (6.2) holds for $t, x, \alpha$ as in (5.27).
Proof: As in the two preceding theorems, we refer to [14, Theorem 3.7] and its proof as regards uniqueness, which follows as in the Stokes case. Concerning existence, the argument is also the same as in the Stokes case. We refer to [44, p. 171-176 and p. 180], in particular [44, p. 175, (1.65)]. The equation $\int_{\bar{\Omega}^{c}} \partial_{1} V \cdot V d x=0$ valid for $V \in W_{0}^{1,2}\left(\bar{\Omega}^{c}\right)^{3}$ is the reason why the Oseen term does not generate a major problem. In order to show that
$u \in C^{0}\left(\left[0, T_{0}\right), L^{2}\left(\bar{\Omega}^{c}\right)^{3}\right)$, we consider $v$ as a weak solution of the time-dependent Stokes system, in an analogous way as stated in (5.1) for the Oseen system, but now with the right-hand side $f-\tau \partial x_{1} u \in L^{2}\left(0, T_{0}, \mathcal{V}^{\prime}\right)$. Then [44, Theorem 3.1.1] yields continuity of $u$ on $\left[0, T_{0}\right)$ with values in $L^{2}\left(\bar{\Omega}^{c}\right)^{3}$. Suppose that $f \in L^{2}\left(0, T_{0}, L_{\sigma}^{2}\left(\bar{\Omega}^{c}\right)\right)$. Since Theorem 2.3 with $\kappa=q=2$ implies $u \in L^{2}\left(0, T_{0}, L^{6}\left(\bar{\Omega}^{c}\right)^{3}\right)$, we thus get that $u$ satisfies all the assumptions imposed on $u$ at the beginning of Section 5 and in Theorem 5.2 if we take $n_{0}=1, p_{1}=2, f^{(1)}=f, m_{0}=1, \varrho_{1}=2, G^{(1)}=u \mid{\overline{B_{S_{0}}}}^{c} \times\left(0, T_{0}\right), q_{0}=q_{1}=\widetilde{p}=\gamma_{j}=2$ for $1 \leq j \leq 3, q=5 / 4$. Therefore inequality (6.2) follows from (5.27).

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