

# Hydrogen sulfide solubility in 50 wt% and 70 wt% aqueous methyldiethanolamine at temperatures from 283 to 393 K and total pressures from 500 to 10000 kPa

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2	Hydrogen Sulfide Solubility in 50 wt.% and 70 wt.% Aqueous
3	Methyldiethanolamine at Temperatures from 283 to 393 K and Total
4	Pressures from 500 to 10000 kPa
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#### 23 Abstract

The hydrogen sulfide (H<sub>2</sub>S) absorption capacity of a 70 wt.% 24 aqueous 25 methyldiethanolamine (MDEA) solution was investigated in a static-analytic apparatus at temperatures of 283, 353 and 393 K and pressures of 2000, 6000 and 10000 kPa in the 26 27 presence of methane. New experimental data were also produced for a 50.1 wt.% aqueous 28 MDEA at 323 K and pressures of 500 and 3000 kPa as part of the apparatus validation procedure. A model based on electrolyte non-random two-liquid (eNRTL) activity 29 30 coefficient model to describe the liquid phase and Peng-Robinson Equation of State to 31 describe the vapor phase non-idealities was developed for the system H<sub>2</sub>S-MDEA-H<sub>2</sub>O, 32 which can potentially be used also for the system in the presence of methane at low 33 pressures. Vapor pressure measurements of pure MDEA were also performed in the range of 405 - 435 K in an ebulliometer and parameters for the Antoine correlation were 34 35 proposed.

36

37 *Keywords:* gas processing; absorption; hydrogen sulfide; methane; MDEA; high

38 pressure; vapor-liquid equilibrium; vapor pressure

#### 39 1. Introduction

Natural and refinery gas streams usually contain acid gases, carbon dioxide and sulfur
compounds, which must be removed in order to ensure trouble-free and safe operations.
Typical sulfur compounds are hydrogen sulfide, carbonyl sulfide, mercaptans, with the
first one being the most important one as it occurs in the largest concentrations [1].
Hydrogen sulfide (H<sub>2</sub>S) gas content is routinely controlled by absorption into aqueous
methyldiethanolamine (MDEA), which can then be thermally regenerated and reused.

46 A 50 wt.% MDEA-H<sub>2</sub>O concentration is considered a benchmark solvent in H<sub>2</sub>S removal, due to its equilibrium behavior and low corrosion. Aqueous MDEA has been 47 48 long established in the industry due to among others, the amine's availability, low cost and energy requirements, resistance to degradation, ability to meet the 4 ppm 49 specification requirement for pipeline gas and to selectively remove H<sub>2</sub>S over CO<sub>2</sub>, which 50 51 often coexist. MDEA owes its latter characteristic to its structure; as a tertiary amine, aqueous MDEA reacts instantaneously with hydrogen sulfide while it requires more time 52 53 to react with CO<sub>2</sub>. Thus, by regulating the contact time between the solvent and the gas, H<sub>2</sub>S removal to specification and minimum co-absorption of CO<sub>2</sub> can be achieved [2], 54 [3]. 55

56 The motivation of this work has been the investigation of highly concentrated MDEA 57 for the combined H<sub>2</sub>S removal and hydrate control for subsea application. Oil and gas 58 reservoirs are turning sour in the course of time [4], [5], which is tackled today by using triazine to control the H<sub>2</sub>S levels [6]. Main disadvantages of employment of triazine are 59 60 related to the non-regeneration of the solvent, weight, space, transportation and disposal requirements. These constraints are of outmost importance, especially as the available 61 production fields are sourer, deeper and in longer distances from the shore [7]. MDEA is 62 already used offshore as a pH stabilizer [8] facilitating its employment subsea, while the 63 64 fact that, as a polar compound, it has affinity for water, renders highly concentrated aqueous MDEA a good candidate for acting both as a hydrate inhibitor and as an H<sub>2</sub>S 65 removal agent. The solvent could be used and regenerated offshore, supported by new 66 technological developments, such as "subsea on a stick" [9]. 67

This work is a first step in the investigation of this multifunctional solvent, with focus on the effect of total pressure in the  $H_2S$  removal capacity of the solvent. The 70 measurements were conducted at high pressures, up to 10000 kPa, with methane as the 71 pressurization medium, since it is the main constituent of natural gas. Few researchers have previously studied the effect of high-pressure methane for the systems CH<sub>4</sub>-CO<sub>2</sub>-72 MDEA-H<sub>2</sub>O [10], [11] and CH<sub>4</sub>-H<sub>2</sub>S-MDEA-H<sub>2</sub>O; a detailed literature review for the 73 74 latter is provided in Section 2.1. The main finding has been that for both CO<sub>2</sub> and H<sub>2</sub>S-75 contained systems, an increase in total pressure leads to increase in the acid gas partial 76 pressure. To our best knowledge, there are no data reported for the system CH<sub>4</sub>-H<sub>2</sub>S-77 MDEA-H<sub>2</sub>O and MDEA solutions with concentrations higher than 50 wt.% MDEA-H<sub>2</sub>O. A 50.1 wt.% MDEA-H<sub>2</sub>O and a 70 wt.% MDEA-H<sub>2</sub>O system were used in this work 78 79 to obtain vapor-liquid equilibrium data (VLE) with hydrogen sulfide and methane. The new VLE data for the system CH<sub>4</sub>-H<sub>2</sub>S-MDEA-H<sub>2</sub>O with 70 wt.% MDEA-H<sub>2</sub>O mixtures 80 81 were obtained at temperatures of approximately 283, 353 and 393 K and pressures of 2000, 6000 and 10000 kPa. The experiments were performed isothermally and the 82 temperature of 283 K was chosen to simulate the low-temperature subsea conditions 83 84 while the temperature of 393 K was chosen to simulate the high regeneration temperature.

## 85 2. Literature Review

#### 86 2.1 H<sub>2</sub>S-MDEA-H<sub>2</sub>O-makeup gas system

An updated list of available VLE data for the system H<sub>2</sub>S-MDEA-H<sub>2</sub>O, including data with makeup gas, is provided in **Table 1**. The amine concentration is expressed in a weight basis for all reference sources to allow for direct comparisons. Concentrations reported in molarities [12]–[14] have been converted to weight fractions using the density correlations presented by Bernal-García et al. [15]. The solution preparation temperature was assumed to be 298.15 K due to lack of this information.

As also other authors working with the system H<sub>2</sub>S-MDEA-H<sub>2</sub>O have observed, the available data in the literature are rather scattered, especially at low loadings. The literature data have been evaluated for self-consistency and mutual-consistency with reported data in similar experimental conditions, following Chunxi and Fürst's approach [16]. This evaluation was performed in order to decide if some data sets would be

98 excluded during our thermodynamic modeling. During the evaluation, the partial 99 pressures for H<sub>2</sub>S from Kuranov et al. [17], Kamps et al. [18] and Sidi-Boumedine et al. [19], who all report total pressures in the absence of makeup gases, were calculated by 100 101 subtracting the vapor pressure of the solvent calculated by Dalton's Law (Eq. 1). The 102 vapor pressure of H<sub>2</sub>O was calculated by the correlations proposed by NIST for the given 103 temperature ranges while the vapor pressure of MDEA was calculated based on the 104 Antoine correlation fitted to existing and new data as presented in Section 5. Results and 105 **Discussion**.

$$P_{solv}^s = P_{MDEA}^s \cdot x_{MDEA} + P_{H_2O}^s \cdot x_{H_2O}$$
 Eq. 1

106

Li and Shen [13] measured H<sub>2</sub>S solubility in 29.9 wt.% aqueous MDEA at temperatures up to 373 K. During the evaluation of the data, a sharp increase of partial pressure at loadings > 0.7 mol H<sub>2</sub>S/mol MDEA was noticed, resulting in a cross-over of literature data reported for 35 wt.% and 50 wt.% MDEA-H<sub>2</sub>O solutions. For this reason, the data from Li and Shen [13] were not included in our database used in the model parametrization, as chosen also by Huttenhuis et al. [20].

113 Jou and coworkers [12], [21] have published experimental data for a 48.9 wt.% and 114 for a 35 wt.% MDEA solution. Two observations can be made for the low loading region: a) the data with a 35 wt.% [21] and a 48.9 wt.% [12] MDEA solution are very similar and 115 116 b) the deviations between the data with a 48.9 wt.% and a 50 wt.% solution look larger than what one would expect with such similar concentrations. Uncertainty information is 117 118 not given in the first publication of Jou et al. [12], while the authors on their second 119 publication report 3% error in liquid loading and 0.1% full scale (FS) error in pressure. 120 Taking this into account, the deviations related to a) and b) are within the experimental 121 uncertainty. Generally, the data from Jou et al. agree with literature values in different 122 concentrations and temperatures besides at low loadings. For example, good agreement 123 is observed between the data Jou et al. [12] for a 23.4 wt.% aqueous MDEA at 313 K and 124 from two other sources [14], [22] at loadings > 0.4 mol H<sub>2</sub>S/mol MDEA. Any small deviations are justified in terms of reported experimental uncertainties provided by 125 MacGregor and Mather [14] (pressure, loading, composition) as well as by Zoghi and 126 Shokouhi [22] (pressure and composition). At lower loadings, significant deviations are 127

seen between the data by Jou et al. [12] and MacGregor and Mather [14] compared to
Huang and Ng [23] as well as Rogers et al. [24]. These differences are difficult to explain
by the reported uncertainties. At higher loadings, some inconsistencies are also seen, for
example, the data from Kuranov and coworkers [17] for a 32.3 wt.% amine solution are
close to the data reported for a 50 wt.% MDEA solution [12], [23].

No pattern was identified between the analysis method and the uncertainty of the 133 results. Unfortunately, often the uncertainty in loading, which could enlighten the reasons 134 for the scatter observed at low loadings, is not reported. The literature sources reporting 135 uncertainties in either pressure or loading are marked in Table 1. In addition, the 136 137 differences observed in the reported data could also be attributed to the purity of the chemicals. Although most of the authors report the use relatively high-purity chemicals 138 (>98-99 wt.% MDEA, >99 vol.% H<sub>2</sub>S), the chemical's aging (contamination, contact with 139 atmospheric humidity, light degradation etc.) could also have contributed to the 140 differences observed. 141

	<b>T</b> (17)	D (1-D-)	D (1-D-)	T	Malanna	Analy	sis Method	<b>S</b>	ND
wt.% aq. MDEA	I (K)	PH2S (KPA)	Ptot (KPa)	Loading	Makeup gas	Vapor Phase	Liquid Phase	Source	NP
11.8, 23.4, 48.9	298.15, 313.15, 323.15, 373.15, 393.15	0.0013-5890	-	0.00129-3.229	Nitrogen (P <sub>H2S</sub> < 200 kPa)	GC	Iodometric back- titration with thiosulfate	Jou et al. [12]	153
11.9, 20	298.15, 310.95, 338.75, 388.75	13.23-1536.6	-	0.18-2.1703	-	Mass balance	Mass balance	Bhairi, Maddox et al. [26], [27] <sup>+*</sup>	49
23.4	313.15	0.52-1600		0.13-1.725	-	GC	Iodometric back- titration with thiosulfate	MacGregor and Mather $[14]^{+*}$	27
35, 50	313.15, 373.15	0.00183-313	-	0.00410-1.077	Nitrogen (P <sub>H2S</sub> < 350 kPa)	GC	Iodometric back- titration with thiosulfate	Jou et al. [21]*	50
29.9	313.15, 333.15, 353.15, 373.15	1.498-445.7	-	0.082-0.902	Nitrogen (P <sub>H2S</sub> < 200 kPa)	$\begin{array}{l} Mass \ balance \\ (P_{H2S} < 200 \ kPa) \ / \\ GC \ (P_{H2S} > 200 \ kPa) \end{array}$	Iodometric back- titration with thiosulfate	Li and Shen [13]*	43
23.1, 50	313.15, 343.15, 373.15, 393.15	0.0033-3673	-	0.00240-1.74	Nitrogen (P <sub>H2S</sub> < P <sub>amb</sub> )	Mass balance	Iodometric back- titration with thiosulfate	Huang and Ng $[23]^*$	42
23, 50	313.15, 323.15	0.00069-5.268	96-110	0.00219-0.313	Nitrogen (P <sub>H2S</sub> < P <sub>amb</sub> )	FTIR	FTIR	Rogers et al. [24] <sup>+*</sup>	30
11.83, 23.63	~298.15, ~313.1	0.023-1.611	-	0.0101-0.2610	-	Mass balance	Mass balance	Lemoine et al. [28] <sup>+*</sup>	29
18.7, 32.2	313.16, 333.15, 373.15, 393.15, 413.15	-	165.2-4895.9	0.48-1.934	-	Mass balance	Mass balance	Kuranov et al. $[17]^{+*}$	71
48.8	313.11, 353.16, 393.15	-	147.9-2783	0.153-1.428	-	Mass balance	Mass balance	Kamps et al. [18] <sup>+*</sup>	26
46.78	~ 313, ~373	-	6.21-1040	0.039-1.116	-	Mass balance	Mass balance	Sidi-Boumedine et al. [19] <sup>+*</sup>	27
23.7	313.2	14-1361	-	0.505-1.639	-	Mass balance	Mass balance	Zoghi and Shokouhi [22] <sup>+*</sup>	12
35.50	283, 298, (313)	0.141-18.892	690-6900	0.028-0.575	Methane	GC	Iodometric back- titration with thiosulfate	Huttenhuis et al. [25]	30
50	323.15	3-278	493-700	0.096-0.889ª	Methane	GC	Mass balance	Dicko et al. [29]+*	5
50	322.95, 343.15	31-974	1480-7090	0.267-1.042	Methane	GC	Titration with silver nitrate	Sadegh et al. $[30]^{+*}$	39

**Table 1.** Literature VLE data for H<sub>2</sub>S-MDEA-H<sub>2</sub>O including data with makeup gas.

<sup>a</sup>:1g48bal loading, <sup>+</sup>: reported uncertainty in pressure, <sup>\*</sup>: reported uncertainty in H<sub>2</sub>S loading/mole fraction

## 144 **2.2 MDEA-H<sub>2</sub>O system**

145 Vapor-liquid equilibrium (VLE), freezing-point depression (FPD) and molar excess 146 enthalpy  $H^{E}$  data for the binary subsystem MDEA-H<sub>2</sub>O are given in **Table 2**. The data were 147 used to model the binary system first in order to reduce the number of parameters to be fitted 148 for the ternary system H<sub>2</sub>S-MDEA-H<sub>2</sub>O onwards, as it will be further explained later in **Section** 149 **4. Thermodynamic modeling**. Eight points from Chang et al. [31] were excluded due to their 150 deviations from the data by Fosbøl et al. [32]

Property	wt.% aq. MDEA	$T / \Delta T_F (\mathbf{K})$	P (kPa)	Source	NP
VLE	3-78.61	313.15-373.15	6.47-100.40	Kim et al. [33]	61
	10-70	326.15-381.15	13.08-101.67	Xu et al. [34]	34
	30-98.9	350.15-458.65	40-66.7	Voutsas et al. [35]	27
FPD	17.4-39.1	(-3.3)-(-13.8)	101.13	Chang et al. [31]	21
	2.6-39.6	(-0.4)-(-14.2)	101.3	Fosbøl et al. [32]	12
$H^E$	9.6-92.5	298.15-342.45	-	Posey [36]	16
	17.5-96.7	298.15-313.15	-	Maham et al. [37]	26
	41.8-98.4	338.15	-	Maham et al. [38]	9

151 **Table 2.** Literature VLE, FPD and  $H^{E}$  data for the binary system MDEA-H<sub>2</sub>O.

## 152 **2.3 Pure MDEA**

A literature review was also performed for the vapor pressure of MDEA. As seen in Table 3,
the data already reported in the literature cover a large range of temperatures, from 293 to 738
K.

**Table 3.** Literature vapor pressure data for pure MDEA.

<i>T</i> (K)	P <sup>s</sup> (kPa)	Source	NP
293.69-401.97	0.0006-1.4776	Noll et al. [39]	26
406.69-435.50	2.48-7.98	Kim et al. [33]	7
420.45-513.85	3.68-90.44	Daubert et al. [40]	14
467.39, 479.39, 488.15	20, 30, 40	Yang et al. [41]	3
519.7-738.4	98.59-3985	VonNiederhausern et al. [42]	9

#### 158 **3. Experimental Work**

#### 159 **3.1 Materials**

Information for the chemicals used are provided in Table 4. MDEA was used as received 160 from the supplier without further purification. Ultra-pure Millipore water was used in this work 161 to prepare the aqueous amine solutions. Both the amine and the water were degassed 162 independently and they were mixed under vacuum to eliminate presence of air during the 163 experiment. The solutions were prepared gravimetrically in a METTLER PM1200 scale with 164 an accuracy of  $1 \cdot 10^{-5}$  kg. The composition uncertainty is the same for each component in a 165 binary mixture as explained in Appendix, and it was found to be u(w)=0.002 for 50.1 wt.% 166 MDEA-H<sub>2</sub>O and u(w)=0.003 for 70 wt.% MDEA-H<sub>2</sub>O. The gases used in this work are 167 hydrogen sulfide and methane as a makeup gas. 168

169 **Table 4**: Chemical Sample Table.

Component	IUPAC name	CAS	Supplier	Purity	Analysis method
N- methyldiethanolamine (MDEA)	2-[2- hydroxyethyl(methyl) amino] ethanol)	105-59-9	Sigma-Aldrich	$\geq$ 99 wt.%	GC
Water	Oxidane	-	-	Ultra-pure	-
Hydrogen sulfide	Sulfane	7783-06-4	Air Liquide	$\geq$ 99.5 vol.%	GC
Methane	Methane	74-82-8	Air Liquide	$\geq$ 99.995 vol.%	GC

170

171

#### **3.2 Experimental set-ups**

**High-pressure VLE.** The high-pressure vapor-liquid equilibrium (VLE) experiments were conducted in an in-house manufacture by ARMINES employing the static-analytic method [43] The apparatus is designed for measurements with acid gases and can be operated in the pressure range from 0.5 to 19.9 MPa and temperatures, from 223 to 473 K. Temperature regulation with an accuracy of  $\pm$  0.01 K is achieved through immersing the cell into an oil bath. The apparatus is similar to the one previously presented by [29] and its schematic is given in **Figure 1**.



Figure 1: High-pressure VLE setup. DTD: Displacement Transducer Display, DLD: Data Logging
Device, DT: Displacement Transducer, GC: Gas Chromatograph, MS: Mobile Sampler for the analysis
of liquid phase, NRV: Non-Rotating valve, PC: Personal Computer for data acquisition, PT: Pressure
Transducer, T: Thermocouple, V: Valve, VS: Vapor Sampler for the analysis of gas phase, VVP:
variable volume press.

The set-up consists of three distinct parts: a) the equipment for filling up the equilibrium 184 cell, i.e. the variable volume press (VVP), the gas bottles and gas tanks, b) the equilibrium cell, 185 including automatic samplers for the gas and the liquid phase(s) and c) the equipment for the 186 analysis of the samples, i.e. the gas chromatograph. Each of these parts consists of various 187 valves and instrumentation. A variable volume press composed by the variable volume pressure 188 cell, a piston and a displacement transducer, was used to introduce the liquid inside the cell, 189 under vacuum. The transducer measures the piston displacement with an accuracy of  $\pm 1.10^{-5}$ 190 m and, by knowing the exact dimensions of the cell, the exact volume of the solvent introduced 191 was determined. Approximately  $6 \cdot 10^{-6}$  m<sup>3</sup> of solvent were introduced in every experiment. An 192 H<sub>2</sub>S bottle was connected to a gas tank with volume of  $1.61 \cdot 10^{-4} \pm 5 \cdot 10^{-8} \text{ m}^3$ , which was further 193 connected to the cell. The presence of a small gas tank between the gas bottle and the cell was 194 195 dictated as an extra safety barrier in case of leakage of the toxic H<sub>2</sub>S. Pressurization of the cell with methane was done directly from the CH<sub>4</sub> bottle. 196

The equilibrium cell is a sapphire tube standing between two Hastelloy flanges. Kalrez O-197 rings are used for sealing the tube. The upper flange accommodates two non-rotating stem 198 loading valves, for H<sub>2</sub>S and for CH<sub>4</sub>, and the lower flange accommodates two more, only one 199 of which was used for the loading of the liquid solution and the discharge of the cell. The 200 201 temperature is monitored and controlled by two platinum probes and two 100  $\Omega$  Platinum resistance temperature detectors (Pt100) with an uncertainty of  $\pm$  0.02 K. Each of the two 202 located in each flange. They are connected to an HP data acquisition unit and are carefully 203 periodically calibrated. The cell is equipped with two Druck<sup>TM</sup> pressure transducers, one 204 calibrated for 0–3 MPa and the other for 0-30 MPa pressure range respectively. The transducers 205 are maintained at the temperature they were calibrated at and the uncertainty is 0.6 kPa. The 206 volume of the cell is  $33.12 \cdot 10^{-6} \pm 5 \cdot 10^{-8}$  m<sup>3</sup> (or  $32.24 \cdot 10^{-6}$  m<sup>3</sup> when the low-pressure transducer 207 is isolated). A stirring system is integrated to the cell in order to reduce the time of equilibration 208 209 and ensure phase homogeneity. The variable-speed stirrer is composed by a rotating axis inside the cell, two propellers mounted on the rotating axis for stirring both the gas phase and the 210 211 liquid phase and a magnetic rod mounted on the rotating axis in order to allow for rotation of the axis by a stirring motor located below the cell. 212

Agilent software BenchLink is used for online monitoring of pressure and temperature, 213 214 enabling the determination of equilibrium. Once the equilibrium is reached, micro samples can be withdrawn and transferred to the GC for analysis. Automatic sampling is allowed through 215 two capillary samplers (ROLSI®) Armines' patent [44]. Two capillaries are fixed in the 216 cylindrical wall of the cell at levels designed to withdraw vapor and liquid phase samples. The 217 samplers are connected to a PERICHROM model PR-2100 gas chromatograph, through a 218 219 heated transfer line. The temperature selected is higher than the boiling point of the heaviest component (MDEA) to avoid any sample condensation. The chromatograph is equipped with a 220 221 thermal conductivity detector (TCD) and a flame ionization detector (FID), and WINILAB III 222 software is used for GC acquisition and treatment.

**Ebulliometer.** A modified Swietoslawski ebulliometer was used, described earlier in detail by Kim et al. [33]. The apparatus can be operated at temperatures up to 473 K and at subatmospheric and atmospheric pressure. The temperatures were measured with calibrated Pt100 resistance thermosensors with an uncertainty of  $\pm 0.05$  K. A DP1520 pressure controller from Druck<sup>TM</sup> was used, calibrated against a BeamexC5 calibrator with an accuracy of  $\pm 0.03$  kPa. The solution is accommodated inside a  $2 \cdot 10^{-4}$  m<sup>3</sup> glass equilibrium still and the set-up allows for the sampling of both the vapor and the liquid phase.

#### **3.3 Experimental Procedure**

231 **High-pressure VLE**. After thorough cleaning with hot deionized (DI) water and ethanol, the cell and tubings were left to dry and set to vacuum during the previous night. The solution was 232 prepared under vacuum directly inside the VVP and the solution preparation temperature was 233 approximately 298 K. Back-pressure of ca. 500 kPa of methane was applied to the VVP. The 234 solution was introduced inside the cell, and the end displacement position was recorded, so as 235 236 the exact amount of solution added could be calculated. The cell was immersed into the bath, the stirrer was turned on, the temperature of the experiment was set and the system was left to 237 238 equilibrate. Temperature stabilization required approximately 30-60 min, after which the vapor pressure of the solution was recorded. 239

The desired global loading, i.e. mol of  $H_2S$  inside the cell per mol of amine, was first decided and based on the PVT conditions of the  $H_2S$  gas tank before and after the filling of the cell, the amount of  $H_2S$  introduced was determined. The calculations were performed using REFPROP software [45] and a Helmholtz energy-based equation of state developed by [46] for pure  $H_2S$ was used. The global loading was, thus, calculated by:

$$n_{H2S} = n_{H2S,tank}^{before} - n_{H2S,tank}^{after}$$
Eq. 2
$$\alpha_{glob} = \frac{n_{H2S}}{n_{MDEA}}$$
Eq. 3

For the experiments with the 50.1 wt.% MDEA aqueous solution, initially a small amount of 245 H<sub>2</sub>S was introduced and it was left to equilibrate. Reaction of H<sub>2</sub>S and MDEA is fast and 246 247 equilibrium was reached within one hour. Because the total pressure was lower than the minimum required pressure of 500 kPa for the ROLSI® samplers and GC to function, methane 248 249 was added up to 500 kPa. Equilibrium was reached in approximately one hour, and the sampling started. In our experiments, sampling and analysis was conducted only for the vapor phase. 250 251 Higher loadings were reached by adding more H<sub>2</sub>S into the cell and repeating the above-252 mentioned procedure.

For the 70 wt.% MDEA solution VLE investigation, two series of experiments were conducted based on the global loading, one for 0.2 and one for 0.5 mol H<sub>2</sub>S/mol MDEA approximately. The experimental procedure varies in the way that after equilibrium was reached, methane was added in 3 stages, up to 2000, 6000 and 10000 kPa. At each pressure level, sampling and analysis of the vapor phase was performed upon equilibrium. The experiments were performed under isothermal conditions, at 283, 353 and 393 K. At the end of the experiment, the cell was depressurized and emptied safely through a caustic solution
(NaOH) in order to neutralize the system. At each temperature, a new experiment was
conducted using fresh solution. We aimed at having the same global loading at all temperatures,
however it was not practically possible to reach exactly the same loadings in every experiment.
The study at each temperature and global loading lasted approximately one week.

264 The analysis of vapor phase concentration was performed in a GC equipped with a Porapak-R column R80/100 mesh (length 2 m, diameter 2 mm) from RESTEK. The carrier gas was 265 helium at a flow rate of 20 ml/min. A constant temperature program at 363 K was used for the 266 quantification of both methane and hydrogen sulfide. Analysis at 383 K was also performed to 267 268 check for water presence in the vapor phase. In order to check the repeatability of the measurements and to perform uncertainty analysis, five samples at least were withdrawn, the 269 first two of them usually were required to saturate the transfer lines in terms of adsorption. 270 Disturbance to equilibrium was considered negligible due to the small volume of each sample. 271

Knowing the pressure, temperature and the composition of the vapor phase, the density of the vapor phase was estimated using REFPROP software [45]. The amount of  $n_{H2S}$  in the vapor and liquid phase and finally the H<sub>2</sub>S loading in the liquid phase, liquid loading  $\alpha$ , were calculated according to Eq. 4 - Eq. 7.

$$n_{tot}^{\nu} = \rho^{\nu} \cdot V^{\nu}$$
 Eq. 4

$$n_i^{\nu} = n_{tot}^{\nu} \cdot y_i$$
 Eq. 5

$$n_i^l = n_{tot} - n_i^{\nu} \qquad \qquad \text{Eq. 6}$$

$$\alpha = \frac{n_{H2S}^l}{n_{MDEA}}$$
 Eq. 7

where  $\rho^{\nu}$  is the molar density of the gas mixture, calculated using REFPPROP and  $V^{\nu}$  is the 276 volume of the vapor phase. The latter is the difference between the volume of the cell, ca.  $33 \cdot 10^{-10}$ 277  $^{6}$  m<sup>3</sup>, which is known from our calibration data and the volume of the liquid which was 278 estimated by the correlations proposed by [15], assuming that the effect of pressure in the liquid 279 280 volume is negligible. Bernal-García and coworkers measured the density of aqueous MDEA in the whole composition range at temperature range of 263.15 - 363.15 K and, based on their 281 data, calculated the excess molar volumes of the binary systems. For our calculations at the 282 temperature of 393 K which was not studied in the afore-mentioned work, the excess molar 283 volume was extrapolated. It is worth mentioning that the deviations in number of moles of H<sub>2</sub>S 284 calculated by the Ideal Gas Law equation and REFPROP employing the most up-to-date 285

Helmholtz energy-based EoS led to deviations in the liquid loading lower than 1.5% at 283 and
353 K, while the deviations were higher at 393 K (max 2.7%). For more accurate results, we
used the results based on the latter.

**Ebulliometer.** Approximately  $0.8 \cdot 10^{-4}$  m<sup>3</sup> of liquid was charged inside the still, preceding purge with nitrogen. The desired temperature was set and equilibrium was assumed after 10 min of stable pressure and temperature. The vapor pressure of MDEA was measured at the temperature range of 405 – 435 K. Validation of the apparatus was performed by measuring the vapor pressure of water and a 1.5% maximum error from the literature was found in equilibrium pressure.

#### 295 4. Thermodynamic modeling

High pressure VLE. An in-house MATLAB-based rigorous model has been developed to 296 describe the chemical and phase equilibrium for the system H<sub>2</sub>S-MDEA-H<sub>2</sub>O. The same 297 algorithm has been previously used to successfully describe CO<sub>2</sub>-amine-H<sub>2</sub>O systems relevant 298 299 to carbon capture processes [47], [48]. Peng-Robinson EoS [49] with the original alpha function 300 was employed to describe the non-idealities of the vapor phase, coupled with the traditional van der Waals one-fluid mixing rules. The binary interaction parameters for Peng-Robinson EoS in 301 302 this work were set to zero. To account for the non-idealities in the liquid phase, the electrolyte non-random two-liquid (eNRTL) model [50] was utilized. The models are presented in the 303 304 Appendix. The required critical parameters and acentric factors for pure components are given in Supplementary Information. 305

The chemical reactions assumed in the liquid phase are the ionization of water, the protonation of MDEA and the dissociation of H<sub>2</sub>S (R. 1-3). The second dissociation reaction of hydrogen sulfide, from bisulfide to sulfide, is not considered in our model due to the low concentration of S<sup>2-</sup> in the solution and in order to reduce the number of parameters in the model [16].

$$2H_2 0 = 0H^- + H_3 0^+$$
 R. 1

$$MDEA + H_3O^+ = MDEAH^+ + H_2O$$
 R. 2

$$H_2S + H_2O = H_3O^+ + HS^-$$
 R. 3

The chemical equilibrium constants as well as Henry's constant for hydrogen sulfide are described by Eq. 8, parametrized according to **Table 5**, where *x* stands for either the chemical equilibrium constant  $K_{eq}$  or Henry's constant  $H_{H2S}$ . Temperature is expressed in K and Henry's constant for hydrogen sulfide in kg·atm.

$$\ln(x) = A + \frac{B}{T} + C\ln(T) + DT$$
 Eq. 8

316

**Table 5**. Mole fraction-based parameters for Eq. 8, reported only with their significant digits.

	А	В	С	D	Reference
Chemical Equilibrium constant for R. 1	132.89	-13445	-22.477	0	Posey [36]
Chemical Equilibrium constant for R. 2	-60.03	-1974	7.533	0	Oscarson et al. [51]
Chemical Equilibrium constant for R. 3	214.58	-12995	-33.547	0	Posey [36]
Henry's constant for H <sub>2</sub> S	342.595	-13237	-55.0551	0.05957	Edwards et al. [52]

318

The vapor pressure for hydrogen sulfide and water is estimated using the Riedel correlation (Eq. 9) where *T* expressed in K and  $P^{sat}$  in Pa. The parameters are presented in **Table 6**. MDEA vapor pressure has been measured in this work and fitted to Antoine correlation. The Antoine parameters used in this work can be found in **Section 1.5.2**.

$$\ln(P^{sat}) = A + \frac{B}{T} + C\ln(T) + DT^{E}$$
 Eq. 9

**Table 6.** Parameters for pure component vapor pressure correlations for Eq. 9.

Component	Model	А	В	С	D	Е	Reference
$H_2S$	Riedel	106.47	-5018	-13.306	-0.09	-0.13	DIPPR [53]
$H_2O$	Riedel	73.649	-7258	-7.304	4.2E-06	2	DIPPR [53]

324

A sensitivity analysis was performed to evaluate the significant numbers in the parameters retrieved from the literature. In **Table 5** and **Table 6** the parameters are provided only with their significant digits.

The adjustable parameters for the eNRTL model are the non-randomness factors,  $\alpha$ , and the energy parameters,  $\tau_{ij}$ . The optimization of the H<sub>2</sub>S-MDEA-H<sub>2</sub>O system requires the regression of a total of 78 parameters. In order to reduce this high number of parameters to be adjusted, the following steps have been taken: 332 I) All non-randomness factors  $\alpha$  have been given fixed values according to **Table 7**.

II) The energy parameters for the subsystem H<sub>2</sub>S-H<sub>2</sub>O have been fixed to the default values
used in Aspen Plus V10 simulation software (Table 7).

III) The energy parameters for the subsystem MDEA-H<sub>2</sub>O have been fixed to the values
obtained by the regression of the literature data presented in Subsection 2.2 MDEA-H<sub>2</sub>O
system.

As a result, the number of parameters is reduced to 36. The temperature dependency of the energy parameters is described by Eq. 10, where  $a_{ij}$  and  $b_{ij}$  were fitted to experimental data.

$$\tau_{ij} = a_{ij} + \frac{b_{ij}}{T}$$
 Eq. 10

The fixed non-randomness factors and fixed energy parameter values are presented in **Table 7**, where *m* denotes molecule and *c*-*a* cation-anion (salt). The non-randomness factors were fixed at 0.2 for molecule-molecule and water-salt interactions, and at 0.1 for the H<sub>2</sub>S-salt and MDEAsalt interactions, according to Hessen and coworkers [54].

344

#### **Table 7:** Fixed parameters of eNRTL model used in this work.

Non-randomness factors, $\alpha$								
Com	ponents							
i	j	ij	ji					
m	m	0.2	0.2					
H <sub>2</sub> O	c-a	0.2	0.2					
$H_2S$	c-a	0.1	0.1					
MDEA	c-a	0.1	0.1					

Energy parameters, $ au_{ij}$									
Co	omponents		a	1	o				
i	j	ij	ji	ij	ji				
$H_2O$	$H_2S$	0	0	0	0				
$H_2O$	H <sub>3</sub> O <sup>+</sup> -OH <sup>-</sup>	8	-4	0	0				
H <sub>2</sub> O	$H_3O^+-HS^-$	8	-4	0	0				

$H_2S$	$H_3O^+-OH^-$	15	-8	0	0
$H_2S$	H <sub>3</sub> O <sup>+</sup> -HS <sup>-</sup>	15	-8	0	0

347 The optimization routine used in this work is Particle Swarm Optimization (PSO), developed by Kennedy and Eberhart [55]. This algorithm allows for the optimization of continuous non-348 linear functions, using particle swarm methodology. The advantage of this optimization routine 349 is that it uses random initialization, thus, unlike other optimization methods, its convergence is 350 not dependent on the first approximations. In order to avoid local minima and find an optimal 351 352 solution, local best topology was used [56], [57]. The PSO parameters are swarm size of 40, maximum number of iterations 600 in 3 loops which terminate once the optimized value 353 deviates more than 10<sup>-4</sup> (tolerance criterion) from its preceding one or if less than 1% 354 improvement is achieved during 60 iterations. The minimization of the absolute average relative 355 356 deviation (AARD) shown in Eq. 11, was chosen as the objective function, where Y was either the partial pressure of  $H_2S$ ,  $P_{H2S}$ , or the total pressure,  $P_{tot}$ . 357

$$F_{obj}(\%) = \frac{1}{N} \sum_{i}^{N} \frac{|Y_i^{exp} - Y_i^{pred}|}{Y_i^{exp}} \cdot 100$$
 Eq. 11

#### 358 5. Results and Discussion

#### 359 **5.1 Experimental results**

High-pressure VLE. The experimental vapor-liquid equilibrium data obtained in this work with 50.1 wt.% and 70 wt.% MDEA solution for the system  $CH_4$ - $H_2S$ -MDEA- $H_2O$  at various pressures and temperatures are presented in **Table 8** and **Table 9**. As mentioned earlier, knowing the experimental uncertainty of reported data could possibly help us understand the scatter observed in the data for the system  $H_2S$ -MDEA- $H_2O$ . Therefore, we performed a thorough investigation of our measurements' uncertainty in order to properly evaluate our data and conclude on the impact of experimental uncertainty on our results.

We have reported the combined uncertainties employing the Law of propagation of uncertainty according to NIST guidelines [58]. The new data are accompanied by the standard uncertainties for total pressure and temperature as well as the combined uncertainties for the

partial pressure of hydrogen sulfide, the global and the liquid loading. It was found that the 370 main contributor to the uncertainty of the partial pressure of H<sub>2</sub>S is the total pressure of the 371 system, as can be observed by the increasing uncertainty of  $P_{\rm H2S}$  for increasing total pressure. 372 The main contribution to the global loading uncertainty is associated with the loading itself, 373 374 while the uncertainty of the liquid loading is mostly affected by the uncertainty of the total moles of H<sub>2</sub>S introduced in the cell. The repeatability of our measurements was taken into 375 account by virtue of the multiple samples analyzed on the GC at each equilibrium pressure and 376 temperature. The complete uncertainty analysis can be found in Supporting Information. 377

378

Table 8. Experimental vapor-liquid equilibrium data and their corresponding combined
uncertainties at total pressure of 500 kPa (and one measurement at total pressure 3000 kPa) and
temperature of 323 K for the system CH<sub>4</sub>-H<sub>2</sub>S-MDEA-H<sub>2</sub>O and 50.1 wt.% aqueous MDEA.
Methane is used as makeup gas.

Т	$P_{ m tot}$	$P_{\rm H2S}$	$u_c(P_{H2S})$	$lpha_{ m glob}$	$u_c(\alpha_{glob})$	$lpha_{ m liq}$	$u_c(\alpha_{liq})$	NS
Κ	kPa	kPa	kPa	$mol H_2S$	mol H <sub>2</sub> S	$mol H_2S$	$mol H_2S$	
				global/	global/	iiquid/	iiquid/	
				mol MDEA	mol MDEA	mol MDEA	mol MDEA	
				Experim	ent 1			
322.98	493.81	2.99	0.03	0.096	0.003	0.095	0.001	9
322.98	480.01	11.27	0.12	0.214	0.005	0.211	0.002	9
322.98	500.72	49.11	0.43	0.490	0.005	0.477	0.002	7
322.98	604.01	177.59	1.20	0.822	0.006	0.775	0.003	10
				Experim	ent 2			
322.98	493.92	2.60	0.02	0.085	0.003	0.084	0.002	6
322.98	493.50	22.33	0.19	0.312	0.004	0.303	0.002	6
322.98	498.13	72.79	0.56	0.588	0.006	0.559	0.003	8
322.98	530.82	139.10	0.93	0.760	0.013	0.703	0.006	5
322.98	545.53	168.46	1.06	0.820	0.039	0.751	0.020	9
322.98	3106.96	179.67	1.52	0.820	0.039	0.745	0.020	8

<sup>a</sup> Standard uncertainties not included above are u(T) = 0.02 K, u(P) = 0.6 kPa.

Table 8 and Figure 2 reveal information regarding both the reproducibility of the measurements in this work as well as their comparison with the literature for the system CH<sub>4</sub>-H<sub>2</sub>S-MDEA-H<sub>2</sub>O with a 50-50.1 wt.% MDEA solution at approximately 323 K. The measured vapor fractions of methane and hydrogen sulfide are reported in Supporting Information together with the uncertainty analysis. Our measurements in the presence of 500 kPa of methane

were performed in two different experiments, and as one can observe in the figure, the same 389 390 behavior is followed and the measurements can be reproduced. The data obtained in this work are in agreement with the data reported by Dicko et al. [29] under similar conditions. These 391 data together with Sadegh et al.'s data [30] at total pressure of 1500 kPa and 7000 kPa show 392 that, for a given liquid loading, an increase in the total pressure of the system leads to an increase 393 in the H<sub>2</sub>S partial pressure. Our single measurement at total pressure of 3000 kPa for this system 394 follows this trend, too. An exception is the last point reported by Dicko et al. at  $\alpha_{liq} = 0.832$  mol 395 H<sub>2</sub>S / mol MDEA, which also differs from the trend in our data. This point is measured at total 396 pressure 700 kPa but lies between the data reported by Sadegh et al. at 1500 kPa and 7000 kPa 397 total pressure. Here it is important to mention that the measurements reported by Dicko et al. 398 399 are global loadings, and the liquid loadings shown in Figure 2 are the ones calculated by the 400 authors.



401

402 **Figure 2:** Equilibrium H<sub>2</sub>S partial pressures as a function of liquid loading and total pressure 403 for 50 wt.% MDEA-H<sub>2</sub>O at 323 K. ●  $P_{tot} = 500$  kPa (This work, Experiment 1),  $\circ P_{tot} = 500$ -404 600 kPa (This work, Experiment 2), ◆  $P_{tot} = 500$ -700 kPa [29], ■  $P_{tot} = 1500$  kPa [30], ×  $P_{tot} =$ 405 3000 kPa (This work), ▲  $P_{tot} = 7000$  kPa [30].

406

Table 9. Experimental vapor-liquid equilibrium data and their corresponding combined
uncertainties as a function of total pressure and temperature for the system CH<sub>4</sub>-H<sub>2</sub>S-MDEAH<sub>2</sub>O and 70 wt.% aqueous MDEA<sup>a</sup>. Methane is used as makeup gas.

	Т	$P_{\rm tot}$	$P_{\rm H2S}$	$u_c(P_{H2S})$	$lpha_{ m glob}$	$u_c(\alpha_{glob})$	$lpha_{ m liq}$	$u_c(\alpha_{liq})$	NS
	K	kPa	kPa	kPa	mol H <sub>2</sub> S				
					mol	mol	mol	mol MDEA	
					MDEA	MDEA	MDEA		
28	3.00	2011.87	3.48	0.03	0.232	0.003	0.231	0.003	9
28	3.00	6030.85	3.85	0.05	0.232	0.003	0.231	0.003	10
28	3.00	10052.50	4.68	0.04	0.232	0.003	0.231	0.003	7
35	2.99	1976.07	106.23	0.92	0.239	0.002	0.211	0.002	5
35	2.99	3954.66	108.30	0.98	0.239	0.002	0.210	0.002	10
35	2.99	5957.76	108.03	1.03	0.239	0.002	0.210	0.002	7
35	2.99	7976.36	111.95	1.04	0.239	0.002	0.209	0.002	6
35	2.99	9988.18	111.42	1.12	0.239	0.002	0.208	0.002	6
39	3.00	2024.40	375.10	1.46	0.246	0.002	0.167	0.002	8
39	2.99	5979.36	376.17	1.72	0.246	0.002	0.165	0.002	10
39	3.00	9925.29	364.18	1.97	0.246	0.002	0.167	0.002	8
28	3.00	1975.74	13.13	0.15	0.488	0.002	0.484	0.002	7
28	3.00	5990.55	17.37	0.28	0.488	0.002	0.482	0.002	5
28	3.00	10045.17	21.56	0.24	0.488	0.002	0.480	0.002	6
35	2.92	2006.00	264.36	1.30	0.478	0.002	0.415	0.002	7
35	2.92	5980.37	281.97	1.57	0.478	0.002	0.408	0.002	8
35	2.92	9975.23	300.30	1.67	0.478	0.002	0.402	0.002	8
39	3.05	974.22	834.43	0.49	0.484	0.002	0.304	0.003	3
39	3.00	2034.17	818.12	2.26	0.484	0.002	0.308	0.003	8
39	3.01	5893.45	806.74	3.34	0.484	0.002	0.309	0.003	7
39	3.00	9915.85	809.32	3.68	0.484	0.002	0.307	0.003	9

410 <sup>a</sup> Standard uncertainties not included above are u(T) = 0.02 K, u(P) = 0.6 kPa.

The observation of increased H<sub>2</sub>S partial pressure upon increase in total pressure can be made also for the 70 wt.% aqueous MDEA system for the temperatures of 283 K and 353 K. The deviations in partial pressure are higher for higher global loadings. On the other hand, the liquid loading remains unchanged at 283 K while the one at 353 K seems to decrease. At 393 K, not clear trends are shown. This behavior is noticed for all global loadings, though the fact that the water present in the vapor phase could not be quantified through the GC analysis, and it was therefore calculated based on the vapor pressure of the solvent under the assumption that it was constant with increasing total pressures, might have its share on the latter. The effect of temperature is the expected one given the exothermic nature of the chemical reactions; the lower the temperature, the higher the absorption of  $H_2S$  in the liquid phase at constant partial pressure of hydrogen sulfide. The features discussed above are illustrated in **Figure 3** and **Figure 4**. It is worth mentioning that error bars representing the uncertainty in pressures and loadings are included in the figures, however uncertainties in pressure are too low to be visible.





425

**Figure 3:** Experimental H<sub>2</sub>S solubility in a 70 wt.% MDEA-H<sub>2</sub>O system with methane as makeup gas at total pressure of 2000 kPa at temperature; • 283 K,  $\blacktriangle$  353 K and = 393 K. Error bars for both H<sub>2</sub>S partial pressure and loading are included.



Figure 4:  $H_2S$  liquid phase loading of a 70 wt.% MDEA- $H_2O$  system with methane as makeup gas as a function of total pressure and temperature; • 283 K,  $\blacktriangle$  353 K and = 393 K. Error bars for both total pressure and liquid loading are included.

434

Although there is a clear trend of the pressure effect on the partial pressure of H<sub>2</sub>S, taking 435 436 into account the uncertainties, it can be seen that the deviations in liquid loading are similar to the experimental uncertainty. In fact, at 283 K and for global loading 0.232, no change at all in 437 liquid loading is observed. The fact that hydrogen sulfide is chemically bound to the amine 438 reinforces the argument that the differences in loading are due to uncertainty in measurements. 439 440 The amount of methane dissolved in the liquid phase is too low to have an impact on the reaction 441 of hydrogen sulfide with the amine solution which is an exothermic reaction whose reversion 442 requires high amounts of energy. Overall, it is observed that the effect of increasing the total pressure from 2000 kPa to 10000 kPa in terms of H<sub>2</sub>S loading in a 70 wt.% aqueous MDEA at 443 temperatures of 283 K, 353 K and 393 K is not significant and, in most cases, it is within or 444 445 very close to experimental uncertainty.

As Sadegh et al. [30] showed, taking into account the gas fugacities is adequate to explain 446 the deviations of the equilibrium H<sub>2</sub>S pressures at different total pressures for a 50 wt.% 447 aqueous MDEA. Indeed, Figure 5 shows how the fugacity exhibits the same behavior for all 448 data obtained in a 50-50.1 wt.% MDEA-H<sub>2</sub>O solution in the presence of methane from different 449 450 literature sources. The figure is similar to one provided by Sadegh et al. [30], this time enriched with our data at total pressure of 500 and 3000 kPa demonstrating the same behavior. The 451 fugacities were calculated using Peng-Robinson EoS with binary interaction parameters set to 452 zero. As far as the data obtained for the 70 wt.% aqueous MDEA are concerned, the fugacity 453

454 can explain the partial pressure trend observed for our data at 283 K and 353 K. At 393 K, the
455 uncertainty in liquid loadings are such that no solid conclusions can be drawn.





**Figure 5:** Equilibrium H<sub>2</sub>S fugacities as a function of liquid loading and total pressure for 50 wt.% MDEA-H<sub>2</sub>O at 323 K. •  $P_{tot} = 500$  kPa (This work, Experiment 1),  $\circ P_{tot} = 500-600$  kPa (This work, Experiment 2), •  $P_{tot} = 500-700$  kPa [29], •  $P_{tot} = 1500$  kPa [30], ×  $P_{tot} = 3000$  kPa (This work, Experiment 2), •  $P_{tot} = 7000$  kPa [30].

461

The effect of amine concentration was also studied by means of comparison with reported 462 data in the literature at 283 K and 393 K, shown in Figure 6 and Figure 7 respectively. Data at 463 353 K in our range of loading and pressure are not available in the literature, therefore no 464 comparison could be performed. A clear effect of increasing molar concentration with 465 increasing amine concentration and constant  $H_2S$  pressure can be seen in the comparison 466 performed at 283 K. This can be expected since the more amine available, the higher the 467 468 capacity of the solvent. It is important to state that the literature data are reported only from one reference source [25] where methane makeup gas was also used. Because of the effect of 469 470 methane presence, we have plotted the available data at similar total pressures; our data only for total pressure of 6000 kPa and the literature data at total pressure of 6900 kPa in order to 471 472 allow for a fairer comparison. The molar concentration of H<sub>2</sub>S is also increasing with amine 473 content in the solution at 393 K, but only up to 50 wt.%. Our data at 70 wt.% overlap with the 474 literature data obtained in a 50 wt.% aqueous MDEA study. This is illustrated in Figure 7 where

- 475 we have only plotted the data with very little methane or with total pressure of 2000 kPa from
- 476 our work.



Figure 6: Hydrogen sulfide molar concentration in the liquid phase for the system CH<sub>4</sub>-H<sub>2</sub>SMDEA-H<sub>2</sub>O as a function of partial pressure and amine concentration at total pressures 60006900 kPa and at 283 K; ◆ 35 wt.% MDEA-H<sub>2</sub>O [25], ■ 50 wt.% MDEA-H<sub>2</sub>O [25], ○ 70 wt.%
MDEA-H<sub>2</sub>O (This work).



Figure 7: Hydrogen sulfide molar concentration in the liquid phase for the system H<sub>2</sub>S-MDEAH<sub>2</sub>O as a function of partial pressure and amine concentration at 393 K; ◆ 11.8 wt.% MDEAH<sub>2</sub>O [12], ■ 18.7 wt.% MDEA-H<sub>2</sub>O [17], + 23.1 wt.% MDEA-H<sub>2</sub>O [23], - 32.2 wt.% MDEA-

486 H<sub>2</sub>O [17], ● 48.8 wt.% MDEA-H<sub>2</sub>O [12], × 48.8 wt.% MDEA-H<sub>2</sub>O, ▲ 50 wt.% MDEA-H<sub>2</sub>O
487 [23], ○ 70 wt.% MDEA-H<sub>2</sub>O (This work).

488

Hydrogen sulfide can react directly with MDEA through a typical acid-base reaction [2]. At 489 490 the same time, the presence of water would enhance the acid gas uptake through the dissolution of hydrogen sulfide as well the protonation of the amine. Therefore, we could identify two 491 possible mechanisms through which H<sub>2</sub>S is absorbed; one directly into the amine and one via 492 water. Moreover, hydrogen sulfide absorption in MDEA-H<sub>2</sub>O is the result of both physical and 493 494 chemical absorption. Therefore, in order to provide a good discussion about the behavior observed in Figure 6 and Figure 7, the physical absorption of hydrogen sulfide into MDEA-495 496 H<sub>2</sub>O systems should be taken into account. To our best knowledge, only Rinker and Sandall 497 [59] have reported such information. They measured H<sub>2</sub>S solubility in protonated aqueous 498 MDEA and their measurements showed that the solubility increases with amine content. 499 Although the available data cover 0-50 wt.% MDEA-H<sub>2</sub>O systems, it can be assumed that the 500 same trends would be followed and the physical absorption of H<sub>2</sub>S in to a 70 wt.% aqueous MDEA is higher than in a 50 wt.% aqueous MDEA. 501

502 Based on the above, the fact that  $x_{H2S}$  is not increased with amine content from 50 to 70 wt.% at 393 K and constant H<sub>2</sub>S pressure indicates that the contribution of the chemical absorption 503 decreases as the amine content increases. This can be also confirmed by observing the slope of 504 indicative tendency curves in Figure 7 (better illustrated in Figure S7 in Supporting 505 506 Information, where non-logarithmic scale is used for the y axis). The slope reveals information 507 about the absorption capacity of the systems. It is observed that as the amine composition increases, the P-x curve has a lower slope (apparent Henry's constant). The lower the slope, the 508 closer to linearity and, thus, higher physical absorption. For example, at 500 kPa, the apparent 509 Henry's constant is 535 kPa·m<sup>3</sup>/kmol for 11.8 wt.% MDEA-H<sub>2</sub>O and 300 kPa·m<sup>3</sup>/kmol for 48.8 510 wt.% MDEA-H<sub>2</sub>O at 393 K. This behavior is followed also at higher pressure; at 3000 kPa, the 511 apparent Henry's constant is 1169 kPa·m<sup>3</sup>/kmol for 11.8 wt.% MDEA-H<sub>2</sub>O and 715 512 kPa·m<sup>3</sup>/kmol for 48.8 wt.% MDEA-H<sub>2</sub>O at 393 K. Unfortunately, our data are too few to assess 513 the P-x linearity for 70 wt.% MDEA-H<sub>2</sub>O, nonetheless it can be said that the chemical 514 contribution in the overall H<sub>2</sub>S uptake is decreased. In the case of low temperatures such as in 515 516 our studied temperature of 283 K, these effects could probably not be visible because the absorption capacity is very high and our data as well as the data reported in the literature are 517 produced for low H<sub>2</sub>S partial pressure. 518

**Ebulliometer.** The measurements conducted in the ebulliometer are shown in **Table 10**. The main limitation of ebulliometric measurements is the absence of stirring. Experimental measurement of the vapor pressure of the binary mixtures used in this work was not possible because two phases formed, associated with the high viscosity of pure MDEA, i.e. ca. 77 mPa·s at 298.15 K [60]–[62]. Therefore, only the vapor pressure of MDEA was measured.

525

527

		P	» (kPa)		
<i>T</i> (K)		DIP	PR	This work (	(Table 11)
	Experimental	Predicted	ARD (%) <sup>b</sup>	Predicted	ARD (%) <sup>b</sup>
405.34	1.79	1.95	9%	1.79	0%
411.00	2.29	2.53	10%	2.34	2%
415.31	2.79	3.06	10%	2.86	2%
418.58	3.29	3.54	7%	3.31	1%
421.73	3.79	4.05	7%	3.80	0%
424.52	4.29	4.55	6%	4.28	0%
427.21	4.79	5.09	6%	4.80	0%
429.49	5.29	5.59	6%	5.28	0%
431.60	5.79	6.08	5%	5.76	0%
433.49	6.29	6.56	4%	6.22	1%
435.34	6.79	7.05	4%	6.71	1%

**Table 10.** Experimental vapor pressure  $P^s$  / kPa for pure MDEA<sup>a</sup>.

<sup>a</sup> Standard uncertainties are u(T)=0.1 K, u(P)=0.1 kPa.

528 
$${}^{b}ARD (\%) = \frac{|P_{s}^{pred} - P_{s}^{exp}|}{P_{s}^{exp}} \cdot 100$$

#### 529 5.2 Modeling results

In this section, we present first the results from the ebulliometer following by the modeling
results for the high-pressure VLE data, since the first ones are used in the model parametrization
for the H<sub>2</sub>S-MDEA-H<sub>2</sub>O equilibrium.

533

Ebulliometer. The Antoine correlation was fitted to available data from the literature (Table 534 3) as well as the newly obtained data of this work, covering a large range of temperatures and 535 pressures. In **Table 10**, our experimental measurements are compared with the predicted vapor 536 537 pressures by our fitted Antoine correlation and the DIPPR equation. At the temperature range of 405-435 K studied in this work, the absolute relative deviation (ARD) between the 538 539 experimental and the estimated value is 7% with DIPPR equation and 1% in our correlation, which has been fitted to available data in the literature covering temperatures from 293 K to 540 541 738 K. The new parameters for Antoine correlation proposed for the estimation of the vapor 542 pressure of MDEA, are shown in **Table 11**. The Average Absolute Relative Deviation (AARD) is 4% for our correlation and 30% for DIPPR. The high deviation for DIPPR equation is mainly 543 544 due to the vapor pressure predictions at temperatures higher than 530 K, which explains the high AARD. In the fitting, we excluded the data from Kim et al. [33] which are slightly higher 545 than the data obtained on the same conditions by Daubert et al. [40] as well as our 546 547 measurements. However, including those data leads to modeled vapor pressures with only the slightly higher AARD of 5%. 548

549

**Table 11.** Parameters for the Antoine correlation<sup>a</sup> for pure MDEA vapor pressure.

		А	В	С
	MDEA	$9.676\pm0.014$	$-1965.6\pm8.9$	$-99.33\pm0.69$
551	$a \log_{10} P^s = A + $	$-\frac{B}{T+C}$ . T in K. P in	Pa. Temperature ran	ge: 294 – 738 K

552	High-pressure VLE. The parameter fitting for the MDEA-H <sub>2</sub> O system returned satisfactory
553	AARDs for all three variables fitted, i.e. VLE ( $P_{tot}$ ), FPD and $H^E$ , as described in Sections 2.2
554	and <b>4</b> . The calculated AARDs for each variable are shown in <b>Table 12</b> .

**Table 12:** AARDs for the fitted  $P_{tot}$ , FPD and  $H^E$  for the MDEA-H<sub>2</sub>O system.

Variable	Source	AARD (%)
$P_{\rm tot}$	Kim et al. [33]	1.1

	Xu et al. [34]	1.9
	Voutsas et al. [35]	6.4
	Overall	2.5
FPD	Chang et al. [31]	10.3
	Fosbøl et al. [32]	4.4
	Overall	6.0
$H^{\rm E}$	Posey [36]	7.6
	Maham et al. [37]	3.1
	Maham et al. [38]	11.5
	Overall	7.4

The model can predict very well the total pressure of the binary system, as witnessed above by the low AARD. The excess enthalpy  $H^{E}$  can be well predicted at temperatures of 298.15 and 313.15 K, though the model yields lower excess enthalpies at 338.15 K for MDEA concentrations lower than 85 wt.%. At this temperature, the model was fitted to experimental data reported by Maham et al. [38] which shows the highest AARD. The corresponding figures for the total pressure, excess enthalpy and the freezing point depression are presented in Supporting Information.

565 The fixed parameters in **Table 7** and the regressed parameters for the binary subsystem MDEA-H<sub>2</sub>O (Supplementary Information) were used for the regression of the ternary system 566 H<sub>2</sub>S-MDEA-H<sub>2</sub>O. Initially, all the data from **Table 1** were used for the parametrization of the 567 model, except for the data from Li and Shen, as well as the data in the presence of methane 568 569 (Case A). The data obtained in the presence of nitrogen were all included. The scatter already 570 discussed earlier at low loadings resulted in high AARD, especially for the data points reported in terms of partial pressure of H<sub>2</sub>S. The high deviations are also attributed to the much lower 571 572 values of partial pressures in comparison with total ones, leading to higher relative numbers. Therefore, we have decided to also perform the data regression excluding all data at loadings 573 lower than 0.05 mol H<sub>2</sub>S/mol MDEA (Case B). This indeed improved substantially the fitting 574 575 of the partial pressures, as one can see in the AARDs in Table 13, from approximately 30% to 18%. The parity plot for the predicted and experimental values is shown in Figure 8 while 576 577 Figure 9 shows the difference between predicted and experimental H<sub>2</sub>S partial pressure as a function of the experimental value. 578



Figure 8: Parity plot for different literature sources; ○ Lemoine et al. [28], □ Huang and Ng
[23], ◊ Rogers et al. [24], × (MacGregor and Mather [14], ⊽ Jou et al. [21], + Jou et al. [12],
Zoghi and Shokouhi [22], □ Maddox et al. [27], (-) y=x.



584

Figure 9: Difference between predicted and experimental H<sub>2</sub>S partial pressure as a function of
the experimental value. ○ Lemoine et al. [28], □ Huang and Ng [23], ◊ Rogers et al. [24], ×
MacGregor and Mather [14], 7 Jou et al. [21], + Jou et al. [12], □ Zoghi and Shokouhi [22], □
Maddox et al. [27].

589

Significant scatter can be seen in the plots above at the lower pressures. At pressures P < 1</li>
kPa, the model underestimates the data from Lemoine et al. [28] while overestimating the data
from Jou et al. [12] and MacGregor and Mather [14]. The accuracy of the model is good for the

data from Rogers et al. [24] and Jou et al. [21], with some data being underpredicted. The visual

observations are depicted on the bias and AARD (%) calculations presented in Table 13. The 594 negative bias whose absolute value is the same as the AARD for Lemoine et al.'s work shows 595 that all data have been underestimated by the model. In addition, the fact that the AAD for this 596 source is 0.2 kPa shows that the high AARD of 27% is due to the low values in partial pressures. 597 At intermediate pressures, the scatter is less pronounced, but still the model overpredicts the 598 data of MacGregor and Mather [14] and Jou et al. [12]. The BIAS and AARDs for these two 599 are -37% and 38% and -27% and 27% respectively, while the rest of the sources show AARDs 600 lower than 17%. At higher pressures, both Figure 8 and Figure 9 show that the model can 601 602 predict well the literature data.

603 Overall, maximum AARD was found for the data from MacGregor and Mather [14] showing an almost 50% AARD in Case A and 38% in Case B. The minimum deviations observed were 604 for the data from Maddox et al. [27] in Case A (13%) and from Huang and Ng [23] in Case B 605 (9%). From the three experimental sets of total pressure, the one reported by Kuranov et al. 606 showed the lowest deviations for both cases. Similar observations were made also by 607 Huttenhuis et al. [20] during the evaluation of their model developed for the H2S-MDEA-H2O 608 system. Although the model framework they used (electrolyte EoS for both phases) differs from 609 ours, their model predictions also showed highest deviations for the data from MacGregor and 610 Mather and lowest for the data from Maddox et al. [27] and Kuranov et al. [17]. Figure 10 611 shows experimental and modelled values for a 50 wt.% aqueous MDEA system as a function 612 of temperature in Case B. 613



**Figure 10**: Hydrogen sulfide loading for 50 wt.% MDEA-H<sub>2</sub>O as a function of partial pressure and temperature; (red) 313 K, (green) 343 K, (blue) 373 K, (orange) 393 K; (–) model,  $\Box$ Huang and Ng [23],  $\Delta$  Rogers et al. [24] (1998),  $\Box$  Jou et al. [21] (1993). Regression in Case B.

The differences in  $H_2S$  partial pressure noticed in the literature data as well as in our data 618 obtained in the presence of methane for relatively low total pressure levels, are comparable to 619 the accuracy of the model. Therefore, since also the effect of methane in the liquid loading has 620 621 been found to be negligible for a 70 wt.% MDEA-H<sub>2</sub>O, we also fitted the model to data available in the presence of methane. However, the code was not modified but, instead, the data for partial 622 pressure of H<sub>2</sub>S and loading were used as if methane was not present. Only data with maximum 623 total pressure of 2000 kPa were considered, due to the more significant  $P_{H2S}$  deviations 624 observed at higher pressures in the literature for a 50 wt.% aqueous MDEA (Case C). To sum 625 up, three cases were studied: 626

627 <u>Case A</u>. Regression of all available data in the absence of methane.

628 <u>Case B</u>. Regression of all available data in the absence of methane and loadings  $\alpha > 0.05$ 629 mol H<sub>2</sub>S/mol MDEA.

630 <u>Case C</u>. Regression of all available data in the absence of methane and loadings  $\alpha > 0.05$ 631 mol H<sub>2</sub>S/mol MDEA, and the data in the presence of methane, loadings  $\alpha > 0.05$  mol H<sub>2</sub>S/mol

632 MDEA and maximum total pressures  $P_{tot}$  of 2000 kPa.



633

**Figure 11**: Hydrogen sulfide loading for 70 wt.% MDEA-H<sub>2</sub>O as a function of partial pressure and temperature; (brown) 283 K, (green) 353 K, (orange) 393 K; (–) model,  $\circ$  This work. Regression in Case C.

		Case A	L			Case B				Case C		
Source	Pressure range	Bias	AAD	AARD	Pressure range	Bias	AAD	AARD	Pressure range	Bias	AAD	AARD
	kPa	(%)	kPa	(%)	kPa	(%)	kPa	(%)	kPa	(%)	kPa	(%)
					Partial	pressur	e, $P_{\rm H2S}$					
Lemoine et al. [28]	0.023-1.611	-42.8	0.1	42.8	0.176-1.611	-27.0	0.2	27.0	0.176-1.611	-23.7	0.1	23.7
Huang and Ng [23]	0.0033-3673	-32.7	80.4	38.8	2.34-3673	-5.0	82.9	8.6	2.34-3673	-2.5	69.0	13.3
Rogers et al. [24]	0.00069-5.268	-25.7	0.1	32.2	0.2-5.268	-12.1	0.2	13.0	0.2-5.268	-6.8	0.2	12.3
MacGregor and Mather [14]	0.52-1600	48.4	19.3	48.7	0.52-1600	36.6	11.5	37.7	0.52-1600	49.0	15.8	49.2
Jou et al. [21]	0.00183-313	-8.1	5.9	23.6	0.295-313	-2.0	3.6	13.0	0.295-313	6.1	5.9	13.4
Jou et al. [12]	0.0013-5890	18.2	125.6	29.7	0.0273-5890	9.5	109.7	17.6	0.0273-5890	13.9	105.7	20.6
Zoghi and Shokouhi [22]	28-1361	24.2	42.4	26.2	14-1361	14.4	32.1	16.4	14-1361	20.4	36.5	22.2
Maddox et al. [27]	13.23-1536.6	-4.9	74.8	15.2	13.23-1536.6	-8.5	78.0	13.4	13.23-1536.6	-6.4	77.1	13.9
Huttenhuis et al. [25]				-				-	0.141-1.495	-35.0	0.3	35.0
Dicko et al. [29]				-				-	3-278	-17.0	11.6	17.0
Sadegh et al. [30]				-				-	53-386	-10.9	11.4	10.9
This work				-				-	2.60-818.12	-20.2	36.5	21.3
Overall		0.9	70.4	30.4		3.5	66.2	17.8		5.1	58.4	20.6
					Total	pressur	re, $P_{\rm tot}$					
Kuranov et al. [17]	165.2-4895.9	-9.6	240.3	12.5	165.2-4895.9	-10.5	241.1	13.6	165.2-4895.9	-10.0	241.5	12.9
Kamps et al. [18]	147.9-2783	-15.2	213.7	16.0	147.9-2783	-20.8	231.5	20.8	147.9-2783	-13.7	173.5	14.2
Sidi-Boumedine et al. [19]	6.21-1040	-10.8	55.4	12.6	6.21-1040	-16.4	70.3	16.7	6.21-1040	-9.9	56.2	11.0
Overall		-11.0	194.4	13.2		-13.9	202.9	15.8		-10.8	187.9	12.7

**Table 13:** BIAS<sup>a</sup>, AADs<sup>b</sup> and AARDs<sup>c</sup> for the fitted total pressures,  $P_{\text{tot}}$ , and H<sub>2</sub>S partial pressures,  $P_{\text{H2S}}$ , for Cases A, B and C.

639 
$${}^{a}\text{BIAS}(\%) = \frac{1}{N} \sum \frac{P_{s}^{pred} - P_{s}^{exp}}{P_{s}^{exp}} \cdot 100, \ {}^{b}\text{AAD} = \frac{1}{N} \sum \left|P_{s}^{pred} - P_{s}^{exp}\right|, \ {}^{c}\text{AARD}(\%) = \frac{1}{N} \sum \left|\frac{P_{s}^{pred} - P_{s}^{exp}}{P_{s}^{exp}}\right| \cdot 100$$

The model parameters obtained from the data regression in each case studied are given in 641 Supplementary Information. Figure 11 shows experimental and modelled values for a 70 wt.% 642 aqueous MDEA system as a function of temperature in Case C while Table 13 contains 643 information about each regression in terms of Bias, AADs and AARDs. The performance of 644 645 the model for a 70 wt% MDEA-H<sub>2</sub>O system is good, especially considering the few data available for this solvent concentration. In Table 13, it can be seen that the accuracy of the 646 model does not significantly change upon the addition of the experimental points with methane 647 in the regression. The overall AARD for the partial pressure is altered from 18% to 21%, which 648 649 is also the AARD calculated for the data published in this work. The data from MacGregor and 650 Mather [14] exhibit again the highest deviations while the measurements reported by Sadegh et 651 al. [30] in total pressure of 1500 kPa with methane as makeup gas show the lowest deviation, 652 11%. The slight deterioration of the fitting for the equilibrium  $H_2S$  can be also attributed to the 653 fact that experimental points obtained for high amine concentrations are used, i.e. 70 wt.% in this work, but it can also be the result of the sensitivity of the algorithm to the numerical method. 654 655 To illustrate the latter, we repeated the data regression for Case A. The resulted AARDs were 29.8% and 30.1%, using the exact same data and fixed parameters. As far as the ability of the 656 657 model to predict the total pressure is concerned, the accuracy has surprisingly improved. This is merely a lucky coincidence due to the fitting of the experimental points for methane-included 658 systems. 659

Speciation information is necessary in the development of process models for the accurate 660 661 design and operation of gas processing plants. Speciation results, calculated with the model presented in this work, are provided in Figure 12 where mole fractions of all the species in the 662 663 liquid phase are plotted against liquid loading for 70 wt.% aqueous MDEA at 353 K. It is shown that as the loading increases, the concentration of MDEA declines and the concentration of 664 665 protonated amine MDEAH<sup>+</sup> increases. At loadings close to 1, most of the amine has been protonated and the mole fractions of MDEA and H<sub>2</sub>S are equal. The curves representing 666 667 MDEAH<sup>+</sup> and HS<sup>-</sup> overlap, a behavior expected since the formation of sulfide was not taken into account due to its low concentration, therefore the amount of HS<sup>-</sup> and MDEAH<sup>+</sup> formed 668 are balanced. H<sub>3</sub>O<sup>+</sup> and OH<sup>-</sup> also overlap and they are practically zero throughout the loading 669 range. No experimental data were found for the speciation distribution in the H<sub>2</sub>S-MDEA-H<sub>2</sub>O 670 system to confirm the model predictions. Speciation graphs for 50.1 wt.% and 70 wt.% MDEA-671 H<sub>2</sub>O at the temperatures studied in this work are provided in Supplementary Information. 672



674

**Figure 12:** Predicted speciation of  $H_2S$ , MDEA and  $H_2O$  in 70 wt.% MDEA- $H_2O$  at 353 K. (-) H<sub>2</sub>S, (-) MDEA, (-) H<sub>2</sub>O, (-) MDEAH<sup>+</sup>, (- -) OH<sup>-</sup>, (- -) HS<sup>-</sup>.

Overall, although the model developed in this work contains MDEA,  $H_2O$ ,  $H_2S$  and the relevant ionic species, it can predict vapor-liquid equilibria for systems containing methane at low total pressures with similar accuracy as the systems in the absence of methane. However, it is recommended to be used only for rough estimations for  $H_2S$ -CH<sub>4</sub>-MDEA-H<sub>2</sub>O system and a model taking into account the methane solubility to be used if available. This model should not be used for systems with total pressure higher than 2000 kPa, where the gas fugacities change substantially.

#### 685 6. Conclusions

686 Experimental vapor-liquid equilibrium data were measured for a 50.1 wt.% aqueous MDEA at temperature of 323 K and pressure up to 3000 kPa as well as a 70 wt.% aqueous MDEA at 687 688 temperature of 283 K, 353 K and 323 K and pressures up to 10000 kPa, due to their relevance for subsea H<sub>2</sub>S removal of natural gas. Therefore, methane was used as makeup gas. The 689 690 experimental data indicate that the effect of total pressure on the liquid loading of the solvent is within the experimental uncertainties, while for the 50 wt.% MDEA-H<sub>2</sub>O system the impact 691 692 on the partial pressure of hydrogen sulfide is attributed to the non-idealities of the vapor phase 693 and it is lower with decreasing total pressure. The system H<sub>2</sub>S-MDEA-H<sub>2</sub>O up to 70 wt.%

MDEA was modeled employing Peng-Robinson EoS to describe the vapor phase and eNRTL 694 activity coefficient model for the liquid phase. The AARD for the partial pressure of H<sub>2</sub>S and 695 for the total system pressure was found to be 18% and 16% respectively. The effect of including 696 data in the presence of methane and maximum total pressure of 2000 kPa in the data regression 697 was studied and found minimal. However, for higher total pressure and different conditions 698 than the studied ones, the use of models taking into account the methane presence was 699 suggested. Last but not least, new parameters for Antoine correlation were proposed for the 700 estimation of the vapor pressure of MDEA based on our new measurements and all available 701 702 literature data covering a wide temperature range.

703

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709

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715

#### 716 List of Symbols

717	а	parameter of Eq. 10
718	$A_{\Phi}$	Debye-Hückel parameter
719	b	parameter of Eq. 10
720	c-a	cation-anion
721	D	Dielectric constant (-)
722	g <sup>ex</sup>	Molar excess Gibbs energy (J/mol)
723	G	eNRTL auxiliary function (-)
724	Н	Henry's constant (kPa m <sup>3</sup> /kmol)

725	I <sub>x</sub>	Ionic strength in mole fraction scale (mol/m <sup>3</sup> )
726	k	Boltzmann constant (J/K)
727	т	molecule
728	М	Molecular weight (kg/kmol)
729	NA	Avogadro number (mol <sup>-1</sup> )
730	Р	Pressure (kPa)
731	<i>r</i> <sub>BORN</sub>	Born radius (m)
732	R	Gas constant (J mol <sup>-1</sup> K <sup>-1</sup> )
733	Т	Absolute temperature (K)
734	v	Molar volume (m <sup>3</sup> /mol)
735	W	Weight fraction (-)
736	x	Mole fraction (-)
737	X	eNRTL mole fraction (-)
738	Z.	Ionic charge (-)
739	Ζ	Absolute value of the ionic charge (-)
740		
741	Greek letters	
742	α	Loading (mol H <sub>2</sub> S/mol MDEA)
743	γ	Activity coefficient (-)
744	3	Permittivity (F/m)
745	ρ	Molar density (mol/cm <sup>3</sup> )
746	$ ho_{ m pdh}$	Closest approach parameter of the Pitzer-Debye-Hückel formulation (-)
747	τ	Energy parameter (-)
748		
749	Superscripts	
750	E	Excess property
751	exp	Experimental value
752	1	Liquid phase
753	lc	Local composition

754	pdh	Pitzer-Debye-Hückel formulation
755	phys	Physical absorption
756	pred	Predicted value
757	v	Vapor phase
758		
759	Subscripts	
760	amb	Ambient
761	app	Apparent
762	aq, MDEA	Aqueous MDEA
763	c	Critical
764	glob	Global, refers to global loading $\alpha_{glob}$ (mol H <sub>2</sub> S in the cell/mol MDEA)
765	i, j, k	Component in a mixture
766	ij	Cross parameter
767	liq	Liquid, refers to liquid loading $\alpha_{liq}$ (mol H <sub>2</sub> S/mol MDEA)
768	S	solvent
769	W	water
770		
771	Abbreviation	IS
772	AAD	Average Absolute Deviation
773	AARD	Average Absolute Relative Deviation
774	eNRTL	electrolyte Non-Random Two Liquids
775	EoS	Equation of State
776	FPD	Freezing Point Depression
777	FTIR	Fourier-Transform infrared
778	GC	Gas Chromatography
779	MDEA	Methyldiethanolamine
780	NP	Number of data points
781	NS	Number of vapor phase samples for GC analysis
782	VLE	Vapor-Liquid Equilibrium

## 784 Appendix A: eNRTL model

The activity coefficients were calculated by the electrolyte Non-Random Two Liquids (eNRTL)model [50]:

The starting point for the description of the liquid phase is the expression of excess Gibbs energy as the sum of two terms; one related to the long-range forces between the ions (first term) and one to the short-range forces between all the species (second term):

$$\frac{g^E}{RT} = \frac{g^{E,pdh}}{RT} + \frac{g^{E,lc}}{RT}$$
Eq. A1

790 This equation lead to:

$$\ln \gamma_i = \ln \gamma_i^{pdh} + \ln \gamma_i^{lc}$$
 Eq. A2

The subscript *pdh* denotes Pitzer-Debye-Hückel formulation for the long-range interactions and
the subscript *lc* denotes Local Composition model. The formulation of the former is:

$$\frac{g^{E,pdh}}{RT} = -\left(\sum_{k} x_{k}\right) \left(\frac{1000}{M_{s}}\right)^{\frac{1}{2}} \left(\frac{4A_{\phi}I_{x}}{\rho_{pdh}}\right) \ln(1+\rho_{pdh}I_{x}^{\frac{1}{2}}) \qquad \text{Eq. A3}$$

793 By derivation, the activity coefficient is expressed according to:

$$\ln \gamma_i^{*pdh} = -(1000/M_s)^{\frac{1}{2}} A_{\phi} \left[ \left( 2 \frac{Z_i^2}{\rho_{pdh}} \right) \ln \left( 1 + \rho_{pdh} I_x^{\frac{1}{2}} \right) + \left( Z_i^2 I_x^{\frac{1}{2}} - 2 I_x^{\frac{3}{2}} \right) / (1 + \rho_{pdh} I_x^{\frac{1}{2}}) \right]$$
Eq. A4

where  $I_x$  is the ionic strength and  $A_{\phi}$  is the Debye-Hückel parameter, expressed as following:

$$I_x = \frac{1}{2} \sum_i Z_i^2 x_i, \qquad A_{\phi} = \frac{1}{3} \left(\frac{2\pi N_o d}{1000}\right)^{\frac{1}{2}} \left(\frac{e^2}{DkT}\right)^{3/2}$$
Eq. A5

The reference state for the *pdh* term is infinite dilution in the mixed solvent while the reference state for the *lc* term is infinite dilution in water. To account for the excess Gibbs energy of transfer from the infinite dilution in the mixed solvent to the infinite dilution in water, a term is added in the long-range interaction expression. This additional term is described by the Born equation:

$$\frac{g^{E,BORN}}{RT} = -\frac{N_A e^2}{8\pi RT r_{BORN} \varepsilon_o} \left(\frac{1}{\varepsilon_s} - \frac{1}{\varepsilon_w}\right) \sum_i x_i z_i^2$$
Eq. A6

800 Further,

$$\ln \gamma_i^{BORN} = \frac{N_A e^2}{8\pi R T r_{BORN} \varepsilon_o} \left(\frac{1}{\varepsilon_s} - \frac{1}{\varepsilon_w}\right) z_i^2$$
 Eq. A7

801

#### 802 The short-range contribution is described by the eNRTL model as following:

$$\frac{g^{E,lc}}{RT} = \sum_{m} X_m \frac{\sum_j X_j G_{jm} \tau_{jm}}{\sum_k X_k G_{km}} + \sum_c X_c \sum_{a'} \frac{X_{a'}}{\sum_k X_k G_{km}} \frac{\sum_j X_j G_{jc,a'c} \tau_{jc,a'c}}{\sum_k X_k G_{kc,a'c}} + \sum_a X_a \sum_{c'} \frac{X_{c'}}{\sum_{c''} X_{c''}} \frac{\sum_j X_j G_{ja,c'a} \tau_{ja,c'a}}{\sum_k X_k G_{ka,c'a}}$$
Eq. A8

where *m*, *c* and *a* denote molecule, cation and anion, while  $X_j = C_j x_j$  with j: *m*, *c*, *a* effective local mole fraction. Equations A9 and A10 are given using the ion-like repulsion assumption and the local electroneutrality assumption.

$$G_{cm} = \frac{\sum_{a} X_a G_{ca,m}}{\sum_{a'} X_{a'}}$$
Eq. A9

$$G_{am} = \frac{\sum_{c} X_{c} G_{ca,m}}{\sum_{c'} X_{c'}}$$
Eq. A10

806 The  $G_{i,j}$  and  $\tau_{i,j}$  parameters are related through the non-randomness parameter,  $a_{i,j}$ :

$$G_{ij} = \exp(-a_{i,j}\tau_{i,j})$$
Eq. A11

807 The equations presented below describe the non-randomness parameters:

$$a_{cm} = \frac{\sum_{a} X_{a} a_{ca,m}}{\sum_{a'} X_{a'}}$$
 Eq. A12

$$a_{am} = \frac{\sum_{c} X_{c} a_{ca,m}}{\sum_{c'} X_{c'}}$$
Eq. A13

808 The energy parameters  $\tau_{mc,ac}$  and  $\tau_{ma,ca}$  are given by:

$$\tau_{mc,ac} = \tau_{cm} - \frac{a_{ca,m}}{a_{mc,ac}} \left( \tau_{ca,m} - \tau_{m,ca} \right)$$
Eq. A14

$$\tau_{ma,ca} = \tau_{am} - \frac{a_{ca,m}}{a_{ma,ac}} \left( \tau_{ca,m} - \tau_{m,ca} \right)$$
Eq. A15

809 where

$$a_{mc,ac} = a_{cm}$$
 Eq. A16

$$a_{am,ac} = a_{am}$$
 Eq. A17

811 The adjustable binary parameters are the non-randomness factors  $a_{ca,m}$ ,  $a_{ca,ca'}$ ,  $a_{ca,c'a}$ ,  $a_{mm'}$ , 812 and the energy parameters  $\tau_{ca,m}$ ,  $\tau_{m,ca}$ ,  $\tau_{ca,ca'}$ ,  $\tau_{ca,ca}$ ,  $\tau_{ca,c'a}$ ,  $\tau_{m,m'}$ ,  $\tau_{m'm}$ .

813

814 From Equation A8, the activity coefficients are calculated:

815 For molecules:

$$\ln \gamma_{m}^{lc} = \frac{\sum_{j} X_{j} G_{jm} \tau_{jm}}{\sum_{k} X_{k} G_{km}} + \sum_{m'} \frac{X_{m'} G_{mm'}}{\sum_{k} X_{k} G_{km'}} \left( \tau_{mm'} - \frac{\sum_{k} X_{k} G_{km'} \tau_{km'}}{\sum_{k} X_{k} G_{km'}} \right) \\ + \sum_{c} \sum_{a'} \frac{X_{a'}}{\sum_{a''} X_{a''}} \frac{X_{c'} G_{mc,a'c}}{\sum_{k} X_{k} G_{kc,a'c}} \left( \tau_{mc,a'c} - \frac{\sum_{k} X_{k} G_{kc,a'c} \tau_{kc,a'c}}{\sum_{k} X_{k} G_{kc,a'c}} \right) \\ + \sum_{a} \sum_{c'} \frac{X_{c'}}{\sum_{c''} X_{c''}} \frac{X_{a'} G_{ma,c'a}}{\sum_{k} X_{k} G_{ka,c'a}} \left( \tau_{ma,c'a} - \frac{\sum_{k} X_{k} G_{ka,c'a} \tau_{ka,c'a}}{\sum_{k} X_{k} G_{ka,c'a}} \right) \\ \text{Eq. A18}$$

816

817 For cations:

$$\frac{1}{Z_{c}}\ln\gamma_{c}^{lc} = \sum_{a'} \frac{X_{a'}}{\sum_{a''} X_{a''}} \frac{\sum_{k} X_{k} G_{kc,a'c} \tau_{kc,a'c}}{\sum_{k} X_{k} G_{kc,a'c}} \\
+ \sum_{m} \frac{X_{m} G_{cm}}{\sum_{k} X_{k} G_{km}} \left(\tau_{cm} - \frac{\sum_{k} X_{k} G_{km} \tau_{km}}{\sum_{k} X_{k} G_{km}}\right) \\
+ \sum_{a} \sum_{c'} \frac{X_{c'}}{\sum_{c''} X_{c''}} \frac{X_{a'} G_{ca,c'a}}{\sum_{k} X_{k} G_{ka,c'a}} \left(\tau_{ca,c'a} - \frac{\sum_{k} X_{k} G_{ka,c'a} \tau_{ka,c'a}}{\sum_{k} X_{k} G_{ka,c'a}}\right)$$

Eq. A19

818

819 For anions:

$$\frac{1}{Z_a} \ln \gamma_a^{lc} = \sum_{c'} \frac{X_{c'}}{\sum_{c''} X_{c''}} \frac{\sum_k X_k G_{ka,c'a} \tau_{ka,c'a}}{\sum_k X_k G_{ka,c'a}} \\
+ \sum_m \frac{X_m G_{am}}{\sum_k X_k G_{km}} \left( \tau_{am} - \frac{\sum_k X_k G_{km} \tau_{km}}{\sum_k X_k G_{km}} \right) \\
+ \sum_c \sum_{a'} \frac{X_{a'}}{\sum_{a''} X_{a''}} \frac{X_c G_{ac,a'c}}{\sum_k X_k G_{kc,a'c}} \left( \tau_{ac,a'c} - \frac{\sum_k X_k G_{kc,a'c} \tau_{kc,a'c}}{\sum_k X_k G_{kc,a'c}} \right) \\$$
Eq. A20

821 The expressions of activity coefficients at infinite dilution are then:

$$\ln \gamma_m^{lc,\infty} = \tau_{wm} + G_{mw} \tau_{mw}$$
 Eq. A21

$$\frac{1}{Z_c} \ln \gamma_c^{lc,\infty} = \frac{X_a}{\sum_{a'} X_{a'}} \tau_{wc,ac} + G_{cw} \tau_{cw}$$
Eq. A22

$$\frac{1}{Z_a} \ln \gamma_a^{lc,\infty} = \frac{X_c}{\sum_{c'} X_{c'}} \tau_{wa,ca} + G_{aw} \tau_{aw}$$
 Eq. A23

By combination of Equations A2, A4, A7, A18 and A23, the activity coefficient for the liquidphase is found by:

$$\gamma_{i} = \gamma_{i}^{pdh} \gamma_{i}^{BORN} \gamma_{i}^{lc} / \gamma_{i}^{lc,\infty}$$
 Eq. A24

where i = m, *c* or *a* for all components, besides the amine in this work. For MDEA, the symmetric reference state for the short-range interactions contribution,  $\gamma_{MDEA}^{lc,\infty}$  is fixed to 1.

## 826 Appendix B: Peng-Robinson Equation of State

827 The fugacity coefficients were calculated by Peng-Robinson equation of state [49]:

$$P = \frac{RT}{v - b} - \frac{\alpha a(T)}{v^2 + 2bv - b^2}$$
 Eq. A25

828 where:

$$a = 0.42724 \frac{R^2 T_c^2}{P_c}, \qquad b = 0.07780 \frac{RT_c}{P_c}, \quad a(T) = \left[1 + m\left(1 - T_R^{\frac{1}{2}}\right)\right]^2 \&$$
$$m = 0.37464 + 1.54226\omega - 2.26992 \:\omega^2$$

830 The traditional van der Waals one-fluid mixing rules were used for the estimation of the gas831 mixture parameters from the pure components' properties.

$$a a(T) = \sum_{i} \sum_{j} x_{i} x_{j} (aa(T))_{i}^{\frac{1}{2}} (aa(T))_{j}^{\frac{1}{2}} (1 - k_{ij}) \qquad \& \qquad b = \sum_{i} x_{i} b_{i}$$

832 In our work, the binary interaction parameter  $k_{ij}$  is set to zero, so as the eNRTL model 833 parameters are the only ones fitted.

834 The critical properties used in this work can be found in Supporting Information.

#### 835 Appendix C: Supplementary Information

- 836 Supplementary information includes:
- 837 A. Critical properties and acentric factors for pure components
- 838 B. Model parametrization
- 839 C. Modeling results
- 840 D. Experimental results
- 841 E. Uncertainty analysis

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