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BENJAMINI-SCHRAMM CONVERGENCE AND SPECTRUM OF RANDOM HYPERBOLIC SURFACES OF HIGH GENUS

LAURA MONK

ABSTRACT. We study geometric and spectral properties of *typical* hyperbolic surfaces of high genus, excluding a set of small measure for the Weil-Petersson probability. We first prove a quantitative estimate of the speed of Benjamini-Schramm convergence to the hyperbolic plane as the genus g goes to infinity. We then use the Selberg pretrace formula to find an estimate for the number of eigenvalues in an interval $[a, b]$, as b and/or g go to infinity. This result proves that, as $g \rightarrow +\infty$, the spectrum approaches the continuous spectrum of the Laplacian on the hyperbolic plane, with the density appearing in the Selberg trace formula. It also leads to a probabilistic uniform Weyl law, when $b \rightarrow +\infty$. We also prove a probabilistic upper bound on the number of eigenvalues in the interval $[a, b]$, and noticeably that the number of eigenvalues below $\frac{1}{4}$ is $\mathcal{O}\left(g(\log g)^{-\frac{3}{4}}\right)$.

1. INTRODUCTION AND MAIN RESULTS

Let X be a compact (oriented, connected, without boundary) hyperbolic surface. It is isometric to a quotient \mathcal{H}/Γ , where $\mathcal{H} = \{x + iy, y > 0\}$ is the upper half-plane equipped with the metric $ds^2 = \frac{dx^2 + dy^2}{y^2}$, and $\Gamma \subset \mathrm{PSL}_2(\mathbb{R})$ is a co-compact Fuchsian group. The metric induces a hyperbolic distance $d_{\mathcal{H}}$ on \mathcal{H} and d_X on X . The hyperbolic structure on X induces a volume form μ_X on X and $\mu_{\mathcal{H}} = \frac{dx dy}{y^2}$ on \mathcal{H} . The total volume of X is $2\pi(2g - 2)$, where g is the genus of X . Therefore, studying surfaces of high genus is equivalent to studying surfaces of large volume.

In the following, the notation $T_1 = \mathcal{O}(T_2)$ means that there exists a universal constant $C > 0$ such that $|T_1| \leq C T_2$ for any choice of parameters. If the constant depends on some parameter x , then we will write $T_1 = \mathcal{O}_x(T_2)$.

1.1. Spectrum of the Laplacian on a hyperbolic surface. Let Δ_X be the (positive) Laplace-Beltrami operator on $L^2(X)$, and let $(\lambda_j)_{j \geq 0}$ be its non-decreasing sequence of eigenvalues (with multiplicities). For any real numbers $0 \leq a \leq b$, let $N_X^\Delta(a, b)$ be the number of eigenvalues of Δ_X in the interval $[a, b]$. This article aims at understanding the behavior of $N_X^\Delta(a, b)$ as the genus g approaches infinity. To put this paper in context, here are the known results about these counting functions.

- $N_X^\Delta(0, \frac{1}{4}) \leq 2g - 2$ [19], and this bound is optimal [20, 9].
- The examples from [9] also prove that, for any $\varepsilon \in (0, \frac{1}{4})$, there are compact hyperbolic surfaces such that $N_X^\Delta(0, \varepsilon) = 2g - 2$.
- On the contrary, for any $\varepsilon > 0$, there cannot be a topological bound on $N_X^\Delta(0, \frac{1}{4} + \varepsilon)$. Indeed, there exist compact hyperbolic surfaces of a given genus $g \geq 2$ with an arbitrarily large number of eigenvalues below $\frac{1}{4} + \varepsilon$ [9].

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- The Weyl law gives the asymptotic behavior of $N_X^\Delta(0, b)$ for a *fixed* surface X as b goes to infinity. In our setting, the best known estimate is the following [4, 21]:

$$\frac{N_X^\Delta(0, b)}{\mu_X(X)} = \frac{b}{4\pi} + \mathcal{O}_X\left(\frac{\sqrt{b}}{\log b}\right)$$

where the implied constant depends on the surface X .

1.2. Random compact hyperbolic surfaces of high genus. The question one may now ask is: what are the spectral properties of a *typical* compact hyperbolic surface? Can we improve the previous results if we allow ourselves to exclude a *small* set of surfaces?

Our approach to this problem involves working with *random* compact hyperbolic surfaces. There are several ways to do this. We chose here to work with the Weil-Petersson volume, following Mirzakhani's approach [18], but there is also, for instance, a construction of Brooks and Makover [8], in which similar work can probably be done.

In the following, our probability space will be the *moduli space* \mathcal{M}_g of compact hyperbolic surfaces of genus g . It is the set of all compact hyperbolic surfaces of genus g , up to isometry. The moduli space is equipped with a natural symplectic form ω_{wp} called the *Weil-Petersson form* [23]. This induces a volume form $\text{Vol}_{\text{wp}}^{(g)} = \frac{\omega_{\text{wp}}^{\wedge(3g-3)}}{(3g-3)!}$ on \mathcal{M}_g , which is of finite volume $V_g = \text{Vol}_{\text{wp}}^{(g)}(\mathcal{M}_g)$. Therefore, it can be renormalized to obtain a probability measure $\mathbb{P}_{\text{wp}}^{(g)} = \frac{1}{V_g} \text{Vol}_{\text{wp}}^{(g)}$.

We say that an event $\mathcal{A}_g \subset \mathcal{M}_g$ occurs *with high probability* when $\mathbb{P}_{\text{wp}}^{(g)}(\mathcal{A}_g) \rightarrow 1$ as $g \rightarrow +\infty$. We will prove results true in that sense, and hence not for all surfaces but most. This probabilistic approach has proven to be a good method in the study of graphs [14]. A significant example of a spectral result true with high probability is Alon Friedman's theorem [15, 6], for random large *regular* graphs. These are known to have similar geometric and spectral properties to the ones of compact hyperbolic surfaces of high genus.

1.3. Geometry of random surfaces. Multiple aspects of the geometry of random surfaces of high genus have already been studied. In [18], Mirzakhani estimated various quantities (injectivity radius, Cheeger constant, diameter ...) with high probability.

Radius of injectivity. The *radius of injectivity* $\text{InjRad}_z(X)$ at a point z of a compact hyperbolic surface X is the supremum of all real numbers $r \geq 0$ such that the ball of radius r centered at z in X is isometric to a ball in the hyperbolic plane \mathcal{H} . The *global radius of injectivity* $\text{InjRad}(X)$ is the infimum over X of the radius of injectivity, or equivalently twice the length of the shortest closed geodesic on X .

In the following, we will work under the assumption that the random surfaces we are considering are a *uniformly discrete* family, that is to say that their radius of injectivity is bounded below by a constant $r(g)$. In order to guarantee this, we will have to exclude some surfaces. We will control the probability measure of the set of excluded surfaces using the following result.

Theorem 1 (Theorem 4.2 in [18]). *There exists a constant $C > 0$ such that, for any large enough g and any small enough $r > 0$,*

$$(1) \quad \frac{1}{C} r^2 \leq \mathbb{P}_{\text{wp}}^{(g)}(\text{InjRad}(X) \leq r) \leq C r^2.$$

The lower bound of this statement motivates the fact that our uniform discreteness bound $r(g)$ will need to go to zero as g approaches infinity in order for the event to occur with high probability.

Benjamini-Schramm convergence. The notion of Benjamini-Schramm convergence has first been introduced by Benjamini and Schramm in the context of sequences of graphs [3], but can naturally be extended to a continuous setting (see [1, 2, 7]). There is a general definition of Benjamini-Schramm convergence for a *deterministic* sequence of hyperbolic surfaces $(X_g)_g$. In the special case when the limit of $(X_g)_g$ is the hyperbolic plane \mathcal{H} , it is equivalent to the following simpler property:

$$(2) \quad \forall L > 0, \lim_{g \rightarrow +\infty} \frac{\mu_{X_g}(\{z \in X_g : \text{InjRad}_z(X_g) < L\})}{\mu_{X_g}(X_g)} = 0$$

which we will therefore use as a definition. The idea behind this characterization is the following. We consider a distance $L > 0$. On the (fixed) surface X_g , one can pick a point z at random, using the normalized measure $\frac{1}{\mu_{X_g}(X_g)} \mu_{X_g}$. Equation (2) means that the probability for the ball of center z and radius L to be isometric to a ball in the hyperbolic plane goes to one as $g \rightarrow +\infty$.

The case where $\text{InjRad}(X_g) \rightarrow +\infty$ as $g \rightarrow +\infty$ is a simple situation which implies Benjamini-Schramm convergence towards \mathcal{H} . However, Theorem 1 proves that this does not occur with high probability for our probabilistic model.

Our first result is a quantitative estimate of the Benjamini-Schramm speed of convergence of random hyperbolic surfaces of high genus towards \mathcal{H} .

Theorem 2. *For any $g \geq 2$ and any $L, M > 0$, there exists a set $\mathcal{A}_{g,L,M} \subset \mathcal{M}_g$ such that for any hyperbolic surface $X \in \mathcal{A}_{g,L,M}$,*

$$(3) \quad \mu_X(\{z \in X : \text{InjRad}_z(X) < L\}) \leq e^L M,$$

$$\text{and } 1 - \mathbb{P}_{\text{wp}}^{(g)}(\mathcal{A}_{g,L,M}) = \mathcal{O}\left(\frac{Le^{2L}}{M}\right).$$

We recall that the implied constant is independent of the genus g and the parameters L, M . Since the total area of a compact hyperbolic surface of genus g is $2\pi(2g - 2)$, this result will only be interesting for $L \leq \log g$. Specifying the parameters to be $L = \frac{1}{6} \log g$, $M = g^{\frac{1}{2}}$ and $r_g = g^{-\frac{1}{24}}(\log g)^{\frac{9}{16}}$, Theorem 1 and 2 together lead to the following Corollary.

Corollary 3 (geometric assumptions). *For large enough g , there exists a subset $\mathcal{A}_g \subset \mathcal{M}_g$ such that, for any hyperbolic surface $X \in \mathcal{A}_g$,*

$$(4) \quad \text{InjRad}(X) \geq g^{-\frac{1}{24}}(\log g)^{\frac{9}{16}}$$

$$(5) \quad \frac{\mu_X(\{z \in X : \text{InjRad}_z(X) < \frac{1}{6} \log g\})}{\mu_X(X)} = \mathcal{O}\left(g^{-\frac{1}{3}}\right)$$

$$\text{and } 1 - \mathbb{P}_{\text{wp}}^{(g)}(\mathcal{A}_g) = \mathcal{O}\left(g^{-\frac{1}{12}}(\log g)^{\frac{9}{8}}\right).$$

The estimate (5) is similar to the one proved in Section 4.4 of [18]. The proof is the same, though an incorrect argument has been modified here.

The rest of this article is devoted to the study of the spectral properties of a given element $X \in \mathcal{A}_g$, where \mathcal{A}_g is the set from Corollary 3. There are no more probabilistic arguments from this point.

1.4. Spectrum of random surfaces. Geometry and spectrum of Riemannian manifolds, and especially of hyperbolic surfaces, are connected in a variety of ways. A straightforward example of this interaction can be found in Cheeger and Buser's inequalities

$$\frac{h^2}{4} \leq \lambda_1 \leq h(1 + 10h),$$

where h is the Cheeger constant of the hyperbolic surface X [12, 10]. Mirzakhani deduced from her probabilistic uniform lower bound on the Cheeger constant that there exists

a constant $c_0 > 0$ such that, with high probability, $\lambda_1 \geq c_0$ [18]. This article aims at providing more information on the distribution of eigenvalues for most random surfaces.

Counting functions. Our main results are the following two estimates on the counting functions $N_X^\Delta(a, b)$, true with high probability.

Theorem 4. *For any large enough g , any $0 \leq a \leq b$ and any $X \in \mathcal{A}_g$ from Corollary 3,*

$$(6) \quad \frac{N_X^\Delta(a, b)}{\mu_X(X)} = \mathcal{O} \left(b - a + \sqrt{\frac{b+1}{\log g}} \right).$$

Whenever $b \leq \frac{1}{4}$, if we set $c = 2^{-15}$, we also have

$$(7) \quad \frac{N_X^\Delta(0, b)}{\mu_X(X)} = \mathcal{O} \left(\frac{g^{-c(\frac{1}{4}-b)^2}}{(\log g)^{\frac{3}{4}}} \right).$$

This Theorem provides us with an upper bound on the number of eigenvalues in an interval $[a, b]$. In equation (6), the term $\sqrt{\frac{b+1}{\log g}}$ corresponds to a minimum scale, below which we can have no additional information by shrinking the spectral window.

The second part of the statement controls the number of small eigenvalues. When we take $b = \frac{1}{4}$, we obtain that $N_X^\Delta(0, \frac{1}{4})$ is $\mathcal{O} \left(g (\log g)^{-\frac{3}{4}} \right)$. This is better than the (optimal) deterministic estimate $N_X^\Delta(0, \frac{1}{4}) \leq 2g - 2$. Furthermore, this bound becomes better the further b is from the bulk spectrum $[\frac{1}{4}, +\infty)$. As a consequence, the examples of surfaces with $2g - 2$ eigenvalues in $[0, \varepsilon]$ (for arbitrarily small ε) are not typical.

The second result is a more precise approximation of the counting functions.

Theorem 5. *There exists a universal constant $C > 0$ such that, for any large enough g , any $0 \leq a \leq b$ and any hyperbolic surface $X \in \mathcal{A}_g$ from Corollary 3, one can write the counting function $N_X^\Delta(a, b)$ as*

$$\frac{N_X^\Delta(a, b)}{\mu_X(X)} = \frac{1}{4\pi} \int_{\frac{1}{4}}^{+\infty} \mathbb{1}_{[a, b]}(\lambda) \tanh \left(\pi \sqrt{\lambda - \frac{1}{4}} \right) d\lambda + R(X, a, b)$$

where

$$-C \sqrt{\frac{b+1}{\log g}} \leq R(X, a, b) \leq C \sqrt{\frac{b+1}{\log g}} \log \left(2 + (b-a) \sqrt{\frac{\log g}{b+1}} \right)^{\frac{1}{2}}.$$

We recall that the sets \mathcal{A}_g from Corollary 3 are fixed subsets of \mathcal{M}_g , of probability going to 1 as $g \rightarrow +\infty$. Therefore, the previous results hold *with high probability*, uniformly with respect to the parameters a and b . The remainder $R(X, a, b)$ is negligible compared to the main term as soon as the size $b - a$ of the spectral window is much larger than the minimal spacing $\sqrt{\frac{b+1}{\log g}}$. There are several limits that can be interesting to study.

- When $[a, b]$ is fixed and $g \rightarrow +\infty$, our result shows that the spectrum of Δ_X approaches the continuous spectrum of the Laplacian on \mathcal{H} . This is the analogous of the fact that, if a sequence of d -regular graphs converges in the sense of Benjamini-Schramm to the d -regular tree, then their spectral measure converges to the Kesten-McKay law.
- By taking $a = 0$ and b going to infinity, one can recover a uniform Weyl law, with remainder of order $\mathcal{O}_g(\sqrt{b \log b})$. The constant is independent of the surface, and explicit in terms of g . We could probably have obtained a better remainder (for instance, $\mathcal{O}_g\left(\frac{\sqrt{b}}{\log b}\right)$ as in [4, 21]) had we allowed the probability set \mathcal{A}_g to depend on the parameter b , which we did not do here in order to make the discussion simpler.

- One can also consider mixed regimes, where both b and g go to infinity.

Our approach in proving the two Theorems is inspired by Part 9 of [17], where Le Masson and Sahlsten prove the convergence of $\frac{N_{X_g}^\Delta(a,b)}{\mu_{X_g}(X_g)}$ to $\frac{1}{4\pi} \int_{\frac{1}{4}}^{+\infty} \mathbf{1}_{[a,b]}(\lambda) \tanh\left(\pi\sqrt{\lambda - \frac{1}{4}}\right) d\lambda$ as $g \rightarrow +\infty$, for a uniformly discrete sequence of compact hyperbolic surfaces $(X_g)_g$ converging to \mathcal{H} in the sense of Benjamini-Schramm. Here, we do not consider a sequence but a fixed surface of high genus. Furthermore, we estimate precisely the error term, which leads us to considering different kernels in the pretrace formula.

Eigenvalue multiplicity. For a compact hyperbolic surface X of genus g , let $m_j(X)$ denote the multiplicity of the eigenvalue λ_j of Δ_X . It is known that $m_j(X) \leq 4g + 2j + 1$ [5].

We can estimate $m_j(X)$ using Theorem 4 provided we have information on the size of λ_j . This can be achieved using a result from Cheng [13]:

$$0 \leq \lambda_j \leq \frac{1}{4} + \frac{16\pi^2 j^2}{\text{diam}(X)^2}$$

where $\text{diam}(X)$ is the diameter of X . Mirzakhani [18] proved that, with high probability, the diameter of a hyperbolic surface is smaller than $40 \log(g)$. We deduce from these observations and Theorem 4 that with high probability, for any integer j ,

$$\frac{m_j(X)}{g} = \mathcal{O}\left(\frac{1}{\sqrt{\log g}} + \frac{j^2}{(\log g)^2}\right).$$

A better probabilistic upper bound $\frac{m_j(X)}{g} = \mathcal{O}\left(\frac{1}{\log g}\right)$ has been proved recently in [16], by estimating the L^p -norms of eigenfunctions on random hyperbolic surfaces of high genus. In the case when $\lambda_j \leq \frac{1}{4} - \varepsilon$, the authors also prove that $\frac{m_j(X)}{g} = \mathcal{O}\left(g^{-c\sqrt{\varepsilon}}\right)$, which is weaker than our bound (7).

Perspectives. Inspired by the similar geometric and spectral properties of large regular graphs and large genus hyperbolic surfaces, Wright conjectured [24] that Alon-Friedman's theorem holds for random hyperbolic surfaces taken with the Weil-Petersson probability measure.

Conjecture 6. For any $\varepsilon > 0$,

$$\lim_{g \rightarrow +\infty} \mathbb{P}_{\text{WP}}^{(g)}\left(N_X^\Delta\left(0, \frac{1}{4} - \varepsilon\right) = 1\right) = 1.$$

In other words, for any fixed $\varepsilon > 0$, with high probability, there is no non-trivial eigenvalue smaller than $\frac{1}{4} - \varepsilon$.

It might be possible to prove this result using the Selberg trace formula with an adequate test function. However, obtaining an estimate as precise as $N_X^\Delta(0, \frac{1}{4} - \varepsilon) = 1$ would be highly technical, and cannot be achieved using the approach presented in this article.

Organization of the paper. We start by proving Benjamini-Schramm convergence for the Weil-Petersson probabilistic model in Section 2. Then, the proof of Theorems 4 and 5 spans over Section 3 and 4, which corresponds respectively to the case where $0 \leq a \leq b \leq 1$ (bottom of the spectrum) and $\frac{1}{2} \leq a \leq b$ (away from small eigenvalues). The way the different parts depend on one another is explained in Figure 1 for clarity.

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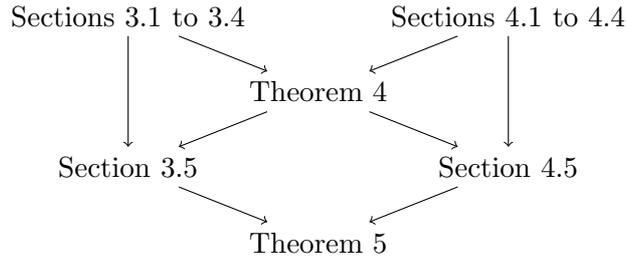


FIGURE 1. The steps of the proofs of Theorems 4 and 5, and the way they depend on one another. The left part corresponds to the case $b \leq 1$, and the right part to the case $a \geq \frac{1}{2}$.

2. BENJAMINI-SCHRAMM CONVERGENCE OF RANDOM SURFACES

We here proceed to the proof of Theorem 2, which states that, with high probability, random surfaces are close to the hyperbolic plane in the Benjamini-Schramm sense. It is a generalization of a result proved in Section 4.4 of [18].

Proof of Theorem 2. Let X be a compact hyperbolic surface of genus g . In order to estimate the volume of

$$X^-(L) = \{z \in X : \text{InjRad}_z(X) < L\},$$

we establish a link between this volume and the number of small geodesics on X . In order to do so, we will work in the universal cover \mathcal{H} of $X = \mathcal{H}/\Gamma$, where we make the geometric construction represented in Figure 2.

Let z be a point in X of radius of injectivity $r < L$. There is a simple geodesic arc c in X starting at z of length $2r$. It corresponds to a geodesic segment \tilde{c} in \mathcal{H} of length $2r$ going from one lift \tilde{z} of z to another, $T\tilde{z}$, for a $T \in \Gamma$. Let $\tilde{\gamma}$ be the axis of the hyperbolic transformation T on \mathcal{H} ; it corresponds to a simple closed geodesic γ on X , of length $\ell \leq 2r$.

Let \tilde{w} be the orthogonal projection of \tilde{z} on $\tilde{\gamma}$. Let \tilde{c}_\perp be the common perpendicular segment of $\tilde{\gamma}$ and \tilde{c} . Then the two trirectangles seen in Figure 2 are isometric. We apply the trigonometry formula 2.3.1(v) of [11] to one of them, and get:

$$\sinh(r) = \cosh(d_{\mathcal{H}}(\tilde{z}, \tilde{\gamma})) \sinh\left(\frac{\ell}{2}\right).$$

As a consequence,

$$\cosh(d_X(z, \gamma)) \leq \frac{\sinh(r)}{\sinh\left(\frac{\ell}{2}\right)} \leq \frac{e^L}{2 \sinh\left(\frac{\ell}{2}\right)} =: \cosh(d_{\max}(\gamma, L)).$$

Then z belongs to the collar of width $d_{\max}(\gamma, L)$ around the closed geodesic γ in X . The volume of this collar is less than the volume of the corresponding collar in the cylinder $\mathcal{H}/\langle T \rangle$, for which γ is the central geodesic. This volume can be computed using the Fermi coordinates:

$$\int_0^\ell \int_{-d_{\max}(\gamma, L)}^{d_{\max}(\gamma, L)} \cosh(\rho) \, d\rho \, dt = 2\ell \sinh(d_{\max}(\gamma, L)) \leq 2\ell \cosh(d_{\max}(\gamma, L)) = \frac{\ell e^L}{\sinh\left(\frac{\ell}{2}\right)}.$$

Since $x \leq \sinh x$, the volume of the collar of width $d_{\max}(\gamma, L)$ around γ in X is smaller than $2e^L$. As a consequence, any point $z \in X^-(L)$ is in a collar of volume less than $2e^L$ around a simple closed geodesic of length at most $2r \leq 2L$. This implies:

$$\mu_X(X^-(L)) \leq 2e^L N_X^\ell(2L),$$

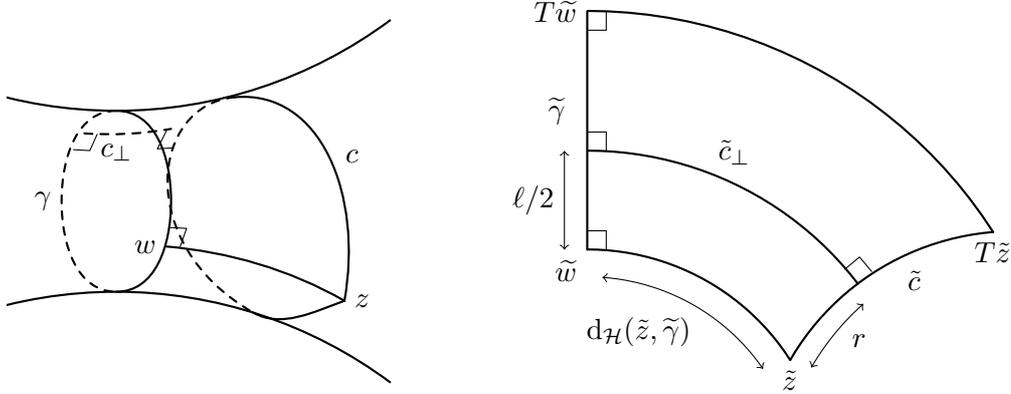
(A) On the surface X (B) On the universal cover \mathcal{H}

FIGURE 2. Illustration of the geometric construction in the proof of Theorem 2.

where, for any positive L , $N_X^\ell(L)$ is the number of simple closed geodesics with length at most L on the hyperbolic surface X . Then, on the set

$$\mathcal{A}_{g,L,M} = \left\{ X \in \mathcal{M}_g : N_X^\ell(2L) \leq \frac{M}{2} \right\}$$

(3) is proved. We estimate the Weil-Petersson probability of this event. In order to do so, let us write

$$\begin{aligned} \mathbb{P}_{\text{wp}}^{(g)}(\mathcal{M}_g \setminus \mathcal{A}_{g,L,M}) &= \frac{1}{V_g} \int_{\mathcal{M}_g} \mathbf{1}_{\{N_X^\ell(2L) > \frac{M}{2}\}} d \text{Vol}_{\text{wp}}^{(g)}(X) \\ &\leq \frac{2}{M} \frac{1}{V_g} \int_{\mathcal{M}_g} N_X^\ell(2L) d \text{Vol}_{\text{wp}}^{(g)}(X) = \frac{2}{M} \mathbb{E}_{\text{wp}}^{(g)}[N_X^\ell(2L)] \end{aligned}$$

where, for a function F on the moduli space, the Weil-Petersson expectation is defined by:

$$\mathbb{E}_{\text{wp}}^{(g)}[F] = \frac{1}{V_g} \int_{\mathcal{M}_g} F(X) d \text{Vol}_{\text{wp}}^{(g)}(X).$$

The expectation of the length counting function $N_X^\ell(L)$ is estimated in Lemma 4.1 in [18]. We sum up the equations (4.1) with $k = 1$ and (4.3) for k between 1 and $\lfloor \frac{g}{2} \rfloor$ in this Lemma, with the length $2L$, and use equation (3.19), to prove that, as $g \rightarrow +\infty$:

$$\mathbb{E}_{\text{wp}}^{(g)}[N_X^\ell(2L)] = \mathcal{O}(Le^{2L}).$$

□

3. PROOF OF THEOREMS 4 AND 5 AT THE BOTTOM OF THE SPECTRUM

Let us prove Theorems 4 and 5 in the case where the spectral window is bounded above, more precisely $0 \leq a \leq b \leq 1$. The reason why we make the assumption $b \leq 1$ is that our

choice of test function behaves poorly for large values of b . The value 1 is arbitrary, and could be replaced by any fixed value larger than $\frac{1}{2}$ (because the proof in the case away from small eigenvalues only works for $a \geq \frac{1}{2}$).

3.1. Pretrace formula, test function and sketch of the proof.

The Selberg pretrace formula. Our main tool in this proof is the Selberg pretrace formula, which can be expressed the following way.

Theorem 7 (Selberg pretrace formula [22]). *Let $X = \mathcal{H}/\Gamma$ be a compact hyperbolic surface. Let $(\psi_j)_{j \geq 0}$ be an orthonormal basis of eigenfunctions of the Laplacian in $L^2(X)$ of associated non decreasing sequence of eigenvalues $(\lambda_j)_{j \geq 0}$. Let $h : \mathbb{C} \rightarrow \mathbb{C}$ be a function satisfying:*

- (1) $h(-r) = h(r)$ for any $r \in \mathbb{C}$;
- (2) h is analytic in the strip $|\operatorname{Im} z| \leq \frac{1}{2} + \varepsilon$ for some $\varepsilon > 0$;
- (3) for all r in that strip,

$$(8) \quad h(r) = \mathcal{O}\left(\frac{1}{(1 + |r|^2)^{1+\varepsilon}}\right).$$

Then the following formula holds (with every term well-defined and converging):

$$(9) \quad \sum_{j=0}^{+\infty} f(\lambda_j) = \frac{\mu_X(X)}{4\pi} \int_{\frac{1}{4}}^{+\infty} f(\lambda) \tanh\left(\pi\sqrt{\lambda - \frac{1}{4}}\right) d\lambda + \int_D \sum_{\gamma \in \Gamma \setminus \{\operatorname{id}\}} K(z, \gamma \cdot z) d\mu_{\mathcal{H}}(z)$$

where $f : [0, +\infty) \rightarrow \mathbb{C}$ is defined by $f\left(\frac{1}{4} + r^2\right) = h(r)$, D is a fundamental domain of the action of Γ on \mathcal{H} and $K(z, w) = K(d_{\mathcal{H}}(z, w))$ is the kernel associated to h . K has the following expression:

$$(10) \quad K(\rho) = -\frac{1}{\sqrt{2}\pi} \int_{\rho}^{+\infty} \frac{g'(u)}{\sqrt{\cosh u - \cosh \rho}} du,$$

where $g(u) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} h(r) e^{iru} dr$ is the inverse Fourier transform of the function h .

Here is a brief description of the terms in equation (9).

- The left hand term is a spectral average with density f . For instance, if f resembles the indicator function of a segment (as will happen in the following), it will be close to the number of eigenvalues in that spectral window.
- The first term on the right hand side does not depend on the metric on X but only on its topology, because the volume of any hyperbolic surface of genus g is $2\pi(2g - 2)$. If f approaches a step function, then it will converge to the integral term in Theorem 5.
- The last term is a geometric term. It can be reformulated in terms of lengths of the closed geodesics on X (this being the Selberg trace formula). It is the one we are going to estimate using our geometric assumptions from Corollary 3. This term is going to behave better with highly regular test functions, because of the Fourier transform. This will limit the speed to which the test function we use can approximate the indicator function of $[a, b]$.

The test function. Let $0 \leq a \leq b \leq 1$. The test function we are going to use in this Section is defined by

$$h_t(r) = f_t\left(\frac{1}{4} + r^2\right),$$

where, for $\lambda \geq 0$,

$$f_t(\lambda) = (\mathbb{1}_{[a,b]} \star v_t)(\lambda) = \frac{t}{\sqrt{\pi}} \int_a^b \exp(-t^2(\lambda - \mu)^2) d\mu.$$

Here, $t > 0$ is a parameter that will grow like $\sqrt{\log g}$, and $v_t(x) = \frac{t}{\sqrt{\pi}} \exp(-t^2 \rho^2)$ is the centered normalized Gaussian of variance $\frac{1}{t}$. As a consequence, f_t is a smooth (pointwise) approximation of the function $\tilde{\mathbb{1}}_{[a,b]}$ defined by

$$\tilde{\mathbb{1}}_{[a,b]}(x) = \begin{cases} 1 & \text{if } x \in (a, b) \\ \frac{1}{2} & \text{if } x = a \text{ or } x = b \\ 0 & \text{otherwise.} \end{cases}$$

It is clear that $h_t : \mathbb{C} \rightarrow \mathbb{C}$ is analytic and even. In order to apply the pretrace formula, we need to prove an estimate like the one in equation (8), which is the aim of the following Lemma.

Lemma 8. *Let $0 \leq a \leq b$ and $t > 0$. There exists a constant $C = C(a, b, t) > 0$ such that*

$$\forall x \in \mathbb{R}, \forall y \in [-1, 1], |h_t(x + iy)| \leq C \exp(-t^2 x^4 + 2t^2(3 + b)x^2).$$

Proof. Let us write

$$h_t(r) = \frac{t}{\sqrt{\pi}} \int_{a-\frac{1}{4}}^{b-\frac{1}{4}} \exp(-t^2(r^2 - \mu)^2) d\mu.$$

The modulus of the integrand for a $r = x + iy$, $-1 \leq y \leq 1$ is

$$\begin{aligned} \left| \exp(-t^2(r^2 - \mu)^2) \right| &= \exp(-t^2(x^4 + y^4 + \mu^2 - 6x^2y^2 - 2\mu x^2 + 2\mu y^2)) \\ &\leq \exp\left(-t^2 x^4 + 6t^2 x^2 + 2t^2\left(b - \frac{1}{4}\right)x^2 + 2t^2\left|\frac{1}{4} - a\right|\right). \end{aligned}$$

We then integrate this inequality between $a - \frac{1}{4}$ and $b - \frac{1}{4}$. □

This exponential decay guarantees a polynomial decay in the strip $\{|\operatorname{Im} r| \leq 1\}$ like the one required in the pretrace formula. Therefore, one can apply it to h_t , and rewrite equation (9) as:

$$(11) \quad \frac{1}{\mu_X(X)} \sum_{j=0}^{+\infty} f_t(\lambda_j) = \frac{1}{4\pi} \int_{\frac{1}{4}}^{+\infty} \mathbb{1}_{[a,b]}(\lambda) \tanh\left(\pi\sqrt{\lambda - \frac{1}{4}}\right) d\lambda + \mathcal{R}_I(t, a, b) + \mathcal{R}_K(X, t, a, b)$$

where

$$(12) \quad \mathcal{R}_I(t, a, b) = \frac{1}{4\pi} \int_{\frac{1}{4}}^{+\infty} (f_t(\lambda) - \mathbb{1}_{[a,b]}(\lambda)) \tanh\left(\pi\sqrt{\lambda - \frac{1}{4}}\right) d\lambda,$$

$$(13) \quad \mathcal{R}_K(X, t, a, b) = \frac{1}{\mu_X(X)} \int_D \sum_{\gamma \in \Gamma \setminus \{\text{id}\}} K_t(z, \gamma \cdot z) d\mu_{\mathcal{H}}(z),$$

K_t is the kernel associated to h_t and D is a fundamental domain of $X = \mathcal{H}/\Gamma$.

Sketch of the proof. There are four steps in the proof of Theorems 4 and 5 at the bottom of the spectrum.

- Control the integral term $\mathcal{R}_I(t, a, b)$ (Section 3.2).
- Estimate the geometric term $\mathcal{R}_K(X, t, a, b)$ using the uniform discreteness and the Benjamini-Schramm convergence assumptions (Section 3.3).
- Control $N_X^\Delta(a, b)$ with $\sum_{j=0}^{+\infty} f_t(\lambda_j)$, and deduce Theorem 4 (Section 3.4).
- Compare more precisely the sum $\sum_{j=0}^{+\infty} f_t(\lambda_j)$ to $N_X^\Delta(a, b)$, and conclude (Section 3.5).

3.2. The integral term. In order to control $\mathcal{R}_I(t, a, b)$, we need to know more about the speed of convergence of f_t towards $\tilde{\mathbb{1}}_{[a,b]}$ as t goes to infinity, which is the aim of the following Lemma.

Lemma 9. *Let $0 \leq a \leq b$. For any $t > 0$ and $\lambda \in [0, +\infty)$,*

$$(14) \quad |f_t(\lambda) - \tilde{\mathbb{1}}_{[a,b]}(\lambda)| \leq \begin{cases} s(t|\lambda - a|) & \text{if } \lambda \in [0, a) \cup \{b\} \\ s(t|\lambda - a|) + s(t|\lambda - b|) & \text{if } \lambda \in (a, b) \\ s(t|\lambda - b|) & \text{if } \lambda \in \{a\} \cup (b, +\infty) \end{cases}$$

where $s : (0, +\infty) \rightarrow \mathbb{R}$ is the (decreasing) function defined by $s(\rho) = \frac{e^{-\rho^2}}{2\sqrt{\pi}\rho}$.

Proof. Let us assume that $\lambda \in (b, +\infty)$. In that case, $\tilde{\mathbb{1}}_{[a,b]}(\lambda) = 0$, and

$$|f_t(\lambda) - \tilde{\mathbb{1}}_{[a,b]}(\lambda)| = f_t(\lambda) \leq \frac{1}{\sqrt{\pi}} \int_{t(\lambda-b)}^{+\infty} e^{-\rho^2} d\rho \leq \frac{e^{-t^2(\lambda-b)^2}}{2\sqrt{\pi} t(\lambda-b)}$$

since for $\rho > t(\lambda-b)$, $1 < \frac{2\rho}{2t(\lambda-b)}$. All the other cases can be proved in the same way, using, when $\lambda \in [a, b]$, the fact that the Gaussian we used in the definition of f_t is normalized. \square

We can now prove the following estimate.

Proposition 10. *Let $0 \leq a \leq b$. For any $t > 0$,*

$$\mathcal{R}_I(t, a, b) = \mathcal{O}\left(\frac{1}{t}\right).$$

Whenever $b \leq \frac{1}{4}$, if $\delta_b = \frac{1}{4} - b$,

$$\mathcal{R}_I(t, a, b) = \mathcal{O}\left(\frac{e^{-\frac{3}{4}t^2\delta_b^2}}{t^{\frac{3}{2}}}\right).$$

Proof. Let us start with the case when $b \leq \frac{1}{4}$. Since $\tanh(x) \leq x$ for any $x \geq 0$,

$$\begin{aligned} \mathcal{R}_I(t, a, b) &= \mathcal{O}\left(\int_{\frac{1}{4}}^{+\infty} \frac{\exp(-t^2(\lambda-b)^2)}{t(\lambda-b)} \sqrt{\lambda - \frac{1}{4}} d\lambda\right) \\ &= \mathcal{O}\left(\frac{1}{t^{\frac{3}{2}}} \int_{t\delta_b}^{+\infty} \frac{\exp(-u^2)}{\sqrt{u}} du\right) = \mathcal{O}\left(\frac{\exp(-\frac{3}{4}t^2\delta_b^2)}{t^{\frac{3}{2}}}\right) \end{aligned}$$

since $\frac{\exp(-\frac{1}{4}u^2)}{\sqrt{u}}$ has finite integral on $(0, +\infty)$.

In the general case, since the estimate (14) diverges around a and b , we cut the integral on a small interval of size $\varepsilon > 0$ around those values. Let C_ε be the set

$$C_\varepsilon = \left\{ \lambda \in \left[\frac{1}{4}, +\infty\right) : |\lambda - a| < \varepsilon \text{ or } |\lambda - b| < \varepsilon \right\}.$$

C_ε has at most two connected components, each of them of length at most 2ε . Since $|f_t - \mathbb{1}_{[a,b]}| \leq 1$,

$$\int_{C_\varepsilon} |f_t(\lambda) - \mathbb{1}_{[a,b]}(\lambda)| d\lambda \leq 4\varepsilon.$$

There are at most three connected components in $[\frac{1}{4}, +\infty) \setminus C_\varepsilon$, and the estimate in every case is the same so we limit ourselves to the study of $[b + \varepsilon, +\infty)$. Lemma 9 implies that

$$\int_{b+\varepsilon}^{+\infty} |f_t(\lambda) - \mathbb{1}_{[a,b]}(\lambda)| d\lambda = \mathcal{O}\left(\int_{b+\varepsilon}^{+\infty} \frac{\exp(t^2(\lambda-b)^2)}{t(\lambda-b)} d\lambda\right) = \mathcal{O}\left(\frac{1}{\varepsilon t^2} \int_{\varepsilon t}^{+\infty} e^{-\rho^2} d\rho\right).$$

Putting the two contributions together, we obtain

$$\mathcal{R}_I(t, a, b) = \mathcal{O} \left(\varepsilon + \frac{1}{\varepsilon t^2} \int_{\varepsilon t}^{+\infty} e^{-\rho^2} d\rho \right),$$

which leads to our claim if we set $\varepsilon = \frac{1}{t}$. \square

3.3. The geometric term. Let us now control the geometric term

$$\mathcal{R}_K(X, t, a, b) = \frac{1}{\mu_X(X)} \int_D \sum_{\gamma \in \Gamma \setminus \{\text{id}\}} K_t(z, \gamma \cdot z) d\mu_{\mathcal{H}}(z)$$

for any compact hyperbolic surface $X \in \mathcal{A}_g$ from Corollary 3. In order to do so, we will estimate the kernel function K_t . We will then regroup the terms in the sum according to the distance between z and $\gamma \cdot z$. This is where we will use the Benjamini-Schramm hypothesis. Indeed, if $z \in D$ has a large radius of injectivity, then the decay of K_t will cause the sum to be small. Otherwise, the sum might not be small, but the volume of the set of such z will.

Fourier estimate. In order to estimate the kernel K_t , we first need to know about the derivative of g_t , the inverse Fourier transform of the test function h_t .

Lemma 11. *Let $0 \leq a \leq b \leq 1$ and $r \in (0, 3)$. For any $u \geq r$, $t \geq \frac{1}{200}$,*

$$(15) \quad g'_t(u) = \mathcal{O} \left(r^{-\frac{2}{3}} \exp \left(-t^2 \delta_b^2 - \frac{7}{32} u^{\frac{4}{3}} t^{-\frac{2}{3}} + \frac{3}{16} u^{\frac{2}{3}} t^{\frac{2}{3}} \right) \right)$$

where $\delta_b = \max(\frac{1}{4} - b, 0)$ is the distance between $[a, b]$ and $[\frac{1}{4}, +\infty)$.

Proof. By definition of g_t ,

$$g_t(u) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} h_t(r) e^{iru} dr = \frac{t}{2\pi^{\frac{3}{2}}} \int_{-\infty}^{+\infty} \int_a^b e^{-t^2(\frac{1}{4}+r^2-\mu)^2} e^{iru} d\mu dr.$$

By the change of variables $\tilde{\mu} = t(\mu - \frac{1}{4})$ and $\tilde{r} = \sqrt{t} r$ and Fubini's theorem, one can rewrite this integral as

$$g_t(u) = \frac{1}{2\pi^{\frac{3}{2}} \sqrt{t}} \int_{t(a-\frac{1}{4})}^{t(b-\frac{1}{4})} \int_{-\infty}^{+\infty} e^{-(r^2-\mu)^2} e^{i\frac{ru}{\sqrt{t}}} dr d\mu.$$

As a consequence, the derivative of g_t is

$$(16) \quad g'_t(u) = \frac{i}{2\pi^{\frac{3}{2}} t} \int_{t(a-\frac{1}{4})}^{t(b-\frac{1}{4})} F_{\frac{u}{\sqrt{t}}}(\mu) d\mu$$

where for $u > 0$ and $\mu \in \mathbb{R}$,

$$F_u(\mu) = \int_{-\infty}^{+\infty} r e^{-(r^2-\mu)^2+iru} dr.$$

Let us estimate this integral using a change of contour. Let $R > 0$ and $R' > 0$ be two real parameters. The function $z \in \mathbb{C} \mapsto z e^{-(z^2-\mu)^2+izu}$ is holomorphic, so its contour integral on the rectangle of vertices R' , $R' + iR$, $-R' + iR$, $-R'$ is equal to zero. We compute the modulus of the integrand for a complex number $z = x + iy$:

$$(17) \quad |z e^{-(z^2-\mu)^2+izu}| = \sqrt{x^2 + y^2} \exp(-x^4 - y^4 + 6x^2y^2 - \mu^2 + 2x^2\mu - 2y^2\mu - yu).$$

It follows directly from this inequality that the integrals on the vertical sides of the rectangle go to zero as R' approaches infinity. As a consequence,

$$F_u(\mu) = \int_{\mathbb{R}+iR} z e^{-(z^2-\mu)^2+izu} dz.$$

We use the triangle inequality and equation (17) to deduce that

$$|F_u(\mu)| \leq 2 \exp(-R^4 - \mu^2 - 2R^2\mu - uR) \int_0^{+\infty} (x+R) \exp(-x^4 + 6x^2R^2 + 2x^2\mu) dx.$$

We now study two distinct cases, depending on the sign of μ .

- If $\mu \geq 0$, then

$$|F_u(\mu)| \leq 2 \exp(8R^4 + 4R^2\mu - uR) \int_0^{+\infty} (x+R) e^{-(x^2-x_0^2)^2} dx$$

where $x_0 = \sqrt{3R^2 + \mu} > 0$. We observe that

$$\begin{aligned} \int_0^{+\infty} (x+R) e^{-(x^2-x_0^2)^2} dx &\leq \int_{-x_0}^{+\infty} (|y| + x_0 + R) e^{-y^4} \underbrace{e^{-4y^2x_0(y+x_0)}}_{\leq 1} dy \\ &= \mathcal{O}(1 + x_0 + R), \end{aligned}$$

and therefore

$$F_u(\mu) = \mathcal{O}\left((1 + R + \mu^{\frac{1}{2}}) \exp(8R^4 + 4R^2\mu - uR)\right).$$

- If $\mu < 0$, we do the same with $x_0 = \sqrt{3}R$.

$$\begin{aligned} |F_u(\mu)| &\leq 2 \exp(-\mu^2 + 8R^4 - 2R^2\mu - uR) \int_0^{+\infty} (x+R) e^{-(x^2-3R^2)^2} dx \\ &= \mathcal{O}\left((1 + R) \exp(-\mu^2 + 8R^4 + 2R^2|\mu| - uR)\right). \end{aligned}$$

As a conclusion, for any $\mu \in \mathbb{R}$,

$$F_u(\mu) = \mathcal{O}\left((1 + R + |\mu|^{\frac{1}{2}}) \exp(-\mu_-^2 + 8R^4 + 4R^2|\mu| - uR)\right)$$

where $\mu_- = \min(\mu, 0)$. We take $R = \frac{1}{4} u^{\frac{1}{3}}$ and obtain that, for any $u > 0$ and $\mu \in \mathbb{R}$

$$(18) \quad F_u(\mu) = \mathcal{O}\left((1 + u^{\frac{1}{3}} + |\mu|^{\frac{1}{2}}) \exp\left(-\mu_-^2 - \frac{7}{32}u^{\frac{4}{3}} + \frac{1}{4}|\mu|u^{\frac{2}{3}}\right)\right).$$

We then integrate the upper bound (18) in equation (16).

$$\begin{aligned} g'_t(u) &= \mathcal{O}\left(\frac{1}{t} \int_{t(a-\frac{1}{4})}^{t(b-\frac{1}{4})} |F_{\frac{u}{\sqrt{t}}}(\mu)| d\mu\right) \\ &= \mathcal{O}\left(\frac{1 + u^{\frac{1}{3}}t^{-\frac{1}{6}} + t^{\frac{1}{2}}}{t} \exp\left(-t^2\delta_b^2 - \frac{7}{32}u^{\frac{4}{3}}t^{-\frac{2}{3}}\right) \int_{-\frac{3}{4}t}^{\frac{3}{4}t} \exp\left(\frac{1}{4}|\mu|u^{\frac{2}{3}}t^{-\frac{1}{3}}\right) d\mu\right) \end{aligned}$$

because $|\mu| \leq \frac{3}{4}t$ for any $\mu \in [t(a-\frac{1}{4}), t(b-\frac{1}{4})]$. Replacing the integral by its value and using the hypotheses about t , u and r concludes the proof. \square

Estimate of the kernel function. Let us now estimate the kernel function K_t , using its expression in terms of g'_t , equation (10).

Lemma 12. *Let $0 \leq a \leq b \leq 1$ and $r \in (0, 3)$. For any $\rho \geq r$, $t \geq \frac{1}{200}$,*

$$(19) \quad K_t(\rho) = \begin{cases} \mathcal{O}\left(\frac{t}{r^2} \exp\left(-t^2\delta_b^2 - \frac{\rho^{\frac{4}{3}}}{8t^{\frac{2}{3}}}\right)\right) & \text{if } \rho \geq 6t^2 \\ \mathcal{O}\left(\frac{t}{r^2} \exp\left(-t^2\delta_b^2 + t^2 - \frac{\rho^{\frac{4}{3}}}{8t^{\frac{2}{3}}}\right)\right) & \text{if } \rho \leq 6t^2 \end{cases}$$

where $\delta_b = \max(\frac{1}{4} - b, 0)$ is the distance between $[a, b]$ and $[\frac{1}{4}, +\infty)$.

Proof. By definition of the kernel associated to h_t ,

$$K_t(\rho) = -\frac{1}{\sqrt{2\pi}} \int_{\rho}^{+\infty} \frac{g'_t(u)}{\sqrt{\cosh u - \cosh \rho}} du.$$

Using equation (15), we obtain

$$(20) \quad K_t(\rho) = \mathcal{O} \left(\frac{\exp(-t^2 \delta_b^2)}{r^{\frac{2}{3}}} \int_{\rho}^{+\infty} \underbrace{\frac{\exp\left(-\frac{7}{32}u^{\frac{4}{3}}t^{-\frac{2}{3}} + \frac{3}{16}u^{\frac{2}{3}}t^{\frac{2}{3}}\right)}{\sqrt{\cosh u - \cosh \rho}}}_{(*)} du \right).$$

Let us cut this integral into two contributions $I_1 = \int_{\rho}^{\rho_{\text{cut}}} (*)$ and $I_2 = \int_{\rho_{\text{cut}}}^{+\infty} (*)$ where

$$\rho_{\text{cut}} = \max(2\rho, 12t^2).$$

This choice of ρ_{cut} allows us deal with the cancellation of the denominator in I_1 only, and to be in the asymptotic regime where the decreasing part of the exponential term is predominant everywhere in I_2 .

In the second integral, since $u \geq \rho_{\text{cut}} \geq 12t^2$, $\frac{3}{16}u^{\frac{2}{3}}t^{\frac{2}{3}} < \frac{2}{32}u^{\frac{4}{3}}t^{-\frac{2}{3}}$. Hence, the quantity in the exponential function is less than $-\frac{5}{32}u^{\frac{4}{3}}t^{-\frac{2}{3}}$. We deal with the denominator by observing that $\cosh u - \cosh \rho \geq \frac{1}{2}(u - \rho)^2 \geq \frac{1}{2}\rho^2$ since $u \geq \rho_{\text{cut}} \geq 2\rho$. As a consequence,

$$I_2 = \mathcal{O} \left(\frac{1}{\rho} \int_{\rho_{\text{cut}}}^{\infty} \exp\left(-\frac{5}{32}u^{\frac{4}{3}}t^{-\frac{2}{3}}\right) du \right).$$

This integral can be controlled by observing that $1 \leq u^{\frac{1}{3}}\rho_{\text{cut}}^{-\frac{1}{3}}$, and then by explicit integration:

$$(21) \quad I_2 = \mathcal{O} \left(\frac{t^{\frac{2}{3}}}{\rho \rho_{\text{cut}}^{\frac{1}{3}}} \exp\left(-\frac{5}{32}\rho_{\text{cut}}^{\frac{4}{3}}t^{-\frac{2}{3}}\right) \right) = \mathcal{O} \left(\frac{t^{\frac{2}{3}}}{r^{\frac{4}{3}}} \exp\left(-\frac{5}{16}\rho_{\text{cut}}^{\frac{4}{3}}t^{-\frac{2}{3}}\right) \right).$$

Now, in the first integral we use the inequality $\cosh u - \cosh \rho \geq (u - \rho) \sinh \rho \geq (u - \rho)\rho$.

$$\begin{aligned} I_1 &\leq \frac{\exp\left(-\frac{7}{32}\rho^{\frac{4}{3}}t^{-\frac{2}{3}} + \frac{3}{16}\rho_{\text{cut}}^{\frac{2}{3}}t^{\frac{2}{3}}\right)}{\sqrt{\rho}} \int_{\rho}^{\rho_{\text{cut}}} \frac{du}{\sqrt{u - \rho}} \\ &= \frac{1}{2} \sqrt{\frac{\rho_{\text{cut}} - \rho}{\rho}} \exp\left(-\frac{7}{32}\rho^{\frac{4}{3}}t^{-\frac{2}{3}} + \frac{3}{16}\rho_{\text{cut}}^{\frac{2}{3}}t^{\frac{2}{3}}\right). \end{aligned}$$

- When $\rho \leq 6t^2$, $\rho_{\text{cut}} = 12t^2$ so

$$(22) \quad I_1 = \mathcal{O} \left(\frac{t}{r^{\frac{1}{2}}} \exp\left(-\frac{7}{32}\rho^{\frac{4}{3}}t^{-\frac{2}{3}} + t^2\right) \right).$$

- Otherwise, $\rho_{\text{cut}} = 2\rho$ so

$$(23) \quad I_1 = \mathcal{O} \left(\exp\left(-\frac{7}{32}\rho^{\frac{4}{3}}t^{-\frac{2}{3}} + \frac{3}{16}2^{\frac{2}{3}}\rho^{\frac{2}{3}}t^{\frac{2}{3}}\right) \right) = \mathcal{O} \left(\exp\left(-\frac{1}{8}\rho^{\frac{4}{3}}t^{-\frac{2}{3}}\right) \right)$$

because the fact that $\rho \geq 6t^2$ implies that $\frac{3}{16}2^{\frac{2}{3}}\rho^{\frac{2}{3}}t^{\frac{2}{3}} < \frac{3}{32}\rho^{\frac{4}{3}}t^{-\frac{2}{3}}$.

Putting together (20), (21), (22) and (23) leads to what was claimed. \square

Kernel summation. We can now proceed to the estimate of the geometric term. In order to do so, we will arrange the terms in the sum depending on the distance between z and $\gamma \cdot z$, and use Lemma 12.

Lemma 13. *Let $0 \leq a \leq b \leq 1$. Let $r \in (0, 3)$ and X be a compact hyperbolic surface of radius of injectivity larger than r . For any $t \geq \frac{1}{200}$, $L \geq 2^{12}t^2$,*

$$(24) \quad \mathcal{R}_K(X, t, a, b) = \mathcal{O} \left(\frac{t^3}{r^4} \exp(-t^2 \delta_b^2) \left[\exp(-L) + \frac{\mu_X(X^-(L))}{\mu_X(X)} \exp(L) \right] \right)$$

where $\delta_b = \max(\frac{1}{4} - b, 0)$ is the distance between $[a, b]$ and $[\frac{1}{4}, +\infty)$ and $X^-(L)$ is the set of points in X of radius of injectivity smaller than L .

Proof. Let us write a fundamental domain D of $X = \mathcal{H}/\Gamma$ as a disjoint union of $D^+(L)$ and $D^-(L)$, respectively the set of points in D of radius of injectivity larger and smaller than L . We cut according to this partition of D the integral in the definition of

$$\mathcal{R}_K(X, t, a, b) = \frac{1}{\mu_X(X)} \int_D \sum_{\gamma \neq \text{id}} K_t(z, \gamma \cdot z) \, d\mu_{\mathcal{H}}(z)$$

into two contributions, $\mathcal{R}_K^+(X, t, a, b, L)$ and $\mathcal{R}_K^-(X, t, a, b, L)$.

Let us first estimate the term $\mathcal{R}_K^+(X, t, a, b, L)$. Lemma 12 allows us to control $K_t(z, \gamma \cdot z)$ in terms of the distance between z and $\gamma \cdot z$. In order to use it, we regroup the terms of the sum according to this quantity:

$$\mathcal{R}_K^+(X, t, a, b, L) = \frac{1}{\mu_X(X)} \int_{D^+(L)} \sum_{j \geq L} \sum_{\substack{\gamma \neq \text{id} \\ j \leq d_{\mathcal{H}}(z, \gamma \cdot z) < j+1}} K_t(z, \gamma \cdot z) \, d\mu_{\mathcal{H}}(z).$$

One should notice that the sum only runs over integers larger than or equal to L as a consequence of the definition of $D^+(L)$. For any $z \in D^+(L)$, $j \geq L$ and $\gamma \in \Gamma \setminus \{\text{id}\}$ such that $j \leq d_{\mathcal{H}}(z, \gamma \cdot z) < j+1$, by Lemma 12 and since $d_{\mathcal{H}}(z, \gamma \cdot z) \geq L > 6t^2$,

$$K_t(z, \gamma \cdot z) = \mathcal{O} \left(\frac{t}{r^2} \exp \left(-t^2 \delta_b^2 - \frac{j^{\frac{4}{3}}}{8t^{\frac{2}{3}}} \right) \right).$$

We then apply the following Lemma, inspired by [17], to control the number of $\gamma \in \Gamma$ contributing to the sum.

Lemma 14. *Let $r \in (0, 3)$ and $X = \mathcal{H}/\Gamma$ be a compact hyperbolic surface of radius of injectivity larger than r . Then, for any $z \in \mathcal{H}$ and any $j > 0$,*

$$(25) \quad \#\{\gamma \in \Gamma : d_{\mathcal{H}}(z, \gamma \cdot z) \leq j\} = \mathcal{O} \left(\frac{e^j}{r^2} \right).$$

Proof. By definition of the radius of injectivity, the balls B_γ of center $\gamma \cdot z$ and radius $\frac{r}{2}$, for $\gamma \in \Gamma$, are disjoint. If γ is such that $d_{\mathcal{H}}(z, \gamma \cdot z) \leq j$, then B_γ is included in the ball of center z and radius $j + \frac{r}{2}$. Since the volume of a hyperbolic ball of radius R is $\cosh(R) - 1$, the number of such γ is smaller than

$$\frac{\cosh(j + \frac{r}{2}) - 1}{\cosh(\frac{r}{2}) - 1} = \mathcal{O} \left(\frac{e^j}{r^2} \right).$$

□

Therefore, and because $\mu_{\mathcal{H}}(D^+(L)) \leq \mu_X(X)$,

$$\mathcal{R}_K^+(X, t, a, b, L) = \mathcal{O} \left(\frac{t}{r^4} \exp(-t^2 \delta_b^2) \sum_{j \geq L} \exp \left(j - \frac{j^{\frac{4}{3}}}{8t^{\frac{2}{3}}} \right) \right).$$

Let $S(t, L)$ denote the sum. The fact that $L \geq 2^{12}t^2$ implies that, for any $j \geq L$, $j \leq \frac{j^{\frac{4}{3}}}{16t^{\frac{2}{3}}}$. As a consequence,

$$S(t, L) \leq \sum_{j \geq L} \exp\left(-\frac{j^{\frac{4}{3}}}{16t^{\frac{2}{3}}}\right)$$

which can be estimated by comparison with an integral:

$$S(t, L) \leq \left(1 + \frac{12t^{\frac{2}{3}}}{L^{\frac{1}{3}}}\right) \exp\left(-\frac{L^{\frac{4}{3}}}{16t^{\frac{2}{3}}}\right) = \mathcal{O}(\exp(-L)) \quad \text{since } L^{\frac{1}{3}} \geq 16t^{\frac{2}{3}}.$$

Therefore,

$$(26) \quad \mathcal{R}_K^+(X, t, a, b, L) = \mathcal{O}\left(\frac{t}{r^4} \exp(-t^2\delta_b^2) \exp(-L)\right).$$

The same method, applied to $\mathcal{R}_K^-(X, t, a, b, L)$, leads to

$$\mathcal{R}_K^-(X, t, a, b, L) = \mathcal{O}\left(\frac{t}{r^4} \exp(-t^2\delta_b^2) \frac{\mu_X(X^-(L))}{\mu_X(X)} \sum_{j \geq 0} (1 + \mathbb{1}_{[0, 6t^2]}(j) e^{t^2}) \exp\left(j - \frac{j^{\frac{4}{3}}}{8t^{\frac{2}{3}}}\right)\right).$$

We cut the sum at $j_{\text{cut}}^{(1)} = \lfloor 6t^2 \rfloor + 1$ and $j_{\text{cut}}^{(2)} = \lfloor 2^{12}t^2 \rfloor + 1$. The term where $j \geq j_{\text{cut}}^{(2)}$ satisfies the same estimate as before since $j_{\text{cut}}^{(2)} \geq 2^{12}t^2$, and therefore is $\mathcal{O}(\exp(-j_{\text{cut}}^{(2)})) = \mathcal{O}(1)$. We control naively the two other terms, which are $\mathcal{O}(t^2 \exp(2^{12}t^2))$. As a consequence,

$$(27) \quad \mathcal{R}_K^-(X, t, a, b, L) = \mathcal{O}\left(\frac{t^3}{r^4} \exp(-t^2\delta_b^2) \frac{\mu_X(X^-(L))}{\mu_X(X)} \exp(2^{12}t^2)\right).$$

Our claim follows directly from equations (26) and (27). \square

Geometric estimate. We are now ready to use the geometric properties of random hyperbolic surfaces, and obtain an estimate of the geometric term in the pretrace formula true with high probability.

Proposition 15. *For any large enough g , any $0 \leq a \leq b \leq 1$ and any hyperbolic surface $X \in \mathcal{A}_g$ defined in Corollary 3, if we set $t = \frac{\sqrt{\log g}}{64\sqrt{6}}$, then*

$$(28) \quad \mathcal{R}_K(X, t, a, b) = \mathcal{O}\left(\frac{g^{-\frac{1}{3 \cdot 2^{13}}\delta_b^2}}{(\log g)^{\frac{3}{4}}}\right)$$

where $\delta_b = \max(\frac{1}{4} - b, 0)$ is the distance between $[a, b]$ and $[\frac{1}{4}, +\infty)$.

Proof. It is a direct consequence of Lemma 13 and the properties of the elements of \mathcal{A}_g , namely that if X is an element of \mathcal{A}_g and $L = \frac{1}{6} \log g = 2^{12}t^2$, then

- the injectivity radius of X is greater than $r = g^{-\frac{1}{24}}(\log g)^{\frac{9}{16}}$;
- $\frac{\mu_X(X^-(L))}{\mu_X(X)} = \mathcal{O}(g^{-\frac{1}{3}})$.

Since $L \geq 2^{12}t^2$, $r < 3$ and $t > \frac{1}{200}$, we can apply Lemma 13:

$$\begin{aligned} \mathcal{R}_K(X, t, a, b) &= \mathcal{O}\left(\frac{t^3}{r^4} \exp(-t^2\delta_b^2) \left[\exp(-L) + \frac{\mu_X(X^-(L))}{\mu_X(X)} \exp(L)\right]\right) \\ &= \mathcal{O}\left(\frac{(\log g)^{\frac{3}{2}}}{g^{-\frac{1}{6}}(\log g)^{\frac{9}{4}}} g^{-\frac{1}{3 \cdot 2^{13}}\delta_b^2} \left[g^{-\frac{1}{6}} + g^{-\frac{1}{3} + \frac{1}{6}}\right]\right). \end{aligned}$$

\square

3.4. Proof of Theorem 4 at the bottom of the spectrum. When we put together equation (11), Proposition 10 and 15, we obtain directly the following statement, which is an estimate of the pretrace formula. Theorems 4 and 5 (when $b \leq 1$) will then follow, since the spectral sum $\sum_{j=0}^{+\infty} f_t(\lambda_j)$ approaches $N_X^\Delta(a, b)$ as $t \rightarrow +\infty$.

Corollary 16. *For any large enough g , any $0 \leq a \leq b \leq 1$ and any hyperbolic surface $X \in \mathcal{A}_g$ defined in Corollary 3, if we set $t = \frac{\sqrt{\log g}}{64\sqrt{6}}$, then*

$$(29) \quad \frac{1}{\mu_X(X)} \sum_{j=0}^{+\infty} f_t(\lambda_j) = \frac{1}{4\pi} \int_{\frac{1}{4}}^{+\infty} \mathbb{1}_{[a,b]}(\lambda) \tanh\left(\pi\sqrt{\lambda - \frac{1}{4}}\right) d\lambda + \mathcal{O}\left(\frac{1}{\sqrt{\log g}}\right).$$

If $b \leq \frac{1}{4}$, then we also have, for $c = 2^{-15}$,

$$(30) \quad \frac{1}{\mu_X(X)} \sum_{j=0}^{+\infty} f_t(\lambda_j) = \mathcal{O}\left(\frac{g^{-c(\frac{1}{4}-b)^2}}{(\log g)^{\frac{3}{4}}}\right).$$

Proof of Theorem 4 when $b \leq 1$. Let $t = \frac{\sqrt{\log g}}{64\sqrt{6}}$. Let us distinguish two cases.

- Whenever $t(b-a) \geq \frac{1}{\sqrt{3}}$, since the function f_t only takes positive values,

$$\frac{N_X^\Delta(a, b)}{\mu_X(X)} \times \inf_{[a,b]} f_t \leq \frac{1}{\mu_X(X)} \sum_{j=0}^{+\infty} f_t(\lambda_j).$$

It follows directly from equation (30) that the right hand term is

$$\mathcal{O}\left(b - a + \frac{1}{\sqrt{\log g}}\right).$$

In order to deal with the infimum, we use Lemma 9 that states that

$$\inf_{[a,b]} f_t \geq \frac{1}{2} - \frac{e^{-t^2(b-a)^2}}{2\sqrt{\pi}t(b-a)} \geq \frac{1}{2} - \frac{\sqrt{3}e^{-\frac{1}{3}}}{2\sqrt{\pi}} \geq \frac{1}{10}$$

since we assumed $t(b-a) \geq \frac{1}{\sqrt{3}}$. Therefore, $(\inf_{[a,b]} f_t)^{-1} = \mathcal{O}(1)$.

- Otherwise, let a' be $b - \frac{1}{\sqrt{3}t}$. For large enough g , $a' > -\frac{1}{2}$, and careful examination of the previous proof implies Corollary 16 still holds. By the first part and since $a' \leq a$,

$$\frac{N_X^\Delta(a, b)}{\mu_X(X)} \leq \frac{N_X^\Delta(a', b)}{\mu_X(X)} = \mathcal{O}\left(b - a' + \frac{1}{\sqrt{\log g}}\right) = \mathcal{O}\left(\frac{1}{\sqrt{\log g}}\right).$$

The other proof is the same, using the small eigenvalue case of Corollary 16. \square

3.5. Proof of Theorem 5 at the bottom of the spectrum. Let us now proceed to the proof of Theorem 5 when $b \leq 1$. Beware that, in the proof of the lower bound, we will need to use Theorem 4 for any $0 \leq a \leq b$. This is not an issue, as was shown in Figure 1.

Proof of the upper bound of Theorem 5 when $b \leq 1$. Let $t = \frac{\sqrt{\log g}}{64\sqrt{6}}$.

If $t(b-a) \leq \sqrt{2e}$, then the integral term is $\mathcal{O}(b-a) = \mathcal{O}\left(\frac{1}{\sqrt{\log g}}\right)$, so the result follows directly from Theorem 4.

Let us assume $t(b-a) \geq \sqrt{2e}$. The issue in the previous estimate was that the convergence of f_t is slow around a and b , and noticeably $f_t(a)$ and $f_t(b)$ go to $\frac{1}{2}$ and not 1 as $t \rightarrow +\infty$. In order to deal with this, we cut a small segment around a and b . Let $\frac{1}{t} \leq \varepsilon \leq \frac{b-a}{2}$, then

$$N_X^\Delta(a, b) = N_X^\Delta(a, a + \varepsilon) + N_X^\Delta(a + \varepsilon, b - \varepsilon) + N_X^\Delta(b - \varepsilon, b).$$

By Theorem 4,

$$\frac{N_X^\Delta(a, a + \varepsilon) + N_X^\Delta(b - \varepsilon, b)}{\mu_X(X)} = \mathcal{O}\left(\varepsilon + \sqrt{\frac{b+1}{\log g}}\right) = \mathcal{O}(\varepsilon).$$

We use the same method as before to control the middle term:

$$\begin{aligned} \frac{N_X^\Delta(a + \varepsilon, b - \varepsilon)}{\mu_X(X)} \times \inf_{[a+\varepsilon, b-\varepsilon]} f_t &\leq \frac{1}{\mu_X(X)} \sum_{j=0}^{+\infty} f_t(\lambda_j) \\ &\leq \frac{1}{4\pi} \int_{\frac{1}{4}}^{+\infty} \mathbb{1}_{[a,b]}(\lambda) \tanh\left(\pi\sqrt{\lambda - \frac{1}{4}}\right) d\lambda + \frac{C'}{\sqrt{\log g}} \end{aligned}$$

for a constant $C' > 0$, given by Corollary 16. By Lemma 9, and because $\varepsilon t \geq 1$,

$$\inf_{[a+\varepsilon, b-\varepsilon]} f_t \geq 1 - \frac{e^{-t^2\varepsilon^2}}{2\sqrt{\pi\varepsilon t}} \geq \frac{1}{1 + e^{-\varepsilon^2 t^2}}.$$

Putting all the contributions together, there exists a constant $C'' > 0$ such that

$$\frac{N_X^\Delta(a, b)}{\mu_X(X)} \leq \frac{1}{4\pi} \int_{\frac{1}{4}}^{+\infty} \mathbb{1}_{[a,b]}(\lambda) \tanh\left(\pi\sqrt{\lambda - \frac{1}{4}}\right) d\lambda + C'' \left(\varepsilon + (b-a) e^{-t^2\varepsilon^2}\right).$$

We can now set

$$\varepsilon = \frac{1}{t} \sqrt{\log\left(\frac{\sqrt{e} t(b-a)}{\sqrt{2}}\right)}.$$

The hypothesis $t(b-a) \geq \sqrt{2e}$ directly implies that $\varepsilon t \geq 1$. Furthermore, the fact that for any $x \geq 1$, $\sqrt{\log x} \leq \frac{x}{\sqrt{2e}}$ implies that $\varepsilon \leq \frac{b-a}{2}$. Direct substitution of ε and t by their values in the previous estimate leads to our claim. \square

Proof of the lower bound of Theorem 5 when $b \leq 1$. Let $t = \frac{\sqrt{\log g}}{64\sqrt{6}}$. Since $0 \leq f_t \leq 1$,

$$N_X^\Delta(a, b) \geq \sum_{a \leq \lambda_j \leq b} f_t(\lambda_j) = \sum_{j=0}^{+\infty} f_t(\lambda_j) - \sum_{0 \leq \lambda_j < a} f_t(\lambda_j) - \sum_{\lambda_j > b} f_t(\lambda_j).$$

By Corollary 16, there exists a $C' > 0$ such that

$$\frac{1}{\mu_X(X)} \sum_{j=0}^{+\infty} f_t(\lambda_j) \geq \frac{1}{4\pi} \int_{\frac{1}{4}}^{+\infty} \mathbb{1}_{[a,b]}(\lambda) \tanh\left(\pi\sqrt{\lambda - \frac{1}{4}}\right) d\lambda - \frac{C'}{\sqrt{\log g}},$$

so it suffices to prove that the two remaining terms are $\mathcal{O}\left(\frac{\mu_X(X)}{\sqrt{\log g}}\right)$.

Both the terms behave the same way, so we only detail the sum over b . Let us divide $(b, +\infty)$ using a subdivision $b_k = b + \frac{k}{t}$, $k \geq 0$. We regroup the terms of the sum according to these numbers.

$$\begin{aligned} \frac{1}{\mu_X(X)} \sum_{\lambda_j > b} f_t(\lambda_j) &= \sum_{k=0}^{+\infty} \frac{1}{\mu_X(X)} \sum_{b_k < \lambda_j \leq b_{k+1}} f_t(\lambda_j) \\ &\leq \sum_{k=0}^{+\infty} \frac{N_X^\Delta(b_k, b_{k+1})}{\mu_X(X)} \times \sup_{[b_k, b_{k+1}]} f_t \\ &= \mathcal{O}\left(\sum_{k=0}^{+\infty} \left(b_{k+1} - b_k + \sqrt{\frac{b_{k+1} + 1}{\log g}}\right) \times \sup_{[b_k, b_{k+1}]} f_t\right) \end{aligned}$$

by Theorem 4. As a consequence,

$$\frac{1}{\mu_X(X)} \sum_{\lambda_j > b} f_t(\lambda_j) = \mathcal{O} \left(\frac{1}{\sqrt{\log g}} + \frac{1}{\sqrt{\log g}} \sum_{k=1}^{+\infty} \sqrt{k} \times \sup_{[b_k, b_{k+1}]} f_t \right).$$

By Lemma 9,

$$\sum_{k=1}^{+\infty} \sqrt{k} \times \sup_{[b_k, b_{k+1}]} f_t = \mathcal{O} \left(\sum_{k=1}^{+\infty} \frac{\exp(-k^2)}{\sqrt{k}} \right) = \mathcal{O}(1).$$

□

4. PROOF OF THEOREMS 4 AND 5 AWAY FROM SMALL EIGENVALUES

We now proceed to the proof of Theorems 4 and 5 in the case when $\frac{1}{2} \leq a \leq b$. It is easy to see that it suffices to prove the results for these two situations, for we can then apply them to $a, \frac{3}{4}$ and $\frac{3}{4}, b$ and add up the two contributions if $a < \frac{1}{2}$ and $b > 1$.

The proof here is very similar to the previous proof, apart from the fact that the test function we use is different. We will not give all the details, and mostly highlight the differences between the two proofs.

The reason why we need to assume $a \geq \frac{1}{2}$ is that the test function we use behaves poorly for small eigenvalues. We will use the fact that there are at most $2g - 2$ of them [19], and that the spectral window is far enough from them, to deal with this situation (see Section 4.4).

4.1. Pretrace formula, test function and sketch of the proof. We will use once again the Selberg pretrace formula, but with a different test function this time.

The test function. Let $a = \frac{1}{4} + \alpha^2$ and $b = \frac{1}{4} + \beta^2$, for some $0 \leq \alpha \leq \beta$. Let us consider the function

$$h_t(r) = (\mathbb{1}_{[\alpha, \beta]} \star v_t)(r) = \frac{t}{\sqrt{\pi}} \int_{\alpha}^{\beta} \exp(-t^2(r - \rho)^2) d\rho = \frac{1}{\sqrt{\pi}} \int_{t(\alpha-r)}^{t(\beta-r)} \exp(-\rho^2) d\rho,$$

where t still grows like $\sqrt{\log g}$. h_t now is a smooth approximation of the function $\tilde{\mathbb{1}}_{[\alpha, \beta]}$. We make h_t into an even test function by setting $H_t(r) = h_t(r) + h_t(-r)$. It is clear that $H_t : \mathbb{C} \rightarrow \mathbb{C}$ is analytic and even. The following Lemma is an estimate on h_t aimed at applying the pretrace formula, but we make it a bit more precise than necessary for later use.

Lemma 17. *Let $0 \leq \alpha \leq \beta$, $a = \frac{1}{4} + \alpha^2$ and $b = \frac{1}{4} + \beta^2$. For any $t > 0$,*

$$\forall r = x + iy, \quad |h_t(r)| \leq \frac{1}{2\sqrt{\pi} \alpha t} \exp(t^2(y^2 - x^2 + 2\beta x - \alpha^2)).$$

Proof. Let $r = x + iy$. The modulus of the integrand in the definition of $h_t(r)$ is

$$|\exp(-t^2(r - \rho)^2)| = \exp(-t^2(x - \rho)^2 + t^2y^2).$$

As a consequence,

$$|h_t(r)| \leq \frac{t}{\sqrt{\pi}} \exp(t^2(y^2 - x^2 + 2\beta x)) \int_{\alpha}^{+\infty} \exp(-t^2\rho^2) d\rho$$

which allows us to conclude, using the Gaussian tail estimate. □

Therefore, one can apply the pretrace formula to H_t :

$$(31) \quad \frac{1}{\mu_X(X)} \sum_{j=0}^{+\infty} (h_t(r_j) + h_t(-r_j)) = \frac{1}{4\pi} \int_a^b \tanh \left(\pi \sqrt{\lambda - \frac{1}{4}} \right) d\lambda + \mathcal{R}_I(t, a, b) + \mathcal{R}_K(X, t, a, b)$$

where

$$(32) \quad \mathcal{R}_I(t, a, b) = \frac{1}{4\pi} \int_0^{+\infty} (h_t(r) + h_t(-r) - \mathbb{1}_{[\alpha, \beta]}(r)) r \tanh(\pi r) dr,$$

$$(33) \quad \mathcal{R}_K(X, t, a, b) = \frac{1}{\mu_X(X)} \int_D \sum_{\gamma \in \Gamma \setminus \{\text{id}\}} K_t(z, \gamma \cdot z) d\mu_{\mathcal{H}}(z),$$

K_t is the kernel associated to H_t and D is a fundamental domain of $X = \mathcal{H}/\Gamma$.

Sketch of the proof. The steps of the proof are exactly the same as before, and are organized the same way. The only additional step is dealing with the contributions of the small eigenvalues to the sum $\sum_{j=0}^{+\infty} (h_t(r_j) + h_t(-r_j))$, and can be found in Section 4.4. This is necessary here and was not before because the function h_t is no longer real valued and small on the imaginary axis. This complication is the reason why this test function does not work whenever $a < \frac{1}{2}$.

4.2. The integral term. The integral estimate is the following.

Proposition 18. *Let $\frac{1}{4} \leq a \leq b$. For any $t \geq \frac{1}{10}$,*

$$\mathcal{R}_I(t, a, b) = \mathcal{O}\left(\frac{\sqrt{b}}{t}\right).$$

The proof uses the same method as before, and the following Lemma to control the speed of convergence of h_t towards $\tilde{\mathbb{1}}_{[\alpha, \beta]}$ as t goes to infinity.

Lemma 19. *Let $0 \leq \alpha \leq \beta$. For any $t > 0$ and $r \in \mathbb{R}$,*

$$(34) \quad |h_t(r) - \tilde{\mathbb{1}}_{[\alpha, \beta]}(r)| \leq \begin{cases} s(t|r - \alpha|) & \text{if } r \in (-\infty, \alpha) \cup \{\beta\} \\ s(t|r - \alpha|) + s(t|r - \beta|) & \text{if } r \in (\alpha, \beta) \\ s(t|r - \beta|) & \text{if } r \in \{\alpha\} \cup (\beta, +\infty) \end{cases}$$

where $s : (0, +\infty) \rightarrow \mathbb{R}$ is the (decreasing) function defined in Lemma 9.

4.3. The geometric term. The control of the geometric term is simpler in this case, because the test function H_t is a convolution of two functions with simple Fourier transforms (a Gaussian and a step function). Therefore, its Fourier transform has a simple expression.

Lemma 20. *Let $\frac{1}{4} \leq a \leq b$ and $r \in (0, 3)$. For any $t \geq \frac{1}{10}$, $u > r$,*

$$(35) \quad g'_t(u) = \mathcal{O}\left(\frac{\sqrt{b}}{r} \exp\left(-\frac{u^2}{4t^2}\right)\right).$$

Proof. Let us write $a = \frac{1}{4} + \alpha^2$ and $b = \frac{1}{4} + \beta^2$, with $0 \leq \alpha \leq \beta$.

We can compute g_t explicitly, knowing the Fourier transform of a Gaussian and a step function:

$$g_t(u) = \frac{\beta \operatorname{sinc}(\beta u) - \alpha \operatorname{sinc}(\alpha u)}{\pi} \exp\left(-\frac{u^2}{4t^2}\right)$$

where $\operatorname{sinc}(x) = \frac{\sin x}{x}$. Therefore, the derivative of g_t is

$$g'_t(u) = -\frac{u}{2t^2} g_t(u) + \frac{\beta^2 \operatorname{sinc}'(\beta u) - \alpha^2 \operatorname{sinc}'(\alpha u)}{\pi} \exp\left(-\frac{u^2}{4t^2}\right).$$

We use the fact that $|\operatorname{sinc}(x)| \leq 1$ and $|x \operatorname{sinc}'(x)| = |\cos x - \operatorname{sinc} x| \leq 2$ to conclude. \square

This leads directly to an estimate on the kernel function, by cutting the integral (10) expressing K_t in terms of g_t at 2ρ and using the same inequalities as before for the denominator.

Lemma 21. *Let $\frac{1}{4} \leq a \leq b$ and $r \in (0, 3)$. For any $\rho \geq r$, $t \geq \frac{1}{10}$,*

$$(36) \quad K_t(\rho) = \mathcal{O} \left(\frac{t\sqrt{b}}{r^2} \exp \left(-\frac{\rho^2}{4t^2} \right) \right).$$

Then, the same summation process leads to the following Lemma.

Lemma 22. *Let $\frac{1}{4} \leq a \leq b$ and $r \in (0, 3)$. Let $X = \mathcal{H}/\Gamma$ be a compact hyperbolic surface of injectivity radius larger than r . For any $t \geq \frac{1}{10}$, $L \geq 8t^2$,*

$$(37) \quad \mathcal{R}_K(X, t, a, b) = \mathcal{O} \left(\frac{t^3\sqrt{b}}{r^4} \left[\exp(-L) + \frac{\mu_X(X^-(L))}{\mu_X(X)} \exp(L) \right] \right)$$

where $X^-(L)$ is the set of points in X of radius of injectivity smaller than L .

We can then conclude using the geometric properties of random surfaces.

Proposition 23. *For any large enough g , any $\frac{1}{4} \leq a \leq b$ and any hyperbolic surface $X \in \mathcal{A}_g$ defined in Corollary 3, if we set $t = \frac{\sqrt{\log g}}{4\sqrt{3}}$, then*

$$(38) \quad \mathcal{R}_K(X, t, a, b) = \mathcal{O} \left(\sqrt{\frac{b}{\log g}} \right).$$

4.4. Small eigenvalues term, and proof of Theorem 4 away from them. The behavior of the function h_t is different on the imaginary and real axes. Noticeably, the function h_t is positive on the real axis, but it is not real valued on the imaginary axis. This will cause some of the inequalities from the previous part to fail. Also, when a is close to $\frac{1}{4}$, the modulus of h_t on the segment $[-\frac{i}{2}, \frac{i}{2}]$ becomes too large, and the remainder we will obtain will be unsatisfactory. This is the reason why this test function is only suitable for values of a greater than $\frac{1}{2}$.

We shall now deal with the small eigenvalues, so that they do not intervene anymore afterwards.

Lemma 24. *Let $\frac{1}{2} \leq a \leq b$. For any compact hyperbolic surface X and any $t > 0$,*

$$(39) \quad \frac{1}{\mu_X(X)} \sum_{r_j \notin \mathbb{R}} (h_t(r_j) + h_t(-r_j)) = \mathcal{O} \left(\frac{1}{t} \right).$$

Proof. Let $\frac{1}{2} \leq \alpha \leq \beta$ such that $a = \frac{1}{4} + \alpha^2$ and $b = \frac{1}{4} + \beta^2$. If $r_j \notin \mathbb{R}$, then $r_j = iy_j$ with $y_j \in [-\frac{1}{2}, \frac{1}{2}]$. By Lemma 17,

$$|h_t(\pm r_j)| \leq \frac{1}{2\sqrt{\pi} \alpha t} \exp(t^2(y_j^2 - \alpha^2)) = \mathcal{O} \left(\frac{1}{t} \right) \quad \text{since } \alpha \geq \frac{1}{2}.$$

The number of such terms is $\leq 2g - 2 = \mathcal{O}(\mu_X(X))$ by [19]. □

When we put together equation (31), Proposition 18, 23 and Lemma 24, we obtain directly the following statement.

Corollary 25. *For any large enough g , any $\frac{1}{2} \leq a \leq b$ and any hyperbolic surface $X \in \mathcal{A}_g$ defined in Corollary 3, if we set $t = \frac{\sqrt{\log g}}{4\sqrt{3}}$, then*

$$(40) \quad \frac{1}{\mu_X(X)} \sum_{r_j \in \mathbb{R}} (h_t(r_j) + h_t(-r_j)) = \frac{1}{4\pi} \int_a^b \tanh \left(\pi \sqrt{\lambda - \frac{1}{4}} \right) d\lambda + \mathcal{O} \left(\sqrt{\frac{b}{\log g}} \right).$$

It is straightforward to deduce Theorem 4 from this result as was done before.

4.5. Proof of Theorem 5 away from small eigenvalues. A version of Theorem 5 in terms of α and β follows directly from the method of Section 3.

Theorem 26 (Theorem 5 away from small eigenvalues). *There exists a universal constant $C > 0$ such that, for any large enough g , any $\frac{1}{2} \leq \alpha \leq \beta$ and any hyperbolic surface $X \in \mathcal{A}_g$ from Corollary 3, if we set $a = \frac{1}{4} + \alpha^2$ and $b = \frac{1}{4} + \beta^2$, then one can write the counting function $N_X^\Delta(a, b)$ as*

$$(41) \quad \frac{N_X^\Delta(a, b)}{\mu_X(X)} = \frac{1}{4\pi} \int_a^b \tanh\left(\pi\sqrt{\lambda - \frac{1}{4}}\right) d\lambda + R(X, a, b)$$

where

$$-C\sqrt{\frac{b}{\log g}} \leq R(X, a, b) \leq C\sqrt{\frac{b}{\log g}} \log\left(2 + (\beta - \alpha)\sqrt{\log g}\right)^{\frac{1}{2}}.$$

We translate this statement in terms of a and b thanks to the fact that $\beta - \alpha = \frac{b-a}{\beta+\alpha}$, and therefore, as soon as $b \geq \frac{1}{2}$,

$$(42) \quad \beta - \alpha \leq \frac{b-a}{\sqrt{b-\frac{1}{4}}} \leq \sqrt{2} \frac{b-a}{\sqrt{b}}.$$

REFERENCES

- [1] Miklos Abert, Nicolas Bergeron, Ian Biringer, Tsachik Gelander, Nikolay Nikolov, Jean Raimbault, and Iddo Samet. On the growth of Betti numbers of locally symmetric spaces. *Comptes Rendus Mathématique. Académie des Sciences. Paris*, 349(15-16):831–835, 2011.
- [2] Miklos Abert, Nicolas Bergeron, Ian Biringer, Tsachik Gelander, Nikolay Nikolov, Jean Raimbault, and Iddo Samet. On the growth of L^2 -invariants for sequences of lattices in Lie groups. *Annals of Mathematics. Second Series*, 185(3):711–790, 2017.
- [3] Itai Benjamini and Oded Schramm. Recurrence of distributional limits of finite planar graphs. *Electronic Journal of Probability*, 6, 2001.
- [4] Pierre H. Bérard. On the wave equation on a compact Riemannian manifold without conjugate points. *Mathematische Zeitschrift*, 155(3):249–276, August 1977.
- [5] Gérard Besson. Sur la multiplicité de la première valeur propre des surfaces riemanniennes. *Université de Grenoble. Annales de l'Institut Fourier*, 30(1):x, 109–128, 1980.
- [6] Charles Bordenave. A new proof of Friedman's second eigenvalue Theorem and its extension to random lifts. *arXiv:1502.04482*, 2015.
- [7] Lewis Bowen. Cheeger constants and L^2 -Betti numbers. *Duke Mathematical Journal*, 164(3):569–615, 2015.
- [8] Robert Brooks and Eran Makover. Random Construction of Riemann Surfaces. *Journal of Differential Geometry*, 68(1):121–157, 2004.
- [9] Peter Buser. Riemannsche Flächen mit Eigenwerten in $(0,1/4)$. *Commentarii Mathematici Helvetici*, 52(1):25–34, 1977.
- [10] Peter Buser. A note on the isoperimetric constant. *Annales Scientifiques de l'École Normale Supérieure. Quatrième Série*, 15(2):213–230, 1982.
- [11] Peter Buser. *Geometry and Spectra of Compact Riemann Surfaces*. Birkhäuser, Boston, 1992.
- [12] Jeff Cheeger. A lower bound for the smallest eigenvalue of the Laplacian. *Problems in analysis (Papers dedicated to Salomon Bochner, 1969)*, pages 195–199, 1970.
- [13] Shiu Yuen Cheng. Eigenvalue comparison theorems and its geometric applications. *Mathematische Zeitschrift*, 143(3):289–297, October 1975.
- [14] P. Erdős and A. Rényi. On the evolution of random graphs. *Magyar Tud. Akad. Mat. Kutató Int. Közl.*, 5:17–61, 1960.
- [15] Joel Friedman. A proof of Alon's second eigenvalue conjecture and related problems. *Memoirs of the American Mathematical Society*, 195(910):viii+100, 2008.
- [16] Clifford Gilmore, Etienne Le Masson, Tuomas Sahlsten, and Joe Thomas. Short geodesic loops and L^p norms of eigenfunctions on large genus random surfaces. *arXiv:1912.09961 [math-ph]*, December 2019.
- [17] Etienne Le Masson and Tuomas Sahlsten. Quantum ergodicity and Benjamini-Schramm convergence of hyperbolic surfaces. *Duke Mathematical Journal*, 166(18):3425–3460, 2017.

- [18] Maryam Mirzakhani. Growth of Weil-Petersson volumes and random hyperbolic surfaces of large genus. *Journal of Differential Geometry*, 94(2):267–300, 2013.
- [19] Jean-Pierre Otal and Eulalio Rosas. Pour toute surface hyperbolique de genre g , $\lambda_{2g-2} > 1/4$. *Duke Mathematical Journal*, 150(1):101–115, 2009.
- [20] Burton Randol. Small eigenvalues of the Laplace operator on compact Riemann surfaces. *Bulletin of the American Mathematical Society*, 80:996–1000, 1974.
- [21] Burton Randol. The Riemann hypothesis for Selberg’s zeta-function and the asymptotic behavior of eigenvalues of the Laplace operator. *Transactions of the American Mathematical Society*, 236:209–223, 1978.
- [22] Atle Selberg. Harmonic analysis and discontinuous groups in weakly symmetric Riemannian spaces with applications to Dirichlet series. *J. Indian Math. Soc. (N.S.)*, 20:47–87, 1956.
- [23] André Weil. On the moduli of Riemann surfaces. In *Œuvres Scientifiques - Collected Papers II:1951-1964*, pages 379–389. Springer-Verlag, Berlin Heidelberg, 1958.
- [24] Alex Wright. A tour through Mirzakhani’s work on moduli spaces of Riemann surfaces. *arXiv:1905.01753 [math]*, May 2019.

UNIVERSITÉ DE STRASBOURG, CNRS, IRMA UMR 7501, F-67000 STRASBOURG, FRANCE
E-mail address: monk@math.unistra.fr