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# A local multiple proper generalized decomposition based on the partition of unity

Rubén Ibáñez<sup>1</sup> | Emmanuelle Abisset-Chavanne<sup>1</sup> | Francisco Chinesta<sup>2</sup> | Antonio Huerta<sup>3</sup> | Elías Cueto<sup>4</sup> 

<sup>1</sup>ICI - High Performance Computing Institute, Ecole Centrale de Nantes, Nantes, France

<sup>2</sup>ESI Group Chair @ PIMM - Procédés et Ingénierie en Mécanique et Matériaux, Ensam ParisTech, Paris, France

<sup>3</sup>Laboratori de Calcul Numeric, Universitat Politècnica de Catalunya, BarcelonaTech, Barcelona, Spain

<sup>4</sup>Aragon Institute of Engineering Research, Universidad de Zaragoza, Zaragoza, Spain

## Correspondence

Francisco Chinesta, ESI Group Chair @ PIMM - Procédés et Ingénierie en Mécanique et Matériaux, Ensam ParisTech, 151 boulevard de l'Hôpital, Paris, France.  
Email: Francisco.Chinesta@ensam.eu

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## Summary

It is well known that model order reduction techniques that project the solution of the problem at hand onto a low-dimensional subspace present difficulties when this solution lies on a nonlinear manifold. To overcome these difficulties (notably, an undesirable increase in the number of required *modes* in the solution), several solutions have been suggested. Among them, we can cite the use of nonlinear dimensionality reduction techniques or, alternatively, the employ of linear local reduced order approaches. These last approaches usually present the difficulty of ensuring continuity between these local models. Here, a new method is presented, which ensures this continuity by resorting to the paradigm of the partition of unity while employing proper generalized decompositions at each local patch.

## KEYWORDS

local reduced order modeling, model order reduction, partition of unity, proper generalized decomposition

## 1 | INTRODUCTION

The finite element method (FEM) is the ubiquitous technique for the approximation of partial differential equations (PDEs). Its generality has led it to succeed in many areas of engineering interest. However, when real-time or many-query applications are envisaged, it is well known to be a technique somewhat slow. In recent years, model order reduction (MOR) techniques have shown that a minimum number of carefully-chosen degrees of freedom may be enough for an accurate solution of these same PDE. Instead of choosing general-purpose piecewise polynomials as basis functions, MOR techniques construct ad hoc basis functions following different techniques. For instance, proper orthogonal decomposition (POD)<sup>1-5</sup> constructs an efficient basis from a set of precomputed snapshots of the full-order PDE solution, ie,

$$u(x, t) = \sum_{i=1}^N \alpha_i(t) \phi_i(x), \quad (1)$$

where, very much like in the finite element context,  $\alpha_i$  are a set of time-dependent coefficients that evolve in time, and  $\phi_i(x)$  are time-independent basis functions obtained by some algebraic treatment of the system snapshots. For instance, many MOR techniques employ the most energetic eigenfunctions of the snapshot autocorrelation matrix to construct these  $\phi_i(x)$ . These play a similar role to the finite element (FE) shape functions, albeit they are global instead of local. Other techniques, such as reduced basis methods, for instance,<sup>6-8</sup> employ some snapshots of the full-order solution as basis for the approximate solution of the system. These snapshots are calculated in a greedy fashion, at time (or parameter) instants at which the error in the approximation is maximal.

It is important to note that Equation (1) constitutes, in fact, a separated expression of the solution (note the space-time separation, which also holds in FE approximations). Assuming this separate or affine decomposition of the solution is also on the origin of another family of MOR techniques coined as proper generalized decomposition (PGD).<sup>9-12</sup> These methods compute the basis function *on the fly*, that is to say, by means of a greedy algorithm that enriches successively the basis until the desired precision is achieved. This methodology has proven to be very effective for a wide variety of high-dimensional problems ranging from the resolution of Fokker-Planck equation<sup>13</sup> to patient-specific liver responses,<sup>14,15</sup> structural dynamics,<sup>16</sup> computational rheology,<sup>17</sup> or, more generally, to any parametric problem that could be written in separate form.<sup>13</sup>

There are situations, however, when the solution is highly nonseparable. In other words, the solution of the problem lives on a nonlinear manifold. In this situation, what MOR techniques do is to project the solution on the tangent space to the manifold at a given (time or parameter) point.<sup>18</sup> This leads to poor approximation properties far from the tangency point, unless special techniques are chosen, as in the work of Niroomandi et al,<sup>19</sup> for instance, where asymptotic expansions were used.

A way to alleviate this problem is to follow the same philosophy than nonlinear-dimensionality reduction methods. For instance, locally-linear-embedding<sup>20</sup> tries to unveil the latent variables by means of imposing a local linear variation on the function, which will change from neighborhood to neighborhood. Another technique that deals with nonlinear problems is the so-called kernel principal component analysis.<sup>21</sup> In this particular case, the definition of an efficient kernel function allows to project the snapshots to high-dimensional spaces (potentially infinite dimensional) in which the solution manifold is flat. In this particular situation, standard interpolation techniques work well.

Badías et al<sup>22</sup> studied the case of a moving source in a transient heat transfer problem. The problem of the nonseparability of the solution was circumvented by means of making a partition of the time domain, dedicating a different PGD for each partition. However, the imposition of the interface conditions will become a tedious task when dealing with partitions involving variables other than time.

The main objective of this paper is to develop a generalized PGD formulation in which continuity between subdomains is guaranteed. This will be achieved by resorting to the partition of unity (PU) paradigm.<sup>23,24</sup> By employing the PU, continuity of the solution is guaranteed if the chosen PU is continuous. This will allow us to glue different PGD approximations defined at particular regions of the space, time, or parameter spaces, thus ensuring a global solution with a minimum of degrees of freedom.

The structure of this paper is organized as follows. The second section depicts the general aspects of the proposed methodology; the third section shows the capability of the method to approximate highly nonlinear functions; the fourth section shows applications for different PDEs from fully diffusive to transient equations; and the fifth section will be devoted for the conclusions. All details related to the variational form will be given in the Appendix.

## 2 | BASICS OF THE METHOD

The method here proposed is based, as mentioned earlier, in the application of the PU paradigm.<sup>23,24</sup> In essence, the PU method states that, given a collection of nonoverlapping patches defined over the domain  $\Omega$ ,  $\Omega_i$ ,  $i = 1, \dots, n_{\text{patch}}$ , a PU  $\varphi_i$  defined on these patches (ie,  $\sum_i \varphi_i = 1$  everywhere in the domain), and function spaces  $V_i$ , associated to each patch, the approximation obtained by

$$V = \sum_{i=1}^{n_{\text{patch}}} \varphi_i V_i$$

For the sake of simplicity, but without losing generality, the local PGD method here presented is illustrated by starting with a weak form related to a general two-dimensional PDE, ie,

$$\mathcal{L}(u^*(x, y), u(x, y); \boldsymbol{\mu}) = F(u^*(x, y), f(x, y)) \text{ in } \Omega.$$

Normally, solving the weak form (2) requires an approximation space for the essential variable  $u(x, y)$  and its associated variation  $u^*(x, y)$ . The classical FE approximation, ie,

$$u(x, y) = \sum_{i \in \mathcal{I}} N_i(x, y) u_i,$$

where  $\mathcal{I}$  is the entire set of shape functions defined over the integration domain,  $\Omega$ , satisfies the PU and linear consistency properties,

$$\begin{aligned} \sum_{i \in \mathcal{I}} N_i(x, y) &= 1, & \forall x, y \in \Omega, \\ \sum_{i \in \mathcal{I}} N_i(x, y) x_i &= x, & \forall x, y \in \Omega, \\ \sum_{i \in \mathcal{I}} N_i(x, y) y_i &= y, & \forall x, y \in \Omega. \end{aligned}$$

The main ingredient of the method here proposed is the combination of the FE shape functions,  $N_i(x, y)$ , as an example of very convenient PU, enriched with low rank PGD approximations. Therefore, each degree of freedom  $u_i$  associated to the FEM shape function  $N_i(x, y)$  will be enriched with a PGD approximation, aiming to capture the details of the solution, which are not captured by the standard low-order FE meshes. Therefore, the approximation of the solution will read

$$u(x, y) = \sum_{i \in \mathcal{I}} N_i(x, y) \sum_{k=1}^M X_k^i(x) Y_k^i(y),$$

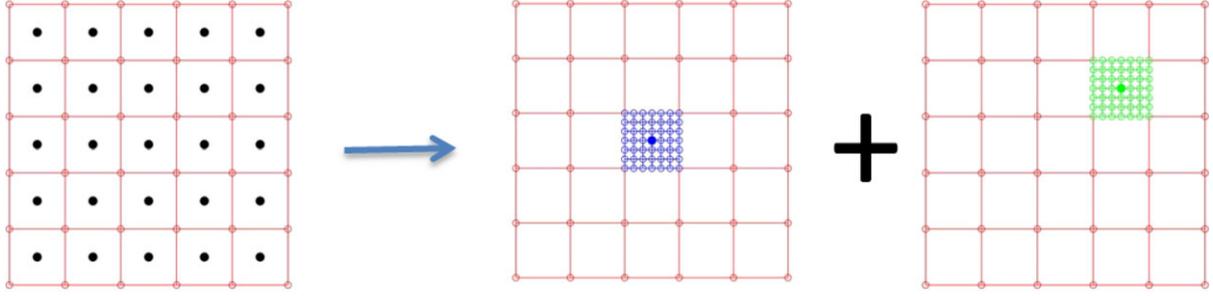
where  $X_k^i(x)$  and  $Y_k^i(y)$  functions are the  $k$ th one-dimensional modes related to the  $i$ th PGD.

Several approaches can be adopted to define the trial function. In this particular case, a Bubnov-Galerkin projection is selected. Hence, the same approximating space is chosen for  $u(x, y)$  and  $u^*(x, y)$ , ie,

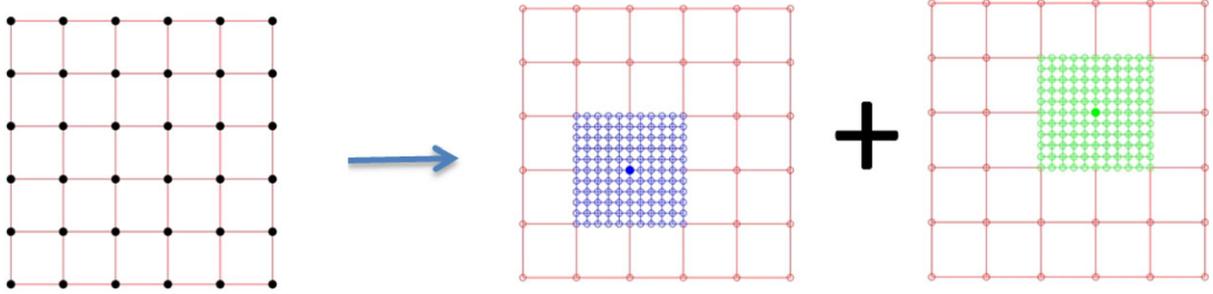
$$u^*(x, y) = \sum_{i \in \mathcal{I}} N_i(x, y) (X_M^{i*}(x) Y_M^i(y) + X_M^i(x) Y_M^{i*}(y)). \quad (2)$$

It is worth noting the greedy nature of the PGD algorithm, since the variation only takes into account the last,  $M$ th, PGD mode. Previous modes are considered as known, and therefore do not appear in Equation (2). Furthermore, a nonlinear problem, which has been created due to the separation of variables (note that we solve for a pair of functions,  $X_M, Y_M$ ), is solved using an alternate direction scheme. Conceptually speaking, it is not necessary to have the same number of modes,  $M$ , for each macroshape function in  $\mathcal{I}$ . A possible strategy would be to stop enriching with new modes in areas where the residual of the equation is low enough. In this first work, all PGDs are enriched with the same number of modes,  $M$ , as the main focus is placed on the methodology itself and proving that combining different PGDs within a continuous framework could be useful to obtain separated solutions in complex scenarios where standard PGD performance is decimated. Indeed, the definition of mode adopted herein is the number of separated functions (1 per element in the set  $\mathcal{I}$ ) that has to be given to the macropartition to reproduce the solution. In other words, the number of modes is equal to  $M$ . It is also important to reckon that the size of the mode will increase according to the number of elements in  $\mathcal{I}$ , since the mode involves as many separated functions per direction as number of degrees of freedom in  $\mathcal{I}$ .

The properties of the method will change depending on the way FE shape functions are defined. Figure 1 exemplifies the partition when piecewise constant shape functions are used. As it can be seen, no overlapping between different PGDs exists. Therefore, extra effort has to be done to set proper interface conditions along the red lines appearing in the same figure. This situation would be equivalent to the method described in the work of Badías et al.<sup>22</sup> This is in sharp contrast with the purely additive approach to couple FE and PGD proposed in the work of Ammar et al.<sup>25</sup> or in the work of Quesada et al.<sup>26</sup> for instance. In that case, the method resembled the  $s$ -version of the FEM<sup>27</sup> or the multiscale FEM proposed by Rank and Krause.<sup>28</sup> In that case, a single PGD approximation was enriched by FE to capture high-gradient



**FIGURE 1** Left, piecewise constant finite element approximation. Center, domain of influence of  $i$ th proper generalized decomposition (PGD). Right, domain of influence of  $k$ th PGD. As can be noticed, no overlapping exists between different PGD approximations



**FIGURE 2** Left, piecewise linear partition of unity. Center, domain of influence of  $i$ th proper generalized decomposition (PGD). Right, domain of influence of  $k$ th PGD

details of the solution, which usually provoke an increase in the number of PGD modes. Here, on the contrary, we develop an approximation by means of multiple PGD approximations, glued together by means of the PU paradigm. We have coined this method as *PU-PGD*.

Figure 2 depicts a piecewise linear partition of the domain. As it can be noticed, a quadrilateral element has contributions coming from four different PGDs. The coupling conditions between different PGDs are automatically taken into account in the elemental contributions. Indeed, this kind of partition imposes a smooth transition between PGDs, which vary continuously along the domain in accordance with the partition of the domain.

It is worth mentioning that the full potential of PGD approximation resides in the capability of writing the integral form in a separated manner. By doing that, all integrals in a high-dimensional space will be split into a product of integrals related to each one of the subspaces. For the sake of simplicity, we will assume that the shape functions  $N_i(x, y)$  are exactly separable in one single mode

$$N_i(x, y) = N_i^x(x)N_i^y(y).$$

It is important to notice that this will be the case, for instance, of 2D quadrilateral elements with straight and parallel sides.

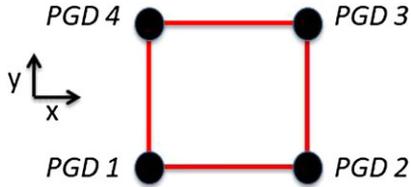
Even though this assumption may be seen as a limitation, it is important to notice that there are some variables, such as time or parameters, that naturally admit cartesian decompositions. Therefore, when dealing with a complex spatial geometry that evolves in time, a smart partition would be to keep the space variable as is. Nevertheless, this very first work is meant to provide an insight of the method, thus, all results presented in the sequel will be related to piecewise linear FE shape functions based on squared elements.

### 3 | *PU-PGD* FOR APPROXIMATION PROBLEMS

#### 3.1 | Basics of the method

In this section, we analyze the capability of the proposed methodology in approximation problems of the form

$$\int_{\Omega} u^*(x, y) (u(x, y) - f(x, y)) d\Omega = 0, \quad (3)$$



**FIGURE 3** In the case of bilinear finite elements, there is one proper generalized decomposition (PGD) enrichment per node [Colour figure can be viewed at [wileyonlinelibrary.com](http://wileyonlinelibrary.com)]

where a function  $f(x, y)$  is to be approximated in an  $PU$ -PGD framework. As it can be noticed, there are no partial derivatives involved in this type of problems. Instead, just a compact expression of the function  $f(x, y)$  is sought. It is important to highlight that continuity is not a requirement, given the fact that no differential equation is involved in this particular case. Indeed, considering a discontinuous approximation space with a much more local procedure will provide better results if the function to be captured presents high gradients or even discontinuities. However, this very first example is included as an initial point to illustrate the continuous approach, which is mandatory to solve PDEs.

Following the spirit of standard FE approximation, the integral over the entire domain  $\Omega$  is split into a sum of integrals for each one of the FEs  $\Omega_e$  appearing in the domain. Recall that piecewise bilinear shape functions present contributions from four different PGDs per element (one for each corner), as shown in Figure 3.

It is important to notice that the  $i$ th PGD, thanks to the PU, affects the support of each FE node, and thus four elements in a regular lattice.

Omitting the dependence with respect to space variables,  $x$  and  $y$ , and thanks to the separability of the FE shape function (if defined in a regular lattice as a tensorial product of one-dimensional functions), the first term in the integral form of Equation (3) particularized for a element  $\Omega_e$  reads

$$\int_{\Omega_e} u^* u dx dy = \sum_{i=1}^4 \sum_{j=1}^4 \int_{\Omega_e} N_i^x N_i^y (X_M^i Y_M^i)^* \left( N_j^x N_j^y \sum_{k=1}^M X_k^j Y_k^j \right) dx dy.$$

The next step is to split the 2D integral form into a tensorial product of 1D integrals as

$$\int_{\Omega_e} u^* u dx dy = \sum_{i=1}^4 \sum_{j=1}^4 \sum_{k=1}^M \int_x X_M^{i*} N_i^x N_j^x X_k^j dx \int_y Y_M^i N_i^y N_j^y Y_k^j dy + \sum_{i=1}^4 \sum_{j=1}^4 \sum_{k=1}^M \int_x X_M^i N_i^x N_j^x X_k^j dx \int_y Y_M^{i*} N_i^y N_j^y Y_k^j dy. \quad (4)$$

As can be noticed, the variation along the  $x$  direction shares all the operators with the variation along the  $y$  direction. Indeed, this property arises naturally from the alternate directions scheme that is used to solve the nonlinear system of the PGD.

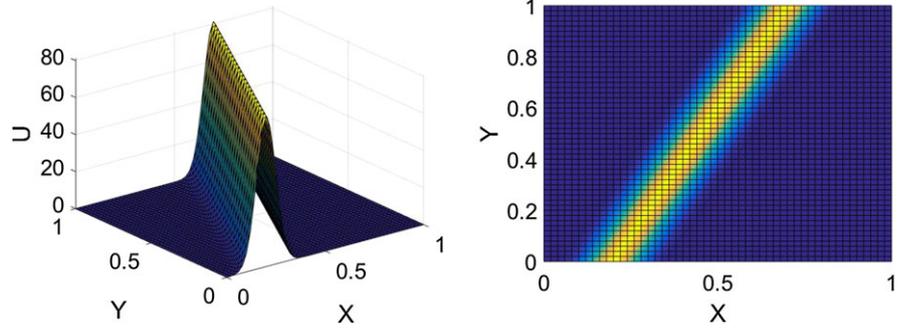
With the source term  $f(x, y)$  expressed in a separated format (eg, by invoking the SVD) as

$$f(x, y) = \sum_{z=1}^Z f_z^x(x) f_z^y(y),$$

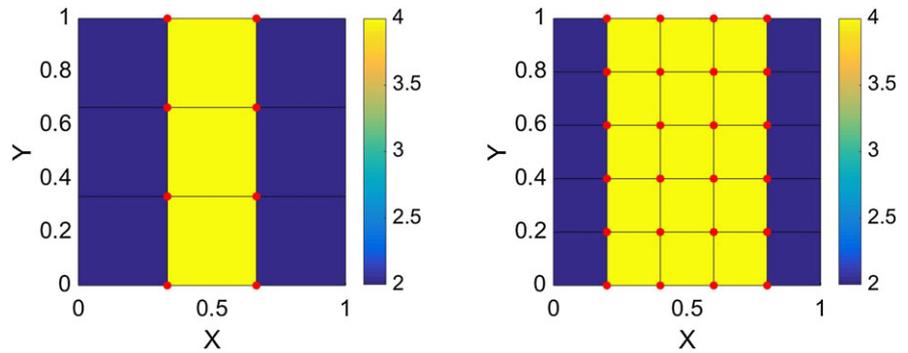
the right-hand side term appearing in Equation (3) reads

$$\int_{\Omega_e} u^* f dx dy = \sum_{i=1}^4 \sum_{z=1}^Z \int_x X_M^{i*} N_i^x f_z^x dx \int_y Y_M^i N_i^y f_z^y dy + \sum_{i=1}^4 \sum_{z=1}^Z \int_x X_M^i N_i^x f_z^x dx \int_y Y_M^{i*} N_i^y f_z^y dy,$$

for  $x$  and  $y$  systems, respectively.



**FIGURE 4** Left, 3D view of  $f_1(x, y)$ . Right, top view of  $f_1(x, y)$  [Colour figure can be viewed at [wileyonlinelibrary.com](http://wileyonlinelibrary.com)]



**FIGURE 5** Left, domain enriched with 8 proper generalized decompositions (PGDs). Right, domain enriched with 24 PGDs. Red points, centroid of each PGD. The legend shows the number of PGD enrichments affecting each element

### 3.2 | A preliminar example

Let us illustrate the methodology by analyzing a highly nonseparable function

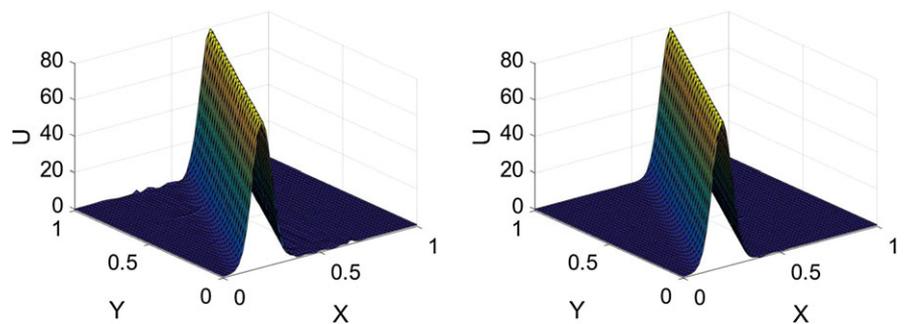
$$f_1(x, y) = \frac{10}{\sigma \sqrt{2\pi}} e^{-\frac{(x-(v+y+x_0))^2}{2\sigma^2}},$$

where  $\sigma = 0.05$ ,  $v = 0.5$ , and  $x_0 = 0.2$ .

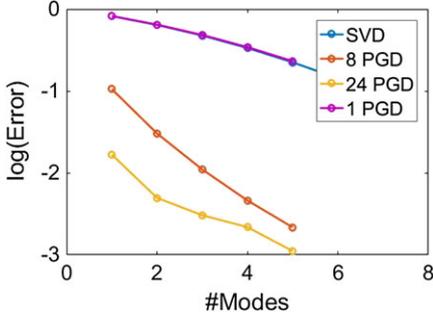
Figure 4 depicts the resulting  $u(x, y) = f_1(x, y)$  scalar field from different perspectives. As it can be seen, the source term is going through the diagonal of the domain, generating a highly nonseparable function. Indeed, a standard PGD algorithm encounters many problems to capture such kind of solution.

Two different partitions of the domain have been tested, as shown in Figure 5. Along the right and left sides of the domain, no enrichment is added to the FE approximation since the sought solution is already zero along this boundary. It is important to reckon that right and left sides of the domain could have been enriched, since no boundary conditions need to be imposed in this approximation problem. However, the fact that compatible boundary conditions (ie, zero along the right and left sides) can be imposed was used as a validation point for the next numerical examples involving PDEs. The elements are colored in accordance with the number of active PGDs acting in each element.

Figure 6 shows the reconstructed solution using 4 modes per local PGD, when the domain has 8 active PGDs (left) and 24 PGDs (right). As it can be seen, even though both approximations capture the main features of the solution, the solution using 24 PGDs is slightly better than the one with 8 PGDs.



**FIGURE 6** Left, solution with 8 proper generalized decomposition (PGDs). Right, solution with 24 PGDs. Every PGD enrichment incorporates four modes [Colour figure can be viewed at [wileyonlinelibrary.com](http://wileyonlinelibrary.com)]



**FIGURE 7** Convergence error for the approximation problem. Comparison between singular value decomposition (SVD), standard proper generalized decomposition (PGD), and two different local strategies

Figure 7 shows the convergence error for the approximation of  $f_1(x, y)$ . The relative error is measured as

$$\mathcal{E}^m = \frac{\|u_{ref} - u^m\|}{\|u_{ref}\|},$$

where  $u^m$  is the reconstructed solution using  $m$  modes.  $u_{ref}$  represents the ground truth. As can be noticed, singular value decomposition (SVD) method suffers from separating this kind of solutions, having a slow decay of the relative error with respect to the number of modes. As it is well known from the literature, standard PGD in two dimensions is exactly equivalent to SVD when an identity operator is employed (the so-called PGD in approximation<sup>29,30</sup>). On the other hand, using the *PU*-PGD algorithm, the relative error decays much faster than the one related to standard SVD or PGD approximations.

It is important to notice that the more a domain is partitioned, the smaller the relative error is. This behavior is expected since the smaller the elements are, the easier it is to capture a local behavior. However, the price to pay is that the cost of computing one mode per PGD will increase just like the storage cost as well. Certainly, systems solved in the fixed point iteration are not local due to the fact that macroshape functions are overlapped. If the macropartition involves  $\#I$  shape functions and each local PGD is discretized using  $n_x$  and  $n_y$  degrees of freedom, the number of unknowns in each PGD fixed point are  $\#I \cdot n_x$  and  $\#I \cdot n_y$  for  $x$  and  $y$  directions, respectively. Therefore, it is important to keep in mind that refining the macromesh would reduce the number of modes  $M$  to represent the solution, but it would also increase the computational cost for obtaining one mode. However, the main focus of this work was to prove that the proposed methodology was able to obtain separated solutions in scenarios where standard PGD performance was not good enough.

## 4 | *PU*-PGD FOR THE SOLUTION OF PARTIAL DIFFERENTIAL EQUATIONS

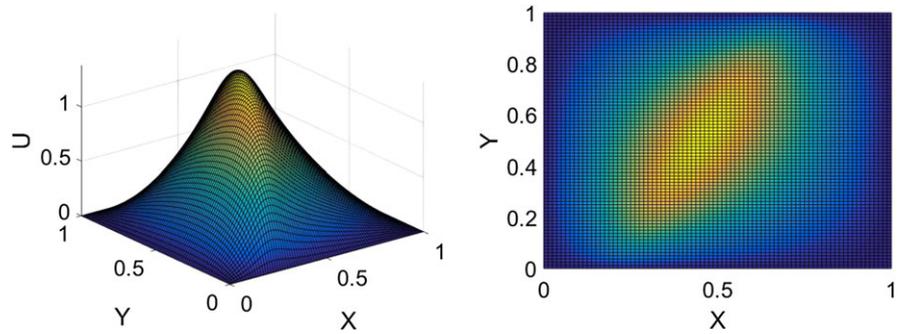
The aim of this section is to show the potential of the proposed methodology when applied to the solution of different PDEs. We consider examples of increasing complexity. Therefore, the first test case is a pure diffusion equation. The second case is a convection-reaction-diffusion equation. Finally, the last example relates to a transient thermal problem with a moving source term. Even though the construction of the operators for each PDE follows the same trend than the ones presented in the *PU*-PGD in approximation, further details are given in the appendix.

### 4.1 | Diffusion PDE

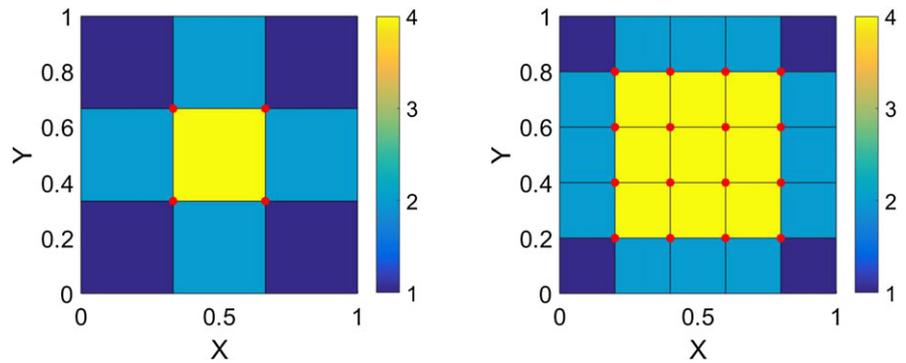
We first consider the following weak form corresponding to a diffusion problem:

$$\int_{\Omega} \eta \nabla u^* \cdot \nabla u dx dy = \int_{\Omega} u^* f_1 dx dy.$$

Vanishing Dirichlet boundary conditions are imposed on the entire boundary. The source term is equal to the function  $f_1(x, y)$  defined in the previous section and the diffusion parameter,  $\eta$ , is set to 1.



**FIGURE 8** Left, 3D view of  $u(x,y)$ . Right, planar view of  $u(x,y)$  for the diffusion equation [Colour figure can be viewed at [wileyonlinelibrary.com](http://wileyonlinelibrary.com)]



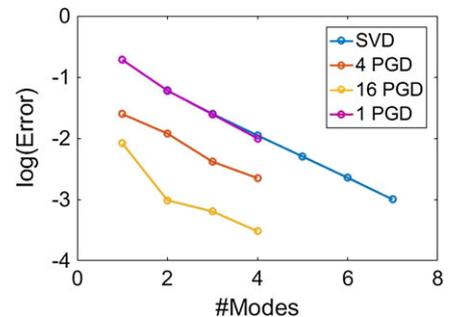
**FIGURE 9** Left, problem solved with four proper generalized decomposition (PGDs). Right, problem solved with 16 PGDs. Red points represent finite element nodes enriched with PGD. The legend represents the number of PGDs per element

Figure 8 depicts the reference  $u(x,y)$  scalar field for the diffusion equation from different perspectives. It has been obtained with a  $80 \times 80$  FE discretization. As it can be seen, the solution diffuses the source term through the diagonal of the domain.

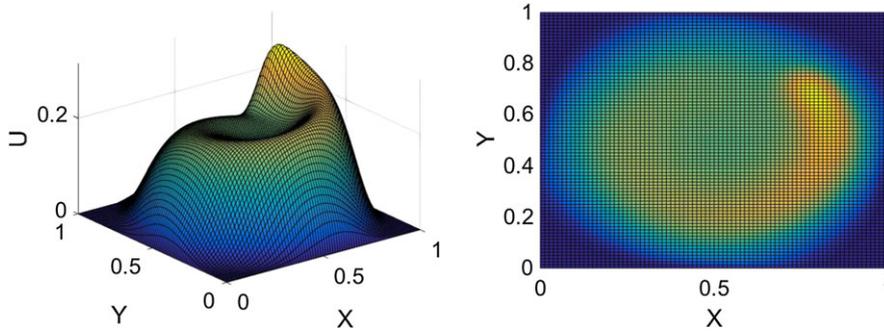
Two different PGD enrichments have been tested, as shown in Figure 9. The PGD enrichments acting on top, bottom, left, and right sides are set to zero to be consistent with the problem statement, where zero Dirichlet boundary conditions are imposed along the entire boundary domain. The red points indicate the enriched nodal supports. The elements are colored in accordance with the number of enriching PGDs. Elements on the corners have only one active PGD, whereas elements on the sides have two active PGDs and interior elements have four active PGDs.

The reconstructed solution using four modes per PGD enrichment, when the domain has four active PGDs (left) and sixteen PGDs (right), is in perfect visual agreement with the reference solution appearing in Figure 9.

Figure 10 shows the convergence error for the diffusive equation. As it can be noticed, all methods converge monotonically, SVD being the one that presents the slower convergence with respect the number of modes. Standard PGD shows a similar convergence rate to SVD, as expected. Even if no formal proof exists of this behavior for operators other than the identity, this is the observed rate for all the experiments considered so far. Finally, using the multi-PGD algorithm presents a faster decay in the relative error.



**FIGURE 10** Convergence error for the diffusion problem. Comparison between the number of modes provided by singular value decomposition (SVD), standard proper generalized decomposition (PGD), and two different enrichment strategies considered so far



**FIGURE 11** Left, 3D view of the reference solution  $u(x, y)$ . Right, planar view of  $u(x, y)$  for convection-reaction-diffusion equation [Colour figure can be viewed at wileyonlinelibrary.com]

## 4.2 | Convection-reaction-diffusion PDE

In this subsection a convection-reaction-diffusion equation is studied. Indeed, the major novelty is the introduction of the convective term since it will involve a nonsymmetric operator. The weak form related to this equation is

$$\int_{\Omega} u^* \mathbf{v} \cdot \nabla u + \eta \nabla u^* \cdot \nabla u + \sigma u^* u \, dx dy = \int_{\Omega} u^* f_2 \, dx dy, \quad (5)$$

where vanishing Dirichlet boundary conditions are considered on the entire boundary. The integration domain is  $\Omega = [0, 1] \times [0, 1]$ . The convective velocity field is assumed to be  $\mathbf{v} = 500(y - 0.5, 0.5 - x)^T$ . The reaction term is set to  $\sigma = 10$  and the diffusion coefficient to  $\eta = 1$ . Regarding the source term, a Gaussian located in  $[x_0, y_0] = [0.75, 0.75]$  as

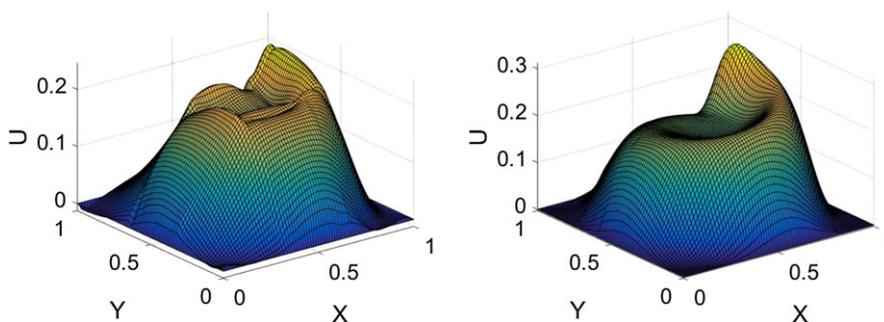
$$f_2(x, y) = \frac{800}{\sqrt{2\pi}} e^{-\frac{(x-x_0)^2 + (y-y_0)^2}{0.005}}$$

is considered.

Figure 11 depicts the reference  $u(x, y)$  scalar field for the solution of Equation (5) (computed by employing an  $80 \times 80$  FE mesh, thus with no need for stabilization) from different perspectives. Note how the Gaussian source term is convected circularly according to the prescribed velocity field. Moreover, the diffusion term makes the Gaussian to be smoothed along its convection path.

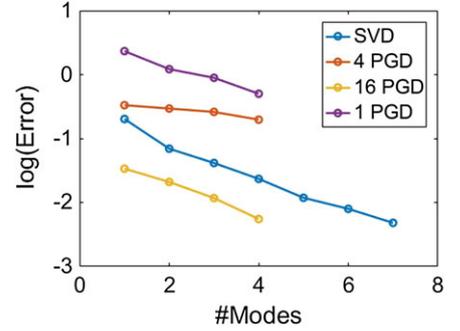
The same partitions than in the diffusive case have been tested, as shown in Figure 9, since also vanishing Dirichlet boundary conditions are imposed along the entire boundary of the domain. Figure 12 shows the reconstructed solution using 4 modes per local PGD, when the problem has been approximated by 4 PGDs (left) and 16 PGDs (right). Notice how the approximation related to 4 PGDs does not capture the solution well. Indeed, a higher number of modes will be needed to converge to the reference solution. On the other hand, the solution involving 16 PGDs is in perfect accordance with the reference solution, capturing all the features of the scalar field without any perceptible oscillation.

Figure 13 shows the convergence of the relative error in logarithmic scale for the SVD and the two different partitions proposed for the CDR equation. It can be clearly seen that the convergence related to the 4 PGD partition is the slowest one. Indeed, this problem is derived from the fact that the first PGD modes have to capture a highly nonlinear behavior, meaning that this partition will require a high amount of modes to converge to the real solution. On the other hand, the 16 PGD partition convergence is quite fast compared with the other two methods.

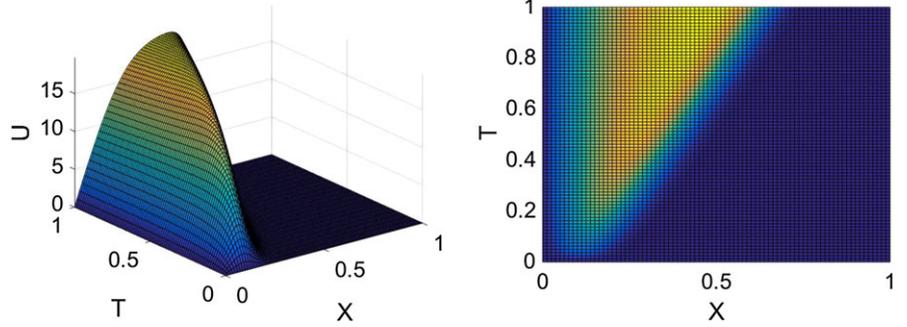


**FIGURE 12** Solution of the convection-diffusion-reaction equation. Left, 4 proper generalized decomposition (PGDs). Right, 16 PGDs. Both solutions have four modes per PGD [Colour figure can be viewed at wileyonlinelibrary.com]

**FIGURE 13** Convergence plot for the convection-diffusion-reaction problem. Comparison between singular value decomposition (SVD), standard proper generalized decomposition (PGD), and two different enrichment strategies



**FIGURE 14** Left, 3D view of  $u(x, t)$ . Right, top view of  $u(x, t)$  for the transient heat transfer equation [Colour figure can be viewed at wileyonlinelibrary.com]



### 4.3 | Transient heat transfer equation

This last example is devoted to the analysis of a transient 1D heat transfer problem, a test proposed by Idelsohn<sup>31</sup> to verify the compactness of different MOR techniques, ie,

$$\int_{\Omega} u^* \frac{\partial u}{\partial t} + \eta \frac{\partial u^*}{\partial x} \frac{\partial u^*}{\partial x} dx dt = \int_{\Omega} u^* f_3 dx dt,$$

where the diffusion coefficient is set to  $\eta = 0.01$ .

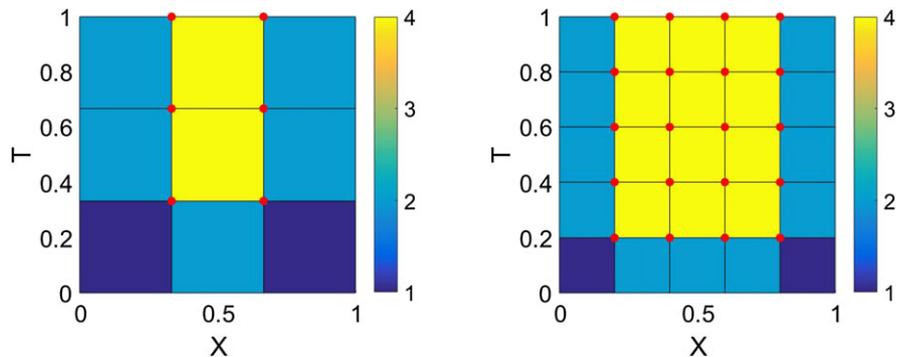
The domain of study is  $\Omega = \Omega_x \times \Omega_t = [0, 1] \times [0, 1]$ . The set of boundary conditions imposed for the field  $u(x, t)$  is  $u(x, 0) = 0$ ,  $u(0, t) = 0$ , and  $u(1, t) = 0$ . The source term  $f_3(x, y)$  is

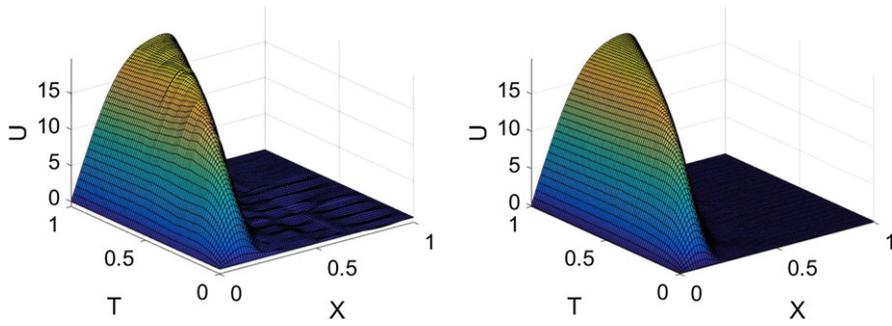
$$f_3(x, t) = \frac{10}{\sigma \sqrt{2\pi}} e^{-\frac{(x-(vt+x_0))^2}{2\sigma^2}},$$

where  $\sigma = 0.05$ ,  $v = 0.5$ , and  $x_0 = 0.1$ .

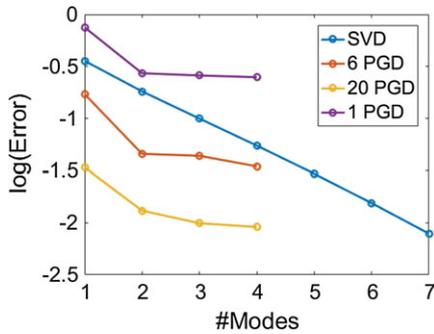
Figure 14 depicts the reference  $u(x, y)$  scalar field for the transient heat equation from different perspectives. Again, an  $80 \times 80$  FE mesh of bilinear elements has been employed. This kind of PDE creates a boundary layer along the diagonal of the space-time domain, which is very hard to capture when using standard separate approximations like POD and classical PGD.

**FIGURE 15** Left, problem solved with 6 proper generalized decompositions (PGDs). Right, problem solved with 20 PGDs. Red points represent the finite element nodes that incorporate an enrichment. The legend indicates the number of PGDs per element





**FIGURE 16** Solution of the transient heat equation. Left, 6 proper generalized decomposition (PGDs). Right, 20 PGDs. Both solutions have four modes per PGD [Colour figure can be viewed at [wileyonlinelibrary.com](http://wileyonlinelibrary.com)]



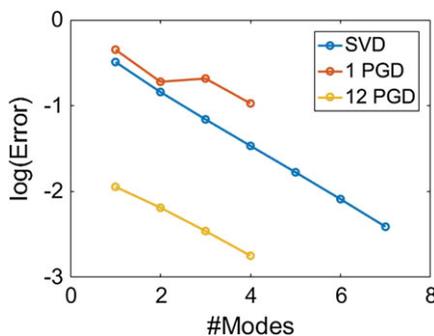
**FIGURE 17** Convergence error for transient heat problem in logarithmic scale for different number of modes. Comparison between singular value decomposition (SVD), standard proper generalized decomposition (PGD), and two different partitions of the domain

Two different enrichment strategies have been tested, as shown in Figure 15. It is important to highlight that PGDs acting on  $x = 0$ ,  $x = 1$ , and  $t = 0$  are set to zero to satisfy the null Dirichlet and initial boundary conditions in this portion of the boundary. However, the PGD acting on the interior nodes of the line  $t = 1$  is not set to zero. This fact is a direct consequence of the pure convective behavior associated to the time variable. Indeed, the line  $t = 1$  acts like an outflow boundary, thus, no condition should be imposed there to ensure the well-posedness of the problem.

Figure 16 shows the reconstructed solution using four modes per local PGD, when the domain has 6 active PGDs (left) and 20 PGDs (right). Notice how both approximations capture the main features of the reference solution. However, the approximation with 6 active PGDs presents small oscillations due to the lack of PGD modes to converge to the reference solution.

Figure 17 shows the convergence plot for the SVD and the two different enrichments proposed for the transient heat equation. It can be stated that both partitions present a faster convergence than the SVD convergence. However, the convergence slope of the PGD partitions tend to be flatter than the SVD one in the last part of the convergence plot.

To establish a fair comparison between the methods presented so far, it is interesting to perform a comparison on a per-degree-of-freedom basis. Figure 18 shows the convergence plot related to the transient thermal problem in which the standard PGD approximation has the same degrees of freedom than the *PU*-PGD approximation. It is worth noting the fast convergence of the *PU*-PGD, which comes from the fact that it takes advantage of the locality of the solution, avoiding problem of highly nonlinear global functions.



**FIGURE 18** Convergence plot for Idelsohn problem using different partitions of the domain. All approximations have the same number of degrees of freedom. PGD, proper generalized decomposition; SVD, singular value decomposition

## 5 | DISCUSSION. FUTURE WORK

The method presented overcomes some of the most important limitations of standard PGD approximations. These, which are indeed common to every MOR technique, are related to the nonlinearity of the solution map. As in previous works, in the literature, an appealing strategy is to develop local approximations. This approach comes from the well-known fact that a manifold is flat in the neighborhood of every point.

The local multiple approximations here developed have shown to reconstruct the solution up to a certain tolerance by using a moderate number of modes. However, as previously stated, the size of the mode will increase according to the macropartition of the domain. Therefore, the choice of the macropartition constitutes one of the central key points of this approach. A macropartition that minimizes the number of macroelements while it reduces the number of modes to represent the solution would be the optimal one. Certainly, a possible future route is to mimic the element refinement in FEM, starting with a coarse macropartition and perform only macrorefinement in areas where the residual of the equation is high enough.

One drawback, however, of the proposed technique is related to the curse of dimensionality in high-dimensional problems. The *PU*-PGD approach introduces a mesh to construct the PU. This may result in prohibitive computational costs if the number of dimensions of the solution grows.

However, as in previous local approaches of PGD,<sup>22</sup> there is no need to partition every dimension into elements. Instead, only those parameter interval of interest that cause the most nonlinear response of the solution should be partitioned. Knowing in advance the nonlinear shape of the solution manifold is by no means straightforward. This constitutes our current effort of research.

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### ORCID

Elias Cueto  <https://orcid.org/0000-0003-1017-4381>

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## APPENDIX

### ELEMENTAL OPERATORS

In this appendix, we give exhaustive details about the construction of the operators required to make the *PU*-PGD approximation for the different PDEs studied in this paper. Detailed information was given for the case of the *PU*-PGD in approximation. Hence, further details will be shown related to the diffusive and convective terms.

#### A.1 | Diffusive operator

In this section, a diffusion of a scalar field  $u(x, y)$  along the  $x$  direction is developed. The derivation for the diffusion along the  $y$  direction follows similar guidelines. For the sake of simplicity, but without losing generality, we will also assume

$$\begin{aligned}
& \int_{\Omega_e} \frac{\partial u^*}{\partial x} \frac{\partial u}{\partial x} dx dy = \\
& \sum_{i=1}^4 \sum_{j=1}^4 \sum_{k=1}^M \int_{\Omega_e} \left( \frac{\partial N_i}{\partial x} X_M^{i*} Y_M^i + N_i \frac{\partial X_M^{i*}}{\partial x} Y_M^i \right) \left( \frac{\partial N_j}{\partial x} X_k^j Y_k^j + N_j \frac{\partial X_k^j}{\partial x} Y_k^j \right) dx dy = \\
& = \sum_{i=1}^4 \sum_{j=1}^4 \sum_{k=1}^M \int_x X_M^{i*} \frac{\partial N_i^x}{\partial x} \frac{\partial N_j^x}{\partial x} X_k^j dx \int_y Y_M^i N_i^y N_j^y Y_k^j dy + \\
& + \sum_{i=1}^4 \sum_{j=1}^4 \sum_{k=1}^M \int_x \frac{\partial X_M^{i*}}{\partial x} N_i^x \frac{\partial N_j^x}{\partial x} X_k^j dx \int_y Y_M^i N_i^y N_j^y Y_k^j dy + \\
& + \sum_{i=1}^4 \sum_{j=1}^4 \sum_{k=1}^M \int_x \frac{\partial X_M^{i*}}{\partial x} N_i^x N_j^x \frac{\partial X_k^j}{\partial x} dx \int_y Y_M^i N_i^y N_j^y Y_k^j dy + \\
& + \sum_{i=1}^4 \sum_{j=1}^4 \sum_{k=1}^M \int_x X_M^{i*} \frac{\partial N_i^x}{\partial x} N_j^x \frac{\partial X_k^j}{\partial x} dx \int_y Y_M^i N_i^y N_j^y Y_k^j dy. \tag{A1}
\end{aligned}$$

It can be highlighted that four different contributions are appearing due to the application of the chain rule. Furthermore, all terms appearing in the last part of Equation (A1) present already a separated format.

## A.2 | Convective operator

In this section, a pure convection of a scalar field  $u(x,y)$  along the  $x$  direction is derived. It is important to notice that this case is a particular case of a general convection given by the velocity field  $\mathbf{v}$ , where a separated representation of the velocity field would be required. The derivation for the convection along the  $y$  direction follows similar guidelines. For the sake of simplicity, but without losing generality, we will also assume that a new mode in the  $x$  direction is desired. Therefore, all variations related to the  $y$  direction are set to zero, ie,

$$\begin{aligned}
& \int_{\Omega_e} u^* \frac{\partial u}{\partial x} dx dy = \\
& \sum_{i=1}^4 \sum_{j=1}^4 \sum_{k=1}^M \int_{\Omega_e} N_i X_M^{i*} Y_M^i \left( \frac{\partial N_j}{\partial x} X_k^j Y_k^j + N_j \frac{\partial X_k^j}{\partial x} Y_k^j \right) dx dy = \\
& = \sum_{i=1}^4 \sum_{j=1}^4 \sum_{k=1}^M \int_x X_M^{i*} N_i^x \frac{\partial N_j^x}{\partial x} X_k^j dx \int_y Y_M^i N_i^y N_j^y Y_k^j dy + \\
& + \sum_{i=1}^4 \sum_{j=1}^4 \sum_{k=1}^M \int_x X_M^{i*} N_i^x N_j^x \frac{\partial X_k^j}{\partial x} dx \int_y Y_M^i N_i^y N_j^y Y_k^j dy. \tag{A2}
\end{aligned}$$

Once again all integrals appearing in the last part of Equation (A2) are written already in a separated way, improving the efficiency of the algorithm.