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Abstract

The parallel finite-element solution of large-scale time-harmonic wave problems is addressed with a non-overlapping optimized Schwarz domain decomposition method (DDM). It is well-known that the efficiency of this kind of method strongly depends on the transmission condition enforced on the interfaces between the subdomains. Local conditions based on high-order absorbing boundary conditions (HABCs) have proved to be well-suited, as a good compromise between basic impedance conditions, which lead to suboptimal convergence, and conditions based on the exact Dirichlet-to-Neumann (DtN) map related to the complementary of the subdomain — which are too expensive to compute. However, a direct application of the approach for configurations with interior cross-points (where more than two subdomains meet) and boundary cross-points (points that belong to both the exterior boundary and at least two subdomains) is suboptimal and, in some cases, can lead to incorrect results.

In this work, we extend a non-overlapping DDM with HABC-based transmission conditions approach to efficiently deal with cross-points for lattice-type partitioning. The proposed cross-point treatment relies on corner conditions developed for Padé-type HABCs. Two-dimensional numerical results with a nodal finite-element discretization are proposed to validate the approach, including convergence studies with respect to the frequency, the mesh size and the number of subdomains. These results demonstrate the efficiency for settings with regular partitions and homogeneous media. Numerical experiments with non-regular partitions and smoothly varying heterogeneous media show the robustness of the approach.

1 Introduction

The efficient and accurate solution of high-frequency time-harmonic wave scattering problems remains a challenging issue in computational engineering. Indeed, first, the unbounded domain must be truncated, which requires a specific treatment to simulate the outward propagation of scattered waves far from the obstacle. Then, the discretization of formulations related to Helmholtz-type problems leads to very large, complex and indefinite (dense or sparse) linear systems, especially in the high-frequency regime, corresponding to the situation where the wavelength is small compared with the characteristic size of the scatterer.

Various approaches can be designed for solving high-frequency scattering problems. Among the most popular ones, the boundary element method (BEM) based on the discretization of an integral equation [18, 60], in conjunction with a preconditioned Krylov subspace iterative solver [1, 19] and fast compression algorithms of integral kernels [14, 15, 64], is a first direction. An alternative for engineering applications is to introduce a boundary condition on a fictitious surface or an absorbing boundary layer enclosing the scatterer to truncate the domain and then to discretize the associated variational formulation in the bounded volume domain *e.g.* using a finite element method (FEM). Many possibilities exist to bound the computational domain. Basically, there is always a trade-off between accuracy and computational effort when choosing one of the truncation strategies. The most basic absorbing boundary condition (*i.e.* Sommerfeld's condition) is easy to use but is not very accurate. Non-local non-reflective boundary conditions [45] and BEM-FEM coupling [23] have been proposed, which give the perfect accuracy but yield expensive additional costs due to their nonlocal nature, leading to dense matrix blocks into the discrete weak formulation. As a good compromise between accuracy and computational cost, local high-order absorbing boundary conditions (HABCs) [3, 31, 48] and perfectly matched layers (PMLs) [8, 9, 72] provide high accuracy (at least for problems with homogeneous media) at the price of a larger number of unknowns and associated computational cost. The accuracy and the cost can be controlled by choosing the order of the HABC or the thickness of the PML. Usually, PMLs are easier to implement than HABCs, but the selection of the tuning parameters of a HABC is simpler to manage. In this article, we consider the HABC developed in [58].

Once the weak formulation is discretized thanks to the FEM with HABCs, the sparse linear system remains to be solved. It is well-known that this issue is still problematic since the sparse complex-valued linear system is very large and highly indefinite, most particularly for high frequencies. A direct solver cannot be used and standard Krylov iterative solvers are extremely difficult to make converge, even with the most advanced preconditioners. A natural alternative solution which has been introduced more than 30 years ago is to use an iterative/direct hybrid approach based on a domain decomposition of the large global computational domain, and iterate between the subdomains where the local subproblems are solved in parallel by a direct solver. Tremendeous efforts have been made to develop efficient domain decomposition methods (DDMs) with good rate of convergence (see *e.g.* [27, 63, 71] for general introductions). For Helmholtz-type problems, we can mention Schwarz methods with overlap [16, 41, 50] or without overlap [6, 20, 37], FETI algorithms [24, 34–36] and the method of polarized traces [75], which are eventually combined with preconditioning techniques (see *e.g.* [22, 43, 46, 66, 67, 73]). A recent overview of these methods has been proposed by Gander and Zhang [42].

In this work, we investigate a domain decomposition approach with non-overlapping subdomains, which minimizes the data transfer between subdomains. It is admitted that the convergence rate of the non-overlapping DDMs strongly depends on the transmission condition enforced on the interfaces between the subdomains. The optimal transmission operator corresponds to the nonlocal Dirichlet-to-Neumann (DtN) map related to the complementary of the subdomain, which is a Schur complement at the discrete level. Since the cost of computing the exact DtN is prohibitive, strategies based on approximate DtN operators started to be investigated in the late 80's and early 90's (see *e.g.* [47, 59]). For Helmholtz-type problems, Després [6, 25] used Robin-type transmission conditions and proved that the iterative DDM converges. The Robin-type operator is a coarse approximation of the exact DtN operator which is sometimes used as a basic absorbing boundary condition (Sommeferld's ABC). Improved Schwarz methods with optimized transmission conditions based on modified second-order transmission operators have next been introduced in [37, 62]. In parallel, FETI methods were adapted to Helmholtz problems as FETI-H [24, 34] and FETI-DPH [36] techniques, which can also be interpreted as optimized Schwarz DDMs. Later, optimized Schwarz methods with HABCs were

developed to improve the convergence rate and robustness of the methods [12, 13, 49, 54], as well as PML-based approaches [4, 65, 66, 73] and nonlocal transmission conditions [21, 51, 68]. As for ABCs, transmission boundary conditions related to HABCs and PML represent a good compromise between the basic impedance conditions (which lead to suboptimal convergence) and the exact Dirichlet-to-Neumann (DtN) map related to the complementary of the subdomain (which is expensive to compute).

In the perspective of large-scale applications, the DDMs must be applicable with domain partitions having interior cross-points (where more than two subdomains meet) and boundary cross-points (that belong to both the exterior boundary and at least two subdomains). The cross-points require special care at both continuous and discrete levels. For Robin-type operators, the convergence is proved for the continuous case and mixed finite element discretizations in [6, 25]. These proofs are extended for more general abstract operators in [20]. Several approaches have been proposed for nodal finite element discretizations. In the FETI framework, dual methods, which [24, 34] are examples, lead to saddle point problems with constraints associated to the cross-points. These constraints can be redundant (see *e.g.* [27, 63, 71] for more details). The direct discretizations of the optimized Schwarz methods described in [37, 62] belong to this category. In the context of elliptic problems, several non-standard discretization at the cross-points have been investigated in [38–40, 53]. With primal-dual methods, such as FETI-DPH [36] and the method proposed in [7, 11], global variables are associated to the physical fields at the cross-points. When combined with a preconditioning technique, this approach improves the convergence, but it requires all-to-all communications to solve a global system, which could deteriorate the parallel scaling of the method. Recently, a cross-point treatment has been investigated for second-order transmission operators [26, 61]. PML-type operators have been tested in configurations with cross-points, but only in the context of DDM preconditioning [4, 52, 65]. The treatment of interior cross-points for optimized Schwarz solvers with HABC-based, PML-based and non-local transmission conditions is a complicated problem which has, to the best of our knowledge, not been carefully addressed.

In the present paper and following [13], we consider an optimized Schwarz DDM with a transmission condition based on a Padé-type HABC operator. Nevertheless, in [13], no specific cross-point treatment was used on both interior and boundary cross-points. In addition, only first- and second-order exterior ABCs were considered. Here, we address the question of the cross-point treatment when the HABC operator is used in the transmission condition, or when it is used in the exterior boundary condition, or both. By contrast with the works reviewed in the previous paragraph, a specific care is required already at the continuous level. Indeed, for a complete definition of the local problems defined on the subdomains, additional conditions are required at the interior corners of the subdomains. Following the recent contribution [58] on the treatment of corners with HABCs, we introduce suited corner conditions into the variational formulation of the subproblems and additional transmission variables at the cross-points. The obtained cross-point treatment accelerates the convergence of the method with a very limited overcost. When a HABC is used as an exterior condition, the cross-point treatment is actually necessary, since the method cannot converge without it. While the approach is designed for regular lattice-type domain partition (*i.e.* with only parallel and perpendicular interfaces) and wave propagation in homogeneous media, it gives very good results with distorted partitions and smoothly varying heterogeneous media.

The paper is organized as follows. In Section 2, we present the Helmholtz boundary-value problem with a HABC and its suitable corner treatment based on adding suitable boundary conditions. The nodal FEM formulation is given next. Section 3 introduces the optimized Schwarz DDM with high-order transmitting boundary conditions. The cross-point treatment

is detailed for two subdomains and then for the multi-subdomain decomposition. The FEM formulation is next stated and some technical aspects about the algorithmic procedure are discussed. In Section 4, we propose some numerical examples to analyze the behavior of the proposed method. Three model configurations with lattice-type partitions are considered for the convergence study. In addition, a sensitivity analysis of the method is reported thanks to the tuning parameters of the high-order transmitting boundary conditions arising in the DDM procedure, the frequency parameter, the mesh refinement and the scalability of the method. After, a numerical investigation with distorted partitions and heterogeneous media is proposed. Finally, we conclude in Section 5.

2 Helmholtz problem with HABC and corner treatment

To describe the method, we consider a two-dimensional Helmholtz problem defined on a rectangular computational domain Ω,

$$\begin{cases} -\Delta u - \kappa^2 u = s, & \text{in } \Omega, \\ \partial_{n_f} u + \mathcal{B}_f u = 0, & \text{on each } \Gamma_f, \end{cases} \tag{1}$$

where κ is the wavenumber, assumed to be a positive constant, and $s(\mathbf{x})$ is a source term. In the numerical simulations, the source term is replaced with a scattering object. The edges of the rectangle are denoted by Γ_f ($f = 1 \dots 4$). For each edge Γ_f, ∂_{n_f} is the (exterior) normal derivative and $\mathcal{B}_f$ is an impedance operator which takes into account the behavior of waves outside the computational domain, that we suppose to be the free-space here. We take the convention that the time-dependence of the fields is $e^{-\imath\omega t}$, where ω is the angular frequency and t is the time.

To simulate wave propagation in free-space, the simplest boundary condition is the Sommerfeld Absorbing Boundary Condition (ABC), which corresponds to using the impedance operator $\mathcal{B}_f = -\imath\kappa$ on the edges. This condition is cheap and easy to use, but the accuracy is known to be poor. In this work, we consider Padé-type high-order absorbing boundary conditions (HABCs) [3, 31, 48, 58], which provide a better accuracy. To preserve the accuracy at the corners of the rectangle, a specific treatment based on compatibility relations derived in [58] is used leading to very low spurious reflections at the boundary. For the HABC, a finite element implementation of the problem is described later. Let us remark that other alternative solutions could be considered for truncating the free-space, like for example by using the well-known Perfectly Matched Layer (PML) approach introduced by Bérenger in [8] and studied *e.g.* in [9, 72] for Helmholtz-type problems. Nevertheless, we do not address this situation here since PMLs are related to the introduction of a surrounding layer which is out of the framework presented here.

2.1 High-order absorbing boundary condition (HABC)

The Padé-type HABC is obtained by approximating an exact non-reflecting boundary condition derived for planar boundaries. Assuming that the exterior medium is homogeneous and free of sources, solving the exterior half-space problem gives the exact (non-local) boundary condition $\partial_n u + \mathcal{B}^{\text{ex}} u = 0$, with the (pseudo-differential) impedance operator $\mathcal{B}^{\text{ex}}$ defined as

$$\mathcal{B}^{\text{ex}} = -\imath\kappa\sqrt{1 + \partial_{\tau\tau}/\kappa^2}, \tag{2}$$

where ∂_n and ∂_τ are respectively the (exterior) normal and tangential derivatives (see *e.g.* [31]). Following [56], this operator is localized by using a Padé approximation of the square-root after

a rotation of the branch-cut. For each face Γ_f, this leads to the HABC impedance operator

$$\mathcal{B}_f = -\imath\kappa\alpha_f \left[1 + \frac{2}{M_f}\sum_{i=1}^{N_f} c_{f,i}\left(1 + \alpha_f^2(c_{f,i}+1)\left[(\alpha_f^2 c_{f,i}+1) + \partial^2_{\tau_f\tau_f}/\kappa^2\right]^{-1}\right)\right], \tag{3}$$

with $\alpha_f = e^{\imath\phi_f/2}$, $c_{f,i} = \tan^2(i\pi/M_f)$ and $M_f = 2N_f + 1$. The accuracy of the Padé-type HABC depends on the number of terms N_f and the angle of rotation ϕ_f (see [48, 58] for further details). In particular, the parameters $N_f = 0$ and $\phi_f = 0$ yield $\mathcal{B}_f = -\imath\kappa$, which corresponds to the Sommerfeld ABC.

For the effective implementation of the HABC, N_f auxiliary fields $\{\varphi_{f,i}\}_{i=1\ldots N_f}$ are defined on Γ_f, and the boundary condition is rewritten as

$$\partial_{n_f} u + B_f\Big(u, \{\varphi_{f,i}\}_{i=1\ldots N_f}\Big) = 0, \quad \text{on } \Gamma_f, \tag{4}$$

with the operator B_f defined as

$$B_f\Big(u, \{\varphi_{f,i}\}_{i=1\ldots N_f}\Big) = -i\kappa\alpha_f\left[u + \frac{2}{M_f}\sum_{i=1}^{N_f} c_{f,i}\,(u + \varphi_{f,i})\right]. \tag{5}$$

The additional fields are governed by the auxiliary equations

$$-\partial_{\tau_f\tau_f}\varphi_{f,i} - \kappa^2\left[(\alpha_f^2 c_{f,i}+1)\varphi_{f,i} + \alpha_f^2(c_{f,i}+1)u\right] = 0, \quad \text{on } \Gamma_f. \quad (i = 1\ldots N_f) \tag{6}$$

The linear multivariate function B_f is introduced to simplify the expressions in the remainder of the paper.

2.2 Corner treatment

When the HABC is prescribed on a boundary with corners, a specific treatment must be used at the corners. Because of the second-order spatial derivative in equation (6), boundary conditions must be added on the auxiliary fields at the extremities of each edge, which are at the corners of the domain. In a previous work [58], we have analyzed several strategies to preserve the accuracy of the solution at the corners. For configurations with right angles, the best approach consists in using a different set of auxiliary fields for each edge, with compatibility relations to couple the auxiliary fields of adjacent edges at the common corner.

Consider two adjacent edges Γ_f and $\Gamma_{f'}$ meeting at the corner $P_{f,f'} = \Gamma_f \cap \Gamma_{f'}$. Two sets of surface fields $\{\varphi_{f,i}\}_{i=1\ldots N_f}$ and $\{\varphi_{f',i'}\}_{i'=1\ldots N_{f'}}$ are defined on Γ_f and $\Gamma_{f'}$, respectively. Globally, a total of $N_f + N_{f'}$ boundary conditions must be written on these auxiliary fields at the corner $P_{f,f'}$. Following the approach detailed in [58], well-suited conditions are such that

$$\partial_{n_{f'}}\varphi_{f,i} + B_{f'}\Big(\varphi_{f,i}, \{\psi_{ff',ii'}\}_{i'=1\ldots N_{f'}}\Big) = 0, \quad \text{on } P_{ff'}, \quad (i = 1\ldots N_f) \tag{7}$$

$$\partial_{n_f}\varphi_{f',i'} + B_f\Big(\varphi_{f',i'}, \{\psi_{f'f,i'i}\}_{i=1\ldots N_f}\Big) = 0, \quad \text{on } P_{ff'}, \quad (i' = 1\ldots N_{f'}) \tag{8}$$

with $N_f \times N_{f'}$ auxiliary variables $\{\psi_{ff',ii'}\}_{i=1\ldots N_f, i'=1\ldots N_{f'}}$ defined as

$$\psi_{ff',ii'} = -\frac{\alpha_{f'}^2(c_{f',i'}+1)\varphi_{f,i} + \alpha_f^2(c_{f,i}+1)\varphi_{f',i'}}{\alpha_f^2 c_{f,i} + \alpha_{f'}^2 c_{f',i'} + 1}, \quad \text{on } P_{ff'}. \quad (i = 1\ldots N_f, i' = 1\ldots N_{f'}) \tag{9}$$

Let us remark that $\psi_{f'f,i'i} = \psi_{ff',ii'}$. In a nutshell, the HABC defined on the field u on one edge is also imposed on the auxiliary fields living on the adjacent edge at the common corner [58], with new auxiliary variables defined at the corner. For instance, the HABC set on $\Gamma_{f'}$ is also forced on the fields $\{\varphi_{f,i}\}_{i=1\dots N_f}$ at $P_{f,f'}$ (equation (7)).

As a particular case, let us consider a configuration with a HABC given on Γ_f and the basic ABC set on the adjacent edge $\Gamma_{f'}$, *i.e.*

$$\partial_{n_f} u + B_f\Big(u, \{\varphi_{f,i}\}_{i=1\dots N_f}\Big) = 0, \quad \text{on } \Gamma_f, \tag{10}$$

$$\partial_{n_{f'}} u - \imath\kappa u = 0, \quad \text{on } \Gamma_{f'}. \tag{11}$$

At the corner $P_{f,f'}$, N_f boundary conditions must be imposed on the auxiliary fields living on Γ_f. Following the approach, the basic ABC must be prescribed

$$\partial_{n_{f'}} \varphi_{f,i} - \imath\kappa\varphi_{f,i} = 0, \quad \text{on } P_{f,f'}, \quad (i = 1\dots N_f) \tag{12}$$

which corresponds to equation (7) with $N_{f'} = 0$ and $\phi_{f'} = 0$.

2.3 Finite element formulation

The problem finally consists in solving the main field u on the rectangular domain with a HABC on each edge by (4). Auxiliary fields defined on the edges are governed by 1D Helmholtz equations through (6) and are coupled at the corners by auxiliary relations by (7)-(8) and auxiliary variables using (9). If the basic ABC is given for the main field on an edge, there is no auxiliary field on that edge, and the basic ABC is prescribed on the auxiliary variables living on the adjacent edges at the common corners.

In order to solve the problem with a finite element scheme, we straightforwardly adapt the bilinear form of the Helmholtz equation. The variational formulation of the problem reads: find $u \in H^1(\Omega)$ and $\varphi_{f,i} \in H^1(\Gamma_f)$, for $f = 1\dots 4$ and $i = 1\dots N_f$, such that

$$\int_\Omega \Big[\nabla u \cdot \nabla v - \kappa^2 uv\Big]\, d\Omega - \sum_{f=1}^{4} \int_{\Gamma_f} B_f\big(u, \{\varphi_{f,i}\}_{i=1\dots N_f}\big) v \, d\Gamma = -\int_\Omega sv \, d\Omega, \quad \forall v \in H^1(\Omega),$$

and

$$\int_{\Gamma_f} \left[\left(\partial_{\tau_f}\varphi_{f,i}\right)\left(\partial_{\tau_f}\rho_f\right) - \kappa^2\left((\alpha_f^2 c_{f,i} + 1)\varphi_{f,i} + \alpha_f^2(c_{f,i}+1)u\right)\rho_f\right] d\Gamma$$
$$- \sum_{f'} \Big[B_{f'}\big(\varphi_{f,i}, \{\psi_{ff',ii'}\}_{i'=1\dots N_{f'}}\big)\rho_f\Big]_{P_{ff'}} = 0, \quad \forall \rho_f \in H^1(\Gamma_f).$$

In the last equation, the index f' corresponds to any edge $\Gamma_{f'}$ adjacent to Γ_f, and the variables $\{\psi_{ff',ii'}\}_{i'=1\dots N_{f'}}$ are defined on $P_{ff'}$ by equation (9). Standard Lagrange finite elements can then be used to discretize the problem.

3 Domain decomposition method with cross-point treatment

In this section, we present a non-overlapping domain decomposition method (DDM) for lattice-type partitions of the domain. The convergence of the method is accelerated by using a Padé-type HABC as a transmission condition with a novel strategy to deal with cross-points. This

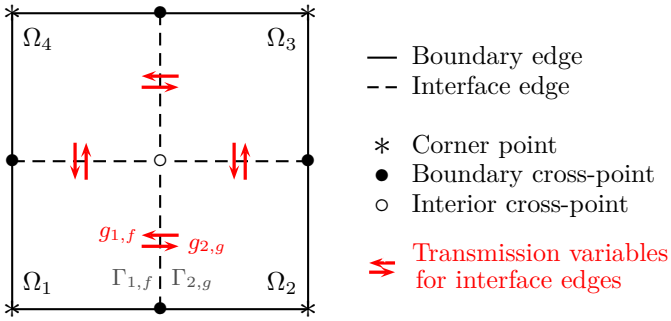


Figure 1: Terminology and transmission variables across the interface edges. In this example, the continuity of the local solution u_1 and u_2 on the interface edge $\Gamma_{1,f} = \Gamma_{2,g}$ is ensured thanks to the transmission variables $g_{1,f}$ and $g_{2,g}$.

strategy relies on the corner treatment derived for the HABCs in the previous section. The DDM and the cross-point strategy are presented in Sections 3.1 and 3.2, respectively. The finite element scheme and the algorithmic procedure are described in Section 3.3.

We consider a partition of the rectangular domain Ω into a grid of rectangular non-overlapping subdomains Ω_I ($I = 1 \dots N^{\text{dom}}$). The edges of each subdomain Ω_I are denoted by $\Gamma_{I,f}$ ($f = 1 \dots 4$). Each edge is either a *boundary edge* if it belongs to the boundary of the global domain ($\Gamma_{I,f} \subset \partial\Omega$), or an *interface edge* if there is a neighboring subdomain beyond the edge ($\Gamma_{I,f} \not\subset \partial\Omega$). In this decomposition, two kinds of points deserve attention: the *boundary cross-points* that belong to two subdomains and that touch the boundary of the global domain, and the *interior cross-points* belonging to four subdomains and that do not touch the boundary of the global domain. These edges and points are illustrated on Figure 1 for a 2×2 partition.

3.1 Optimized Schwarz-type domain decomposition method (DDM)

Following the standard optimized Schwarz-type method, the global problem (1) is decomposed into local subproblems defined on the subdomains. The solution u_I for the subdomain Ω_I is obtained by solving

$$\begin{cases} -\Delta u_I - \kappa^2 u_I = s, & \text{in } \Omega_I, \\ \partial_{n_{I,f}} u_I + \mathcal{B}_{I,f} u_I = g_{I,f}, & \text{on each } \Gamma_{I,f}, \end{cases} \tag{13}$$

where $\mathcal{B}_{I,f}$ is an impedance operator and $g_{I,f}$ is a transmission variable which is set to zero if $\Gamma_{I,f}$ is a boundary edge, while it depends on the local solution belonging of the neighboring subdomain if $\Gamma_{I,f}$ is an interface edge. In the latter case, to ensure the compatibility to the global problem, the transmission variable is defined as

$$g_{I,f} = \partial_{n_{I,f}} u_J + \mathcal{B}_{I,f} u_J, \tag{14}$$

where u_J is the solution of the neighboring subdomain Ω_J. If the impedance operators used on both sides of the interface are the same, which means $\mathcal{B}_{I,f} = \mathcal{B}_{J,g}$, the transmission variables defined on the shared interface edge $\Gamma_{I,f} = \Gamma_{J,g}$ verify

$$g_{I,f} = -g_{J,g} + 2\mathcal{B}_{J,g} u_J, \tag{15}$$

where $g_{J,g}$ is the transmission variable defined on the edge $\Gamma_{J,g}$ of Ω_J.

In the global iterative DDM procedure, each iteration consists in solving concurrently all the subproblems (13) and updating the transmission variables using equation (15) (see Section 3.3 for further details). For a fast convergence of this procedure, the impedance operators used at the interface edges must be chosen wisely. Ideally, for a given subdomain, the operators should correspond to the DtN map related to the complementary of the subdomain. Approximations of this DtN map are also used to define ABCs. Indeed, if there is no source outside the subdomain, the transmission variables in the local system (13) are cancelled, and the transmission conditions should be non-reflecting boundary conditions.

Following [13], the impedance operators for the transmission conditions are based on Padé-type HABCs. For each subdomain Ω_I, the local solution u_I verifies

$$\begin{cases} -\Delta u_I - \kappa^2 u_I = s, & \text{in } \Omega_I, \\ \partial_{n_{I,f}} u_I + B_{I,f}\Big(u_I, \{\varphi_{I,f,i}\}_{i=1\dots N_{I,f}}\Big) = g_{I,f}, & \text{on each } \Gamma_{I,f}, \end{cases} \tag{16}$$

with the transmission variable $g_{I,f}$ that verifies

$$g_{I,f} = \begin{cases} 0, & \text{if } \Gamma_{I,f} \subset \partial\Omega, \\ -g_{J,g} + 2B_{J,g}\Big(u_J, \{\varphi_{J,g,j}\}_{j=1\dots N_{J,g}}\Big), & \text{if } \Gamma_{I,f} \not\subset \partial\Omega. \end{cases} \tag{17}$$

The second equation of system (16) is a boundary condition if $\Gamma_{I,f}$ is a boundary edge, or a transmission condition if $\Gamma_{I,f}$ is an interface edge. In both cases, if $N_{I,f} > 0$, auxiliary fields $\{\varphi_{I,f,i}\}_{i=1\dots N_{I,f}}$ are defined on the edge, and are governed by

$$-\partial_{\tau_{I,f}\tau_{I,f}}\varphi_{I,f,i} - \kappa^2\left[(\alpha_{I,f}^2 c_{I,f,i} + 1)\varphi_{I,f,i} + \alpha_{I,f}^2(c_{I,f,i}+1)u_I\right] = 0, \quad \text{on } \Gamma_{I,f}, \tag{18}$$

with $i = 1 \dots N_{I,f}$. The parameters of the transmission conditions used on both sides of an interface edge must be the same (*i.e.* $N_{I,f} = N_{J,g}$ and $\phi_{I,f} = \phi_{J,g}$, with $\Gamma_{I,f} = \Gamma_{J,g}$), since we assumed that the impedance operators are the same on a shared interface. For consistency, the same boundary condition must be prescribed on the boundary edges of the subdomains and on the corresponding edges of the global domain.

Boundary conditions must be set on the auxiliary fields at the extremities of the edges because of the second-order partial derivative in the governing equation (18). The extremities of an edge are at corners of a subdomain, and correspond to interior cross-points, boundary cross-points or corners of the global domain. The cross-point treatment, described in the next section, actually provides the missing boundary conditions at the cross-points.

3.2 Dealing with cross-points

The cross-point treatment relies on the corner treatment described in Section 2. It is applied at the corners of the subdomains. Depending on the configuration, it provides boundary conditions or transmission conditions for the auxiliary fields at the cross-points. In the latter case, new transmission variables are defined at the cross-points.

3.2.1 Two-subdomain case

To describe the approach, we first consider a partition of the rectangular domain Ω into two rectangular subdomains with an interface Γ and two boundary cross-points. Three configurations, represented on Figure 2, are studied: the basic ABC prescribed on $\partial\Omega$ with a HABC-based

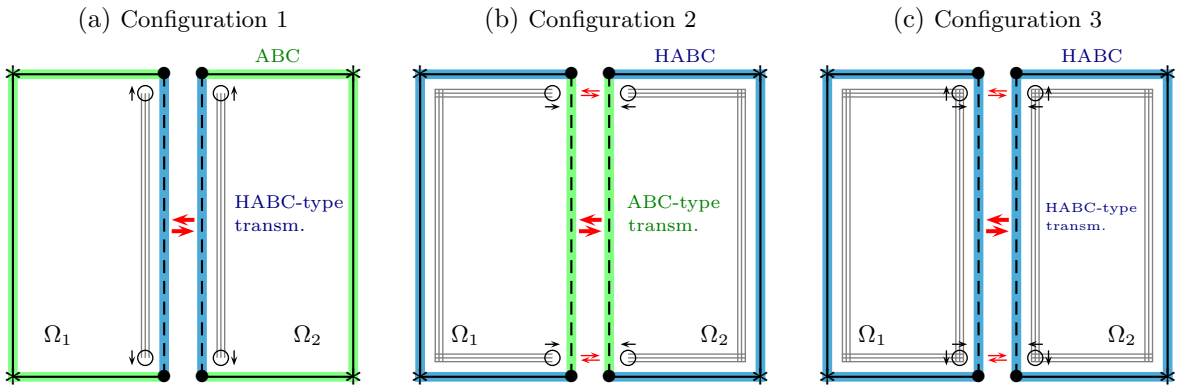


Figure 2: Three configurations for two-subdomain case. The exterior boundary condition is a basic ABC or a HABC, and the transmission condition is based on the HABC operator or the basic ABC operator. The thin gray lines illustrate the position of auxiliary fields. The black arrows indicate where boundary conditions are required for auxiliary fields. The red arrows indicate transmission conditions on the edge or at the cross-points.

transmission condition on Γ (Configuration 1), a HABC on $\partial\Omega$ with a transmission condition based on the basic ABC on Γ (Configuration 2), and the HABC operator used both for $\partial\Omega$ and Γ (Configuration 3). Because the HABC is used on the exterior boundary and/or the interface, a specific treatment must be used at the boundary cross-points.

In the first configuration (Figure 2a), auxiliary fields are defined on both sides of the interface. These fields require boundary conditions at the extremities of the interface, which are corners of the subdomains. The basic ABC is set on the adjacent edges (*i.e.* the upper and lower boundary edges). Following the strategy of Section 2, the basic ABC is also given on the auxiliary fields at the boundary cross-points.

In the second configuration (Figure 2b), a HABC is given on each gobal edge Γ_f in the global problem, auxiliary fields are defined on each edge and the corner treatment is used. After the domain partition, a HABC is imposed on each boundary edge $\Gamma_{I,f}$ of each subdomain Ω_I, and a set of auxiliary fields is defined on each of these edges. For the consistency of the global problem, the parameters of the HABC on $\Gamma_{I,f}$ must be the same as the parameters of the HABC given on the global edge $\Gamma_f \supset \Gamma_{i,f}$. For a global edge Γ_f that has been divided by the partitioning (upper and lower edges on Figure 2b), the continuity of the auxiliary fields must be enforced at the cross-points. As the ABC-based transmission condition is used on the main field on the interface, this transmission condition is also used on each auxiliary field at the boundary cross-points and auxiliary transmission variables are defined at these points.

The last configuration combines the difficulties. The exterior boundary condition and the transmission condition are based on HABCs (Figure 2c). The auxiliary fields living on every edge require boundary conditions at the boundary cross-points. To deal with this case, we recall that the operators used on the edges of each subdomain should approximate the DtN map of the free-space if the exterior medium relative to each subdomain is homogeneous. If the transmission variables are canceled, it corresponds to forcing a HABC on every edge. Therefore, we apply the corner treatment described in Section 2 to all the corners of the subdomains, which gives boundary conditions for the auxiliary fields. If the transmission variables are not canceled, the continuity of the fields $\{u_I\}_{I=1,2}$ is enforced at the interface thanks to the right-hand side of the second equation of system (13). For the auxiliary fields living on the boundary edges, the boundary conditions at the cross-point become transmission conditions by adding transmission

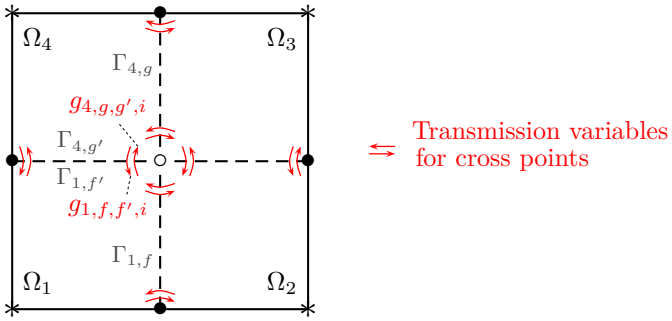


Figure 3: Transmission variables across the boundary and interior cross-points, if the HABC operator is used both in the exterior boundary condition and in the interface conditions. In the example, the continuity of the auxiliary fields $\varphi_{1,f,i}$ and $\varphi_{4,g,i}$ (defined on the aligned edges $\Gamma_{1,f}$ and $\Gamma_{4,g}$) at the interior cross-point $P_{1,ff'} = P_{4,gg'}$ is ensured thanks to the transmission variables $g_{1,ff',i}$ and $g_{4,gg',i}$. These variables verify equation (20).

variables in the right-hand sides, as for the second configuration. These transmission variables verify relations similar to equation (17) at the cross-points.

3.2.2 Multi-subdomains case

In the general case, the rectangular domain Ω is partitioned into a grid of rectangular subdomains, with interior and boundary cross-points. The strategy relies on the following principles, which generalize the approaches used for the three previous configurations:

- The same transmission condition is used on both sides of each interface edge. The boundary condition used on each boundary edge is the same as the one prescribed on the corresponding edge of the global domain. In the domain partition, the HABC operators used on edges that are on a same line have the same parameters (*e.g.* on Figure 3: the top edges of Ω_1 and Ω_2, the top edges of Ω_3 and Ω_4, the left edges of Ω_1 and Ω_4, ...).
- If auxiliary fields are defined on an edge $\Gamma_{I,f}$ of a subdomain Ω_I, boundary conditions or transmission conditions must be set on these fields at the extremities of this edge (which can be interior cross-points, boundary cross-points, or corners of Ω). These conditions are given by the condition already used for u_I on the adjacent edges. If a transmission condition is used on u_I on an adjacent edge, transmission conditions are considered on the auxiliary fields at the cross-point, and new transmission variables are introduced.
- The corner treatment described in Section 2 is used at the corners of each subdomain, which gives boundary conditions to the auxiliary fields living on the edges. At the cross-points, these conditions can become transmission conditions by adding transmission variables in the right-hand sides, which are similar to equation (17).

Following these principles, the description of the problem with domain decomposition can be completed.

For each subdomain Ω_I, the local solution u_I verifies equations (16). For each edge $\Gamma_{I,f}$, the transmission variable $g_{I,f}$ satisfies equation (17). Each auxiliary field $\varphi_{I,f,i}$ $(i = 1 \dots N_{I,f})$

defined on a boundary or interface edge $\Gamma_{I,f}$ is such that

$$\begin{cases} -\partial_{\tau_{I,f}\tau_{I,f}}\varphi_{I,f,i} - \kappa^2\big((\alpha_{I,f}^2 c_{I,f,i}+1)\varphi_{I,f,i} + \alpha_{I,f}^2(c_{I,f,i}+1)u_I\big) = 0, & \text{on } \Gamma_{I,f}, \\ \partial_{n_{I,f'}}\varphi_{I,f,i} + B_{I,f'}\left(\varphi_{I,f,i}, \{\psi_{I,ff',ii'}\}_{i'=1\dots N_{I,f'}}\right) = g_{I,ff',i}, & \text{on each } P_{I,ff'}, \end{cases} \tag{19}$$

with the transmission variable $g_{I,ff',i}$

$$g_{I,ff',i} = \begin{cases} 0, & \text{if } \Gamma_{I,f'} \subset \partial\Omega, \\ -\,g_{J,gg',i} + 2B_{J,g}\left(\varphi_{J,g,i}, \{\psi_{J,gg',ii'}\}_{i'=1\dots N_{J,g'}}\right), & \text{if } \Gamma_{I,f'} \not\subset \partial\Omega. \end{cases} \tag{20}$$

In these relations, $\Gamma_{I,f'}$ is any edge that is adjacent to $\Gamma_{I,f}$, and $P_{I,ff'} = \Gamma_{I,f} \cap \Gamma_{I,f'}$ is the corner that is shared by these edges. The second equation of system (19) is a boundary condition if $\Gamma_{I,f'}$ is a boundary edge, or a transmission condition if $\Gamma_{I,f'}$ is an interface edge. The transmission variable is set to zero in the former case, and it depends on the solution of the other side of $\Gamma_{I,f'}$ in the latter case. The variables $\psi_{I,ff',ii'}$ are defined using equation (9).

In equation (20), the indices are chosen in such a way that Ω_J is the neighboring subdomain on the other side of $\Gamma_{I,f'}$, the edge $\Gamma_{J,g'}$ is shared by the subdomains (*i.e.* $\Gamma_{I,f'} = \Gamma_{J,g'}$), and the edge $\Gamma_{J,g}$ is aligned with $\Gamma_{I,f}$ (*i.e.* $f = g$), as illustrated on Figure 3. The variable $g_{J,gg',i}$ is used in a transmission condition for an auxiliary field $\varphi_{J,g,i}$ living on $\Gamma_{J,g}$. Therefore, the transmission conditions enforce the continuity of the auxiliary fields $\varphi_{I,f,i}$ and $\varphi_{J,g,i}$, which live on edges that are on the same line. Let us note that, since the HABC parameters are the same for edges that are aligned, $N_{I,f} = N_{J,g}$ and $\phi_{I,f} = \phi_{J,g}$.

3.3 Finite element scheme and algorithmic procedure

Each step of the DDM iterative procedure consists in solving a local subproblem on each subdomain, and updating the transmission variables both on the interface edges and at the cross-points. The numerical solution of the subproblems is performed with a standard nodal finite element scheme built on a conformal mesh made of triangles or quadrangles. For each subdomain Ω_I, the variational formulation of the subproblem reads: find $u_I \in H^1(\Omega_I)$ and $\varphi_{I,f,i} \in H^1(\Gamma_{I,f})$, with $i = 1 \dots N_{I,f}$ and $f = 1 \dots 4$, such that

$$\begin{aligned} \int_{\Omega_I} \Big[\nabla u_I \cdot \nabla v_I - \kappa^2 u_I v_I\Big]\, d\Omega_I + \sum_{f=1}^{4} \int_{\Gamma_{I,f}} B_{I,f}\Big(u_I, \{\varphi_{I,f,i}\}_{i=1\dots N_{I,f}}\Big) v_I \, d\Gamma_{I,f} \\ = \int_{\Omega_I} s v_I \, d\Omega_I + \sum_{f=1}^{4} \int_{\Gamma_{I,f}} g_{I,f} v_I \, d\Gamma_{I,f}, \qquad \forall v_I \in H^1(\Omega_I), \end{aligned} \tag{21}$$

and

$$\begin{aligned} \int_{\Gamma_{I,f}} \Big[(\partial_{\tau_{I,f}}\varphi_{I,f,i})\,(\partial_{\tau_{I,f}}\rho_{I,f}) - \kappa^2\Big((\alpha_{I,f}^2 c_{I,f,i}+1)\varphi_{I,f,i} + \alpha_{I,f}^2(c_{I,f,i}+1)u_I\Big)\,\rho_{I,f}\Big]\, d\Gamma_{I,f} \\ + \sum_{f'} \Big[B_{I,f'}\left(\varphi_{I,f,i}, \{\psi_{I,ff',ii'}\}_{i'=1\dots N_{I,f'}}\right)\rho_{I,f}\Big]_{P_{I,ff'}} \\ = \sum_{f'} \Big[g_{I,ff',i}\,\rho_{I,f}\Big]_{P_{I,ff'}}, \qquad \forall \rho_{I,f} \in H^1(\Gamma_{I,f}). \end{aligned} \tag{22}$$

In the last equation, the index f' corresponds to any edge $\Gamma_{I,f'}$ adjacent to $\Gamma_{I,f}$, and $P_{I,ff'} = \Gamma_{I,f} \cap \Gamma_{I,f'}$ is the shared corner. The variables $\psi_{I,ff',ii'}$ are defined using (9). This variational formulation is an extension of the one used in [13] (see equation (62) in that reference). In that work, there is only one set of auxiliary fields and equations on the subdomain boundary $\partial\Omega_I$. Here, there is one set for each edge $\Gamma_{I,f}$ of the subdomain, and new terms appear in (22) to deal with the corners of the subdomain.

In the DDM iterative procedure, the transmission variables computed at an iteration n are used in the right-hand side of equations (21)-(22) to compute the local fields of the iteration $n+1$. The transmission variables are then updated using equations (17)-(20). Therefore, at each iteration, the *interface* transmission variables are computed using

$$g_{I,f}^{(n+1)} = -g_{J,g}^{(n)} + 2B_{J,g}\Big(u_J^{(n+1)}, \{\varphi_{J,g,j}^{(n+1)}\}_{j=1\ldots N_{J,g}}\Big), \tag{23}$$

for each interface edge $\Gamma_{I,f} \not\subset \partial\Omega$. Similarly, the *cross-point* transmission variables are updated through

$$g_{I,ff',i}^{(n+1)} = -g_{J,gg',i}^{(n)} + 2B_{J,g}\left(\varphi_{J,g,i}^{(n+1)}, \{\psi_{J,gg',ii'}^{(n+1)}\}_{i'=1\ldots N_{J,g'}}\right), \tag{24}$$

at each cross-point $P_{I,ff'}$, with $\Gamma_{I,f'} \not\subset \partial\Omega$.

The global process can be recast as one application of an iterative operator $\mathcal{A} : \mathcal{G} \to \mathcal{G}$ defined by

$$g^{(n+1)} = \mathcal{A}g^{(n)} + b, \tag{25}$$

where $g^{(n)} \subset \mathcal{G}$ is the set of transmission data, and b depends on the source term s. This can be seen as one iteration of the Jacobi method to solve the linear system $(\mathcal{I}-\mathcal{A})g = b$, where $\mathcal{I}$ is the identity operator. Following a well-known strategy (see *e.g.* [7, 13]), a GMRES Krylov subspace iterative solver is used on the top of the DDM procedure to solve this linear system efficiently. Here, by contrast with most of the works, the transmission data contains transmission variables associated to both interfaces and cross-points.

4 Numerical results

This section reports some finite element simulations to study the HABC-based domain decomposition method with cross-point treatment. After a description of three benchmarks in Section 4.1, we analyze the convergence history (Section 4.2), the sensitivity to the HABC parameters (Section 4.3) and the influence of the wavenumber, the mesh density and the number of subdomains on the convergence rate (Section 4.4). In section, configurations with non-regular domains partitions and smoothly varying heterogemeous media are investigated in Section 4.5.

4.1 Description of the benchmarks

The reference benchmark used through this section is the scattering of an incident plane wave $u^{\text{inc}}(\mathbf{x}) = e^{\imath\kappa x}$ by a sound-soft circular scatterer. For a circle of radius R centered at the origin, the scattered field is given by

$$u^{\text{ref}}(r,\theta) = -\sum_{m=0}^{\infty} \epsilon_m \imath^m \frac{J_m(\kappa R)}{H_m^{(1)}(\kappa R)} H_m^{(1)}(\kappa r)\cos(m\theta), \quad r \geq R, \tag{26}$$

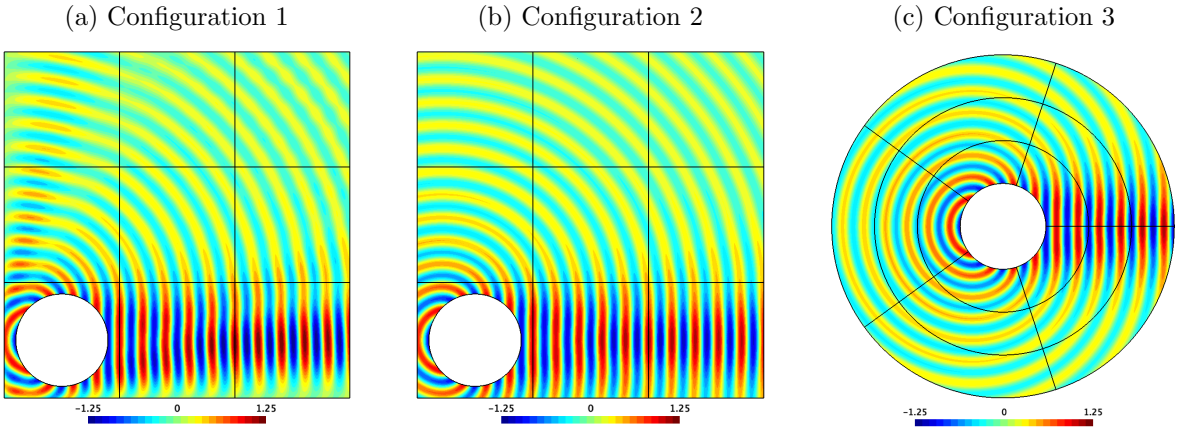


Figure 4: Scattering benchmarks: real part of the scattered field of the reference numerical solution for the three configurations with $\kappa = 4\pi$. The basic ABC, a HABC and the second-order Bayliss-Turkel ABC are set on the exterior border, respectively.

where (r, θ) are the polar coordinates, J_m is the m^{th}-order Bessel's function, $H_m^{(1)}$ is the m^{th}-order first-kind Hankel function, and ϵ_m is the Neumann function which is equal to 1 for $m = 0$ and 2 otherwise.

Three configurations are considered. For the first configuration (Figure 4a), the finite element simulations are performed on the square computational domain $[-1.25, 6.25] \times [-1.25, 6.25]$ with checkerboard partitions. The scatterer is the unit disk centered at the origin. The basic ABC, *i.e.* $\partial_n u - \imath\kappa u = 0$, is set on the exterior boundary of the domain. Because this boundary condition is a rather inaccurate non-reflecting boundary treatment, the numerical solution contains both the scattered field and spurious waves reflected on the exterior boundary. For the second configuration (Figure 4b), the HABC is used on the edges of the square domain with the suited treatment at the corners. The HABC parameters $N = 6$ and $\phi = 0.3\pi$ have been selected to avoid any visible modelling error in the numerical solution (*i.e.* the numerical error due to the finite element scheme is significantly larger than the modelling error due to the approximate boundary condition — see [58]). For the third configuration (Figure 4c), we have considered a circular domain of radius $R_1 = 4$ with radial/longitudinal partitions. The second-order Bayliss-Turkel ABC [2, 5]

$$\partial_r u - \imath\kappa u + \frac{1}{2R_1} u - \frac{1}{8R_1^2(R_1^{-1} - \imath\kappa)} u - \frac{1}{2(R_1^{-1} - \imath\kappa)} \partial_{\tau\tau} u = 0, \tag{27}$$

is given on the exterior circular boundary, where $\partial_\tau = R_1^{-1}\partial_\theta$ is the tangential derivative over the circle of radius R_1 in polar coordinates (r, θ).

For all the configurations, the finite element scheme is based on meshes made of second-order curvilinear triangular elements and quadratic polynomial basis functions (P2). The Dirichlet BC $u = -u^{\text{inc}}$ is set at the boundary of the (sound-soft) scatterer. By default, the wavenumber is $\kappa = 4\pi$ and the characteristic number of vertices per wave length is $n_\lambda = 10$. The meshes of the squared domain and the circular domain are made of 56 538 and 49 718 triangles, respectively. For the three configurations, the relative L^2-errors of the finite element solutions compared to the reference solution (26) are 2.2×10^{-1}, 2.4×10^{-3} and 1.7×10^{-3}, respectively. We used the GetDDM framework [70] which combines the mesh generator Gmsh [44] and the finite element solver GetDP [28].

4.2 Convergence analysis

We begin by analyzing the convergence of the DDM procedure with cross-point treatment for the different configurations. The relative L^2-errors and the relative residuals are plotted as functions of the number of the GMRES iterations on Figure 5 for the three configurations. The L^2-error is calculated by comparing the solution obtained in each subdomain to the reference numerical solution computed on the same mesh without domain decomposition. In every case, HABC-based transmission conditions with different numbers of auxiliary fields are tested ($N = 0$, 2, 4 and 6 with $\phi = 0.3\pi$). The effect of the cross-point treatment is analyzed by keeping or removing the corresponding terms in the finite element scheme. The latter case consists in setting a homogeneous Neumann BC on the auxiliary fields at the cross-points. On all the figures, the dotted lines are associated to results without the cross-point treatment.

For the first configuration (*i.e.* squared domain with basic ABC), the relative residual and the relative error have the same order of magnitude in all the cases (Figures 5a-5b) and decrease during the iterations. Using the cross-point treatment clearly accelerates the convergence, especially for transmission conditions with large values of N. The number of iterations to reach a relative error of magnitude 10^{-6} is reduced by 20% to 40% thanks to the treatment. When the cross-point treatment is enabled, the decay of residual and error can be accelerated further, up to a certain point, by taking a number of auxiliary fields N sufficiently large. Higher values for N does not change the results, while, without the cross-point treatment, increasing N slightly slows down the decays.

The good results obtained with the cross-point treatment and N sufficiently large can be interpreted by looking at the numerical solution after each iteration (Figure 6). At the initialization, the right-hand side term of the iteration system is computed by solving each subproblem with source terms only (see Section 3). Here, only the subdomain containing the scattering disk has a source, and then non-zero solution (Figure 6a). The numerical solution in this subdomain is already rather accurate since the transmission condition acts as a HABC, and the cross-point treatment behaves as the suited corner treatment. Since there is neither source nor very significant reflected waves generated outside the subdomain, the HABC and the corner treatment constitute a very good boundary treatment for the subdomain. During the iterations, the signal is propagated from subdomain to subdomain. At the fourth iteration, the signal reaches the last subdomain. This coincides with a sharp reduction of both the residual and error by an order one in magnitude.

For the second configuration (*i.e.* squared domain with HABC), the impact of the cross-point treatment is more important. When the cross-point treatment is not enabled, the residuals decrease with the iterations (Figure 5c), but the relative errors reach a plateau and stagnate at 10^{-1} (Figure 5d). This can be explained by noting the only difference with the previous configuration: a HABC is prescribed on the exterior boundary instead of a basic ABC, and auxiliary fields defined on the edges of the domain Ω. Without the cross-point treatment, the derivative of these auxiliary fields is set to zero at the boundary cross-points. Then, the problem with domain decomposition is not compatible with the original problem, and the iterative schemes converge towards a wrong solution. To fix this, the continuity of the auxiliary fields living on the boundary edges must be enforced at the boundary cross-points. With the cross-point treatment, transmission conditions are set at the boundary cross-points, and the error decays together with the residual, as it should be. It is worth to note that the absence of cross-point treatment in the first configuration does not break the convergence of the error because no auxiliary fields are defined on the boundary edges. In that case, the problem with domain partition is compatible with the original problem.

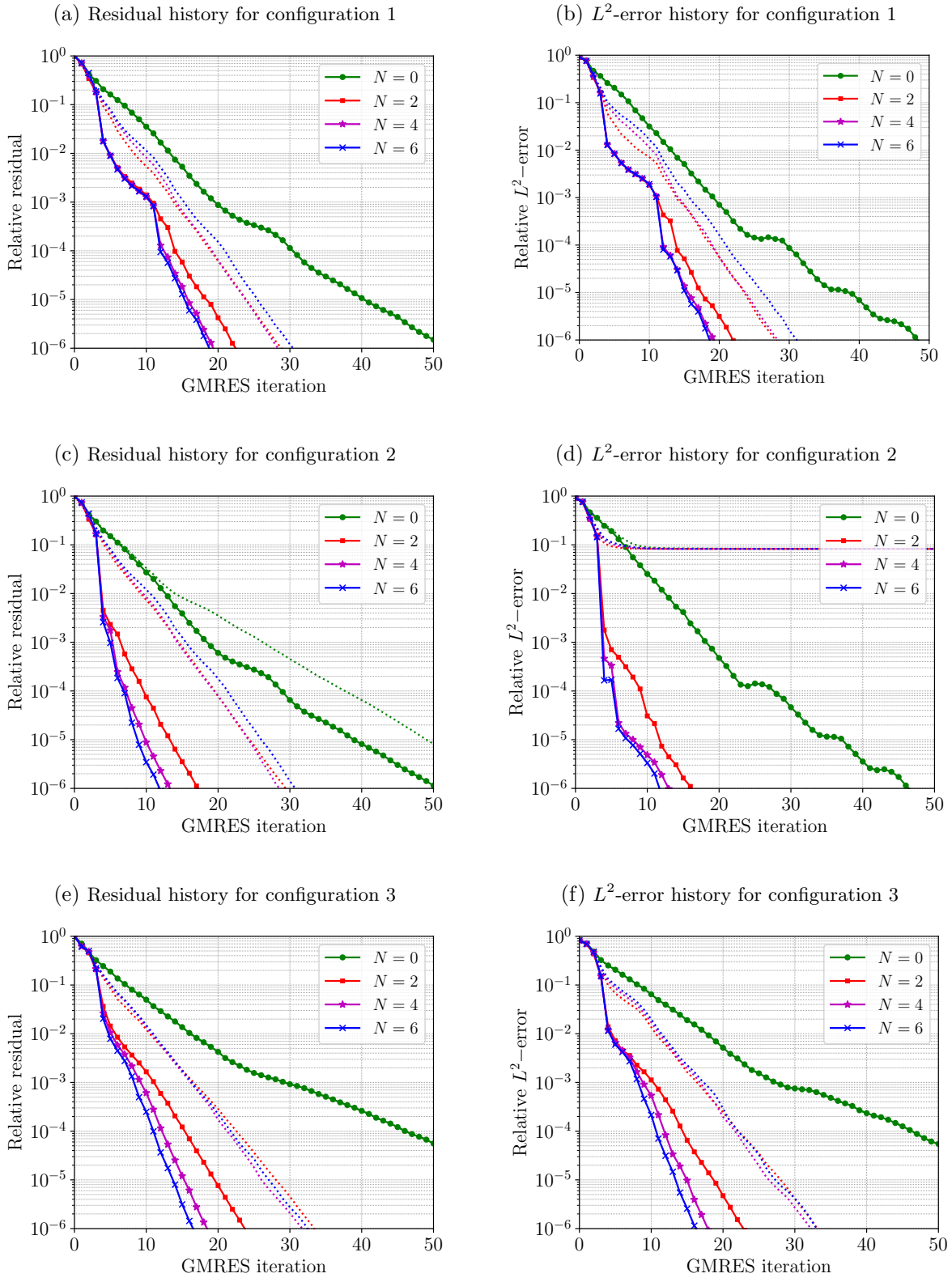


Figure 5: Evolution of relative residual *(left)* and relative L^2-error *(right)* in the course of the GMRES iterations for the three configurations represented on Figure 4. HABC-based transmission conditions with $N = 0, 2, 4, 6$ auxiliary fields and $\phi = 0.3\pi$ are used. The dotted lines correspond to the results obtained when the cross-point treatment is not used. Handling the cross-point procedure is represented by continuous lines.

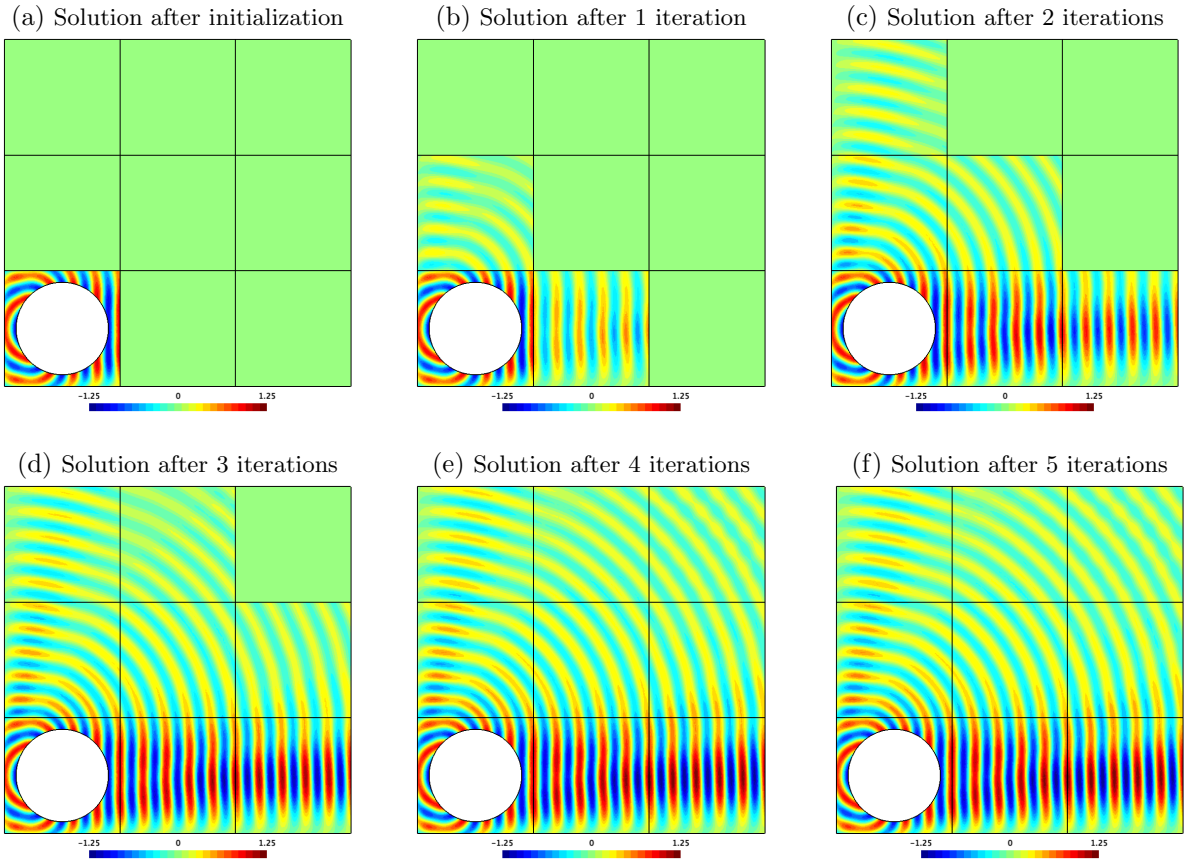


Figure 6: Evolution of the solution during the GMRES iterations for configuration 1 and the HABC-based transmission condition with $N = 4$ and $\phi = 0.3\pi$. The first picture is obtained after initialization of the RHS of the transmission system.

When comparing the results of the two first configurations for high values of N, we observe that the error decays faster for the second configuration, especially between the iterations 3 and 4, where the error drops by at least 3 orders of magnitudes. This tremendous result is likely due to the specificity of the benchmark: the exact scattering solution verifies the exact free-space boundary condition on the boundary and the interfaces. Since the HABC is used both as exterior BC and transmission condition, the exact behavior of the solution is captured with a few iterations. By contrast, when the basic ABC is used as exterior BC, small waves reflected on the ABC must travel towards the subdomains.

For the third configuration (*i.e.* circular domain with Bayliss-Turkel ABC), both the relative residual and the relative error decrease in all the cases. Again, the decay is faster with the cross-point treatment. Let us highlight that no treatment is used at the boundary cross-points, which corresponds to the junction of interface edges with the exterior boundary (with a Bayliss-Turkel ABC) or the border of the circular scatterer (with an inhomogeneous Dirichlet BC). The method can then be applied to settings with such boundary conditions without any issue.

4.3 Sensitivity to the HABC parameters

The efficiency of the transmission condition depends on the number N of auxiliary fields and the rotating angle ϕ. To study the sensitivity of the convergence to these parameters, we perform the DDM procedure with several values of N and ϕ for the three configurations. The number

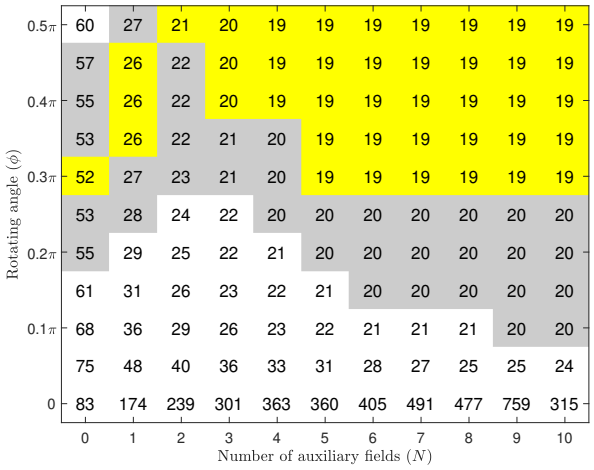


Figure 7: Number of GMRES iterations to reach the relative residual 10^{-6} in configuration 1 for different values of the number of auxiliary fields N and rotating angle ϕ. For each column (*i.e.* each value of N), cells in yellow correspond to the minimal number of iterations, while cells in gray are up to 10% from the minimal number of iterations.

of GMRES iterations to reach the relative residual 10^{-6} are reported on Figure 7 for the first configuration.

For any given ϕ, increasing the number of auxiliary fields N accelerates the convergence, up to a certain limit, as already mentioned in the previous section. The only exception is for $\phi = 0$. Nevertheless, increasing N leads to a higher computational cost and the amount of data to exchange at the cross-points. It is then advantageous to take the smallest N yielding the best convergence. For practical applications, the optimal N would likely depends on the configuration.

The selection of the parameter ϕ is an important matter, because it accelerates the convergence of the iterative process at no additional cost. We observe first that the Padé case ($\phi = 0$) gives the worst result in all the cases, and it should be avoided. The optimal value for ϕ, represented for each N by yellow cells on Figure 7, depends on the number N of auxiliary fields. This can make the parameter selection rather tricky. Fortunately, the number of iterations is not very sensitive to ϕ as soon as it is sufficiently large (*i.e.* larger than $\pi/4$ here). The range of the nearly-optimal values of ϕ, represented by the gray zone on Figure 7, is indeed rather wide.

The results for the other configurations lead to similar conclusions. They are not reported here for the sake of conciseness. In the remainder of the paper, we always use $\phi = 0.3\pi$, which is a nearly-optimal value for all the configurations.

4.4 Influence of the wavenumber, the mesh density and the number of subdomains

In this section, we study the sensitivity of the method thanks to the wavenumber κ, the mesh density n_λ and the number of subdomains. High frequency simulations are challenging because they require fine meshes with high mesh densities to avoid the pollution effect. The efficiency of the method for large values of κ and n_λ is therefore an important issue.

Figure 8 shows the number of iterations to reach the relative residual 10^{-6} with respect to κ and n_λ for the various configurations and several values of N. For configurations 1 and 3, the

dotted lines correspond to cases where the cross-point treatment is not used.

We first analyze the influence of κ on the convergence. For $N = 0$, the number of iterations increases with respect to κ in all the cases (Figures 8a, 8c and 8e). The increase is very slow for the second configuration, and faster for the third one. For higher values of N, the convergence does not change significantly with κ when the cross-point treatment is used. As already observed, higher values of N accelerate the convergence, and the convergence is slower if the cross-point treatment is not used.

For the first and third configurations, the number of iterations increases with the mesh density n_λ for all the values of N (Figures 8b and 8f). Fortunately, the number of iterations can be kept constant when increasing n_λ by taking N larger: the number of iterations then remains approximately 20 for the first configuration and 17 for the third configuration. Therefore, a convergence independent of the mesh density can be achieved provided that N is sufficiently large. This was already observed in [13] on benchmarks without cross-points treatment. The results are slightly different for the second configuration (Figure 8d): the number of iterations increases very slowly for $N = 0$ and 2, while it decreases until a plateau for $N = 6$ and 12. The plateau is lower for $N = 12$ than for $N = 6$. This is likely due to the fact that the numerical solution is closer to the exact free-space scattering solution, and that the HABC-based transmission condition is perfectly suited to this specific case.

These results then indicate that the method is well-adapted to high-frequency problems with high density meshes, provided that N is sufficiently large.

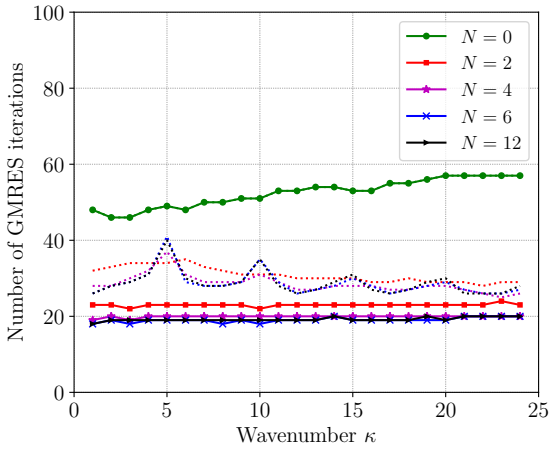
Figure 9 shows the evolution of the number of GMRES iterations with respect to the number of subdomains for the first and third configurations. The simulations have been performed with increased numbers of subdomains in the x- and y-directions for the squared domain (resp. $N_{\text{dom},x}$ and $N_{\text{dom},y}$) and in the r- and θ-directions for the circular domain (resp. $N_{\text{dom},r}$ and $N_{\text{dom},\theta}$). The size of the domains increases with the number of subdomains: the squared domain is $[-1.25, 2.5N_{\text{dom},x} - 1.25] \times [-1.25, 2.5N_{\text{dom},y} - 1.25]$ and the circular domain is $\{(r, \theta) : r \in [1, 1 + N_{\text{dom},r}], \theta \in [0, 2\pi]\}$. The results for the second configuration are similar to those for the first one. They are not reported for the sake of shortness. The strong scaling analysis for the third configuration (*i.e.* increasing the number of subdomains without increasing the size of the domain) also leads to similar results.

We directly see that the method scales in both cases (Figures 9a and 9b). The number of iterations increases linearly with the number of subdomains in each direction, which is the expected behavior. Indeed, since the transmission of propagating waves from subdomain to subdomain is local with the transmission conditions, a larger number of iterations is required to allow the propagation of waves across a larger number of subdomains. Preconditioning techniques based on sweeps (*e.g.* [32, 66, 69, 73, 74]) and coarse spaces (*e.g.* [4, 10, 22, 34]) allow for global transmissions of information between the subdomains with improved convergences. The combination of our approach with preconditioning techniques is currently under investigation.

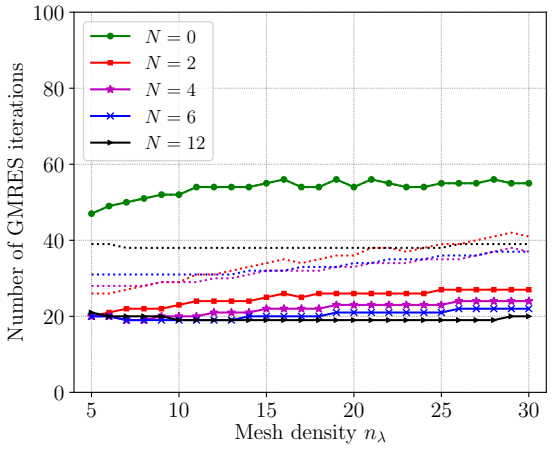
4.5 Experiments with non-right angles and heterogeneous media

The proposed DDM is *a priori* suited only to wave propagation in homogeneous media and lattice-type domain partitions with right angles. Indeed, the HABC operator used in the transmission condition is built under the hypothesis of a homogeneous medium, and the compatibility relations used in the cross-point treatment are derived for corners with right angles (see Section 2). Nevertheless, the HABC can be used as a good approximation with smoothly-varying heterogeneous media, since it can represent locally the transmission of waves at the interface (see *e.g.* [57]). The compatibility relations derived for right-angle corners can be used as an

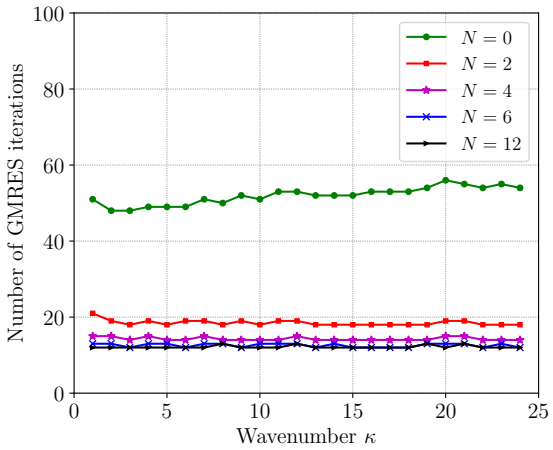
(a) Number of iterations *vs* κ for configuration 1



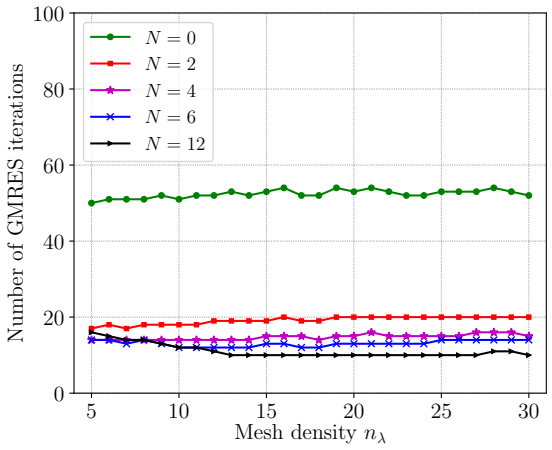
(b) Number of iterations *vs* n_λ for configuration 1



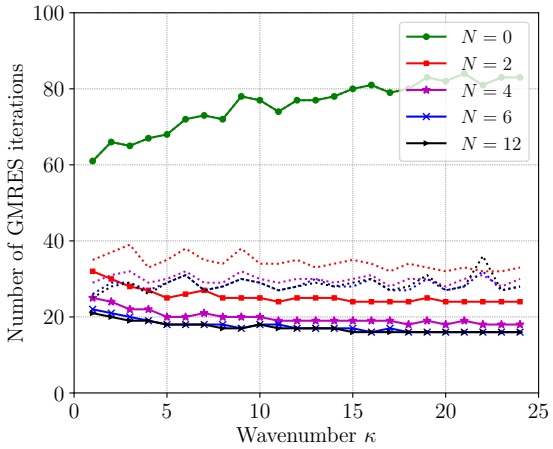
(c) Number of iterations *vs* κ for configuration 2



(d) Number of iterations *vs* n_λ for configuration 2



(e) Number of iterations *vs* κ for configuration 3



(f) Number of iterations *vs* n_λ for configuration 3

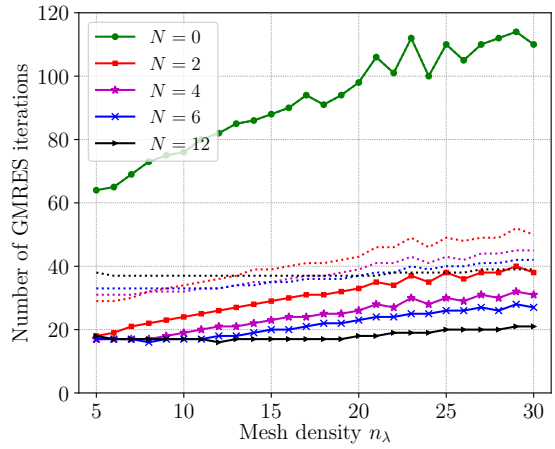
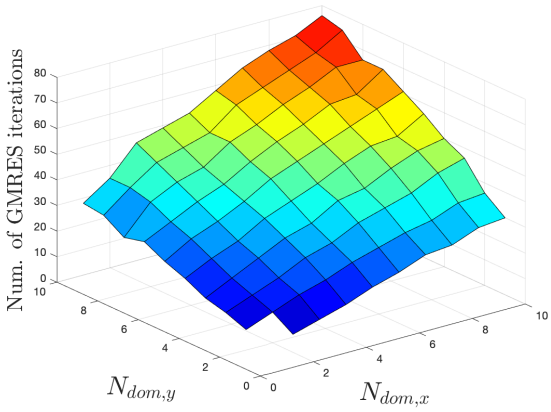


Figure 8: Number of GMRES iterations to reach the relative residual 10^{-6} as a function of the wavenumber κ with a fixed mesh density $n_\lambda = 10$ *(left)* or as a function of the mesh density n_λ with a fixed wavenumber $\kappa = 4\pi$ *(right)* to assert the scaling of the solution with κ and n_λ.

(a) Scaling for configuration 1



(b) Scaling for configuration 3

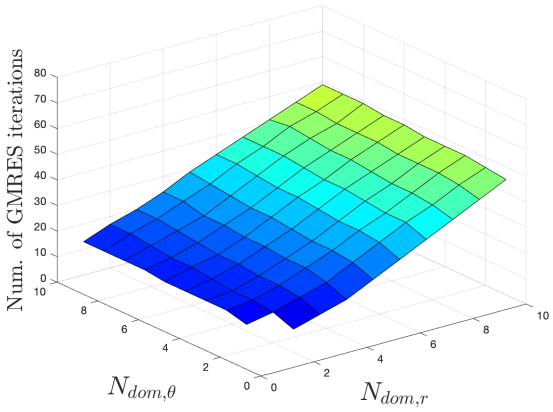


Figure 9: Number of GMRES iterations to reach the relative residual 10^{-6} for different number of subdomains to assert the scaling of the procedure. The size of the main domain increases with the number of subdomains in the x-, y- and r-directions.

approximate treatment with non-right angles [58].

In this section, our approach is tested for configurations where the corner treatment is not exact: first, for domain partitions with non-right interior angles (Section 4.5.1) and second for wave propagation in heterogeneous media (Section 4.5.2).

4.5.1 Configurations with distorted partitions

To analyze the method for partitions with non-right angles, we consider the scattering benchmark and the three configurations described in Section 4.1. The partitions are deformed by moving the cross-points, which create acute and obtuse angles, as shown on Figure 10. The points are shifted for the two first configurations (on distances 0.5, 1 and 1.5) and twisted for the third one (by angles 0.1π, 0.2π and 0.3π). In every case, HABC-based transmission conditions with different numbers of auxiliary fields are tested ($N = 0$, 2, 4 and 6 with $\phi = 0.3\pi$). The effect of the cross-point treatment is analyzed by keeping or removing the corresponding terms in the finite element scheme. The terms implemented for the right-angle case are used without modification for non-right angles.

Table 1 shows the number of GMRES iterations to reach the relative residual 10^{-6} for each case. The relative L^2-error (not shown for the sake of shortness) is always close to 10^{-6}, except for the second configuration (*i.e.* squared domain with a HABC on the exterior border) without cross-point treatment. As discussed in Section 4.2, the compatibility is not ensure at the boundary cross-points for that case. We have observed that, when using the cross-point treatment, the method converges towards the correct solution, even with an important distortion of the partition. In that case, several interfaces starting from boundary cross-points are not perpendicular to the exterior border.

In nearly all the cases, the number of GMRES iterations increases when the distortion of the partitions is amplified. For the first configuration, the increase is rather small, with and without cross-point treatment. For the two other configurations, the number of iterations increases more rapidly when the cross-point treatment is used. Nevertheless, in all the cases,

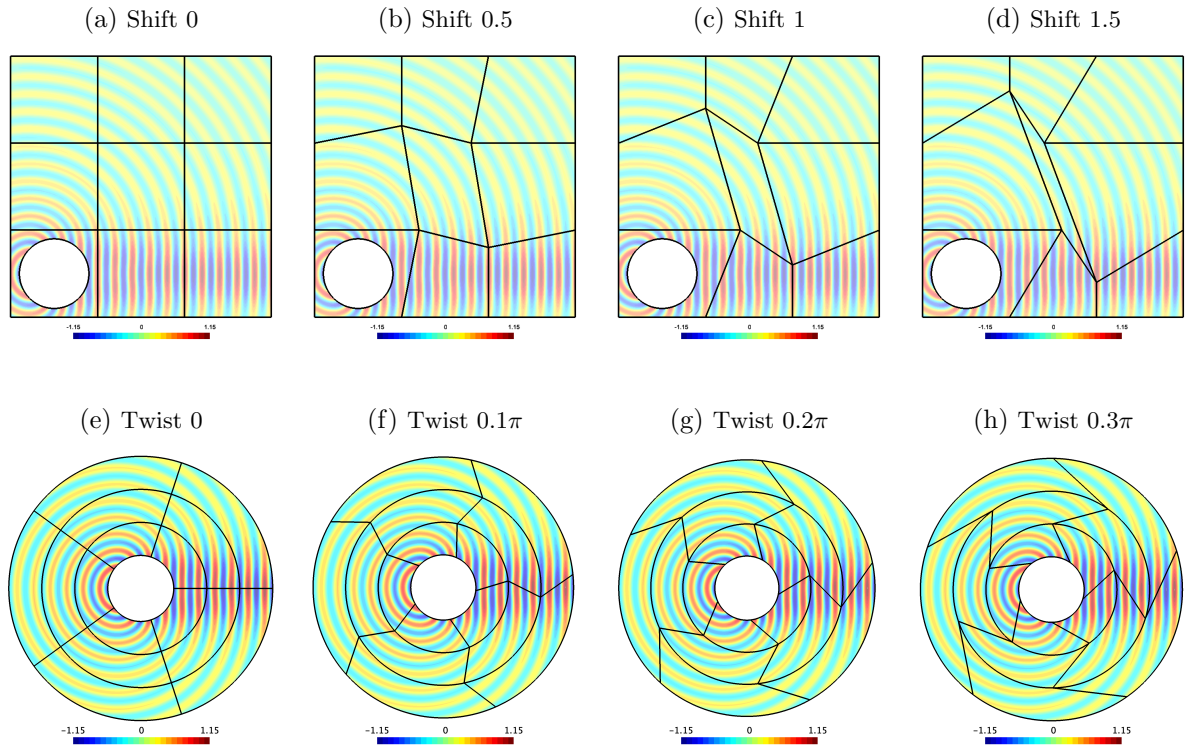


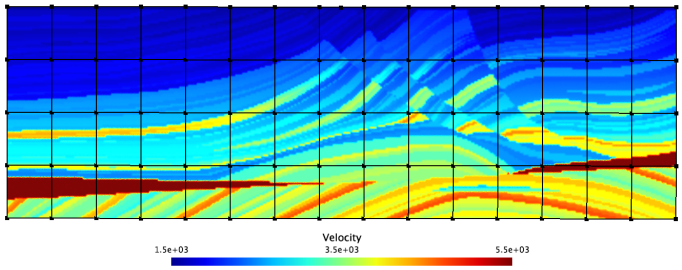
Figure 10: Snapshot of the distorted partitions for the squared domain (configurations 1 and 2, Figure 10a-10d) and the circular domain (configuration 3, figures 10e-10h).

Table 1: Number of GMRES iterations to reach the relative residual 10^{-6} for the different configurations with distorted domain partitions. The final relative L^2-error is also approximately 10^{-6} for every case, except for the second configuration without cross-point treatment (results not shown) where the method is not consistent.

		Configuration 1				Configuration 2				Configuration 3			
	Shift/Twist →	0	0.5	1	1.5	0	0.5	1	1.5	0	0.1π	0.2π	0.3π
No cross-point treatment	$N=0$	52	55	58	66	-	-	-	-	76	81	85	91
	$N=2$	29	32	32	34	-	-	-	-	34	38	38	38
	$N=4$	29	31	32	33	-	-	-	-	32	34	33	34
	$N=6$	31	32	33	34	-	-	-	-	33	32	31	32
With cross-point treatment	$N=0$	52	55	58	66	51	57	60	67	76	81	85	91
	$N=2$	23	24	25	28	18	22	25	27	24	29	32	34
	$N=4$	20	20	21	24	14	18	20	23	19	24	28	30
	$N=6$	19	19	20	22	12	17	20	21	17	23	26	28

using the cross-point treatment accelerates the convergence. The speedup is smaller for the third configuration, but it is still significant. For the most distorted configurations (*i.e.* shift with 1.5 and twist with 0.3π), the smallest numbers of iterations always correspond to the cases with both the largest N and the cross-point treatment. These results show the robustness of the approach with non-right angles.

(a) Velocity model



(b) Real part of the numerical solution u

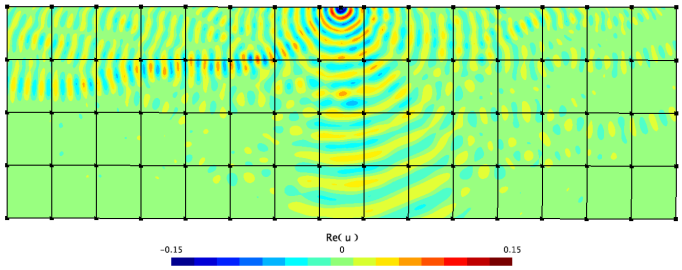


Figure 11: Marmousi benchmark: Velocity model (a) and reference numerical solution (b). The HABC is used as boundary condition.

4.5.2 Benchmark with a smoothly varying heterogeneous medium

As a preliminary study to analyze the effectiveness of the method with heterogeneous media, we consider the Marmousi model (Figure 11), a velocity map $c(\mathbf{x})$ that represents a geological structure. This model is frequently used to evaluate modeling and imaging techniques. Here, the Helmholtz equation with spatially varying wavenumber is solved on the domain,

$$-\Delta u - \kappa(\mathbf{x})^2 u = \delta(\mathbf{x} - \mathbf{x}_{\text{sou}}), \quad \text{on } \Omega,$$

with $k(\mathbf{x}) = 2\pi\nu/c(\mathbf{x})$ and $\Omega = [0, 9.192\text{m}] \times [-2.904\text{m}, 0]$. A Dirac source point is placed at the coordinate $\mathbf{x}_{\text{sou}} = (4585\text{m}, -10\text{m})$, closed to the surface. The homogeneous Neumann boundary condition is prescribed on the upper side of the domain, and the basic ABC is set on the other sides. The domain is partitioned into 4×15 rectangular subdomains. The spatially varying $\kappa(\mathbf{x})$ is used *as is* in the equations of the DDM initially derived for a constant wavenumber. The simulations are performed for the frequency $\nu = 10\ \text{s}^{-1}$ and the characteristic mesh size $h \approx 30\text{m}$. The mesh is made of 84.022 triangles and 170.539 second-order nodes.

Because the parameter analysis performed in Section 4.3 is not longer valid, the HABC-type transmission condition is tested with several number of auxiliary fields and rotating angles. The number of GMRES iterations to reach the relative residual 10^{-6} with several values for N and ϕ are shown on Table 2. The final relative L^2-error, not shown, is about 10^{-6} in all the cases.

We observe that using the HABC-based transmission condition (with $N \geq 1$) instead of the basic impedance condition (corresponding to $N = 0$) still accelerates the convergence for every value of $\phi \geq 0.1\pi$. However, the speedup is limited and stagnates rather rapidly when increasing N: using $N = 2$ is sufficient. This is expected since the HABC operator is not designed for heterogeneous media. Nevertheless, it provides a significant speedup in comparison with the

Table 2: Marmousi. Number of GMRES iteration.

	$\phi \rightarrow$	0	0.1π	0.2π	0.3π	0.4π	0.5π
No cross-point treatment	$N=0$	168	134	123	125	135	155
	$N=1$	>500	106	90	83	80	80
	$N=2$	>500	96	83	78	77	78
	$N=3$	>500	95	82	78	77	77
	$N=4$	>500	89	80	78	77	77
With cross-point treatment	$N=0$	168	134	123	125	135	155
	$N=1$	>500	89	79	74	72	72
	$N=2$	>500	82	73	71	70	69
	$N=3$	>500	76	71	70	70	70
	$N=4$	>500	80	74	72	72	73

basic impedance condition, with a moderate supplementary computational cost. In addition, the result is not very sensitive to the value of the parameter ϕ: the values in the range $[0.3\pi, 0.5\pi]$ give similar results.

Finally, we observe that, in all the cases, using the cross-point treatment accelerates the convergence. The speedup is rather small, but this is also expected since the issue here is related to the quality of the transmission condition, more than the treatment of the cross-points.

5 Conclusion

In this work, we presented a suitable way to address the cross-point problem for the efficient parallel finite element solution of high-frequency scattering problems with an optimized Schwarz DDM. We considered cases where a Padé-type HABC operator is used for the transmission condition (to accelerate the convergence of the procedure), for the exterior boundary condition (to improve the accuracy of the solution) or for both conditions.

To handle the cross-points, suitable relations and additional transmission variables were introduced at the points. Numerical results have shown that the convergence rate of the obtained DDM is improved. We systematically analyzed the way the convergence depends on the tuning parameters of the method as well as the frequency, the mesh refinement and the number of subdomains. Configurations with distorted partitions and heterogeneous media were tested. While the method was conceived for lattice-type partitions with right angles, it also performed very well with partitions having non-right angles. As expected, the efficiency of the approach for configurations with heterogeneous media was not as performant. Current approaches to tackle this problem are based on non-local methods (see *e.g.* [21, 68, 69, 75]) and preconditioners (see *e.g.* [4, 22, 33, 43]).

The extension to the 3D Helmholtz equation can be obtained by adapting the developments of the present paper and the technical details given in [58]. Even if the DDM gains in efficiency thanks to the cross-point treatment, the reported scalability results prove that the method is intrinsically dependent of the number of subdomains since the iterated wave field needs to be propagated through the subdomains, translating hence the nonlocal nature of wave-like problems, whatever is the optimized local transmission condition. Furthermore improvements to avoid the problem are currently being developed by using fast sweeping preconditioners [69, 73, 74] and coarse space approximations [4, 22]. Finally, extensions to other time-harmonic wave problems, in particular for electromagnetic [30] and elastic waves, is still needed [17, 29, 55]. This is under study but the problem is technically much more complicated.

Acknowledgments

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