

# Numerical Scheme Impacts on Time Domain Full Waveform Inversion

Pierre Jacquet, Andreas Atle, Hélène Barucq, Henri Calandra, Julien Diaz

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## Numerical Scheme Impacts on Time Domain Full Waveform Inversion

#### Mathias 2019

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Pau, FRANCE









### Seismic Acquisition



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### Seismic Acquisition



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 $Pierre\ Jacquet\ \texttt{pierre}\ \texttt{,}\ \texttt{jacquet}\ \texttt{@inria.fr}\ |\ \textbf{Numerical\ Scheme\ Impacts\ on\ Time\ Domain\ FWI}$ 



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Cost function to minimize :

$$\mathcal{J}(\boldsymbol{m}) = rac{1}{2} ||\boldsymbol{d}_{obs} - \mathcal{F}(\boldsymbol{m})||^2$$

- F(*m*) is the restriction on the receivers of the simulated waves in the medium *m*. (With  $m = c, \rho, \kappa$ ...)
- FWI iterates until  $\mathcal{J}(\boldsymbol{m}) \longrightarrow 0$

#### [1] Patrick Lailly

The seismic inverse problem as a sequence of before stack migrations Conference on Inverse Scattering

#### [2] Albert Tarantola

Inversion of seismic reflection data in the acoustic approximation Geophysics, Vol. 49, 1984

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### **Continuous Forward Model**

First order acoustic wave equation

$$\begin{cases} \frac{1}{\rho c^2} \frac{\partial \boldsymbol{p}}{\partial t} + \nabla \cdot \boldsymbol{v} = f_{\rho} & \text{on } \Omega \\ \rho \frac{\partial \boldsymbol{v}}{\partial t} + \nabla \boldsymbol{p} = 0 & \text{on } \Omega \\ \boldsymbol{p} = 0 & \text{on } \boldsymbol{\Gamma}_1 \\ \frac{\partial \boldsymbol{p}}{\partial t} + \boldsymbol{c} \nabla \boldsymbol{p} \cdot \boldsymbol{n} = 0 & \text{on } \boldsymbol{\Gamma}_2 \\ \boldsymbol{p}(0) = 0, \quad \boldsymbol{v}(0) = 0 \end{cases}$$



Conditions

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### **Discrete Forward Model**

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Space Discretization : Discontinuous Galerkin Elements

- Nodal (Lagrangian / Jacobian)
- Modal (Bernstein-Bézier)

### **Discrete Forward Model**

Contraction 5

Space Discretization : Discontinuous Galerkin Elements

- Nodal (Lagrangian / Jacobian)
- Modal (Bernstein-Bézier)

Semi-discretized model :

$$\frac{\partial}{\partial t} \bar{\boldsymbol{U}}(t) = A \bar{\boldsymbol{U}}(t) + \bar{\boldsymbol{F}}(t)$$

with :

$$\mathbf{\bar{U}}(t) = \begin{pmatrix} \mathbf{\bar{P}}(t) \\ \mathbf{\bar{V}}(t) \end{pmatrix}$$

### **Discrete Forward Model**

Contraction 5

Space Discretization : Discontinuous Galerkin Elements

- Nodal (Lagrangian / Jacobian)
- Modal (Bernstein-Bézier)

Time schemes :

- Runge Kutta 2/4
- Adams Bashforth 3

Semi-discretized model :

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with :

$$\bar{\boldsymbol{U}}(t) = \begin{pmatrix} \bar{\boldsymbol{P}}(t) \\ \bar{\boldsymbol{V}}(t) \end{pmatrix}$$



Discontinuous Galerkin Method



Assets of Discontinuous Galerkin Methods :

- Unstructured grid (enable to match the topography and medium irregularities)
- Robust to physical discontinuities
- hp-adaptivity
- Massively parallel performance properties



h-adaptivity



p-adaptivity with P1, P2, P3 elements

#### Outline



#### Time Domain Full Waveform Inversion

Seismic Acquisition FWI Workflow Forward Discretization

#### Some Results

1D Results 2D Time Domain FWI Results 2D Multiscale Reconstruction 1D Preliminary tests



### FWI 1D Results



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### 2D Time Domain Reconstruction

#### 2D FWI :

- Developped in Total environmement (DIP<sup>1</sup>)
- Nodal Space Operators (Lagrangian/Jacobian)
- Modal Space Operators (Bernstein-Bézier)
- Runge Kutta 2/4 and Adams Bashforth 3 time-schemes

Gradient expression :

$$\nabla_{\frac{1}{\kappa}} \mathcal{J} = \int_0^T \int_\Omega \frac{\partial \boldsymbol{\rho}}{\partial t} \boldsymbol{\lambda}_1 d\Omega dt \quad \text{with} : \boldsymbol{\kappa} = \boldsymbol{\rho} \boldsymbol{c}^2$$

 $m{c},\,
ho$  and  $m{\kappa}$  Constant per elements

<sup>&</sup>lt;sup>1</sup>http://dip.inria.fr/



# 2D Time Domain FWI Reconstructions

#### **RK2** Reconstructed *c* Model (30 iterations)

5,000 4,000 3,000 2,000





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 $m \cdot s^{-1}$ 



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depth (km)

#### 2D Time Domain FWI Reconstructions Time-schemes comparison

#### RK4 Reconstructed c Model (30 iterations)

5,000 4,000 3,000 2,000







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 $m \cdot s^{-1}$ 



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depth (km)

#### 2D Time Domain FWI Reconstructions Time-schemes comparison

#### AB3 Reconstructed c Model (30 iterations)









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depth (km)

# 2D Time Domain FWI Reconstructions



- 47k P1 elements
- Constant  $\rho$  model ( $\rho = 1$ )
- 19 sources / 181 Receivers
- Noise : SNR=10
- 30 iterations
- 120 cores
- Polynomial basis : Nodal

#### **Computational time :**

- RK2 : 3h15
- RK4 : 4h30
- AB3 : 5h10



# 2D Time Domain FWI Reconstructions



- 47k P1 elements
- Time Scheme : RK2
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- ▶ 19 sources / 181 Receivers
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- 30 iterations
- 120 cores

#### **Computational time :**

- Nodal : 3h15
- Modal: 4h30<sup>[1]</sup>

Cost function evolution :



#### Chan J. and Warburton T. GPU-Accelerated Bernstein Bézier Discontinuous Galerkin Methods for Wave Problems SIAM Journal on Scientific Computing 2017



Reconstruction with an initial smooth model

#### Reconstructed model *c* Model (30 iterations RK2)









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 $m \cdot s^{-1}$ 



#### 2D Multiscale Reconstructions Multiscale Principle [1]



 C. Bunks, F. M. Saleck, S. Zaleski, and G. Chavent Multiscale seismic waveform inversion GEOPHYSICS, Vol. 60, No. 5, 1995

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depth (km)







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Reconstruction with an initial smooth model

#### Reconstructed *c* Model with 1.0-7.5Hz filter









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 $m \cdot s^{-1}$ 





Reconstruction with an initial smooth model

depth (km)

#### Reconstructed *c* Model with 1.0-15Hz filter







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 $m \cdot s^{-1}$ 











- 47k P1 elements
  Time Scheme : RK2
- Constant  $\rho$  model ( $\rho = 1$ )
- 19 sources / 181 Receivers
- Noise : SNR=10
- 120 cores
- 20 FWI iterations per filter
- Computation time : 10h
- Frequencies : 1-2.5Hz, 1-5.0Hz





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### Conclusion



#### Main Results :

- 2D Acoustic Reconstruction performed with different discretization
- Multiscale FWI implemented and working on Marmousi

### Conclusion



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- 2D Acoustic Reconstruction performed with different discretization
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#### Perspectives :

- Perform reconstruction on other test cases (2D/3D)
- Develop enhanced optimizers (NLCG, Limited BFGS)
- Adapt the code to use High order Model
- Extend the code to elastic and elasto-acoustic propagator
- Exploit coupled numerical method (SEM/DG) (Aurélien Citrain Thesis)

### Conclusion



#### Main Results :

- 2D Acoustic Reconstruction performed with different discretization
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Thank you.





### Adjoint State Method

Lagrangian fonctional [1] :

$$\mathcal{L}(\widehat{\boldsymbol{u}},\widehat{\boldsymbol{\lambda}},\boldsymbol{m}) = rac{1}{2} ||\boldsymbol{d}_{obs} - \mathcal{R}(\widehat{\boldsymbol{u}})||^2 + < \textit{Forward}_{\boldsymbol{m}}(\widehat{\boldsymbol{u}}) - f_{p},\widehat{\boldsymbol{\lambda}} >$$

If  $\hat{\boldsymbol{u}} = \boldsymbol{u}$  Solution of the Direct Problem  $\iff (Forward_{\boldsymbol{m}}(\boldsymbol{u}) - f_{p} = 0)$ :

 $\mathcal{J}(\boldsymbol{m}) = \mathcal{L}(\boldsymbol{u}, \widehat{\boldsymbol{\lambda}}, \boldsymbol{m})$ 

 Plessix R-E A review of the adjoint-state method for computing the gradient of a functional with geophysical applications Geophysical Journal International, Volume 167, Issue 2, 2006

### Adjoint State Method

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$$\mathcal{J}(\textbf{\textit{m}}) = \mathcal{L}(\textbf{\textit{u}}, \widehat{oldsymbol{\lambda}}, \textbf{\textit{m}})$$

Let us choose  $\widehat{\boldsymbol{\lambda}} = \boldsymbol{\lambda}$  such as  $\frac{\partial \mathcal{L}}{\partial \boldsymbol{u}} = 0$ 

$$(\mathcal{R}^* d_{obs} - u) + \mathit{Forward}^*_{m}(oldsymbol{\lambda}) = 0$$

#### Plessix R-E

A review of the adjoint-state method for computing the gradient of a functional with geophysical applications Geophysical Journal International, Volume 167, Issue 2, 2006

### Adjoint State Method

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Let us choose  $\widehat{\boldsymbol{\lambda}} = \boldsymbol{\lambda}$  such as  $\frac{\partial \mathcal{L}}{\partial \boldsymbol{u}} = 0$ 

$$(\mathcal{R}^* \textit{d}_{obs} - \textit{u}) + \textit{Forward}^*_{m}(m{\lambda}) = 0$$

For *Forward*<sub>*m*</sub>( $\boldsymbol{u}$ ) -  $f_p = 0$ :

$$\partial_{\boldsymbol{m}_i} \mathcal{J}(\boldsymbol{m}) = \partial_{\boldsymbol{m}_i} \mathcal{L}(\boldsymbol{u}, \boldsymbol{\lambda}, \boldsymbol{m}) = \partial_{\boldsymbol{m}_i} < \textit{Forward}_{\boldsymbol{m}}(\boldsymbol{u}), \boldsymbol{\lambda} >$$

#### [1] Plessix R-E

A review of the adjoint-state method for computing the gradient of a functional with geophysical applications Geophysical Journal International, Volume 167, Issue 2, 2006

### **Adjoint Formulation**



Continuous Direct Problem





### Adjoint Formulation



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### **Adjoint Formulation**



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#### AtD : Adjoint then Discretized Strategy

$$\mathcal{J}(\boldsymbol{p}) = rac{1}{2} ||\boldsymbol{d}_{obs} - R\boldsymbol{p}||^2$$

1

$$\begin{cases} \frac{1}{\rho \boldsymbol{c}^2} \frac{\partial \boldsymbol{p}}{\partial t} + \nabla \cdot \boldsymbol{v} = f_p & \text{on } \Omega \\ \rho \frac{\partial \boldsymbol{v}}{\partial t} + \nabla \boldsymbol{p} = 0 & \text{on } \Omega \\ \boldsymbol{p} = 0 & \text{on } \boldsymbol{\Gamma}_1 \\ \frac{\partial \boldsymbol{p}}{\partial t} + \boldsymbol{c} \nabla \boldsymbol{p} \cdot \boldsymbol{n} = 0 & \text{on } \boldsymbol{\Gamma}_2 \\ \boldsymbol{p}(0) = 0, \quad \boldsymbol{v}(0) = 0 \end{cases}$$

$$\frac{1}{\rho c^2} \frac{\partial \lambda_1}{\partial t} + \nabla \cdot \lambda_2 = \frac{\partial \mathcal{J}}{\partial \rho} \quad \text{on } \Omega$$
$$\rho \frac{\partial \lambda_2}{\partial t} + \nabla \lambda_1 = 0 \quad \text{on } \Omega$$
$$\lambda_1 = 0 \quad \text{on } \Gamma_1$$
$$\frac{\partial \lambda_1}{\partial t} - c \nabla \lambda_1 \cdot n = 0 \quad \text{on } \Gamma_2$$
$$\lambda_1(T) = 0, \quad \lambda_2(T) = 0$$

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 $t \in [0, T]$   $t \in [T, 0]$ 

### AtD : Adjoint then Discretized Strategy

$$\mathcal{J}(\boldsymbol{p}) = \frac{1}{2} ||\boldsymbol{d}_{obs} - \boldsymbol{R}\boldsymbol{p}||^{2}$$

$$\begin{cases} \frac{\partial \boldsymbol{\bar{U}}}{\partial t}^{n} = \boldsymbol{A}\boldsymbol{\bar{U}}^{n} + \boldsymbol{\bar{F}}^{n} \\ \text{With} : \quad \boldsymbol{\bar{U}}^{n} = \begin{pmatrix} \boldsymbol{\bar{P}}^{n} \\ \boldsymbol{\bar{V}}^{n} \end{pmatrix} \qquad \begin{cases} \frac{\partial \boldsymbol{\bar{\Lambda}}}{\partial t}^{n} = \boldsymbol{A}\boldsymbol{\bar{\Lambda}}^{n} + \boldsymbol{R}^{*}(\boldsymbol{R}\boldsymbol{\bar{U}}^{n} - \boldsymbol{d}_{obs}) \\ \text{With} : \quad \boldsymbol{\bar{\Lambda}}^{n} = \begin{pmatrix} \boldsymbol{\bar{\Lambda}}_{1}^{n} \\ \boldsymbol{\bar{\Lambda}}_{2}^{n} \end{pmatrix} \end{cases}$$

$$0 \underbrace{\boldsymbol{1}}_{\text{Time-steps going Forward}} T \qquad 0 \underbrace{\boldsymbol{1}}_{\text{Time-steps going Backward}} T$$

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All time scheme can be summed-up such as :

 $L\overline{U} = E\overline{F}$ 

RK4 time-scheme leads to :

$$\bar{\bm{U}}^{n+1} = B\bar{\bm{U}}^n + \bm{C_0}\bar{\bm{F}}^n + \bm{C_1}\bar{\bm{F}}^{n+\frac{1}{2}} + \bm{C_1}\bar{\bm{F}}^{n+1}$$



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# DtA : Discretize then Adjoint Strategy



All time scheme can be summed-up such as :

#### $L\bar{U} = E\bar{F}$

We are looking for a Discrete Adjoint state satisfying :

 $\boldsymbol{L^*} \boldsymbol{\bar{\Lambda}} = -R^* (\boldsymbol{d_{obs}} - R \boldsymbol{\bar{U}})$ 

With the adjoint operator L\* satisfying :

 $< Lar{U},ar{\Lambda}> = <ar{U},L^*ar{\Lambda}>$ 



All time scheme can be summed-up such as :

 $L\bar{U}=E\bar{F}=\bar{G}$ 

We are looking for a Discrete Adjoint state satisfying :

 $\boldsymbol{L^*} \boldsymbol{\bar{\Lambda}} = -R^* (\boldsymbol{d_{obs}} - R \boldsymbol{\bar{U}}) = \boldsymbol{\bar{D}}$ 

With the adjoint operator L\* satisfying :

 $< Lar{U},ar{\Lambda}> = <ar{U},L^*ar{\Lambda}>$ 

 $<ar{m{G}},ar{m{\Lambda}}>=<ar{m{U}},ar{m{D}}>$  (Adjoint Test)

Adjoint test succeeds  $\iff$  operator  $L^*$  well established



#### DtA : Discretize then Adjoint Strategy Example with RK4

RK4 time-scheme leads to :

$$\bar{\bm{U}}^{n+1} = B\bar{\bm{U}}^n + \bm{C_0}\bar{\bm{F}}^n + \bm{C_1}\bar{\bm{F}}^{n+\frac{1}{2}} + \bm{C_1}\bar{\bm{F}}^{n+1}$$

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$$\boldsymbol{L}\boldsymbol{\bar{U}} = \boldsymbol{E}\boldsymbol{\bar{F}} = \boldsymbol{\bar{G}}$$

$$\begin{pmatrix} I & & \\ -B & I & \\ & -B & I & \\ & & \ddots & \ddots & \\ & & & -B & I \end{pmatrix} \begin{pmatrix} \boldsymbol{\bar{U}}^{0} \\ \boldsymbol{\bar{U}}^{1} \\ \boldsymbol{\bar{U}}^{2} \\ \vdots \\ \boldsymbol{\bar{U}}^{n} \end{pmatrix} = \begin{pmatrix} \boldsymbol{\bar{G}}^{0} \\ \boldsymbol{\bar{G}}^{1} \\ \boldsymbol{\bar{G}}^{2} \\ \vdots \\ \boldsymbol{\bar{G}}^{n} \end{pmatrix}$$

So :

$$\boldsymbol{L^*} = \begin{pmatrix} I & -B^* & & \\ & I & -B^* & \\ & \ddots & \ddots & \\ & & & I & -B^* & \\ & & & & I & \end{pmatrix}$$

## Adjoint Strategies Comparison

#### **Adjoint Then Discretize**

- + Physical approach
- + Same discrete operators for Forward and Backward
- -- Approximate gradient [1]
  - Consistent with the discretization

#### **Discretize then Adjoint**

- + Numerical approach
- + Has an Adjoint Test
- Tremendous work to develop the adjoint operators

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• Non-consistency of the adjoint state [2]

- Sirkes, Ziv and Tziperman, Eli Finite Difference of Adjoint or Adjoint of Finite Difference ? 1997
- [2] Sei Alain and Symes William A Note on Consistency and Adjointness for Numerical Schemes 1997

## Adjoint Strategies Comparison

#### **Adjoint Then Discretize**

- + Physical approach
- + Same discrete operators for Forward and Backward
- - Approximate gradient [1]
  - Consistent with the discretization

#### **Discretize then Adjoint**

- + Numerical approach
- + / Has an Adjoint Test (in theory)

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- Tremendous work to develop the adjoint operators
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### 1D Preliminary tests :

#### 1D FWI :

Gradient expression :

- Lagrange / B-Bézier Operators
- RK4 / AB3 time-schemes

 $\nabla_{\boldsymbol{c}} \mathcal{J} = -\int_{0}^{T} \int_{\Omega} \frac{2}{\rho \boldsymbol{c}^{3}} \frac{\partial \boldsymbol{p}}{\partial t} \boldsymbol{\lambda}_{1} d\Omega dt$ 

### 1D Preliminary tests :

#### 1D FWI :

- Lagrange / B-Bézier Operators
- RK4 / AB3 time-schemes

Adjoint test passed with :

- With a canonical space inner-product (< u, v ><sub>X</sub> = ∑<sub>i</sub> u<sub>i</sub>v<sub>i</sub>)
- ► With a M-space inner product (< u, v ><sup>M</sup><sub>X</sub> =< Mu, v ><sub>X</sub>)

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Gradient expression :

$$\nabla_{\boldsymbol{c}} \mathcal{J} = -\int_{0}^{T} \int_{\Omega} \frac{2}{\boldsymbol{\rho} \boldsymbol{c}^{3}} \frac{\partial \boldsymbol{\rho}}{\partial t} \boldsymbol{\lambda}_{1} d\Omega dt$$

./run --- Adjoint test ---inner product U/D 553123.57586755091 inner product G/Q 553123.57586756046



### 1D Preliminary tests :

#### 1D FWI :

- Lagrange / B-Bézier Operators
- RK4 / AB3 time-schemes

#### Adjoint test passed with :

- With a canonical space inner-product (< u, v ><sub>X</sub> = ∑<sub>i</sub> u<sub>i</sub>v<sub>i</sub>)
- ► With a M-space inner product (< u, v ><sup>M</sup><sub>X</sub> = < Mu, v ><sub>X</sub>)

Gradient expression :

$$\nabla_{\boldsymbol{c}} \mathcal{J} = -\int_{0}^{T} \int_{\Omega} \frac{2}{\boldsymbol{\rho} \boldsymbol{c}^{3}} \frac{\partial \boldsymbol{\rho}}{\partial t} \boldsymbol{\lambda}_{1} d\Omega dt$$

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```
./run
--- Adjoint test -----
inner product U/D 553123.57586755091
inner product G/Q 553123.57586756046
./run
--- Adjoint test -----
inner product U/D -75077.332007383695
inner product G/Q -75077.332007386358
./run
---- Adjoint test -----
inner product U/D 125669.89223600870
inner product U/D 125669.89223600870
```



c Model at the 100th FWI iteration

<sup>2</sup>With Bernstein-Bézier elements and AB3 time scheme <sup>3</sup>With canonical scalar product





- For AB3 scheme : DtA is slighly better than AtD
- No predominant behaviour