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# Recovering underlying graph for networks of 1D waveguides by reflectometry and transferometry

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## Abstract

We present a method for blind recovery of network made out of a tree of 1D homogeneous waveguides with the same physical characteristics using reflectogram and transferogram(s).

**Keywords:** Inverse problem, Topology recovering, Quantum graph, Reflectometry, Wire analysis

## 1 Introduction

We consider an unknown quantum graph  $\mathcal{G}$  (see [2]) equipped with a wave operator along its branches and some transmission conditions on its nodes connecting together the quantities evaluated on the branches. Our graph is a rooted tree graph, where all branches are oriented from a root-point Inp to end-points Out<sub>*k*</sub> ( $k = 1 \dots K$ ). We will consider a maximum of two consecutive nodes between Inp and any Out<sub>*k*</sub> and at least a node in  $\mathcal{G}$ .

We will now explain how a wave  $V$  propagates along the graph  $\mathcal{G}$  :

- On each branch of  $\mathcal{G}$  the wave satisfies a homogeneous wave's equation

$$\partial_{tt}^2 V - c^2 \partial_{xx}^2 V = 0,$$

where  $t$  denotes the time and  $x$  the abscissa along the considered branch. The celerity  $c$  of the waves is supposed to be a known constant.

- Following Kirchhoff's rules,  $V$  is continuous on  $\mathcal{G}$  and at each node  $J$

$$\partial_x V|_{\mathbf{e}_{j_0}}(J) = \sum_{j_k=j_1}^{j_K} \partial_x V|_{\mathbf{e}_{j_k}}(J),$$

where  $\mathbf{e}_{j_k}$  are the branches connected to  $J$ , with  $\mathbf{e}_{j_0}$  the branch closest to Inp.

- On Inp, we have an impedance boundary condition

$$\partial_t V(\text{Inp}, t) - c \frac{Z_u}{Z_c} \partial_x V(\text{Inp}, t) = (\partial_t u)(t)$$

where the constant  $Z_u$  and  $u \in H_{\text{loc}}^1(\mathbb{R})$  are known. The unknown characteristic impedance  $Z_c$  is supposed to be constant.

- At each Out<sub>*k*</sub> we have an impedance condition

$$\partial_t V + c \frac{Z_k}{Z_c} \partial_x V = 0,$$

where  $Z_k$  is an unknown constant.

## 2 Graph recovery problem

Reflectometry and transferometry methods can be applied to any practical electrical or acoustic network. The reflectogram is the following Steklov operator :

$$u(t) \mapsto \mathfrak{R}(t) := V(\text{Inp}, t),$$

whereas transferograms are operators for  $k = 1 \dots K$ :

$$u(t) \mapsto \mathfrak{T}_k(t) := V(\text{Out}_k, t).$$

We suppose that we can control  $u$ . With a known celerity  $c$  and the input load  $Z_u$ , the reflectogram and optionally some transferograms, we want to recover  $\mathcal{G}$  that is to say to determine

- the number of nodes and end-points, and their ordering (topology),
- the length  $\ell_j$  of all branches,
- the end-points load  $Z_k$ ,
- the characteristic impedance  $Z_c$ .

## 3 Injectivity

There can exist several quantum graphs with the same reflectogram, so we will make two hypotheses. Firstly, no scatterer (node or end-points) have the same distance from Inp, to ensure they can be dissociated. Secondly, no  $Z_k$  is equal to  $Z_c$  to ensure waves are reflected on the end-point. We will suppose that  $Z_k > Z_c$  is always satisfied.

## 4 Scattering

We can choose the excitation signal  $u$  such it is a peak function (sole local maximum). It propagates at celerity  $c$  along a branch until it meets a scatterer.  $\mathfrak{R}$  is the sum of attenuated  $u$ -shaped peaks, i-e of the form  $\sum_p S_p u(t - t_p)$  where where each  $(t_p, S_p)$  - called echo - corresponds to the duration and the amplitude attenuation of a propagation of  $u$  in  $\mathcal{G}$  through transmissions  $T$  and reflections  $\Gamma$  looping on Inp.  $T_k$  is similarly generated, with propagations from Inp to Out $_k$ . An echo amplitude  $S_p$  gives the nature of its contained scatterers (order for a node, load for an end-point), its abscissa the path length between the observation point (Inp for  $\mathfrak{R}$  and Out $_k$  for  $\mathfrak{T}_k$ ) and the scatterer.

The algorithm presented in [1] identifies echoes in a complex reflectogram and associates them with unknown scatterers in  $\mathcal{G}$ , giving their nature and location. It runs iteratively, dispelling ambiguities from peaks overlapping and accumulated reflections.

But this method requires knowledge of  $Z_c$  and supposes that  $Z_c = Z_u$  (no reflections at Inp). It can be enhanced by the use of transferograms.

## 5 Algorithm

### 5.1 Recovering $Z_c$

We simply recover  $Z_c$  from the reflectogram at origin where we see an echo (called mismatch echo) of amplitude  $T_u = (1 - \Gamma_u)$  with  $\Gamma_u = (Z_u - Z_c)/(Z_u + Z_c)$ . Of course if  $Z_c = Z_u$  then the mismatch peak is null.

### 5.2 Recovering the first node

The first echo observe in  $\mathfrak{R}$  after the mismatch echo have for abscissa  $2\ell_0/c$  and for amplitude  $T_u(1 + \Gamma_u)\Gamma_0$  with  $\Gamma_0 = (2/m_0 - 1)$  where  $m_0$  is the order of the first node  $J_0$ . We thus recover  $\ell_0$  and  $m_0$ .

### 5.3 Using the transferograms

If the amplitude of the first echo  $(t_1, S_1)$  of  $\mathfrak{T}_k$  is above  $4T_u(\Gamma_0 + 1)/3$ , then Out $_k$  is directly connected to  $J_0$ . Thus we have  $S_1 = T_u(2/m_0)(1 + \Gamma_k)$  with  $\Gamma_k = (Z_k - Z_c)/(Z_k + Z_c)$ , so we recover  $Z_k$ . The length of the  $J_0$  to Out $_k$  branch is  $(ct_1 - \ell_0)$ . If  $S_1$  is under  $T_u(\Gamma_0 + 1)$ , a node exists between  $J_0$  and Out $_k$ . This recovered topology

can remove branch location ambiguities for an unknown scatterer with the reflectogram.

## 5.4 Using the reflectogram

We use the algorithm developed in [1] to continue the analysis of  $\mathfrak{R}$ . We changed the procedure to use informations from the transferograms, and adapt to  $Z_u \neq Z_c$ . Indeed, we need to apply a  $T_u(1 - \Gamma_u)$  factor to all  $\mathfrak{R}$  echoes amplitude and to consider reflexions on Inp when discriminating between echoes.

This method achieves an error-free topology reconstruction if some technical hypothesis on  $u$  are fulfilled.  $\ell_j$  are retrieved with an accuracy decreasing when farther from Inp (relative error under 5% from 350 simulations), as are  $Z_c$  (under 0.1%) and  $Z_k$  (under 10%). Better determination is possible by optimizing the all lengths  $\tilde{\ell}$  and all loads  $\tilde{\mathbf{Z}}$  vectors such that they minimize the functional

$$J(\ell, \mathbf{Z}) = \int_0^{8\ell_{max}/c} \frac{|\mathfrak{R}(t) - \mathfrak{R}_{\ell, \mathbf{Z}}(t)|^2}{3\ell_{max}/c + t^2} dt$$

where  $\ell_{max}$  is the maximum  $\ell_j$  from previous steps and  $\mathfrak{R}_{\ell, \mathbf{Z}}$  the simulated reflectogram using the recovered topology with  $\ell$  and  $\mathbf{Z}$ . We look for the minimum of  $J$  with Newton's algorithm, initializing  $\tilde{\ell}$  and  $\tilde{\mathbf{Z}}$  with the previously recovered values.

## 6 Applications

This algorithm can be used for recovering unknown electrical networks, with one reflectometry device and optionnal transferometry transceivers on end-points. Removing the condition on  $Z_u$  makes the algorithm more resilient to real-life implementation limitations.

## References

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