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Focusing large spectral bandwidths through scattering media

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Abstract: Wavefront shaping is a powerful method to refocus light through a scattering medium. Its application to large spectral bandwidths or multiple wavelengths refocusing for nonlinear bio-imaging in-depth is however limited by spectral decorrelations. In this work, we demonstrate ways to access a large spectral memory of a refocus in thin scattering media and thick forward-scattering biological tissues. First, we show that the accessible spectral bandwidth through a scattering medium involves an axial spatio-spectral coupling, which can be minimized when working in a confocal geometry. Second, we show that this bandwidth can be further enlarged when working in a broadband excitation regime. These results open important prospects for multispectral nonlinear imaging through scattering media.

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1. Introduction

Nonlinear label-free bio-imaging is strongly developing as a promising tool for in situ biomedical applications. Reaching large penetration depths that surpass several hundreds of micrometers in tissues is however still a major challenge [1,2]. Adaptive optics tools are able to generate nonlinear signals at shallow depths only, since they primarily make use of ballistic photons that rapidly vanish when reaching several scattering mean free paths [3]. When highly scattering regimes dominate, other imaging strategies are required. Wavefront shaping is an interesting approach that makes use of all scattered paths to modify the initial random intensity speckle pattern into a refocus, using optimization [4] or measurement of the transmission matrix of the medium [5]. Initially developed for imaging through scattering media, this method is potentially applicable in reflection [6] which is of high importance for biological imaging. Recently, it has been used to generate efficient second harmonic generation (SHG) signals from narrow band pulses through anisotropic biological scattering media [7].

Many nonlinear optical interactions used for tissue imaging are however more complex than SHG since they involve polychromatic configurations with more than one beam, sometimes in a broadband ultra-short pulsed regime. The most frequent examples are multicolour multiphoton fluorescence imaging [8,9], four wave mixing, coherent anti stokes Raman scattering (CARS) [10], and sum frequency generation (SFG) [11]. These situations are very challenging for wavefront shaping since coherent manipulation of waves through a scattering medium is only applicable within a wavelength range below the spectral bandwidth of the medium [12–15]. This bandwidth, which represents the number of independent spectral modes of the medium, is

generally assessed by its speckle spectral correlation bandwidth [14–16] and does not surpass a few tens of nanometers in millimetre-thick biological media [17]. As a consequence, the spectral memory of the refocus formed by wavefront shaping, e.g. its robustness to spectral detuning, is limited by the medium spectral bandwidth (in what follows, we use refocus the term for a focus formed by wavefront shaping, to distinguish it from the free-space geometrical focus).

In this work, we show that two additional parameters strongly influence the spectral memory of a refocus in a non-trivial way. First, the axial spatio-spectral coupling of the waves propagation in thin or forward-scattering media can strongly reduce the accessible refocus spectral bandwidth. We evaluate the role of the axial position of the refocus plane along the propagation direction on the refocus spectral bandwidth and show that when the imaging system is in the confocal geometry (e.g. the focal points of the excitation and collection optics in free space coincide), large spectral bandwidths can be obtained with minimal axial expansion. Second, we investigate the effect of the spectral bandwidth of the incident laser beam and demonstrate that broadband refocussing is able to enlarge the accessible refocus spectral bandwidth beyond the spectral bandwidth of the medium, as a direct consequence of the coherent regime of interferences occurring when the refocus is formed.

2. 3D speckle spectral correlations

The monochromatic illumination of a scattering medium leads to the formation of a speckle, which results from interferences of spatially random nature, between the different paths undergone by light within the medium. Under monochromatic illumination with a tuneable incident wavelength, the spectral bandwidth of the refocus obtained by wavefront shaping is expected to be the same as the speckle spectral bandwidth [12]. However since in thin scattering media the spectral properties of a speckle are also strongly correlated to its spatial axial expansion [18], the axial dimension can have strong consequences on the measured speckle and refocus spectral bandwidths. The importance of this effect is studied in what follows.

The experimental set up depicted in Fig. 1a permits to measure speckle as well as refocussing spectral decorrelations in a microscopy imaging geometry. The incident beam is produced by a tunable OPO (Coherent) pumped by a picosecond laser (pump 1031 nm, 3ps/pulse, 0.1 nm spectral width, APE) expanded with a telescope. This source can be considered as close to monochromatic even for thin scattering media. Light is sent on a 256×256 pixel reflective spatial light modulator (SLM, Boulder Nonlinear Systems), conjugated to the back focal plane of the excitation objective O1 (Olympus LCPLN-IR 20x, 0.45 NA) (Fig. 1a). Note that the SLM can be considered as achromatic in the spectral ranges used in the experiment. The beam is focused on the sample plane, which is imaged by the objective O2 (Olympus PlanFLN 40x, 0.6 NA) on a CMOS camera (Blackfly U3-23S6M-C, FLIR). The sample is placed between the objectives with a fixed holder. O2 can be translated along the optical axis using a translation stage to explore the axial properties of both speckles and refocussed points. For refocussing behind the scattering medium, the transmission matrix (TM) of the medium is measured as described in [5]. The SLM is used to phase-tune part of the incident wave front, the other part serving as a reference beam. Self-reference interferometry using a wave front decomposition on Hadamard basis components is used to extract the TM matrix relation between the incident and out-going field modes [5]. We use 1024 (32×32 macropixels on the SLM) Hadamard basis components to scan the TM. Note that this TM also contains possible contributions from the illumination and detection optics. The long focal length of objective O1 (8.3 mm) allows to place the scattering medium at a chosen position from its focal plane F within millimetre distances (Fig. 1b). In addition the refocus plane R, where the TM is measured, is not necessarily the focal plane F (Fig. 1b).

Speckles measured at a few mm from a polycarbonate diffuser (Newport, 10°, Light Shaping Diffuser) are displayed in Fig. 1c for different wavelengths separated by 6 nm, starting from $\lambda_0 = 828$ nm. This wavelength shift induces visible decorrelations between the obtained speckles.

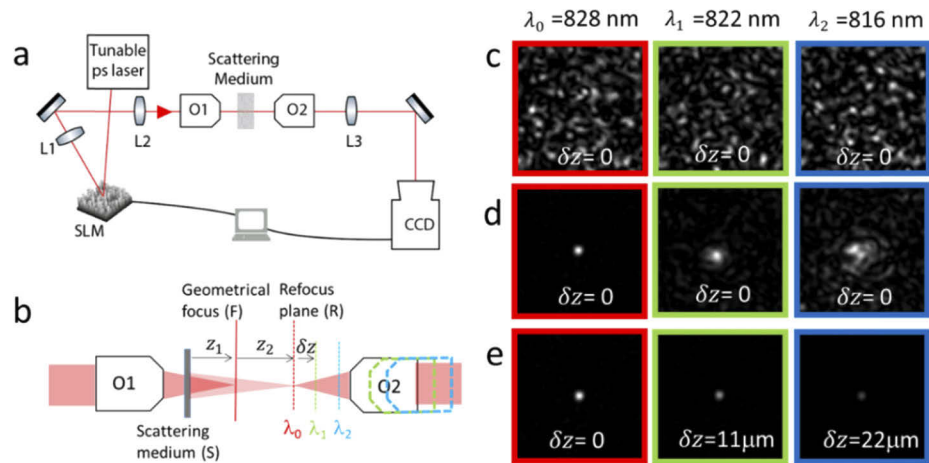


Fig. 1. (a) Experimental layout. O1, O2: objectives. SLM: spatial light modulator. L1, L2 (lenses) are used to image the SLM on the back focal plane of O1. L3 (lens) is used to image the refocus plane on the CCD camera. (b) Positioning of the sample and objectives. $z_1 = SF$ is the distance between the scattering medium (here a diffuser) and the geometrical focus (F). $z_2 = FR$ is the distance between F and the refocus plane R at λ_0 . δz is the distance between R and the new position of the refocus plane when the wavelength λ is shifted from its initial value λ_0 . (c) Speckle images obtained at the fixed plane R for several wavelengths separated by 6 nm, for $z_1 = 2$ mm and $z_2 = 1$ mm. (d) Wavefront shaping experiment producing a refocus in R. Once the TM is measured for λ_0 , the incident wavelength is tuned to λ_1 and λ_2 , producing distorted focus images at $\delta z = 0$. (e) Translating the image planes by δz allows to recover a focus at λ_1 and λ_2 .

We evaluated how this speckle decorrelation affects a refocussing experiment by wavefront shaping. After the TM of the scattering medium is measured at λ_0 to produce a refocus in the reference plane R, the wavelength is tuned. At $\lambda_1 = 822$ nm and in a more pronounced way at $\lambda_2 = 816$ nm, the refocus is seen to quickly degrade (Fig. 1d). This is expected from the observed speckle spectral decorrelation (Fig. 1c), which indicates a similar degree of decorrelation in the measured TM [12]. Importantly however, moving the O2 objective away from the R plane along the axial propagation direction allows recovering a better refocus quality (e.g. a tighter refocus), even though the enhancement performances are slightly degraded (the enhancement being defined as the ratio between the refocus maximal intensity divided by the averaged of its surrounding speckle) (Fig. 1e). Although chromatic aberrations of lenses contribute to this effect, they cannot entirely explain the large axial shifts observed. Indeed the observed axial shift is $22 \mu\text{m}$ for a wavelength shift of 12 nm, while pure chromatic effects, measured in a similar geometry without the scattering medium, lead to a spatial displacement of less than $6 \mu\text{m}$. As a reference for chromatic aberrations in the system, an axial shift dependence of $0.45 \mu\text{m}/\text{nm}$ is measured when focussing light at the same R plane without the scattering medium.

We verified that the recovery of the focus quality when shifting the image plane by δz is also accompanied by an increase of speckle correlations. Figure 2 represents the speckle spectral decorrelation curve measured at the reference refocus plane R, obtained by tuning λ away from the reference wavelength $\lambda_0 = 828$ nm. In this experiment, the scattering medium is placed a few mm's before the geometrical focus F and the plane R is close to F (Fig. 2a). Note that a limited amount of spectral points are measured per curve to guaranty the stability of the set-up during wavefront shaping and spectral/axial tuning. In the plane R, a speckle decorrelation bandwidth $\Delta\lambda$ of about 18 nm is obtained from a rough estimate of the full width half maximum of the

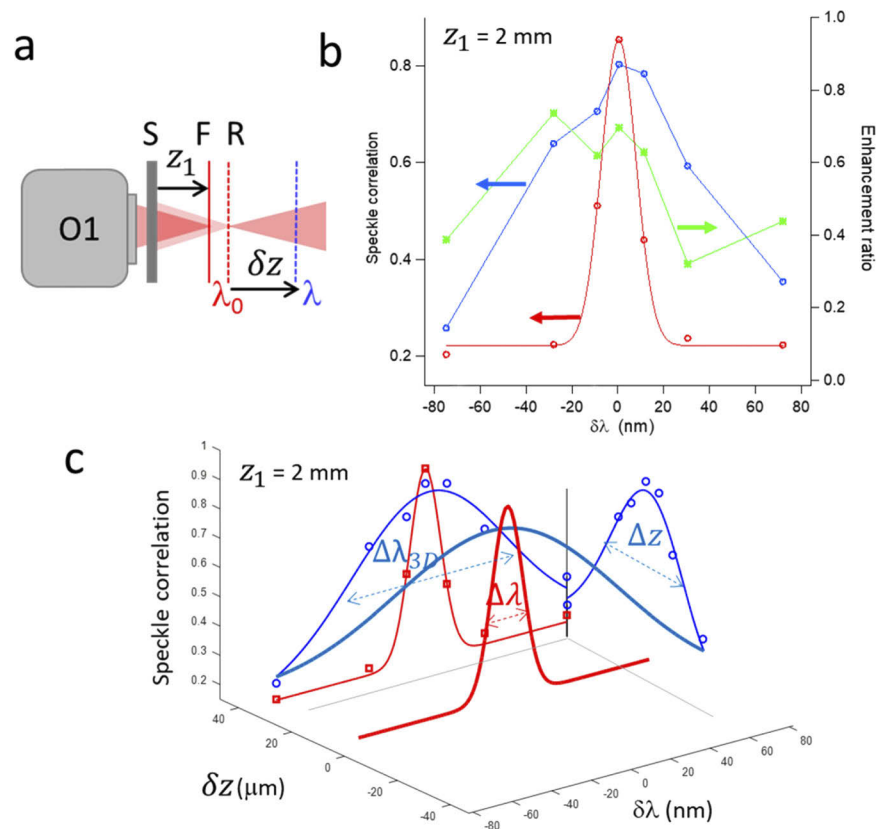


Fig. 2. Measured speckle correlations in 2D, 3D, and enhancement ratio of the refocus at the best focus plane, in a diffuser. (a) The reference R plane position is set close to the geometrical focus plane ($z_2 \sim 120 \mu\text{m}$). The scattering medium output surface is positioned at $z_1 = 2 \text{ mm}$ from the geometrical focus F. (b) Measured speckle correlations obtained when tuning the incident wavelength away from $\lambda_0 = 828 \text{ nm}$: in 2D (red markers), 3D (blue markers) and enhancement ratio at the best focus plane (green markers). is the shift of the measured wavelength with respect to the reference wavelength λ_0 . The continuous line is a fit of the 2D experimental data by a Gaussian function used here to estimate the width of the curve. The non-zero background of the curve is attributed to the used Speckle field of view. Here, the refocus 2D bandwidth is not represented due to the fast degradation of the refocus quality when changing the incident wavelength, making the estimation of the enhancement ratio imprecise. (c) Same configuration represented in a 3D plot where the translation distance δz , necessary to find the best focus at wavelengths λ , is represented as an additional axis. The speckle correlation 2D (red) and 3D (blue) experimental points (markers) and fits (continuous lines) are represented both in the λ -space and in the z -space, as well as in 3D (continuous thick lines crossing the graph). In this situation the measured spectral decorrelation bandwidths are respectively $\Delta\lambda = 18 \text{ nm}$ in 2D and $\Delta\lambda_{3D} = 87 \text{ nm}$ in 3D, while the spatial extent of the 3D bandwidth is $\Delta z = 30.5 \mu\text{m}$. Note that this size is larger than the speckle lateral and axial grain sizes (which are less than 1 and 8 μm respectively).

curve (Fig. 2b, red). When varying λ , an axial shift of the image plane away from the reference plane R permits however to recover a higher degree of correlation with the reference speckle (Fig. 2b, blue curve). We define as ‘3D bandwidth’ $\Delta\lambda_{3D}$ the speckle spectral width obtained when correlation measurements are performed at the planes with highest correlation for each tuned incident wavelength. $\Delta\lambda_{3D}$ is about 87 nm in the configuration described in Fig. 2 (Fig. 2b, blue curve), almost a factor 5 above the initial 2D bandwidth. Interestingly, the 3D spectral decorrelations curves of both speckle and refocus (Fig. 2b, green curve), are relatively similar, and it is found that the highest correlation plane for the speckle is also the plane for recovering an optimized focus when the wavelength is tuned away from λ_0 . Figure 2c displays the axial position δz of the refocus when it reaches its highest quality as the wavelength is tuned, together with its corresponding normalized enhancement. This leads to a $\delta z/\delta\lambda$ representation that shows the optimal position of a refocus for a given incident wavelength. In the present configuration, it is seen that the axial translation necessary to cover the full 3D spectral width of an optimized focus is of the order of tens of μm (Fig. 2c).

Overall, this experiment shows that both speckle and refocus spectral properties are affected by a coupling with the axial dimension z . Note that the chromatic nature of spatial speckle correlations can take complex forms, since axial and lateral dimensions contribute differently to speckle propagation. In addition to axial correlations which are the center of this work, a lateral dilatation of the speckle is also expected when λ and z are tuned, which originates from the convergent illumination of the scattering medium [18,19]. This lateral dilatation of the speckle may explain the slight discrepancy between the blue and green curves plotted in Fig. 2b. Nevertheless, lateral effects along the (x, y) transverse dimensions are non-negligible only away from the propagation axis, as previously observed in nonlinear refocus-scanning imaging [20].

3. Spatial extent of the refocus 3D bandwidth

To get further insight in the spatio-spectral characteristics of the refocus along the axial direction, we quantified the spatial extent along z of the 3D refocus spectral bandwidth, for different imaging geometries. Figure 3 shows how this spatial extent varies with the distance z_2 , for a scattering medium placed at a few mm from the geometrical focus F ($z_1 = 5$ mm) (Fig. 3a). When z_2 is a few millimeters, the full spatial extent Δz necessary to scan the full 3D spectral bandwidth $\Delta\lambda_{3D}$ reaches hundreds of μm (Fig. 3b). This spatio-spectral dependence is symmetric with respect to the reference plane position (Fig. 3c). Removing the scattering medium in similar convergent illumination shows that less than 30% of the measured displacement values is attributed to chromatic aberrations. Note that whatever the axial position of the produced refocus, its 3D spectral extent is limited by the intrinsic spectral bandwidth of the diffuser, which is measured to be around 87 nm from Fig. 2.

Figure 3 shows that increasing the distance z_2 (and similarly z_1) leads to an axial expansion of the spectral bandwidth. In contrast when the scattering medium and observation planes are at the geometric focus position ($z_1 = z_2 = 0$), e.g. when the system is set in the confocal geometry for both excitation and detection objectives, the spatial extent of the 3D spectral bandwidth is minimized. Note that removing the excitation objective to produce a planar wave illumination leads to a 3D bandwidth that cannot be confined in one given plane as expected from free space speckle propagation.

The dependence observed in Fig. 3 is attributed to the presence of spatio-spectral coupling along the axial direction of the wave propagation after the scattering medium. This coupling naturally takes the form of phase contributions in the wave propagation taking place at the exit of the medium, which contains in particular a parabolic form when a thin medium is illuminated under convergence illumination [18]. We model below the simple case of scattering by a thin diffuser to emphasize the role of spatio-spectral coupling. Under coherent convergent illumination, the formation of a speckle is the result of diffraction which takes a form that depends both on λ

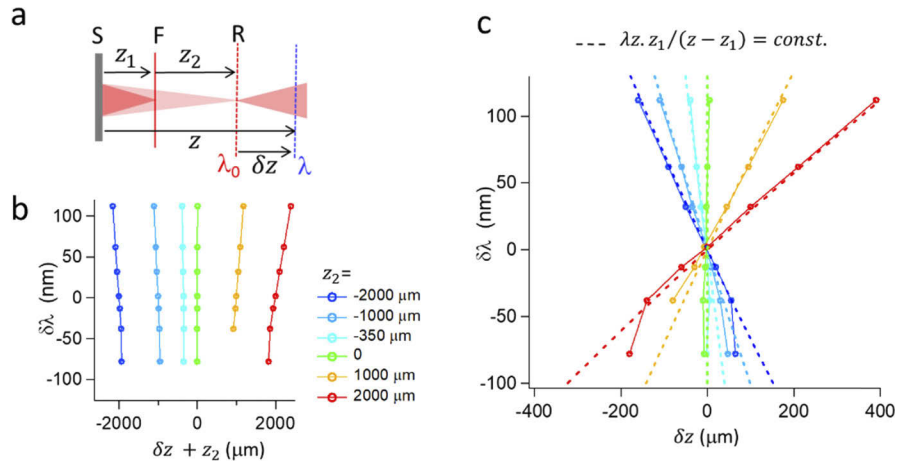


Fig. 3. Relation between wavelength detuning $\delta\lambda$ and axial distance between foci. (a) Experimental configuration showing the different planes involved, δz is the translation required to recover an optimized focus at λ (ranging from 750 nm to 940 nm), relative to the reference plane R of refocus at $\lambda_0 = 828$ nm. The distance from the scattering medium to the geometrical focus is $z_1 = 5$ mm, and the refocus plane position R is varied between $z_2 = -2$ mm and 2 mm. (b) Displacement relative to the plane R ($\delta z + z_2$) required for wavelengths shifts $\delta\lambda$ defined by ($\lambda = \lambda_0 + \delta\lambda$), for different values z_2 . (c) Same data shown as a function of δz . The dashed lines represent a function that follows the law $\lambda z z_1 / (z - z_1) = const.$

and axial propagation distance z in an intrinsically coupled manner [18]. In the Fresnel regime, the observed intensity follows the Fresnel-Kirchhoff equation for a convergent illumination [18]:

$$I(u, v) = \left| \iint G(x, y) \exp\left(\frac{-i\pi}{\lambda z} \frac{(z - z_1)}{z_1} (x^2 + y^2)\right) \exp\left(\frac{-i2\pi}{\lambda z} (ux + vy)\right) dx dy \right|^2 \quad (1)$$

With $I(u, v)$ the intensity measured in the image plane of coordinates (u, v) , placed at a distance z from the scattering thin diffuser made of a phase-only transmission mask $G(x, y)$. $z = (\delta z + z_1 + z_2)$ is the distance between the scattering medium and the observation plane for the wavelength λ (Fig. 3a).

In Eq. (1), simple spatio-spectral correlation properties can be found in two extreme cases. First, under planar wave illumination, the term in z/z_1 in the first complex exponential in Eq. (1) vanishes and the remaining term is invariant with respect to the dependence $\lambda z = const.$ [18] (where $const.$ is a constant number). This situation has permitted depth access in speckle imaging using spatio-spectral correlations [21,22]. Second, under convergent illumination and observation close to the optical axis ($(u, v) \approx (0, 0)$), the second term vanishes and a parabolic phase remains, leaving an invariance with respect to the dependence $\lambda z z_1 / (z - z_1) = const.$ In the context of this work we follow this second situation, the refocus being measured on the optical axis. Note that far from the optical axis, this equation emphasizes extra spatio-spectral dilatations in the lateral dimension (x, y) as mentioned above [18,19].

When λ or z are varied, the case $(u, v) \approx (0, 0)$ permits to evidence planes of constant spatio-spectral correlations which follow $\lambda z z_1 / (z - z_1) = const.$ This dependence is displayed in Fig. 3c together with the experimental data, fixing (z_1, z_2) as the distances used in the experiment. The interpretation of our results with the Fresnel propagation model provides a good qualitative agreement with the experimentally observed spatio-spectral coupling in thin scattering medium. In particular, it displays a symmetry with respect to $\delta z = 0$, and a minimal spatial extent of the

spectral variables when $z_2 = 0$, as observed experimentally. Note however that we assumed in the model that the transmission coefficient of the diffuser is achromatic, which is an approximation.

The results of Fig. 3 lead to two important conclusions. First, despite the presence of the diffuser, the light beam keeps the ‘memory’ of the geometrical focus of the objective lens (manifested as a parabolic phase) when the wavelength is shifted. This is expected by the fact that the scattering regime explored here is far from a multiple scattering situation, which would randomly scramble propagation phase contributions. Second, in a thin scattering medium the spectral memory of the refocus formed by wavefront shaping at a given plane depends on the axial position of this plane. It is in particular maximized when this plane coincides with the geometrical focus of the excitation optics, e.g. when the set-up is set in a confocal geometry for both excitation and detection objectives.

4. Spatial extent of the refocus 3D bandwidth in thin and anisotropic scattering media

The previous measurements were performed using a diffuser as a scattering medium, which is close to a thin scattering medium, manifested also by a large angular memory effect [23,24]. We evaluated in Fig. 4 the 2D and 3D spectral bandwidths as well as the axial expansion of the 3D bandwidth in different media, either thin or exhibiting an anisotropic scattering nature. In such cases, we expect the memory of the excitation geometrical focus to be present to some degree, which should preserve the spatio-spectral coupling effect observed above. First, a fixed 1 mm thick opaque brain slice from a mouse (coronal cross section) is used as biological sample. This sample exhibits 2D and 3D spectral bandwidths that are similar to those of a diffuser. The enhancement of the spectral bandwidth when moving from a 2D measurement to 3D is about 5 to 10 fold. The scattering mean free path of a biological tissue is about 100 μm , but its anisotropic factor g close to 1 makes its transport mean free path larger than 1 mm [25], which is the thickness explored here. We compared these values to scattering media exhibiting a lower anisotropy factor, but thin with respect to their transport mean free path. These samples are made of diluted TiO_2 particles of 500 nm size deposited on a coverslip with thicknesses 20 μm and 40 μm , embedded in a glue of refraction index close to 1.5. This permits to reduce their scattering power, and in particular to generate measurable 2D bandwidths, above the 0.1 nm spectral width of the laser used. Even though this regime is thus far from being diffusive, both the 2D and 3D bandwidths measured in the TiO_2 samples are lower than in the biological tissue, which is attributed to their lower thickness-to-transport mean free path ratio. In all reported samples, scattering is originating either from a thin or anisotropic medium. As a control a very thick TiO_2 sample is also reported, where spectral bandwidths are measured to be below 2 nm. In this case, no spatio-spectral coupling effect was measurable, as expected from a multiple scattering regime where the geometrical phase memory of the excitation wavefront is lost.

A common slope $\Delta\lambda_{3D}/\Delta z$ can be drawn out of the measurements performed in the thin and anisotropic media (Fig. 4). This slope is expected to depend only on pure geometrical factors (in particular z_1) and is dictated by Eq. (1) in thin media. The axial expansion of the spectral bandwidth therefore does not reflect the medium but rather propagation properties of the speckle generated by the medium. Figure 4 also shows that in the anisotropic scattering media studied here, which thickness is similar or below their transport mean free path, very large 3D bandwidth can be obtained. This increased capacity for spectral memory can be related to the previously reported large geometrical memory effect in such media [26], which has been generalized to a shift/tilt memory effect [27] combining angular [28] and spatial [29] memory effects.

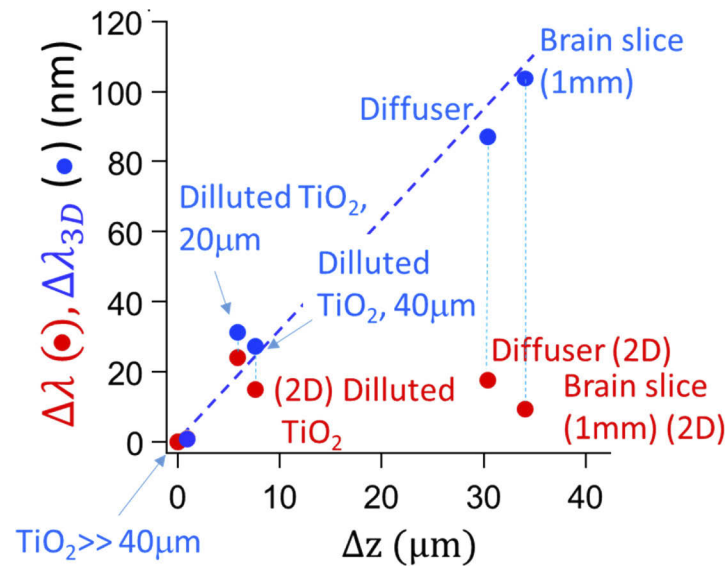


Fig. 4. 2D (red) and 3D (blue) spectral bandwidths reported as a function of the spatial extent Δz of $\Delta\lambda_{3D}$ for different media (see text) in the geometrical configuration $z_1 = 5$ mm. The refocus (reference) R plane position is set close to the scattering medium surface ($z_2 \sim 120 \mu\text{m}$). The blue dashed line is added as a guide line.

5. Enlarging the 3D refocus bandwidth using broadband wavefront shaping

All experiments above were performed close to a monochromatic regime. We explore in what follows the consequence of use of a broadband excitation pulse, which is generally required for nonlinear optical generation. A short pulse is expected to be strongly distorted at the exit of a scattering media, as a consequence of the randomization of propagation lengths in the medium. To control short pulsed laser light with large bandwidths in scattering media, several strategies have been developed, based on the direct optimization of a nonlinear signal [30,31], the control of the incident pulse in the time domain [32,33], or on the independent manipulation of different spectral channels [34]. A faster and more direct approach has recently been achieved by measuring the transmission matrix of the medium in the broadband regime [35,36]. The measured ‘broadband TM’, performed using a laser which bandwidth is close to the medium spectral bandwidth, has to be distinguished from the ‘multispectral TM’ reconstructed from a collection of individual monochromatic TM’s. The broadband TM has been shown in particular to preserve the short pulse nature of the refocus, due to an intrinsic coherent time gating effect produced by self-reference interferometry [35], provided that the medium bandwidth is comparable with the laser bandwidth. As a consequence a broadband refocus exploits favourably short propagation paths in a scattering medium [37,35]. This effect, which has been exploited for nonlinear imaging [7] and polarization-preserved imaging [37], should also have consequences on the spectral bandwidth of the refocus.

To assess the spectral memory of such a broadband refocus, we use the confocal geometry mentioned above ($z_2 = z_1 = 0$). and measure the TM of the medium under short pulsed, broadband conditions. The laser is now a tunable OPO (Coherent) pumped by a Ti:Sa oscillator (150 fs/pulse, 80 MHz) (Coherent) with a spectral width of about 8 nm. Experiments were first performed at various z_1 distances, the reference plane being at the geometrical focus of the objective O1 (Fig. 5a). The measured 3D bandwidth is considerably enhanced (Fig. 5b) and reaches $\Delta\lambda_{3D}$ values of 260 nm, about a factor 3 above the bandwidth obtained in the ps (close to

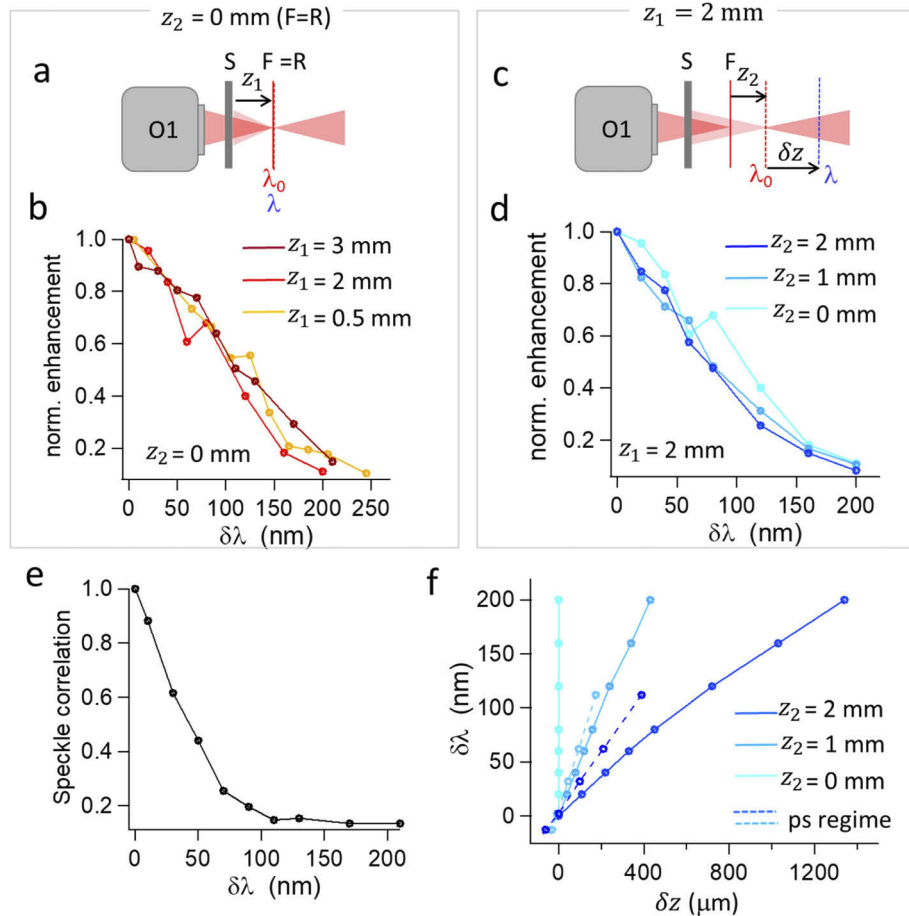


Fig. 5. Broadband refocussing in the multispectral regime, using a 150 fs source centered on $\lambda_0 = 785$ nm, and a diffuser as the scattering medium. (a) Scheme for evaluating the dependence of $\Delta\lambda_{3D}$ as a function of the z_1 (= SF) distance. The reference refocus plane R for λ_0 is at the geometrical focus position. (b) 3D spectral correlation of refocussed beams, for different SF distances. (c) Scheme for evaluating the dependence of $\Delta\lambda_{3D}$ as a function of the FR distance. (d) $\Delta\lambda_{3D}$ obtained for different FR distances ($z_1 = 2$ mm). (e) 2D speckle correlation bandwidth obtained in the confocal geometry (this bandwidth is similar to 3D speckle correlation bandwidths when the system is not in the confocal geometry). (f) Relation between spectral detuning and the axial displacement of the best refocus plane, for different z_2 (= FR) distances. The data are compared to the values obtained in Fig. 3.

monochromatic) regime. This increase is attributed to the coherent gating effect produced by the measurement of the broadband TM transmission matrix [35], which intrinsically enlarges spectral capabilities as recently demonstrated by time gating [38]. When tuning the z_2 distance (z_1 is fixed at 2 mm from the diffuser) (Fig. 5c), the obtained 3D bandwidth is very similar for all z_2 distances (Fig. 5d), confirming that the obtained 3D bandwidth value is an intrinsic property of the medium. Note that this spectral bandwidth is twice as large as the speckle correlation bandwidth (about 80 nm in the confocal geometry, Fig. 5e), which keeps a similar magnitude as the one measured in the close-to-monochromatic ps regime. This is due to the fact that while a refocus is constructed in coherent conditions, the speckle measurement is performed on intensities. At last, the spatial dependence of the spectrally large refocus (Fig. 5f) shows that the axial expansion of $\Delta\lambda_{3D}$ is minimized around F when the reference plane is at the geometrical focus, while it extends hundreds of μm away when z_2 increases. This behaviour is similar to what was observed in the ps, close-to-monochromatic regime (Fig. 5f).

6. Conclusion

We have shown that spatio-spectral coupling in thin and anisotropic scattering media has important consequences on the spectral memory of a refocus formed by wavefront shaping. The axial expansion of refocus spectral correlations is in particular minimized when working in the confocal excitation/detection geometry, leading to large spectral bandwidth limited only by the medium. This work demonstrates moreover that under broadband illumination conditions, it is possible to refocus wavelengths that are distant from each other by a large amount. This situation is typically met in four wave mixing and coherent anti Stokes Raman microscopy, or in sum frequency nonlinear excitations. In non-confocal situations, it should be possible to use a single TM for the refocussing of shifted wavelengths, making the medium characterization for further imaging much faster than in sequential-optimization based methods. That could be done by compensating for the axial shift induced by speckle correlation properties using additional correction wavefronts.

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