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Carleman estimate for a 1D linear elastic problem involving interfaces. Application to an inverse problem.

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Abstract

The aim of this work is to study the stability of the reconstruction of some mechanical parameters that characterize interfaces within linear elastic bodies. The model considered consists of two bonded elastic solids. The adhesive elastic layer between them is a thin interphase that is approximated, by asymptotic analysis, as an interface. The resulting interface model is in turn characterized by a typical spring-type linear elastic behavior. In this context and on a 1D configuration, we investigate a Carleman-type estimate that relates to a unique continuation principle. The proof is based on the construction of suitable weight functions: their gradients are non-zero, the jumps of the derivatives are positive across the interface and the averages of the derivatives vanish. The design of such weight functions enables the control of the interface terms in the estimates. With this at hand, we establish a Lipschitz stability estimate for the inverse problem of identifying an interface stiffness parameter from measurements that are available on both sides of the external boundary.

Keywords: interface stiffness parameter, 1D elasticity, Carleman's estimate, Lipschitz stability

1 Introduction

The study of solid interfaces in structural problems has garnered much scientific attention due to its meaningful role in several engineering fields (e.g. civil engineering, mechanical engineering, biomechanics, geomechanics,...). Within the framework of those studies on mechanics of composite materials, a common assumption is adopted, i.e. the continuity of the displacement and the stresses along with the interface. Nevertheless, the assumption of perfect interfaces is generally inadequate in describing real-life engineering problems. These limitations originate the consideration of imperfect interfaces and enhance their effect on the determination of the effective elastic properties in micro-cracked media [6, 13].

Aim of this study is the identification of some mechanical parameters that characterize linear elastic interfaces in mechanical structures. More precisely, the studied model corresponds to the simple problem of two bonded elastic solids. The adhesive elastic layer between them is thin material interphase that is modelled as an interface using appropriate asymptotic analysis. The resulting interface model is in turn characterized by a typical spring-type linear behaviour. Namely, the interface tractions are assumed to be continuous, but the displacement can be discontinuous across the debonded region and, as the case may be, proportional to the interface traction. This parametric inverse problem is treated from two aspects: a numerical one by developing an algorithm aiming at the quantification reconstruction of the interfaces elastic stiffnesses, which is developed in a previous work [11] and a theoretical one by quantifying the stability of the parameters, which is presented in this paper.

Various forms of stability estimates for the Robin coefficients are developed in the literature: Both local and monotone global Lipschitz stability for regular Robin coefficient are provided by Chaabane et al. [7] under the assumption that the flux is non negative, Lipschitz stability estimate for piecewise constant Robin coefficient is proved by Sincich [14], monotone global Lipschitz stability is developed by Alessandrini et al. [1] using the three-spheres-inequalities as the main tool and also one can find global logarithmic type [3, 8] based on the general Carleman inequality proven by Phung [12]. We refer to the following review [15] for more details about this kind of inequality. Although, Belhachmi et al. [2] considered a Laplace equation with a Robin transmission condition, i.e. the potential is continuous and the flux is discontinuous, a dual problem for the one considered here. The proof of their directional Lipschitz stability has the same lines as the one developed in [7]. Jaoua et al. [10] proved Local Lipschitz stability for an impedance defined on a crack. Up to our knowledge, there are no studies in the literature concerning the stability issue of the interfaces elastic stiffnesses.

In this framework, we focus on the toy problem of the imperfect interface problem as a first step, avoiding the difficulties encountered with the complexity of the mathematical resolution. Therefore, we are interested in the parametric inverse problem in the case of one-dimensional elasticity equation with a spring-type linear law, more precisely, in the stability issue of an interface elastic stiffness defined on a known point under the assumption that the data are available on both sides of the external boundaries. Roughly speaking, we seek some information on the stiffness parameters from the measurements. The reason for asking this question is that the experiment's measurements are naturally subjected to errors, so that, we are looking for a deviation of the measurements from the variations of the parameters. Thus, we focus on establishing a Lipschitz stability estimate, under a priori assumption on the measurements, based on Carleman's inequalities. To do so, we prove a Carleman estimate for the problem under study by reconstructing an appropriate weight function. Carleman estimates for parabolic equations with standard transmission conditions in the one-dimensional case were proven by Benabdallah et al. [4, 5]. The proof is based on the construction of suitable weight functions, whose gradient is non-zero in the observation region and the jump of the derivative is positive across the interface. Here, we impose to these constraints that the average of the derivative of the weight function vanishes as a new parameter which enables to control the interface terms in the derivation of the Carleman-type estimate. Carleman's estimates in the 2D case is still an open question.

This article is outlined as follows. In the forthcoming, the forward and inverse problems of interest are reviewed and the main results are stated. Carleman inequality is established in section 3. Subsection 3.3 details the choice of the weight function by writing the quadratic form in terms of the jump and the average of the weight function gradient. The fourth section presents the stability result, based on the established Carleman inequality, for the interface stiffness parameter. The closing section is devoted to some comment.

2 Problem setting and main results

2.1 The forward problem

Consider $\Omega = (0, 1) \subset \mathbb{R}$ be a bounded open set and a be a point in $(0, 1)$. The point a split the reference domain Ω into two sub-domains Ω^\pm being defined as $\Omega^- = (0, a)$, $\Omega^+ = (a, 1)$. These sub-domains are characterized by the corresponding Young's moduli E^\pm which are elements of

$$\{E^\pm \in C^1(\Omega^\pm); 0 < e_1 \leq E^\pm(x) \leq e_2, \forall x \in \Omega^\pm\}.$$

Also, we denote $\Omega' = \Omega^+ \cup \Omega^-$, $S = \{a\}$ and $\Gamma = \{0, 1\}$.

The one-dimensional elasticity problem considered is represented by the following elliptic boundary value problem that involves a spring-type linear interface law. More precisely, the field q satisfies

$$-\partial_x(E\partial_x q)(x) = f(x), \quad x \in \Omega'. \quad (1)$$

Moreover, the interface flux is assumed to be continuous, but the field can be discontinuous across the singular point a and, as the case may be, proportional to the interface flux. These conditions are represented as follows

$$\llbracket E\partial_x q \rrbracket = 0, \quad \llbracket q \rrbracket = k\langle E\partial_x q \rangle, \quad (2)$$

where the jump of a quantity q across a point a being defined as $\llbracket q \rrbracket = q(a^+) - q(a^-)$, also the average of q is defined by $\langle q \rangle = \frac{1}{2}(q(a^+) + q(a^-))$ and k is a positive real-valued parameter. Besides, we impose Neumann boundary conditions on the boundary Γ .

2.2 The inverse problem

We are interested in the inverse problem of recovering a stiffness parameter k defined on the singular point a from available data on the boundary Γ in the one-dimensional elasticity problem. Supposing q_i as a solution satisfying (1,2) associated to $k = k_i$, for $i = 1, 2$ and also $f(x) = 0$ for all $x \in \Omega'$. Here to achieve a stability result, we have to derive a Carleman-type estimate for the difference between two solutions, q_1 and q_2 . This difference is the solution of a non-homogeneous elliptic equation, because of the difference of the stiffness constant it does not satisfy the appropriate spring-type linear interface law, on the singular point a , defined above.

For this reason, we derive a peculiar Carleman-type estimate which includes additional interface terms.

$$\llbracket E\partial_x q \rrbracket = 0, \quad \llbracket q \rrbracket = k_2\langle E\partial_x q \rangle + g, \quad (3)$$

where $g = (k_1 - k_2)\langle E\partial_x q_1 \rangle$.

2.3 Statement of the main results

Here, we state our main results developed in this work. In first time, we present the following lemma specifying the suitable weight function which enables to control the interface terms in the derivation of Carleman's estimate.

Lemma 1. *There exists a continuous weight function β satisfying*

$$\llbracket \beta' \rrbracket > 0, \quad \langle \beta' \rangle = 0,$$

for which the matrix \mathcal{M} of the quadratic form is definite and positive, more precisely

$$\det \mathcal{M} \simeq_{s\lambda \gg} c_1^2 c^4 \left(-4r(er - 1)^2 + 8r(er + 1)^2 - 4r \frac{(er + 1)^4}{(er - 1)^2} \right) = 0,$$

where $\llbracket \beta' \rrbracket = c(r - 1)$, $\langle \beta' \rangle = \frac{c}{2}(r + 1)$, $c = \beta'(a^-)$, $r = \frac{\beta'(a^+)}{\beta'(a^-)}$, $c_1 = E^-$ and $e = \frac{E^+}{E^-}$. For example, β has the following shape (see Figure 1)

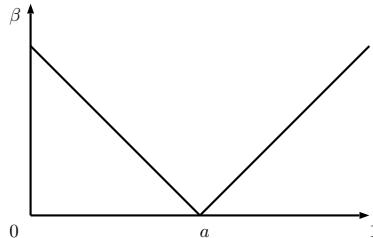


Figure 1: Typical shape's sketch for β .

A Carleman-type estimate for linear elastic interfaces is established. This Carleman estimate would be used later for the uniqueness and stability of the inverse solution of the parameter identification problem under study.

Theorem 1. *There exist $\lambda_0 > 0$, $s_0 > 0$ and a continuous weight function β satisfying the following assumptions*

$$\llbracket \beta' \rrbracket > 0, \quad \langle \beta' \rangle = 0,$$

so that the following estimate holds

$$\begin{aligned} & \|\mathcal{P}_1\psi\|_{L^2(\Omega')}^2 + \|\mathcal{P}_2\psi\|_{L^2(\Omega')}^2 + s\lambda^2 \int_{\Omega'} \varphi e^{2s\varphi} |\partial_x q|^2 dx + s^3\lambda^4 \int_{\Omega'} \varphi^3 e^{2s\varphi} |q|^2 dx \\ & \leq C \|e^{s\varphi} f\|_{L^2(\Omega')}^2 + \frac{s\lambda\varphi(a)}{k^2} e^{2s\varphi(a)} |g|^2 + s\lambda\varphi(1) e^{2s\varphi(1)} |E\partial_x q|^2(1) + s\lambda\varphi(0) e^{2s\varphi(0)} |E\partial_x q|^2(0) \\ & + s^3\lambda^3 \varphi^3(1) e^{2s\varphi(1)} |q|^2(1) + s^3\lambda^3 \varphi^3(0) e^{2s\varphi(0)} |q|^2(0), \end{aligned}$$

for all $\lambda \geq \lambda_0$, $s \geq s_0$ and for all q piece-wise smooth field satisfying (1,2).

The stability estimate is a fundamental result in the inverse problem which immediately yields the uniqueness result. Stability estimates are so meaningful from the theoretical aspects and also form the numerical one, more precisely this type of estimates specifies the property of the numerical algorithms. The stability estimate obtained here is based on the main global Carleman inequality which is valid only in one-dimensional.

Theorem 2. *We assume that Ω' is a bounded open subset of \mathbb{R} with a boundary Γ . We denote by q_i the solution of (1,2) associated with $k = k_i$ such that $q_i \in H^1(\Omega')$, for $i = 1, 2$. Let $m > 0$ be a constant such that $\langle E\partial_x q_1 \rangle \geq m$ on the point a . Then, there exists a positive constant \mathcal{C} such that*

$$|k_1 - k_2|^2 \leq \mathcal{C} (|q_1 - q_2|^2(1) + |q_1 - q_2|^2(0)).$$

This result is a very important step to numerically solve the parameters identification problem in elastic interfaces because it provides the condition in which this problem is stable numerically.

3 Carleman inequality

For simplicity, we divide the proof into five steps.

3.1 Change of variables and outlines

In this step, we set the differential equation satisfied by a new function ψ , which will be q up to a weight function. Thus, we introduce

$$\psi(x) = e^{s\varphi(x)} q(x), \quad \mathcal{P}\psi(x) = -e^{s\varphi(x)} \partial_x (E\partial_x (e^{-s\varphi(x)} \psi)),$$

where s is a positive constant and φ is a positive function defined as a weight function as follows

$$\varphi(x) = e^{\lambda\beta(x)}.$$

Let us assume that we already find a weight function $\beta(x)$. So that, we have the following relations in Ω' holds

$$\partial_x \varphi = \lambda\beta' \varphi.$$

We can rewrite the right-hand side of (1) by

$$\int_{\Omega'} f^2 e^{2s\varphi} dx = \int_{\Omega'} |\mathcal{P}\psi(x)|^2 dx.$$

Therefore, our task is a lower estimate of $\|\mathcal{P}\psi\|_{L^2(\Omega')}^2$. Direct calculations yield

$$\mathcal{P}\psi = -\partial_x(E\partial_x\psi) + 2s\lambda\varphi E\beta' \partial_x\psi - s^2\lambda^2\varphi^2 (\beta')^2 E\psi + s\lambda\varphi (E\beta')' \psi + s\lambda^2\varphi E(\beta')^2 \psi.$$

For obtaining a lower estimate for $\|\mathcal{P}\psi\|_{L^2(\Omega')}^2$, we decompose the operator \mathcal{P} into the symmetric part \mathcal{P}_+ and the anti-symmetric part \mathcal{P}_- such that $\mathcal{P} = \mathcal{P}_+ + \mathcal{P}_-$.

We define the adjoint operator \mathcal{P}^* as follows

$$\mathcal{P}^*\psi = -\partial_x(E\partial_x\psi) - 2s\lambda\varphi E\beta' \partial_x\psi - s^2\lambda^2\varphi^2 (\beta')^2 E\psi - s\lambda\varphi (E\beta')' \psi - s\lambda^2\varphi E(\beta')^2 \psi.$$

We define the symmetric part \mathcal{P}_+ and the anti-symmetric part \mathcal{P}_- of \mathcal{P} by

$$\mathcal{P}_+ = \frac{1}{2}(\mathcal{P} + \mathcal{P}^*), \quad \mathcal{P}_- = \frac{1}{2}(\mathcal{P} - \mathcal{P}^*).$$

Then, we have $\mathcal{P}\psi = \mathcal{P}_+\psi + \mathcal{P}_-\psi$ and

$$\mathcal{P}_+\psi = -\partial_x(E\partial_x\psi) - s^2\lambda^2\varphi^2 (\beta')^2 E\psi$$

and

$$\mathcal{P}_-\psi = 2s\lambda\varphi E\beta' \partial_x\psi + s\lambda\varphi (E\beta')' \psi + s\lambda^2\varphi E(\beta')^2 \psi.$$

To get the desired estimate, we take into consideration other changes in the above decomposition as follows

$$\begin{aligned} \mathcal{P}_1\psi &= \mathcal{P}_+\psi = -\partial_x(E\partial_x\psi) - s^2\lambda^2\varphi^2 (\beta')^2 E\psi, \\ \mathcal{P}_2\psi &= 2s\lambda\varphi E\beta' \partial_x\psi + 2s\lambda^2\varphi E(\beta')^2 \psi, \\ f_s &= e^{s\varphi} f - s\lambda\varphi (E\beta')' \psi + s\lambda^2\varphi E(\beta')^2 \psi. \end{aligned}$$

Hence

$$\begin{aligned} \|f_s\|_{L^2(\Omega')}^2 &= \|\mathcal{P}_1\psi + \mathcal{P}_2\psi\|_{L^2(\Omega')}^2 = \|\mathcal{P}_1\psi\|_{L^2(\Omega')}^2 + \|\mathcal{P}_2\psi\|_{L^2(\Omega')}^2 + 2(\mathcal{P}_1\psi, \mathcal{P}_2\psi)_{L^2(\Omega')} \\ &\geq 2(\mathcal{P}_1\psi, \mathcal{P}_2\psi)_{L^2(\Omega')}. \end{aligned}$$

In the following steps, we see that the definition we have made of φ makes $(\mathcal{P}_1\psi, \mathcal{P}_2\psi)_{L^2(\Omega')}$ positive up to several terms that can be controlled whenever we make an appropriate choice of the parameters s and λ .

More precisely, in the next step, we make the computations of the double products $(\mathcal{P}_1\psi, \mathcal{P}_2\psi)_{L^2(\Omega')}$. This will give an inequality with two global terms of $|\psi|^2$ and $|\partial_x\psi|^2$ and also an interfacial term on the left-hand side, while the observations terms will appear on the right-hand side. Finally, we will turn back to the original function q and deduce the Carleman inequality.

3.2 First estimates

In this step, we estimate the right-hand side of the desired estimate from below by means of $(\mathcal{P}_+\psi, \mathcal{P}_-\psi)_{L^2(\Omega')}$. For this, we develop the four terms appearing in $(\mathcal{P}_1\psi, \mathcal{P}_2\psi)_{L^2(\Omega')}$ and we have

$$\begin{aligned} 2(\mathcal{P}_1\psi, \mathcal{P}_2\psi)_{L^2(\Omega')} &= 2(-\partial_x(E\partial_x\psi) - s^2\lambda^2\varphi^2 (\beta')^2 E\psi, 2s\lambda\varphi E\beta' \partial_x\psi + 2s\lambda^2\varphi E(\beta')^2 \psi)_{L^2(\Omega')} \\ &= 2(-\partial_x(E\partial_x\psi), 2s\lambda\varphi E\beta' \partial_x\psi)_{L^2(\Omega')} + 2(-\partial_x(E\partial_x\psi), 2s\lambda^2\varphi E(\beta')^2 \psi)_{L^2(\Omega')} \\ &\quad + 2(-s^2\lambda^2\varphi^2 (\beta')^2 E\psi, 2s\lambda\varphi E\beta' \partial_x\psi)_{L^2(\Omega')} \\ &\quad + 2(-s^2\lambda^2\varphi^2 (\beta')^2 E\psi, 2s\lambda^2\varphi E(\beta')^2 \psi)_{L^2(\Omega')} \\ &= 2 \sum_{i,j=1}^2 \mathcal{I}_{ij}. \end{aligned}$$

We reduce the orders of derivatives of ψ by the integration by parts, so that we have the term \mathcal{I}_{11} , with the consideration $\partial_x(E\partial_x\psi)E\partial_x\psi = \frac{1}{2}\partial_x(|E\partial_x\psi|^2)$, as follows

$$\begin{aligned}\mathcal{I}_{11} &= -2s\lambda \int_{\Omega'} \varphi \beta' \partial_x(E\partial_x\psi) E \partial_x\psi \, dx \\ &= -s\lambda \int_{\Omega'} \varphi \beta' \partial_x(|E\partial_x\psi|^2) \, dx \\ &= -s\lambda [\varphi \beta' |E\partial_x\psi|^2]_{S\cup\Gamma} + s\lambda^2 \int_{\Omega'} \varphi (\beta')^2 |E\partial_x\psi|^2 \, dx + s\lambda \int_{\Omega'} \varphi \beta'' |E\partial_x\psi|^2 \, dx,\end{aligned}$$

where, for a quantity p ,

$$[p]_{S\cup\Gamma} = p(1) - p(0) - \llbracket p \rrbracket.$$

The term \mathcal{I}_{11} is given by

$$\mathcal{I}_{11} = -s\lambda [\varphi \beta' |E\partial_x\psi|^2]_{S\cup\Gamma} + s\lambda^2 \int_{\Omega'} \varphi (\beta')^2 |E\partial_x\psi|^2 \, dx + \mathcal{X}_{11},$$

where

$$\mathcal{X}_{11} = s\lambda \int_{\Omega'} \varphi \beta'' |E\partial_x\psi|^2 \, dx.$$

Using the fact that $\partial_x\varphi = \lambda\beta'\varphi$, the term \mathcal{I}_{12} is given by

$$\begin{aligned}\mathcal{I}_{12} &= -2s\lambda^2 \int_{\Omega'} \varphi E(\beta')^2 \psi \partial_x(E\partial_x\psi) \, dx \\ &= -2s\lambda^2 [\varphi E(\beta')^2 \psi E\partial_x\psi]_{S\cup\Gamma} + 2s\lambda^3 \int_{\Omega'} \varphi E^2(\beta')^3 \psi \partial_x\psi \, dx + 2s\lambda^2 \int_{\Omega'} \varphi (E(\beta')^2)' \psi E\partial_x\psi \, dx \\ &\quad + 2s\lambda^2 \int_{\Omega'} \varphi (\beta')^2 |E\partial_x\psi|^2 \, dx \\ &= 2s\lambda^2 \int_{\Omega'} \varphi (\beta')^2 |E\partial_x\psi|^2 \, dx + \mathcal{X}_{12},\end{aligned}$$

where

$$\mathcal{X}_{12} = -2s\lambda^2 [\varphi E(\beta')^2 \psi E\partial_x\psi]_{S\cup\Gamma} + 2s\lambda^3 \int_{\Omega'} \varphi E^2(\beta')^3 \psi \partial_x\psi \, dx + 2s\lambda^2 \int_{\Omega'} \varphi (E(\beta')^2)' \psi E\partial_x\psi \, dx.$$

Next, noting that $2\psi \partial_x\psi = \partial_x(|\psi|^2)$ and by an integration by parts, the term \mathcal{I}_{21} is given by

$$\begin{aligned}\mathcal{I}_{21} &= -2s^3\lambda^3 \int_{\Omega'} \varphi^3 E^2(\beta')^3 \psi \partial_x\psi \, dx \\ &= -s^3\lambda^3 \int_{\Omega'} \varphi^3 E^2(\beta')^3 \partial_x(|\psi|^2) \, dx \\ &= -s^3\lambda^3 [\varphi^3 E^2(\beta')^3 |\psi|^2]_{S\cup\Gamma} + 3s^3\lambda^4 \int_{\Omega'} \varphi^3 E^2(\beta')^4 |\psi|^2 \, dx + s^3\lambda^3 \int_{\Omega'} \varphi^3 (E^2(\beta')^3)' |\psi|^2 \, dx \\ &= -s^3\lambda^3 [\varphi^3 E^2(\beta')^3 |\psi|^2]_{S\cup\Gamma} + 3s^3\lambda^4 \int_{\Omega'} \varphi^3 E^2(\beta')^4 |\psi|^2 \, dx + \mathcal{X}_{21},\end{aligned}$$

where

$$\mathcal{X}_{21} = s^3\lambda^3 \int_{\Omega'} \varphi^3 (E^2(\beta')^3)' |\psi|^2 \, dx.$$

The term \mathcal{I}_{22} is given by

$$\mathcal{I}_{22} = -2s^3\lambda^4 \int_{\Omega'} \varphi^3 E^2(\beta')^4 |\psi|^2 \, dx.$$

Therefore, noting that s^3 and s are the maximum orders respectively among all the terms ψ^2 and $|\partial_x \psi|^2$ and estimating any lower order terms as non-principal terms, we have

$$\begin{aligned} (\mathcal{P}_1 \psi, \mathcal{P}_2 \psi)_{L^2(\Omega')} &= -s\lambda [\varphi \beta' |E\partial_x \psi|^2]_{S \cup \Gamma} + s\lambda^2 \int_{\Omega'} \varphi (\beta')^2 |E\partial_x \psi|^2 dx + \mathcal{X}_{11} \\ &\quad + 2s\lambda^2 \int_{\Omega'} \varphi (\beta')^2 |E\partial_x \psi|^2 dx + \mathcal{X}_{12} \\ &\quad - s^3 \lambda^3 [\varphi^3 E^2(\beta')^3 |\psi|^2]_{S \cup \Gamma} + 3s^3 \lambda^4 \int_{\Omega'} \varphi^3 E^2(\beta')^4 |\psi|^2 dx + \mathcal{X}_{21} \\ &\quad - 2s^3 \lambda^4 \int_{\Omega'} \varphi^3 E^2(\beta')^4 |\psi|^2 dx. \end{aligned}$$

Adding all the terms together to the form $(\mathcal{P}_1 \psi, \mathcal{P}_2 \psi)_{L^2(\Omega')}$ in $\|f_s\|_{L^2(\Omega')}^2 = \|\mathcal{P}_1 \psi\|_{L^2(\Omega')}^2 + \|\mathcal{P}_2 \psi\|_{L^2(\Omega')}^2 + 2(\mathcal{P}_1 \psi, \mathcal{P}_2 \psi)_{L^2(\Omega')}$, we have

$$\begin{aligned} &\|\mathcal{P}_1 \psi\|_{L^2(\Omega')}^2 + \|\mathcal{P}_2 \psi\|_{L^2(\Omega')}^2 + 6s\lambda^2 \int_{\Omega'} \varphi (\beta')^2 |E\partial_x \psi|^2 dx + 2s^3 \lambda^4 \int_{\Omega'} \varphi^3 E^2(\beta')^4 |\psi|^2 dx \\ &- 2s\lambda [\varphi \beta' |E\partial_x \psi|^2]_{S \cup \Gamma} - 2s^3 \lambda^3 [\varphi^3 E^2(\beta')^3 |\psi|^2]_{S \cup \Gamma} = \|f_s\|_{L^2(\Omega')}^2 - 2\{\mathcal{X}_{11} + \mathcal{X}_{12} + \mathcal{X}_{21}\}. \end{aligned}$$

3.3 Choice of a weight function

In this step, we focus our attention on the jump term at the singular point a on the left-hand side of the above equality to find the appropriate weight function β for μ as a positive definite quadratic form

$$\mu := s\lambda [\varphi \beta' |E\partial_x \psi|^2] + s^3 \lambda^3 [\varphi^3 E^2(\beta')^3 |\psi|^2].$$

We can write the quadratic form μ in a matrix form as follows

$$\mu = \Psi^T \cdot \mathcal{M} \cdot \Psi,$$

where Ψ is the following vector

$$\Psi = ((E\partial_x \psi)(a^-), s\lambda(\varphi\psi)(a^-))^T.$$

We develop the jump term

$$\begin{aligned} \mu &= s\lambda \varphi(a^+) \beta'(a^+) |E\partial_x \psi|^2(a^+) - s\lambda \varphi(a^-) \beta'(a^-) |E\partial_x \psi|^2(a^-) \\ &\quad + s^3 \lambda^3 \varphi^3(a^+) (E^2(\beta')^3)(a^+) |\psi|^2(a^+) - s^3 \lambda^3 \varphi^3(a^-) (E^2(\beta')^3)(a^-) |\psi|^2(a^-). \end{aligned}$$

From this expression, we remark that the interface information is available from one side, not from both sides of the singularity, as not the case of the interface condition. To avoid this difficulty, we propose to write the quadratic form μ in terms of the average of ψ and $\partial_x \psi$.

Remark 1. *First, we remark that from this point of calculus, we are not in the same case as the work of Benabdallah et al. [5]. Second, we adopt our proposed idea for the standard transmission conditions as [5] and we get the following forms*

$$[\psi] = 0, \quad [E\partial_x \psi] = s\lambda \varphi [E\beta'] \langle \psi \rangle$$

and

$$\mu := s\lambda \varphi [E\beta'] \langle E\partial_x \psi \rangle^2 + 2s^2 \lambda^2 \varphi^2 \langle \beta' \rangle [E\beta'] \langle \psi \rangle \langle E\partial_x \psi \rangle + s^3 \lambda^3 \varphi^3 \left\{ [E^2(\beta')^3] + \frac{1}{4} [E\beta'] [E\beta']^2 \right\} \langle \psi \rangle^2,$$

which gives the same choice of weight's function as [5].

Remark 2. We did the calculations for β continuous and even discontinuous and we remark the same behaviour for both cases. To simplify the calculations, we consider β as a continuous weight function at the point a .

Now, we adopt this idea for our case and for that we develop these formulas below

$$k\langle E\partial_x q \rangle = -s\lambda k\varphi e^{-s\varphi} \langle E\beta' \psi \rangle + e^{-s\varphi} k \langle E\partial_x \psi \rangle, \quad \llbracket E\partial_x q \rrbracket = -s\lambda\varphi e^{-s\varphi} \llbracket E\beta' \psi \rrbracket + e^{-s\varphi} \llbracket E\partial_x \psi \rrbracket.$$

So that, we write the spring-type linear interface law verified by ψ at the point a using (3) under the condition $\llbracket E\beta' \rrbracket \neq 0$ and denoting $\mathcal{Q} = (1 + \frac{1}{4}s\lambda k\varphi \llbracket E\beta' \rrbracket)^{-1}$

$$\llbracket \psi \rrbracket = -s\lambda k\varphi \mathcal{Q} \langle E\beta' \rangle \langle \psi \rangle + k\mathcal{Q} \langle E\partial_x \psi \rangle + \mathcal{Q}g_s$$

and

$$\llbracket E\partial_x \psi \rrbracket = (s\lambda\varphi \llbracket E\beta' \rrbracket - s^2\lambda^2 k\varphi^2 \mathcal{Q} \langle E\beta' \rangle^2) \langle \psi \rangle + s\lambda k\varphi \mathcal{Q} \langle E\beta' \rangle \langle E\partial_x \psi \rangle + s\lambda\varphi \mathcal{Q} \langle E\beta' \rangle g_s,$$

where $g_s = e^{s\varphi} g$ and we use the fact that

$$\llbracket qp \rrbracket = \llbracket q \rrbracket \langle p \rangle + \langle q \rangle \llbracket p \rrbracket, \quad \langle qp \rangle = \frac{1}{4} \llbracket q \rrbracket \llbracket p \rrbracket + \langle q \rangle \langle p \rangle.$$

After, we need the following formula in the calculation of the form μ

$$\begin{aligned} \llbracket \psi \rrbracket^2 &= s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2 \langle E\beta' \rangle^2 \langle \psi \rangle^2 + k^2 \mathcal{Q}^2 \langle E\partial_x \psi \rangle^2 - 2s\lambda k^2\varphi \mathcal{Q}^2 \langle E\beta' \rangle \langle \psi \rangle \langle E\partial_x \psi \rangle \\ &\quad + \mathcal{Q}^2 |g_s|^2 - 2s\lambda k\varphi \mathcal{Q}^2 \langle E\beta' \rangle \langle \psi \rangle g_s + 2k\mathcal{Q}^2 \langle E\partial_x \psi \rangle g_s \end{aligned}$$

and

$$\begin{aligned} \llbracket E\partial_x \psi \rrbracket^2 &= (s^2\lambda^2\varphi^2 \llbracket E\beta' \rrbracket^2 + s^4\lambda^4 k^2\varphi^4 \mathcal{Q}^2 \langle E\beta' \rangle^4 - 2s^3\lambda^3 k\varphi^3 \mathcal{Q} \llbracket E\beta' \rrbracket \langle E\beta' \rangle^2) \langle \psi \rangle^2 \\ &\quad + s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2 \langle E\beta' \rangle^2 \langle E\partial_x \psi \rangle^2 \\ &\quad + 2(s^2\lambda^2 k\varphi^2 \mathcal{Q} \llbracket E\beta' \rrbracket \langle E\beta' \rangle - s^3\lambda^3 k^2\varphi^3 \mathcal{Q}^2 \langle E\beta' \rangle^3) \langle \psi \rangle \langle E\partial_x \psi \rangle + s^2\lambda^2\varphi^2 \mathcal{Q}^2 \langle E\beta' \rangle^2 |g_s|^2 \\ &\quad + 2(s^2\lambda^2\varphi^2 \mathcal{Q} \llbracket E\beta' \rrbracket \langle E\beta' \rangle - s^3\lambda^3 k\varphi^3 \mathcal{Q}^2 \langle E\beta' \rangle^3) \langle \psi \rangle g_s + 2s^2\lambda^2 k\varphi^2 \mathcal{Q}^2 \langle E\beta' \rangle^2 \langle E\partial_x \psi \rangle g_s. \end{aligned}$$

Below, we write the form μ as a quadratic form with the fact that

$$\llbracket p^2 \rrbracket = 2\llbracket p \rrbracket \langle p \rangle, \quad \langle p^2 \rangle = \frac{1}{4} \llbracket p \rrbracket^2 + \langle p \rangle^2.$$

So, the form μ is as following

$$\begin{aligned} \mu &= s\lambda\varphi \llbracket \beta' \rrbracket |E\partial_x \psi|^2 + s^3\lambda^3\varphi^3 \llbracket E^2(\beta')^3 \rrbracket |\psi|^2 \\ &= s\lambda\varphi \llbracket \beta' \rrbracket \langle |E\partial_x \psi|^2 \rangle + s\lambda\varphi \langle \beta' \rangle \llbracket |E\partial_x \psi|^2 \rrbracket + s^3\lambda^3\varphi^3 \llbracket E^2(\beta')^3 \rrbracket \langle |\psi|^2 \rangle + s^3\lambda^3\varphi^3 \langle E^2(\beta')^3 \rangle \llbracket |\psi|^2 \rrbracket \\ &= \frac{1}{4} s\lambda\varphi \llbracket \beta' \rrbracket \llbracket |E\partial_x \psi|^2 \rrbracket + s\lambda\varphi \langle \beta' \rangle \langle |E\partial_x \psi|^2 \rangle + 2s\lambda\varphi \langle \beta' \rangle \llbracket |E\partial_x \psi|^2 \rrbracket \langle |E\partial_x \psi|^2 \rangle \\ &\quad + \frac{1}{4} s^3\lambda^3\varphi^3 \llbracket E^2(\beta')^3 \rrbracket \llbracket |\psi|^2 \rrbracket + s^3\lambda^3\varphi^3 \llbracket E^2(\beta')^3 \rrbracket \langle |\psi|^2 \rangle + 2s^3\lambda^3\varphi^3 \langle E^2(\beta')^3 \rangle \llbracket |\psi|^2 \rrbracket \langle |\psi|^2 \rangle. \end{aligned}$$

Using the conditions verified by ψ at the singular point a , we get the quadratic form μ_1 in function of the average of ψ and $E\partial_x \psi$ as follows and the form μ_2 in function of g_s

$$\mu = \mu_1 + \mu_2,$$

where

$$\begin{aligned}
\mu_1 = & s\lambda\varphi \left\{ \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[\beta']\langle E\beta' \rangle^2 + [\beta'] + 2s\lambda k\varphi \mathcal{Q}\langle\beta'\rangle\langle E\beta' \rangle + \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[E^2(\beta')^3] \right\} \langle E\partial_x\psi \rangle^2 \\
& + s^3\lambda^3\varphi^3 \left\{ \frac{1}{4}[\beta']\langle E\beta' \rangle^2 + \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[\beta']\langle E\beta' \rangle^4 - \frac{1}{2}s\lambda k\varphi[\beta']\langle E\beta' \rangle\langle E\beta' \rangle^2 \right\} \langle\psi\rangle^2 \\
& + s^3\lambda^3\varphi^3 \left\{ \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[E^2(\beta')^3]\langle E\beta' \rangle^2 + [E^2(\beta')^3] - 2s\lambda k\varphi \mathcal{Q}\langle E^2(\beta')^3 \rangle\langle E\beta' \rangle \right\} \langle\psi\rangle^2 \\
& + 2s^2\lambda^2\varphi^2 \left\{ \frac{1}{4}s\lambda k\varphi \mathcal{Q}[\beta']\langle E\beta' \rangle\langle E\beta' \rangle - \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[\beta']\langle E\beta' \rangle^3 \right\} \langle\psi\rangle\langle E\partial_x\psi \rangle \\
& + 2s^2\lambda^2\varphi^2 \left\{ \langle\beta'\rangle[\beta']\langle E\beta' \rangle - s\lambda k\varphi \mathcal{Q}\langle\beta'\rangle\langle E\beta' \rangle^2 - \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[E^2(\beta')^3]\langle E\beta' \rangle \right. \\
& \left. + s\lambda k\varphi \mathcal{Q}\langle E^2(\beta')^3 \rangle \right\} \langle\psi\rangle\langle E\partial_x\psi \rangle
\end{aligned}$$

and

$$\begin{aligned}
\mu_2 = & \frac{1}{4}s^3\lambda^3\varphi^3 \mathcal{Q}^2([\beta']\langle E\beta' \rangle^2 + [E^2(\beta')^3]) |g_s|^2 + 2s^2\lambda^2\varphi^2 \mathcal{Q}\langle\beta'\rangle\langle E\beta' \rangle\langle E\partial_x\psi \rangle g_s \\
& + \frac{1}{2}s^3\lambda^3 k\varphi^3 \mathcal{Q}^2([\beta']\langle E\beta' \rangle^2 + [E^2(\beta')^3]) \langle E\partial_x\psi \rangle g_s \\
& + s^3\lambda^3\varphi^3 \mathcal{Q} \left(\frac{1}{2}[\beta']\langle E\beta' \rangle\langle E\beta' \rangle + 2\langle E^2(\beta')^3 \rangle \right) \langle\psi\rangle g_s \\
& - \frac{1}{2}s^4\lambda^4 k\varphi^4 \mathcal{Q}^2([\beta']\langle E\beta' \rangle^3 + [E^2(\beta')^3]\langle E\beta' \rangle) \langle\psi\rangle g_s.
\end{aligned}$$

We note that we can write the quadratic form μ_1 in matrix form as follows

$$\mu_1 = \Psi^T \cdot \mathcal{M} \cdot \Psi,$$

where Ψ is the following vector

$$\Psi = (\langle E\partial_x\psi \rangle, s\lambda\varphi\langle\psi\rangle)^T$$

and the coefficient of the symmetric matrix \mathcal{M} are given by

$$\begin{aligned}
\mathcal{M}_{11} = & \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[\beta']\langle E\beta' \rangle^2 + [\beta'] + 2s\lambda k\varphi \mathcal{Q}\langle\beta'\rangle\langle E\beta' \rangle + \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[E^2(\beta')^3], \\
\mathcal{M}_{22} = & \frac{1}{4}[\beta']\langle E\beta' \rangle^2 + \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[\beta']\langle E\beta' \rangle^4 - \frac{1}{2}s\lambda k\varphi \mathcal{Q}[\beta']\langle E\beta' \rangle\langle E\beta' \rangle^2 + \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[E^2(\beta')^3]\langle E\beta' \rangle^2 \\
& + [E^2(\beta')^3] - 2s\lambda k\varphi \mathcal{Q}\langle E^2(\beta')^3 \rangle\langle E\beta' \rangle, \\
\mathcal{M}_{12,21} = & \frac{1}{4}s\lambda k\varphi \mathcal{Q}[\beta']\langle E\beta' \rangle\langle E\beta' \rangle - \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[\beta']\langle E\beta' \rangle^3 + \langle\beta'\rangle[\beta']\langle E\beta' \rangle - s\lambda k\varphi \mathcal{Q}\langle\beta'\rangle\langle E\beta' \rangle^2 \\
& - \frac{1}{4}s^2\lambda^2 k^2\varphi^2 \mathcal{Q}^2[E^2(\beta')^3]\langle E\beta' \rangle + s\lambda k\varphi \langle E^2(\beta')^3 \rangle.
\end{aligned}$$

So, the determinant of the matrix \mathcal{M} is

$$\begin{aligned}
\det\mathcal{M} \simeq_{s\lambda\gg} & [\beta']^2\langle E\beta' \rangle^2 - 4[\beta']\langle E\beta' \rangle^2\langle\beta'\rangle^2 - 8[\beta']\langle E\beta' \rangle^2\langle E\beta' \rangle^2 + 32\langle\beta'\rangle^2\langle E\beta' \rangle^2 + 16\frac{[\beta']^2\langle E\beta' \rangle^4}{[\beta']^2} \\
& - 64\frac{\langle\beta'\rangle^2\langle E\beta' \rangle^4}{[\beta']^2}.
\end{aligned}$$

We remark that the determinant of the matrix \mathcal{M} is a polynomial of fourth order with the consideration of the following changes

$$\llbracket \beta' \rrbracket = c(r-1), \quad \langle \beta' \rangle = \frac{c}{2}(r+1), \quad c = \beta'(a^-), \quad r = \frac{\beta'(a^+)}{\beta'(a^-)}.$$

So, the determinant is as follows

$$\det \mathcal{M} \simeq_{s\lambda \gg} c^4 \left((r-1)^4 + 3(r+1)^4 - \frac{(r+1)^6}{(r-1)^2} - 3(r-1)^2(r+1)^2 \right).$$

For the case of non homogeneous domain, it has the following form

$$\det \mathcal{M} \simeq_{s\lambda \gg} c_1^2 c^4 \left(-4r(er-1)^2 + 8r(er+1)^2 - 4r \frac{(er+1)^4}{(er-1)^2} \right),$$

where $c_1 = E^-$ and $e = \frac{E^+}{E^-}$. From Figure 2, we choose the derivative of the weight function as follows

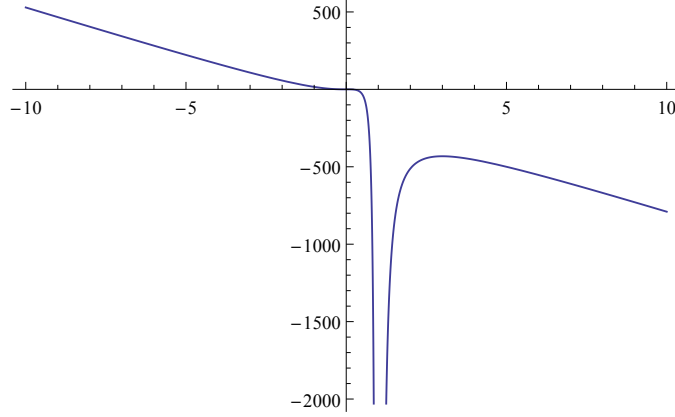


Figure 2: The determinant curve.

$$\llbracket \beta' \rrbracket = \alpha > 0, \quad \langle \beta' \rangle = 0.$$

So, for this choice, the form μ is a definite positive form as follows

$$\mu_1 \geq Cs\lambda\varphi(a)\langle E\partial_x\psi \rangle^2 + Cs^3\lambda^3\varphi^3(a)\langle \psi \rangle^2.$$

3.4 Final estimates

The final step will be to estimate the terms on the right-hand side of the following equation. So, we obtain after the consideration of the weight function choice

$$\begin{aligned} & \|\mathcal{P}_1\psi\|_{L^2(\Omega')}^2 + \|\mathcal{P}_2\psi\|_{L^2(\Omega')}^2 + 6s\lambda^2 \int_{\Omega'} \varphi(\beta')^2 |E\partial_x\psi|^2 dx + 2s^3\lambda^4 \int_{\Omega'} \varphi^3 E^2(\beta')^4 |\psi|^2 dx \\ & + Cs\lambda\varphi(a)\langle E\partial_x\psi \rangle^2 + Cs^3\lambda^3\varphi^3(a)\langle \psi \rangle^2 + \frac{1}{2}s^3\lambda^3\varphi^3\mathcal{Q}^2 \left(2\llbracket \beta' \rrbracket \langle E\beta' \rangle^2 + \frac{1}{4}\llbracket \beta' \rrbracket \llbracket E\beta' \rrbracket^2 \right) |g_s|^2 \\ & \leq \|f_s\|_{L^2(\Omega')}^2 - 2\{\mathcal{X}_{11} + \mathcal{X}_{12} + \mathcal{X}_{21} + \mathcal{Y}_1 + \mathcal{Y}_2\} + 2s\lambda\varphi(1)\beta'(1)|E\partial_x\psi|^2(1) - 2s\lambda\varphi(0)\beta'(0)|E\partial_x\psi|^2(0) \\ & + 2s^3\lambda^3\varphi^3(1)(E^2(\beta')^3)(1)|\psi|^2(1) - 2s^3\lambda^3\varphi^3(0)(E^2(\beta')^3)(0)|\psi|^2(0), \end{aligned} \tag{4}$$

where

$$\mathcal{Y}_1 = \frac{1}{2}s^3\lambda^3 k\varphi^3 \mathcal{Q}^2 \left(2\llbracket\beta'\rrbracket\langle E\beta'\rangle^2 + \frac{1}{4}\llbracket\beta'\rrbracket\llbracket E\beta'\rrbracket^2 \right) \langle E\partial_x\psi\rangle_{g_s}$$

and

$$\mathcal{Y}_2 = \frac{3}{2}s^3\lambda^3\varphi^3 \mathcal{Q}\llbracket\beta'\rrbracket\llbracket E\beta'\rrbracket\langle E\beta'\rangle\langle\psi\rangle_{g_s} - \frac{1}{2}s^4\lambda^4 k\varphi^4 \mathcal{Q}^2\langle E\beta'\rangle \left(2\llbracket\beta'\rrbracket\langle E\beta'\rangle^2 + \frac{1}{4}\llbracket\beta'\rrbracket\llbracket E\beta'\rrbracket^2 \right) \langle\psi\rangle_{g_s}.$$

Remark 3. Observe that the coefficients in front of the trace terms at 0 and 1 on the r.h.s. in the inequality (4) are positive because of the properties of the function β .

Now, we estimate the right-hand side terms in the above inequality (4) using the properties of the derivative of β and of the Young moduli E as follows

$$|\mathcal{X}_{11}| = \left| s\lambda \int_{\Omega'} \varphi\beta'' |E\partial_x\psi|^2 dx \right| \leq C(\alpha_1, e_2)s\lambda \int_{\Omega'} \varphi|\partial_x\psi|^2 dx, \quad C(\alpha_1, e_2) = \alpha_1 e_2^2,$$

$$|\mathcal{X}_{21}| = s^3\lambda^3 \left| \int_{\Omega'} \varphi^3 (E^2(\beta')^3)' |\psi|^2 dx \right| \leq C(\alpha_1, e_2)s^3\lambda^3 \int_{\Omega'} \varphi^3 |\psi|^2 dx, \quad C(\alpha_1, e_2) = 3e_2^2\alpha_1\alpha^2.$$

We remark that the second term of the following equation absorb the third one

$$|\mathcal{X}_{12}| = \left| -2s\lambda^2 [\varphi E(\beta')^2 \psi E\partial_x\psi]_{S\cup\Gamma} + 2s\lambda^3 \int_{\Omega'} \varphi E^2(\beta')^3 \psi \partial_x\psi dx + 2s\lambda^2 \int_{\Omega'} \varphi (E(\beta')^2)' \psi E\partial_x\psi dx \right|$$

and we estimate it by the Young inequality $mn \leq \frac{m^2}{2\epsilon} + \frac{\epsilon n^2}{2}$ as follows

$$2s\lambda^3 \int_{\Omega'} \varphi |E^2(\beta')^3 \psi \partial_x\psi| dx \leq C_\epsilon(\alpha, e_2, \epsilon)s\lambda^4 \int_{\Omega'} \varphi|\psi|^2 dx + \epsilon s\lambda^2 \int_{\Omega'} \varphi|\partial_x\psi|^2 dx,$$

where

$$C_\epsilon(\alpha, e_2, \epsilon) = \frac{e_2^2\alpha^3}{\epsilon}.$$

We develop the first term in \mathcal{X}_{12} using the formula for the jump of function's product

$$\begin{aligned} & -2s\lambda^2 [\varphi E(\beta')^2 \psi E\partial_x\psi]_{S\cup\Gamma} \\ &= -2s\lambda^2\varphi(1)(E(\beta')^2)(1)\psi(1)(E\partial_x\psi)(1) + 2s\lambda^2\varphi(0)(E(\beta')^2)(0)\psi(0)(E\partial_x\psi)(0) + 2s\lambda^2\varphi(a)\llbracket E(\beta')^2 \psi E\partial_x\psi \rrbracket \\ &= -2s\lambda^2\varphi(1)(E(\beta')^2)(1)\psi(1)(E\partial_x\psi)(1) + 2s\lambda^2\varphi(0)(E(\beta')^2)(0)\psi(0)(E\partial_x\psi)(0) + \frac{1}{2}s\lambda^2\varphi(a)\llbracket E(\beta')^2 \rrbracket\llbracket \psi \rrbracket\llbracket E\partial_x\psi \rrbracket \\ &+ 2s\lambda^2\varphi(a)\llbracket E(\beta')^2 \rrbracket\langle\psi\rangle\langle E\partial_x\psi\rangle + 2s\lambda^2\varphi(a)\langle E(\beta')^2 \rrbracket\llbracket \psi \rrbracket\langle E\partial_x\psi\rangle + 2s\lambda^2\varphi(a)\langle E(\beta')^2 \rangle\langle\psi\rangle\llbracket E\partial_x\psi \rrbracket. \end{aligned}$$

We use the conditions verified by ψ and $E\partial_x\psi$ to obtain the following expression

$$\begin{aligned} & 2s\lambda^2\varphi(a)\llbracket E(\beta')^2 \psi E\partial_x\psi \rrbracket \\ &= \left(\frac{1}{2} - s^3\lambda^4 k\varphi^3(a)\mathcal{Q}\llbracket E(\beta')^2 \rrbracket\langle E\beta'\rangle\llbracket E\beta'\rrbracket + \frac{1}{2}s^4\lambda^5 k^2\varphi^4\mathcal{Q}^2\llbracket E(\beta')^2 \rrbracket\langle E\beta'\rangle^3 + 2s^2\lambda^3\varphi^2(a)\langle E(\beta')^2 \rangle\llbracket E\beta'\rrbracket \right. \\ & \left. - 2s^3\lambda^4 k\varphi^3(a)\mathcal{Q}\langle E(\beta')^2 \rangle\langle E\beta'\rangle^2\langle\psi\rangle^2 + \left(\frac{1}{2}s^2\lambda^3 k^2\varphi^2(a)\mathcal{Q}^2\llbracket E(\beta')^2 \rrbracket\langle E\beta'\rangle + 2s\lambda^2 k\varphi(a)\mathcal{Q}\langle E(\beta')^2 \rangle \right) \langle E\partial_x\psi \rangle^2 \right. \\ & \left. + \left(\frac{1}{2}s^2\lambda^3 k\varphi^2(a)\mathcal{Q}\llbracket E(\beta')^2 \rrbracket\llbracket E\beta'\rrbracket - s^3\lambda^4 k^2\varphi^3(a)\mathcal{Q}^2\llbracket E(\beta')^2 \rrbracket\langle E\beta'\rangle^2 + 2s\lambda^2\varphi(a)\llbracket E(\beta')^2 \rrbracket \right) \langle\psi\rangle\langle E\partial_x\psi \rangle \right. \\ & \left. + \frac{1}{2}s^2\lambda^3\varphi^2(a)\mathcal{Q}^2\langle E\beta'\rangle\llbracket E(\beta')^2 \rrbracket|g_s|^2 - s^3\lambda^4 k\varphi^3(a)\mathcal{Q}^2\llbracket E(\beta')^2 \rrbracket\langle E\beta'\rangle^2\langle\psi\rangle_{g_s} + \frac{1}{2}s^2\lambda^3\varphi^2(a)\mathcal{Q}\llbracket E(\beta')^2 \rrbracket\llbracket E\beta'\rrbracket\langle\psi\rangle_{g_s} \right. \\ & \left. + 2s^2\lambda^3\varphi^2(a)\mathcal{Q}\langle E(\beta')^2 \rangle\langle E\beta'\rangle\langle\psi\rangle_{g_s} + s^2\lambda^3 k\varphi^2(a)\mathcal{Q}^2\langle E\beta'\rangle\llbracket E(\beta')^2 \rrbracket\langle E\partial_x\psi \rangle_{g_s} + 2s\lambda^2\varphi(a)\mathcal{Q}\langle E(\beta')^2 \rangle\langle E\partial_x\psi \rangle_{g_s}. \right. \end{aligned}$$

We estimate the above term with the consideration that $\frac{1}{1+\frac{1}{4}s\lambda k\varphi\llbracket E\beta'\rrbracket} \leq \frac{4}{s\lambda k\varphi\llbracket E\beta'\rrbracket}$ for $s\lambda$ sufficiently large

$$\begin{aligned}
& 2s\lambda^2\varphi(a)\llbracket E(\beta')^2\psi E\partial_x\psi\rrbracket \\
& \leq (2s^2\lambda^3\varphi^2(a)\llbracket E(\beta')^2\rrbracket\langle E\beta'\rangle + 8s^2\lambda^3\varphi^2(a)\frac{\llbracket E(\beta')^3\rrbracket\langle E\beta'\rangle^3}{\llbracket E\beta'\rrbracket^2} + 2s^2\lambda^3\varphi^2(a)\langle E(\beta')^2\rangle\llbracket E\beta'\rrbracket \\
& + 8s^2\lambda^3\varphi^2(a)\frac{\langle E(\beta')^2\rangle\langle E\beta'\rangle^2}{\llbracket E\beta'\rrbracket}\langle\psi\rangle^2 + \left(8\lambda\frac{\llbracket E(\beta')^2\rrbracket\langle E\beta'\rangle}{\llbracket E\beta'\rrbracket^2} + 8\lambda\frac{\langle E(\beta')^2\rangle}{\llbracket E\beta'\rrbracket}\right)\langle E\partial_x\psi\rangle^2 \\
& + \left(4s\lambda^2\varphi(a)\llbracket E(\beta')^2\rrbracket + 16s\lambda^2\varphi(a)\frac{\llbracket E(\beta')^2\rrbracket\langle E\beta'\rangle^2}{\llbracket E\beta'\rrbracket^2}\right)\langle\psi\rangle\langle E\partial_x\psi\rangle \\
& + \frac{8\lambda\langle E\beta'\rangle\llbracket E(\beta')^2\rrbracket}{k^2\llbracket E\beta'\rrbracket^2}|g_s|^2 + \frac{s\lambda^2\varphi(a)}{k}\left(\frac{16\llbracket E(\beta')^2\rrbracket\langle E\beta'\rangle^2}{\llbracket E\beta'\rrbracket^2} + 2\llbracket E(\beta')^2\rrbracket + \frac{8\langle E(\beta')^2\rangle\langle E\beta'\rangle}{\llbracket E\beta'\rrbracket}\right)\langle\psi\rangle g_s \\
& + \left(\frac{16\lambda\langle E\beta'\rangle\llbracket E(\beta')^2\rrbracket}{k\llbracket E\beta'\rrbracket^2} + \frac{8\lambda\langle E(\beta')^2\rangle}{k\llbracket E\beta'\rrbracket}\right)\langle E\partial_x\psi\rangle g_s.
\end{aligned}$$

Using $\varphi \leq \varphi^3$, we obtain

$$\begin{aligned}
|\mathcal{X}_{12}| & \leq C_\epsilon(\alpha, e_2, \epsilon)s\lambda^4 \int_{\Omega'} \varphi^3|\psi|^2 dx + \epsilon s\lambda^2 \int_{\Omega'} \varphi|\partial_x\psi|^2 dx + (Cs^2\lambda^3 + C_\epsilon s\lambda^3 + \epsilon s\lambda^3)\varphi^3(a)\langle\psi\rangle^2 \\
& + (C\lambda + C\epsilon s\lambda + \epsilon\lambda)\varphi(a)\langle E\partial_x\psi\rangle^2 + (C\lambda + C_\epsilon s\lambda\varphi(a) + C_\epsilon\lambda)\frac{1}{k^2}|g_s|^2 \\
& + s^2\lambda^3\varphi^2(1)|\psi|^2(1) + s^2\lambda^3\varphi^2(0)|\psi|^2(0) + \lambda|E\partial_x\psi|^2(1) + \lambda|E\partial_x\psi|^2(0),
\end{aligned}$$

where C is a constant depending only on the Young moduli and the weight function β .

The terms $\mathcal{Y}_{j=1,2}$ are estimated as follows

$$|\mathcal{Y}_1| \leq C_\epsilon \frac{s\lambda\varphi(a)}{k^2}|g_s|^2 + \epsilon s\lambda\varphi(a)\langle E\partial_x\psi\rangle^2$$

and

$$|\mathcal{Y}_2| \leq \frac{C_\epsilon s\lambda\varphi(a)}{k^2}|g_s|^2 + \epsilon s^3\lambda^3\varphi^3(a)\langle\psi\rangle^2.$$

Similarly, we have

$$\|f_s\|_{L^2(\Omega')}^2 \leq C\|e^{s\varphi}f\|_{L^2(\Omega')}^2 + Cs^2\lambda^4 \int_{\Omega'} \varphi^3|\psi|^2 dx.$$

By gathering all the estimates, we get the following inequality

$$\begin{aligned}
& \|\mathcal{P}_1\psi\|_{L^2(\Omega')}^2 + \|\mathcal{P}_2\psi\|_{L^2(\Omega')}^2 + s\lambda^2 \int_{\Omega'} \varphi(\beta')^2|E\partial_x\psi|^2 dx + s^3\lambda^4 \int_{\Omega'} \varphi^3 E^2(\beta')^4|\psi|^2 dx \\
& + Cs\lambda\varphi(a)\langle E\partial_x\psi\rangle^2 + Cs^3\lambda^3\varphi^3(a)\langle\psi\rangle^2 + \frac{1}{2}s^3\lambda^3\varphi^3\mathcal{Q}^2\left(2\llbracket\beta'\rrbracket\langle E\beta'\rangle^2 + \frac{1}{4}\llbracket\beta'\rrbracket\llbracket E\beta'\rrbracket^2\right)|g_s|^2 \\
& \leq C\|e^{s\varphi}f\|_{L^2(\Omega')}^2 + (s^2\lambda^4 + s^3\lambda^3 + C_\epsilon s\lambda^4) \int_{\Omega'} \varphi^3|\psi|^2 dx + (s\lambda + \epsilon s\lambda^2) \int_{\Omega'} \varphi|\partial_x\psi|^2 dx \\
& (Cs^2\lambda^3 + C_\epsilon s\lambda^3 + \epsilon s\lambda^3 + \epsilon s\lambda^3 + \epsilon s^3\lambda^3)\varphi^3(a)\langle\psi\rangle^2 + (C\lambda + C\epsilon s\lambda + \epsilon\lambda + \epsilon\lambda + \epsilon s\lambda)\varphi(a)\langle E\partial_x\psi\rangle^2 \\
& + (C\lambda + C_\epsilon s\lambda + C_\epsilon\lambda + C\lambda + C_\epsilon s\lambda + C_\epsilon\lambda + C_\epsilon s\lambda)\frac{1}{k^2}\varphi(a)|g_s|^2 \\
& + (s\lambda + \lambda)\varphi(1)|E\partial_x\psi|^2(1) + (s\lambda + \lambda)\varphi(0)|E\partial_x\psi|^2(0) \\
& + (s^3\lambda^3 + s^2\lambda^3)\varphi^3(1)|\psi|^2(1) + (s^3\lambda^3 + s^2\lambda^3)\varphi^3(0)|\psi|^2(0).
\end{aligned}$$

By choosing ϵ sufficiently small and taking $\lambda \geq \lambda_0$, $s \geq s_0$, we obtain

$$\begin{aligned}
& \|\mathcal{P}_1\psi\|_{L^2(\Omega')}^2 + \|\mathcal{P}_2\psi\|_{L^2(\Omega')}^2 + s\lambda^2 \int_{\Omega'} \varphi |\partial_x \psi|^2 dx + s^3 \lambda^4 \int_{\Omega'} \varphi^3 |\psi|^2 dx + s^3 \lambda^3 \varphi^3(a) \langle \psi \rangle^2 \\
& + s\lambda \varphi(a) \langle E \partial_x \psi \rangle^2 + \frac{1}{2} s^3 \lambda^3 \varphi^3 \mathcal{Q}^2 (\llbracket \beta' \rrbracket \langle E \beta' \rangle^2 + \llbracket E^2(\beta')^3 \rrbracket) |g_s|^2 \\
& \leq C \|e^{s\varphi} f\|_{L^2(\Omega')}^2 + \frac{s\lambda \varphi(a)}{k^2} |g_s|^2 + s\lambda \varphi(1) |E \partial_x \psi|^2(1) + s\lambda \varphi(0) |E \partial_x \psi|^2(0) \\
& + s^3 \lambda^3 \varphi^3(1) |\psi|^2(1) + s^3 \lambda^3 \varphi^3(0) |\psi|^2(0).
\end{aligned} \tag{5}$$

3.5 Carleman's inequality

Recalling that $\psi = e^{s\varphi} q$, we have

$$e^{s\varphi} \partial_x q = \partial_x \psi - s\lambda \varphi \beta' \psi \quad \text{in } \Omega',$$

which yields

$$s\lambda^2 \varphi e^{2s\varphi} |\partial_x q|^2 \leq C s\lambda^2 \varphi |\partial_x \psi|^2 + C s^3 \lambda^4 \varphi^3 |\psi|^2 \quad \text{in } \Omega',$$

to be used on the left-hand side of (5). Consequently, we obtain

$$\begin{aligned}
& \|\mathcal{P}_1\psi\|_{L^2(\Omega')}^2 + \|\mathcal{P}_2\psi\|_{L^2(\Omega')}^2 + s\lambda^2 \int_{\Omega'} \varphi e^{2s\varphi} |\partial_x q|^2 dx + s^3 \lambda^4 \int_{\Omega'} \varphi^3 e^{2s\varphi} |q|^2 dx + s^3 \lambda^3 \varphi^3(a) e^{2s\varphi} \langle q \rangle^2 \\
& + s\lambda \varphi(a) e^{2s\varphi} \langle E \partial_x q \rangle^2 + \frac{1}{4} s\lambda \varphi(a) e^{2s\varphi} \llbracket q \rrbracket^2 + \frac{1}{2} s^3 \lambda^3 \varphi^3 \mathcal{Q}^2 (\llbracket \beta' \rrbracket \langle E \beta' \rangle^2 + \llbracket E^2(\beta')^3 \rrbracket) e^{2s\varphi} |g|^2 \\
& \leq C \|e^{s\varphi} f\|_{L^2(\Omega')}^2 + \frac{s\lambda \varphi(a)}{k^2} e^{2s\varphi(a)} |g|^2 + s\lambda \varphi(1) e^{2s\varphi(1)} |E \partial_x q|^2(1) + s\lambda \varphi(0) e^{2s\varphi(0)} |E \partial_x q|^2(0) \\
& + s^3 \lambda^3 \varphi^3(1) e^{2s\varphi(1)} |q|^2(1) + s^3 \lambda^3 \varphi^3(0) e^{2s\varphi(0)} |q|^2(0).
\end{aligned}$$

Remark 4. *The Carleman estimate, established here, gives the unique continuation for the 1D elasticity equation with an imperfect interface from the knowledge of the boundary data on both sides.*

4 Stability of the inverse problem

Supposing q_i as a solution satisfying (1,2) associated to $k = k_i$, for $i = 1, 2$ and also $f(x) = 0$ for all $x \in \Omega'$. Let q be the difference between both solutions q_1 and q_2 which satisfies the following boundary condition at the point a

$$\llbracket E \partial_x q \rrbracket = 0, \quad \llbracket q \rrbracket = k_2 \langle E \partial_x q \rangle + (k_1 - k_2) \langle E \partial_x q_1 \rangle.$$

We apply the Carleman inequality, Theorem (1), to the inverse problem presented above with the interface term $g = (k_1 - k_2) \langle E \partial_x q_1 \rangle$ and we get

$$\begin{aligned}
& \|\mathcal{P}_1\psi\|_{L^2(\Omega')}^2 + s^3 \lambda^4 \int_{\Omega'} e^{2s\varphi} \varphi^2 |q|^2 dx \\
& \leq \|\mathcal{P}_1\psi\|_{L^2(\Omega')}^2 + \|\mathcal{P}_2\psi\|_{L^2(\Omega')}^2 + s\lambda^2 \int_{\Omega'} e^{2s\varphi} \varphi |E \partial_x q|^2 dx + s^3 \lambda^4 \int_{\Omega'} e^{2s\varphi} \varphi^3 |q|^2 dx \\
& \leq \text{obser} + C \frac{s\lambda \varphi}{k^2} e^{2s\varphi} |g|^2.
\end{aligned} \tag{6}$$

Using the fact that $\psi = e^{s\varphi} q$, one develops the first term of the above inequality as follows

$$\begin{aligned}
\mathcal{P}_1\psi & = -\partial_x (E \partial_x \psi) - s^2 \lambda^2 \varphi^2 (\beta')^2 E \psi \\
& = -\partial_x (E \partial_x (e^{s\varphi} q)) - s^2 \lambda^2 \varphi^2 (\beta')^2 E e^{s\varphi} q \\
& = - (s\lambda (E \beta')' \varphi + s\lambda^2 E (\beta')^2 \varphi^2 + 2s^2 \lambda^2 E (\beta')^2 \varphi^2) e^{s\varphi} q - 2s\lambda \beta' \varphi e^{s\varphi} E \partial_x q,
\end{aligned} \tag{7}$$

with the fact that $\partial_x(E\partial_x q) = 0$.

From (7), we get the following estimate

$$-(s\lambda^2 E(\beta')^2 \varphi^2 + 2s^2 \lambda^2 E(\beta')^2 \varphi^2) e^{s\varphi} q - 2s\lambda\beta' \varphi e^{s\varphi} E\partial_x q = \mathcal{P}_1 \psi + s\lambda(E\beta')' \varphi e^{s\varphi} q. \quad (8)$$

Now, we estimate the last equality (8) using Cauchy's inequality as follows

$$\begin{aligned} & \| (s\lambda^2 E(\beta')^2 \varphi^2 + 2s^2 \lambda^2 E(\beta')^2 \varphi^2) e^{s\varphi} q + 2s\lambda\beta' \varphi e^{s\varphi} E\partial_x q \|_{L^2(\Omega')}^2 + s^3 \lambda^4 \int_{\Omega'} e^{2s\varphi} \varphi^2 |q|^2 dx \\ & \leq \| \mathcal{P}_1 \psi + s\lambda(E\beta')' \varphi e^{s\varphi} q \|_{L^2(\Omega')}^2 + s^3 \lambda^4 \int_{\Omega'} e^{2s\varphi} \varphi^2 |q|^2 dx \\ & \leq \| \mathcal{P}_1 \psi \|_{L^2(\Omega')}^2 + C s^2 \lambda^2 \int_{\Omega'} e^{2s\varphi} \varphi^2 |q|^2 dx + s^3 \lambda^4 \int_{\Omega'} e^{2s\varphi} \varphi^2 |q|^2 dx. \end{aligned} \quad (9)$$

Then, we remark that the third term in the last inequality (9) absorbs the second one and we get by Carleman's estimate (6) the following estimate

$$\begin{aligned} & \| (s\lambda^2 E(\beta')^2 \varphi^2 + 2s^2 \lambda^2 E(\beta')^2 \varphi^2) e^{s\varphi} q - 2s\lambda\beta' \varphi e^{s\varphi} E\partial_x q \|_{L^2(\Omega')}^2 + s^3 \lambda^4 \int_{\Omega'} e^{2s\varphi} \varphi^2 |q|^2 dx \\ & \leq \| \mathcal{P}_1 \psi \|_{L^2(\Omega')}^2 + s^3 \lambda^4 \int_{\Omega'} e^{2s\varphi} \varphi^2 |q|^2 dx \\ & \leq \text{obser} + C \frac{s\lambda\varphi}{k_2^2} e^{2s\varphi} |g|^2. \end{aligned} \quad (10)$$

Next, we need to estimate the boundary integral in terms of the domain integral and for that we present the following theorem.

Theorem 3. (for more details see [9], Theorem 1.5.1.10) Let Ω be a bounded open subset of \mathbb{R}^n with a Lipschitz boundary Γ . Then, there exists a constant C such that

$$\int_{\Gamma} |u|^p ds \leq C \left(\epsilon^{1-\frac{1}{p}} \int_{\Omega} |\nabla u|^p dx + \epsilon^{-\frac{1}{p}} \int_{\Omega} |u|^p dx \right),$$

for all $u \in W_p^1(\Omega)$ and $\epsilon \in]0, 1[$. In addition, C depends only on the norm of μ in $C^1(\bar{\Omega})$ and on δ (defined in Lemma 1.5.1.9 [9]).

We present the above estimation in the following form

$$\epsilon^{\frac{1}{p}-1} \int_{\Gamma} |u|^p ds \leq C \left(\int_{\Omega} |\nabla u|^p dx + \epsilon^{-1} \int_{\Omega} |u|^p dx \right)$$

and to apply it in our case, we make the following choices

$$\begin{aligned} u &= s\lambda\varphi e^{s\varphi} q, \\ \nabla u &= s\lambda^2 \beta' \varphi e^{s\varphi} q + s^2 \lambda^2 \beta' \varphi^2 e^{s\varphi} q + s\lambda\varphi e^{s\varphi} \partial_x q, \\ \epsilon &= (s\lambda^2)^{-1}. \end{aligned}$$

By Theorem 3, we get then the trace inequality

$$\begin{aligned} s^{5/2} \lambda^3 e^{2s\varphi} (|q(a^+)|^2 + |q(a^-)|^2) & \leq s^{5/2} \lambda^3 e^{2s\varphi} (|q(1)|^2 + |q(0)|^2 + |q(a^+)|^2 + |q(a^-)|^2) \\ & \leq \text{obser} + C \frac{s\lambda\varphi}{k_2^2} e^{2s\varphi} |g|^2. \end{aligned} \quad (11)$$

We use the expression $[[q]] = k_2 \langle E \partial_x q \rangle + g$ to get the following inequality

$$\begin{aligned} |g|^2 &= |[q] - k_2 \langle E \partial_x q \rangle|^2 \\ &\leq 2 \left(|[q]|^2 + |k_2 \langle E \partial_x q \rangle|^2 \right) \\ &\leq 2 |q(a^+) - q(a^-)|^2 + 2 |k_2 \langle E \partial_x q \rangle|^2 \\ &\leq 4 \left(|q(a^+)|^2 + |q(a^-)|^2 \right) + 2 |k_2 \langle E \partial_x q \rangle|^2. \end{aligned}$$

Then, inserting the above inequality in (11), it yields

$$\begin{aligned} s^{5/2} \lambda^3 e^{2s\varphi(a)} |g|^2 &\leq s\lambda \varphi(1) e^{2s\varphi(1)} |E \partial_x q|^2(1) + s\lambda \varphi(0) e^{2s\varphi(0)} |E \partial_x q|^2(0) + s^3 \lambda^3 \varphi^3(1) e^{2s\varphi(1)} |q|^2(1) \\ &\quad + s^3 \lambda^3 \varphi^3(0) e^{2s\varphi(0)} |q|^2(0) + C \frac{s\lambda\varphi}{k^2} e^{2s\varphi(a)} |k_1 - k_2|^2 \langle \partial_x q_1 \rangle^2 + s^{5/2} \lambda^3 e^{2s\varphi(a)} |k_2 \langle E \partial_x q \rangle|^2. \end{aligned}$$

We divide all the terms by the factor $s^{5/2} \lambda^3$

$$\begin{aligned} &e^{2s\varphi(a)} |k_1 - k_2|^2 \langle E \partial_x q_1 \rangle^2 \\ &\leq \frac{1}{\sqrt{s}} \varphi(1) e^{2s\varphi(1)} |E \partial_x q|^2(1) + \frac{1}{\sqrt{s}} \varphi(0) e^{2s\varphi(0)} |E \partial_x q|^2(0) + \sqrt{s} \varphi^3(1) e^{2s\varphi(1)} |q|^2(1) \\ &\quad + \sqrt{s} \varphi^3(0) e^{2s\varphi(0)} |q|^2(0) + C \frac{1}{\sqrt{s}} \frac{\varphi(a)}{k^2} e^{2s\varphi(a)} |k_1 - k_2|^2 \langle \partial_x q_1 \rangle^2 + e^{2s\varphi(a)} |k_2 \langle E \partial_x q \rangle|^2. \end{aligned} \tag{12}$$

For s large enough, the left-hand side in (12) can absorb the last term of the right-hand side and the first terms of the right-hand side are sufficiently small. Therefore, since φ is a bounded function, we obtain

$$|k_1 - k_2|^2 \langle E \partial_x q_1 \rangle^2 \leq C (|q|^2(1) + |q|^2(0)) + e^{2s\varphi(a)} |k_2 \langle E \partial_x q \rangle|^2.$$

For $\langle E \partial_x q_1 \rangle \geq m > 0$, we have the stability estimate

$$|k_1 - k_2|^2 \leq C (|q_1 - q_2|^2(1) + |q_1 - q_2|^2(0)) + e^{2s\varphi(a)} |k_2 \langle E \partial_x q \rangle|^2.$$

Now, we estimate the second term by using a trace inequality, q as a harmonic function and the properties of the function φ

$$\|v\|_{L^2(\partial\Omega)}^2 \leq C' \|v\|_{H^1(\Omega)}^2 \leq C' \|v\|_{L^2(\Omega)}^2.$$

Here, we choose

$$v = \varphi^{1/2} e^{s\varphi} \partial_x q.$$

So, one gets

$$s^{5/2} \lambda^3 \varphi(a) e^{2s\varphi(a)} |k_2 \langle E \partial_x q \rangle|^2 \leq C' s^{5/2} \lambda^3 \|\varphi^{1/2} e^{s\varphi} \partial_x q\|_{L^2(\Omega)}^2$$

and we estimate the right-hand side of the above inequality by Carleman's inequality (1) as following

$$\begin{aligned} &s^{5/2} \lambda^3 \|\varphi^{1/2} e^{s\varphi} \partial_x q\|_{L^2(\Omega)}^2 \\ &= s^{3/2} \lambda \left\{ s\lambda^2 \|\varphi^{1/2} e^{s\varphi} \partial_x q\|_{L^2(\Omega)}^2 \right\} \\ &\leq s^{3/2} \lambda \left\{ s\lambda \varphi(1) e^{2s\varphi(1)} |E \partial_x q|^2(1) + s\lambda \varphi(0) e^{2s\varphi(0)} |E \partial_x q|^2(0) + s^3 \lambda^3 \varphi^3(1) e^{2s\varphi(1)} |q|^2(1) \right\} \\ &\quad + s^{3/2} \lambda \left\{ s^3 \lambda^3 \varphi^3(0) e^{2s\varphi(0)} |q|^2(0) + C \frac{s\lambda\varphi}{k^2} e^{2s\varphi(a)} |k_1 - k_2|^2 \langle \partial_x q_1 \rangle^2 \right\}. \end{aligned}$$

We divide all the terms by $s^{5/2}\lambda^3$

$$\begin{aligned} & \varphi(a)e^{2s\varphi(a)}|k_2\langle E\partial_x q \rangle|^2 \\ & \leq \frac{1}{\lambda}\varphi(1)e^{2s\varphi(1)}|E\partial_x q|^2(1) + \frac{1}{\lambda}\varphi(0)e^{2s\varphi(0)}|E\partial_x q|^2(0) + s^2\lambda\varphi^3(1)e^{2s\varphi(1)}|q|^2(1) \\ & \quad + s^2\lambda\varphi^3(0)e^{2s\varphi(0)}|q|^2(0) + C\frac{1}{\lambda}\varphi(a)e^{2s\varphi(a)}|k_1 - k_2|^2|\partial_x q_1|^2. \end{aligned}$$

For λ large enough, the left-hand side can absorb the last term of the right-hand side and the first terms of the right-hand side are sufficiently small.

Therefore, since φ is a bounded function, we obtain

$$|k_2\langle E\partial_x q \rangle|^2 \leq C\{|q|^2(1) + |q|^2(0)\}.$$

Remark 5. *The main idea for the above proof is presented by two step, the first one by estimating the interfacial term $\langle E\partial_x q \rangle$ by the parameter s and the second one by estimating the interface additional term $\langle E\partial_x q_1 \rangle$ by the parameter λ . It is indispensable to write the Carleman estimate in terms of the two parameters s and λ .*

5 Conclusion

By borrowing the approach developed by Benabdallah et al. [5], we prove a Carleman-type estimate that quantifies the unique continuation result for the case of one-dimensional elasticity equation with a spring-type linear law. The proof is based on the construction of suitable weight functions, whose gradient is non-zero, the jump of the derivative is positive across the interface and the average of the derivative vanishes which is a new parameter imposed to enable the control of the interface terms. We have proved a stability estimate for the inverse problem of identifying a stiffness parameter from measurements available on both sides of the external boundary. The extendibility of this stability estimate to the case of the non-stationary problem will be pursued in the future. An extension to higher dimensions will be a continuation of this work such that the challenge will be in the reconstruction of the weight function to get a higher dimensional Carleman-type estimate but the stability proof will have the same lines as the one presented in this work.

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