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# GRADIENT ESTIMATES AND MAXIMAL DISSIPATIVITY FOR THE KOLMOGOROV OPERATOR IN $\Phi_{2}^{4}$ 

GIUSEPPE DA PRATO AND ARNAUD DEBUSSCHE

Abstract. We consider the transition semigroup $P_{t}$ of the $\Phi_{2}^{4}$ stochastic quantisation on the torus $\mathbb{T}^{2}$ and prove the following new estimate (Theorem 3.9)

$$
\left|D P_{t} \varphi(x) \cdot h\right| \leq c t^{-\beta}|h|_{C^{-s}}\|\varphi\|_{0}\left(1+|x|_{C^{-\alpha}}\right)^{\gamma},
$$

for some $\alpha, \beta, \gamma, s$ positive. Thanks to this estimate, we show that cylindrical functions are a core for the corresponding Kolmogorov equation. Some consequences of this fact are discussed in a final remark.

## 1. Introduction

We consider the equation of the $\Phi_{2}^{4}$ stochastic quantisation on the torus $\mathbb{T}^{2}$ :

$$
\left\{\begin{array}{l}
d X=\left(A X-: X^{3}:\right) d t+d W  \tag{1.1}\\
X(0)=x
\end{array}\right.
$$

where

$$
A x=\Delta x-x, \quad D(A)=H^{2}\left(\mathbb{T}^{2}\right)
$$

This equation has been the object of several studies. In particular, in [DaDe03-2] existence and uniqueness of a mild solution to (1.1) has been proved for almost all initial datas $x$ with respect to the Gibbs measure $\nu$ defined below; later in [MoWe15] well-posedness was proved for all $x$ in a negative Besov space. Moreover, in [HaMa16] the strong Feller property of the corresponding transition semigroup $\left(P_{t}\right)_{t \geq 0}$ was proved. This latter result is improved in [TsWe18] where it is proved that, for $t>0, P_{t}$ maps bounded Borel functions to $\alpha$-Hölder functions for some $\alpha \in(0,1)$.

The main result of the present paper (Theorem 3.9 below) goes further in this direction: we prove that $P_{t}$ maps bounded borelian functions to Lipschitz functions and give an estimate of the gradient of $P_{t} \varphi$. We use a method similar to the one used in [ $\left.\mathrm{DaDe} 03-1\right]$ and $[\mathrm{DaDe} 07]$ as well as similar estimates as in [TsWe18]. In the second part we use this result for showing that cylindrical functions are a core for the infinitesimal generator $\mathcal{K}$ of $P_{t}$ and present some consequences of this fact.

## 2. NOTATIONS

We use the classical $L^{p}=L^{p}\left(\mathbb{T}^{2}\right)$ spaces for $p \in[1, \infty]$. For $p=2$, we write $H=L^{2}$. We use also the $L^{2}$ based Sobolev spaces $H^{s}=H^{s}\left(\mathbb{T}^{2}\right)$ for $s \in \mathbb{R}$ and the Besov spaces $B_{p, q}^{s}=B_{p, q}^{s}\left(\mathbb{T}^{2}\right)$, for $s \in \mathbb{R}, p, q \in[1, \infty]$. These are defined in terms of Fourier series and Littlewood-Paley theory (see [BaChDa11]). Note that $H^{s}=B_{2,2}^{s}$. When $p=q=\infty$, we write $B_{\infty, \infty}^{s}=C^{s}$ and for $s>0$ non integer these spaces coincide with the classical Hölder spaces.

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Besov spaces are convenient when multiplying two functions when one has negative regularity. Recall that for $p_{1}, p_{2} \in[1, \infty]$ and $\alpha_{1}, \alpha_{2} \in \mathbb{R}$ with $\alpha_{1}+\alpha_{2}>0$ we have for $\alpha=\alpha_{1} \wedge \alpha_{2} \wedge\left(\alpha_{1}+\alpha_{2}\right)$, $\frac{1}{p}=\frac{1}{p_{1}}+\frac{1}{p_{2}}, \tilde{p}=p_{1} \vee p_{2}$

$$
\|u v\|_{B_{p, \tilde{p}}^{\alpha}}^{\alpha} \leq C\|u\|_{B_{p_{1}, p_{1}}^{\alpha_{1}}}\|v\|_{B_{p_{2}, p_{2}}^{\alpha_{2}}}
$$

and for any $\kappa>0$

$$
\|u v\|_{B_{p, p}^{\alpha-\kappa}}^{\alpha-1} \leq C\|u\|_{B_{p_{1}, p_{1}}^{\alpha_{1}}}\|v\|_{B_{p_{2}, p_{2}}^{\alpha_{2}}} .
$$

We use the eigenprojections $\Pi_{N}$ corresponding to the eigenvalues $k^{2} \leq N, k \in \mathbb{Z}^{2}$, of $A$.
Define the Gaussian measure

$$
\mu=\prod_{k \in \mathbb{Z}^{2}} N\left(0, \frac{1}{1+|k|^{2}}\right), \text { on } \mathcal{H}=\mathbb{R}^{\mathbb{Z}^{2}},
$$

and for $x_{N} \in \Pi_{N} H$

$$
: x_{N}^{2}:=x_{N}^{2}-\rho_{N}^{2}, \quad: x_{N}^{3}:=x_{N}^{3}-3 \rho_{N}^{2} x_{N},
$$

where

$$
\rho_{N}=\frac{1}{2 \pi}\left[\sum_{|k| \leq N} \frac{1}{1+|k|^{2}}\right]^{1 / 2}
$$

More generally, the $\mathrm{n}^{\text {th }}$ Wick power of $x_{N}$ is defined by : $x_{N}^{n}:=\sqrt{n!} \rho_{N}^{n} H_{n}\left(\frac{1}{\rho_{N}} x_{N}\right)$ where $H_{n}$ is the $\mathrm{n}^{\text {th }}$ Hermite polynomial.

As well known, there exists the limit

$$
\begin{equation*}
\lim _{N \rightarrow \infty}:\left(\Pi_{N} x\right)^{n}:=: x^{n}:, \text { in } L^{r}\left(\mathcal{H}, \mu ; B_{p, q}^{s}\right), \tag{2.1}
\end{equation*}
$$

for any $r \in[1, \infty), p, q \in[1, \infty], s<0$.
Then, we define

$$
Z_{\infty}(t)=\int_{-\infty}^{t} e^{(t-s) A} d W(s)
$$

Since $Z_{\infty}(t)$ has law $\mu$, we know that : $\Pi_{N} Z_{\infty}^{n}$ : converges to : $Z_{\infty}^{n}$ : in $L^{r}\left(\mathcal{H}, \mu ; B_{p, q}^{s}\right)$ for any $r \in[1, \infty), p, q \in[1, \infty], s<0$

We set moreover $X=Y_{\infty}+Z_{\infty}=Y+Z$ where

$$
Z(t)=\int_{0}^{t} e^{(t-s) A} d W(s)=Z_{\infty}(t)-e^{A t} Z_{\infty}(0)
$$

Clearly:

$$
:\left(\Pi_{N} X\right)^{3}:=\left(\Pi_{N} Y_{\infty}\right)^{3}+3\left(\Pi_{N} Y_{\infty}\right)^{2} \Pi_{N} Z_{\infty}+3\left(\Pi_{N} Y_{\infty}\right):\left(\Pi_{N} Z_{\infty}\right)^{2}:+:\left(\Pi_{N} Z_{\infty}\right)^{3}: .
$$

It is therefore natural to consider the following definition for : $X^{3}$ : in (1.1):

$$
: X^{3}:=Y_{\infty}^{3}+3 Y_{\infty}^{2} Z_{\infty}+3 Y_{\infty}: Z_{\infty}^{2}:+: Z_{\infty}^{3}: .
$$

We also have

$$
: X^{3}:=Y^{3}+3 Y^{2} Z+3 Y: Z^{2}:+: Z^{3}:
$$

if we set

$$
: Z^{3}:=-\left(e^{A t} Z_{\infty}(0)\right)^{3}+3\left(e^{A t} Z_{\infty}(0)\right)^{2} Z_{\infty}-3\left(e^{A t} Z_{\infty}(0)\right): Z_{\infty}^{2}:+: Z_{\infty}^{3}: .
$$

We will see that $Y$ and $Y_{\infty}$ has positive regularity whereas $Z$ has the same smoothness as $Z_{\infty}$. Using the product rules above, it can be seen that all products above make sense.

Finally, we define $Z_{N}=\Pi_{N} Z$ and introduce the following approximation of (1.1):

$$
\left\{\begin{array}{l}
\frac{d Y_{N}}{d t}=A Y_{N}-\left(Y_{N}^{3}+3 Y_{N}^{2} Z_{N}+3 Y_{N}: Z_{N}^{2}:+: Z_{N}^{3}:\right)  \tag{2.2}\\
Y_{N}(0)=x
\end{array}\right.
$$

Equivalently, setting $X_{N}=Y_{N}+Z_{N}$, we may write with a slight abuse of notation

$$
: X_{N}^{3}:=Y_{N}^{3}+3 Y_{N}^{2} Z_{N}+3 Y_{N}: Z_{N}^{2}:+: Z_{N}^{3}:
$$

and

$$
\left\{\begin{array}{l}
d X_{N}=\left(A X_{N}+: X_{N}^{3}:\right) d t+\Pi_{N} d W  \tag{2.3}\\
X_{N}(0)=x
\end{array}\right.
$$

Note that this is not a Galerkin approximation. However, since : $X_{N}^{3}:=X_{N}^{3}-3 \rho_{N}^{2} X_{N}$ and $\Pi_{N} d W$ is a finite dimensional white noise, it is classical to prove that $X_{N}$ exists and is unique. Moreover, it is smooth in space when the initial datum is smooth; in this case $Y_{N}$ is also spatially smooth and has moreover regularity $3 / 2^{-}$in time. All the computations done in the sections below are rigorous when the initial datum is smooth, for initial data in a negative Besov space, they are obtained thanks to a preliminary smoothing of the initial datum.

It is not difficult to prove that $\left(X_{N}\right)$ converges to the unique solution of (1.1) (see [TsWe18]) in convenient topologies, $C\left([0, T] ; C^{-\alpha}\right)$ for $\alpha>0$ for instance.

Let us define the Gibbs measure $\nu$ on $\mathcal{H}$

$$
\begin{equation*}
\nu(d x)=z^{-1} e^{-2 U(x)} \mu(d x), \tag{2.4}
\end{equation*}
$$

where

$$
\begin{equation*}
U(x)=\frac{1}{4} \int_{H}: x^{4}:(\xi) d \xi \tag{2.5}
\end{equation*}
$$

and

$$
z:=\int_{H} e^{-2 U(x)} \mu(d x) .
$$

Then

$$
e^{-2 U(x)} \in L^{r}(\mathcal{H}, \mu),
$$

for any $r \in[1, \infty)$ thanks to the Nelson inequality.
We can write equation (1.1) as

$$
\left\{\begin{array}{l}
\frac{d Y}{d t}=A Y-Y^{3}-3 Y^{2} Z-3 Y: Z^{2}:-: Z^{3}:  \tag{2.6}\\
Y(0)=x
\end{array}\right.
$$

or, in mild form,

$$
\begin{equation*}
Y(t)=e^{A t} x-\int_{0}^{t} e^{A(t-s)}\left[Y^{3}(s)+3 Y^{2}(s) Z(s)+3 Y(s): Z^{2}(s):+: Z^{3}(s):\right] d s \tag{2.7}
\end{equation*}
$$

We define the transition semigroups $P_{t}$ and $P_{t}^{N}$ associated to (1.1) and (2.3) respectively and introduce the following modified semigroup:

$$
\begin{equation*}
S_{t}^{N} \varphi(x)=\mathbb{E}\left(e^{-K \int_{0}^{t} V\left(X_{N}(s)\right) d s} \varphi\left(X_{N}(t, x)\right)\right) \tag{2.8}
\end{equation*}
$$

where the potential $V$ is given by

$$
V\left(X_{N}\right)=\left|: X_{N}^{2}:\left.\right|_{H^{-\alpha}} ^{p}=\left|X_{N}^{2}-\rho_{N}^{2}\right|_{H^{-\alpha}}^{p}\right.
$$

for some $K>0, p>0, \alpha>0$ to be chosen. This quantity is well defined for $\varphi \in \mathcal{B}_{b}\left(C^{-\alpha}\right)$, the space of borelian bounded functions on $C^{-\alpha}$.

The following relation formally follows from Duhamel formula since $P_{N} \varphi$ satisfies the Kolmogorov equation associated to (2.3) and $S_{N} \varphi$ satisfies the same equation with an extra term due to the potential $V$ :

$$
\begin{equation*}
P_{t}^{N} \varphi(x)=S_{t}^{N} \varphi(x)+K \int_{0}^{t} S_{t-s}^{N}\left(V P_{s} \varphi\right)(x) d s \tag{2.9}
\end{equation*}
$$

It can be proved rigorously by an approximation argument.

## 3. Estimates

Let us consider the equation:

$$
\left\{\begin{array}{l}
\frac{d \eta^{h, N}}{d t}=\left(A \eta^{h, N}-3: X_{N}^{2}: \eta^{h, N}\right),  \tag{3.1}\\
\eta^{h, N}(0)=h
\end{array}\right.
$$

It is classical that $\eta^{h, N}=D_{x} X_{N} \cdot h$ is the differential of $X_{N}$ with respect to the intial data $x \in C^{-\alpha}$ in the direction $h \in C^{-\alpha}$.

In this section, we derive several estimates on $\eta_{N}$ and $X_{N}$. These hold only on finite time intervals and we restrict to $t \in[0,1]$.
Lemma 3.1. Let $\alpha<\frac{1}{3}, \epsilon>0$ such that $\alpha+\epsilon<\frac{1}{3}$ and $p \geq 1$, $s>0$ such that $\alpha+\epsilon+\frac{2}{3 p}+\frac{s}{3}<\frac{1}{3}$ then there exists $c$ depending on $\alpha, \epsilon, p$ such that for all $t \in(0,1], x \in C^{-\alpha}, h \in C^{-s}$

$$
\left|\eta^{h, N}(t)\right|_{C^{\alpha+\epsilon}} \leq c t^{-(\alpha+\epsilon+s) / 2}|h|_{C^{-s}} e^{c \int_{0}^{t}\left|: X_{N}^{2}(s):\right|_{H^{-\alpha}}^{p} d s}
$$

and

$$
\left|\eta^{h, N}(t)\right|_{L^{\infty}} \leq c t^{-s / 2}|h|_{C^{-s}} e^{c \int_{0}^{t}\left|: X_{N}^{2}(s):\right|_{H^{-\alpha}}^{p} d s} .
$$

Proof. Write (3.1) in mild form

$$
\eta^{h, N}(t)=e^{A t} h-3 \int_{0}^{t} e^{A(t-s)}\left(: X_{N}^{2}(s): \eta^{h, N}(s)\right) d s
$$

Below we omit the dependence on $N$. By the smoothing of the heat semigroup $e^{A t}$, the product rule in Besov spaces and $t \in[0,1]$, we have

$$
\begin{aligned}
& \left|\eta^{h}(t)\right|_{C^{\alpha+\epsilon}} \leq c t^{-(\alpha+\epsilon+s) / 2}|h|_{C^{-s}}+c \int_{0}^{t}|t-r|^{-(\alpha+\epsilon+1 / 2)}\left|: X^{2}(r): \eta^{h}(r)\right|_{H^{-\alpha-\epsilon / 2}} d r \\
& \leq c t^{-(\alpha+\epsilon+s) / 2}|h|_{C^{-s}}+\left.c \int_{0}^{t}|t-r|^{-(\alpha+\epsilon+1 / 2)}\left|: X^{2}(r):\left.\right|_{H^{-\alpha}}\right| \eta^{h}(r)\right|_{C^{\alpha+\epsilon}} d r .
\end{aligned}
$$

Let $\lambda(t)=t^{(\alpha+\epsilon+s) / 2}\left|\eta^{h}(t)\right|_{C^{\alpha+\epsilon}}$, then since $t \in[0,1]$,

$$
\lambda(t) \leq c|h|_{C^{-s}}+c \int_{0}^{t}|t-r|^{-(\alpha+\epsilon+1 / 2)} r^{-(\alpha+\epsilon+s) / 2}\left|: X^{2}(r):\right|_{H^{-\alpha}} \lambda(r) d r .
$$

Let $\frac{1}{p}+\frac{1}{q}=1$ and

$$
\frac{(3 \alpha+3 \epsilon+1+s) q}{2}<1, \quad \text { i.e. } \frac{3 \alpha+3 \epsilon+1+s}{2}<1-\frac{1}{p} .
$$

Then

$$
\begin{aligned}
\lambda^{p}(t) & \leq c|h|_{C^{-s}}^{p}+c\left(\int_{0}^{t}|t-r|^{-q(\alpha+\epsilon+1 / 2)} r^{-q(\alpha+\epsilon+s) / 2} d r\right)^{p / q} \int_{0}^{t}\left|: X^{2}(r):\right|_{H^{-\alpha}}^{p} \lambda^{p}(r) d r \\
& \leq c|h|_{C^{-s}}^{p}+c \int_{0}^{t}\left|: X^{2}(r):\right|_{H^{-\alpha}}^{p} \lambda^{p}(r) d r
\end{aligned}
$$

The first inequality follows by Gronwall's lemma.
For the second inequality, we write:

$$
\begin{aligned}
& \left|\eta^{h}(t)\right|_{L^{\infty}} \leq c t^{-s / 2}|h|_{C^{-s}}+c \int_{0}^{t}|t-r|^{-(\alpha+\epsilon+1 / 2)}\left|: X^{2}(r): \eta^{h}(r)\right|_{H^{-\alpha-\epsilon / 2}} d r \\
& \leq c t^{-s / 2}|h|_{C^{-s}}+\left.c \int_{0}^{t}|t-r|^{-(\alpha+\epsilon+1 / 2)}\left|: X^{2}(r):\left.\right|_{H^{-\alpha}}\right| \eta^{h}(r)\right|_{C^{\alpha+\epsilon}} d r
\end{aligned}
$$

and with $p, q$ as above:

$$
\begin{aligned}
\left|\eta^{h}(t)\right|_{L^{\infty}}^{p} & \leq c t^{-s p / 2}|h|_{C^{-s}}^{p}+\left.c \int_{0}^{t} r^{p(\alpha+\epsilon+s) / 2}\left|: X^{2}(r):\left.\right|_{H^{-\alpha}} ^{p}\right| \eta^{h}(r)\right|_{C^{\alpha+\epsilon}} ^{p} d r \\
& \leq c t^{-s p / 2}|h|_{C^{-s}}^{p}+\left.c \int_{0}^{t}\left|: X^{2}(r):\left.\right|_{H^{-\alpha}} ^{p} e^{c p \int_{0}^{t} r\left|: X_{N}^{2}(s):\right|_{H^{-\alpha}}^{p} d s}\right| h\right|_{C^{-s}} ^{p} d r \\
& \leq c|h|_{C^{-s}}^{p}\left(t^{-s p / 2}+e^{c p \int_{0}^{t}\left|: X_{N}^{2}(s):\right|_{H^{-\alpha}}^{p} d s}\right)
\end{aligned}
$$

The conclusion follows.
We now estimate moments of $Y_{N}$ solution of (2.2). We estimate $Y_{N}$ in $L^{r}\left(\mathbb{T}^{2}\right), r \geq 2$, using similar arguments as in [TsWe18] and [MoWe15], but in a simplified way since we do not need such refined estimates for large times.

Lemma 3.2. Let $\alpha \in[0,1)$ and $r \geq 2$, there exist $c>0$ depending on $\alpha$ and $r$ such that

$$
\frac{1}{r} \frac{d}{d t}\left|Y_{N}\right|_{L^{r}}^{r}+\frac{r-1}{r^{2}}\left|\nabla\left(\left(Y_{N}\right)^{r / 2}\right)\right|_{L^{2}}^{2}+\frac{1}{2}\left|Y_{N}\right|_{L^{r+2}}^{r+2} \leq c \mathcal{L}_{\alpha}\left(Z_{N}\right)^{k_{r}}
$$

where $k_{r}=\frac{r+2}{1-\alpha}$ and

$$
\mathcal{L}_{\alpha}\left(Z_{N}\right)=\left|Z_{N}\right|_{C^{-\alpha}}+\left|: Z_{N}^{2}:\left.\right|_{C^{-\alpha}}+\left|: Z_{N}^{3}:\right|_{C^{-\alpha}}+1\right.
$$

Proof. Again we omit the dependence on $N$. Multiplying equation (2.2) by $Y^{r-1}$ and integrating on $\mathbb{T}^{2}$ we obtain:

$$
\frac{1}{r} \frac{d}{d t}|Y|_{L^{r}}^{r}+\frac{4(r-1)}{r^{2}}\left|\nabla\left(Y^{r / 2}\right)\right|_{L^{2}}^{2}+|Y|_{L^{r+2}}^{r+2}=-\int_{\mathbb{T}^{2}}\left(3 Y^{r+1} Z+3 Y^{r}: Z^{2}:+Y^{r-1}: Z^{3}:\right) d \xi
$$

We estimate each of the three terms in the right hand side. By Proposition A. 8 and Proposition A. 9 in [TsWe18], we have

$$
3\left|\int_{\mathbb{T}^{2}} Y^{r+1} Z d \xi\right| \leq 3\left|Y^{r+1}\right|_{B_{1,1}^{\alpha}}|Z|_{C^{-\alpha}} \leq c\left(\left|Y^{r+1}\right|_{L^{1}}^{1-\alpha}\left|\nabla Y^{r+1}\right|_{L^{1}}^{\alpha}+\left|Y^{r+1}\right|_{L^{1}}\right)|Z|_{C^{-\alpha}}
$$

Clearly

$$
\left|Y^{r+1}\right|_{L^{1}}=|Y|_{5}^{L^{r+1}} \underset{r+1}{s} \leq|Y|_{L^{r+2}}^{r+1}
$$

Moreover

$$
\left|\nabla Y^{r+1}\right|_{L^{1}}=(r+1)\left|Y^{r} \nabla Y\right|_{L^{1}} \leq(r+1)\left|Y^{\frac{r+2}{2}}\right|_{L^{2}}\left|Y^{\frac{r-2}{2}} \nabla Y\right|_{L^{2}}=\frac{2(r+1)}{r}|Y|_{L^{r+2}}^{\frac{r+2}{2}}\left|\nabla Y^{r / 2}\right|_{L^{2}} .
$$

It follows that

$$
\begin{aligned}
& 3\left|\int_{\mathbb{T}^{2}} Y^{r+1} Z d \xi\right| \leq c\left(|Y|_{L^{r+2}}^{(1-\alpha)(r+1)+\alpha(r+2) / 2}\left|\nabla Y^{r / 2}\right|_{L^{2}}^{\alpha}+|Y|_{L^{r+2}}^{r+1}\right)|Z|_{C^{-\alpha}} \\
& =c\left(|Y|_{L^{r+2}}^{((1-\alpha / 2) r+1}\left|\nabla Y^{r / 2}\right|_{L^{2}}^{\alpha}+|Y|_{L^{r+2}}^{r+1}\right)|Z|_{C^{-\alpha}} \\
& \leq \frac{1}{6}|Y|_{L^{r+2}}^{r+2}+\frac{r-1}{r^{2}}\left|\nabla Y^{r / 2}\right|_{L^{2}}^{2}+c\left(|Z|_{C^{-\alpha}}^{p_{1}}+|Z|_{C^{-\alpha}}^{p_{2}}\right)
\end{aligned}
$$

with

$$
\frac{(1-\alpha / 2) r+1}{r+2}+\frac{\alpha}{2}+\frac{1}{p_{1}}=1
$$

and $p_{2}=r+2$. Note that since $\alpha \in(0,1), p_{1}$ exists and is equal to $\frac{r+2}{1-\alpha}$. We proceed similarly for the second one and obtain:

$$
3\left|\int_{\mathbb{T}^{2}} Y^{r}: Z^{2}: d \xi\right| \leq \frac{1}{6}|Y|_{L^{r+2}}^{r+2}+\frac{r-1}{r^{2}}\left|\nabla Y^{r / 2}\right|_{L^{2}}^{2}+c\left(\left|: Z^{2}:\left.\right|_{C^{-\alpha}} ^{p_{3}}+\left|: Z^{2}:\right|_{C^{-\alpha}}^{p_{4}}\right)\right.
$$

with

$$
\frac{\alpha}{2}+\frac{r(1-\alpha / 2)}{r+2}+\frac{1}{p_{3}}=1
$$

and $p_{4}=\frac{r+2}{2}$ and for the third one

$$
\left|\int_{\mathbb{T}^{2}} Y^{r-1}: Z^{3}: d \xi\right| \leq \frac{1}{6}|Y|_{L^{r+2}}^{r+2}+\frac{r-1}{r^{2}}\left|\nabla Y^{r / 2}\right|_{L^{2}}^{2}+c\left(\left|: Z^{3}:\left.\right|_{C^{-\alpha}} ^{p_{5}}+\left|: Z^{3}:\right|_{C^{-\alpha}}^{p_{6}}\right)\right.
$$

with

$$
\frac{\alpha}{2}+\frac{r(1-\alpha / 2)-1}{r+2}+\frac{1}{p_{5}}=1
$$

and $p_{6}=\frac{r+2}{3}$.
Using these three inequalities in the energy equality and noting that $p_{1}=\frac{r+2}{1-\alpha} \geq p_{i}, i=2, \ldots, 5$ we obtain

$$
\frac{1}{2} \frac{d}{d t}|Y|_{L^{r}}^{r}+\frac{r-1}{r^{2}}\left|\nabla\left(Y^{r / 2}\right)\right|_{L^{2}}^{2}+\frac{1}{2}|Y|_{L^{r+2}}^{r+2} \leq c\left(|Z|_{C^{-\alpha}}+\left|: Z^{2}:\left.\right|_{C^{-\alpha}}+\left|: Z^{3}:\right|_{C^{-\alpha}}+1\right)^{p_{1}}\right.
$$

This gives the result with $k_{r}=p_{1}$.
Lemma 3.3. Let $\alpha<\frac{1}{3}, k \geq 1$ there exists $C_{\alpha, k}>0$ and $\kappa_{\alpha, k}>0$ such that

$$
\mathbb{E}\left(\sup _{t \in[0,1]}\left|Y_{N}(t)\right|_{C^{-\alpha}}^{k}\right) \leq C_{\alpha, k}\left(|x|_{C^{-\alpha}}+1\right)^{\kappa_{\alpha, k}}
$$

Proof. Since

$$
\left|Y_{N}\right|_{L^{r}}^{r+2} \leq c\left|Y_{N}\right|_{L^{r+2}}^{r+2}
$$

it follows from Lemma 3.2 that

$$
\frac{d}{d t}\left|Y_{N}(t)\right|_{L^{r}}^{r}+c\left|Y_{N}(t)\right|_{L^{r}}^{r+2} \leq c \mathcal{L}_{\alpha}\left(Z_{N}(t)\right)^{k_{r}}
$$

and by [TsWe18, Lemma 3.8] it follows for $t \in[0,1]$

$$
\left|Y_{N}(t)\right|_{L^{r}} \leq c \max \left(t^{-r / 2}, \sup _{t \in[0,1]} \mathcal{L}_{\alpha}\left(Z_{N}(t)\right)^{\frac{r}{1-\alpha}}\right)
$$

We choose $r>\frac{2}{\alpha}$ and use the embedding $L^{r} \subset C^{-\alpha}$ to deduce

$$
\begin{equation*}
\left|Y_{N}(t)\right|_{C^{-\alpha}} \leq c \max \left(t^{-r / 2}, \sup _{t \in[0,1]} \mathcal{L}_{\alpha}\left(Z_{N}(t)\right)^{\frac{r}{1-\alpha}}\right) . \tag{3.2}
\end{equation*}
$$

It remain to obtain a bound for $\left|Y_{N}(t)\right|_{C^{-\alpha}}$ for small time. Write the mild form for the equation satisfied by $Y_{N}$ :

$$
Y_{N}(t)=e^{A t} x-\int_{0}^{t} e^{A(t-s)}\left(Y_{N}^{3}(s)+3 Y_{N}^{2}(s) Z_{N}(s)+3 Y_{N}(s): Z_{N}^{2}(s):+: Z_{N}^{3}(s):\right) d s
$$

Let $\alpha<\frac{1}{3}, \beta>\alpha$ such that

$$
\gamma:=\frac{\alpha+\beta}{2}<\frac{1}{3}
$$

and define

$$
M(t)=\sup _{s \in[0, t]}\left(\left|Y_{N}(s)\right|_{C^{-\alpha}}+s^{\gamma}\left|Y_{N}(s)\right|_{C^{-\beta}}\right) .
$$

Then similar computations as in the proof of Theorem 3.3 in [TsWe18] give for $t \in[0,1]$ :

$$
M(t) \leq c_{1}|x|_{C^{-\alpha}}+c_{2} t^{1-3 \gamma} M(t)^{3}+c_{3} t^{1-3 \gamma} \sup _{t \in[0,1]} \mathcal{L}_{\alpha}\left(Z_{N}(t)\right)^{3},
$$

for some constants $c_{1}, c_{2}$ which can be chosen larger than 1 .
Therefore if we set

$$
t^{*}=\min \left\{\left(c_{3} \sup _{t \in[0,1]} \mathcal{L}_{\alpha}\left(Z_{N}(t)\right)^{3}\right)^{-\frac{1}{1-3 \gamma}},\left(8 c_{2}\left(c_{1}|x|_{C^{-\alpha}}+1\right)^{2}\right)^{-\frac{1}{1-3 \gamma}}, 1\right\}
$$

we have

$$
\begin{equation*}
\sup _{t \in\left[0, t^{*}\right]} M(t) \leq 2\left(c_{1}|x|_{C^{-\alpha}}+1\right) \tag{3.3}
\end{equation*}
$$

It follows from (3.2) and (3.3):

$$
\begin{aligned}
\sup _{t \in[0,1]}\left|Y_{N}(t)\right|_{C^{-\alpha}} & \leq \sup _{t \in\left[0, t^{*}\right]}\left|Y_{N}(t)\right|_{C^{-\alpha}}+\sup _{t \in\left[t^{*}, 1\right]}\left|Y_{N}(t)\right|_{C^{-\alpha}} \\
& \leq 2\left(c_{1}|x|_{C^{-\alpha}}+1\right)+\max \left(\left(t^{*}\right)^{-2 / r}, \sup _{t \in[0,1]} \mathcal{L}_{\alpha}\left(Z_{N}(t)\right)^{\frac{r k_{r}}{r+2}}\right)
\end{aligned}
$$

The result follows taking the expectation and using finiteness of the moments of

$$
\sup _{t \in[0,1]} \mathcal{L}_{\alpha}\left(Z_{N}(t)\right) .
$$

We set

$$
\widetilde{Z}_{N}(t)=e^{A t} x+\Pi_{N} \int_{0}^{t} e^{A(t-s)} d W(s)=e^{A t} x+Z_{N}(t)
$$

and

$$
\widetilde{Y}_{N}(t)=X_{N}(t)-\widetilde{Z}_{N}(t)
$$

Then $\widetilde{Y}_{N}$ satisfies the same equation as $Y_{N}$, see (2.2), with $Z_{N}$ replaced by $\widetilde{Z}_{N}$ and 0 initial datum.

Lemma 3.4. Let $\alpha \in(0,1 / 9), k \geq 1$ then

$$
\mathbb{E}\left(\sup _{t \in[0,1]}|\widetilde{Y}|_{L^{2}}^{2}+\int_{0}^{t}\left|\nabla \widetilde{Y}_{N}(s)\right|_{L^{2}}^{2} d s\right)^{k} \leq C_{\alpha, k}\left(|x|_{C^{-\alpha}}^{12 k /(1-\alpha)}+1\right)
$$

for any $t>0$.
Proof. Lemma 3.2 clearly also holds for $\widetilde{Y}_{N}$ provided $Z_{N}$ is replaced by $\widetilde{Z}_{N}$. We use it with $r=2$ and again omit to indicate the dependency on $N$ :

$$
\frac{1}{2} \frac{d}{d t}|\widetilde{Y}|_{L^{2}}^{2}+\frac{1}{4}|\nabla \widetilde{Y}|_{L^{2}}^{2}+\frac{1}{2}|\widetilde{Y}|_{L^{4}}^{4} \leq c \mathcal{L}_{\alpha}(\widetilde{Z})^{k_{2}}
$$

$c$ an $k_{2}=\frac{4}{1-\alpha}$ are given in Lemma 3.2. We use the product rules in Besov spaces and the smoothing effect of $e^{A t}$ to prove:

$$
\mathcal{L}_{\alpha}(\widetilde{Z}(s)) \leq C_{\alpha, \epsilon}\left(\left(s^{-(2 \alpha+\epsilon)}+1\right)|x|_{C^{-\alpha}}^{3}+1\right) \mathcal{L}_{\alpha}(Z(s))
$$

the result follows by integrating on $[0, T]$ and taking expectation of the $k^{\text {th }}$ power. Note that since $\alpha<1 / 9$, it is possible to choose $\epsilon>0$ small enough $(2 \alpha+\epsilon) k_{2}<1$.

Lemma 3.5. Let $\alpha \in(0,1 / 9), k \geq 1$, there exists $\gamma_{\alpha, k}, C_{\alpha, k}$ such that

$$
\mathbb{E}\left(\sup _{t \in[0,1]}\left|\widetilde{Y}_{N}(t)\right|_{H^{\alpha+\epsilon}}^{k}\right) \leq C_{\alpha, k}\left(|x|_{C^{-\alpha}}+1\right)^{\gamma_{\alpha, k}} .
$$

Proof. We write, omitting again $N$

$$
\widetilde{Y}(t)=\int_{0}^{t} e^{A(t-s)}\left(\widetilde{Y}^{3}(s)+3 \widetilde{Y}^{2}(s) \widetilde{Z}(s)+3 \widetilde{Y}(s): \widetilde{Z}^{2}(s):+: \widetilde{Z}^{3}(s):\right) d s
$$

and get

$$
|\widetilde{Y}(t)|_{H^{\alpha+\epsilon}} \leq c \int_{0}^{t}(t-s)^{-(\alpha+\epsilon / 2)}\left(\left|\widetilde{Y}^{3}\right|_{H^{-\alpha}}+\left|\widetilde{Y}^{2} \widetilde{Z}\right|_{H^{-\alpha}}+\left|\widetilde{Y}: \widetilde{Z}^{2}:\left.\right|_{H^{-\alpha}}+\left|: \widetilde{Z}^{3}:\right|_{H^{-\alpha}}\right) d s\right.
$$

We have

$$
\begin{aligned}
\left|\widetilde{Y}^{3}\right|_{H^{-\alpha}} & \leq c\left|\widetilde{Y}^{2}\right|_{H^{\alpha+\epsilon}}|\widetilde{Y}|_{C^{-\alpha}} \leq c|\widetilde{Y}|_{B_{4,4}^{\alpha+\epsilon}}^{2}|\widetilde{Y}|_{C^{-\alpha}} \leq c|\widetilde{Y}|_{H^{\alpha+\epsilon+1 / 2}}^{2}|\widetilde{Y}|_{C^{-\alpha}} \\
& \leq c\left(|\widetilde{Y}|_{L^{2}}^{2(1 / 2-\alpha-\epsilon)}|\nabla \widetilde{Y}|_{L^{2}}^{2(1 / 2+\alpha+\epsilon)}+|\widetilde{Y}|_{L^{2}}^{2}\right)|\widetilde{Y}|_{C^{-\alpha}} .
\end{aligned}
$$

Similarly

$$
\left|\widetilde{Y}^{2} \widetilde{Z}\right|_{H^{-\alpha}} \leq c\left|\widetilde{Y}^{2}\right|_{H^{\alpha+\epsilon}}|\widetilde{Z}|_{C^{-\alpha}} \leq c\left(|\widetilde{Y}|_{L^{2}}^{2(1 / 2-\alpha-\epsilon)}|\nabla \widetilde{Y}|_{L^{2}}^{2(1 / 2+\alpha+\epsilon)}+|\widetilde{Y}|_{L^{2}}^{2}\right)|\widetilde{Z}|_{C^{-\alpha}}
$$

and

$$
\left.\left|\widetilde{Y}: \widetilde{Z}^{2}:\left.\right|_{H^{-\alpha}} \leq c\right| \widetilde{Y}\right|_{H^{\alpha+\epsilon}}\left|: \widetilde{Z}^{2}:\left.\right|_{C^{-\alpha}} \leq c\left(|\widetilde{Y}|_{L^{2}}^{1-(\alpha+\epsilon)}|\nabla \widetilde{Y}|_{L^{2}}^{\alpha+\epsilon}+|\widetilde{Y}|_{L^{2}}\right)\right|: \widetilde{Z}^{2}:\left.\right|_{C^{-\alpha}}
$$

Gathering these results gives

$$
\begin{aligned}
|\widetilde{Y}(t)|_{H^{\alpha+\epsilon}} & \leq c \int_{0}^{t}(t-s)^{-(\alpha+\epsilon / 2)}\left[\left(|\widetilde{Y}|_{L^{2}}^{2(1 / 2-\alpha-\epsilon)}|\nabla \widetilde{Y}|_{L^{2}}^{2(1 / 2+\alpha+\epsilon)}+|\widetilde{Y}|_{L^{2}}^{2}\right)\left(|\widetilde{Y}|_{C^{-\alpha}}+|\widetilde{Z}|_{C^{-\alpha}}\right)\right. \\
& +\left(|\widetilde{Y}|_{L^{2}}^{1-(\alpha+\epsilon)}|\nabla \widetilde{Y}|_{L^{2}}^{\alpha+\epsilon}+|\widetilde{Y}|_{L^{2}}\right)\left|: \widetilde{Z}^{2}:\left.\right|_{C^{-\alpha}}+\left|: \widetilde{Z}^{3}:\right|_{C^{-\alpha}}\right] d s=:(a)+(b)+(c) .
\end{aligned}
$$

The first term is bounded thanks to Hölder inequality

$$
\begin{aligned}
& (a)=c \int_{0}^{t}(t-s)^{-(\alpha+\epsilon / 2)}\left[\left(|\widetilde{Y}|_{L^{2}}^{2(1 / 2-\alpha-\epsilon)}|\nabla \widetilde{Y}|_{L^{2}}^{2(1 / 2+\alpha+\epsilon)}+|\widetilde{Y}|_{L^{2}}^{2}\right)\left(|\widetilde{Y}|_{C^{-\alpha}}+|\widetilde{Z}|_{C^{-\alpha}}\right) d s\right. \\
& \leq c\left(|\widetilde{Y}|_{L^{p_{1}}\left(0, t, L^{2}\right)}^{2(1 / 2-\alpha-\epsilon)}\left(|\nabla \widetilde{Y}|_{L^{2}\left(0, t, L^{2}\right)}^{2(1 / 2+\alpha+\epsilon)}+|\widetilde{Y}|_{L^{p_{2}}\left(0, t, L^{2}\right)}\right)\left(|Y|_{L^{p_{3}\left(0, t, C^{-\alpha}\right)}}+|x|_{C^{-\alpha}}+|Z|_{L^{p_{3}\left(0, t, C^{-\alpha}\right)}}\right)\right.
\end{aligned}
$$

with

$$
\frac{2(1 / 2-\alpha-\epsilon)}{p_{1}}+(1 / 2+\alpha+\epsilon)+\alpha+\frac{\epsilon}{2}<1+\frac{1}{p_{3}}
$$

and

$$
\frac{2}{p_{2}}+\alpha+\frac{\epsilon}{2}+\frac{1}{p_{3}}<1
$$

Since $\alpha<\frac{1}{4}, \epsilon, p_{1}, p_{2}, p_{3}$ can be chosen such this is satisfied. Then

$$
\begin{aligned}
(b) & \leq c \int_{0}^{t}(t-s)^{-(\alpha+\epsilon / 2)}\left(|\widetilde{Y}|_{L^{2}}^{1-(\alpha+\epsilon)}|\nabla \widetilde{Y}|_{L^{2}}^{\alpha+\epsilon}+|\widetilde{Y}|_{L^{2}}\right) \\
& \times\left[s^{-(\alpha+\epsilon / 2)}\left(|x|_{C^{-\alpha}}^{2}+|x|_{C^{-\alpha}}|Z|_{C^{-\alpha}}\right)+\left|: Z^{2}:\right|_{C^{-\alpha}}\right] d s \\
& \leq c\left(|\widetilde{Y}|_{L^{p_{4}\left(0, t, L^{2}\right)}}^{1-(\alpha+\epsilon)}|\nabla \widetilde{Y}|_{L^{2}\left(0, t, L^{2}\right)}^{\alpha+\epsilon}+|\widetilde{Y}|_{L^{p_{5}\left(0, t, L^{2}\right)}}\right) \\
& \times\left(|x|_{C^{-\alpha}}^{2}+|x|_{C^{-\alpha}}|Z|_{L^{p_{6}\left(0, t, C^{-\alpha}\right)}}+\left|: Z^{2}:\right|_{L^{p_{6}}\left(0, t, C^{-\alpha}\right)}\right)
\end{aligned}
$$

with

$$
\alpha+\frac{\epsilon}{2}+\frac{1-(\alpha+\epsilon)}{p_{4}}+\frac{\alpha+\epsilon}{2}+\alpha+\frac{\epsilon}{2}+\frac{1}{p_{6}}<1
$$

and

$$
2\left(\alpha+\frac{\epsilon}{2}\right)+\frac{1}{p_{5}}+\frac{1}{p_{6}}<1
$$

which again is possible since $\alpha<\frac{1}{4}$. Similarly

$$
\begin{aligned}
& (c) \leq c \int_{0}^{t}(t-s)^{-(\alpha+\epsilon / 2)}\left(s^{-(2 \alpha+\epsilon)}\left(|x|_{C^{-\alpha}}^{3}+|x|_{C^{-\alpha}}^{2}|Z|_{C^{-\alpha}}\right)\right. \\
& +s^{-(\alpha+\epsilon / 2)}|x|_{C^{-\alpha}}\left|: Z^{2}:\left.\right|_{C^{-\alpha}}+\left|: Z^{3}:\right|_{C^{-\alpha}}\right) d s
\end{aligned}
$$

Taking expectation we find thanks to Lemma 3.3, Lemma 3.4, Hölder inequality and finiteness of the moments of $\sup _{t \in[0,1]}|Z|_{C^{-\alpha}}, \sup _{t \in[0,1]}\left|: Z^{2}:\right|_{C^{-\alpha}}$ and $\sup _{t \in[0,1]}\left|: Z^{3}:\right|_{C^{-\alpha}}:$

$$
\mathbb{E}\left(\sup _{t \in[0,1]}|\widetilde{Y}(t)|_{H^{\alpha+\epsilon}}^{k}\right) \leq c\left(|x|_{C^{-\alpha}}^{\kappa_{\alpha, k}}+1\right)
$$

for some $\kappa_{\alpha, k}$ which could be written explicitly.
Lemma 3.6. Let $\alpha \in(0,1 / 9)$, for any $k \geq 1$, there exists $\bar{\gamma}_{\alpha, k}, C_{\alpha, k}$ such that

$$
\mathbb{E}\left(\left|: X_{N}(t)^{2}:\right|_{H^{-\alpha}}^{k} d s\right) \leq C_{\alpha, k}\left(t^{-\alpha k}+1\right)\left(|x|_{C^{-\alpha}}+1\right)^{\bar{\gamma}_{\alpha, k}} .
$$

Proof. We have, omitting the dependence in $N$,

$$
\begin{aligned}
\left|: X(t)^{2}:\right|_{H^{-\alpha}} & \leq c\left(\left|Y^{2}(t)\right|_{H^{-\alpha}}+|Y(t)|_{H^{\alpha+\epsilon}}|Z(t)|_{C^{-\alpha}}+\left|: Z(t)^{2}:\right|_{C^{-\alpha}}\right) \\
& \leq c\left(|Y(t)|_{H^{\alpha+\epsilon}}\left(|Y(t)|_{C^{-\alpha}}+|Z(t)|_{C^{-\alpha}}\right)+\left|: Z(t)^{2}:\right|_{C^{-\alpha}}\right) \\
& \leq c\left(\left(t^{-\alpha}|x|_{C^{-\alpha}}+|\widetilde{Y}(t)|_{H^{\alpha+\epsilon}}\right)\left(|Y(t)|_{C^{-\alpha}}+|Z(t)|_{C^{-\alpha}}\right)+\left|: Z(t)^{2}:\right|_{C^{-\alpha}}\right)
\end{aligned}
$$

The result follows by Lemma 3.3 and Lemma 3.5.
3.1. Estimates on $S_{t}^{N}$. We have

$$
S_{t}^{N} \varphi(x)=\mathbb{E}\left(e^{-K \int_{0}^{t} V\left(X_{N}(s)\right) d s} \varphi\left(X_{N}(t, x)\right)\right)
$$

where the potential $V$ is given by

$$
V(x)=\left|: x^{2}:\right|_{H^{-\alpha}}^{p} .
$$

Lemma 3.7. Let $\alpha \in\left(0, \frac{1}{9}\right), s>0, p \geq 1$ such that $\alpha+\frac{1}{3 p}+\frac{s}{3}<\frac{1}{3}$ and $\alpha(2 p-1)+s<2$. Then there exists $\gamma$ such that for any $\varphi \in \mathcal{B}_{b}\left(C^{-\alpha}\right)$ we have for $t \in[0,1], x \in C^{-\alpha}, h \in H$

$$
\left|D S_{t}^{N} \varphi(x) \cdot h\right| \leq c\left(t^{-(1+s+2 \alpha p) / 2}+1\right)|h|_{C^{-s}}\|\varphi\|_{0}\left(1+|x|_{C^{-\alpha}}\right)^{\gamma}
$$

and

$$
\left|D S_{t}^{N}(V \varphi)(x) \cdot h\right| \leq c\left(t^{-(1+s+2 \alpha p) / 2}+1\right)|h|_{C^{-s}}\|\varphi\|_{0}\left(1+|x|_{C^{-\alpha}}\right)^{\gamma} .
$$

Remark 3.8. $V \varphi$ is not bounded but it is clear from Lemma 3.6 that $S_{t}^{N}$ can be extended to borelian functions with polynomial growth.

Proof. We only prove the second inequality, the first one is similar.
We omit to mention the dependence of $N$. From [DaZa97] we know that $S_{t}^{N} \varphi$ is differentiable in any direction $h \in H$ and we have the following modified Bismut-Elworthy-Li formula:

$$
\begin{aligned}
D S_{t}(V \varphi)(x) \cdot h= & \frac{1}{t} \mathbb{E}\left(e^{-K \int_{0}^{t} V(X(s, x)) d s} V(X(t, x)) \varphi(X(t, x)) \int_{0}^{t}\left(\eta^{h}(s), d W(s)\right)\right) \\
& \left.+2 K \mathbb{E}\left(e^{-K \int_{0}^{t} V(X(s, x)) d s} V(X(t, x)) \varphi(X(t, x)) \int_{0}^{t}\left(1-\frac{s}{t}\right) V^{\prime}(X(s, x)) \cdot \eta^{h}(s) d s\right)\right) \\
= & (a)+(b) .
\end{aligned}
$$

Clearly

$$
\begin{aligned}
&(a) \leq \frac{1}{t}\|\varphi\|_{0}\left[\mathbb{E}(V(X(t, x)))^{2}\right]^{1 / 2} \\
& \times\left[\mathbb { E } \left(e^{-2 K} \int_{0}^{t} V(X(s, x) d s\right.\right. \\
&\left.\left.\left|\int_{0}^{t}\left(\eta^{h}(s), d W(s)\right)\right|^{2}\right)\right]^{1 / 2}
\end{aligned}
$$

By Lemma 3.6 we have for $K$ large enough

$$
\left[\mathbb{E}(V(X(t, x)))^{2}\right]^{1 / 2} \leq c\left(t^{-\alpha p}+1\right)\left(1+|x|_{C^{-\alpha}}\right)^{\bar{\gamma}_{\alpha, 2 p} / 2}
$$

To estimate the second factor, we proceed as in the proof of [DaDe03-2, Lemma 4.1]. Writing Itô's formula for

$$
d\left|e^{-K \int_{0}^{t} V(X(s, x) d s} \int_{10}^{t}\left(\eta^{h}(s), d W(s)\right)\right|^{2}
$$

we deduce for $K$ large enough

$$
\begin{aligned}
& \mathbb{E}\left(e^{-2 K \int_{0}^{t} V(X(s, x) d s}\left|\int_{0}^{t}\left(\eta^{h}(s), d W(s)\right)\right|^{2}\right)^{1 / 2} \\
& \leq \mathbb{E}\left(\int_{0}^{t} e^{-2 K \int_{0}^{s} V(X(r, x) d r}\left|\eta^{h}(s)\right|^{2} d s\right)^{1 / 2} \\
& \leq c \mathbb{E}\left(\int_{0}^{t} \sigma^{-s} d s\right)^{1 / 2}|h|_{C^{-s}}
\end{aligned}
$$

thanks to Lemma 3.1 and the embedding $L^{\infty} \subset L^{2}$. In Lemma 3.1, we choose $\epsilon>0$ such that $\alpha+\epsilon+\frac{1}{3 p}+\frac{s}{3}<\frac{1}{3}$.

It follows that

$$
(a) \leq c\left(t^{-(1+s) / 2}+1\right)\left(t^{-\alpha p}+1\right)|h|_{C^{-s}}\|\varphi\|_{0}\left(1+|x|_{C^{-\alpha}}\right)^{\gamma_{\alpha, 2 p} / 2}
$$

Similarly, the other term is bounded by

$$
(b) \leq c\|\varphi\|_{0}\left[\mathbb{E}(V(X(t, x)))^{2}\right]^{1 / 2} \mathbb{E}\left(e^{-2 K \int_{0}^{t} V(X(s, x) d s}\left|\int_{0}^{t} V^{\prime}(X(s, x)) \cdot \eta^{h}(s) d s\right|^{2}\right)^{1 / 2}
$$

We have

$$
\begin{aligned}
\left|V^{\prime}(X(s, x)) \cdot \eta^{h}(s)\right| & =2 p \mid\left(: X^{2}(s):\left.\right|_{H^{-\alpha}} ^{p-2}\left|\left(: X^{2}(s):, X(s) \eta^{h}(s)\right)\right|_{H^{-\alpha}}\right. \\
& \leq\left. 2 p\left|: X^{2}(s):\left.\right|_{H^{-\alpha}} ^{p-1}\right| X(s) \eta^{h}(s)\right|_{H^{-\alpha}} \\
& \leq 2 p\left|: X^{2}(s):\left|H_{H^{-\alpha}}^{p-1}\right| X(s)\right|_{H^{-\alpha}}\left|\eta^{h}(s)\right|_{C^{\alpha+\epsilon}} .
\end{aligned}
$$

By Lemma 3.1, Minkowski inequality, Lemma 3.3 and Lemma 3.5, we thus can write for $K$ large enough:

$$
\begin{aligned}
& E\left(e^{-2 K \int_{0}^{t} V(X(s, x) d s}\left|\int_{0}^{t} V^{\prime}(X(s, x)) \cdot \eta^{h}(s) d s\right|^{2}\right)^{1 / 2} \\
& \leq c \mathbb{E}\left(\left(\left.\int_{0}^{t}\left|: X^{2}(s):\left.\right|_{H^{-\alpha}} ^{p-1}\right| X(s)\right|_{H^{-\alpha}} s^{-(\alpha+\epsilon+s) / 2} d s\right)^{2}\right)^{1 / 2}|h|_{C^{-s}} \\
& \leq c \int_{0}^{t} \mathbb{E}\left(\left(\left.\left|: X^{2}(s):\left.\right|_{H^{-\alpha}} ^{p-1}\right| X(s)\right|_{H^{-\alpha}}\right)^{2}\right)^{1 / 2} s^{-(\alpha+\epsilon+s) / 2} d s|h|_{C^{-s}} \\
& \leq c\left(\int_{0}^{t}\left(1+s^{-(\alpha+\epsilon+s+2 \alpha(p-1)) / 2}\right) d s\right)\left(1+|x|_{C^{-\alpha}}\right)^{\left(\bar{\gamma}_{\alpha, 2 p} / 2+\bar{\gamma}_{\alpha, 4(p-1)}+\kappa_{\alpha, 4}\right) / 4}|h|_{C^{-s}}
\end{aligned}
$$

We deduce for $\epsilon>0$ small enough so that $\alpha+\epsilon+s+2 \alpha(p-1))<2$ :

$$
(b) \leq c\|\varphi\|_{0}\left(1+t^{-\alpha p}\right)\left(1+|x|_{C^{-\alpha}}\right)^{\left(\bar{\gamma}_{\alpha, 4(p-1)}+\kappa_{\alpha, 4}\right) / 4}|h|_{C^{-s}}
$$

The conclusion follows when gathering the estimates on $(a)$ and $(b)$.
3.2. Estimate on $P_{t}^{N}$. Recall that

$$
P_{t}^{N} \varphi(x)=S_{t}^{N} \varphi(x)+K \int_{0}^{t} S_{t-s}^{N}\left(V P_{s}^{N} \varphi\right) d s
$$

Theorem 3.9. Let $s \in(0,1)$ and $\alpha \in(0,1 / 9)$ such that $\frac{2 \alpha}{1-s}<1-s-3 \alpha$. Let $p \geq 1$ such that $\frac{2 \alpha}{1-s}<\frac{1}{p}<1-s-3 \alpha$, then there exists $\gamma>0$, and $c>0$ depending on $s, \alpha$ and $p$ such that for any $\varphi \in \mathcal{B}_{b}\left(C^{-\alpha}\right)$ we have for $t \in[0,1], x \in C^{-\alpha}, h \in H$ :

$$
\left|D P_{t}^{N} \varphi(x) \cdot h\right| \leq c\left(t^{-(1+s+2 \alpha p) / 2}+1\right)|h|_{C^{-s}}\|\varphi\|_{0}\left(1+|x|_{C^{-\alpha}}\right)^{\gamma} .
$$

Remark 3.10. It is easy to find $\alpha, s$ satisfying the assumptions of Theorem 3.9. For instance, choose any $s \in(0,1), \delta \in\left(0, \frac{1-s}{2}\right)$ and $\alpha<\min \left\{\frac{\delta(1-s)}{1+3 \delta} ; s\right\}$. Then

$$
\frac{\alpha}{\delta}<1-s-3 \alpha .
$$

Moreover

$$
\frac{\alpha}{\delta}>\frac{2 \alpha}{1-s} .
$$

It follows that one may choose $p$ such that $\frac{2 \alpha}{1-s}<\frac{\alpha}{\delta}<\frac{1}{p}<1-s-3 \alpha$ and Theorem 3.9 implies

$$
\left|D P_{t}^{N} \varphi(x) \cdot h\right| \leq c\left(t^{-(1+s+\delta) / 2}+1\right)|h|_{C^{-s}}\|\varphi\|_{0}\left(1+|x|_{C^{-\alpha}}\right)^{\gamma} .
$$

Remark 3.11. Since $H$ is dense in $C^{-s}$, this result shows that $D P_{t}^{N} \varphi(x)$ belongs to the dual of $C^{-s}$ and the bound extends to $h \in C^{-s}$.

Proof. Since for $t \geq 1,\left\|P_{t} \varphi\right\|_{0} \leq\|\varphi\|_{0}$, it suffices to prove the result for $t \in[0,1]$ and use the Markov property.

Since $\frac{2 \alpha}{1-s}>\frac{2 \alpha}{2-s+\alpha}$, we have $\frac{2 \alpha}{2-s+\alpha}<\frac{1}{p}<1-s-3 \alpha$ and this is equivalent to the assumptions of Lemma 3.7, therefore we may write:

$$
\begin{aligned}
\left|D P_{t}^{N} \varphi(x) \cdot h\right| \leq & \left|D S_{t}^{N} \varphi(x) \cdot h\right|+K \int_{0}^{t}\left|D S_{t-s}^{N}\left(V P_{s} \varphi\right)(x)\right| d s \\
\leq & c\left(t^{-(1+s+2 \alpha p) / 2}+1\right)|h|_{C^{-s}}\|\varphi\|_{0}\left(1+|x|_{C^{-\alpha}}\right)^{\gamma} \\
& +K \int_{0}^{t} c(t-s)^{-(1+s+2 \alpha p) / 2}|h|_{C^{-s}}\|\varphi\|_{0}\left(1+|x|_{C^{-\alpha}}\right)^{\gamma} d s .
\end{aligned}
$$

This implies the result since $1+s+2 \alpha p<2$.
Letting $N \rightarrow \infty$ in Theorem 3.9, we deduce:
Corollary 3.12. Let $s \in(0,1)$ and $\alpha \in(0,1 / 9)$ such that $\frac{2 \alpha}{1-s}<1-s-3 \alpha$. Let $p \geq 1$ such that $\frac{2 \alpha}{1-s}<\frac{1}{p}<1-s-3 \alpha$, then there exists $\gamma>0$, and $c>0$ depending on $s, \alpha$ and $p$ such that for any $\varphi \in \mathcal{B}_{b}(H)$ we have for $t \in[0,1], x, y \in C^{-\alpha}$ :

$$
\left|P_{t} \varphi(x)-P_{t} \varphi(y)\right| \leq c\left(t^{-(1+s+2 \alpha p) / 2}+1\right)|x-y|_{C^{-s}}\|\varphi\|_{0}\left(1+|x|_{C^{-\alpha}}+|y|_{C^{-\alpha}}\right)^{\gamma} .
$$

## 4. The Kolmogorov operator

Given $N \in \mathbb{N}$, we denote by $\mathcal{F}_{N} C_{b}^{\infty}$ the set of all functions $\varphi: H \rightarrow \mathbb{R}$ of the form

$$
\begin{equation*}
\varphi(x)=g\left(x_{j},|j| \leq N\right), \tag{4.1}
\end{equation*}
$$

where $g: \mathbb{R}^{(2 N+1)^{2}} \rightarrow \mathbb{R},\left(\xi_{j},|j| \leq N\right) \rightarrow g\left(\xi_{j},|j| \leq N\right)$ is of class $C_{b}^{\infty}$. Moreover, we set

$$
\mathcal{F} C_{b}^{\infty}=\bigcap_{N=1}^{\infty} \mathcal{F}_{N} C_{b}^{\infty} .
$$

Note that if $\varphi \in \mathcal{F}_{N} C_{b}^{\infty}$ is of the form (4.1) we have

$$
\begin{equation*}
D \varphi(x)=\sum_{|k| \leq N} D_{k} g\left(\xi_{j},|j| \leq N\right) e_{k} \tag{4.2}
\end{equation*}
$$

and

$$
\begin{equation*}
D^{2} \varphi(x)=\sum_{|k| \leq N,|h| \leq N} D_{h} D_{k} g\left(\xi_{j},|j| \leq N\right) e_{h} \otimes e_{k} . \tag{4.3}
\end{equation*}
$$

Let us introduce the Kolmogorov operator on $\mathcal{F} C_{b}^{\infty}$, setting

$$
\begin{equation*}
\mathcal{K} \varphi:=\frac{1}{2} \operatorname{Tr}\left[D^{2} \varphi\right]+\left\langle A x-: x^{3}:, D \varphi\right\rangle, \quad \varphi \in \mathcal{F} C_{b}^{\infty} . \tag{4.4}
\end{equation*}
$$

This definition is meaningful because if $\varphi$ is given by (4.2) we have

$$
\left\langle: x^{3}:, D \varphi\right\rangle=\sum_{|k| \leq N}\left\langle D_{k} g\left(\xi_{j},|j| \leq N\right) e_{k},: x^{3}:\right\rangle
$$

We set

$$
\rho(x)=Z^{-1} e^{-2 U(x)},
$$

so that

$$
\begin{equation*}
D \log \rho(x)=-2 D U(x)=-2: x^{3}: \tag{4.5}
\end{equation*}
$$

We also consider the approximate measure $\nu_{N}$ in $H$

$$
\begin{equation*}
\nu_{N}(d x)=Z_{N}^{-1} e^{-U_{N}(x)} \mu_{N}(d x) \tag{4.6}
\end{equation*}
$$

where

$$
\begin{equation*}
U_{N}(x)=\frac{1}{4} \int_{H}: x_{N}^{4}:(\xi) d \xi \tag{4.7}
\end{equation*}
$$

and

$$
Z_{N}:=\int_{H} e^{-U_{N}(x)} \mu_{N}(d x)
$$

Note that

$$
\begin{equation*}
\mathcal{K}\left(\varphi^{2}\right)=2 \varphi \mathcal{K} \varphi+|D \varphi|^{2}, \quad \forall \varphi \in \mathcal{F}_{N} C_{b}^{\infty} . \tag{4.8}
\end{equation*}
$$

Integrating with respect to $\nu$ over $\mathcal{H}$, yields

$$
\begin{equation*}
\int_{\mathcal{H}} \mathcal{K} \varphi \varphi d \nu=-\frac{1}{2} \int_{\mathcal{H}}|D \varphi|^{2} d \nu, \quad \forall \varphi \in \mathcal{F} C_{b}^{\infty} . \tag{4.9}
\end{equation*}
$$

We also introduce an approximating operator $\mathcal{K}_{N}$ on $\mathcal{F} C_{b}^{\infty}$ setting, for $\varphi(x)=g\left(x_{j},|j| \leq M\right)$

$$
\begin{align*}
& \mathcal{K}_{N} \varphi:=\frac{1}{2} \operatorname{Tr}\left[\Pi_{N} D^{2} \varphi\right]+\left\langle A x-: x_{N}^{3}:, D \varphi\right\rangle \\
& =\frac{1}{2} \sum_{|h| \leq N \wedge M} D_{h}^{2} g\left(\xi_{j},|j| \leq M\right)-\sum_{|h| \leq M}\left(\alpha_{h}-\rho_{N}\right) x_{h} D_{h} g\left(\xi_{j},|j| \leq M\right)  \tag{4.10}\\
& -\sum_{|h| \leq M} \sum_{\left|j_{1}\right|,\left|j_{2}\right|,\left|j_{3}\right| \leq M} x_{j_{1}} x_{j_{2}} x_{j_{3}} \delta_{h, j_{1}+j_{2}+j_{3}} D_{h} g\left(\xi_{j},|j| \leq M\right),
\end{align*}
$$

where $\alpha_{h}=\left(1+|h|^{2}\right)$.

Similarly as above, we have

$$
\begin{equation*}
\int_{\mathcal{H}} \mathcal{K}_{N} \varphi \varphi d \nu_{N}=-\frac{1}{2} \int_{\mathscr{H}}|D \varphi|^{2} d \nu_{N}, \quad \forall \varphi \in \mathcal{F}_{N} C_{b}^{\infty} . \tag{4.11}
\end{equation*}
$$

4.1. $m$-dissipativity of $\mathcal{K}$. Fix $N_{0} \in \mathbb{N}$ and let $f \in \mathcal{F}_{N_{0}} C_{b}^{\infty}$. For $N \in \mathbb{N}$ we consider the solution $\varphi_{N}$ of the approximating equation

$$
\begin{equation*}
\lambda \varphi_{N}-\frac{1}{2} \operatorname{Tr}\left[D^{2} \Pi_{N} \varphi_{N}\right]+\left\langle A x_{N}-: x_{N}^{3}:, D \varphi_{N}\right\rangle=f . \tag{4.12}
\end{equation*}
$$

Then, taking into account (4.11) we find the estimate

$$
\begin{equation*}
\int_{H}\left|D \varphi_{N}\right|^{2} d \nu_{N} \leq \frac{2}{\lambda} \int_{H} f^{2} d \nu_{N} . \tag{4.13}
\end{equation*}
$$

Theorem 4.1. $\mathcal{K}$ is extendible to an $m$-dissipative operator in $L^{1}(\mathcal{H}, \nu)$ whose extension we denote by $\mathcal{K}^{(1)}$. Moreover, $\mathcal{F} C_{b}^{\infty}$ is a core for $\mathcal{K}^{(1)}$.

Proof. We have only to show

$$
\begin{equation*}
\lim _{N \rightarrow \infty} \int_{\mathscr{H}}\left\langle: x_{N}^{3}:-: x^{3}:, D \varphi_{N}\right\rangle d \nu=0 \tag{4.14}
\end{equation*}
$$

and then to apply the Lumer-Phillips theorem in some Besov spaces.
We take $\alpha, s$ and $p$ satisfying the assumptions of Theorem 3.9 and write

$$
\left.\left|D P_{t}^{N} f(x) \cdot h\right| \leq c t^{-(1+s+\alpha p) / 2}+1\right)|h|_{C^{-s}}\|f\|_{0}\left(1+|x|_{C^{-\alpha}}\right)^{\gamma}, \quad t \geq 0
$$

for some $\alpha<\frac{1}{8}, \epsilon>0$ and $c, \gamma$ depending on $s, \alpha$, $\epsilon$. Moreover, we can choose $1+\alpha+\epsilon+s<2$.
We write

$$
\varphi_{N}(x)=\int_{0}^{\infty} e^{-\lambda t} P_{t}^{N} f(x) d s
$$

and deduce

$$
\left|D \varphi_{N}(x) \cdot h\right| \leq c|h|_{C^{-s}}\|f\|_{0}\left(1+|x|_{C^{-\alpha}}\right)^{\gamma} .
$$

It then suffices to write

$$
\left|\int_{\mathcal{H}}\left\langle: x_{N}^{3}:-: x^{3}:\right\rangle D \varphi_{N}\right| d \nu \leq c \int_{\mathcal{H}}\left|: x_{N}^{3}:-: x^{3}:\right|_{C^{-s}}\|f\|_{0}\left(1+|x|_{C^{-\alpha}}\right)^{\gamma} d \nu
$$

By the Hölder inequality this converges to 0 .

Remark 4.2. Arguing as in [DaDe07] one can show that the closure $\mathcal{K}^{(p)}$ of $\mathcal{K}$ in $L^{p}(\mathcal{H}, \nu), p \geq 1$ is $m$-dissipative. Moreover the gradient operator

$$
D: \mathcal{F} C_{b}^{1} \subset L^{p}(\mathcal{H}, \nu) \rightarrow L^{p}(\mathcal{H}, \nu ; H)
$$

is closable. We denote by $\bar{D}_{p}$ its closure and by $W^{1, p}(\mathcal{H}, \nu)$ its domain.
Finally,

$$
D\left(\mathcal{K}^{(p)}\right) \subset W^{1,2}(\mathcal{H}, \nu)
$$

and it results

$$
\begin{equation*}
\int_{\mathcal{H}} \mathcal{K}^{(2)} \varphi \varphi d \nu=-\frac{1}{2} \int_{\mathscr{H}}|D \varphi|^{2} d \nu, \quad \forall \varphi \in D\left(\mathcal{K}^{(2)}\right) . \tag{4.15}
\end{equation*}
$$

Moreover since, as easily checked

$$
\int_{\mathscr{H}} \mathcal{K}^{(2)} \varphi \psi d \nu=-\frac{1}{2} \int_{\mathscr{H}}\langle D \varphi, D \psi\rangle d \nu, \quad \forall \varphi, \psi \in \mathcal{F} C_{b}^{2}(H),
$$

it results

$$
\begin{equation*}
\int_{\mathcal{H}} \mathcal{K}^{(2)} \varphi \psi d \nu=-\frac{1}{2} \int_{\mathcal{H}}\langle D \varphi, D \psi\rangle d \nu, \quad \forall \varphi, \psi \in D\left(\mathcal{K}^{(2)}\right) \tag{4.16}
\end{equation*}
$$

so that $\mathcal{K}^{(2)}$ is symmetric. Finally, let us recall the classical integration by parts formula for the Gaussian measure $\mu$

$$
\begin{equation*}
\int_{\mathcal{H}}\langle D \varphi, z\rangle \psi d \mu=-\int_{\mathcal{H}}\langle D \psi, z\rangle \varphi d \mu+\int_{\mathcal{H}} W_{z} \psi \varphi d \mu, \quad h \in \mathbb{Z}^{2}, \psi, \varphi \in \mathcal{F} C_{b}^{1}(\mathcal{H}) \tag{4.17}
\end{equation*}
$$

where

$$
W_{z}(x)=\left\langle Q^{-1 / 2} x, z\right\rangle=\sum_{k \in \mathbb{Z}^{2}} \sqrt{1+|k|^{2}} z_{k} x_{k}, \quad \phi \in \mathcal{H}
$$

and $Q$ is the covariance of $\mu$. Since $\nu(d x)=\rho(x) \mu(d x)$, setting $\psi=\rho$ in (4.17) yields

$$
\int_{\mathcal{H}}\langle D \varphi, z\rangle \rho d \mu=-\int_{\mathcal{H}}\langle D \rho, z\rangle \varphi d \mu+\int_{\mathcal{H}} W_{z} \rho \varphi d \mu, \quad h \in \mathbb{Z}^{2}, \varphi \in \mathcal{F} C_{b}^{1}(\mathcal{H}) .
$$

which is equivalent to

$$
\int_{\mathcal{H}}\langle D \varphi, z\rangle d \nu=-\int_{\mathcal{H}}\langle D \log \rho, z\rangle \varphi d \nu+\int_{\mathcal{H}} W_{z} \varphi d \nu, \quad h \in \mathbb{Z}^{2}, \varphi \in \mathcal{F} C_{b}^{1}(\mathcal{H}) .
$$

Since $\langle D \log \rho, z\rangle=-\left\langle: x^{3}:, z\right\rangle$ we have

$$
\int_{\mathcal{H}}\langle D \varphi, z\rangle d \nu=\int_{\mathcal{H}}\left\langle: x^{3}:, z\right\rangle \varphi d \nu+\int_{\mathcal{H}} W_{z} \varphi d \nu, \quad h \in \mathbb{Z}^{2}, \varphi \in \mathcal{F} C_{b}^{1}(\mathcal{H})
$$

Therefore, for any $p>1$ and any $z \in \mathbb{Z}^{2}$ there is $C_{p}>0$ such that

$$
\begin{equation*}
\left|\int_{\mathcal{H}}\langle D \varphi, z\rangle d \nu\right| \leq C_{p}\|\varphi\|_{L^{p}(\mathcal{H}, \nu)}|z| . \tag{4.18}
\end{equation*}
$$

This inequality is useful to develop a geometrical theory for the measure $\nu$, see $[\mathrm{DaDe} 17],[\mathrm{DaDe} 16]$ and [DaLuTu18].

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