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# A physically based method for soil evaporation estimation by revisiting the soil drying process

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#### Abstract 15

16 While numerous models exist for soil evaporation estimation, they are more or less empirically based either in the model structure or in the determination of introduced 17 parameters. The main difficulty lies in representing the water stress factor, which is 18 19 usually thought to be limited by capillarity-supported water supply or by vapor 20 diffusion flux. Recent progress in understanding soil hydraulic properties, however, 21 have found that the film flow, which is often neglected, is the dominant process under 22 low moisture conditions. By including the impact of film flow, a reexamination on the typical evaporation process found that this usually neglected film flow might be the 23 24 dominant process for supporting the stage-II evaporation (i.e. the fast falling rate stage), besides the generally accepted capillary-flow-supported stage- I evaporation 25 26 and the vapor-diffusion-controlled stage-III evaporation. A physically based model for

2

estimating the evaporation rate was then developed by parameterizing the 27 Buckingham-Darcy's law. Interestingly, the empirical Bucket model was found to be a 28 specific form of the proposed model. The proposed model requires the in-equilibrium 29 relative humidity as the sole input for representing water stress and introduces no 30 31 adjustable parameter in relation to soil texture. The impact of vapor diffusion was also 32 discussed. Model testing with laboratory data yielded an excellent agreement with observations for both thin soil and thick soil column evaporation experiments. Model 33 34 evaluation at 15 field sites generally showed a close agreement with observations, with a great improvement in the lower range of evaporation rates in comparison with 35 the widely applied PT-JPL model. 36

# 37 **1. Introduction**

38 Soil evaporation plays an important role in the mass and energy exchange between the 39 land surface and the atmosphere. Together with plant transpiration, they return about 40 60% of precipitation into the air (Oki & Kanae, 2006) and consume nearly 25% of the incoming solar radiation globally (Trenberth et al., 2009). Estimating soil evaporation 41 42 accurately and separating its contribution from the plant transpiration are crucial to understand the water and energy cycles (Seneviratne et al., 2010; Wang & Dickinson, 43 2012; Maxwell & Condon, 2016; Gu et al., 2018) and to quantify the carbon cycle 44 process that is highly related to plant behavior (Sutanto et al., 2012; Schlesinger & 45 46 Jasechko, 2014; Fisher et al., 2017). Physically, soil evaporation is the transition of 47 soil water from the liquid phase to the vapor phase and the escape of water vapor to the above atmosphere. The phase transition requires an energy supply while the vapor 48 escape is mainly a molecule diffusion process (Haghighi et al., 2013). By focusing 49 50 mainly on energy supply or on vapor transport process, the existing evaporation estimation methods can generally be summarized into two series, the energy budget 51 one and the mass transfer one (Brutsaert, 2005). 52

53 The energy budget methods vary in different forms, among them, the classic one 54 relies on the concepts of atmospheric demand and water supply. The atmospheric 55 demand is represented by the potential evaporation rate, often estimated by the Penman (1948) equation or by the Priestley & Taylor (1972) equation. The water 56 supply, also known as the water stress factor, is usually expressed empirically as a 57 linear function of soil water content after the identification of a "critical water content" 58 59 below which soil water supply is limited (Seneviratne et al., 2010). This kind of method, also termed as the Bucket model (Budyko, 1974; Brutsaert, 2005), is 60 probably the oldest method of evaporation estimation. Due to its simplicity (which 61 62 may vary in different forms), it is still extensively applied, in particular within remote sensing-based evaporation models (e.g., Fisher et al., 2008; Miralles et al., 2011; 63 Martens et al., 2017). 64

The mass transfer- or resistance-based methods consider the vapor transfer process in 65 a direct way. The evaporation process is described as vapor diffusion from either the 66 soil surface or a depth below. The former is the so-called  $\alpha$  formulation (e.g., Barton, 67 1979; Noilhan & Planton, 1989). An aerodynamic resistance term is introduced to 68 represent the vapor transfer efficiency from soil surface to the atmosphere. When the 69 70 evaporation process is considered from a plane below the soil surface, where vapor pressure is saturated, an additional soil surface resistance term is included, leading to 71 the so-called  $\beta$  formulation (e.g., Deardorff, 1977; Dorman & Sellers, 1989). 72

The Bucket model and the resistance-based methods, although widely applied, are 73 74 somehow empirically based, especially in determining the water stress and the resistance factors. The linear relationship in the Bucket model is empirical (Brutsaert, 75 2005; Seneviratne et al., 2010) and might not be unique even in the same location 76 (Haghighi et al., 2018). The resistance methods, although physically based, rely on 77 empirical estimation of the resistance terms (e.g., Mahfouf & Noilhan, 1991; Merlin 78 79 et al., 2016, 2018). The ambiguities in parameter estimation methods, as well as the difficulty in accurate data acquisition (e.g., the surface water content and the soil 80 hydraulic properties), therefore, would introduce high uncertainty in evaporation rate 81 82 estimation.

Alternatively, the soil evaporation can be estimated from the perspective of soil water
transport ability. Instead of being represented as a stress factor or a resistance term

85 (usually expressed empirically) as in the former methods, the soil water supply can be calculated directly and the actual evaporation rate is regarded as the minimum term 86 between the soil water flux and the potential evaporation rate. This method is termed 87 as the threshold formulation (e.g., Mahrt & Pan, 1984; Dickinson, 1984). With the soil 88 water flux described physically by the Buckingham-Darcy's law, which can be 89 calculated by solving the Richards equation (e.g., Sutanto et al., 2012) or by 90 parameterizing over a thin soil surface layer (e.g., Mahrt & Pan, 1984), this method 91 92 provides a much more solid basis for evaporation estimation. In the literature, for a drying process, capillary flow is regarded as the main soil water flow form and vapor 93 diffusion is included when soil becomes very dry (Philip & de Vries, 1957; Saito et al., 94 2006; Lehmann et al., 2008; Or et al., 2013; Or & Lehmann, 2019). 95

However, such capillary flow- and vapor diffusion-supported evaporation process 96 may be problematic for not considering the potential impact of film flow. This thin 97 film form water, held by the adsorption forces on the soil particle surface, is usually 98 treated as unremovable and hence be unimportant in water flow process (e.g., Philip 99 100 & de Vries, 1957; Idso et al., 1974). Recent progress in soil hydraulic properties 101 modeling however has found that this film flow, is in fact the predominant water flow form under low moisture conditions (e.g., Tuller & Or, 2001; Wang et al., 2013, 2018). 102 By including this film flow, the newly proposed soil hydraulic models greatly 103 104 improved the performance under low moisture conditions in comparison with capillary-based models (Tuller & Or, 2001; Lebeau & Konrad, 2010; Peters, 2013; 105 Wang et al., 2016, 2017, 2018). The re-recognition of the importance of this thin 106 film, therefore, requires us to reconsider the soil evaporation process, which is 107 108 commonly assumed to be supported only by capillary flow and vapor diffusion. 109 Although early research in the engineering area has taken into account of the impact of the so-called bound water (e.g., Chen & Pei, 1989), the analysis was problematic 110 due to the unclear definition of the hydraulic conductivity properties in relation to the 111 112 bound water.

Hence, there are three main objectives in this study: 1) to reexamine the typical evaporation process by including the impact of film flow, besides the commonly

- recognized capillary flow and vapor diffusion; 2) to develop a theoretical evaporation
- estimation model based on the detailed analysis of the soil drying process; and 3) to
- 117 evaluate the model performance with both laboratory and field observations.

# 118 **2. Theoretical Development**

# 119 2.1. The Soil Drying Process Revisited

120 [Figure 1 near here]



121 122

As shown in Figure 1 (a), the typical evaporation process from an initially saturated 123 124 soil column can be generally distinguished into the energy-limited stage and the moisture-limited stage, respectively (e.g., Lehmann et al., 2008; Shokri et al., 2011; 125 Or et al., 2013). The energy-limited stage is also termed as stage- I evaporation. 126 When the atmospheric demand keeps constant, this stage would have a constant 127 evaporation rate (assuming the impact of capillary limitations and nonlinear boundary 128 layer interactions are not important), so it is also known as the constant-rate period 129 130 (e.g., Yiotis et al., 2007). The moisture-limited stage can be subdivided into two

131 stages: the stage-II or the transition stage marked by a fast falling evaporation rate 132 and the stage-III where evaporation rate keeps low and changes smoothly (e.g., Idso et 133 al., 1974; Merz et al., 2015).

Correspondingly, the soil water flow supporting the drying process can also be 134 divided into three forms, including the commonly recognized capillary flow in full 135 pores and in corners and the vapor diffusion in void pores, as well as the usually 136 neglected thin film flow on soil particle surfaces (Tuller & Or, 2001; Peters, 2013; 137 138 Wang et al., 2018). Notably, this thin film flow held by adsorption forces is different from the so-called "thick film flow" presented in the literature (e.g., Yiotis et al., 2003, 139 2007, 2012; Lehmann et al., 2008), where it represents the water flow supported 140 mostly by capillarity and controls the water transport between the saturated zone and 141 142 the evaporation surface. In these referred work, this thin film flow was assumed to be unimportant and then be neglected (Yiotis et al., 2007). 143

# 144 **2.1.1. The Stage- I Evaporation**

For the drying process, the stage- I evaporation is extensively studied (e.g., Lehmann 145 146 et al., 2008, 2018; Shahraeeni et al., 2012; Shokri et al., 2008; Or et al., 2013 and references therein). In this period, an efficient water supply between the drying front 147 and the soil surface is maintained by capillary flow. The drying process is 148 accompanied by a decrease in soil surface moisture and an increase in drying front 149 150 depth. The end of this period should depend on both the capillary flow ability and the atmospheric demand. However, because the capillarity-supported water supply is 151 usually higher than the atmospheric demand even under low water saturation 152 conditions (Shahraeeni et al., 2012), the stage- I evaporation generally ends when a 153 154 critical surface water content or a characteristic depth is reached, where the capillary 155 water potential gradient between the drying front and the soil surface cannot overcome the gravitational forces and viscous dissipation (Lehmann et al., 2008). This 156 157 critical water content then can be seen as air begins to invade the finest pores at soil 158 surface and it is roughly corresponding to the so-called residual water content (Or et 159 al., 2013). The duration of the stage- I evaporation is therefore determined by soil pore size distribution properties. In this period, the evaporation rate is generally equal 160

to the potential evaporation rate. However, it should be noted that for very fine
textured soil and for high atmospheric demand, a higher critical water content and a
lower evaporation rate are expected due to the resistance induced by viscous effects
(Haghighi et al., 2013; Lehmann et al., 2018; Or & Lehmann, 2019) and/or due to the
vapor exchange limitations across the air boundary layer (Shahraeeni et al., 2012;
Haghighi & Or, 2013). In the present study, these effects are not included.

# 167 **2.1.2. The Stage- I Evaporation**

168 While extensive research on stage- I evaporation exists in the literature, little work has been done on the stage-II evaporation (Shokri et al., 2011). In the classic soil 169 170 evaporation theory, when the capillary flow is interrupted at the end of stage-I evaporation, the vaporization plane begins to recede into the inside soil and the vapor 171 172 diffusion becomes be dominant at the soil surface (Philip & de Vries, 1957; Saito et al., 2006; Lehmann et al., 2008; Or et al., 2013). Notably at this point, liquid water 173 still exists at soil surface, mainly in the form of thin film adsorbed by soil particle 174 (Tuller & Or, 2011; Wang et al., 2018). In the capillarity-based theory (e.g., van 175 176 Genuchten, 1980; Corey & Brooks, 1999), this left liquid water is represented by the so-called residual water content and means the low limitation of the free soil water. 177 Accordingly, it has nearly no contribution to the evaporation process. 178

This left and unremovable soil surface water content, however, is inconsistent with 179 180 recent laboratory observations. . . For example, the magnetic resonance imaging of a sand column presented by Merz et al. (2015) demonstrated clearly a gradual 181 decreasing soil surface moisture during the stage- II evaporation (see also Figure 1 b). 182 The vapor-diffusion-dominated assumption is also inconsistent with laboratory 183 observations. According to the Fick's law, a vapor-diffusion-controlled stage-II 184 185 evaporation requires a gradually increased diffusion layer thickness to explain the fast falling evaporation rate (Shokri et al., 2009, 2011). The dye experiment by Shokri et 186 al. (2011) however observed an abrupt jump of the vaporization plane at the end of 187 the stage- I evaporation (the so-called transition period in Shokri et al., 2011). This 188 189 abrupt jump indicated that the vaporization plane should remain in the soil surface during the stage-II evaporation. In other words, the water transport near the soil 190

191 surface is still dominated by liquid flow otherwise a gradually receding vaporization 192 plane would be observed. The gradually decreased soil surface moisture and the dominant liquid flow observed in stage- I evaporation happens to be consistent with 193 recent progress in soil hydraulic properties modeling. Commonly used soil hydraulic 194 models conceptualize pore space as a bundle of cylindrical capillaries (e.g., van 195 Genuchten, 1980; Mualem, 1976), neglecting the flow in liquid films held by 196 adsorption forces. This kind of capillary models often underestimate the hydraulic 197 198 conductivities under dry conditions (e.g., Tuller & Or, 2001; Wang et al., 2013, 2018). 199 a lot of research confirms that the liquid film, also explained as in Recently. corresponding to the residual water content in capillary model (Corey & Brooks, 200 1999), is actually flowable (Tuller & Or, 2001; Tokunaga, 2009; Wang et al., 2013; 201 202 Peters, 2013). By including both capillary flow and film flow, recently proposed models significantly improve the model performance under dry conditions (Lebeau & 203 Konrad, 2010; Peters, 2013; Tuller & Or, 2001; Tuller et al., 1999; Wang et al., 2016, 204 2017, 2018). An illustration in Figure 1 (c) shows clearly that this neglected film 205 206 conductivity is dominating after the soil moisture reaches the so-called residual water content while the vapor diffusion is only important under extremely dry conditions. 207

At pore scale, when air invades the finest pore at soil surface -marking the end of 208 stage- I evaporation-, liquid flow still exists in film form held by the adsorption force 209 210 as well as in corner form retained by capillarity. This is consistent with the following statement made by Scherer (1990) for the fast falling rate evaporation period: "The 211 liquid in the pores near the surface remains in the funicular condition, so there are 212 contiguous pathways along which flow can occur" (see Figure 4 in Scherer, 1990). 213 Tuller & Or (2001) showed that the hydraulic conductivity resulting from corner flow 214 215 was generally negligible in comparison with film flow under dry conditions. Therefore, the usually neglected film flow, rather than vapor diffusion, might be the 216 supporting mechanism that limits the water supply during stage- I evaporation. The 217 vapor diffusion also contributes to the evaporation process, however, in a magnitude 218 219 generally less than liquid film flow.

220 Notably this film dominant zone is restricted in a depth of several millimeters, below

221 which water flow is still supported by capillarity. During this stage, the evaporation process is also accompanied by a decrease in soil surface moisture while the drying 222 front depth keeps almost constant due to the discontinuity of capillary flow (Lehmann 223 et al., 2008;). Because the film flow flux is controlled by the soil specific surface area 224 and the film thickness that depends on matric potential (Bird, 1960; Tokunaga, 2009), 225 when considering film flow limitations solely, the soils with finer texture should 226 generally have a higher evaporation rate (at the same matric potential) and a longer 227 228 decreasing period during stage- I evaporation.

# 229 2.1.3. The Stage-III Evaporation

The stage-III evaporation is therefore the period when vapor diffusion actually 230 controls (Shokri et al., 2011). As shown in Figure 1 (c), the vapor conductivity only 231 232 exceeds the film conductivity when soil water content is extremely low. In this stage, the thin soil surface layer is (almost) completely dried (Figure 1 a and 1 b), with some 233 tightly-bounded thin liquid films left. The soil surface water content is equal to the 234 air-dry value. The evaporation process is accompanied by the receding of vaporization 235 236 plane (also known as the secondary drying front) into deeper soil. The vapor diffusion-controlled evaporation rate is very low (Figure 1 a) and depends on the 237 diffusion length between vaporization plane and soil surface (Shorkri et al., 2011). 238

### 239 **2.1.4. The Complete Evaporation Process**

By including film flow, the evaporation process can then be summarized into three 240 typical stages as shown in Figure 1 (d), with however some different explanations 241 with the classic theory: (1) The stage- I evaporation is supported by capillary flow 242 from drying front to soil surface. This stage ends when the driving capillary water 243 244 potential difference between drying front and soil surface cannot overcome the 245 gravitational forces and viscous dissipation, marked by a critical surface water content or a characteristic drying depth (Lehmann et al., 2008). The duration of this stage 246 depends on the width of pore size distribution and is also impacted by the nonlinear 247 boundary layer interactions (Shahraeeni et al., 2012; Haghighi & Or, 2013) as well as 248 249 the soil texture-dependent capillary flow limitations (Haghighi et al., 2013; Lehmann et al., 2018). When neglecting capillary flow limitations, the evaporation rate is 250

mainly controlled by the atmospheric demand. (2) The stage-II evaporation is also 251 252 supported by liquid flow, however, in the form of liquid film within a very thin soil surface layer. The capillarity-driven water can no longer achieve the soil surface 253 directly. This stage ends when vapor conductivity becomes dominant over film 254 conductivity, marked with the jump of vaporization plane from surface to deeper soil. 255 The duration of this stage depends mainly on the specific surface area of porous 256 media. The evaporation rate is limited by the soil water transport ability mainly in 257 258 film form. (3) The stage-III evaporation is the vapor diffusion stage. An almost 259 completely dried layer is developed at the soil surface and it grows deeper and deeper when the drying process keeps going. The evaporation rate depends on the depth of 260 the vaporization plane, which is also known as the secondary drying front. 261

At pore scale (Figure 1 d), the end of stage- I evaporation occurs when the air begins to invade the finest pores at the soil surface. During the stage- II evaporation, almost no saturated pores exist at the soil surface, liquid water is mainly retained in the form of film and can move along the surface of porous media. In the stage-III evaporation, almost no liquid water exists in pores and on the soil particle surface, except some extremely thin films tightly bounded by adsorption force (depends on the air humidity).

# 269 2.2. Model Development

## 270 **2.2.1.** Accounting for Film Flow in the Evaporation Estimation

Under most field atmospheric conditions, the evaporation rate is approximately equal
to potential evaporation during stage- I period. In stage-II period, the evaporation rate
is very low. Therefore, the key is to define the evaporation rate in stage- II evaporation.
because,

During stage- II period, the evaporation rate is limited by the soil water transport ability, which can be expressed by the Buckingham-Darcy's law, written as

277 
$$LE = -K\left(\theta\right) \frac{dh}{d\theta} \frac{d\theta}{dz}\Big|_{z=0}$$
(1)

where *LE* (L T<sup>-1</sup>) is the actual evaporation rate, and *K* (L T<sup>-1</sup>) and *h* (L) are the hydraulic conductivity and water potential at the surface water content  $\theta$ , respectively. When ignoring capillary flow limitations and nonlinear boundary layer interactions, the evaporation rate in stage- I period is equal to the potential evaporation rate,  $LE_p$ . Therefore, at the onset of stage- II evaporation, Eq. (1) becomes

283 
$$LE_{c} = -K\left(\theta_{c}\right) \frac{dh}{d\theta} \frac{d\theta}{dz}\Big|_{z=0,\theta=\theta_{c}} = LE_{p}$$
(2)

with  $\theta_c$  being the critical surface water content marking the discontinuity of capillary 284 flow. It is roughly close to the so-called residual water content (Lehmann et al., 2008; 285 286 Or et al., 2013). It should be noted that when capillary flow limitations in relation to fine-textured soils (Haghighi et al., 2013) and/or vapor diffusion through above thin 287 boundary layer in relation to high atmospheric demand (Shahraeeni et al., 2012) 288 become be important in stage- I period, no constant evaporation rate stage would be 289 290 observed and the  $LE_c$  might have a value less than  $LE_p$  (Lehmann et al., 2018; Or & Lehmann, 2019). 291

To solve Eqs. (1) and (2), information on the soil hydraulic properties are required. In 292 the literature, the capillary based models, the van Genuchten (1980)-Mualem (1976) 293 294 model for instance, are applied (e.g., Saito et al., 2006). This is questionable since the film flow, as discussed previously, is thought to be dominant during stage-II 295 evaporation. According to Campbell & Shiozawa (1992), a linear relationship exists 296 between water content and log-scale water potential under dry conditions, so that the 297 soil water retention curve that accounts for film flow can be expressed as (Wang et al., 298 299 2016)

$$S_f = \frac{\theta}{\theta_c} = 1 - \frac{\ln(h/h_c)}{\ln(h_0/h_c)}$$
(3)

with  $S_f$  being the saturation degree that accounts for film water, and  $h_c$  and  $h_0$  the water potential at  $\theta_c$  and zero, respectively. As suggested by Schneider & Goss (2012),  $h_0$  can be approximately set to  $6.3 \times 10^4$  m for soils with different texture properties.

The hydraulic conductivity that accounts for film flow depends on the soil specific surface area *SA* and the film thickness f (Bird, 1960), in the form of

306 
$$\frac{K(\theta)}{K(\theta_c)} = \frac{SA \times f^3}{SA_c \times f_c^3}$$
(4)

307 In general, the specific surface area is kept as constant in film conductivity models (Tuller & Or, 2001; Lebeau & Konrad, 2010; Wang et al., 2017). However, this 308 assumption is not appropriate when film thickness is thick than the width of parallel 309 plates or when the water film does not cover the entire soil surface under extremely 310 dry conditions. When considering the sole contribution of film flow, an effective 311 surface area can be applied. It was expressed as the water content to film thickness 312 ratio, that is,  $SA = \theta / f$ . Since the soil surface water content is very low under stage- II 313 314 evaporation, only van der Waals forces were considered. The film thickness formula proposed in Iwamatsu & Horii (1996) was applied, written as 315

316 
$$\frac{f}{f_c} = \left(\frac{h}{h_c}\right)^{-1/3}$$
(5)

Substituting Eq. (5) and  $SA = \theta / f$  into Eq. (4) gives

318 
$$K(\theta) = K(\theta_c) S_f \left(\frac{h}{h_c}\right)^{-2/3}$$
(6)

Notably the impact of modified viscosity on film conductivity under very thin film
thickness conditions (Tuller & Or, 2001; Lebeau & Konrad, 2010) was not considered
here.

The water content gradient  $d\theta/dz$  at soil surface is also required in Eqs. (1) and (2). 322 The stage-II evaporation can be seen as a secondary drying process of a very thin soil 323 layer (with several millimeters in length, see Figure 1 b) with uniform initial water 324 content  $\theta_c$ . Similar to the classic evaporation process as described in Gardner (1959) 325 and in Brutsaert (2014), the upper boundary can be set to a fixed water content,  $\theta_m$ , 326 which is actually the soil air-dry value. The value depends on the soil texture 327 328 properties and the environmental conditions. The lower boundary can be set at a fixed 329 depth defined as  $z_d$ , with the water content assigned as the surface water content, varying from  $\theta_c$  to  $\theta_m$ . When the water potential profile in the thin soil surface layer 330 shows high nonlinearity, the water content profile can be seen as linear under dry 331 conditions (Figure 1 b). Therefore, this gradient  $d\theta/dz$  is approximately equal to  $(\theta -$ 332  $\theta_m)/z_d$ . 333

Dividing Eq. (1) by Eq. (2) and with the substitution of  $d\theta/dz \approx (\theta - \theta_m)/z_d$  and Eqs. (3) and (6), one obtains

336 
$$\frac{LE}{LE_p} = S_f \left(\frac{h}{h_c}\right)^{1/3} \frac{\left(S_f - S_m\right)}{\left(1 - S_m\right)}$$
(7)

Eq. (7) provides a theoretical formula for scaling evaporation rate under stage- II evaporation, with the assumption that film flow dominates the soil surface water flow process. The normalized evaporation rate ranges from 1 at the beginning of stage- II evaporation to 0 when the air-dry soil water content is reached at the end of stage- II evaporation. Note that the impact of vapor diffusion is not considered in Eq. (7). This film-flow-based evaporation rate estimation method is simply termed as the E\_FILM model.

The E\_FILM model requires the input of soil surface water potential h, which is highly sensitive to water content and depends on soil texture properties. Alternatively, h can be calculated from the well-known Kelvin equation under dry conditions, assuming a thermodynamic equilibrium between liquid and vapor phase. Therefore, hcan be written as:

349 
$$h = \frac{RT}{Mg} \ln \left( RH \right) \tag{8}$$

where *R* is the universal gas constant (8.314 J mol<sup>-1</sup>K<sup>-1</sup>), *T* (K) is the absolute temperature, *M* is the molecular weight of water (0.018015 kg mol<sup>-1</sup>), *g* is the gravitational acceleration (9.81 m s<sup>-2</sup>) and *RH* is the in-equilibrium relative humidity at soil surface. An illustration in Figure 2 (a) g shows that the temperature effect on water potential estimation is not significant. With neglecting the temperature effect, Eq. (7), therefore, can be written in the form of *RH*, as

356 
$$\frac{LE}{LE_{p}} = \left[\frac{\ln(RH)}{\ln(RH_{c})}\right]^{1/3} \frac{\ln\left[\ln(RH_{0})/\ln(RH)\right]}{\ln\left[\ln(RH_{0})/\ln(RH_{c})\right]} \frac{\ln\left[\ln(RH_{m})/\ln(RH)\right]}{\ln\left[\ln(RH_{m})/\ln(RH_{c})\right]}$$
(9)

where  $RH_o$  is the in-equilibrium relative humidity at the water potential of  $h_o$ , being 1.04% at 20 °C, while  $RH_c$  and  $RH_m$  are in corresponding to the water potential  $h_c$  and  $h_m$ , respectively. 360 The merit of using *RH* is that when vapor pressure is in equilibrium with soil water potential, the RH keeps almost constant as 100% during stage- I evaporation since the 361 soil moisture supply is sufficient. It begins to decrease rapidly in stage- II evaporation 362 then change slowly in stage-III evaporation, reflecting a moisture limited 363 and condition. This changing is consistent with the evaporation rate dynamic. Hence, the 364 critical relative humidity  $RH_c$  and  $RH_m$  can be chosen as the maximum (except the 365 value of 1) and the minimum value of the observed RH, respectively. Tuller & Or 366 (2005) suggested an empirical threshold potential around  $-10^3$  m (corresponding to a 367 RH of 93% at 20 °C) where the capillary condensation becomes negligible. When 368 continues observation of RH is unavailable, this empirical value is suggested as the 369 critical water potential value for different soils. 370

371 Eq. (9) provides a scaling method for soil evaporation estimation, requiring *RH* as the sole input. As shown in Figure 2 (b), a nonlinear relationship exists between the 372 evaporation rate and the surface water content. It depends on the critical soil water 373 potential where stage-II evaporation begins. The more negative critical water 374 potential is expected for soils with finer texture under film flow limitations as 375 illustrated in Figure 2 (c). Notably here we didn't consider the possible impact of 376 capillary flow limitations. We emphasize here again that the E FILM model deals 377 with the water flux in a very thin soil surface layer, with only several millimeters in 378 379 depth. Therefore, much negative matric potential values are expected compared to those observed at much deeper layers in the literature. 380

381 [Figure 2 near here]



382

It is interesting to note that if the surface area *SA* was kept constant in Eq. (4), the
E\_FILM model would have the form of

$$\frac{LE}{LE_p} = \frac{S_f - S_m}{1 - S_m} \tag{10}$$

386 Eq. (10) yields a linear relationship between the scaled evaporation rate and the soil surface water content. It is in the same form as the widely applied but empirically 387 based Bucket model (Budyko, 1974). Hence, Eq. (10) provides a physical explanation 388 for the Bucket method. The slight difference between Eq. (10) and the Bucket model 389 is that the soil water content in Eq. (10) is observed in a depth of several millimeters, 390 391 while in the Bucket model, it generally represents a depth of several to tens of centimeters. However, the assumption that the water content shows a similar changing 392 trend at different depths (near the soil surface) seems to be appropriate in the drying 393 394 process.

The difference between Eqs. (7) and (10) comes from the different conductivity functions associated with film flow (Eq. 4). In deriving Eq. (7), only film-form water flow is considered, assuming no contribution comes from capillary flow when water content is less than the critical value  $\theta_c$ . It represents a very dry condition under which the soil water potential dynamic can usually be captured by changes in *RH* (Tuller & Or, 2005). When stage- I evaporation ends at a more positive critical water potential (in a magnitude cannot be captured by *RH* variation), water flow in the soil surface
might be supplied by both capillary flow (retained in very fine pores) and film flow
(along soil particle surface). Under this situation, Eq. (10) is preferred with the input
of soil surface water potential or water content. This is because the constant specific
surface area assumption in Eq. (4) generally yielded a close agreement with observed
conductivities in relatively high water potential range (Tuller & Or, 2001; Lebeau &
Konrad, 2010; Wang et al., 2017).

# 408 2.2.2. The Impact of Capillary Flow Limitations in Stage- I Evaporation

409 In the former section, it is assumed that the capillarity-driven water supply is sufficient to meet the atmospheric water demand in stage- I, therefore, LE is equal to 410  $LE_p$  for surface water content higher than  $\theta_c$  (Eq. 2). However, when the capillary 411 412 flow limitations become be important in relation to fine-textured soils (Haghighi et al., 2013; Lehmann et al., 2018; Or & Lehmann, 2019) and/or the vapor diffusion 413 limitations through the above thin boundary layer is significant under high 414 atmospheric demand (Shahraeeni et al., 2012), LE may also show a falling trend and 415 416 be less than  $LE_p$  during stage- I period.

Under this situation, the resistances come from capillary flow, film flow as well as 417 vapor diffusion should be taken into account together to describe a complete soil 418 evaporation process. Haghighi et al. (2013) and more recently, Lehmann et al. (2018) 419 420 provided a formula for describing the capillary limitations on evaporation rate. By estimating  $LE_c$  in Eq. (2) with this formula, the evaporation rate dynamic during both 421 stage- I and stage- I period may be described. Such combination method, however, 422 requires a detail parameterization of soil hydraulic properties and involves a lot of 423 424 uncertainty. For example, due to the high specific surface area, the film flow may be 425 important even in stage-I period for fine textured soils. This possible impact however was not taken into account in the existing theory that dealing with capillary 426 limitations (e.g., Haghighi et al., 2013; Lehmann et al., 2018). In the present study, 427 therefore, we focus solely on the limitations come from film flow in stage- I period 428 429 and ignore the possible impact of capillary flow limitations in stage- I period.

430 **2.2.3. The Impact of Vapor Flow** 

Eq. (7) considers film flow only and ignores the potential impact of vapor transport that is dominant during stage-III evaporation. Here, the isothermal vapor flow is included to evaluate its influence on the evaporation rate. The vapor diffusion contribution is calculated following the classical theory as provided in Philip & de Vries (1957) and in Saito et al. (2006). The detail is presented in Appendix A. It should be noted that only isothermal vapor diffusion is considered, so that the impact of temperature gradient is not included.

Figure 2 (c) shows that the contribution of vapor flow is, in general, much less than that of film flow. It is only important for soils with coarse texture and for very dry conditions. Considering vapor flow in soil evaporation estimation, however, will produce a very complex formula and requires additional information related to specific soil properties. Here, we suggest a simple equation to include the impact of vapor flow, written as:

$$LE = S_{f} \left(\frac{h}{h_{c}}\right)^{1/3} \frac{\left(S_{f} - S_{m}\right)}{\left(1 - S_{m}\right)} \left(LE_{p} - LE_{v}\right) + LE_{v}$$
(11)

where  $LE_{\nu}$  represents the contribution from vapor flow. Shokri et al. (2011) showed that the vapor diffusion flux was quite similar for soils with different texture, often in a range from 0.5 to 2.5 mm d<sup>-1</sup> at the onset of stage-III evaporation. The mean vapor diffusion rate is about 1.5 mm d<sup>-1</sup> for different soils. The accurate estimation of vapor diffusion rate, however, requires the length between the vaporization plane and the soil surface and requires also the atmospheric vapor pressure. Eq. (11) is only suggested when vapor flow is believed to be important.

452 **2.2.4. Evaporation Rate Estimation in Field** 

444

When it comes to field scale, the soil evaporation estimation becomes much more complicated. In contrast with the monotonically increase of the film dominated soil layer thickness as in laboratory experiment, the thin layer thickness under real field conditions would grow and reduce during daytime and nighttime due to the soil moisture redistribution (Brutsaert, 2014). As a result, a distinct diurnal pattern of evaporation rate would be observed in field (e.g., Jackson et al, 1976; Idso et al., 1974, 1979). However, when considering the evaporation process at the daily scale, this diurnal patter can be avoided (Brutsaert, 2014). The vapor transport can also be
neglected for its minor effect on the daily evaporation rate (e.g., Milly, 1984a, 1984b;
Saravanapavan & Salvucci, 2000).

Another difficulty with field data is that the relative humidity is often observed at the 463 height of 2 m above the soil surface, which might not be in equilibrium with soil 464 surface. However, the relative humidity at different heights generally shows a similar 465 trend during the evaporation process. Eq. (9) shows that it is the  $log(RH) / log(RH_c)$ 466 467 ratio that controls the value of LE. Therefore, it might be appropriate to assume that the ratio of log-scale relative humidity at soil surface can be represented by that 468 observed at the height of 2 m. However, unlike laboratory experiments, it is hard to 469 define the critical  $RH_c$  where the evaporation rate begins to decrease as the 470 471 (atmospheric) boundary condition is an open system. In this study, a maximum value of 0.85 derived by trial and error is suggested as the upper boundary for the critical 472  $RH_c$ . That is,  $RH_c$  is the minimum value between 0.85 and the observed maximum RH. 473  $RH_m$  is simply set to the minimum value of the observed RH. This requires the 474 475 observations cover a complete drying period. The influence of the soil heterogeneity here is simplified by assuming that the applied air relative humidity represents a 476 mixture of soil surface moisture conditions. 477

With these assumptions and by ignoring vapor flow, Eq. (9) can now be used to estimate the field soil evaporation rate, requiring only the meteorological data as input. In this study, the potential evaporation rate  $LE_p$  is calculated by the Priestley & Taylor (1972) equation, written as:

$$LE_{p} = \alpha_{PT} \frac{\Delta}{\Delta + \gamma} (R_{n} - G)$$
(12)

where  $\alpha_{PT}$  is the Priestley-Taylor constant, with the value of 1.26,  $\Delta$  (kPa K<sup>-1</sup>) the slope of the vapor pressure curve at air temperature,  $\gamma$  (kPa K<sup>-1</sup>) the psychrometric constant,  $R_n$  (W m<sup>-2</sup>) the net radiation and G (W m<sup>-2</sup>) the soil heat flux.

482

# 486 **3. Data Description**

#### 487 **3.1. Data from Laboratory Experiment**

The proposed E\_FILM model is evaluated at both laboratory and field scales. The laboratory data include nine thin (less than 1 mm in length) soil evaporation experiments provided in Wilson (1990) and one thick soil column (50 cm in length) evaporation experiment presented in Zhang et al. (2015).

#### 492 [Table 1 near here]

493 The thin soil evaporation test includes three soil types, Beaver Creek Sand, Custom 494 Silt and Regina Clay. For each soil type, three tests were presented (see Table 1). In 495 all tests except for tests silt\_1 and clay\_1, the thin soil samples were prepared by 496 gently dusting a layer of dry soil onto a sheet of aluminum foil. The sample was then 497 saturated with distilled water using a mist applicator. For tests silt\_1 and clay\_1, the soil sample was prepared as slurry and poured into the evaporation pan to achieve a 498 499 thicker soil layer. As described in Wilson (1990), some difficulties were encountered when using the slurried soil method. For example, non-uniform drying was found for 500 silt\_1 test and for clay\_1 test. Besides, shrinking and deformation was frequently 501 502 observed for clay\_1 test, which resulted in problems such as curling and irregular drying. The evaporation was preceded in room temperature and the relative humidity 503 504 of air was kept almost constant. For the details please refer to Wilson (1990).

The thick soil evaporation experiment was provided in Zhang et al. (2015). The soil column was 50 cm in length. Different from the thin soil evaporation, an infrared lamp was used for supplying heat to soil surface. The relative humidity near soil surface was recorded during the drying process. The experiment details can be seen in Zhang et al. (2015).

# 510 **3.2. Data Collected from Field Sites**

The field sites are chosen from the collection presented in Merlin et al. (2016), including 15 sites representing bare soil conditions across different countries (see table 2). Note that only sites with more than 40 days of observation are selected. These sites were selected mostly from the national and international flux station networks (OZnet, European Flux Database and AmeriFlux) while two sites were
chosen from short term intensive field campaigns such as the HAPEX
(Hydrology-Atmosphere Pilot Experiment) and IHOP (International H2O Project).
Most of these sites, however, are not under true bare soil conditions. Merlin et al.
(2016) provided a principle for choosing the "bare soil" period where plant
transpiration is thought to be negligible. For more details please refer to Merlin et al.
(2016).

522 [Table 2 near here]

# 523 4. Results and Discussion

# 524 **4.1. Model Testing with Thin Soil Evaporation**

The thin soil evaporation presented in Wilson (1990) represents exactly the physical 525 process described by the E\_FILM model, that is, a drying process of extremely thin 526 527 (less than 1 mm) soil layer. Hence there exists no impact of moisture supply from 528 below soil as in thick soil column evaporation. The vapor flow within soil can be neglected because there is no moisture supply during stage-III evaporation. The soil 529 water potential was calculated from observed water content by the measured soil 530 water retention curve. The in-equilibrium relative humidity at soil surface is then 531 provided according to the Kelvin equation. Therefore, these thin soil evaporation 532 experiments provide the perfect data for model testing. 533

534 [Figure 3 near here]



535

Model testing results as shown in Figure 3 demonstrated clearly that the provided E\_FILM model was in excellent agreement with observations for almost all nine experiments. Two exceptions are for tests silt\_1 and clay\_1, where the evaporation rate began to decrease while the observed *RH* kept constant as 1. As demonstrated in section 3.1, these two tests applied a different method for sample preparation and yielded a non-uniform drying process. This may explain the mismatch between model predictions and observations.

Another mismatch was observed during stage- I evaporation, where the observed soil evaporation was generally larger than the potential rate. In this experiment, the potential evaporation rate was observed from free water evaporation under the same environment. Therefore, this mismatch may be due to the different vapor diffusion through the thin atmospheric layer above soil and water surface, respectively. This issue is not in the scope of the E\_FILM model.

549 When using E\_FILM model to predict the evaporation rate, only in-equilibrium *RH* 

550 and potential evaporation rate  $LE_p$  are needed, requiring no adjustable parameter in 551 relation to soil texture. The different evaporation rate change among different soils have already been captured by the in-equilibrium RH. According to Figure 2 (c), the 552 critical matric potential is expected to be smaller (more negative) for soils with finer 553 texture under film flow limitations situation. On the contrary, the observed  $RH_c$  was 554 close to 0.99 (corresponding to a water potential of -138 m at 20°C) for all soils. This 555 may be explained by that the critical matric potential that marks the beginning of 556 557 stage-II evaporation is too large (higher than -138 m for all soils) to be captured by RH observations. This is consistent with the description of the evaporation process in 558 section 2.1, where the critical water potential is thought to be close to the value in 559 corresponding to the residual water content. In the literature (e.g., Wang et al., 2016; 560 561 2017), this value was generally higher than -138 m, especially for soils with coarse texture. Nevertheless, as soon as the soil water supply cannot meet the atmospheric 562 demand, the soil surface water potential would decrease dramatically to very negative 563 values, resulting in a decrease of observed RH. 564

# 565 **4.2. Model Testing with Soil Column Evaporation**

566 [Figure 4 near here]



567

The sand column evaporation experiment presented in Zhang et al. (2015) was 568 undertaken with an external heat supply, resulting in much low RH observations near 569 570 soil surface, changing from 18% to 58%. Therefore, the soil water potential may not 571 be in equilibrium with the observed RH. However, as discussed previously, it is the h/ $h_c$  ratio that controls the change of evaporation rate and the observed RH can then be 572 used for predicting the evaporation rate. Besides, with external heat supply, the 573 574 atmospheric demand was very high through the drying process. A decreasing evaporation rate was observed during stage- I period (Figure 4). This might be 575 attributed to the vapor diffusion limitations through above thin boundary layer 576 (Shahraeeni et al., 2012). This effect is not included in the E\_FILM model 577 578 development. Therefore, we only considered the drying process after stage- I period and the critical  $RH_c$  was chosen at day 2 from RH observations. 579

Figure 4 showed that the estimated evaporation rate was generally in good agreement with observations. An underestimation however was found in the low evaporation rate range. Different from the thin soil evaporation, the vapor flow can be important in the soil column evaporation case, especially during stage-III evaporation. By including vapor flow and by setting  $LE_v$  to 1.5 mm d<sup>-1</sup> (the mean vapor diffusion rate from Shokri et al. 2011) in Eq. (11), the model estimation was in excellent agreement with observations (after days 2).

Here, no enhancement factor in relation to vapor diffusion is included (Shokri et al. 2011). The close agreement with observations when considering both film flow and isothermal vapor diffusion might indicate that the unclearly defined enhancement factor (e.g., Saito et al., 2006) actually represents the impact of film flow. However, more testing with different soil columns are needed.

592 The excellent agreement with the observations of both thin soil and thick soil column 593 evaporation revealed that the E\_FILM model (without adjustable parameter) might 594 have captured the actual physical mechanism during stage- II evaporation.

# 595 **4.3. Model Testing with Field Observations**

When testing the E\_FILM model with field observations, significant uncertainties are 596 visible. The main reason is the lack of RH observations at the soil surface. When 597 using the RH observed at 2 m instead, it is difficult to define the critical  $RH_c$  as the 598 599 humidity is not only impacted by soil wetness but also by other atmospheric factors. Under wet conditions, the *RH* at soil surface is generally higher than that at 2 m above, 600 and therefore, the critical  $RH_c$  when using observations at 2 m should have a value 601 smaller than that (close to 0.99) observed in laboratory. In this study, a maximum 602 value of 0.85 derived by trial and error is suggested as the upper boundary for the 603 critical RHc. For model illustration, the soil evaporation module from the widely 604 applied PT-JPL model (Fisher et al., 2008) was chosen for comparison. This module 605 uses an empirical equation to express the water stress, also as a function of RH (see 606 607 Appendix B).

608 [Figure 5 near here]



609

Model testing with 15 sites as shown in Figures 5 and 6 demonstrated that the 610 611 E\_FILM model generally captured the dynamics of evaporation observations. Especially, the E\_FILM model improved the prediction under low evaporation 612 conditions in comparison with the PT-JPL model, as shown at sites DEGeb, FRAur, 613 FRAvi, FRLam and NIHAP. The overestimation of evapotranspiration (ET) rates 614 under arid and semi-arid sites was known to be a common problem in almost all ET 615 models (e.g., Michel et al., 2016). Therefore, the developed E\_FILM model may 616 provide a possible solution for this problem. 617

618 [Figure 6 near here]



619

620 The E\_FILM model underestimated evaporation rate at sites BElon, USDK1 and 621 USFwf. This underestimation can be improved by choosing a lower  $RH_c$ , indicating 622 that the critical value is actually different under different conditions.

In summarize, the E\_FILM model generally yielded lower RMSEs in comparison with the PT-JPL model, with a mean value of 31.67 W m<sup>-2</sup> and 35.43 W m<sup>-2</sup>, respectively (Figures 5 and 6). For  $R^2$ , however, the PT-JPL model generally presented higher values, with a mean value of 0.45 compared to the 0.38 of the E\_FILM model. This much lower  $R^2$  with the E\_FILM model was mainly due to the poor performance at sites BELon and USDK1.

The relatively poor performance with field observations might be due to two main 629 reasons. Firstly, most testing sites were not in a true bare soil situation. As shown in 630 table 2, only two sites (NIHAP and USIHO) were under bare soil condition while the 631 observations from other sites were chosen from a "bare soil" period assuming plant 632 transpiration was "negligible or small compared to soil evaporation" (Merlin et al., 633 2016). However, without direct measurement of the soil evaporation, it is hard to 634 evaluate the data quality. Secondly, high uncertainty existed in relation to the RH 635 input. In this study, the RH data used were observed at 2 m above the soil surface and 636 should be different from the in-equilibrium RH at the soil surface. Besides, an upper 637

638 boundary of 0.85 was set for the critical  $RH_c$  in this study. However, as discussed previously, this value might be different at different sites. For example, the 639 underestimation at sites BELon and USDK1 can be improved by setting a lower 640 critical  $RH_c$ . However, it was hard to define the actual  $RH_c$  since this value was not 641 only related to soil texture but also to atmospheric conditions. Uncertainty also 642 existed in defining  $RH_m$ . In this study,  $RH_m$  was simply set as the lowest value of the 643 observed RH. This principle was appropriate only when the observations covered a 644 645 complete drying period. Note that the impacts of  $RH_c$  and  $RH_m$  are highly correlated in the E\_FILM model. 646

In spite of this disadvantage, the E\_FILM model was attractive for its solid physical 647 base, which was evidenced by laboratory observations. Moreover, the estimation of 648 soil evaporation in field showed an improvement under dry conditions. The film 649 dominant assumption in the E\_FILM model was also consistent with the field drying 650 experiment in Goss & Madliger (2007), where they found the film flow has a 651 significant contribution in soil evaporation. However, further work is needed to 652 653 derive a complete evaporation rate estimation model by including both capillary and film limitations, and to improve the model performance in field, for example, 654 considering the input of soil surface water content. 655

#### 656 **5. Conclusions**

657 This study provided a new interpretation of the typical soil evaporation process by including the impact of film flow along the soil particle surface. This film flow was 658 usually thought to be unimportant and then be neglected in soil drying process (e.g., 659 Philip & de Vries, 1957; Idso et al., 1974). However, recent progress in soil hydraulic 660 661 modeling development confirmed this film flow as a dominant process under low moisture conditions (e.g., Tuller & Or, 2001; Peters, 2013; Wang et al., 2016, 2017, 662 2018). By including this film flow, the typical drying process was revisited. The result 663 664 found that this film flow might be the dominant process for limiting the evaporation loss under stage- II evaporation, in which the vapor flow was usually regarded as the 665

limitation factor (e.g., Philip & de Vries, 1957; Saito et al., 2006; Lehmann et al.,
2008; Or et al., 2013). The typical evaporation process was then interpreted as the
capillary-flow-supported stage- I evaporation, the film-flow-controlled stage- II
evaporation and the vapor-diffusion-dominant stage- II evaporation.

Based on the assumption that film flow controlled the stage-II evaporation, a physically based model was developed by parameterizing the Buckingham-Darcy's law. The model provided a solid basis for describing the moisture limitation on evaporation rate, requiring only meteorological data as input and introducing no adjustable parameter. The impact of vapor flow was also discussed. It was found to be only important for soils with coarse texture under very dry conditions.

Model testing with laboratory data, including nine thin soil thickness evaporation tests 676 and one thick sand column evaporation, yielded excellent agreement with 677 observations. The model evaluation with 15 field sites however introduced some 678 uncertainty. The main reason is the lack of in-equilibrium relative humidity 679 observations near the soil surface. Nevertheless, the proposed E FILM model 680 681 significantly improves the performance under dry conditions in comparison with the widely applied PT-JPL model. Since the evaporation overestimation in arid and 682 semi-arid regions has been found in almost all evapotranspiration estimation models 683 (Michel et al., 2016), this E\_FILM model provides an opportunity to improve the 684 evaporation estimation under such dry conditions. However, the present E\_FILM 685 model didn't consider the impact of capillary limitations which might be important in 686 fine-textured soils (Haghighi et al., 2013; Lehmann et al., 2018) and further work is 687 needed to test and improve the model performance in field, especially in relation to 688 689 the determination of the critical  $RH_c$  that marks the beginning of stage- I evaporation.

# 690 Appendix A

691 When both film flow and vapor diffusion are included, the evaporation rate can be 692 expressed by parameterizing the Buckingham-Darcy's law (Eq. 7), as

$$LE \approx \left(K_f + K_v\right) h \ln\left(\frac{h_0}{h_c}\right) \frac{\left(S_f - S_m\right)}{z_d}$$
(A1)

694 where  $K_f$  and  $K_v$  are the hydraulic conductivity that accounts for film flow and 695 isothermal vapor flow, respectively.

696 By applying the effective surface area, the film conductivity is expressed as

693

$$K_f = \frac{2\rho g \theta_f}{3\pi\eta} f^2$$
(A2)

698 where  $\rho$  is the water density (9.98×10<sup>2</sup> kg m<sup>-3</sup>), g is the acceleration of gravity (9.81 699 m s<sup>-2</sup>),  $\eta$  is fluid viscosity (1.005×10<sup>-3</sup> Pa s at 293 K) and the film thickness f is 700 expressed following Tuller & Or (2001), as

$$f = \sqrt[3]{\frac{-A_{svl}}{6\pi\rho gh}}$$
(A3)

where  $A_{svl}$  is the Hamaker constant for solid-vapor interactions through the intervening liquid and is set as  $-6.0 \times 10^{-20}$  J following Tuller & Or (2001).

The isothermal vapor conductivity is given in Saito et al. (2006), as

705 
$$K_{\nu} = \frac{\rho_{\nu}}{\rho} D \frac{Mg}{RT} RH$$
(A4)

where  $\rho_{\nu}$  (kg m<sup>-3</sup>) is the saturated vapor density, D (m<sup>2</sup> s<sup>-1</sup>) is the vapor diffusivity, written as

708  $D = \tau \theta_a D_a \tag{A5}$ 

with  $\theta_a$  being the air-filled porosity,  $\tau$  the tortuosity factor calculated according to Millington and Quirk (1961), as

- 711  $\tau = \frac{\theta_a^{7/3}}{\theta_s^2}$ (A6)
- 712 with  $\theta_s$  the saturated water content.

And  $D_a$  is the vapor diffusivity in air, written as

714 
$$D_a = 2.12 \times 10^{-5} \left(\frac{T}{273.15}\right)^2$$
(A7)

The soil evaporation in the PT-JPL model (Fisher et al., 2008) is estimated by

717 
$$LE = RH^{VPD/\beta}LE_{p}$$
(B1)

with VPD (kPa) being vapor pressure deficit,  $\beta$  (1.0 kPa) representing the relative sensitivity to VPD.

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Figure 1. (a) Illustration of the normalized evaporation rate and the predicted soil 988 surface water content from an initial saturated soil column. Three stages are identified, 989 including the near constant rate stage- I evaporation, the fast falling rate stage- II 990 991 evaporation and the smoothly changing stage-III evaporation. (b) The near surface soil 992 water content profiles observed by magnetic resonance imaging at different evaporation stages after Merz et al. (2015). (c) Illustration of the soil hydraulic 993 994 conductivity curve over the complete moisture range with the EMFX model proposed in Wang et al. (2016, 2017), including capillary flow  $(K_c)$ , film flow  $(K_f)$  and vapor 995 996 diffusion  $(K_v)$ . The vapor diffusion is calculated following Saito et al. (2006). (d) 997 Illustration of the typical soil water flow process under three stages, the soil particle is 998 in yellow and the water is in blue. The stage- I evaporation is supported by capillary flow retained in soil pores, the drying process is accompanied by an increase in the 999 1000 drying front depth; During stage-II evaporation, capillary flow can no longer reach 1001 soil surface directly, the drying process is limited by film flow along soil particle surface, dominated in a thin -several millimeters in depth- soil surface layer; Under 1002 1003 stage-III evaporation, this thin soil surface layer is almost completely dried (with 1004 water content be equal to the air-dry water content), vapor diffusion becomes the 1005 dominant process in this layer.

Figure 2. (a) The relationship between water potential and relative humidity under different temperatures. (b) Illustration of the E\_FILM model with different critical water potential values; the surface water content is calculated with Eq. (3). (c) Illustration of the evaporation rate calculated with different critical water content values; the solid line accounts for both film flow and vapor diffusion while the dashed line accounts for film flow only (see Appendix A). The air-dry water content  $\theta_m$  is set as zero and the depth of the thin layer  $z_d$  is set as 5 mm for illustration.

Figure 3. E\_FILM model testing with thin soil evaporation experiments from Wilson(1990), including sand, silt and clay.

1015 **Figure 4**. E\_FILM model testing with thick sand column evaporation experiment 1016 from Zhang et al. (2015)

1017 **Figure 5**. Predicted evaporation flux with the E\_FILM model and the PT-JPL model

1018 in the first nine field sites,  $LE_o$  and  $LE_s$  are the observed and the simulated 1019 evaporation rate, respectively.

- 1020 **Figure 6**. Predicted evaporation flux with the E\_FILM model and the PT-JPL model
- 1021 in last six field sites,  $LE_o$  and  $LE_s$  are the observed and the simulated evaporation rate,

1022 respectively.

1023

# 1024 Table captions

#### 1025

1026 **Table 1**. The laboratory evaporation tests from Wilson (1990) and Zhang et al. (2015)

Test No.	Soil Type	Sample Thickness (mm)	<i>RH</i> of air	Reference
Sand_1	Beaver Creek Sand	0.7	0.53	Wilson (1990)
Sand_2		0.5	0.44	
Sand_3		0.5	0.58	
Silt_1	Custom Silt	0.5	0.21	Wilson (1990)
Silt_2		0.3	0.39	
Silt_3		0.3	0.62	
Clay_1	Regina Clay	0.7	0.39	Wilson (1990)
Clay_2		0.3	0.35	
Clay_3		0.2	0.50	
Soil column	Medium Sand	500	0.50	Zhang et al. (2015)

# 1027

1028 **Table 2.** Flux Sites including bare soil periods, modified from Merlin et al. (2016)

Site	Exp./Net.	Lat;lon	Land cover	Soil Texture	Reference
AUStu	OzFlux	-17.15;133.35	grass	Silt loam	Beringer et al. (2011)

BELon	GHGEurope	50.55;4.74	crop	Silt loam	Papale et al. (2006)		
CHOe2	GHGEurope	47.29;7.73	crop	Silty clay	Alaoui and Goetz (2008)		
DEGeb	GHGEurope	51.10;10.91	crop	Silty clay loam	Kutsch et al. (2010)		
ESES2	GHGEurope	39.28; -0.32	crop	Silty clay	Kutsch et al. (2010)		
FRAur	GHGEurope	43.55;1.11	crop	Clay loam	Béziat et al. (2009)		
FRAvi	GHGEurope	43.92;4.88	crop	Silty clay loam	Garrigues et al. (2015)		
FRLam	GHGEurope	43.50;1.24	crop	Clay	Béziat et al. (2009)		
ITBCi	GHGEurope	40.52;14.96	crop	Clay	Denef et al. (2013)		
NIHAP	HAPEX	2.24;13.20	bare	Sand	Wallace et al. (1993)		
USArm	AmeriFlux	36.61; -97.49	crop	Clay	Fischer et al. (2007)		
USDk1	AmeriFlux	35.97; -79.09	grass	Loam	Novick et al. (2004)		
	AmeriFlux		35.45;			Dave et al. (2012)	
USFWI		-111.77	grass	Silt loam	Dore et al. (2012)		
USIb1	AmeriFlux	41.86; -88.22	crop	Silty clay loam	Wu et al. (2012)		
USIHO	IHOP	36.47; 100.62	bare	Sandy clay loam	Lemone et al. (2007)		