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PARTICLE MATERIAL RECOGNITION USING DIGITAL IN LINE HOLOGRAPHY AND VIRTUAL/REAL IMAGE CONCEPTS

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Abstract

The authors have shown recently [1, 2] that digital in-line holography (DIH) can recover some features of the electromagnetic near field of large spherical transparent particles. Among these features, a special attention was paid to the photonic jet whose intensity maxima positioning allows to estimate the particle refractive index. In the present communication, we show that, in the framework of the Fresnel diffraction approximation, the reconstructed photonic jet can be seen as the Airy's spot of a spherical lens. More surprisingly, we also demonstrate that DIH can recover an image of this near field structure that can be real or virtual, referring to geometrical optics terminology. Besides the theoretical interest, this concept provides a simple and numerically effective solution to characterize the composition of solid, liquid and gas particles and opens up huge perspectives for multiphase flow investigations.

1 Introduction

Digital in line holography (DIH) is a powerful volumetric imaging technique for the characterization of moving particle's properties: position, shape, etc. (e.g. [3, 4]). For obvious reasons, the particle composition is also an important property, but very few studies are devoted to this kind of measurement, especially in combination with the previous features. Considering DIH, the particle material recognition can be achieved either by means of morphological [5] or refractive index considerations [6]. Shape considerations are limited by the necessity of an a priori knowledge of the particles morphology [7] or the use of geometrical criteria [5]. On the other hand, diagnostics based on refractive index measurement are classically built on a best fitting approach [8] between experimental and directly simulated holograms. For small particle (e.g., diameter $D < 10 \mu\text{m}$), the hologram patterns recorded in the forward direction are quite sensitive to the particle composition and the direct simulation of the corresponding holograms remains numerically affordable [9]. This is not the case anymore for large particles, as the sensitivity to the refractive index of the fitting approach decreases rapidly while its computation time increases drastically.

Recently the present authors have shown that, for refracting particles, DIH enables to recover some features of their photonic jet (maximum intensity position) useful to determine the particle refractive index. In the scope of geometrical optics, the photonic jet was seen as the apex of

an exterior cusp caustic (e.g. [10]). In this communication, this near field structure is interpreted in the framework of the scalar diffraction theory (e.g. [11]) resulting in more information about the particle's composition.

2 Photonic jet interpretation in wave optics scope

In the following discussion, we consider an unpolarised harmonic plane wave, with a wavelength in vacuum λ , which propagates along the z -axis in a free medium with real refractive index m_2 . It illuminates a spherical particle of radius a , with real refractive index m_1 and relative refractive index $m = m_1/m_2$.

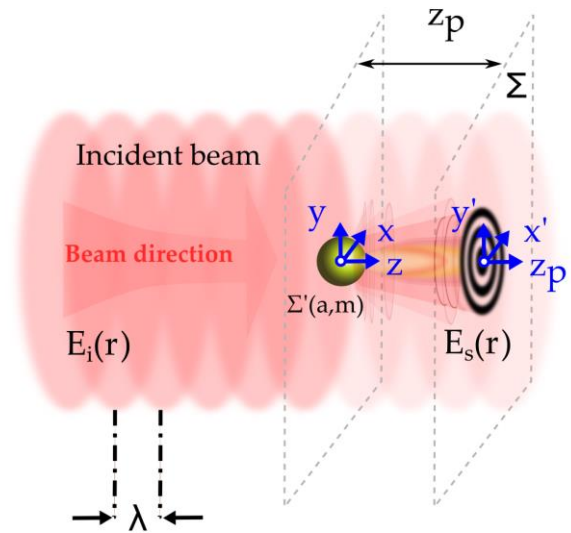


Figure 1: Schematics of DIH principle

For the sake of simplicity, the particle is supposed to be positioned at the centre of the Cartesian coordinate system ($Oxyz$), and the sensor (a quadratic imaging surface $\Sigma(x, y; z_p)$) is set to be perpendicular to the z -axis at an axial distance $z = z_p$, as illustrated in Figure 1. In DIH, the particle hologram is the intensity image of the scattering pattern recorded by the sensor, when it is exposed both to the incident field $E_i(r)$ (constant) and the field that is scattered $E_s(r)$ in the forward direction by the particle. The corresponding total time-averaged intensity can be expressed as:

$$I_t(r_p) = \frac{\varepsilon c}{2} |E_i(r_p) + E_s(r_p)|^2 \quad (1)$$

where $r_p(x, y, z_p)$ is a coordinate vector describing the surface Σ , and ε and c are respectively the electrical permittivity and light speed in the medium. By introducing the other intensities, $I_\chi(r_p) = \varepsilon c |E_\chi(r_p)|^2 / 2$ where $\chi = i, s, t$ stand respectively for the incident, scattered and total fields, Eq. (1) can be rewritten as:

$$I_t(r_p) \propto I_i(r_p) + I_s(r_p) + 2\Re\{E_i(r_p)\bar{E}_s(r_p)\} \quad (2)$$

The last term in Eq. (2) depicts the interference pattern corresponding to the hologram. Classically, to increase the contrast and eliminate hidden normalization and gain constants, the hologram is conveniently divided by the intensity of the incident field (without particles). Additionally, it is usually assumed that the intensity of the scattered field is negligible compared to the incident field, $I_s(r_p)/I_i(r_p) \ll 1$. Consequently, DIH, which uses the incident wave as a reference wave, is limited to diluted particle systems. The normalized hologram $I_h(r_p)$ can be simply approximated as follows:

$$I_h(r_p) = \frac{2\Re\{E_i(r_p)\bar{E}_s(r_p)\}}{I_i(r_p)} \approx \frac{I_t(r_p)}{I_i(r_p)} - 1 \quad (3)$$

To simulate $I_t(r_p)$, electromagnetic theories like the Lorenz-Mie theory (LMT) or the Debye theory can be used. Several authors have demonstrated that this approach is applicable to small particles of micrometer size but becomes quickly prohibitive from a computational point of view for the characterization of large particles with variable composition. For the latter the scalar diffraction theory is preferred especially considering the Fresnel's paraxial approximation:

$$E_s(r) = \frac{e^{jkz}}{j\lambda z} e^{j\frac{k}{2z}(x^2+y^2)} \iint_{\Sigma'} \left\{ E_p(r') e^{j\frac{k}{2z}(x'^2+y'^2)} \right\} e^{-j\frac{jk}{z}(xx'+yy')} dx' dy' \quad (4)$$

where $r'(x', y'; z=0)$ is a coordinate vector describing the surface Σ' and $E_p(r')$ is characteristic field of the particle to be propagated to the sensor. The Fresnel approximation is valid only when $ka \gg 1$ and $z \gg \sqrt{ka^4/2}$.

In Eq. (4) the field $E_p(r')$ is the product of the incident field, by the particle transmission function $T(x', y') = t(x', y')e^{j\phi(x', y')}$, where $t(x', y')$ is an amplitude term and $\phi(x', y')$ is a pure phase term. Usually, considering the diffraction of a circular object [11], the term t is modelled by a 2D pupil function. In a similar way, in DIH, a spherical particle with radius a is simply assimilated to a circular hole (or an opaque disk according to Babinet's principle), with the following transmission function:

$$t(x', y') = \begin{cases} 1, & \text{if } x'^2 + y'^2 \leq a^2 \\ 0 & \text{elsewhere} \end{cases} \quad (5)$$

where $\phi(x', y') = 0$.

For large particles and far field measurement, the models rely on the several order of magnitude deviation prevailing between the pure diffraction process (partial waves of order $p=-1$ [13]) and all other scattering mechanisms: specular reflection ($p=0$), single refraction ($p=1$), etc. [12, 13], allowing to neglect the latters.

Nevertheless, this assumption is not valid in near-field. Indeed, if we consider only the dominant effect after the diffraction process, i.e. single refraction ($p=1$) for transparent refracting particles ($m>1$, e.g. droplets or glass beads in water and $m<1$, e.g. gas bubbles in water) the term $e^{j\phi(x', y')}$ accounts for the angular deviation of the incident rays [14] and cannot be neglected.

In the paraxial approximation, simple ray optics gives for the effective focal length, f , of a spherical object of radius a and relative refractive index m [11]:

$$f = ma / (2(m-1)) \quad (6)$$

with $f > 0$ for refracting particles and $f < 0$ for reflecting ones. This allows to retain the following phase term for a spherical particle [11]:

$$\phi(x, y) = 2ka - \frac{k}{2f}(x'^2 + y'^2) \quad (7)$$

The combination Eqs. (4) and (7) give for the scattered field:

$$E_s(r) = E_i \frac{e^{jk(z+2a)}}{j\lambda z} e^{j\frac{jk}{2z}(x^2+y^2)} \times \iint_{x'^2+y'^2 \leq a^2} e^{-j\frac{jk}{2f}(x'^2+y'^2)} e^{j\frac{jk}{2z}(x'^2+y'^2)} e^{-j\frac{jk}{z}(xx'+yy')} dx' dy' \quad (8)$$

For $z=f$, the two quadratic phase exponential terms of Eq. (8) cancel each other and E_s becomes:

$$E_s(x, y; f) = E_i \frac{e^{jk(z+2a)}}{j\lambda f} e^{j\frac{jk}{2f}(x^2+y^2)} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} t(x', y') e^{-j\frac{jk}{f}(xx'+yy')} dx' dy' \quad (9)$$

Note that in Eq. (9), the integration boundaries have been extended to infinite, and the pupil function reintroduced in order to retrieve a scaled Fourier transform of the particle transmission function. For $t=1$ and a plane wave illumination $E_i = 1$, E_s can be written as:

$$E_s(x, y; z_p = f) \approx \frac{e^{jk(z+2a)}}{j\lambda f} e^{j\frac{jk}{2z}(x^2+y^2)} \left[2\pi a^2 \frac{J_1(ka\rho/f)}{ka\rho/f} \right] \quad (10)$$

with for the radial intensity profile:

$$I_s(\rho; z_p = f) = 4 \left(\frac{\pi a}{\lambda f} \right)^2 \left[\frac{J_1(ka\rho/f)}{ka\rho/f} \right]^2 \quad (11)$$

where $\rho = \sqrt{x^2 + y^2}$ and J_1 is the first order spherical Bessel function. The function in Eq. (11) is known as the Airy distribution.

Properties of Airy's spot produced by a thin lens are widely reported in the literature [15], especially its intensity profile that takes the following form:

$$I_s(0,0,z) = \left(\frac{\pi a}{\lambda f} \right)^2 \left[\frac{\sin(km_2 a^2 / f^2)}{km_2 a^2 / f^2} \right]^2 \quad (12)$$

Generally, the transverse and axial characteristic dimensions of the Airy spot are defined by two radii $a_x \approx 0.61\lambda a / f$ and $a_z \approx 2z_R$ respectively, where $z_R = 4m_2\lambda a^2 / f^2$ stands for the Rayleigh distance.

From the foregoing, several conclusions must be highlighted:

- The photonic jet predicted by electromagnetic theories [16] and observed experimentally for droplets in liquid-liquid systems [1, 2] can be seen as the Airy's spot of transparent focusing particle when considering wave optics.
- It can be recovered by DIH near-field reconstruction but only roughly evaluated considering geometrical optics approximations as we observe the Point Spread function (PSF) of the lens for $z_p = f$ [11].
- The particle acts as converging $f > 0$ or diverging $f < 0$ lens depending on the relative refractive index ($m > 1$ or $m < 1$). As the above interpretation is valid for both lenses, the existence of a "virtual" photonic jet (or Airy's spot) is predictable for the diverging ones.

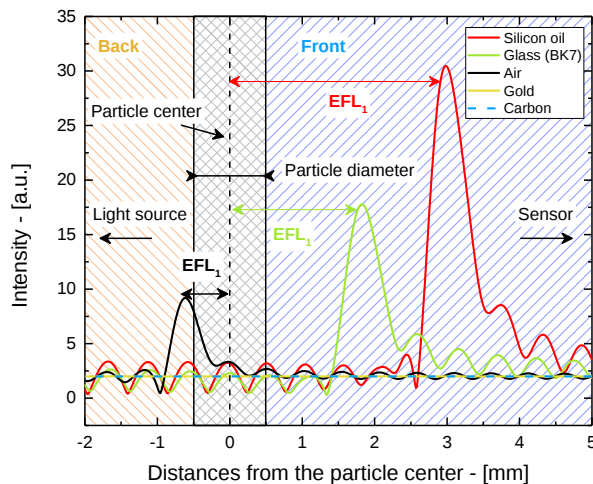


Figure 2: Axial profiles in the near-field, calculated by back-propagating LMT hologram's simulation, for five particles having the same radius (500μm) and different compositions. DIH parameters: HeNe laser ($\lambda=0.6328\mu\text{m}$), square sensor (with 2000x2000 square pixels of 5.0μm on a side) located at 100mm from the particle center.

3 Numerical and experimental illustrations

The axial profiles, plotted in Figure 2, have been calculated by back-propagating holograms I_h simulated with LMT. The back-propagation was performed in the Fresnel approximation [1]. Five types of particles dispersed in water are considered in these simulations. They all have the same radius (500μm) but different compositions: two opaque particles (gold and carbon beads) and three transparent particles (air bubbles, silicon oil droplets and glass beads). The first thing to note on this figure is the lack of focalisation peak for the opaque particles (gold and carbon). The second remarkable element is the presence of a maximum of intensity corresponding to the axial profile of the Airy's spot for the transparent particles. The position of this maximum, compared to the particle's center, is quite close to the effective focal length f of the considered lens (respectively -0.7mm, 3.3mm and 2.1mm), confirming that geometrical optics can be used as a first rough but fast approximation. Moreover, when considering particles with a relative refractive index <1 , this maximum appears behind the particle. It can be noticed that the axial profiles exhibit an asymmetry not predicted by Eq. (12) probably due to the simplicity of the thin lens model used.

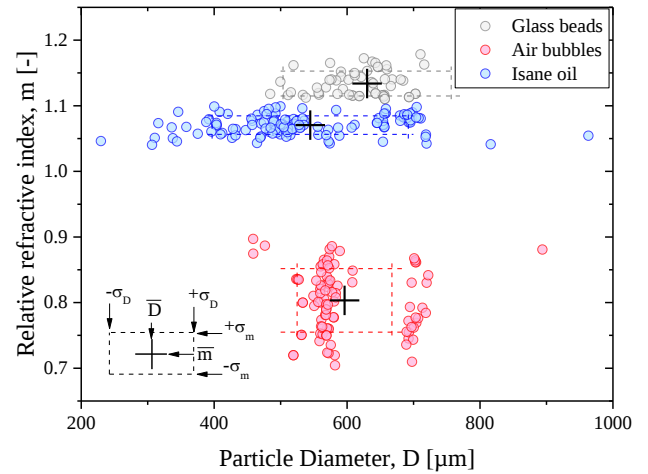


Figure 3: Simultaneous measurement of the relative refractive index and diameter of rising air bubbles (red), rising Isane droplets (blue) and falling glass beads (gray). Black crosses indicate the mean values (diameters and refractive indices), while the dashed lines show the standard deviations at $\pm 1\sigma$.

Finally, a classical digital in line holography experimental setup has been used to demonstrate the capability of the method to determine simultaneously the 3D position, velocity, size and composition of flowing particles. The flow-cell consists of a rectangular tank filled with water. Three dispersed phases were considered: free rising droplets of Isane_oil ($m=1.0826$) and air bubbles ($m=0.7505$) generated at the bottom of the tank by two drop on demand generators, and a stream of falling glass beads ($m=1.1370$). The DIH optical setup was simply composed of a fiber-injected HeNe laser ($\lambda = 632.8\text{nm}$, 10mW),

collimation optics and a S-CMOS camera (2560x2160 square pixels of 6.5 μ m on a side) placed in Gabor configuration at $z_r=100$ mm. The recorded holograms were back-propagated using a generalized Fresnel propagator allowing to account for the water-glass-air interfaces as well as astigmatic effects [1,2]. Figure 3 shows some preliminary experimental results; the three populations are clearly recognized.

4 Conclusion and outlooks

In this communication we report that the photonic jet produced by large transparent particles corresponds to the Airy's spot of a thin lens and can be reconstructed by back-propagation of DIH holograms. This phenomenon is directly linked to refracted rays and cannot be neglected in the near-field. DIH reconstructions display a real image of this spot for refracting particles ($m>1$), and a virtual image in the case of the reflecting ones ($m<1$). Furthermore, the spot disappears logically when considering opaque particles. This approach opens up very promising perspectives for multiphase flow extensive characterization, including material recognition.

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