

Dynamic force microscopy in non-contact mode Laurent Nony

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Institut Matériaux Microélectronique Nanosciences Provence

Dynamic force microscopy in non-contact mode

Laurent Nony laurent.nony@im2np.fr

Aix-Marseille Université CNRS, Institut Matériaux Microélectronique et Nanosciences de Provence, UMR 7334, 13397, Marseille, France

> M2 Nanosciences et Nanotechnologies UE Nano-objects « High-resolution imaging » 2019-2020









Documents & references

 Many documents and references related to the lecture "High-resolution imaging" are available on the AMUBox server:

https://amubox.univ-amu.fr/s/T678xB289tXCJas

 The references cited in the present document are given in the above link within the folder "Dynamic force microscopy"













Atomic Force Microscopy

• Microscopie à Force Atomique = Atomic Force Microscopy : AFM

• Microscopy technique developed in 1986 by the inventors of the Scanning Tunnelling Microscope (STM, 1982 at IBM Zurich, Physics Nobel Price) \rightarrow Images of the surface of any kind of sample in real space with atomic vertical and lateral resolution by using interaction forces between a nano-tip and the surface

STM

Source: Applied Physics Letters 40, p178 (1982)

Tunneling through a controllable vacuum gap

G. Binnig, H. Rohrer, Ch. Gerber, and E. Weibel IBM Zurich Research Laboratory, 8803 Rüschlikon-ZH, Switzerland

(Received 30 September 1981; accepted for publication 4 November 1981)

We report on the first successful tunneling experiment with an externally and reproducibly adjustable vacuum gap. The observation of vacuum tunneling is established by the exponential dependence of the tunneling resistance on the width of the gap. The experimental setup allows for simultaneous investigation and treatment of the tunnel electrode surfaces.

PACS numbers: 73.40.Gk

Source: Physical Review Letters 49, p57 (1982)

Surface Studies by Scanning Tunneling Microscopy

G. Binning, H. Rohrer, Ch. Gerber, and E. Weibel *IBM Zurich Research Laboratory*, 8803 Rüschlikon-ZH, Switzerland [Received 30 April 1982]

Surface microscopy using vacuum tunneling is demonstrated for the first time. Topographic pictures of surfaces on an *atomic scale* have been obtained. Examples of resolved monoatomic steps and surface reconstructions are shown for (110) surfaces of $CaIrSn_4$ and Au.

PACS numbers: 68.20.+t, 73.40.Gk

AFM

Source: Physical Review Letters 56, p930 (1986)

Atomic Force Microscope

G. Binnig^(a) and C. F. Quate^(b) Edward L. Ginzton Laboratory, Stanford University, Stanford, California 94305

and

Ch. Gerber^(c)

IBM San Jose Research Laboratory, San Jose, California 95193 (Received 5 December 1985)

The scanning tunneling microscope is proposed as a method to measure forces as small as 10^{-18} N. As one application for this concept, we introduce a new type of microscope capable of investigating surfaces of insulators on an atomic scale. The atomic force microscope is a combination of the principles of the scanning tunneling microscope and the stylus profilometer. It incorporates a probe that does not damage the surface. Our preliminary results *in air* demonstrate a lateral resolution of 30 Å and a vertical resolution less than 1 Å.

PACS numbers: 68.35.Gy











Outlook

A. AFM in 3 steps (contact mode concept)

B. Dynamic AFM

C. Dynamic AFM in non-contact mode













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A – AFM in 3 steps (contact mode concept)

- 1. Tip-cantilever system
- 2. Optical detection setup
- 3. Surface scanning and feedback















A – AFM in 3 steps (contact mode concept)

- 1. Tip-cantilever system
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• **AFM = scanning probe technique** (nano-tip close to a surface): tip is located at the end of a micro-lever = **cantilever**



• AFM = scanning probe technique (nano-tip close to a surface): tip is located at the end of a micro-lever = cantilever

Mechanical analogy: recessed beam

 Further simplification: spring with stiffness k subject to an interaction force F_{int} with the surface, resonance frequency of the fundamental bending mode f₀

Orders of magnitude cantilevers Size: (100 x 20 x 3) μm³ Stiffness: 0.1 N/m ... 1800 N/m Frequency: 10 kHz ... 1 MHz

Orders of magnitude for tips Aspect ratio upon resalers Height: 10 μm Curvature radius: 5nm (at best)





um

100 µm













• AFM = scanning probe technique (nano-tip close to a surface): tip is located at the end of a micro-lever = cantilever

- Mechanical analogy: recessed beam
 - Further simplification: spring with stiffness k subject to an interaction force F_{int} with the surface, resonance frequency of the fundamental bending mode f₀
- Measure of the lever deflection by means of an optical setup ⇒ indirect measure of forces (1 pN..1 nN):
 - *Electrostatic* (due to charges) : attractive, long-range (> 100 nm)
 - Van der Waals (London, Debye, Keesom) : attractive, long-range
 (> 1 nm)
 - Chemical (interatomic) : attractive or repulsive, short-range (few Å)
 - Magnetic (long-range)
- For a detailed work on forces involved in AFM:
 - "Scanning probe microscopy: the lab on a tip" (E.Meyer, H.J. Hug, R. Bennewitz, Springer-Verlag Berlin Heidelberg 2004
 - "Noncontact Atomic Force Microscopy" (S.Morita, R. Wiesendanger, E. Meyer, Springer-Verlag Berlin Heidelberg 2002
 - "Intermolecular and surface forces, 3rd edition" (J.N.Israelachvili, Elsevier 2011)

















• Tip-surface interaction force¹:











• AFM = scanning probe technique (nano-tip close to a surface): tip is located at the end of a micro-lever = cantilever

Mechanical analogy: recessed beam

 Further simplification: spring with stiffness k subject to an interaction force F_{int} with the surface, resonance frequency of the fundamental bending mode f₀

100 μm 20 μm 10 μm 10 μm



Scaling illustration:



G.Binnig et al., Phys. Rev. Lett. 56, p930 (1986)









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- 4-quadrants PhotoSensitive Detector: PSD
 - 4 PN junctions side-by-side
 - 4 anodes + 1 common cathode



В



sociation







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Movement detection:





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Movement detection:





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Movement detection:





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Movement detection:



Cantilever deflection = « Normal force $(F_N) = (A + B) - (C + D)$

Cantilever torsion = « Lateral force $(F_L) = (A + D) - (B + C)$









A - AFM in 3 steps (contact mode concept)

1. Tip-cantilever system

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A3- Surface scanning & feedback





G.Binnig et al., Phys. Rev. Lett. 56, p930 (1986)













A3- Surface scanning & feedback



In AFM, the vertical feedback can be performed on various signals, not only on the deflection F_N (see hereafter)!

















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A3- Surface scanning & feedback

How AFM Works

www.ParkAFM.com

Park Systems, Inc.













Outlook

A. AFM in 3 steps (contact mode concept)

B. Dynamic AFM

C. Dynamic AFM in non-contact mode













B – Dynamic AFM

- 1. Dynamic AFM specificities
- 2. Force sensitivity
- 3. Non-linear dynamics and modelization
- 4. Tapping mode AFM vs. non-contact AFM













B – Dynamic AFM

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B1- Dynamic AFM specificities

Problem: What is the pressure that applies on a 1 $Å^2$ surface upon a 1 nN force?



$$P = \frac{F}{s} = \frac{10^{-9}N}{(10^{-10}m)^2} = 100 \ GPa$$

NB: Steel's Young modulus: 200 GPa

Consequence: Tip and/or surface damaging.















B1- Dynamic AFM specificities

- Cantilever excitation at a frequency equal or close to its resonance frequency :
 - Advantages:
 - 1. Lowering the shearing force between tip and surface
 - 2. Increasing force sensitivity w.r.t. the contact mode
 - 3. Non-contact (attractive force mapping only)

Difficulties:

- 1. Understanding the cantilever non-linear dynamics
- Detection of the oscillation (system detection bandwidth> MHz)
- 3. Complexity of the electronics scheme
- 4. Origin of the image contrast?

Measures with the tip-surface distance:

- 1. Oscillation frequency
- 2. Oscillation ampltidue
- 3. Excitation amplitude
- 4. Phase shift between excitation and oscillation
- The mechanical system is strictly similar to a forced damped harmonic oscillator (to be checked experimentally however)













Mode normal (fondamental)



Orders of magnitude: Amplitude : 50 pm ... 10 nm Frequency: 32 kHz ... 1 MHz

B1- Dynamic AFM specificities

□ 5 variables : <u>Excitation amplitude</u>, <u>Excitation frequency</u>, <u>Oscillation amplitude</u>, <u>Phase</u> (excitation / oscillation) and <u>Tip-surface distance</u>



Cantilever quality factor and intrinsic dissipation:

$$Q = 2\pi \frac{\langle E_{stored} \rangle_T}{\langle E_{diss} \rangle_T} = \frac{2\pi f_0}{\Gamma_0}$$

Differential equation for
$$\mathbf{D} \to \infty$$

 $\ddot{z}(t) + \frac{\omega_0}{Q} \dot{z}(t) + \omega_0^2 z(t) = \omega_0^2 \mathbf{A}_{\text{exc}}(t)$

Harmonic solution:

$$z(t) = A(\omega)\cos(\omega t + \varphi(\omega))$$

¹ Polesel et al., Nanotechnology **14**,1036 (2003)



$$u\left(u=\frac{\nu}{\nu_0}\right) = \frac{A(u)}{A_0} = \frac{1}{\sqrt{Q^2(1-u^2)^2+u^2}}$$
$$\varphi(u) = \arctan\left\{\frac{u}{Q(u^2-1)}\right\}$$









Exercice

Starting from the differential equation for the forced damped harmonic oscillator (Quality factor Q, pulsation at resonance $\omega_0 = 2\pi v_0$):

$$\ddot{z}(t) + \frac{\omega_0}{Q}\dot{z}(t) + \omega_0^2 z(t) = \omega_0^2 A_{\text{exc}}(t)$$

with a forcing term of the form:

$$A_{\text{exc}}(t) = A_x \cos(\omega t) = \frac{A_0}{Q} \cos(\omega t)$$

and assuming a harmonic solution:

$$z\left(t
ight) = A\left(\omega
ight)\cos\left(\omega t + \varphi\left(\omega
ight)
ight)$$

1- Prove that the normalized amplitude $a(u) = A(u)/A_0$, with $u = \omega/\omega_0$ being the normalized frequency, and the phase of the oscillator write :

$$a\left(u=\frac{\nu}{\nu_0}\right) = \frac{A(u)}{A_0} = \frac{1}{\sqrt{Q^2(1-u^2)^2+u^2}}$$
$$\varphi(u) = \arctan\left\{\frac{u}{Q(u^2-1)}\right\}$$

2- Plot these functions (with Igor Pro for instance) with Q = 100, then 100000.











B – Dynamic AFM

1. Dynamic AFM specificities

2. Force sensitivity

3. Non-linear dynamics and modelization

4. Tapping mode AFM vs. non-contact AFM













B2-Force sensitivity (for cantilevers)

Equipartition theorem on a cantilever subject to thermal fluctuations yields a minimum detectable force :

$$F_{\min} = \sqrt{\frac{4kk_BT}{\omega_0 Q}B}$$

- Orders of magnitude:
 - k = 30 N/m
 - B ~ 250 Hz
 - $\omega_0 = 2\pi f_0 = 2\pi 150 \text{ kHz}$
 - Q = 40 000 (sous vide)
 - k_BT = 25 meV (298 K)

$$F_{min} \sim 60 \ pN$$

- Increasing the sensitivity :
 - k : risk for "jump to contact"
 - ✓ ➤k_BT : not easy

 - Q : depend upon the intrinsic dissipation of the material











B2-Force sensitivity (for cantilevers)

• Tip-surface interaction force:











B2-Force sensitivity (for cantilevers)

• « Jump to contact »:


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^{1.}Pérez et al., Phys Rev. B **58**, 10835 (1998)











Experimental illustration¹:



¹.Polesel et al., Nanotechnology **14**:1036 (2003)











Variational method based on the least action principle¹:

$$\delta S = \delta \left[\int_{0}^{1/f} L(z, \dot{z}, t) \right] = 0$$

$$\left[z(t), z(\dot{t}) \right] = T - V + W$$

$$= \frac{1}{2} \frac{k_c}{(2\pi f_0)^2} z(\dot{t})^2 - \left[\frac{1}{2} k_c z(t)^2 - z(t) k_c A_x \cos(2\pi f t) + V_c^{int}(z(t)) \right]$$

$$- \left[\frac{\Gamma_0 k_c}{(2\pi f_0)^2} z(t) \underline{z(t)} + V_d^{int}(z(t), \underline{z(t)}) \right]$$

Harmonic trial function: $z(t) = A \cos (2\pi f t + \varphi)$

$$\delta A \leqslant \cos(\varphi) = \frac{A}{A_x} \left(1 - \frac{f^2}{f_0^2} \right) + \frac{2}{k A_x} \langle F_c \rangle$$

$$\delta \varphi \leqslant \sin(\varphi) = -\frac{f}{f_0} \times \frac{A}{QA_x} - \frac{2}{k_c A_x} \langle F_d \rangle$$
Conservative & dissipative coupling terms

¹L.Nony et al. J.Chem.Phys.**111**,1610 (1999); "Analyse de la microscopie de force dynamique: application à l'étude de l'ADN", L.Nony, Thèse de l'univ. Bordeaux I (2000).

imenp

L









Coupling terms (Fourier components of the interaction force @ f):

$$\begin{cases} \langle \mathsf{F}_{\text{int}}^{c} \rangle = -\mathsf{f} \times \int_{0}^{1/\mathsf{f}} \frac{\partial}{\partial A} V_{\text{int}}^{c}(z(\mathsf{t})) d\mathsf{t} = \frac{1}{2\pi} \int_{0}^{2\pi} \mathsf{F}_{\text{int}}^{c}(\mathsf{D} - A\cos(\mathsf{u})) \cos(\mathsf{u}) d\mathsf{u} \\\\ \langle \mathsf{F}_{\text{int}}^{d} \rangle = -\frac{\mathsf{f}}{A} \int_{0}^{1/\mathsf{f}} \frac{\partial}{\partial \varphi} V_{\text{int}}^{d}(z(\mathsf{t}); \dot{z}(\mathsf{t})) d\mathsf{t} = \frac{1}{2\pi} \int_{0}^{2\pi} \mathsf{F}_{\text{int}}^{d}(\mathsf{D} - A\cos(\mathsf{u})) \sin(\mathsf{u}) d\mathsf{u} \\\\ \mathbf{V}_{\text{int}}^{c} = V_{\text{Lr}} + V_{\text{Sr}} + V_{\text{Elec}} + V_{\text{Magn}} \\\\ V_{VdW}(t) = -\frac{HR}{6(D - z(t))} \\\\ \mathsf{Long-range:} \\ \end{cases} V_{Vare } (t) = -U_0 \left[2e^{-\frac{D - z(t) - r_c}{\lambda}} - e^{-\frac{2(D - z(t) - r_c)}{\lambda}} \right] \\\\ \mathsf{Long-range:} \\ \end{cases}$$

Dispersive Van der Waals L.Nony *et al.* J.Chem.Phys.**111**,1610 (1999) Morse L.Nony *et al.*, Phys.Rev.B **74**, 235409 (2006)











B – Dynamic AFM

- 1. Dynamic AFM specificities
- 2. Force sensitivity
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- 4. Tapping mode AFM vs. non-contact AFM













B4- Tapping mode AFM vs. non-contact AFM:

5 variables : Excitation <u>amplitude</u>, <u>Excitation frequency</u>, <u>Oscillation amplitude</u>, <u>phase</u> (excitation / oscillation) and <u>tip-surface distance</u>

The variables that are kept constant set the mode :

Tapping (Amplitude Modulation-AFM)

Excitation amplitude and Excitation frequency are constant

 Measures : Oscillation amplitude and phase vs. tip-surface distance
 Probed interaction: most generally, attractive+repulsive interactions (intermittent contact)

non-contact AFM (nc-AFM, or Frequency Mod.-AFM¹)

Phase (=- $\pi/2$: resonance) **and** Oscillation amplitude are constant

• *Measures* : Resonance frequency shift and excitation amplitude required to keep the oscillation amplitude constant vs. tip-surface distance

 Probed interaction: most generally, attractive interaction (non-contact)

¹ T.Albrecht *et al.*, J.Appl.Phys.**69**, 668 (1991)













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C – Dynamic AFM in non-contact mode

- 1. Introduction
- 2. Connection with the modelization
- 3. Experimental illustrations
- 4. Instrumentation: PLL & amplitude regulation

5. Important results & current trends











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C1-Introduction

- Method developed¹ in 1991
- Concept:

Phase (=- $\pi/2$: resonance) and oscillation amplitude are constant

Measures : resonance frequency shift and excitation amplitude required to keep the oscillation amplitude constant (DISSIPATION image) vs. tip-surface distance

z feedback (TOPOGRAPHY image): performed to maintain $\Delta f = \tilde{f_0} - f_0$ constant (equivalent to tunneling current in STM)

¹ T.Albrecht *et al.*, J.Appl.Phys.**69**, 668 (1991)

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C2- Connection with the modelization (1):

Oscillation amplitude and phase (=-π/2) are fixed: Solving the system {cos(φ), sin(φ)} w.r.t. A and φ. Normalized frequency shift :

$$\frac{\Delta f}{f_0} = -\frac{1}{2\pi k_c A} \int_0^{2\pi} F_{\text{int}}^{\,\mathsf{c}}(D + A\cos(u))\cos(u)du$$

F.Giessibl (1997, 2000); U.Dürig (1999); L.Nony (1999, 2006)

Conservative interaction force (Attractive = non-contact):

$$\mathsf{F}_{\mathrm{int}}^{\mathrm{c}}(z_{-}) = -\frac{\partial}{\partial z_{-}} V_{\mathrm{int}}^{\mathrm{c}}(z_{-}) \text{ with } \begin{cases} V_{VdW}(t) = -\frac{HR}{6(D-z(t))} \\ V_{Morse}(t) = -U_0 \left[2e^{-\frac{D-z(t)-r_c}{\lambda}} - e^{-\frac{2(D-z(t)-r_c}{\lambda}} \right] \end{cases}$$

Hence the analytic expression for ∆f:

$$\frac{\Delta f}{f_0} = -\frac{1}{k_c A} \left[\frac{HRA}{6(D^2 - A^2)^{3/2}} + \frac{2U_0}{\lambda} (\Upsilon_{1,1/\lambda} - \Upsilon_{1,2/\lambda}) \right].$$
Long-range Short-range

C2- Connection with the modelization (1):



Exercice

 The calculation of the frequency shift can be simplified upon linearization of the interaction force by assuming small oscillation amplitudes (A~50 pm << D).

$$F_{\rm int}(t) = F_{\rm int}(s(t)) = F_{\rm int}(D - z(t)) \simeq \left. \frac{\partial F_{\rm int}(s)}{\partial s} \right|_{s_{\rm min} = D - A_0} s(t)$$

 In this case, reconsider the differential equation of motion of the cantilever and prove that the shifted resonance frequency then writes :

$$\widetilde{f}_0 = f_0 \sqrt{1 - \frac{1}{k} \left. \frac{\partial F_{\text{int}}}{\partial s} \right|_{s_{\min}}}$$

Deduce out of this the linearized formula for the frequency shift :

$$\frac{\widetilde{f}_0 - f_0}{f_0} = \frac{\Delta f}{f_0} = -\frac{1}{2k} \left. \frac{\partial F_{\text{int}}}{\partial s} \right|_{s_{\min}}$$













C2- Connection with the modelization (2):

Vertical sensitivity in the short-range regime (order of magnitude):

 $\delta z/\delta \Delta f \sim 10 \text{ pm} / \text{Hz}$

- The experimental control on Δf is better than 0,1 Hz (around $f_0 \sim 150$ kHz), i.e. 1 pm vertical resolution !
- This is what is achieved at low temperature (liquid Helium : 4K).
- At room temperature, the noise level of the instrument (z, amplitude detection...) exceeds that value. The resolution onset then is rather 10 pm.
- This sensitivity is enough to achieve atomic resolution at room temperature













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C3- Experimental illustrations



C3- Experimental illustrations

CaF₂



C. Barth *et al.* J. Phys.: Condens. Matter **13**, 2061(2001)

NaCl islands on Cu(111)



R. Bennewitz *et al.* Phys. Rev. B **13**, 2061 (2001)

KBr + e⁻



R. Bennewitz *et al.* Surf. Sci. **474**, L197 (2001)





C. Loppacher *et al.* Phys. Rev. B **62**, 16944 (2000)

NiO(100) @7K



R. Hoffmann *et al.* Phys. Rev. B **67**, 085402 (2003)

InAs(110) @14K



A. Schwarz *et al.* Phys. Rev. B **61**, 2837 (2000)

Si(111)7x7 @RT & 7K



M. Lantz et al. Phys. Rev. Lett. **84**, 2642 (2000)



F. Giessibl Science **267**, 68 (1995)











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- Phase Locked Loop (PLL):
- ✓ A PLL ensures the detection of Δf w.r.t. a reference frequency which is fixed equal to the **resonance frequency of the free cantilever** f_0 , therefore matching the phase condition $\varphi = -\pi/2$



✓ Frequency demodulation around f₀!

In this excitation scheme of the cantilever, there are no bifurcations structures as seen in AM-AFM. Qualitatively, the interpretation of the frequency shift can be made out of the linear approach, but it is to be reminded that the measured shift is the one that derives from the non-linear approach...











• Principle for the retention of the excitation at resonance:





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TITUTS

Principle for the retention of the oscillation amplitude :















Overview of the control electronics "PLL-excitation"¹ type



¹Ch. Loppacher Appl. Phys. A 66, 215 (1998); L. Nony *et al.*, Phys. Rev. B 74, 235439 (2006)













nc-AFM: keypoints

 Demodulation technic of the frequency of the cantilever around a central frequency which is the resonance frequency of the free cantilever:

$$\Delta \mathbf{f} = \widetilde{f}_0 - f_0$$

- The oscillation amplitude and the phase are kept constant : resonance values
- True non-contact method :

□ Attractive interaction regime with the surface

 $\Box \quad \Delta f < 0$

• Very large z sensitivity upon accurate detection of Δf , which allows for large vertical and lateral contrasts (atomic resolution)











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C5- Important results : high-resolution imaging of surfaces

CaF₂



C. Barth et al. J. Phys.: Condens. Matter 13, 2061(2001)

NiO(100) @7K





R. Bennewitz et al. Phys. Rev. B 13, 2061 (2001)

$KBr + e^{-}$



R. Bennewitz et al. Surf. Sci. 474, L197 (2001)





C. Loppacher et al. Phys. Rev. B 62, 16944 (2000)



R. Hoffmann et al. Phys. Rev. B 67, 085402 (2003)

InAs(110) @14K an irection



A. Schwarz et al. Phys. Rev. B 61, 2837 (2000)



M. Lantz et al. Phys. Rev. Lett. 84, 2642 (2000)



F. Giessibl Science 267, 68 (1995)



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1nm





Si(111)7x7 @RT & 7K



C5- Important results : high-resolution imaging of molecular assemblies adsorbed on surfaces



R. Pawlak *et al.* J. Phys. Chem. C **114**, 9290–9295 (2010)



S. Maier *et al.* Small **4**, 1115–1118 (2008)



A. Amrous *et al.* Adv. Mater. Interfaces 1400414 (2014)



S. Kawai *et al.* Phys. Rev. B **84**, 085429 (2011)













C5- Important results : chemical identification at the atomic scale

Atomic resolution imaging and force versus distance measurements on KBr (0 0 1) using low temperature scanning force microscopy

R. Hoffmann^{*}, M.A. Lantz¹, H.J. Hug, P.J.A. van Schendel, P. Kappenberger, S. Martin, A. Baratoff, H.-J. Güntherodt

Applied Surface Science 188 (2002) 238-244



Chemical identification of individual surface atoms by atomic force microscopy

Yoshiaki Sugimoto¹, Pablo Pou², Masayuki Abe^{1,3}, Pavel Jelinek⁴, Rubén Pérez², Seizo Morita¹ & Óscar Custance¹

VOLUME 92, NUMBER 14

PHYSICAL REVIEW LETTERS

week ending 9 APRIL 2004

Sublattice Identification in Scanning Force Microscopy on Alkali Halide Surfaces

R. Hoffmann,^{1,2,*} L. N. Kantorovich,³ A. Baratoff,¹ H. J. Hug,^{1,4} and H.-J. Güntherodt¹ ¹National Center of Competence in Research on Nanoscale Science, Institute of Physics, University of Basel, Klingelbergstrasse 82, CH-4056 Basel, Switzerland ²Physikalisches Institut, Universität Karlsruhe, 76128 Karlsruhe, Germany ³Department of Physics, School of Physical Sciences and Engineering, King's College London, London WC2R 2LS, United Kingdom ⁴Swiss Federal Laboratory for Materials Testing and Research, CH-8600 Dübendorf, Switzerland (Received 17 March 2003; published 9 April 2004)



Vol 446 |1 March 2007 | doi:10.1038/nature05530

Y. Sugimoto *et al.* Nature **446**, 64 (2007)













C5- Important results : manipulations of atoms, molecules and nano-objects

week ending PHYSICAL REVIEW LETTERS VOLUME 90, NUMBER 17 2 MAY 2003 Mechanical Vertical Manipulation of Selected Single Atoms by Soft Nanoindentation Using Near Contact Atomic Force Microscopy Noriaki Oyabu,¹ Óscar Custance,^{2,*} Insook Yi,² Yasuhiro Sugawara,^{1,2} and Seizo Morita^{1,2} ¹Graduate School of Engineering, Osaka University, 2-1 Yamada-Oka, Suita, 565-0871, Japan 4K ²Handai Frontier Research Center (FRC), 2-1 Yamada-Oka, Suita, 565-0871, Japan (Received 27 December 2002; published 2 May 2003) week ending PHYSICAL REVIEW LETTERS PRL 98, 106104 (2007) 9 MARCH 2007 Mechanism for Room-Temperature Single-Atom Lateral Manipulations on Semiconductors RT using Dynamic Force Microscopy

Yoshiaki Sugimoto,¹ Pavel Jelinek,² Pablo Pou,³ Masayuki Abe,^{1,4} Seizo Morita,¹ Ruben Perez,³ and Oscar Custance¹

VOLUME 90, NUMBER 6

PHYSICAL REVIEW LETTERS

week ending 14 FEBRUARY 2003

Direct Determination of the Energy Required to Operate a Single Molecule Switch

Ch. Loppacher,^{1,*} M. Guggisberg,¹ O. Pfeiffer,¹ E. Meyer,¹ M. Bammerlin,² R. Lüthi,² R. Schlittler,² J. K. Gimzewski,^{2,†} H. Tang,³ and C. Joachim³
 ¹Institute of Physics, University of Basel, Klingelbergstrasse 82, CH-4056 Basel, Switzerland
 ²IBM Research Division, Zurich Research Laboratory, CH-8803 Rüschlikon, Switzerland
 ³CEMES/CNRS, 29, rue Jeanne Marvig, 31055-Toulouse Cedex, France (Received 11 October 2002; published 14 February 2003)













C5- Current trends :

- Small amplitudes (50 pm, sensitivity to the SR-regime) : piezoelectric probes (tuning fork, "qPlus" Omicron, "Colibri" Specs). Usually used at low temperature:
 - \rightarrow imaging + spectroscopy: quantitative interpretation (forces, charges...)
- Time resolution: high-frequency cantilevers or higher harmonics...
- High-resolution in Kelvin Probe Force Microscopy (KPFM): quantitative interpretation of the contrast...
- Time-resolved KPFM by pump-probe methods...













C5- Current trends: quartz tuning fork¹



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C5- Current trends: quartz tuning fork

Measuring the Charge State of an Adatom with Noncontact Atomic Force Microscopy

Leo Gross,¹* Fabian Mohn,¹ Peter Liljeroth,^{1,2} Jascha Repp,^{1,3} Franz J. Giessibl,³ Gerhard Meyer¹ 12 JUNE 2009 VOL 324 **SCIENCE**



Topographic sensitivity to the charge state on the atomic scale





Spectroscopic sensitivity to the charge state on the atomic scale

















C5- Current trends: quartz tuning fork

The Chemical Structure of a Molecule Resolved by Atomic Force Microscopy

Leo Gross,¹* Fabian Mohn,¹ Nikolaj Moll,¹ Peter Liljeroth,^{1,2} Gerhard Meyer¹

28 AUGUST 2009 VOL 325 SCIENCE

Pentacene / Cu(111)



Pentacene / 2ML NaCl(001) / Cu(111)



Organic structure determination using atomicresolution scanning probe microscopy

Leo Gross^{1*}, Fabian Mohn¹, Nikolaj Moll¹, Gerhard Meyer¹, Rainer Ebel², Wael M. Abdel-Mageed^{2,3} and Marcel Jaspars²















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Thanks for your attention



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