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1	Development of reflective back contacts for high-efficiency ultrathin Cu(In,Ga)Se ₂ solar
2	cells
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14 Abstract

Because of poor light absorption, Cu(In,Ga)Se₂-based (CIGS) solar cells with an ultrathin 15 absorber layer (<500 nm) require the development of reflective back contacts. To enhance 16 17 rear reflectance in CIGS ultrathin devices, we investigate novel back contact architectures based on a silver metallic mirror covered with a thin layer of In₂O₃:Sn (ITO), which is fully 18 compatible with nanopatterning for further light trapping improvements. First, numerical 19 electromagnetic simulations of complete solar cells have been performed for a 490 nm thick 20 CIGS absorber with various back contacts. We predict a short-circuit current density of J_{SC} = 21 34.0 mA/cm² for a 490 nm thick CIGS absorber with a silver nanostructured mirror. Second, 22 23 we have fabricated and characterized 490 nm thick CIGS solar cells with transparent back contacts made of ITO, and reflective back contacts made of silver covered with ITO. Solar 24 cells with a transparent ITO back contact exhibit an average efficiency of 10.0 %, compared 25 to 9.3 % for standard molybdenum back contacts. A 5 nm thick Ga₂O₃ layer is revealed at the 26 ITO/CIGS interface by transmission electron microscopy and energy dispersive X-ray 27 spectroscopy. When silver is added, the reflective back mirror leads to a J_{SC} improvement of 28 4.6 mA/cm² (from 22.4 to 27.0 mA/cm²). These results pave the way for efficient ultrathin 29 CIGS solar cells on reflective back contacts. 30

31 Keywords

Solar cells, Copper indium gallium selenide, Ultrathin films, Transparent back contact,
Reflective back contact, Nanostructured back mirror

34 1. Introduction

35 Cu(In,Ga)Se₂-based (CIGS) solar cells are one of the most promising thin-film technologies, 36 with a record efficiency of 22.9 % achieved with a 2–3 μm thick absorber [1], [2]. However, 37 the scarcity and high cost of Indium are a drawback for industrial production of competitive 38 modules. With a CIGS absorber thickness lower than 500 nm it is possible to reduce 39 deposition time and materials consumption, resulting in a decreased manufacturing cost [3], 40 [4].

41 Ultrathin CIGS-based solar cells exhibit lower short-circuit current densities (J_{SC}) and 42 efficiencies mainly due to lower light absorption, enhanced recombination [5] and low 43 reflectivity of the conventional Mo back contact [6], [7]. These loss mechanisms can be 44 overcome by substituting the Mo back contact with passivating and more reflective back

contact. Up to now, there are few studies of reflective back contacts for CIGS solar cells, such 45 as: direct use of reflective or metallic back contacts [8]–[10], a combination of transparent 46 conducting oxide and metallic reflector [7], passivating reflective back contact [11], or 47 nanostructured back contacts [12], [13]. Nanostructured back reflectors are a promising 48 strategy for enhanced absorption in ultrathin CIGS layers, as they lead to calculated J_{SC} up to 49 36.3 mA/cm² in a 150 nm thick CIGS layer compared to 23.5 mA/cm² with a standard Mo 50 back contact [12]. Nanostructured mirrors have also yielded significant J_{SC} improvements in 51 52 other types of solar cells such as amorphous Si:H [14] and ultrathin GaAs [15], [16].

In this paper we present the first steps towards the development of a reflective back contact 53 54 based on a silver metallic mirror encapsulated with a thin layer of In₂O₃:Sn (ITO). Reflective back contacts for CIGS solar cells need not only to provide enhanced rear reflectance but also 55 56 to form an ohmic contact with the absorber. Several studies report on the formation of a resistive interfacial layer of Ga₂O₃ after CIGS deposition on oxide layers such as hydrogen-57 doped In₂O₃, ITO, MoO₃, SnO₂:F, ZnO or ZnO:Al [17]–[22]. In particular, ITO is known to 58 prevent Na diffusion from the glass to the CIGS absorber and to form a detrimental layer of 59 Ga₂O₃ for CIGS deposition temperatures above 520 °C [22], [23]. As a result, the stability and 60 electrical properties of ultrathin devices on transparent ITO-based substrates were first 61 investigated by co-evaporation of CIGS at low temperature (450°C) with incorporation of Na 62 from a NaF post-deposition treatment. 63

64 In this study, the improvement of light absorption in a 490 nm thick CIGS absorber is first simulated using a nanostructured and a flat mirror encapsulated in transparent conducting 65 66 oxides such as ITO. Then, the optical properties of glass substrates covered with Mo, Ag and Ag/ITO are determined and compared. 490 nm thick CIGS solar cells are fabricated on ITO 67 68 back contacts, and the ITO/CIGS interface is studied in-depth with transmission electron microscopy (TEM) and energy dispersive X-ray spectroscopy (EDX). The performances of 69 70 cells on ITO and Mo back contacts are compared and discussed, and preliminary results on Ag/ITO back contacts are presented. 71

72 2. Experimental methods

CIGS solar cells were fabricated on soda-lime glass (SLG) substrates with three different back contacts. The first stack is made of a conventional 800 nm thick molybdenum (Mo) back contact deposited on a 300 nm Al₂O₃ layer that acts as a diffusion barrier for alkali elements present in SLG. Al₂O₃ is deposited by atomic layer deposition. The second stack is based on a

transparent 300 nm thick ITO layer sputtered on SLG. The third stack contains a 30 nm thick 77 ITO sputtered on a 150 nm thick Ag layer deposited by electron beam evaporation. A 490 nm 78 thick CIGS layer was deposited by thermal co-evaporation in a one-stage process, at a fixed 79 substrate temperature of 450 °C. This substrate temperature was calibrated in a previous 80 work, using an infrared camera [24]. Average composition of absorbers was determined from 81 X-Ray Fluorescence (XRF) signal of CIGS deposited on Mo, leading to atomic ratios of 82 $[Cu]/([In]+[Ga]) = 0.81 \pm 0.01$ (CGI) and $[Ga]/([Ga]+[In]) = 0.32 \pm 0.01$ (GGI). A NaF post-83 deposition treatment was applied with an evaporation rate of 1 nm/min under Se flux and at 84 fixed substrate temperature of 350 °C. Solar cells were completed with the standard front 85 layer stack consisting in a chemical bath deposited CdS/rf-sputtered ZnO/rf-sputtered ZnO:Al 86 87 with respective thicknesses 50nm/50nm/400nm determined by XRF analysis on control samples. 88

Opto-electrical properties of ITO were determined from a 200 nm thick ITO layer sputtered 89 on SLG. As annealing of ITO improves its optical transparency and carrier mobility, its 90 optical indices and resistivity were measured before and after a 10-minute-annealing 91 performed in air on a hotplate set at 540 °C. Refractive index n and extinction coefficient k92 were extracted from ellipsometric measurements on a UVISEL ellipsometer from Horiba 93 Jobin-Yvon. The ITO resistivity was measured with a four-point probe system before ($\rho =$ 94 9.0x10⁻⁴ Ω .cm) and after annealing ($\rho = 7.0x10^{-4} \Omega$.cm). Finally, a Sentech reflectometer was 95 used to determine the spectral reflectance of four different substrates: SLG/600-nm-Mo, 96 SLG/150-nm-Ag, SLG/150-nm-Ag/30-nm-ITO and SLG/150-nm-Ag/200-nm-ITO substrates. 97

Current-voltage I(V) measurements of 0.1 cm² solar cells were carried out with a Newport 98 class AAA solar simulator under AM1.5G illumination and in the dark. Devices on 99 transparent substrates were placed on a black paper before light I(V) measurements to avoid 100 additional rear-reflection from the substrate holder. External Quantum Efficiency (EQE) 101 measurements were performed with a IQE200 Newport Instrument. A slab of a 490 nm thick 102 CIGS layer deposited on an ITO-based transparent back contact was prepared by a focused 103 ion beam, and analysed with TEM in a FEI Titan Themis XFEG instrument at an acceleration 104 voltage of 200 kV. Composition mappings were determined by EDX. 105

106 **3.** Results and discussion

107 3.1 Absorption simulation

We first investigate numerically the optical effects of a nanostructured back mirror, a flat back 108 mirror, and a transparent back contact as compared to a conventional Mo back contact. The 109 spectral absorption is simulated with a rigorous coupled wave analysis method, and the 110 optical indices are calculated from ellipsometric measurements for CIGS and annealed ITO, 111 and taken from the literature for other materials. More details can be found in reference [12]. 112 The architecture of flat ultrathin devices is sketched in Fig. 1a), while the optimized geometry 113 of a complete device including a nanostructured back mirror is presented in Fig. 1b) and c). 114 This simulated nanostructured back mirror consists of square-shaped dielectric nanogrid on 115 top of a flat ITO-coated Ag stack. The dielectric nanogrid has a refractive index of n = 1.5, a 116 height of 200 nm, a period of 600 nm and hole width of 350 nm. The experimental fabrication 117 of such a dielectric nanogrid is under development, using a scalable process based on direct 118 nanoimprint of a sol-gel derived metal oxide [25], [26]. The simulated absorptions of CIGS 119 120 devices on Mo, ITO, as well as flat and nanostructured reflective back contacts are plotted in Fig. 2. 121

As shown in Fig. 2a, using a standard Mo back contact leads to substantial absorption losses 122 for long wavelengths (> 600 nm), because of the low reflectivity at CIGS/Mo interface. 123 Moreover, optical losses are also observed at short wavelengths (< 500 nm) due to the low 124 bandgap of CdS (2.4 eV) buffer layer. Assuming a perfect collection of photogenerated 125 carriers, a theoretical short-circuit current density J_{SC} can be deduced from the simulated 126 CIGS absorption [12]. It is limited to 28.8 mA/cm² with a Mo back contact. When Mo is 127 replaced by ITO (Fig. 2b) a similar spectral absorption with a J_{SC} of 28.5 mA/cm² is 128 calculated. The introduction of an Ag layer results in enhanced back reflectance and 129 significant improvement of CIGS light absorption at long wavelengths (> 800 nm), (Fig. 2c), 130 which leads to an increased J_{SC} of 32.2 mA/cm². Finally, if the flat mirror is replaced by an 131 optimised nanostructured back mirror made of a nanogrid, additional absorption resonances 132 occur close to the bandgap (Fig. 2d). Multi-resonant absorption in the thin CIGS layer leads to 133 a J_{SC} of 34.0 mA/cm², which represents an additional 1.8 mA/cm² gain as compared to the flat 134 back mirror. 135

136 3.2 Optical characterization of back contacts

Before deposition of CIGS on transparent and reflective back contacts, the spectral reflectance 137 of 4 different back contacts deposited on SLG were first measured and compared: a 600 nm 138 139 thick Mo, a 150 nm thick Ag layer, and 150 nm thick Ag layers with 30 nm and 200 nm ITO coatings (Fig. 3a). As expected, Mo has a low reflectivity and cannot act as a back mirror, 140 resulting in low CIGS absorption for ultrathin CIGS solar cells (Fig. 2a). On the contrary, the 141 silver layer provides a high reflectivity (R>90 % for λ >500 nm), but it cannot be used as a 142 back contact due to diffusion of Ag in CIGS during the growth. When the silver layer is 143 covered with ITO coatings, the reflectivity is significantly reduced due to parasitic absorption. 144 Reflectivity dips are attributed to Fabry-Perot resonances in the ITO layer. After a 10-minute-145 annealing in air at 540 °C, the reflectivity is strongly enhanced and exceeds 90 % for 146 wavelengths above 550 nm. This annealing experiment performed in air proves that the ITO 147 layer efficiently encapsulates the Ag mirror, preventing its oxidation and morphological 148 changes that negatively impact its reflectivity; the stability of such reflective back contact 149 under CIGS deposition was only tested after the complete solar cell fabrication. To determine 150 the origin of the enhanced reflectivity of Ag/ITO stacks after annealing, ellipsometric 151 measurements were carried out on a 200 nm thick ITO layer sputtered on SLG before and 152 after annealing in air at 540 °C during 10 minutes. The refractive index n and extinction 153 coefficient k are displayed in Fig. 3b). The annealing results in an increase of n and a decrease 154 of k in the visible and infrared domains (λ >500 nm). The origin of these effects was not 155 investigated in detail, though it is commonly accepted that annealing of amorphous ITO 156 generally results in an increased layer density, an improved carrier mobility and a reduced 157 158 free-carrier absorption. It leads to a significant improvement of the Ag/ITO reflectivity and a spectral shift of the reflectivity dips toward shorter wavelengths (Fig. 3a). 159

160 3.3 TEM-EDX back contact characterization

A TEM-EDX analysis was performed in order to investigate the ITO/CIGS interface. Fig. 4a) and b) show a TEM image of an ultrathin 490-nm-CIGS layer deposited on an ITO back contact along with the average composition profile of absorber elements deduced from EDX analysis. An average GGI of 0.28 was calculated from the average EDX signal of the CIGS layer, which is close to the GGI value of 0.32 determined by XRF. It is worth mentioning that the EDX signal of Cu is overestimated as Cu from the substrate holder is also detected. Hence CGI cannot be deduced from this analysis. 168 This EDX analysis confirms that the CIGS composition is homogeneous from the front to the 169 back interface, as expected from a one-stage CIGS deposition process. Moreover it reveals 170 that Ga segregates at the back contact.

171 Closer views are shown in Fig. 4. with (c) an additional dark field TEM image and its 172 corresponding compositional EDX mappings of O, Ga, In, and Se elements. They reveal a 173 segregation of both Ga and O elements at the ITO/CIGS interface, which suggests that a thin 174 Ga₂O₃ layer is formed during CIGS deposition. This layer should be highly resistive and 175 possibly n-doped [17]. However, in this study the low substrate temperature (450 °C) during 176 CIGS deposition results in a very thin Ga₂O₃ layer of only 5 nm that does not seem to 177 deteriorate solar cell performances.

178 3.4 Solar cell performances

The performances of ultrathin CIGS solar cells are summarized in Table 1 for Mo and ITO 179 back contacts. Best efficiencies are given together with average I(V) parameters and standard 180 181 deviation for the 10 best cells. Dark I(V) parameters are extracted from the fit of a 2-diode model: saturation currents for diodes with ideality factors of 1 (J₀₁) and 2 (J₀₂), shunt 182 183 resistance (R_{SH}) and series resistance (R_S). I(V) and EQE curves of best cells are displayed on Fig. 5, together with the experimental EQE of a 490 nm thick CIGS cell on a reflective back 184 contact (Ag). The reference solar cells on a SLG/Al₂O₃/Mo substrate exhibit an average 185 186 efficiency of 9.3 %.

Ultrathin CIGS solar cells with an ITO back contact present an average efficiency of 10.0 %. 187 Mo and ITO-based substrates lead to very close dark currents (J₀₁, J₀₂), R_{SH} and fill factors 188 (FF). A slight increase of R_s is observed with an ITO back contact, from 0.1 for Mo to 1.2 189 $\Omega.cm^2$. Importantly, the increase of both open-circuit voltage and J_{SC} results in an absolute 190 efficiency enhancement of 0.7 %. It indicates that ITO has suitable electrical properties to be 191 used as a back contact in CIGS solar cells. Ultrathin CIGS devices on transparent SLG/ITO 192 193 substrates could be used for bifacial solar cells. Alternatively, solar cell efficiency could be enhanced by adding a mirror on the backside of the glass substrate. 194

According to simulation results (Fig. 2), the short-circuit current could also be increased by adding a flat silver back mirror between the ITO layer and the glass substrate, with an expected gain of 4.6 mA/cm² over the Mo reference. We have performed preliminary experiments of 490 nm thick CIGS solar cells deposited on reflective multilayer stacks (MLS) based on silver covered by ITO. They led to poor electrical performances due to process issues that should be overcome by further optimizations. However, the reflective MLS resulted in a strong EQE enhancement. The EQE of ultrathin solar cells with a Mo back contact, an ITO back contact, and a reflective MLS are plotted and compared in Fig. 5b. The reflective back contact leads to an improvement for wavelengths above 500 nm and results in a promising $J_{SC} = 27.0 \text{ mA/cm}^2$. The same EQE measurements are also plotted and compared to numerical simulation in Fig. 2. A very good agreement between experiments and simulations is obtained.

207 **5.** Conclusion

Ultrathin CIGS solar cells with transparent and reflective back contacts were first numerically 208 investigated. A J_{SC} of 32.2 mA/cm² is calculated for 490 nm thick CIGS solar cells with a flat 209 back mirror, corresponding to a 3.4 mA/cm² gain as compared to standard molybdenum back 210 contacts. An additional optical gain is predicted with a nanostructured reflective back contact, 211 leading to a J_{SC} of 34.0 mA/cm². Flat reflective stacks of Ag/ITO were fabricated and are 212 stable under a 10-minute-annealing in air at 540 °C. Their spectral reflectance is over 90 % 213 for wavelengths above 550 nm, which should substantially increase absorption in ultrathin 214 CIGS absorbers. Complete ultrathin CIGS solar cells were fabricated on molybdenum and 215 216 transparent ITO back contacts, with respective average efficiencies of 9.3 % and 10.0 %. ITO forms an ohmic back contact with CIGS, as confirmed by an average FF of 70.3 %. A TEM-217 EDX investigation shows that the thickness of the Ga₂O₃ formed at ITO/CIGS interface 218 during CIGS deposition is limited to 5 nm and has no impact on solar cell performances. 219 Preliminary experiments on reflective back contacts made of Ag and ITO led to a $J_{SC} = 27.0$ 220 mA/cm². These results provide a promising route toward efficient, ultrathin CIGS solar cells 221 222 deposited on reflective back contacts.

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- 311

312 List of figures

- Fig. 1 Sketch of an ultrathin CIGS solar cell a) on planar back contacts such as molybdenum, TCO, or reflective multi-layer stack (MLS) and b) on a nanostructured back mirror consisting of a periodically patterned dielectric on top of reflective MLS. c) Three dimensional view of the ultrathin device on a nanostructured back mirror with optimized geometry. CIGS is deposited on top of a square-shaped dielectric nanogrid with a height of 200 nm, a period of 600 nm and hole width of 350 nm.
- Fig. 2 Simulated absorption within each layer of 490 nm thick CIGS solar cells under
 AM1.5G illumination, with back contacts made of a) molybdenum, b) ITO, c)
 reflective multi-layer stack (MLS) and d) reflective MLS with nanopatterns.
 Respective experimental EQEs are given in a), b) and c) for comparison.

- Fig. 3. a) Reflectance of 600 nm thick molybdenum (grey), 150 nm thick silver (black) and stacks of 30 nm thick (red) and 200 nm thick (blue) ITO on silver before and after annealing in air at 540 °C for 10 minutes (dashed and solid lines, respectively). b) Refractive index n (solid lines) and extinction coefficient k (dashed lines) of a 200 nm thick sputtered ITO layer measured by ellipsometry before (black) and after (red) annealing.
- Fig. 4 a) TEM dark field image of ultrathin CIGS layer deposited on ITO back contact
 with b) corresponding composition profile from average EDX signal, c) TEM dark
 field image of ultrathin CIGS layer close to the ITO back contact with associated
 EDX mappings of O, Ga, In, and Se elements.
- Fig. 5 Experimental I(V) and External Quantum Efficiency (EQE) curves of 490 nm
 thick CIGS solar cells on molybdenum (black), ITO (red) and reflective multilayer stack (blue) back contacts.

337 Table captions

338Table 1.Summary of I(V) parameters for 490 nm thick CIGS solar cells with various back339contacts. A two-diode model is used to fit dark I(V) curves (J_{01} and J_{02} : saturation340currents for respective ideality factors of 1 and 2, R_{SH} : shunt resistance, R_S : series341resistance).





















